

GOLDSTONE MODE & SOFT-MODE AT FINITE ISOSPIN DENSITY

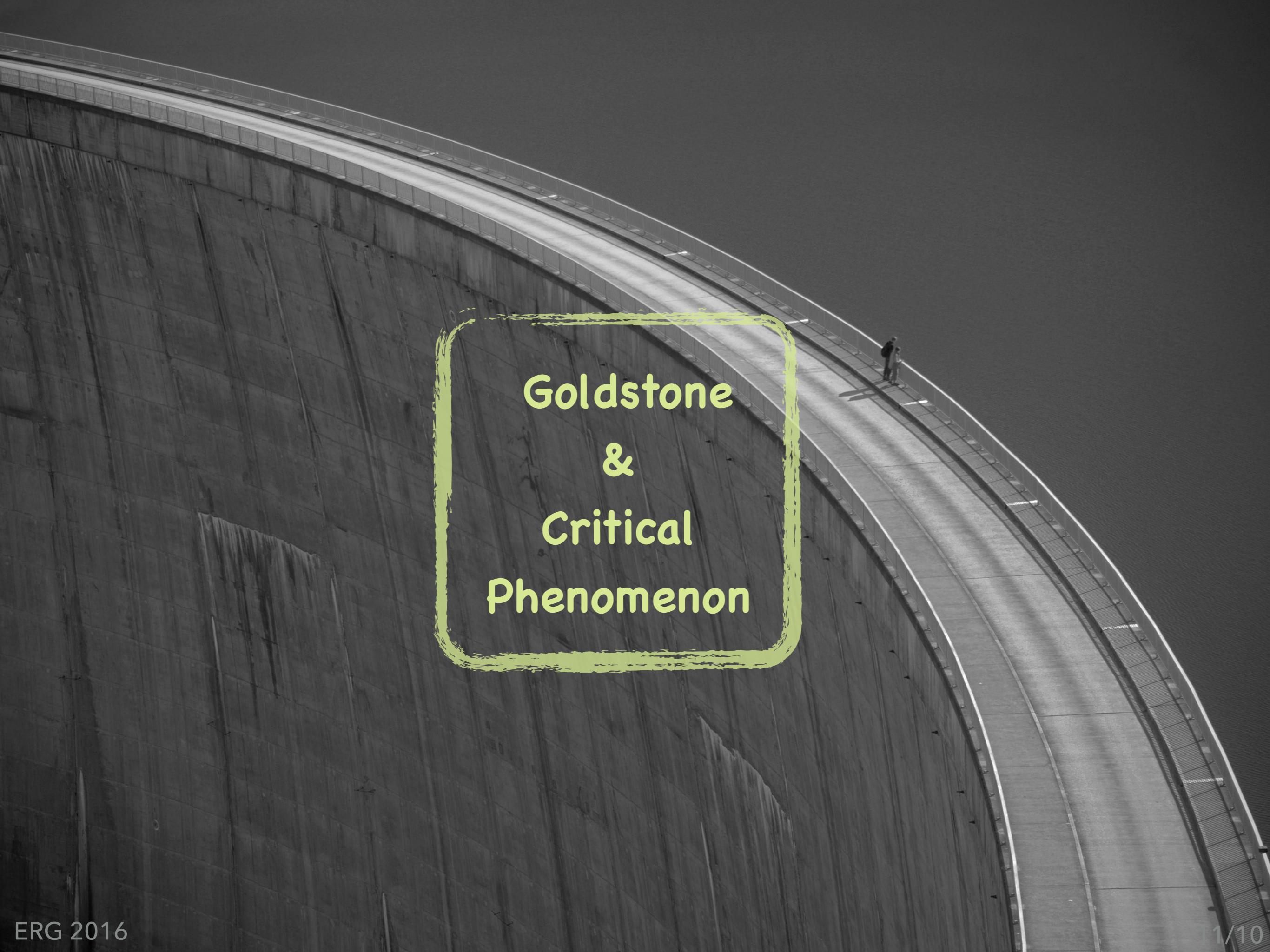
Ziyue WANG and Pengfei ZHUANG

Tsinghua University, Beijing

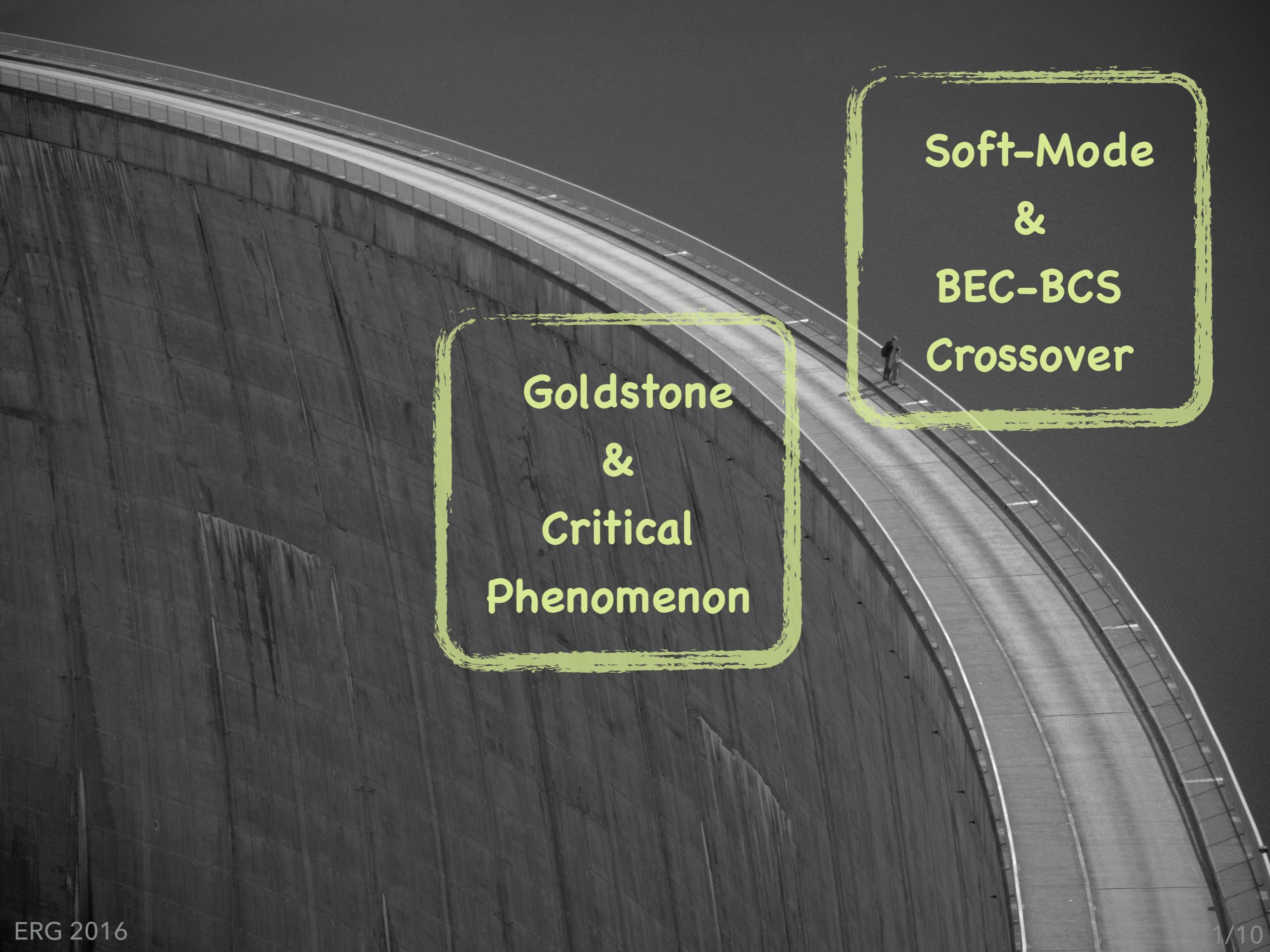
@ ERG 2016, Trieste, Italy

September 19, 2016



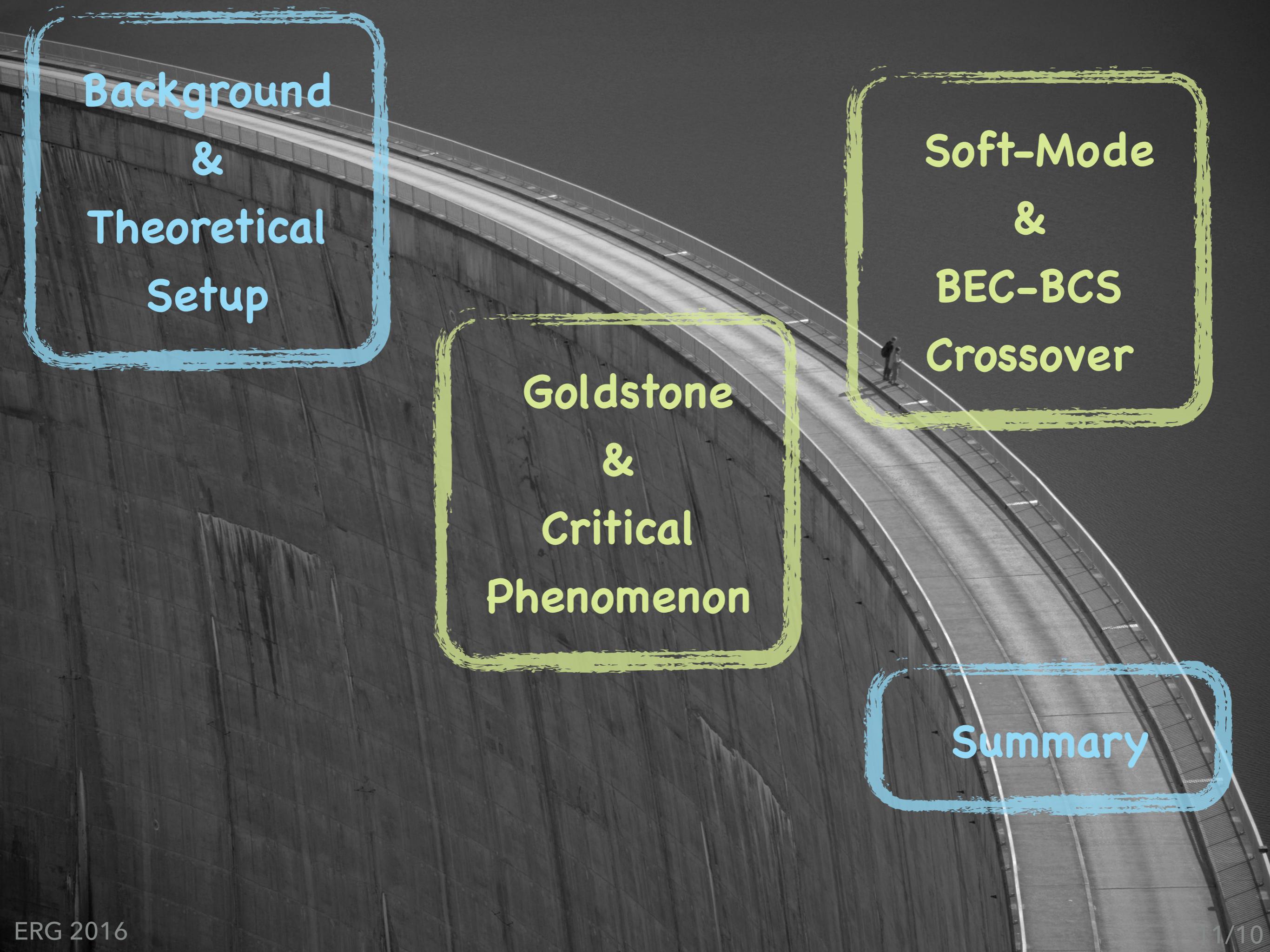


Goldstone
&
Critical
Phenomenon



Goldstone
&
Critical
Phenomenon

Soft-Mode
&
BEC-BCS
Crossover



Background
&
Theoretical
Setup

Soft-Mode
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BEC-BCS
Crossover

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Summary

Pion Superfluidity @ Finite Isospin Density

Background & Theoretical Setup

► Motivation:

- isospin imbalance in the interior of neutron stars

[S. Barshay et al, Phys. Lett. B 47, 107 (1973)] [V. A. Khodel et al, Phys. Rev. Lett. 93, 151101 (2004)]

- free from lattice sign problem – model comparison
- relativistic BEC-BCS crossover

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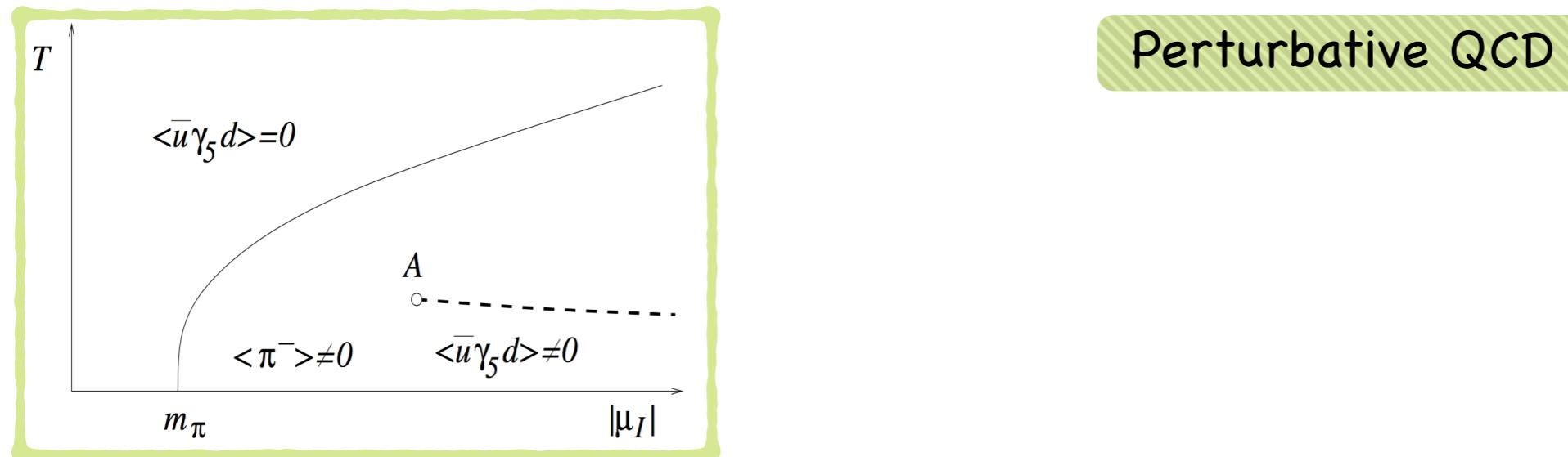
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Thermodynamics in theoretical methods:



[D.T. Son and M.A. Stephanov, PRL.86,592]

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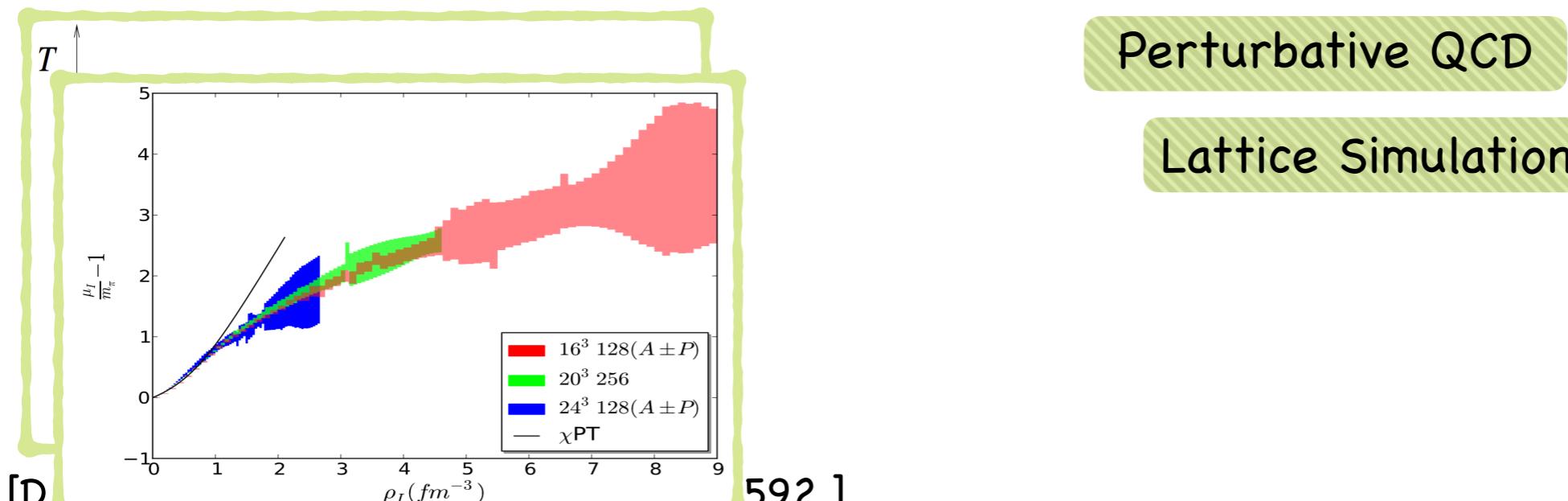
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Thermodynamics in theoretical methods:



[W. Detmold, et al., Phys. Rev. D 86, 054507]

