Critical Casimir forces from the equation of state of quantum critical systems

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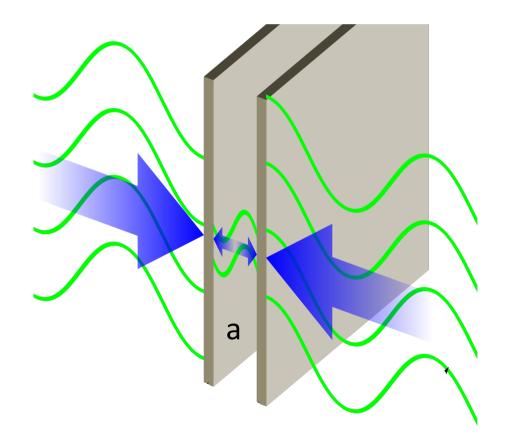


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### Introduction

Casimir 1948 : Casimir effect



## **Critical Casimir force**

Fisher and de Gennes 1978 :

fluctuating (classical) medium between two plates creates a force (if correlation length of the order of distance between the plates)

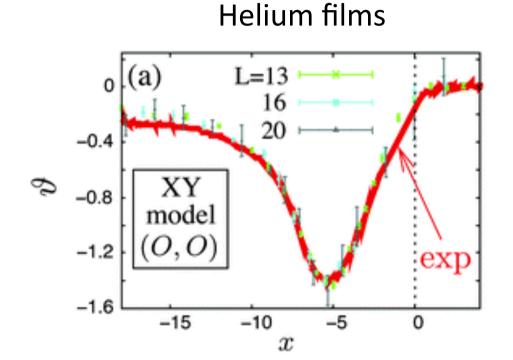
Large correlation length : second order phase transition  $\xi \propto |T - T_c|^{-\nu}$ Also implies universality of the force close to criticality :

Pressure of fluid on the plates : 
$$P = -\frac{1}{A} \frac{\partial F}{\partial a}$$

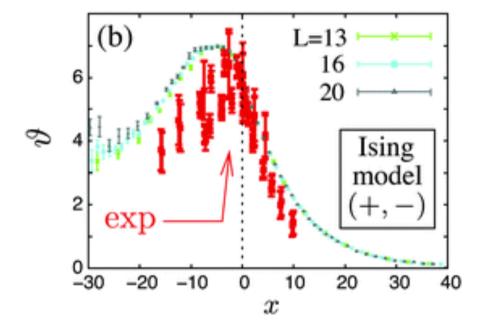
$$F = aA(f_{\text{bulk}} + f_{\text{ex}}) \qquad f_{\text{ex}} = a^{-3}\tilde{f}_s(a/\xi)$$

$$P_c = -\partial_a(a\tilde{f}_s) = a^{-3}\vartheta(a/\xi)$$

### **Experimental realizations**



**Binary mixtures** 



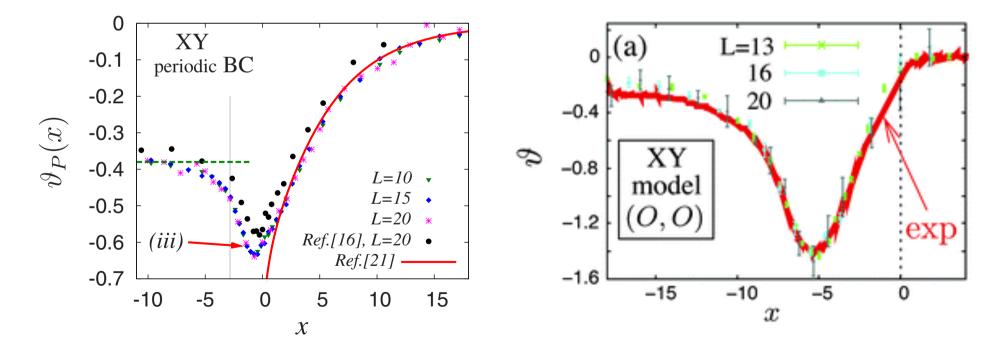
MC simulations : Vasilyev et al. 2007 He experiments : Garcia et al. 1994 Ganshin et al. 2006 Binary mixture : Fukuto et al. 2005

