

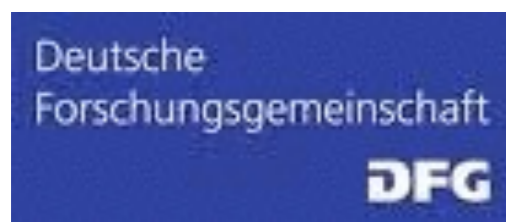
# Generation of structure in spin- and mass-imbalanced Fermi gases

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Jens Braun  
TU Darmstadt

8th International Conference on the Exact Renormalization Group

20/09/2016



# From balanced to imbalanced gases - why?

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- Condensed-matter point of view:

Superconductors form  
when spin-up and spin-down electrons pair up

PHYSICAL REVIEW

VOLUME 108, NUMBER 5

DECEMBER 1, 1957

## Theory of Superconductivity\*

J. BARDEEN, L. N. COOPER,† AND J. R. SCHRIEFFER‡  
*Department of Physics, University of Illinois, Urbana, Illinois*  
(Received July 8, 1957)



- What happens when a magnetic field is applied?

Superconductivity is destroyed  
for a large enough field

A NOTE ON THE MAXIMUM CRITICAL FIELD OF  
HIGH-FIELD SUPERCONDUCTORS

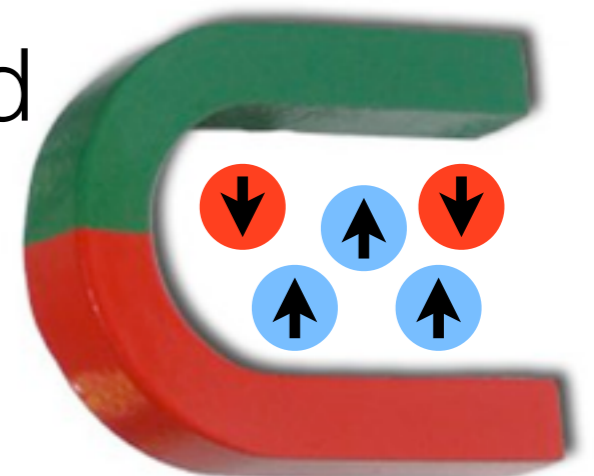
B. S. Chandrasekhar

Westinghouse Research Laboratories, Pittsburgh, Pennsylvania  
(Received July 23, 1962)

UPPER LIMIT FOR THE CRITICAL FIELD IN HARD SUPERCONDUCTORS

A. M. Clogston

Bell Telephone Laboratories, Murray Hill, New Jersey  
(Received August 9, 1962)



# From balanced to imbalanced gases - why?

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- *Related problem:*

*Superfluids* form

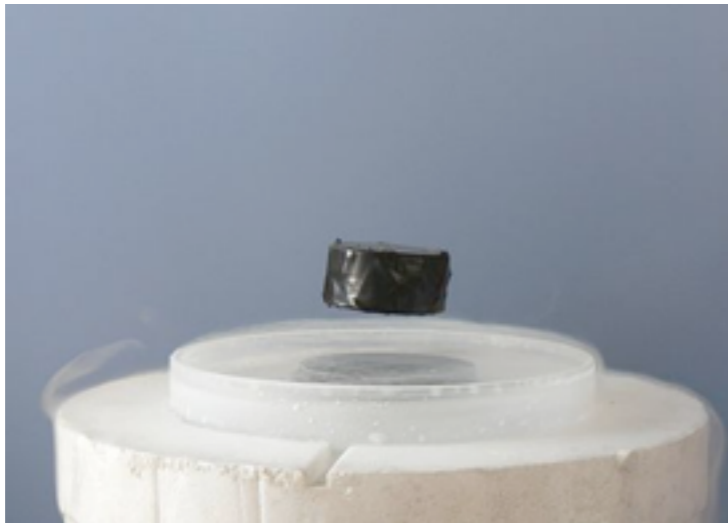
when spin-up and spin-down *fermions* pair up

- What happens when a *spin imbalance* is introduced?

# Relevance for various research fields

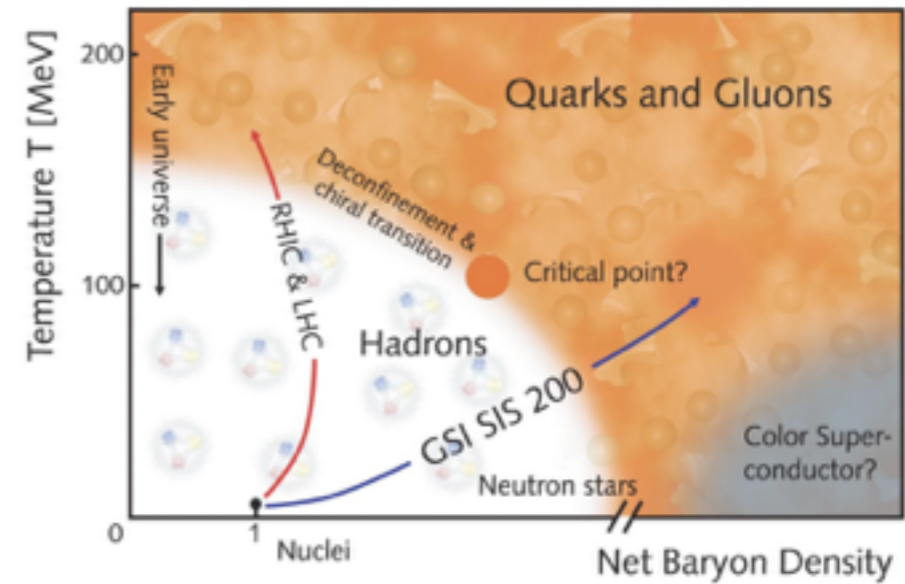
[clearly incomprehensive compilation, guided by the speaker's personal interests]

## Superconductivity



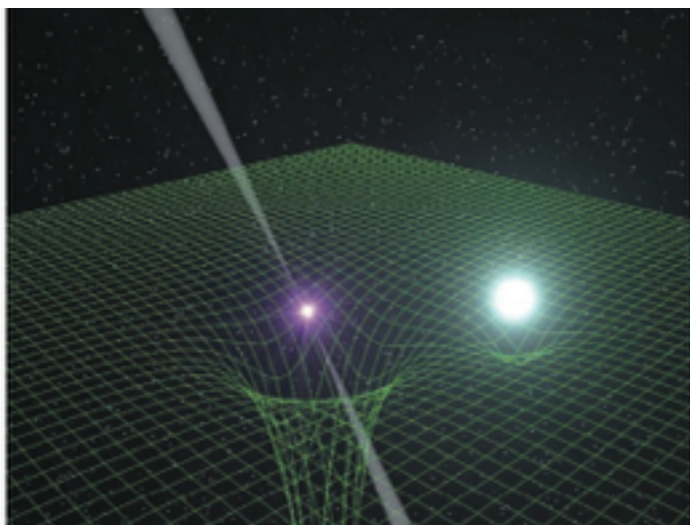
[high-T superconductor, [wikipedia.org](https://en.wikipedia.org)]

## Quantum Chromodynamics



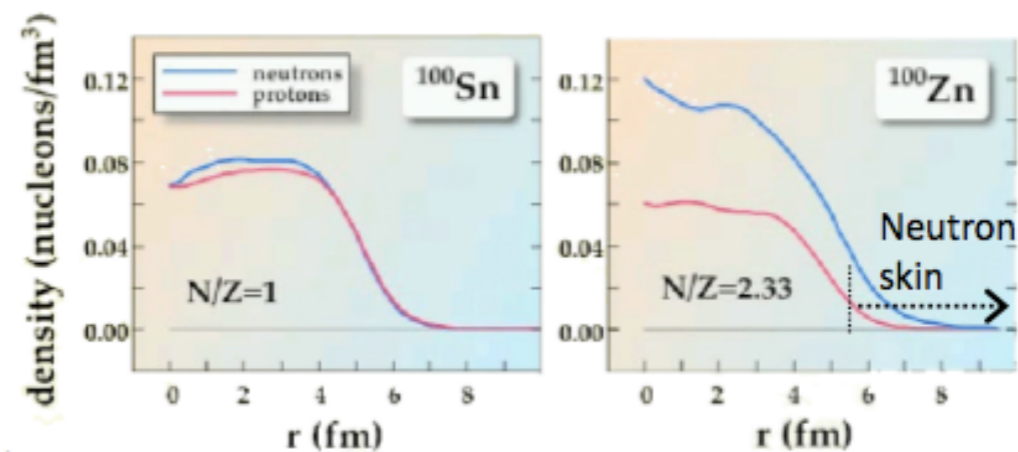
[FAIR, [gsi.de](https://www.gsi.de)]

## Astrophysics



[Demorest *et al.*, Nature (2010); Antoniadis *et al.*, Science (2013)]

## Nuclear Physics

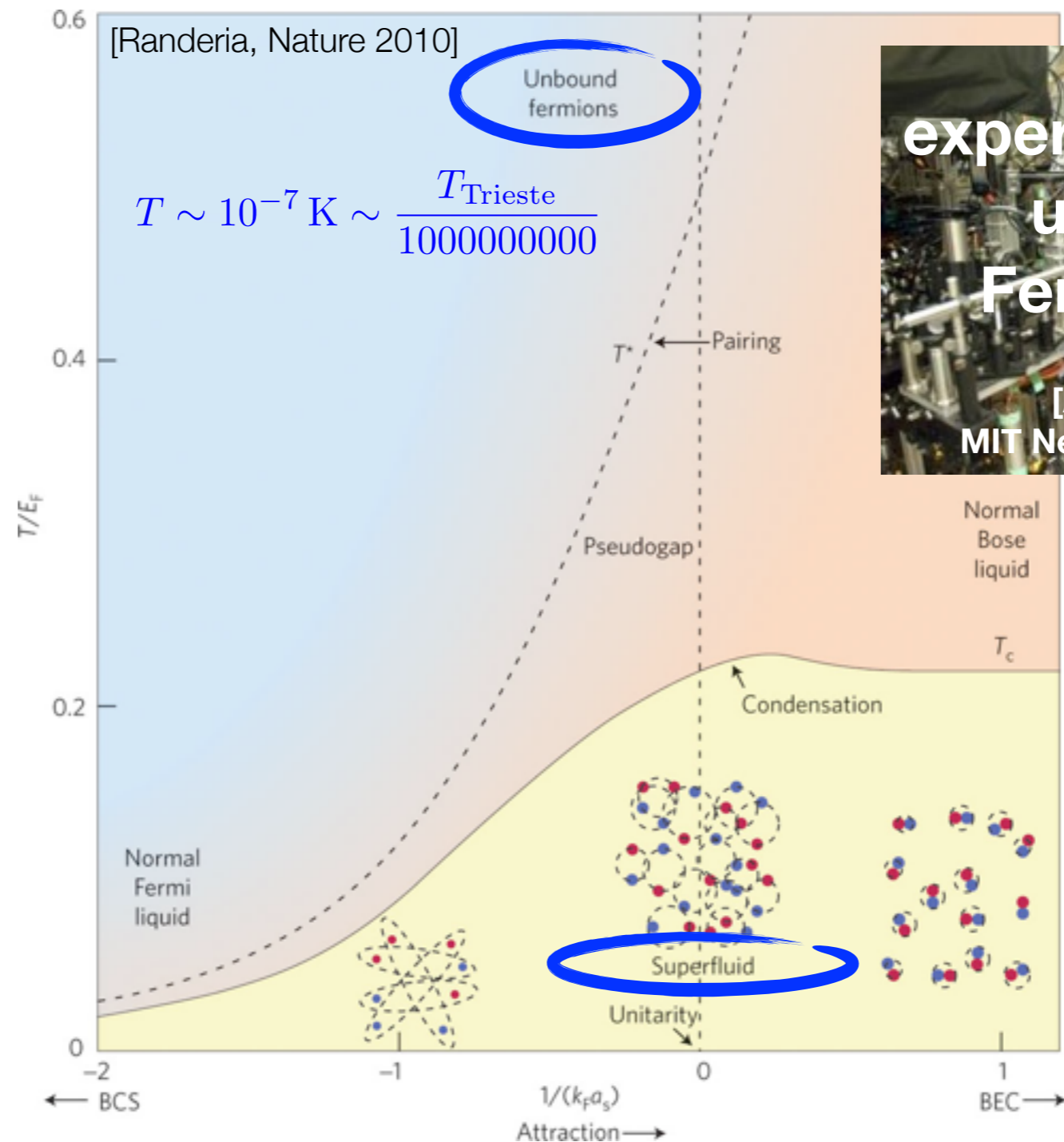


[Dobaczewski *et al.* '94]

Lab for controlled studies of imbalanced systems?

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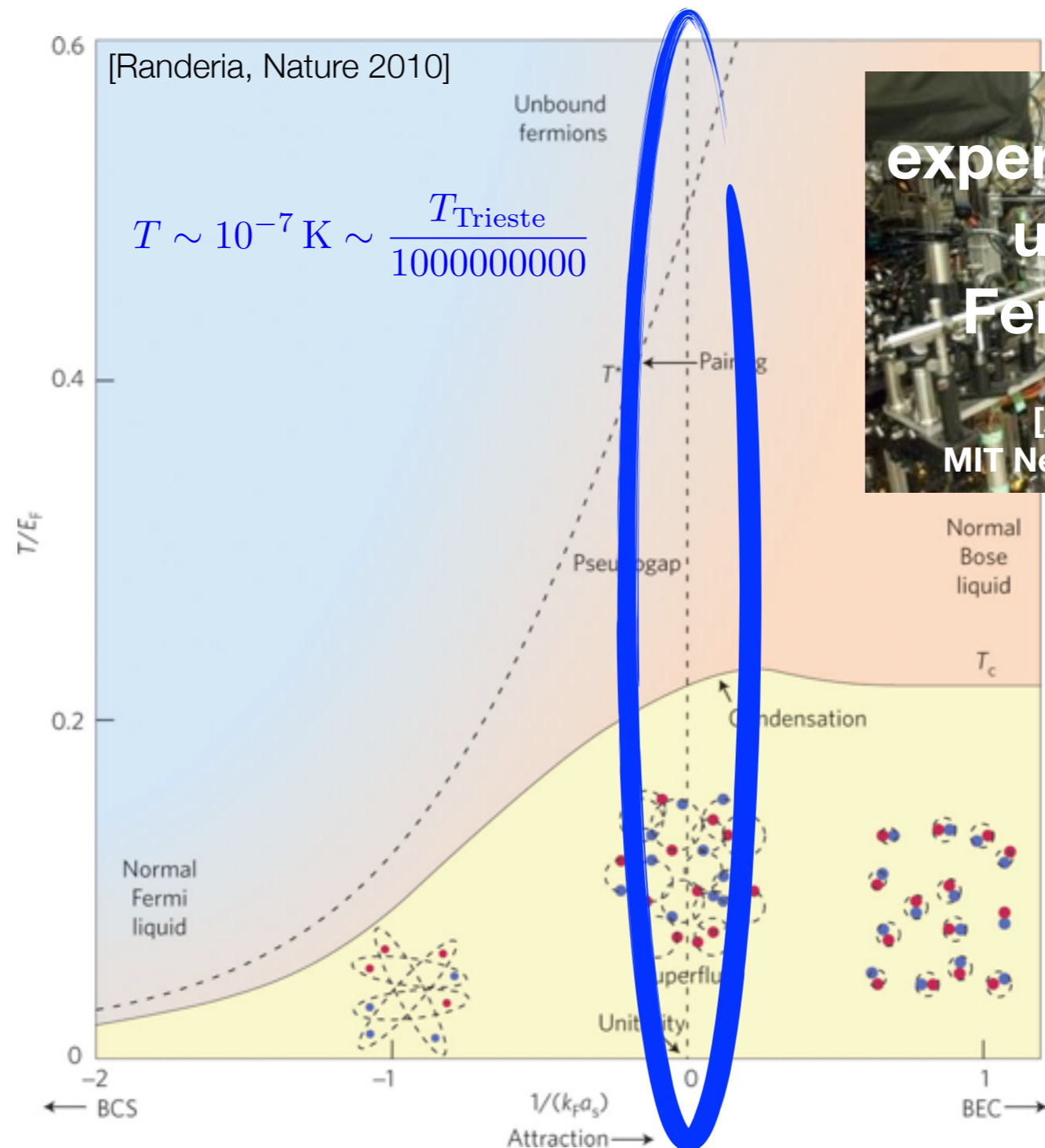
# Lab for controlled studies of imbalanced systems?



experiments with  
ultracold  
Fermi gases

[Anne Trafton  
MIT News Office 09/2009]

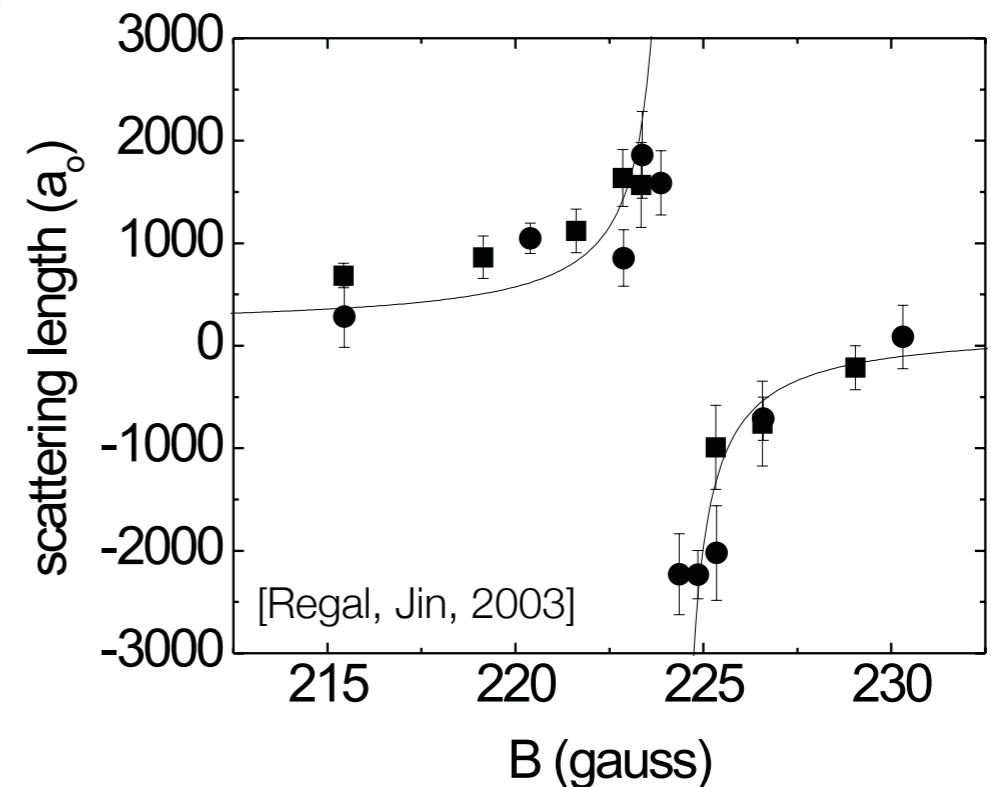
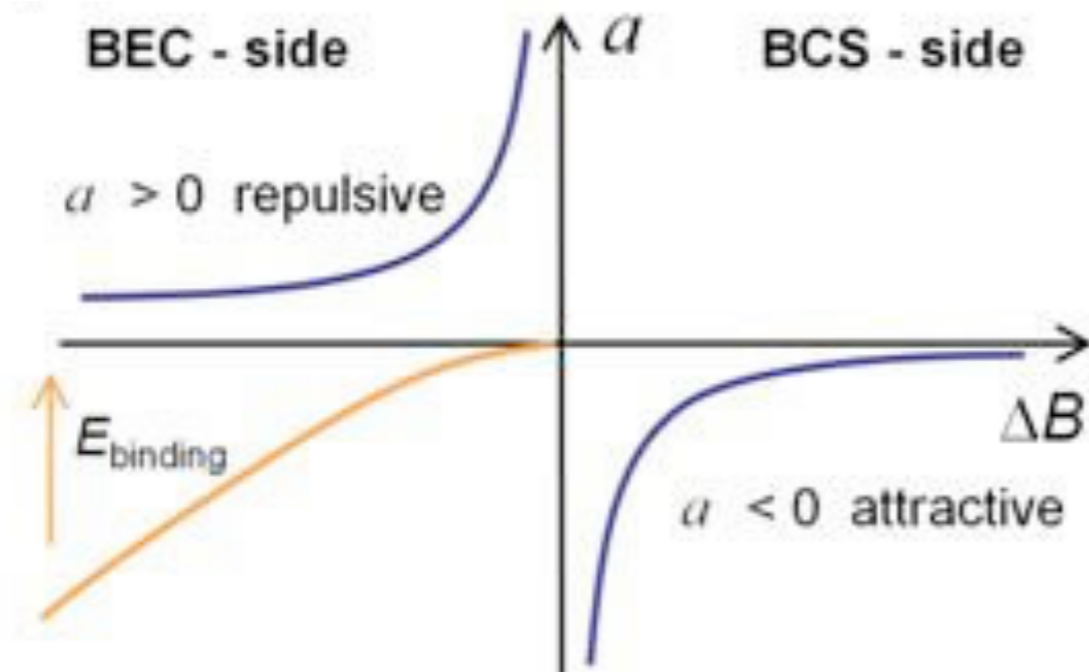
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experiments with  
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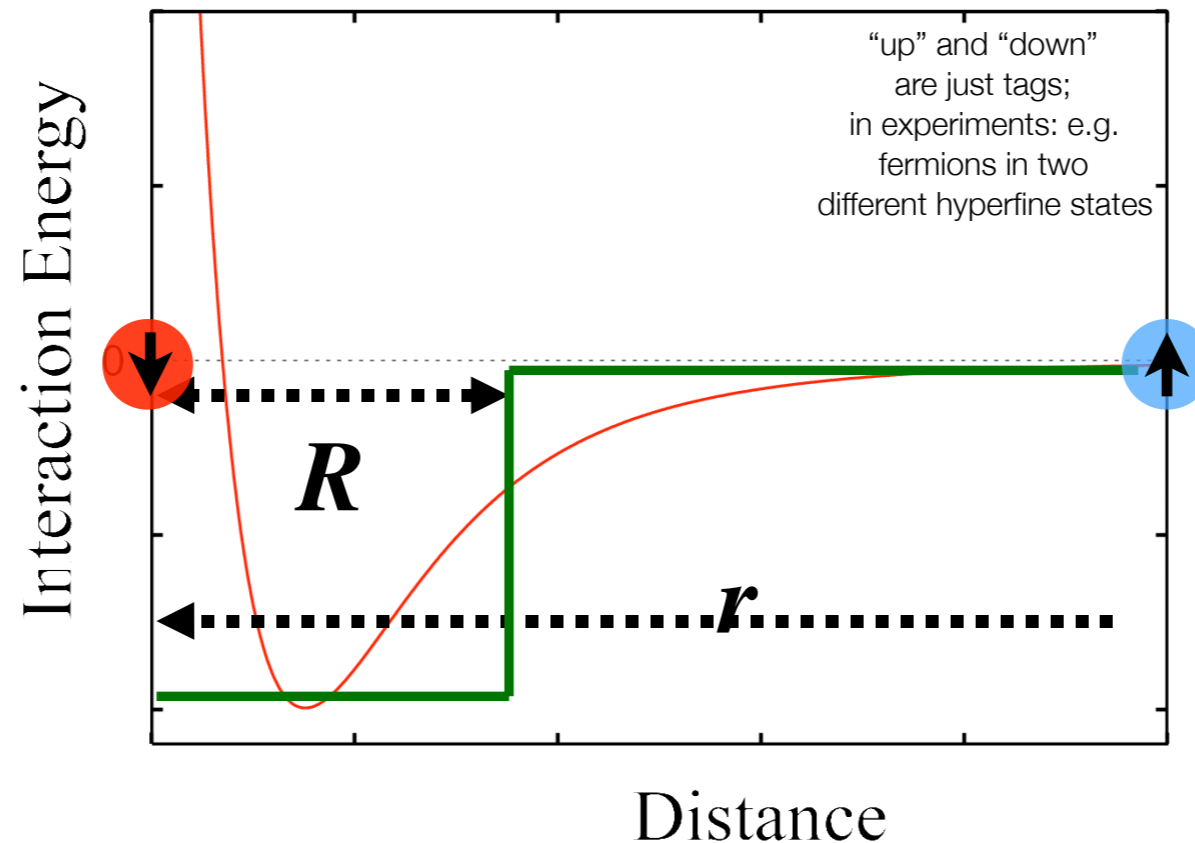
# Unitary Regime



