

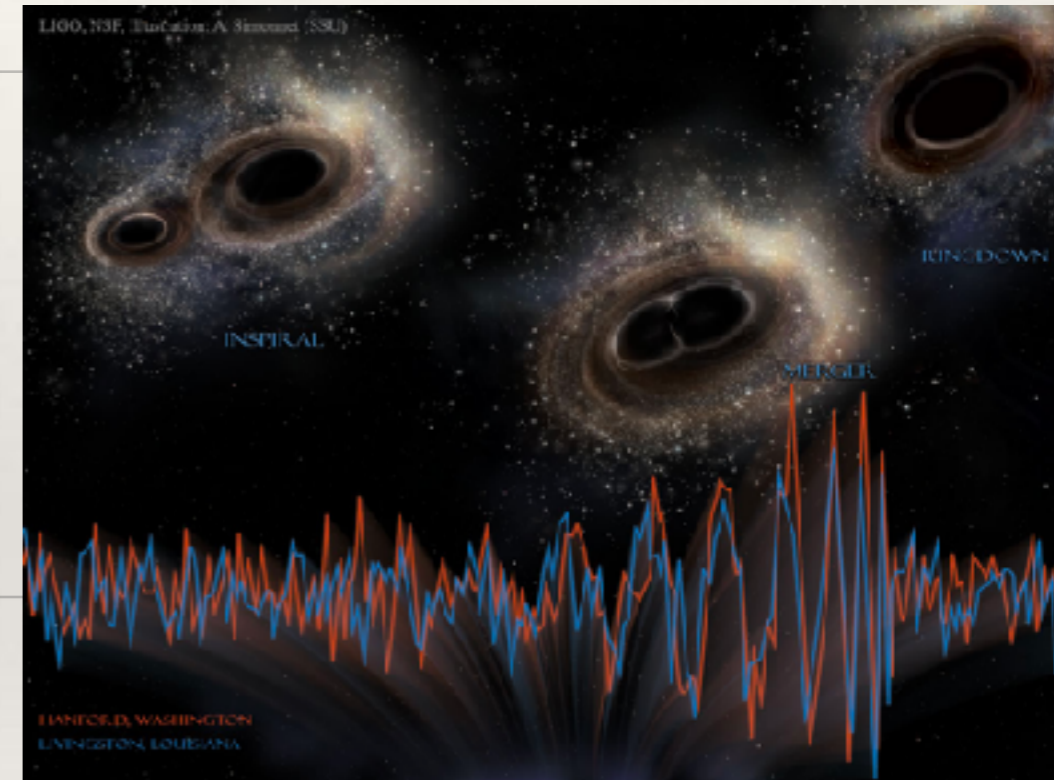
*Stanislav (Stas) Babak.*

*AstroParticule et Cosmologie, CNRS (Paris)*



# Gravitational waves

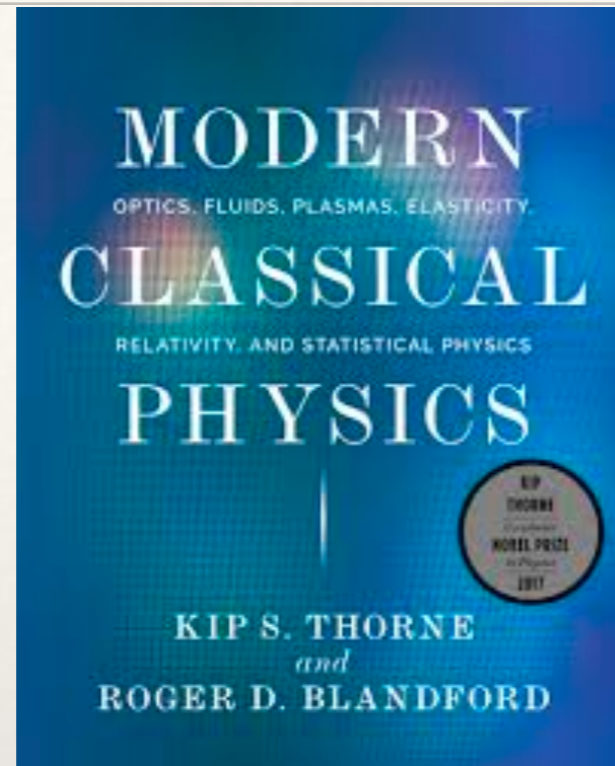
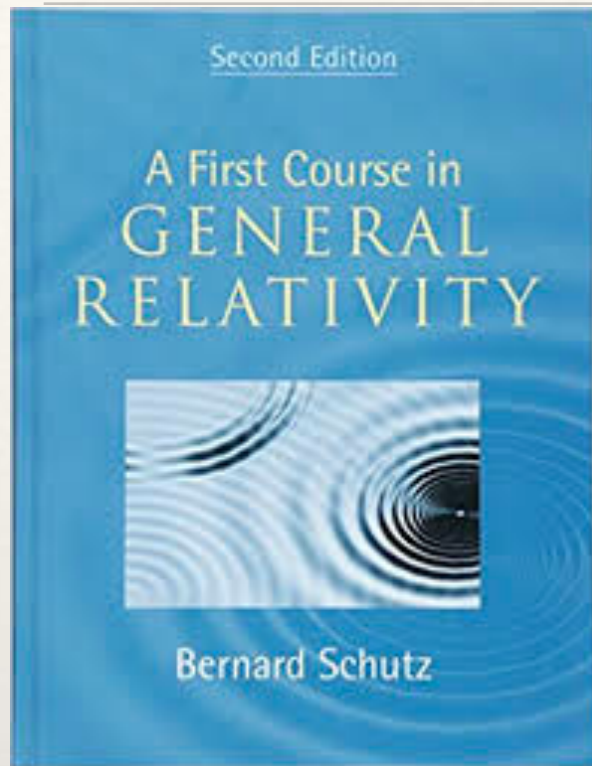
## Part I



ICTP, 18-22 June 2018

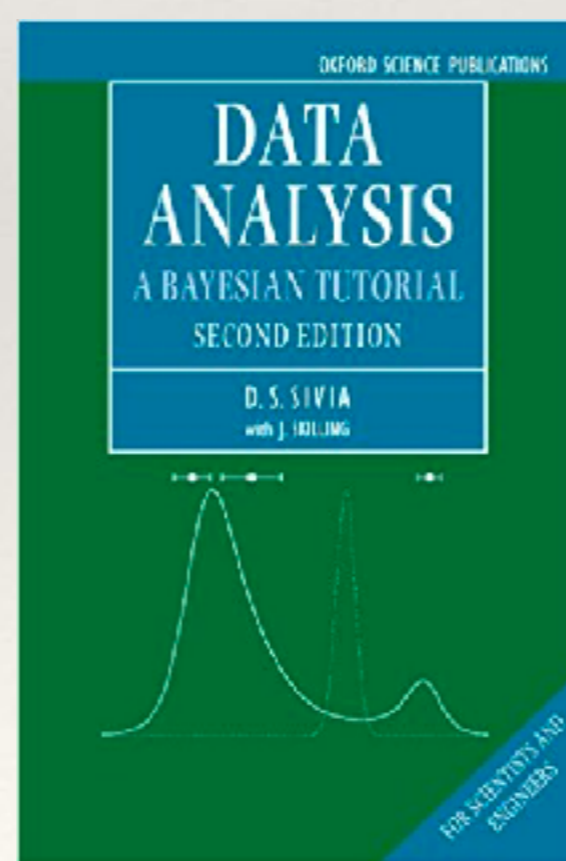
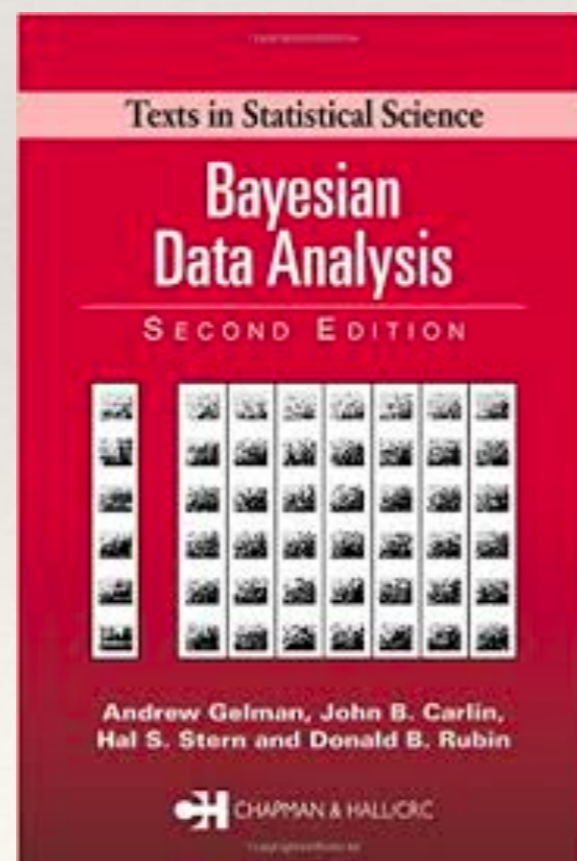
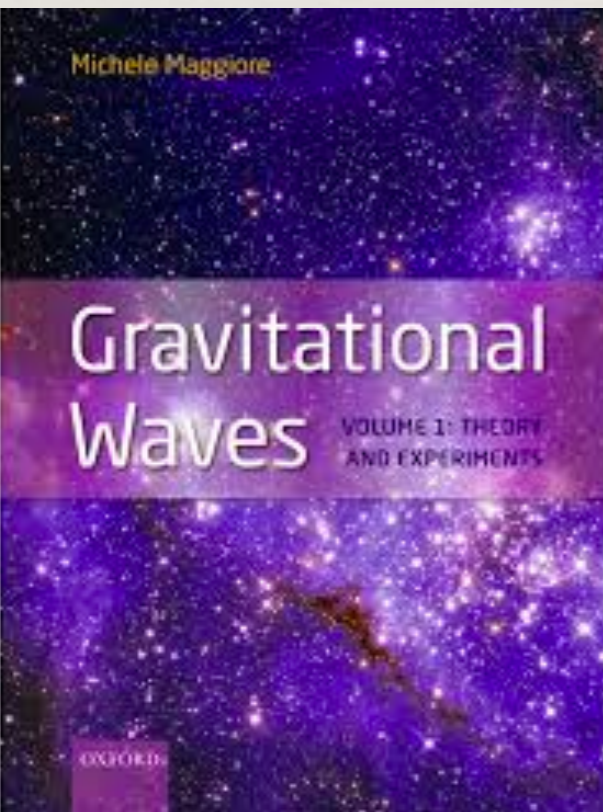


