

Strongly interacting Fermi gas in two dimensions

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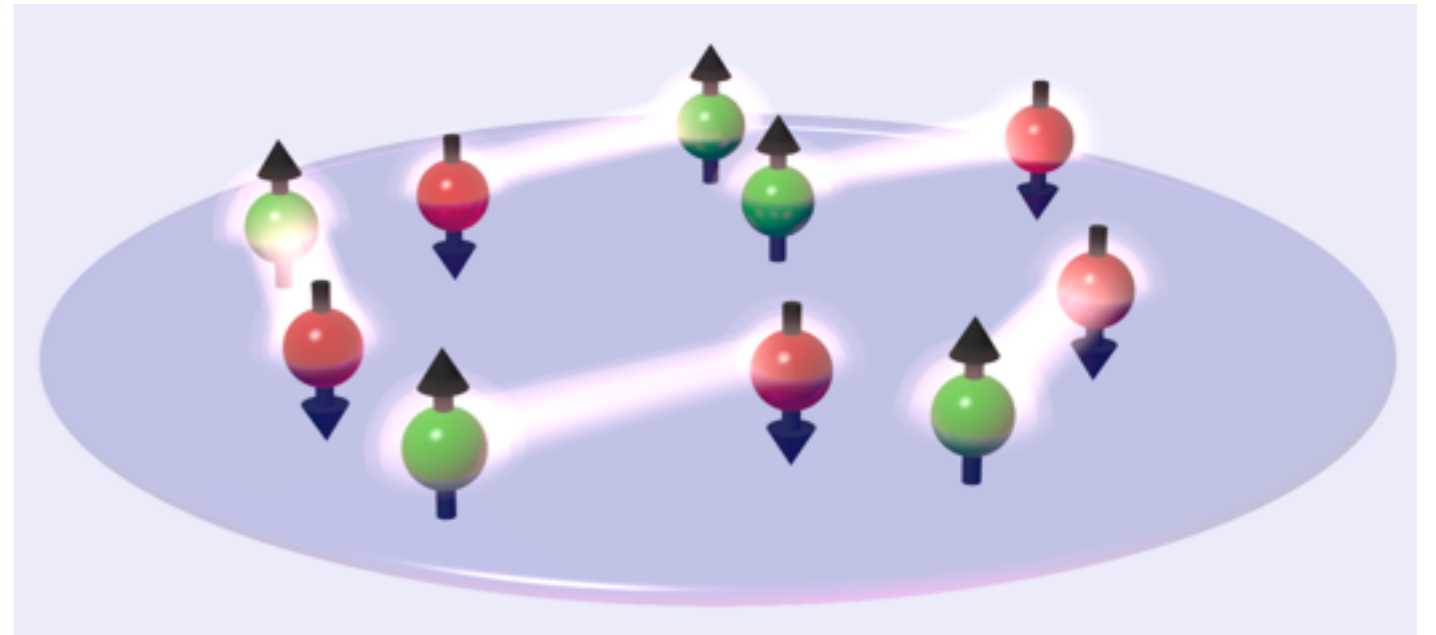
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Motivation

- understand properties of strongly interacting Fermi gas



- **thermodynamics:** equation of state (EoS), density $n(\mu, T, a)$, pressure $P(\mu, T, a)$
- **transport & dynamical properties...**
- dilute gas of nonrelativistic \uparrow and \downarrow fermions with contact interaction

$$\mathcal{H} = \int d\mathbf{x} \sum_{\sigma=\uparrow,\downarrow} \psi_{\sigma}^{\dagger} \left(-\frac{\hbar^2 \nabla^2}{2m} - \mu_{\sigma} \right) \psi_{\sigma} + g_0 \psi_{\uparrow}^{\dagger} \psi_{\downarrow}^{\dagger} \psi_{\downarrow} \psi_{\uparrow}$$

2D Fermi gas

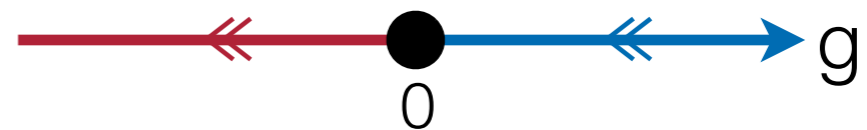
$$\mathcal{L} = \sum_{\sigma} \bar{\psi}_{\sigma} \left(\partial_{\tau} - \frac{\nabla^2}{2m} - \mu \right) \psi_{\sigma} + g_0 \bar{\psi}_{\uparrow} \bar{\psi}_{\downarrow} \psi_{\downarrow} \psi_{\uparrow}$$

- vacuum ($\mu=0$) classically scale invariant with $z=2$, $\dim[g_0]=0$
- exact beta function:

$$\frac{dg}{d\ell} = -\frac{g^2}{2}$$

attractive: strong binding

repulsive: asympt. free



- log. running coupling, energy scale dependence breaks scale invariance

Holstein 1993; Pitaevskii & Rosch 1997

exact 2D scattering amplitude:

always bound state!

$$\text{3D: } f(k) = \frac{1}{-1/a_{3D} - ik}$$



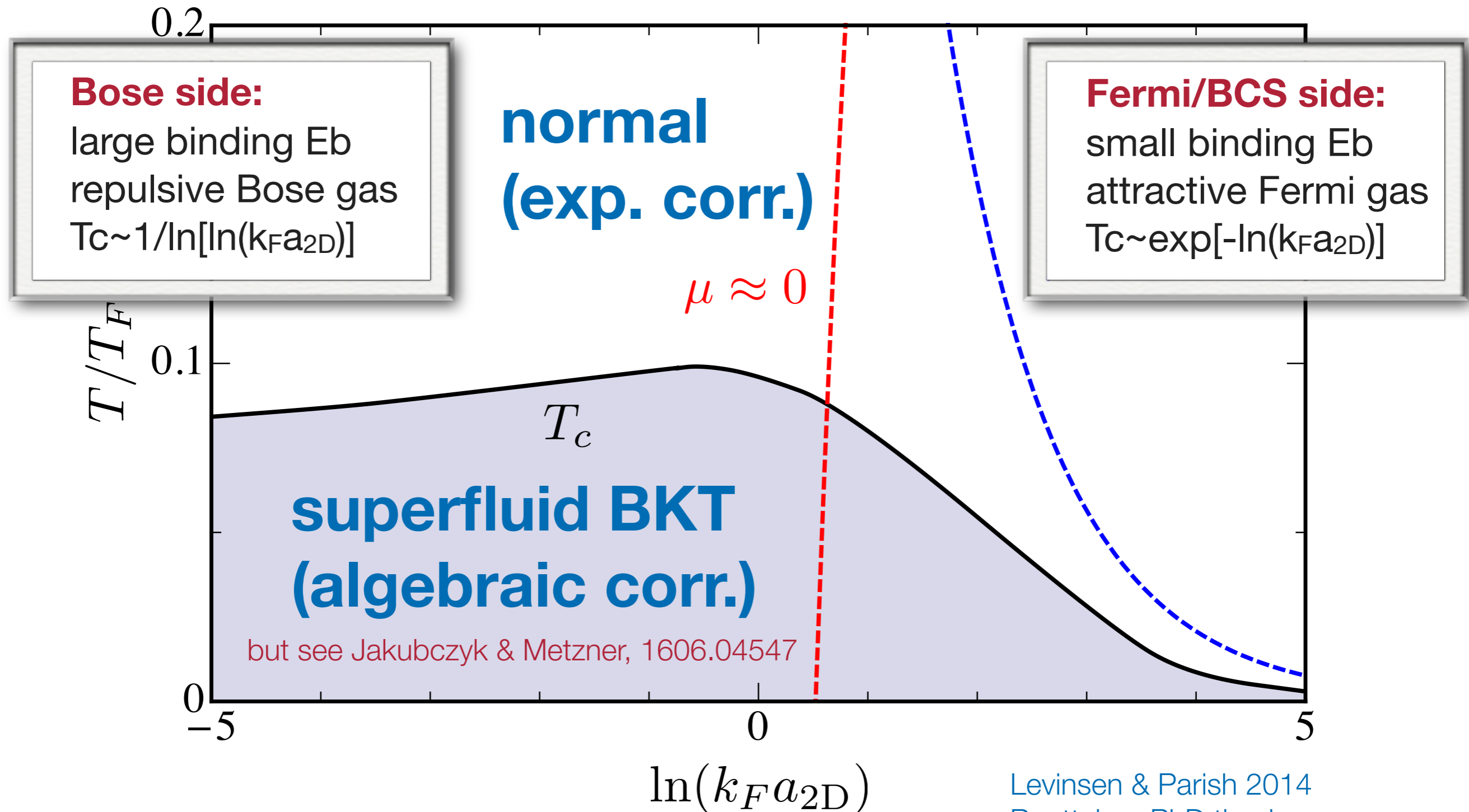
$$\text{2D: } f(k) = \frac{1}{\ln(1/k^2 a_{2D}^2) + i\pi}$$

$$\varepsilon_B = \frac{\hbar^2}{ma_{2D}^2}$$

- typical scale $k=k_F$: expansion parameter $g=-1/\ln(k_F a_{2D})$

Adhikari 1986

Phase diagram of 2D Fermi gas



Thermodynamics

- dilute gas ($k_F r_0 \ll 1$): **universal properties** depend only on T/T_F and $k_F a_{2D}$
- **Contact density:** probability to find up and down in same place [S. Tan 2008]

$$C = m^2 g_0^2 \langle \bar{\psi}_\uparrow \bar{\psi}_\downarrow \psi_\downarrow \psi_\uparrow(r) \rangle \qquad \frac{dE}{d \ln a_{2D}} = \frac{C}{2\pi m}$$

