

# **Searching eV-scale sterile neutrino in Atmospheric Neutrino Experiments**

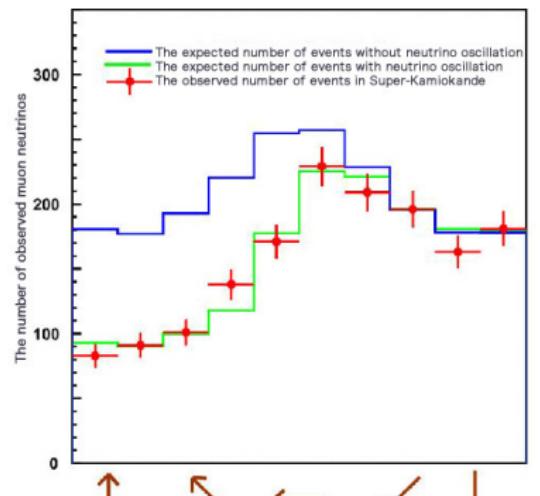
**Jordi Salvado**

Advanced Workshop on Physics of Atmospheric Neutrinos - PANE



UNIVERSITAT DE  
BARCELONA

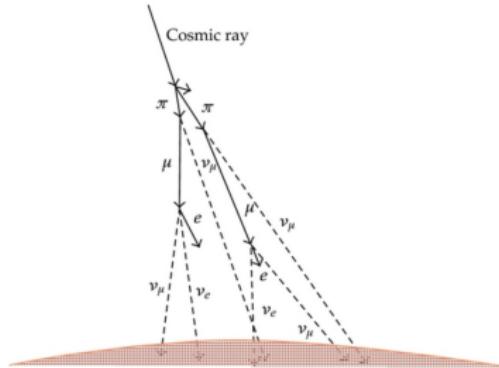
# Atmospheric neutrinos were crucial (announced 1998)



Upward going Neutrinos  
Flight length: 12800km  
Only a half of the expected number (blue line) was observed.

Horizontal going Neutrinos  
Flight length: 500km  
Only 80% of the expected number was observed.

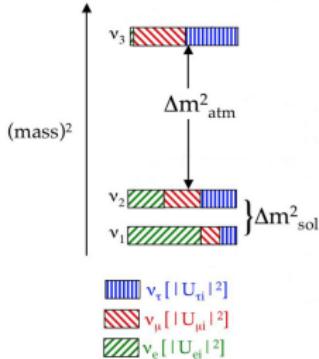
Downward going Neutrinos  
Flight length: 15km  
Consistent with the expected number.



- ▶ Proton decay experiments see neutrinos missing in some zenith directions:  
Takaaki Kajita
- ▶ Oscillation with small matter effect was the answer!

$$|\Delta m_{\text{atm}}^2| = 2.534 \times 10^{-3} \text{ eV}^2$$

# Neutrinos are massive!



$$\Delta m_{\text{sol}}^2 = 7.50 \times 10^{-5} \text{ eV}^2$$

$$|\Delta m_{\text{atm}}^2| = 2.534 \times 10^{-3} \text{ eV}^2$$

$$|U|_{3\sigma} = \begin{pmatrix} 0.799 \rightarrow 0.844 & 0.516 \rightarrow 0.582 & 0.141 \rightarrow 0.156 \\ 0.242 \rightarrow 0.494 & 0.467 \rightarrow 0.678 & 0.639 \rightarrow 0.774 \\ 0.284 \rightarrow 0.521 & 0.490 \rightarrow 0.695 & 0.615 \rightarrow 0.754 \end{pmatrix}$$

[B. Kayser, hep-ph/0506165 (2004)]

[Fig. from: Ivan Esteban et. al. JHEP 01 (2017) 087 www.nu-fit.org]

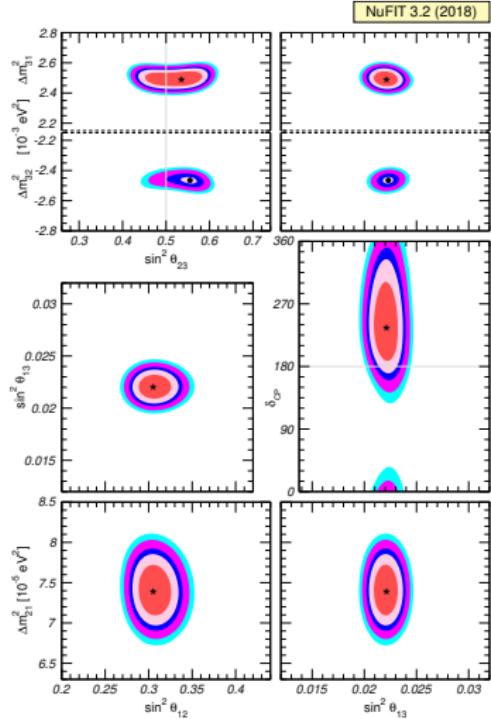
[F. Capozzi et al. Arxiv:1804.09678]

2015 NOBEL PRIZE  
in Physics



NEUTRINO OSCILLATIONS

The discovery of these oscillations shows that neutrinos have mass.



