

REMEDIOS

- LHC precision tests (di/tri-boson): what is being searched for? -



LHC

$L=27\text{km}$

$E_{\text{beam}}=362\text{MJ}$

Precision=??



AlpTransit

$L=35\text{km}$

$E_{\text{train}}=362\text{MJ}$

Precision= 10^{-6}

1cm
8cm
8cm
35 km

Francesco Riva
(CERN)

In collaboration with
Liu, Pomarol, Rattazzi 1603.03064,
Azatov, Contino, Machado 1607.05236

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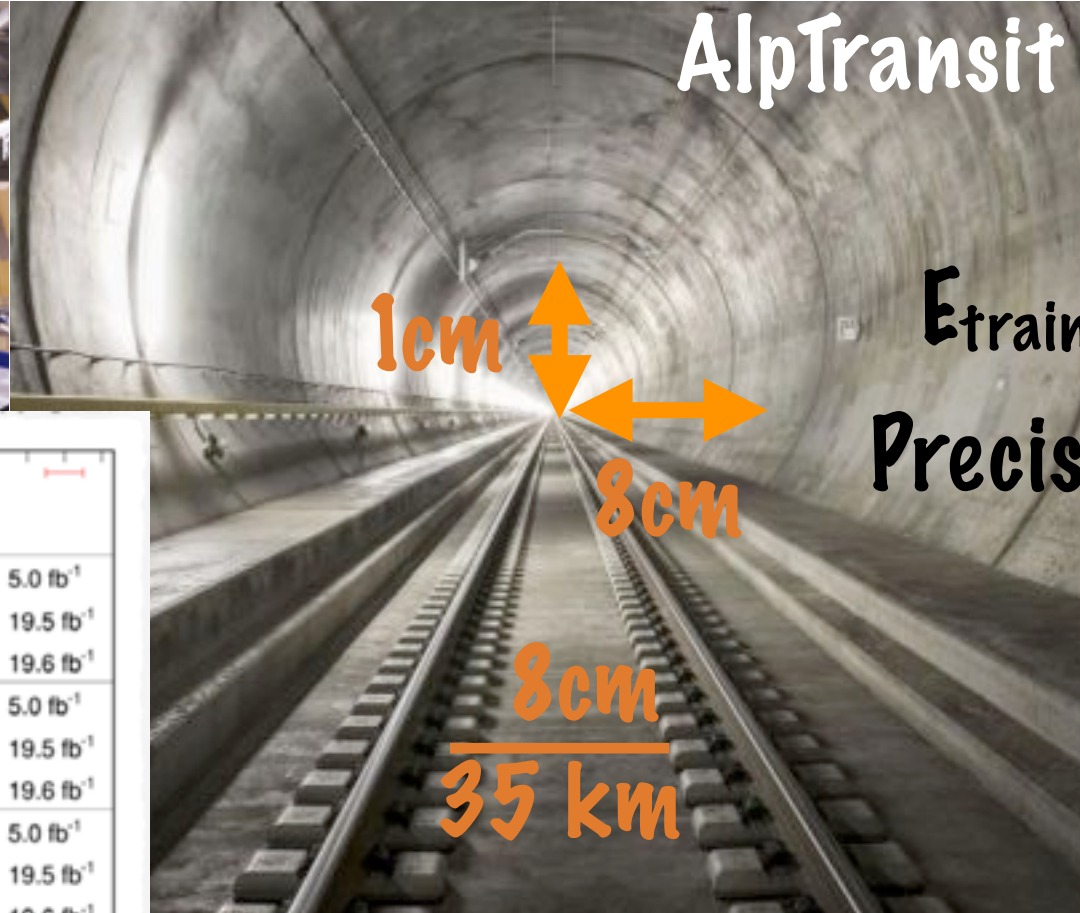


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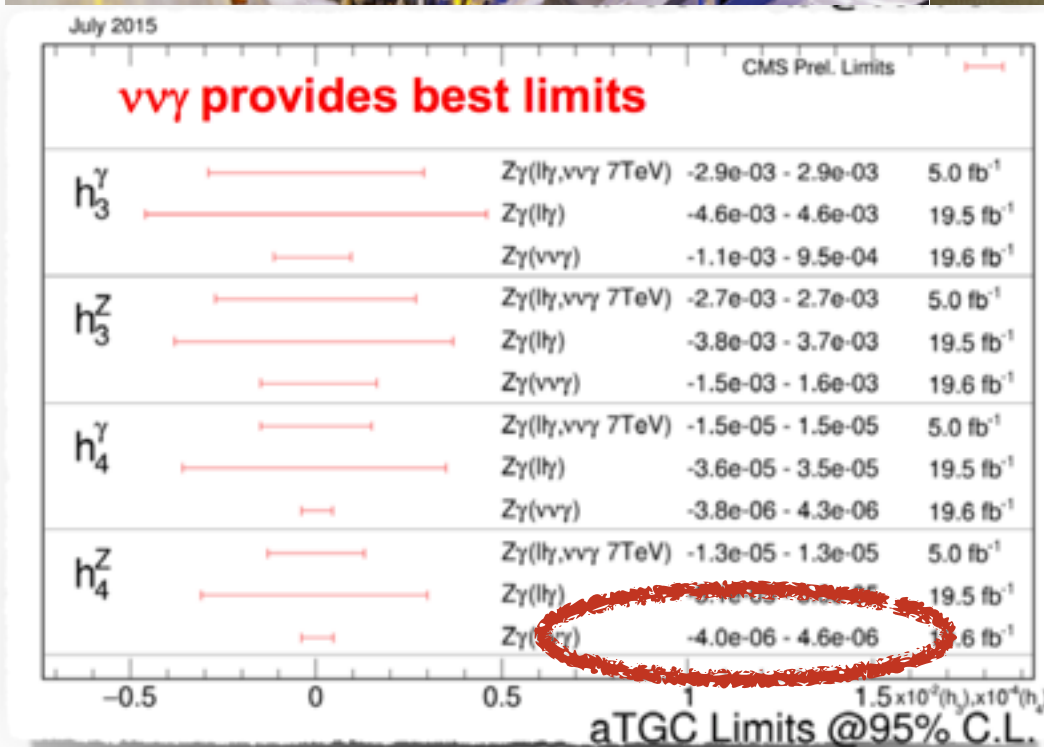


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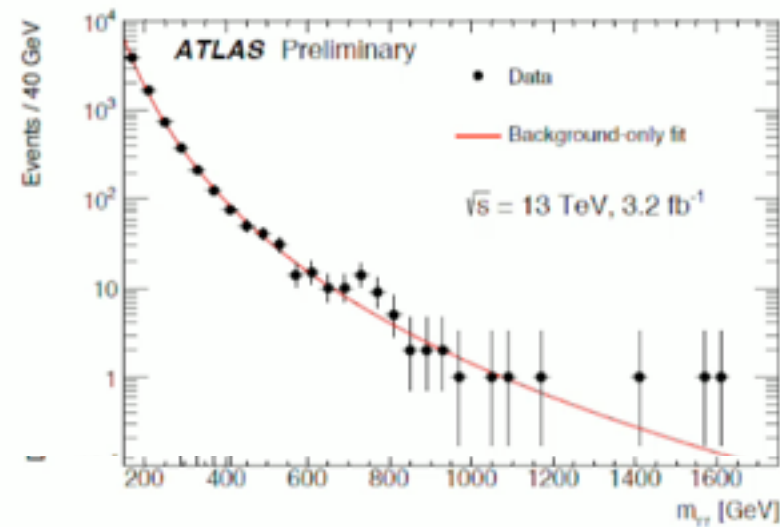


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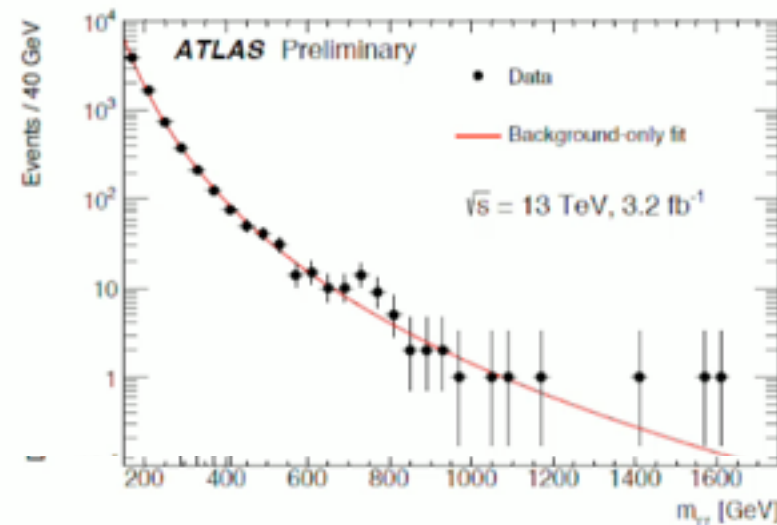
Two modes of exploration at LHC:

A) Direct Searches:

