Universal short-time dynamics: FRG for a temperature quench



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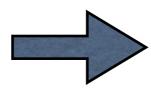
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Trieste, 19th September 2016





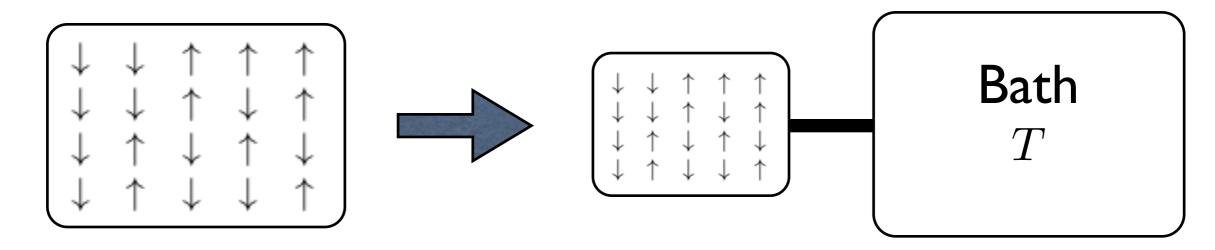




Quenching

From Wikipedia, the free encyclopedia

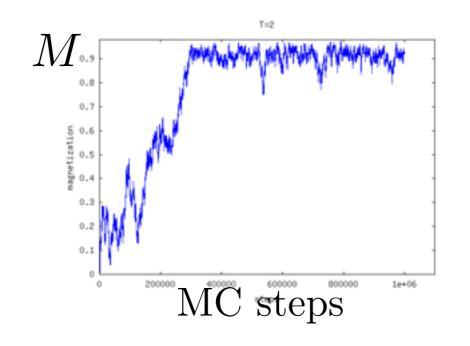
In materials science, **quenching**, a type of heat treating, is the rapid cooling of a workpiece to obtain certain material properties.