- **s-wave scattering length** is tunable by an external magnetic field (Feshbach resonance)
- interaction strength is directly related to the **s-wave scattering length**  $a_s$



# Unitary Regime

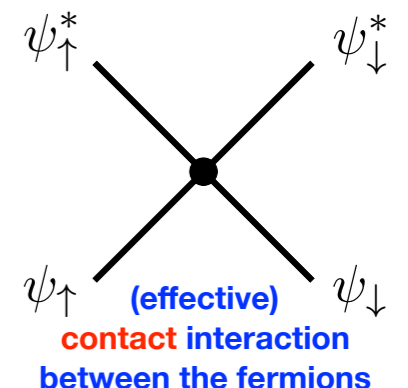


- limit of infinite scattering length  $a_s$  defines a universal regime:

$$0 \approx \frac{1}{|a_s|} \ll \sqrt{\epsilon_F} \sim k_F \sim \frac{1}{r} \ll \frac{1}{R} \approx \infty$$

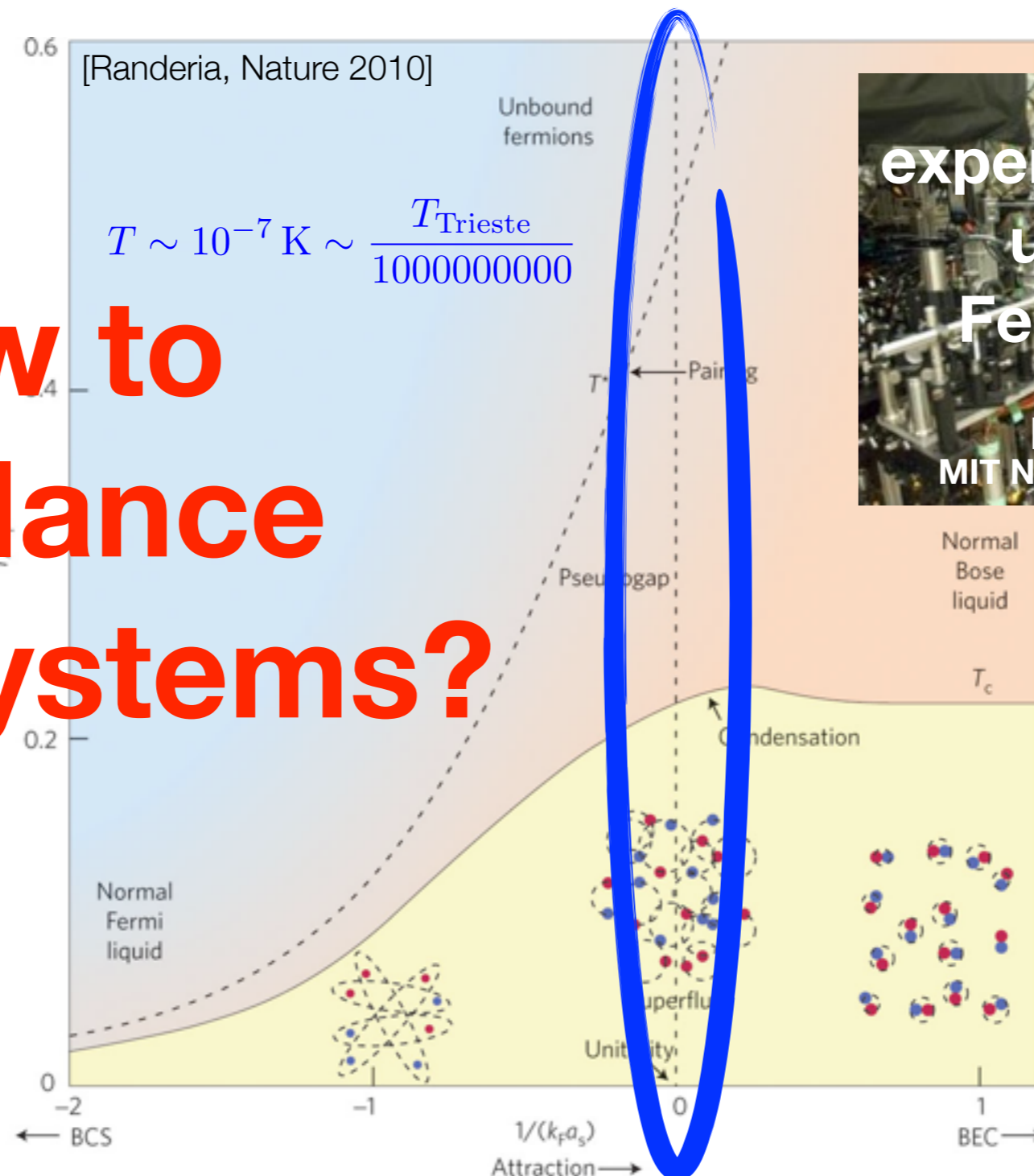
- universal properties:

$$E/N, T_c, \dots \propto \text{universal const(s)} \cdot \epsilon_F$$



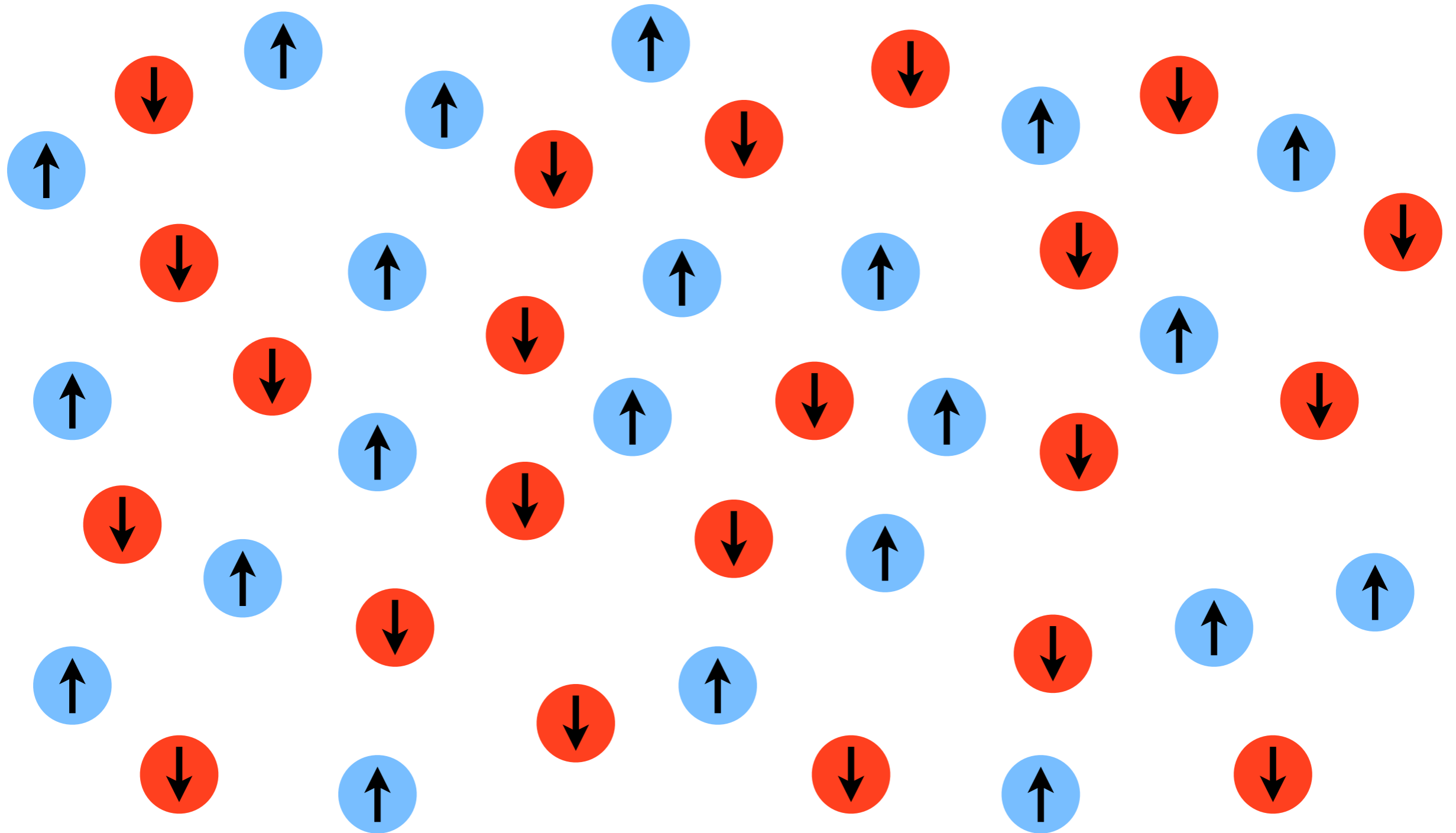
# Lab for controlled studies of imbalanced systems?

**How to  
imbalance  
these systems?**



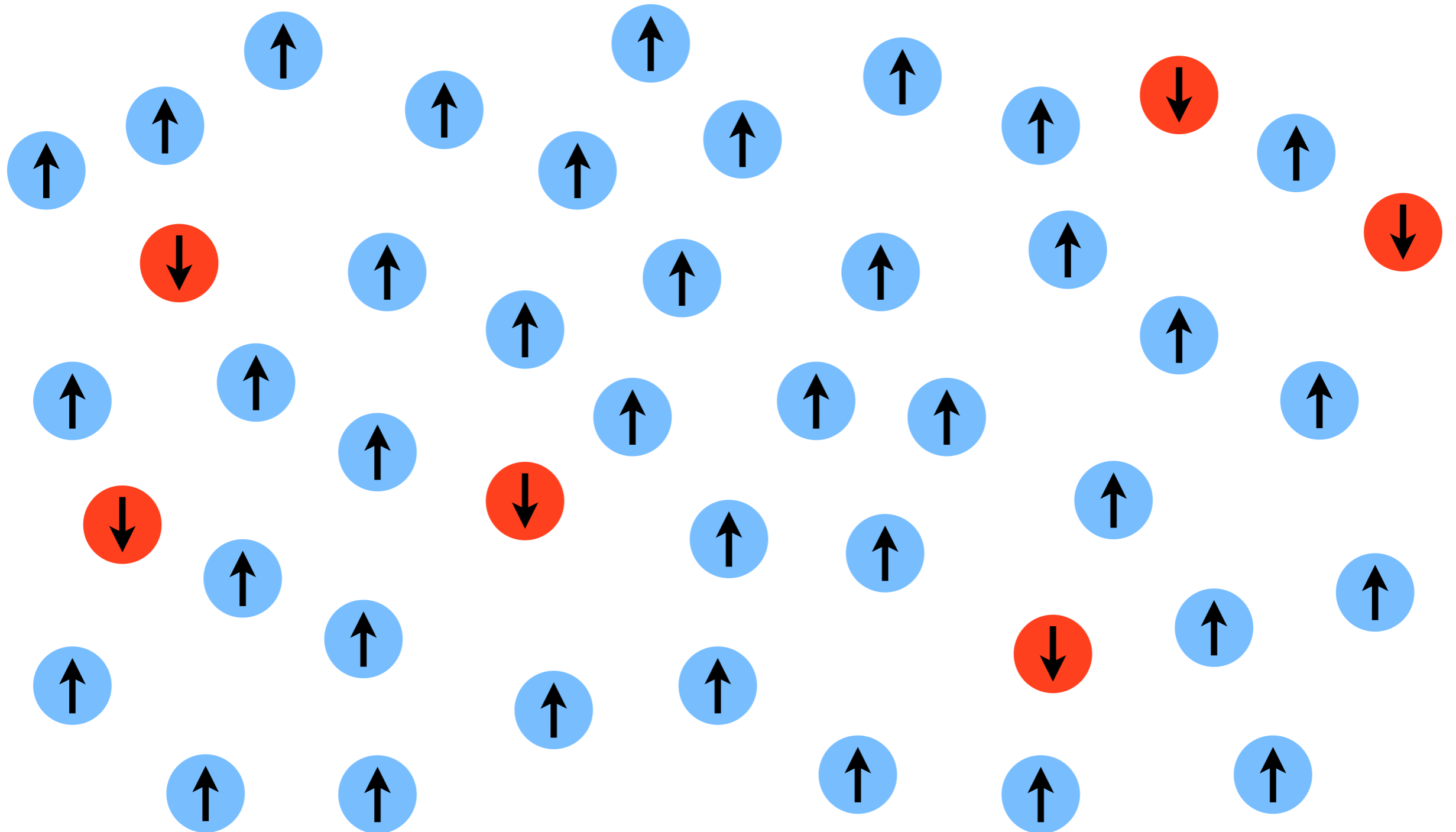
# “Balanced” unitary Fermi gas

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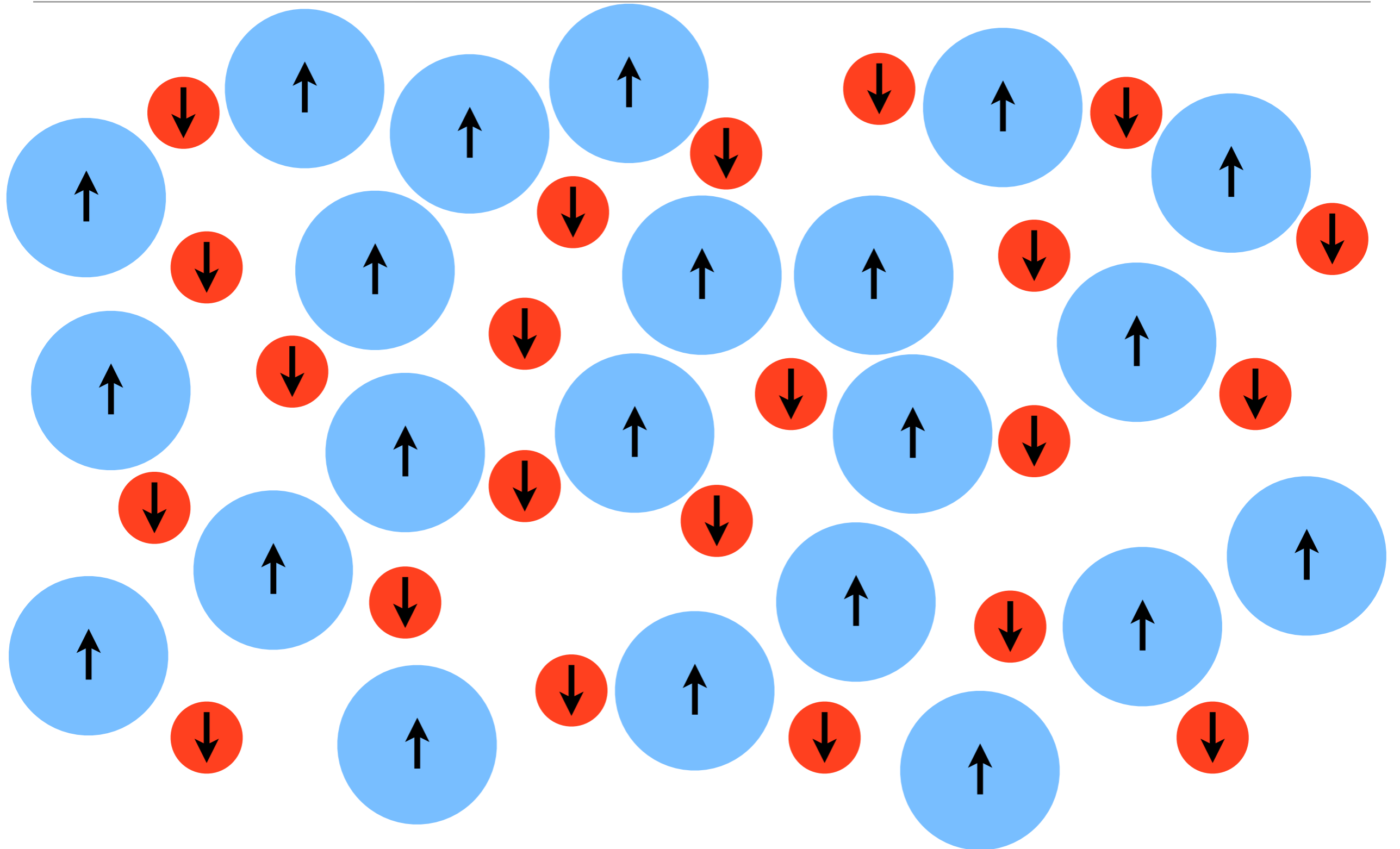
# Deforming Fermi gases: spin imbalance