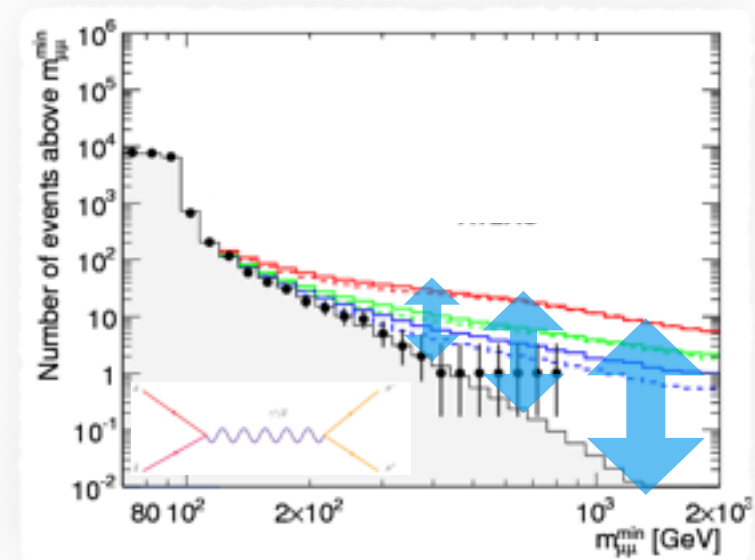
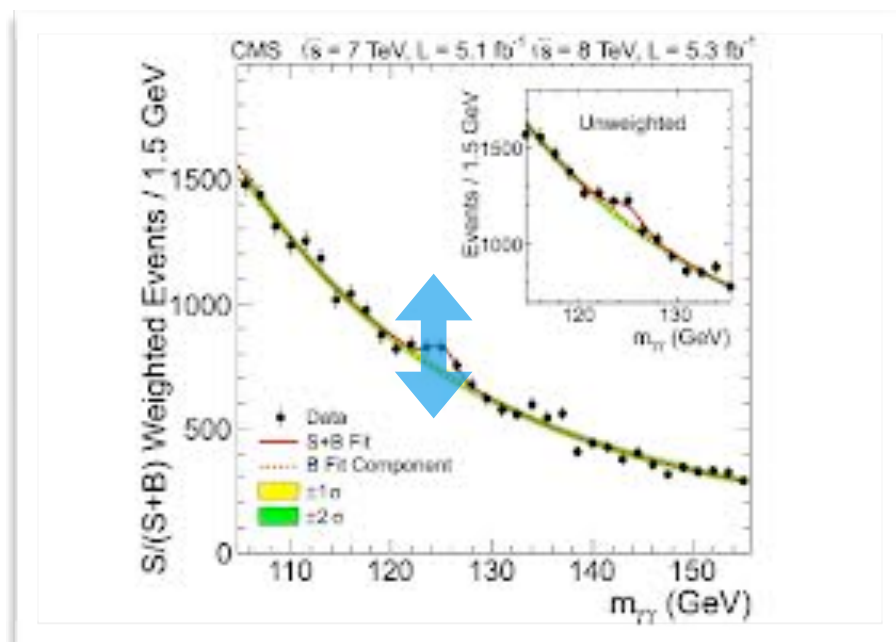


Two modes of exploration at LHC:

A) Direct Searches:



B) Indirect Searches:

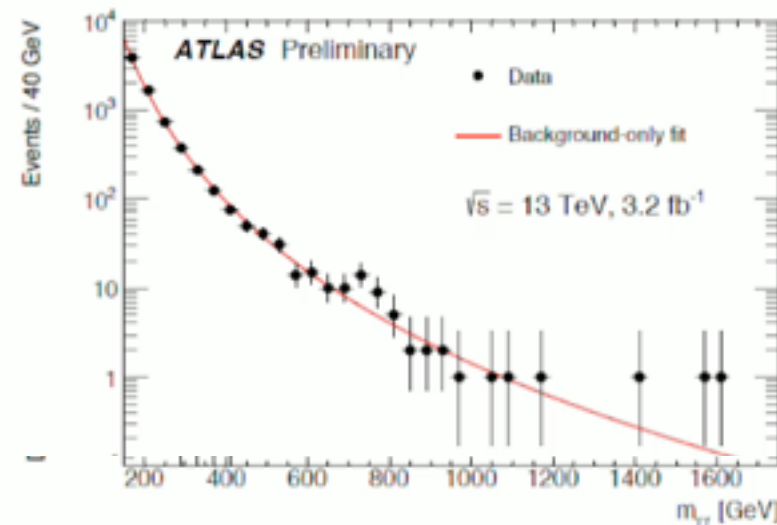


1) On SM resonance $E = m_Z, m_h$

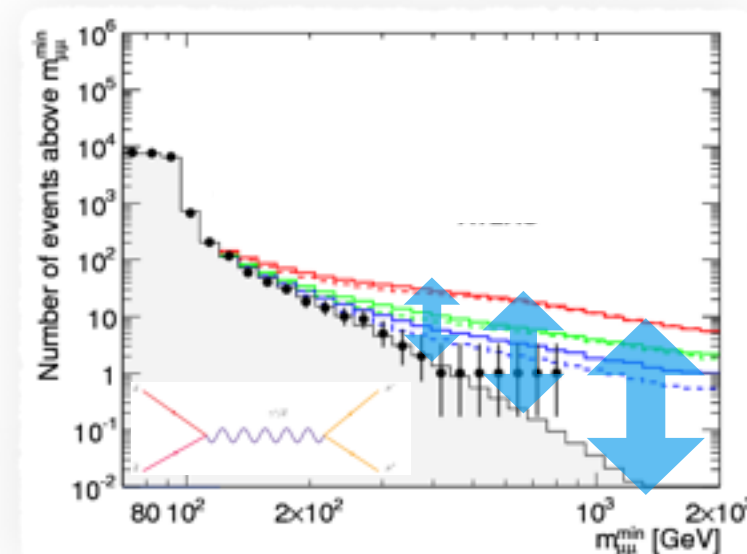
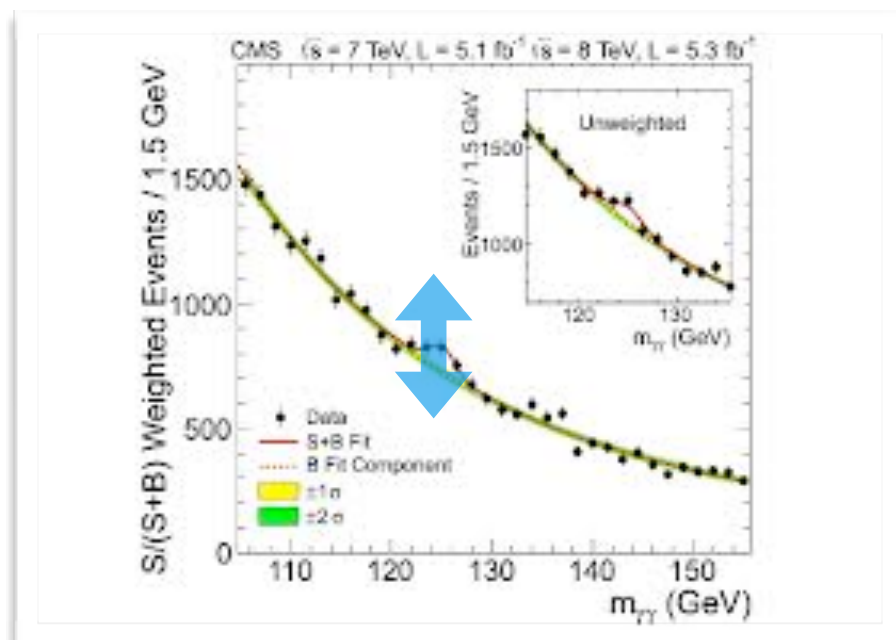
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1) **On SM resonance** $E = m_Z, m_h$

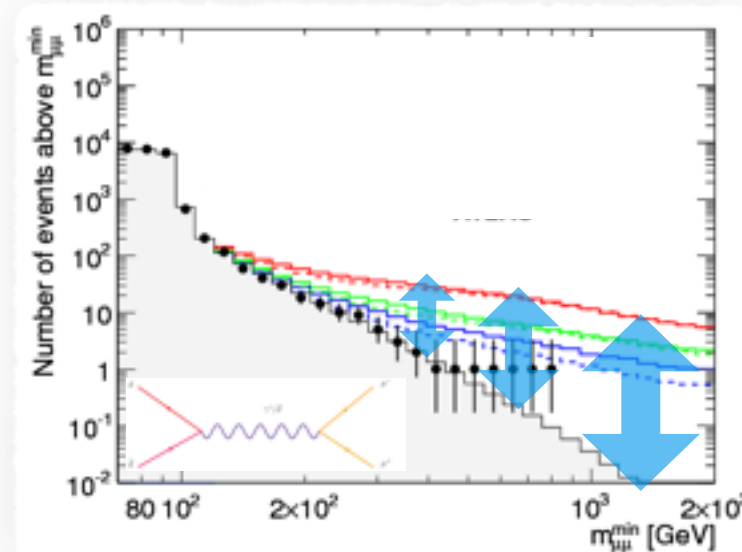
2) **Off SM resonance** $E \gg m_Z, m_h$

► Less precise (small statistics)

► Relies on **large signals!**

Two modes of exploration at LHC:

This talk: $\bar{q}q \rightarrow VV' (V'')$
 $V^{(\prime)} = W^\pm, Z, \gamma$



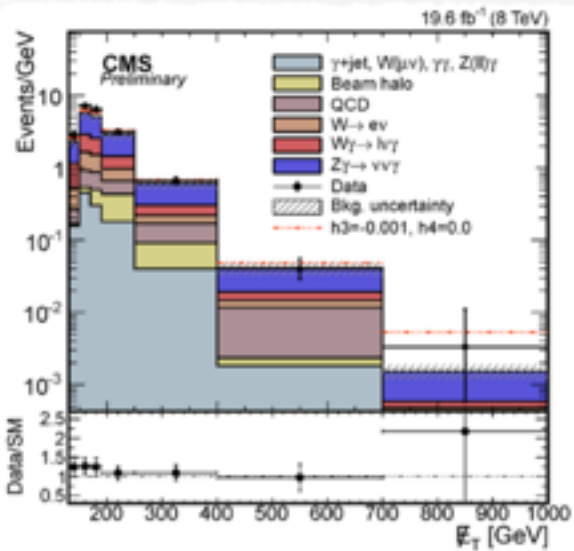
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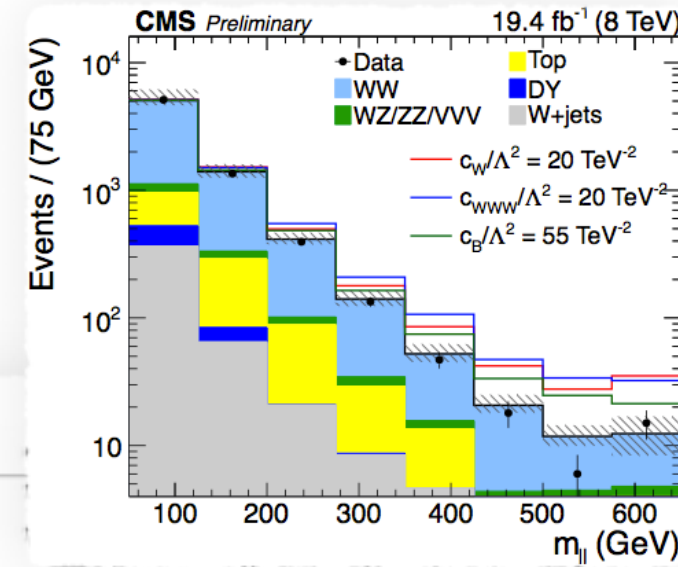
(tri)Dibosons at LHC $\bar{q}q \rightarrow VV'(V'')$