Perturbative QCD

Lattice Simulation

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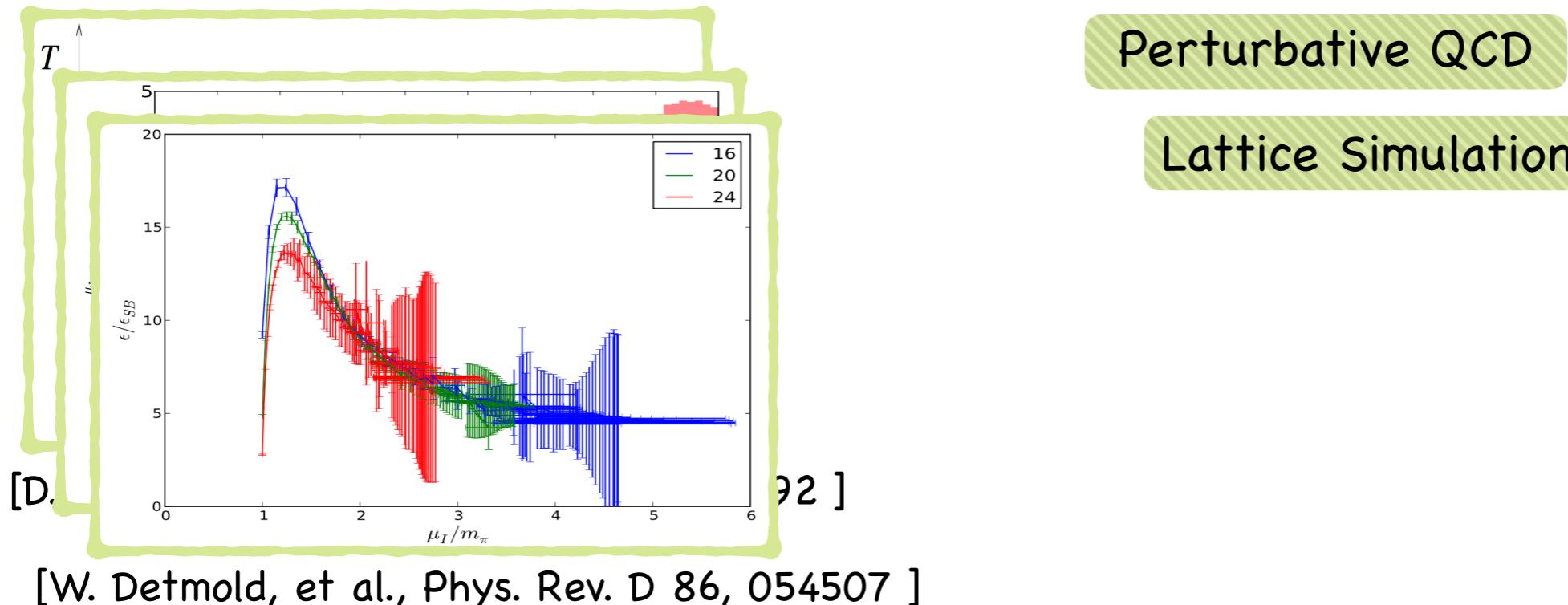
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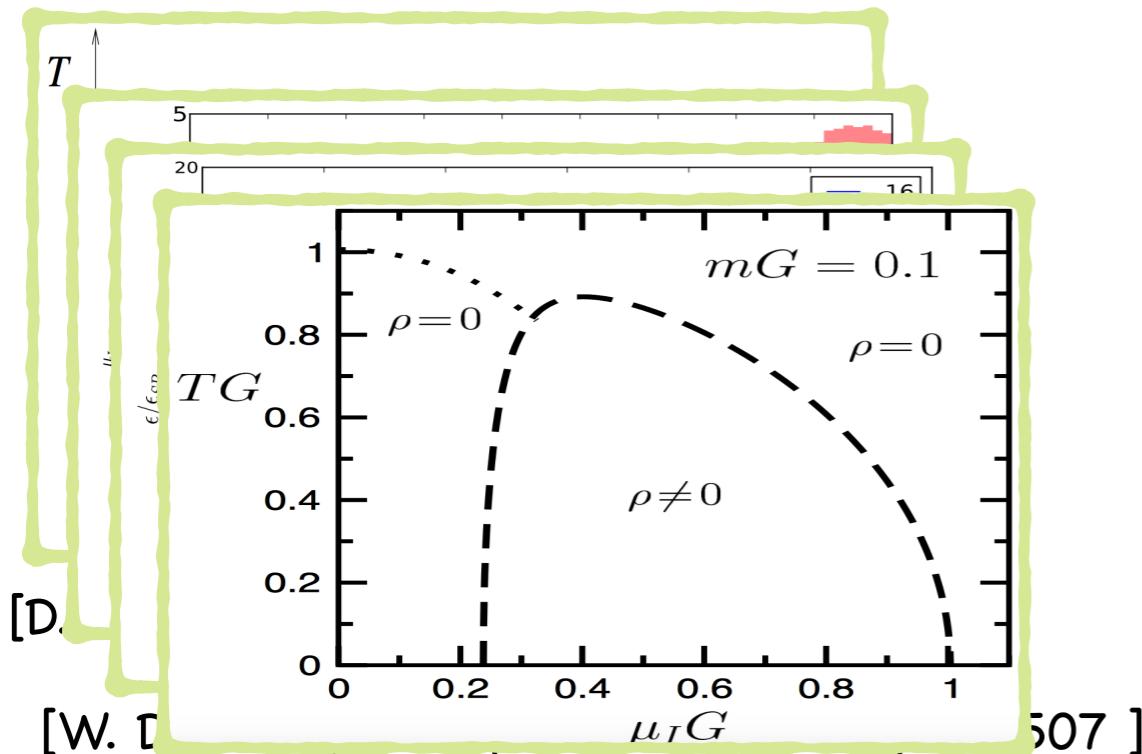
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Thermodynamics in theoretical methods:



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Random Matrix

[B. Klein et al., Phys. Rev. D 72, 015007]

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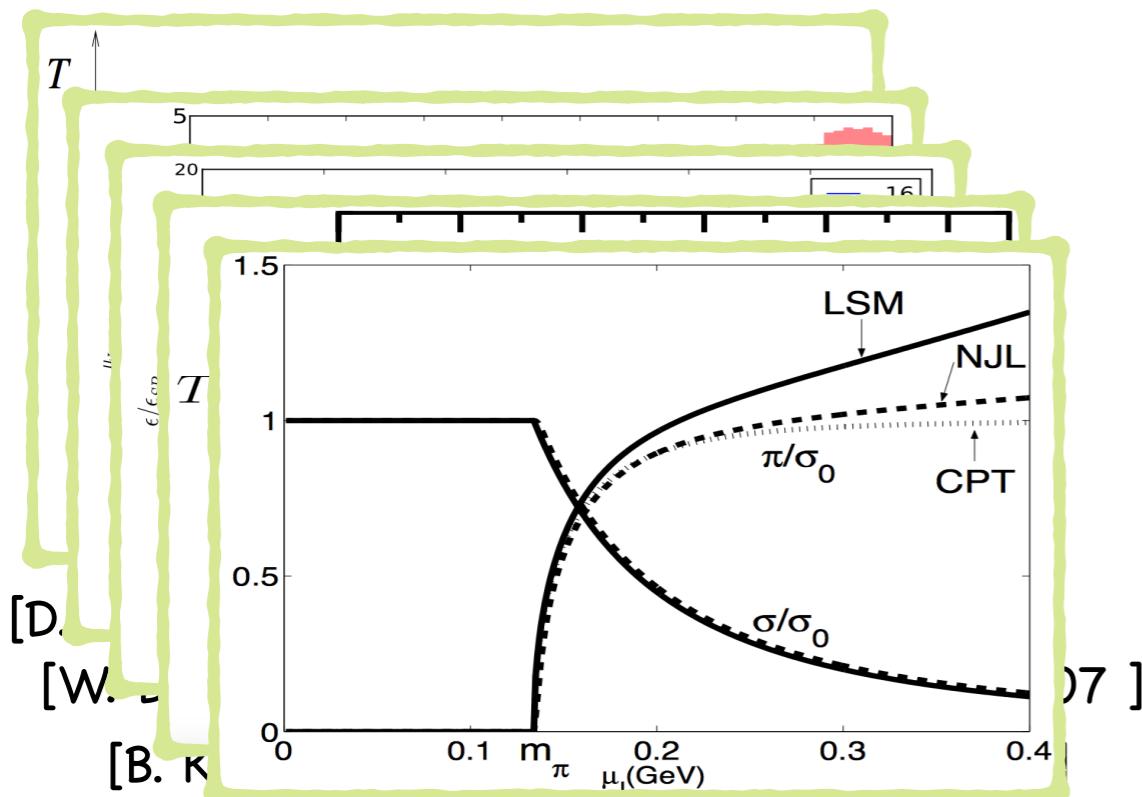
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Thermodynamics in theoretical methods:



[L. He, et al., Phys. Rev. D 71, 116001]

Perturbative QCD

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Nambu- Jona-Lasinio model

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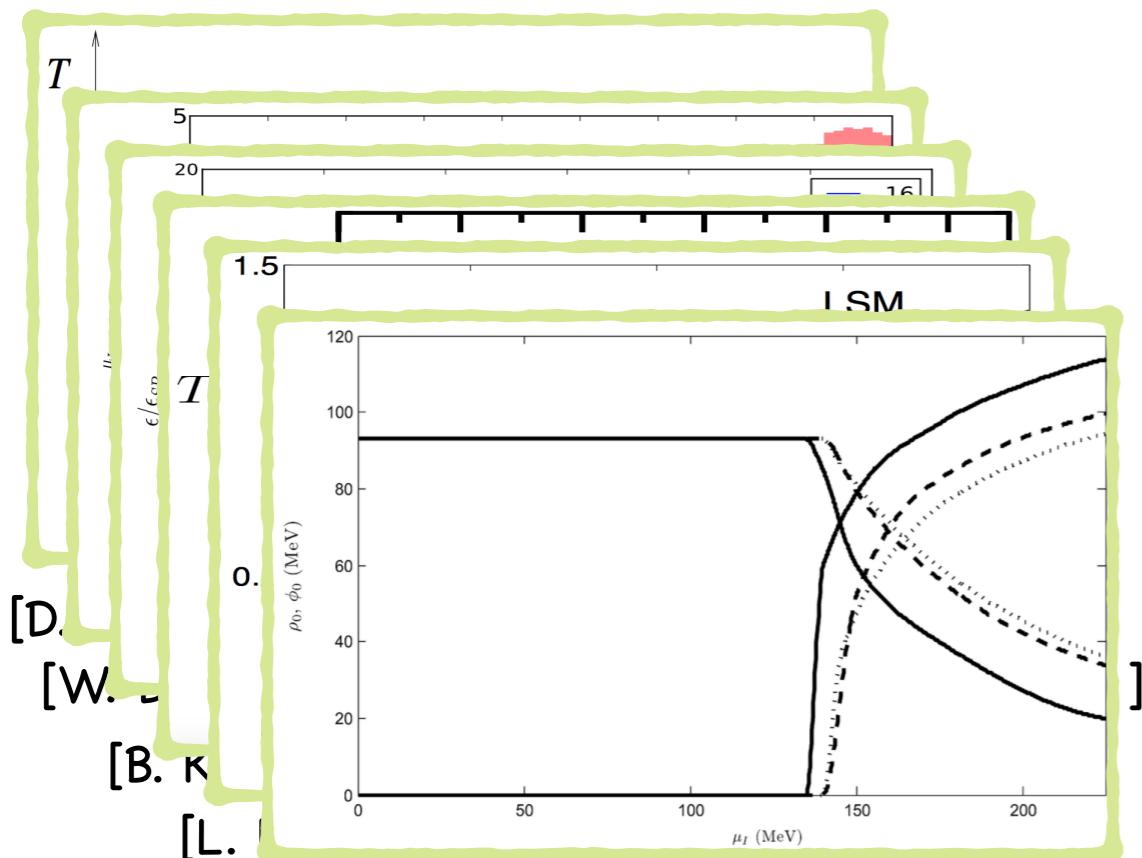
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Linear-Sigma model

[E. E. Svanes et al., Nucl. Phys. A 857, 16]

Pion Superfluidity @ Finite Isospin Density

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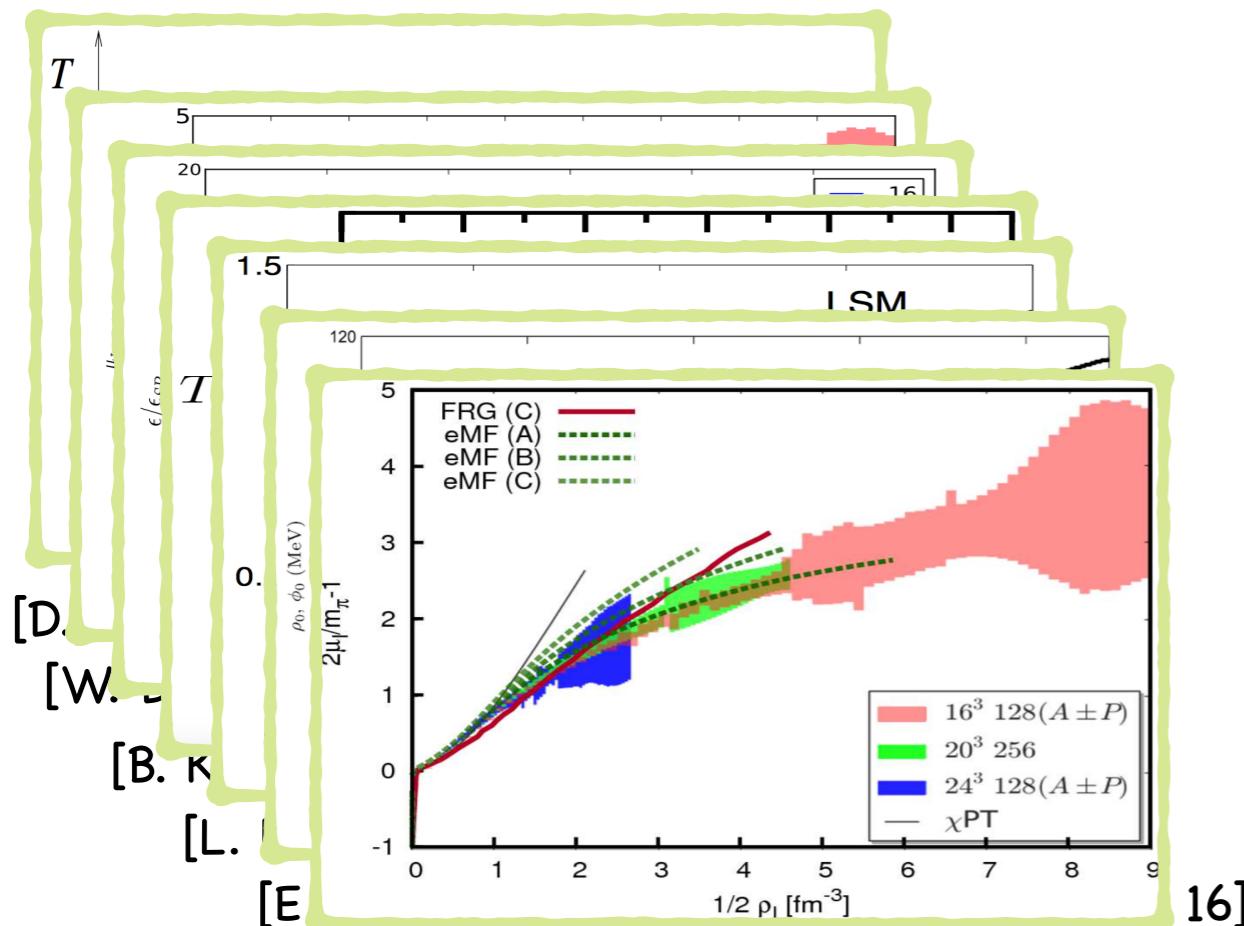
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Thermodynamics in theoretical methods:



[K. Kamikado, et al., Phys. Lett. B 718, 1044]

