

# Chiral symmetry breaking in continuum QCD

Mario Mitter

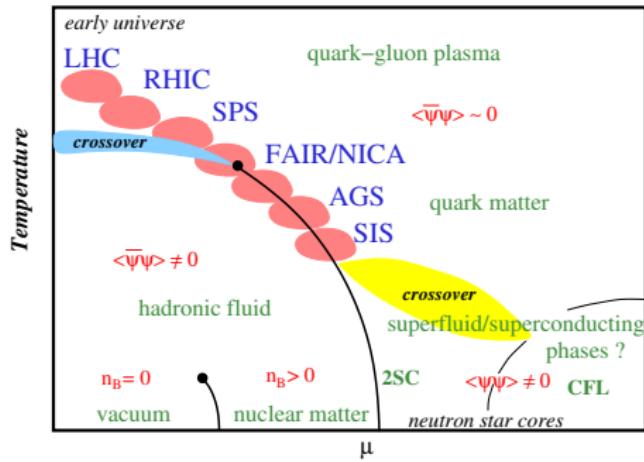
Ruprecht-Karls-Universität Heidelberg

Trieste, September 2016



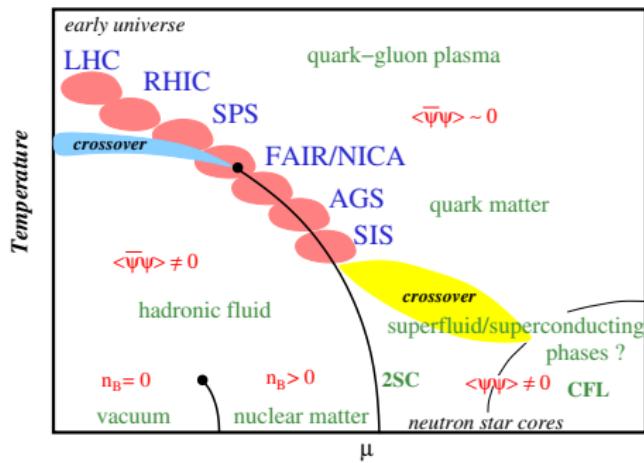
# fQCD collaboration - QCD (phase diagram) with FRG:

J. Braun, L. Corell, A. K. Cyrol, L. Fister, W. J. Fu, M. Leonhardt, MM,  
J. M. Pawłowski, M. Pospiech, F. Rennecke, N. Strodthoff, N. Wink, ...



# fQCD collaboration - QCD (phase diagram) with FRG:

J. Braun, L. Corell, A. K. Cyrol, L. Fister, W. J. Fu, M. Leonhardt, MM,  
J. M. Pawłowski, M. Pospiech, F. Rennecke, N. Strodthoff, N. Wink, ...



large part of this effort: vacuum YM-theory and QCD

# QCD with the FRG

- $S[\Phi] = \Gamma_\Lambda[\Phi]$ : use only perturbative QCD input
  - ▶  $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$
  - ▶  $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$

# QCD with the FRG

- $S[\Phi] = \Gamma_\Lambda[\Phi]$ : use only perturbative QCD input
  - ▶  $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$
  - ▶  $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
- gauge-fixed approach (Landau gauge): ghosts appear

$$\partial_k \Gamma_k[A, \bar{c}, c, \bar{q}, q] = \frac{1}{2}$$



# QCD with the FRG

- $S[\Phi] = \Gamma_\Lambda[\Phi]$ : use only perturbative QCD input
  - ▶  $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$
  - ▶  $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
- gauge-fixed approach (Landau gauge): ghosts appear

$$\partial_k \Gamma_k[A, \bar{c}, c, \bar{q}, q] = \frac{1}{2}$$



- vertex expansion:

$$\begin{aligned} \Gamma[\Phi] = & \\ \sum_n \int_{p_1, \dots, p_{n-1}} & \Gamma_{\phi_1 \dots \phi_n}^{(n)}(p_1, \dots, p_{n-1}) \phi^1(p_1) \dots \phi^n(-p_1 - \dots - p_{n-1}) \end{aligned}$$

# QCD with the FRG

- $S[\Phi] = \Gamma_\Lambda[\Phi]$ : use only perturbative QCD input
  - ▶  $\alpha_s(\Lambda = \mathcal{O}(10) \text{ GeV})$
  - ▶  $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
- gauge-fixed approach (Landau gauge): ghosts appear

$$\partial_k \Gamma_k[A, \bar{c}, c, \bar{q}, q] = \frac{1}{2}$$



- vertex expansion:

$$\begin{aligned} \Gamma[\Phi] &= \\ &\sum_n \int_{p_1, \dots, p_{n-1}} \Gamma_{\phi_1 \dots \phi_n}^{(n)}(p_1, \dots, p_{n-1}) \phi^1(p_1) \dots \phi^n(-p_1 - \dots - p_{n-1}) \end{aligned}$$

- aim for “apparent convergence” of  $\Gamma[\Phi] = \lim_{k \rightarrow 0} \Gamma_k[\Phi]$

# Landau gauge QCD

- two crucial phenomena:  $S\chi$ SB and confinement
- similar scales - hard to disentangle see e.g. [Williams, Fischer, Heupel, 2015]
- quenched QCD: allows separate investigation:

# Landau gauge QCD

- two crucial phenomena:  $S\chi$ SB and confinement
- similar scales - hard to disentangle see e.g. [Williams, Fischer, Heupel, 2015]
- quenched QCD: allows separate investigation:
  - quenched matter part [MM, Strodthoff, Pawłowski, 2014]
  - pure YM-theory [Cyrol, Fister, MM, Pawłowski, Strodthoff, 2016]
  - outlook:
    - ▶ unquenching
    - ▶ YM-theory at finite temperature  $T > 0$

# Landau gauge QCD

- two crucial phenomena:  $S\chi$ SB and confinement
- similar scales - hard to disentangle see e.g. [Williams, Fischer, Heupel, 2015]
- quenched QCD: allows separate investigation:
  - quenched matter part [MM, Strodthoff, Pawłowski, 2014]
  - pure YM-theory [Cyrol, Fister, MM, Pawłowski, Strodthoff, 2016]
  - outlook:
    - ▶ unquenching
    - ▶ YM-theory at finite temperature  $T > 0$
- use results from lattice QCD to gauge truncation

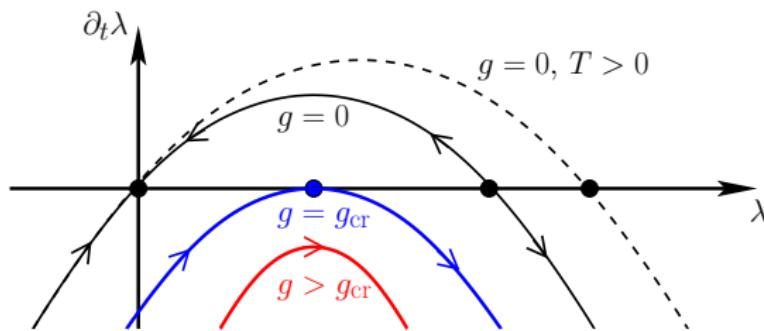
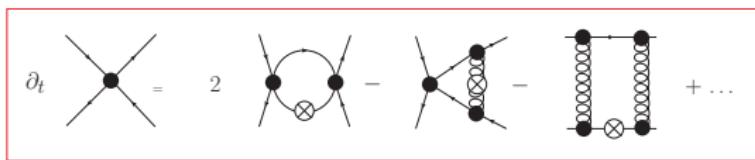
# Chiral symmetry breaking

- $\chi$ SB  $\Leftrightarrow$  resonance in 4-Fermi interaction  $\lambda$  (pion pole):

# Chiral symmetry breaking

- $\chi$ SB  $\Leftrightarrow$  resonance in 4-Fermi interaction  $\lambda$  (pion pole):
- resonance  $\Rightarrow$  singularity without momentum dependency

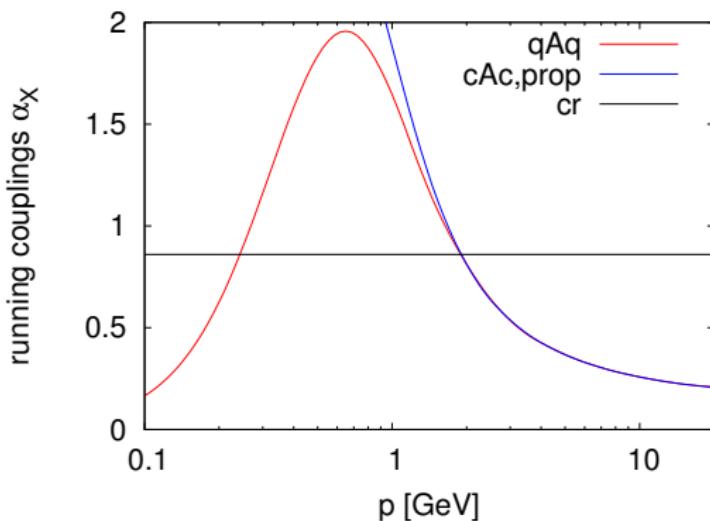
$$\partial_t \lambda = a \lambda^2 + b \lambda \alpha + c \alpha^2, \quad b > 0, \quad a, c \leq 0$$



