

Extra Dimensions in Atmospheric Neutrinos

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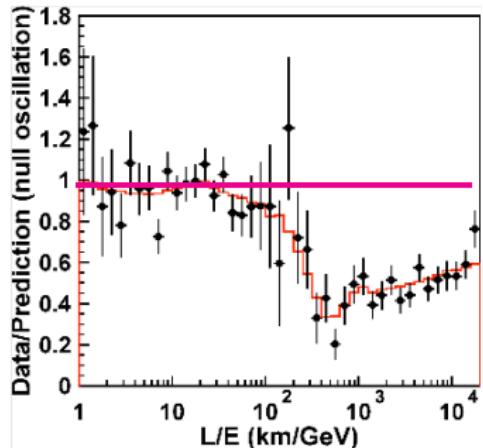
**Advanced Workshop on Physics of Atmospheric
Neutrinos - PANE 2018**

¹In collaboration with Arman Esmaili and Zahra Tabrizi

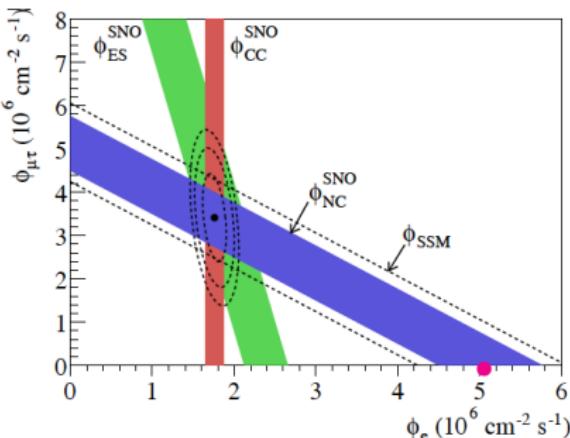


Neutrino Oscillations

- In the latest years there are stronger evidences for neutrino oscillation



Super-Kamiokande Experiment

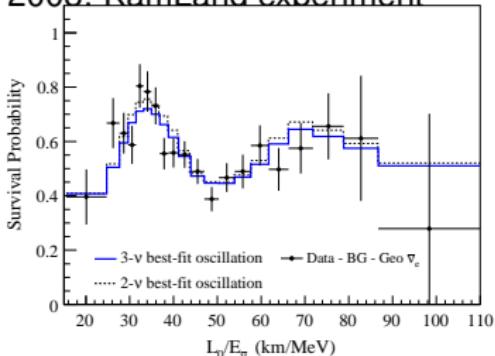


SNO experiment

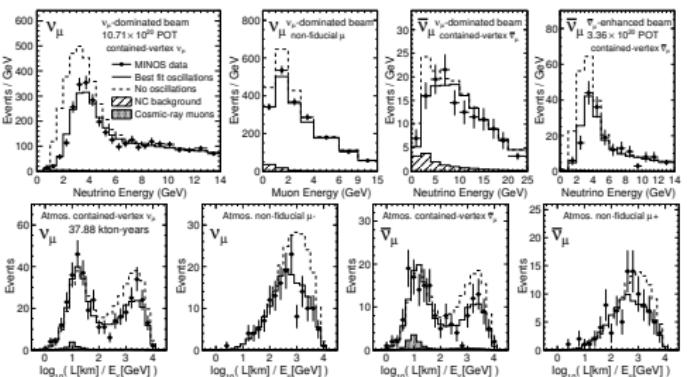
All evidences point to non-zero neutrino masses: Nobel prize in 2015!!

Oscillation discover in other expts

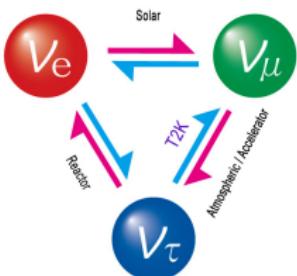
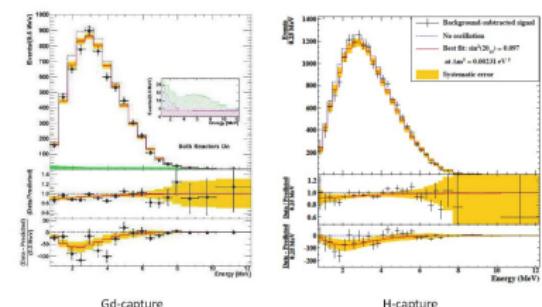
2008: KamLand experiment



MINOS experiment: UFG/USP/UNICAMP



Double CHOOZ: UNICAMP/UFABC/CBPF

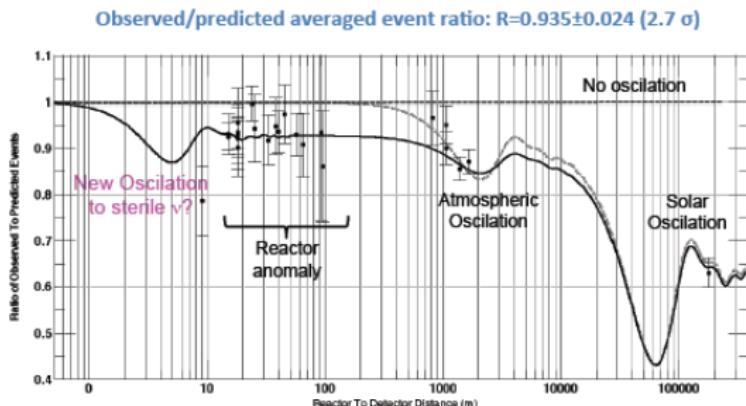
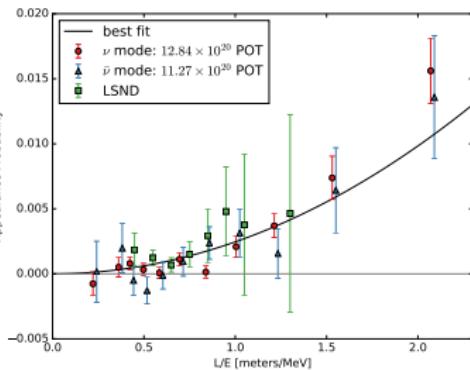
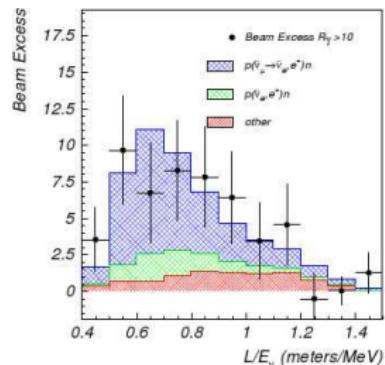