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Linear-Sigma model

Quark-Meson model

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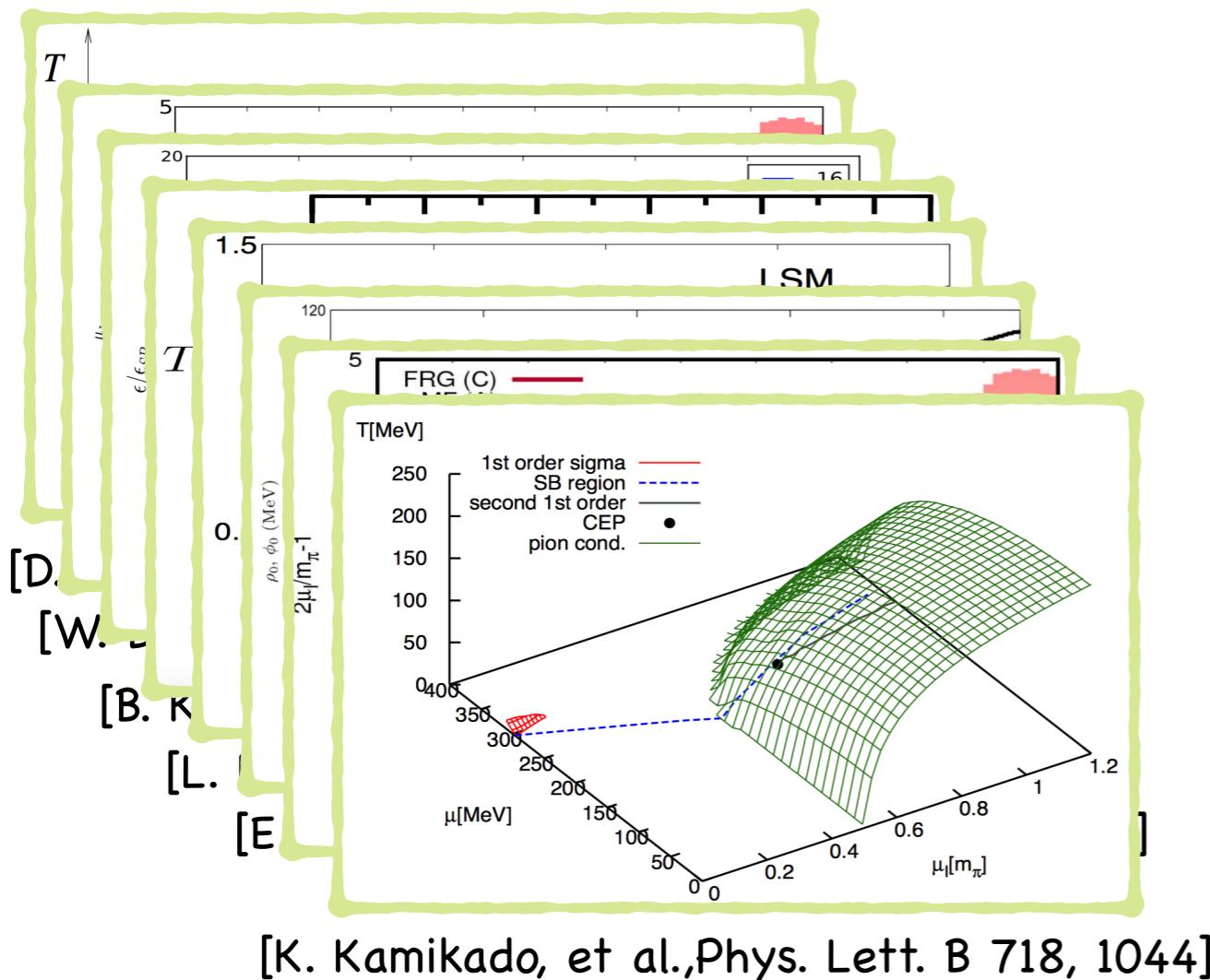
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Second order phase transition
on phase boundary

Background & Theoretical Setup

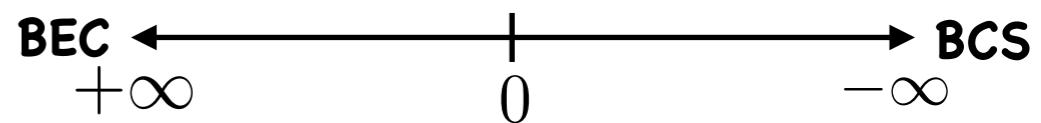
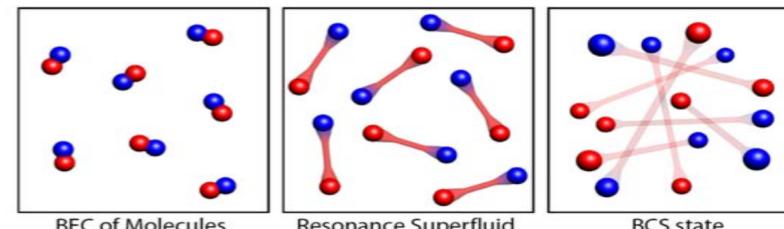
Pion Superfluidity @ Finite Isospin Density

- ▶ BEC-BCS crossover in pion superfluidity

- ultra cold atom
& condensed matter :

attractive strength

$$1/k_F a$$



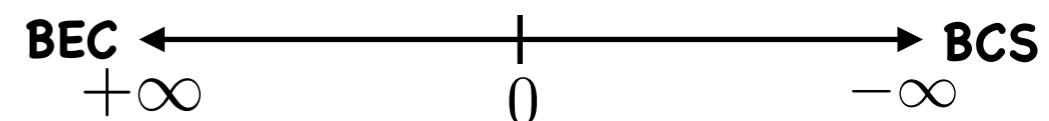
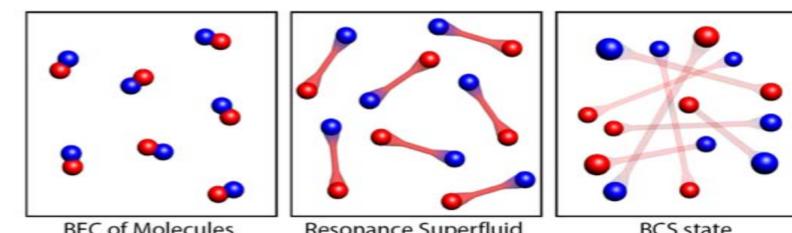
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- QCD system (relativistic) : color-superconductivity & pion superfluidity...
density induced: $T_c < T^*$

$T < T_c :$



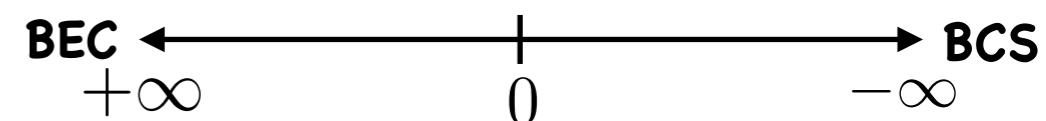
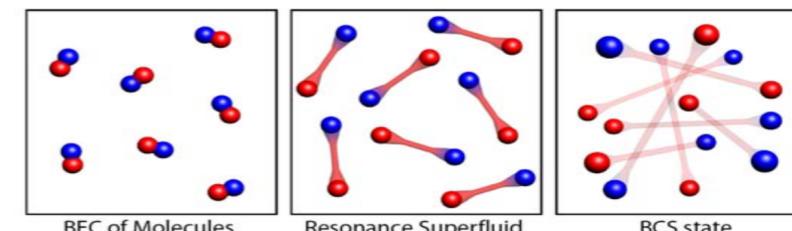
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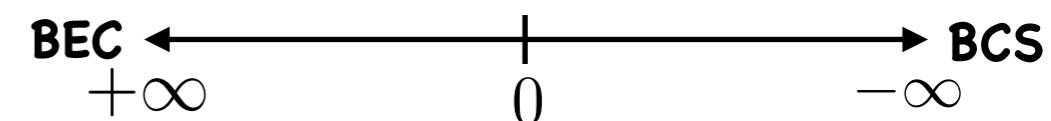
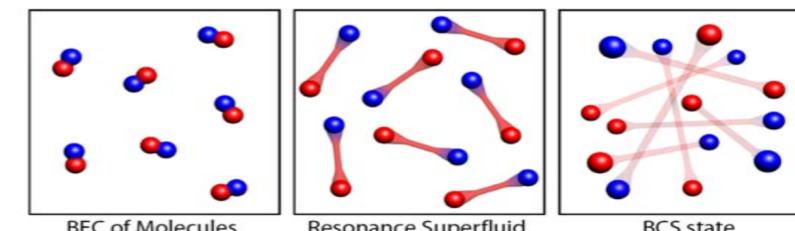
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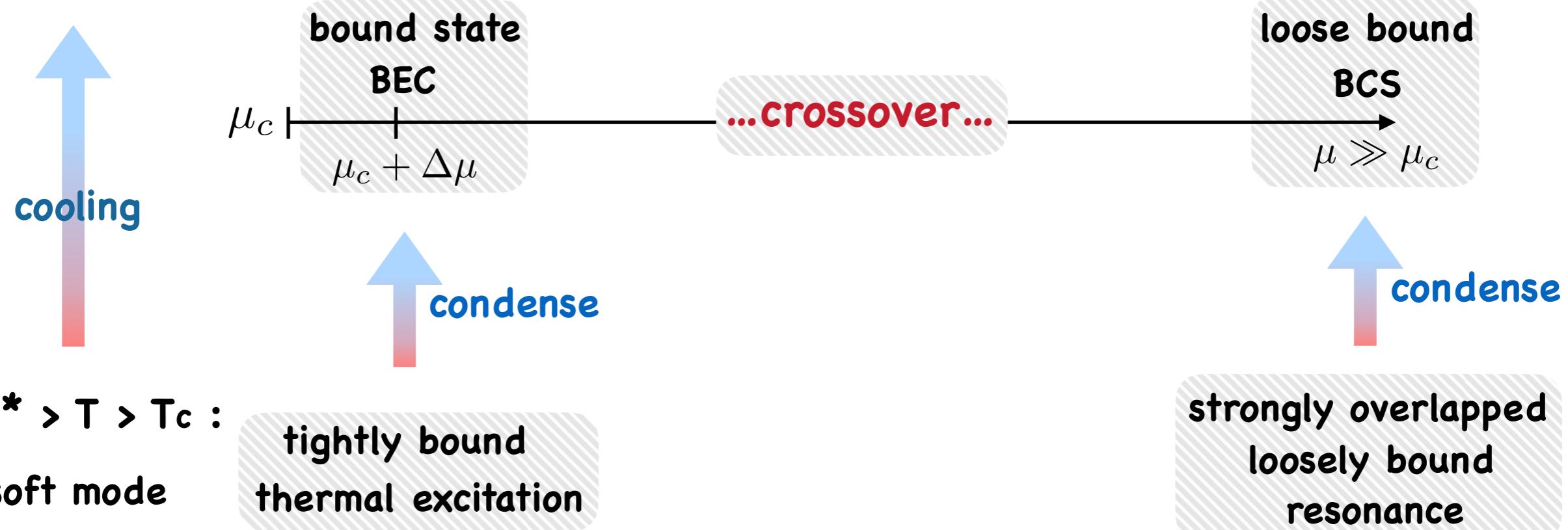


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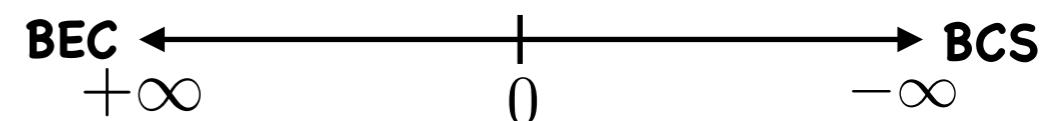
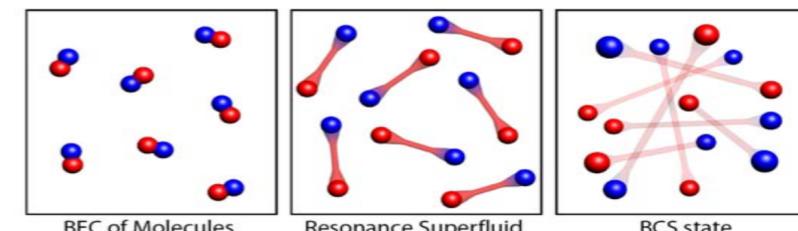
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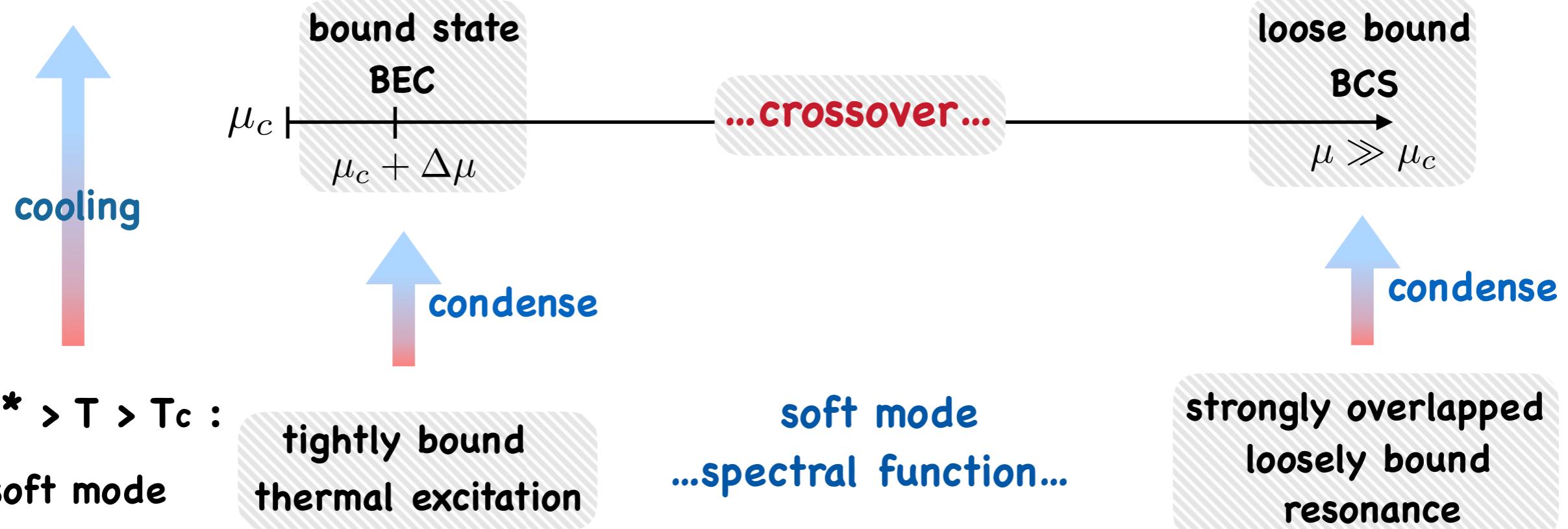


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& Theoretical Setup

► Functional Renormalization Group (FRG)

UV  IR

- strongly coupled many-body system
 - scale transformation (fixed point)
 - based on effective action
 - Analytical continuation
- phase transition
critical phenomenon
 $\Gamma \rightarrow \Gamma^{(2)} \dots \Gamma^{(n)}$
real-time observables

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 $\Gamma \rightarrow \Gamma^{(2)} \dots \Gamma^{(n)}$
real-time observables

► Linear Sigma Model & Quark Meson Model

$$\Gamma_k[\Phi] = \int_x \left\{ \bar{\Psi} S(i\partial) \Psi + \frac{1}{2}(\partial_\mu \sigma)^2 + \frac{1}{2}(\partial_\mu \pi_0)^2 + U_k(\vec{\varphi}^2) - c\sigma \right. \\ \left. + [(\partial_\mu - 2\delta_{\mu 4}\mu_I)\pi_-][(\partial_\mu + 2\delta_{\mu 4}\mu_I)\pi_+] \right\}$$

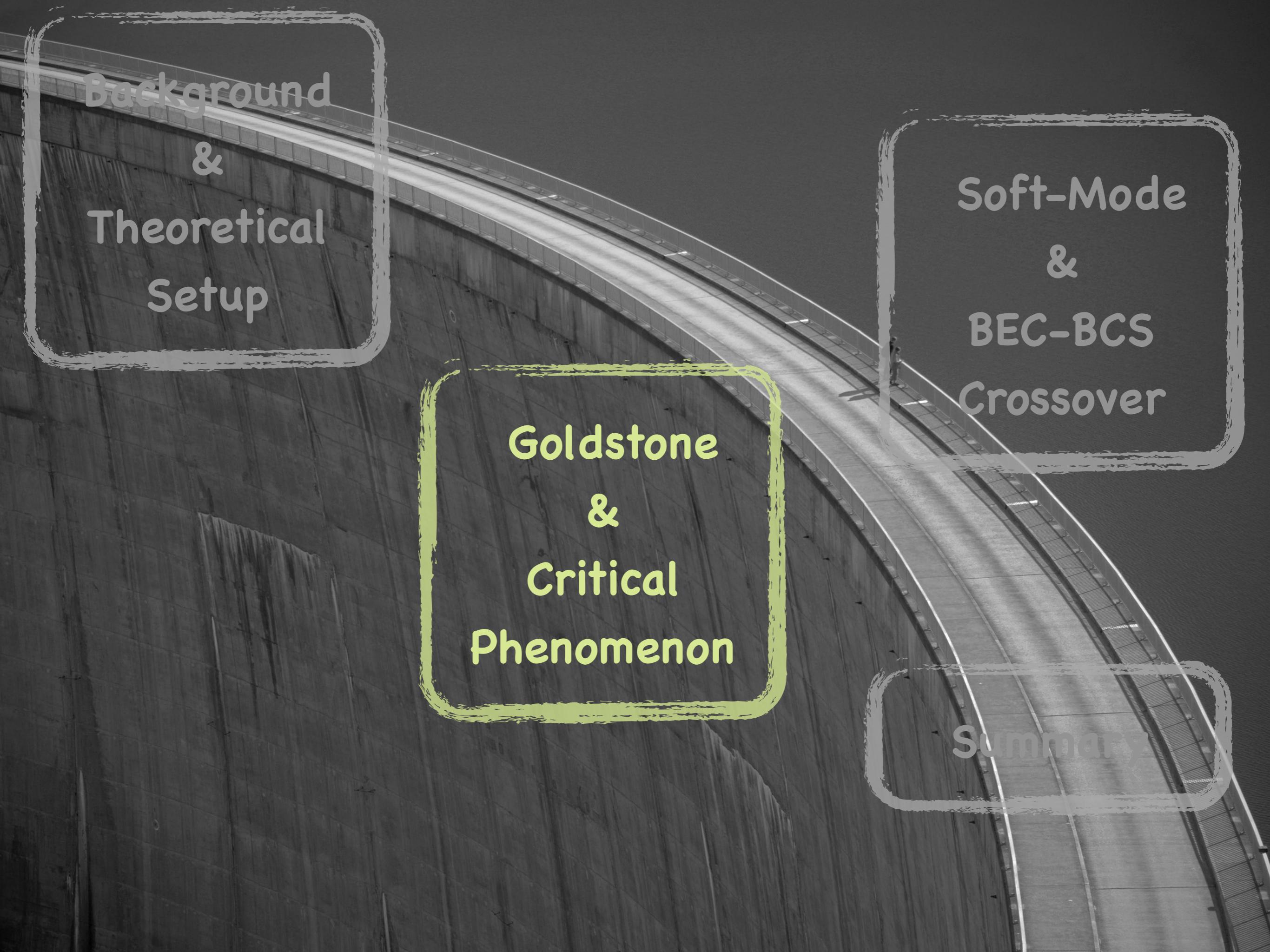
$$S(i\partial) = \begin{pmatrix} i(\partial + \gamma_4\mu_I) + ih(\sigma + i\gamma_5\pi_0) & -\sqrt{2}h\gamma_5\pi_- \\ -\sqrt{2}h\gamma_5\pi_+ & i(\partial - \gamma_4\mu_I) + ih(\sigma - i\gamma_5\pi_0) \end{pmatrix}.$$

