

Measurement of the Multi-TeV Neutrino-Nucleon Cross Section Via Earth Absorption

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PANE

ICTP, Trieste, May 27, 2018

UNIVERSITY OF
COPENHAGEN



Two **seemingly** unrelated questions –

- 1 Where are the most energetic particles coming from?
- 2 What is the structure of matter at the smallest scales?



Neutrinos interactions are weak ...

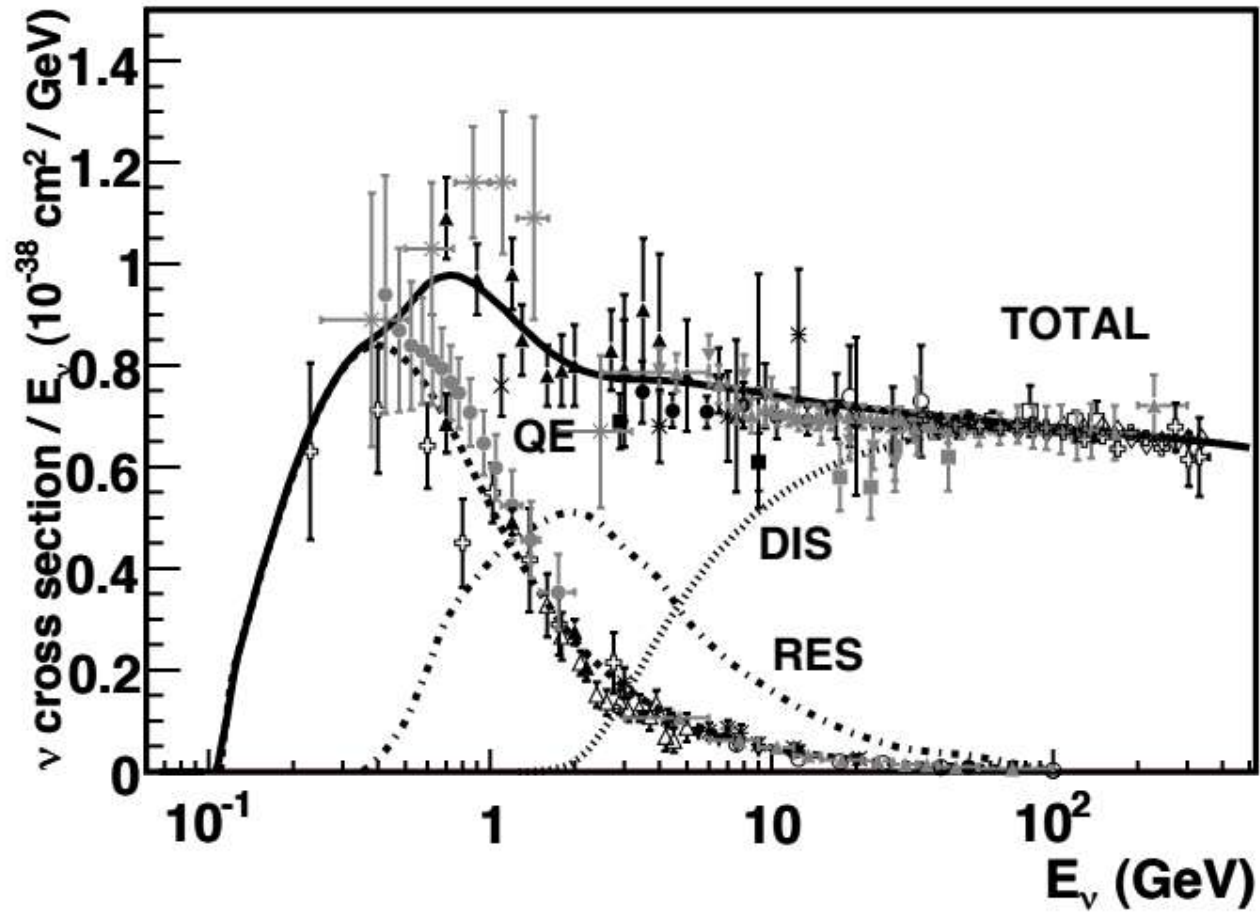
... but we *are* persistent

At center-of-mass energy of 1 GeV:

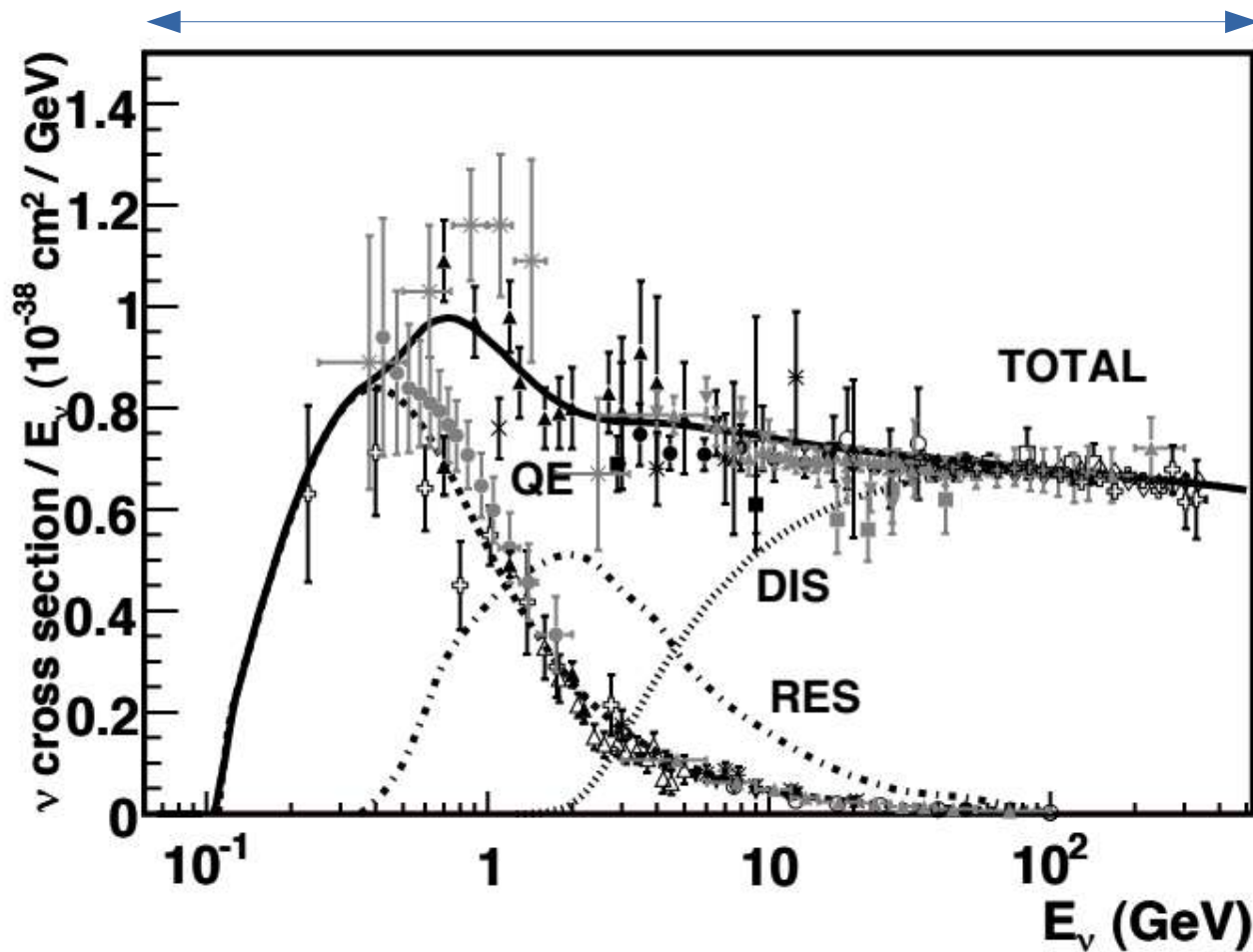
$$\sigma_{pp} \sim 10^{-28} \text{ cm}^2$$

$$\sigma_{\gamma p} \sim 10^{-29} \text{ cm}^2$$

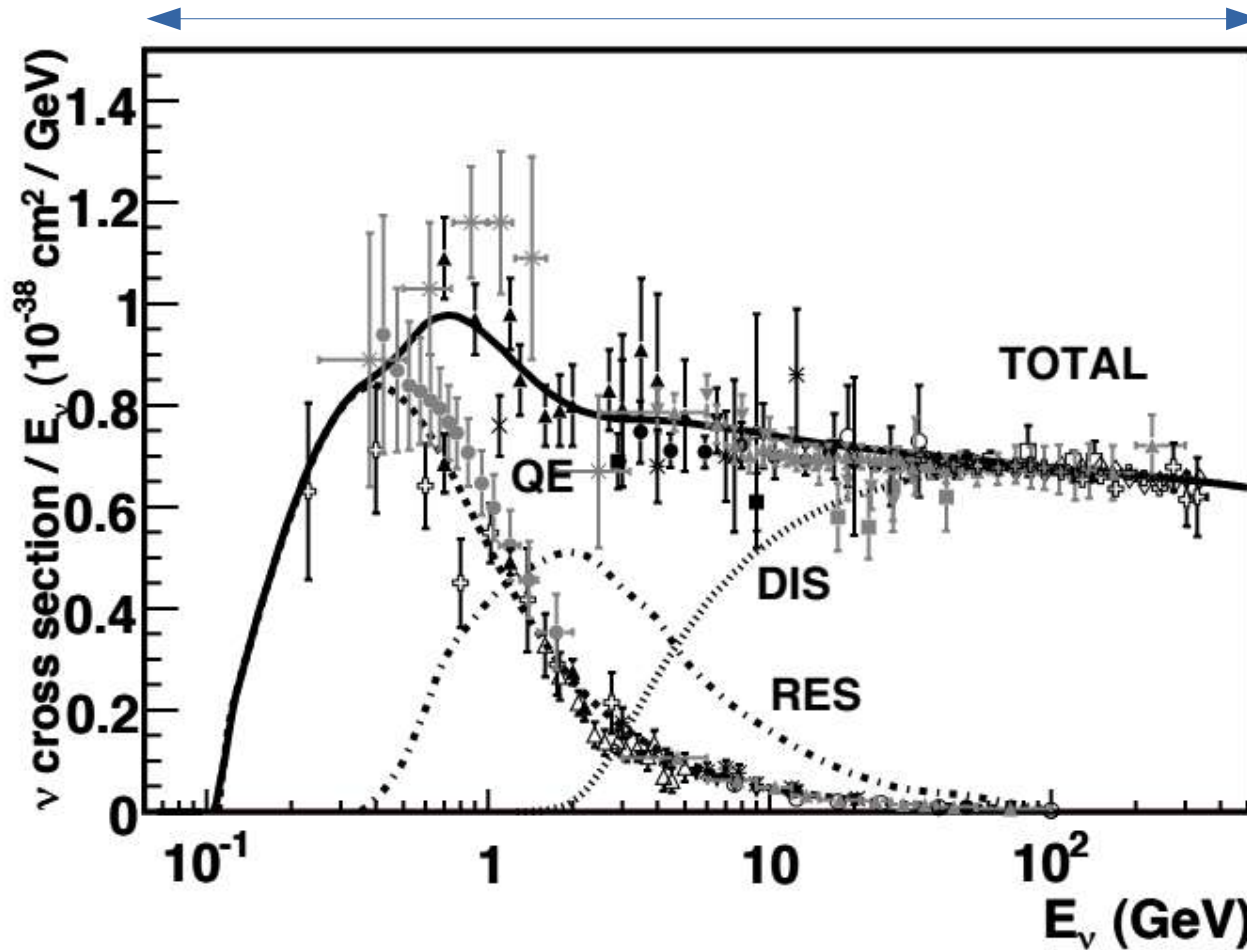
$$\sigma_{\nu p} \sim 10^{-38} \text{ cm}^2$$



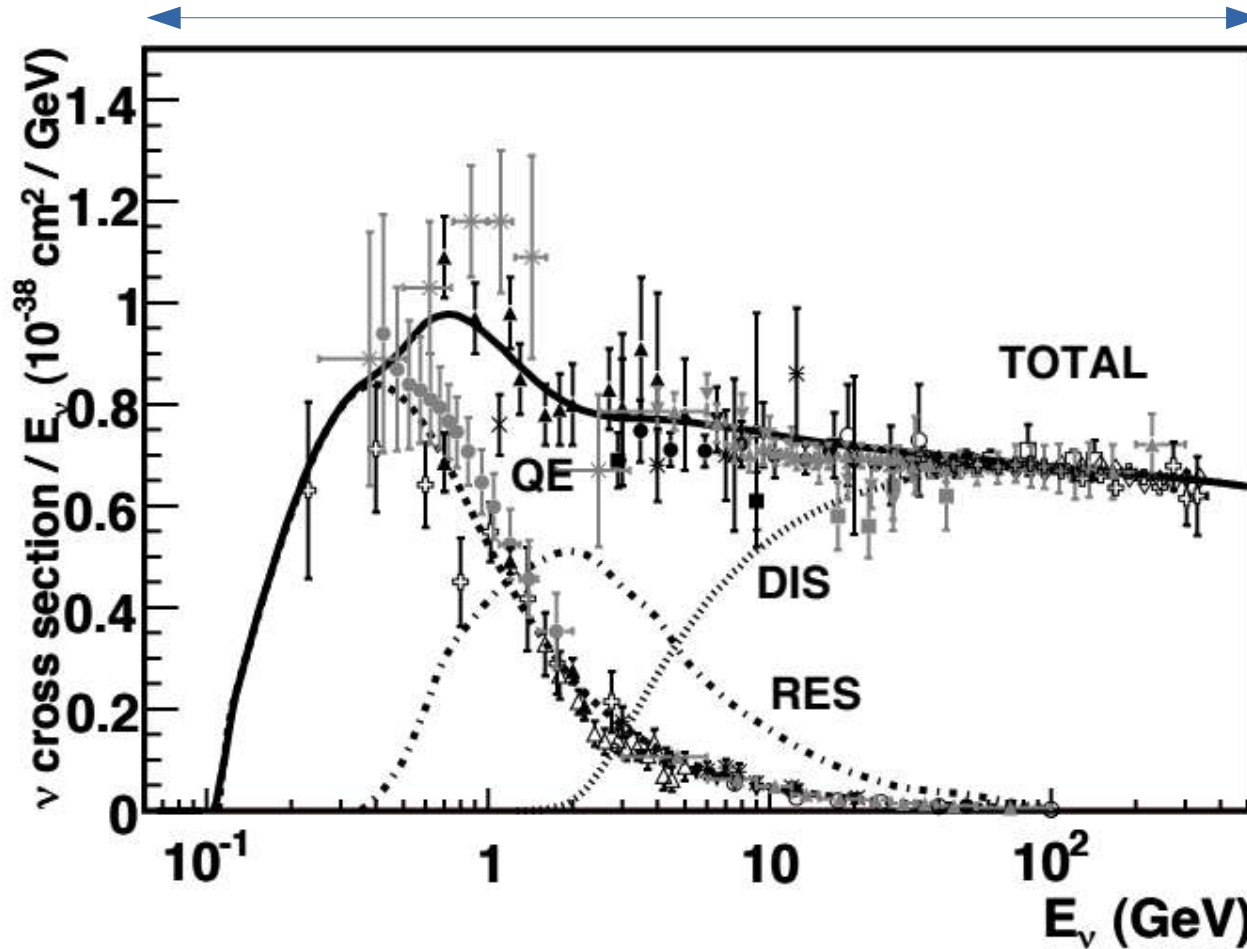
Accelerator experiments



Accelerator experiments



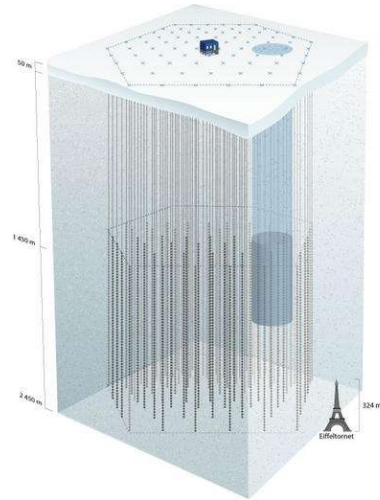
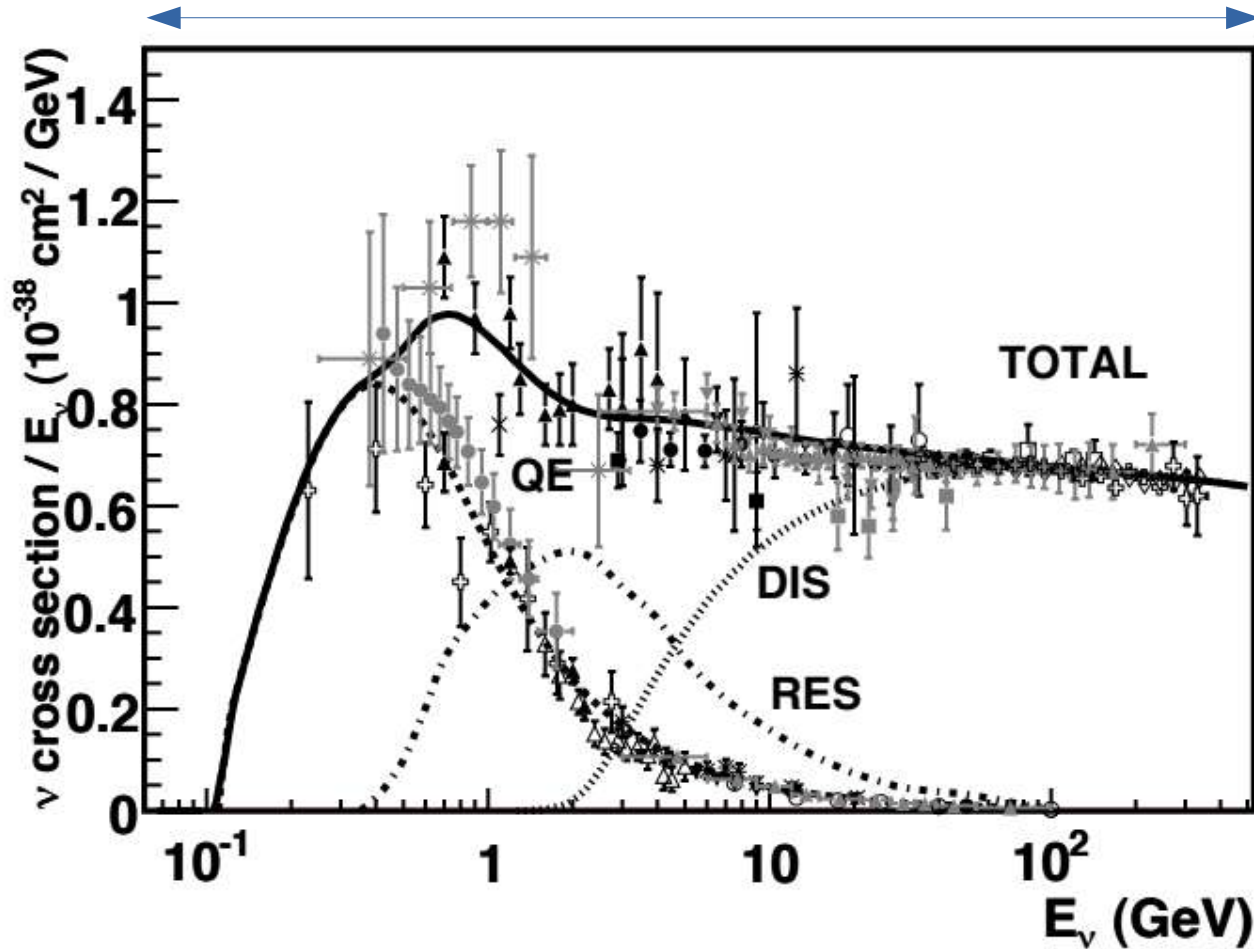
Accelerator experiments

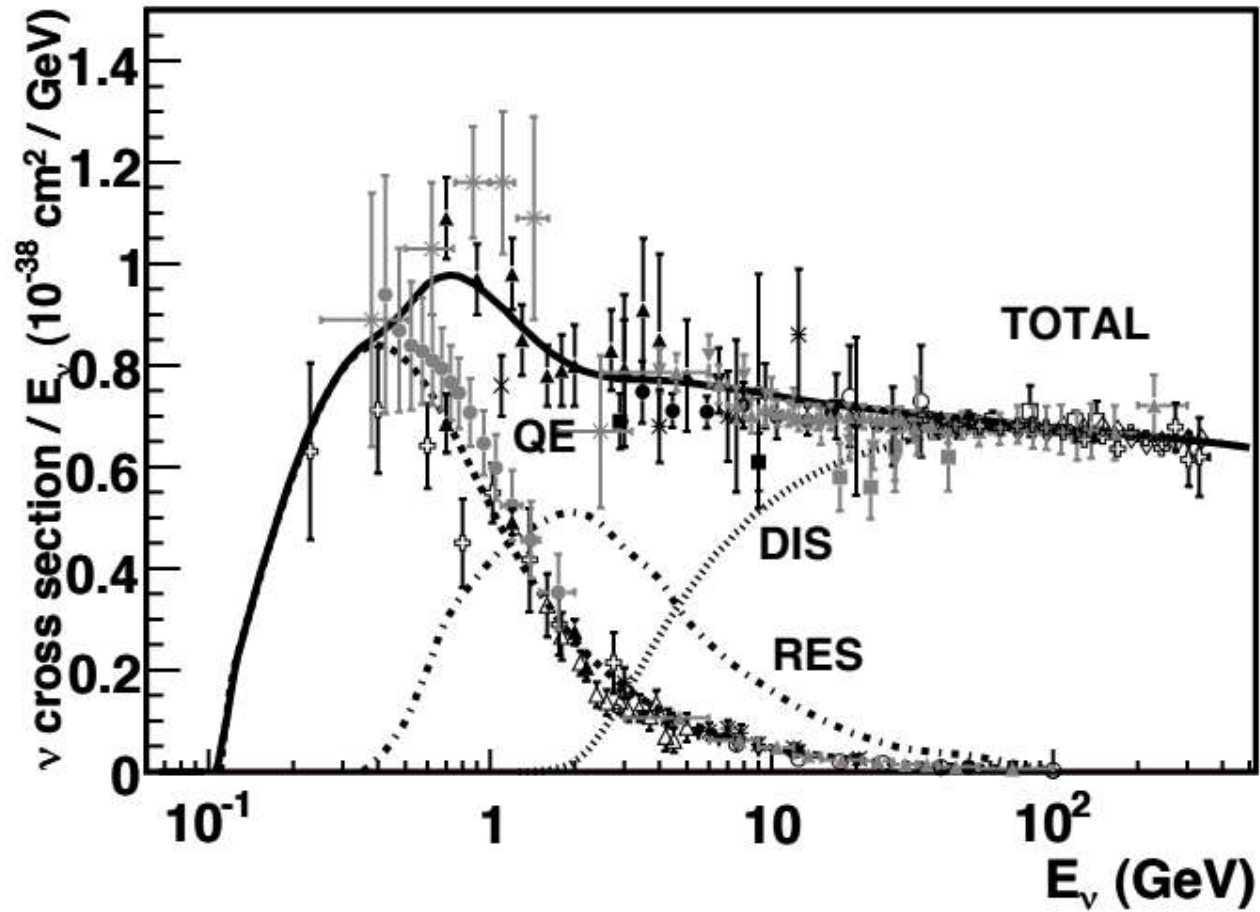


←
One recent
measurement
(COHERENT)

→
No
measurements
... until now!

Accelerator experiments



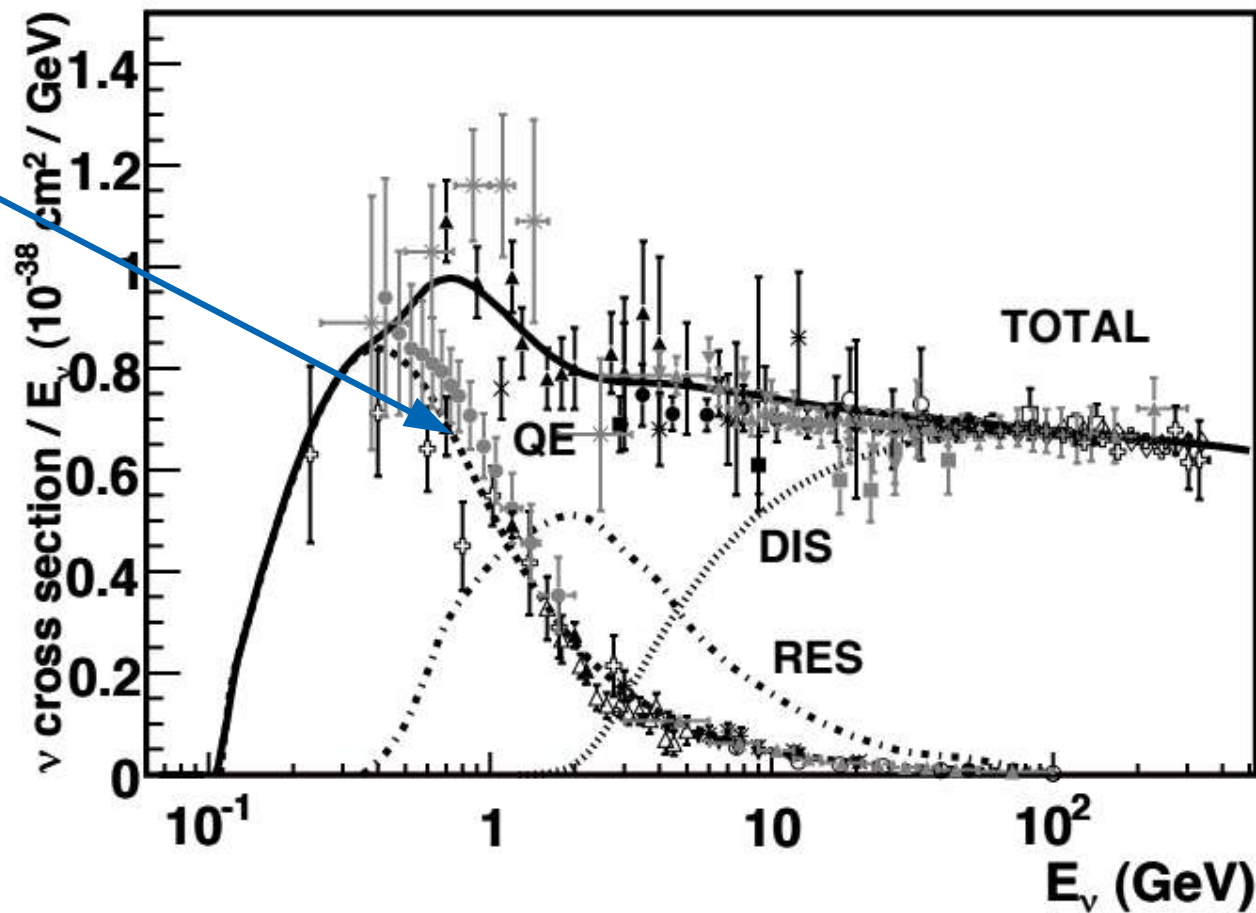


Quasi-elastic

scattering:

$$\nu_l + n \rightarrow l^- + p$$

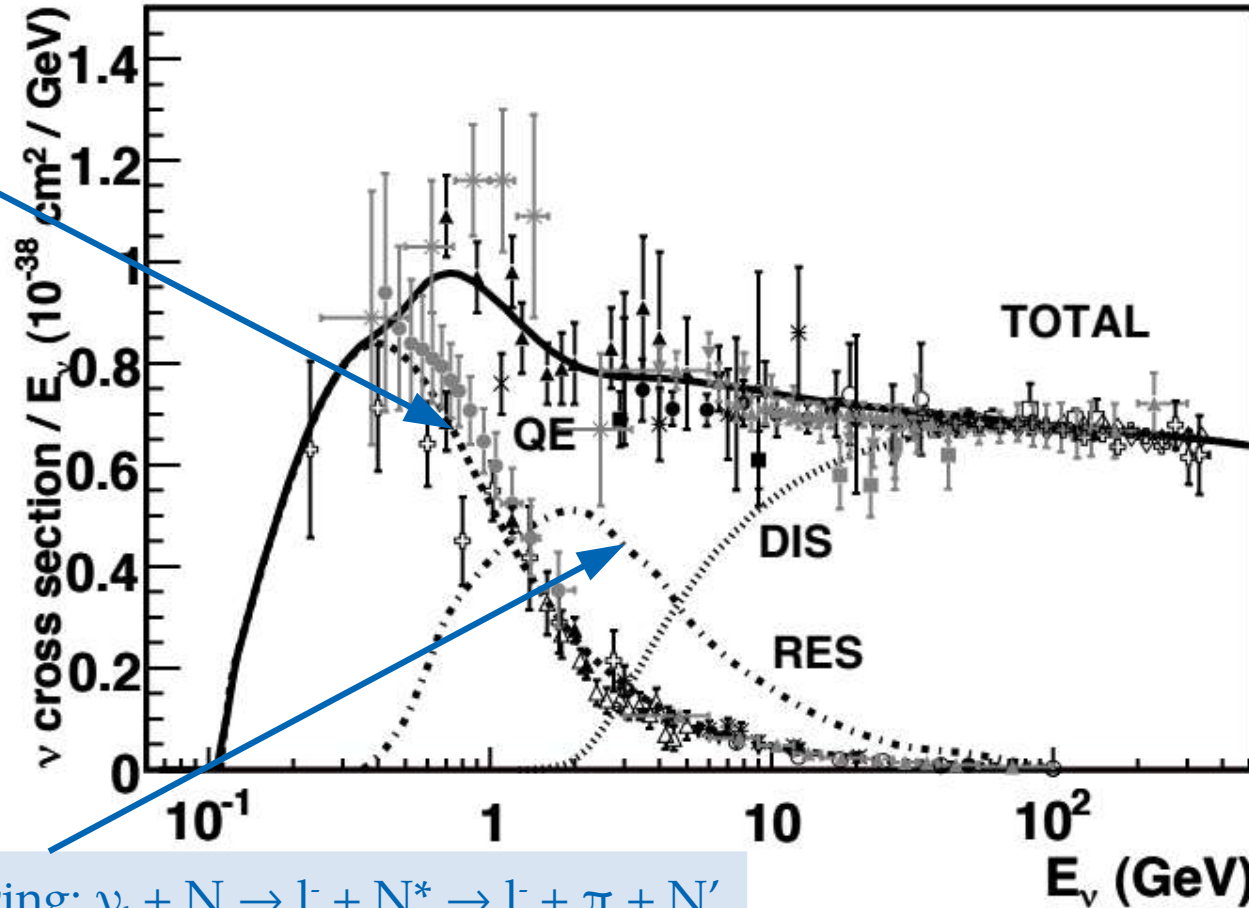
$$\bar{\nu}_l + p \rightarrow l^+ + n$$



Quasi-elastic
scattering:

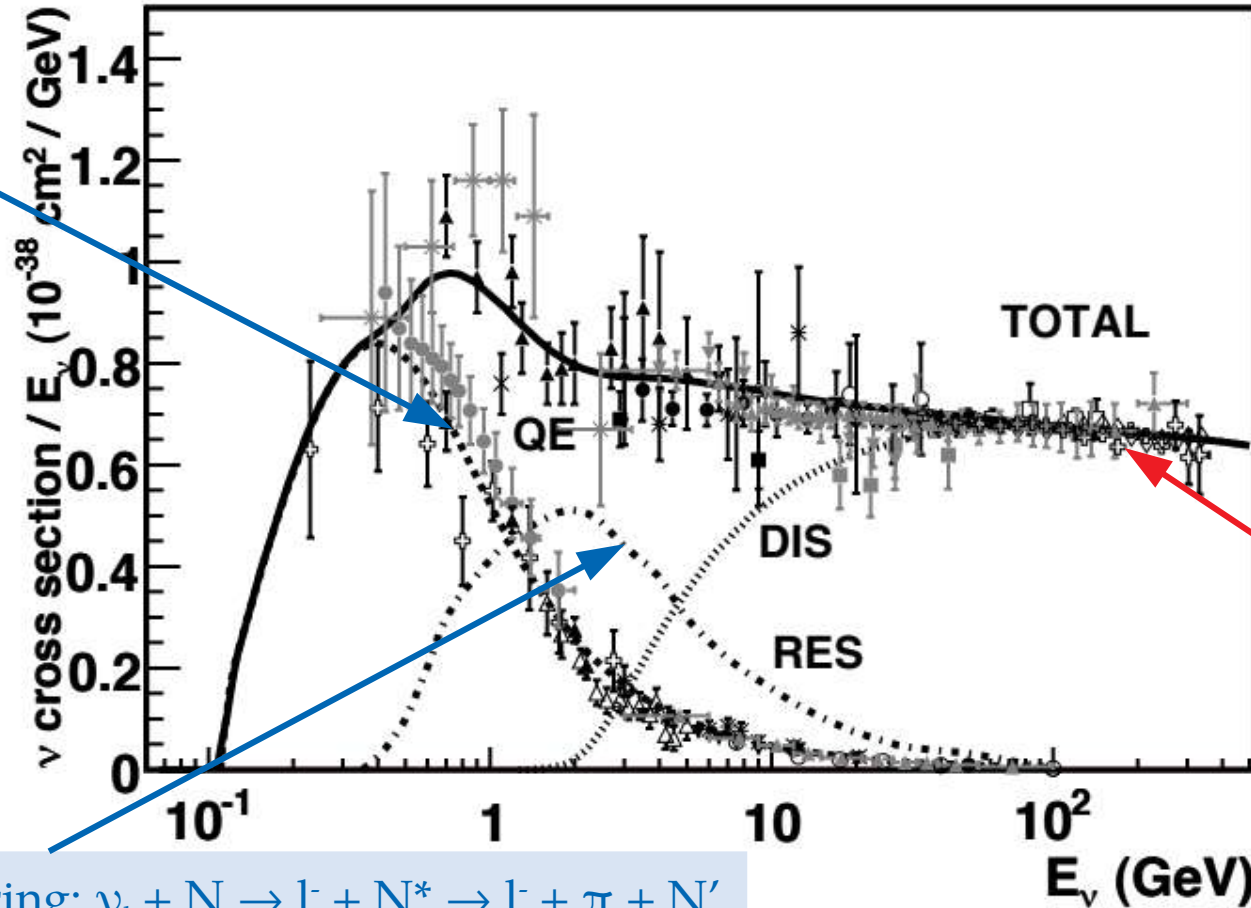
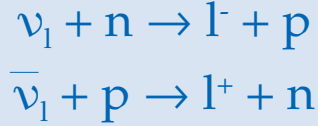
$$\nu_l + n \rightarrow l^- + p$$

$$\bar{\nu}_l + p \rightarrow l^+ + n$$

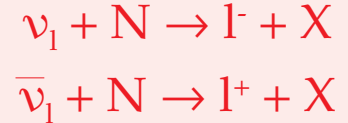


Resonant scattering: $\nu_l + N \rightarrow l^- + N^* \rightarrow l^- + \pi + N'$

Quasi-elastic
scattering:



Deep inelastic
scattering:

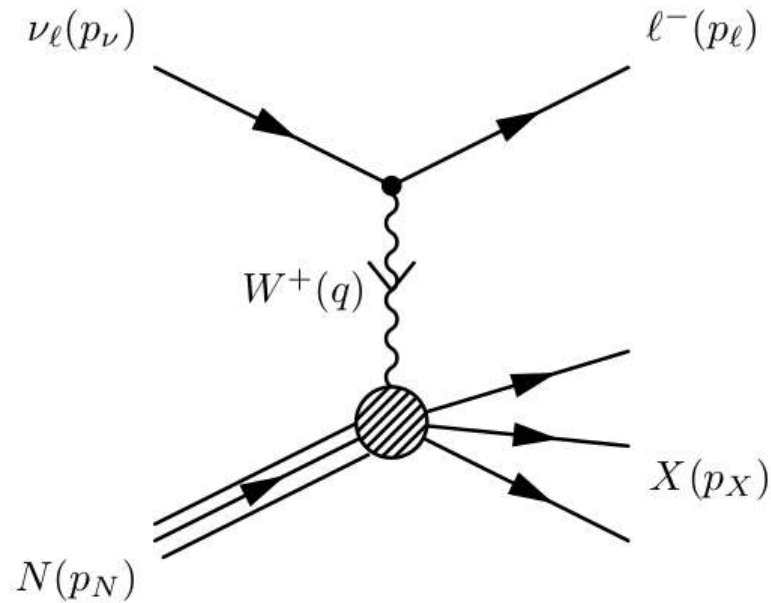


Resonant scattering: $\nu_l + N \rightarrow l^- + N^* \rightarrow l^- + \pi + N'$

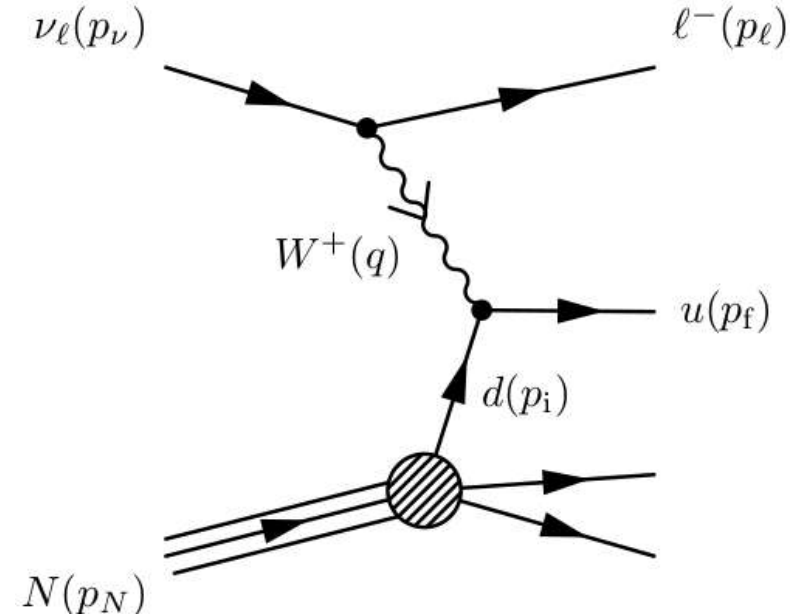
Particle Data Group

How does DIS probe nucleon structure?

What you see

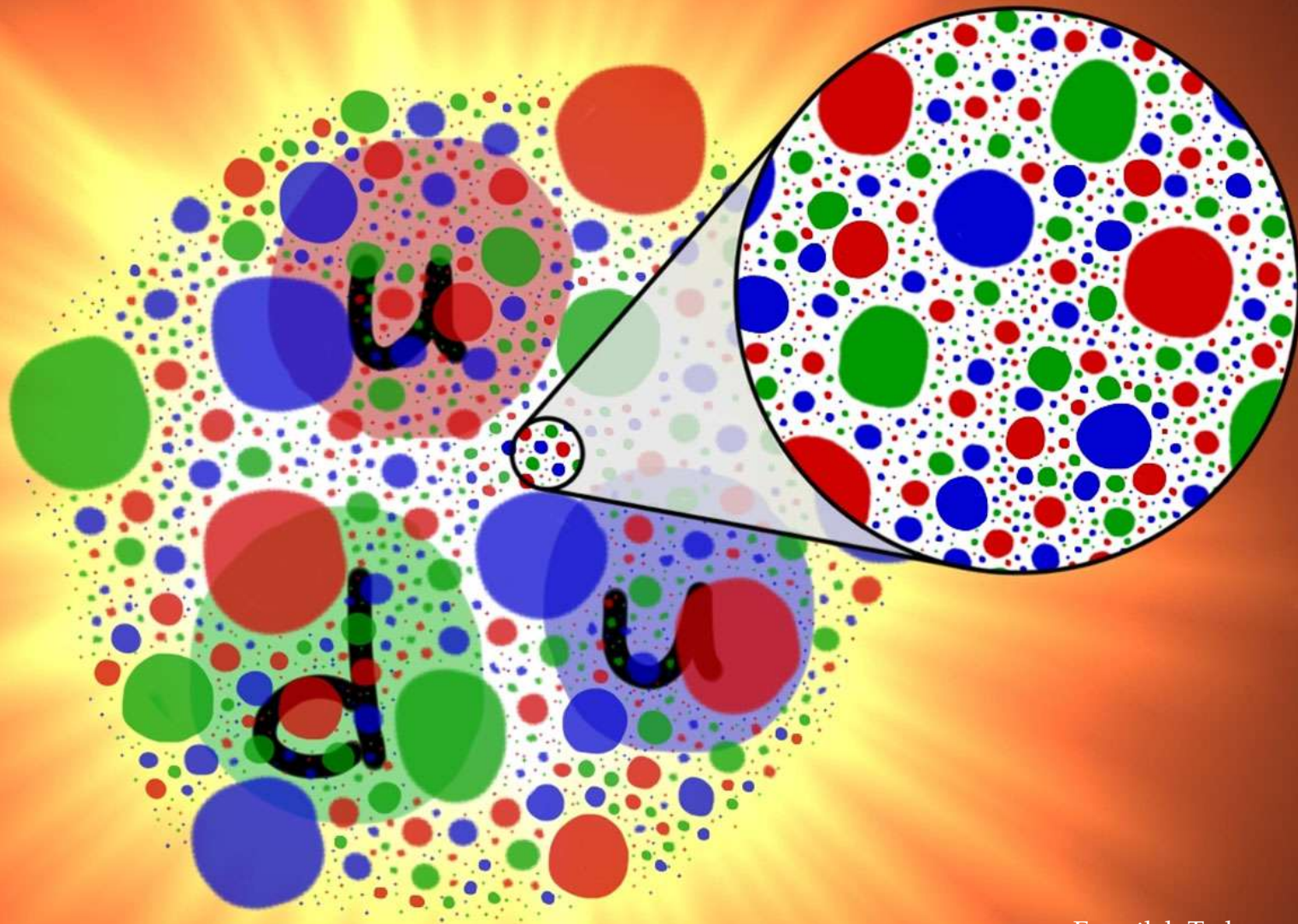


Beneath the hood



(Plus the equivalent neutral-current process (Z-exchange))

Giunti & Kim, *Fundamentals of Neutrino Physics & Astrophysics*



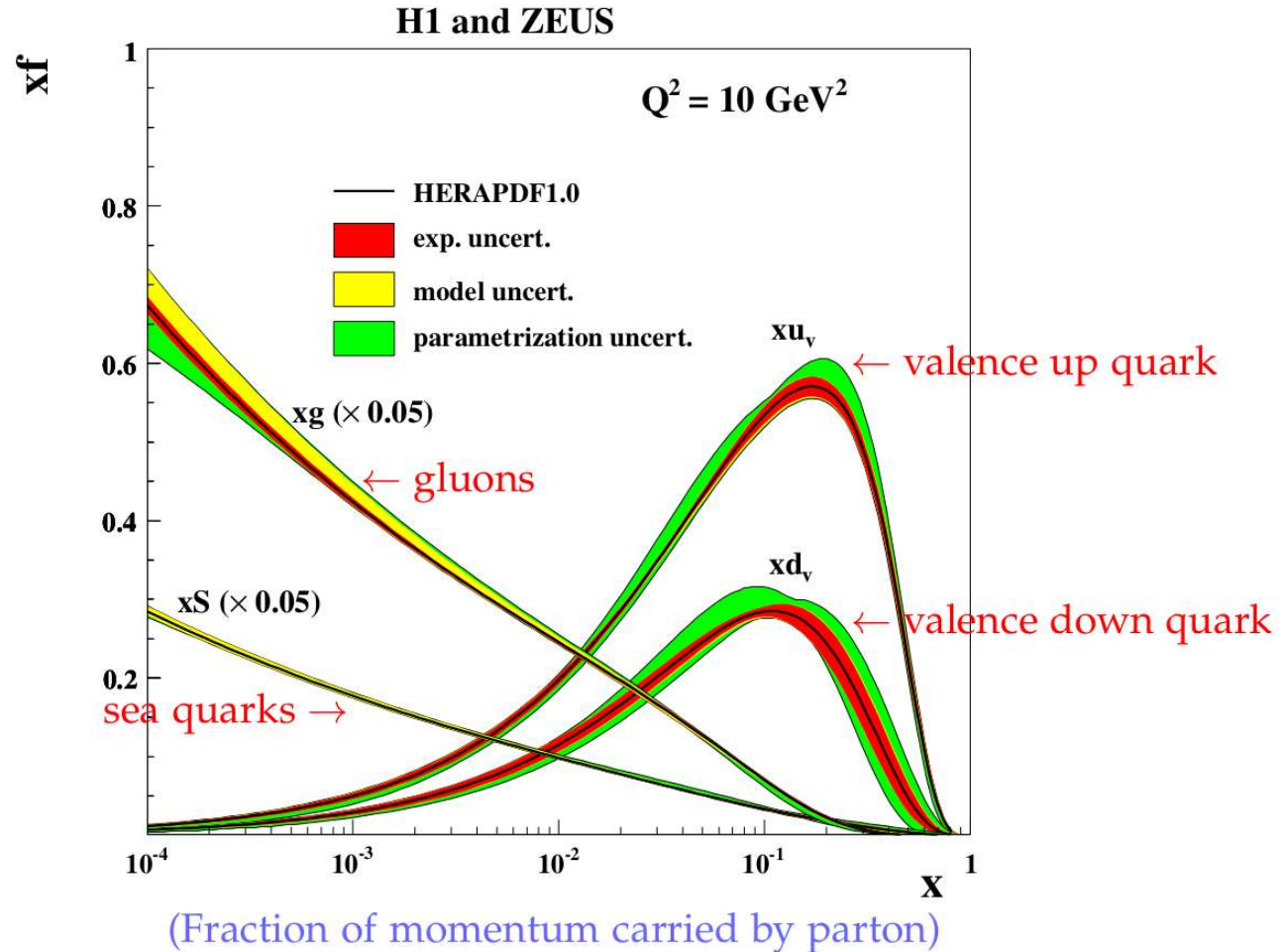
Peeking inside a proton

$$x \sim \frac{m_W}{E_\nu}$$

$$\sim 10^{-4} \text{ for } E_\nu = \text{PeV}$$

← Extrapolation

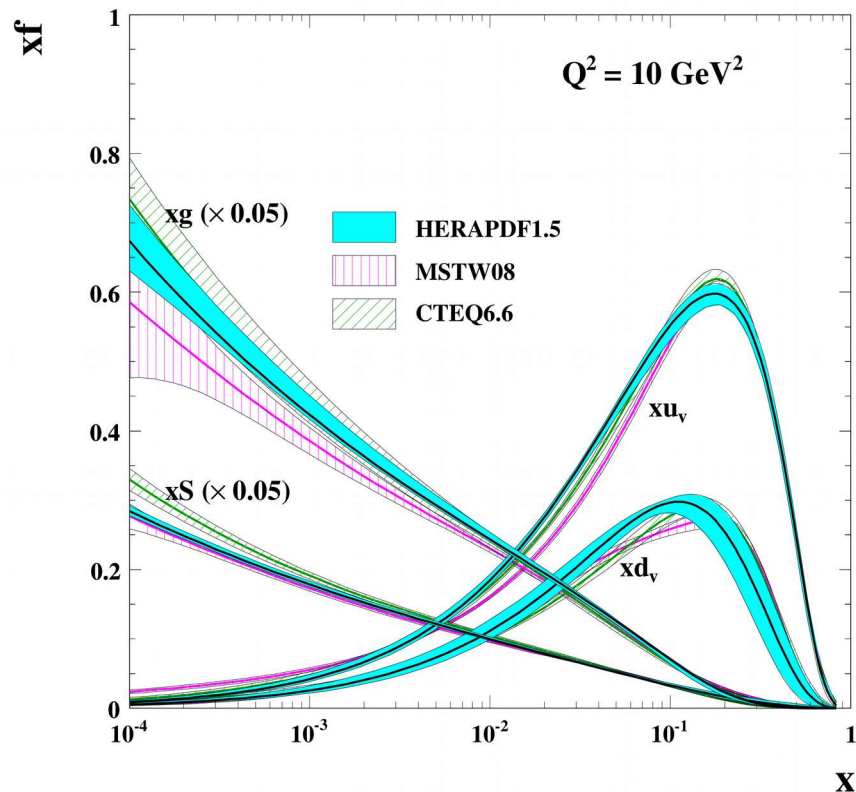
(Parton distribution function)



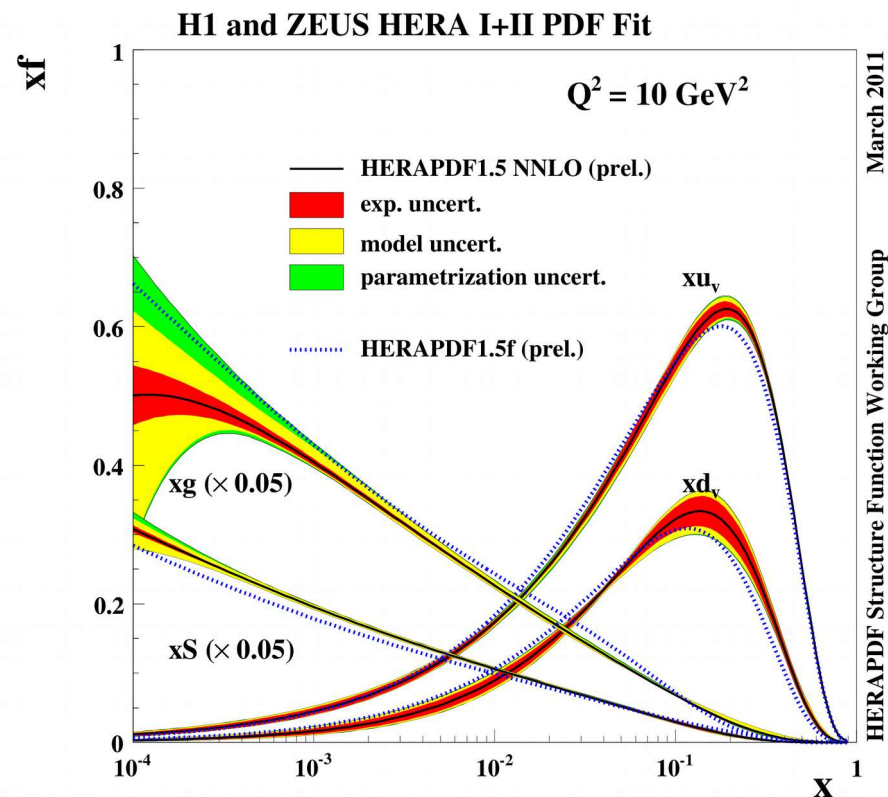
A. COOPER-SARKAR 2012

The world of PDFs is complex

Different fitting groups



Different QCD prescriptions



March 2011

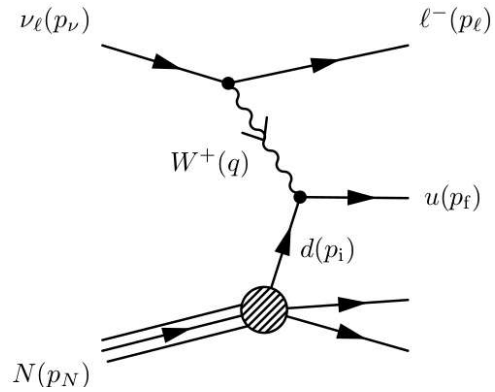
HERAPDF Structure Function Working Group

A. Cooper-Sarkar 2012

Extrapolating the cross section to high energies

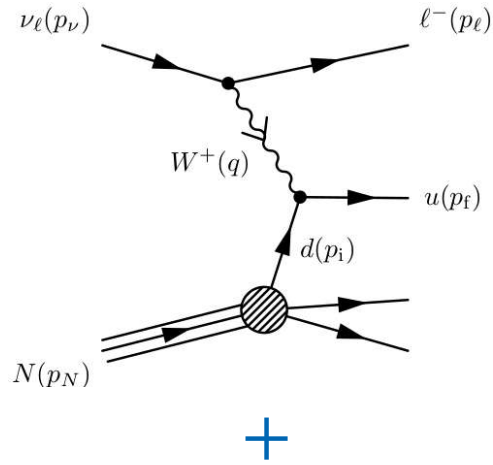
Extrapolating the cross section to high energies

SM



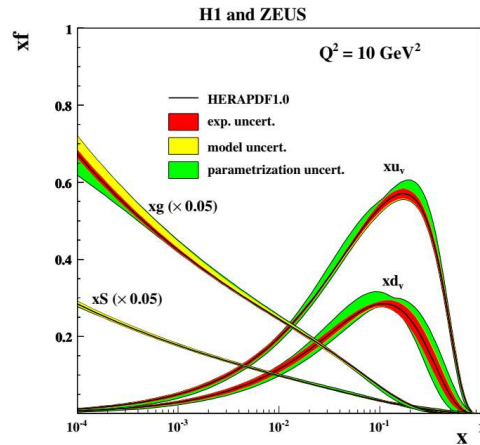
Extrapolating the cross section to high energies

SM



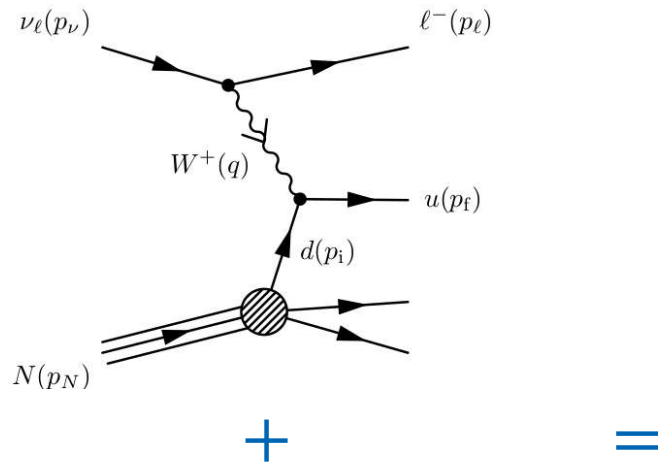
+

PDFs

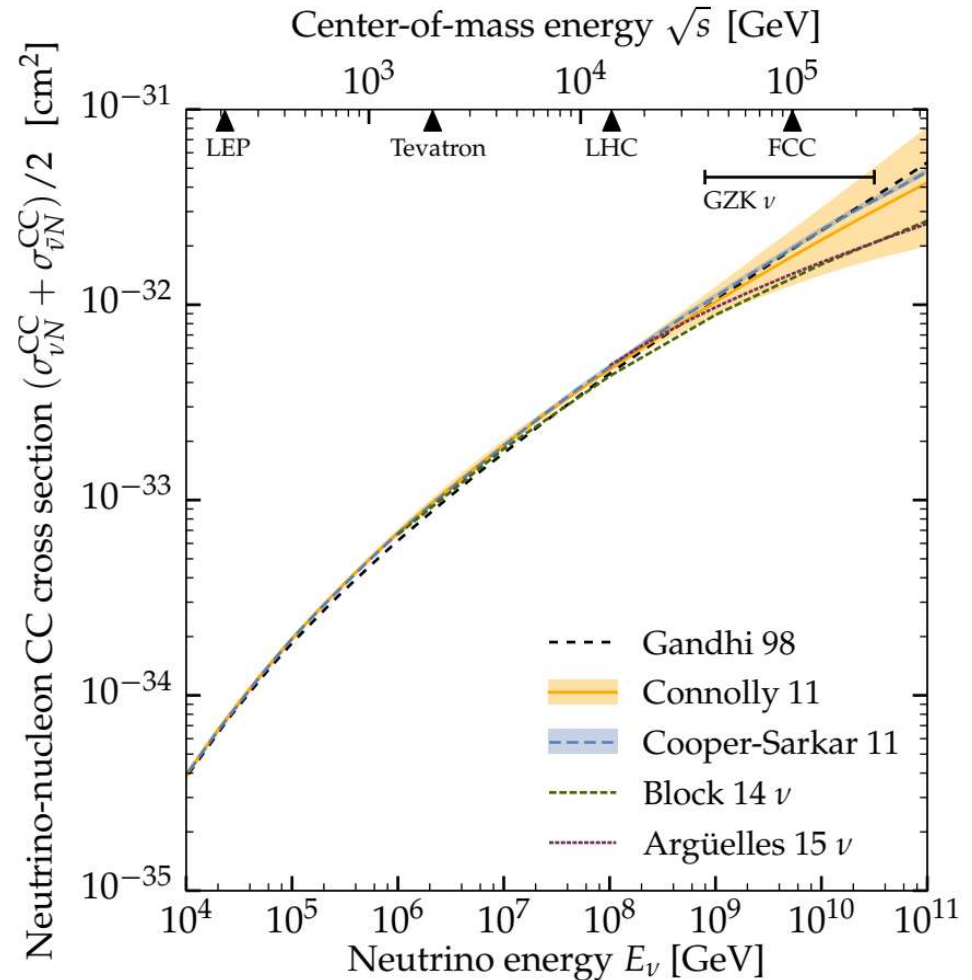
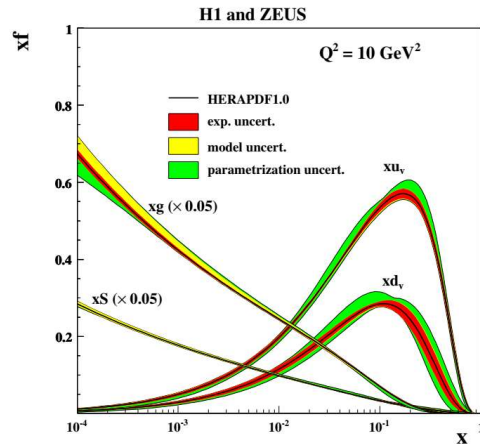


Extrapolating the cross section to high energies

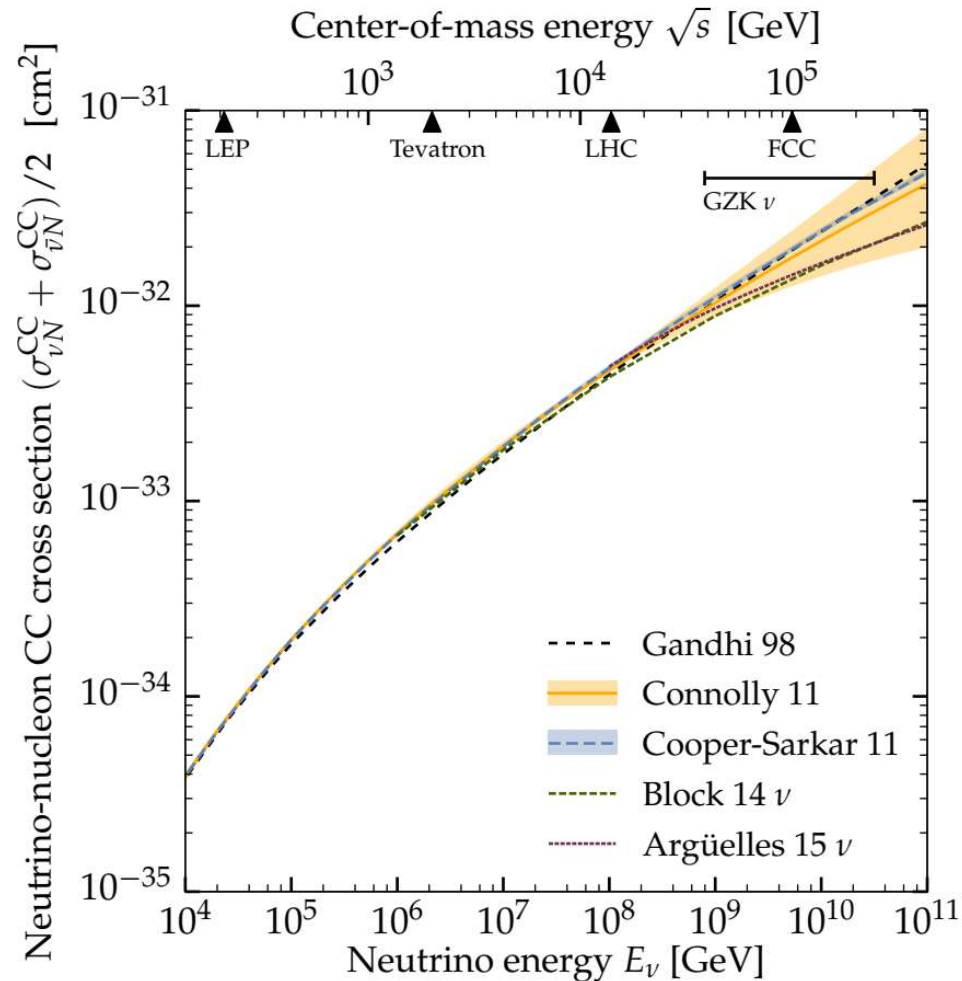
SM



PDFs

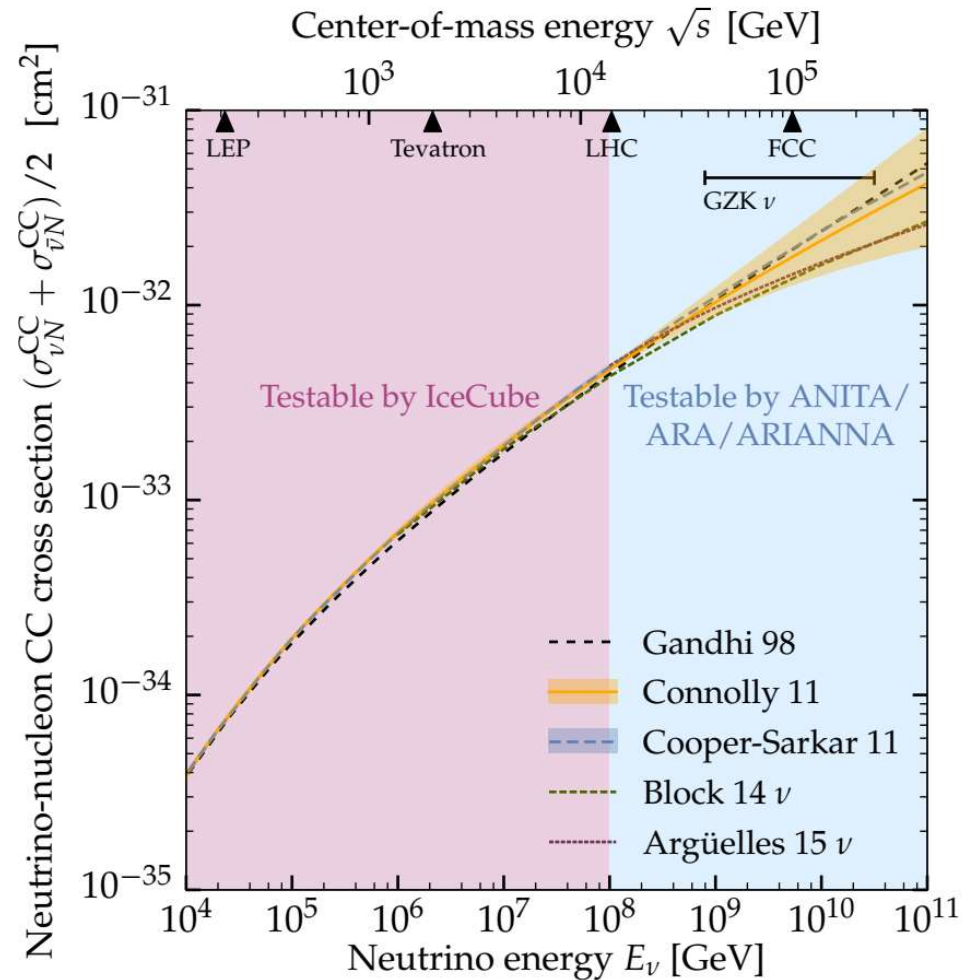


What can we measure *now* and later?