# Literature



Some figures in these lectures are borrowed from these books

LSC+Virgo, Phys. Rev. Lett. 116, 221101 (2016), LSC+Virgo, Phys. Rev. X 6, 041015 (2016), LSC+Virgo, Phys.Rev.Lett. 116 241102 (2016), LSC+Virgo, Phys.Rev. X6 041014 (2016), LSC+Virgo, Phys. Rev. Lett. 118, 221101 (2017), Berti et al., Class.Quantum Grav. 32, 243001 (2015), LSC+VIRGO, ArXiv: 1805.11579, S. Khan+, Phys.Rev. D93 (2016) 044007, LSC+Virgo [arXiv: 1805.11579](https://arxiv.org/abs/1805.11579), LIGO\_Virgo, Astrophys.J. 848 (2017) L12, LIGO+Virgo Phys.Rev.Lett. 119 (2017) 161101, Babak+ Phys.Rev. D95 (2017) 103012, LISA consortium [arXiv:1702.00786](https://arxiv.org/abs/1702.00786), Klein+ Phys.Rev. D93 (2016) 024003, LISA consortium, [arXiv:1305.5720](https://arxiv.org/abs/1305.5720), Amaro-Seoane+ Class.Quant.Grav. 29 (2012) 124016 .



# Lecture 3

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- 📌 Laser interferometer in space: LISA
- 📌 Sources and event rate
- 📌 Data analysis: Part 1

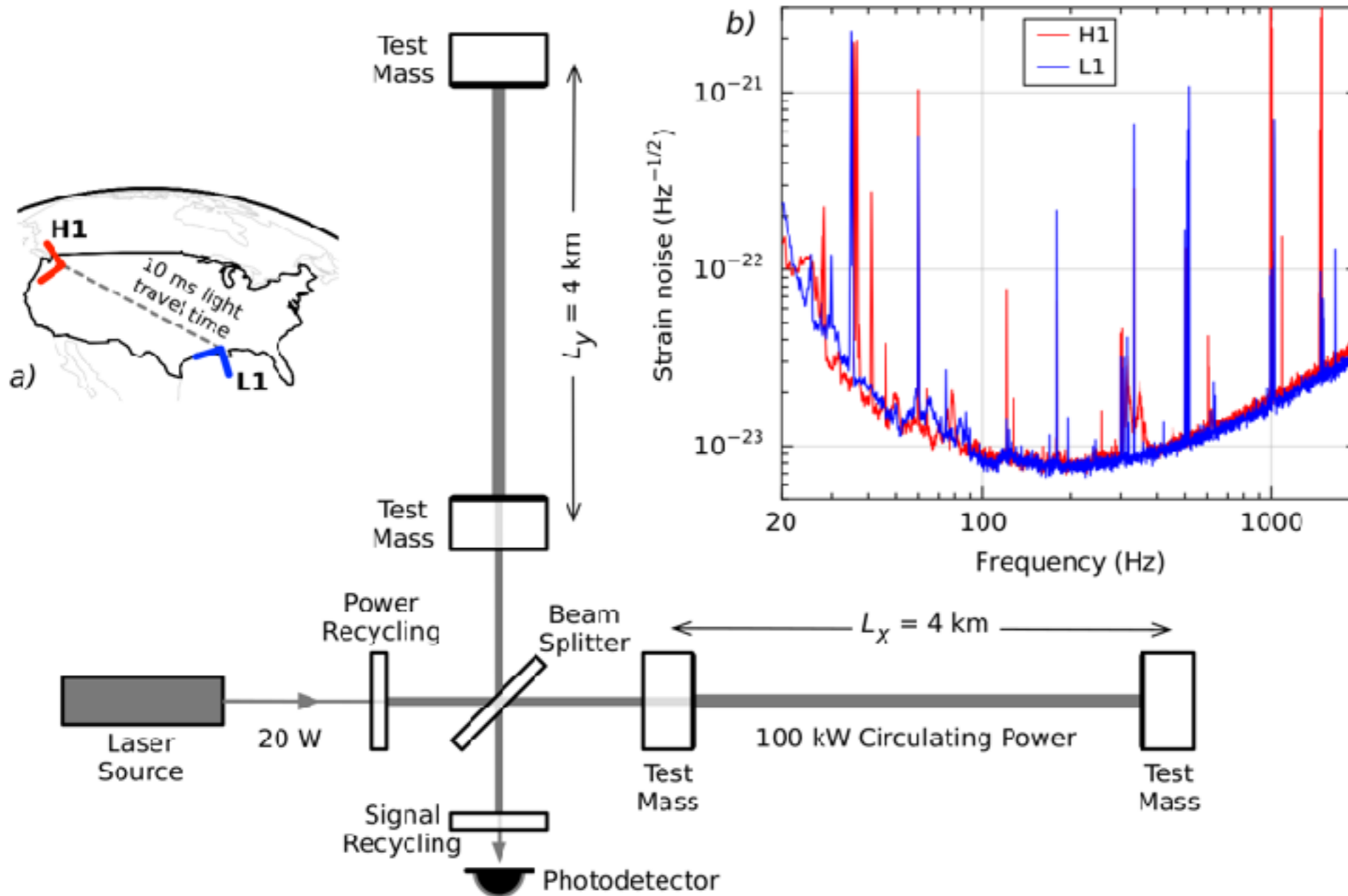




# LIGO/VIRGO

$$\Delta L = \frac{1}{2} h^{GW} L$$

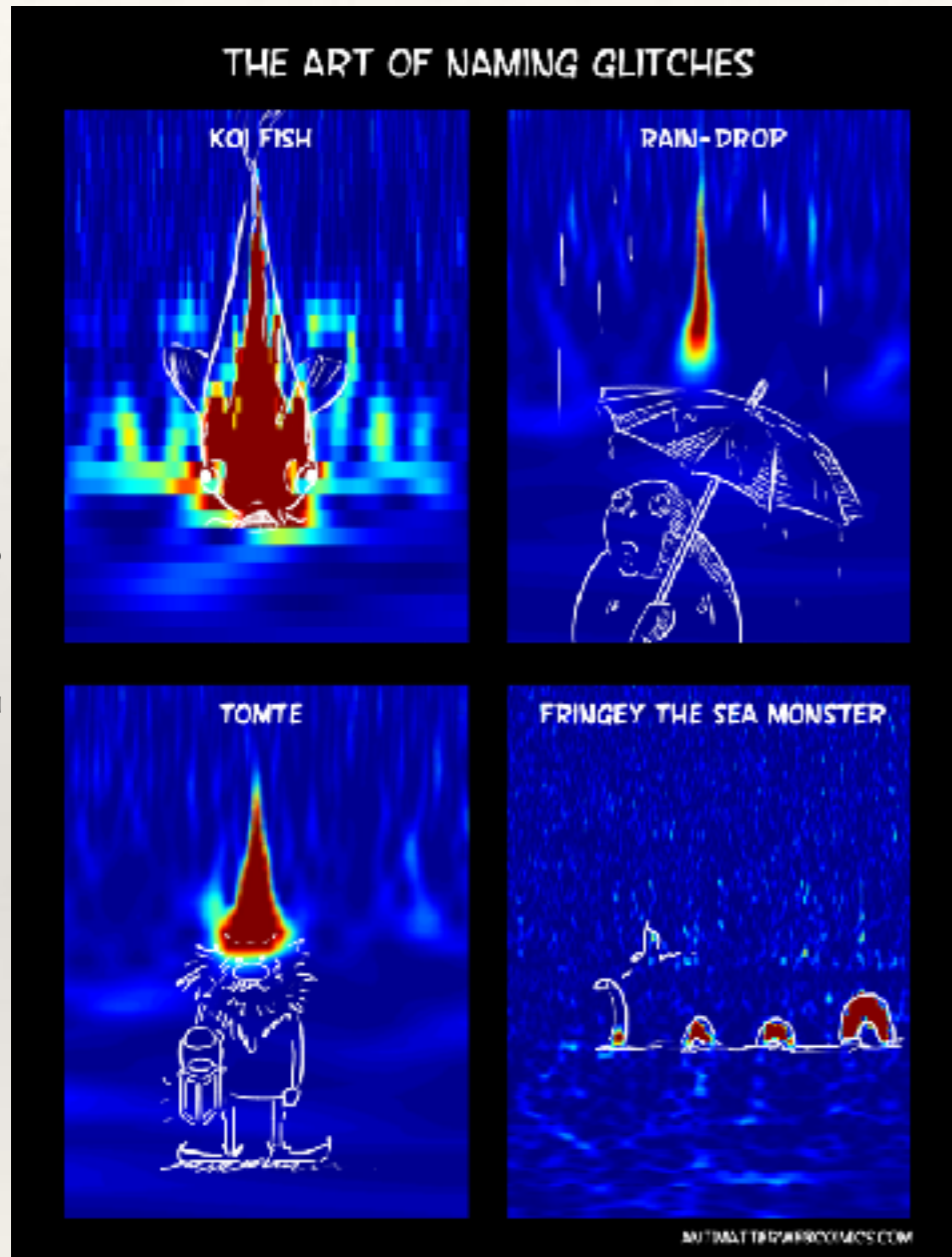
$$h \approx 10^{-22}, \rightarrow L = 4\text{km}, \rightarrow \Delta L \approx 10^{-16}\text{cm}$$