Two flavor chiral effective model $\Psi = \begin{pmatrix} u \\ d \end{pmatrix}$

Symmetry breaking pattern: $O(4) - O(3) - O(2) - Z(2)$

► RG Flow Equation

$$\partial_k \Gamma_k = \frac{1}{2} \text{ (dashed circle with blue cross)} - \text{ (solid circle with red cross)}$$



Background
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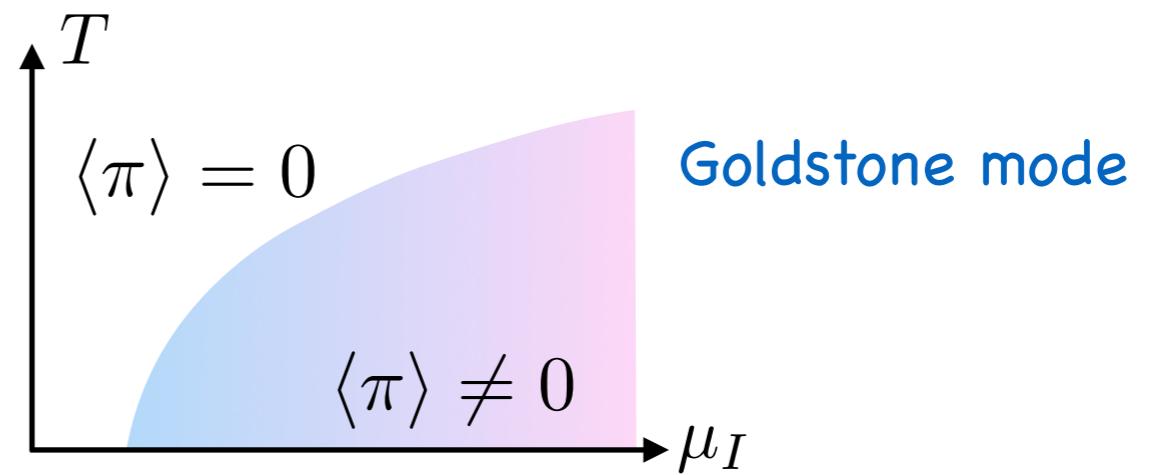
Summary

Goldstone Mode & Critical Phenomenon

- ▶ Symmetry breaking pattern: $O(4) - O(3) - O(2) - Z(2)$
- ▶ Second order phase transition
- ▶ Critical exponents & Universality class

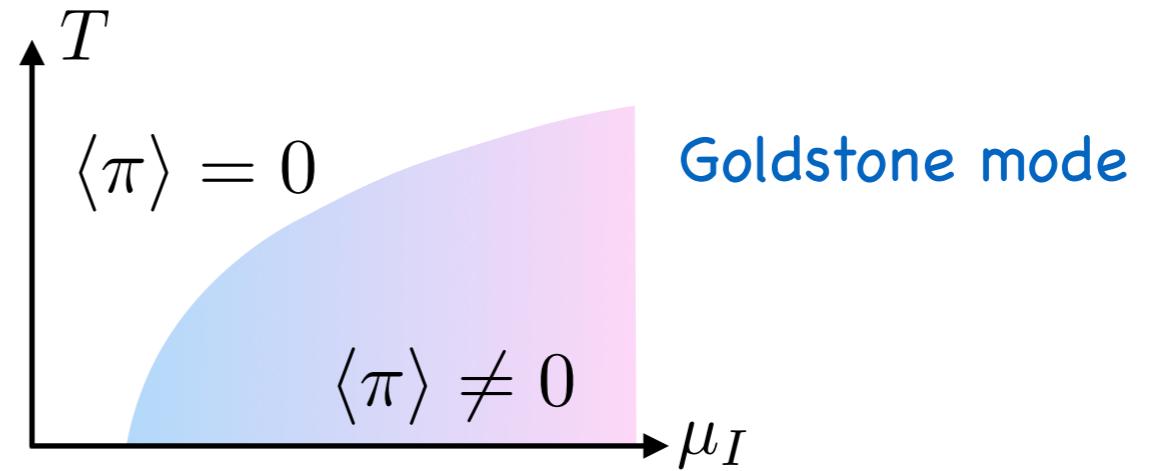


Goldstone mode



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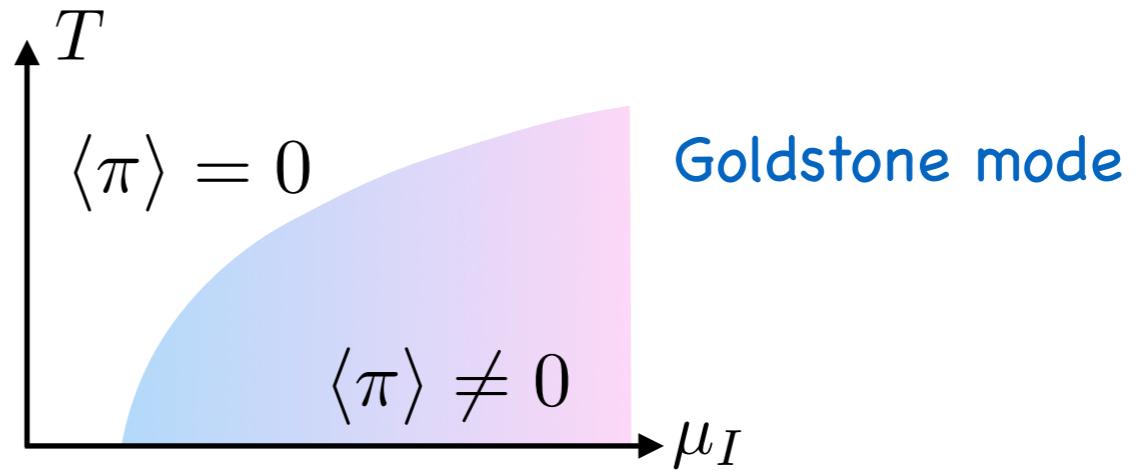


- ▶ LINEAR SIGMA Model + ISOSPIN
FRG & MEAN FIELD
- ▶ Fit critical exponents β, ν
along the phase boundary @ $T_c - \Delta T$

ZyW, Pengfei ZHUANG,
arXiv:1511.05279

Goldstone Mode & Critical Phenomenon

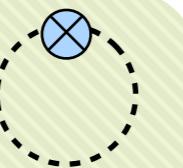
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FRG & MEAN FIELD
- ▶ Fit critical exponents β, ν
along the phase boundary @ $T_c - \Delta T$
- ▶ Compare with a general $O(N)$ MODEL, $N=2$, FIXED POINT
- ▶ $O(2)$ Universality Class with a dimension crossover

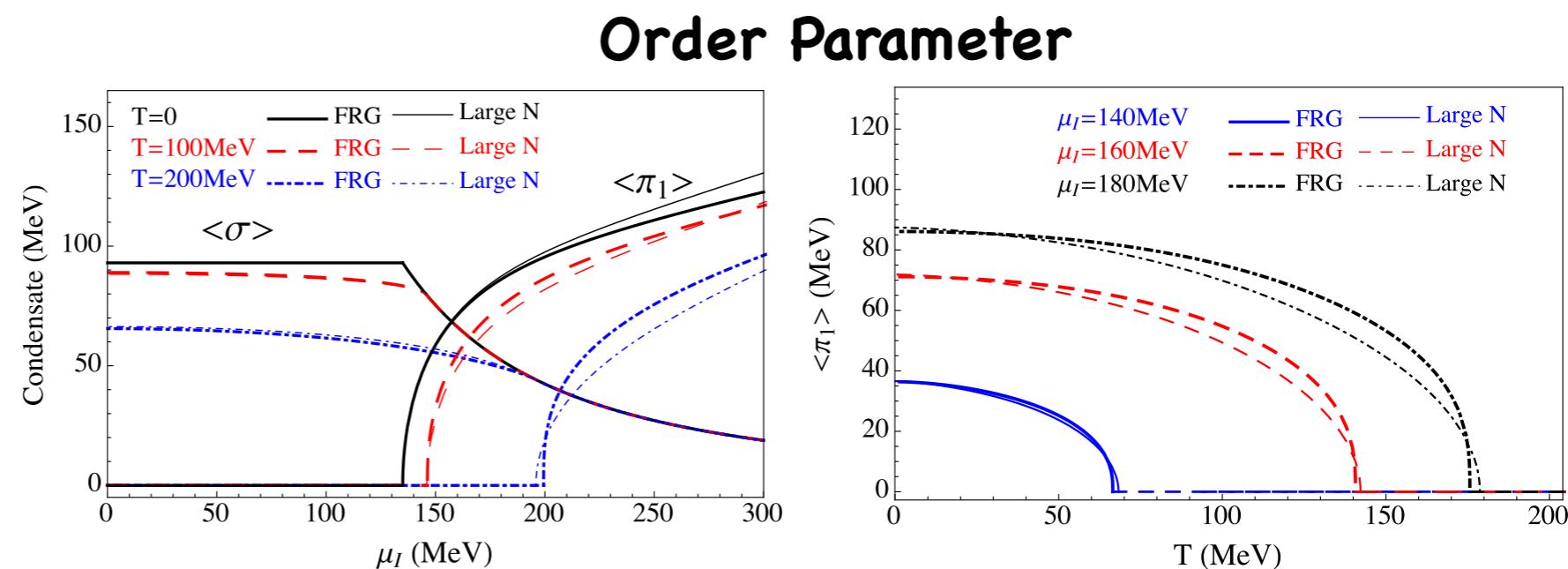
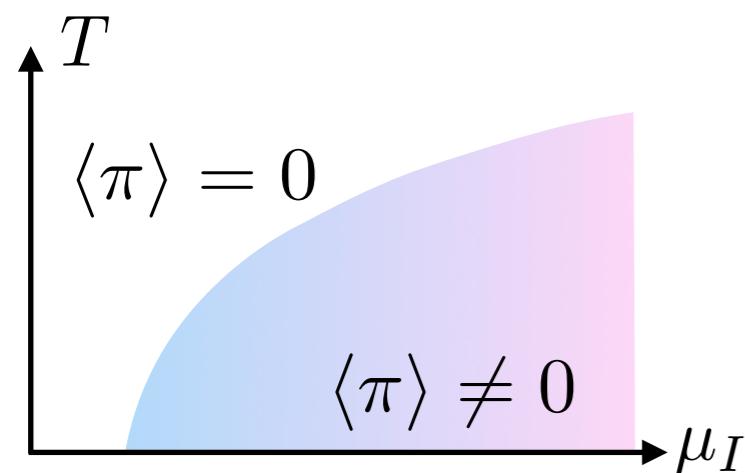
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- FRG : $\partial_k \Gamma_k = \frac{1}{2}$

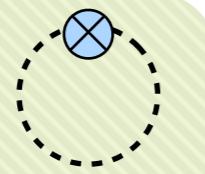


Taylor expansion

- MEAN FIELD :
Large-N expansion



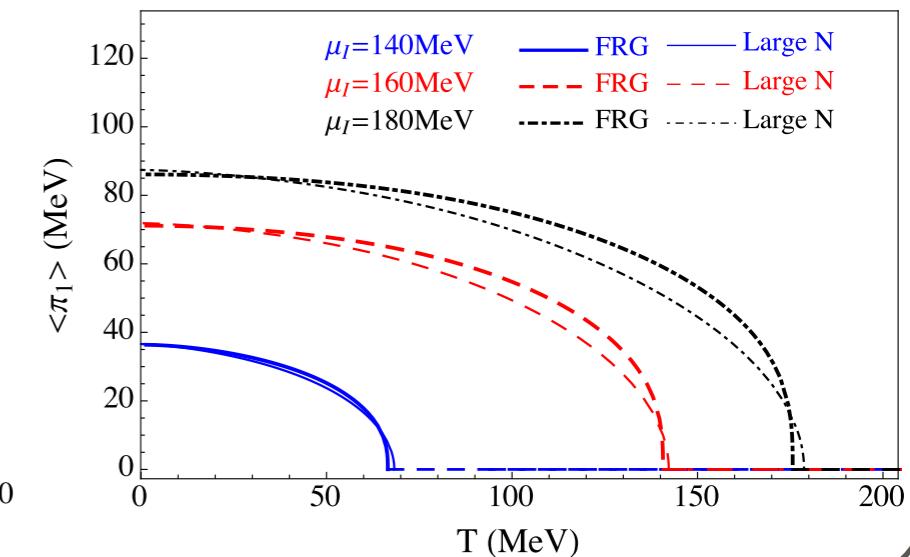
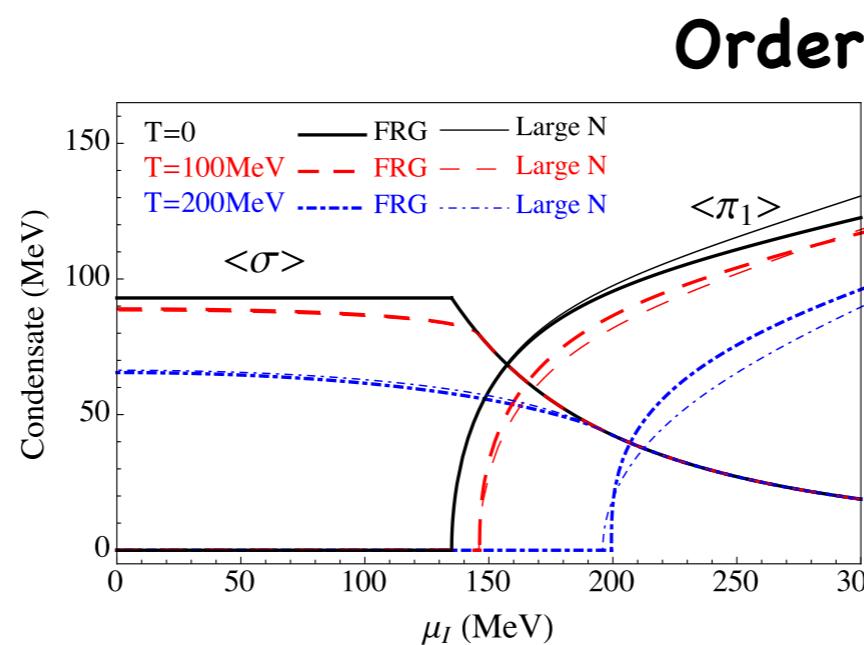
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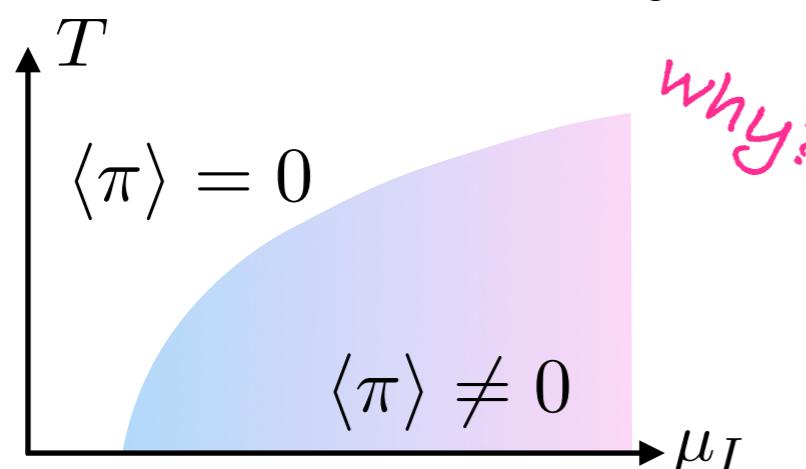


- Fit critical exponent:

$$\langle \pi \rangle \sim \left(\frac{\mu_I - \mu_I^c}{\mu_I^c} \right)^\beta \quad \text{or} \quad \langle \pi \rangle \sim \left(\frac{T_c - T}{T_c} \right)^\beta$$

- FRG:

β drops fast, then saturate at high T



Critical Exponent

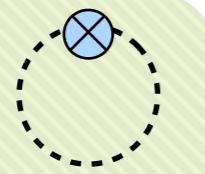
- FRG: 0.5 – 0.3

T (MeV)	0	10	50	100	150	200	250
μ_I^c (MeV)	135.0	135.2	138.0	146.3	164.5	199.4	248.9
β	0.5	0.445	0.380	0.347	0.328	0.318	0.314

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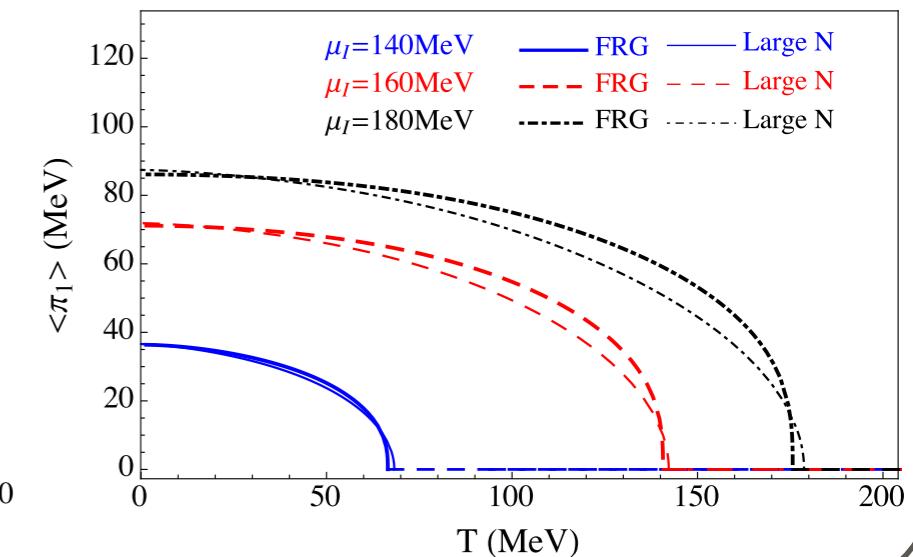
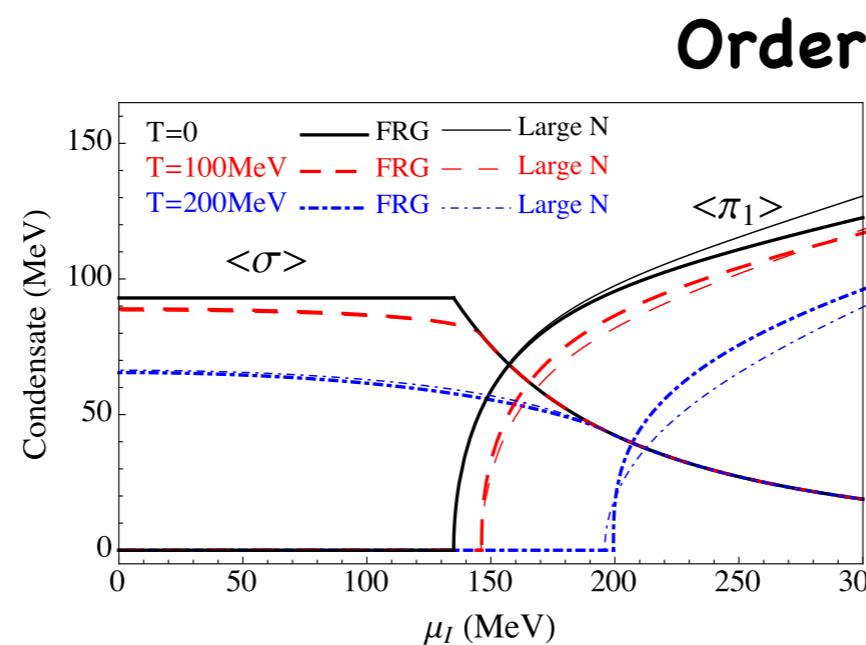
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- Dimension Reduction

$$\int_0^\infty d^d x \rightarrow \int_0^{T^{-1}} dt \int_0^\infty d^{d-1} \mathbf{x}$$

T=0 to high T limit: $S^1 \times R^{d-1} \rightarrow R^{d-1}$

$$d_{\text{eff}} = d \rightarrow d_{\text{eff}} = d - 1$$

Critical Exponent

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► Dimension Reduction

$T=0$ to high T limit: $d_{\text{eff}} = d \rightarrow d_{\text{eff}} = d - 1$

► PION SUPERFLUIDITY

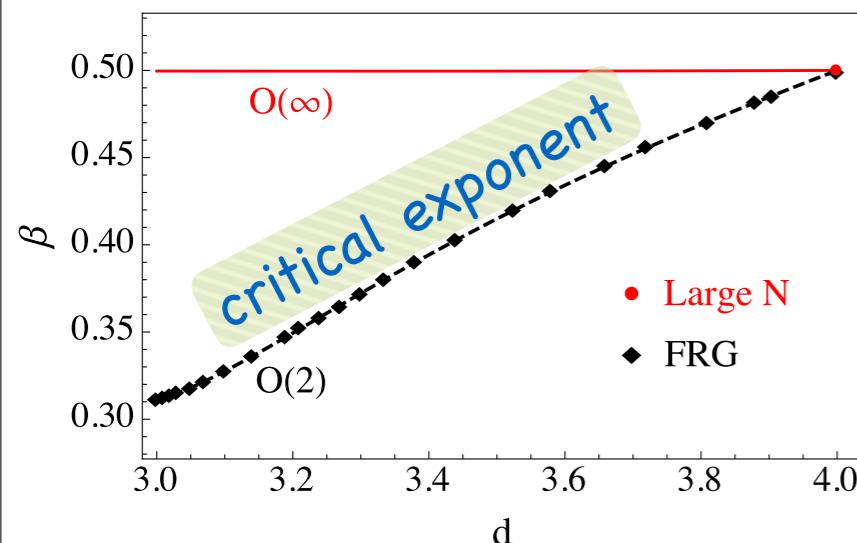
@ finite T

► General O(2) MODEL

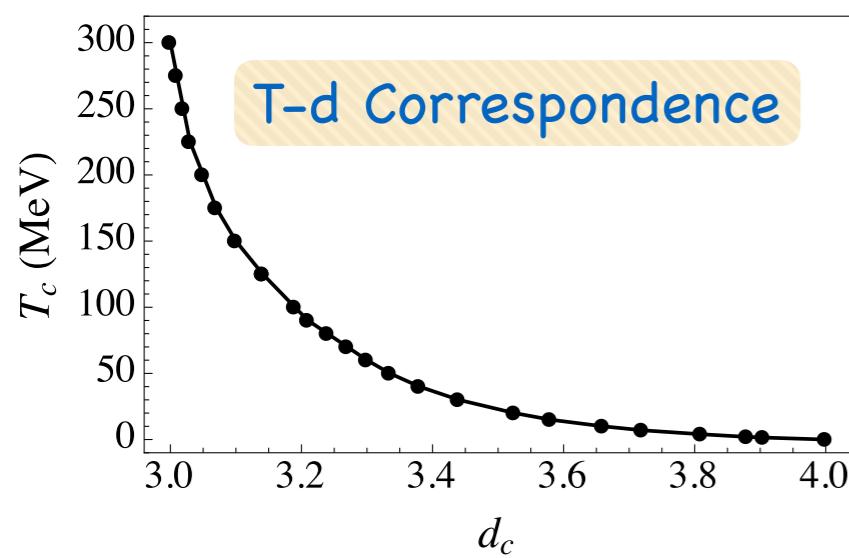
@ continuous d

Same symmetry

T-d Correspondence



O(2) universality class
with a
dimension crossover

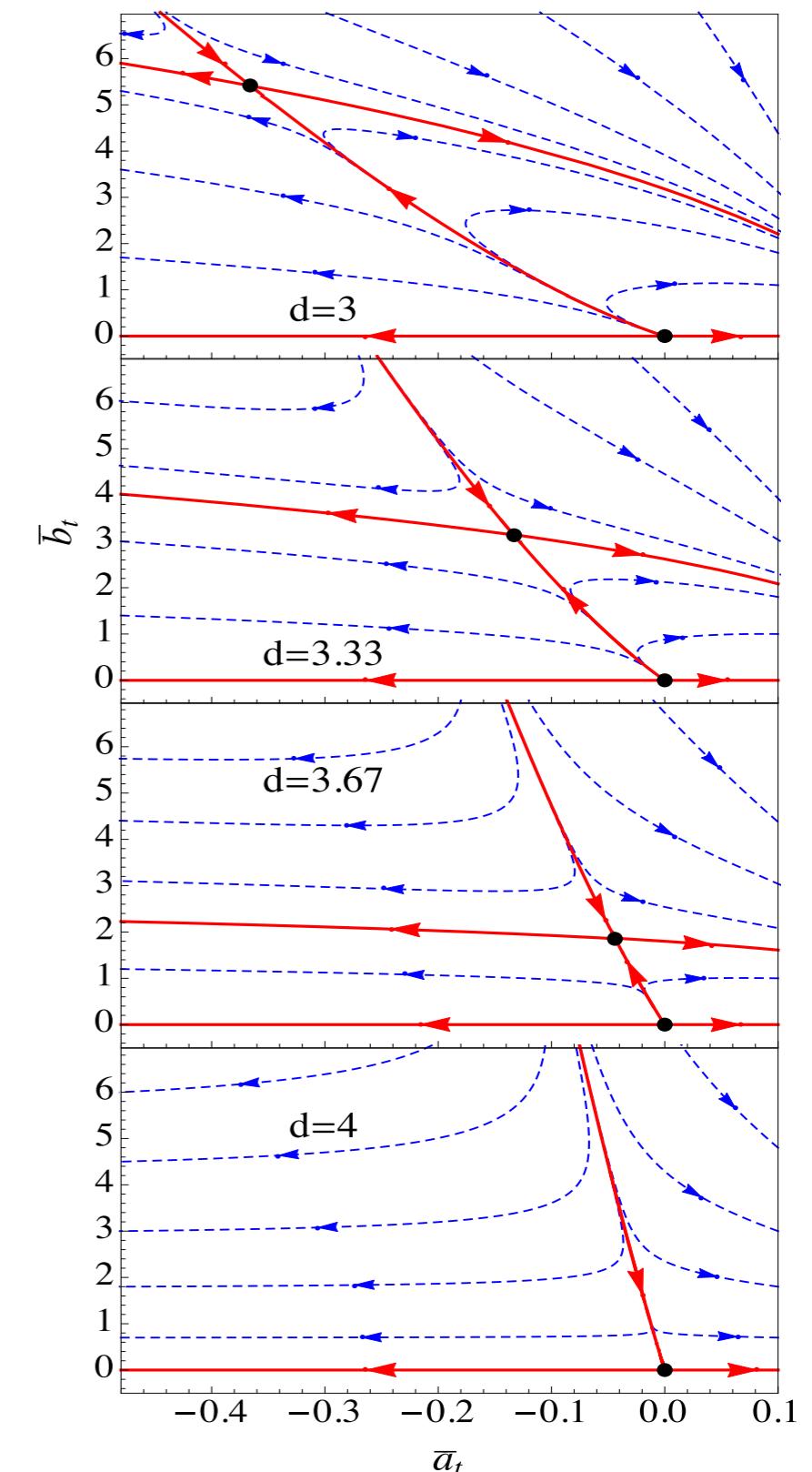


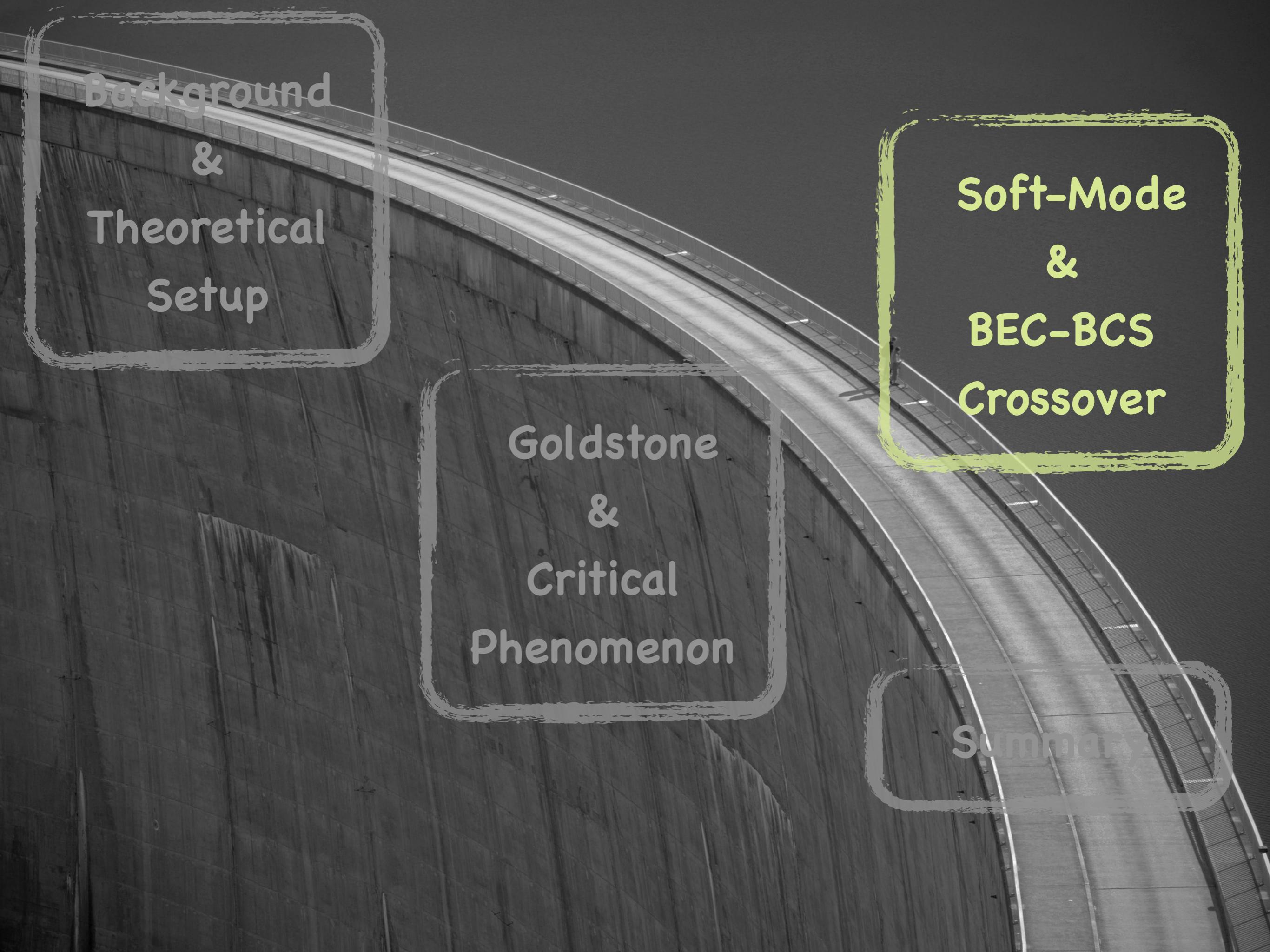
T-d relation
not universal

$$T = 0, d = 4, \beta = 0.5$$

General O(2) MODEL

@ continuous d





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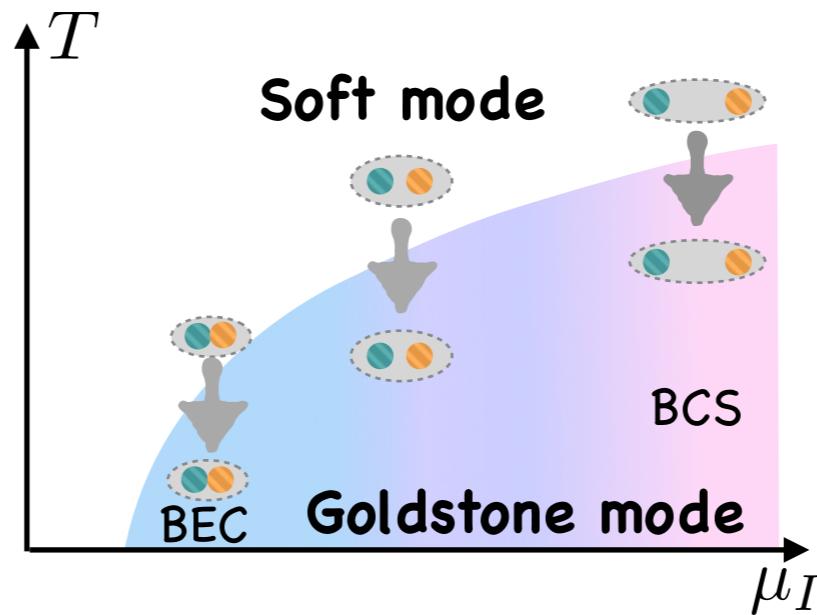
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Summary