[Braun, 2011]

# (transverse) running couplings

[Cyrol, MM, Pawłowski, Strodthoff, in preparation]



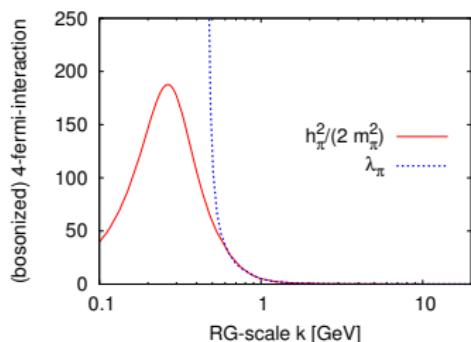
- agreement in perturbative regime required by gauge symmetry
- non-degenerate in nonperturbative regime: reflects gluon mass gap
- $\alpha_{\bar{q}Aq} > \alpha_{cr}$ : necessary for chiral symmetry breaking
- area above  $\alpha_{cr}$  very sensitive to errors

# 4-Fermi vertex via dynamical hadronization

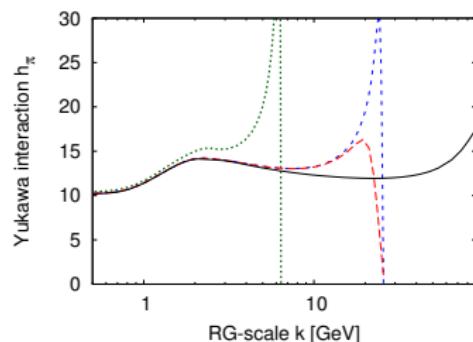
[Gies, Wetterich, 2002]

- change of variables: particular 4-Fermi channels  $\rightarrow$  meson exchange
- efficient inclusion of momentum dependence  $\Rightarrow$  no singularities
- calculation of model parameters from QCD

$$\partial_k \Gamma_k = \frac{1}{2} \left( \text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} + \frac{1}{2} \text{Diagram 4} \right)$$



[MM, Strodthoff, Pawlowski, 2014]



[Braun, Fister, Haas, Pawlowski, Rennecke, 2014]

[MM, Strodthoff, Pawlowski, 2014]

## Vertex Expansion

[MM, Strodtboff, Pawlowski, 2014],

[Cyrol, Fister, MM, Strodtboff, Pawlowski, 2016]

$$\partial_t \quad \text{---}^{-1} = \quad + \quad + \frac{1}{z} \quad -$$

+

$$\partial_t = - \left( \text{sum of diagrams} \right) - \frac{1}{2} \text{ (diagram)} + 2 \text{ (diagram)} - \text{ (perm.)}$$

$$\partial_t = -2 \left( \text{Diagram 1} \right) - \left( \text{Diagram 2} \right) - \left( \text{Diagram 3} \right) - \left( \text{Diagram 4} \right) - \left( \text{Diagram 5} \right) - \left( \text{Diagram 6} \right) - \left( \text{Diagram 7} \right) + \text{perm.}$$

# Vertex Expansion

[MM, Strodthoff, Pawłowski, 2014],

[Cyrol, Fister, MM, Strodthoff, Pawłowski, 2016]

$$\partial_t \text{---}^{-1} = \begin{array}{c} \text{Diagram 1} \\ + \end{array} \begin{array}{c} \text{Diagram 2} \\ + \frac{1}{2} \end{array} \begin{array}{c} \text{Diagram 3} \\ - \end{array} \begin{array}{c} \text{Diagram 4} \\ + \end{array} \begin{array}{c} \text{Diagram 5} \\ - \end{array} \begin{array}{c} \text{Diagram 6} \\ - \end{array}$$

$$\partial_t \text{---}^{-1} = \begin{array}{c} \text{Diagram 1} \\ + \end{array}$$

$$\partial_t \text{---}^{-1} = \begin{array}{c} \text{Diagram 1} \\ - 2 \end{array} \begin{array}{c} \text{Diagram 2} \\ + \frac{1}{2} \end{array}$$

$$\partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - \end{array} \begin{array}{c} \text{Diagram 2} \\ - \end{array} \begin{array}{c} \text{Diagram 3} \\ + \text{perm.} \end{array}$$

$$\begin{array}{c} \partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - \end{array} \begin{array}{c} \text{Diagram 2} \\ - \end{array} \begin{array}{c} \text{Diagram 3} \\ - \end{array} \begin{array}{c} \text{Diagram 4} \\ - \end{array} \begin{array}{c} \text{Diagram 5} \\ - \end{array} \begin{array}{c} \text{Diagram 6} \\ + 2 \end{array} \begin{array}{c} \text{Diagram 7} \\ - \end{array} \begin{array}{c} \text{Diagram 8} \\ + \text{perm.} \end{array} \end{array}$$

$$\partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - \end{array} \begin{array}{c} \text{Diagram 2} \\ + 2 \end{array} \begin{array}{c} \text{Diagram 3} \\ - \end{array} \begin{array}{c} \text{Diagram 4} \\ + \text{perm.} \end{array}$$

$$\partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - \end{array} \begin{array}{c} \text{Diagram 2} \\ - \end{array} \begin{array}{c} \text{Diagram 3} \\ + 2 \end{array} \begin{array}{c} \text{Diagram 4} \\ - \end{array} \begin{array}{c} \text{Diagram 5} \\ + \text{perm.} \end{array}$$

$$\begin{array}{c} \partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - 2 \end{array} \begin{array}{c} \text{Diagram 2} \\ - \end{array} \begin{array}{c} \text{Diagram 3} \\ - \end{array} \begin{array}{c} \text{Diagram 4} \\ - \end{array} \begin{array}{c} \text{Diagram 5} \\ - \end{array} \begin{array}{c} \text{Diagram 6} \\ - \end{array} \begin{array}{c} \text{Diagram 7} \\ - \end{array} \begin{array}{c} \text{Diagram 8} \\ + \text{perm.} \end{array} \end{array}$$

# Vertex Expansion

[MM, Strodthoff, Pawłowski, 2014],

[Cyrol, Fister, MM, Strodthoff, Pawłowski, 2016]

$$\partial_t \text{---}^{-1} = \begin{array}{c} \text{Diagram 1} \\ + \end{array} \begin{array}{c} \text{Diagram 2} \\ + \frac{1}{2} \end{array} \begin{array}{c} \text{Diagram 3} \\ - \end{array} \begin{array}{c} \text{Diagram 4} \\ + \end{array} \begin{array}{c} \text{Diagram 5} \\ - \end{array} \begin{array}{c} \text{Diagram 6} \\ - \end{array}$$

$$\partial_t \text{---}^{-1} = \begin{array}{c} \text{Diagram 1} \\ + \end{array}$$

$$\partial_t \text{---}^{-1} = \begin{array}{c} \text{Diagram 1} \\ - 2 \end{array} \begin{array}{c} \text{Diagram 2} \\ + \frac{1}{2} \end{array}$$

$$\partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - \end{array} \begin{array}{c} \text{Diagram 2} \\ - \end{array} \begin{array}{c} \text{Diagram 3} \\ + \text{perm.} \end{array}$$

$$\begin{array}{c} \partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - \end{array} \begin{array}{c} \text{Diagram 2} \\ - \end{array} \begin{array}{c} \text{Diagram 3} \\ - \end{array} \begin{array}{c} \text{Diagram 4} \\ - \end{array} \begin{array}{c} \text{Diagram 5} \\ - \end{array} \begin{array}{c} \text{Diagram 6} \\ + 2 \end{array} \begin{array}{c} \text{Diagram 7} \\ - \end{array} \begin{array}{c} \text{Diagram 8} \\ + \text{perm.} \end{array} \end{array}$$

$$\partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - \end{array} \begin{array}{c} \text{Diagram 2} \\ + 2 \end{array} \begin{array}{c} \text{Diagram 3} \\ - \end{array} \begin{array}{c} \text{Diagram 4} \\ + \text{perm.} \end{array}$$

$$\partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - \end{array} \begin{array}{c} \text{Diagram 2} \\ - \end{array} \begin{array}{c} \text{Diagram 3} \\ + 2 \end{array} \begin{array}{c} \text{Diagram 4} \\ - \end{array} \begin{array}{c} \text{Diagram 5} \\ + \text{perm.} \end{array}$$

$$\begin{array}{c} \partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - 2 \end{array} \begin{array}{c} \text{Diagram 2} \\ - \end{array} \begin{array}{c} \text{Diagram 3} \\ - \end{array} \begin{array}{c} \text{Diagram 4} \\ - \end{array} \begin{array}{c} \text{Diagram 5} \\ - \end{array} \begin{array}{c} \text{Diagram 6} \\ - \end{array} \begin{array}{c} \text{Diagram 7} \\ - \end{array} \begin{array}{c} \text{Diagram 8} \\ + \text{perm.} \end{array} \end{array}$$

