Equilibrium States are determined by their unstable conditionals.

Pablo D. Carrasco

(joint w/ Federico Rodriguez-Hertz)

ICMC-USP visiting ICTP

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ICTP







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$$P_{top}(\varphi) = \sup_{\nu \in \mathcal{P}_{f}(M)} \{h_{\nu}(f) + \int \varphi d\nu\}.$$





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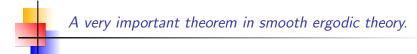
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- Uniqueness of equilibrium states General (but more restrictive) methods (Bowen, or more recently, Climenhaga-Thompson).
- Properties/Description of equilibrium states. (Mixing. Bernoulli. Decay of correlations.)

A very important theorem in smooth ergodic theory.





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• $\varphi \equiv \mathbf{0} \Rightarrow \mu$ is the entropy maximizing measure.



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• $\varphi(x) = -\log |\det df| E_x^u| \Rightarrow \mu$ is the SRB measure (need to assume f is $C^{1+\theta}$)





Main idea of the proof:

Several generalizations of the above theorem exists.

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We discuss now one important example where the available methods fail.

Diagonal actions (on locally homogeneous spaces).

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such that (for some metric) $df|E^u$ expansion, $df|E^s$ contraction and E^c is tangent to the orbits of the action. P. Carrasco, ICMC-USP visiting ICTP Equilibrium States are determined by their unstable conditionals.





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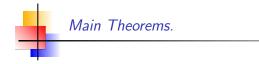


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For a center isometry all bundles $E^s, E^u, E^c, E^{cs} = E^c \oplus E^s, E^{cu} = E^c \oplus E^u$ are integrable to *f*-invariant foliations W^* .





$f: M \to M$ center isometry of class C^2 s.t





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Theorem A [P.C., F. Rodriguez-Hertz] There exist $\mu_{\varphi} \in \mathcal{P}_f(M)$ and families of measures $\mu^u = \{\mu_x^u\}_{x \in M}, \mu^s = \{\mu_x^s\}_{x \in M}, \mu^{cu} = \{\mu_x^{cu}\}_{x \in M}, \mu^{cs} = \{\mu_x^{cs}\}_{x \in M}$ satisfying the following.



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- 1. The probability μ_{φ} is an equilibrium state for the potential φ .
- For every x ∈ M the measure μ^σ, σ ∈ {u, s, cu, cs} is a Radon measure on W^σ(x) which is positive on relatively open sets, and y ∈ W^σ(x) implies μ^σ_x = μ^σ_y.





3. If ξ is a measurable partition that refines the partition by unstable (stable) leaves then the conditionals $(\mu_{\varphi})_x^{\xi}$ of μ_{φ} are equivalent to μ_x^u (resp. μ_x^s) for $\mu_{\varphi} - a.e.(x)$.



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- 4. For every $\epsilon > 0$ sufficiently small, for every $x \in M$ the measure $\mu_{\varphi}|D(x;\epsilon)$ has product structure with respect to the pair μ_x^u, μ_x^{cs} , i.e. its equivalent to $\mu_x^u \times \mu_x^{cs}$.



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- 5. Given $\epsilon > 0$ there exist $a(\epsilon), b(\epsilon) > 0$ such that if

$$U(x,\epsilon,n) = \{y \in W^u(x,\epsilon) : d(f^jx,f^jy) < \epsilon, j = 0, \ldots, n-1\}$$

then

$$a(\epsilon) \leq rac{\mu_x^u(U(x,\epsilon,n)))}{e^{S_n arphi(x) - n P_{top}(arphi)}} \leq b(\epsilon).$$



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Work in progress: Uniqueness also holds in the homogeneous examples (Weyl Chambers' flow).





In the general setting, an SRB measure is a invariant measure whose unstable conditionals are absolutely continuous with respect to Lebesgue.





• SRB measures exist (Sinai-Pesin).



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- \bullet Ledrappier-Young: μ is an SRB if and only if

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Implicit: Unstable manifolds coincide with Pesin's unstable manifolds.

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The families μ^{u},μ^{s} provide the reference measures to which one can compare.