$V^{(\prime)} = W^\pm, Z, \gamma$

Plenty of Data:

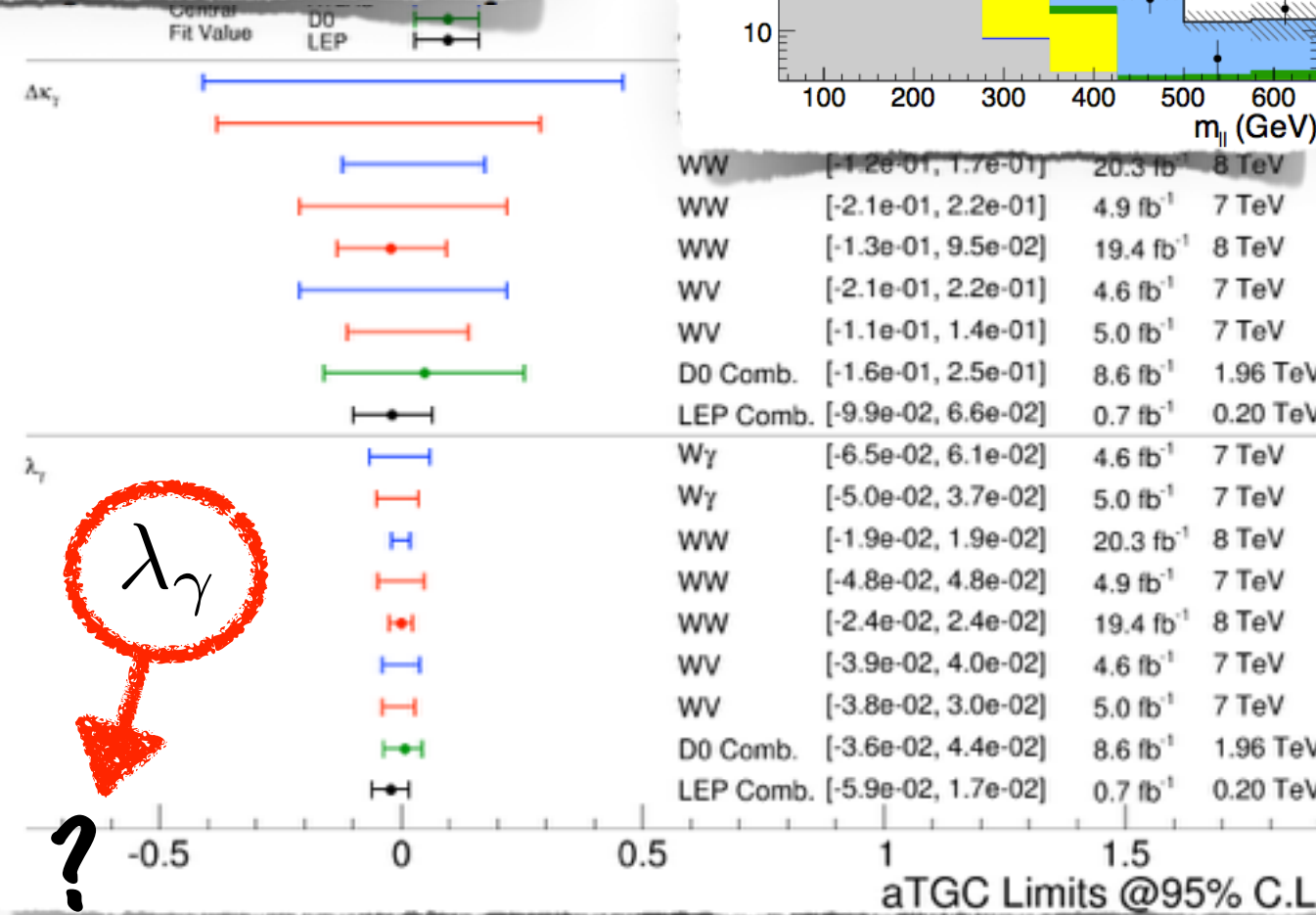


		Observed [TeV ⁻⁴]	Expected [TeV ⁻⁴]
$n = 0$	f_{T0}/Λ^4	$[-0.9, 0.9] \times 10^2$	$[-1.2, 1.2] \times 10^2$
	f_{M2}/Λ^4	$[-0.8, 0.8] \times 10^4$	$[-1.1, 1.1] \times 10^4$
	f_{M3}/Λ^4	$[-1.5, 1.4] \times 10^4$	$[-1.9, 1.8] \times 10^4$



			CMS Prel. Limits
h_3^γ	$Z\gamma(l\bar{l}, \nu\nu\gamma)$ 7TeV	$-2.9\text{e-}03 - 2.9\text{e-}03$	5.0 fb^{-1}
	$Z\gamma(l\bar{l})$	$-4.6\text{e-}03 - 4.6\text{e-}03$	19.5 fb^{-1}
	$Z\gamma(\nu\nu\gamma)$	$-1.1\text{e-}03 - 9.5\text{e-}04$	19.6 fb^{-1}
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h_4^γ	$Z\gamma(l\bar{l}, \nu\nu\gamma)$ 7TeV	$-1.5\text{e-}05 - 1.5\text{e-}05$	5.0 fb^{-1}
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h_4^Z	$Z\gamma(l\bar{l}, \nu\nu\gamma)$ 7TeV	$-1.3\text{e-}05 - 1.3\text{e-}05$	5.0 fb^{-1}
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aTGC Limits @95% C.L.



Why nobody cares?

Part 1

why nobody cares

1 Smallness

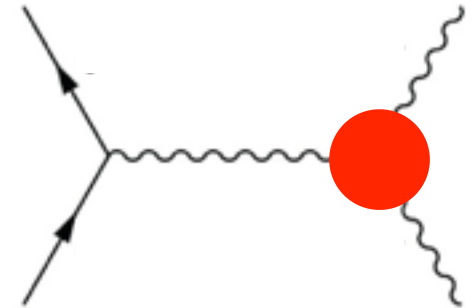
why nobody cares

These parameters $\leftrightarrow$ EFT coefficients:

$$\lambda_\gamma \leftrightarrow \quad \text{dim-6} \quad \epsilon_{abc} W_\mu^{a\,\nu} W_{\nu\rho}^b W^{c\,\rho\mu}$$

$$h_3^Z \leftrightarrow \quad \text{dim-8} \quad H^\dagger D_\mu H D_\rho B_\nu^\rho \tilde{B}^{\mu\nu}$$

$$f_{T,0} \leftrightarrow \quad W_{\mu\nu}^4$$



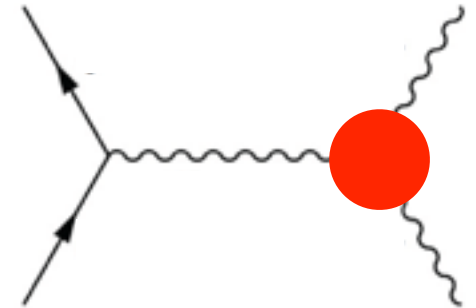
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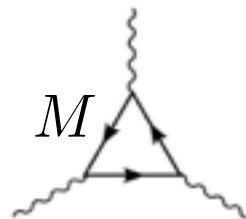
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$$f_{T,0} \leftrightarrow W_{\mu\nu}^4$$



In popular models (e.g. SUSY, CH), V_T elementary, these are **tiny**:



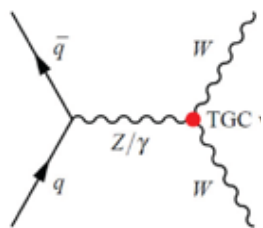
$$g\lambda_\gamma \sim \frac{g^3}{16\pi^2} \frac{m_W^2}{M^2}$$

SILH: Giudice, Grojean, Pomarol, Rattazzi' 2007

2

Inconsistent

These couplings are tested in high-energy processes:



LEP2: $E = 130 - 209 \text{ GeV}$

$$\lambda_\gamma \in [-0.059, 0.017] \sim \frac{g^2}{g_*^2} \frac{m_W^2}{M^2} \Rightarrow M \gtrsim 30 \text{ GeV}$$

95% C.L. (LEP EW: 1302.3415)

(see Contino's talk)

LHC: $E \gtrsim 1 \text{ TeV}$

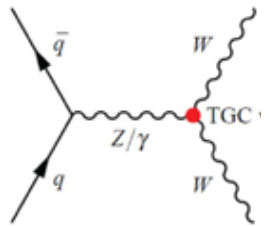
$$\lambda_\gamma \in [-0.019, +0.019]$$

...

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- There is in fact a **fundamental obstruction** to measure weakly coupled theories consistently with hypothesis of scale separation...