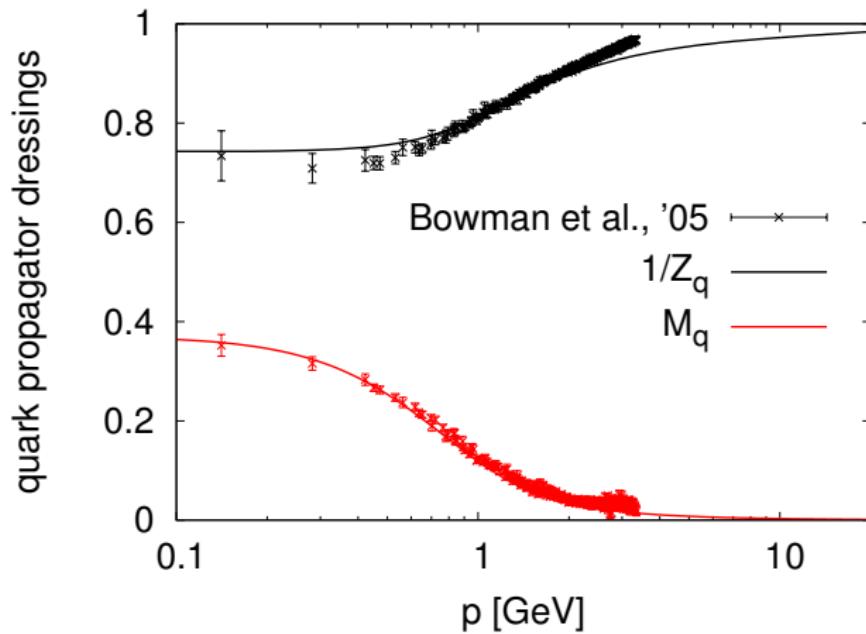
$$\partial_t \text{---}^{-1} = -2 \begin{array}{c} \text{Diagram 1} \\ - \end{array} \begin{array}{c} \text{Diagram 2} \\ + \end{array} \begin{array}{c} \text{Diagram 3} \\ + \frac{1}{2} \end{array}$$

$$\partial_t \text{---} = \begin{array}{c} \text{Diagram 1} \\ - \end{array} \begin{array}{c} \text{Diagram 2} \\ - \end{array} \begin{array}{c} \text{Diagram 3} \\ - \end{array} \begin{array}{c} \text{Diagram 4} \\ + 2 \end{array} \begin{array}{c} \text{Diagram 5} \\ + \text{perm.} \end{array}$$

# Quark propagator

[MM, Pawłowski, Strodthoff, 2014]

- $\Gamma_{\bar{q}q}^{(2)}(p) \propto Z_q(p) (\not{p} + M(p))$

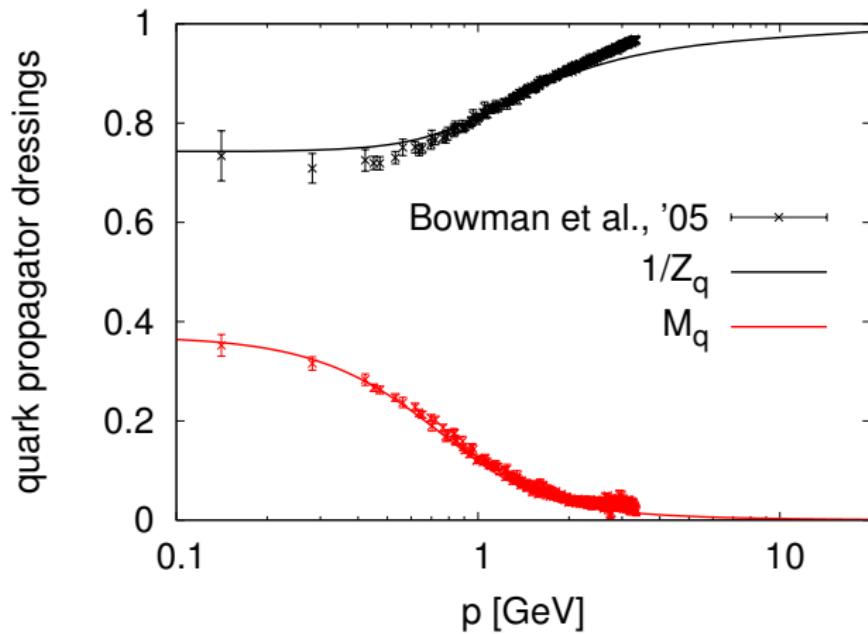


- FRG vs. lattice: bare mass, quenched, scale set via gluon propagator

# Quark propagator

[MM, Pawłowski, Strodthoff, 2014]

- $\Gamma_{\bar{q}q}^{(2)}(p) \propto Z_q(p) (\not{p} + M(p))$

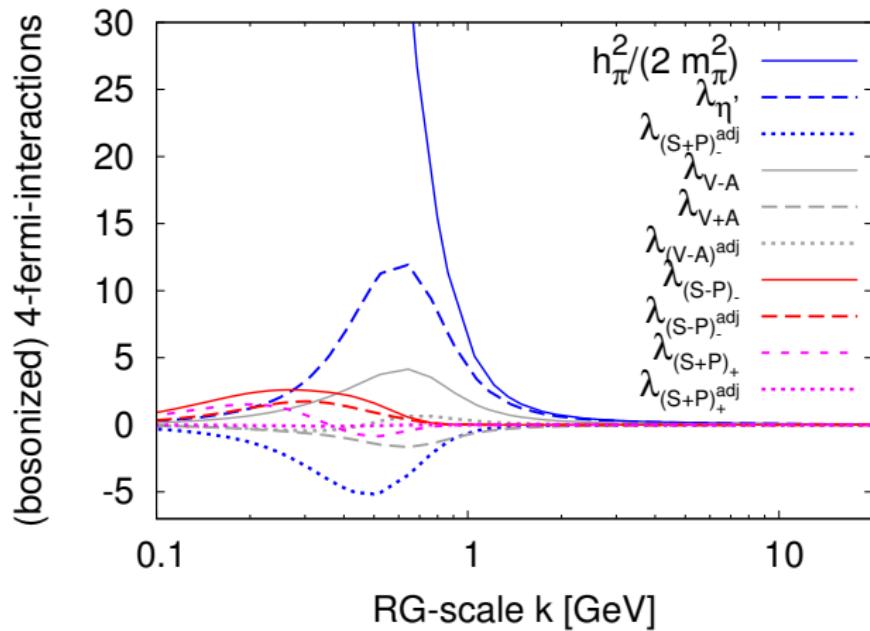


- FRG vs. lattice: bare mass, quenched, scale set via gluon propagator
- agreement not sufficient: need apparent convergence at  $\mu \neq 0$

lattice data: Bowman, Heller, Leinweber, Parappilly, Williams, Zhang , Phys. Rev. D71, 054507 (2005).