Soft-Mode BEC-BCS Crossover

tightly bound
thermal excitation

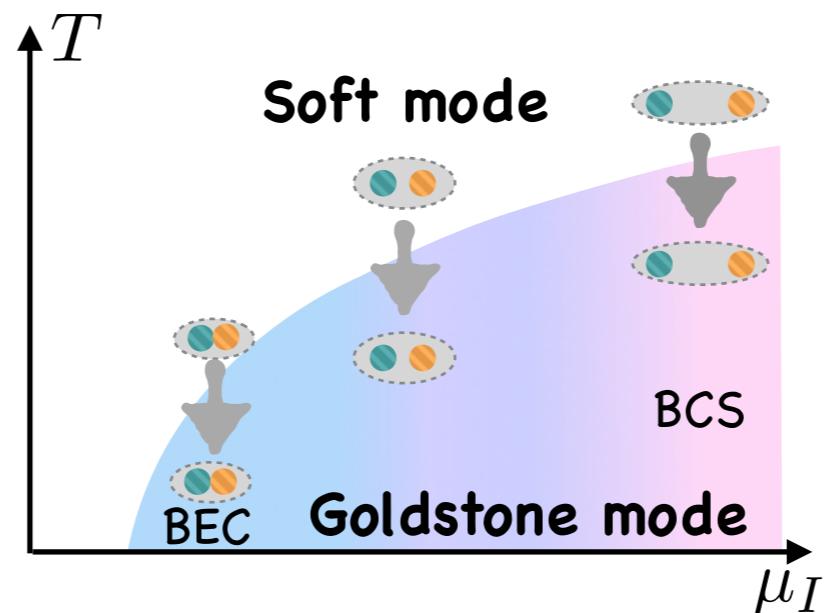
- BEC-BCS Crossover: density & chiral symmetry restoration
- Soft mode @ $T^* > T > T_c$



strongly overlapped
loosely bound
resonance

Soft-Mode BEC-BCS Crossover

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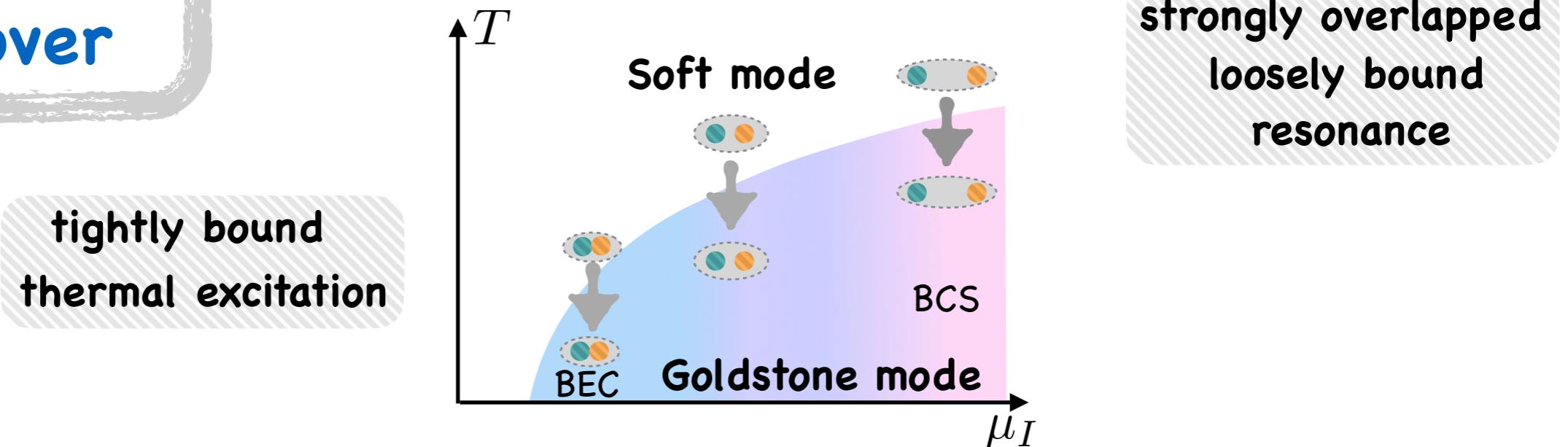
- Meson SPECTRAL function in QUARK MESON model @ $T^* > T > T_c$

Free particle: pole \rightarrow mass

Interaction: pole, branch-cut \rightarrow mass, stability, decay channel

Soft-Mode BEC-BCS Crossover

- BEC-BCS Crossover: density & chiral symmetry restoration
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- Meson SPECTRAL function in QUARK MESON model @ $T^* > T > T_c$

Free particle: pole \rightarrow mass

Interaction: pole, branch-cut \rightarrow mass, stability, decay channel

- Flow equation of meson 2-point function

[K. Kamikado, et al., Eur. Phys. J. C 74, 2806 (2014)]

[R. Tripolt, et al, Phys. Rev. D 89, 034010 (2014)]

[R. Tripolt, et al, Phys. Rev. D 90, 074031 (2014)]

- Analytical continuation $\Gamma^{(2),R}(\omega, \vec{p}) = \lim_{\epsilon \rightarrow 0} \Gamma^{(2),E}(p_0 = -i(\omega + i\epsilon), \vec{p})$

- Spectral function

$$\rho(\omega) = -\frac{1}{\pi} \frac{\text{Im}\Gamma^{(2),R}(\omega)}{\left(\text{Re}\Gamma^{(2),R}(\omega)\right)^2 + \left(\text{Im}\Gamma^{(2),R}(\omega)\right)^2}$$

Flow equation of meson 2-point function

Truncation: LPA, neglect momentum dependence of $\Gamma^{(3)}, \Gamma^{(4)}$

$$\partial_k \Gamma_{\sigma\sigma}^{(2)} = \begin{aligned} & \sigma \text{---} \textcircled{\times} \sigma \text{---} \sigma - \frac{1}{2} \text{---} \textcircled{\times} \sigma \text{---} \sigma + \sigma \text{---} \textcircled{\times} \pi_0 \text{---} \sigma - \frac{1}{2} \text{---} \textcircled{\times} \pi_0 \text{---} \sigma \\ & + \sigma \text{---} \textcircled{\times} \pi_+ \text{---} \sigma - \frac{1}{2} \text{---} \textcircled{\times} \pi_+ \text{---} \sigma + \sigma \text{---} \textcircled{\times} \pi_- \text{---} \sigma - \frac{1}{2} \text{---} \textcircled{\times} \pi_- \text{---} \sigma \\ & -2 \sigma \text{---} \textcircled{\times} u \bar{u} \text{---} \sigma -2 \sigma \text{---} \textcircled{\times} d \bar{d} \text{---} \sigma \end{aligned}$$

$$\partial_k \Gamma_{\pi_0\pi_0}^{(2)} = \begin{aligned} & \pi_0 \text{---} \textcircled{\times} \pi_0 \text{---} \pi_0 + \pi_0 \text{---} \textcircled{\times} \pi_0 \text{---} \pi_0 \\ & - \frac{1}{2} \text{---} \textcircled{\times} \sigma \text{---} \pi_0 \pi_0 - \frac{1}{2} \text{---} \textcircled{\times} \pi_+ \text{---} \pi_0 \pi_0 - \frac{1}{2} \text{---} \textcircled{\times} \pi_- \text{---} \pi_0 \pi_0 - \frac{1}{2} \text{---} \textcircled{\times} \pi_0 \text{---} \pi_0 \pi_0 \\ & -2 \pi_0 \text{---} \textcircled{\times} u \bar{u} \text{---} \pi_0 -2 \pi_0 \text{---} \textcircled{\times} d \bar{d} \text{---} \pi_0 \end{aligned}$$

$$\partial_k \Gamma_{\pi_+\pi_-}^{(2)} = \begin{aligned} & \pi_+ \text{---} \textcircled{\times} \sigma \text{---} \pi_- + \pi_+ \text{---} \textcircled{\times} \pi_- \text{---} \pi_- \\ & - \frac{1}{2} \text{---} \textcircled{\times} \sigma \text{---} \pi_+ \pi_- - \frac{1}{2} \text{---} \textcircled{\times} \pi_+ \text{---} \pi_+ \pi_- - \frac{1}{2} \text{---} \textcircled{\times} \pi_- \text{---} \pi_+ \pi_- - \frac{1}{2} \text{---} \textcircled{\times} \pi_0 \text{---} \pi_+ \pi_- \\ & -4 \pi_+ \text{---} \textcircled{\times} u \bar{d} \text{---} \pi_- \end{aligned}$$

$$\partial_k \Gamma_{\pi_-\pi_+}^{(2)} = \begin{aligned} & \pi_- \text{---} \textcircled{\times} \sigma \text{---} \pi_+ + \pi_- \text{---} \textcircled{\times} \pi_+ \text{---} \pi_+ \\ & - \frac{1}{2} \text{---} \textcircled{\times} \sigma \text{---} \pi_- \pi_+ - \frac{1}{2} \text{---} \textcircled{\times} \pi_+ \text{---} \pi_- \pi_+ - \frac{1}{2} \text{---} \textcircled{\times} \pi_- \text{---} \pi_- \pi_+ - \frac{1}{2} \text{---} \textcircled{\times} \pi_0 \text{---} \pi_- \pi_+ \\ & -4 \pi_- \text{---} \textcircled{\times} \bar{u} d \text{---} \pi_+ \end{aligned}$$

► Flow equation of meson 2-point function

Truncation: LPA, neglect momentum dependence of $\Gamma^{(3)}, \Gamma^{(4)}$

$$\partial_k \Gamma_{\sigma\sigma}^{(2)} = \sigma \circlearrowleft \sigma - \frac{1}{2} \sigma \circlearrowleft \sigma + \sigma \circlearrowleft \pi_0 - \frac{1}{2} \pi_0 \circlearrowleft \sigma \\ + \sigma \circlearrowleft \pi_+ - \frac{1}{2} \pi_+ \circlearrowleft \sigma + \sigma \circlearrowleft \pi_- - \frac{1}{2} \pi_- \circlearrowleft \sigma \\ - 2 \sigma \circlearrowleft u \bar{u}$$

$$\partial_k \Gamma_{\pi_0\pi_0}^{(2)} = \pi_0 \circlearrowleft \pi_0 - \frac{1}{2} \pi_0 \circlearrowleft \pi_0 - \frac{1}{2} \pi_0 \circlearrowleft \pi_+ - \frac{1}{2} \pi_0 \circlearrowleft \pi_- \\ - 2 \pi_0 \circlearrowleft u \bar{u} - 2 \pi_0 \circlearrowleft d \bar{d}$$

Chiral Symmetry Restoration

$$\partial_k \Gamma_{\pi_+\pi_-}^{(2)} = \pi_+ \circlearrowleft \sigma \pi_- + \pi_+ \circlearrowleft \pi_- \\ - \frac{1}{2} \pi_+ \circlearrowleft \sigma \pi_- - \frac{1}{2} \pi_+ \circlearrowleft \pi_+ \pi_- - \frac{1}{2} \pi_+ \circlearrowleft \pi_- \pi_- - \frac{1}{2} \pi_+ \circlearrowleft \pi_0 \pi_- \\ - 4 \pi_+ \circlearrowleft u \bar{d} \pi_-$$

$$\partial_k \Gamma_{\pi_-\pi_+}^{(2)} = \pi_- \circlearrowleft \sigma \pi_+ + \pi_- \circlearrowleft \pi_+ \pi_+ \\ - \frac{1}{2} \pi_- \circlearrowleft \sigma \pi_+ - \frac{1}{2} \pi_- \circlearrowleft \pi_- \pi_+ - \frac{1}{2} \pi_- \circlearrowleft \pi_0 \pi_+ \\ - 4 \pi_- \circlearrowleft \bar{u} d \pi_+$$

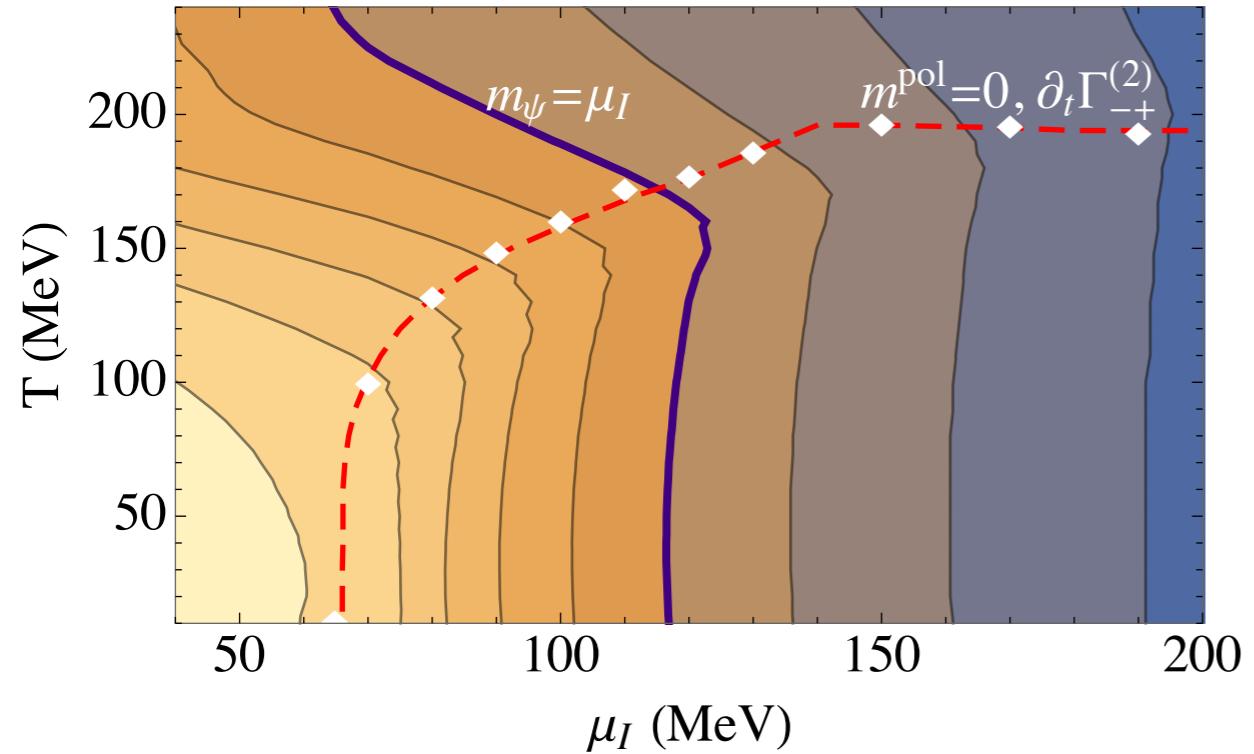
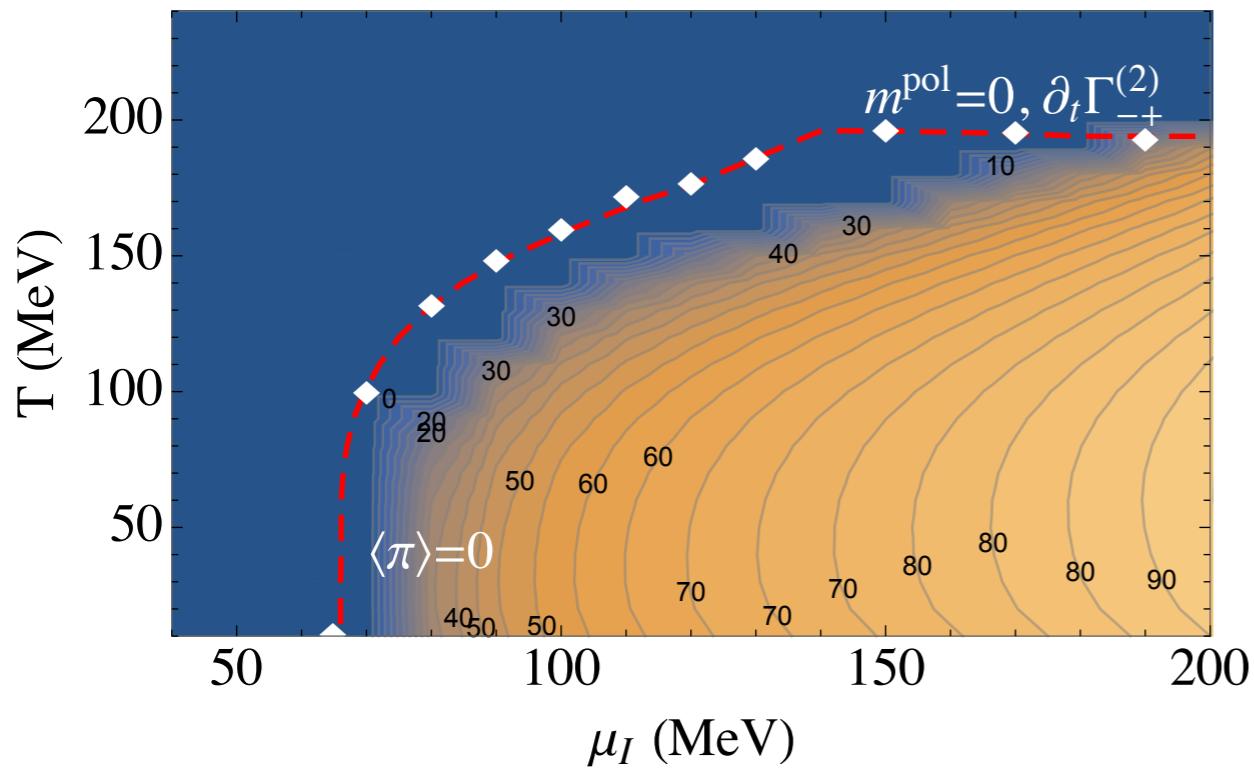
► Initial condition & Input

$$U_\Lambda(\rho) = \frac{1}{2} m_\Lambda^2 \rho + \frac{1}{4} \lambda_\Lambda \rho^2$$