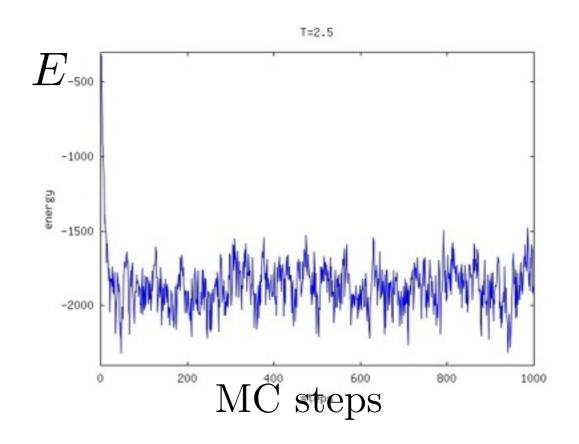


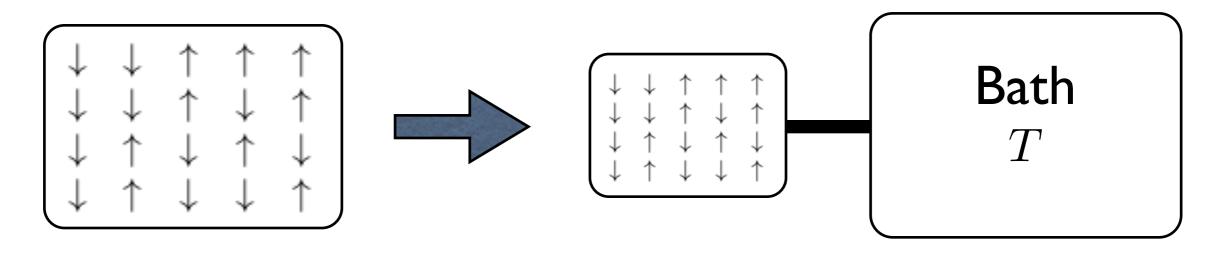
initial distribution $P_0(\{\sigma\})$

equilibrium driven by reservoir $P_{\rm eq}(\{\sigma\},T)$

Example: Metropolis algorithm

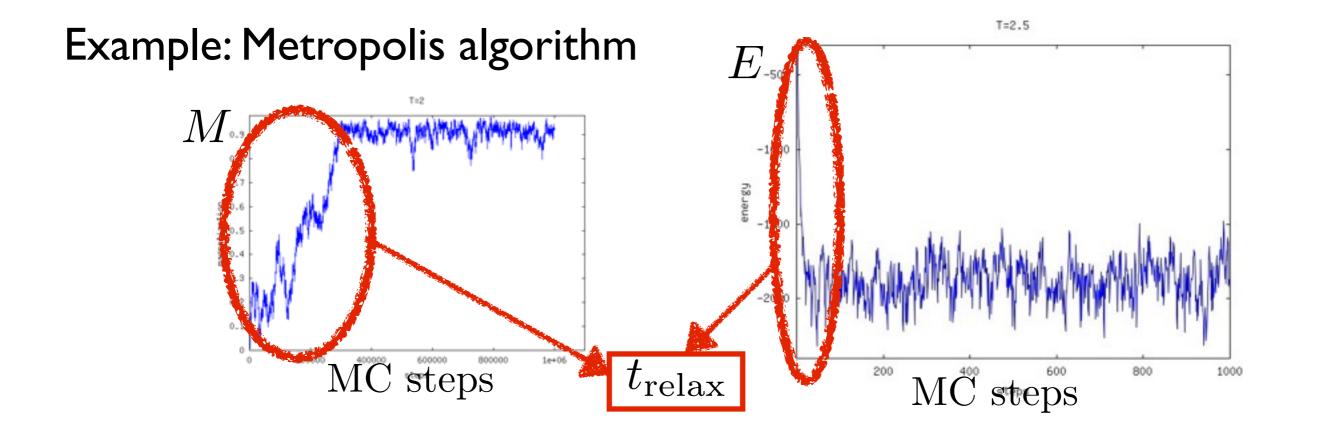


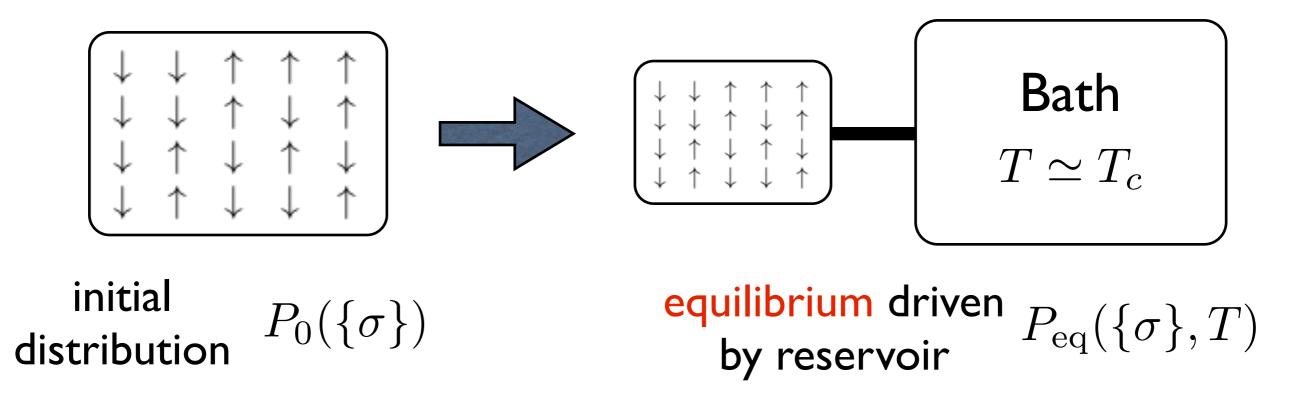




 $\begin{array}{c} \text{initial} \\ \text{distribution} \end{array} P_0(\{\sigma\})$

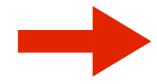
equilibrium driven by reservoir $P_{\rm eq}(\{\sigma\},T)$





 $T \simeq T_c$ equilibrium: correlation length and time diverges quench: $t_{
m relax}$ too! Equilibrium is attained in infinite time

Is there some universal behaviour?