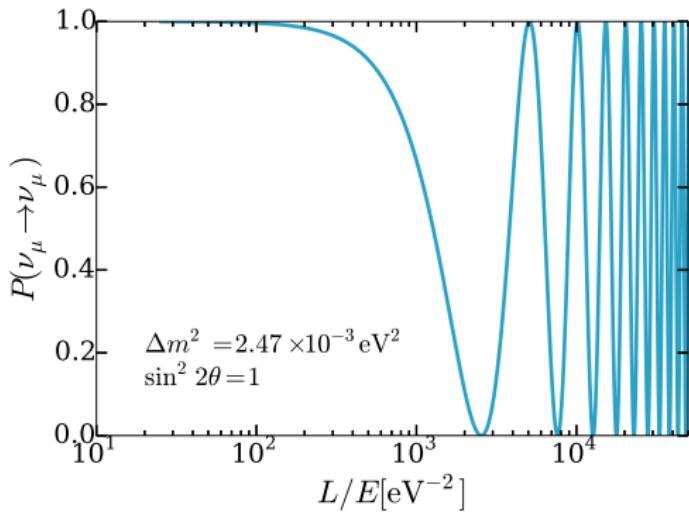
# Neutrino Oscillations

$$H = \frac{1}{2E} U M^2 U^\dagger + V_m$$

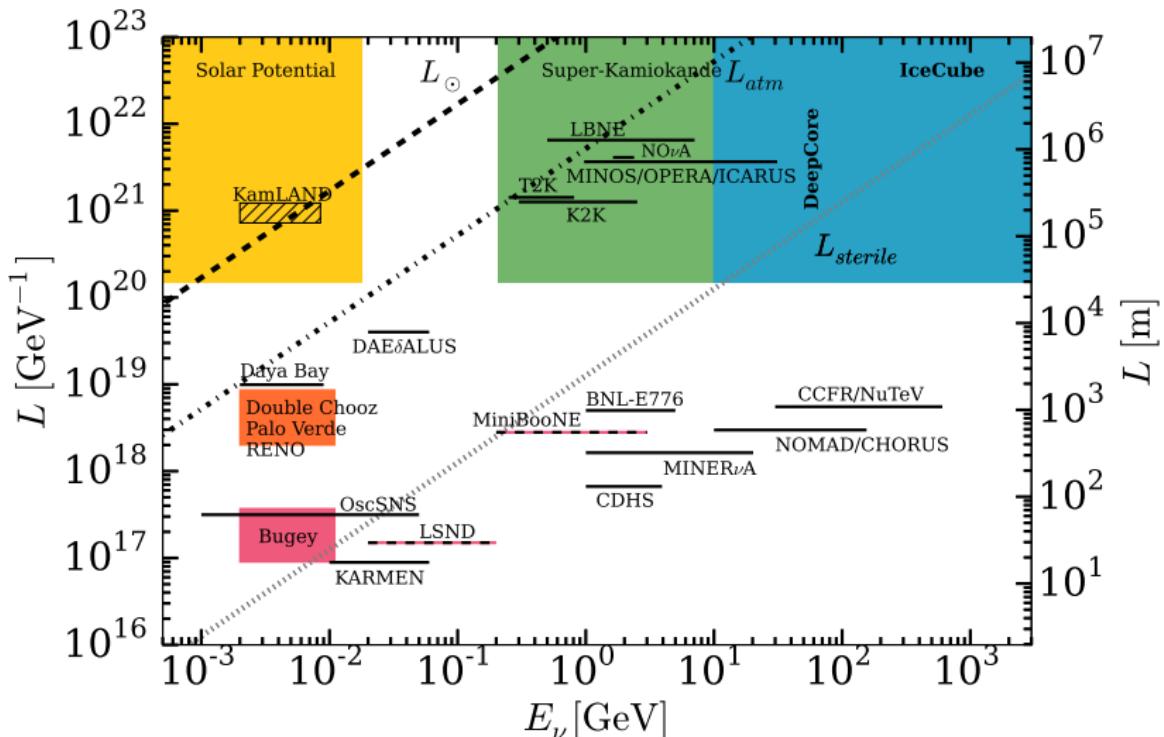
$M$ ,  $V$  and  $U$  are  $3 \times 3$  matrices. In two generations the oscillation probability at a given distance  $L$  and energy  $E$  in vacuum

$$P_{\nu_\alpha \rightarrow \nu_\alpha} \left( \frac{L}{E} \right) = 1 - \sin^2 2\theta \sin^2 \left( \frac{\Delta m^2 L}{4E} \right)$$

- ▶  $\sin^2 2\theta$  : oscillation amplitude
- ▶  $\Delta m^2$ : oscillation frequency
  - ▶  $L/E \ll 1/\Delta m^2 \rightarrow$  no oscillations
  - ▶  $L/E \sim 1/\Delta m^2 \rightarrow$  oscillations
  - ▶  $L/E \gg 1/\Delta m^2 \rightarrow$  fast oscillations ("averaged")



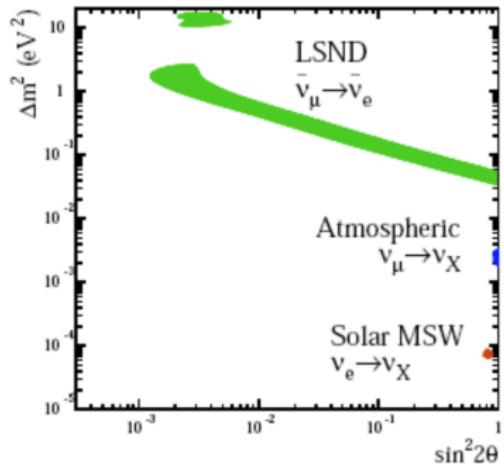
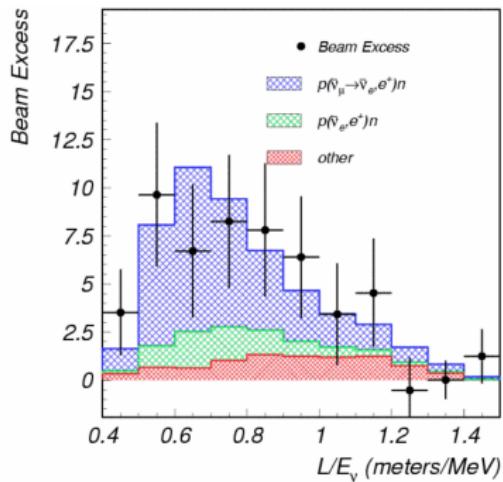
$$\text{Experiments: } L_{\text{osc}} = 2\pi \frac{E}{\Delta m^2} \mid \Delta m^2_{\text{LSND}} = 1 \text{eV}^2$$



[modified from J.S. Diaz and V.A. Kostelecky, Phys.Lett. B700, 25 (2011)]

# Why eV? The LSND experiment (90's)

- ▶ The LSND experiment saw an excess of  $\bar{\nu}_e$  over background.
- ▶  $3.8\sigma$  signal.