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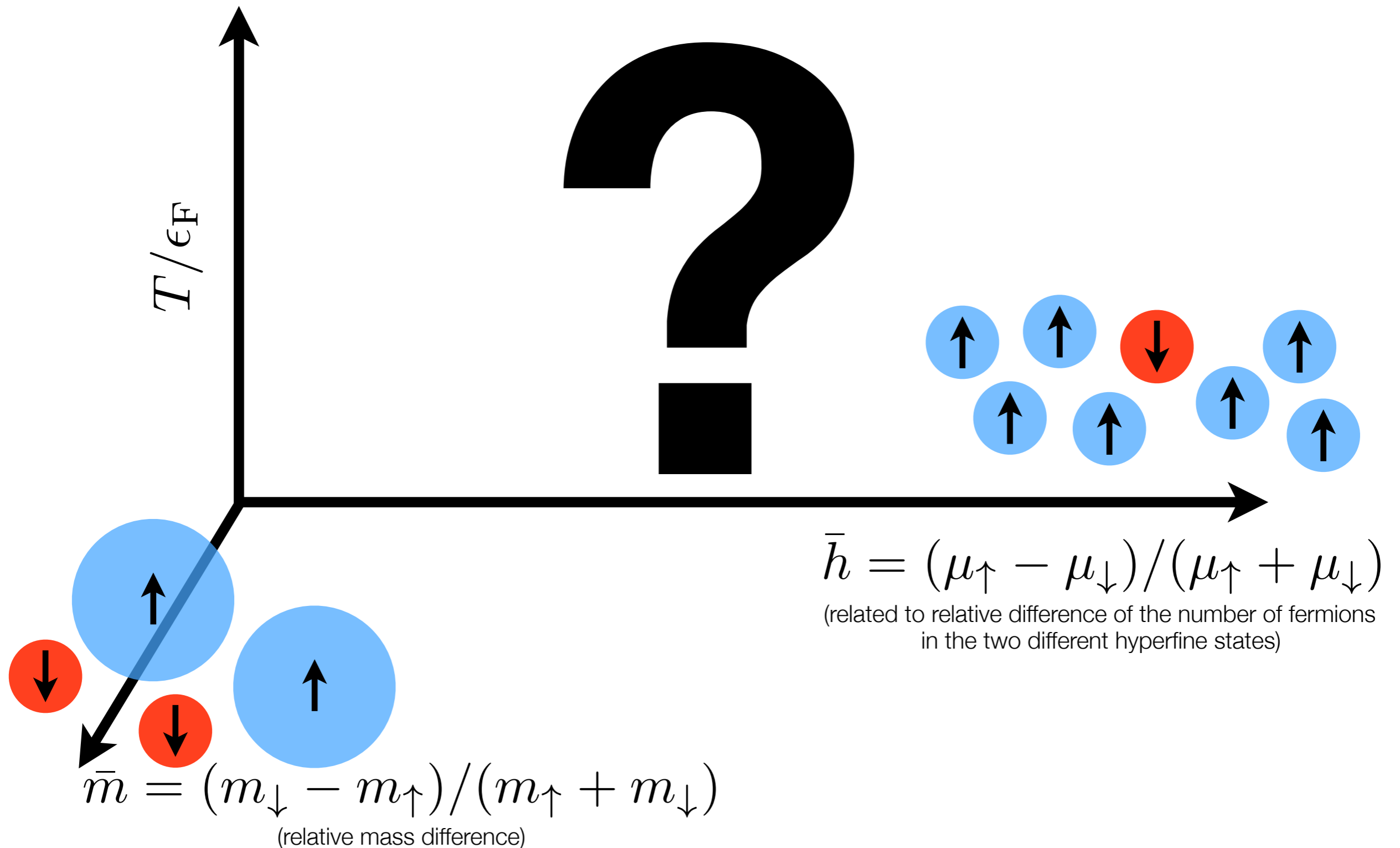


# Deforming Fermi gases: mass imbalance

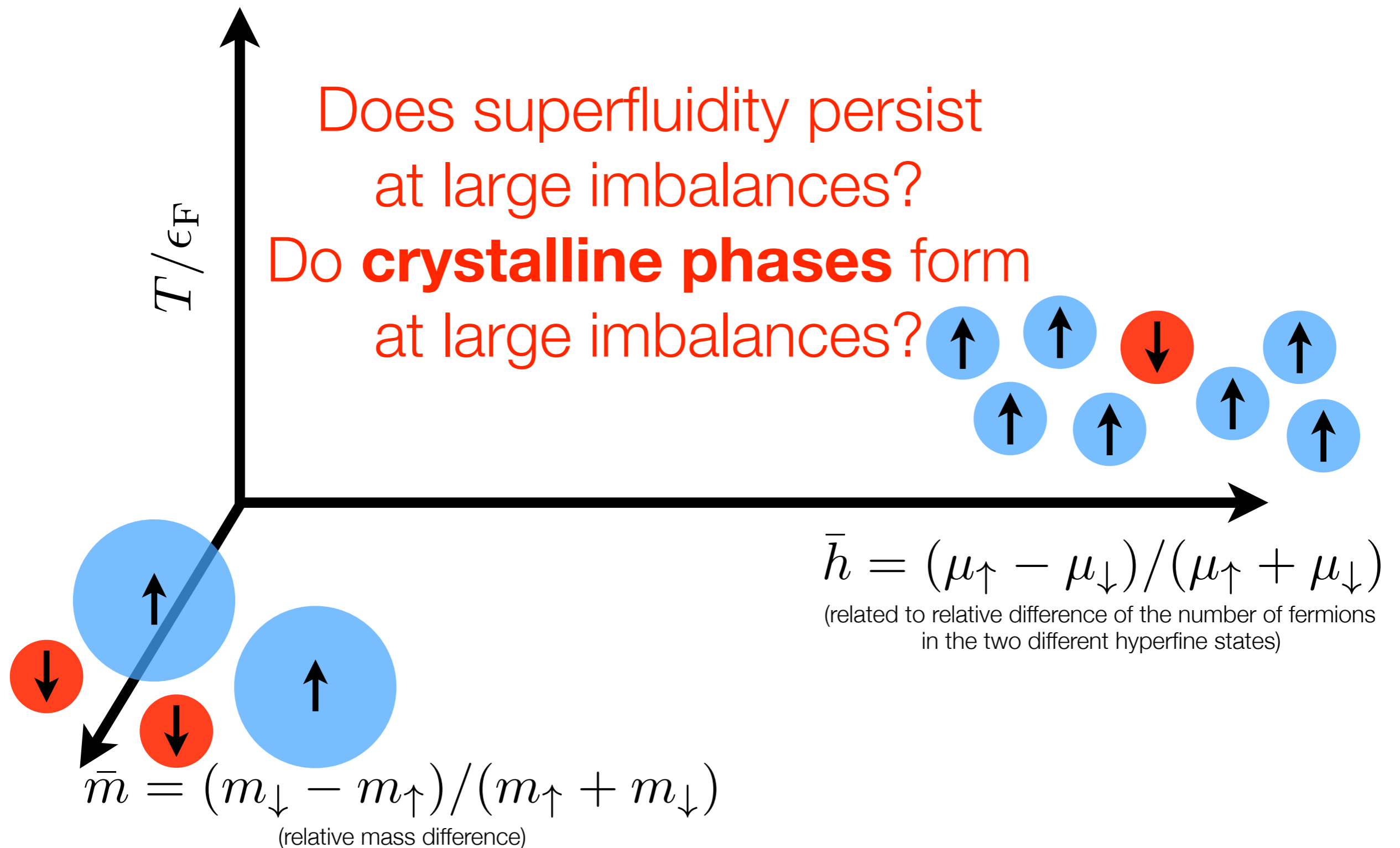
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# Phase diagram of imbalanced unitary Fermi gases



# Phase diagram of imbalanced unitary Fermi gases

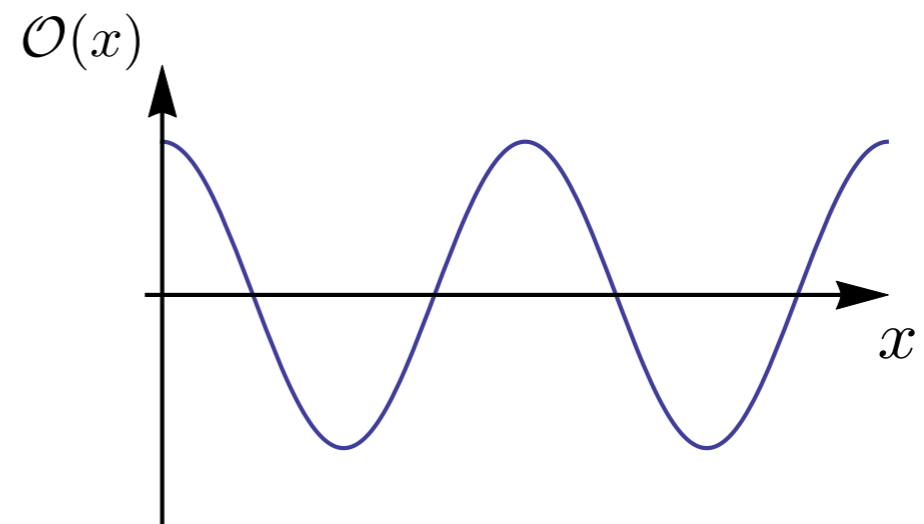
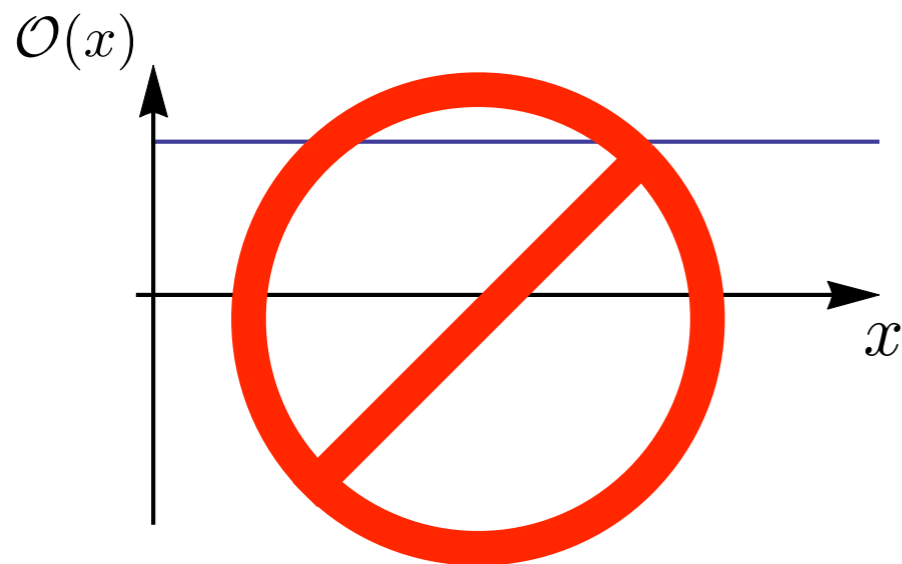


# “Crystal”-type phases in Fermi gases

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- what does “crystal” mean in this case?

Ground-state configuration (and observables) are **not** homogeneous but, e.g., periodic functions of the space-like coordinates







# How to study imbalanced Fermi gases?

Schrödinger equation

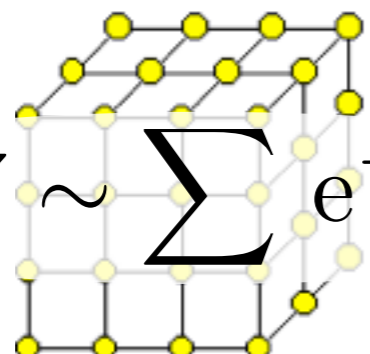
$$\hat{H}|\psi\rangle = i\partial_t|\psi\rangle$$

Wetterich equation

$$\partial_t\Gamma_k = \frac{1}{2}\text{STr}\frac{\partial_t R_k}{\Gamma_k^{(2)} + R_k}$$

[Wetterich '92]

(lattice) Monte Carlo


$$Z \sim \sum e^{-S}$$

Density Functional  
Theory

$$\psi_{\text{gs}}(\vec{x}_1, \dots, \vec{x}_N) \Leftrightarrow n_{\text{gs}}(\vec{x})$$

[Hohenberg & Kohn '65]

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$$Z \sim \sum e^{-S}$$
A diagram of a 3D lattice structure, likely representing a crystal or a lattice of particles. The lattice is composed of yellow spheres connected by lines, forming a cubic grid. The equation  $Z \sim \sum e^{-S}$  is written to the left of the lattice.

Density Functional Theory

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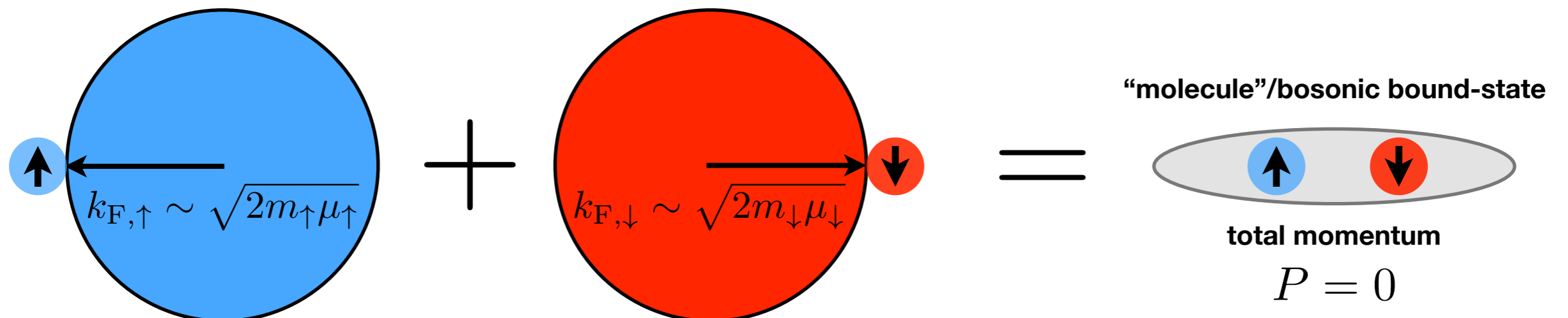
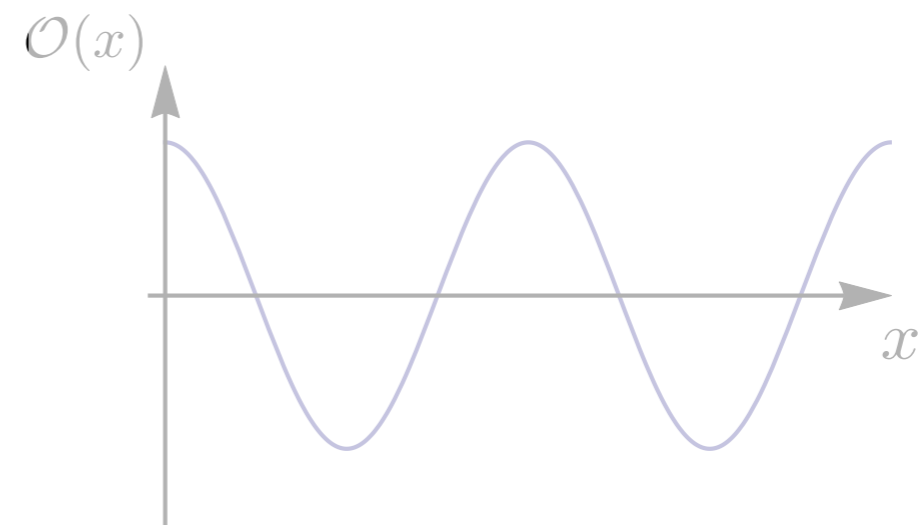
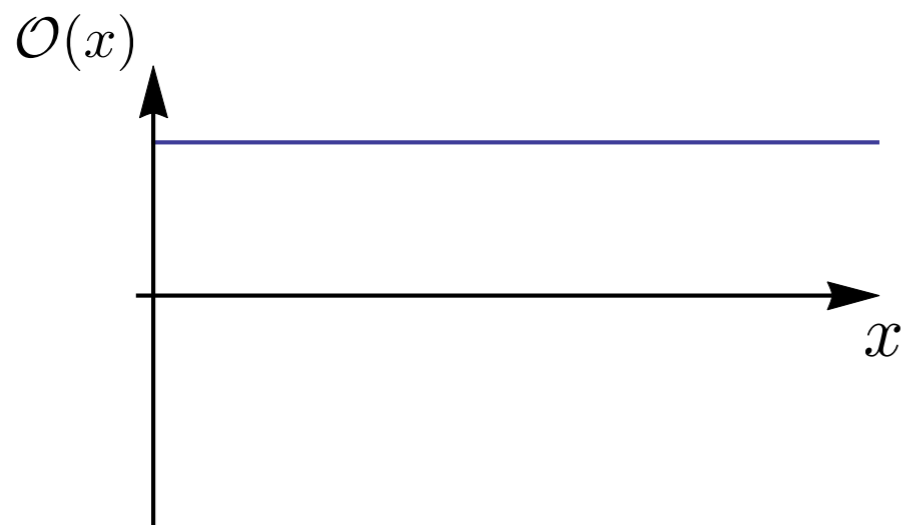
$$\psi_{\text{gs}}(\vec{x}_1, \dots, \vec{x}_N) \Leftrightarrow n_{\text{gs}}(\vec{x})$$

[Hohenberg & Kohn '65]

# “Crystal”-type phases in Fermi gases

- pairing in a nutshell (**BCS**):

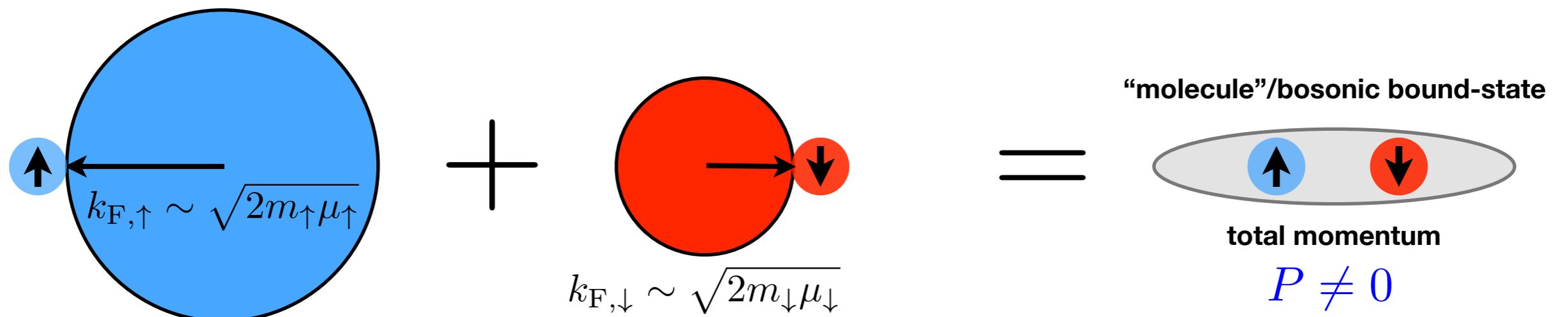
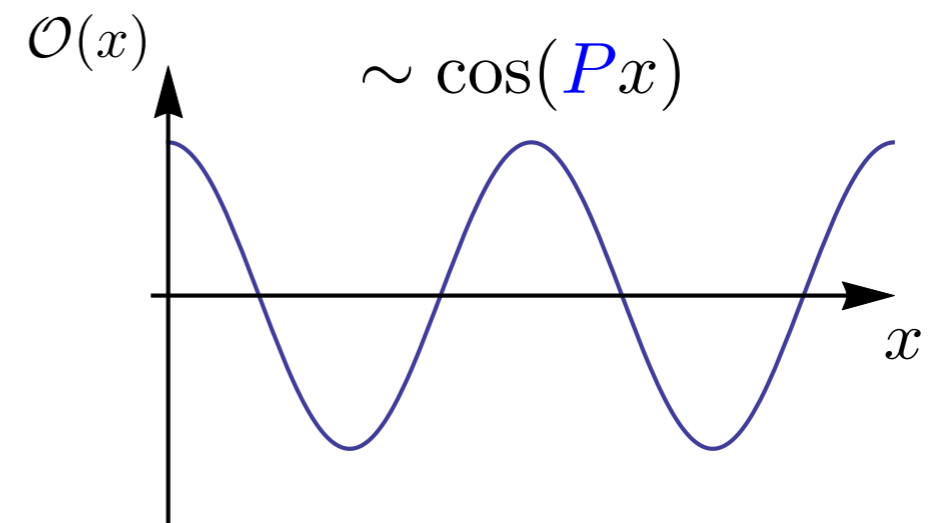
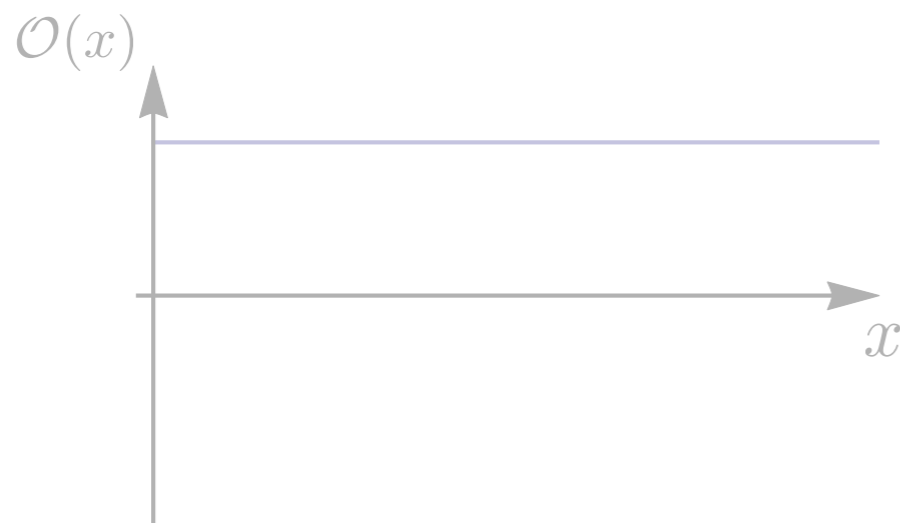
[Bardeen, Cooper, Schrieffer '57]