AGING

Janssen, Schaub, Schmittman '89 rev: Calabrese & Gambassi '05

Aging dynamics

New critical exponents? at criticality: $(\xi=\infty)$ Janssen, Schaub, Schmittman '89 rev: Calabrese & Gambassi '05

I) two-time correlation functions

(response function)
$$G_R(q,t,t') = q^{-2+\eta+z} \mathcal{G}_R(q^zt,q^zt') \qquad \begin{array}{l} \eta,z \\ \text{equilibrium} \\ \text{(correlation function)} \end{array}$$

Aging dynamics

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1) two-time correlation functions

(response function)
$$G_R(q,t,t') = q^{-2+\eta+z} \mathcal{G}_R(q^zt,q^zt') \qquad \begin{array}{l} \eta,z \\ \text{equilibrium} \\ \text{(correlation function)} \end{array}$$

if $t' \ll t$

$$G_C(q, t, t') = q^{-2+\eta} (t/t')^{\theta-1} \widetilde{\mathcal{G}}_C(q^z t, q^z t')$$

$$G_R(q, t, t') = q^{-2+\eta+z} (t/t')^{\theta} \widetilde{\mathcal{G}}_R(q^z t, q^z t')$$

new non-equilibrium exponent θ !

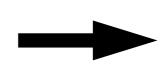
Aging dynamics

New critical exponents? at criticality: $(\xi = \infty)$ Janssen, Schaub, Schmittman '89 rev: Calabrese & Gambassi '05

2) magnetization

Modify quench protocol:

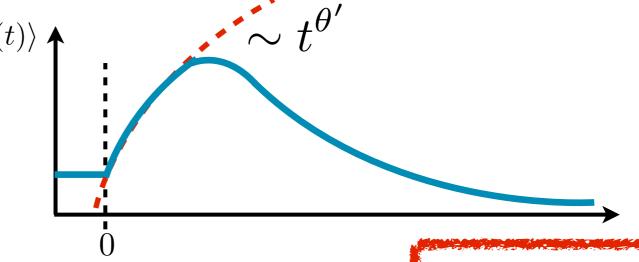
ext. magnetic field for t<0



no ext. magnetic field for t>0

$$M(t) = M_0 t^{\theta'} \mathcal{M}(M_0 t^{\theta' + \beta/z\nu})$$

$$\theta' = \theta + (2 - z - \eta)/z$$



Take home message: breaking of time-translational invariance (or fluctuation-dissipation theorem) generates new exponents Sieberer et al. '13, '14

Marino & Diehl '16

cf. equilibrium systems with spatial boundaries

Theoretical description



coarse graining: Hohenberg-Halperin models

Hohenberg & Halperin '77

e.g., Ising model:

Glauber dynamics of Ising model: Model A

Initial condition ϕ_0 is stochastic: $P[\phi_0] \propto \mathrm{e}^{-\int \tau_0 \phi_0^2/2}$

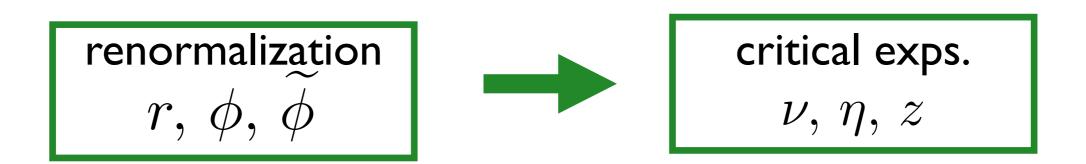
Initial correlation length: $\tau_0^{-1/2} \to 0$

$$\langle O(t) \rangle = \int O(\phi) \langle \delta[\phi - \phi(t)] \rangle_{\zeta,\phi_0} = \int O(\phi) e^{-S[\phi,\widetilde{\phi}]}$$