## More motivation: short baseline anomalies

- ▶ LSND found  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  oscillation with  $\Delta m^2 \sim 1\text{eV}^2$  and  $\sin^2 2\theta \sim 0.003$
- ▶ MiniBoone  $\nu_\mu \rightarrow \nu_e$  and  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  appearance
  - ▶ No significant excess at high energies ( $E > 475 \text{ MeV}$ )
  - ▶ Unexplained events at low energies, interpretation as oscillations similar to LSND:  $\Delta m^2 \sim 1\text{eV}^2$
  - ▶ **New result today!** arxiv:1805.12028  $4.5\sigma$   $6.1\sigma$  with LSND
- ▶ Gallium Anomaly, SAGE and GALLEX event rates lower than expected, can be explained by  $\nu_e$  disappearance with  $\Delta m^2 \geq 1\text{eV}^2$

For a current global status:

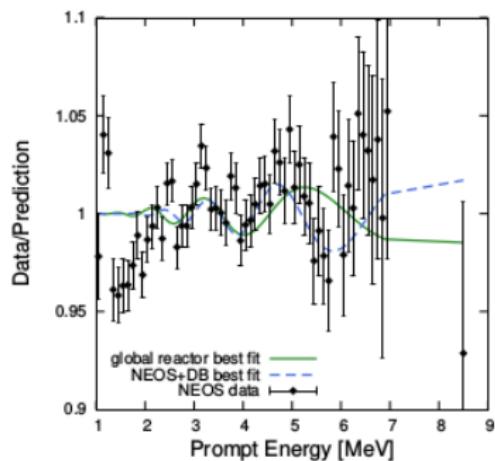
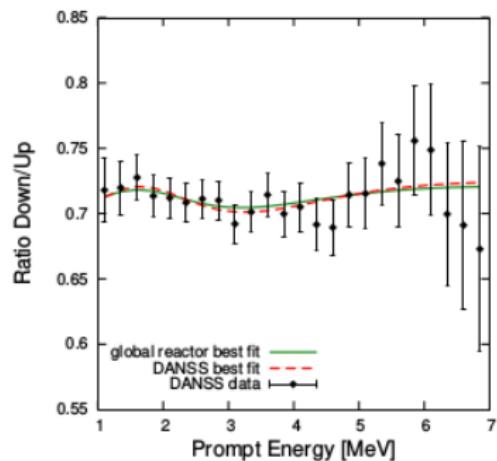
Mona Dentler et al. arXiv:1803.10661

S. Gariazzo, et.al., JHEP 06 (2017) 135

G. H. Collin, et al., Phys. Rev. Lett. 117, 221801 (2016)

## More motivation: New data in reactors

- ▶ New reactor flux calculation (Mueller et al., 1101.2663, P. Huber, 1106.0687) 3% higher, tension in short-baseline ( $L \leq 100m$ ) experiments.
- ▶ After the 5MeV bump we have new DANSS and NEOS results more independent of the flux calculation!



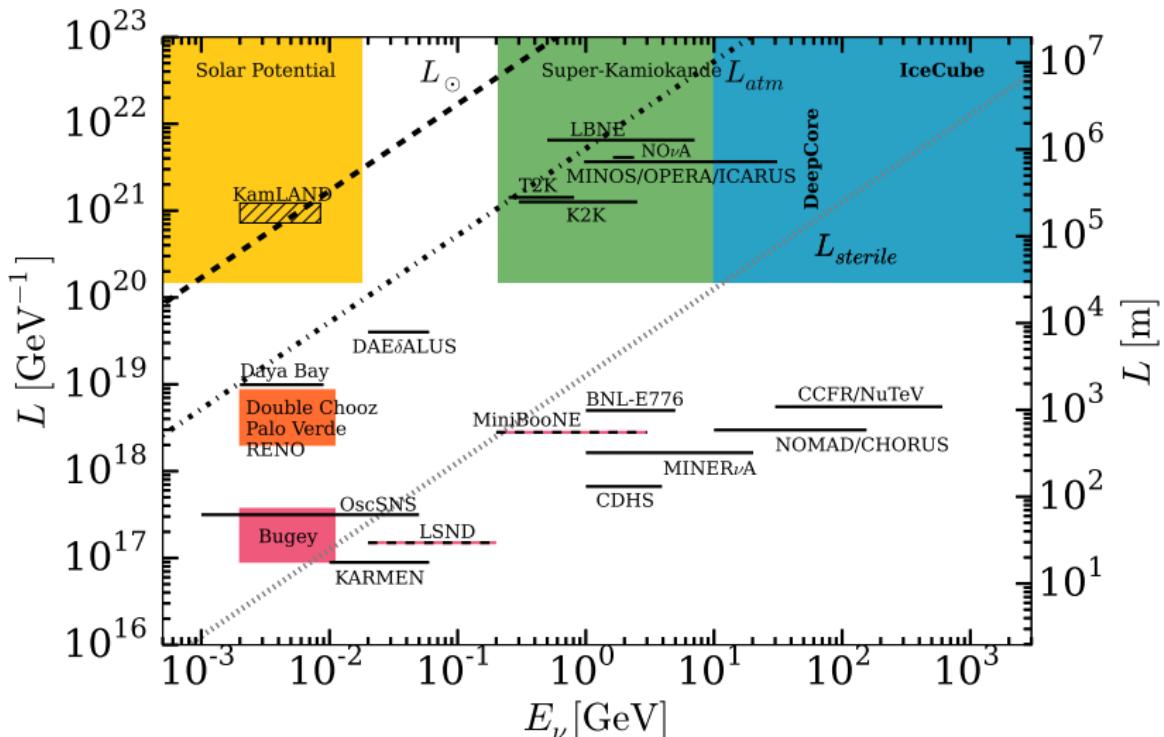
Y. Ko et al., Phys. Rev. Lett. 118 (2017)

M. Danilov, Moriond EW 2017

Mona Dentler et al. JHEP 1711 (2017) 099

C.Giunti, et al. JHEP 10 (2017) 143.

$$\text{Experiments: } L_{\text{osc}} = 2\pi \frac{E}{\Delta m^2} \mid \Delta m_{\text{LSND}}^2 = 1 \text{eV}^2$$



[modified from J.S. Diaz and V.A. Kostelecky, Phys.Lett. B700, 25 (2011)]

## Averaged ATM neutrinos

$$|\nu_\alpha\rangle = \sum_i U_{\alpha i}^* |\nu_i\rangle$$

The mixing matrix for  $n$  neutrino flavours can be decomposed as the product of  $n(n - 1)/2$  rotations.

$$U = V_{3n} V_{2n} V_{1n} V_{3(n-1)} V_{2(n-1)} V_{1(n-1)} \cdots V_{34} V_{24} V_{14} \underbrace{V_{23} V_{13} V_{12}}_{=U_0}$$

with

$$(V_{ij})_{ab} = \begin{cases} \cos(\theta_{ij}), & a = b \in \{i, j\} \\ \sin(\theta_{ij})e^{i\delta_{ij}}, & a = i, b = j \\ -\sin(\theta_{ij})e^{-i\delta_{ij}}, & a = j, b = i \\ 1, & a = b \notin \{i, j\} \\ 0, & \text{otherwise} \end{cases}$$