Neutrino oscillation between three generations

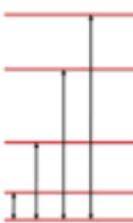
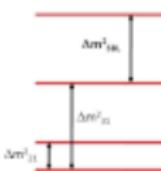


Sterile neutrino :hiccups or new particle?

- LSND, Mini-BooNE (2018) and Reactor anomaly

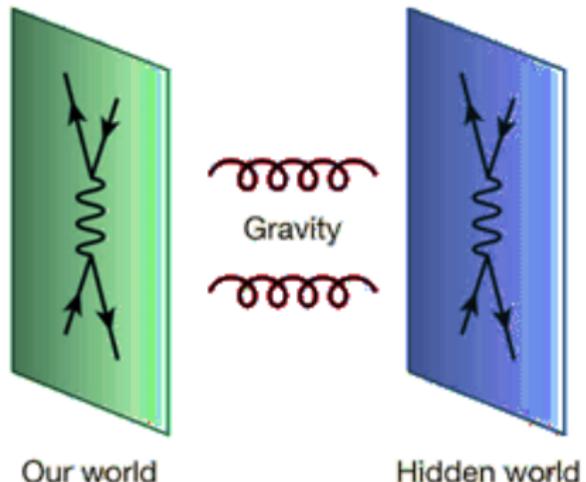


- 3+1 mass scheme
- 3+2 mass scheme



Large extra dimension model

- Proposed by Arkani-Hamed, Savas Dimopoulos and Dvali (ADD)

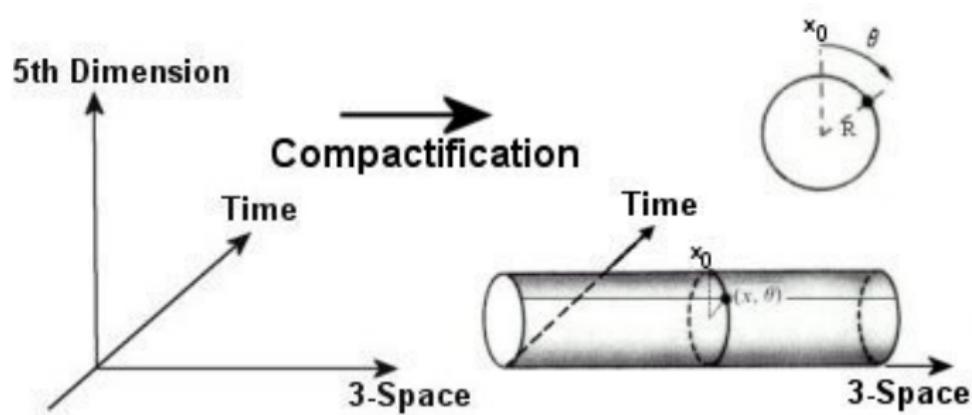


- Gauss law: $V \sim \frac{m_1 m_2}{M_D^{\delta+2} R_{\text{ED}}^\delta r} = \frac{m_1 m_2}{M_{\text{Pl}} r}$

Constrains from $1/r$ tests $R_{\text{ED}} > 1 \text{ mm}$

Large extra dimension model

Compactification: explanation of smallness of gravity: $\propto \frac{1}{V_{\text{ED}}} \propto \frac{1}{R_{\text{ED}}}$



One of extra dimensions have R_{ED} much greater then the other dimensions (5D) .

Neutrino in Large Extra Dimension Models

$$S_i = \int d^4x \, dy \, i\bar{\Psi}_i \Gamma^A \partial_A \Psi_i + \int d^4x \, [i\bar{\nu}_{iL} \not{d} \nu_{iL} + \lambda_{ij} h \bar{\nu}_{iL} \psi_{jR}(x, y=0)]$$

Smallness of neutrino mass: $m_\nu \propto \frac{\lambda_{ij}}{R_{\text{ED}}} \sim 1\text{eV}$

- Fourier modes: infinite number of states



Neutrino in Large Extra Dimension Models

$$S_i = \int d^4x \text{ dy } i\bar{\Psi}_i \Gamma^A \partial_A \Psi_i + \int d^4x [i\bar{\nu}_{iL} \not{\partial} \nu_{iL} + \lambda_{ij} h \bar{\nu}_{iL} \psi_{jR}(x, y=0)]$$

Mass terms

$$\sum_{n=-\infty}^{\infty} m_i^D \bar{\nu}_{iL} \psi_{iR}^{(n)} + \sum_{n=1}^{\infty} \frac{n}{R_{\text{ED}}} \left(\overline{\psi_{iL}^{(n)}} \psi_{iR}^{(n)} - \overline{\psi_{iL}^{(-n)}} \psi_{iR}^{(-n)} \right) + \text{h.c.},$$