# LIGO/VIRGO noise

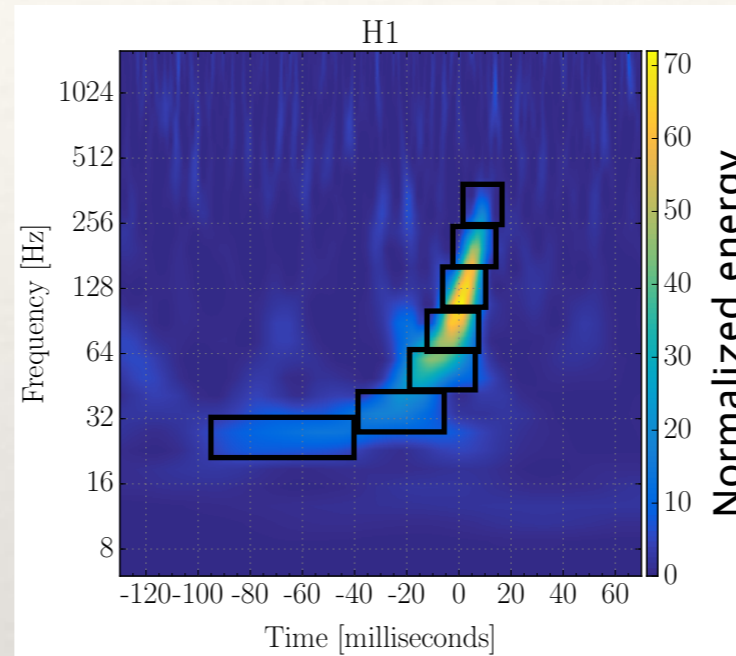
The noise in LIGO/VIRGO is stationary only on relatively short time scale



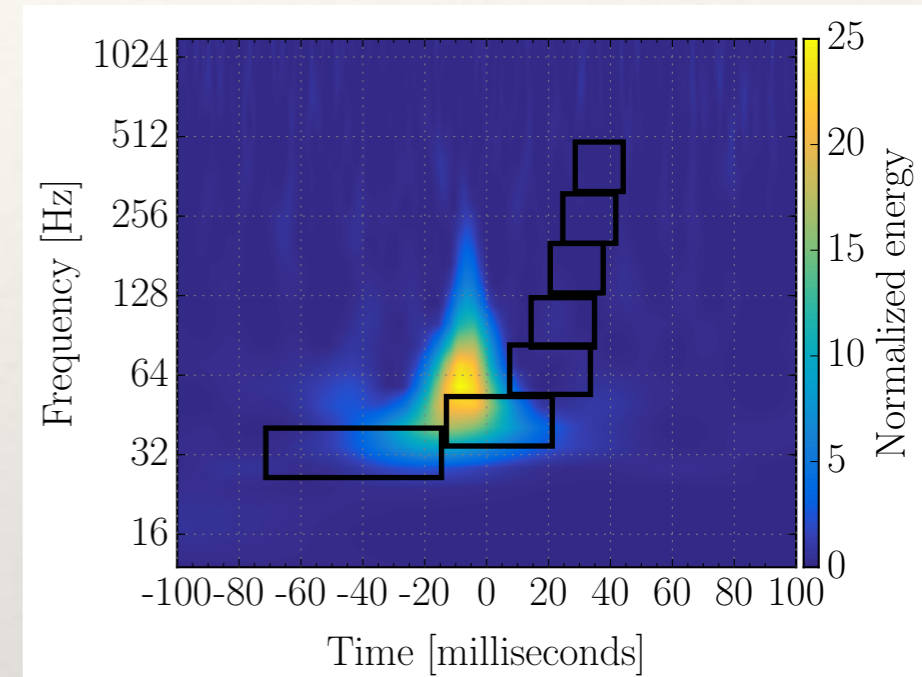
frequency

time

Real



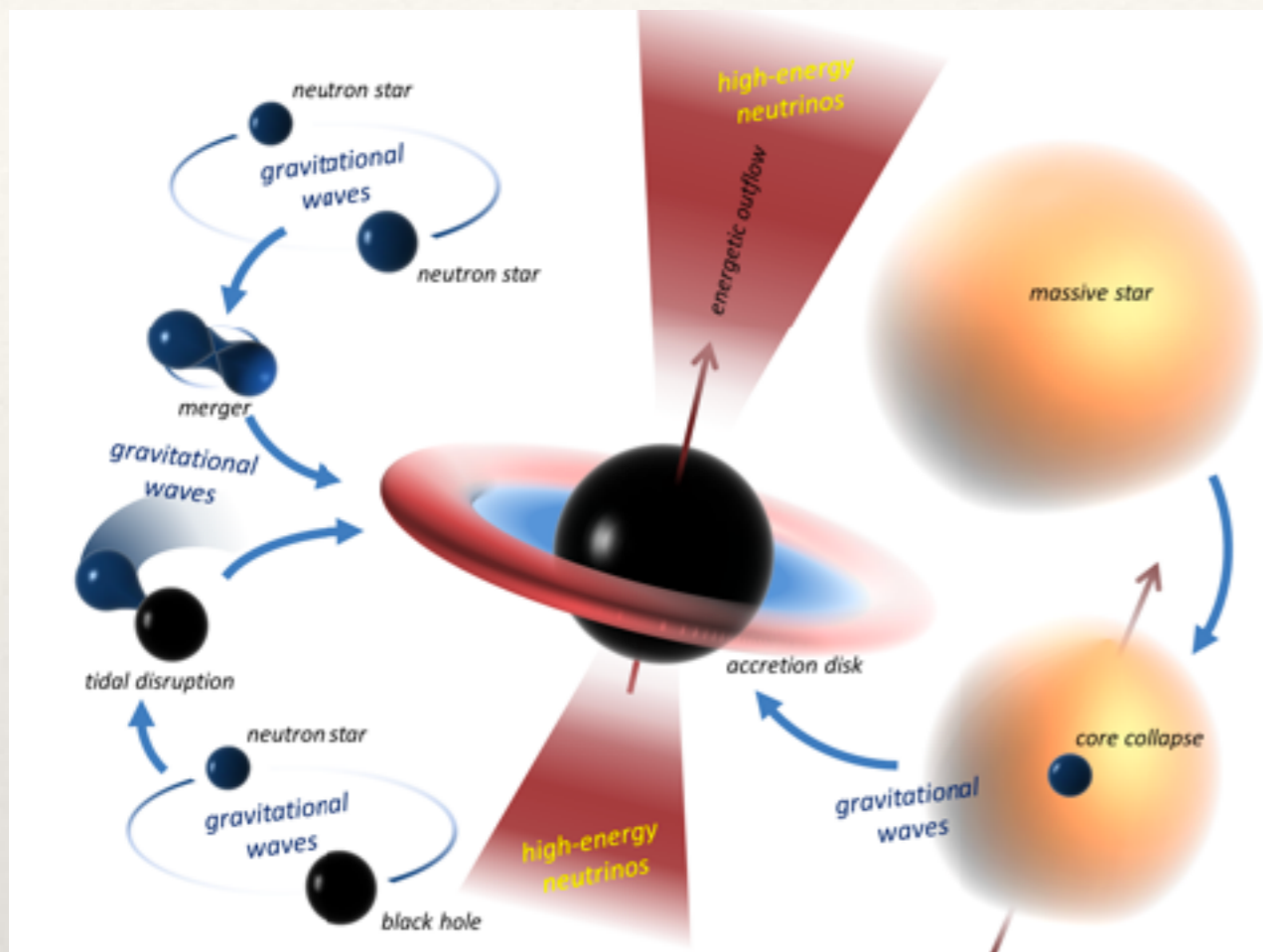
Instrumental



The noise is not Gaussian: need to introduce additional consistency checks into the detection statistic