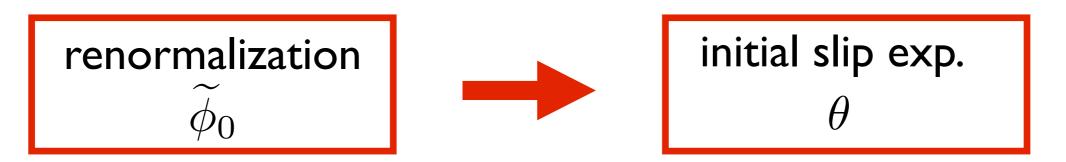
- ullet New field: response field $\widetilde{\phi}$
- ullet $S[\phi,\widetilde{\phi}]$ contain info also on initial conditions
- This action is viable for renormalization

Renormalization and critical exponents

Equilibrium: some divergencies, renormalization required



Quench: new divergences, additional renormalization required



- second order in ϵ expansion

- value of θ : large-N limit of O(N)
 - Monte Carlo

rev: Calabrese & Gambassi '05

Modified action:
$$S_k[\Psi] = S[\Psi] + \Delta S_k[\Psi]$$
 $\Psi = (\varphi, \widetilde{\varphi})$

 $\Delta S_k[\Psi]$ imposes a IR cutoff parametrized by k

$$S_k[\Psi]$$
 coarse-grained over $S_k[\Psi]$ effective action $\Gamma_k[\Phi]$ volume $\sim k^{-d}$)

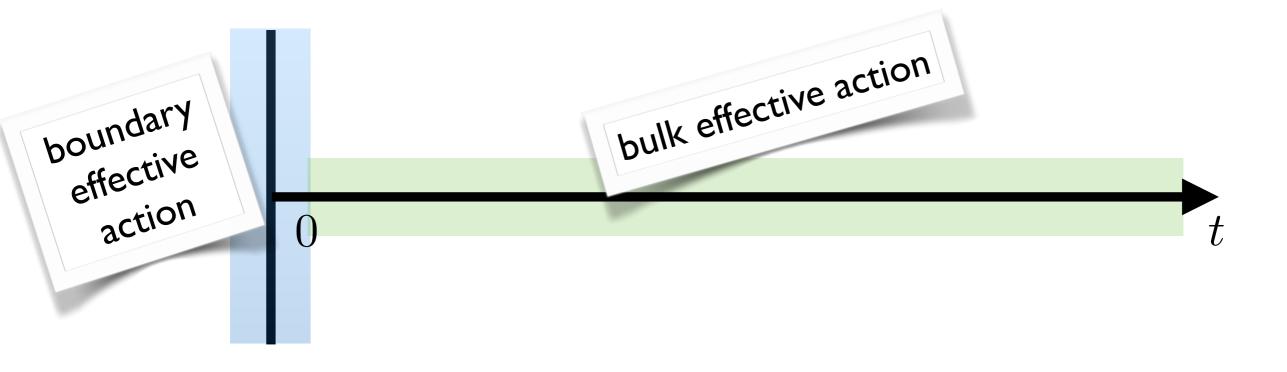
(coarse-grained over

$$\frac{\mathrm{d}\Gamma_k}{\mathrm{d}k} = \frac{1}{2} \int \mathrm{tr} \left[\left(\frac{\delta^2 \Gamma_k}{\delta \Phi^{\mathrm{t}} \delta \Phi} + R_k \right)^{-1} \frac{\mathrm{d}R_k}{\mathrm{d}k} \right] \quad \text{FRG equation}$$
 Wetterich '93

How to include the quench? An ansatz for $\Gamma_k[\Phi]$ has to be used.

$$\Gamma[\phi,\widetilde{\phi}] = \frac{\Gamma_0[\phi_0,\widetilde{\phi}_0]}{\Gamma_0[\phi_0,\widetilde{\phi}_0]} + \int_x \vartheta(t-t_0)\widetilde{\phi} \left(Z\dot{\phi} + K\nabla^2\phi + \frac{\partial \mathcal{U}}{\partial\phi} - D\widetilde{\phi}\right)$$
 initial quench conditions