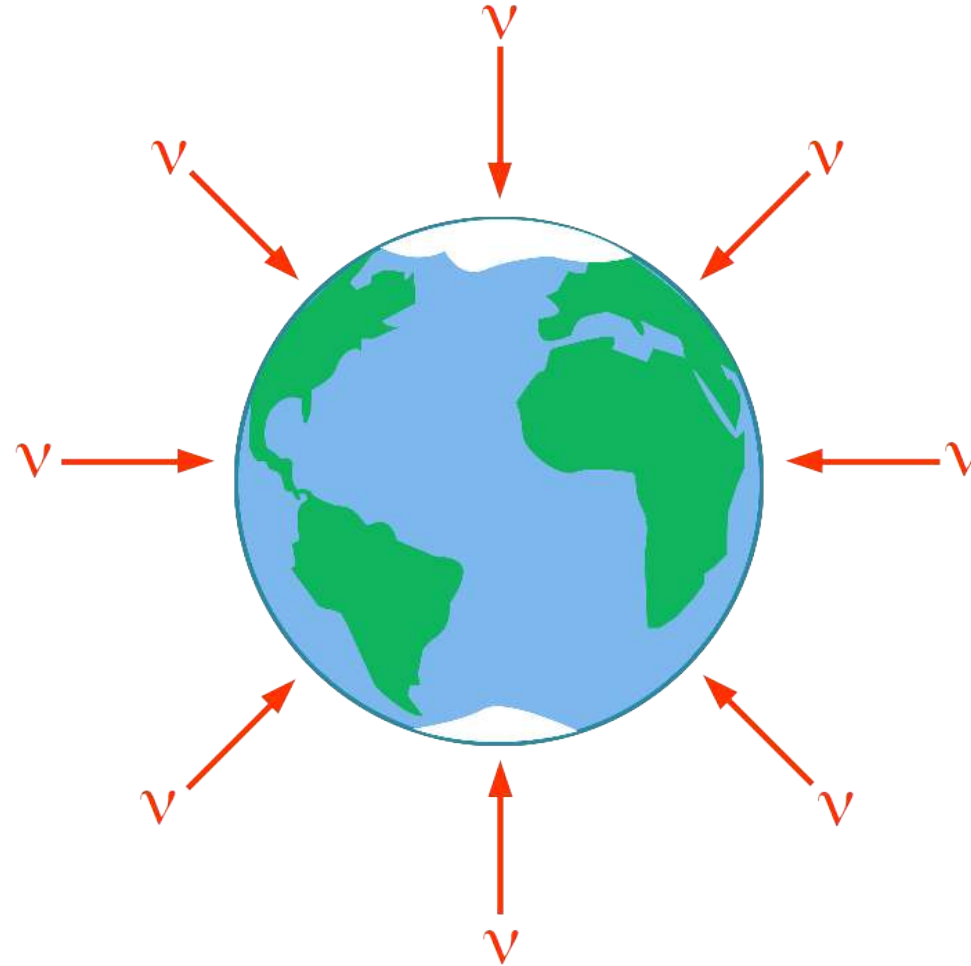
MB & A. Connolly, 1711.11043

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MB & A. Connolly, 1711.11043

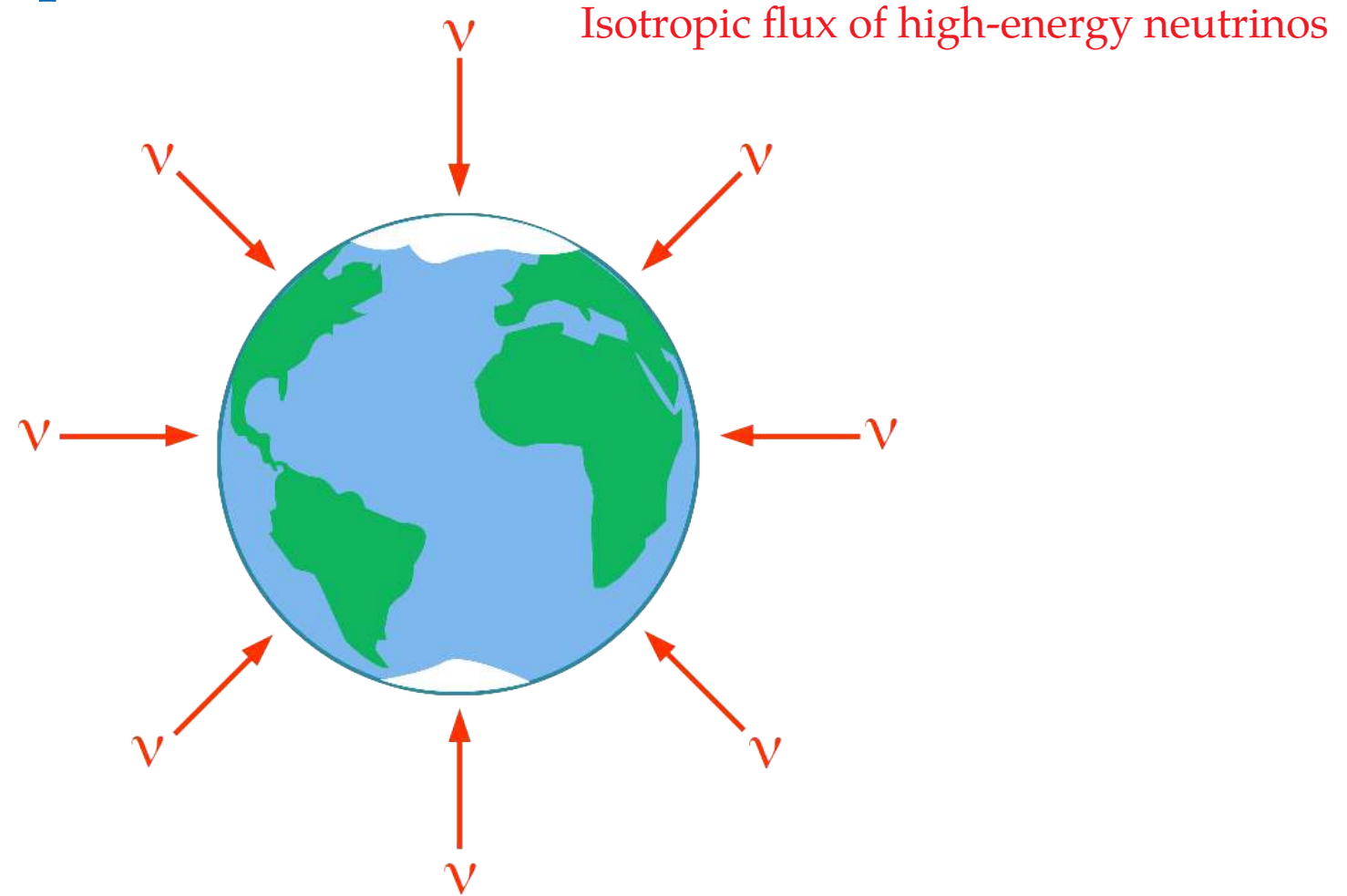
Neutrino, interrupted



Marfatia, McKay, Weiler, *PLB* 2015
Connolly, Thorne, Waters, *PRD* 2011
Hussain, Marfatia, McKay, *PRD* 2008
Borriello *et al.*, *PRD* 2008
Hussain, Marfatia, McKay, Seckel, *PRL* 2006
Hooper, *PRD* 2002

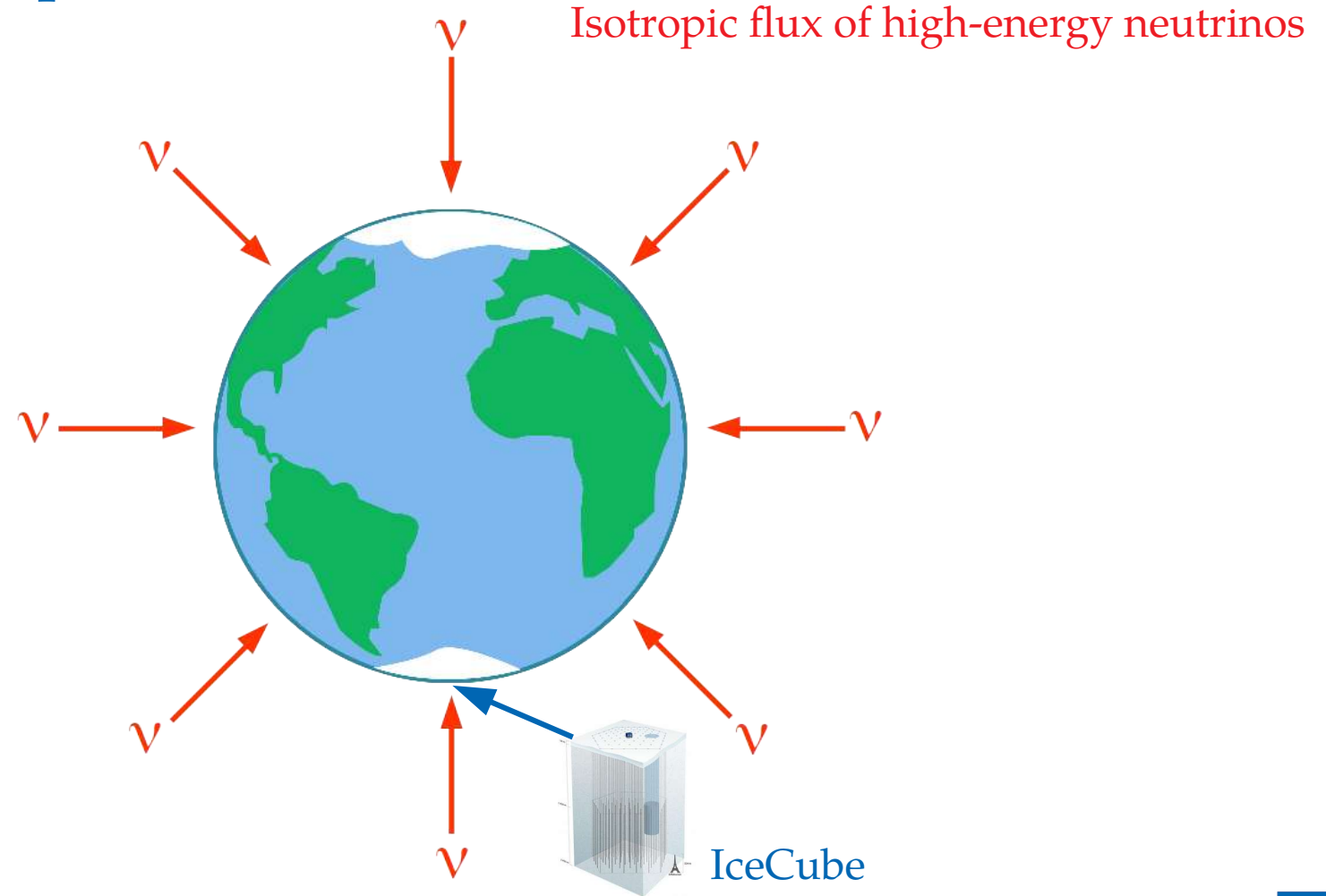
Mauricio Bustamante (Niels Bohr Institute)

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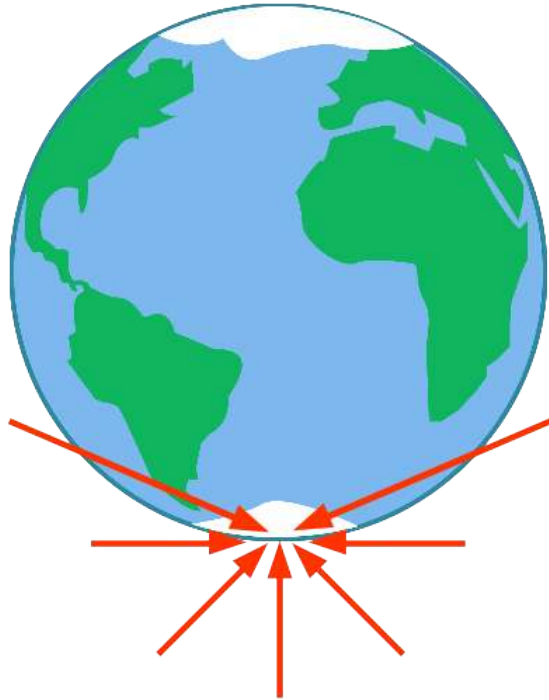
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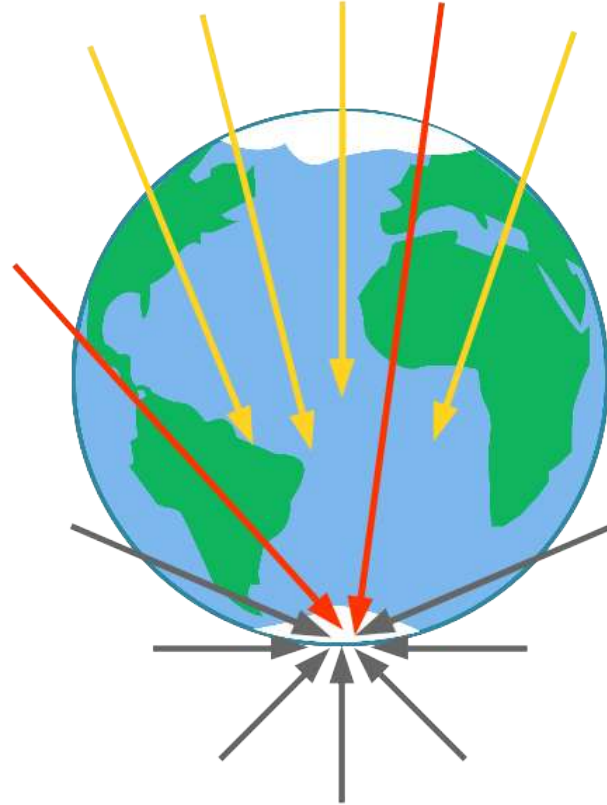


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Hooper, *PRD* 2002

Most of these neutrinos reach IceCube

Neutrino, interrupted

Many of these neutrinos are stopped by the Earth



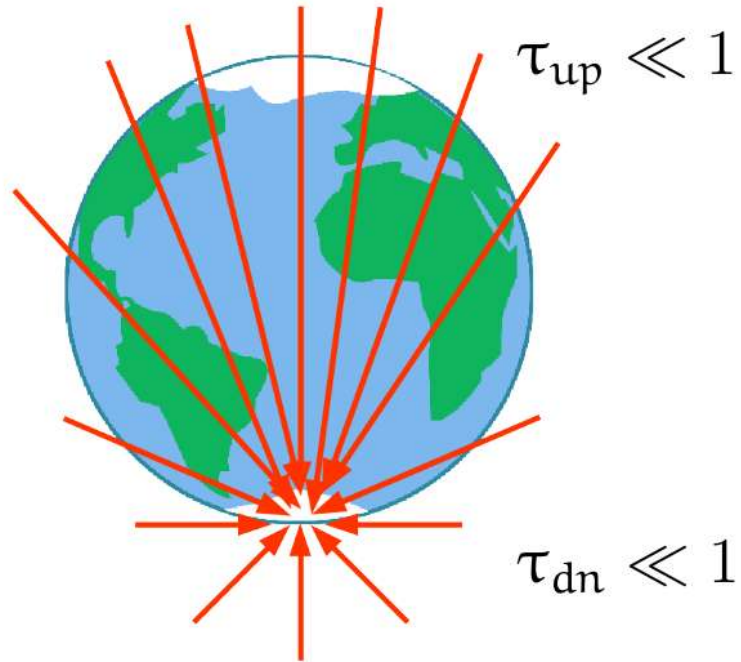
Marfatia, McKay, Weiler, *PLB* 2015
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Hussain, Marfatia, McKay, *PRD* 2008
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Hussain, Marfatia, McKay, Seckel, *PRL* 2006
Hooper, *PRD* 2002

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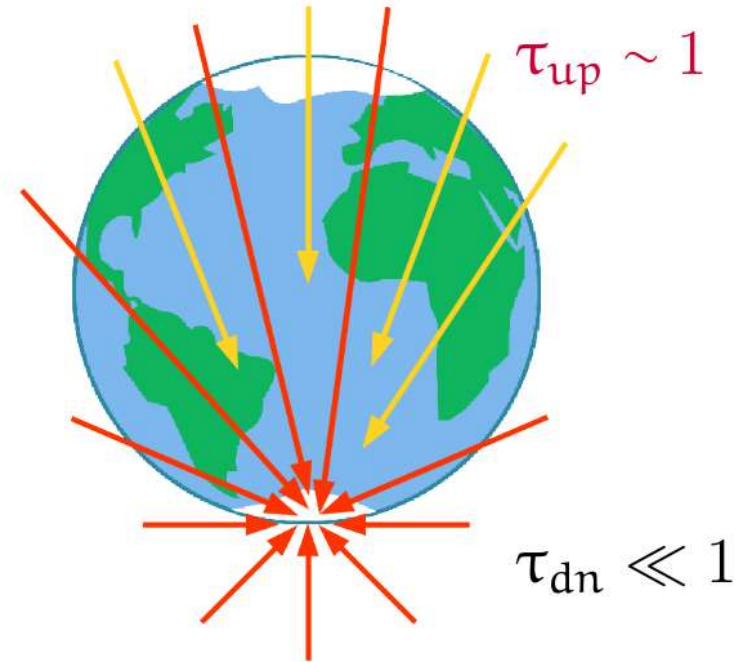
Measuring the high-energy cross section

$$\text{Optical depth to } \nu N \text{ int's} = \frac{\text{Distance from Earth's surface to IceCube}}{\text{Mean free path inside Earth}} \equiv \tau(E_\nu, \theta_z) \propto \sigma_{\nu N}$$

Below ~ 10 TeV: Earth is transparent



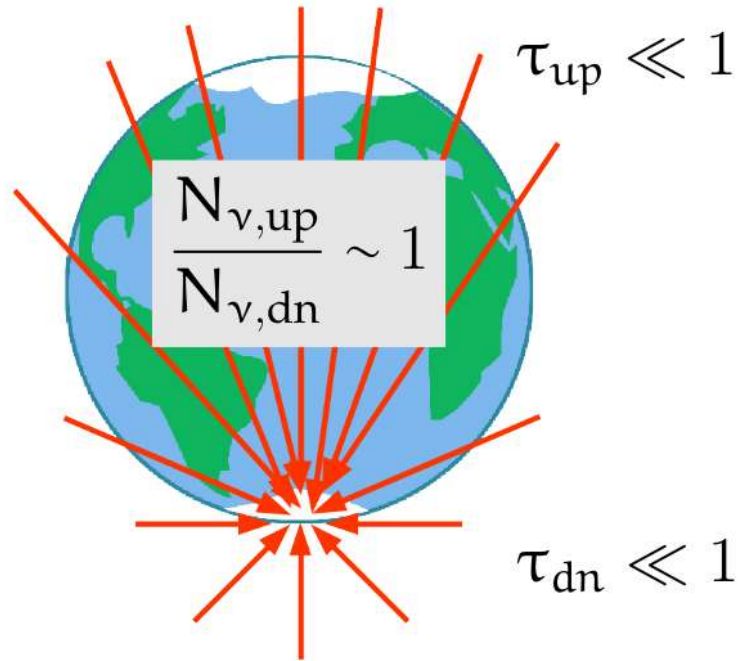
Above ~ 10 TeV: Earth is opaque



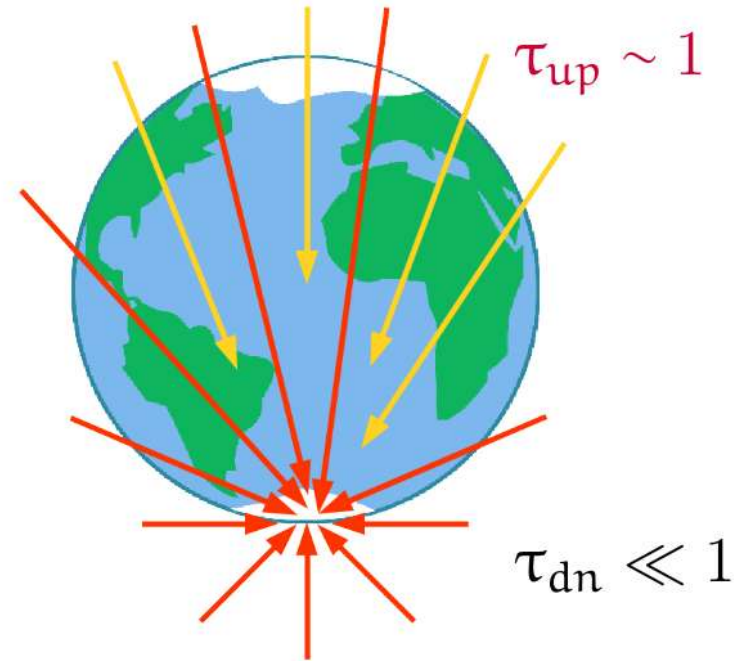
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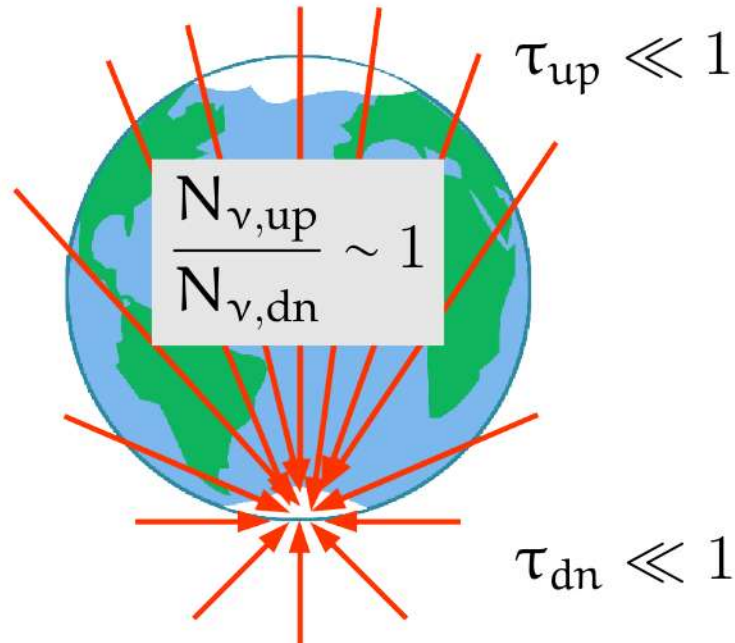
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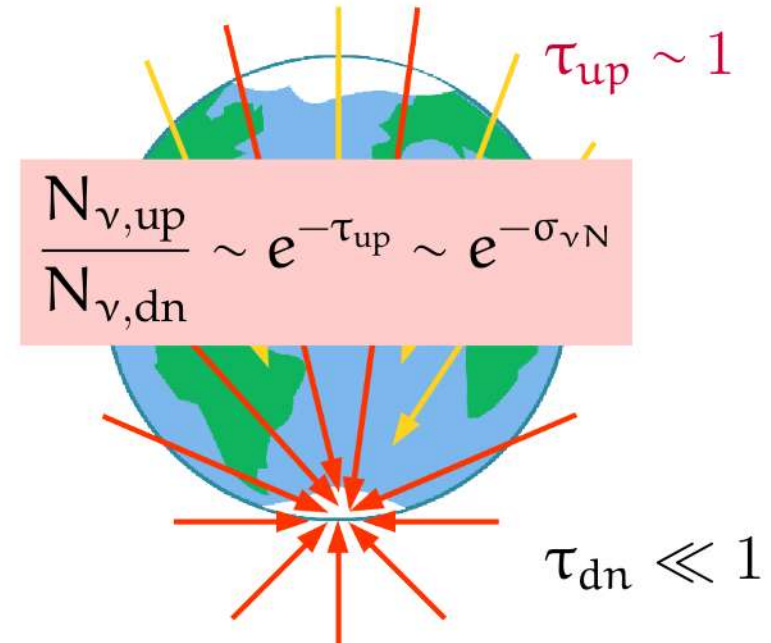
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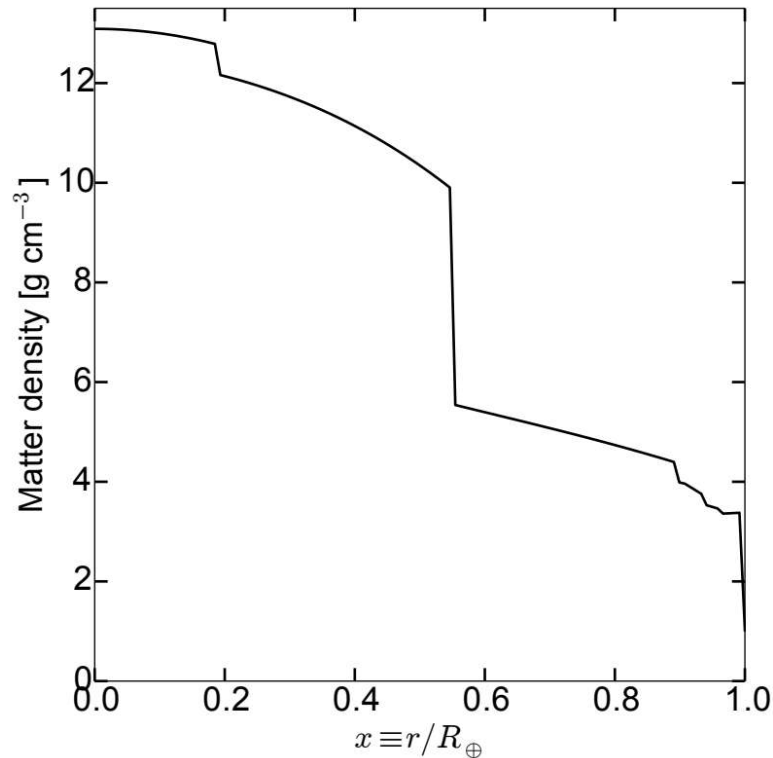
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A feel for the in-Earth attenuation