## other 4-Fermi channels

[MM, Pawlowski, Strodthoff, 2014]



- bosonized only  $\sigma$ - $\pi$ -channel momentum dependently  $\Rightarrow$  sufficient
- other channels: quantitatively not important in loops

# Quark-gluon interactions

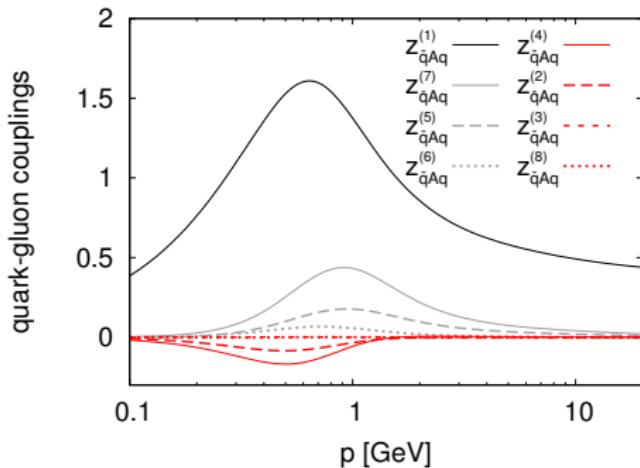
[MM, Pawłowski, Strodthoff, 2014]

- quark-gluon interaction most crucial for chiral symmetry breaking
- full tensor basis, e.g.  $\gamma^\mu$ ,  $i(\not{p} + \not{q})\gamma^\mu$ ,  $\frac{1}{2} [\not{p}, \not{q}] \gamma^\mu$

# Quark-gluon interactions

[MM, Pawlowski, Strodthoff, 2014]

- quark-gluon interaction most crucial for chiral symmetry breaking
- full tensor basis, e.g.  $\gamma^\mu$ ,  $i(\not{p} + \not{q})\gamma^\mu$ ,  $\frac{1}{2}[\not{p}, \not{q}]\gamma^\mu$

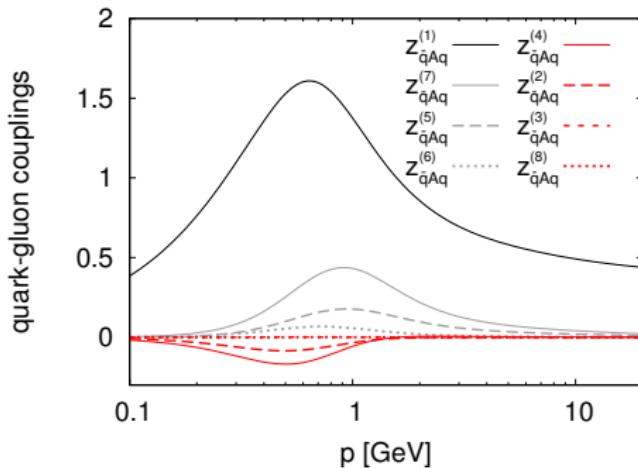


- vertex strength:  
reflects gluon gap
- 8 tensors (transversally projected):
  - ▶ classical tensor
  - ▶ chirally symmetric
  - ▶ **break chiral symmetry**

# Quark-gluon interactions

[MM, Pawlowski, Strodthoff, 2014]

- quark-gluon interaction most crucial for chiral symmetry breaking
- full tensor basis, e.g.  $\gamma^\mu$ ,  $i(\not{p} + \not{q})\gamma^\mu$ ,  $\frac{1}{2}[\not{p}, \not{q}]\gamma^\mu$



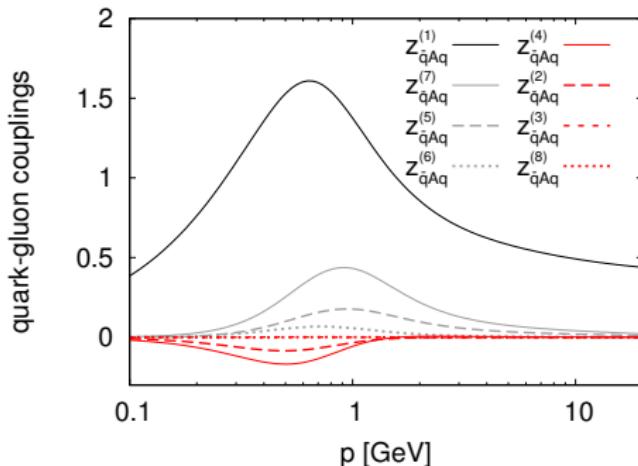
- vertex strength:  
reflects gluon gap
- 8 tensors (transversally projected):
  - ▶ classical tensor
  - ▶ chirally symmetric
  - ▶ **break chiral symmetry**

- chirally symmetric tensors from operator  $\bar{q}\not{\partial}^3 q$  worsen result

# Quark-gluon interactions

[MM, Pawlowski, Strodthoff, 2014]

- quark-gluon interaction most crucial for chiral symmetry breaking
- full tensor basis, e.g.  $\gamma^\mu$ ,  $i(\not{p} + \not{q})\gamma^\mu$ ,  $\frac{1}{2}[\not{p}, \not{q}]\gamma^\mu$



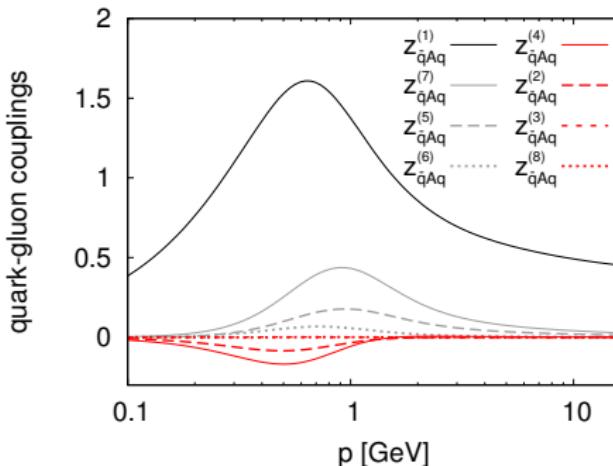
- vertex strength:  
reflects gluon gap
- 8 tensors (transversally projected):
  - ▶ classical tensor
  - ▶ chirally symmetric
  - ▶ **break chiral symmetry**

- chirally symmetric tensors from operator  $\bar{q}\not{\partial}^3 q$  worsen result
- counteracted by tensor structures in  $\Gamma_{AA\bar{q}q}^{(4)}$  and  $\Gamma_{A^3\bar{q}q}^{(5)}$  from  $\bar{q}\not{\partial}^3 q$

# Quark-gluon interactions

[MM, Pawlowski, Strodthoff, 2014]

- quark-gluon interaction most crucial for chiral symmetry breaking
- full tensor basis, e.g.  $\gamma^\mu$ ,  $i(\not{p} + \not{q})\gamma^\mu$ ,  $\frac{1}{2}[\not{p}, \not{q}]\gamma^\mu$



- vertex strength:  
reflects gluon gap
- 8 tensors (transversally projected):
  - ▶ classical tensor
  - ▶ chirally symmetric
  - ▶ **break chiral symmetry**

- chirally symmetric tensors from operator  $\bar{q}\not{\partial}^3 q$  worsen result
- counteracted by tensor structures in  $\Gamma_{AA\bar{q}q}^{(4)}$  and  $\Gamma_{A^3\bar{q}q}^{(5)}$  from  $\bar{q}\not{\partial}^3 q$

⇒ expansion in BRST-invariant operators improves convergence?

## YM theory: modified STIs (mSTIs)

- regulator breaks BRST-symmetry of gauge-fixed action  
⇒ modification to Slavnov-Taylor identities  $\propto R_k$  ( $\rightarrow 0$  for  $k \rightarrow 0$ )

[Ellwanger, Hirsch, Weber, 1996]

## YM theory: modified STIs (mSTIs)

- regulator breaks BRST-symmetry of gauge-fixed action  
⇒ modification to Slavnov-Taylor identities  $\propto R_k$  ( $\rightarrow 0$  for  $k \rightarrow 0$ )

[Ellwanger, Hirsch, Weber, 1996]

- most immediate consequence for gluon propagator at  $\Lambda$ :

$$[\Gamma_{AA}^{(2)}]_{\mu\nu}(p) = Z_\Lambda p^2 \left( \delta_{\mu\nu} - \frac{p^\mu p^\nu}{p^2} \right) + m_\Lambda^2 \delta_{\mu\nu} \quad m_\Lambda^2 \propto \Lambda^2$$

## YM theory: modified STIs (mSTIs)

- regulator breaks BRST-symmetry of gauge-fixed action  
⇒ modification to Slavnov-Taylor identities  $\propto R_k$  ( $\rightarrow 0$  for  $k \rightarrow 0$ )

[Ellwanger, Hirsch, Weber, 1996]

- most immediate consequence for gluon propagator at  $\Lambda$ :

$$[\Gamma_{AA}^{(2)}]_{\mu\nu}(p) = Z_\Lambda p^2 \left( \delta_{\mu\nu} - \frac{p^\mu p^\nu}{p^2} \right) + m_\Lambda^2 \delta_{\mu\nu} \quad m_\Lambda^2 \propto \Lambda^2$$

- mSTI determines  $m_\Lambda^2$ , but hard numerical problem due to  $\Lambda^2$