3 Non-Interference for BSM_6 amplitudes

Azatov, Contino, Machado, FR'16

For $E \gg m_W$ states have well defined helicity h

Amplitudes for $2 \rightarrow 2$ with different h_{tot} don't interfere

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Theorem:

A_4	$ h(A_4^{\text{SM}}) $	$ h(A_4^{\text{BSM}}) $
$VVVV$	0	4,2
$VV\phi\phi$	0	2
$VV\psi\psi$	0	2
$V\psi\psi\phi$	0	2
$\psi\psi\psi\psi$	2,0	2,0
$\psi\psi\phi\phi$	0	0
$\phi\phi\phi\phi$	0	0

Any BSM dim-6 operator

Massless limit + tree level + at least one transverse vector

- ▶ SM and BSM_6 contribute to different helicity amplitudes
- ▶ No interference

3 Non-Interference for $\mathcal{BSM}_6$ amplitudes

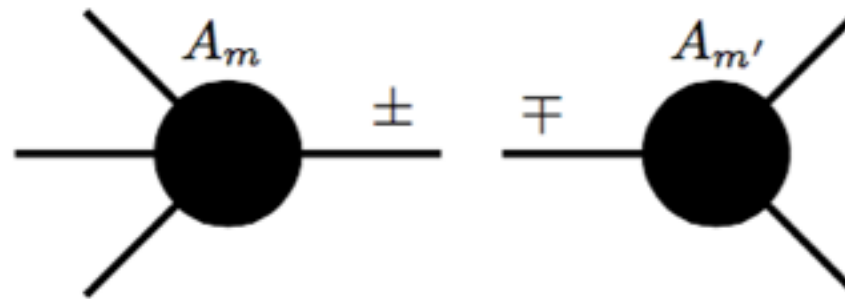
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How?

i) Helicity sums:

$$h(A_n) = h(A_m) + h(A_{m'})$$

$\nearrow n=m+m'-2$ legs



3 Non-Interference for $\mathcal{BSM}_6$ amplitudes

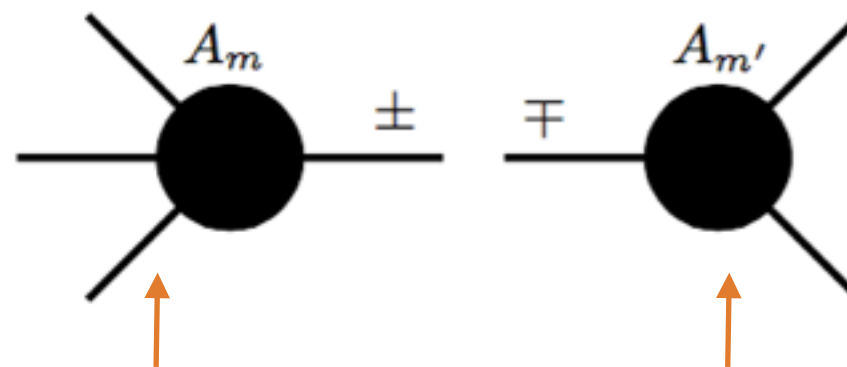
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$n = m + m' - 2$ legs



$p \in \mathbb{C}$ so that on-shell
condition $p^2 = 0$
satisfied also for A_3

SM or $\mathcal{BSM}_6$ (*)

(*)=In a basis where all effects proportional to the Equations of Motion have been eliminated

3 Non-Interference for $\mathcal{BSM}_6$ amplitudes

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ii) Helicity of 3-point $\leftrightarrow$ coupling dimension:

$$|h(A_3)| = 1 - [g]$$

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Spinor Helicity Formalism:

$$v_+(p) = (|p]_{\dot{a}}, 0) \quad \epsilon_+^\mu(p; q) = \frac{\langle q | \gamma^\mu | p]}{\sqrt{2} \langle qp \rangle}$$

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irreps of $SU(2) \times SU(2)$
 $\simeq SO(3, 1)$
 (familiar in SUSY)

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Little Group scaling (Poincaré)

$$|p_i\rangle \rightarrow t_i |p_i\rangle \quad |p_i] \rightarrow t_i^{-1} |p_i] \quad \rightarrow p_i^\mu \text{ invariant} \quad \rightarrow A \sim t_i^{-2h_i}$$

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Dimensional Analysis $[|p\rangle] \sim [p] \sim \text{GeV}^{1/2}$ $[A_{3\text{point}}] \sim \text{GeV}$

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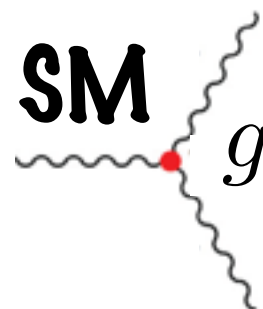
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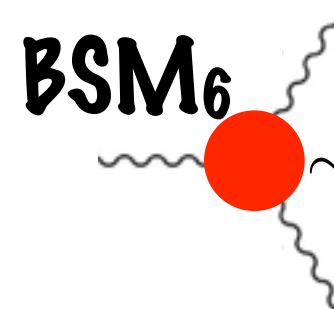
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$\rightarrow A_3(1^{h_1} 2^{h_2} 3^{h_3}) = g \langle 12 \rangle^{h_3 - h_1 - h_2} \langle 13 \rangle^{h_2 - h_1 - h_3} \langle 23 \rangle^{h_1 - h_2 - h_3}$ uniquely fixed

SM  $g \Rightarrow |h^{SM}| = 1$

$\mathcal{BSM}_6$  $\sim \frac{1}{M^2} \Rightarrow |h^{BSM}| = 3$

3 Non-Interference for $\mathcal{BSM}_6$ amplitudes

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How?

iii) SUSY* Ward Identities: $|h(A_4^{SM})| < 2$ (except $\psi^+\psi^+\psi^+\psi^+$)

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SM upliftable to SUSY+R-parity (with 1 Higgs doublet)

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Grisaru, Pendleton, van Nieuwenhuizen'77

→ e.g. $0 = \langle 0 | [Q, \Psi^+ V^+ V^+ V^+] | 0 \rangle = \sum_i \langle 0 | \Psi^+ \dots [Q, V^+] \dots V^+ | 0 \rangle \propto \langle 0 | V^+ V^+ V^+ V^+ | 0 \rangle$

$[Q, \psi] \sim V$ $[Q, V] \sim \psi$ $\psi^+ \psi^+ \dots = 0$

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$[Q, \psi] \sim V$ $[Q, V] \sim \psi$ $\psi^+ \psi^+ \dots = 0$

SM : $A(V^+ V^+ V^+ V^-) = A(V^+ V^+ \psi^+ \psi^-) = A(V^+ V^+ \phi \phi) = A(V^+ \psi^+ \psi^+ \phi) = 0$

BSM : Operators with transverse V not supersymmetrizable

Elias-Miro, Espinosa, Pomarol'14

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- i) Helicity sums $h(A_n) = h(A_m) + h(A_{m'})$
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3 Non-Interference for BSM_6 amplitudes

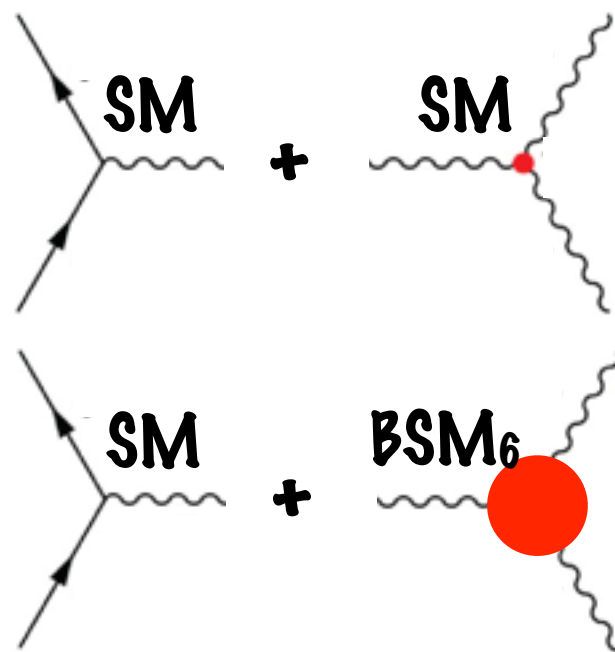
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$$h^{tot} = 1 \pm 3 = 2, 4$$

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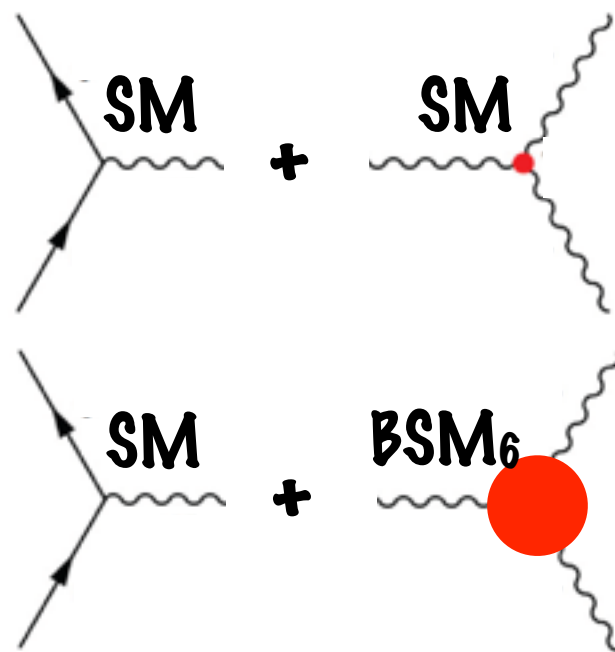
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No Interference
(dim-6, 4-point)

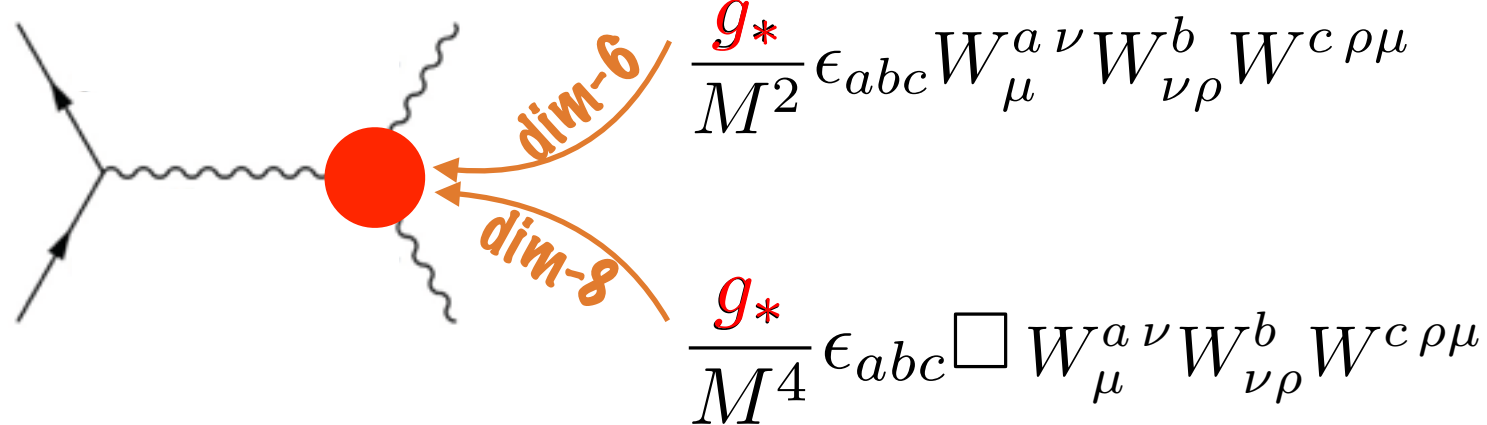
3 Non-Interference for BSM_6 amplitudes

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Implications?