# “Crystal”-type phases in Fermi gases

- pairing in a nutshell (**FFLO**):

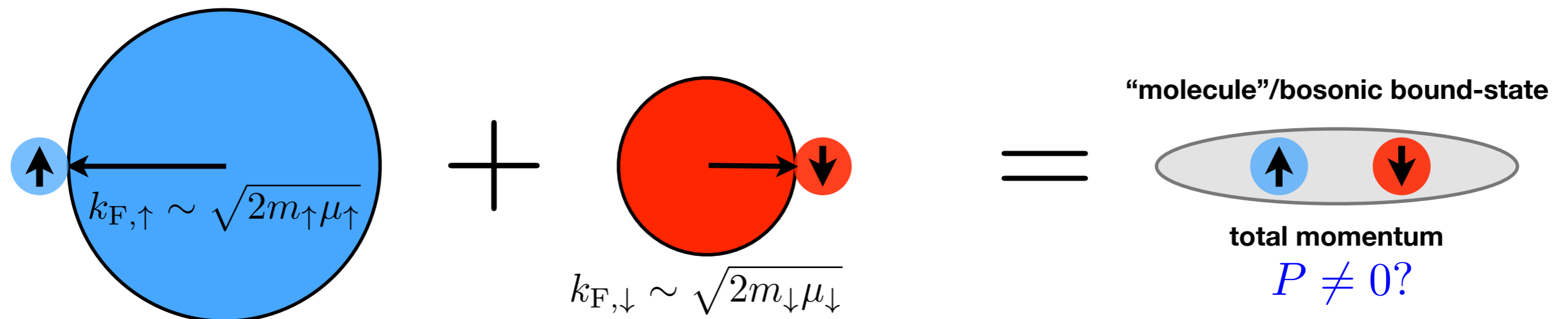
[Fulde, Ferrell '64; Larkin, Ovchinnikov '64]



# “In-medium” two-body Schrödinger equation

[Roscher, JB, Drut '14, '15]

- study pairing in the presence of “inert” Fermi spheres:



- Schrödinger equation:

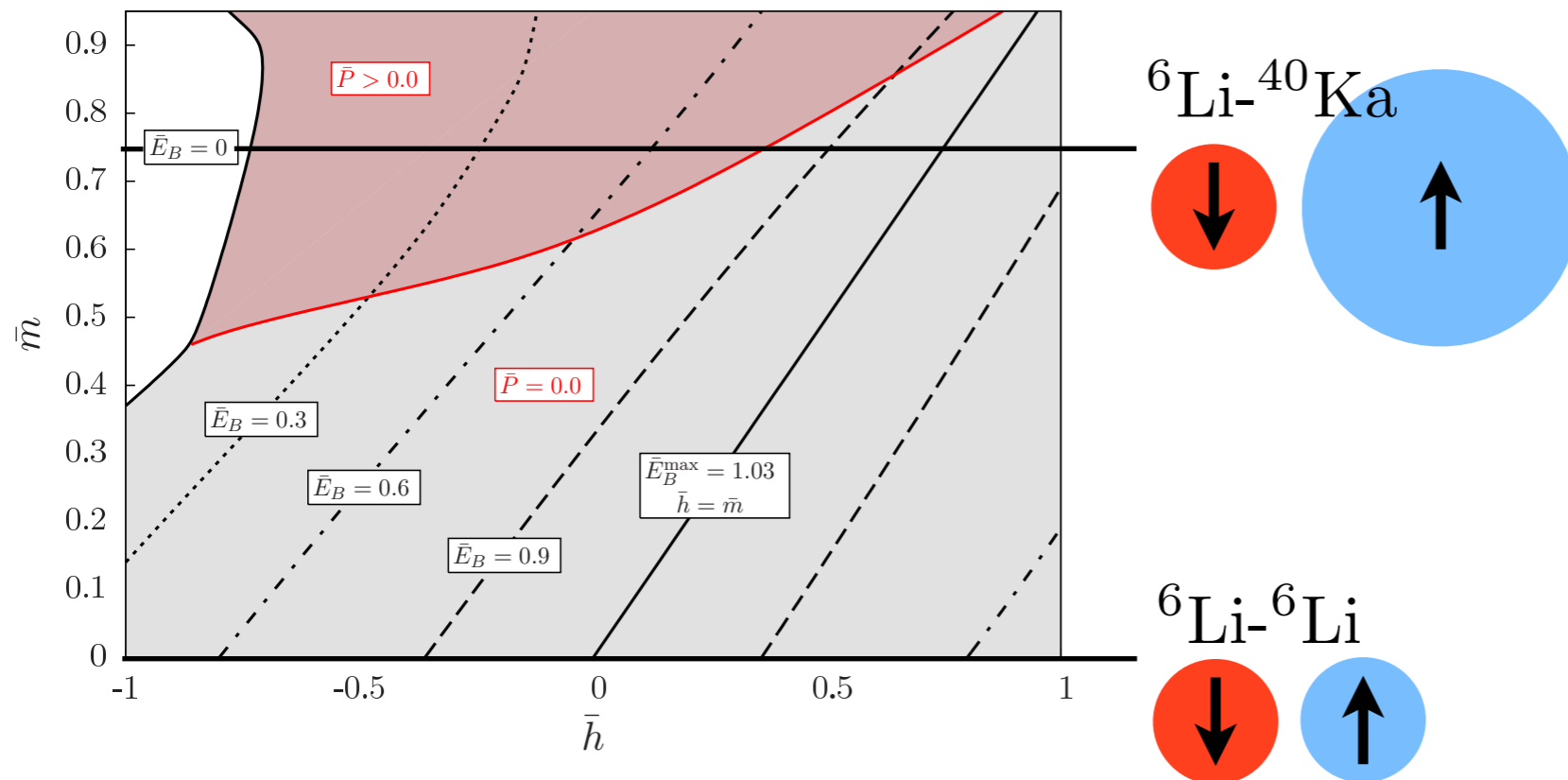
$$\left[ \sum_{\sigma=\uparrow,\downarrow} \epsilon_\sigma(\partial_{x_\sigma}) - g\delta(x_\uparrow - x_\downarrow) + E_B \right] \Psi(x_\uparrow, x_\downarrow) = 0$$

- dispersion relation:

$$\epsilon_\sigma(\partial_{x_\sigma}) = \left| - (2m_\sigma)^{-1} \partial_{x_\sigma}^2 - \epsilon_{F,\sigma} \right|$$

# Phase diagram of imbalanced Fermi gases I

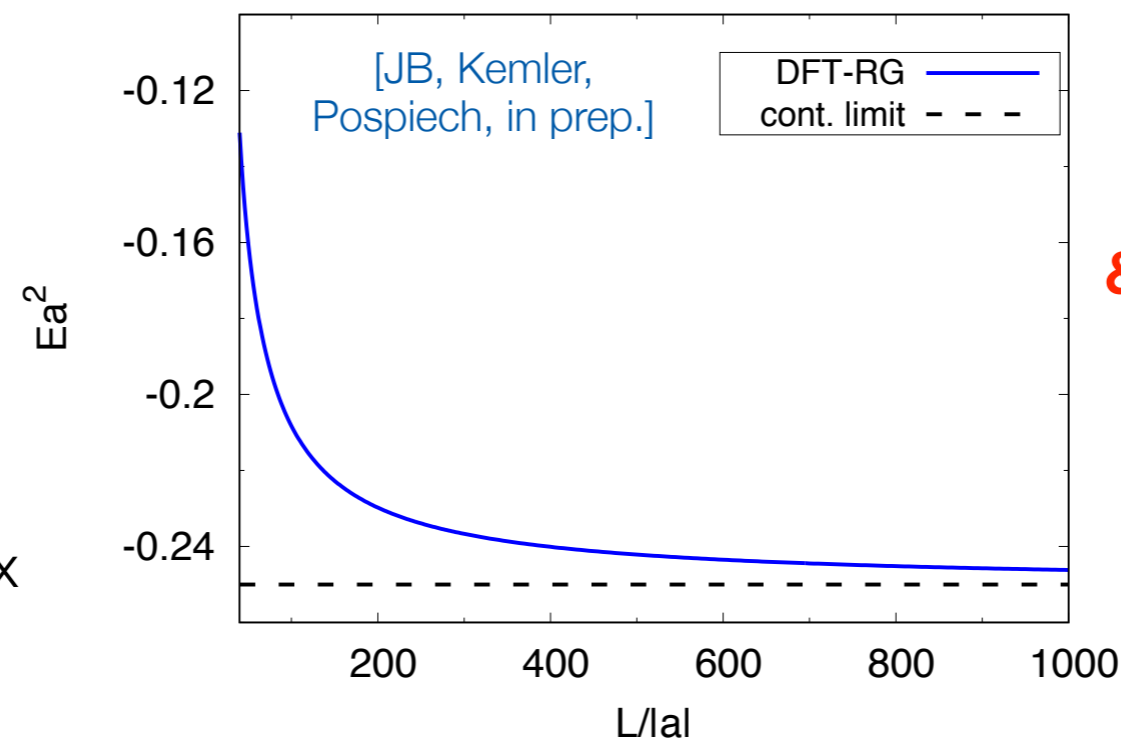
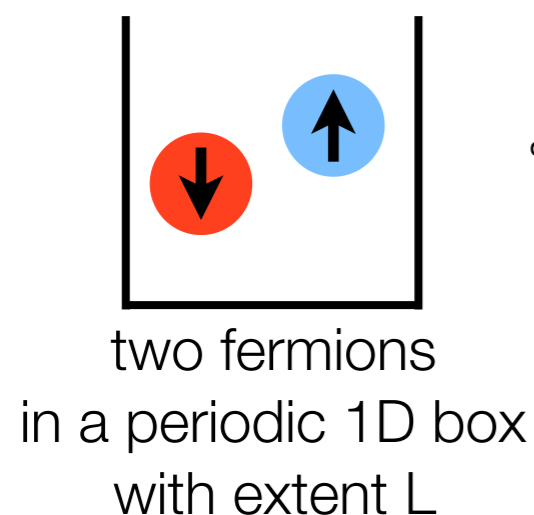
[Roscher, JB, Drut '15]



- **bound-state formation** in the presence of inert Fermi spheres
- necessary condition for macroscopic occupation of the ground state ( $\sim$ condensation): bound-state formation

# From few to many?

- approach via Schrödinger equation becomes essentially impossible for large fermion numbers to study the fully interacting problem [“exponential barrier”, Kohn, Nobel lecture '99]
- functional methods allow to bridge the gap from “few to many” efficiently, such as an **RG approach to Density Functional Theory** [Polonyi & Sailer '02, Polonyi & Schwenk '04; JB '11; Kemler & JB '13; Kemler, Pospiech, JB '16; JB, Kemler, Pospiech, in prep.]



**cf. talk by Sandra Kemler  
& poster by Martin Pospiech  
on DFT-RG study of a  
nuclear physics model**



# How to study imbalanced Fermi gases?

Schrödinger equation

$$\hat{H}|\psi\rangle = i\partial_t|\psi\rangle$$

Wetterich equation

$$\partial_t\Gamma_k = \frac{1}{2}\text{STr}\frac{\partial_t R_k}{\Gamma_k^{(2)} + R_k}$$

[Wetterich '92]



(lattice) Monte Carlo

$$Z \sim \sum e^{-S}$$

Density Functional Theory

$$\psi_{\text{gs}}(\vec{x}_1, \dots, \vec{x}_N) \Leftrightarrow n_{\text{gs}}(\vec{x})$$

[Hohenberg & Kohn '65]

# Imbalanced Fermi gases & Wetterich equation

[Roscher, JB, Drut '15]

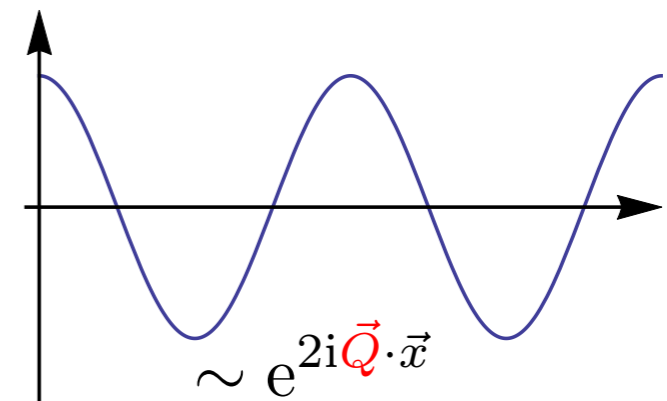
- classical action (partially bosonized version):

$$S[\{\psi_\sigma\}, \bar{\varphi}] = \int_{\tau, \vec{x}} \left[ \sum_{\sigma=\uparrow, \downarrow} \psi_\sigma^* \left( \partial_\tau - \frac{\nabla^2}{2m_\sigma} - \mu_\sigma \right) \psi_\sigma + \bar{m}_\varphi^2 \bar{\varphi}^* \bar{\varphi} - \bar{h}_\varphi \left( \bar{\varphi}^* \psi_\uparrow \psi_\downarrow - \bar{\varphi} \psi_\uparrow^* \psi_\downarrow^* \right) \right].$$

- consider **space-dependent** background/mean field:

$$\bar{\varphi}(\vec{x}) = \bar{\varphi}_0 e^{2i\vec{Q}\cdot\vec{x}}$$

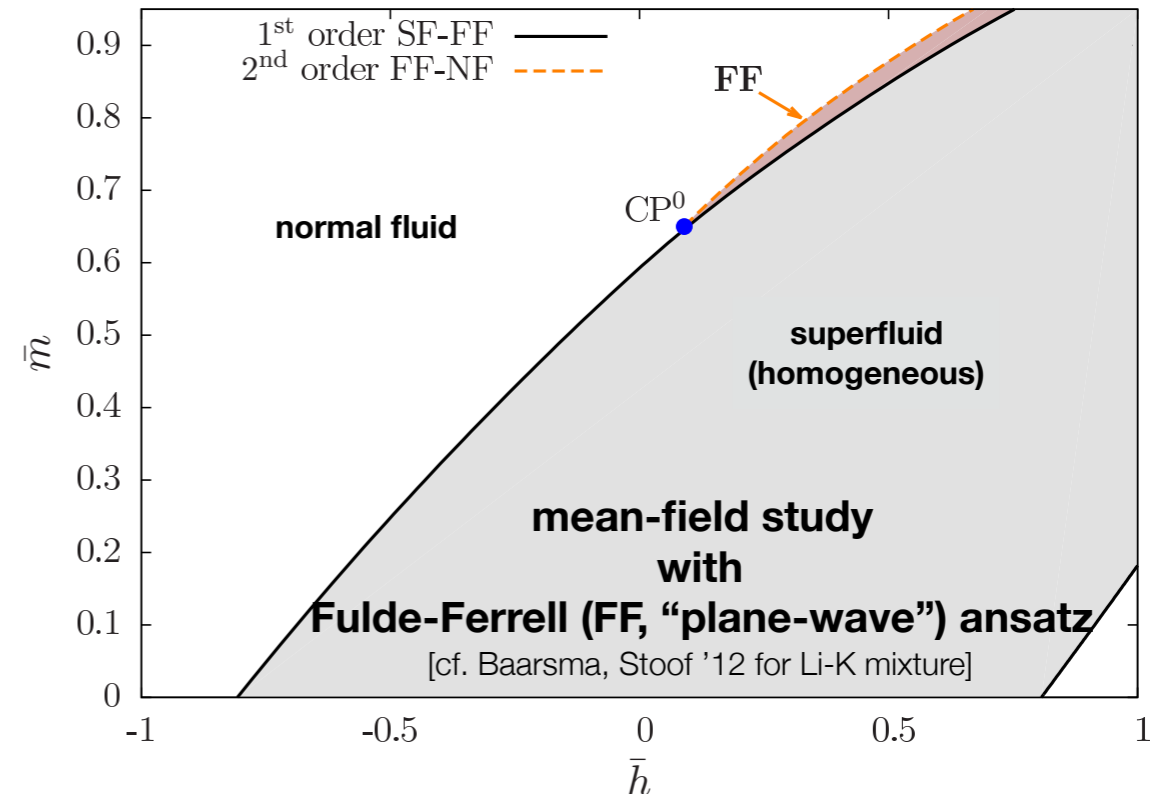
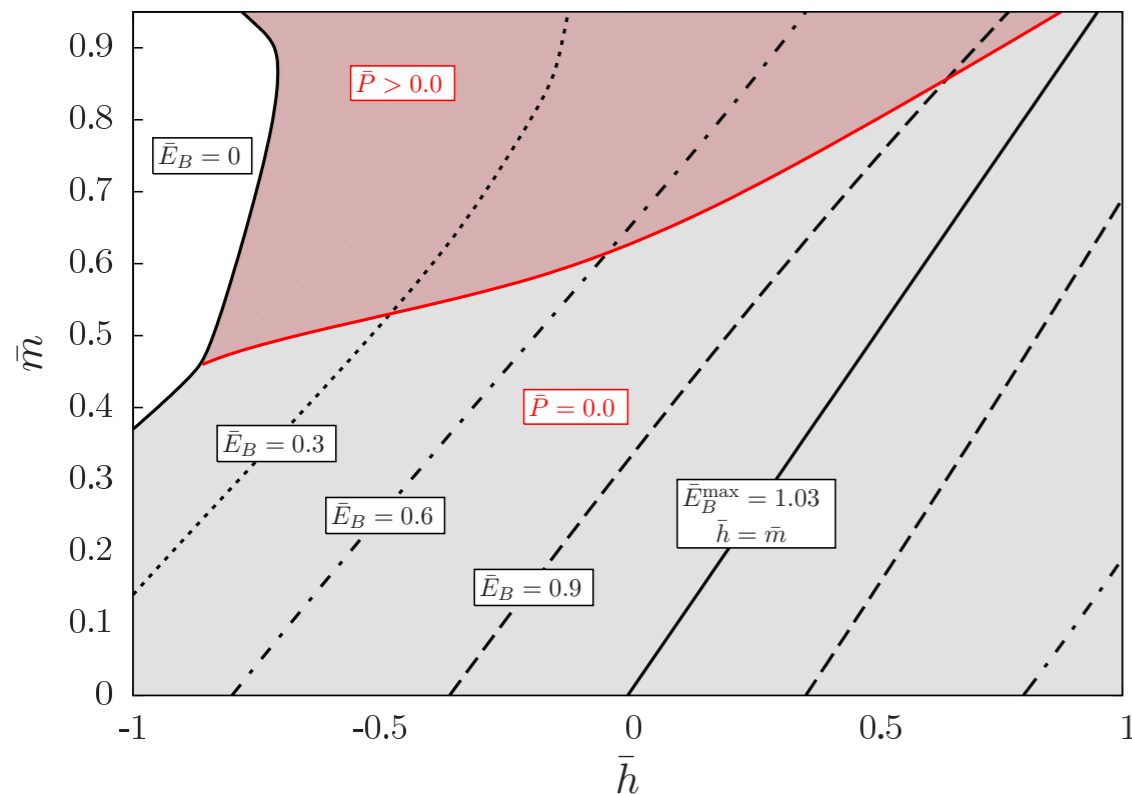
[Fulde-Ferrell ansatz]