Formally similar to a boundary problem:



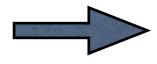
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 initial quench conditions quench

$$\Gamma_0=\int\left(-\frac{Z_0^2}{2\tau_0}\widetilde{\phi}_0^2+Z_0\widetilde{\phi}_0\phi_0\right) \qquad \begin{array}{c} \text{from } P[\phi_0] \text{ and } \\ \text{causality} \end{array}$$

determined causality arguments

Breaking of TTI



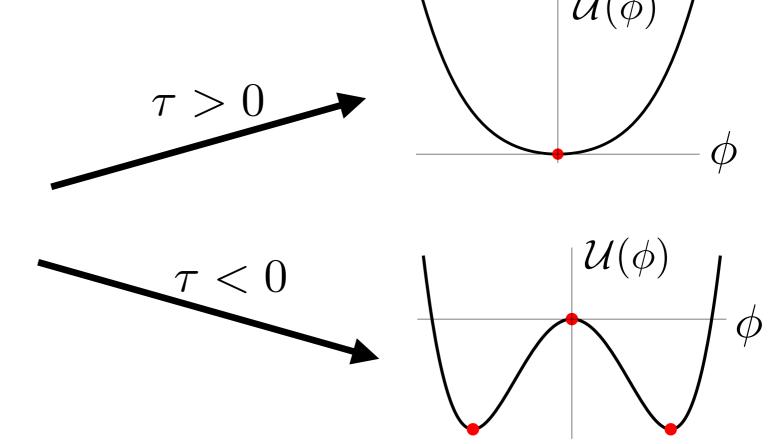
Fourier transform not useful

Our approach: 3 steps

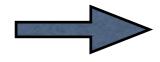
Step I: write an ansatz for the potential $\mathcal{U}(\phi)$

Simplest ansatz:

$$\mathcal{U}(\phi) = \frac{\tau}{2}\phi^2 + \frac{g}{4!}\phi^4$$



Breaking of TTI



Fourier transform not useful

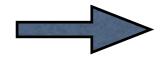
Our approach: 3 steps

Step 2: rewrite FRG eq. as follows:

$$\frac{\mathrm{d}\Gamma}{\mathrm{d}k} = \frac{1}{2} \int_x \mathrm{tr} \left[\vartheta(t-t_0) G(x,x) \frac{\mathrm{d}R}{\mathrm{d}k} \sigma\right]$$

$$G = \left(\Gamma^{(2)} + R\right)^{-1} = (G_0^{-1} - V)^{-1}$$
 field field independent dependent

Breaking of TTI



Fourier transform not useful

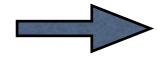
Our approach: 3 steps

$$\frac{\mathrm{d}\Gamma}{\mathrm{d}k} = \frac{1}{2} \int_{x} \mathrm{tr} \left[\vartheta(t - t_0) G(x, x) \frac{\mathrm{d}R}{\mathrm{d}k} \sigma \right]$$

$$G(x,x') = G_0(x,x') + \int_{y} G_0(x,y) V(y) G(y,x')$$
 Dyson-like equation

$$=G_0(x,x')+\sum_{n=1}^{+\infty}G_n(x,x')$$
 Each term can be evaluated!
$$\sim \phi^{2n}$$

Breaking of TTI



Fourier transform not useful

Our approach: 3 steps

Step 3: replace these terms into FRG eq.

constant value: renormalizes bulk

$$\frac{\mathrm{d}\Gamma_k}{\mathrm{d}k} \propto \int_{\mathbf{r}}^{\infty} \int_0^{\infty} \mathrm{d}t \, \widetilde{\phi} \phi \left[1 + f(t)\right]$$
 vanishing function

