

Nicolai Christiansen (ITP Heidelberg)

Higher Derivative Quantum Gravity and Vertex Functions

Trieste, ERG 2016

September 22, 2016

Outline

- (Introduction: Quantum Gravity and Asymptotic Safety)
- Quantum Gravity and Vertex Expansions
- Higher Derivative Interactions
- Fixed Points
- Outlook

Perturbative Quantization

- Einstein-Hilbert action: quadratic in derivatives $[G_N] = \text{energy}^{-2}$

$$S_{EH} = \frac{1}{16\pi} \int d^4x \sqrt{g} \left(-\frac{1}{G} R + 2\frac{\Lambda}{G} \right)$$

- expansion parameter for n-point Greens functions:

dimensionless $\longrightarrow g \equiv G_N E^2 = \frac{E^2}{M_{\text{Pl}}^2} \quad , \quad M_{\text{Pl}} \approx 10^{19} \text{GeV}$


- \longrightarrow higher loop orders require higher derivative counterterms !
- \longrightarrow full theory is either divergent or includes infinitely many free parameters

(perturbatively) non-renormalizable

Asymptotic Safety in a Nutshell

- Non-perturbative renormalization in Quantum Gravity

- (i) d.o.f. carried by the metric field
- (ii) diffeomorphism invariance
- (iii) quantum field theory of point particles

- Quantum Fluctuations  scale dependent couplings

$$g_i \longrightarrow g_i(k)$$

k = energy scale

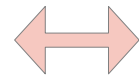
g dimensionless
couplings

UV fixed point:

$$\lim_{k \rightarrow \infty} g_i(k) = g_{i,*} < \infty$$

+ finite number of free parameters (predictive)

Asymptotic Safety



UV completion

S. Weinberg (1979)

- example: Asymptotic Freedom : $g_* = 0$ (perturbative)

A Challenge for Asymptotic Safety

- Technical tool: Functional Renormalization (Wetterich equation)

$$k \frac{d}{dk} \Gamma_k[\phi] = \text{STr} \left(\frac{\delta^2}{\delta \phi^2} \Gamma_k + R_k \right)^{-1} [\phi] k \frac{d}{dk} R_k$$

Wetterich (1993)

- Quantum Gravity: UV physics unknown!
 - systematic expansions and truncation enhancements!
Relevant subsets of theory space?
(Falls, Litim, Nikolakopoulos, Rahmede 2013)
 - Convergence?
 - ...find a deeper, underlying guiding principle...

Approaches to Asymptotic Safety with the FRG

- The background flow equation:

1
$$\dot{\Gamma}_k[\bar{g}, g] \Big|_{\bar{g}=g} = \text{STr} \frac{\dot{R}_k}{\Gamma_k^{(2)} + R_k} \Big|_{\bar{g}=g}$$

- Drawbacks:
 - equation is not closed as it is evaluated at

$$\bar{g} = g \quad \text{but} \quad \Gamma_k^{(2)} = \frac{\delta^2 \Gamma_k}{\delta g^2}$$

- no direct access to vertex functions

- Different approach: vertex expansion \longrightarrow write $g = \bar{g} + h$

2
$$\longrightarrow \Gamma[\bar{g}, h] = \sum \frac{1}{n!} \frac{\delta^n \Gamma}{\delta h^n} \Big|_{h=0} h^n \quad (\text{schematically})$$

\longrightarrow Flow equations for fluctuation field correlators

- Functional derivatives of Wetterich equation

 Scale dependence of full vertex functions $\Gamma^{(n)} = \frac{\delta^n \Gamma[\bar{g}, h]}{\delta h^n}$

$$k \frac{d}{dk} \Gamma^{(2)} = -\frac{1}{2} \text{ (diagram: bubble with cross on top, double lines on bottom)} + \text{ (diagram: bubble with cross on top, double lines on left and right)} - 2 \text{ (diagram: bubble with cross on top, dotted lines on left and right)}$$

$$k \frac{d}{dk} \Gamma^{(3)} = -\frac{1}{2} \text{ (diagram: bubble with cross on top, triple lines on bottom)} + 3 \text{ (diagram: bubble with cross on top, double lines on left, triple lines on right)} - 3 \text{ (diagram: bubble with cross on top, double lines on left, triple lines on right, cross on bottom)} + 6 \text{ (diagram: bubble with cross on top, dotted lines on left and right, triple lines on bottom)}$$

$$k \frac{d}{dk} \Gamma^{(n)} = \text{Flow}[\Gamma^{(2)}, \dots, \Gamma^{(n+2)}] \longrightarrow \text{infinite hierarchy of flow equations}$$