## YM theory: modified STIs (mSTIs)

- regulator breaks BRST-symmetry of gauge-fixed action  
⇒ modification to Slavnov-Taylor identities  $\propto R_k$  ( $\rightarrow 0$  for  $k \rightarrow 0$ )

[Ellwanger, Hirsch, Weber, 1996]

- most immediate consequence for gluon propagator at  $\Lambda$ :

$$[\Gamma_{AA}^{(2)}]_{\mu\nu}(p) = Z_\Lambda p^2 \left( \delta_{\mu\nu} - \frac{p^\mu p^\nu}{p^2} \right) + m_\Lambda^2 \delta_{\mu\nu} \quad m_\Lambda^2 \propto \Lambda^2$$

- mSTI determines  $m_\Lambda^2$ , but hard numerical problem due to  $\Lambda^2$
- BRST-symmetry at  $k = 0$  requires propagators of scaling type:

$$\lim_{p \rightarrow 0} [\Gamma_{AA}^{(2)}]_{\mu\nu}(p) \propto (p^2)^{1-2\kappa} \quad \lim_{p \rightarrow 0} [\Gamma_{\bar{c}c}^{(2)}]_{\mu\nu}(p) \propto (p^2)^{1+\kappa} \quad \kappa \in (0.5, 1)$$

[Lerche, von Smekal, 2002]

- scaling solution (BRST-symmetry) fixes  $m_\Lambda^2$  uniquely

## YM theory: modified STIs (mSTIs)

- regulator breaks BRST-symmetry of gauge-fixed action  
⇒ modification to Slavnov-Taylor identities  $\propto R_k$  ( $\rightarrow 0$  for  $k \rightarrow 0$ )

[Ellwanger, Hirsch, Weber, 1996]

- most immediate consequence for gluon propagator at  $\Lambda$ :

$$[\Gamma_{AA}^{(2)}]_{\mu\nu}(p) = Z_\Lambda p^2 \left( \delta_{\mu\nu} - \frac{p^\mu p^\nu}{p^2} \right) + m_\Lambda^2 \delta_{\mu\nu} \quad m_\Lambda^2 \propto \Lambda^2$$

- mSTI determines  $m_\Lambda^2$ , but hard numerical problem due to  $\Lambda^2$
- BRST-symmetry at  $k = 0$  requires propagators of scaling type:

$$\lim_{p \rightarrow 0} [\Gamma_{AA}^{(2)}]_{\mu\nu}(p) \propto (p^2)^{1-2\kappa} \quad \lim_{p \rightarrow 0} [\Gamma_{cc}^{(2)}]_{\mu\nu}(p) \propto (p^2)^{1+\kappa} \quad \kappa \in (0.5, 1)$$

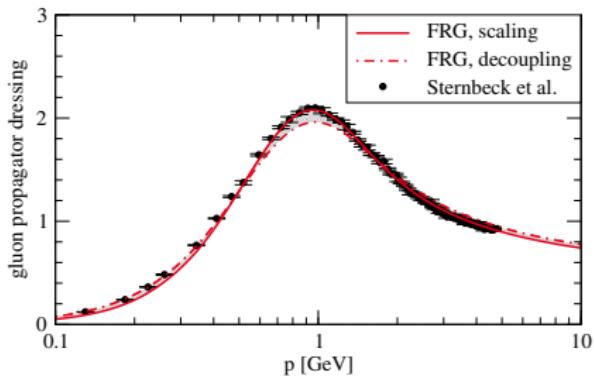
[Lerche, von Smekal, 2002]

- scaling solution (BRST-symmetry) fixes  $m_\Lambda^2$  uniquely
- lattice simulations: no scaling (decoupling solutions) [Cucchieri, Mendes, 2008]
- non-perturbative gauge-fixing on the lattice? [Maas, 2009]

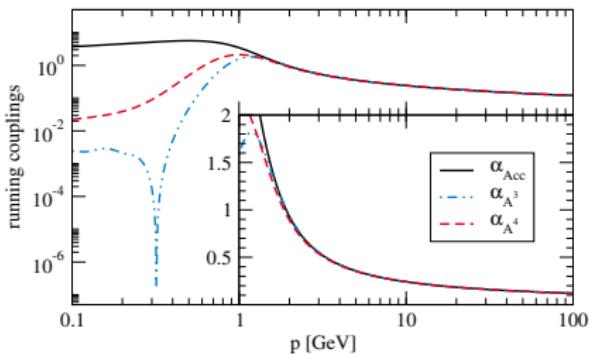
# YM theory: results

[Cyrol, Fister, MM, Pawłowski, Strodthoff, 2016]

- $\Gamma_{AA}^{(2)}(p) \propto Z_A(p) p^2 (\delta^{\mu\nu} - p^\mu p^\nu / p^2)$



- running couplings



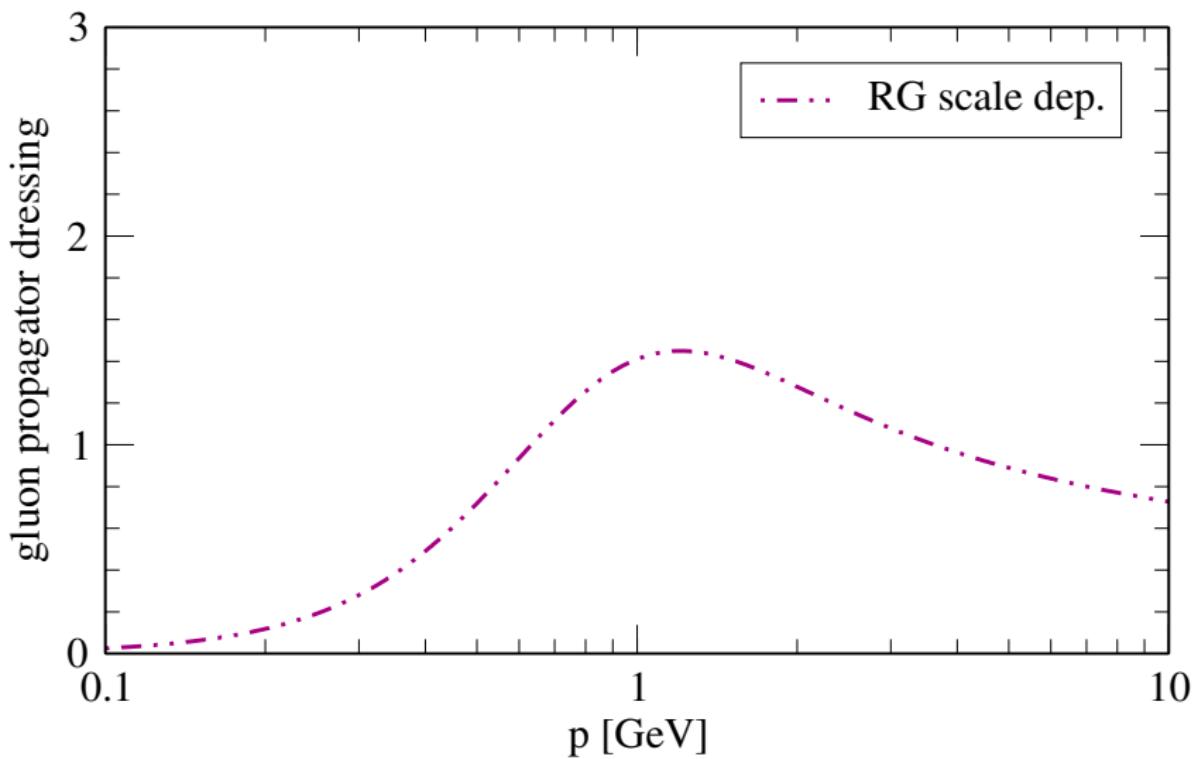
- decoupling solutions:

- ▶ vary  $m_\Lambda^2$  slightly away from scaling value
- ▶ dashed-dot line: largest variation of  $m_\Lambda^2$  that leads to “back-bending”

lattice data: A. Sternbeck, E. M. Ilgenfritz, M. Müller-Preussker, A. Schiller, and I. L. Bogolubsky, PoS LAT2006, 076.