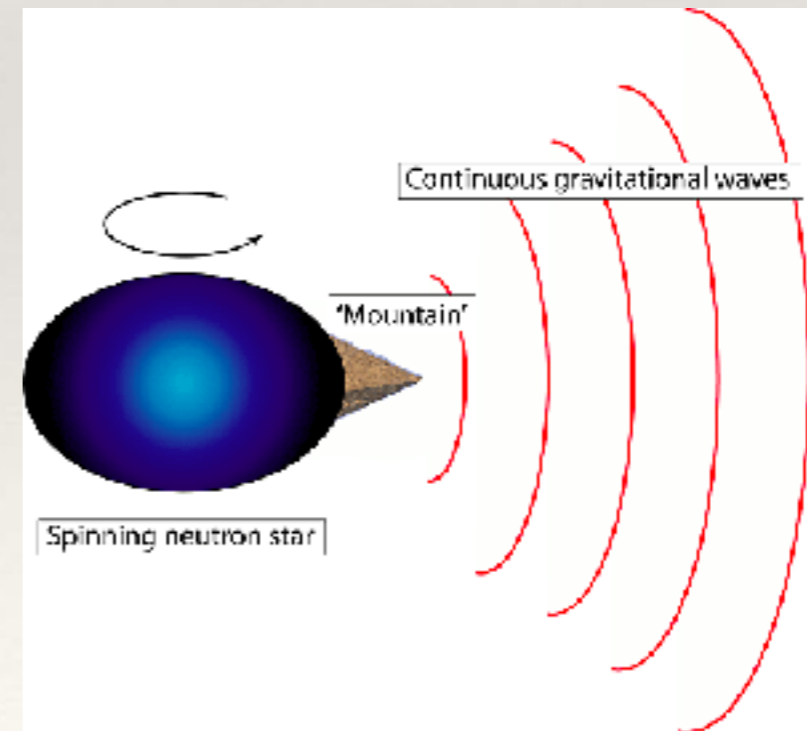


# Other GW sources in LIGO/VIRGO frequency band



Credit: I. Bartos/Based on arXiv:1212.2289

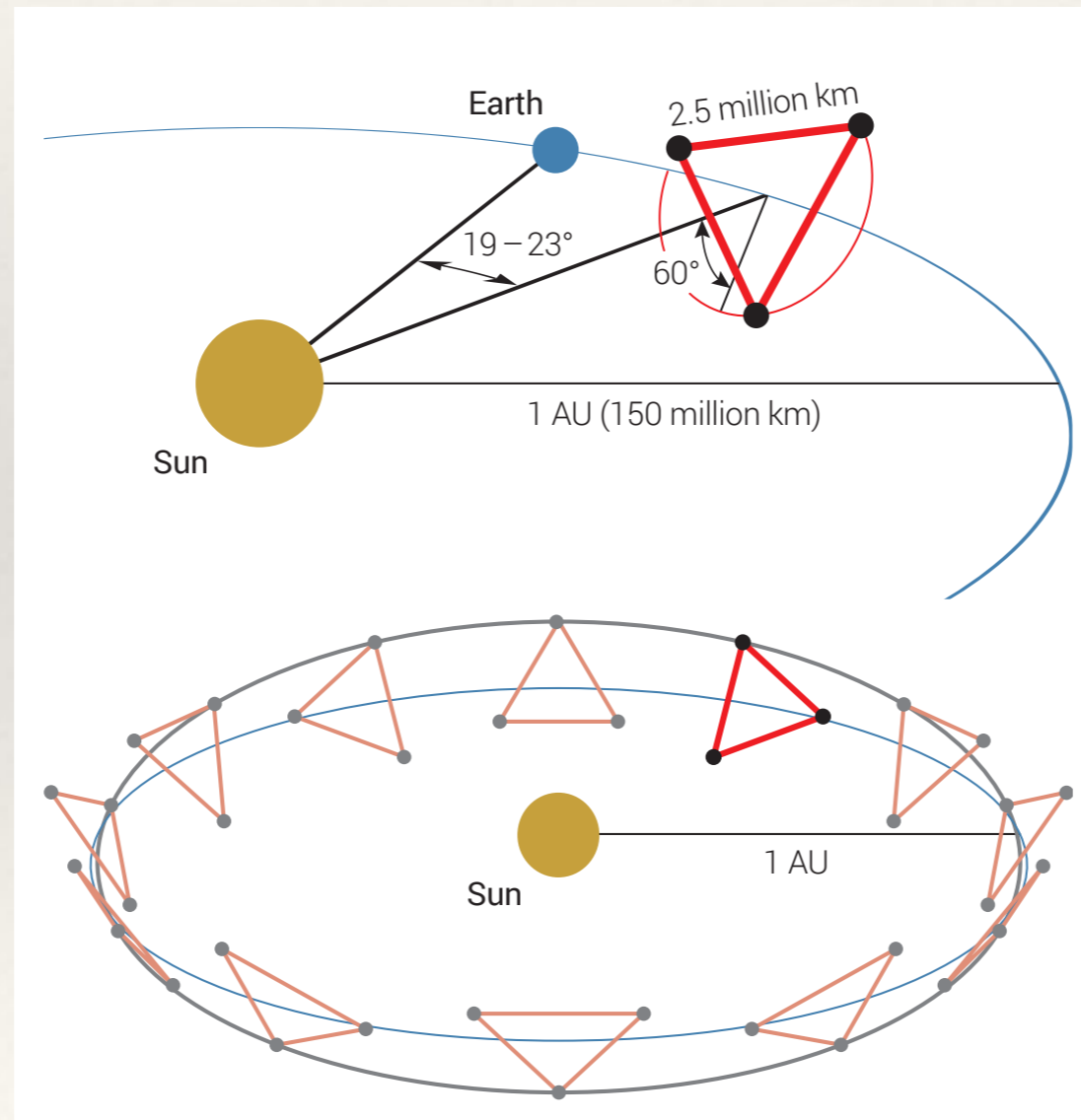
- Core collapse
- Merging NS binary systems
- Binary: NS-BH
- Monochromatic signal from deformed single NSs
- Stochastic GW signal from the energetic processes in early Universe





# LISA: Laser Interferometric Space Antenna

- LISA: GW observatory in space: The launch date 2032 - 2034. Leading by European Space Agency.
- LISAPathfinder - Technological mission to prove the technical readiness of LISA - fantastic results, order of magnitude better than minimum requirement



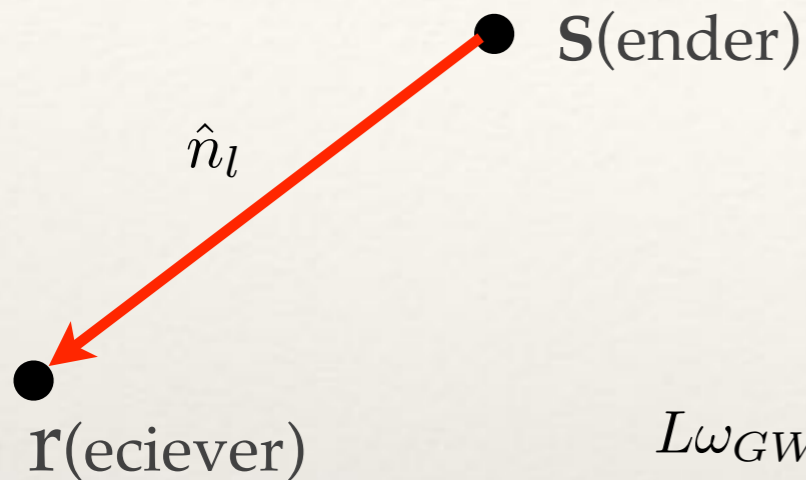
# LISA (cartoon)





# LISA

## Principle of measurement



LISA: three satellites in free falling orbits around the sun, constellation forms equilateral triangle

L1~L2~L3= 2.5 mln. km.

Operates in freq. range 0.1 mHz - 0.1 Hz.

Exchange laser light - measurement of the proper distance between satellites.

$L\omega_{GW} = 1 \rightarrow f_{GW} \approx 20mHz$  Long wavelength is not applicable

We cannot cover the detector by LIF, use “TT” frame: in this frame GW can be seen as affecting the phase (or frequency) of the laser light.