$$\bar{q}q \rightarrow VV'$$

$$V^{(\prime)} = W^\pm, Z, \gamma$$



(3-point vertex has dimension of coupling $\rightarrow$ weight with g^* for illustration)

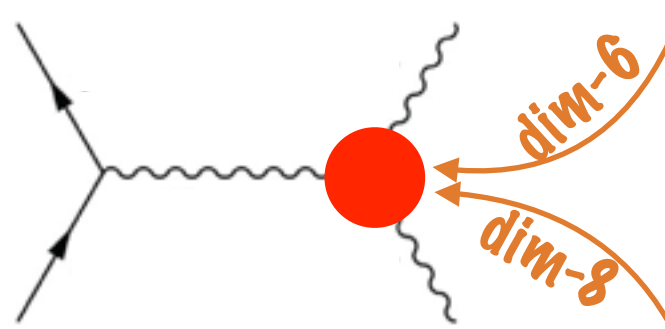
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$$\bar{q}q \rightarrow VV'$$

$$V^{(\prime)} = W^\pm, Z, \gamma$$

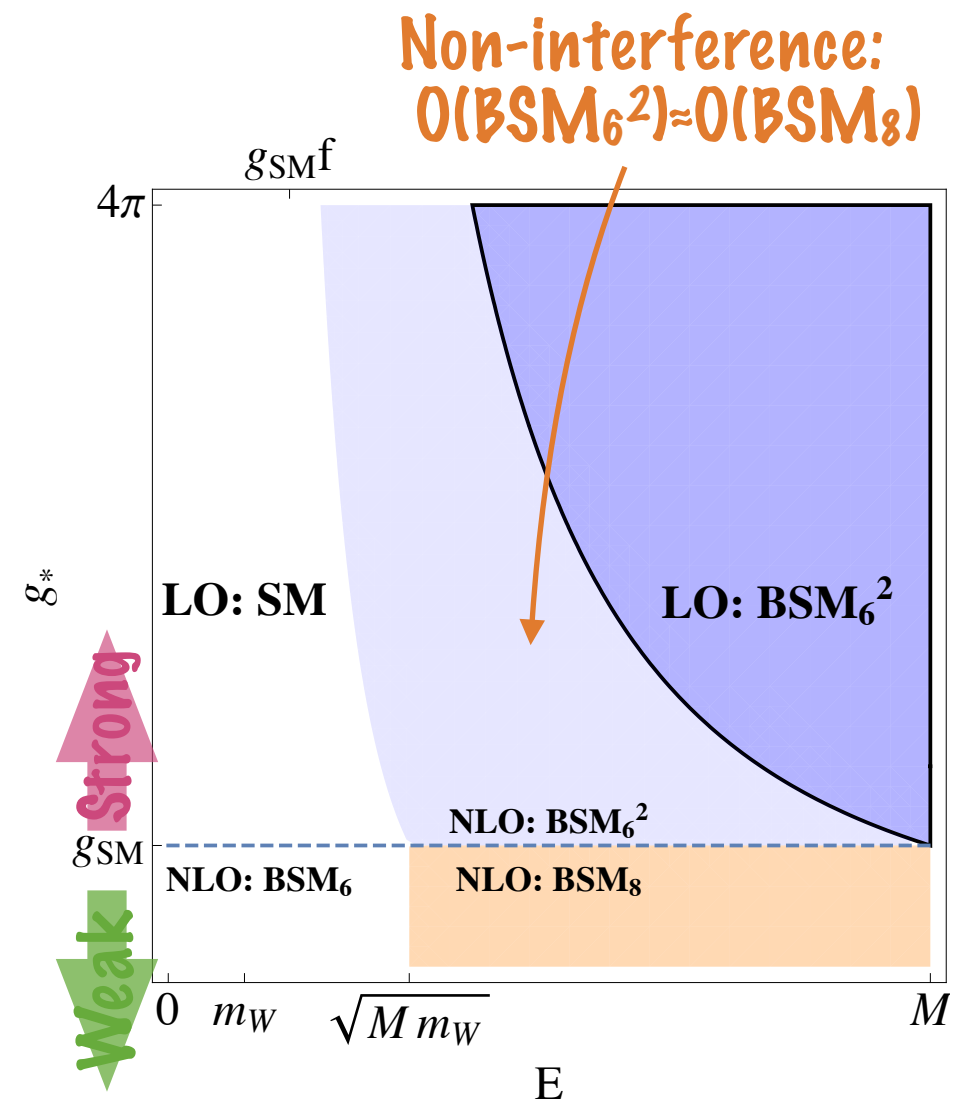


$$\frac{g_*}{M^2} \epsilon_{abc} W_\mu^a{}^\nu W_{\nu\rho}^b W^{c\rho\mu}$$

$$\frac{g_*}{M^4} \epsilon_{abc} \square W_\mu^a{}^\nu W_{\nu\rho}^b W^{c\rho\mu}$$

(3-point vertex has dimension of coupling $\rightarrow$ weight with g_* for illustration)

$$\sigma_T \sim \frac{g_{\text{SM}}^4}{E^2} \left[1 + \underbrace{\frac{g_*}{g_{\text{SM}}} \frac{m_W^2}{\Lambda^2}}_{\text{BSM}_6 \times \text{SM}} + \underbrace{\frac{g_*^2}{g_{\text{SM}}^2} \frac{E^4}{\Lambda^4}}_{\text{BSM}_6^2} + \underbrace{\frac{g_*}{g_{\text{SM}}} \frac{E^4}{\Lambda^4}}_{\text{BSM}_8 \times \text{SM}} + \dots \right]$$



why nobody cares?

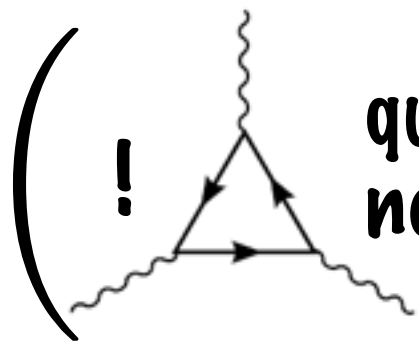
Because there is **no structured scenario**
where these searches are self-consistent
(need **strong coupling**)

Part 2

why caring (Remedios)

Strongly Coupled BSM?

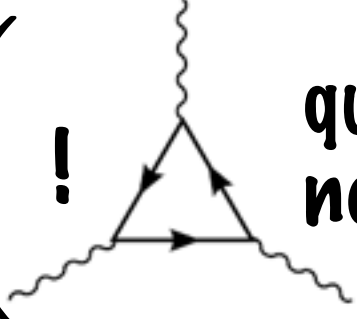
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 quantum effects generally propagate
new couplings to the whole SM !



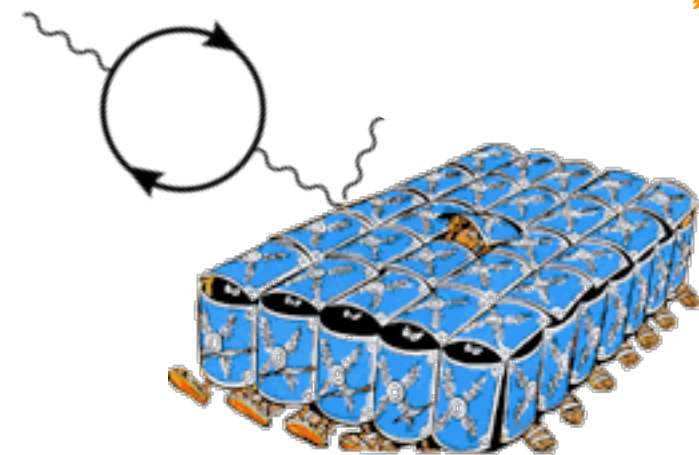
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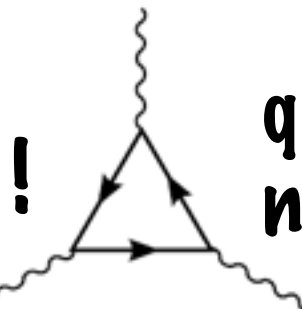
Need shielding:

Approximate Symmetries broken in SM



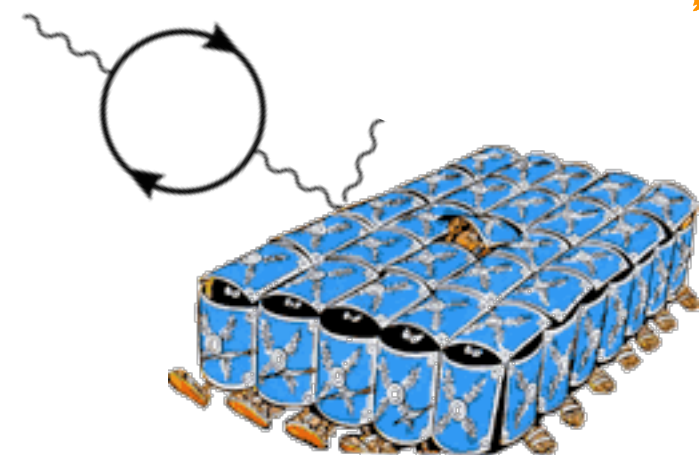
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SM Lagrangian
(dim-4)

$$A \simeq g_{SM}^2 \left(1 + \frac{g_*^2}{g_{SM}^2} \frac{E^2}{M^2} \right) \equiv g^2(E)$$

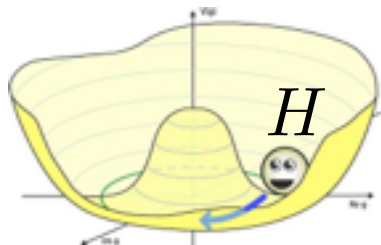


Higher-dim
operators
(dim-6, dim-8...)

Strongly Coupled BSM and Approximate Symmetries

Scalars: Composite Higgs

Higgs is a Pseudo Goldstone boson from new strong sector (symm=SO(5)/SO(4))
Georgi, Kaplan '84; Agashe, Contino, Nomura, Pomarol '04; Giudice, Grojean, Pomarol, Rattazzi '07; ...