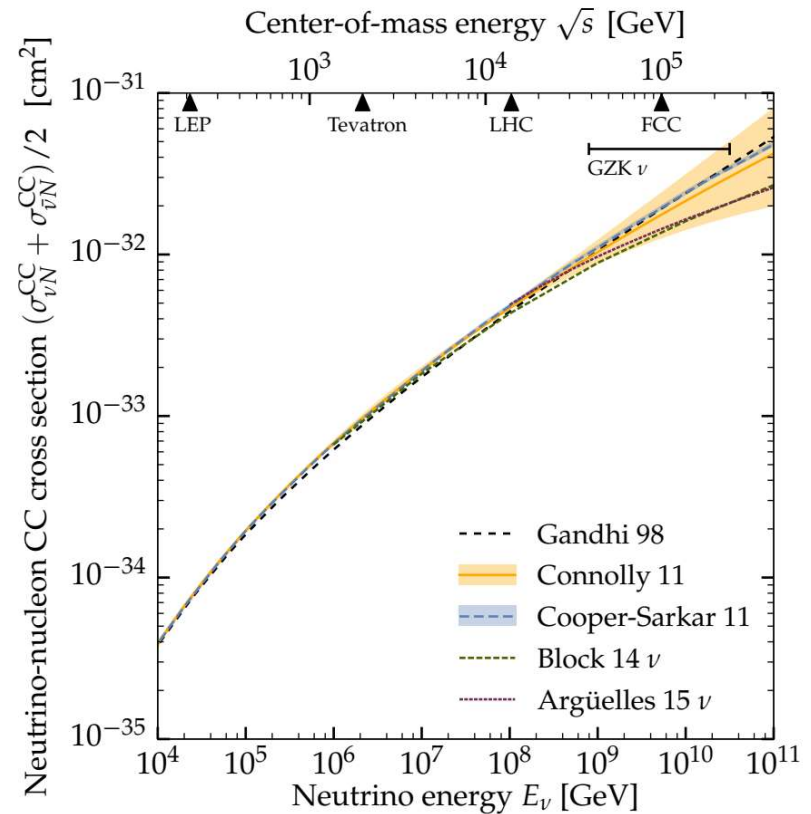
Earth matter density

(Preliminary Reference Earth Model)

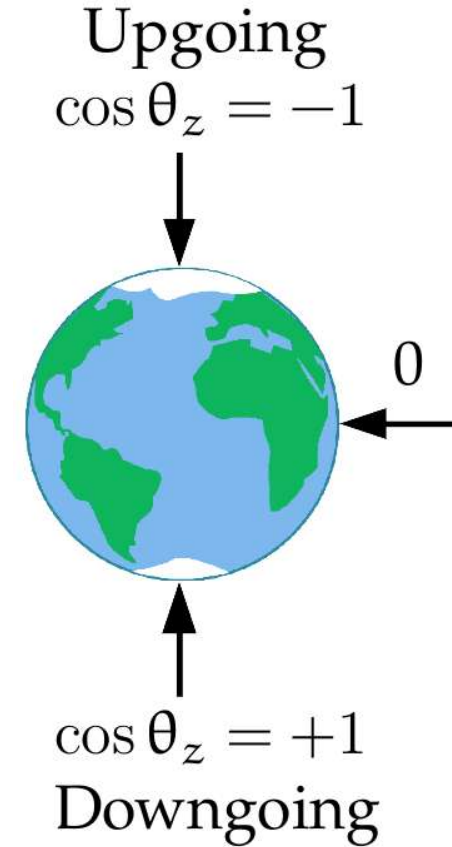
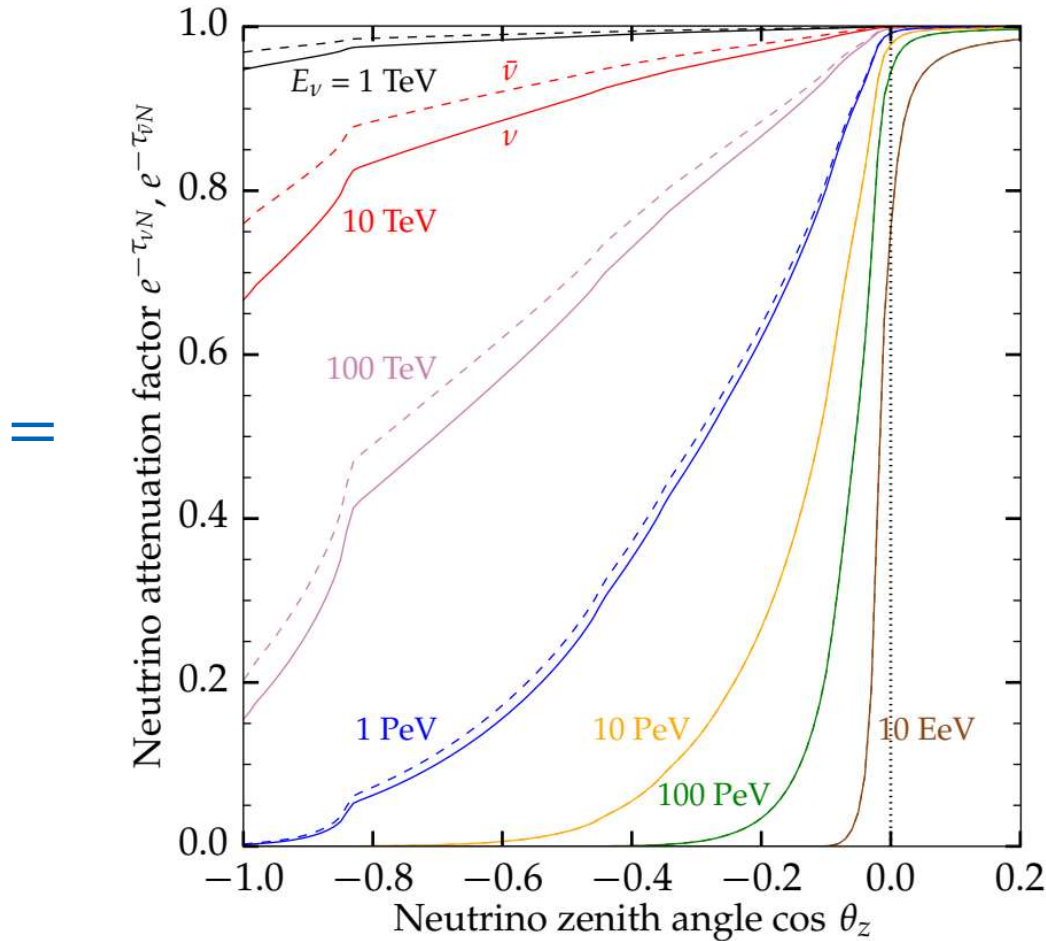


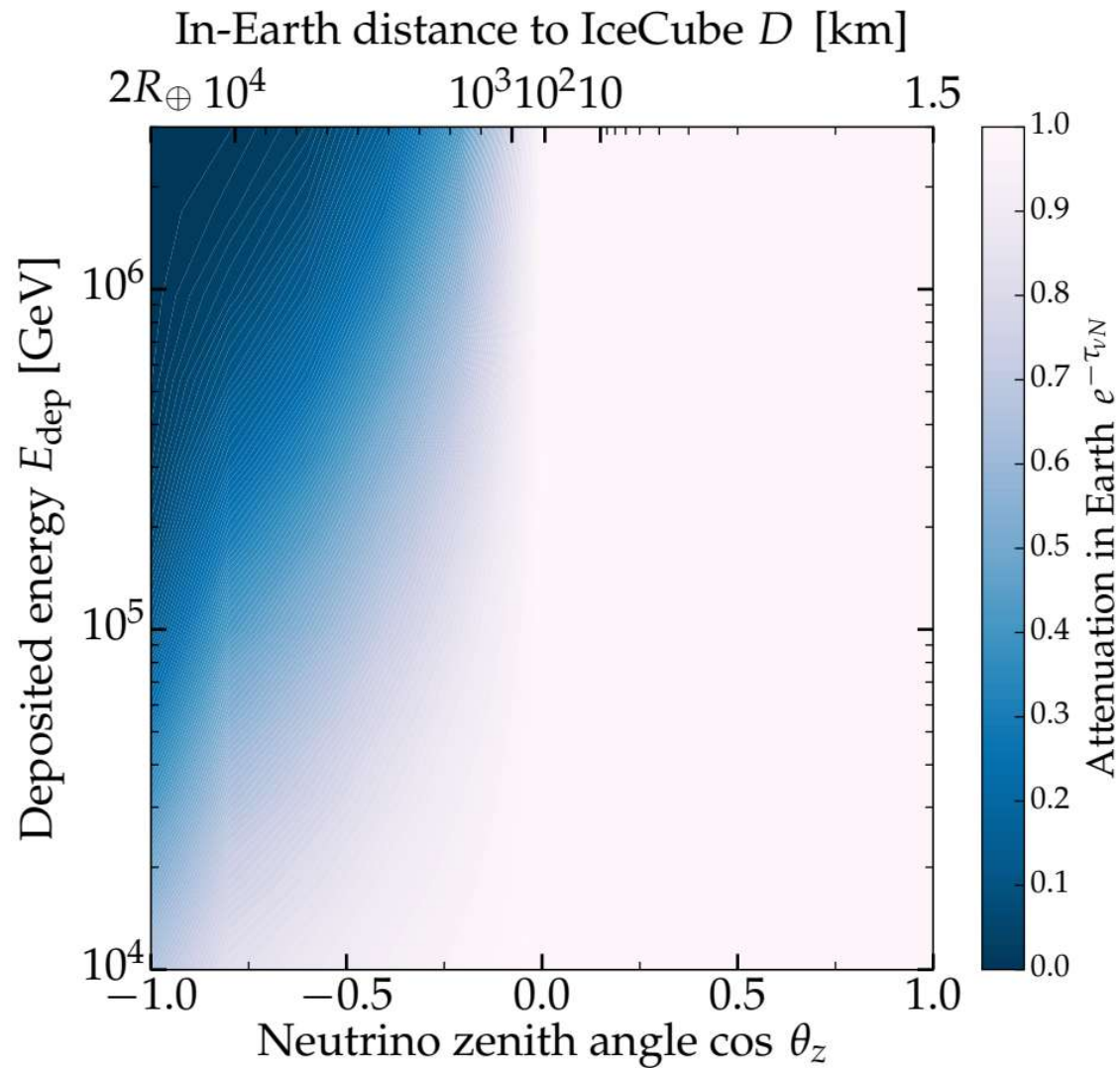
+

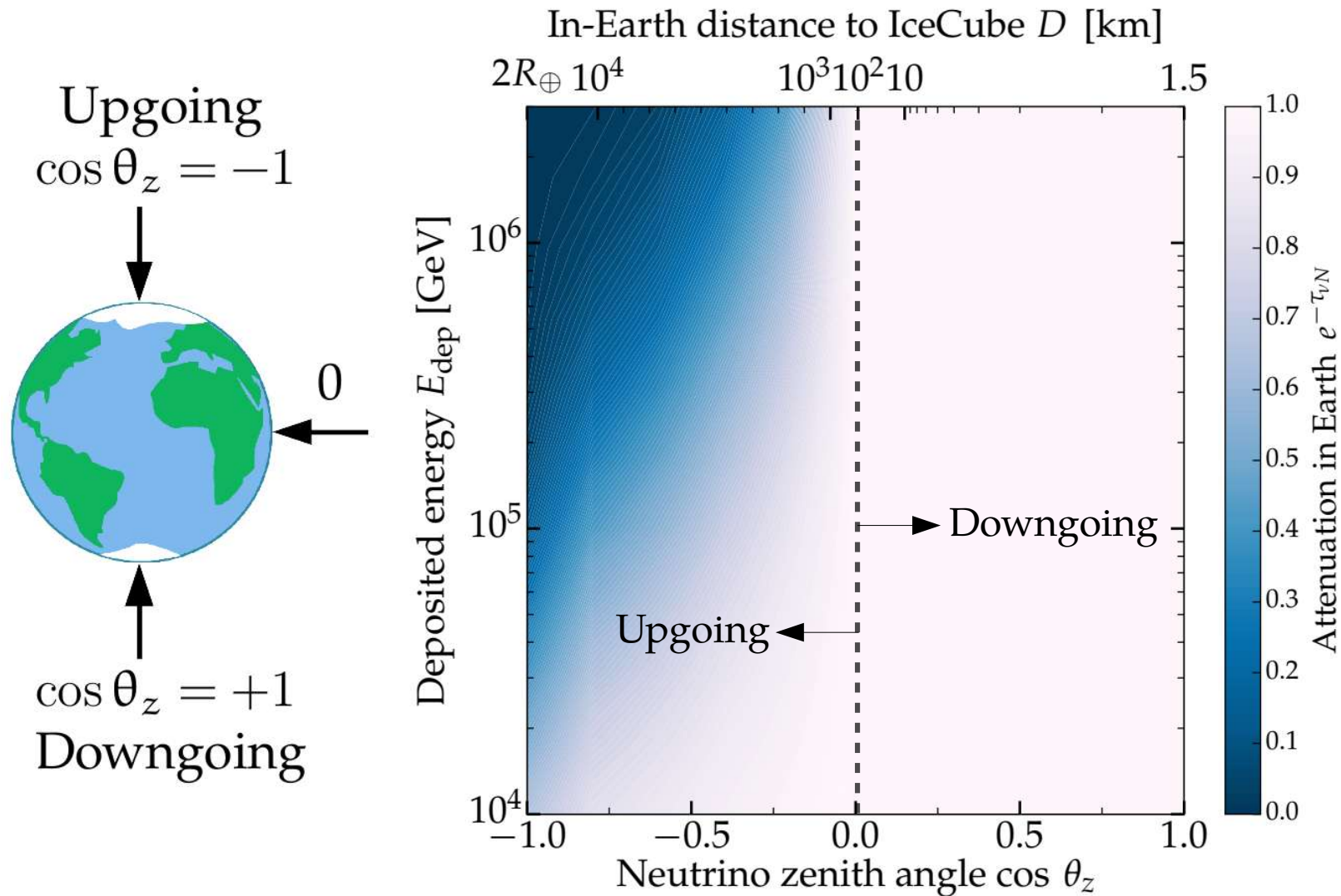
Neutrino-nucleon cross section

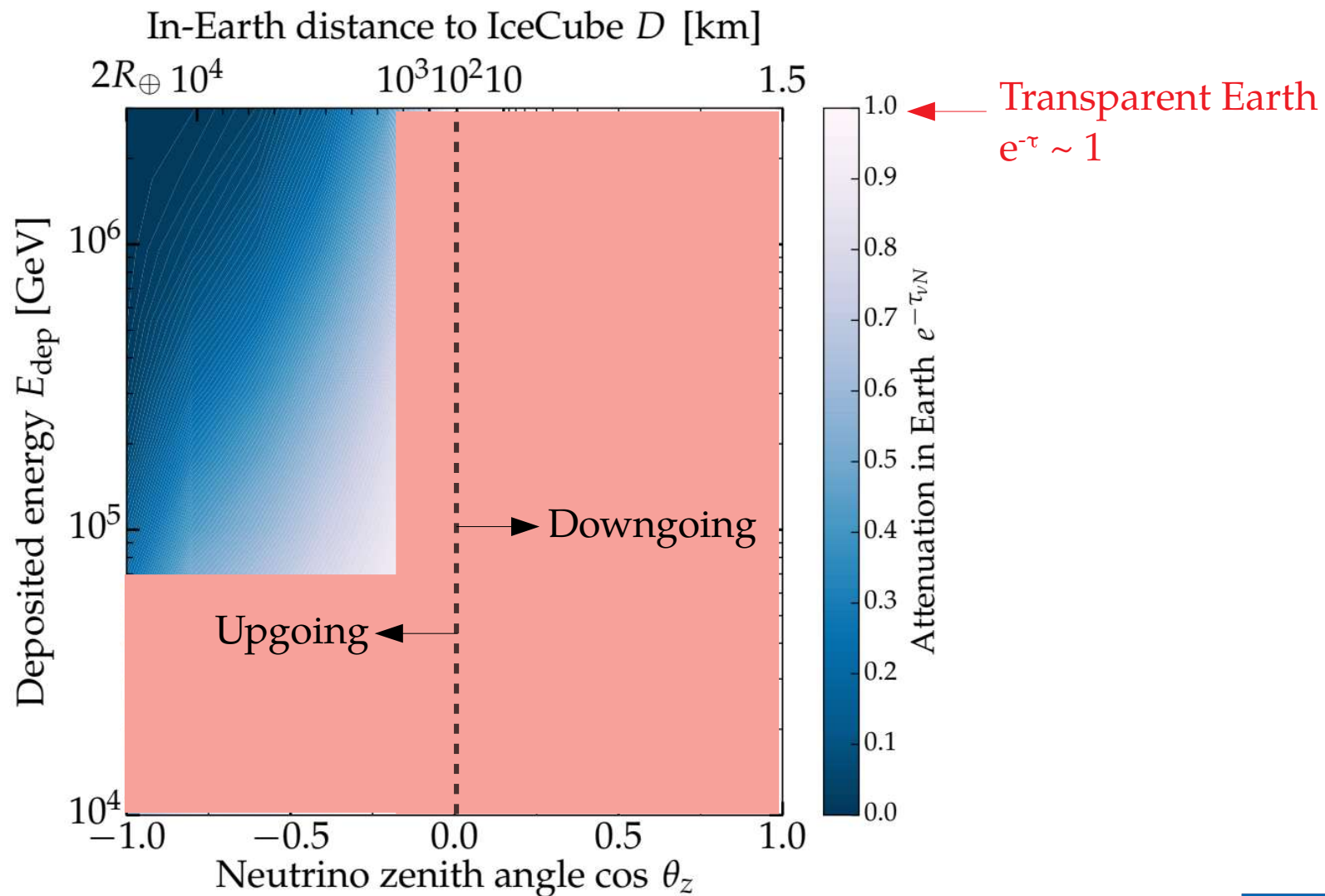
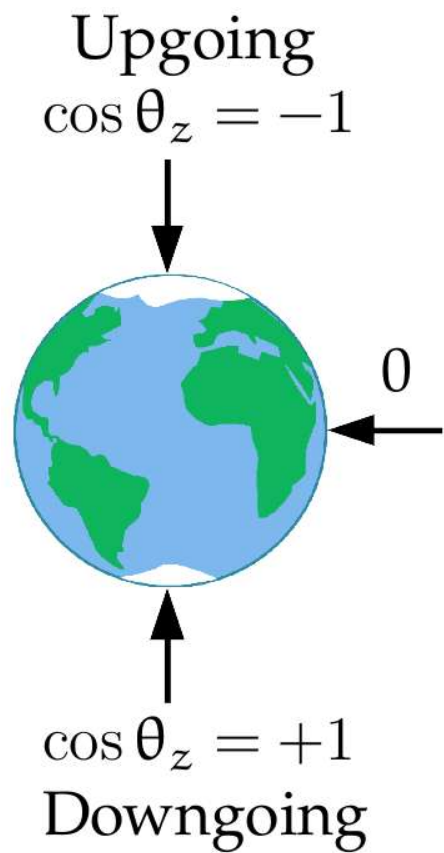


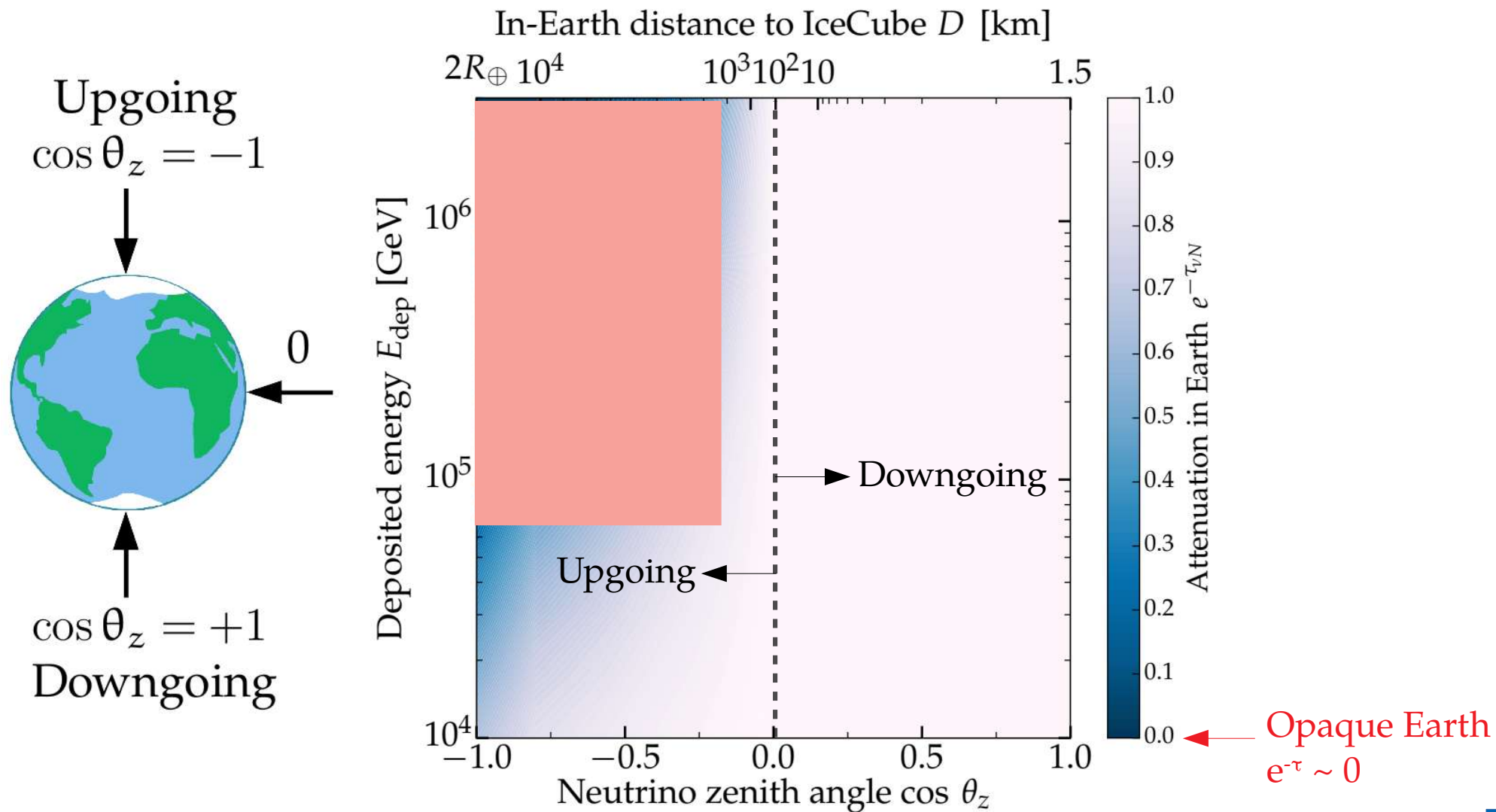
A feel for the in-Earth attenuation



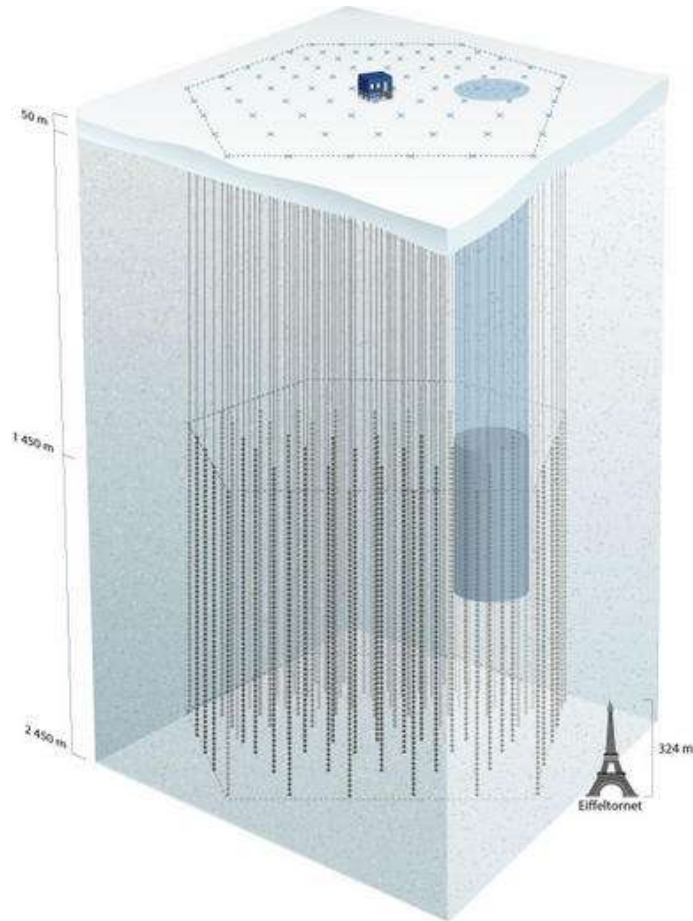








IceCube – What is it?



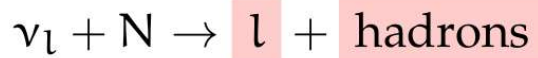
- ▶ Km^3 in-ice Cherenkov detector in Antarctica
- ▶ >5000 PMTs at 1.5–2.5 km of depth
- ▶ Sensitive to neutrino energies $> 10 \text{ GeV}$



How does IceCube see neutrinos?

Two types of fundamental interactions ...

Charged-current (CC)



Neutral-current (NC)



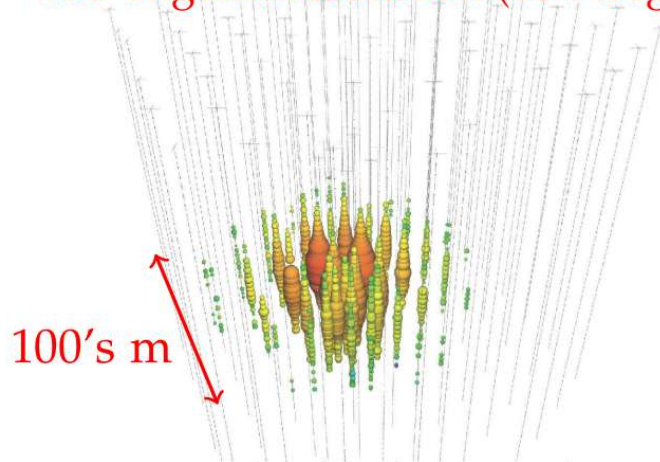
...create two event topologies ...