Change in laser freq. due to GWs  $\rightarrow \frac{\Delta\nu}{\nu_0} = \frac{n_l^i n_l^j \Delta h_{ij}}{2(1 - \hat{k} \cdot \hat{n}_l)}$

unperturbed freq. of a laser  $\rightarrow \nu_0$

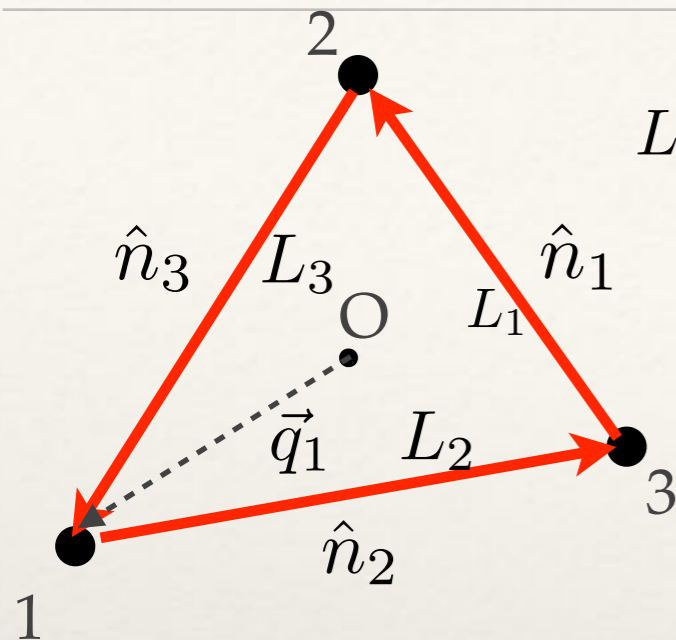
direction of GW propagation  $\rightarrow \hat{k}$

$\Delta h^{ij} = h^{ij}(t_s) - h^{ij}(t_r)$

$$t = t_r, \quad t_s = t - |\vec{R}_r(t) - \vec{R}_s(t_s)| \approx t - |\vec{R}_r(t) - \vec{R}_s(t)| \approx L_l$$



# LISA



$L_1 \approx L_2 \approx L_3 \approx L$  — but not exactly!

$$\frac{\Delta\nu}{\nu_0} \rightarrow y_{slr} \quad \text{— for a single link}$$

Consider GW signal from a binary system

$$y_{slr}^{GW} = -i \frac{\omega L}{2} A_l(\iota, \psi, \theta_s, \phi_s) e^{i\Phi(t - \hat{k} \cdot \vec{R}_0)} e^{-i\omega \hat{k} \cdot \vec{q}_r} \text{Sinc} \left[ \frac{\omega L}{2} (1 - \hat{k} \cdot \hat{n}_l) \right] e^{-i \frac{\omega L}{2} (1 - \hat{k} \cdot \hat{n}_l)}$$

- $A_l$  Amplitude of GW times antenna beam function
- Sinc — zero of sinc function gives freq. of GW signal which cannot be measured (f-n of sky position) - wiggles in the sensitivity at high frequencies
- Phase:  $t - \hat{k} \cdot \vec{R}_0(t)$  Doppler modulation (dominant) due to relative motion of the detector and the source

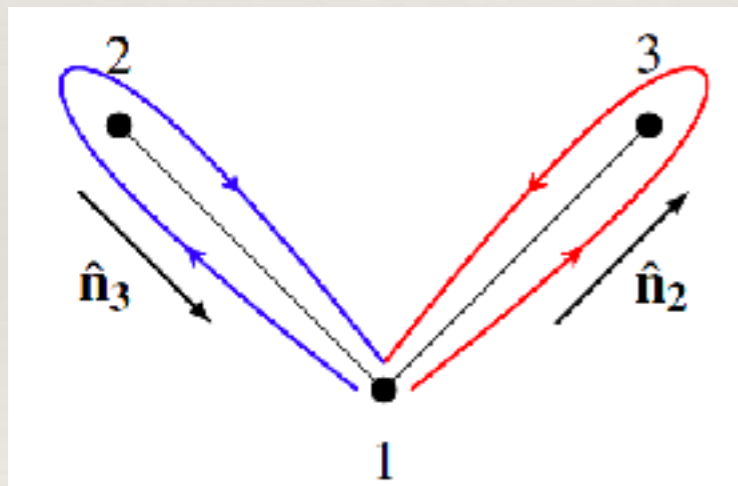




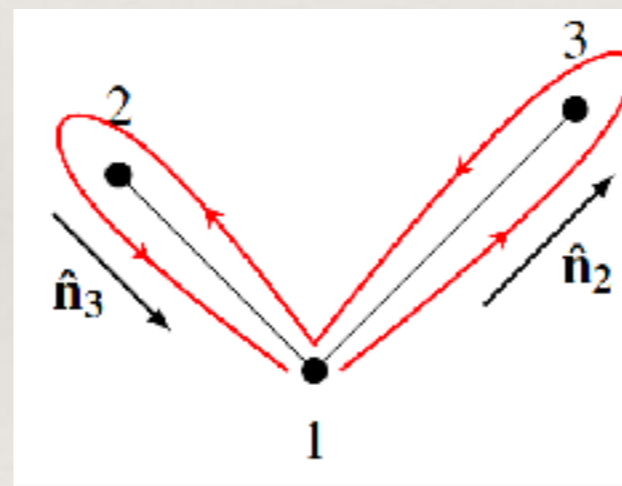
# LISA

- The localization of the sources in the sky comes from the Doppler modulation of the phase and from the amplitude modulation (time dependent antenna beam f-n)
- The term(s) dependent on the position of each spacecraft explicitly (q-vectors): important at very high frequency: constellation “feels” GW propagation

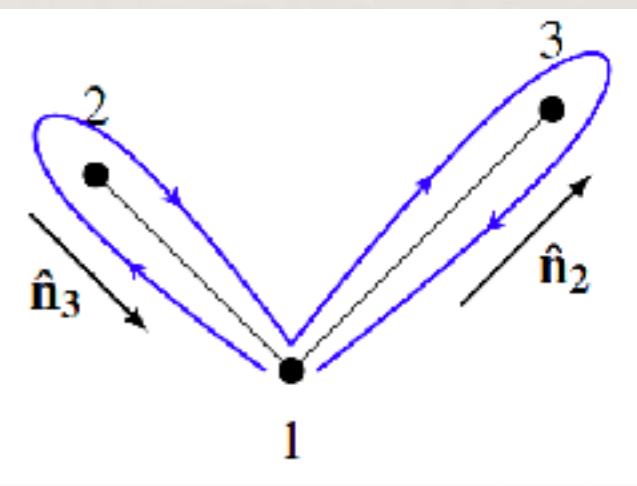
**TDI** — Time Delay Interferometry: technique which we apply to cancel the laser noise (in post-processing the data)