## Averaged ATM neutrinos

$$U = \mathcal{U} U_0 = \begin{pmatrix} 1 - \alpha & \Theta \\ X & Y \end{pmatrix} \begin{pmatrix} \mathcal{U}_\nu & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} (1 - \alpha) \mathcal{U}_\nu & \Theta \\ X \mathcal{U}_\nu & Y \end{pmatrix}$$

$$\alpha = \begin{pmatrix} \alpha_{ee} & 0 & 0 \\ \alpha_{\mu e} & \alpha_{\mu\mu} & 0 \\ \alpha_{\tau e} & \alpha_{\tau\mu} & \alpha_{\tau\tau} \end{pmatrix}$$

whose components to leading order in the active-heavy mixing elements are given by

$$\alpha_{\beta\gamma} \simeq \begin{cases} \frac{1}{2} \sum_{i=4}^n |U_{\beta i}|^2, & \beta = \gamma \\ \sum_{i=4}^n U_{\beta i} U_{\gamma i}^*, & \beta > \gamma \\ 0, & \gamma > \beta \end{cases}$$

## Averaged ATM neutrinos

$$\tilde{H} = \begin{pmatrix} H_0 & 0 \\ 0 & H_1 \end{pmatrix} + V_{\text{NC}} \mathcal{U}^\dagger \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \mathcal{U}$$

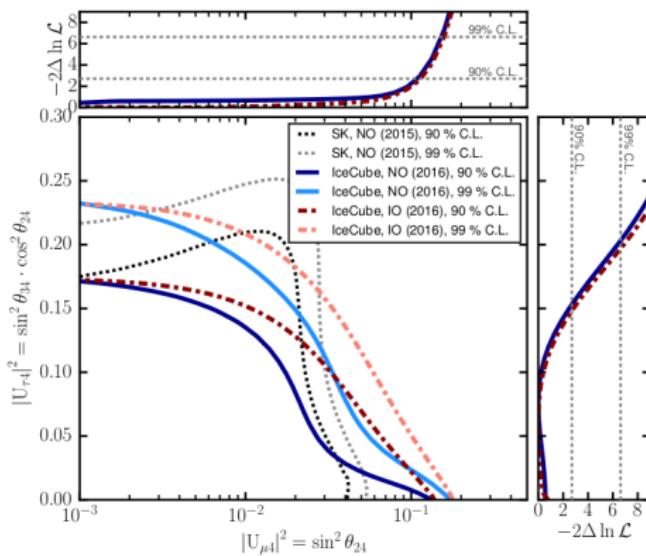
$$\tilde{H}_0 = H_0 + V_{\text{NC}}(1 - \alpha^\dagger)(1 - \alpha) \simeq H_0 - V_{\text{NC}}(\alpha + \alpha^\dagger)$$

$$\tilde{H}_0 = \frac{\Delta m_{31}^2}{4E} \begin{pmatrix} -\cos(2\theta_{23}) & \sin(2\theta_{23}) \\ \sin(2\theta_{23}) & \cos(2\theta_{23}) \end{pmatrix} - V_{\text{NC}} \begin{pmatrix} 2\alpha_{\mu\mu} & \alpha_{\tau\mu}^* \\ \alpha_{\tau\mu} & 2\alpha_{\tau\tau} \end{pmatrix}$$

- ▶ The averaged effect of sterile neutrinos is essentially a modification of the matter potential (NSI Arman Esmaili Talk)

# eV steriles by SuperKamiokande and deepcore

- ▶ SuperKamikande and DC they both measure atmospheric neutrino oscillations and give a bound due to matter effects.
- ▶ eV is too heavy for this data samples to see oscillaitons, esencially a reparameterizaiton of the NSI result.



K. Abe et al. (Super-Kamiokande Collaboration), Phys. Rev. D91, 052019 (2015).  
IceCube Collaboration, Phys. Rev. Lett. 117 no. 7, (2016) 071801

# Matter effects with the Sterile Neutrino at Earth

$$P_{\nu_\alpha \rightarrow \nu_\alpha} = 1 - \sin^2 2\theta_M \sin^2 \left( \frac{\Delta m_M^2 L}{4E_\nu} \right)$$

where  $\theta_M$  and  $\Delta m_M^2$  satisfy

$$\Delta m_M^2 = \sqrt{(\Delta m^2 \cos 2\theta - A)^2 + (\Delta m^2 \sin 2\theta)^2}$$
$$\tan 2\theta_M = \frac{\tan 2\theta}{1 - \frac{A}{\Delta m^2 \cos 2\theta}}$$

and  $A = \pm 2\sqrt{2}EG_F N$ ,  $N$  number density. Resonant flavor transition can happen if

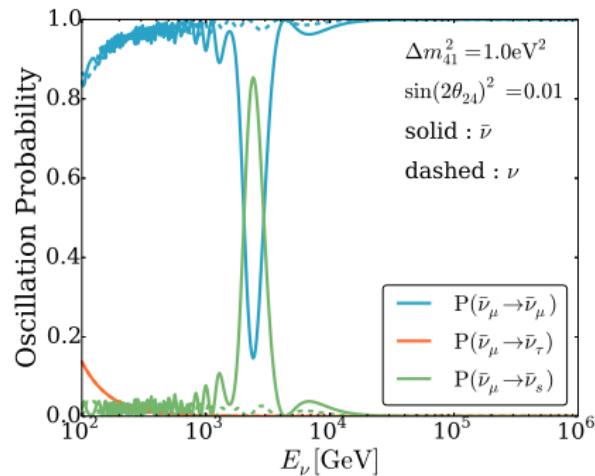
$$E_\nu^{res} = \mp \cos 2\theta \frac{\Delta m^2}{2N} \frac{1}{\sqrt{2G_F}}$$

this resonance can enhance the transition between active and sterile neutrinos. Talk by Alexei Smirnov

# Matter effects with the Sterile Neutrino at Earth

In the **Earth**, sterile neutrino with **small mixing** and  
 $\Delta m^2 = O(1\text{eV}^2)$  the resonance happens when