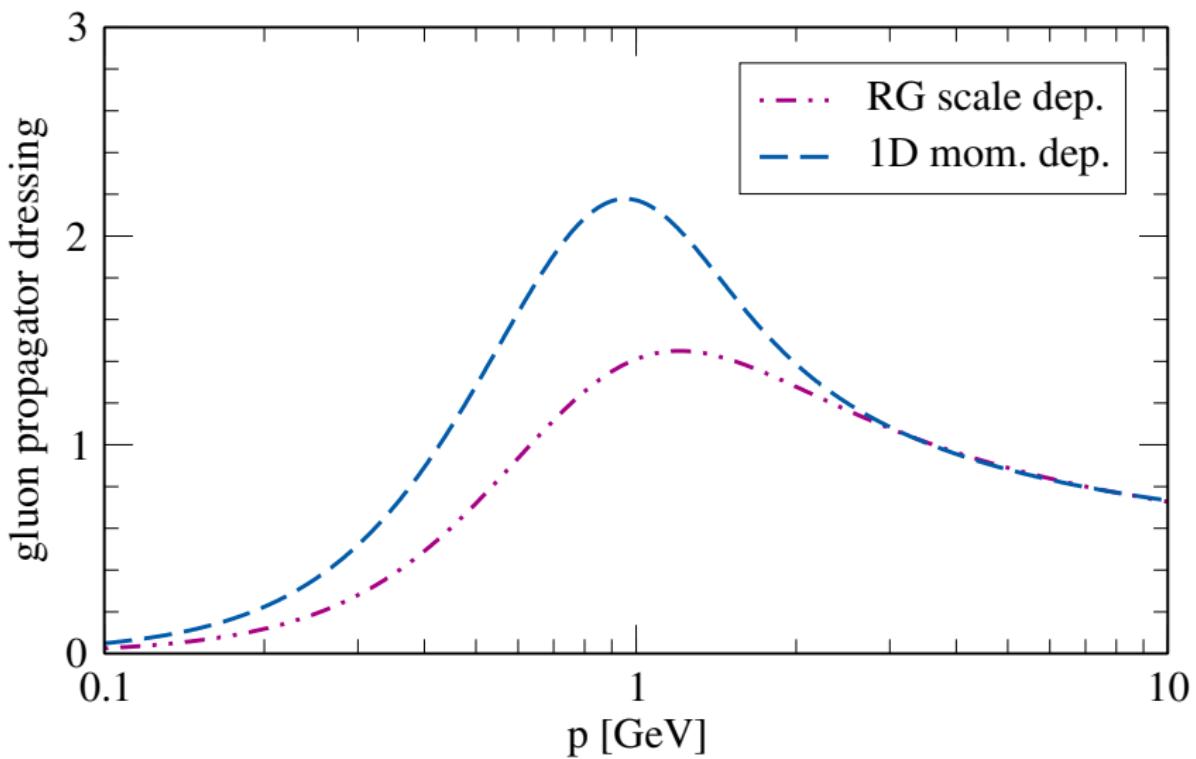
# YM-theory: Truncation dependence

[Cyrol, Fister, MM, Pawłowski, Strodthoff, 2016]



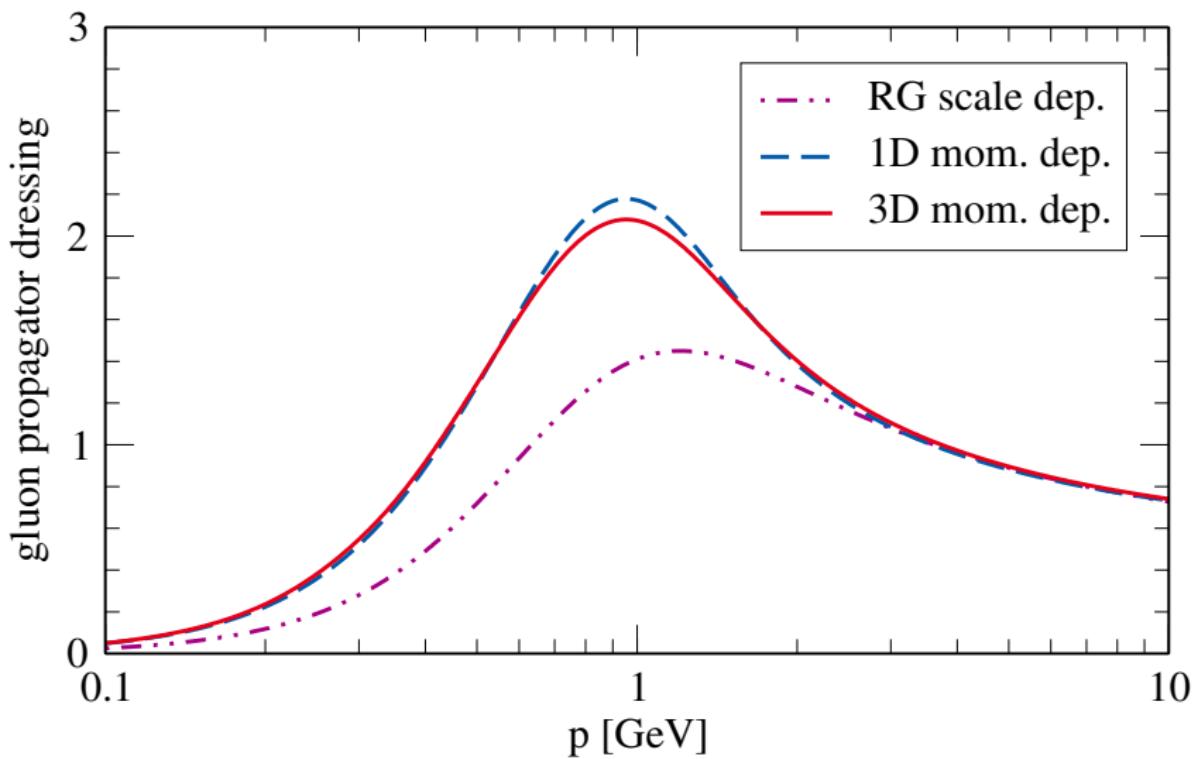
# YM-theory: Truncation dependence

[Cyrol, Fister, MM, Pawłowski, Strodthoff, 2016]

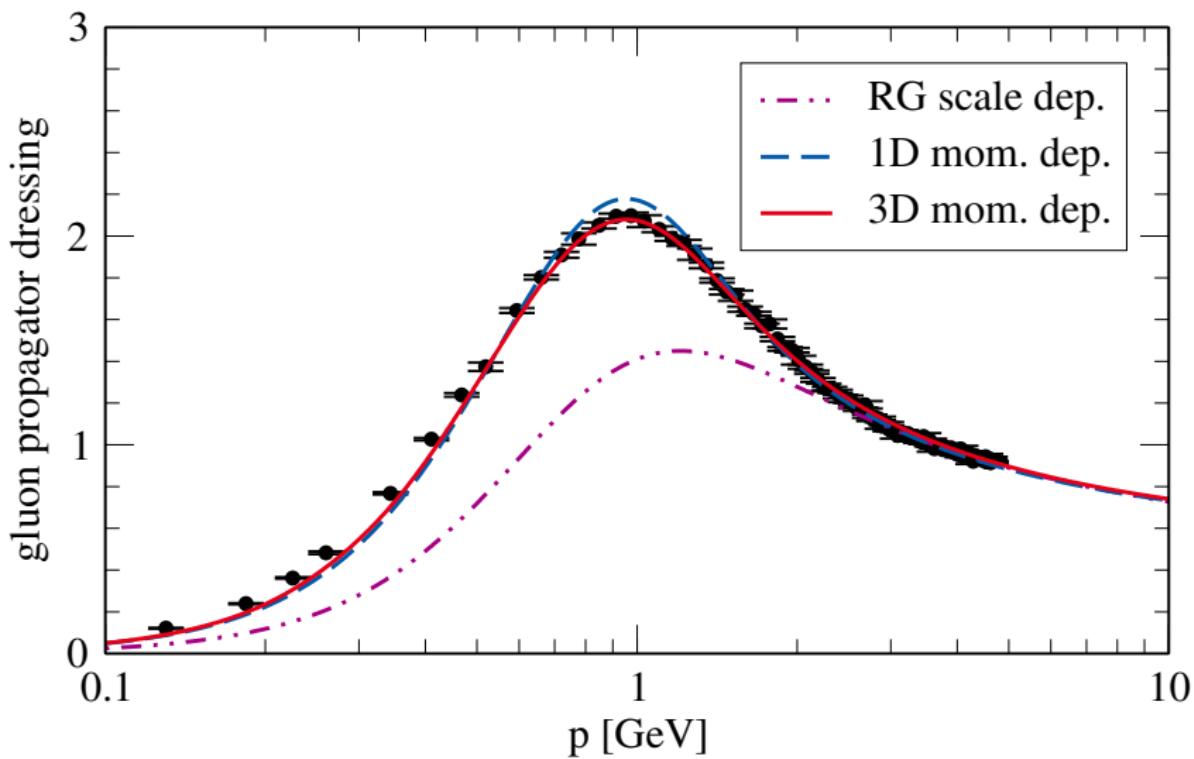


# YM-theory: Truncation dependence

[Cyrol, Fister, MM, Pawłowski, Strodthoff, 2016]