adiabatic theorem for internal energy E via Hellmann-Feynman

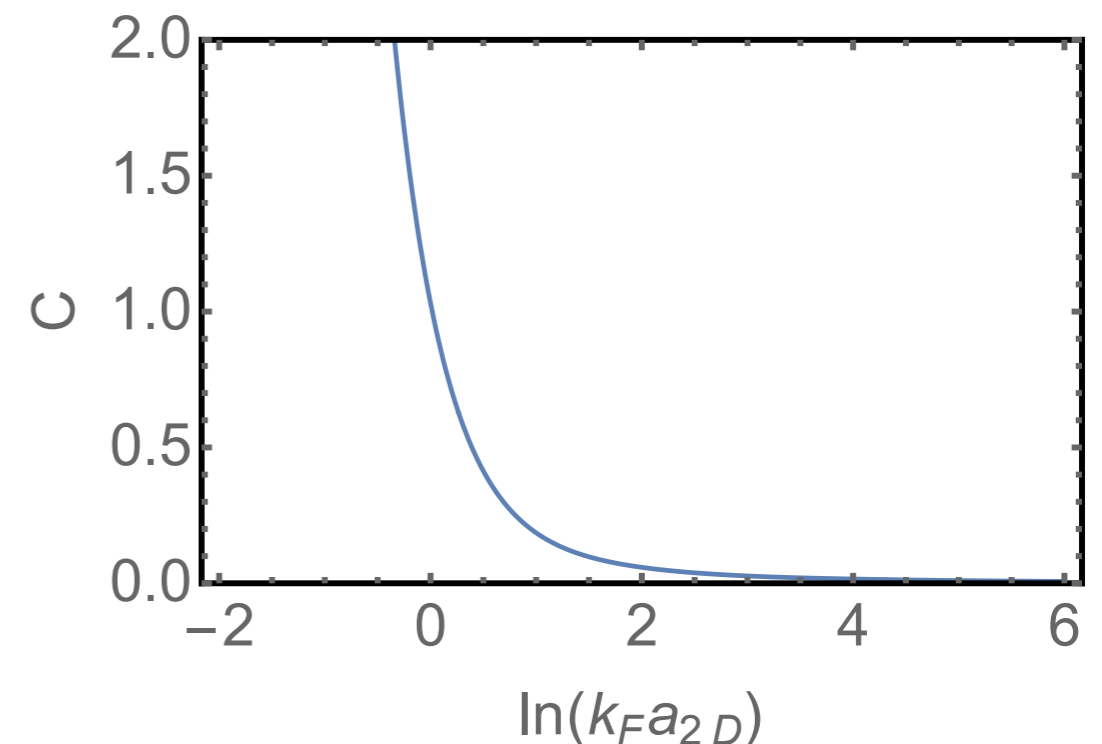
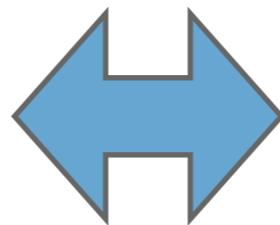
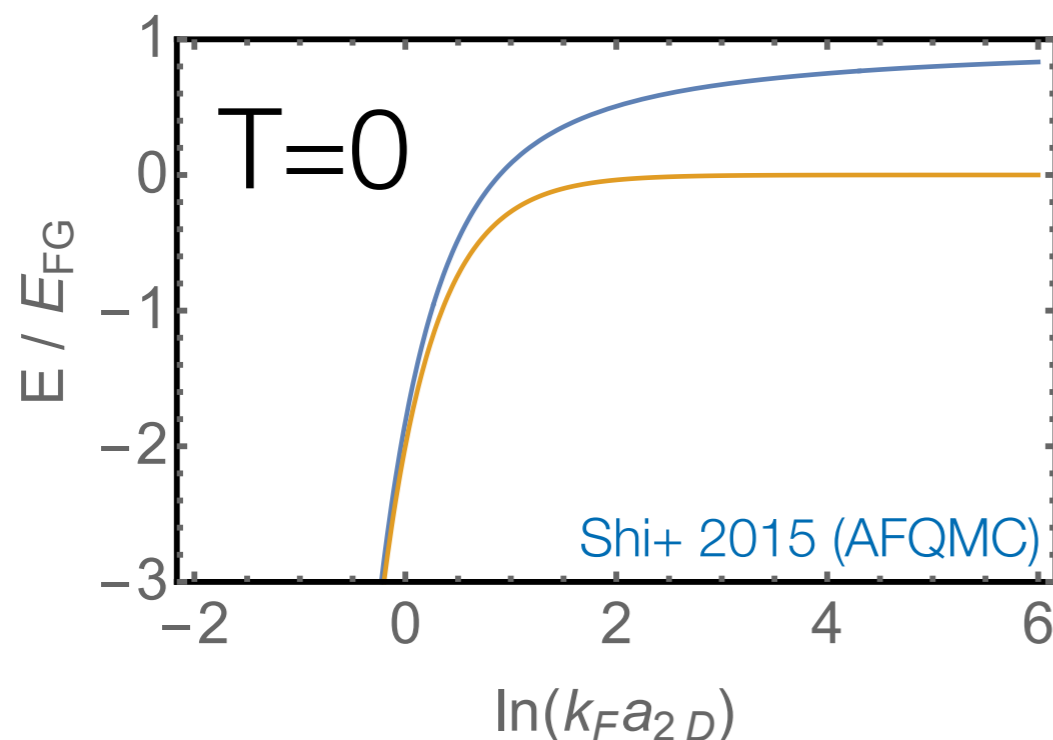
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Universal relations

- dilute gas: Contact determines **UV limit of correlation functions**

- momentum distribution

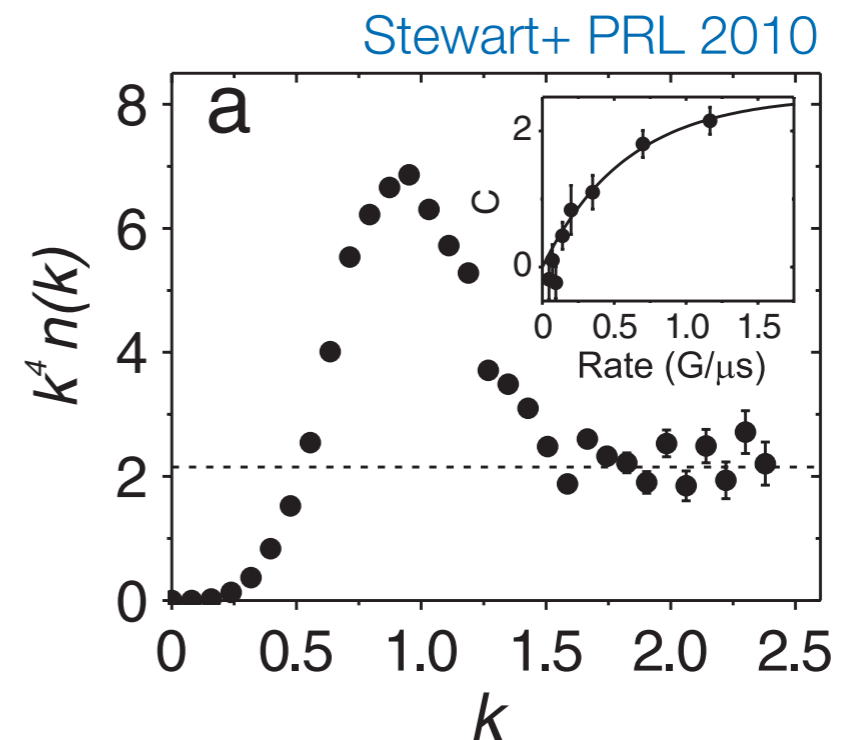
$$n(k) \xrightarrow{k \gg k_F} \frac{C}{k^4}$$

- internal energy

$$E = \sum_{\sigma} \int^{\Lambda} \frac{d^2 k}{(2\pi)^2} \frac{k^2}{2m} n_{\sigma}(k) - \ln(\Lambda a_{2D}) \frac{C}{2\pi m}$$

kinetic

interaction energy



Pressure

scale invariant:

$$P = E$$

interacting 2D Fermi gas:

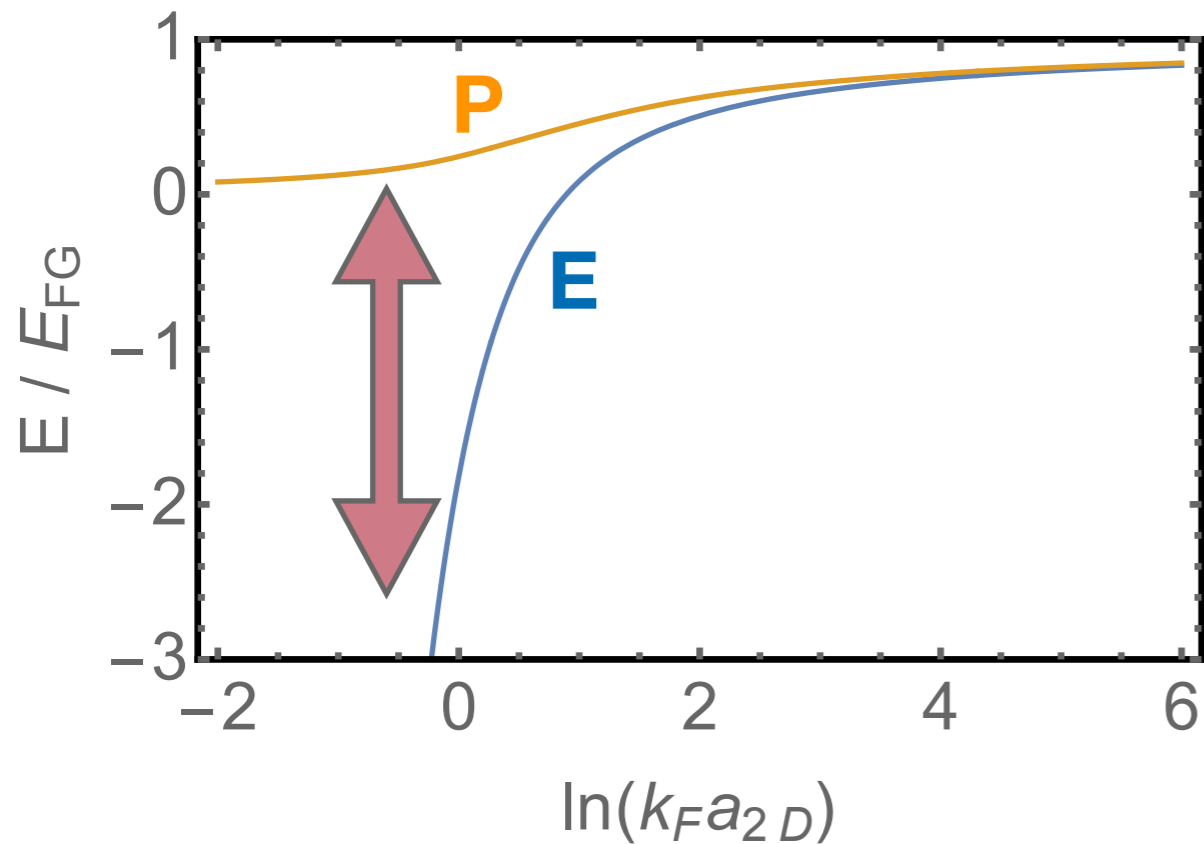
$$P = E + \frac{C}{4\pi m}$$

breaks scale invariance!