These shower and make light

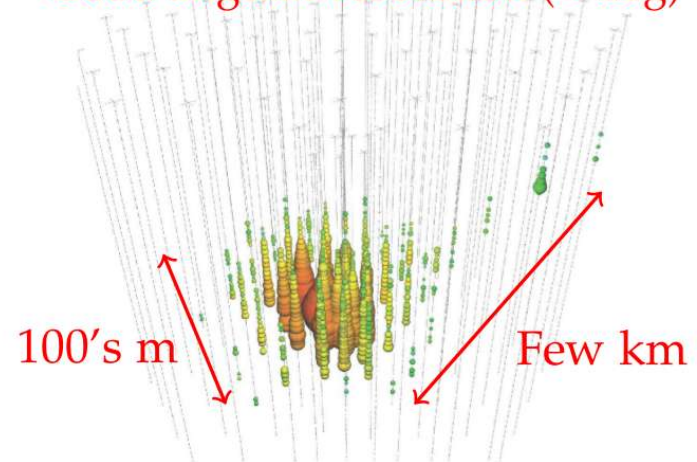
Showers — From CC ν_e or ν_τ , or NC ν_χ

Tracks — From CC ν_μ mainly

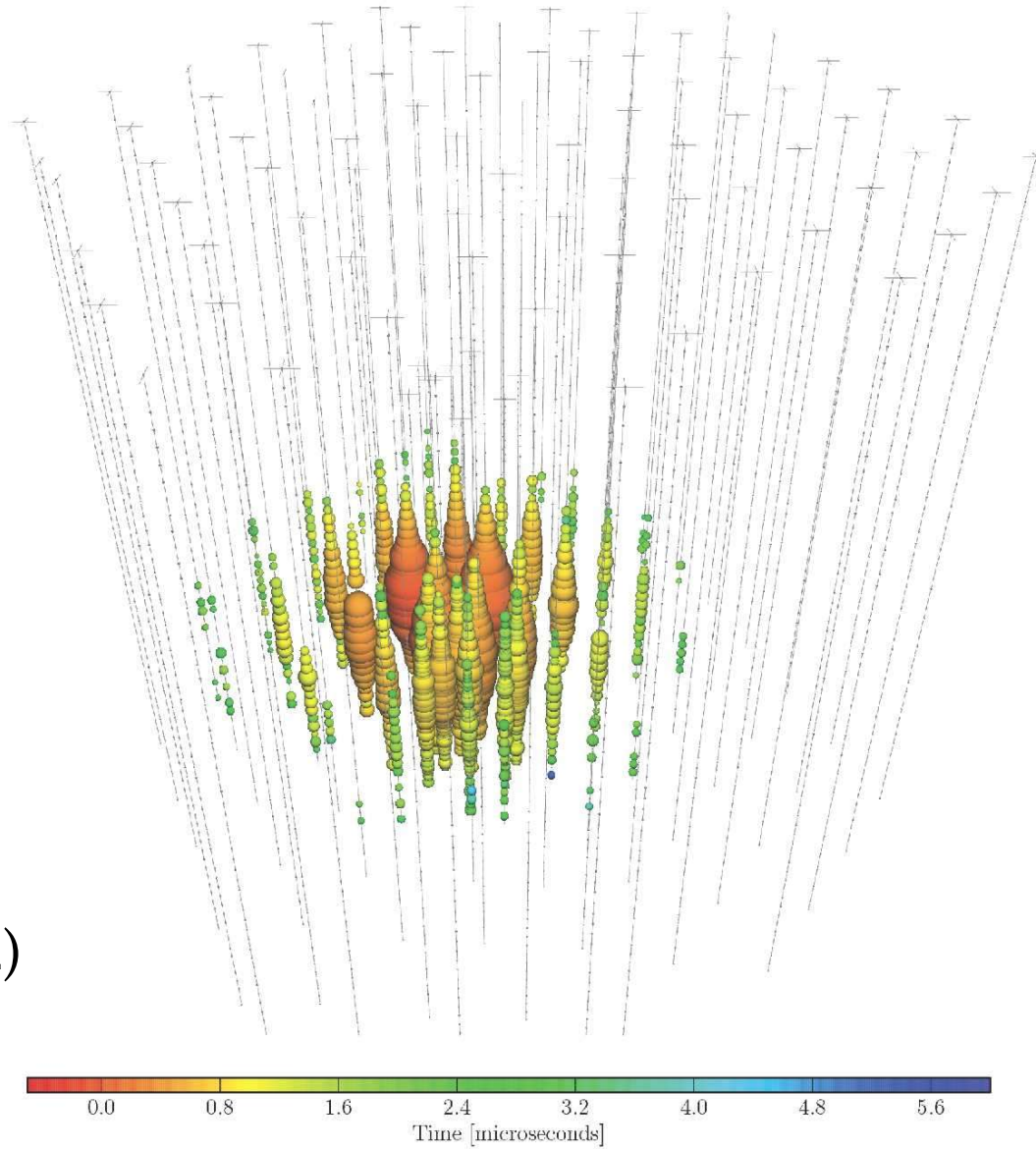
Bad angular resolution (10's deg)



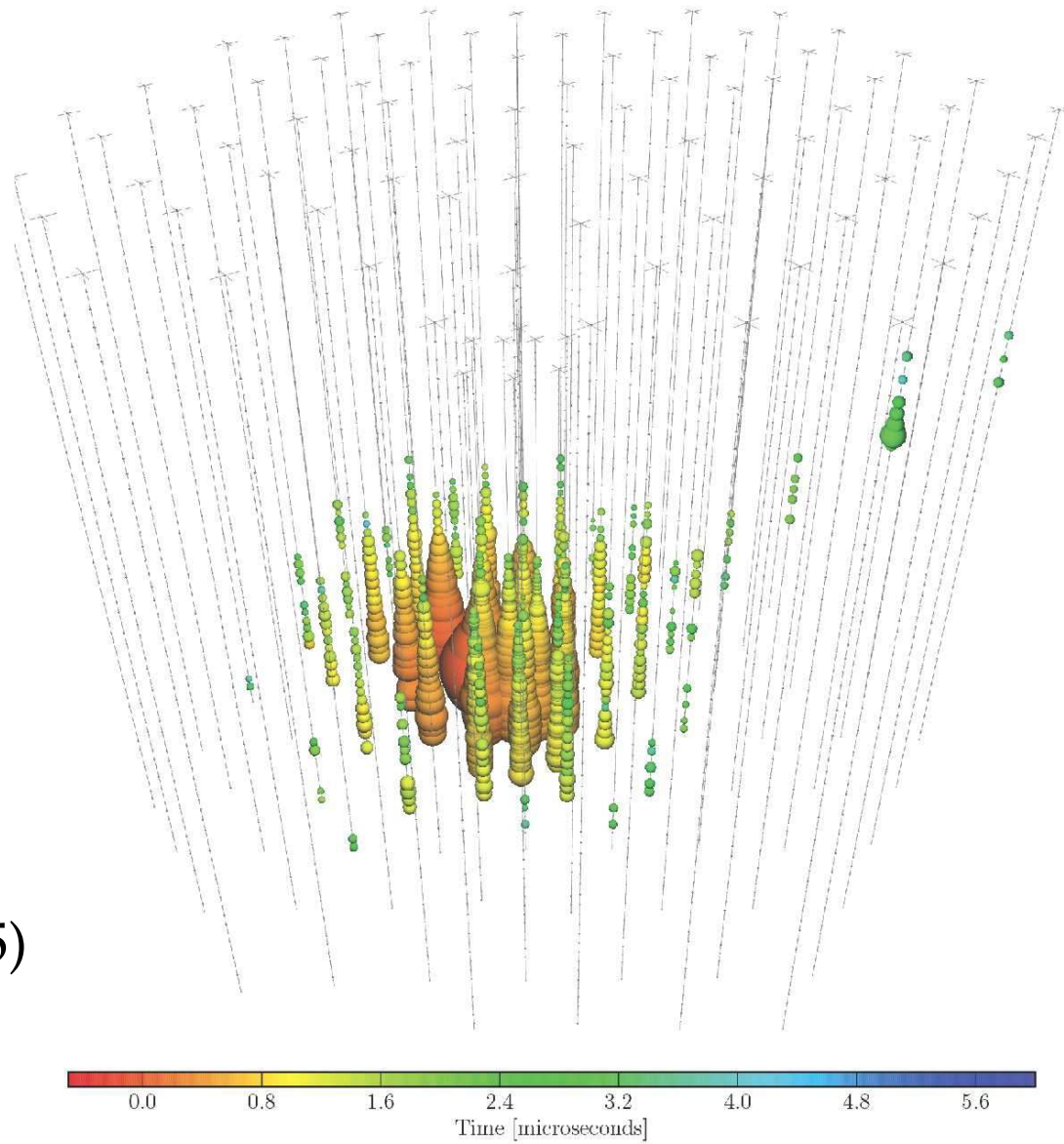
Good angular resolution (< deg)



Shower
(IceCube event #22)



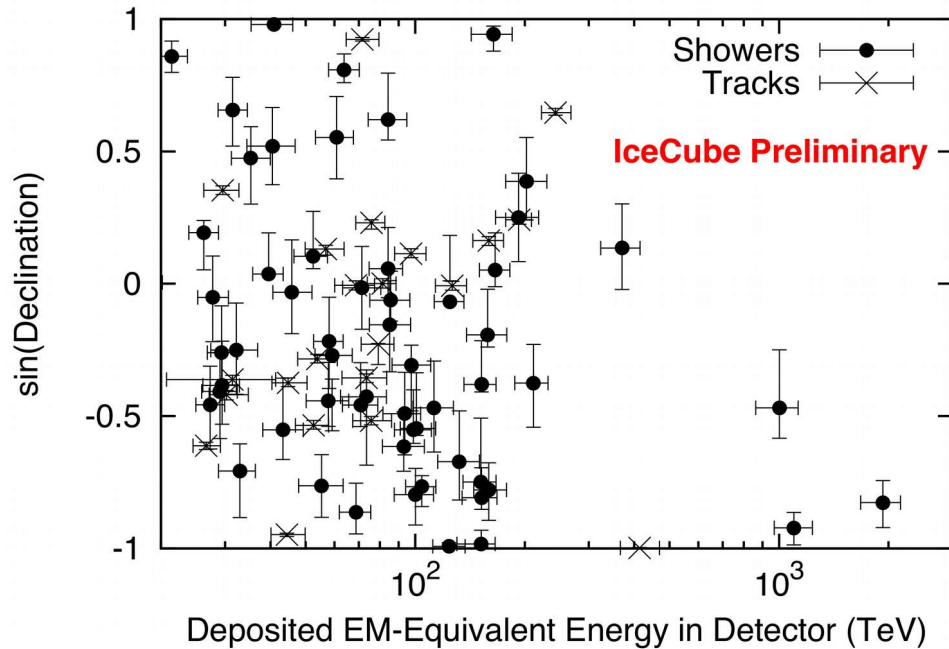
Track
(IceCube event #15)



What has IceCube found so far (6 years)?

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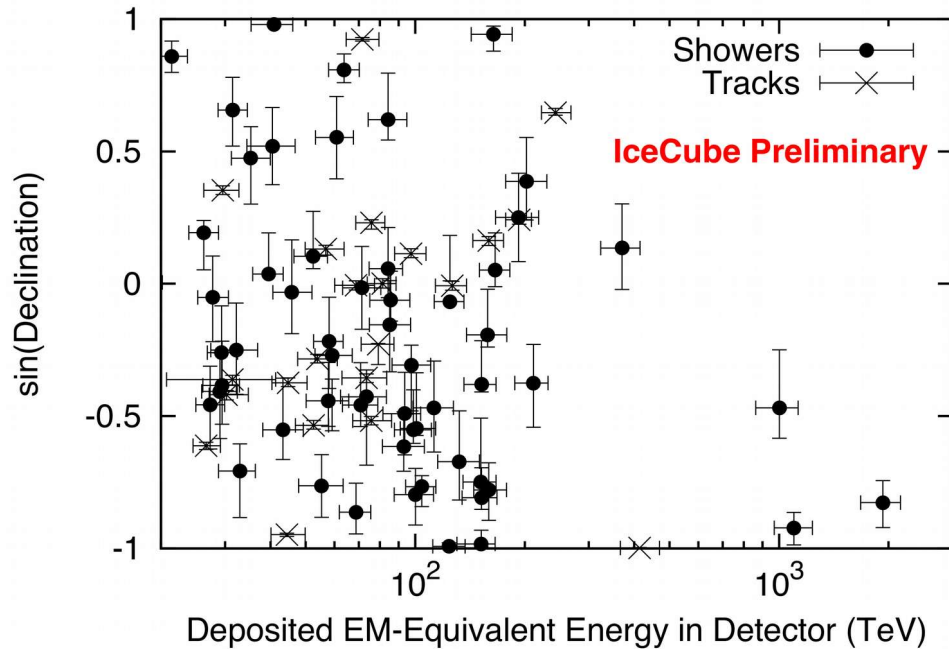
80 contained events between 18 TeV – 2 PeV
(16 atm. neutrinos, 25 atm. muons)



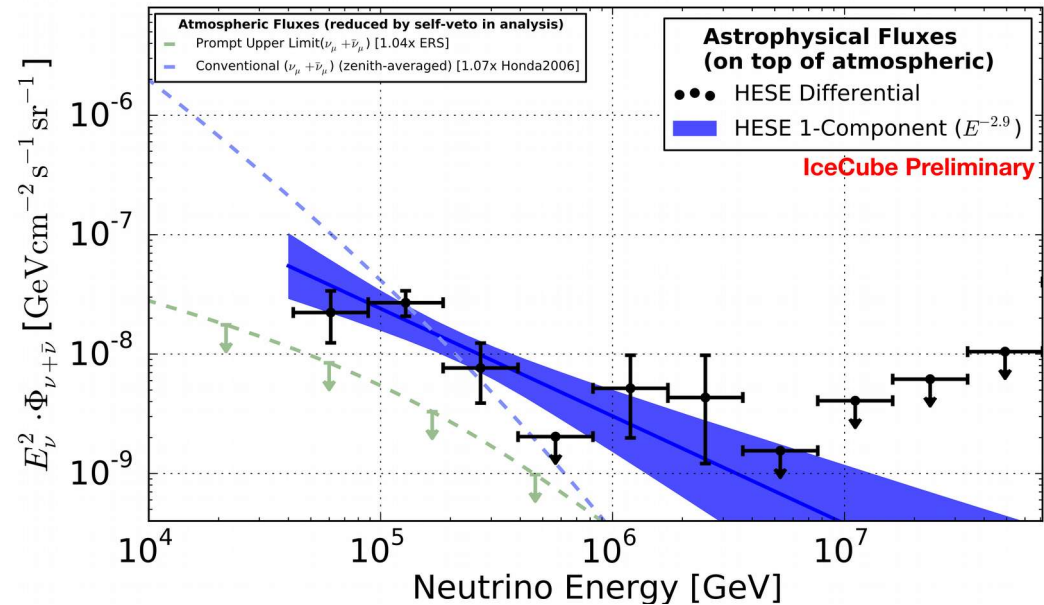
C. Kopper, ICRC 2017

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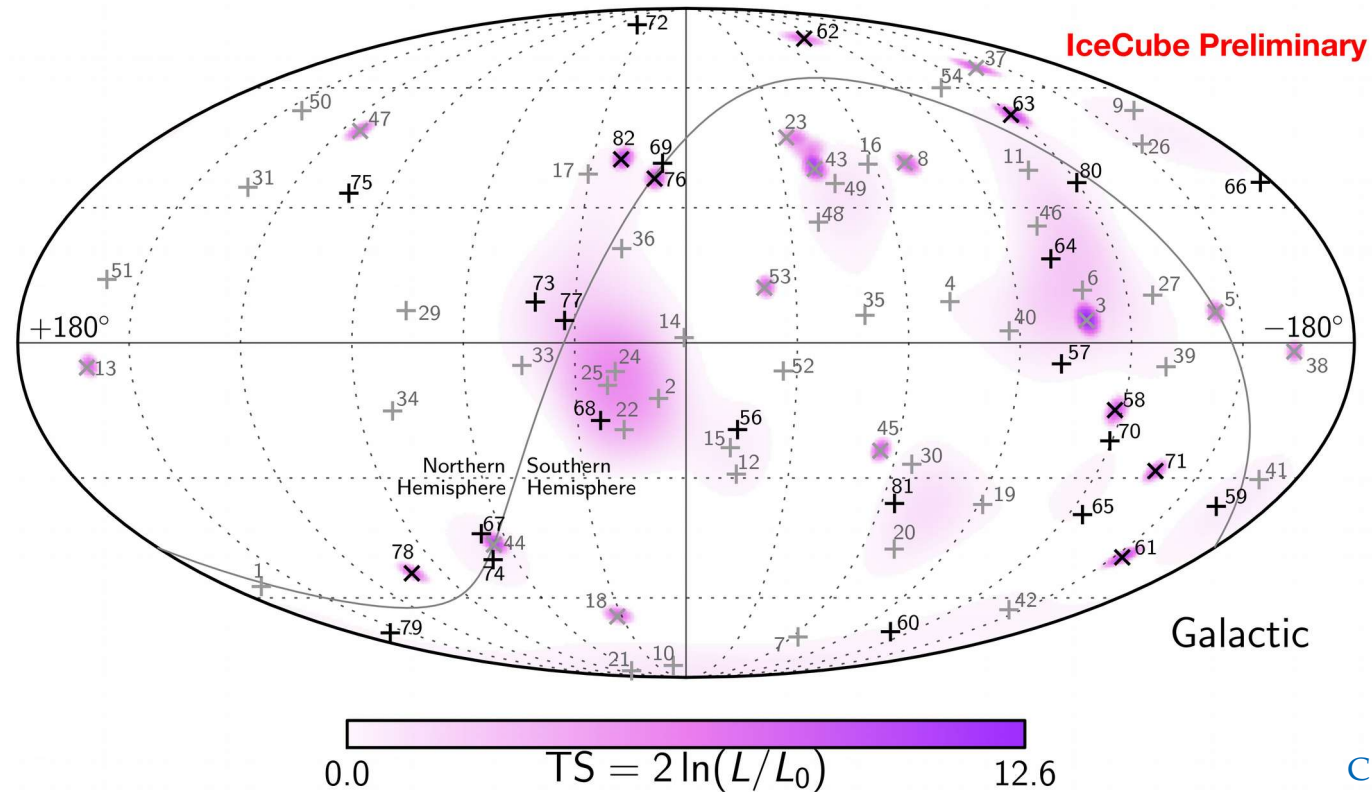
Astrophysical ν flux detected at $> 7\sigma$
(Normalization ok, but steep spectrum)



C. Kopper, ICRC 2017

What has IceCube found so far (6 years)?

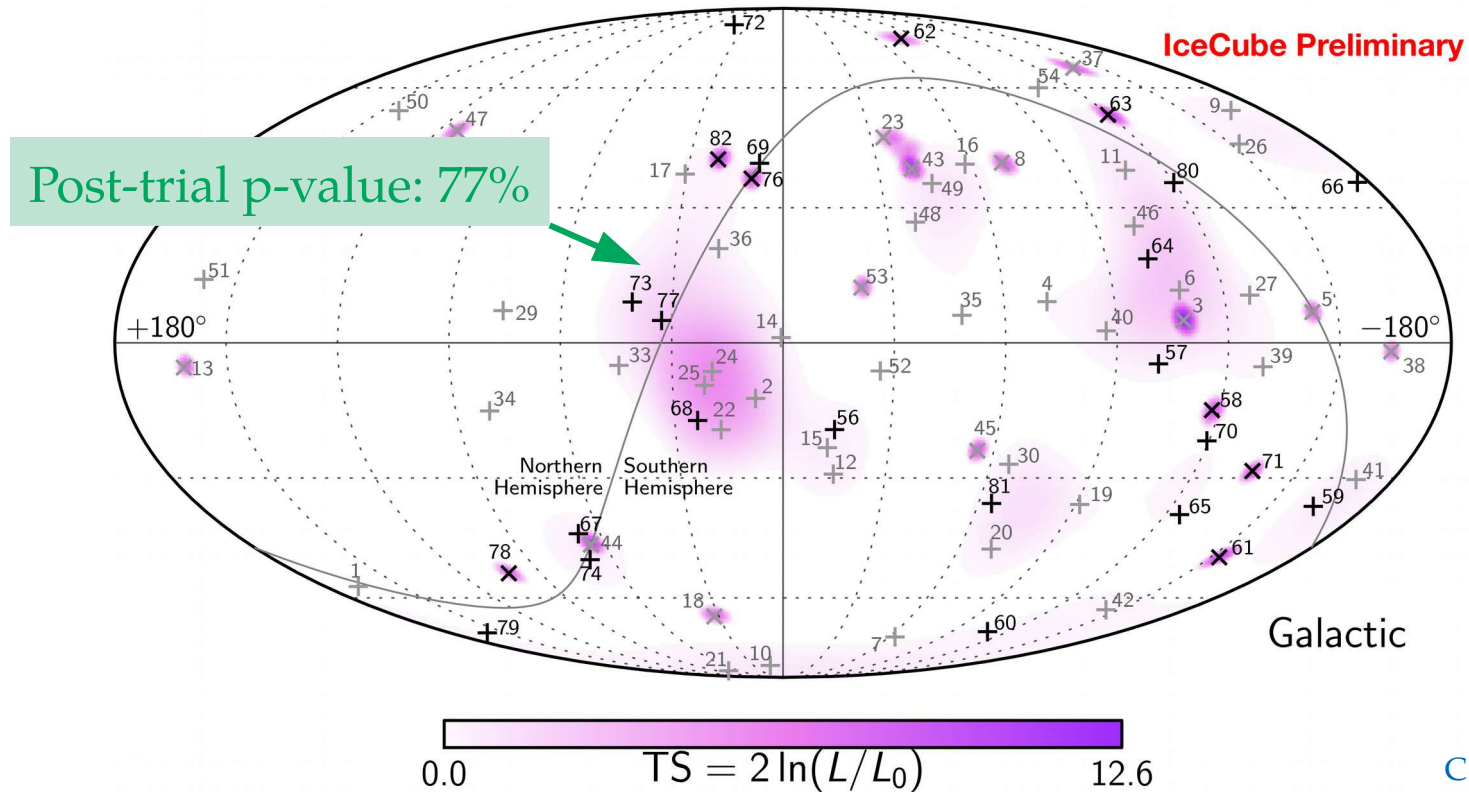
Arrival directions compatible with isotropy



C. Kopper, ICRC 2017

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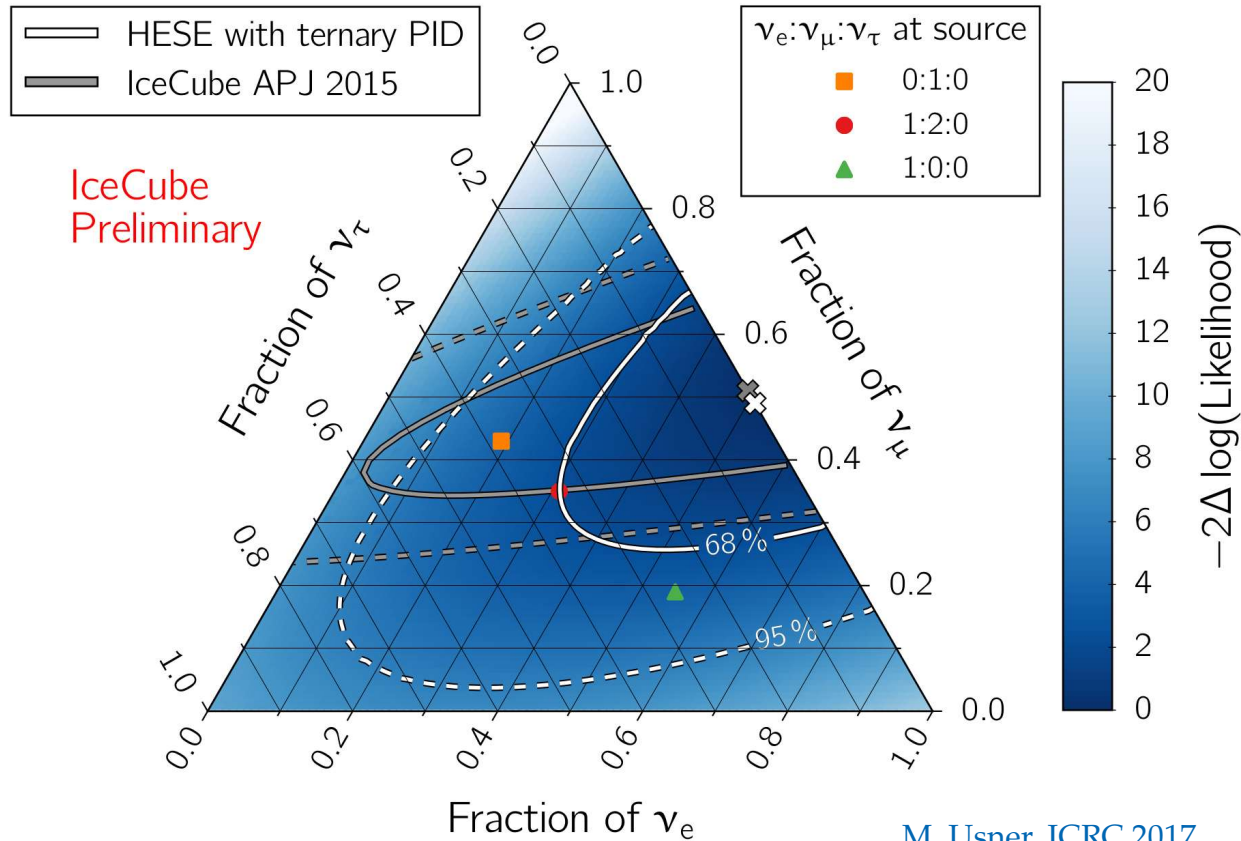
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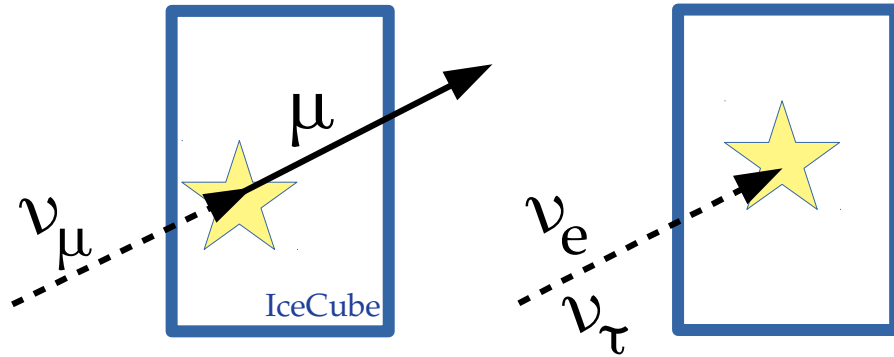
Flavor composition compatible with equal proportion of each flavor



M. Usner, ICRC 2017

Contained vs. uncontained νN interactions

Contained events



Starting track

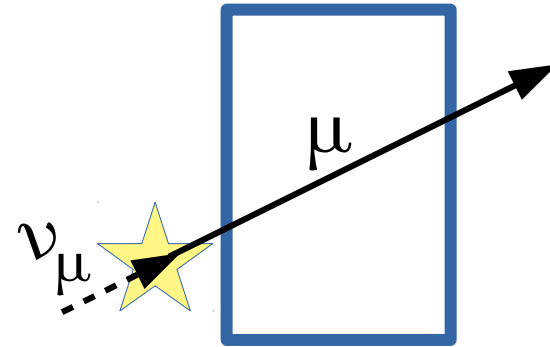
Shower

Pro: Clean determination of E_ν

Con: Few events (<100)

Ref.: MB & A. Connolly, 1711.11043

Uncontained events



Through-going muon

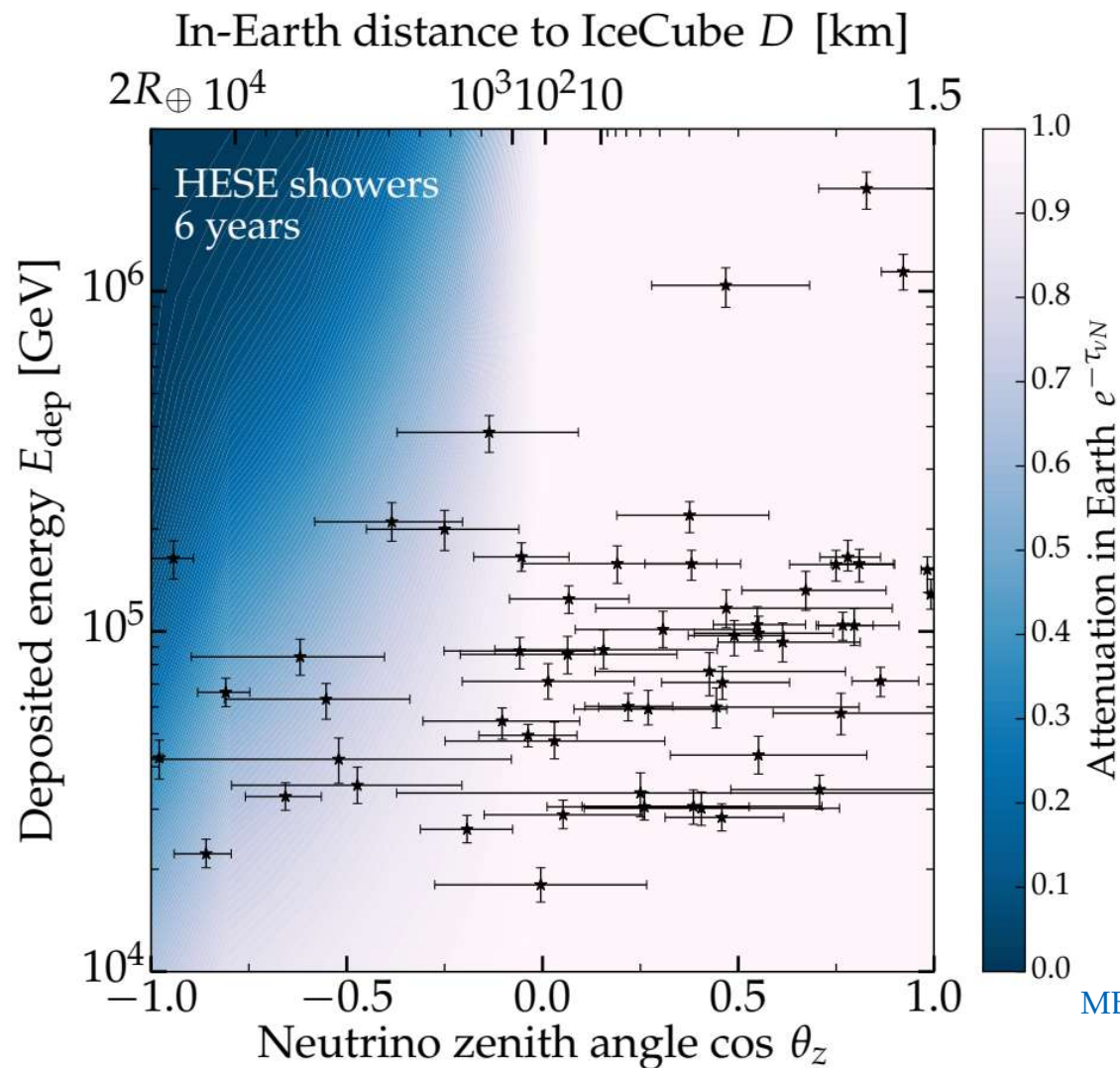
Pro: Lots of events ($\sim 10k$ used)