#### Strict boundary conditions

### **Boundary conditions**

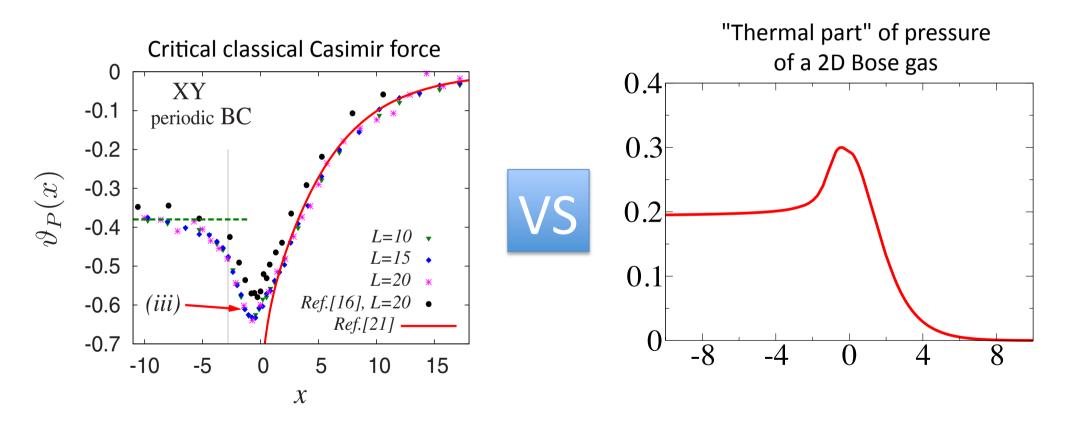
Casimir scaling function universal, depends on :

- dimension
- symmetry of order parameter
- geometry / boundary conditions



Vasilyev et al. 2009

#### From critical Casimir to quantum critical systems...

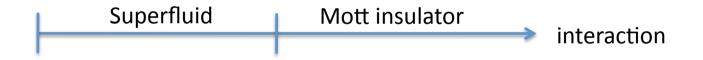


Functional renormalization group study : AR, Kodio, Dupuis, Lecheminant (2013)

### Quantum phase transitions

QPT : transition at zero temperature (change of ground-state) when changing non-thermal parameter  $\delta$ 

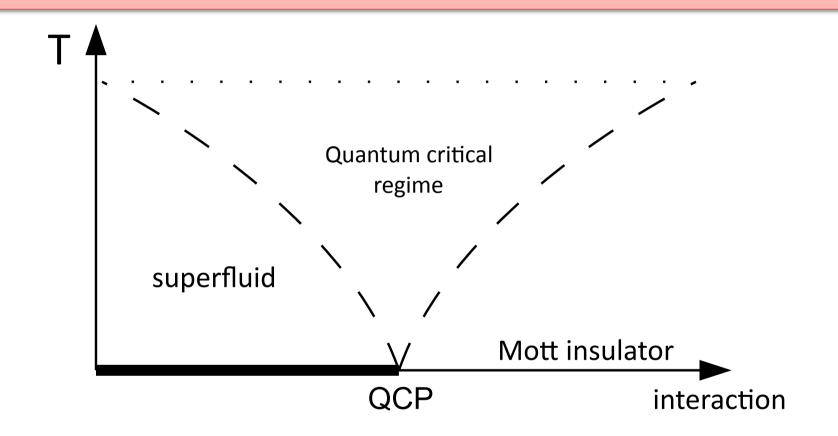
Examples : Bosonic Mott transition at constant density (XY universality class)



Ferro-paramagnetic transition in quantum Ising model in transverse field



## Phase diagram and critical scaling



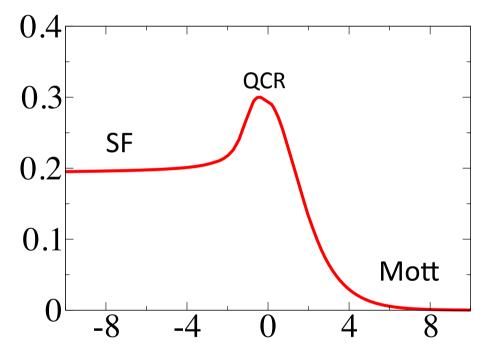
Close to QCP: 
$$f(\delta,T) = \epsilon_0(\delta) + T^3/c^2 \tilde{f}_s(\beta c/\xi)$$