- Direct access to fluctuation correlation functions/couplings
- Access to momentum dependence

Vertex Expansion II

- Essential ingredient: parameterization/truncation of vertices
- Vertex construction:

$$\Gamma^{(n)} = \prod_{i=1}^n \sqrt{Z(p_i)} G_n^{\frac{n}{2}-1} \mathcal{T}^n(p_1, \dots, p_n; \lambda_j)$$

- Construction from „defining action“ $S[\bar{g}, h]$

→ Linear split $g = \bar{g} + h = \mathbb{I} + h$

→ Expand in powers of h : i.e. Tensor structures from $\frac{\delta^n S}{\delta h^n}$

→ Rescaling: $h \longrightarrow \sqrt{GZ}h$

→ Dressing: $g_i \longrightarrow g_{i,k}$ more general: $g_{i,k}(p)$

Systematics in the Vertex Expansion

- Vertex expansions in quantum gravity:

$$\Gamma^{(2)}$$

$$\Gamma^{(3)}$$

$$\Gamma^{(4)}$$

NC,Litim,Pawlowski,Rodigast 2012

NC,Knorr,Pawlowski,Rodigast 2014

NC,Knorr,Meibohm,Pawlowski,Reichert 2015

Denz,Pawlowski,Reichert in prep

→ expansion around curved background with $g(R)$ couplings:

NC,Falls,Pawlowski,Reichert in prep

→ expansion is consistent so far

→ support for Asymptotic Safety: Fixed Point

→ based on Einstein-Hilbert: $S = S_{\text{EH}} + \text{gauge}$

→ Higher order operators ???

Higher Derivative Gravity I

- General local action with curvature invariants up to 4-th order in derivatives: Stelle (1977)

$$S_{HD} = \frac{1}{16\pi} \int d^4x \sqrt{g} \left(\underbrace{a R_{\mu\nu} R^{\mu\nu} + b R^2}_{\text{fourth order}} - \underbrace{\frac{1}{G} R + 2 \frac{\Lambda}{G}}_{\text{Einstein-Hilbert (2nd order)}} \right)$$

- No need for (Riemann-tensor)² term due to topological Gauss-Bonnet term in d=4
- (Ric)² can be traded for (Weyl)²
- dimensionless couplings a,b
- Power-counting perturbatively renormalizable but:

non-unitary

Higher Derivative Gravity II

- Higher Derivative Operators and Asymptotic Safety



Background Field Flows for Γ_k

$R + R^2$: (Lauscher, Reuter 2002)

$R + R^2 + C^2$: (Codello, Peracci 2006),
(Benedetti, Machado, Saueressig 2009/10),

$f(R)$: (Codello, Peracci, Rahmede 2008), (Machado, Saueressig 2008), (Falls, Litim, Nikolakopoulos, Rahmede 2013/2014), (Benedetti 2013),
(Dietz, Morris 2013), (Eichhorn 2015)

$R + C^3$: (Gies, Knorr, Lippoldt, Saueressig 2016)



Flows for Vertex Functions/
Fluctuation correlators/
dynamical couplings



NC, in prep 2016

Higher Derivative Gravity III

- Start with general action quartic in derivatives

$$S_{HD} = \frac{1}{16\pi} \int d^4x \sqrt{-\det g} \left(a R_{\mu\nu} R^{\mu\nu} + b R^2 - \frac{1}{G} R + 2 \frac{\Lambda}{G} \right)$$

- Tensor structures of $\Gamma^{(n)}$ from $\frac{\delta^n S_{HD}}{\delta h^n}$
- Rescaling: $h \longrightarrow \sqrt{GZ} h$
- Dressing: $(Z \longrightarrow Z_k, G \longrightarrow G_k, a \longrightarrow a_k, b \longrightarrow b_k, -2\Lambda \longrightarrow M_k^2)$