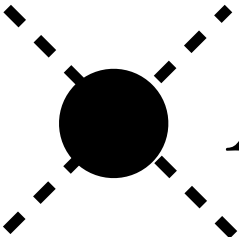
Shift Symm: $H \rightarrow H + c_{\text{n.l.}}$

$$g_* \partial_\mu H_{\text{n.l.}}$$

ϵH ~~symmetry~~

▶ $\frac{g_*^2}{M^2} (\partial_\mu |H|^2)^2$ **large**

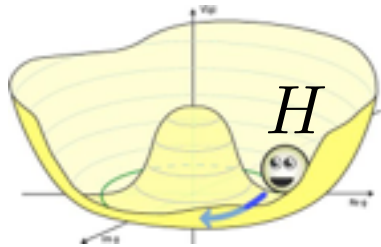
▶ $\lambda |H|^4$ **small**

▶ Large effects $V_L V_L \rightarrow V_L V_L$ at high-E:  $A \sim \lambda \left(1 + \frac{g_*^2}{\lambda} \frac{E^2}{M^2} \right)$

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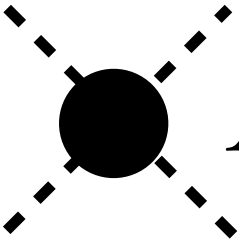
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Fermions: (partially) Composite fermions

Symmetry: Chiral (or Non-linear SUSY)

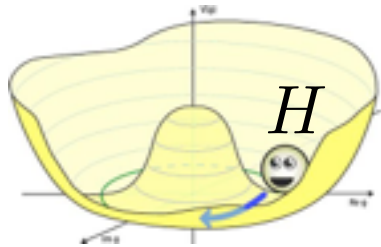
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▶ Large effects in, e.g. dijets at LHC

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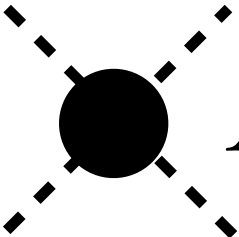
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Vectors: ?

Strong transverse vectors at high-E

Problem: Gauge bosons associated with **weak** SM coupling ($\partial_\mu + igA_\mu$)
how compatible with $g^* \gg g$?

(Even if composite, $g = \frac{4\pi}{\sqrt{N}}$ it would never appear strong)

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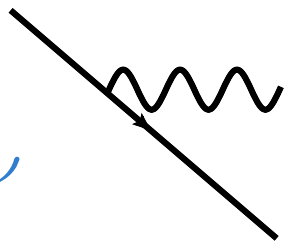
However:

Two ways a particle can couple to gauge boson:

$$g \bar{\psi}_{\text{new}} A_\mu \gamma^\mu \psi_{\text{new}}$$

monopole/dipole

$$g_* \bar{\psi}_{\text{new}} \sigma^{\mu\nu} \psi_{\text{new}} F_{\mu\nu}$$



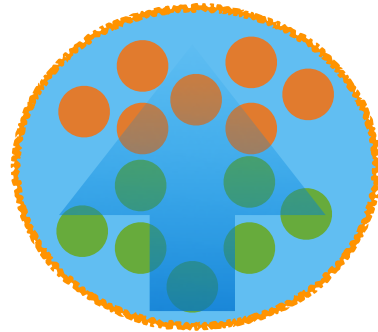
Strong transverse vectors at high-E

$$g = 0$$

- ▶ A_μ composite, g^*
- ▶ No light charged d.o.f

$$\mathcal{L} = \frac{M^4}{g_*^2} L \left(\frac{\partial_\mu}{M}, g_* \frac{\hat{F}_{\mu\nu}}{M^2}, \Phi \right)$$

(Euler-Heisenberg)



- ▶ Strong irrelevant interactions (multipole)

Strong transverse vectors at high-E

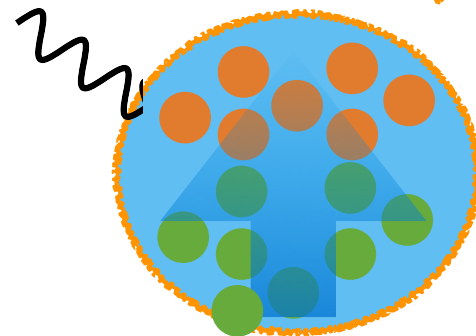
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$$q = \frac{g}{g_*}$$

- ▶ Light d.o.f. with small charge $\frac{g}{g_*}$

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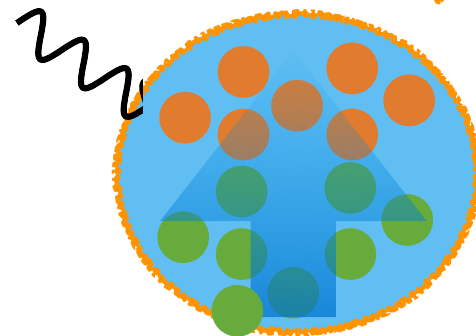
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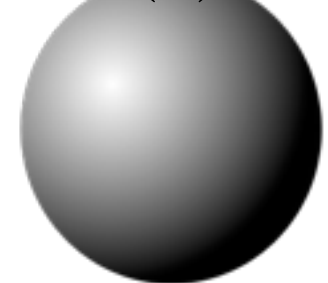
$$U(1)_{local}^3 \times SU(2)_{global}$$



Non-abelian

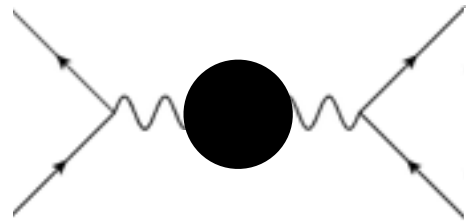
Inonu-Wigner Contraction ('53)
#generators invariant
(Like Poincaré → Galilei)

$$SU(2)_{local}$$



Strong transverse vectors: Implications

▶ $\frac{1}{M^2} (D_\rho W_\mu^{a,\nu})^2$



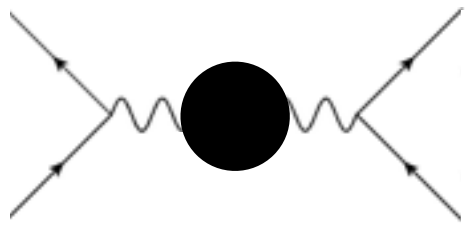
$$W, Y \simeq \frac{m_W^2}{M^2} \lesssim 10^{-3}$$

Barbieri, Pomarol, Rattazzi, Strumia'04

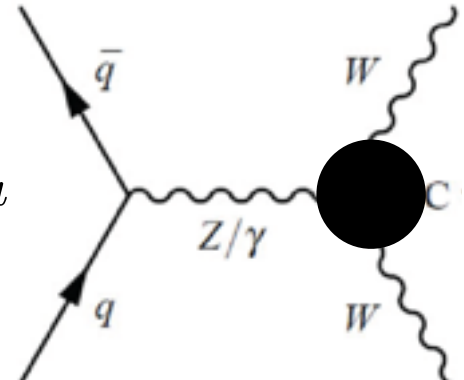


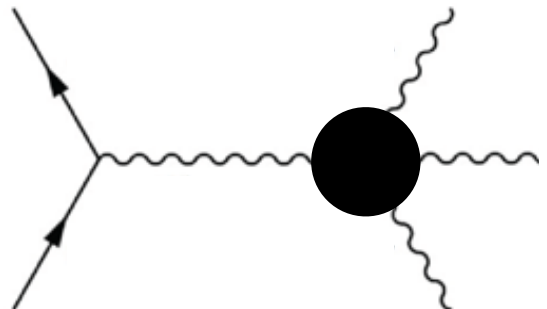
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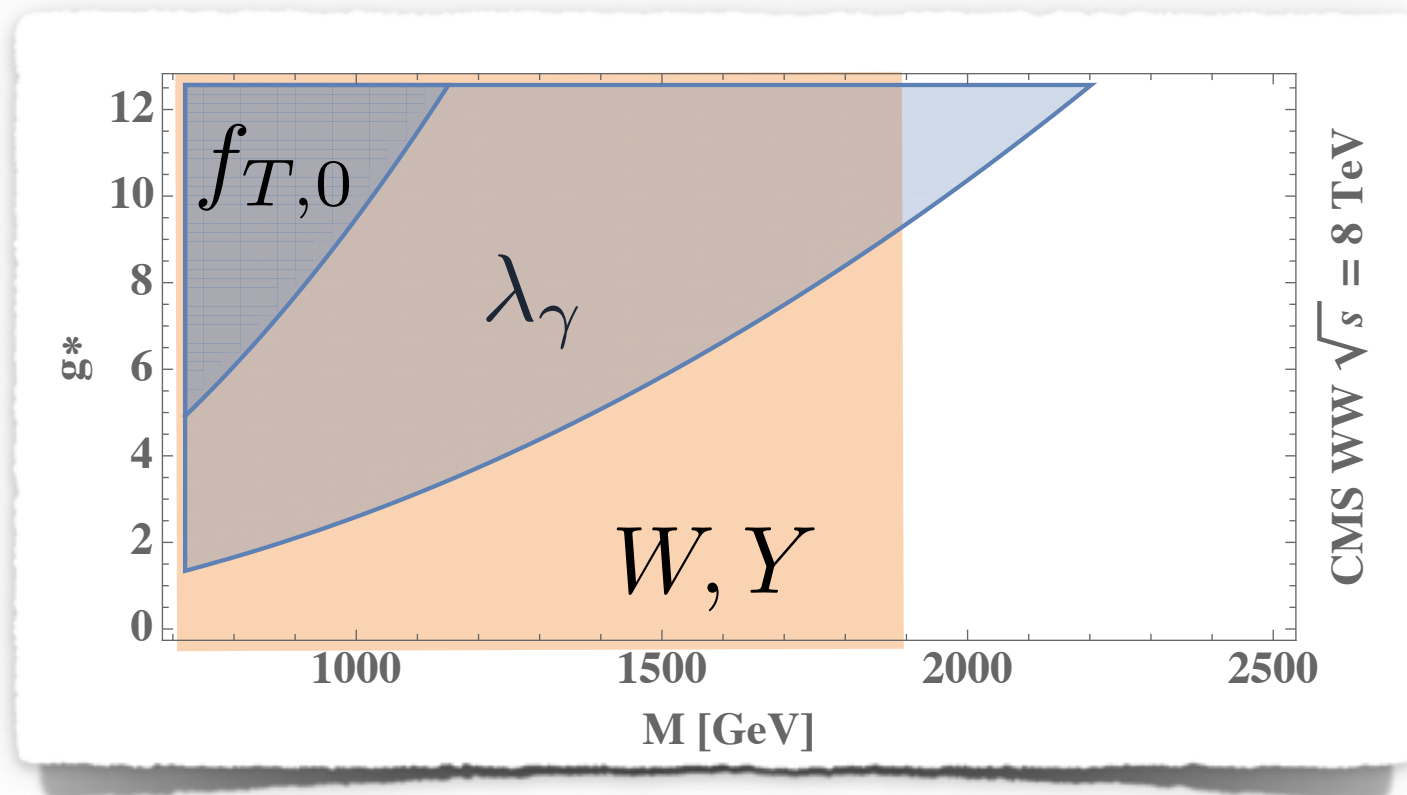
Barbieri, Pomarol, Rattazzi, Strumia'04