Ground-state energy density

 $\xi \propto |\delta - \delta_c|^{-\nu}$ 

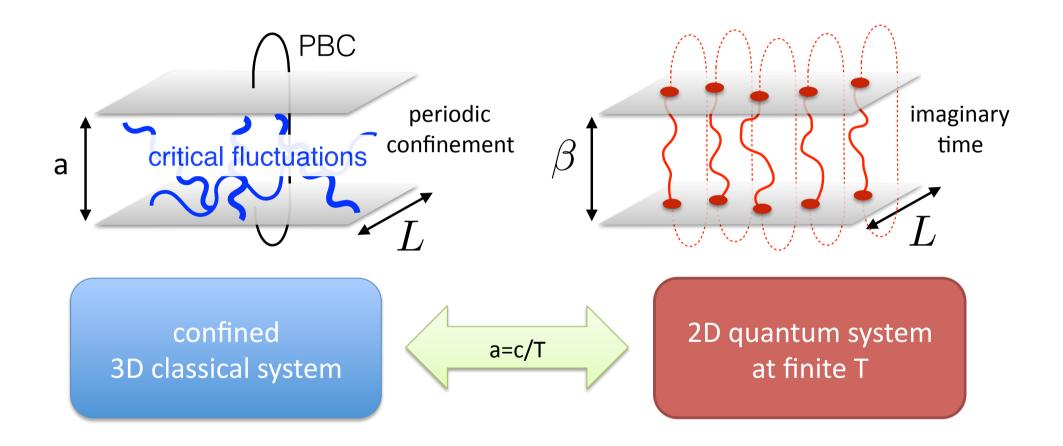
#### Phase diagram and critical scaling

Thermal part of pressure/free energy



Close to QCP:  $f(\delta,T) = \epsilon_0(\delta) + T^3/c^2 \tilde{f}_s(\beta c/\xi)$ 

### Quantum classical correspondence



Close to a critical point, universality implies that scaling functions are the same !

From critical Casimir to quantum critical systems... and back

$$f(\delta, T) = \epsilon_0(\delta) + T^3/c^2 \tilde{f}_s(\beta c/\xi)$$

Universality : same as confined 3D classical system

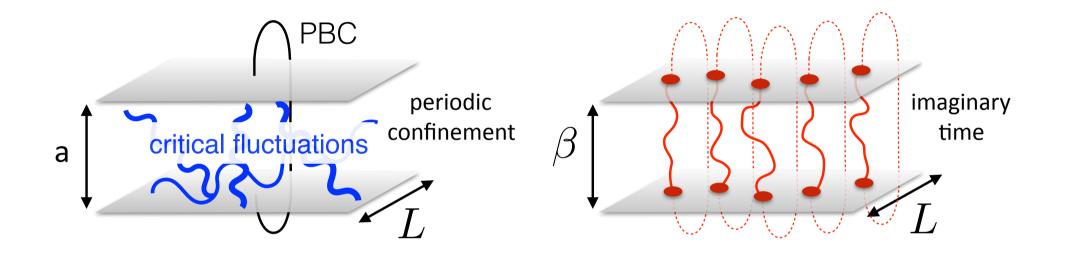
$$\epsilon = L^{-2} \langle \hat{H} 
angle = \partial_eta (eta f)$$
  
Average energy of a quantum critical system :  $\epsilon = \epsilon_0 - rac{T^3}{c^2} artheta_P (eta c / \xi)$ 

Thermodynamic stability of quantum systems implies attractive Casimir force for periodic BC.

$$P_c = -\partial_a(a\tilde{f}_s) = a^{-3}\vartheta(a/\xi)$$

For free bosons : Coleman et al. Am. J. Phys. 2009

### **Critical Casimir vs Equation of State**



Quantum to classical correspondence :

average energy interpreted as (universal) entropic "Casimir" force

Classical to quantum correspondence :

critical Casimir force with PBC can be quantum simulated experimentally!

### **FRG** calculation

Theoretical approaches are scarce to compute scaling functions (large N or epsilon expansion fail here).

AR, Kodio, Lecheminant and Dupuis (2014) : LPA'+ expansion up to  $\varphi^4$  Not good for Ising.