- lowest order: **mean-field truncation** (loops only with internal fermion lines); minimize with respect to  $\bar{\varphi}_0$  &  $\vec{Q}$

# Phase diagram of imbalanced Fermi gases II

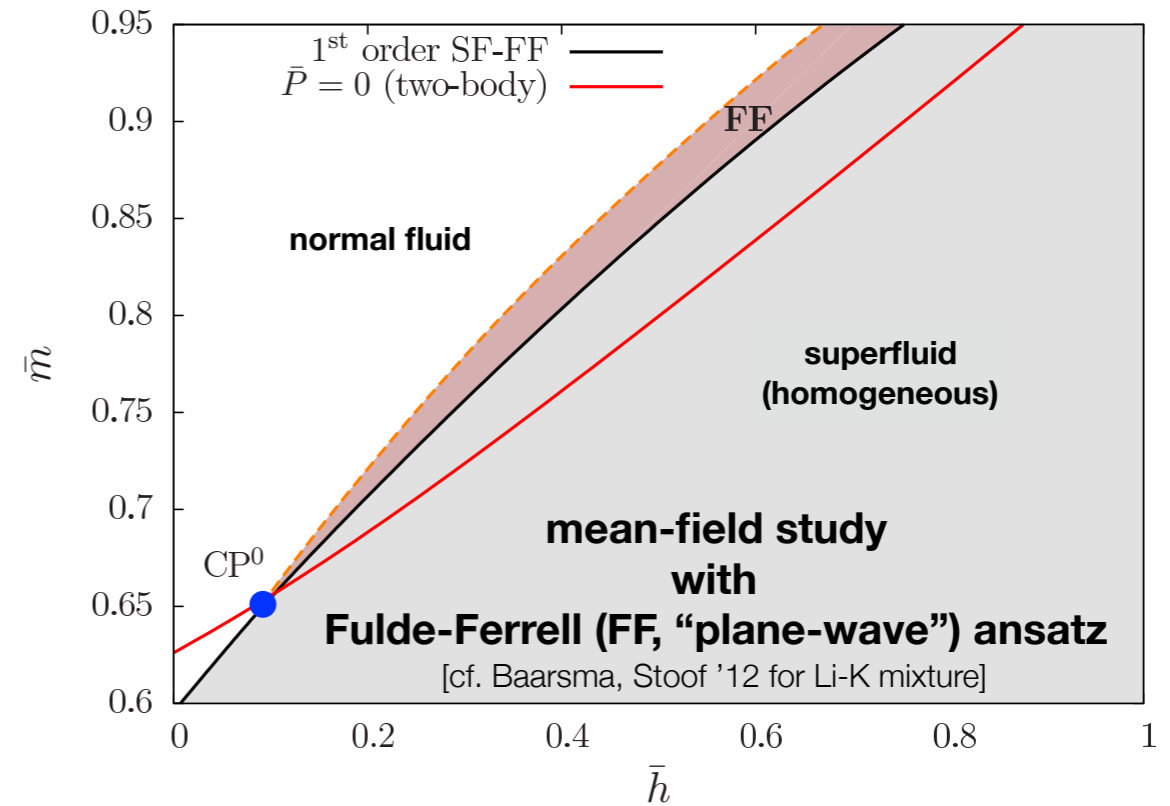
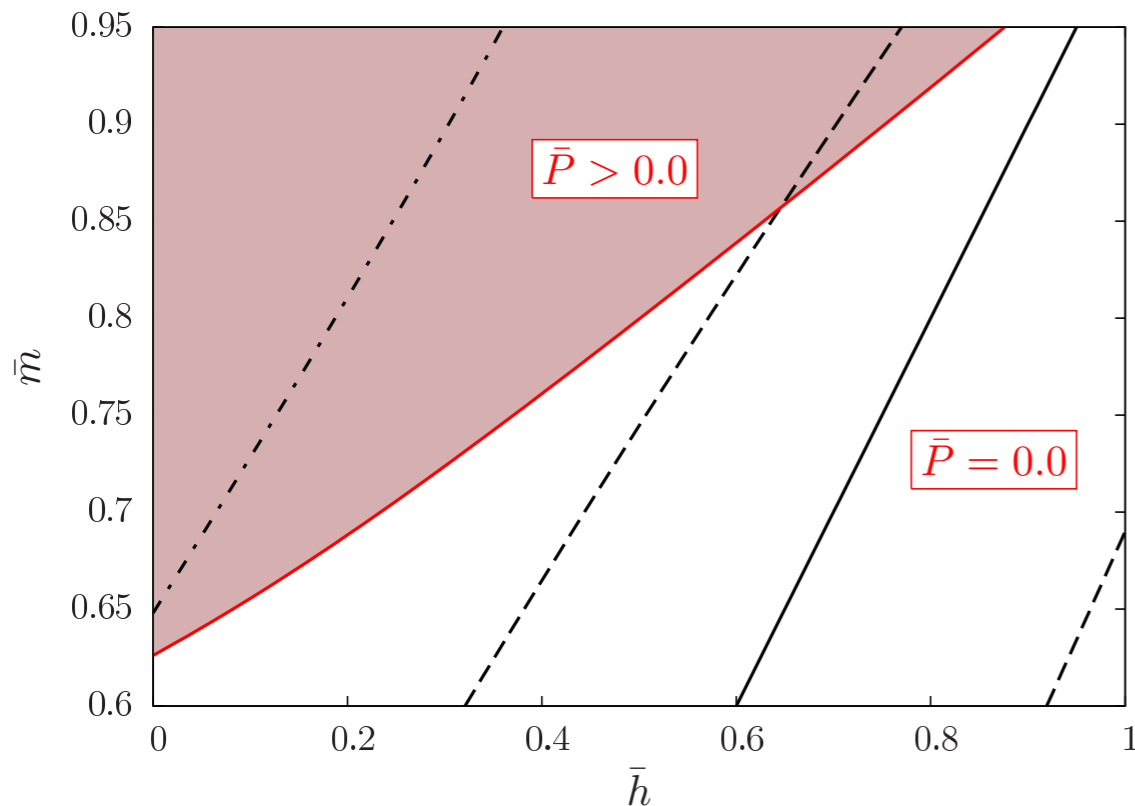
[Roscher, JB, Drut '15]



- **bound-state formation** in the presence of inert Fermi spheres [left panel]
- necessary condition for macroscopic occupation of the ground state ( $\sim$ condensation): bound-state formation
- **many-body phase diagram** [right panel]

# Phase diagram of imbalanced Fermi gases II

[Roscher, JB, Drut '15]



- **bound-state formation** in the presence of inert Fermi spheres [left panel]
- necessary condition for macroscopic occupation of the ground state ( $\sim$ condensation): bound-state formation
- **many-body phase diagram** [right panel]

# Mean-field study “beyond” the Fulde-Ferrell ansatz

[Roscher, JB, Drut '15]

- use **spatially constant** background/mean field
- inverse propagator of the order-parameter field  $\varphi$ :

$$\bar{P}_\varphi(q_0, \vec{q}^2) = \left( iZ_{\varphi,k}q_0 + \frac{1-\bar{m}^2}{2} A_{\varphi,k} \vec{q}^2 + \frac{\partial^2 U}{\partial \bar{\varphi} \partial \bar{\varphi}^*} \Big|_{\bar{\varphi}_0} \right)$$

- derivative expansion:  $A_{\varphi,k}(\vec{q}^2) = A_{\varphi,k}^{(1)} + A_{\varphi,k}^{(2)} \vec{q}^2 + \dots$

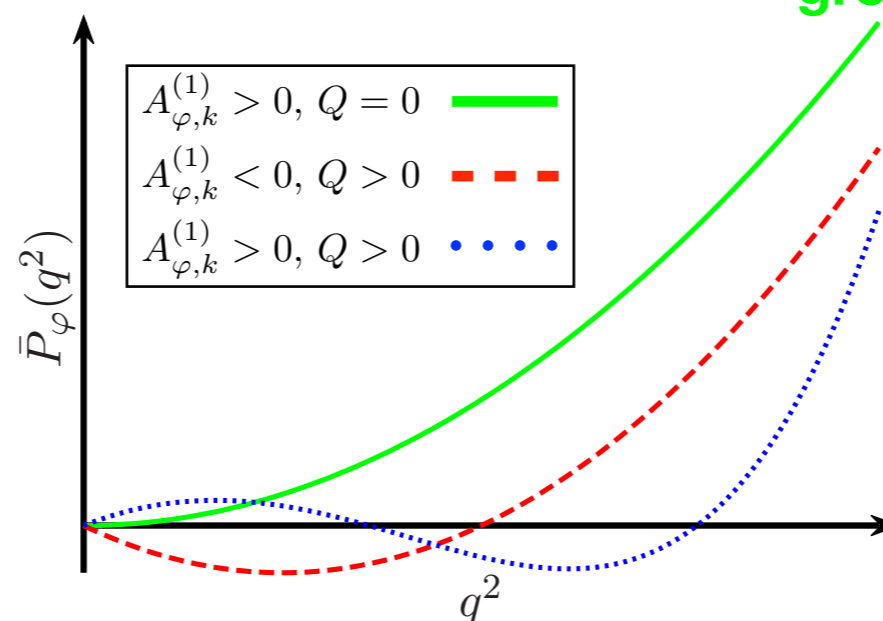
- different scenarios:

**green:** homogenous ground state

**red/blue:** “dominance” of modes favoring a “crystalline” phase

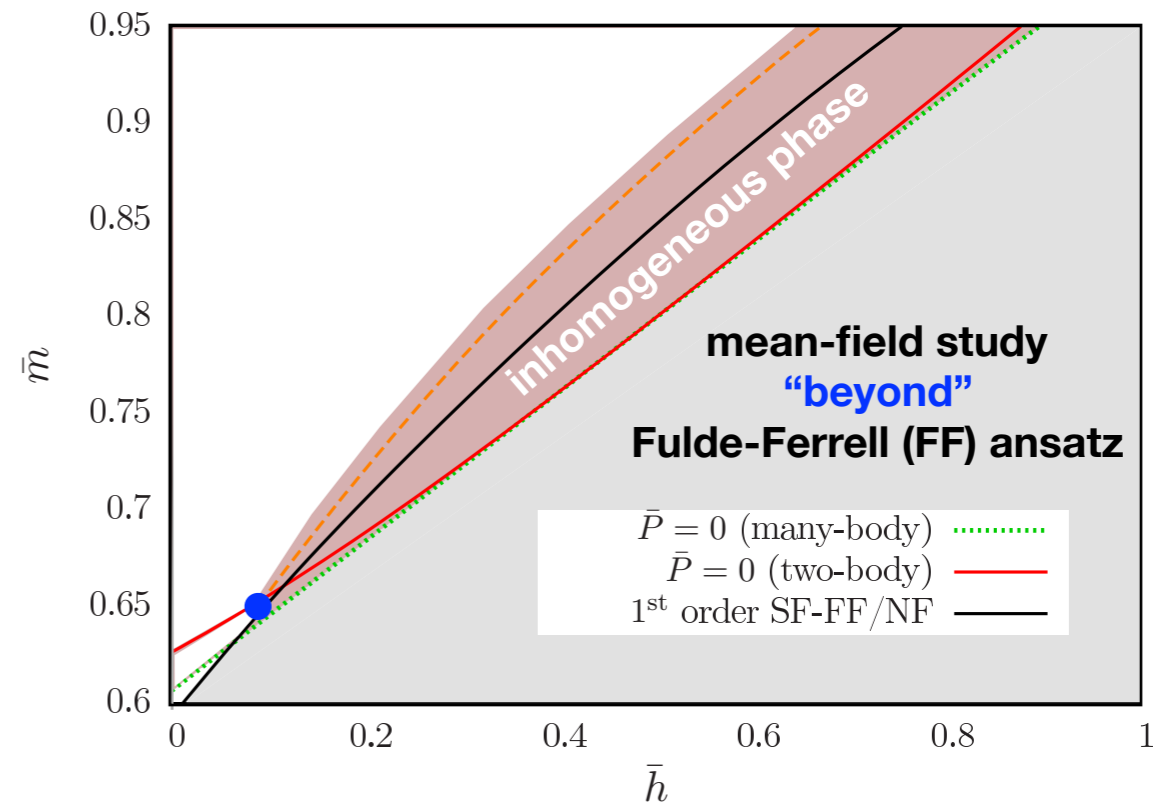
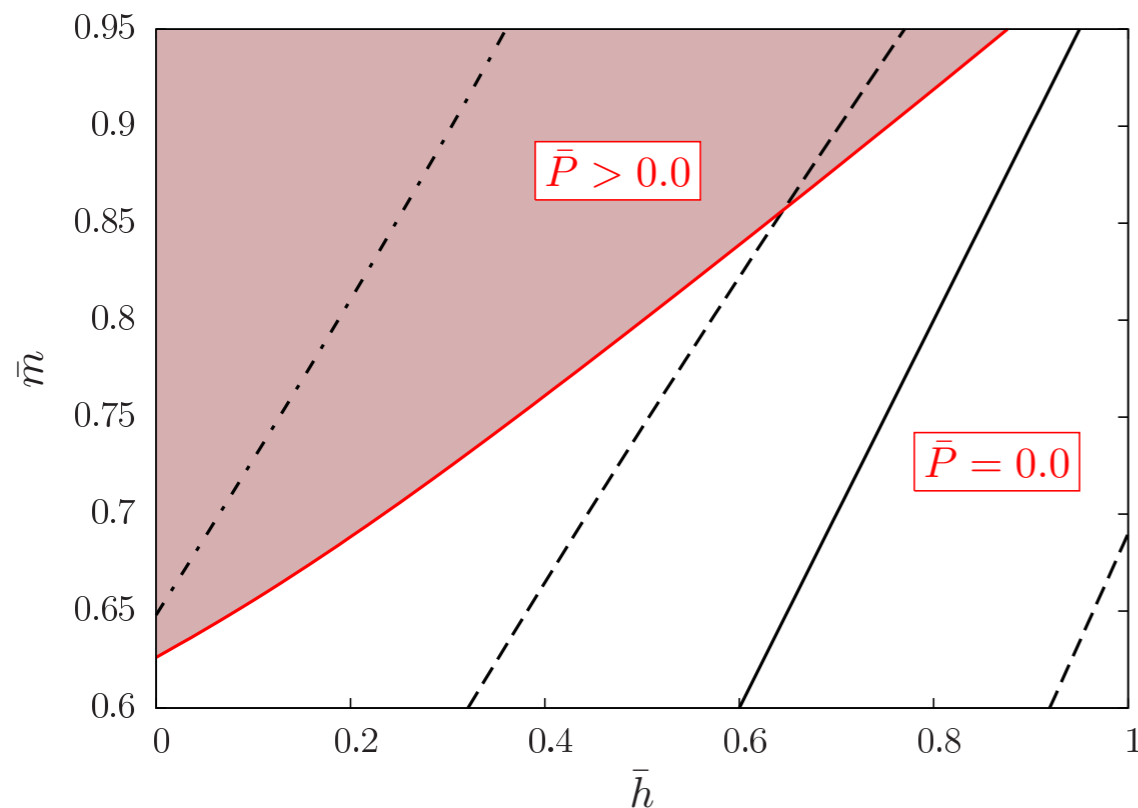
use **sign change of**  $A_{\varphi,k}^{(1)}$   
at LO in the  
derivative expansion  
as **indicator for** the  
onset of

**“crystalline” condensation**



# Phase diagram of imbalanced Fermi gases II

[Roscher, JB, Drut '15]



- **bound-state formation** in the presence of inert Fermi spheres [left panel]
- necessary condition for macroscopic occupation of the ground state ( $\sim$ condensation): bound-state formation
- **many-body phase diagram** [right panel]

# Inclusion of order-parameter fluctuations

[Roscher, JB, Drut '15]

- ansatz for the effective action:

$$\Gamma_k[\{\psi_\sigma\}, \varphi] = \int_{\tau, \vec{x}} \left[ \sum_{\sigma=\uparrow, \downarrow} \psi_\sigma^* \left( \partial_\tau - \frac{\nabla^2}{2m_\sigma} - \mu_\sigma \right) \psi_\sigma \right. \\ \left. + \varphi^* \left( \frac{Z_{\varphi, k}}{A_{\varphi, k}} \partial_\tau - \frac{1 - \bar{m}^2}{2} \nabla^2 \right) \varphi + U_k(\rho) - h_\varphi \left( \varphi^* \psi_\uparrow \psi_\downarrow - \varphi \psi_\uparrow^* \psi_\downarrow^* \right) \right]$$

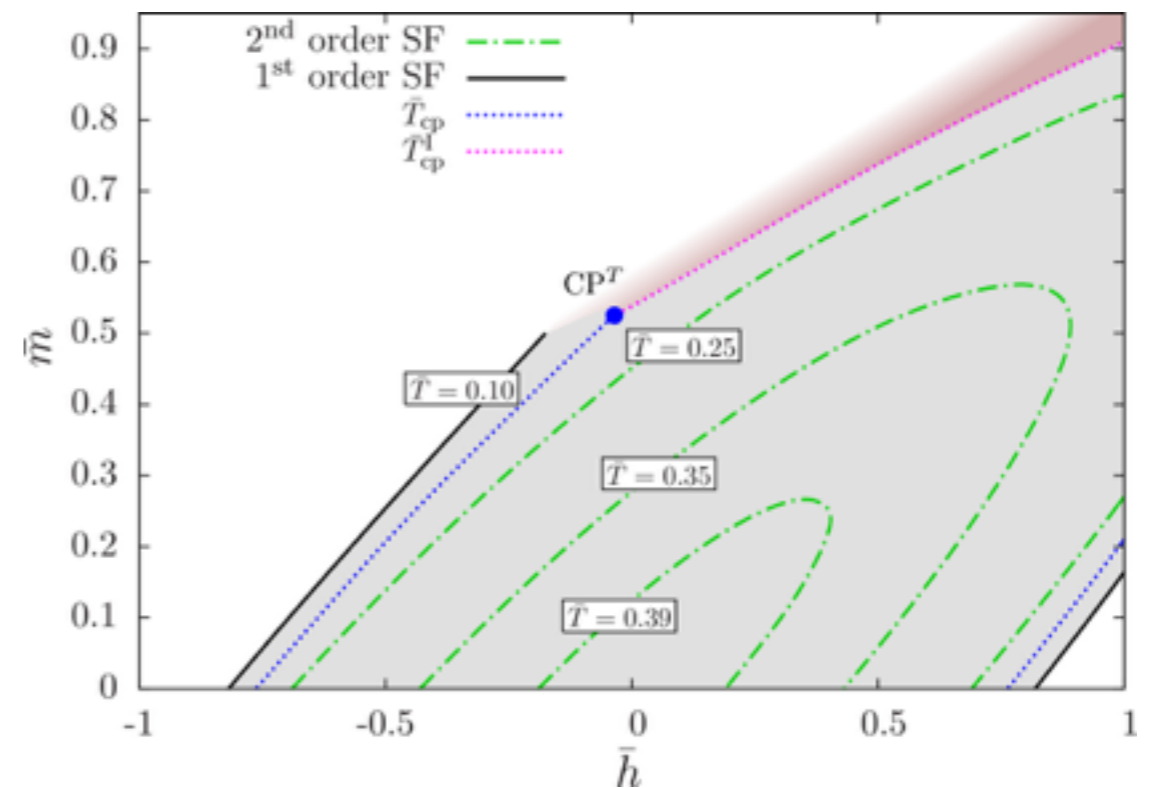
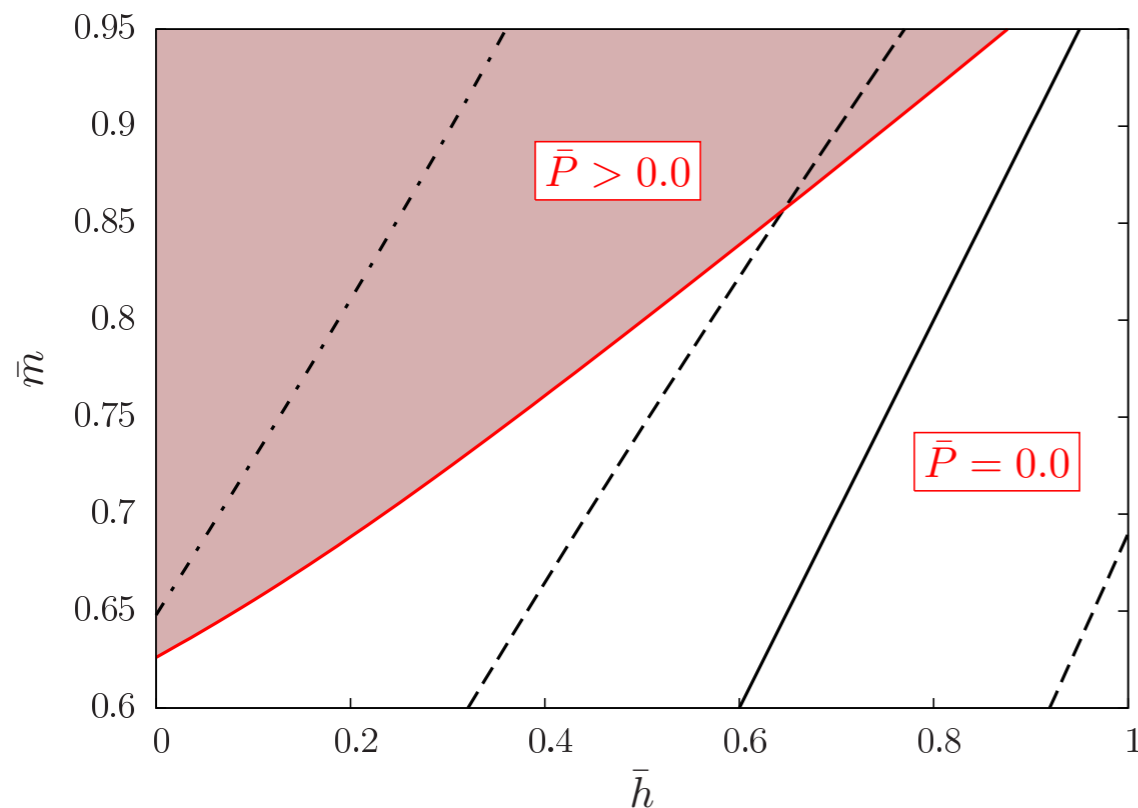
- use **spatially constant** background/mean field
- inverse propagator of the order-parameter field  $\varphi$ :

$$\bar{P}_\varphi(q_0, \vec{q}^2) = \left( i Z_{\varphi, k} q_0 + \frac{1 - \bar{m}^2}{2} A_{\varphi, k} \vec{q}^2 + \frac{\partial^2 U}{\partial \bar{\varphi} \partial \bar{\varphi}^*} \Big|_{\bar{\varphi}_0} \right)$$