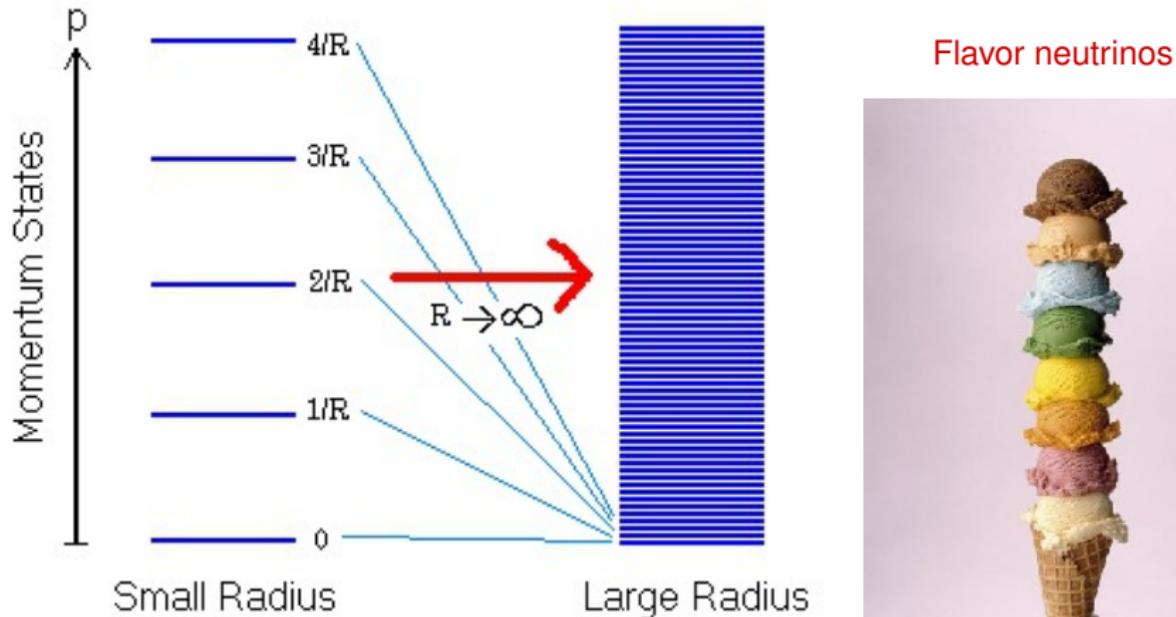
$$M_i = \lim_{n \rightarrow \infty} \begin{pmatrix} m_i^D & \sqrt{2}m_i^D & \sqrt{2}m_i^D & \sqrt{2}m_i^D & \dots & \sqrt{2}m_i^D \\ 0 & 1/R_{\text{ED}} & 0 & 0 & \dots & 0 \\ 0 & 0 & 2/R_{\text{ED}} & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & 0 & \dots & n/R_{\text{ED}} \end{pmatrix}.$$

Only two free parameters m_i^D and R_{ED} .



Neutrino in Large Extra Dimension Models

Depending on radius we can have oscillation on short and long scale



$$\frac{1}{R^2} \gg \Delta m_{31}^2$$

$$\frac{1}{R^2} \sim \Delta m_{31}^2$$



Neutrino in Large Extra Dimension Models

The flavor eigenstates are written as

$$\nu_{\alpha L} = \sum_{i=1}^3 U_{\alpha i} \nu_{iL}^{(0)} = \sum_{i=1}^3 U_{\alpha i} \sum_{n=0}^{\infty} S_i^{0n} \nu_{iL}'^{(n)}, \quad (\alpha = e, \mu, \tau),$$

The eigenvalues are solution of

$$\lambda_i^{(n)} - \pi (m_i^D R_{\text{ED}})^2 \cot(\pi \lambda_i^{(n)}) = 0.$$

with solution

$$m_i^{(n)} = \frac{\lambda_i^{(n)}}{R_{\text{ED}}} \simeq \frac{n}{R_{\text{ED}}}, \quad (n = 0, 1, \dots),$$

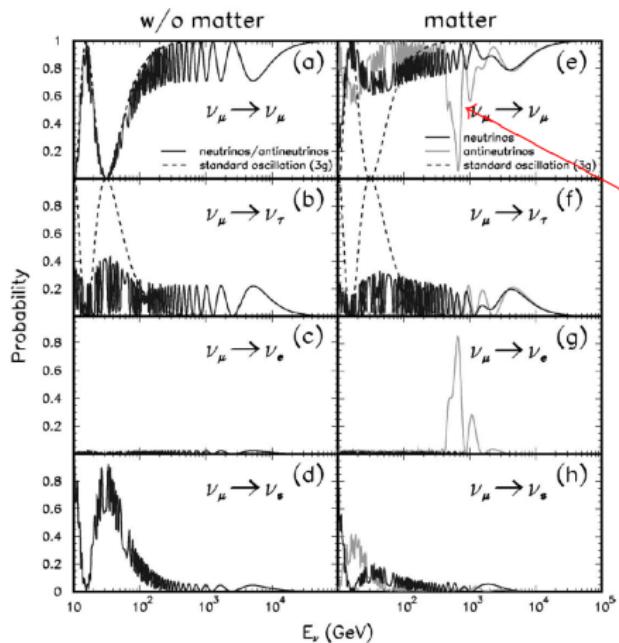
- If $\Delta m_{31}^2 = m_3^2 - m_1^2 \sim \frac{n^2}{R_{\text{ED}}^2}$: Stronger changes in neutrino oscillation
- If the radius R_{ED} is larger also stronger non-unitarity effects

$$U_{\alpha i} \rightarrow U_{\alpha i} S_i^{0n}$$



Matter effects for sterile ν

H. Nunokawa, O.L.G.P. and R. Zukanovich Funchal. Phys.Lett. B562 (2003)



Stronger muon disappearance.