Improved calculation : 2<sup>nd</sup> order of Derivative Expansion

$$\Gamma_{k}[\phi] = \int_{0}^{\hbar\beta} d\tau \int d^{d}r \left\{ \frac{Z_{k}^{x}(\rho)}{2} (\nabla\phi)^{2} + \frac{Z_{k}^{\tau}(\rho)}{2} (\partial_{\tau}\phi)^{2} \right. \\ \left. + \frac{Y_{k}^{x}(\rho)}{4} (\nabla\rho)^{2} + \frac{Y_{k}^{\tau}(\rho)}{4} (\partial_{\tau}\rho)^{2} + U_{k}(\rho) \right\},$$
$$\left. \int f = \lim_{k \to 0} U_{k}(\rho_{0,k}) \right\}$$

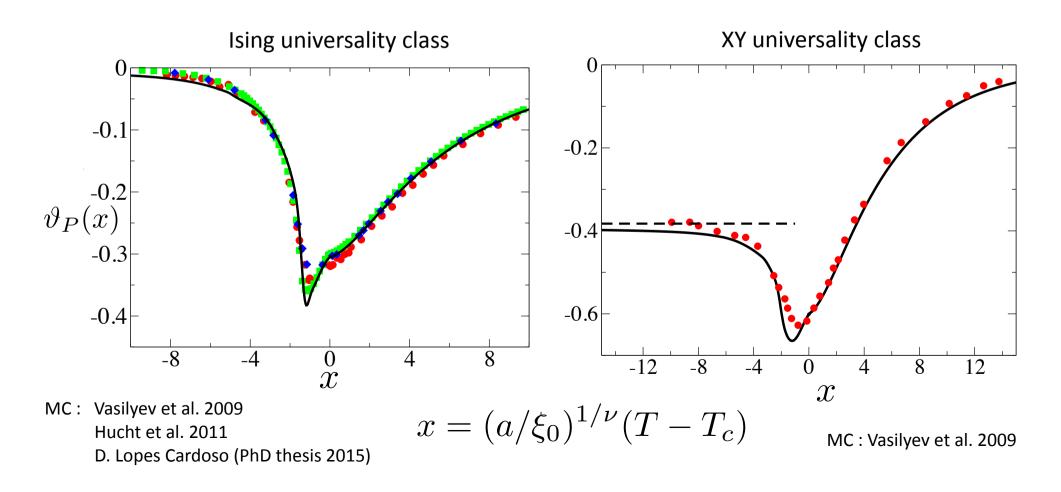
 $\rho = \phi^2/2$ 

See also : Jakubczyk and Napiorkowski 2013

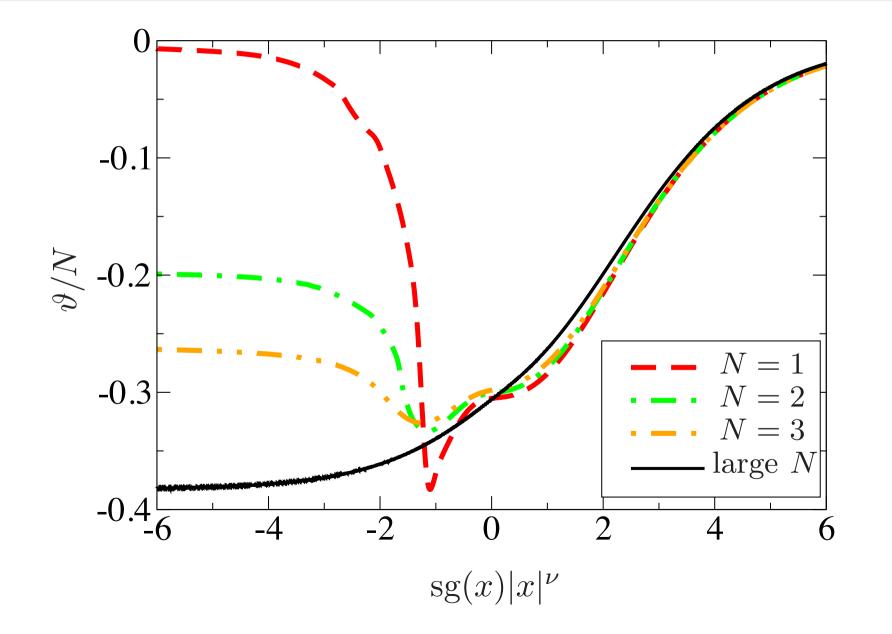
#### Casimir critical force from the FRG