Pressure

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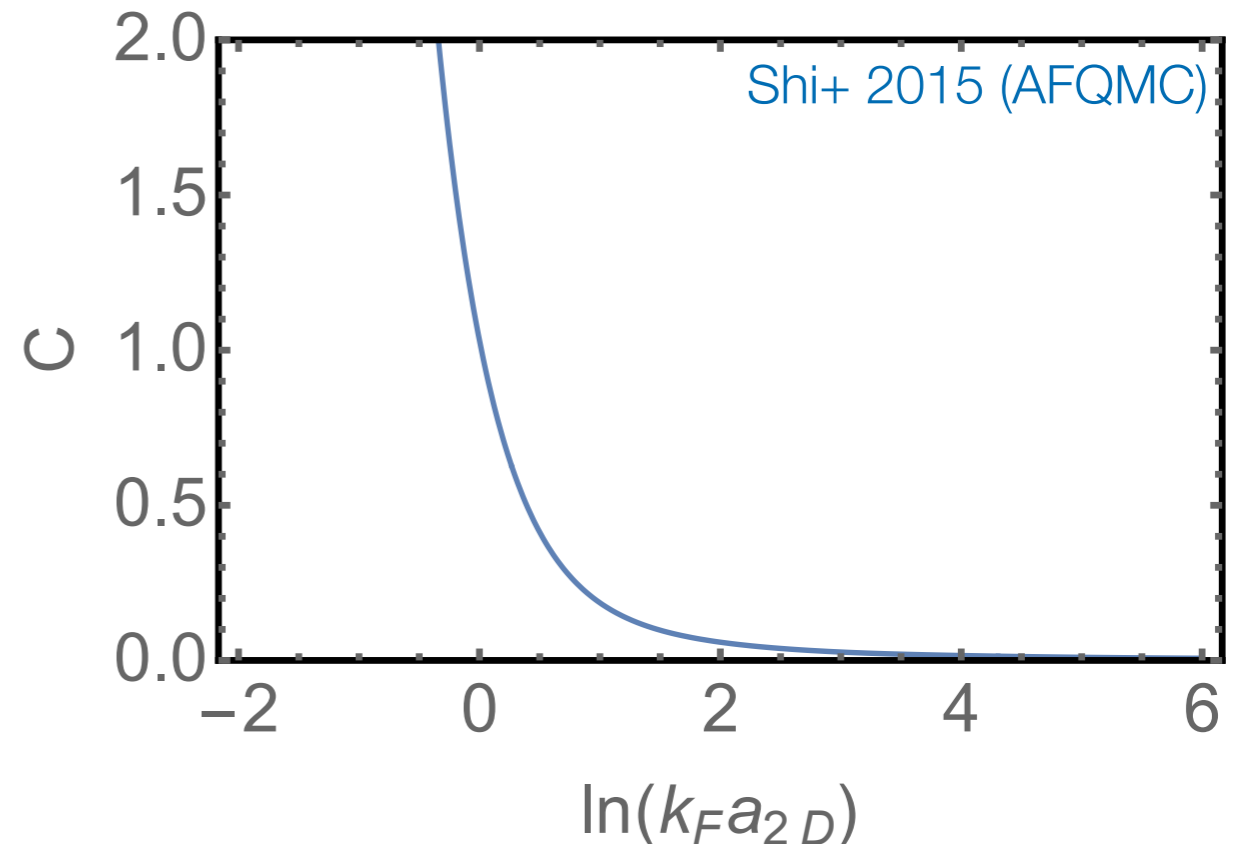
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interacting 2D Fermi gas:

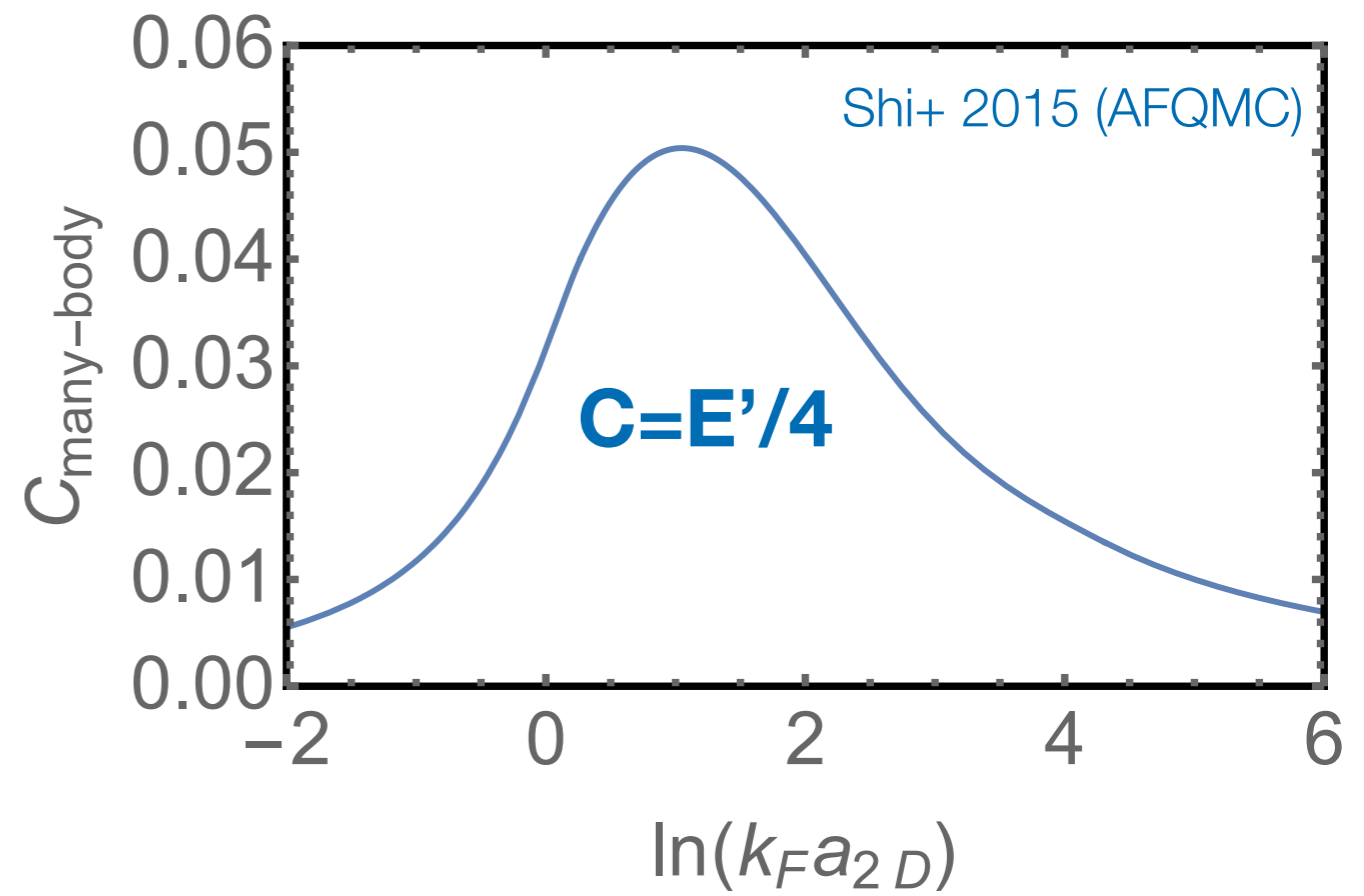
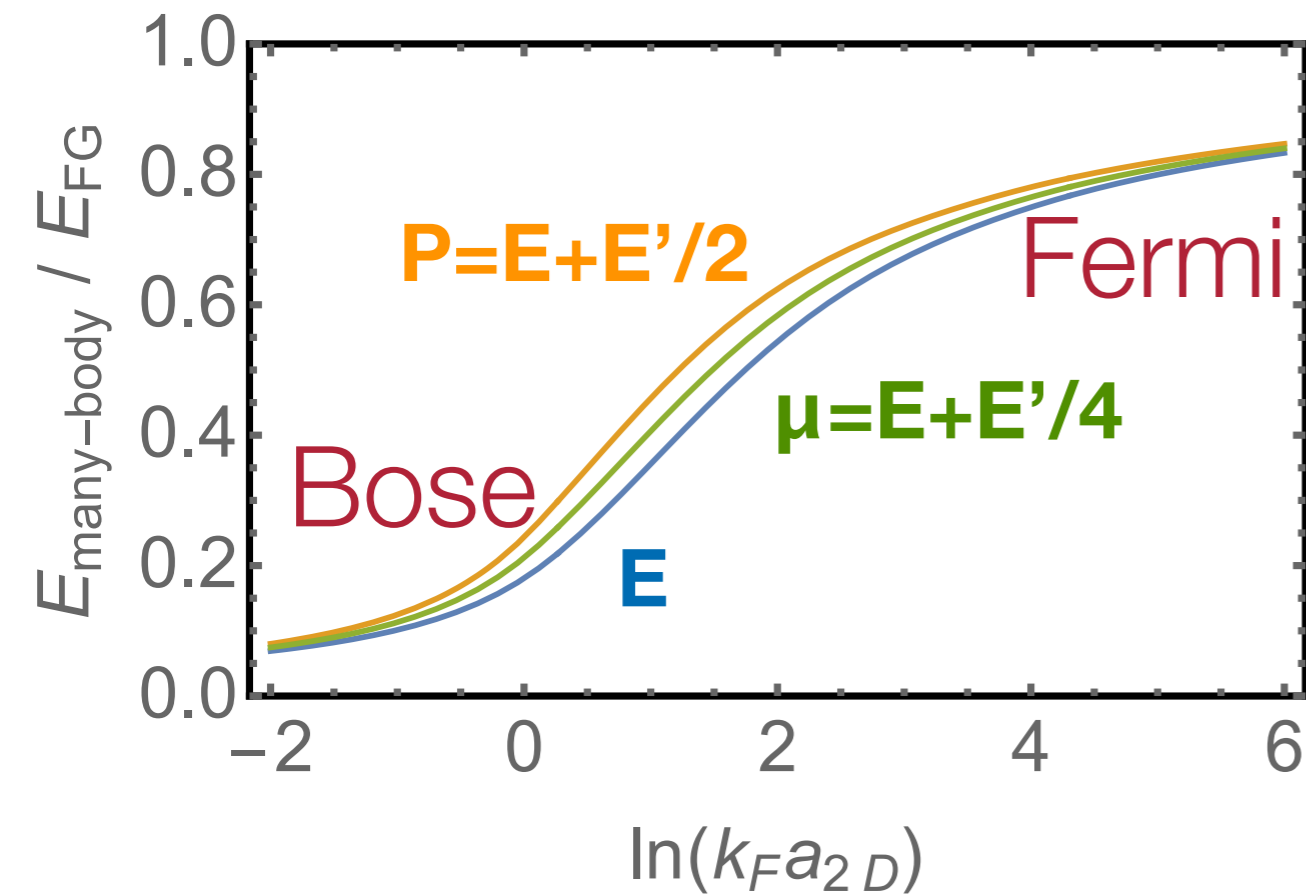
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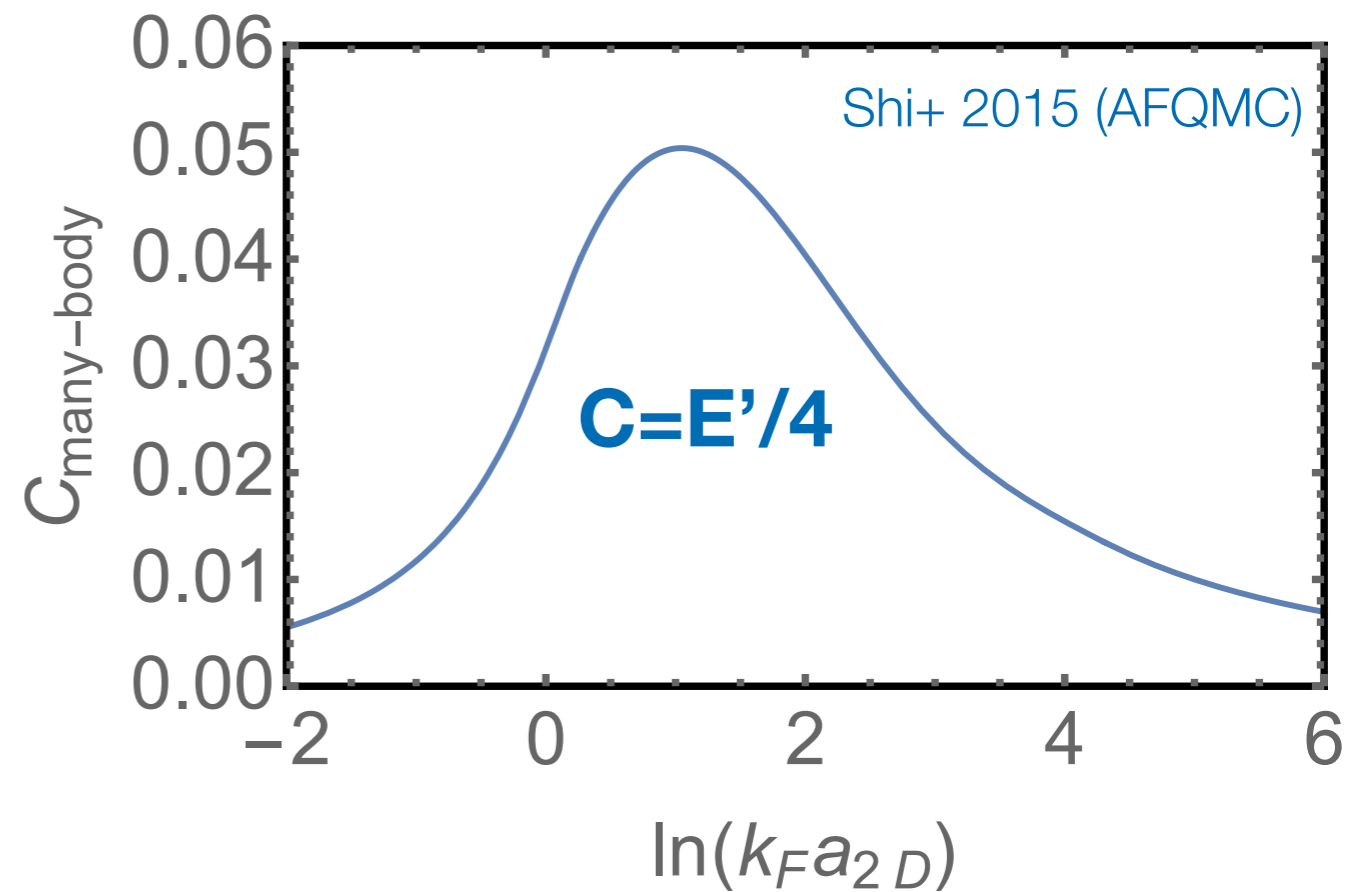
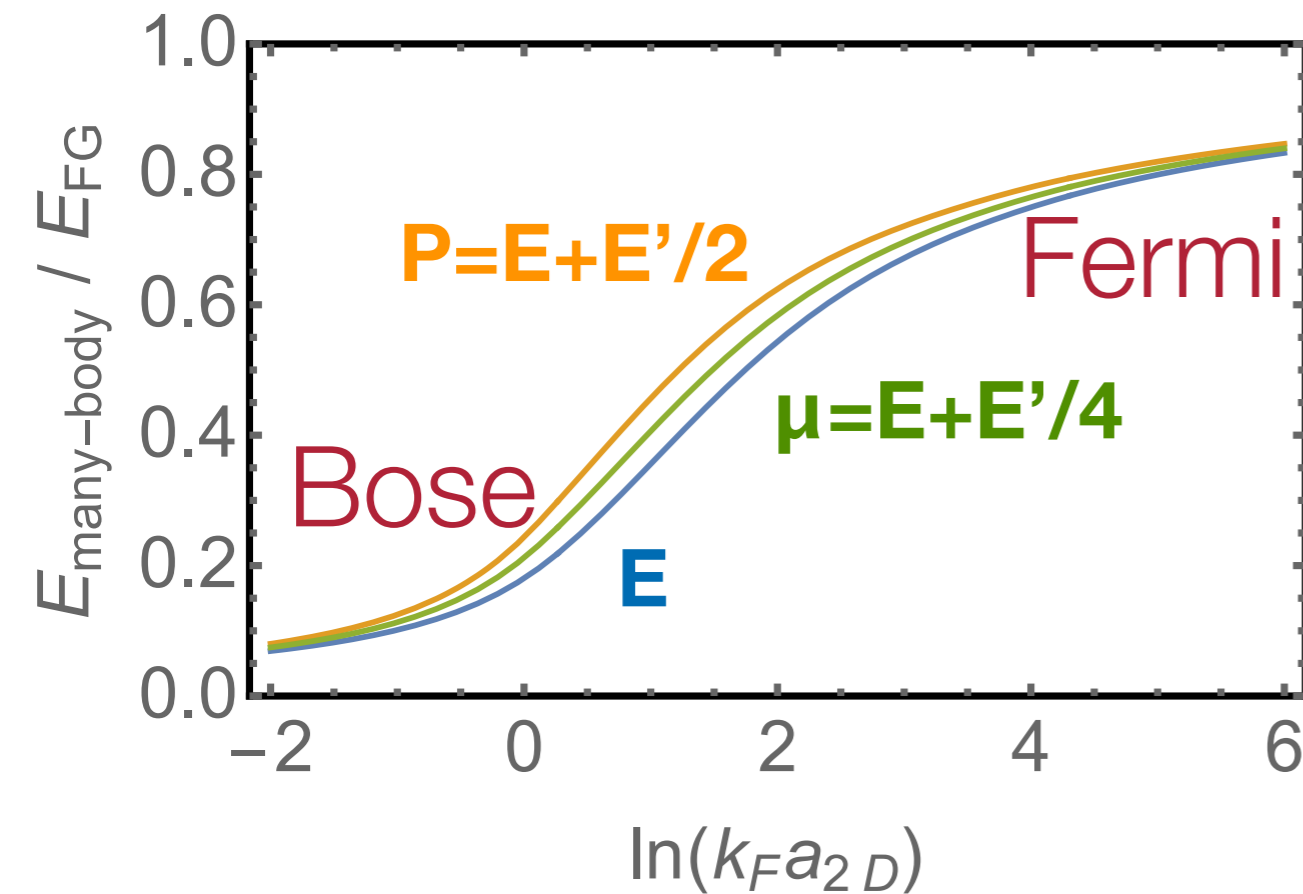
Ground-state energy

subtract two-body binding energy:



Ground-state energy

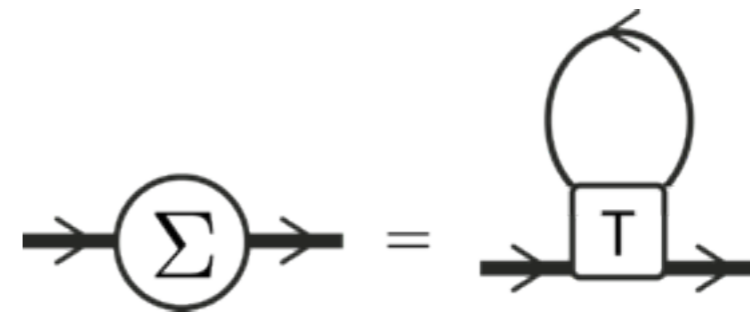
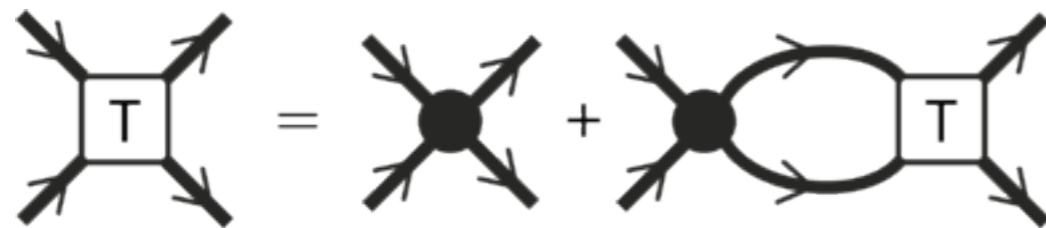
subtract two-body binding energy:



Fermi liquid theory:
Engelbrecht & Randeria 1992

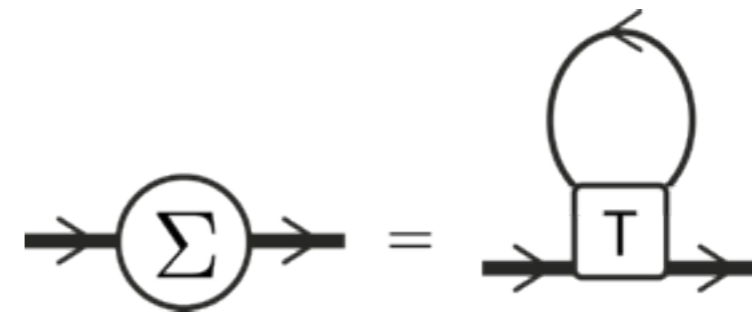
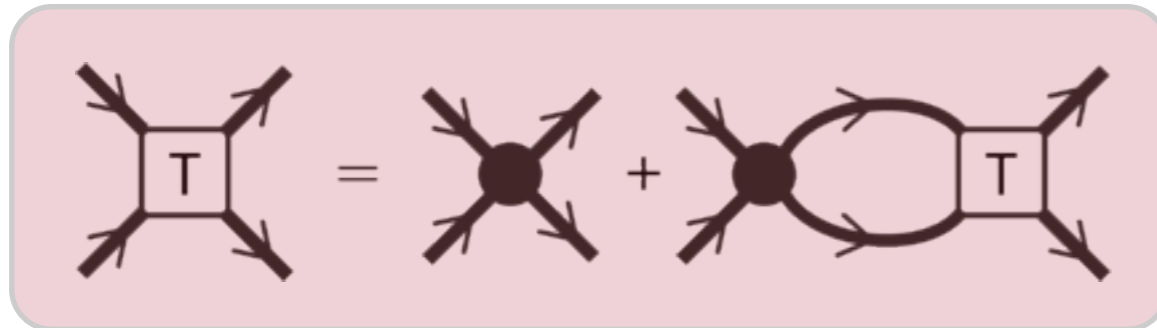
$$\frac{E}{E_{FG}} = 1 + g + \left(\frac{3}{4} - \ln 2\right)g^2 + \dots \quad \left[g = -\frac{1}{\ln(k_F a_{2D})}\right]$$

Many-body T-matrix



Nozières & Schmitt-Rink 1985;
2D: Engelbrecht & Randeria 1990

Many-body T-matrix



Nozières & Schmitt-Rink 1985;
2D: Engelbrecht & Randeria 1990

step 1: compute many-body T-matrix

two-body T-matrix:
$$T_0(E) = \frac{4\pi/m}{\ln(\varepsilon_B/E) + i\pi}$$

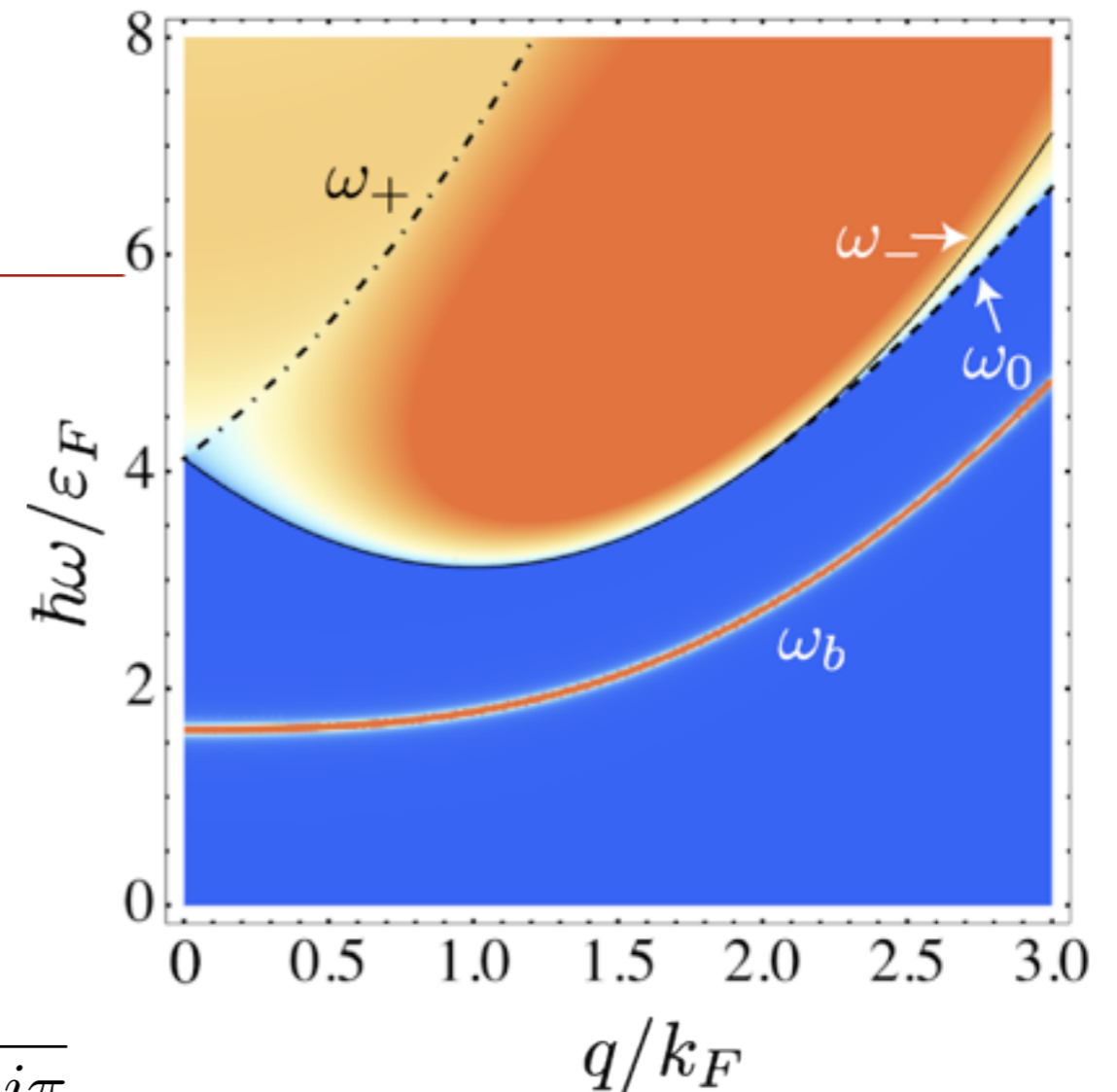
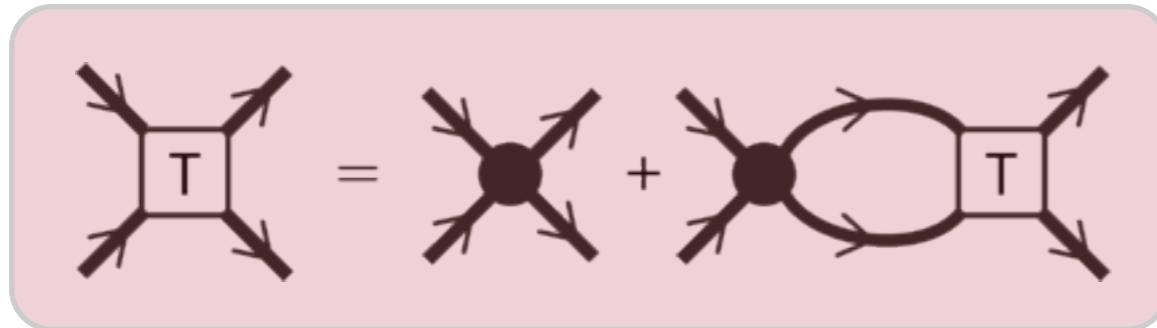
many-body: finite density medium scattering Schmidt, Enss, Pietilä & Demler 2012

$$T^{-1}(\mathbf{q}, \omega) = T_0^{-1}(\omega + i0 + \mu_\uparrow + \mu_\downarrow - \varepsilon_{\mathbf{q}}/2) + \int \frac{d^2k}{(2\pi)^2} \frac{n_F(\varepsilon_{\mathbf{k}} - \mu_\uparrow) + n_F(\varepsilon_{\mathbf{k}+\mathbf{q}} - \mu_\downarrow)}{\omega + i0 + \mu_\uparrow + \mu_\downarrow - \varepsilon_{\mathbf{k}} - \varepsilon_{\mathbf{k}+\mathbf{q}}}$$

we find compact solution

$$T(\mathbf{q}, \omega) = T_0 \left(\frac{1}{2}z \pm \frac{1}{2} \sqrt{(z - \varepsilon_{\mathbf{q}})^2 - 4\varepsilon_F \varepsilon_{\mathbf{q}}} \right) \quad z = \omega + i0 - \varepsilon_F + \mu_\downarrow$$