Also: A. Nicolaidis, *et al.* hep-ph/9904415
and Osamu Yasuda, hep-ph/0102166

Stronger bound on $|U_{\mu 4}|^2$, A. Esmaili, F. Halzen and O.L.G.P. JCAP 1211, 041(2012)

Matter effects in Large extra dimensions

Matter potential in mass basis

$$i \frac{d}{dt} \begin{pmatrix} \nu_1^{(0)} \\ \nu_2^{(0)} \\ \nu_3^{(0)} \\ \nu_1^{(1)} \\ \nu_2^{(1)} \\ \nu_3^{(1)} \\ \nu_1^{(2)} \\ \nu_2^{(2)} \\ \nu_3^{(2)} \\ \vdots \\ \nu_1^{(N)} \\ \nu_2^{(N)} \\ \nu_3^{(N)} \end{pmatrix} = \frac{1}{2ER_{\text{ED}}^2} \begin{pmatrix} \eta_1 + V_{11} & V_{12} & V_{13} & \xi_1 & 0 & 0 & 2\xi_1 & 0 & 0 & \dots & N\xi_1 & 0 & 0 \\ V_{21} & \eta_2 + V_{22} & V_{23} & 0 & \xi_2 & 0 & 0 & 2\xi_2 & 0 & \dots & 0 & N\xi_2 & 0 \\ V_{31} & V_{32} & \eta_3 + V_{33} & 0 & 0 & \xi_3 & 0 & 0 & 2\xi_3 & \dots & 0 & 0 & N\xi_3 \\ \xi_1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & \dots & 0 & 0 & 0 \\ 0 & \xi_2 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & \dots & 0 & 0 & 0 \\ 0 & 0 & \xi_3 & 0 & 0 & 1 & 0 & 0 & 0 & \dots & 0 & 0 & 0 \\ 2\xi_1 & 0 & 0 & 0 & 0 & 0 & 4 & 0 & 0 & \dots & 0 & 0 & 0 \\ 0 & 2\xi_2 & 0 & 0 & 0 & 0 & 0 & 4 & 0 & \dots & 0 & 0 & 0 \\ 0 & 0 & 2\xi_3 & 0 & 0 & 0 & 0 & 0 & 4 & \dots & 0 & 0 & 0 \\ \vdots & \ddots & \vdots & \vdots & \vdots \\ N\xi_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \dots & N^2 & 0 & 0 \\ 0 & N\xi_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \dots & 0 & N^2 & 0 \\ 0 & 0 & N\xi_3 & 0 & 0 & 0 & 0 & 0 & 0 & \dots & 0 & 0 & N^2 \end{pmatrix} \begin{pmatrix} \nu_1^{(0)} \\ \nu_2^{(0)} \\ \nu_3^{(0)} \\ \nu_1^{(1)} \\ \nu_2^{(1)} \\ \nu_3^{(1)} \\ \nu_1^{(2)} \\ \nu_2^{(2)} \\ \nu_3^{(2)} \\ \vdots \\ \nu_1^{(N)} \\ \nu_2^{(N)} \\ \nu_3^{(N)} \end{pmatrix}$$

For neutrinos

$$\eta_i = (N + 1/2) \xi_i^2, \quad V_{ij} = 2ER_{\text{ED}}^2 \sum_{\alpha=e,\mu,\tau} U_{\alpha i}^* U_{\alpha j} V_\alpha$$

$$\xi_i = \sqrt{2} m_i^D R_{\text{ED}}$$

Mass eigenstates

For anti-neutrinos: $\bar{V}_\mu < 0$ then $\bar{\nu}_\mu \rightarrow \nu_s$ transition

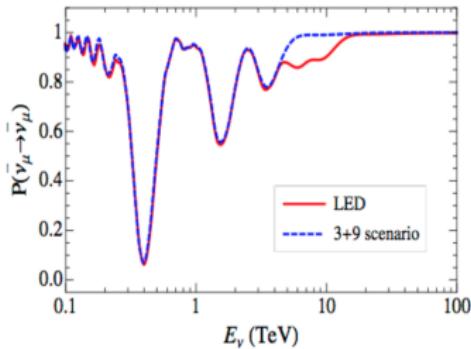
The Oscillation Probabilities

$$\Delta m^2 \sim \frac{1}{R^2} \sim 0.16 \text{ eV}^2$$

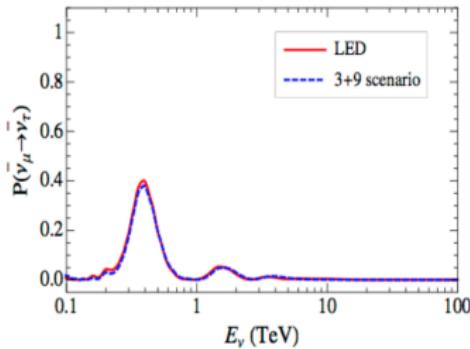
$$\sin^2 \theta \sim (mR)^2 \sim 2.5 \times 10^{-2}$$

- For 3 KK modes:

$$m_1^D = 0.01 \text{ eV} \text{ and } R_{\text{ED}} = 5 \times 10^{-5} \text{ cm}$$