- discretization of the effective potential  $U_k$  on a grid

# Phase diagram of imbalanced Fermi gases III

[Roscher, JB, Drut '15]



- **bound-state formation** in the presence of inert Fermi spheres [left panel]
- necessary condition for macroscopic occupation of the ground state ( $\sim$ condensation): bound-state formation
- **many-body phase diagram** [right panel]



# How to study imbalanced Fermi gases?

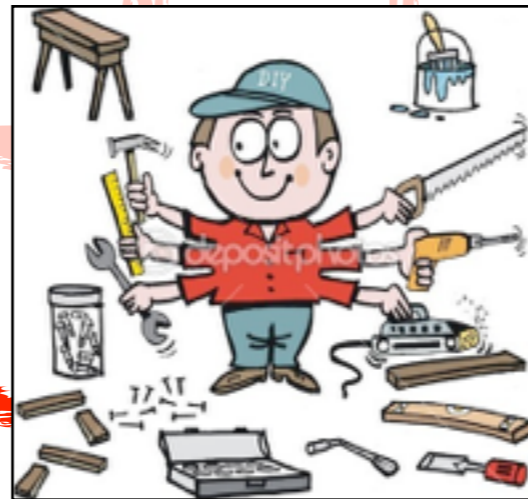
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[Wetterich '92]



(lattice) Monte Carlo

$$Z \sim \sum e^{-S}$$
A diagram of a 3D lattice structure, likely representing a crystal or a lattice of particles. The lattice is composed of yellow spheres connected by lines, forming a cubic grid. The equation  $Z \sim \sum e^{-S}$  is overlaid on the lattice.

Density Functional Theory

$$\psi_{\text{gs}}(\vec{x}_1, \dots, \vec{x}_N) \Leftrightarrow n_{\text{gs}}(\vec{x})$$

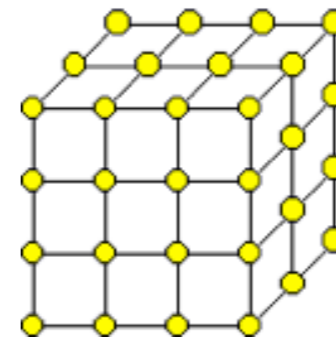
[Hohenberg & Kohn '65]

# Fermi gases and lattice MC simulations

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- Monte Carlo (MC) simulations:

partition sum  $Z \sim \sum_{\text{conf.}} e^{-S}$



(discretization  
of space-time)



- physical observables:

$$\langle \mathcal{O} \rangle \sim \frac{1}{Z} \sum_{\text{conf.}} \mathcal{O} e^{-S} \equiv \frac{1}{Z} \sum_i \mathcal{O}_i \rho_i$$

( $\rho_i$  : probability measure)

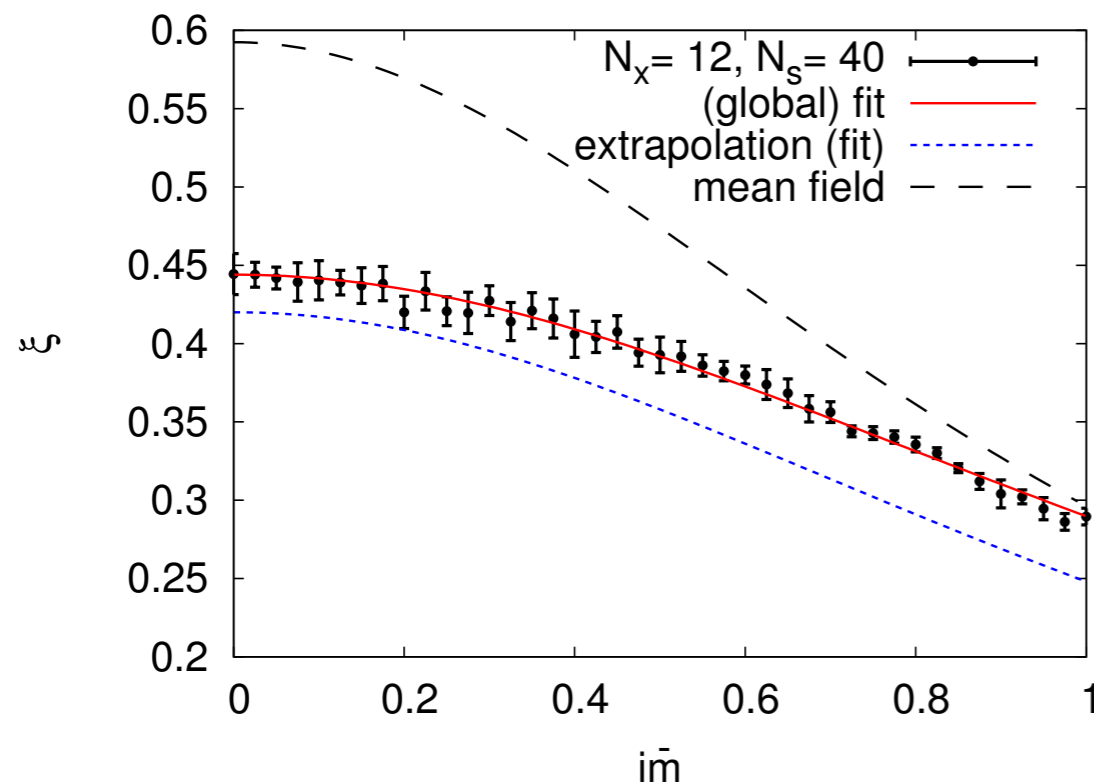
# Imbalanced Fermi gases and lattice MC

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- sum/integral can be evaluated straightforwardly with Monte Carlo techniques for  $\rho_i > 0$
- “**sign**” problem: the probability measure  $\rho_i$  is negative for “some” choices for the physical parameters
- strategy for **ultracold gases**:
  - ▶ Taylor expansion [JB, Drut, work in progress]
  - ▶ probability measure is rendered positive semi-definite by using imaginary-valued polarizations and imaginary mass-imbalance [JB et al. '13; Roscher, JB, Chen, Drut '14]

# Equation of state (EOS) of ultracold Fermi gases

[JB, Drut, Roscher '14]



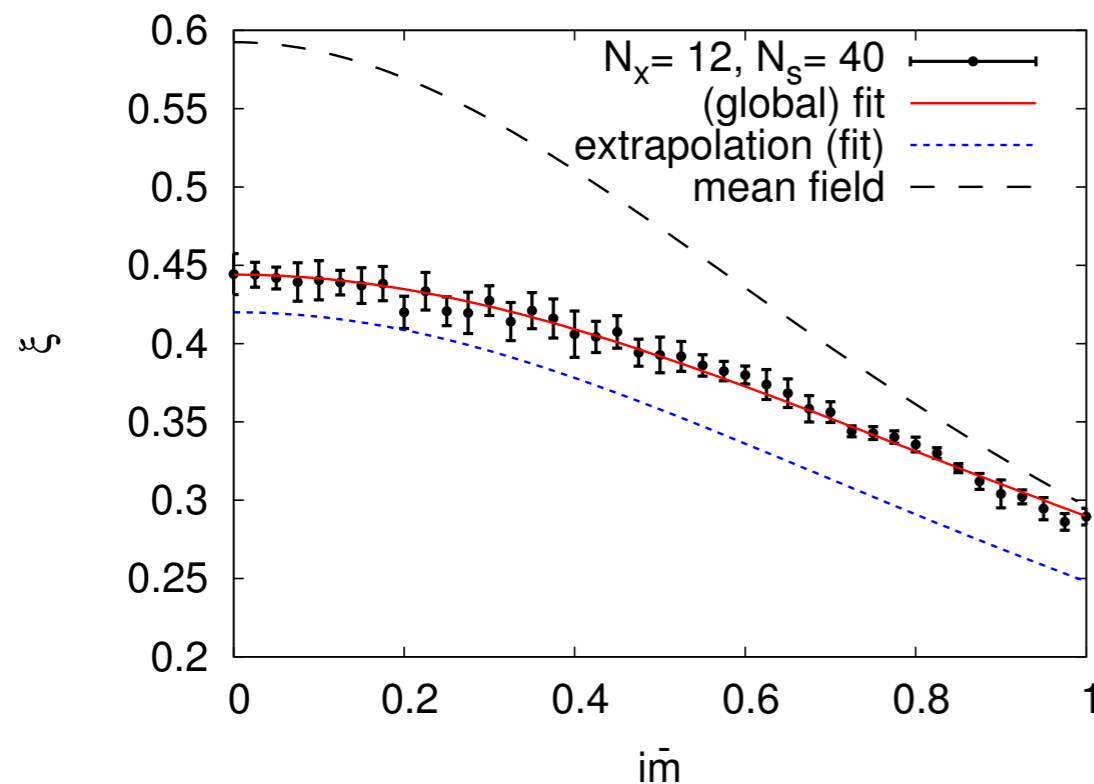
- *Bertsch* parameter & EOS (spin-balanced):

$$E_N(\bar{m}) = \xi E_N^{\text{free}}(\bar{m} = 0)$$

- Monte Carlo simulations for imaginary mass imbalances

# Equation of state (EOS) of ultracold Fermi gases

[JB, Drut, Roscher '14]



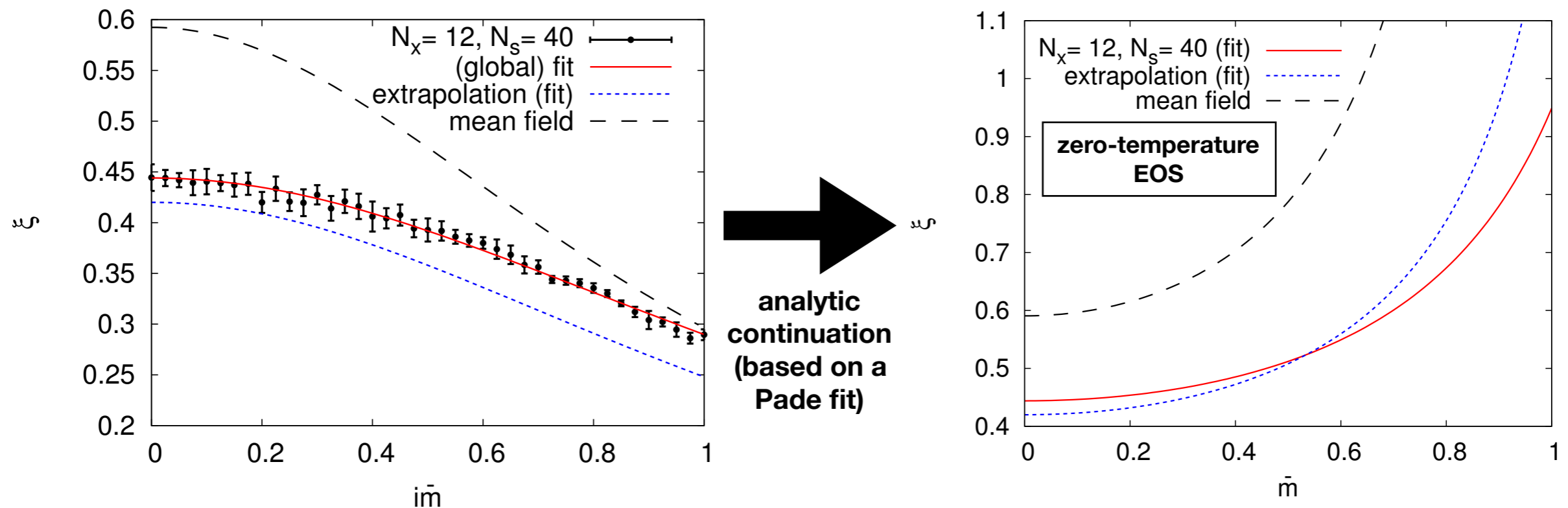
- Hellmann-Feynman theorem:

$$\left. \frac{d\xi(\bar{m})}{d\bar{m}^2} \right|_{\bar{m}=0} = \xi(\bar{m} = 0)$$

- mean-field theory:  $\xi(\bar{m}) = \frac{\xi(\bar{m} = 0)}{1 + (i\bar{m})^2}$

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[JB, Drut, Roscher '14]



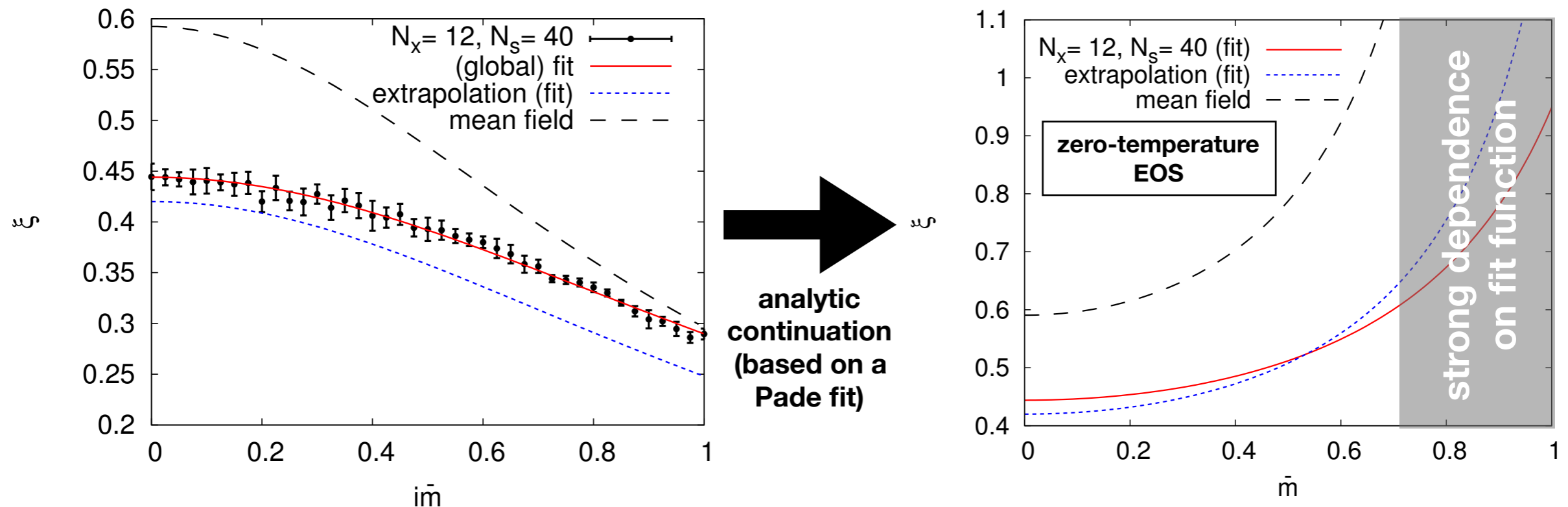
- *Bertsch* parameter & EOS (spin-balanced):

$$E_N(\bar{m}) = \xi E_N^{\text{free}}(\bar{m} = 0)$$

- Monte Carlo simulations for imaginary mass imbalances yield physical equation of state via analytic continuation

# Equation of state (EOS) of ultracold Fermi gases

[JB, Drut, Roscher '14]



- *Bertsch* parameter & EOS (spin-balanced):

$$E_N(\bar{m}) = \xi E_N^{\text{free}}(\bar{m} = 0)$$

- **Phase transition**/emergence of “crystalline” phase at large imbalances? More involved studies are required ...

# How to study imbalanced Fermi gases?

Schrödinger equation

$$\hat{H}|\psi\rangle = E|\psi\rangle$$

Wetterich equation

$$\frac{\partial_t R_k}{(2)} + R_k$$

**Comparison with experimental data?**

[picture taken from Hitchcock's 'Psycho' (1960)]

(lattice) Imc


$$Z \sim \sum e^{-S}$$

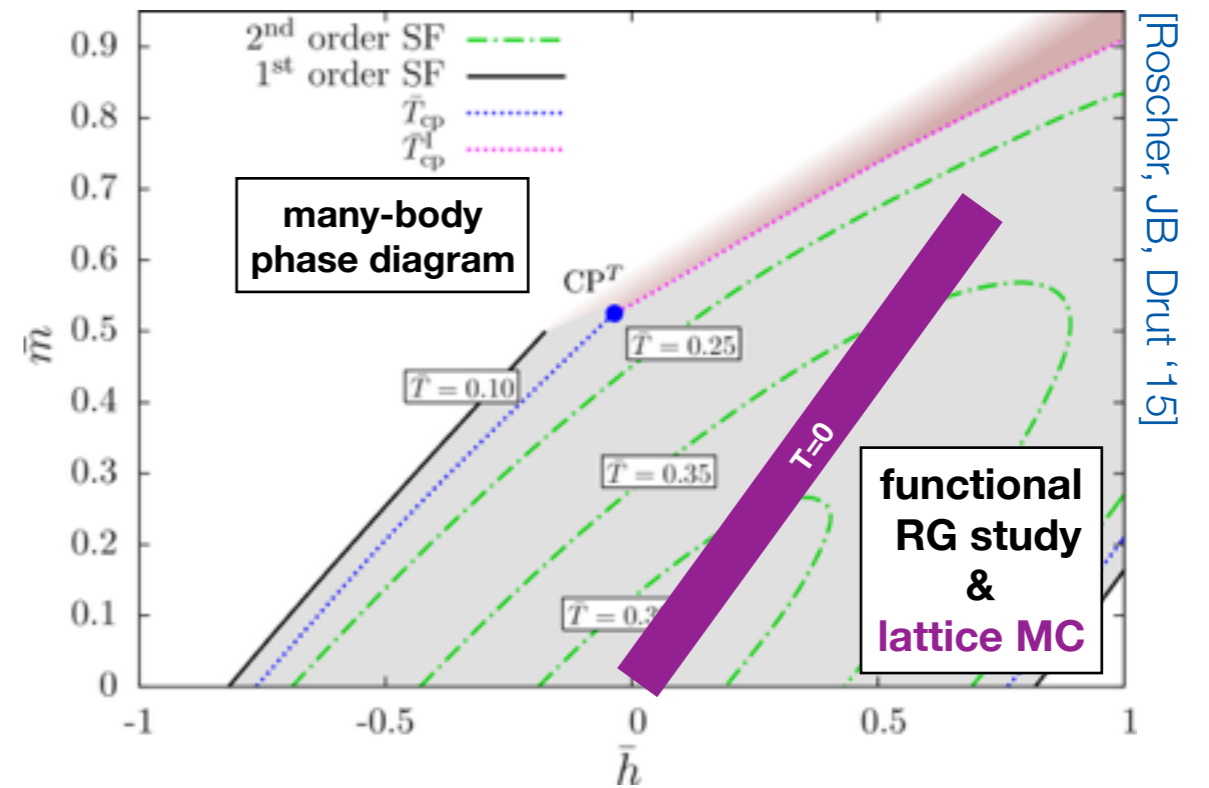
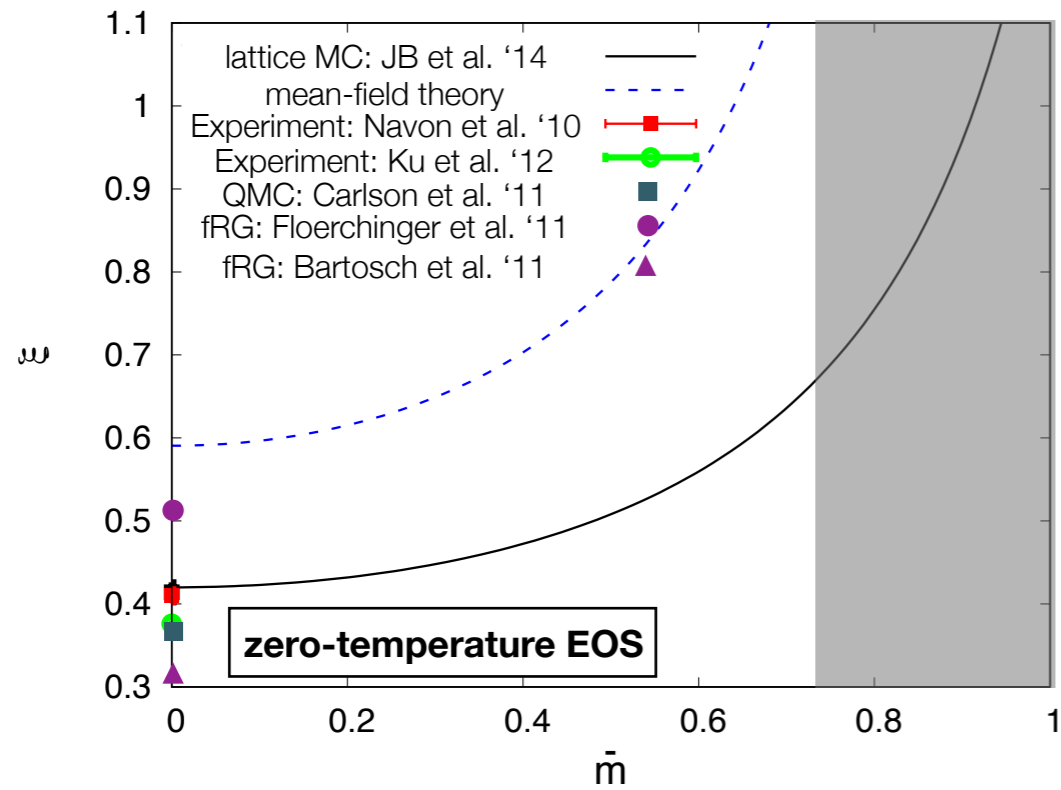
Functional Theory

$$\psi_{\text{gs}}(\vec{x}_1, \dots, \vec{x}_N) \Leftrightarrow n_{\text{gs}}(\vec{x})$$