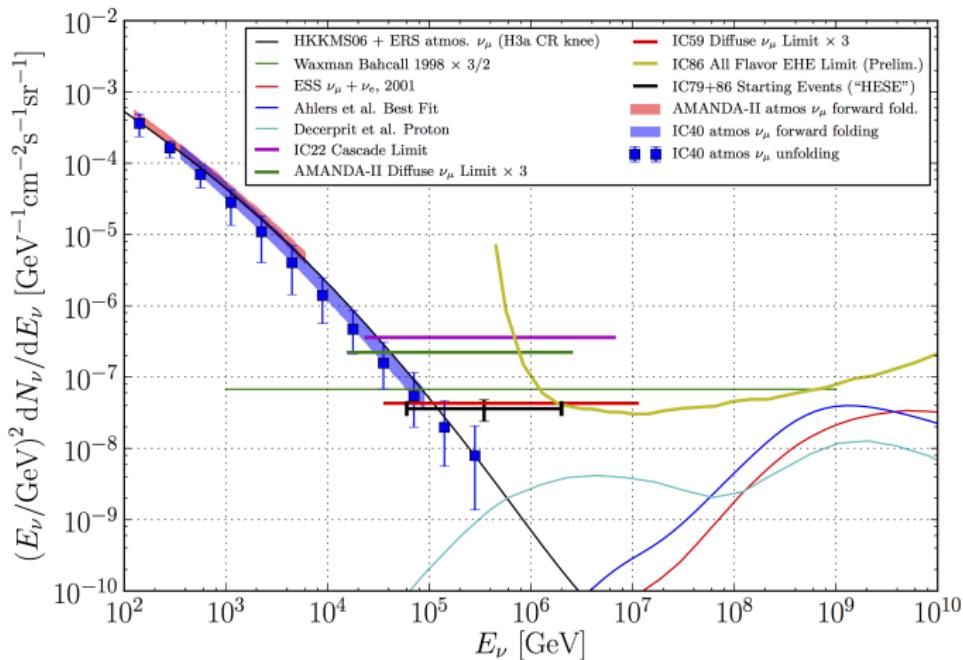
$$E_\nu^{\text{res}} = \frac{\Delta m^2 \cos 2\theta}{2\sqrt{2}G_F N} \sim O(\text{TeV})$$



- M.V. Chizhov, S.T. Petcov. Phys.Rev. D63 (2001) 073003  
H. Nunokawa et al. Phys.Lett. B562 (2003) 279-290  
Sandhya Choubey JHEP 0712 (2007) 014  
Barger et al., Phys.Rev.D85:011302,(2012)  
Arman Esmaili et al. JCAP 1211 (2012) 041

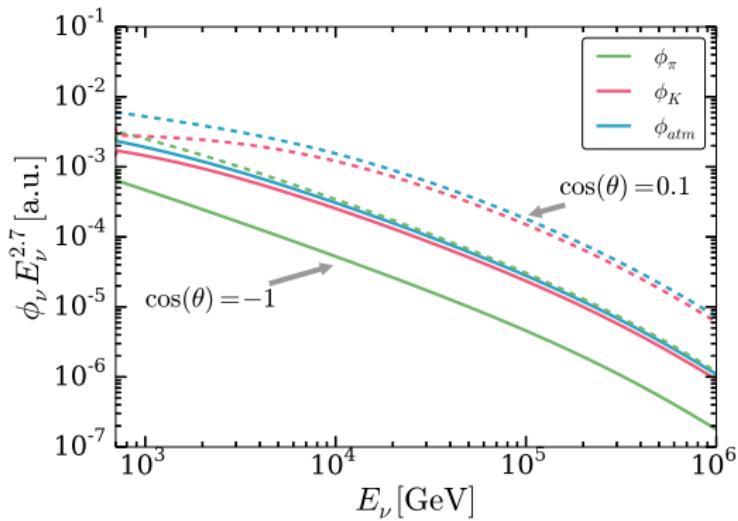
# Matter effects with the Sterile Neutrino at Earth

- ▶ TeV is in the center of the atmospheric data in IceCube.
- ▶ Other experiments are not sensitive at this energies.



# The initial atmospheric neutrino flux

The conventional atmospheric neutrino (muon) flux originates from the decay of  $\pi^\pm$  and  $K^\pm$  in the atmosphere.

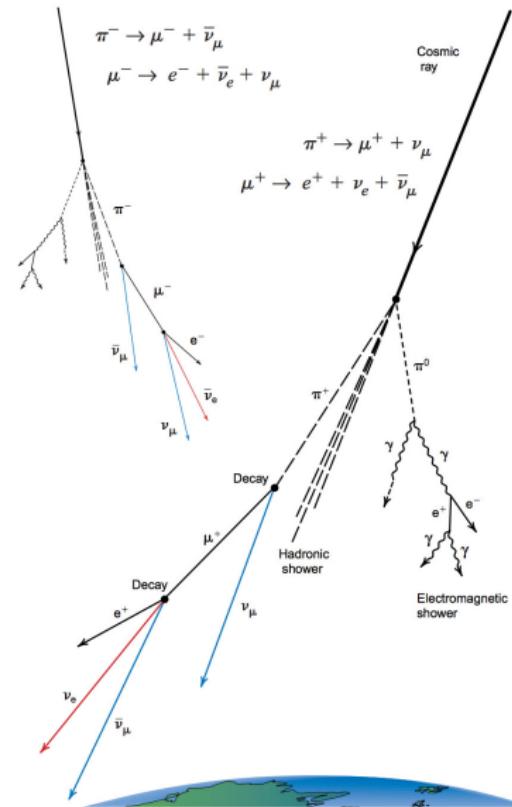


[Honda et al., Phys.Rev.D75:043006 (2007)]

[Louis et al., Los Alamos Science Number 25 (1997)]

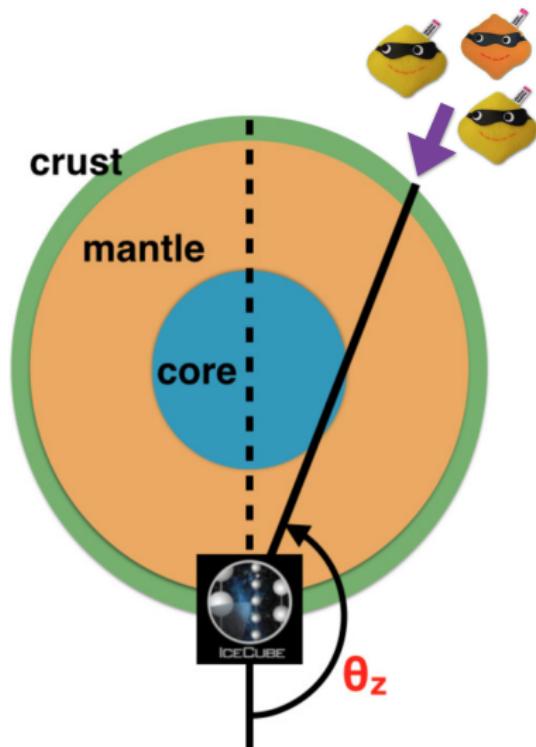
we are improving, talks by Thomas K. Gaisser, Anatoli Fedynitch,