$\triangleright \frac{g_*}{M^2} \epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^b W^{\rho\mu}_c$

 $\lambda_\gamma \approx g_* \frac{m_W^2}{M^2}$
 $M \gtrsim \sqrt{\frac{g_*}{4\pi}} 2.2 \text{ TeV}$

$\triangleright \frac{g_*^2}{M^4} W_{\mu\nu}^4$

 $f_{T,0} \approx \frac{g_*^2}{M^4}$
 $M \gtrsim \sqrt{\frac{g_*}{4\pi}} 1.1 \text{ TeV}$

So far the only models that motivate these searches consistently

Strong transverse vectors: Implications

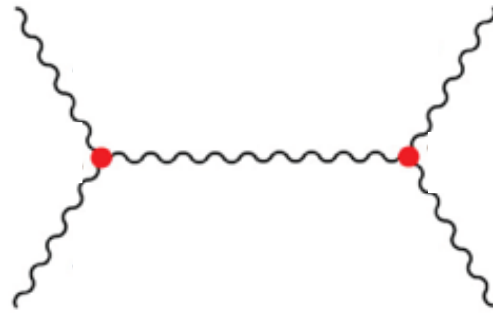


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Strong transverse vectors: Implications

W_T scattering

SM

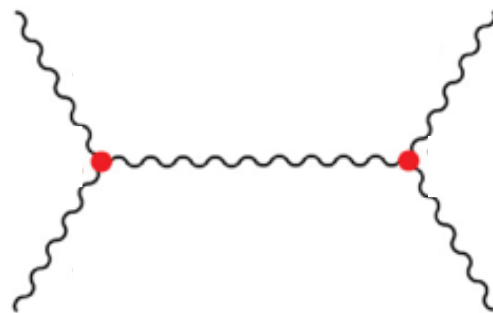


$$\sim g^2$$

Strong transverse vectors: Implications

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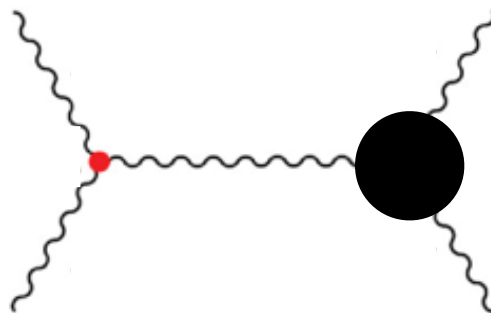
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Strong vectors:

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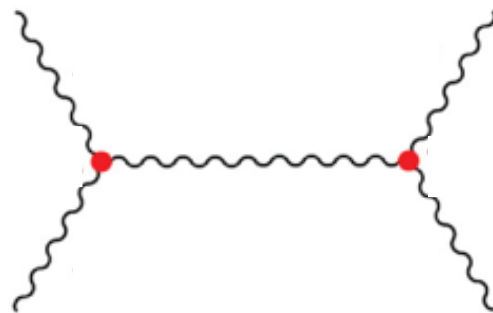
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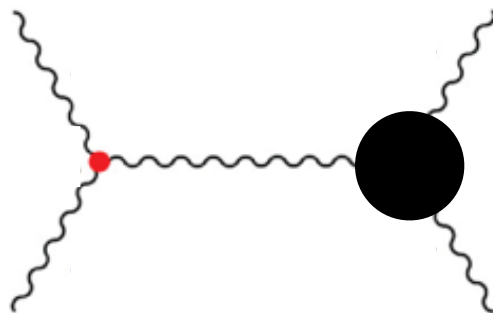
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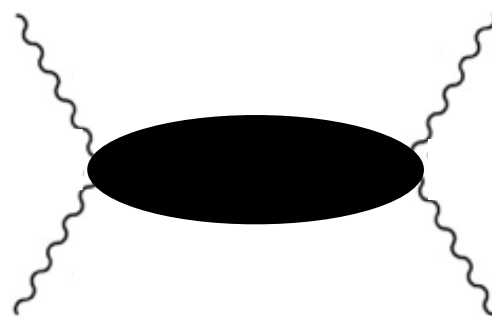
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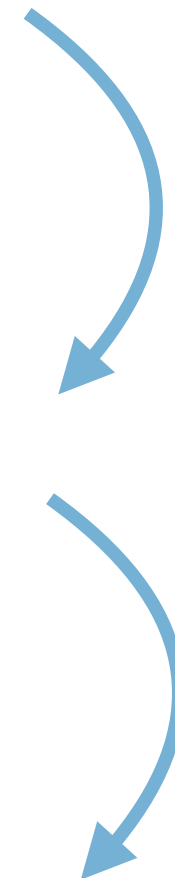
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(NDA: 4-point vertex $\approx$ coupling²)



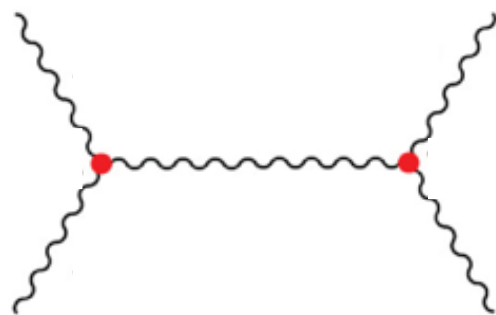
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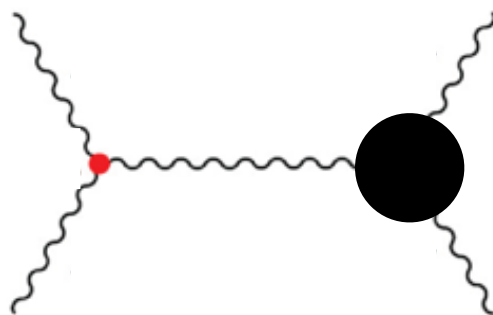
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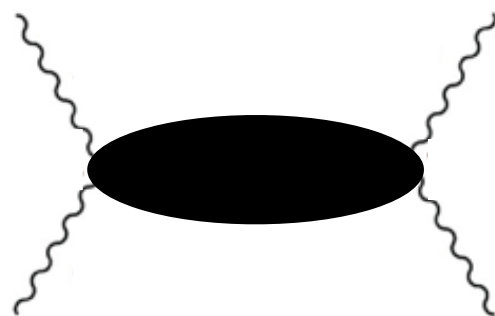
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▶ $\frac{\delta \mathcal{A}_{BSM}}{\mathcal{A}_{SM}} \lesssim 1$ **dimension-6 analysis ok**