Many-body T-matrix



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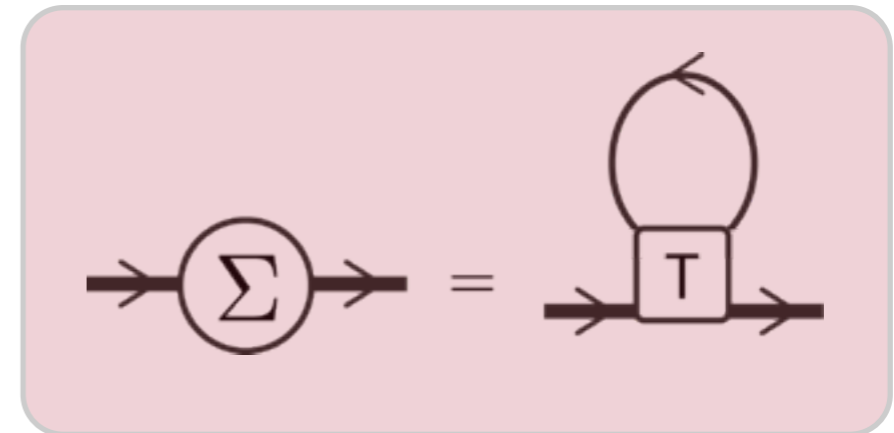
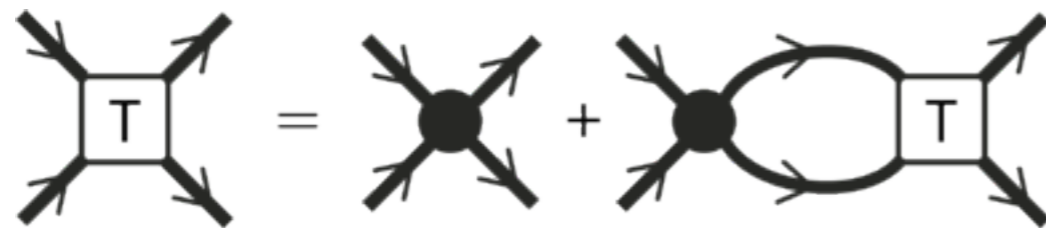
many-body: finite density medium scattering [Schmidt, Enss, Pietilä & Demler 2012](#)

$$T^{-1}(\mathbf{q}, \omega) = T_0^{-1}(\omega + i0 + \mu_\uparrow + \mu_\downarrow - \varepsilon_{\mathbf{q}}/2) + \int \frac{d^2k}{(2\pi)^2} \frac{n_F(\varepsilon_{\mathbf{k}} - \mu_\uparrow) + n_F(\varepsilon_{\mathbf{k}+\mathbf{q}} - \mu_\downarrow)}{\omega + i0 + \mu_\uparrow + \mu_\downarrow - \varepsilon_{\mathbf{k}} - \varepsilon_{\mathbf{k}+\mathbf{q}}}$$

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Fermion spectral function



step 2: fermion self-energy

$$\Sigma_{\downarrow}(\mathbf{p}, \omega) = \int_{k < k_F} \frac{d^2 k}{(2\pi)^2} T(\mathbf{k} + \mathbf{p}, \varepsilon_{\mathbf{k}} - \mu_{\uparrow} + \omega)$$

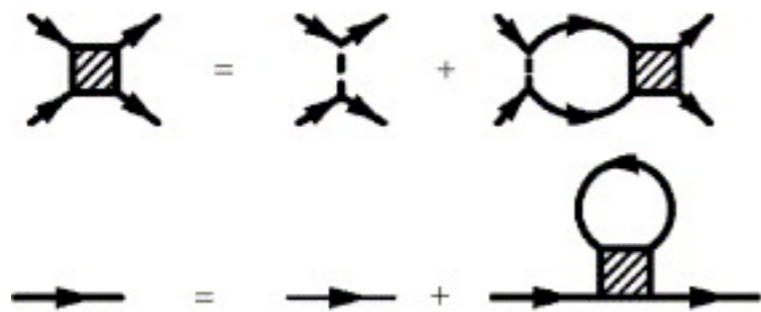
step 3: spectral function

$$\mathcal{A}_{\downarrow}(\mathbf{p}, \omega) = -2 \operatorname{Im} \frac{1}{\omega + i0 + \mu_{\downarrow} - \varepsilon_{\mathbf{p}} - \Sigma_{\downarrow}(\mathbf{p}, \omega)}$$

contains full information about
energy spectrum, quasiparticle weights, decay rates...

Luttinger-Ward approach

- repeated particle-particle scattering dominant in dilute gas:



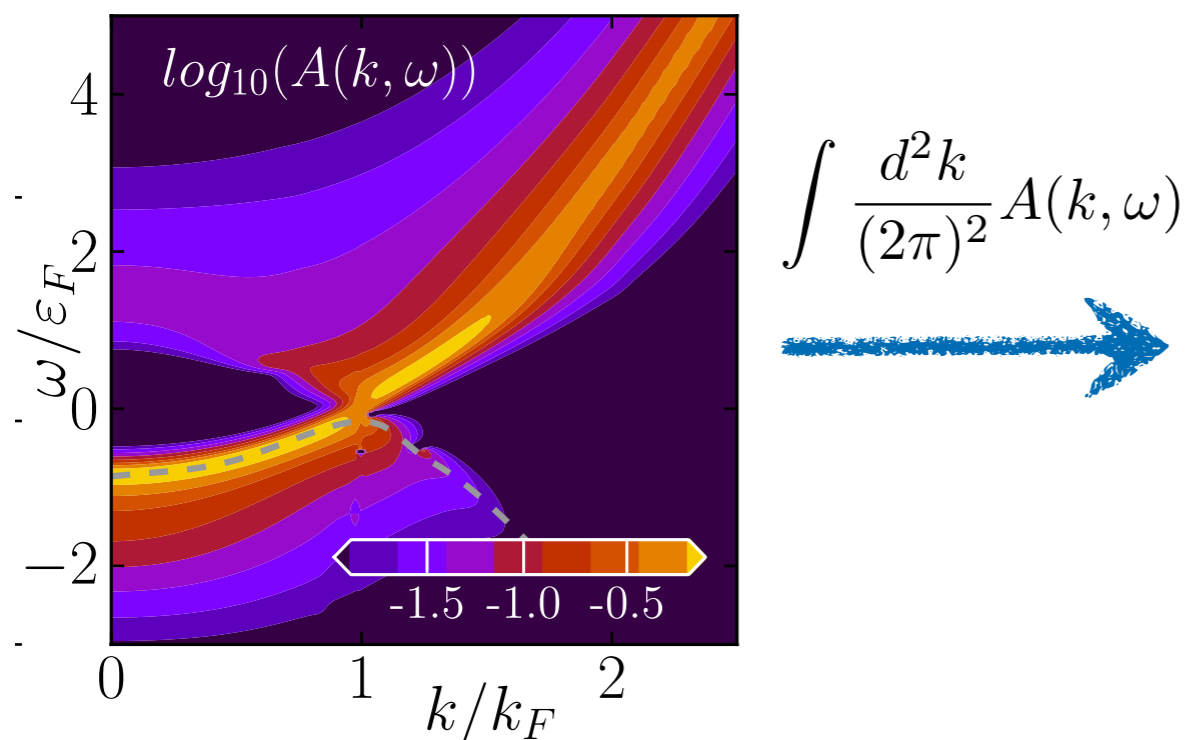
self-consistent T-matrix

Hausmann 1993, 1994;
Hausmann et al. 2007

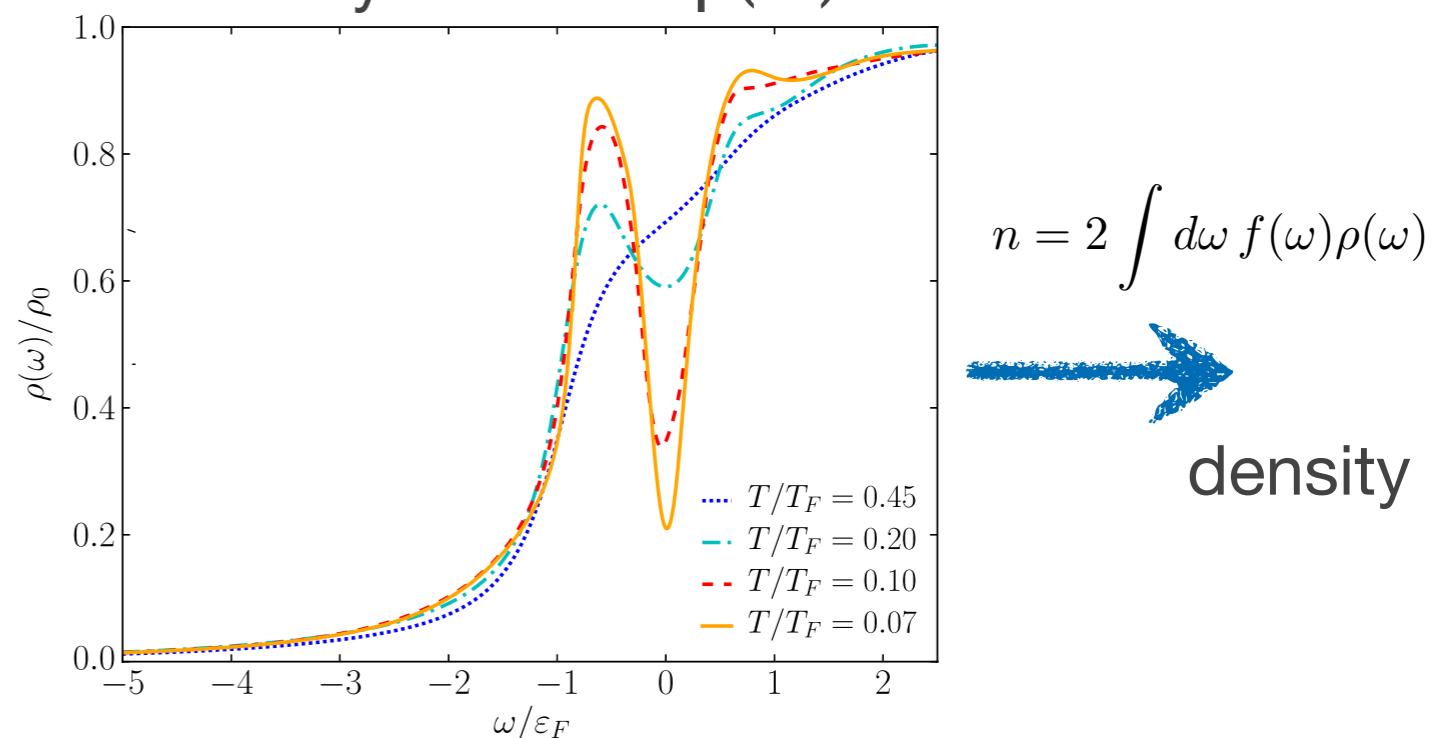
self-consistent fermion propagator
(400 momenta / 400 Matsubara frequencies)