→ Flow equation for the inverse propagator:

$$k \partial_k \Gamma^{(2)}(Z_k, a_k, b_k, M_k^2) = \text{Flow}^{(2)}(G_k, Z_k, a_k, b_k, M_k^2)$$

Vertex-coupling: free parameter or from EH-flows

Higher Derivative Gravity IV

- Gauge fixing: 2nd order in derivatives!

$$S_{\text{GF}} = \frac{1}{2\alpha} \int d^4x \sqrt{g} \bar{g}^{\mu\nu} F_\mu F_\nu \quad , \quad \text{in general:} \quad F_\mu = \bar{\nabla}^\nu h_{\mu\nu} - \frac{1+\beta}{4} \bar{\nabla}_\mu h^\nu_\nu$$

This work: $\alpha = 0 \quad , \quad \beta = -1$  See also: (C. Wetterich 2016)

- Propagator: Projector representation $\left(\Gamma^{(2)}\right)^{-1} = \sum_i \left(\Gamma^{(2)}\right)^{-1}_i \mathcal{P}_i$

TT component

$$\left(\Gamma^{(2)}\right)^{-1} = \frac{16\pi}{Z_k} \begin{pmatrix} (p^2 + a_k G_k p^4 + M_k^2)^{-1} & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & -2(4p^2 - 8p^4(3a_k G_k + b_k G_k) + M_k^2)^{-1} & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

Tr component

Higher Derivative Gravity V

- Projection procedure:

$$\lim_{p \rightarrow 0} \mathcal{P}_{TT} \dot{\Gamma}^{(2)} \longleftrightarrow k \partial_k M_k^2$$

$$\lim_{p \rightarrow 0} \frac{\partial^2}{\partial p^2} \mathcal{P}_{TT} \dot{\Gamma}^{(2)} \longleftrightarrow k \partial_k Z_k$$

$$\lim_{p \rightarrow 0} \frac{\partial^4}{\partial p^4} \mathcal{P}_{TT} \dot{\Gamma}^{(2)} \longleftrightarrow k \partial_k a_k$$

$$\lim_{p \rightarrow 0} \frac{\partial^4}{\partial p^4} \mathcal{P}_{Tr} \dot{\Gamma}^{(2)} \longleftrightarrow k \partial_k b_k$$

—→ clean separation of the four-derivative couplings via tensor projection

Higher Derivative Gravity VI

- Fixed point solutions

1 Solve $(\beta_M, \eta, \beta_a, \beta_b)$ and assume fixed point in g :

→ parametric dependence on Newtons coupling, e.g. $g_* = 0.6$

$$(\mu_* = -0.17, \eta_* = 0.26, a_* = -0.94, b_* = 0.4)$$

Critical exponents: $(\theta_1 = -5.8, \theta_2 = -21.5, \theta_3 = 6.3)$

Irrelevant direction

2 Solve $(\beta_M, \eta, \beta_a, \beta_b, \beta_g)$, β_g from EH calculation

$$(g_* = 0.98, \mu_* = -0.2, \eta_* = 1.06, a_* = -0.56, b_* = 0.25)$$

$$(\theta_1 = -1.67, \theta_2 = -14.5, \theta_3 = -18.7, \theta_4 = 11.1)$$

Higher Derivative Gravity VII

- Discussion:

- system seems to be stable („simulation“ of changes in truncation)

- reasonable fixed point values,
shares features with EH vertex expansion: e.g. $\eta > 0$

- Three relevant, one irrelevant direction

Asymptotic Safety:



...in agreement (Benedetti, Machado, Saueressig 2009)

- critical exponents a bit too large....

- stabilization by higher order operators?

(Falls, Litim, Nikolakopoulos, Rahmede 2013/2014)

Summary and Outlook

- Vertex expansion and flow of the propagator with all four-derivative operators
- Fixed point with three relevant and one irrelevant direction



Asymptotic Safety

Outlook

- Including more momentum dependence
- Including R^3 operators
- Three-point/four-point function with higher-derivative operators
- Unitarity

Thank You!!!