Equal arm Michelson

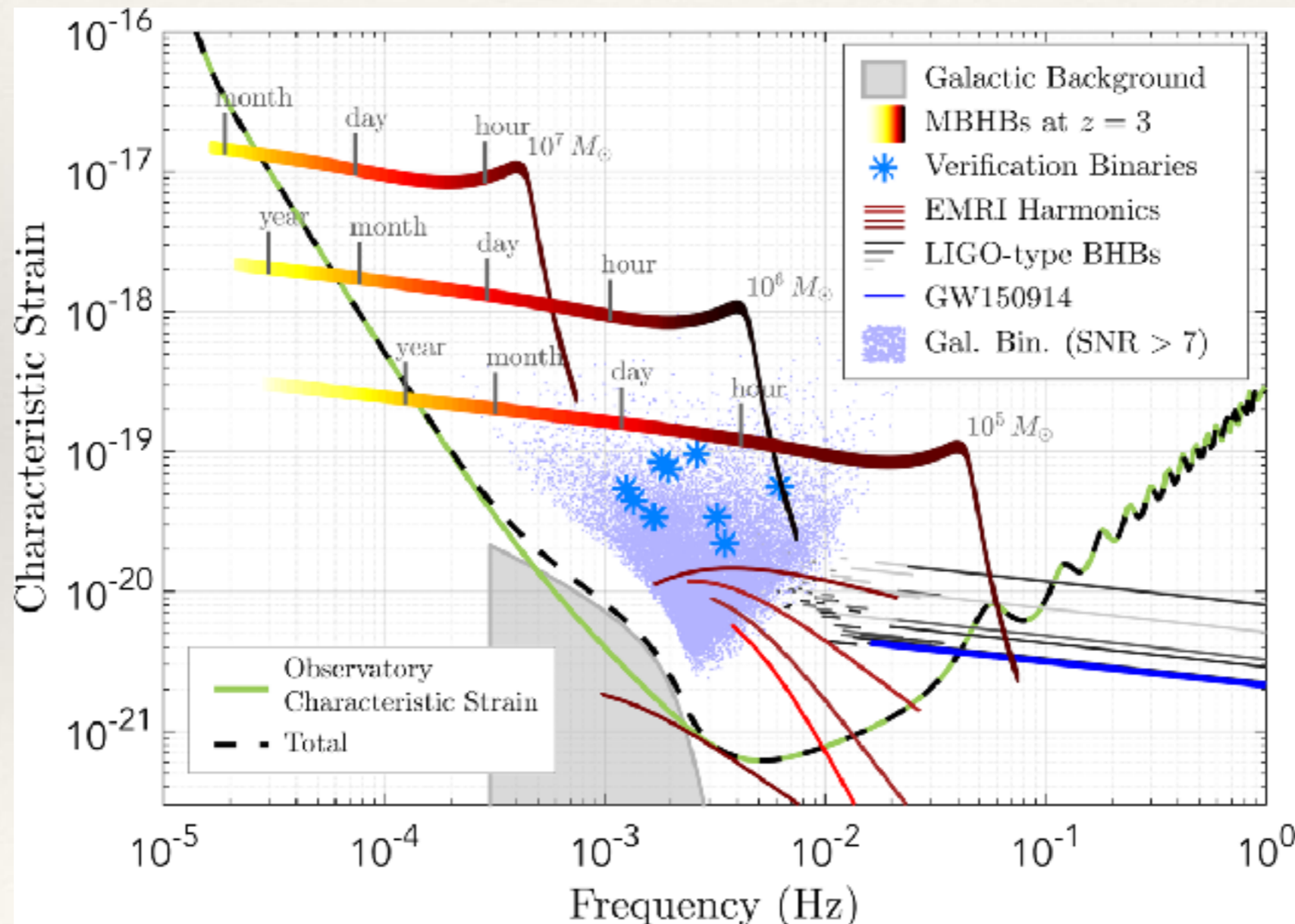


Unequal arm Michelson



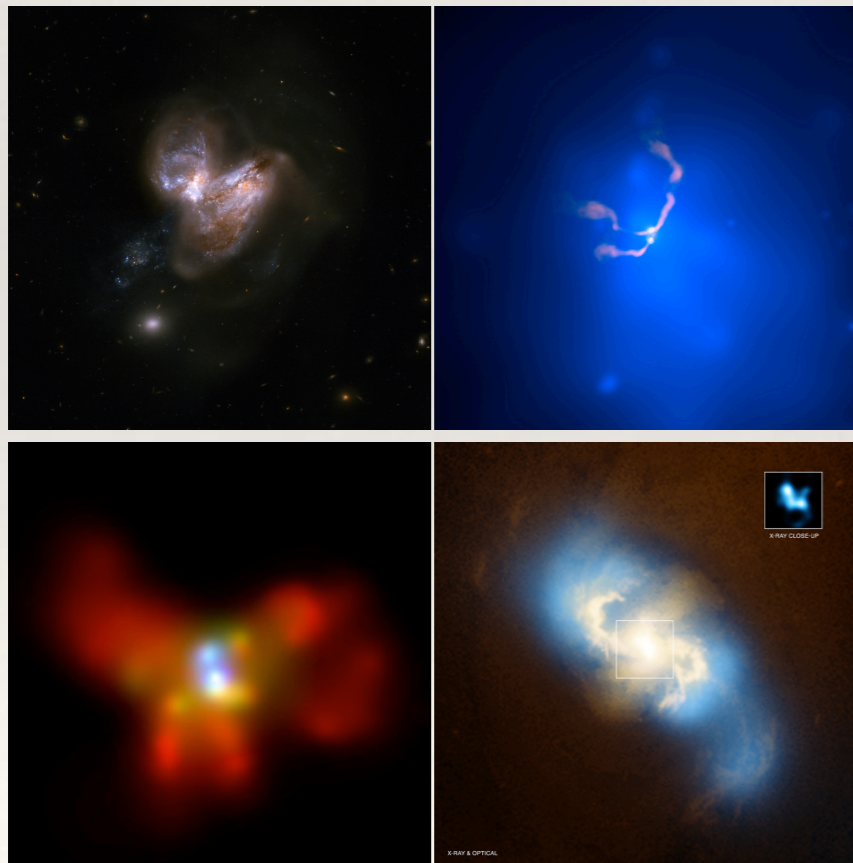
# GW sources in LISA band

- GW signals in LISA are strong and long-lived.
- LISA data will contain thousands of GW signals simultaneously: need to separate and characterize them
- Non-stationary noise



# LISA sources

- We believe that all galactic nuclei host Massive Black Holes: Milky Way has 4 mln. solar mass BH
- Galaxies merge: we can form Massive Black Hole Binary (MBHB) system
- We need stars and gas to bring MBHs close together for GW to be efficient (binary is merging within Hubble time)



[Credits: Hassinger+, VLA, Chandra, NASA]



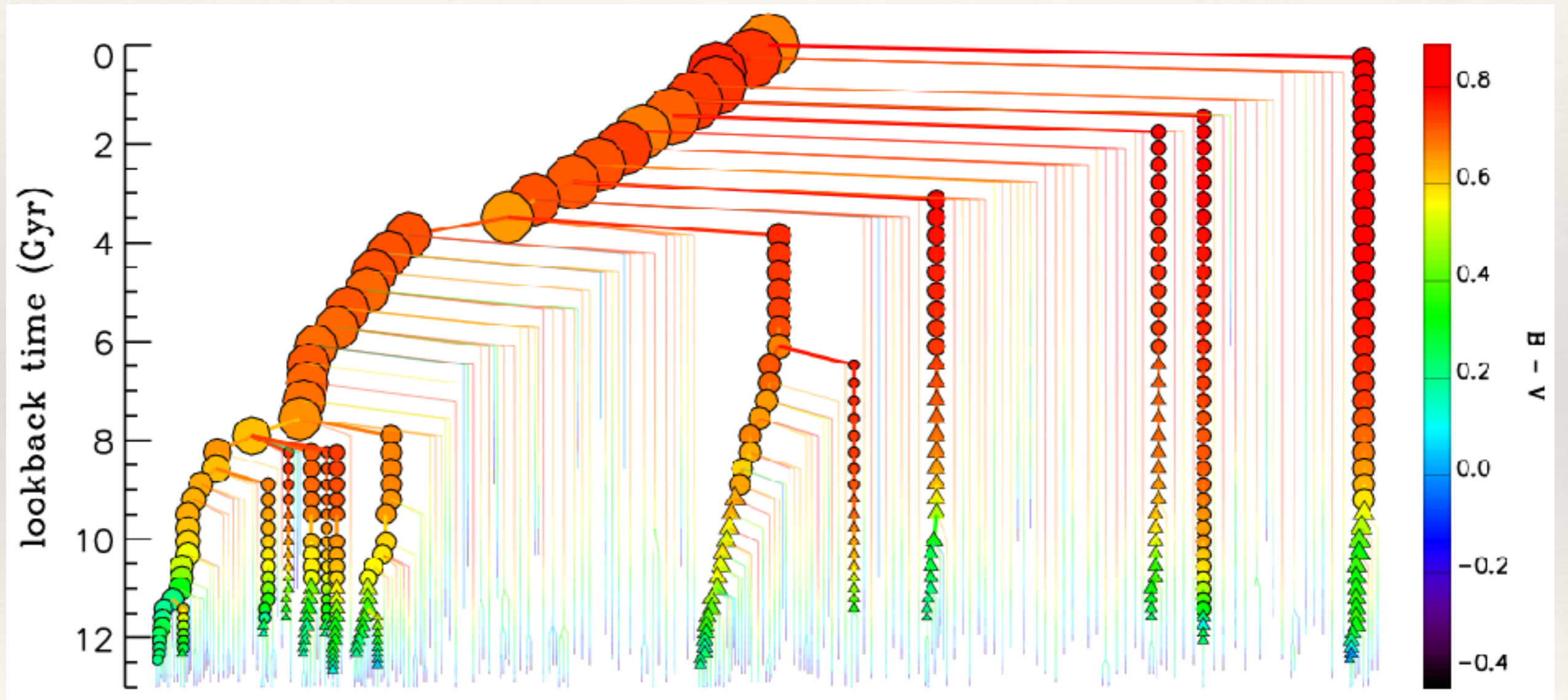
[Image: Hubble telescope]





# LISA sources: MBHB

MBHBs are formed from the initial BH seed. Those seeds could be “light” remnant of the first generation of stars or “heavy” from the direct collapse of a giant gas cloud. BHs accumulated the mass through gas accretion and merging.

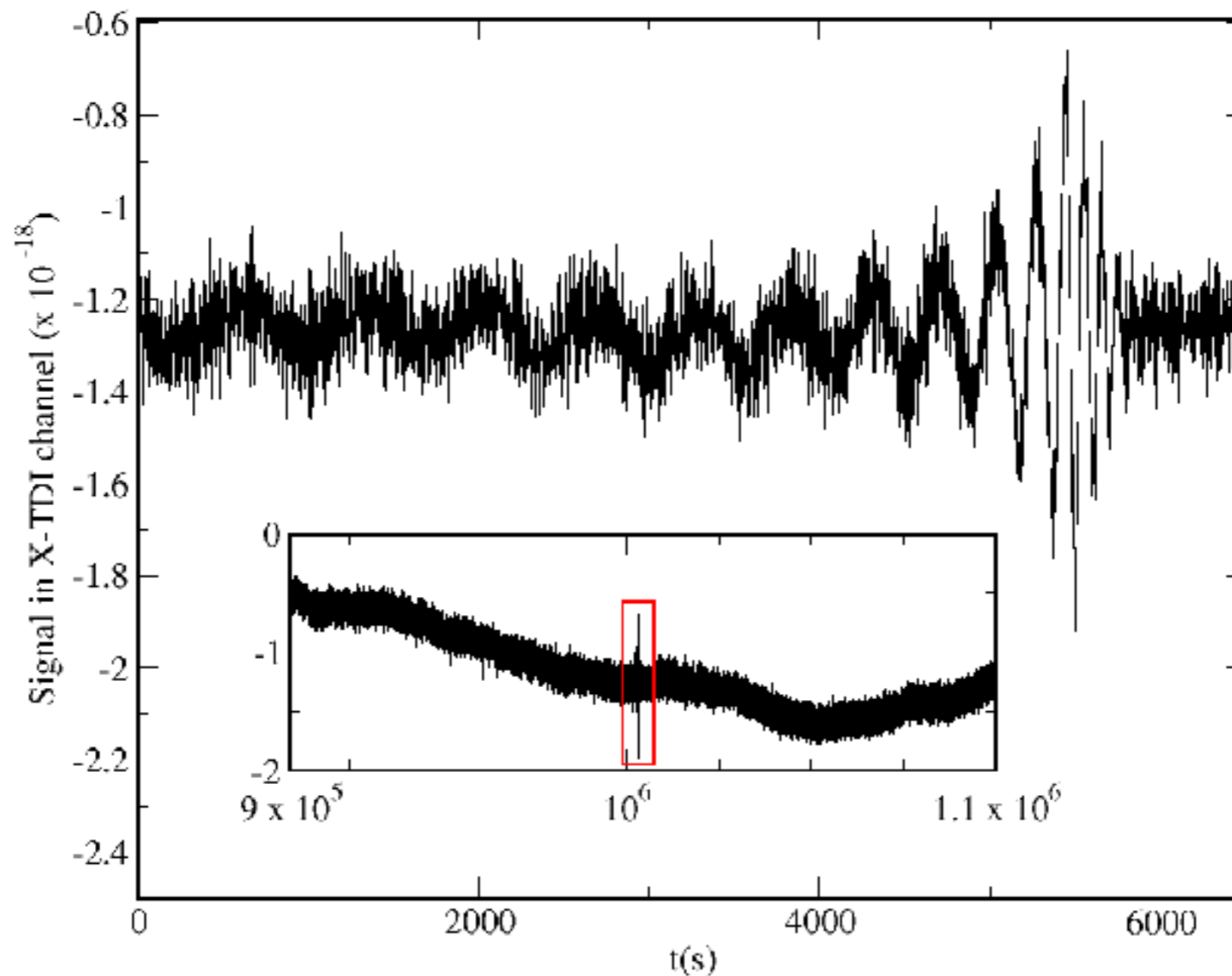


[Credits: Gabriella De Lucia]