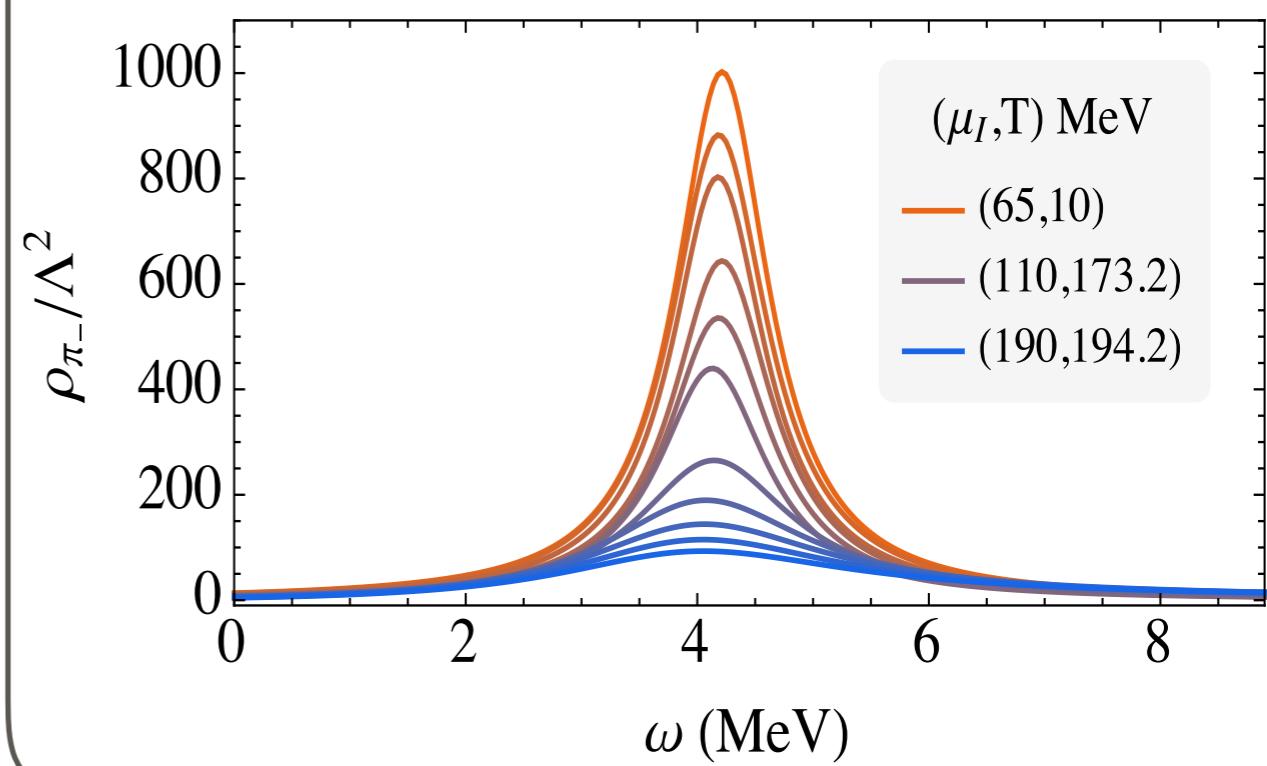
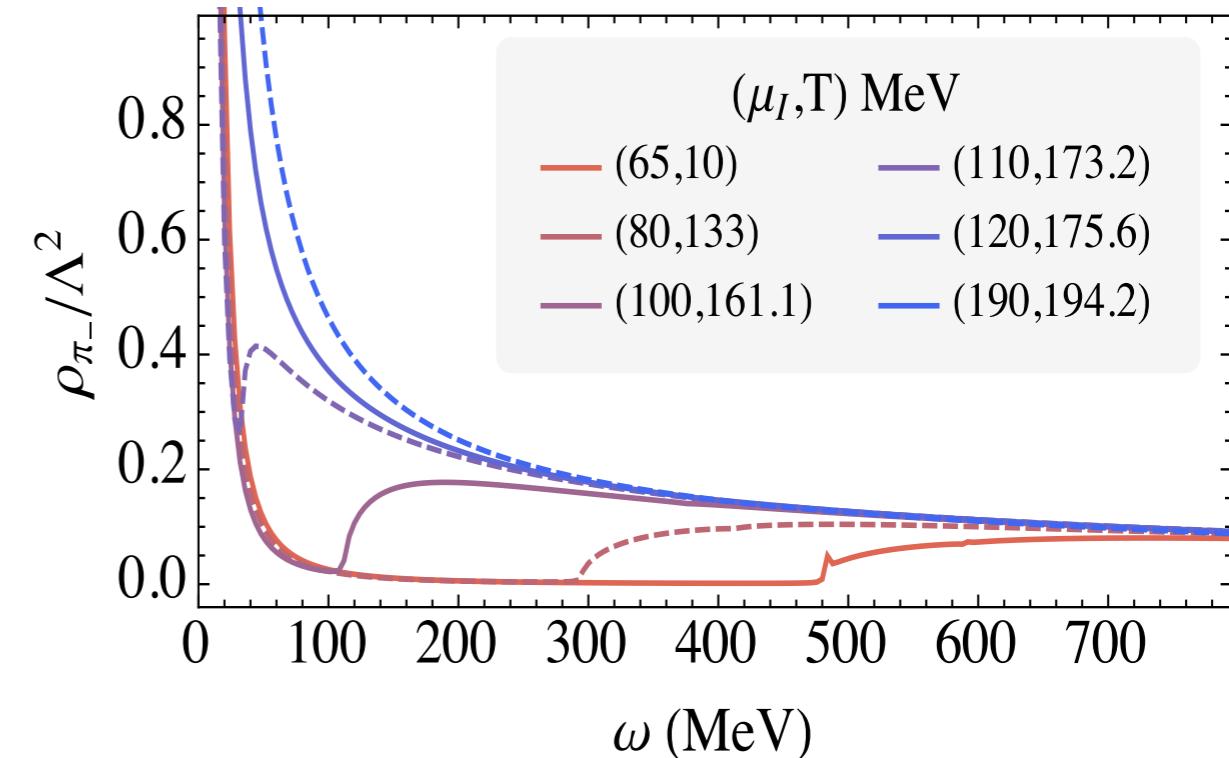
$$\Gamma_{\sigma\sigma,\Lambda}^{(2),R} = -\omega^2 + 2U'_\Lambda + 4\phi^2 U''_\Lambda, \\ \Gamma_{\pi_+\pi_-,\Lambda}^{(2),R} = -(\omega - 2\mu_I)^2 + 2U'_\Lambda,$$

$$\Gamma_{\pi_0\pi_0,\Lambda}^{(2),R} = -\omega^2 + 2U'_\Lambda \\ \Gamma_{\pi_-\pi_+,\Lambda}^{(2),R} = -(\omega + 2\mu_I)^2 + 2U'_\Lambda$$

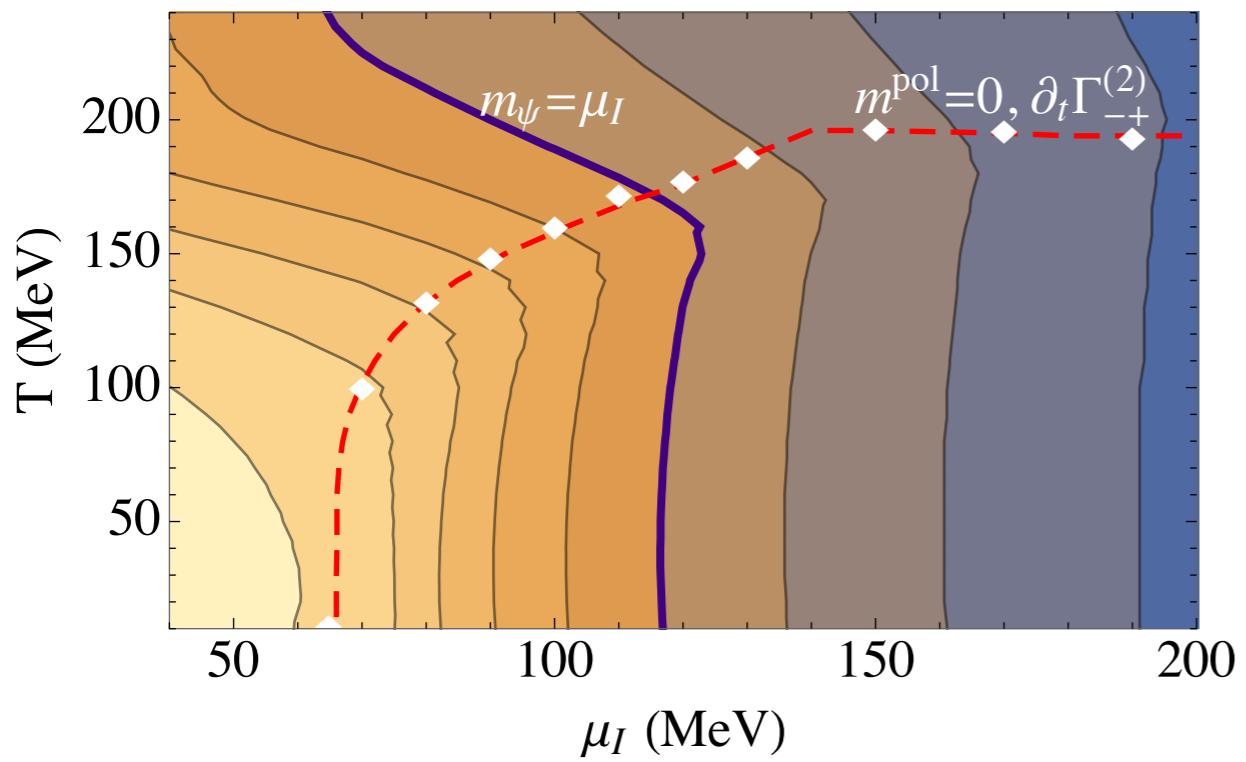
Phase diagram



Spectral function of soft mode



Phase diagram

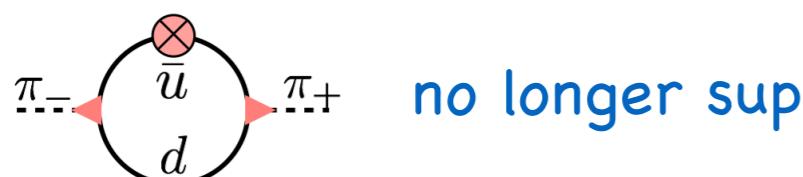


- ▶ Sharp peak in BEC limit



is suppressed

- ▶ Broad resonance in BCS limit

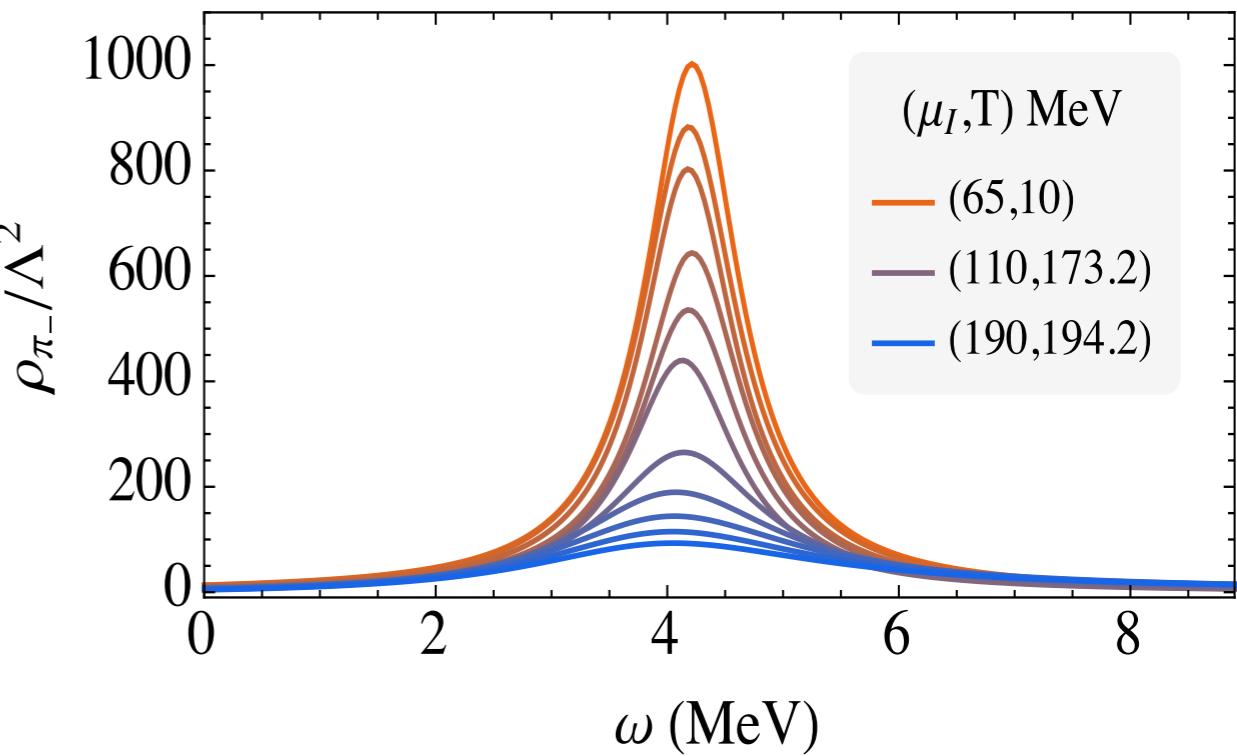
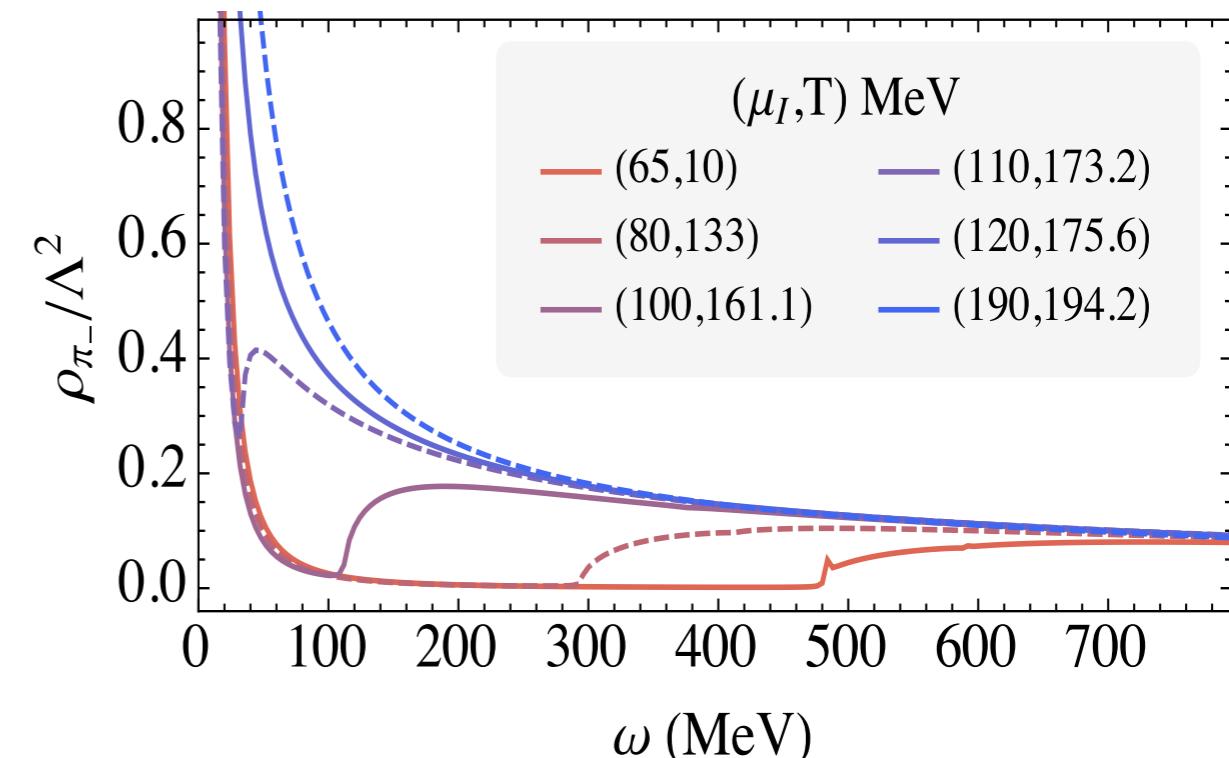


no longer suppressed

- ▶ 2-quark channel opened already at $\omega > 2(m_q - \mu_I)$

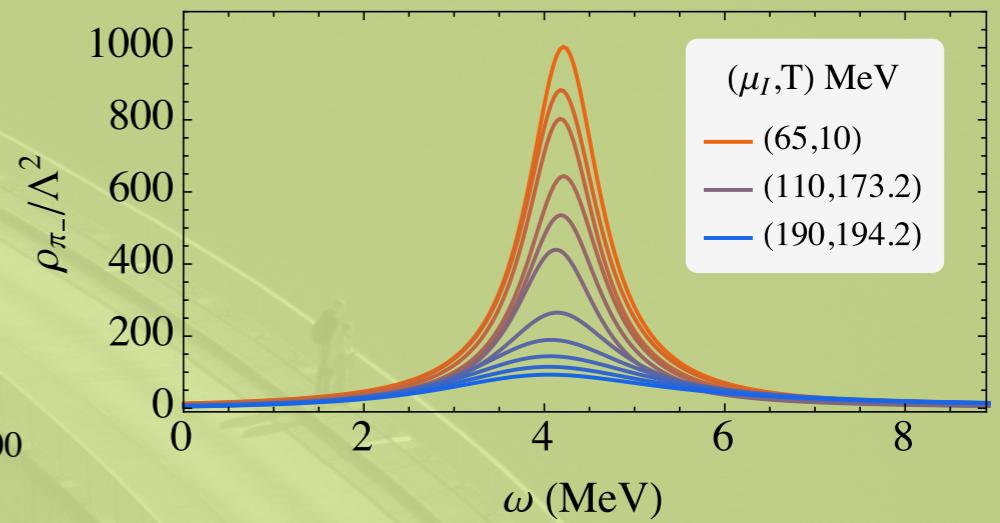
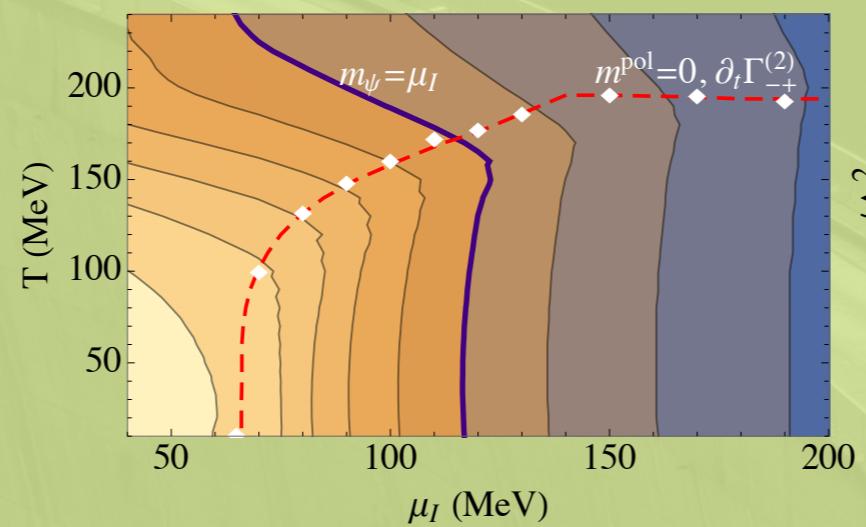
$$\omega = 0$$