Morihiro Honda



# Neutrinos through the Earth

The muon neutrinos come from different zenith angles ( $\theta_z$ ) crossing different Earth layers



**core :**

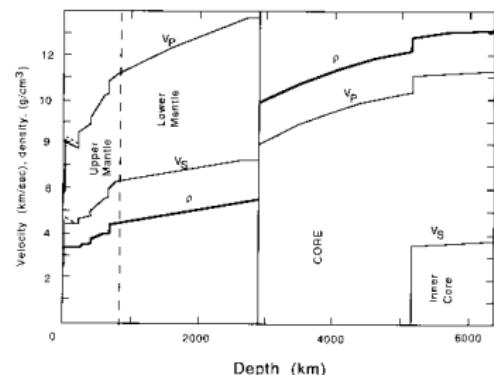
$$\cos \theta_z \sim [-1, -0.8]$$

**mantle :**

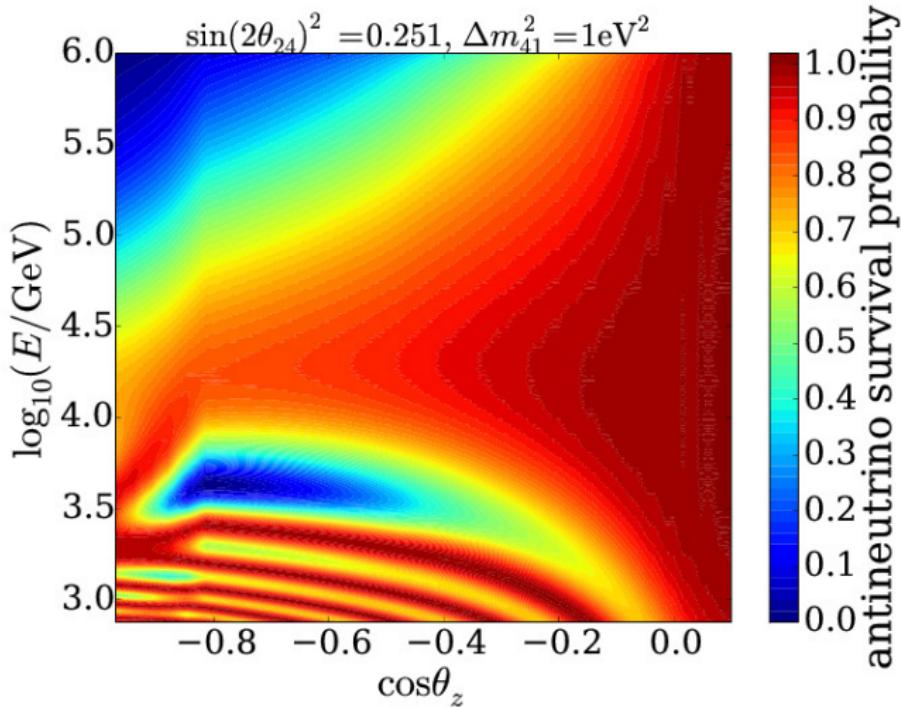
$$\cos \theta_z \sim [-0.8, -0.1]$$

**crust :**

$$\cos \theta_z > -0.1$$



# 3+1 Oscillogram



[Carlos Argüelles, J.S., C. Weaver. *SQuIDS*, CPC 2015.06.022.]

<https://github.com/jsalvado/SQuIDS>

<https://github.com/arguelles/nuSQuIDS>



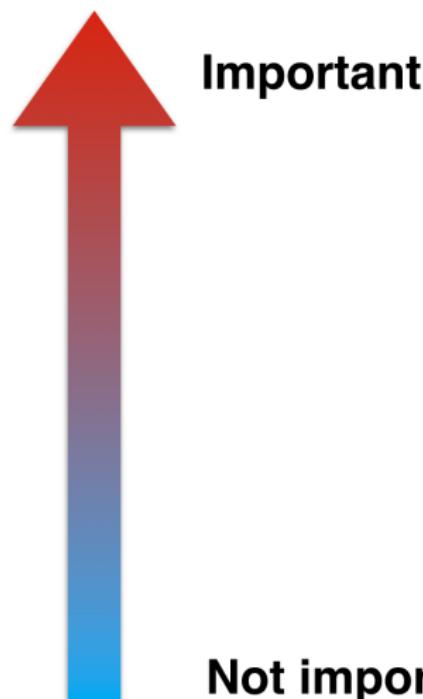
# Systematic errors

Systematics are *very* important; *some more than others*. This are the systematics we considered:

- ▶ *DOM efficiency*
- ▶ *Flux continuous parameters*
  - ▶ *spectral index*
  - ▶  *$\pi/K$  ratio*
  - ▶  *$\nu/\bar{\nu}$  ratio*
- ▶ Air shower hadronic models
- ▶ Primary cosmic ray fluxes
- ▶ Hole Ice
- ▶ Neutrino cross sections
- ▶ Bulk ice scattering/absorption
- ▶ Earth model