Con: Uncertain estimates of E_ν

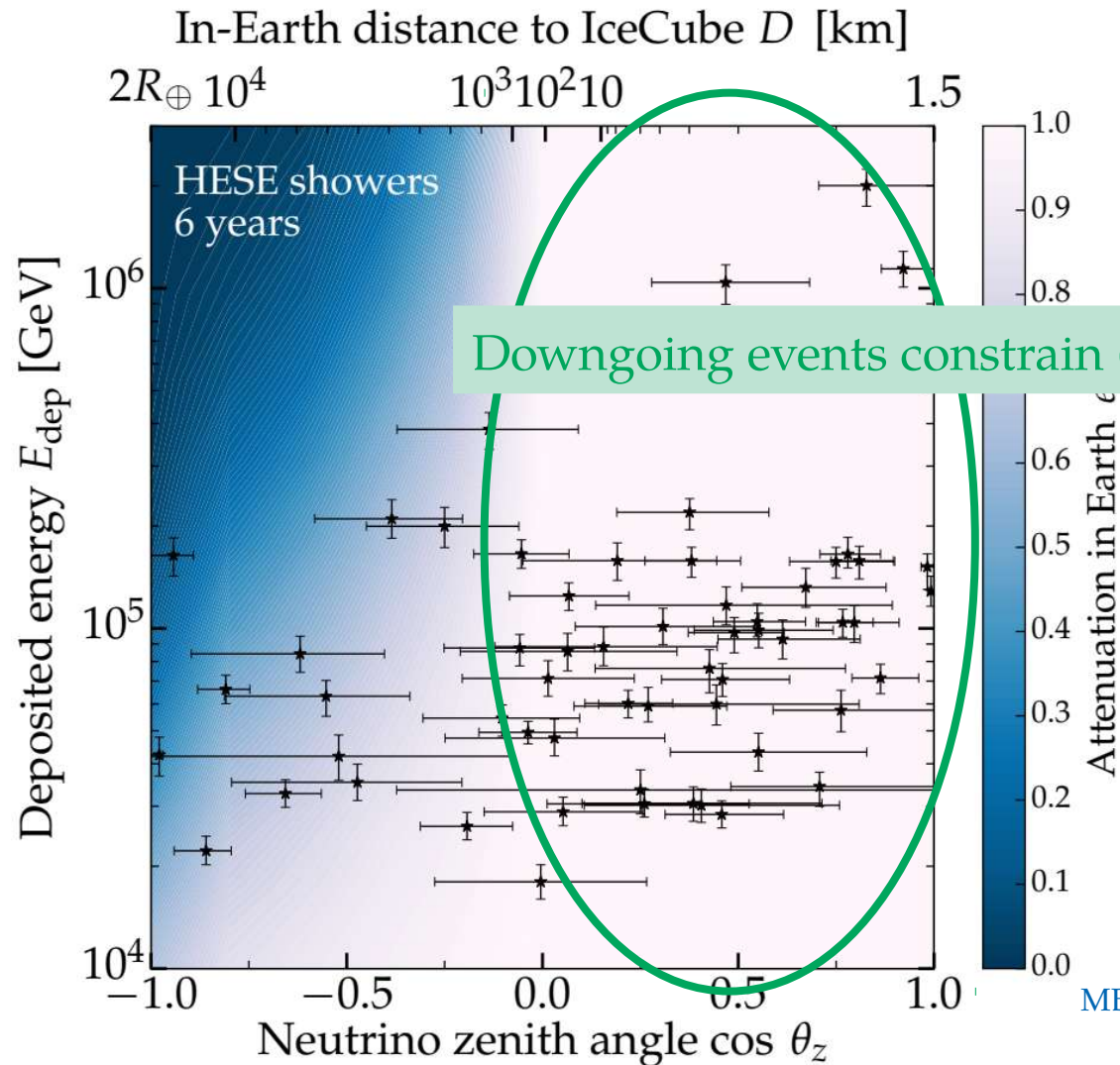
Ref.: IceCube, Nature 2017, 1711.08119

Cross section from contained events

- ▶ $\sigma_{\nu N}$ varies with neutrino energy $\Rightarrow$ use events where E_ν is well-reconstructed
- ▶ These are IceCube High-Energy Starting Events (HESE):
 - ▶ νN interaction occurs inside the detector
 - ▶ ☒ **Showers:** completely contained in the detector ($E_{\text{dep}} \approx E_\nu$)
 - ▶ ☐ **Tracks:** partially contained ($E_{\text{dep}} < E_\nu$)
- ▶ We use the 58 publicly available HESE showers (6-year sample)
- ▶ HESE tracks *could* be used
 - but we would need non-public data to reconstruct E_ν without bias

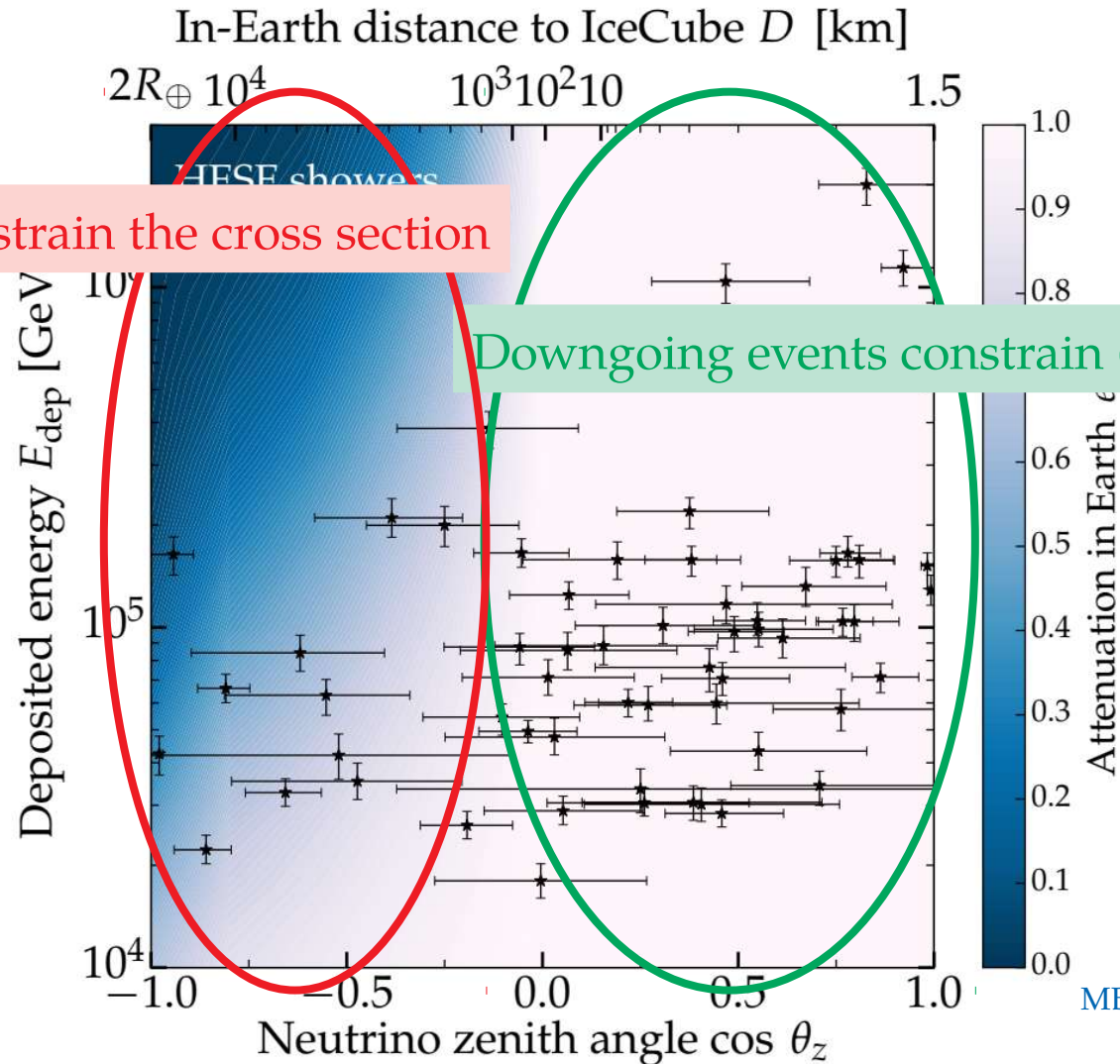


MB & A. Connolly, 1711.11043



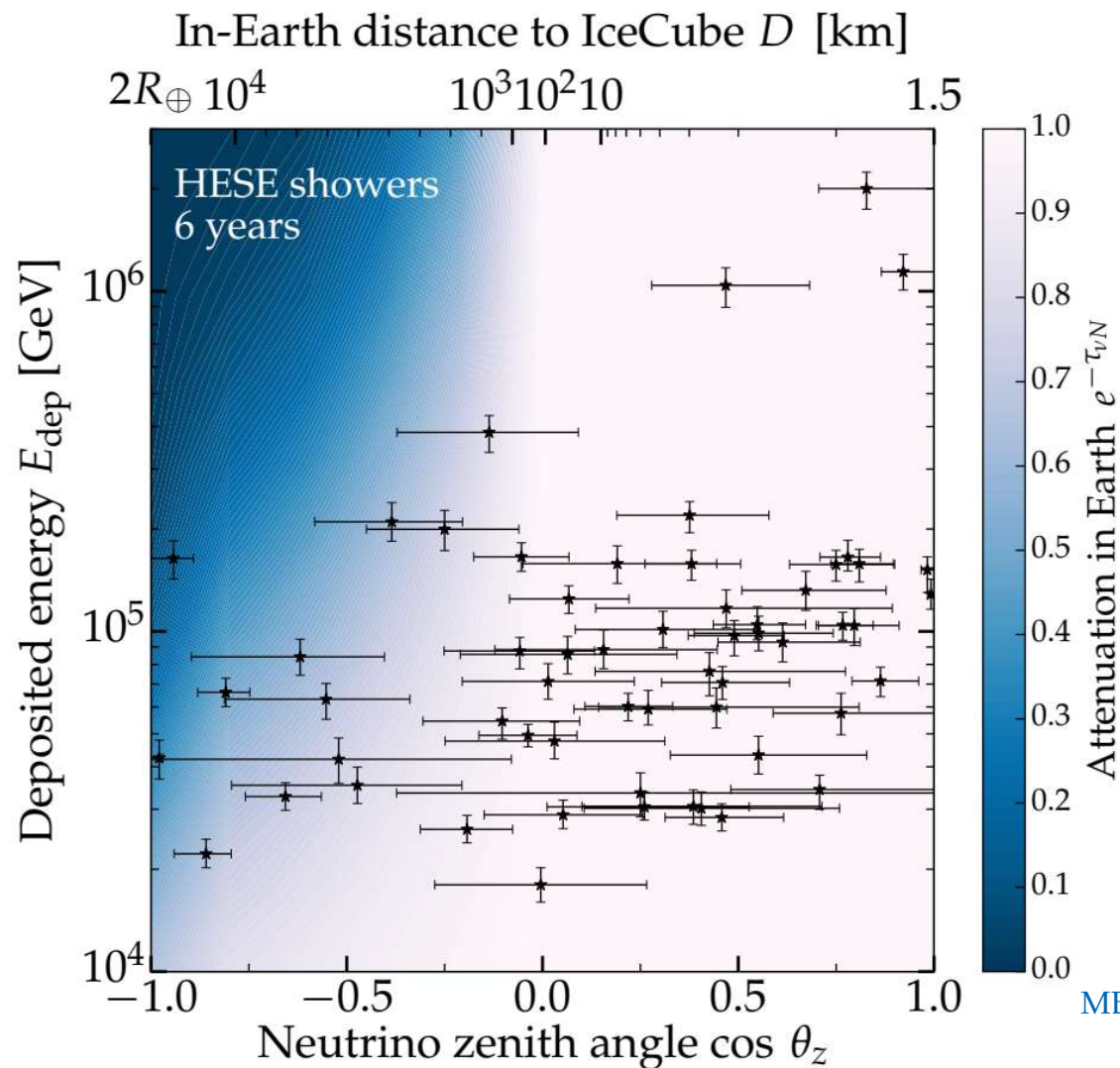
MB & A. Connolly, 1711.11043

Upgoing events constrain the cross section

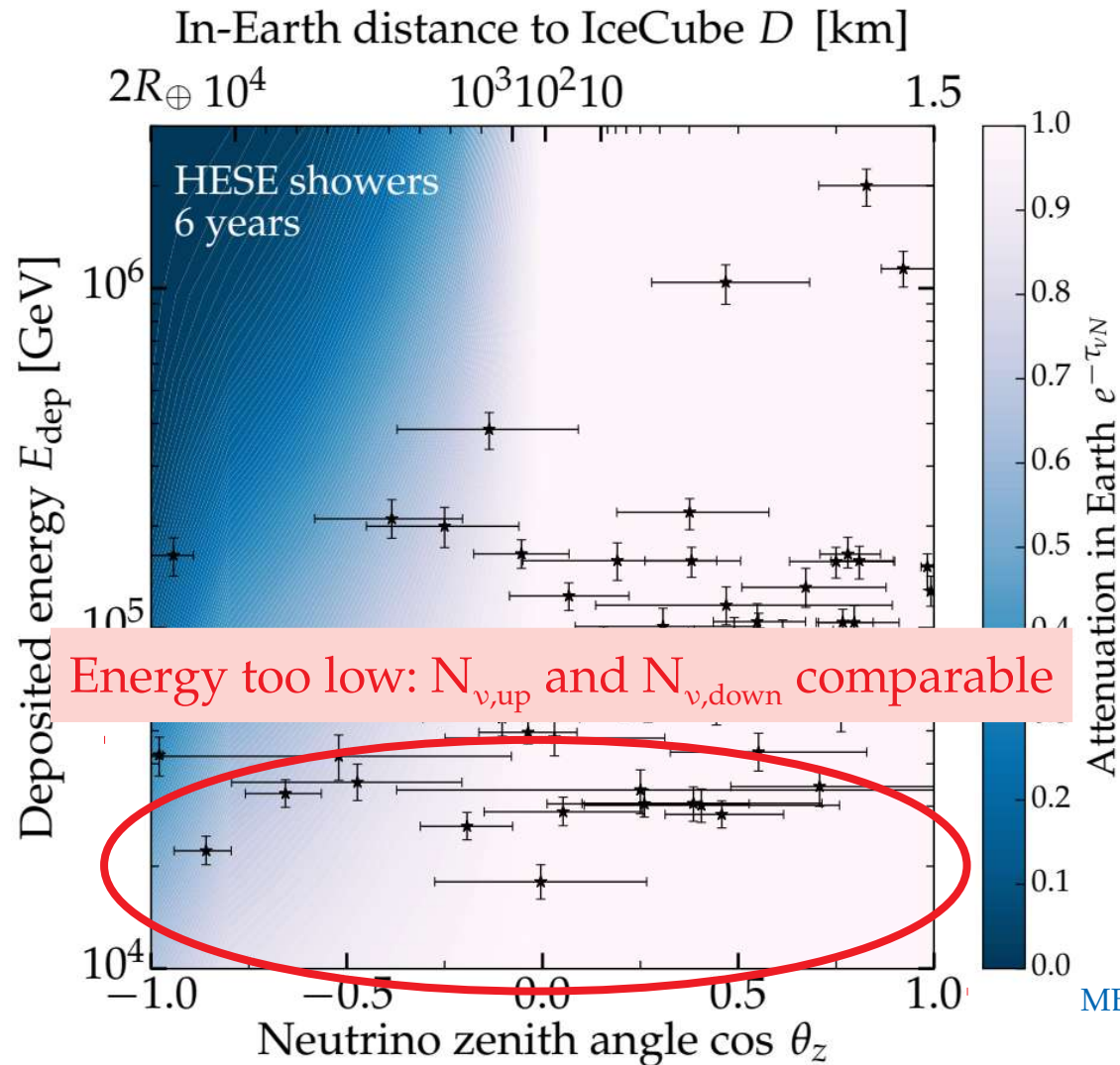


Downgoing events constrain (flux x cross section)

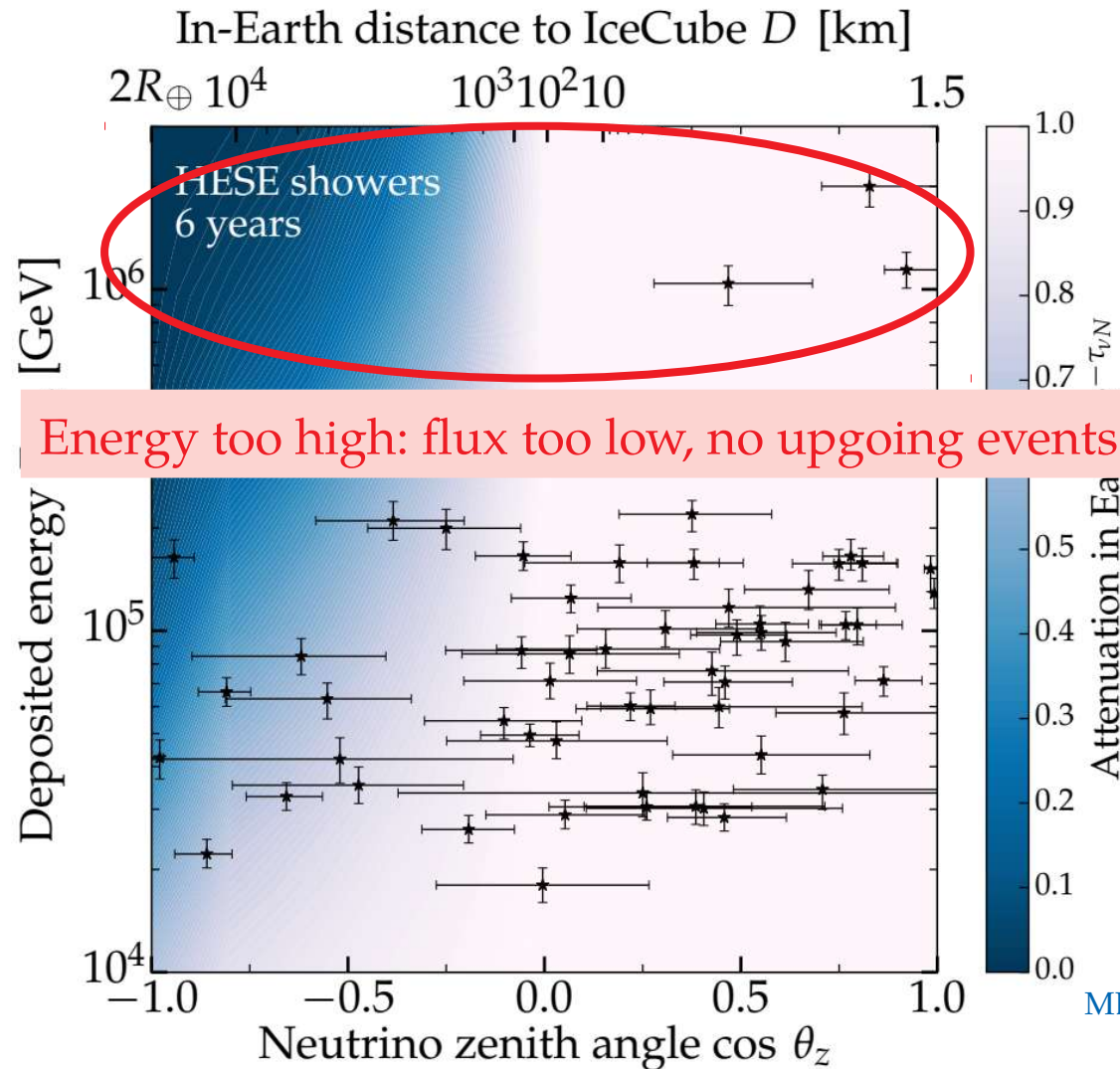
MB & A. Connolly, 1711.11043



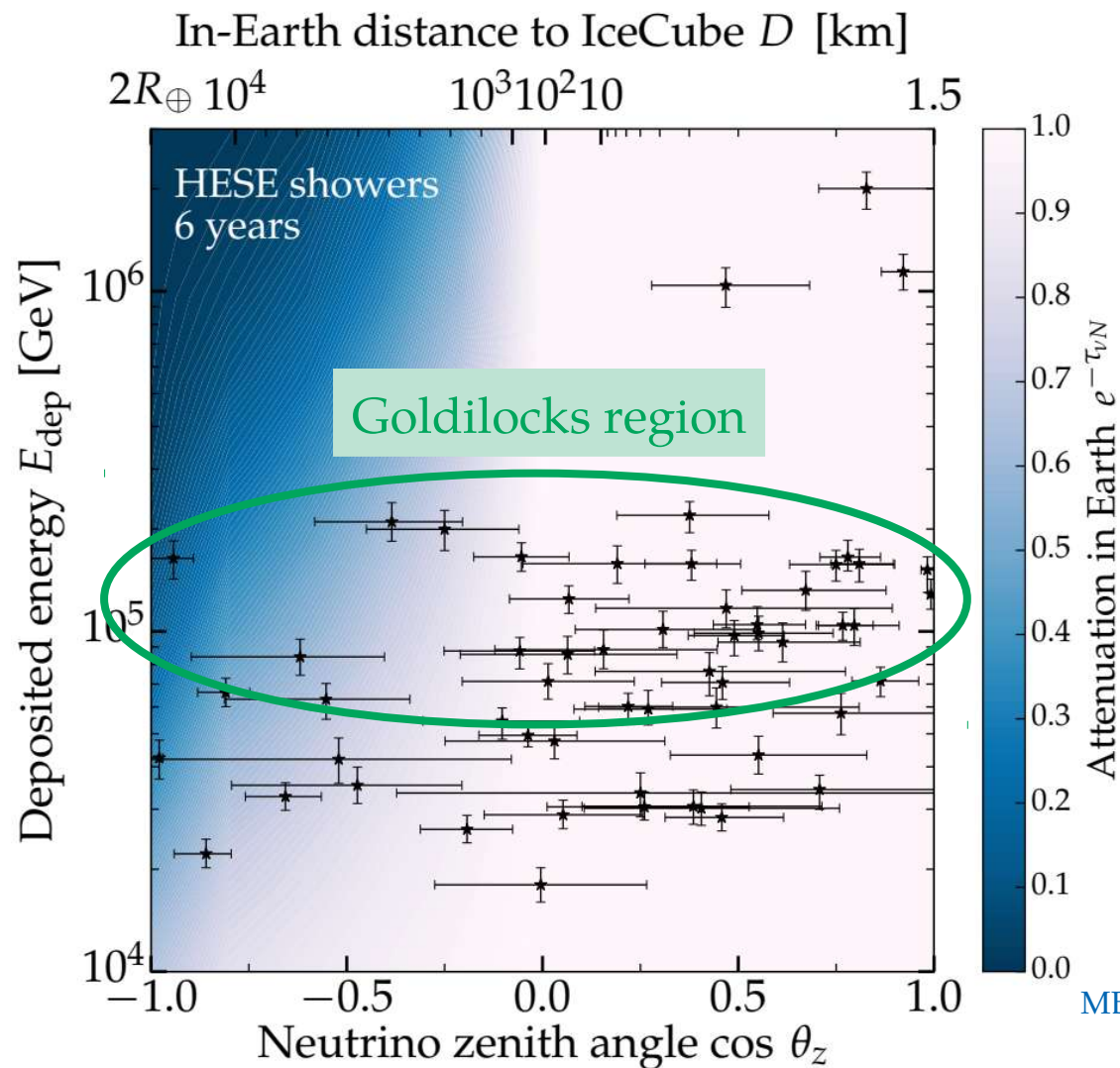
MB & A. Connolly, 1711.11043



MB & A. Connolly, 1711.11043

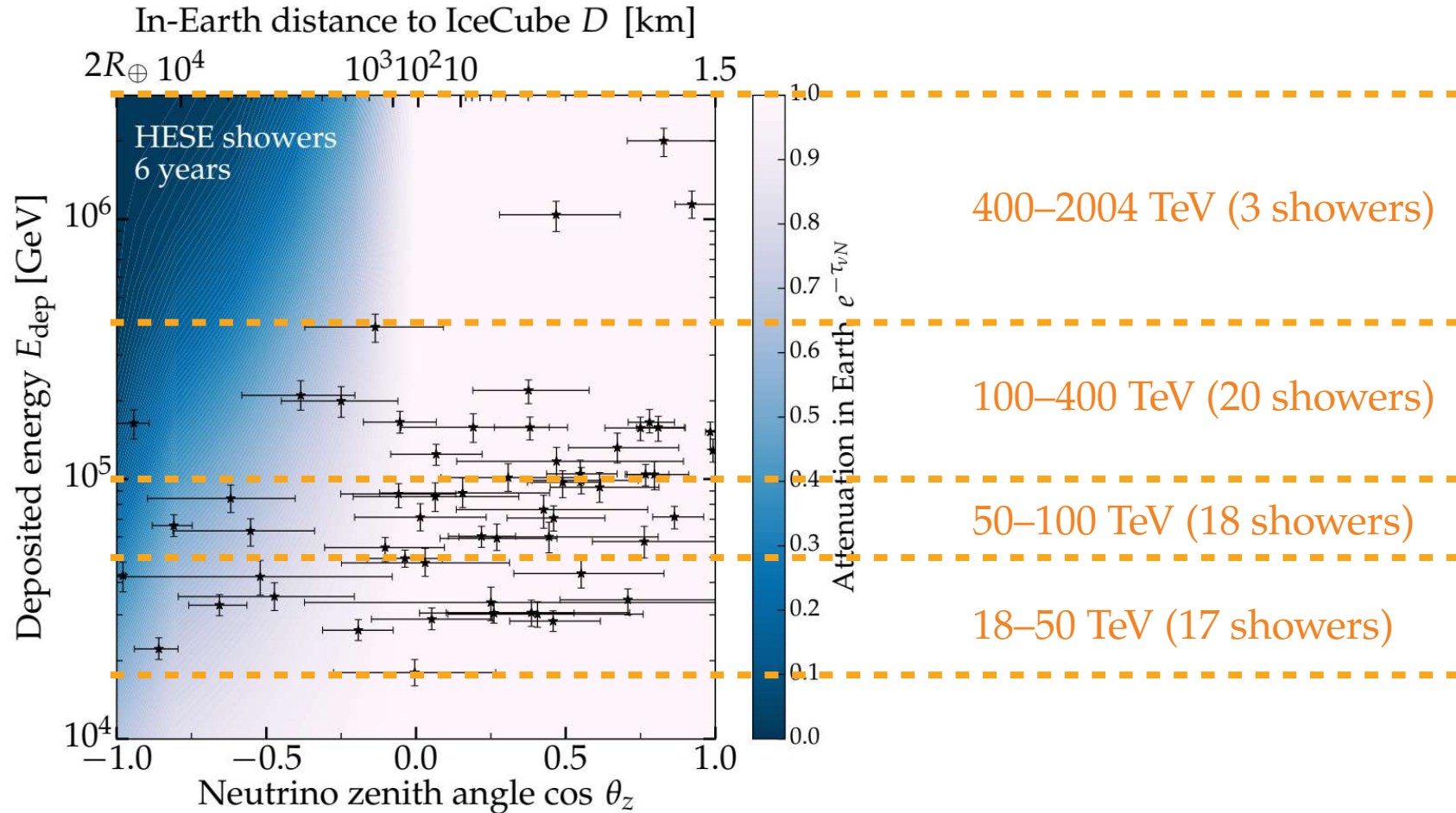


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Bin-by-bin analysis



Sensitivity to σ in each bin

Number of contained events in an energy bin:

$$N_\nu \sim \Phi_\nu \cdot \sigma_{\nu N} \cdot e^{-\tau} = \Phi_\nu \cdot \sigma_{\nu N} \cdot e^{-L\sigma_{\nu N}n_N}$$

Downgoing (no matter)

$$N_{\nu,\text{dn}} \sim \Phi_\nu \cdot \sigma_{\nu N}$$

Downgoing events fix the product $\Phi_\nu \cdot \sigma_{\nu N}$

Upgoing (lots of matter)

$$N_{\nu,\text{up}} \sim N_{\nu,\text{dn}} \cdot e^{-\tau}$$

Upgoing events measure $\sigma_{\nu N}$ via τ

Reality check:

Few events (per energy bin), so we are statistics-limited

The fine print

- ▶ High-energy ν 's: astrophysical (isotropic) + atmospheric (**anisotropic**)
 - ↳ We take into account the shape of the atmospheric contribution
- ▶ The shape of the astrophysical ν **energy spectrum** is still uncertain
 - ↳ We take a $E^{-\gamma}$ spectrum in *narrow* energy bins
- ▶ **NC showers** are sub-dominant to **CC showers**, but they are indistinguishable
 - ↳ Following Standard-Model predictions, we take $\sigma_{\text{NC}} = \sigma_{\text{CC}}/3$
- ▶ IceCube does not **distinguish ν from $\bar{\nu}$** , and their cross-sections are different
 - ↳ We assume equal fluxes, expected from production via pp collisions
 - ↳ We assume the avg. ratio $\langle \sigma_{\bar{\nu}\text{N}} / \sigma_{\nu\text{N}} \rangle$ in each bin known, from SM predictions
- ▶ The **flavor composition** of astrophysical neutrinos is still uncertain
 - ↳ We assume equal flux of each flavor, compatible with theory and observations

What goes into the (likelihood) mix?