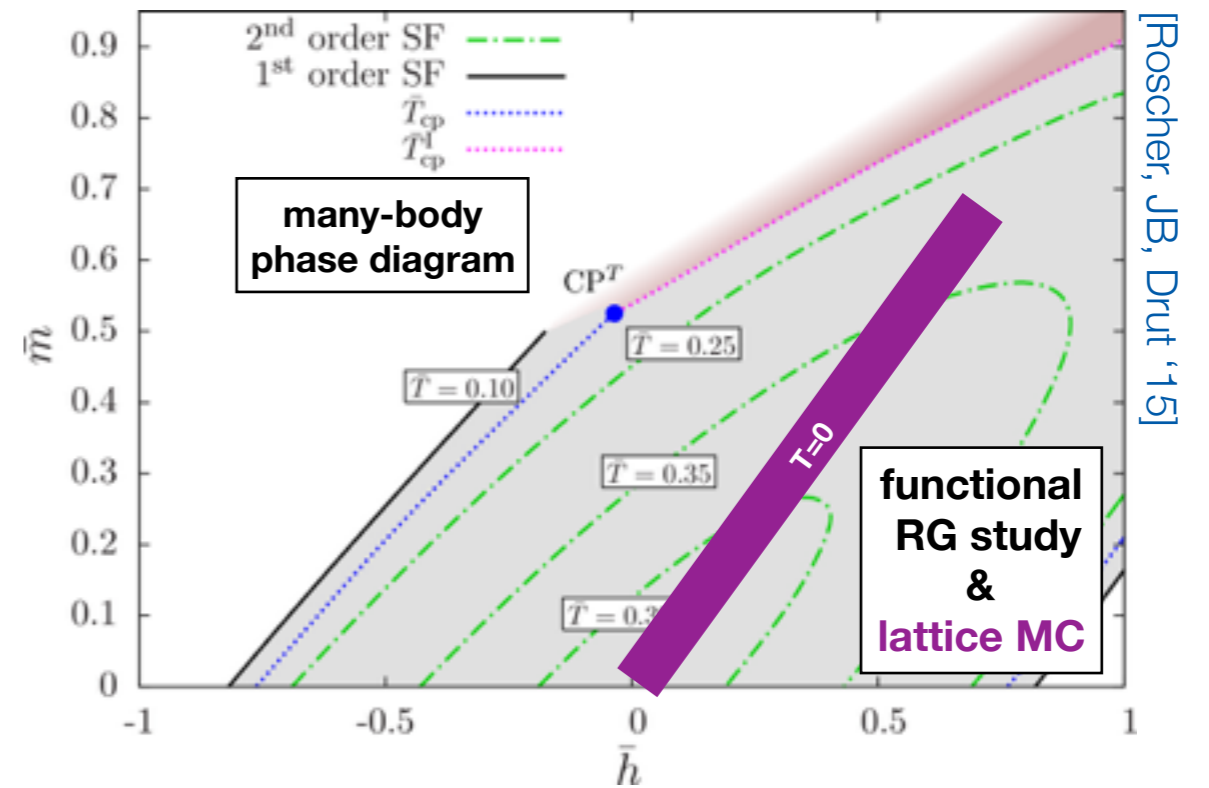
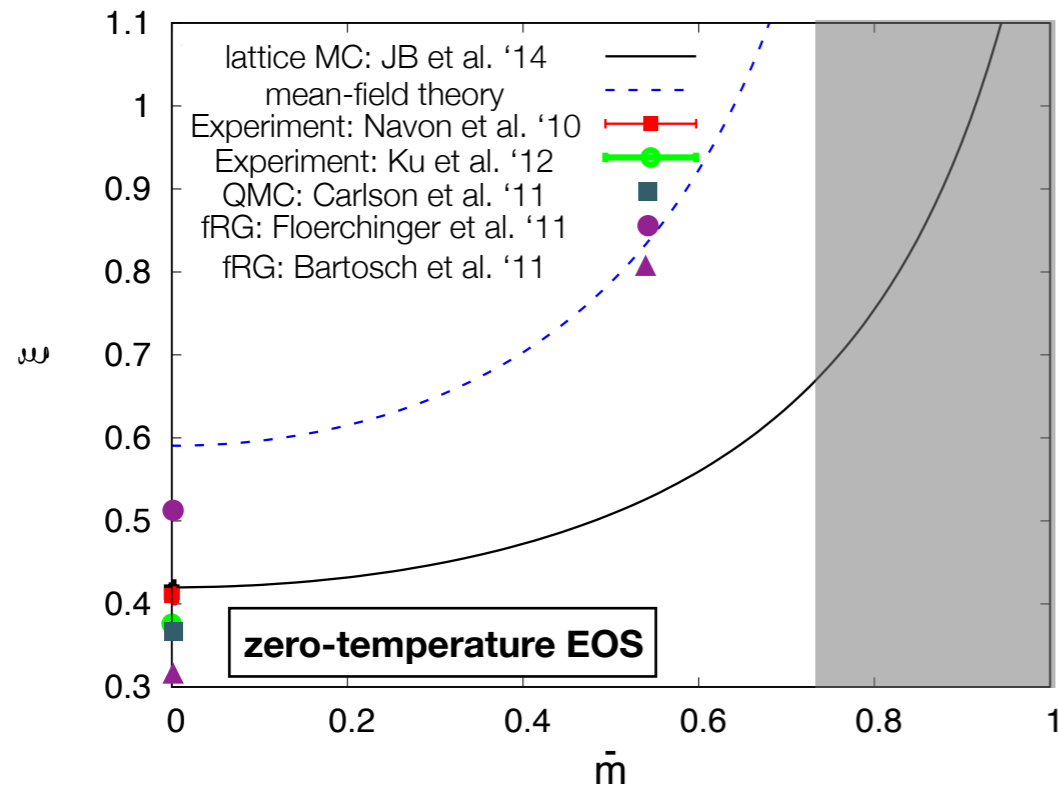
[Hohenberg & Kohn '65]



# Theory “versus” Experiment



# Theory “versus” Experiment



## ● Fermi polaron (“N+1-body problem”)

[MC: Carlson & Reddy '05; Lobo et al. '06; Prokof'ev & Svistunov '08; Pilati & Giorgini '08;

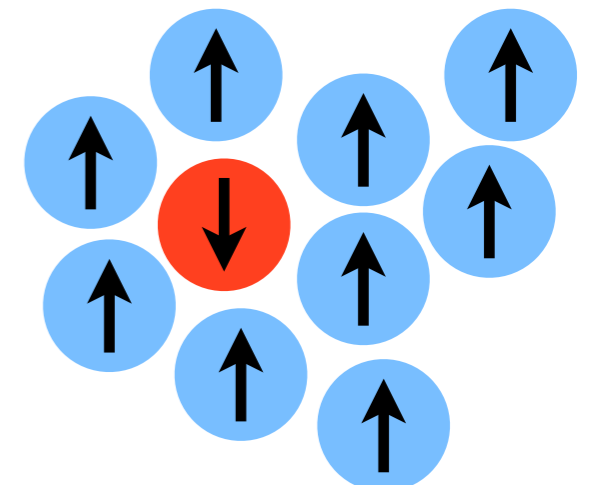
Variational: Chevy '07; Bulgac & Forbes '07; Ku, Braun & Schwenk '09;

functional RG: Enss & Schmidt '11]

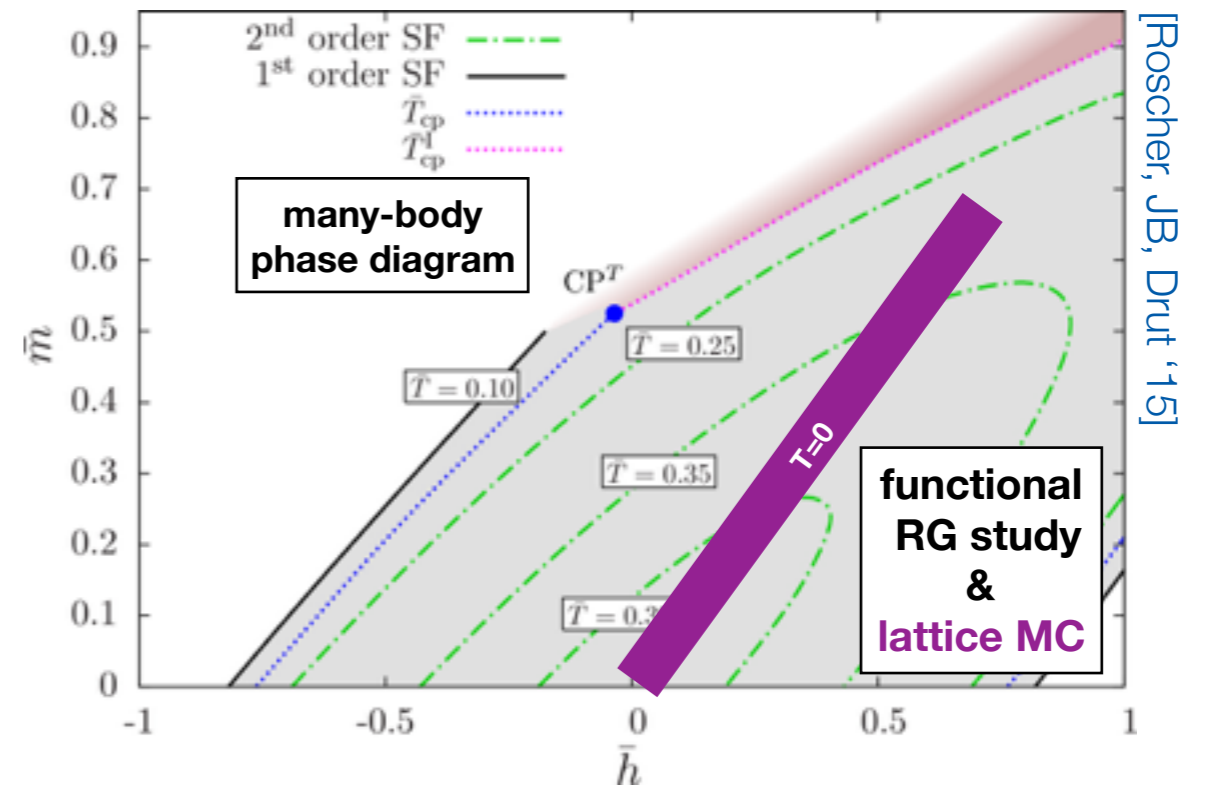
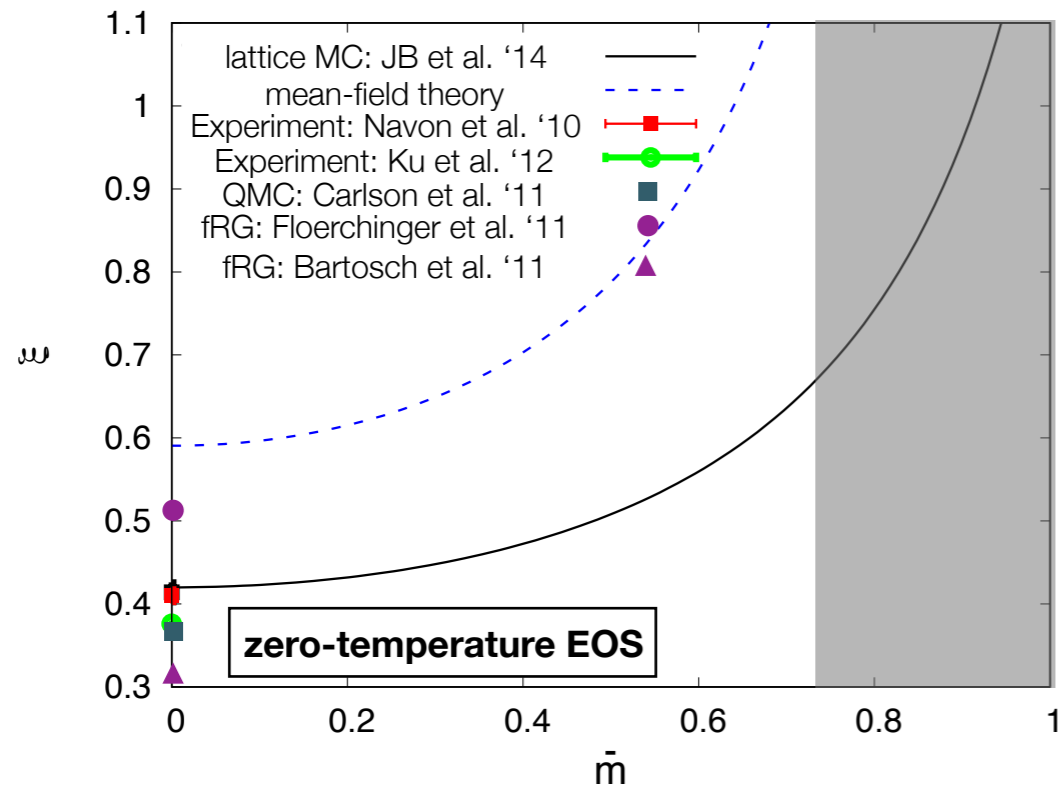
[not discussed  
 in this talk]

$$\eta = \mu_{\downarrow} / \epsilon_F \approx -0.615 \dots -0.57$$

[Note: consistent with experimental constraints, Chevy '07]



# Theory “versus” Experiment

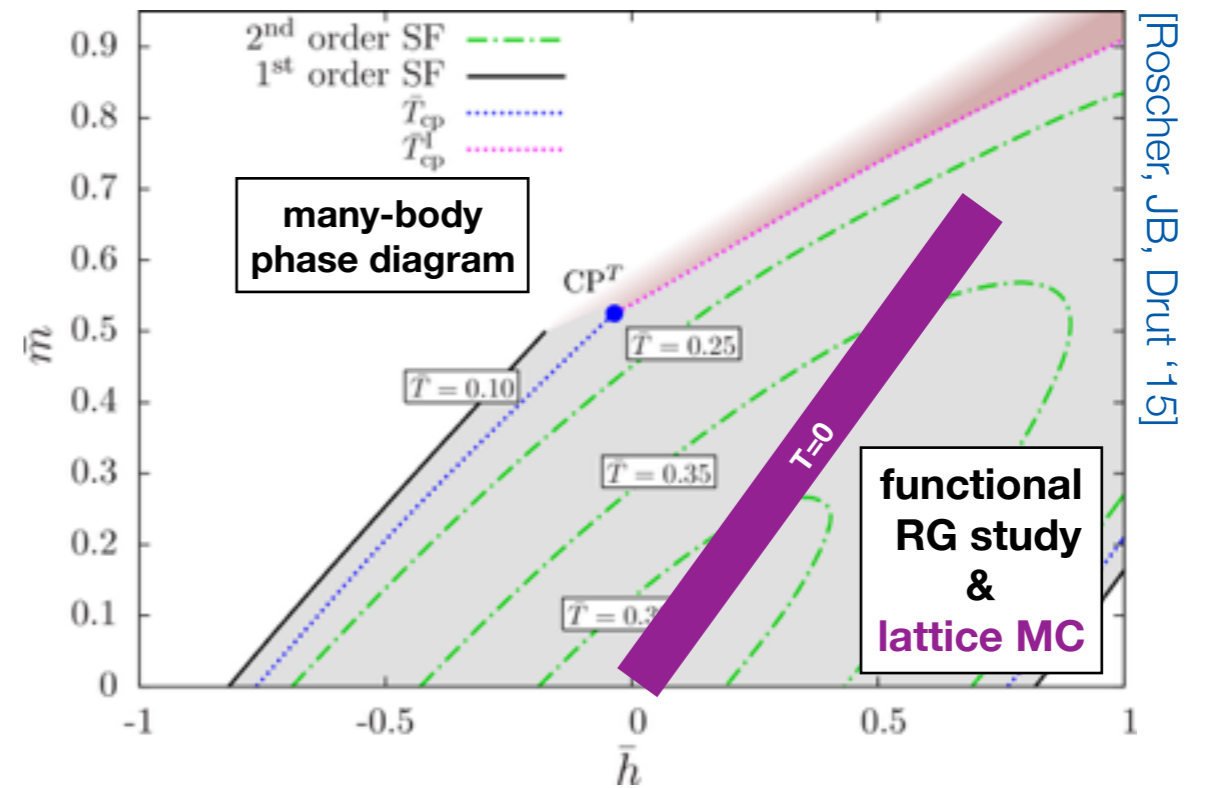
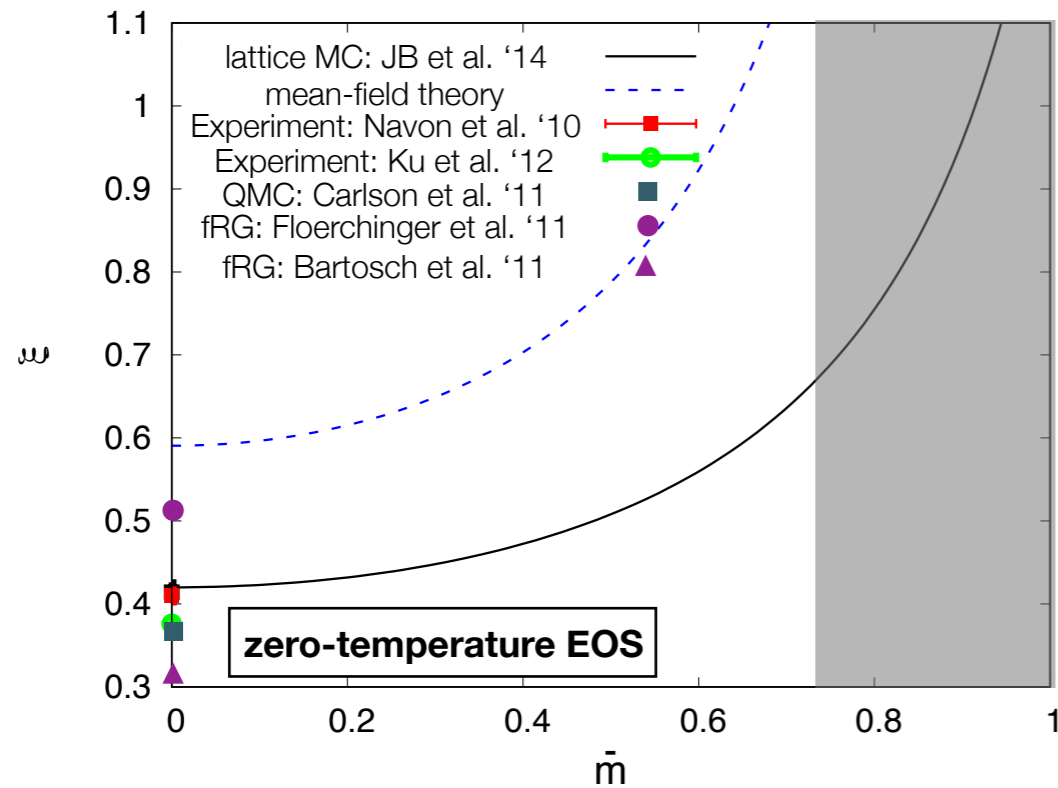