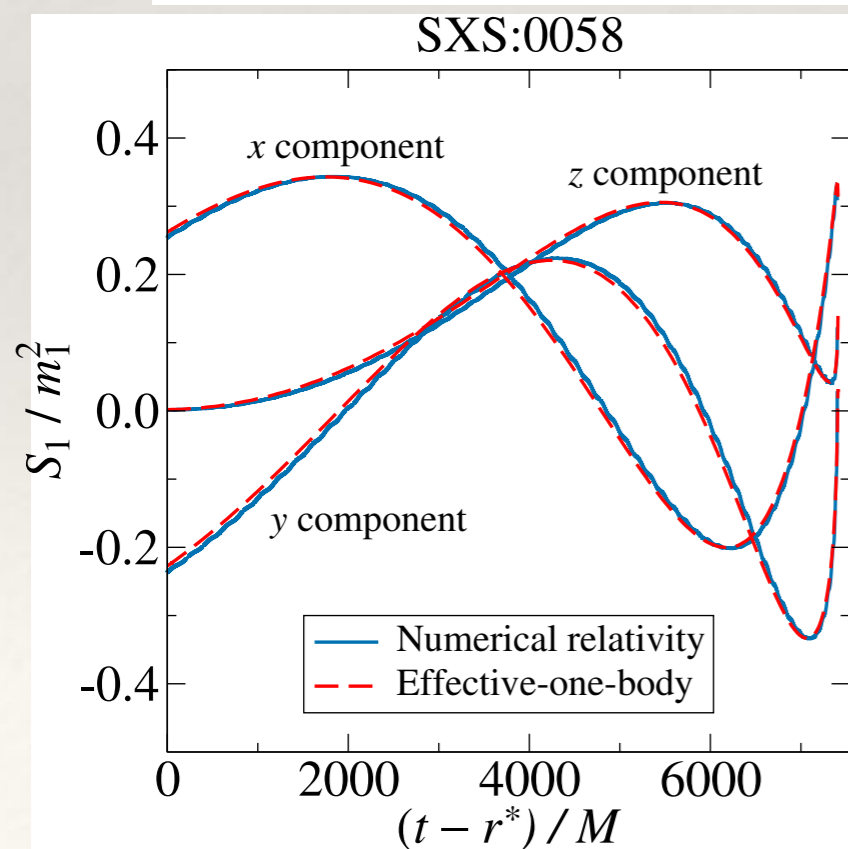
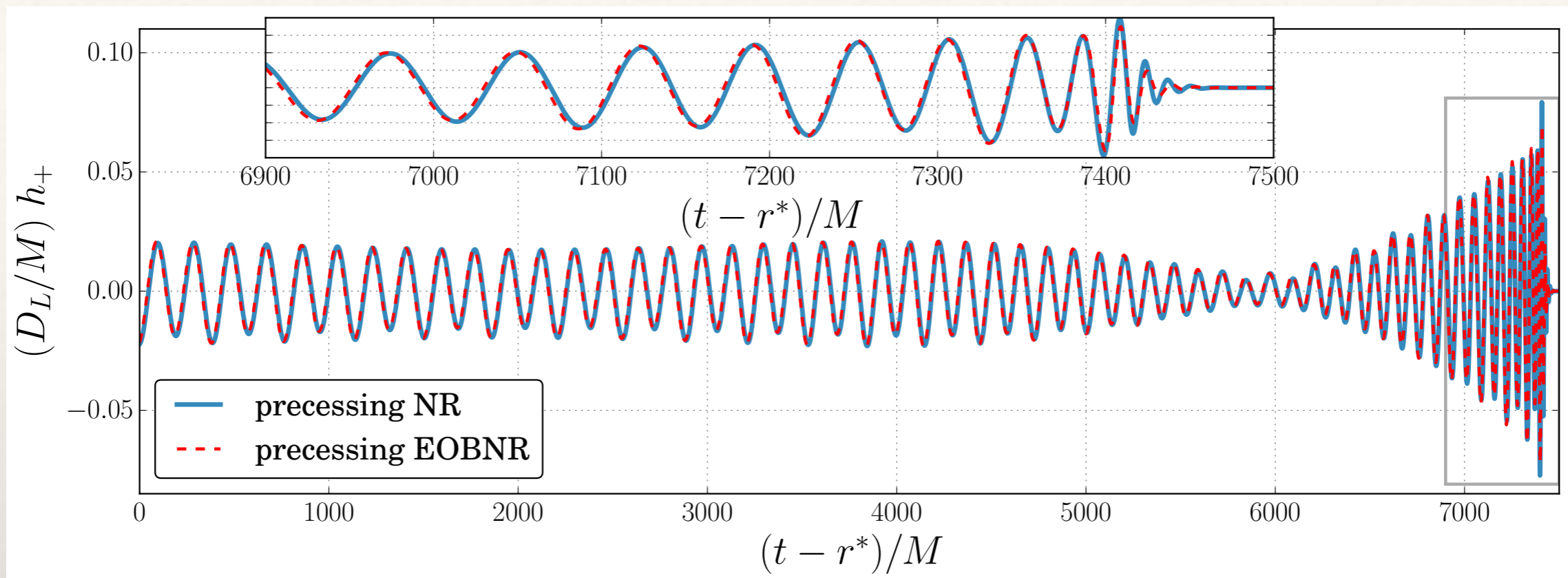


# LISA: GW signal from MBHB

The signal from MBHB is similar to what we have observed in LIGO (scaled up in the amplitude and stretched in time). GW signal from MBHB is expected to be the strongest signal (seen by eye in the simulated data). Imposes stringent demands on the accuracy of GW signal modelling



# Precessing binary





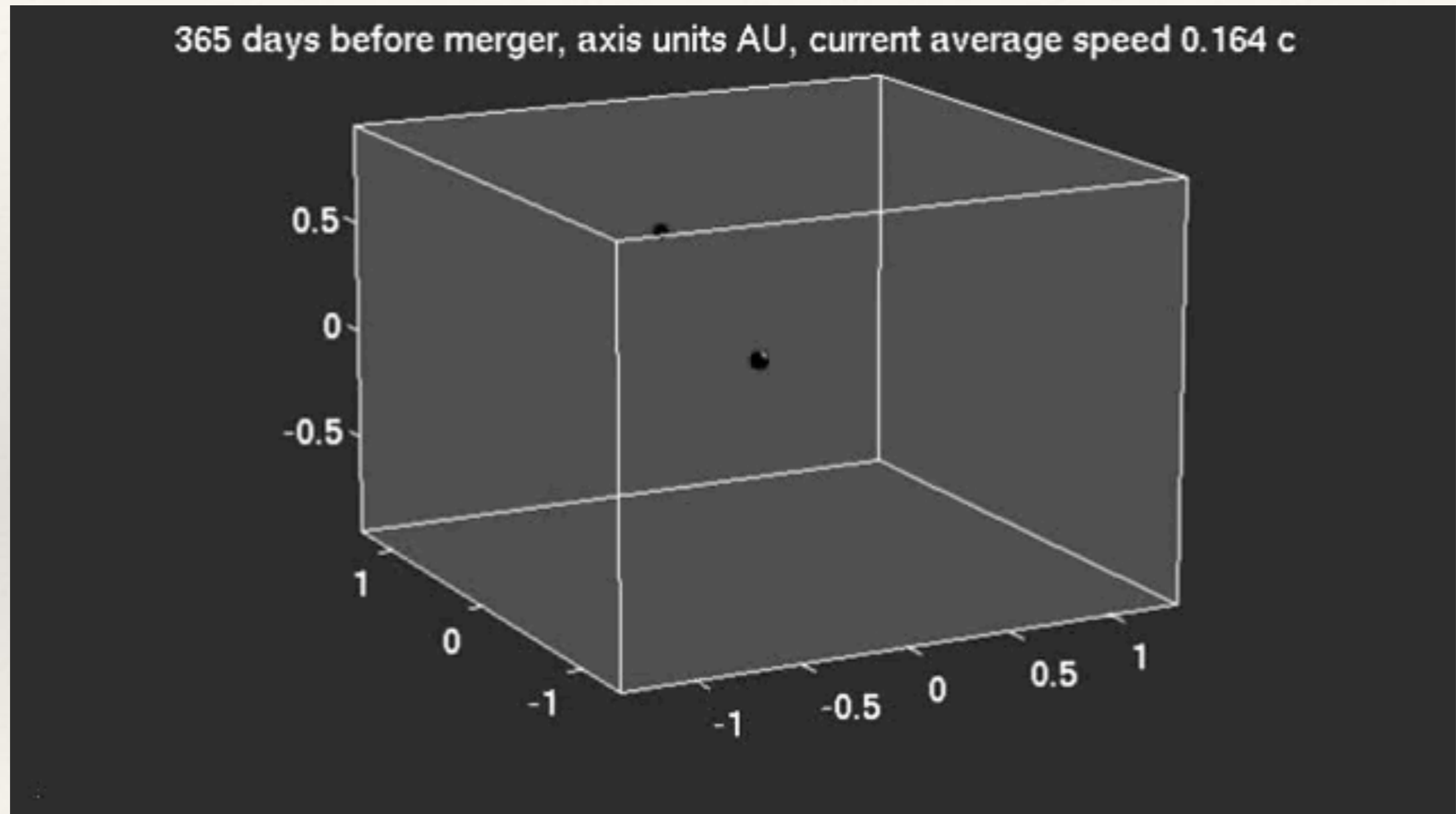
# EMRIs (extreme mass ratio inspirals)