- ▶ Inside each energy bin, we freely vary
 - ▶ N_{ast} (showers from astrophysical neutrinos)
 - ▶ N_{atm} (showers from atmospheric neutrinos)
 - ▶ γ (astrophysical spectral index)
 - ▶ σ_{CC} (neutrino-nucleon charged-current cross section)
- ▶ For each combination, we generate the angular and energy shower spectrum...
- ▶ ... and compare it to the observed HESE spectrum via a likelihood
- ▶ Maximum likelihood yields σ_{CC} (marginalized over nuisance parameters)
- ▶ Bins are independent of each other – there are no (significant) cross-bin correlations

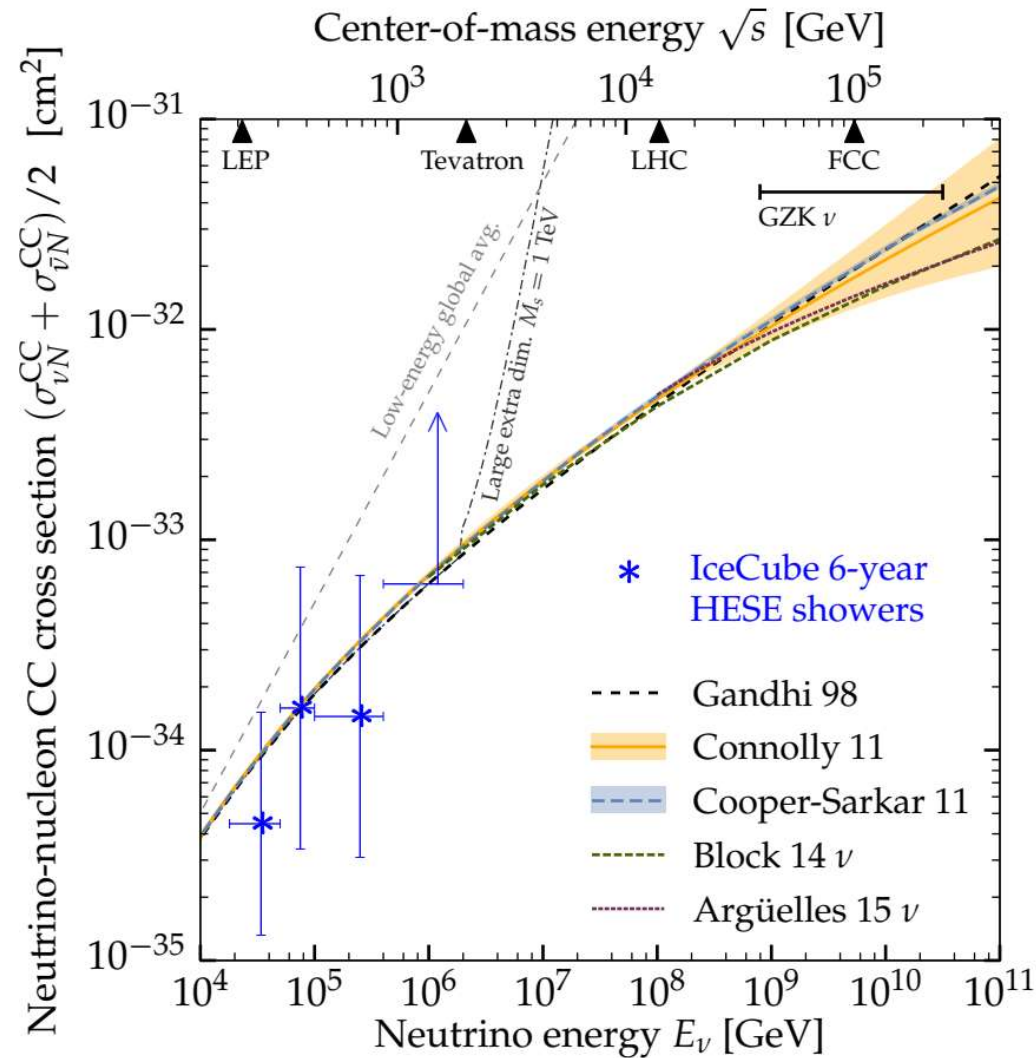
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Including detector resolution
(10% in energy, 15° in direction)

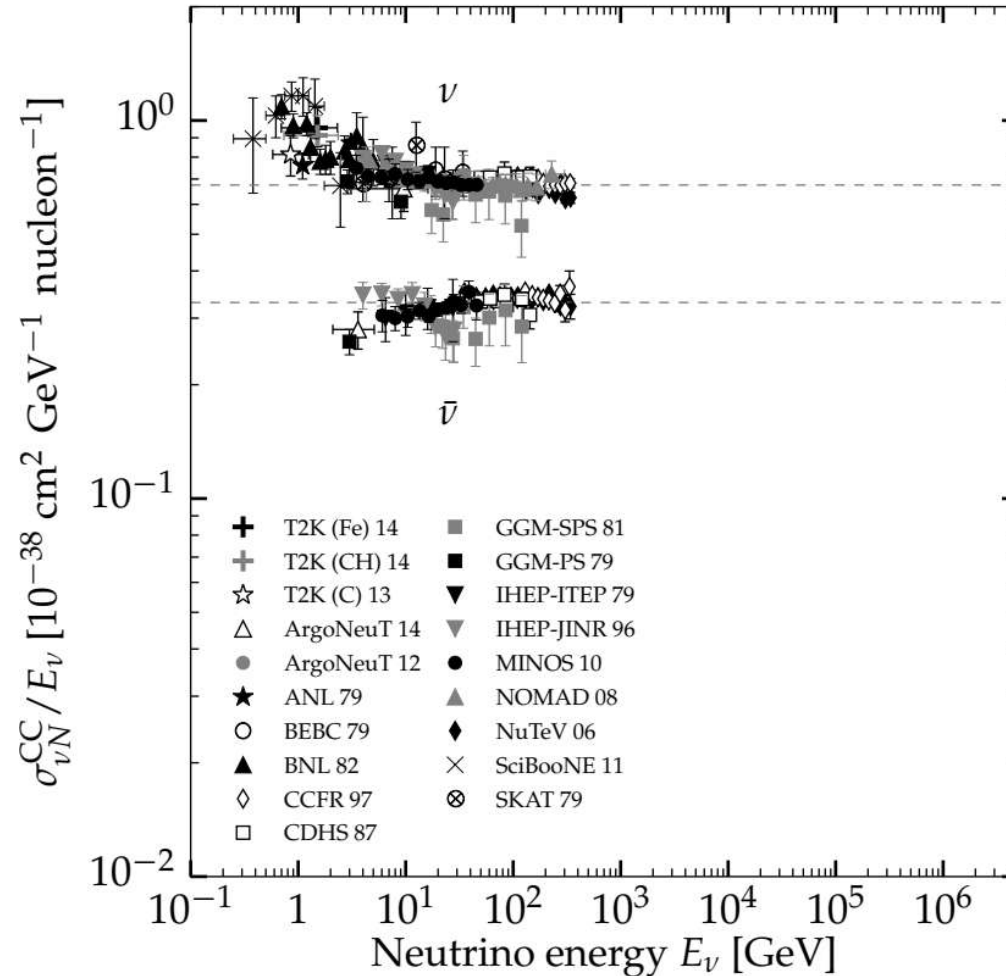


Our result



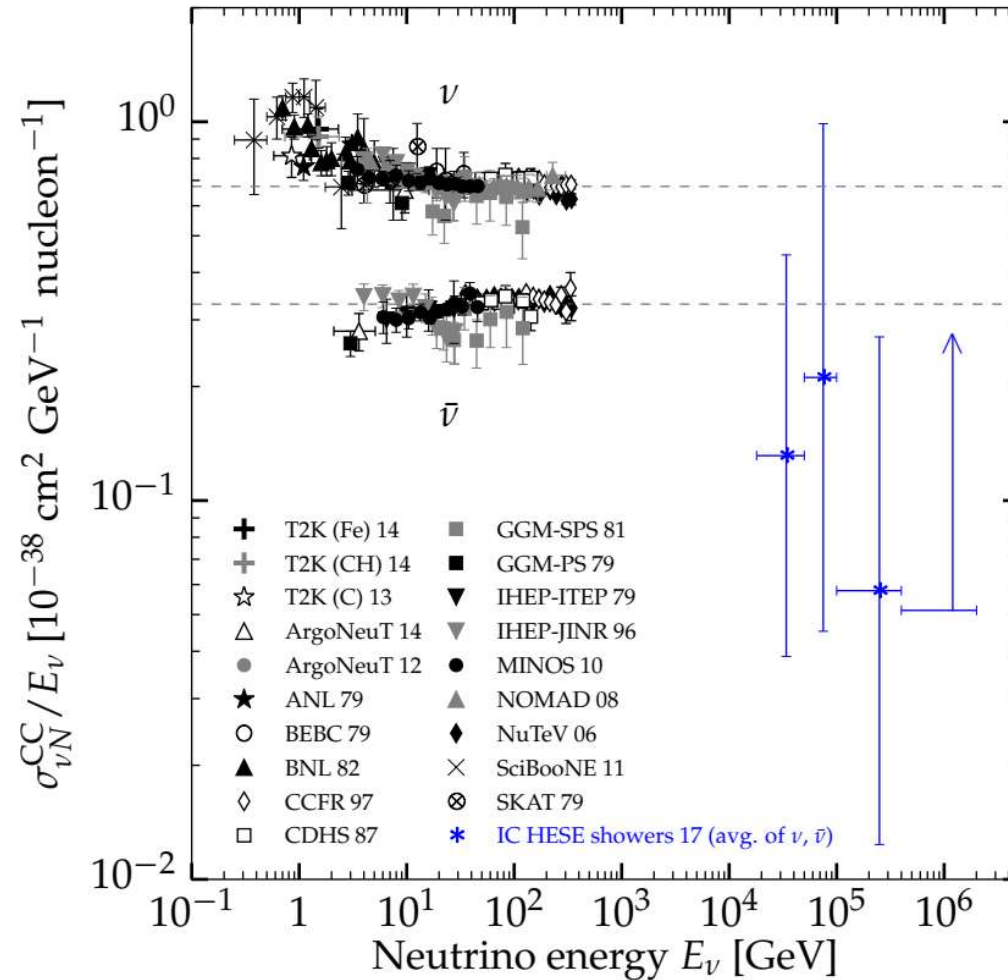
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Extending cross section measurements



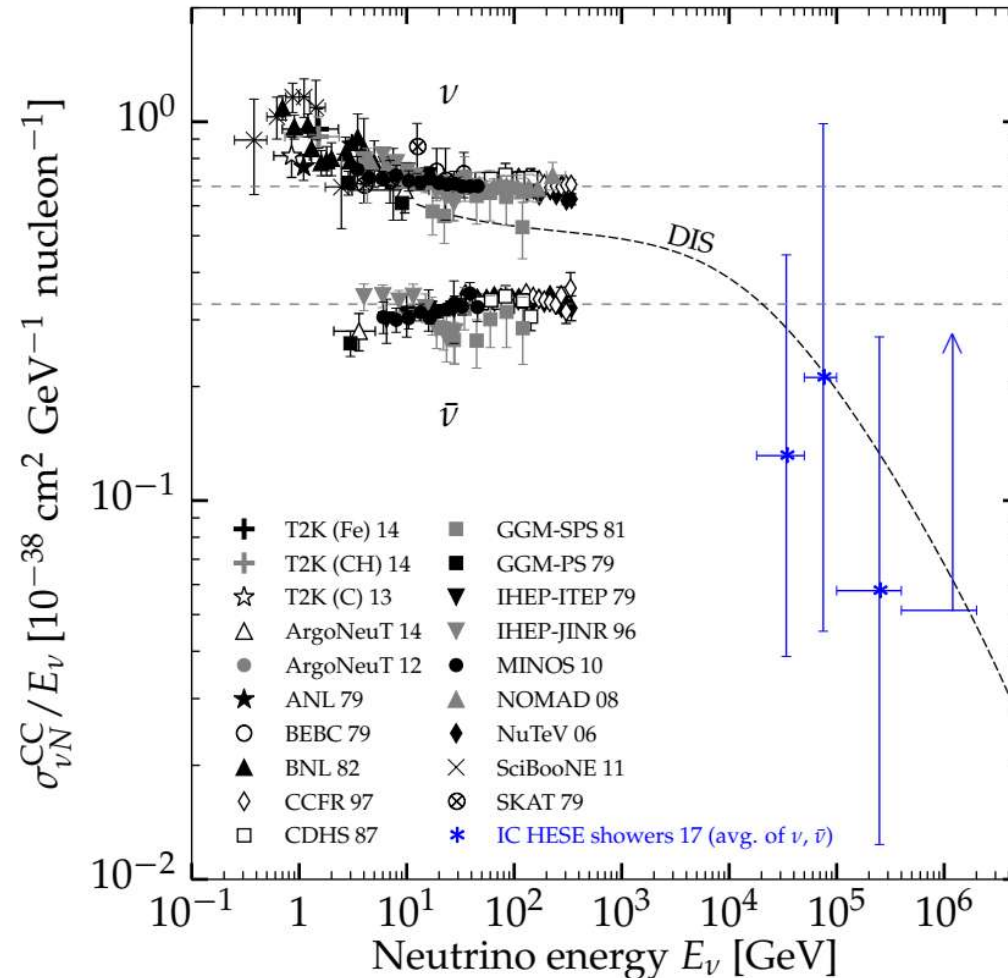
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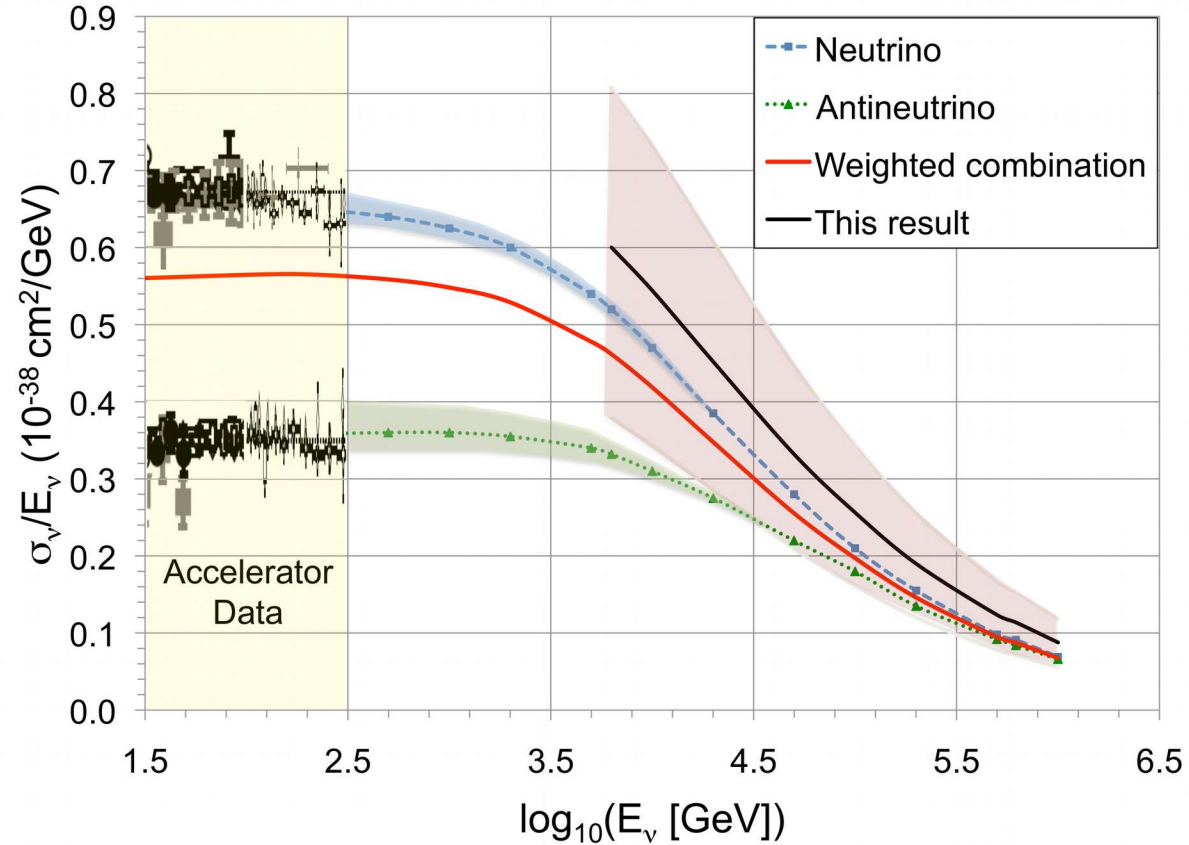
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How to do better / more?

- ▶ Currently, we are statistics-limited
 - ↳ Solvable with more data from IceCube, IceCube-Gen2, KM3NeT
- ▶ Large errors in arrival direction ($\sim 10^\circ$) give errors in attenuation
 - ↳ Solvable with ongoing IceCube improvements + KM3NeT
- ▶ Charged-current + neutral-current cross sections are indistinguishable
 - ↳ Solvable (?) with muon and neutron echoes (Li, MB, Beacom 16)
- ▶ Cannot separate ν from $\bar{\nu}$
 - ↳ Wait to detect Glashow resonance (~ 6.3 PeV), sensitive only to $\bar{\nu}_e$
- ▶ Use starting tracks / through-going muons
 - ↳ Doable / done by IceCube (more next)

Using through-going muons instead

- ▶ Use $\sim 10^4$ through-going muons
- ▶ Measured: dE_μ/dx
- ▶ Inferred: $E_\mu \approx dE_\mu/dx$
- ▶ From simulations (uncertain):
most likely E_ν given E_μ
- ▶ Fit the ratio $\sigma_{\text{obs}}/\sigma_{\text{SM}}$
 $1.30^{+0.21}_{-0.19} \text{ (stat.) } ^{+0.39}_{-0.43} \text{ (syst.)}$
- ▶ All events grouped in a single
energy bin 6–980 TeV



IceCube, Nature 2017

Bonus: Measuring the inelasticity $\langle y \rangle$

- ▶ Inelasticity in CC ν_μ interaction $\nu_\mu + N \rightarrow \mu + X$:

$$E_X = y E_\nu \quad \text{and} \quad E_\mu = (1-y) E_\nu \Rightarrow y = (1 + E_\mu/E_X)^{-1}$$

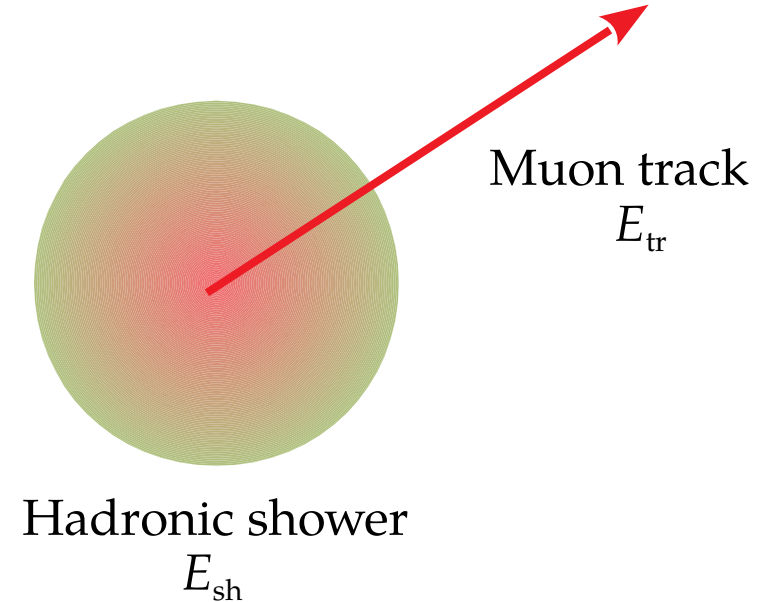
- ▶ The value of y follows a distribution $d\sigma/dy$

- ▶ In a HESE starting track:

$$\left. \begin{array}{l} E_X = E_{\text{sh}} \text{ (energy of shower)} \\ E_\mu = E_{\text{tr}} \text{ (energy of track)} \end{array} \right\} y = (1 + E_{\text{tr}}/E_{\text{sh}})^{-1}$$

- ▶ New IceCube analysis:

- ▶ 5 years of starting-track data (2650 tracks)
- ▶ Machine learning separates shower from track
- ▶ Different y distributions for ν and $\bar{\nu}$



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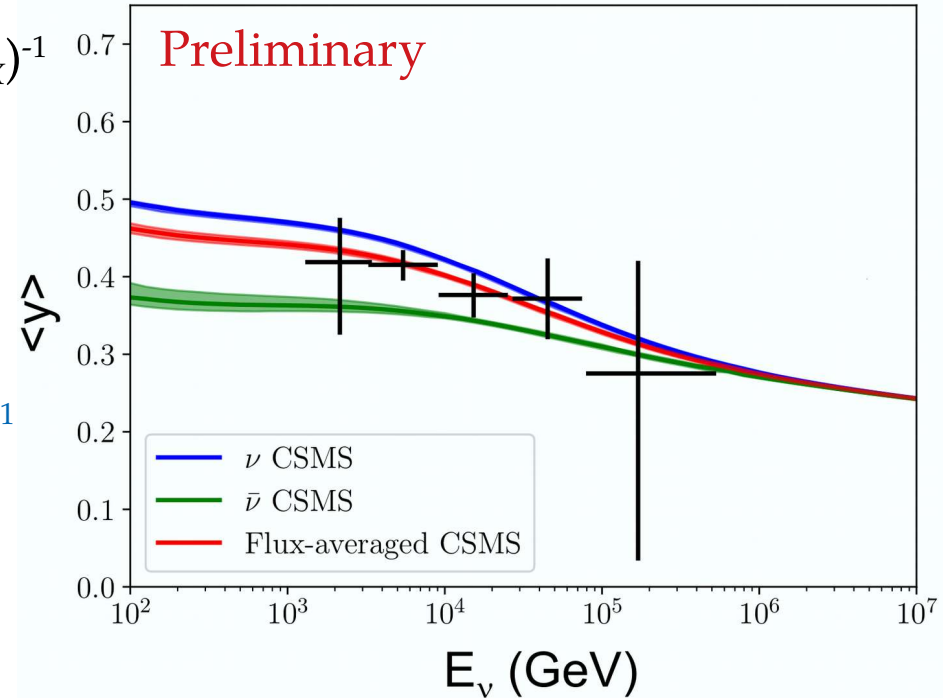
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See presentation by Spencer Klein, ISVHECRI 2018

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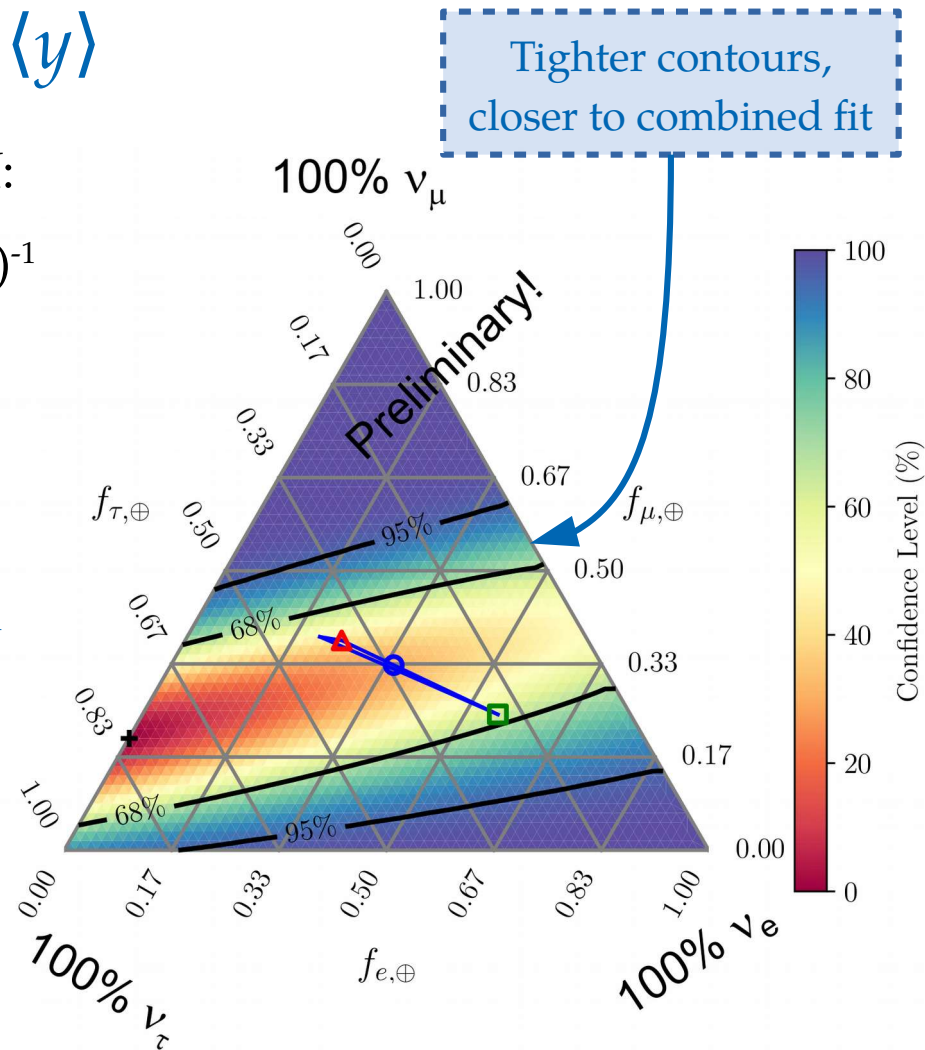
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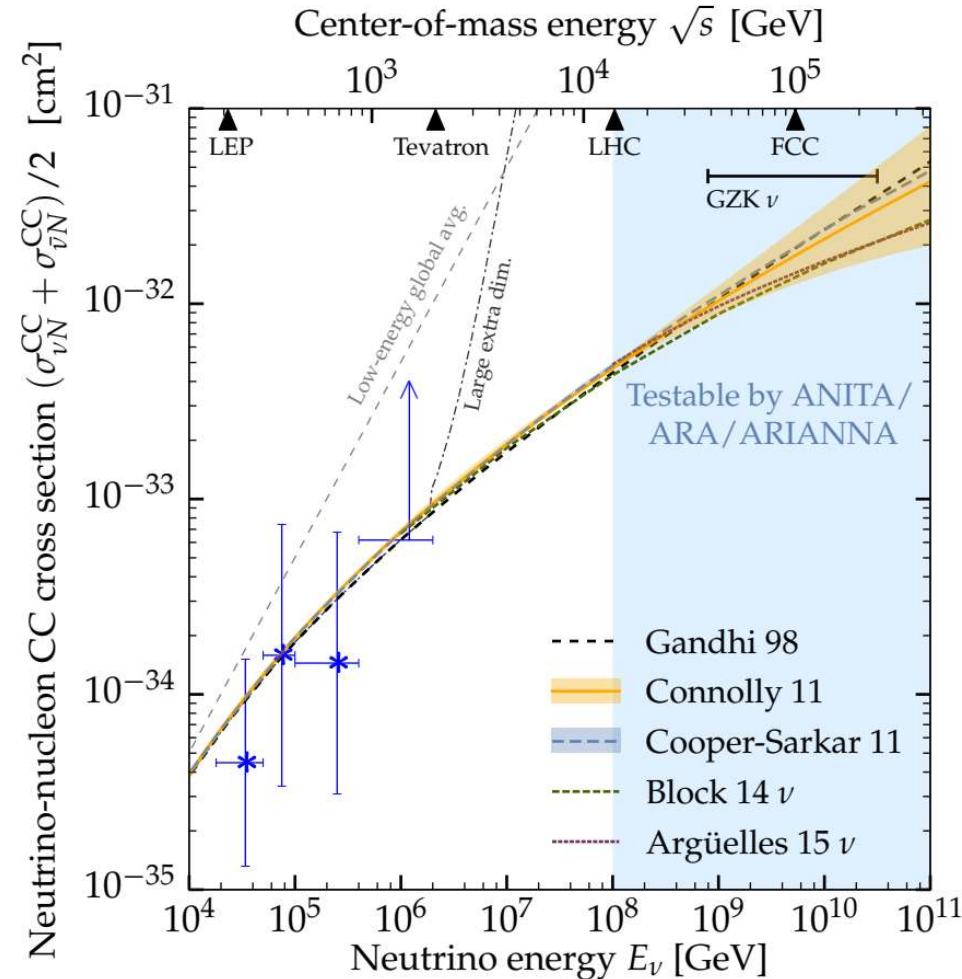
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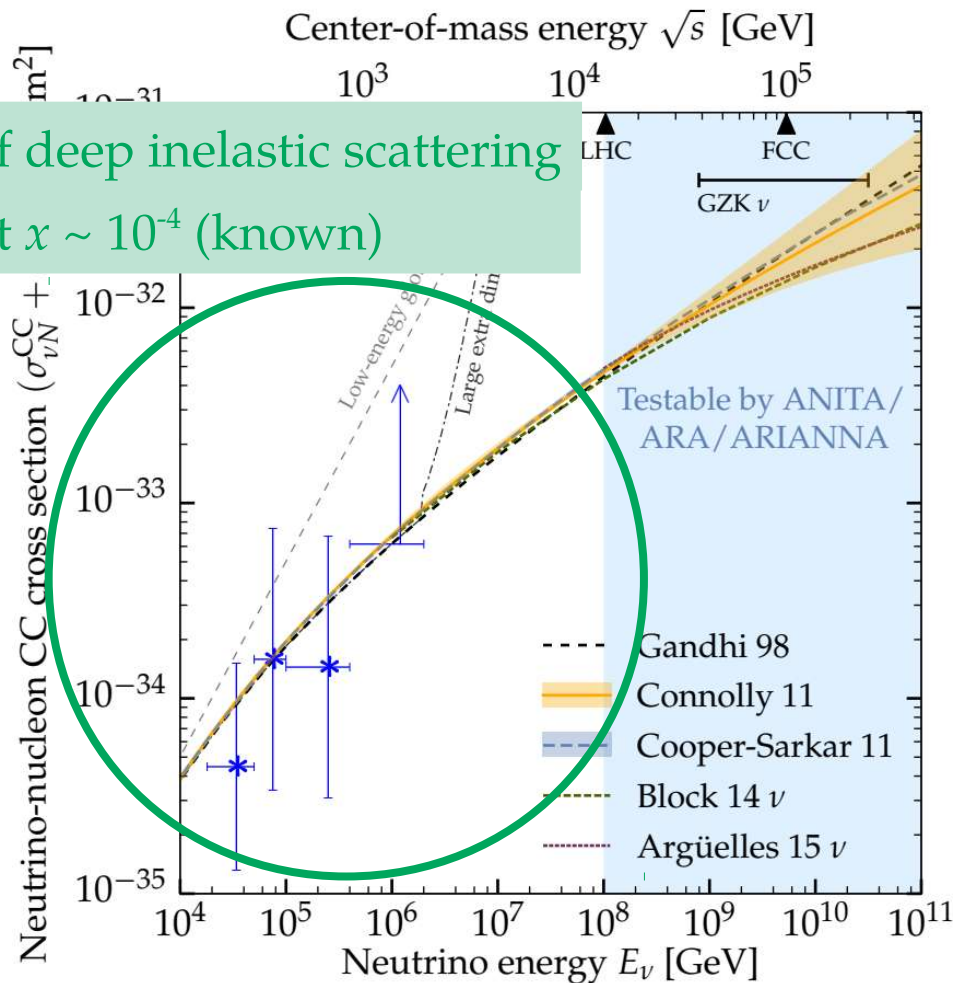
See presentation by Spencer Klein, ISVHECRI 2018