- phase transition temperature  $T_c/\mu$  at zero imbalances:

functional RG	0.40
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[Boettcher, JB, Herbst, Roscher, Pawłowski, Wetterich '14]  
 [Roscher, JB, Drut '15]

# Theory “versus” Experiment



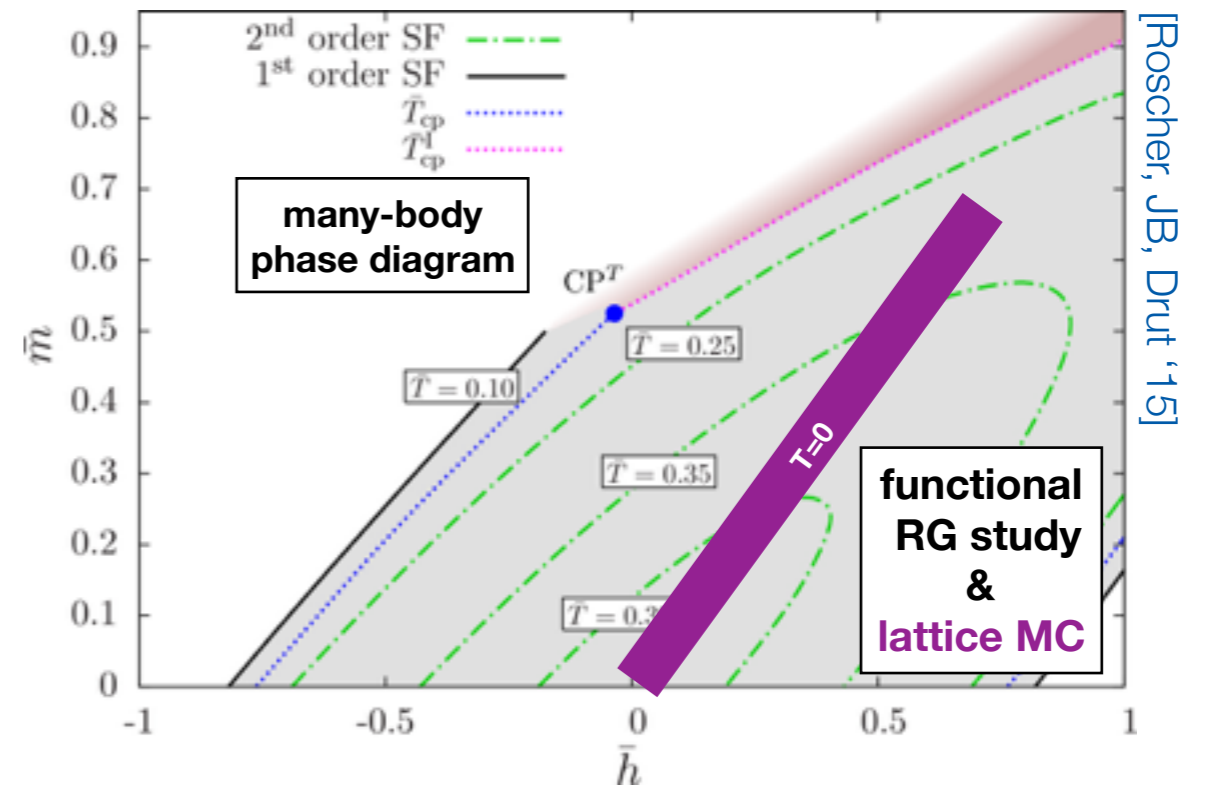
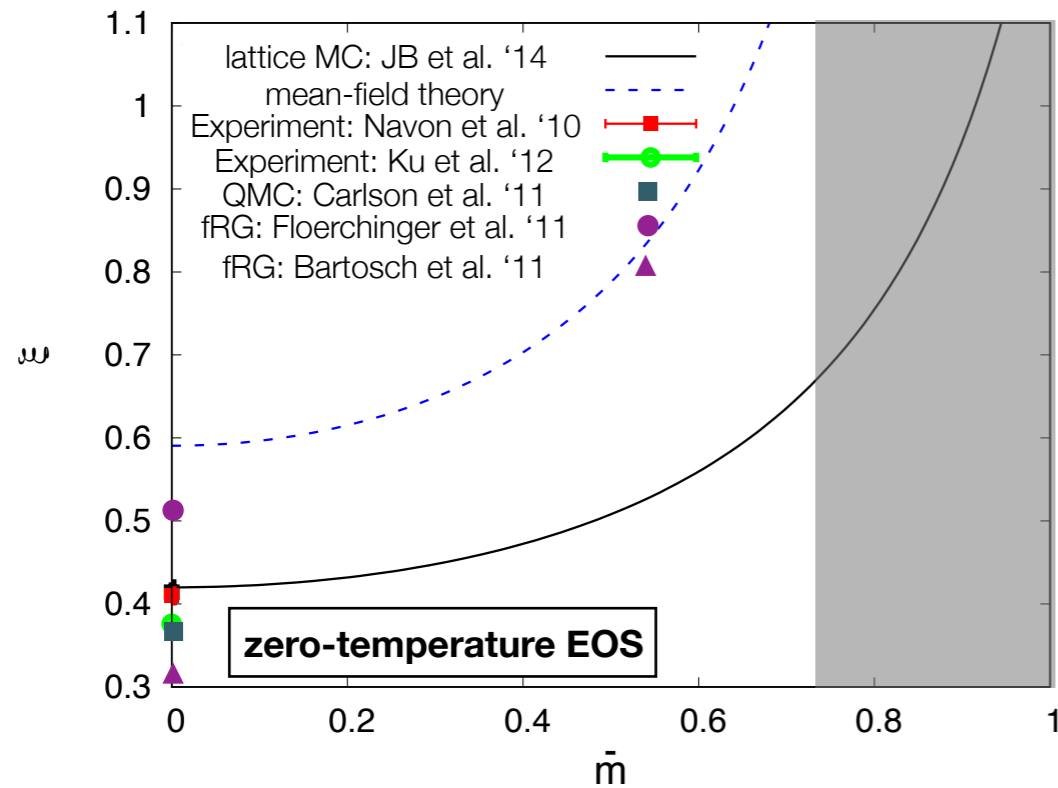
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QMC/lattice MC	0.35

[Boettcher, JB, Herbst,  
 Roscher, Pawłowski, Wetterich '14]  
 [Roscher, JB, Drut '15]

[Bulgac, Drut, Magierski '08]

# Theory “versus” Experiment



- phase transition temperature  $T_c/\mu$  at zero imbalances:

functional RG	0.40
QMC/lattice MC	0.35
DDMC	0.399

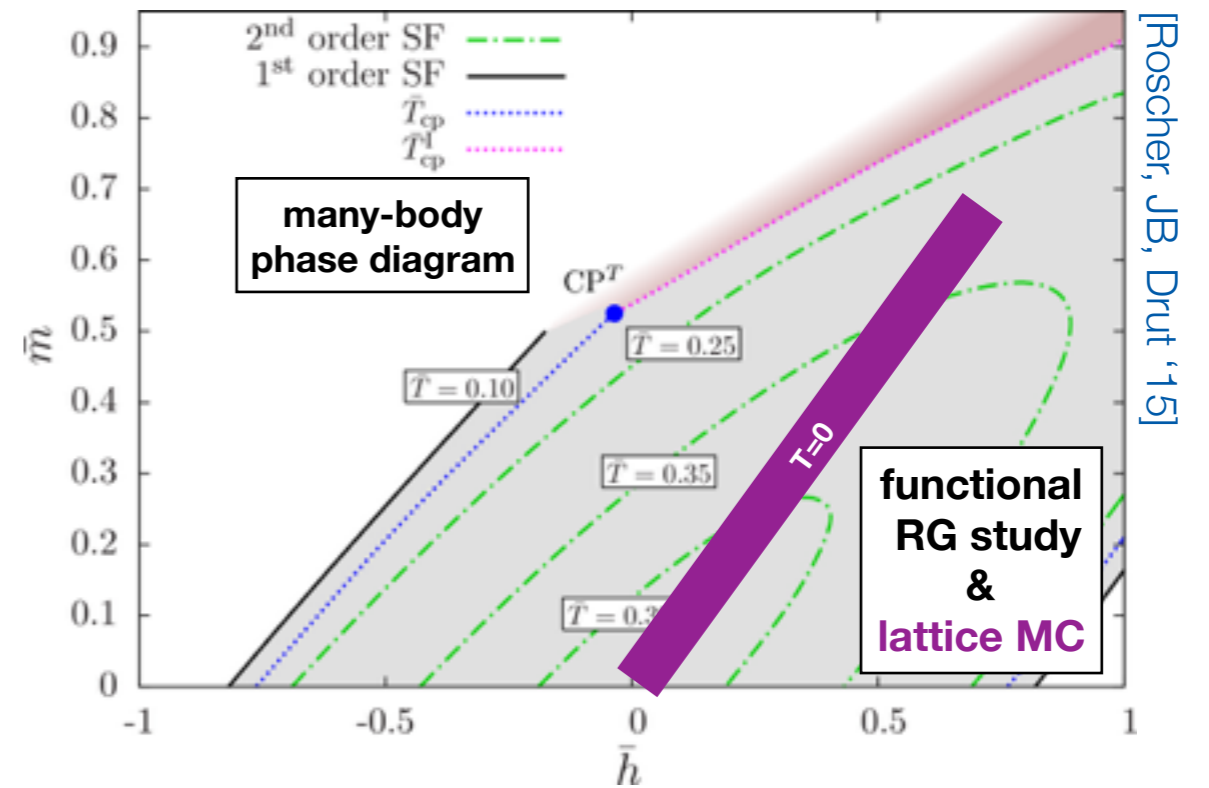
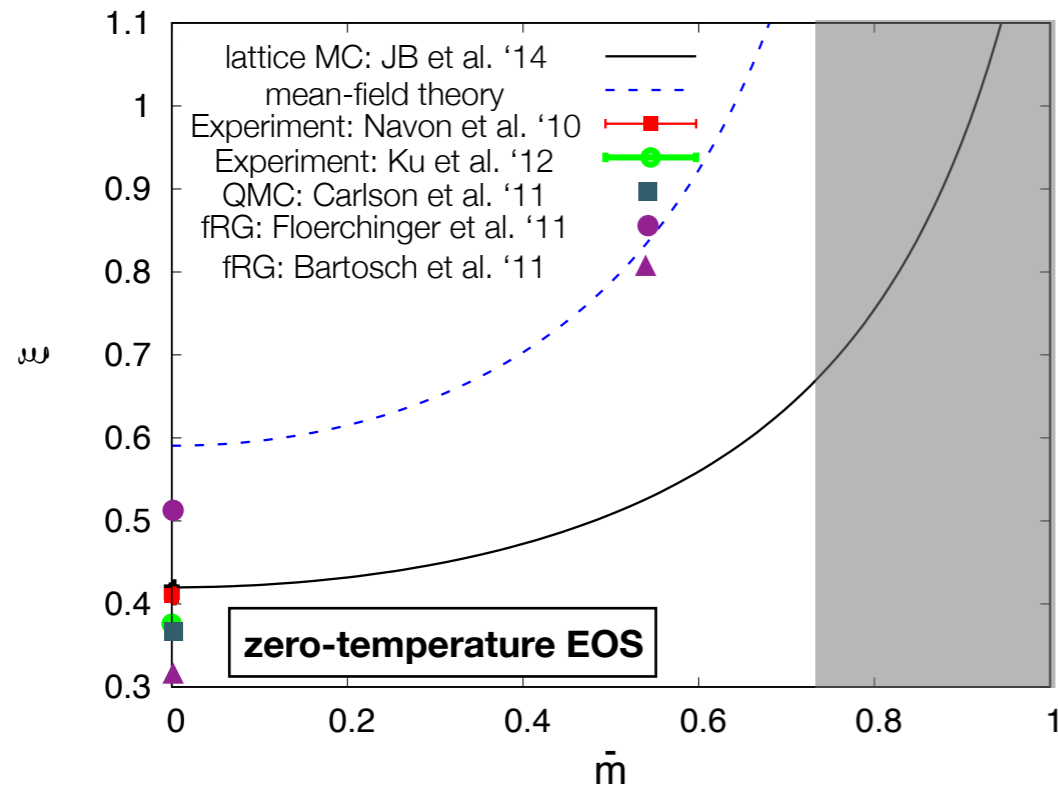
[Boettcher, JB, Herbst, Roscher, Pawłowski, Wetterich '14]

[Roscher, JB, Drut '15]

[Bulgac, Drut, Magierski '08]

[Goulko, Wingate '10]

# Theory “versus” Experiment



- phase transition temperature  $T_c/\mu$  at zero imbalances:

functional RG	0.40
QMC/lattice MC	0.35
DDMC	0.399
Experiment	0.40

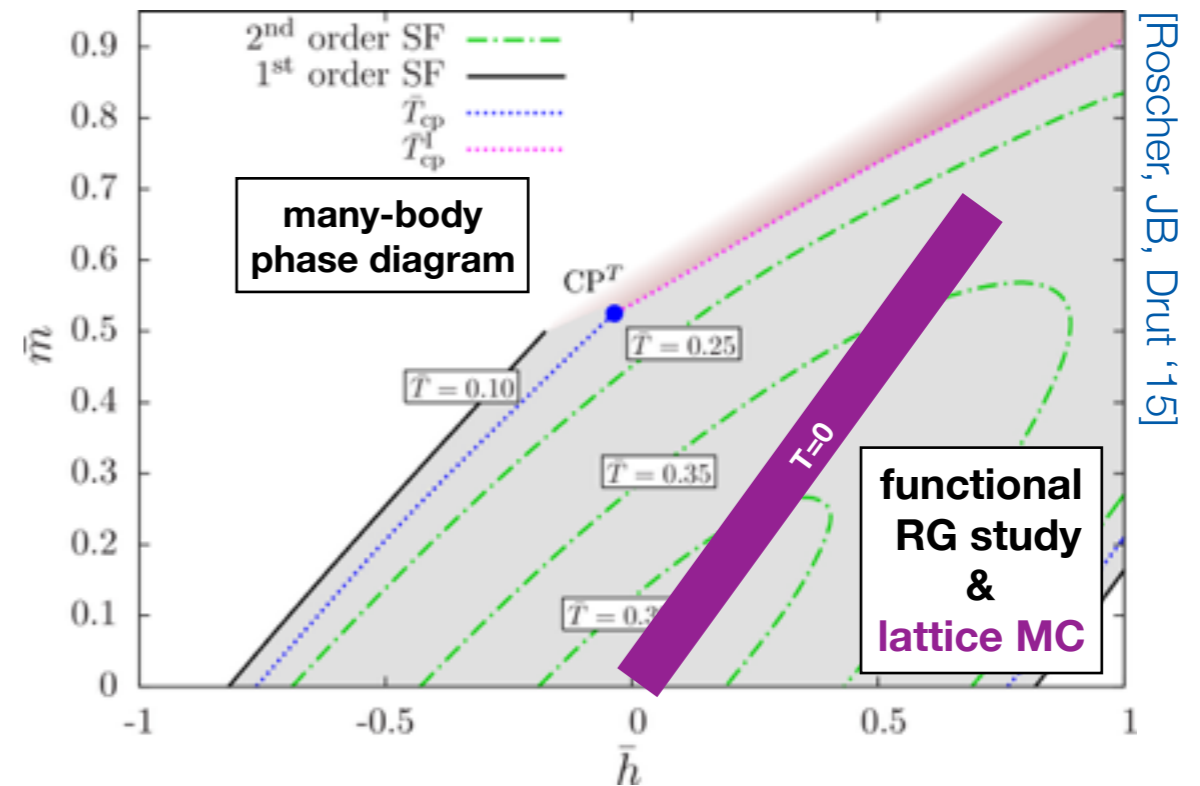
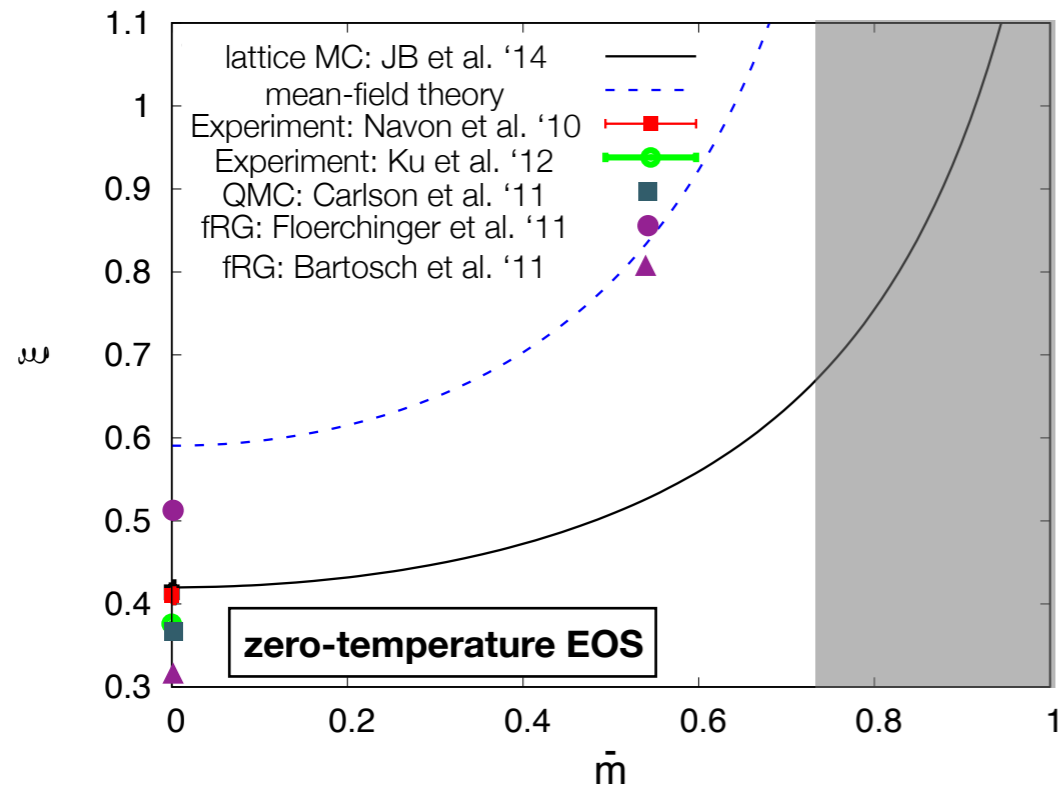
[Boettcher, JB, Herbst,  
 Roscher, Pawłowski, Wetterich '14]  
 [Roscher, JB, Drut '15]

[Bulgac, Drut, Magierski '08]

[Goulko, Wingate '10]

[Ku et al. '12]

# Theory “versus” Experiment



- very good agreement between theory and experiment with respect to the phase transition in the balanced limit
- **experiments** with different **mass imbalances** are now in reach (6Li, 40K, 161Dy, 163Dy and 167Er)

# Summary

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So

- overall: very good agreement between theoretical predictions and experiments for balanced unitary Fermi gases
- EOS of mass-imbalanced Fermi gases
  - testable prediction for future experiments
- strong indications for the existence of a crystalline phase at large mass imbalances
  - (very) exact coordinates?
  - what is the “structure” of the “crystal”?

(k

$\rho_k$

$\vec{x}$ )



# Summary

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So

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**STAY TUNED**