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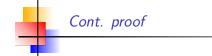
Theorem

For every $x \in M$

1.
$$\mu_{f_X}^{\sigma} = e^{P_{top}(\varphi) - \varphi} f_* \mu_X^{\sigma} \quad \sigma \in \{u, cu\}.$$

2. $\mu_{f_X}^{\sigma} = e^{\varphi - P_{top}(\varphi)} f_* \mu_X^{\sigma} \quad \sigma \in \{s, cs\}.$

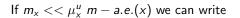




If $m_x << \mu^u_x \ m-a.e.(x)$ we can write

 $dm_x = \rho d\mu_x^u$

where ρ is measurable on M.



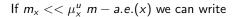
Cont. proof

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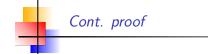
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is constant on the atoms of ξ . From here one deduces that

$$y \in \xi(x) \Rightarrow rac{
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ho(x)} = \prod_{k=1}^{\infty} rac{e^{arphi \circ f^{-k}(y)}}{e^{arphi \circ f^{-k}(x)}} = \Delta_x(y).$$

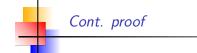
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$$ho(y) = rac{\Delta_x(y)}{L(x)}, \quad y \in \xi(x)$$

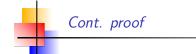
with $L(x) = \int_{\xi(x)} \Delta_x(y) d\mu_x^u(y)$.



$$\rho(y) = rac{\Delta_x(y)}{L(x)}, \quad y \in \xi(x)$$

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Define a measure ν by requiring $\nu = m$ on \mathcal{B}_{ξ} and such its conditionals on ξ are given by $d\nu_x = \rho d\mu_x$.



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It is enough to show $m = \nu$ on every $\mathcal{B}_{f^{-n}\xi}, n \ge 0$.



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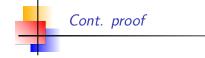
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It is enough to show $m = \nu$ on every $\mathcal{B}_{f^{-n}\xi}$, $n \ge 0$. An induction argument shows that after proving $m = \nu$ on $\mathcal{B}_{f^{-1}\xi}$ we are done.

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$$q(x) = \nu_x(f^{-1}\xi(x)) \Rightarrow$$

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Cont. proof

$$q(x) = \frac{1}{L(x)} \int_{f^{-1}(\xi(fx))} \Delta_x(f^{-1}fy) \frac{e^{P-\varphi(y)}}{e^{P-\varphi(f^{-1}fy)}} d\mu_x^u$$

= $\frac{1}{L(x)} \int_{\xi(fx)} \Delta_x(f^{-1}z) e^{\varphi(f^{-1}z)-P} d\mu_{fx}^u(z) = \frac{L(fx)}{L(x)} e^{\varphi(x)-P} \le 1$

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$$\frac{L(fx)}{L(x)} \leq e^{\varphi(x)-P} \in L^1(M,m)$$

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$$\frac{L(fx)}{L(x)} \le e^{\varphi(x)-P} \in L^1(M,m)$$
$$\int \log \frac{L \circ f}{L} dm = 0 \Rightarrow \int -\log q(x) dm(x) = P - \int \varphi dm.$$

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Since ξ is a SPLY and *m* equilibrium state, Ledrappier, Young tell us that

$$\Rightarrow P - \int \varphi dm = h_m(f) = H(f^{-1}\xi|\xi) = \int -\log m_x(f^{-1}\xi(x))dm(x)$$

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$$\Rightarrow \int -\log \frac{\nu_x(f^{-1}\xi(x))}{m_x(f^{-1}\xi(x))} dm(x) = 0$$

Using strict convexity of the logarithm function plus working (carefully) with the partitions $f^{-1}\xi(x)|\xi(x)$ we deduce

$$\left.\frac{d\nu}{dm}\right|_{f^{-1}\xi}(x)=1 \quad m-a.e.x \in M.$$

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Since ξ is a SPLY and *m* equilibrium state, Ledrappier, Young tell us that

$$\Rightarrow P - \int \varphi dm = h_m(f) = H(f^{-1}\xi|\xi) = \int -\log m_x(f^{-1}\xi(x))dm(x)$$

$$\Rightarrow \int -\log \frac{\nu_x(f^{-1}\xi(x))}{m_x(f^{-1}\xi(x))} dm(x) = 0$$

Using strict convexity of the logarithm function plus working (carefully) with the partitions $f^{-1}\xi(x)|\xi(x)$ we deduce

$$\left.\frac{d\nu}{dm}\right|_{f^{-1}\xi}(x)=1 \quad m-a.e.x \in M.$$

and afterwards, $m_x = \nu_x$ on $\xi(x)$ for ALL atoms, as we wanted to show.

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Grazie!!!