TABLE II. Universal Casimir amplitude  $\vartheta(0)/2$ 

N	1	2	3
NPRG	-0.1527	-0.3006	-0.4472
Monte Carlo [5]	-0.1520(2)	-0.2993(7)	



#### Scaling functions for different N



## **Conclusion and perspectives**

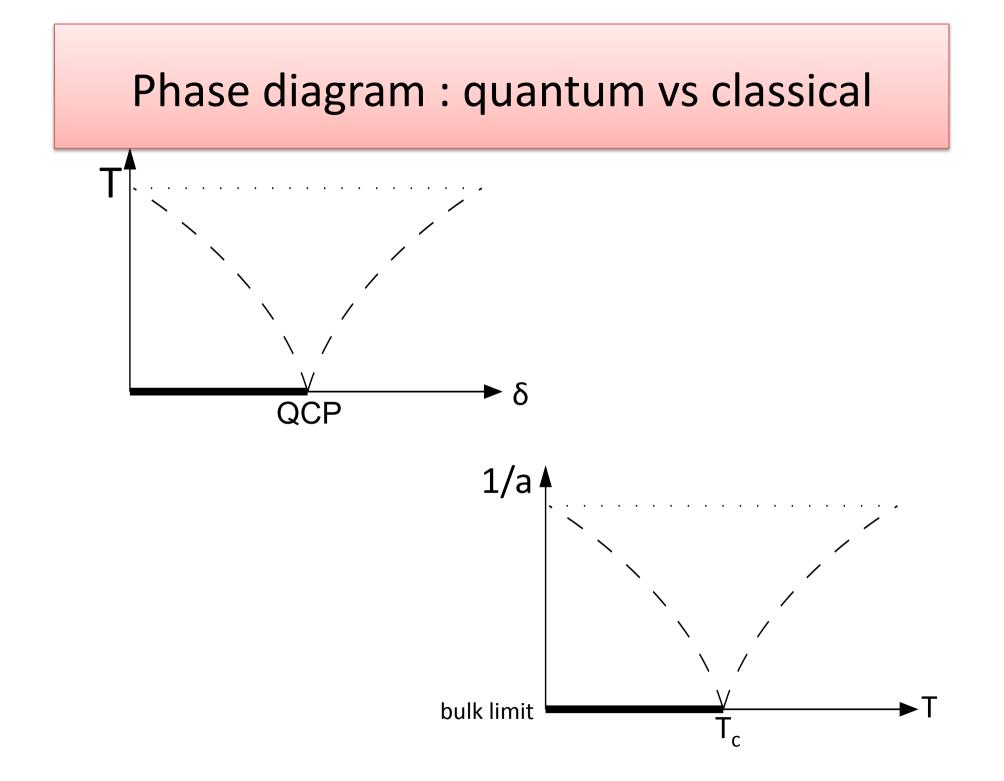
- Critical Casimir forces with periodic BC is the equation of state of a quantum critical system.

- Corresponding scaling functions could be measured in state of the art experiments on quantum systems.

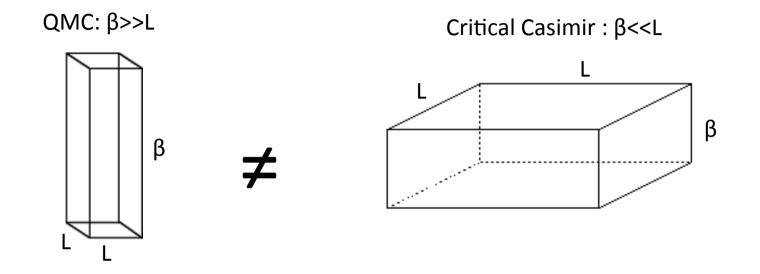
- Tools of quantum many-body problem can be used to study critical Casimir forces (Quantum Monte Carlo).