*continuous systematics*

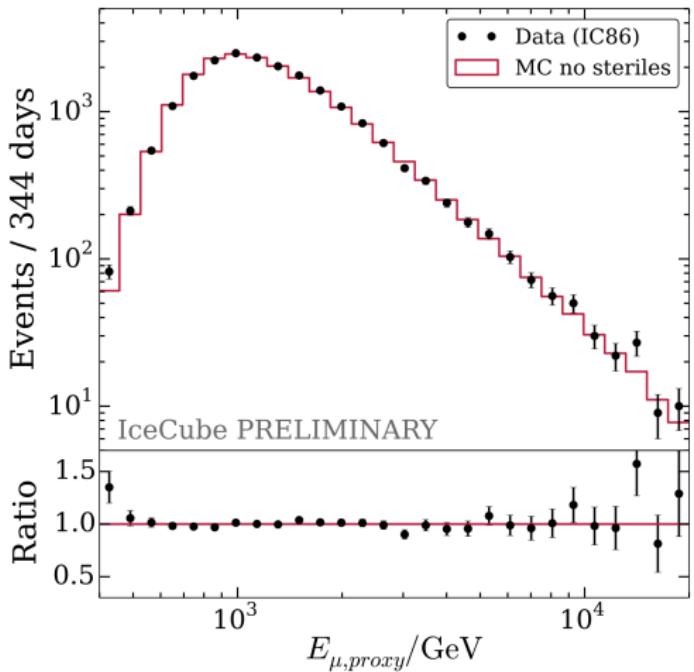
discrete systematic



**Not important**

## How the fit looks

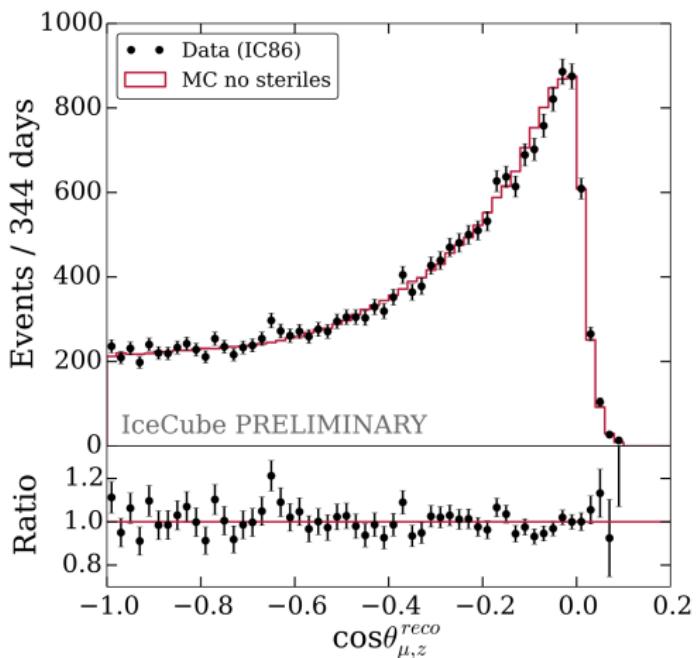
- ▶ We fitted the null hypothesis (no steriles) using the central sets (no variants) on the full 2D sample space.
- ▶ We recover a good fit and sensible nuisance parameters.



Parameter	Value	Prior
Normalization	1.02	No Prior
$\Delta\gamma$	0.05	$G(0.,0.05)$
DOEff	0.985	No Prior
$\pi/K$	1.10	$G(1.,0.1)$
$\nu/\bar{\nu}$	1.0	$G(1.,0.05)$
$\delta$	0.001	$G(0.,0.05)$

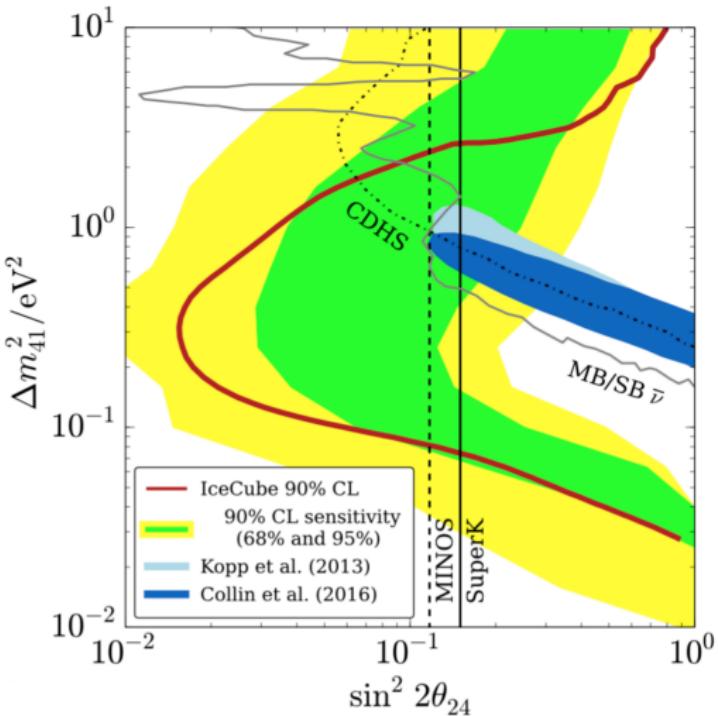
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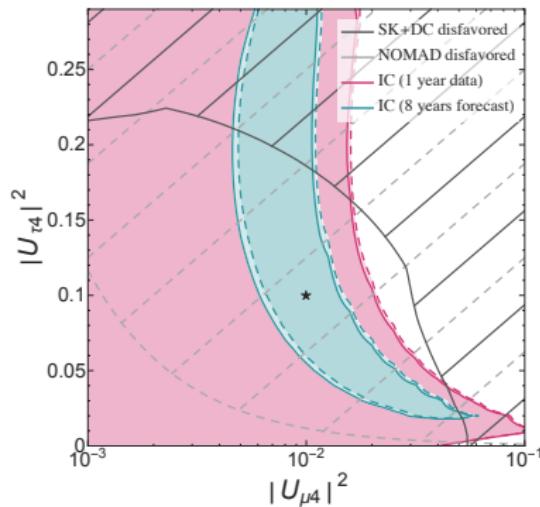
# Main result!



And if they are heavier?  $> 10\text{eV}$

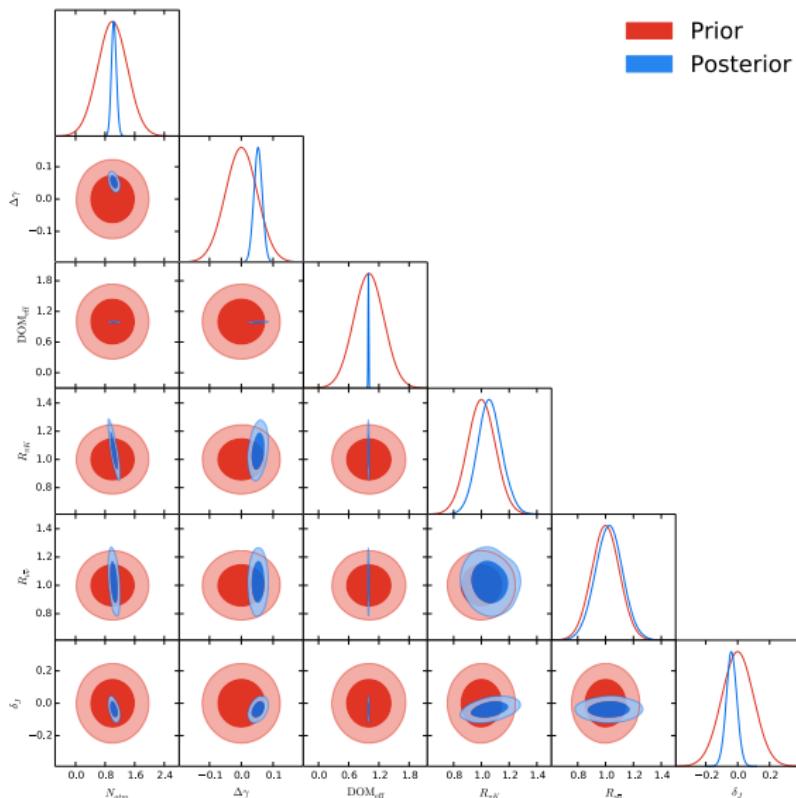
- ▶ Effectively a re-definition of the matter potential, a la NSI.