Spectral function of soft mode

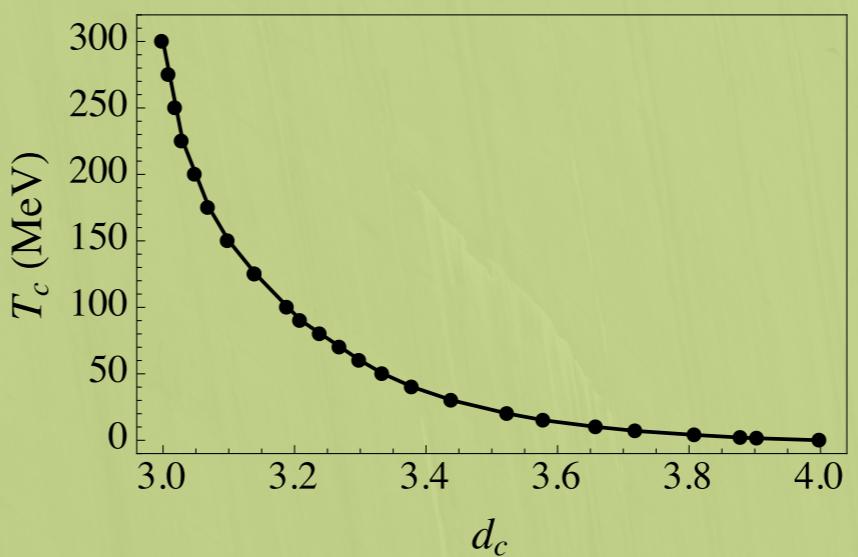
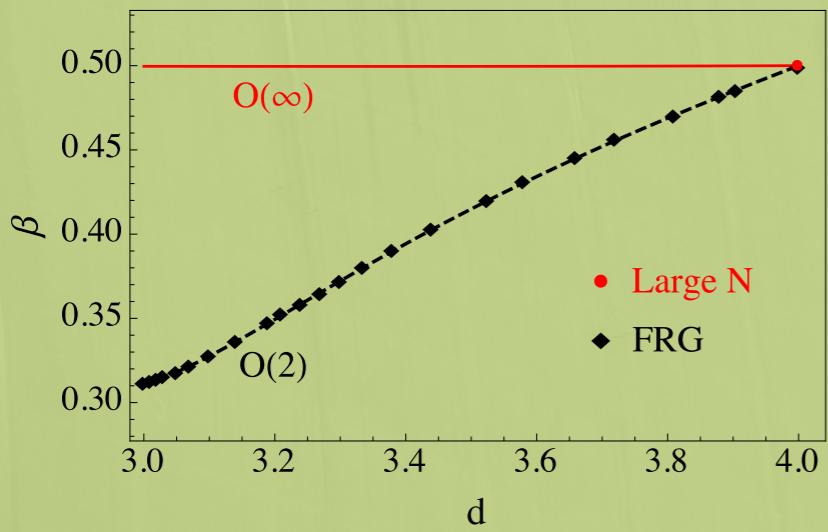


summary

Soft Mode – BEC-BCS Crossover



Goldstone Mode – Critical Phenomenon





THANK YOU!

$$\begin{aligned}\partial_k \Gamma_{\sigma\sigma}^{(2)}(p_4) = & -\frac{1}{2} \Gamma_{\sigma\sigma\sigma\sigma}^{(4)} I_\sigma^{(2)} - \frac{1}{2} \Gamma_{\sigma\sigma\pi\pi}^{(4)} I_{\pi_0}^{(2)} - \frac{1}{2} \Gamma_{\sigma\sigma\pi_+\pi_-}^{(4)} \left(I_{\pi_+}^{(2)} + I_{\pi_-}^{(2)} \right) \\ & + (\Gamma_{\sigma\sigma\sigma}^{(3)})^2 J_{\sigma\sigma}(p_4) + (\Gamma_{\sigma\pi\pi}^{(3)})^2 \left(J_{\pi_0\pi_0}(p_4) + J_{++}(p_4) + J_{--}(p_4) \right) \\ & - 2N_c h^2 \left(F_{++}^\sigma(p_4, \mu_I) + F_{--}^\sigma(p_4, \mu_I) \right)\end{aligned}$$

$$\begin{aligned}\partial_k \Gamma_{\pi_0\pi_0}^{(2)}(p_4) = & -\frac{1}{2} \Gamma_{\sigma\sigma\pi\pi}^{(4)} I_\sigma^{(2)} - \frac{1}{2} \Gamma_{\pi_0\pi_0\pi_0\pi_0}^{(4)} I_{\pi_0}^{(2)} - \frac{1}{2} \Gamma_{\pi_0\pi_0\pi_+\pi_-}^{(4)} \left(I_{\pi_+}^{(2)} + I_{\pi_-}^{(2)} \right) \\ & + \frac{1}{2} (\Gamma_{\sigma\pi\pi}^{(3)})^2 \left(J_{\sigma\pi_0}(p_4) + J_{\pi_0\sigma}(p_4) \right) \\ & - 2N_c h^2 \left(F_{++}^\pi(p_4, \mu_I) + F_{--}^\pi(p_4, \mu_I) \right)\end{aligned}$$

$$\begin{aligned}\partial_k \Gamma_{\pi_+\pi_-}^{(2)}(p_4) = & -\frac{1}{2} \Gamma_{\sigma\sigma\pi_+\pi_-}^{(4)} I_\sigma^{(2)} - \frac{1}{2} \Gamma_{\pi_0\pi_0\pi_+\pi_-}^{(4)} I_{\pi_0}^{(2)} - \frac{1}{2} \Gamma_{\pi_+\pi_-\pi_+\pi_-}^{(4)} \left(I_{\pi_+}^{(2)} + I_{\pi_-}^{(2)} \right) \\ & + (\Gamma_{\sigma\pi\pi}^{(3)})^2 \left(J_{\sigma-}(p_4) + J_{+\sigma}(p_4) \right) \\ & - 4N_c h^2 F_{-+}^\pi(p_4, \mu_I)\end{aligned}$$

$$\begin{aligned}\partial_k \Gamma_{\pi_-\pi_+}^{(2)}(p_4) = & -\frac{1}{2} \Gamma_{\sigma\sigma\pi_+\pi_-}^{(4)} I_\sigma^{(2)} - \frac{1}{2} \Gamma_{\pi_0\pi_0\pi_+\pi_-}^{(4)} I_{\pi_0}^{(2)} - \frac{1}{2} \Gamma_{\pi_+\pi_-\pi_+\pi_-}^{(4)} \left(I_{\pi_+}^{(2)} + I_{\pi_-}^{(2)} \right) \\ & + (\Gamma_{\sigma\pi\pi}^{(3)})^2 \left(J_{-\sigma}(p_4) + J_x \sigma+(p_4) \right) \\ & - 4N_c h^2 F_{+-}^\pi(p_4, \mu_I)\end{aligned}$$

$$I_{\alpha}^{(2)}(\mu_I) = \frac{k^4}{3\pi^2} \frac{1}{4} \left\{ \frac{1 + n_B(E_{\alpha} - 2\mu_I) + n_B(E_{\alpha} + 2\mu_I)}{E_{\alpha}^3} - \frac{n'_B(E_{\alpha} - 2\mu_I) + n'_B(E_{\alpha} + 2\mu_I)}{E_{\alpha}^2} \right\}$$

$$\begin{aligned} J_{\alpha\alpha}(p_4, \mu_I) = & \frac{k^4}{3\pi^2} \frac{1}{4} \left\{ \frac{12E_{\alpha}^2 + p_4^2}{E_{\alpha}^3(4E_{\alpha}^2 + p_4^2)^2} (1 + n_B(E_{\alpha} + 2\mu_I) + n_B(E_{\alpha} - 2\mu_I)) \right. \\ & \left. - \frac{1}{E_{\alpha}^2(2E_{\alpha} - ip_4)ip_4} n'_B(E_{\alpha} - 2\mu_I) + \frac{1}{E_{\alpha}^2(2E_{\alpha} + ip_4)ip_4} n'_B(E_{\alpha} + 2\mu_I) \right\} \end{aligned}$$

$$\begin{aligned} J_{\alpha\beta}(p_4) = & \frac{k^4}{3\pi^2} \frac{1}{4} \left\{ \frac{E_{\beta}^2 - (E_{\alpha} - ip_4)(3E_{\alpha} - ip_4)}{E_{\alpha}^3((E_{\alpha} - ip_4)^2 - E_{\beta}^2)^2} (1 + n_B(E_{\alpha})) + \frac{E_{\beta}^2 - (E_{\alpha} + ip_4)(3E_{\alpha} + ip_4)}{E_{\alpha}^3((E_{\alpha} + ip_4)^2 - E_{\beta}^2)^2} n_B(E_{\alpha}) \right. \\ & + \frac{1}{E_{\alpha}^2((E_{\alpha} - ip_4)^2 - E_{\beta}^2)} n'_B(E_{\alpha}) + \frac{1}{E_{\alpha}^2((E_{\alpha} + ip_4)^2 - E_{\beta}^2)} n'_B(E_{\alpha}) \\ & \left. + \frac{2}{E_{\beta}((E_{\beta} - ip_4)^2 - E_{\alpha}^2)^2} n_B(E_{\beta}) + \frac{2}{E_{\beta}((E_{\beta} + ip_4)^2 - E_{\alpha}^2)^2} (1 + n_B(E_{\beta})) \right\} \end{aligned}$$

$$\begin{aligned}
J_{+\sigma}(p_4, \mu_I) = & \frac{k^4}{3\pi^2} \frac{1}{4} \left\{ \frac{E_\sigma^2 - (E_\pi - ip_4 - 2\mu_I)(3E_\pi - ip_4 - 2\mu_I)}{E_\pi^3((E_\pi - ip_4 - 2\mu_I)^2 - E_\sigma^2)^2} (1 + n_B(E_\pi - 2\mu_I)) + \frac{n'_B(E_\pi - 2\mu_I)}{E_\pi^2((E_\pi - ip_4 - 2\mu_I)^2 - E_\sigma^2)} \right. \\
& + \frac{E_\sigma^2 - (E_\pi + ip_4 + 2\mu_I)(3E_\pi + ip_4 + 2\mu_I)}{E_\pi^3((E_\pi + ip_4 + 2\mu_I)^2 - E_\sigma^2)^2} n_B(E_\pi + 2\mu_I) + \frac{n'_B(E_\pi + 2\mu_I)}{E_\pi^2((E_\pi + ip_4 + 2\mu_I)^2 - E_\sigma^2)} \\
& \left. + \frac{2n_B(E_\sigma)}{E_\sigma((E_\sigma - ip_4 - 2\mu_I)^2 - E_\pi^2)^2} + \frac{2(1 + n_B(E_\sigma))}{E_\sigma((E_\sigma + ip_4 + 2\mu_I)^2 - E_\pi^2)^2} \right\}
\end{aligned}$$

$$\begin{aligned}
J_{+\sigma}(p_4, \mu_I) = & \frac{k^4}{3\pi^2} \frac{1}{4} \left\{ \frac{E_\sigma^2 - (E_\pi - ip_4 - 2\mu_I)(3E_\pi - ip_4 - 2\mu_I)}{E_\pi^3((E_\pi - ip_4 - 2\mu_I)^2 - E_\sigma^2)^2} (1 + n_B(E_\pi - 2\mu_I)) + \frac{n'_B(E_\pi - 2\mu_I)}{E_\pi^2((E_\pi - ip_4 - 2\mu_I)^2 - E_\sigma^2)} \right. \\
& + \frac{E_\sigma^2 - (E_\pi + ip_4 + 2\mu_I)(3E_\pi + ip_4 + 2\mu_I)}{E_\pi^3((E_\pi + ip_4 + 2\mu_I)^2 - E_\sigma^2)^2} n_B(E_\pi + 2\mu_I) + \frac{n'_B(E_\pi + 2\mu_I)}{E_\pi^2((E_\pi + ip_4 + 2\mu_I)^2 - E_\sigma^2)} \\
& \left. + \frac{2n_B(E_\sigma)}{E_\sigma((E_\sigma - ip_4 - 2\mu_I)^2 - E_\pi^2)^2} + \frac{2(1 + n_B(E_\sigma))}{E_\sigma((E_\sigma + ip_4 + 2\mu_I)^2 - E_\pi^2)^2} \right\}
\end{aligned}$$

$$\begin{aligned}
F_{++}^0(p_4, \mu_I) &= \sum_q \left[-m^2 - \vec{q}_r^2 - (q_4 + i\mu_I)(q_4 + 2p_4 + i\mu_I) \right] D_+^2(q) D_+(q+p) 4k\Theta(k^2 - \vec{q}^2) \\
&= \frac{4k^4}{6\pi^2} \left\{ \frac{4E_q^2 - p_4^2}{E_q(4E_q^2 + p_4^2)^2} (n_F(E_q + \mu_I) + n_F(E_q - \mu_I) - 1) \right. \\
&\quad \left. - \frac{1}{2E_q(2E_q - ip_4)} n'_F(E_q - \mu_I) - \frac{1}{2E_q(2E_q + ip_4)} n'_F(E_q + \mu_I) \right\}
\end{aligned}$$

$$\begin{aligned}
F_{++}^0(p_4, \mu_I) &= \sum_q D_+^2(q) D_+(q+p) 4k\Theta(k^2 - \vec{q}^2) \\
&= \frac{4k^4}{6\pi^2} \left\{ \frac{12E_q^2 + p_4^2}{4E_q^3(4E_q^2 + p_4^2)^2} (1 - n_F(E_q + \mu_I) - n_F(E_q - \mu_I)) \right. \\
&\quad \left. + \frac{1}{4E_q^2(2E_q - ip_4)ip_4} n'_F(E_q - \mu_I) - \frac{1}{4E_q^2(2E_q + ip_4)ip_4} n'_F(E_q + \mu_I) \right\}
\end{aligned}$$

$$\begin{aligned}
F_{+-}^\pi(p_4, \mu_I) &= \sum_q \left[-m^2 - \vec{q}_r^2 - (q_4 + i\mu_I)(q_4 + 2p_4 - 3i\mu_I) \right] D_+^2(q) D_-(q+p) 4k\Theta(k^2 - \vec{q}^2) \\
&= \frac{4k^4}{6\pi^2} \left\{ \frac{1}{2E_q(2E_q + ip_4 + 2\mu_I)^2} (2n_F(E_q + \mu_I) - 1) + \frac{1}{2E_q(2E_q - ip_4 - 2\mu_I)^2} (2n_F(E_q - \mu_I) - 1) \right. \\
&\quad \left. - \frac{1}{2E_q(2E_q + ip_4 + 2\mu_I)} n'_F(E_q + \mu_I) - \frac{1}{2E_q(2E_q - ip_4 - 2\mu_I)} n'_F(E_q - \mu_I) \right\}
\end{aligned}$$