Bauer, Parish, Enss PRL 2014

- spectral function $A(k, \omega)$



density of states $\rho(\omega)$

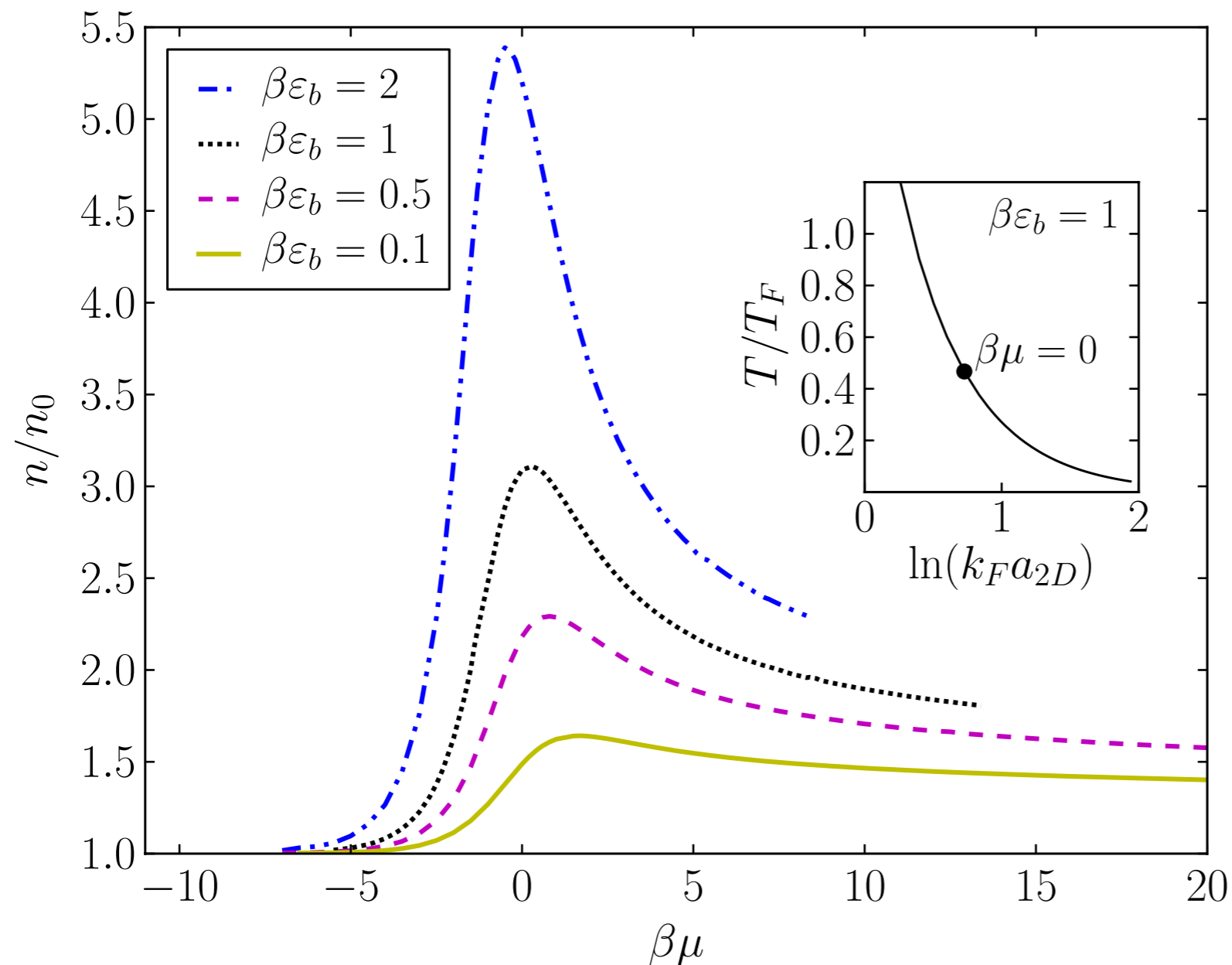


$$n = 2 \int d\omega f(\omega) \rho(\omega)$$

density

Density equation of state: theory

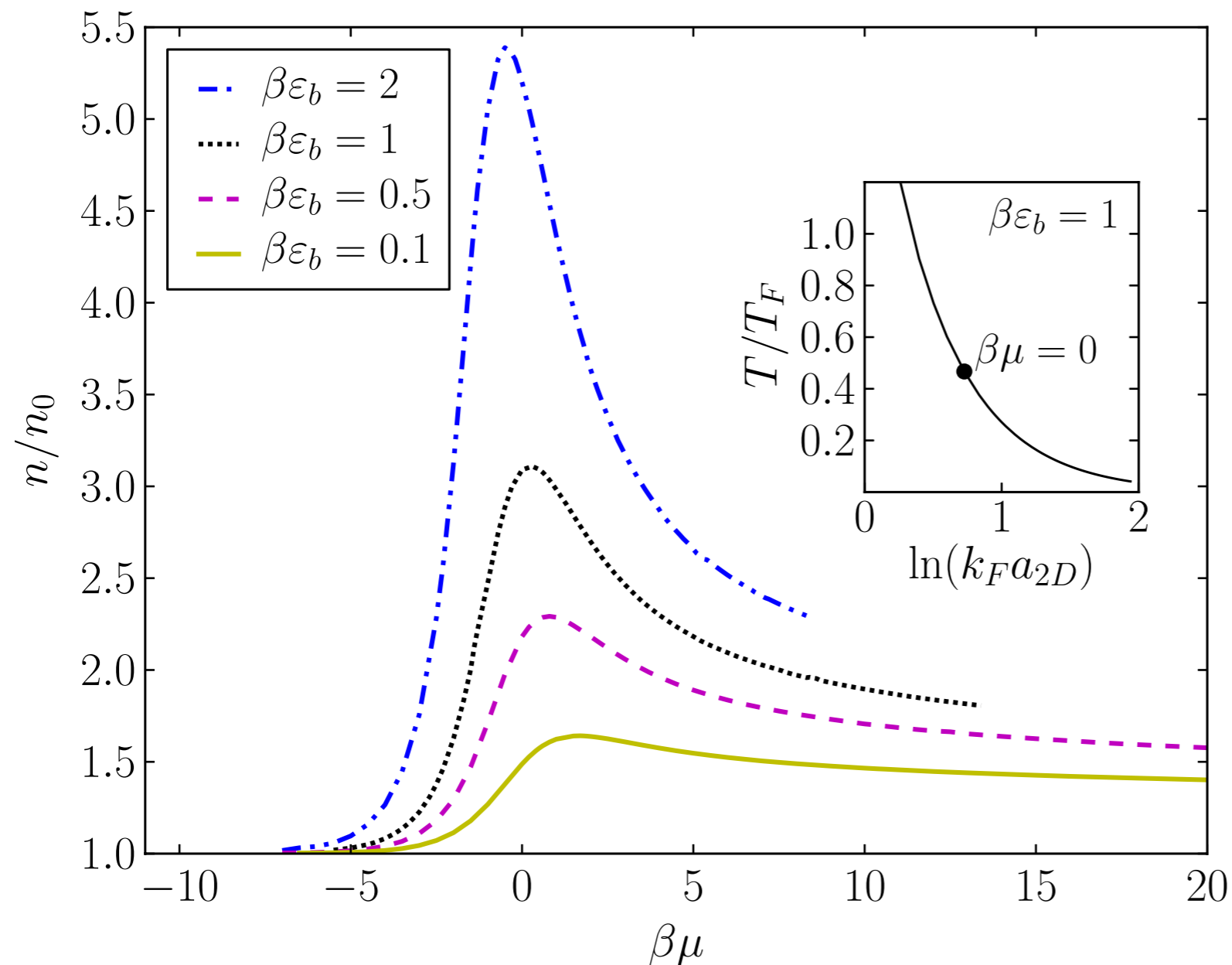
- **maximum** & **density driven crossover**



$$n = 2 \int d\omega f(\omega) \rho(\omega)$$
$$n_0 = 2 \ln(1 + e^{\beta\mu}) / \lambda_T^2$$