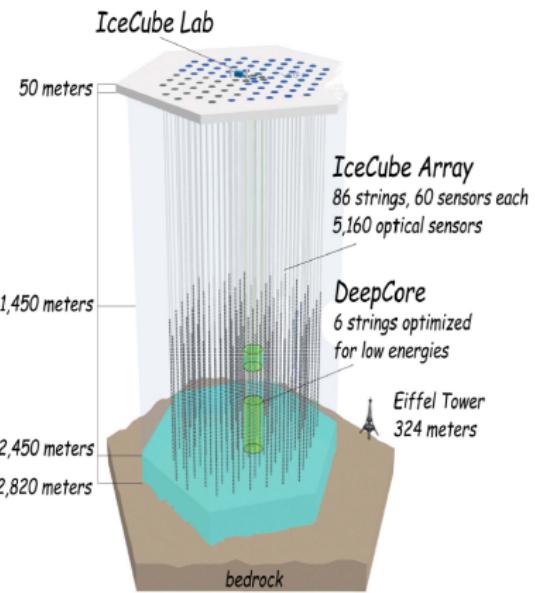
(a) $P(\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu)$



(b) $P(\bar{\nu}_\mu \rightarrow \bar{\nu}_\tau)$

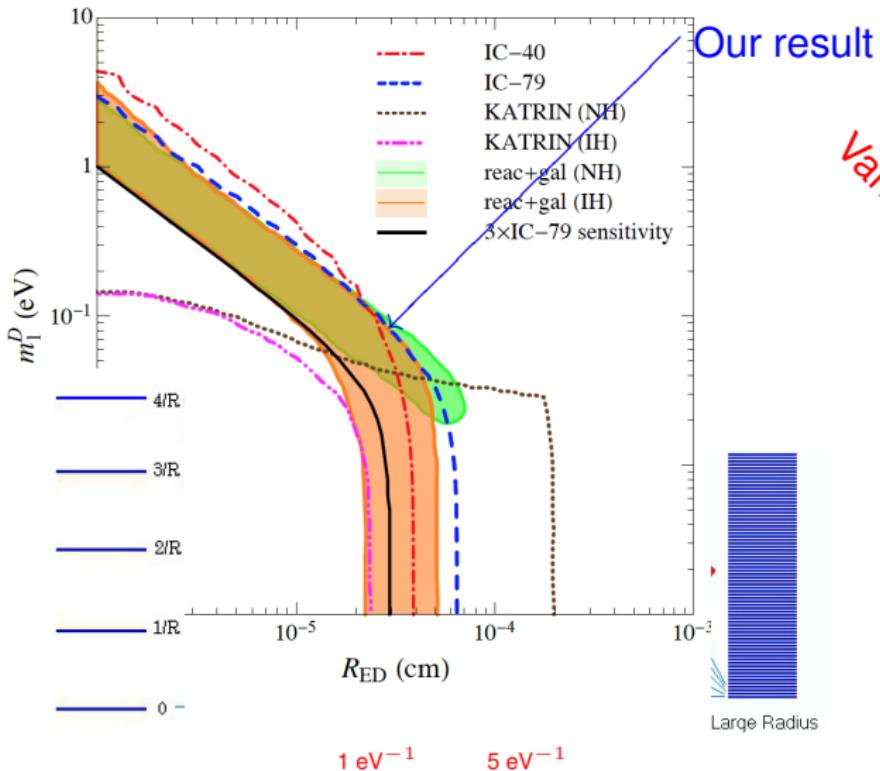
ICECUBE experiment

- ▶ 5160 PMTs
- ▶ 1 km³ volume
- ▶ 86 strings
- ▶ 17 m PMT-PMT spacing per string
- ▶ 120 m string spacing
- ▶ Angular resolution ~ 1°
- ▶ Completed 2010