Quo vadis: IceCube vs. ANITA/ARA/ARIANNA

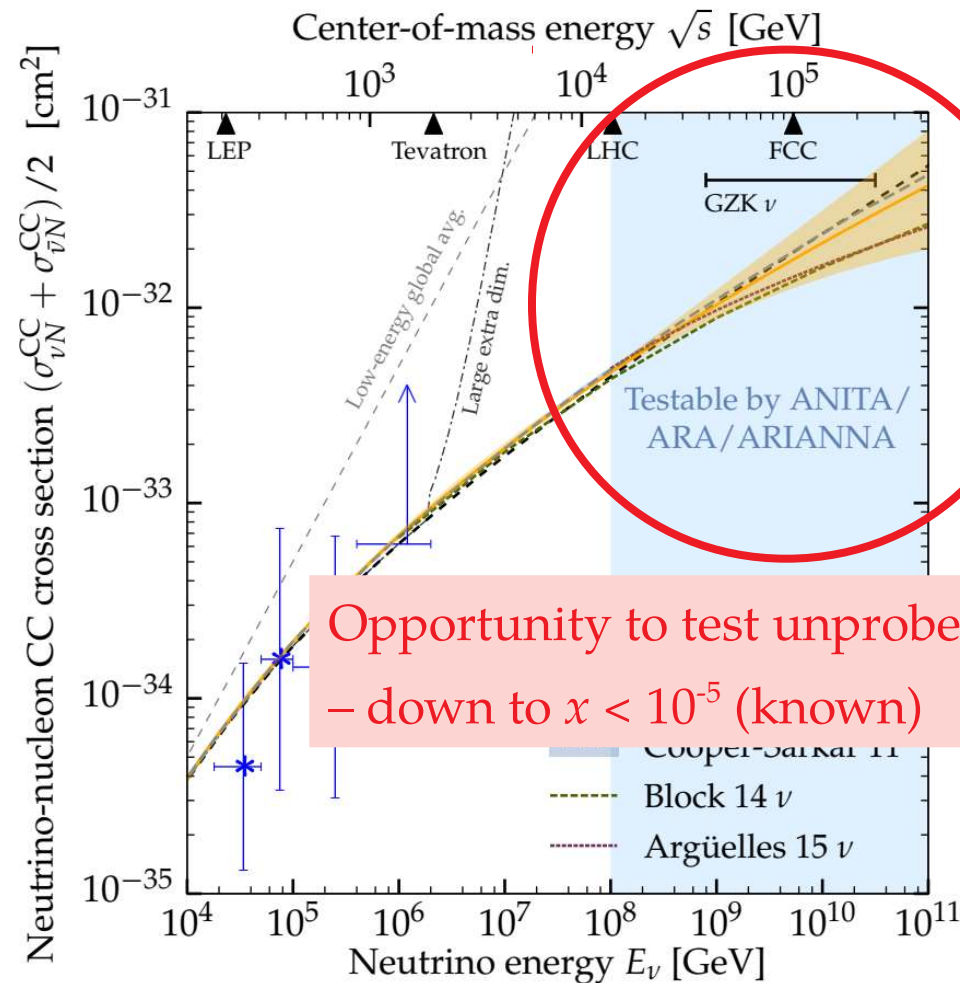


Quo vadis: IceCube vs. ANITA/ARA/ARIANNA

Test predictions of deep inelastic scattering
– down to PDFs at $x \sim 10^{-4}$ (known)



Quo vadis: IceCube vs. ANITA/ARA/ARIANNA



Opportunity to test unprobed nucleon structure
– down to $x < 10^{-5}$ (known)

Summary

- ▶ We extracted the neutrino-nucleon cross section from 18 TeV to 2 PeV
 - ↳ Previously known up to 350 GeV
- ▶ Found consistency with Standard-Model predictions
- ▶ Errors are still large due to statistics and astrophysical unknowns
- ▶ But both will be improved in the future

Already today, neutrino telescopes
are probing fundamental particle physics

Backup slides

Marginalized cross section in each bin

TABLE I. Neutrino-nucleon charged-current inclusive cross sections, averaged between neutrinos ($\sigma_{\nu N}^{\text{CC}}$) and anti-neutrinos ($\sigma_{\bar{\nu} N}^{\text{CC}}$), extracted from 6 years of IceCube HESE showers. To obtain these results, we fixed $\sigma_{\bar{\nu} N}^{\text{CC}} = \langle \sigma_{\bar{\nu} N}^{\text{CC}} / \sigma_{\nu N}^{\text{CC}} \rangle \cdot \sigma_{\nu N}^{\text{CC}}$ — where $\langle \sigma_{\bar{\nu} N}^{\text{CC}} / \sigma_{\nu N}^{\text{CC}} \rangle$ is the average ratio of $\bar{\nu}$ to ν cross sections calculated using the standard prediction from Ref. [60] — and $\sigma_{\nu N}^{\text{NC}} = \sigma_{\nu N}^{\text{CC}}/3$, $\sigma_{\bar{\nu} N}^{\text{NC}} = \sigma_{\bar{\nu} N}^{\text{CC}}/3$. Uncertainties are statistical plus systematic, added in quadrature.

E_ν [TeV]	$\langle E_\nu \rangle$ [TeV]	$\langle \sigma_{\bar{\nu} N}^{\text{CC}} / \sigma_{\nu N}^{\text{CC}} \rangle$	$\log_{10}[\frac{1}{2}(\sigma_{\nu N}^{\text{CC}} + \sigma_{\bar{\nu} N}^{\text{CC}})/\text{cm}^2]$
18–50	32	0.752	-34.35 ± 0.53
50–100	75	0.825	-33.80 ± 0.67
100–400	250	0.888	-33.84 ± 0.67
400–2004	1202	0.957	$> -33.21 (1\sigma)$

Neutrino zenith angle distribution

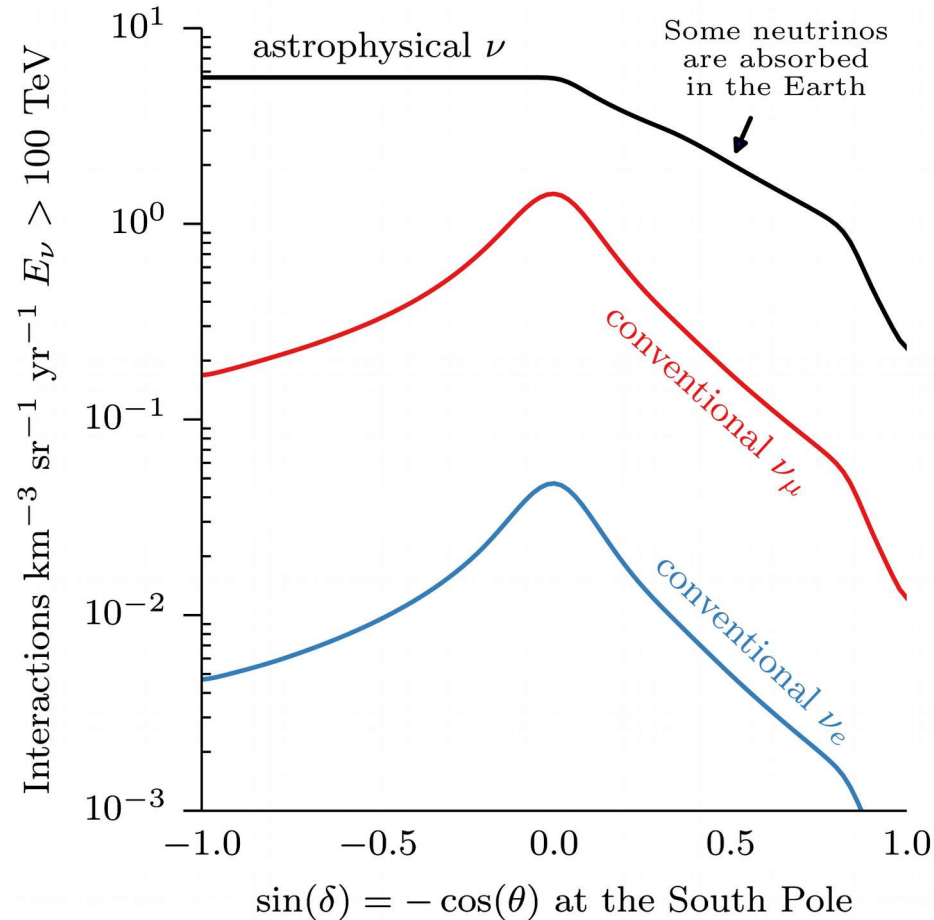
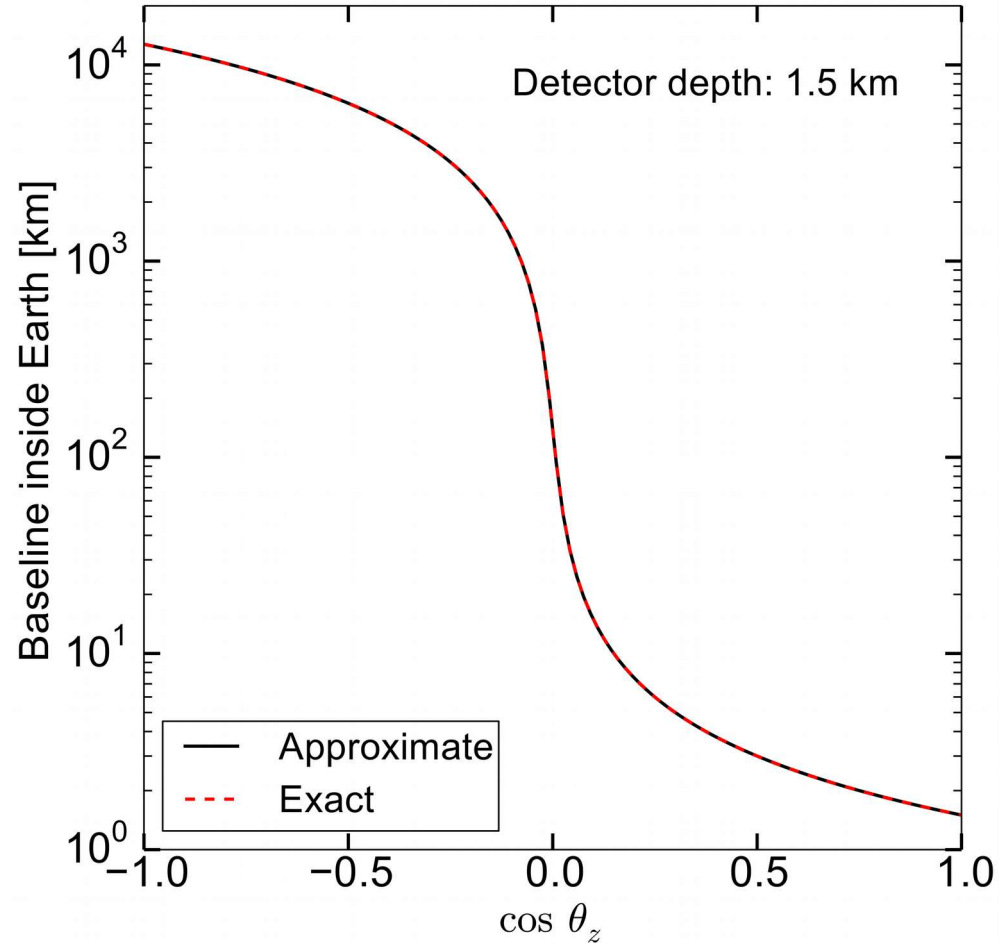
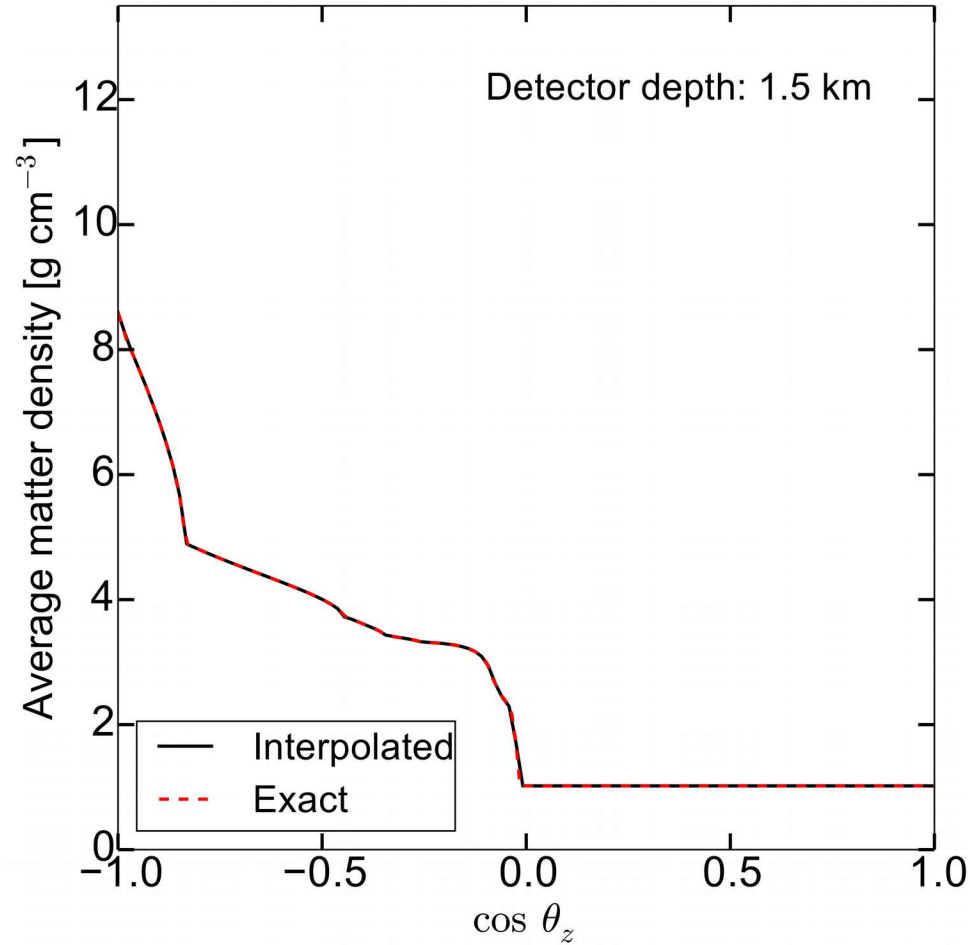


Figure by
Jakob Van Santen
ICRC 2017

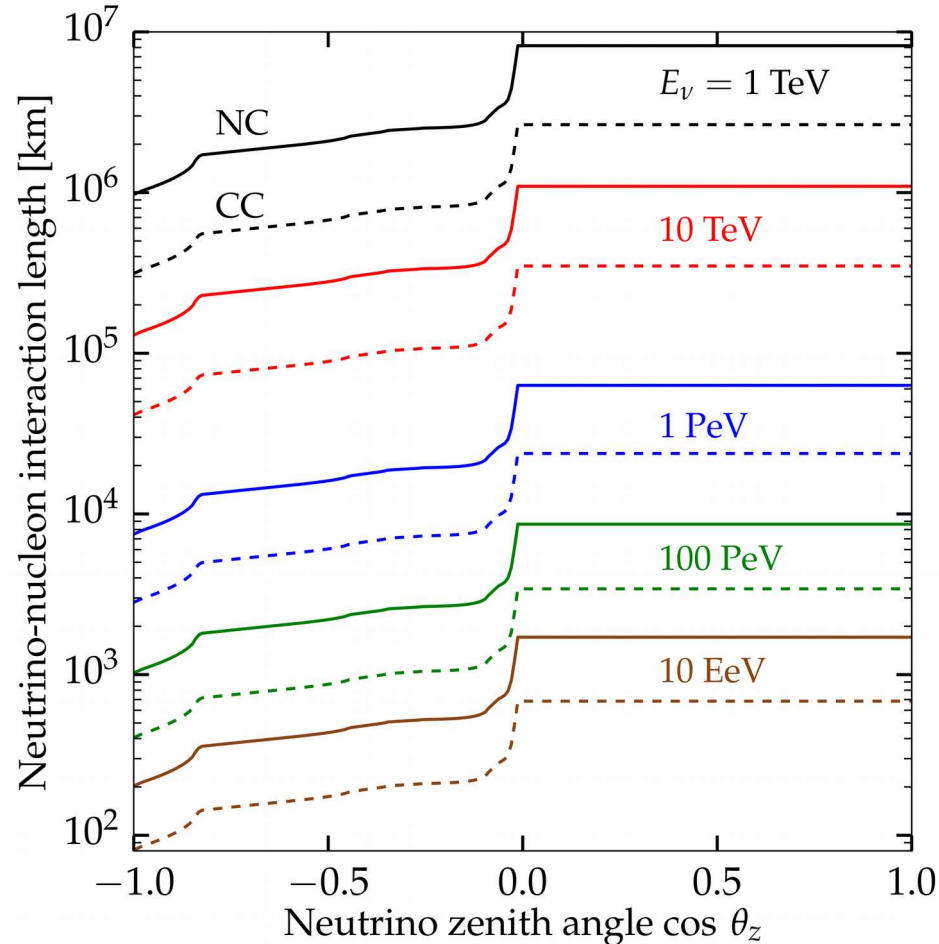
Neutrino baseline inside the Earth



Average Earth matter density



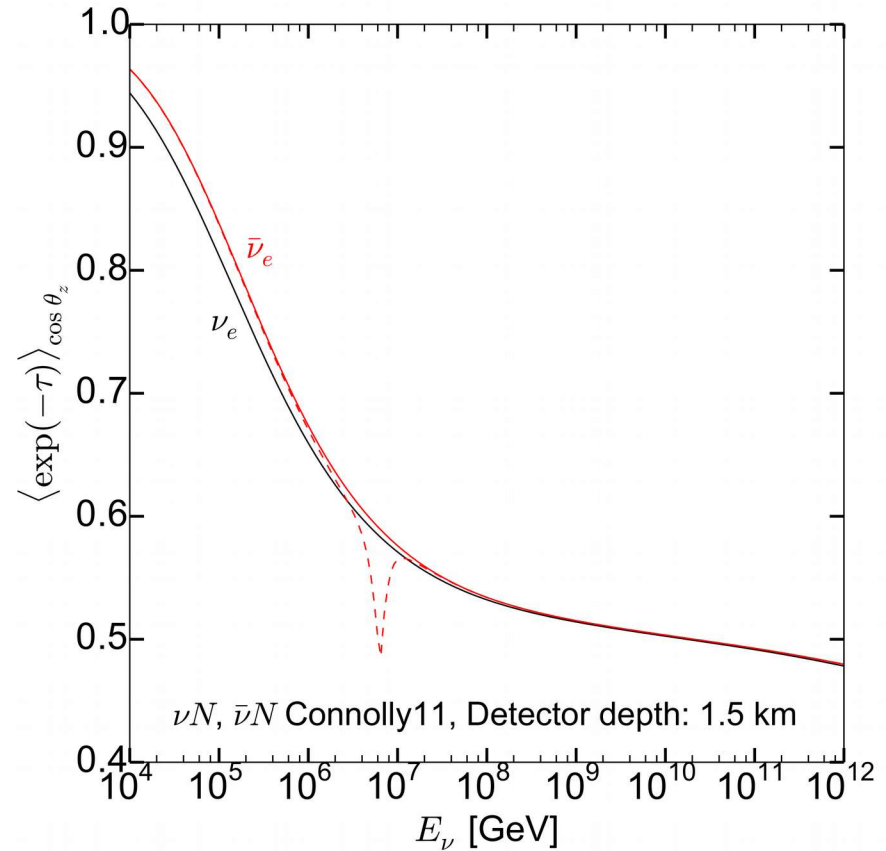
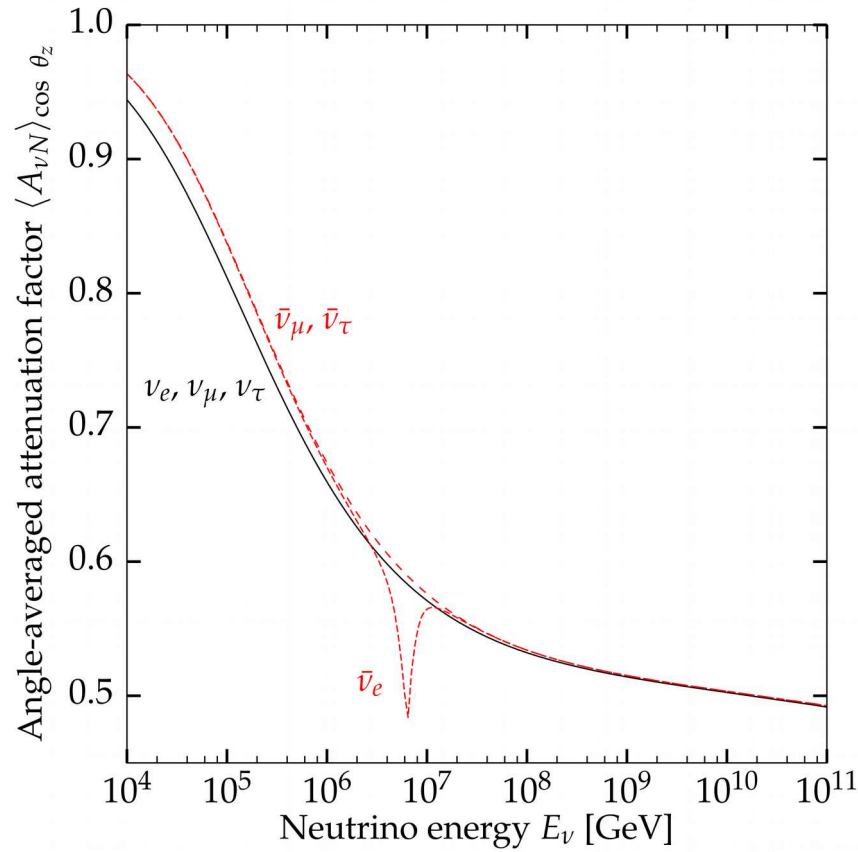
$\bar{\nu}N$ interaction length inside the Earth



The $\bar{\nu}N$ interaction length
is $\sim 60\%$ larger

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Angle-averaged absorption inside the Earth



Energy and angular shower spectra

Rate from all flavors, CC + NC:

$$\frac{d^2 N_{\text{sh}}}{dE_{\text{sh}} d \cos \theta_z} = \frac{d^2 N_{\text{sh},e}^{\text{CC}}}{dE_{\text{sh}} d \cos \theta_z} + \text{Br}_{\tau \rightarrow \text{sh}} \frac{d^2 N_{\text{sh},\tau}^{\text{CC}}}{dE_{\text{sh}} d \cos \theta_z} + \sum_{l=e,\mu,\tau} \frac{d^2 N_{\text{sh},l}^{\text{NC}}}{dE_{\text{sh}} d \cos \theta_z}$$

$= 0.83$

Contribution from one flavor CC:

$$\frac{d^2 N_{\text{sh},l}^{\text{CC}}}{dE_{\text{sh}} d \cos \theta_z}(E_{\text{sh}}, \cos \theta_z) \simeq -2\pi \rho_{\text{ice}} N_A V T \left\{ \Phi_l(E_\nu) \sigma_{\nu N}^{\text{CC}}(E_\nu) e^{-\tau_{\nu N}(E_\nu, \theta_z)} + \Phi_{\bar{l}}(E_\nu) \sigma_{\bar{\nu} N}^{\text{CC}}(E_\nu) e^{-\tau_{\bar{\nu} N}(E_\nu, \theta_z)} \right\} \Big|_{E_\nu = E_{\text{sh}}/f_{l,\text{CC}}}$$

Conversion between shower energy and neutrino energy:

$$f_{l,t} \equiv \frac{E_{\text{sh}}}{E_\nu} \simeq \begin{cases} 1 & \text{for } l = e \text{ and } t = \text{CC} \\ [\langle y \rangle + 0.7(1 - \langle y \rangle)] \simeq 0.8 & \text{for } l = \tau \text{ and } t = \text{CC} \\ \langle y \rangle \simeq 0.25 & \text{for } l = e, \mu, \tau \text{ and } t = \text{NC} \end{cases}$$

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Detector resolution

Number of contained showers:

$$\frac{d^2 N_{\text{sh}}}{dE_{\text{dep}} d\cos\theta_z} = \int dE_{\text{sh}} \int d\cos\theta'_z \frac{d^2 N_{\text{sh}}}{dE_{\text{sh}} d\cos\theta'_z} R_E(E_{\text{sh}}, E_{\text{dep}}, \sigma_E(E_{\text{sh}})) R_\theta(\cos\theta'_z, \cos\theta_z, \sigma_{\cos\theta_z})$$

Energy resolution: [Palomares-Ruiz, Vincent, Mena *PRD* 2015; Vincent, Palomares-Ruiz, Mena *PRD* 2016; MB, Beacom, Murase, *PRD* 2016]

$$R_E(E_{\text{sh}}, E_{\text{dep}}, \sigma_E(E_{\text{sh}})) = \frac{1}{\sqrt{2\pi\sigma_E^2(E_{\text{sh}})}} \exp\left[-\frac{(E_{\text{sh}} - E_{\text{dep}})^2}{2\sigma_E^2(E_{\text{sh}})}\right] \quad \text{with } \sigma_E(E_{\text{sh}}) = 0.1E_{\text{sh}}$$

IceCube, *JINST* 2014

Angular resolution:

$$R_\theta(\cos\theta'_z, \cos\theta_z, \sigma_{\cos\theta_z}) = \frac{1}{\sqrt{2\pi\sigma_{\cos\theta_z}^2}} \exp\left[-\frac{(\cos\theta'_z - \cos\theta_z)^2}{2\sigma_{\cos\theta_z}^2}\right]$$

$$\text{with } \sigma_{\cos\theta_z} \equiv \frac{1}{2} [|\cos(\theta_z + \sigma_{\theta_z}) - \cos\theta_z| + |\cos(\theta_z - \sigma_{\theta_z}) - \cos\theta_z|] \quad \text{and } \sigma_{\theta_z} = 15^\circ$$

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Likelihood

In an energy bin containing $N_{\text{sh}}^{\text{obs}}$ observed showers, the likelihood is

Each energy bin is independent

$$\mathcal{L} = \frac{e^{-(N_{\text{sh}}^{\text{atm}} + N_{\text{sh}}^{\text{ast}})}}{N_{\text{sh}}^{\text{obs}}!} \prod_{i=1}^{N_{\text{sh}}^{\text{obs}}} \mathcal{L}_i$$

Partial likelihood, *i.e.*, relative probability of the i -th shower being from an atmospheric neutrino or an astrophysical neutrino:

Depends on $\sigma_{\nu N}$

$$\mathcal{L}_i = N_{\text{sh}}^{\text{atm}} \mathcal{P}_i^{\text{atm}} + N_{\text{sh}}^{\text{ast}} \mathcal{P}_i^{\text{ast}}$$

$$\mathcal{P}_i^{\text{atm}} = \left(\int_{E_{\text{dep}}^{\text{min}}}^{E_{\text{dep}}^{\text{max}}} dE_{\text{dep}} \int_{-1}^1 d \cos \theta_z \frac{d^2 N_{\text{sh}}^{\text{atm}}}{dE_{\text{dep}} d \cos \theta_z} \right)^{-1} \left(\frac{d^2 N_{\text{sh}}^{\text{atm}}}{dE_{\text{dep}} d \cos \theta_z} \Big|_{E_{\text{dep},i}, \cos \theta_{z,i}} \right)$$

PDF for this shower to be made by an atmospheric ν

$$\mathcal{P}_i^{\text{ast}} = \left(\int_{E_{\text{dep}}^{\text{min}}}^{E_{\text{dep}}^{\text{max}}} dE_{\text{dep}} \int_{-1}^1 d \cos \theta_z \frac{d^2 N_{\text{sh}}^{\text{ast}}}{dE_{\text{dep}} d \cos \theta_z} \right)^{-1} \left(\frac{d^2 N_{\text{sh}}^{\text{ast}}}{dE_{\text{dep}} d \cos \theta_z} \Big|_{E_{\text{dep},i}, \cos \theta_{z,i}} \right)$$

PDF for this shower to be made by an astrophysical ν

Depends on γ and $\sigma_{\nu N}$

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See also: Palomares-Ruiz, Vincent, Mena *PRD* 2015; Vincent, Palomares-Ruiz, Mena *PRD* 2016

Mauricio Bustamante (Niels Bohr Institute)

Best-fit values and uncertainties

TABLE II. Best-fit values and 1σ uncertainties of the nuisance parameters in each energy bin: number of showers due to atmospheric neutrinos $N_{\text{sh}}^{\text{atm}}$, number of showers due to astrophysical neutrinos $N_{\text{sh}}^{\text{ast}}$, and astrophysical spectral index γ .

E_ν [TeV]	$N_{\text{sh}}^{\text{atm}}$	$N_{\text{sh}}^{\text{ast}}$	γ
18–50	4.2 ± 4.9	11.4 ± 3.5	2.38 ± 0.31
50–100	6.3 ± 5.3	11.7 ± 4.5	2.43 ± 0.31
100–400	6.4 ± 6.0	12.9 ± 5.2	2.49 ± 0.31
400–2004	1.2 ± 1.0	1.73 ± 0.89	2.37 ± 0.32

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