# YM-theory: Truncation dependence [Cyrol, Fister, MM, Pawłowski, Strodthoff, 2016]

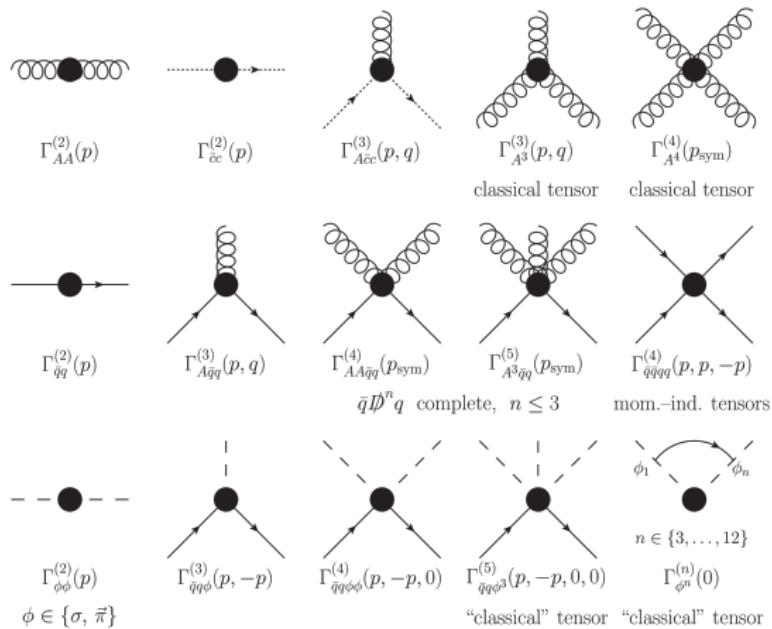


lattice data: A. Sternbeck, E. M. Ilgenfritz, M. Müller-Preussker, A. Schiller, and I. L. Bogolubsky, PoS LAT2006, 076.

# Outlook: unquenching

[Cyrol, MM, Pawłowski, Strodthoff, in prep.]

extended truncation:

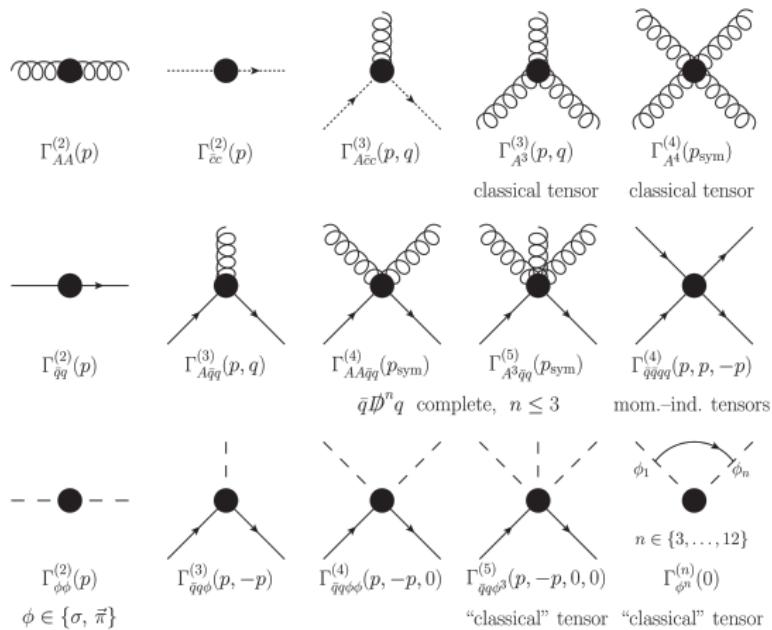


systematics of improving the truncation?

# Outlook: unquenching

[Cyrol, MM, Pawłowski, Strodthoff, in prep.]

extended truncation:



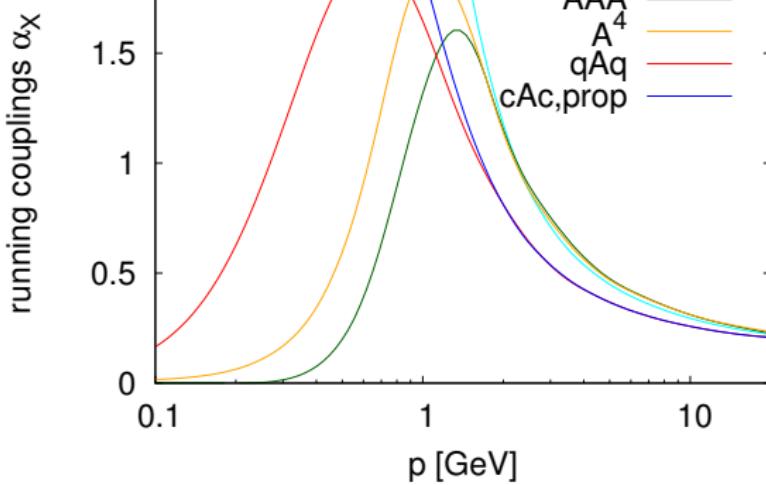
systematics of improving the truncation?

⇒ BRST-invariant operators, e.g.  $\bar{\psi} \not{D}^n \psi$

# Outlook: running couplings

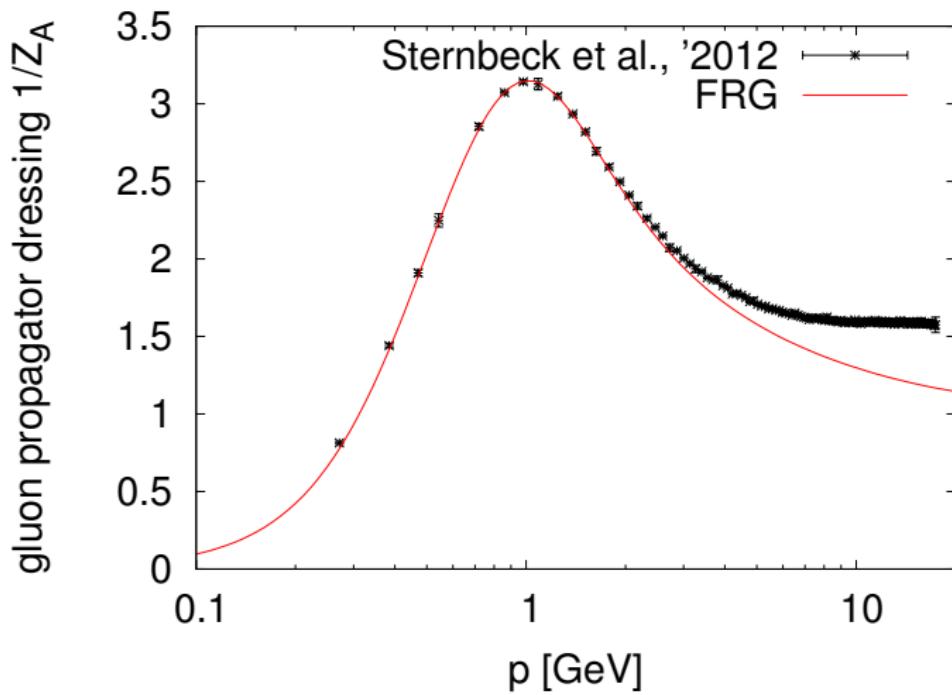
[Cyrol, MM, Pawłowski, Strodthoff, in preparation]

- $\alpha_{cAc} = \frac{\left(\Gamma_{cAc}^{(3)}(p)\right)^2}{4\pi Z_A(p)Z_c(p)^2}$
- $\alpha_{AAA} = \frac{\left(\Gamma_{AAA}^{(3)}(p)\right)^2}{4\pi Z_A(p)^3}$
- $\alpha_{A^4} = \frac{\left(\Gamma_{A^4}^{(4)}(p)\right)}{4\pi Z_A(p)^2}$
- $\alpha_{qAq} = \frac{\left(\Gamma_{qAq}^{(3)}(p)\right)^2}{4\pi Z_A(p)Z_q(p)^2}$
- $\alpha_{cAc,\text{prop}} = \frac{1}{4\pi Z_A(p)Z_c(p)^2}$