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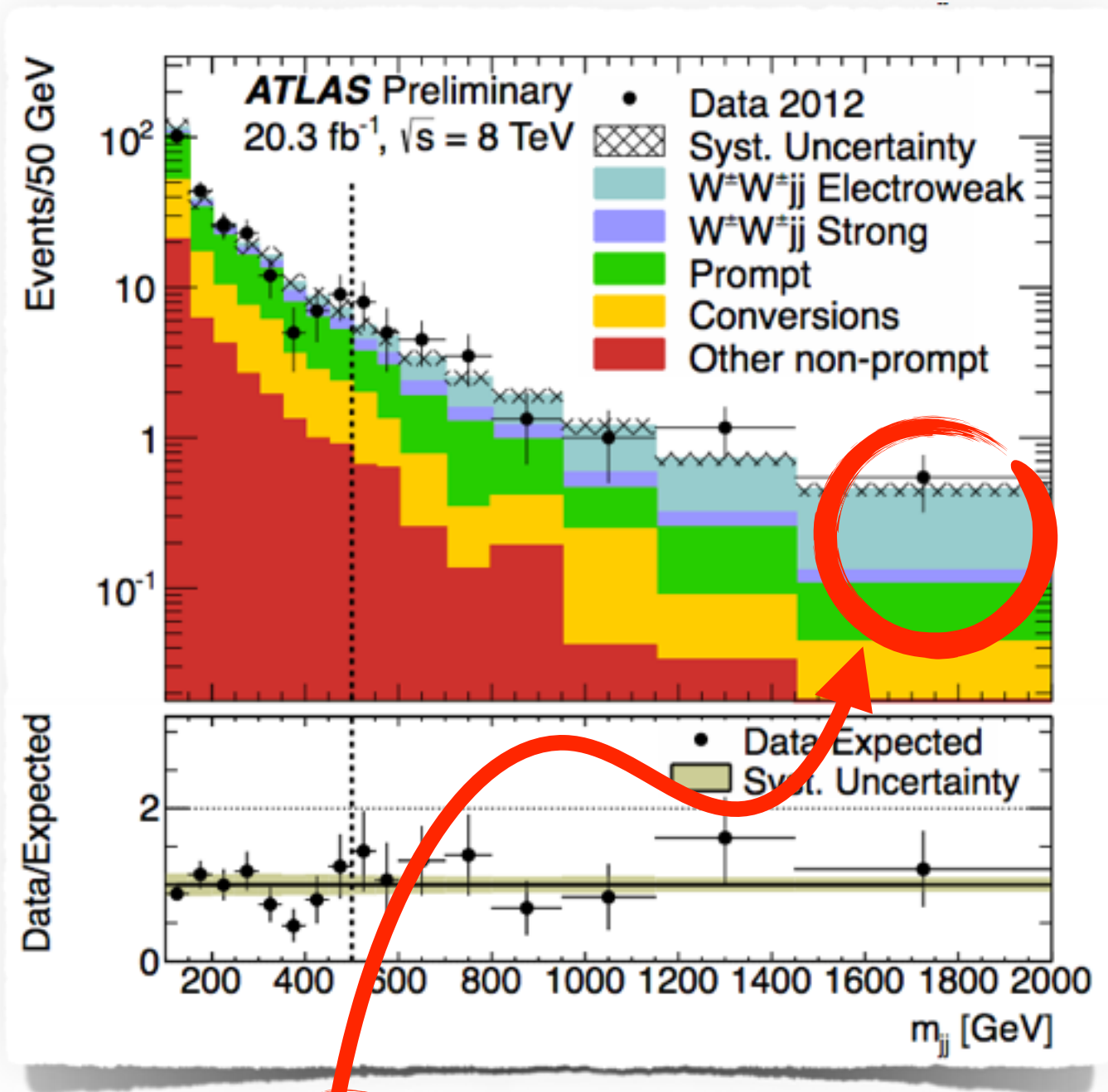
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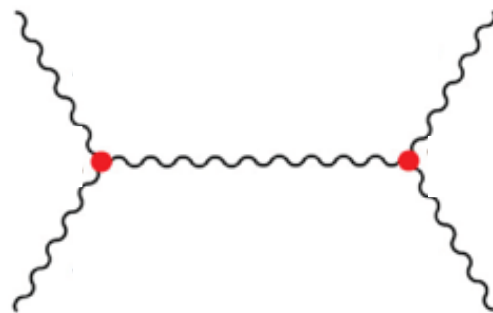
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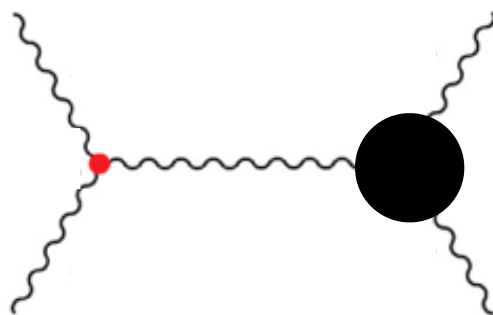
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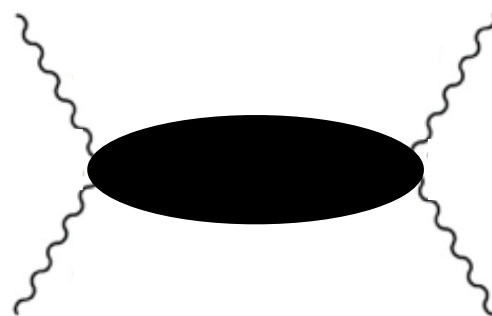
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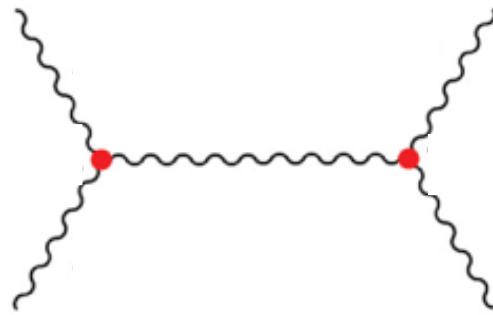
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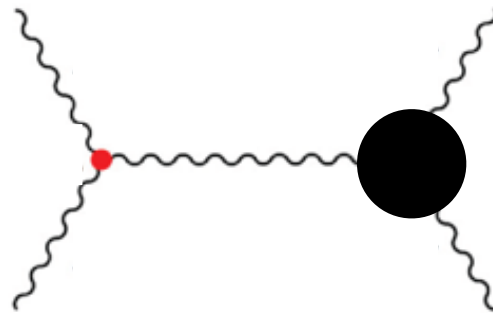
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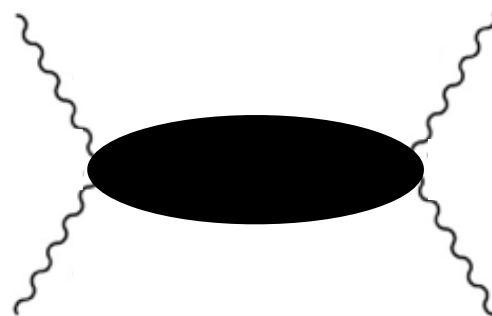
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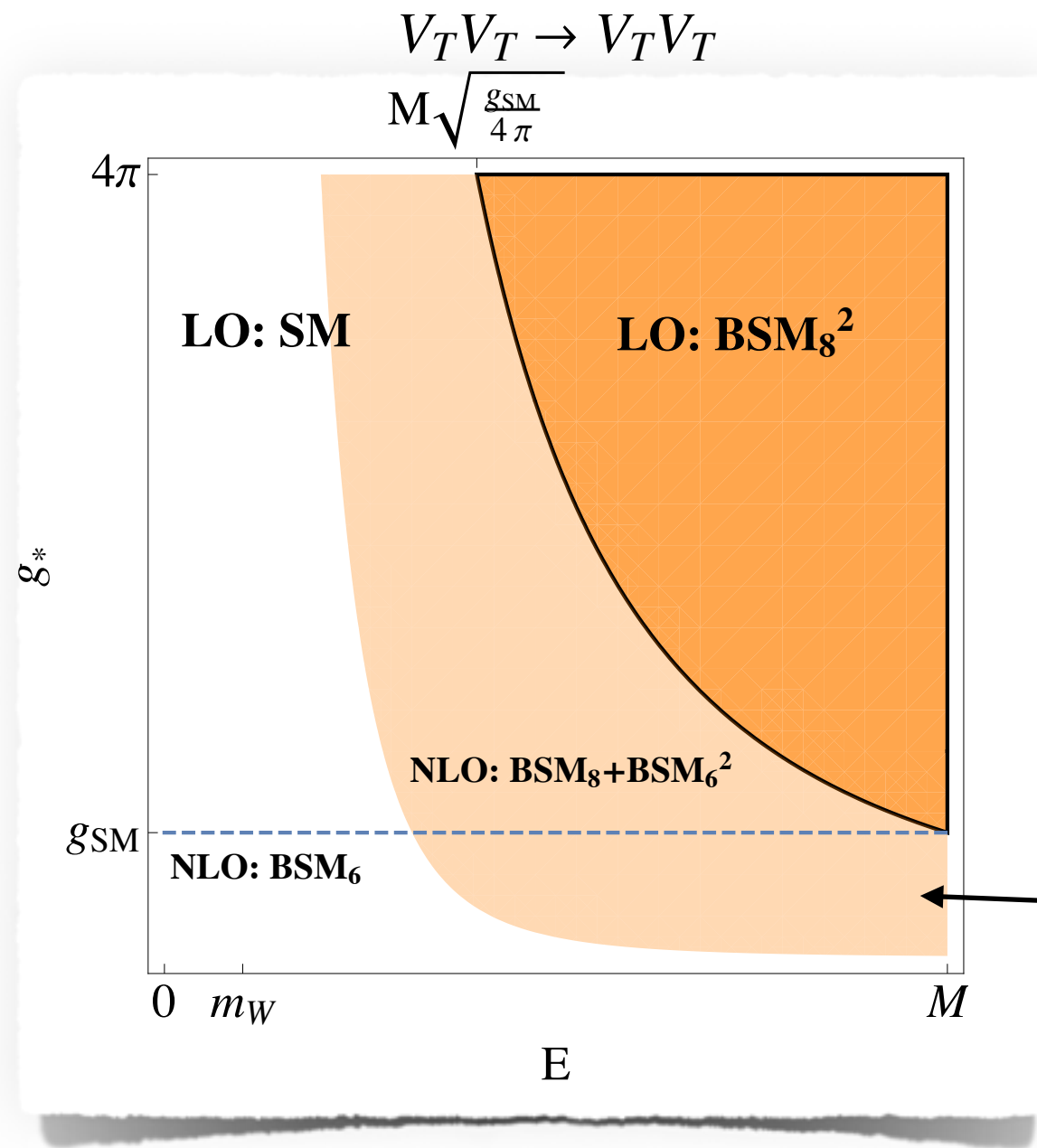
$$\frac{g_*}{g} \frac{E^2}{M^2}$$

► $\frac{\delta \mathcal{A}_{BSM}}{\mathcal{A}_{SM}} \gtrsim 1$ dim-6 not ok! **dim-8 dominate!**

(EFT E/M expansion still valid: dimension-10 small)

Strong transverse vectors: Implications

W_T scattering



Non-interference:
even at low- E $D_8 > D_6$

dim-8 always large for WW scattering

(EFT E/M expansion still valid: dimension-10 small)

Strong transverse vectors: UV?

Through Partial Compositeness?

$$\mathcal{L}_{mix} = \epsilon_A A_\mu J^\mu + \epsilon_F F_{\mu\nu} \mathcal{O}^{\mu\nu}$$

Marginal $\dim[J] = 3$

Relevant if $\dim[\mathcal{O}] < 2$

Unitarity $\dim[\mathcal{O}] \geq 2$

Ferrara,Gatto,Grillo'74,Mack'79

▶ **Either free or irrelevant**

▶ **No CFT or Warped Extra-D model exists**

Strong transverse vectors: UV?

Through P

Strong transverse vectors=**Remedios**

Remedy to motivate (some) LHC searches

No warped UV model

(Remedios the Beauty was not a creature of this world-
Gabriel Garcia Marquez)

► No C

...

Conclusions

LHC/LEP: many €€resources€€ to test transverse vectors

- ▶ Not BSM motivated (non-interference, smallness in ordinary models)

Need structurally robust scenario for strong coupling

- ▶ Remedios:
 - large multipoles/small monopoles
 - deformed symmetry

- ▶ Approximate symmetries can lead to dim-8 domination in $2 \rightarrow 2$ processes



Motivation for some QGC searches

Motivation for dim-8 operator classification

Henning,Lu,Melia,Murayama'15; Lehman,Martin'15

Conclusions

Precision Tests at LHC