Constraints on large extra dimensions

χ^2 analysis of muon disappearance data of ICECUBE

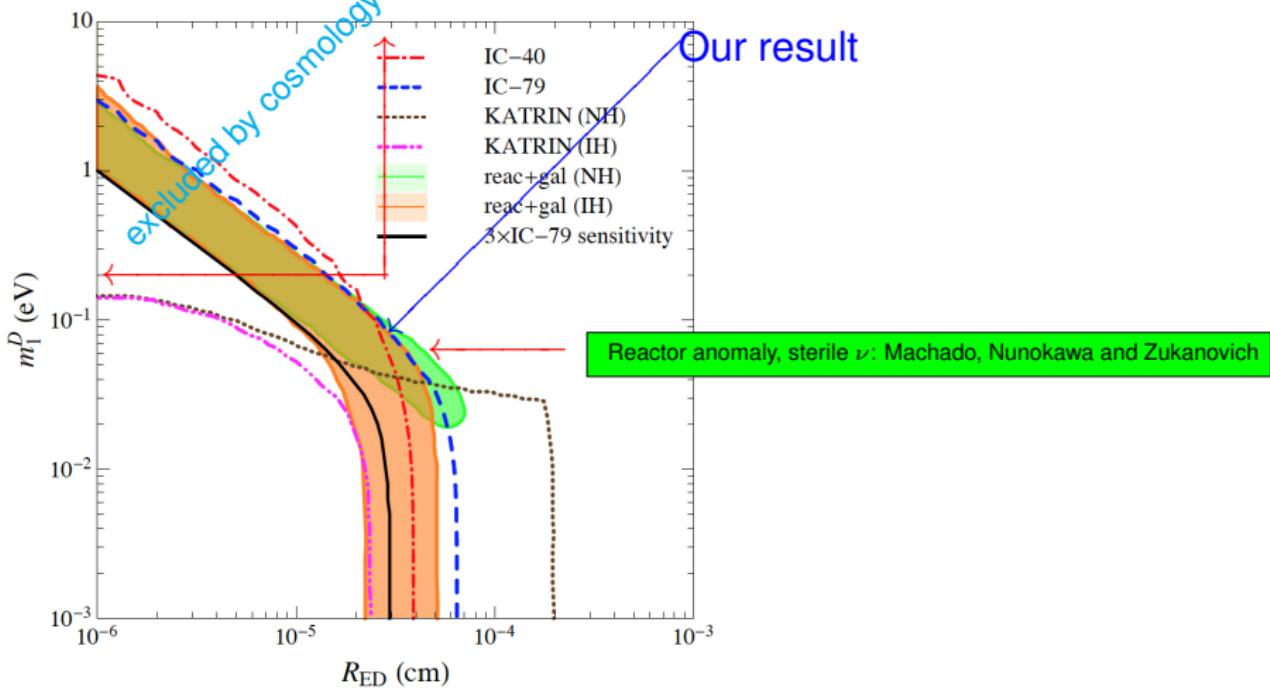


Variables: m_{lightest} and R_{ED}

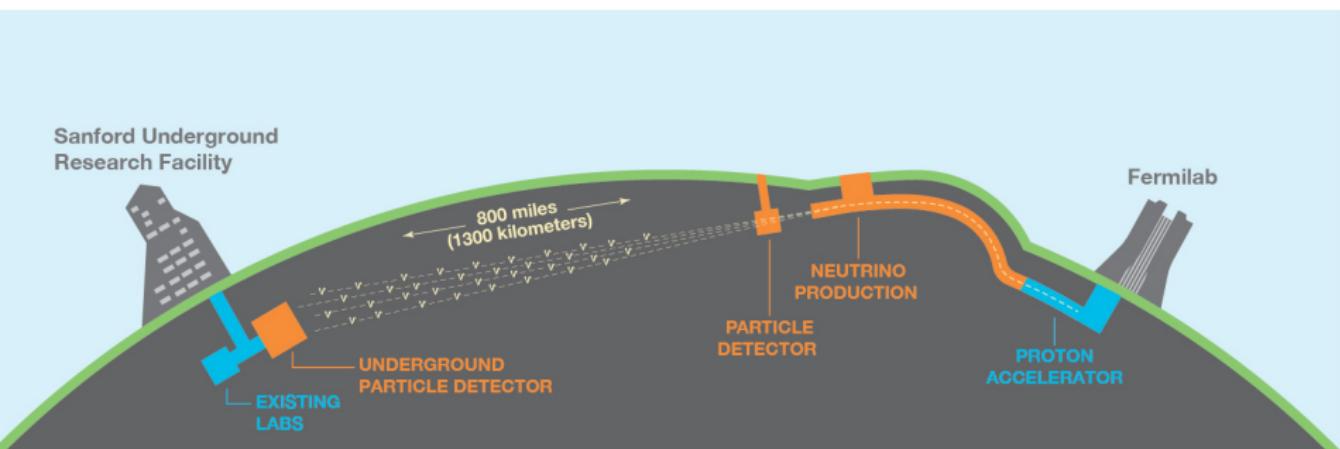


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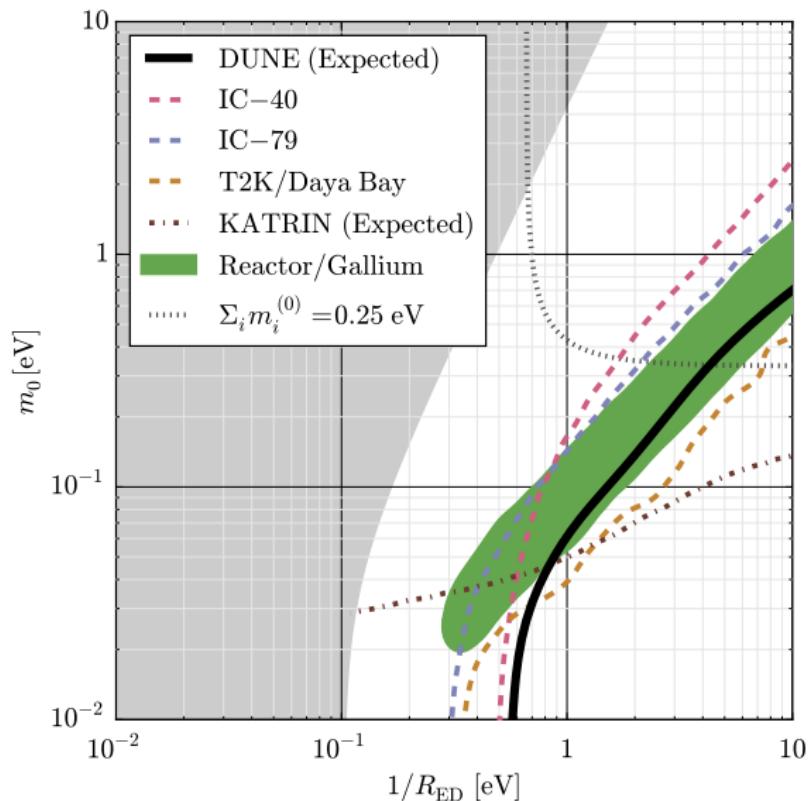
Test in longbaseline experiments: DUNE



Measure μ (e) events: Test of $P(\nu_\mu \rightarrow \nu_\mu)(P(\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu)$ and $P(\nu_\mu \rightarrow \nu_e)(P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e))$