# Outlook: gluon propagator

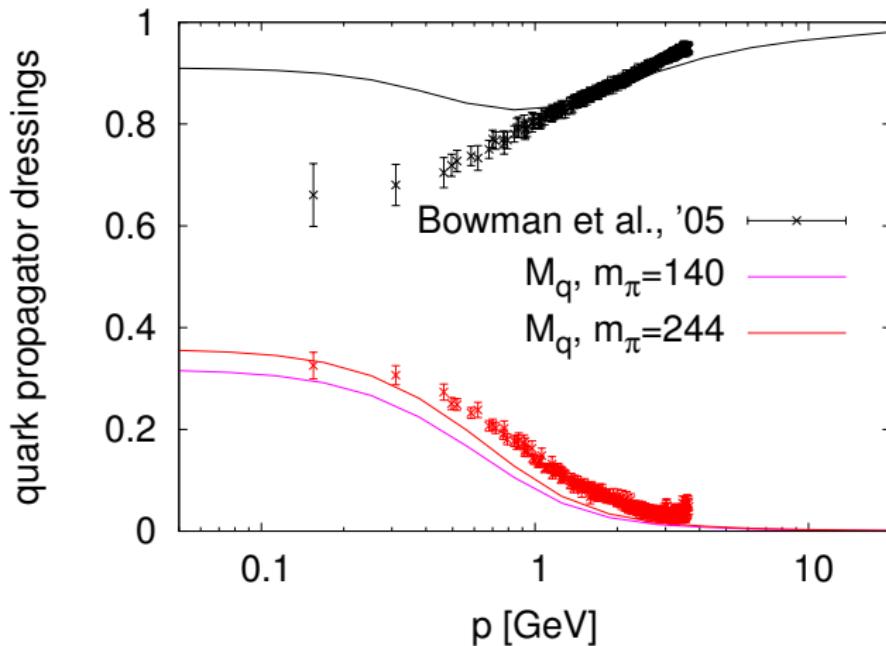
[Cyrol, MM, Pawłowski, Strodthoff, in preparation]



lattice data: A. Sternbeck, K. Maltman, M. Müller-Preussker, L. von Smekal, PoS LATTICE2012 (2012) 243.

# Outlook: quark propagator

[Cyrol, MM, Pawłowski, Strodthoff, in preparation]

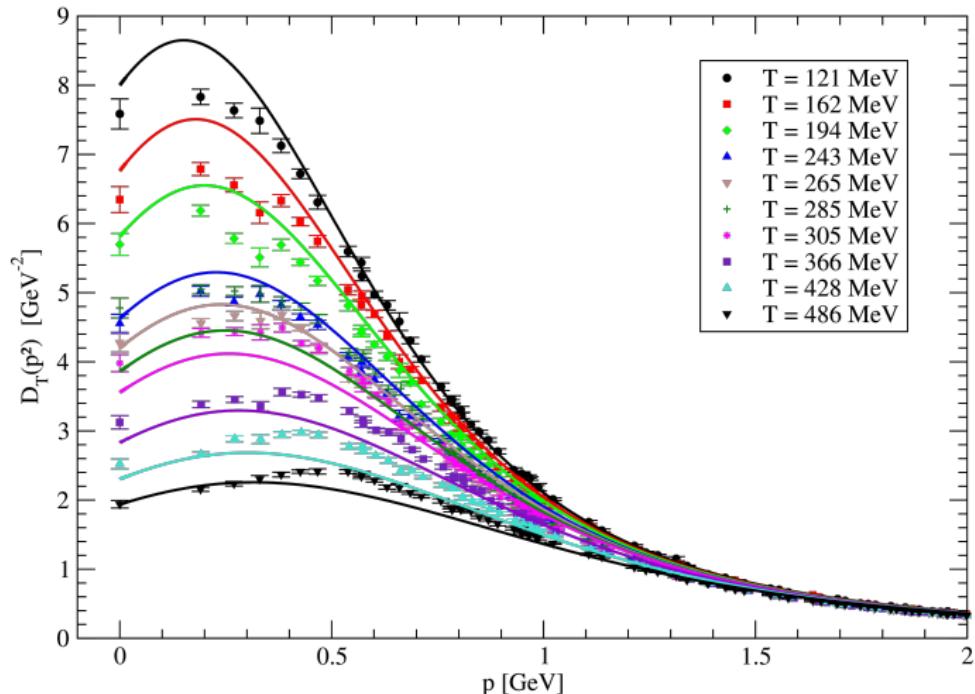


- comparison FRG with lattice: bare mass and scale setting

lattice data: Bowman, Heller, Leinweber, Parappilly, Williams, Zhang , Phys. Rev. D71, 054507 (2005).

# Outlook: (averaged) gluon propagator at $T \neq 0$

[Cyrol, MM, Pawłowski, Strodthoff, in preparation]



lattice data: magnetic component of gluon propagator from Silva et. al. Phys.Rev. D89 (2014)

# Summary and Outlook

## QCD with functional RG

- vertex expansion
- sole input  $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$  and  $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
- good agreement with lattice correlators

# Summary and Outlook

## QCD with functional RG

- vertex expansion
- sole input  $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$  and  $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
- good agreement with lattice correlators

## Outlook

- QCD phase diagram:  
order parameters, equation of state and fluct. of cons. charges
- bound-state properties (form factors, PDA...)
- more checks on convergence of vertex expansion

# Summary and Outlook

## QCD with functional RG

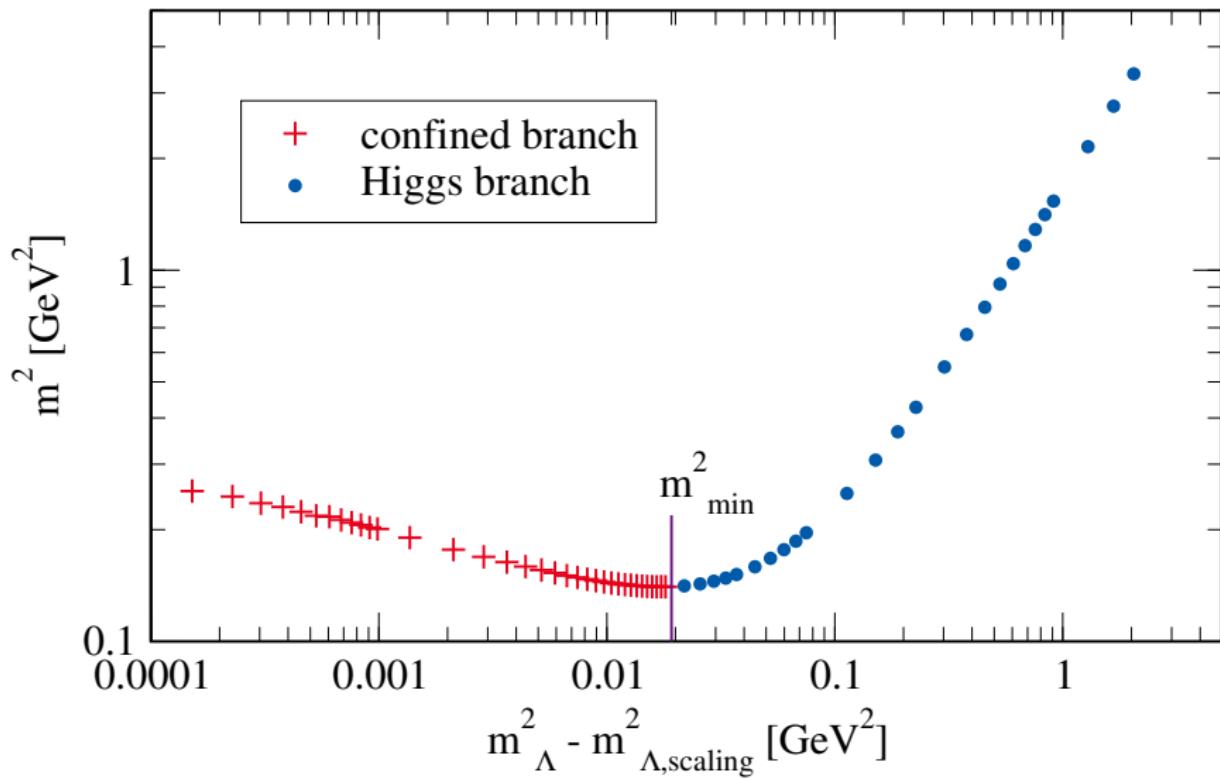
- vertex expansion
- sole input  $\alpha_s(\Lambda = \mathcal{O}(10) \text{ GeV})$  and  $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
- good agreement with lattice correlators

## Outlook

- QCD phase diagram:  
order parameters, equation of state and fluct. of cons. charges
- bound-state properties (form factors, PDA...)
- more checks on convergence of vertex expansion

Poster by Anton K. Cyrol: “fQCD: QCD with the Functional RG”

# Dynamical mass generation



# Gluon propagator maximum over UV mass parameter