Open question : how to tackle other boundary conditions ?
 Ex : Free BC, order parameter depends on position, flow equation much harder to solve.

arXiv:1606.03205



# Finite size scaling for quantum systems



We can thus expect that the universal coefficient of FSS depends on the ratio  $\rho=\beta c/L$ . At the critical point  $\delta=0$ :  $u = E_0 - \frac{T^3}{c^2} \Phi(\rho)$ 

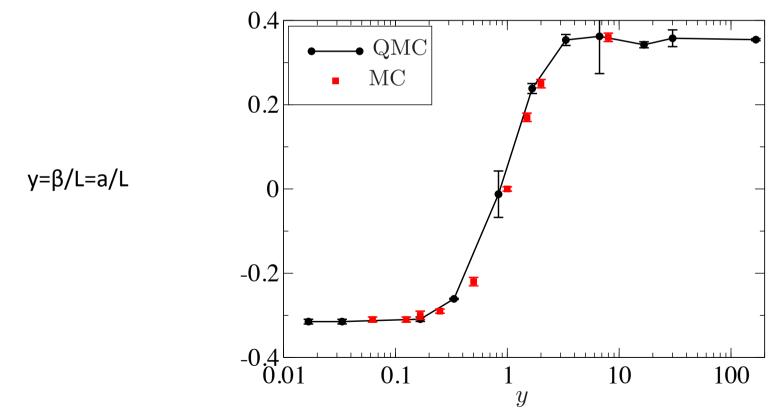
 $\Phi(0)$  Casimir amplitude (=-0.32 for Ising)

 $\lim_{\rho\to\infty} \tilde{\Phi}(\rho) = \alpha \rho^3 \quad \text{with $\alpha$ a universal (non-standard Casimir) amplitude} \\ \quad \ \ (= 0.37 \text{ for Ising})$ 

### Aspect ratio and Finite Size Scaling

Dependence on aspect ratio known in context of Casimir forces.

Continuous imaginary time QMC for quantum Ising in transverse field



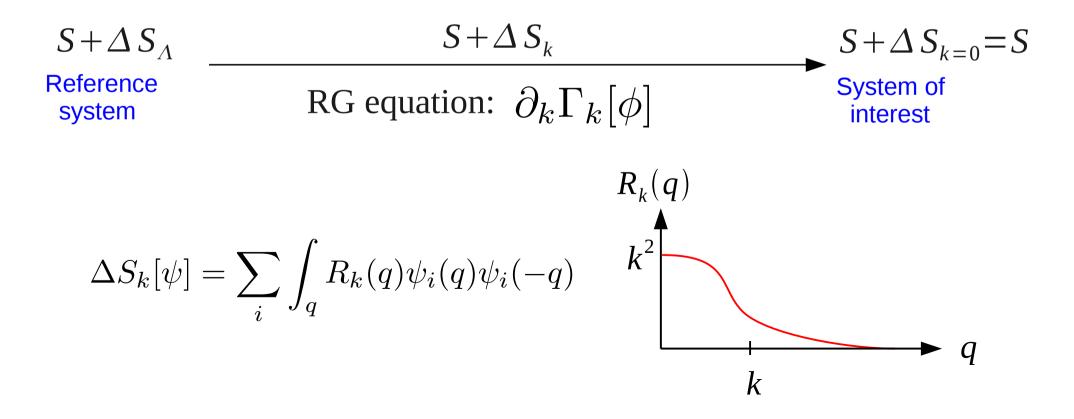
MC: Hucht et al. 2011

#### **Non-Pertubative Renormalization Group**

Subtle calculation in 4- $\epsilon$  and Large N for periodic BC.

Here : Non-Perturbative Renormalization Group (Wetterich 1993)

Family of actions indexed by momentum scale k



## Effective Action ("Gibbs" free energy)

Quantum O(N) model

$$S = \int_0^{\hbar\beta} d\tau \int d^2r \left\{ \frac{(\nabla \varphi)^2}{2} + \frac{(\partial_\tau \varphi)^2}{2c^2} + \frac{r\varphi^2}{2} + \frac{u(\varphi^2)^2}{4!} \right\}$$

Effective action : Legendre transform of Free energy with respect to magnetic field. Depends on the order parameter :  $m{\phi}=\langlem{arphi}
angle$ 

Ansatz : Derivative expansion (low energy fluctuations most important close to QCP)

$$\Gamma_{k}[\phi] = \int_{0}^{\hbar\beta} d\tau \int d^{d}r \left\{ \frac{Z_{k}^{x}(\rho)}{2} (\nabla\phi)^{2} + \frac{Z_{k}^{\tau}(\rho)}{2} (\partial_{\tau}\phi)^{2} + \frac{Y_{k}^{x}(\rho)}{4} (\nabla\rho)^{2} + \frac{Y_{k}^{\tau}(\rho)}{4} (\partial_{\tau}\rho)^{2} + U_{k}(\rho) \right\}$$

 $\rho = \phi^2/2$