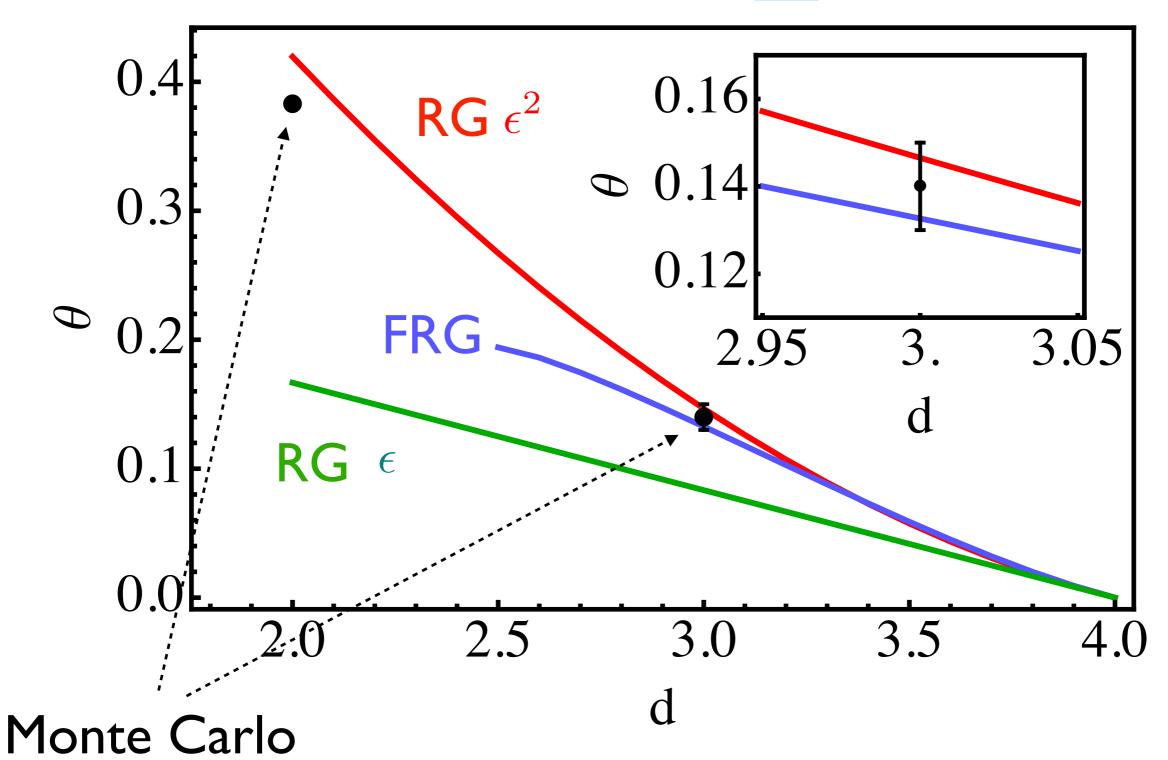
$$\int_0^\infty dt \, \widetilde{\phi}(t)\phi(t)f(t) = \sum_{n=0}^\infty c_n \frac{d^n}{dt^n} \left[\widetilde{\phi}(t)\phi(t) \right]_{t=0}^n$$

Bulk renormalizes the boundary!

$$\mathcal{U} = \frac{g}{4!} (\phi^2 - \phi_m^2)^2 + \frac{\lambda}{6!} (\phi^2 - \phi_m^2)^3$$

more refined ansatz:

$$\mathcal{U} = \frac{g}{4!} (\phi^2 - \phi_m^2)^2 + \frac{\lambda}{6!} (\phi^2 - \phi_m^2)^3$$



Conclusion & outlook

Summary

- Critical quench generates new critical exponents
- FRG can be used to compute non-equilibrium exponents
- Good agreement with numerics and perturbative RG

Near future...

- Full potential
- Application to quenches in isolated quantum systems
- O(N) and Potts models
- Other HH models (conserved dynamics)

... far future

- non-thermal fixed points?
- coarsening dynamics
- towards new exp. platforms (cold atoms, exciton-polaritons...)



Thank you for your attention!