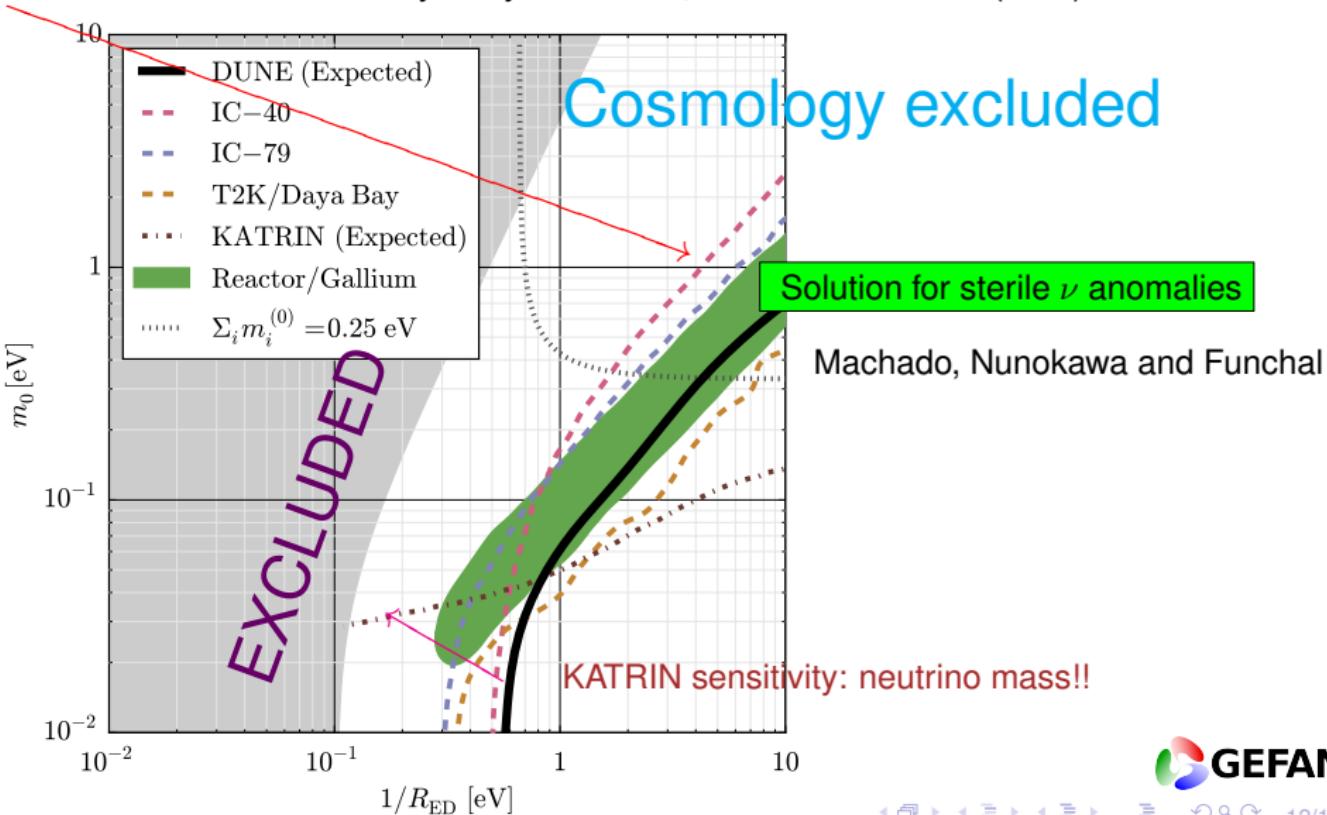


Constrains on large extra dimensions



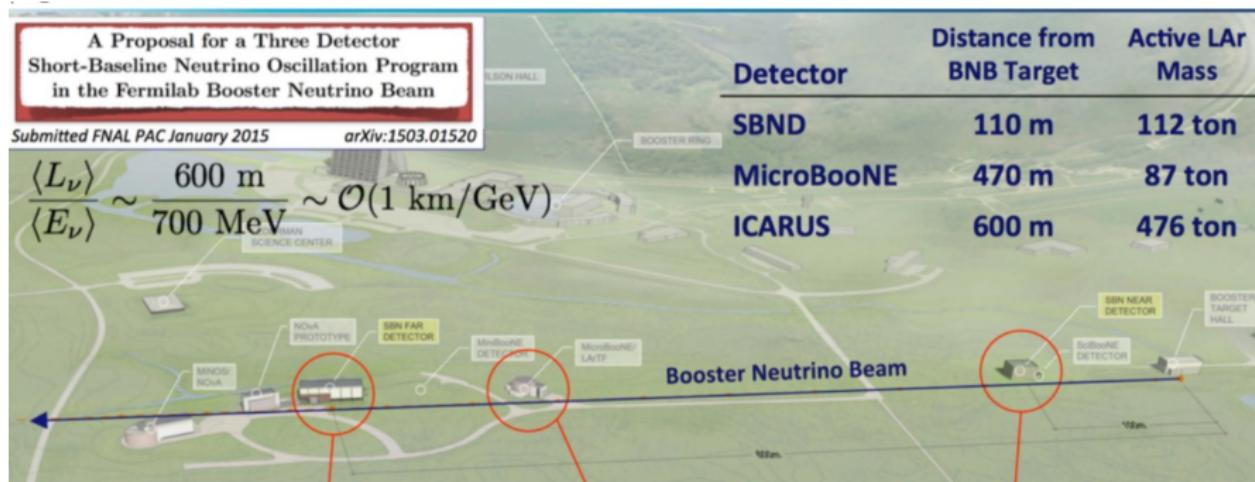
Constrains on large extra dimensions

IC-40/IC-79: ICECUBE analysis by O. L. G. P., Esmaili and Tabrizi (USP)