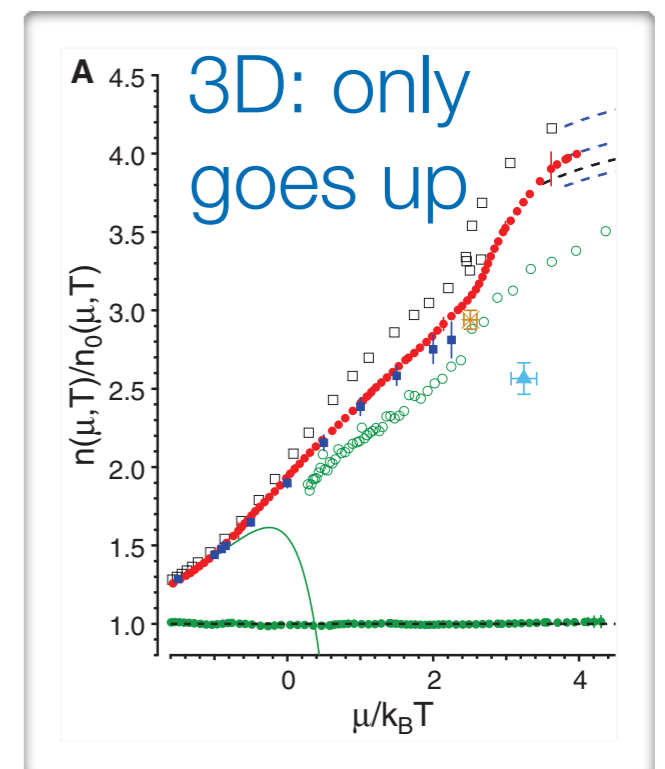
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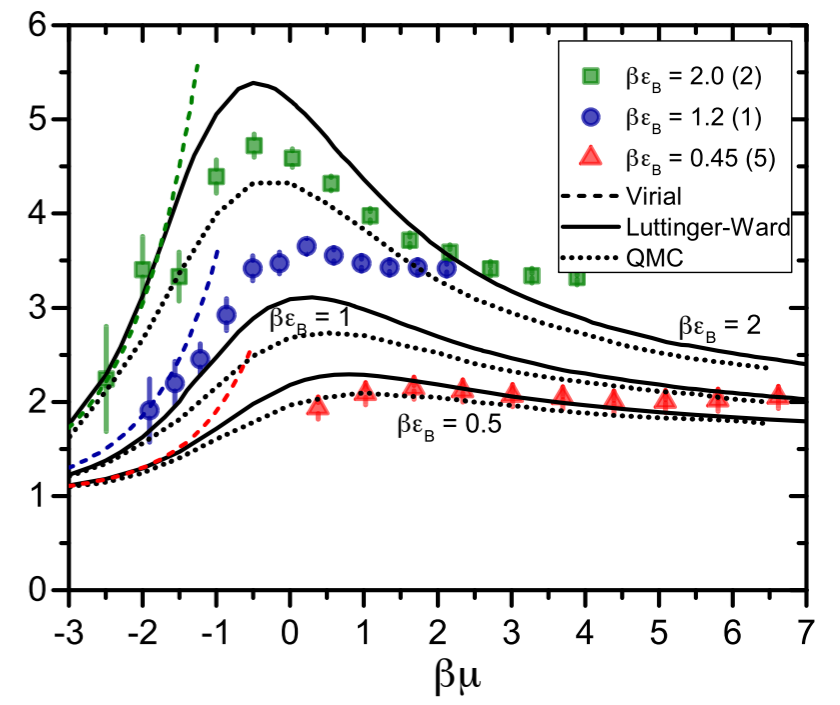
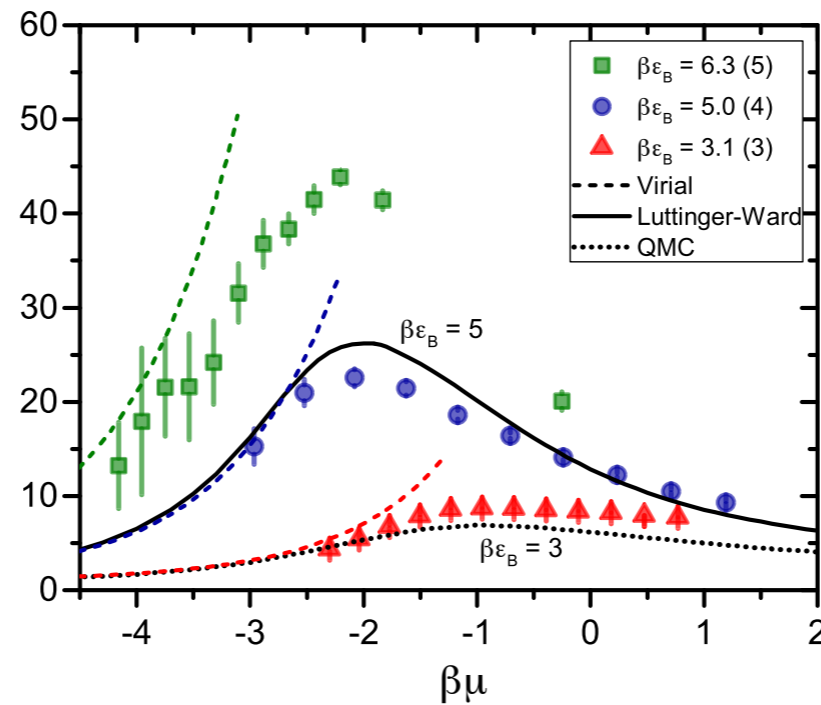
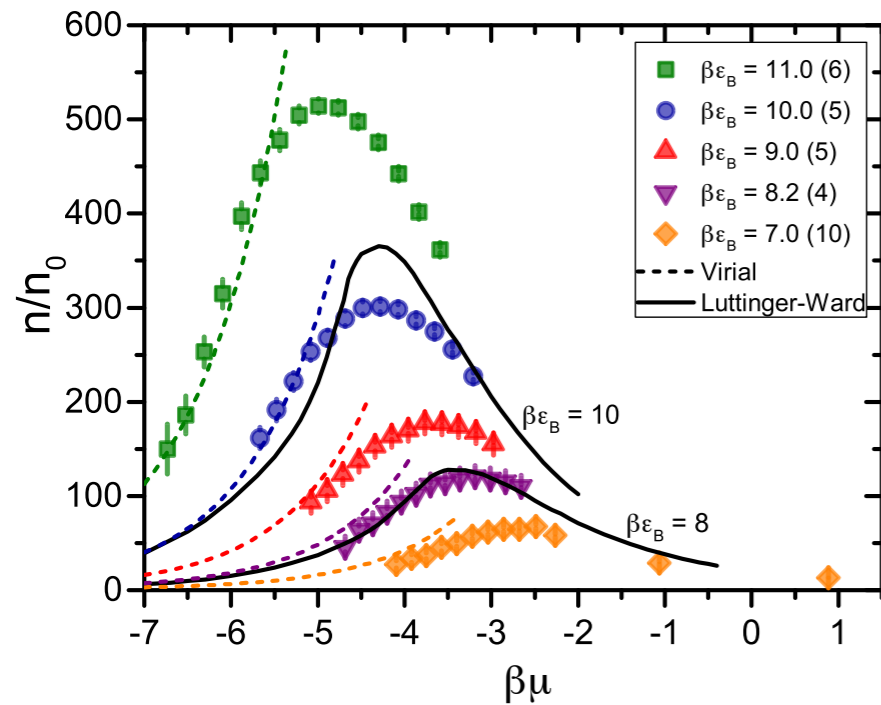


$$n = 2 \int d\omega f(\omega) \rho(\omega)$$

$$n_0 = 2 \ln(1 + e^{\beta\mu}) / \lambda_T^2$$



Equation of state: cold atom experiment



Boettcher, ..., Jochim, Enss PRL 2016

High temperature: virial expansion

- virial expansion

$$n_{\sigma} \lambda_T^2 = \ln(1 + e^{\beta\mu}) + 2\Delta b_2 e^{2\beta\mu} + 3\Delta b_3 e^{3\beta\mu} + \dots$$

$$\Delta b_2 = e^{\beta\varepsilon_B} - \int_{-\infty}^{\infty} \frac{\exp[-e^s / (2\pi)] ds}{\pi^2 + (s - \ln(2\pi\beta\varepsilon_B))^2}$$

Barth & Hofmann PRA 2014

- Bose limit ($\Delta b_2 \approx e^{\beta\varepsilon_B}$):

$$n_{\text{bos}} \approx 2e^{2\beta\mu + \beta\varepsilon_B} \lambda_T^{-2} = e^{\beta\mu_{\text{bos}}} \lambda_{\text{bos}}^{-2}$$

- good variable between Fermi and Bose limits:

$$2\tilde{\mu} = 2\mu + \varepsilon_B = \mu_{\text{bos}}$$

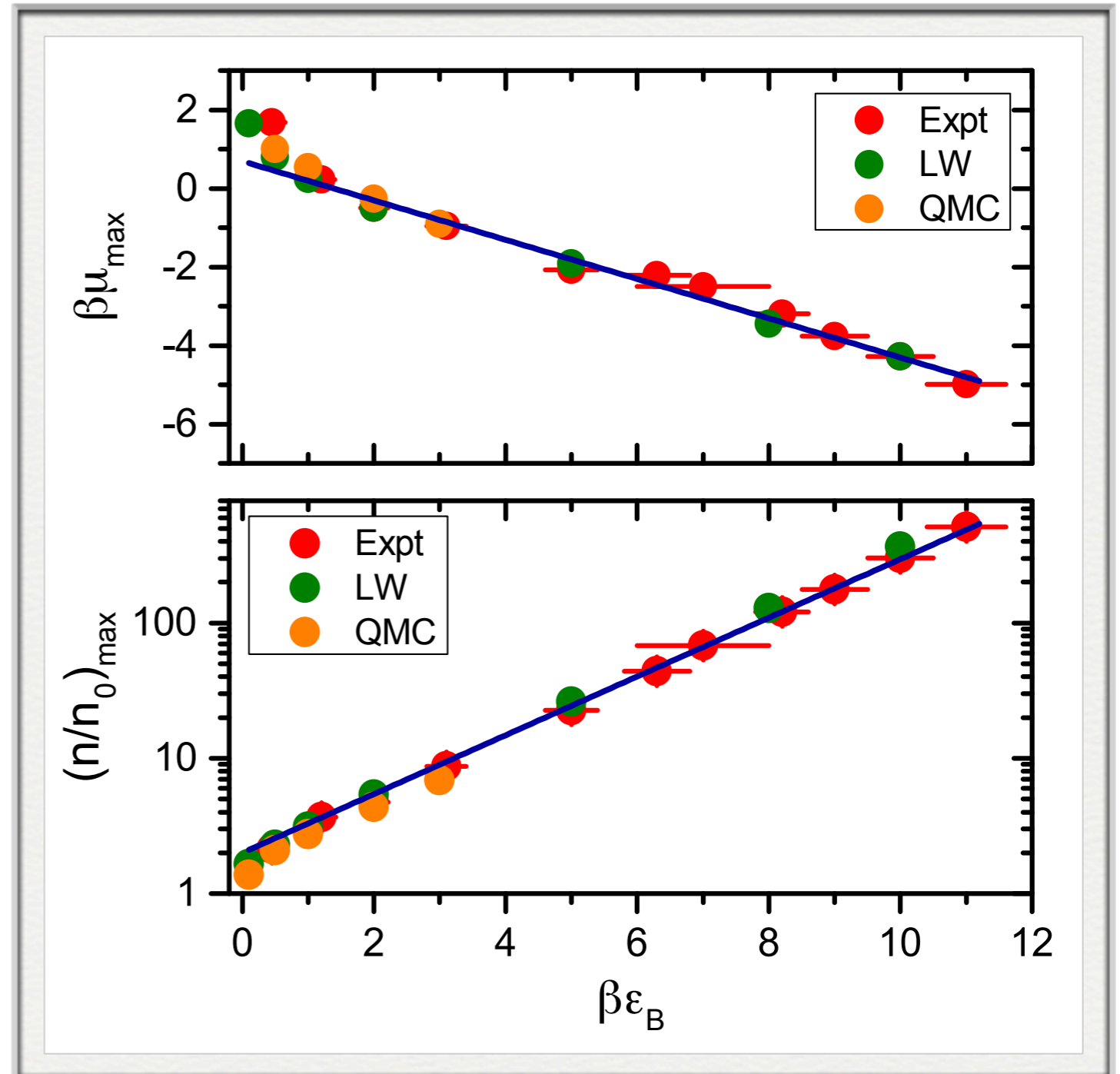
Scaling of density maximum

- **maximum** where $\tilde{\mu} \simeq 0$:

$$(\beta\mu)_{\max} \simeq -\frac{\beta\epsilon_B}{2} + \ln 2$$

at density

$$(n/n_0)_{\max} \simeq 2e^{\beta\epsilon_B/2}$$



Summary & Outlook

- **2D Fermi gas:**
 - scale invariance broken
 - exact universal relations for dilute gas
 - large density renormalization
 - density driven crossover**

Bauer, Parish & Enss 2014

Boettcher, ..., Jochim, Enss 2016

- **outlook:**
 - pseudogap (from fRG...)
 - (transverse) spin diffusion $D_0 \sim \hbar v_F/m$
 - extension to low-temperature phase