$$H = \frac{1}{2E_\nu} U \textcolor{red}{M}^2 U^\dagger + V_m$$



# Continuous Systematics

- ▶ Some of them like the DOMeff are well measured by the data.

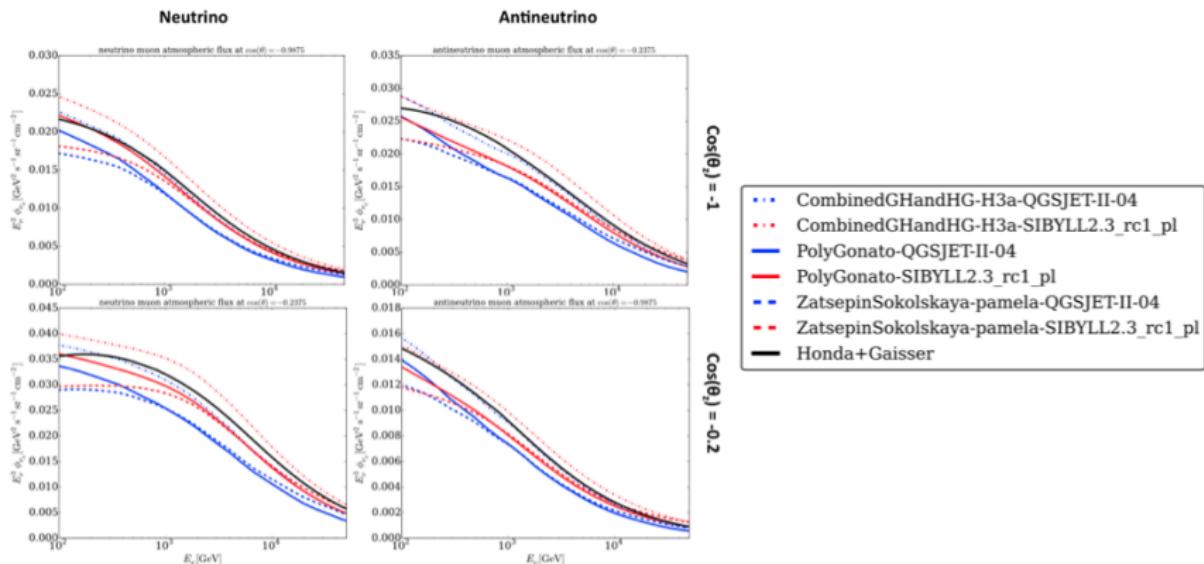


From C.A.

# More on Flux Systematics!

Main uncertainties in the neutrino flux are capture by the *continuous* parameters.

CR flux uncertainties and Hadronic models are discrete systematics

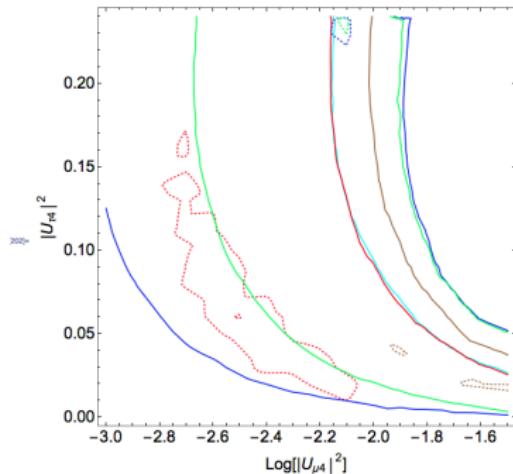
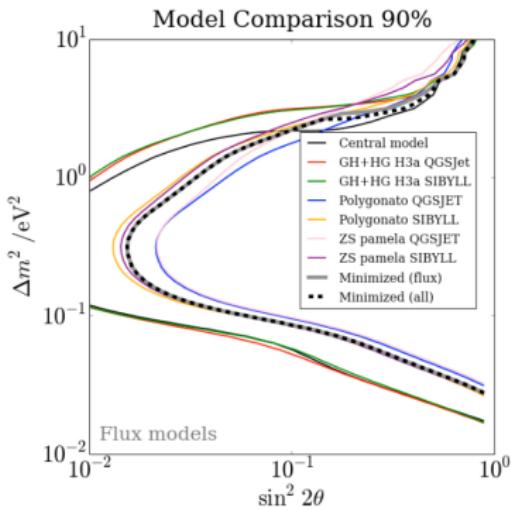


[Fedynitch et al. arXiv:1504.06639]

[Collins et al. URL: <http://dspace.mit.edu/handle/1721.1/98078>]

# More on Flux Systematics!

- ▶ Discrete flux systematics have a big impact.
- ▶ Still OK treatment for 1 year but an improvement is needed for future analysis.



More in C.Arguelles and B.J.P.Jones thesis

## Conclusions

- ▶ Atm neutrinos are great! They constrain eV and heavier sterile neutrinos, specially in the high energy region due to the parametric resonance.
- ▶ Matter effect are essential in both cases(**lowE and highE**) therefore any modification of the matter effects change this bounds.(Yasaman Farzan and Danny Marfatia Talks)
- ▶ In the next updates  $O(10\text{yrs})$  systematic errors are becoming important.
- ▶ Our knowledge of **the flux** is going to be **limiting all the BSM searches** in the future, better knowledge or treatment is very important (talks by Thomas K. Gaisser, Anatoli Fedynitch, Morihiro Honda)

iThanks!