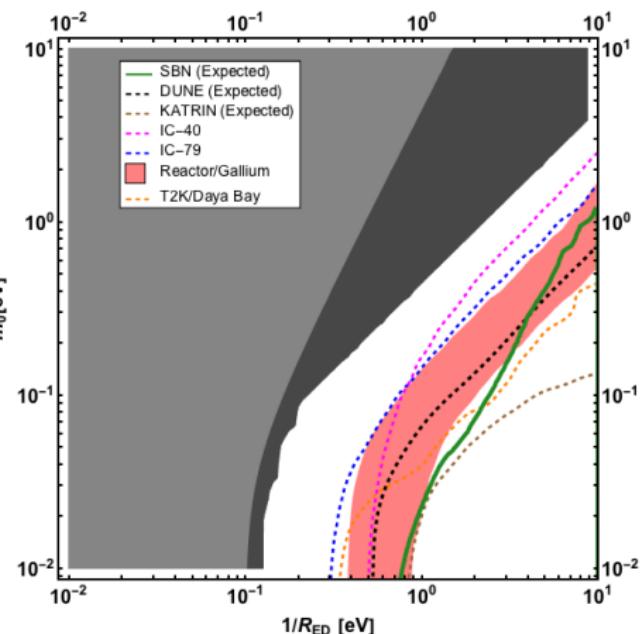
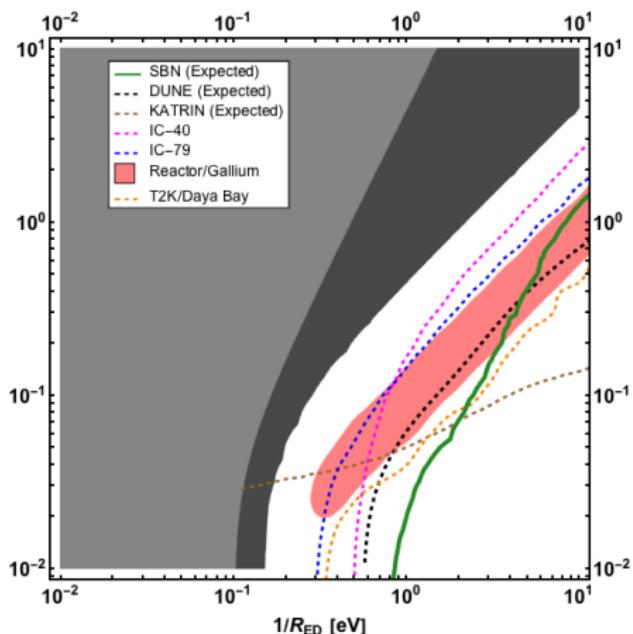
Constraints on short-baseline experiments: SBN

Short Baseline program for FERMILAB



Constraints on short-baseline experiments: SBN

Sensitivity region for SBN

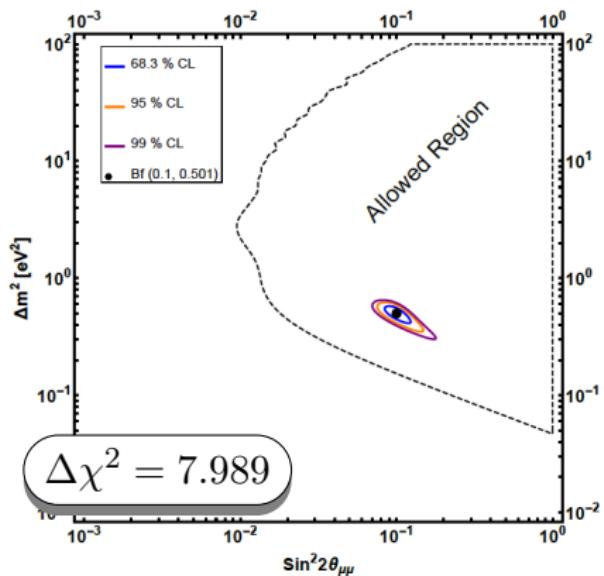


Work in progress, G. Stenico, D. Vanegas Forero and O.L.G.P.

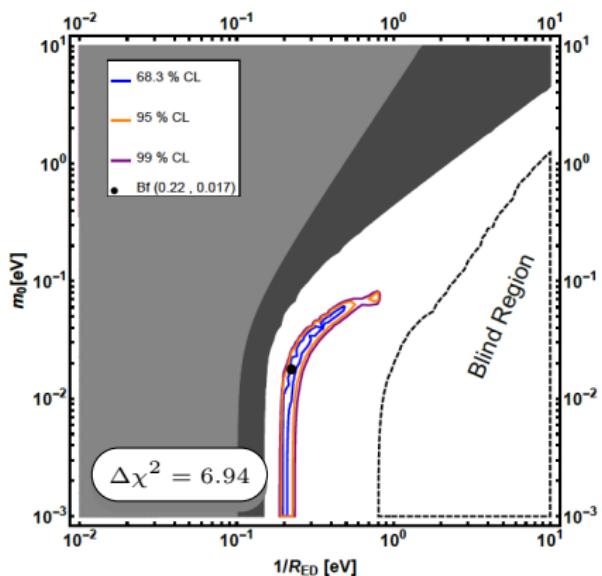


Discrimination power

TRUE LED: → 3+1 scenario



TRUE 3+1 model : → LED scenario



Conclusions

- ① Smallness of neutrino mass: motivated by large extra dimension models?
- ② Neutrinos in large extra dimension models: oscillation pattern and non-unitarity
- ③ Future experiments can provide very good constrains on size of extra dimension
- ④ Full discrimination between the extra dimension oscillation model and the usual three neutrino theory
- ⑤ ICECUBE (3× IC-70) $R_{\text{ED}} > 2,5(3,0) \times 10^{-5}$ cm for small m_0

Support : Projeto Temático FAPESP 2014/19164-6.



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