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- Massive BHs in galactic nuclei surrounded by stars and gas with quite high density
- MBH could capture a compact object (BH, NS, WD) which is thrown on a very eccentric orbit (due to N-body interaction). The orbit shrinks and circularizes due to grav. radiation.
- EMRI: binary system with extreme mass ratio of component  $10^{-7}$  -  $10^{-5}$
- Compact object revolves  $10^6$  orbits in the proximity of MBH before the plunge.



# EMRI

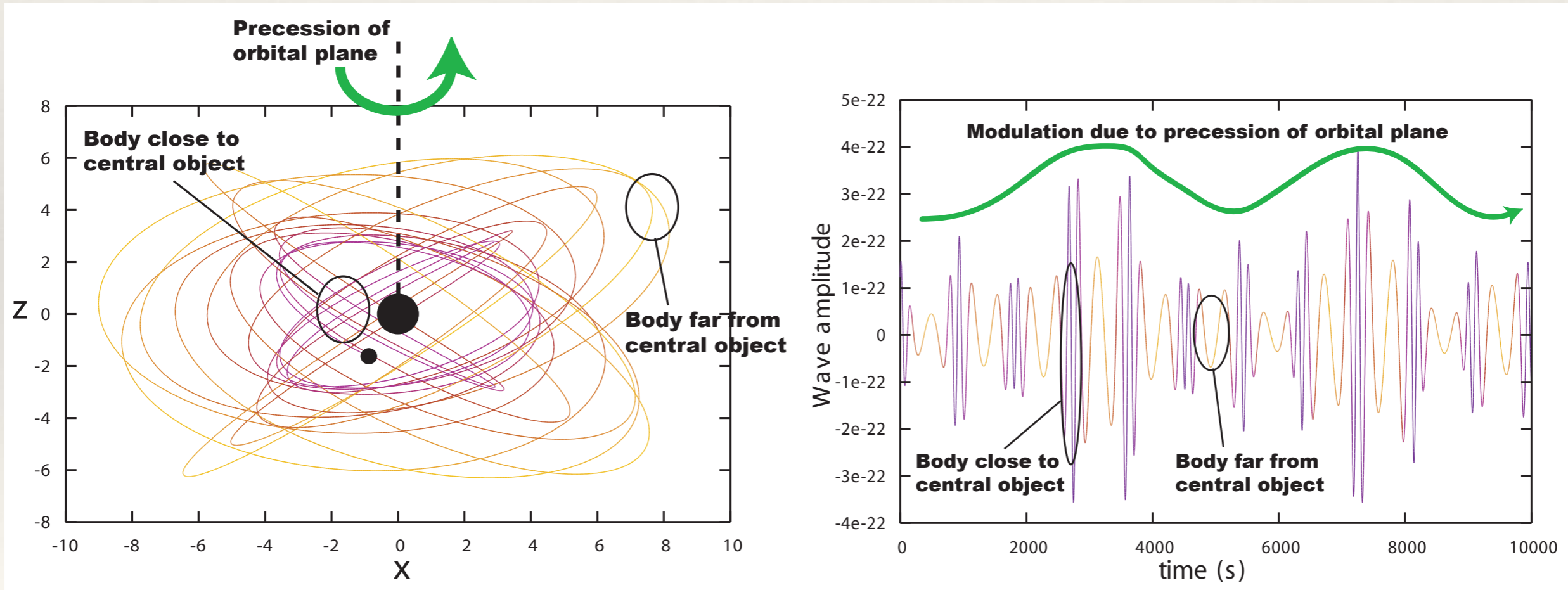


[Credits: S Draco, CalTech]



# EMRI

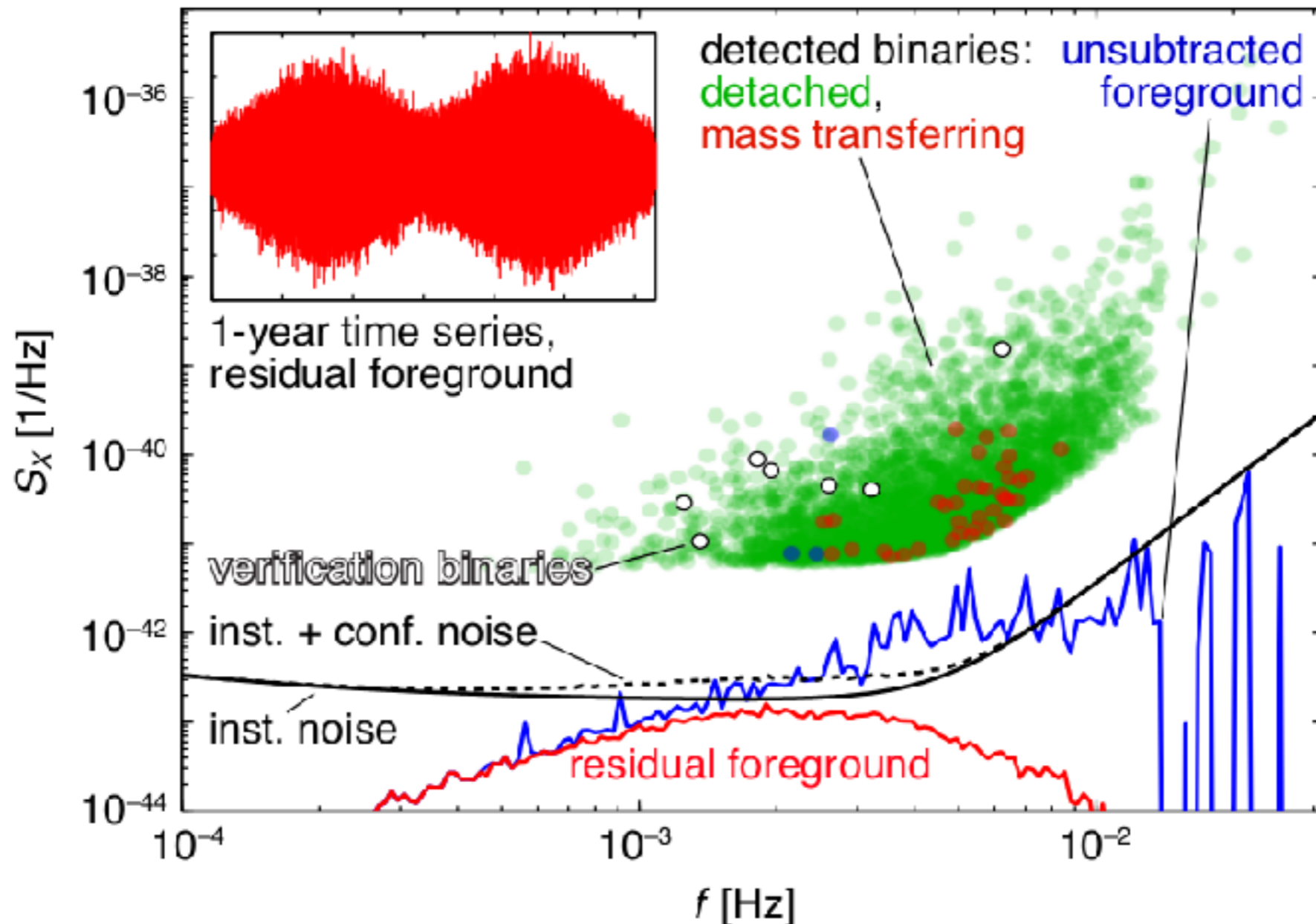
- Orbital motion: (almost) elliptical with a strong relativistic precession + orbital precession due to spin-orbital coupling
- Signal is very rich in structure (hard to detect but gives a lot of information)
- Ultra-precise parameter determination (if detected). Can map spacetime of a heavy object: holidodesy





# Galactic white dwarf binaries

- We expect to have  $10^7$  WD binaries all emitting GWs in the LISA band, only  $10^4$  can be resolved individually, other form stochastic GW signal (foreground)
- GW signal is almost monochromatic
- Verification binaries: known from current e/m observations (+GAIA,+ LSST)



# Expected event rate in LISA

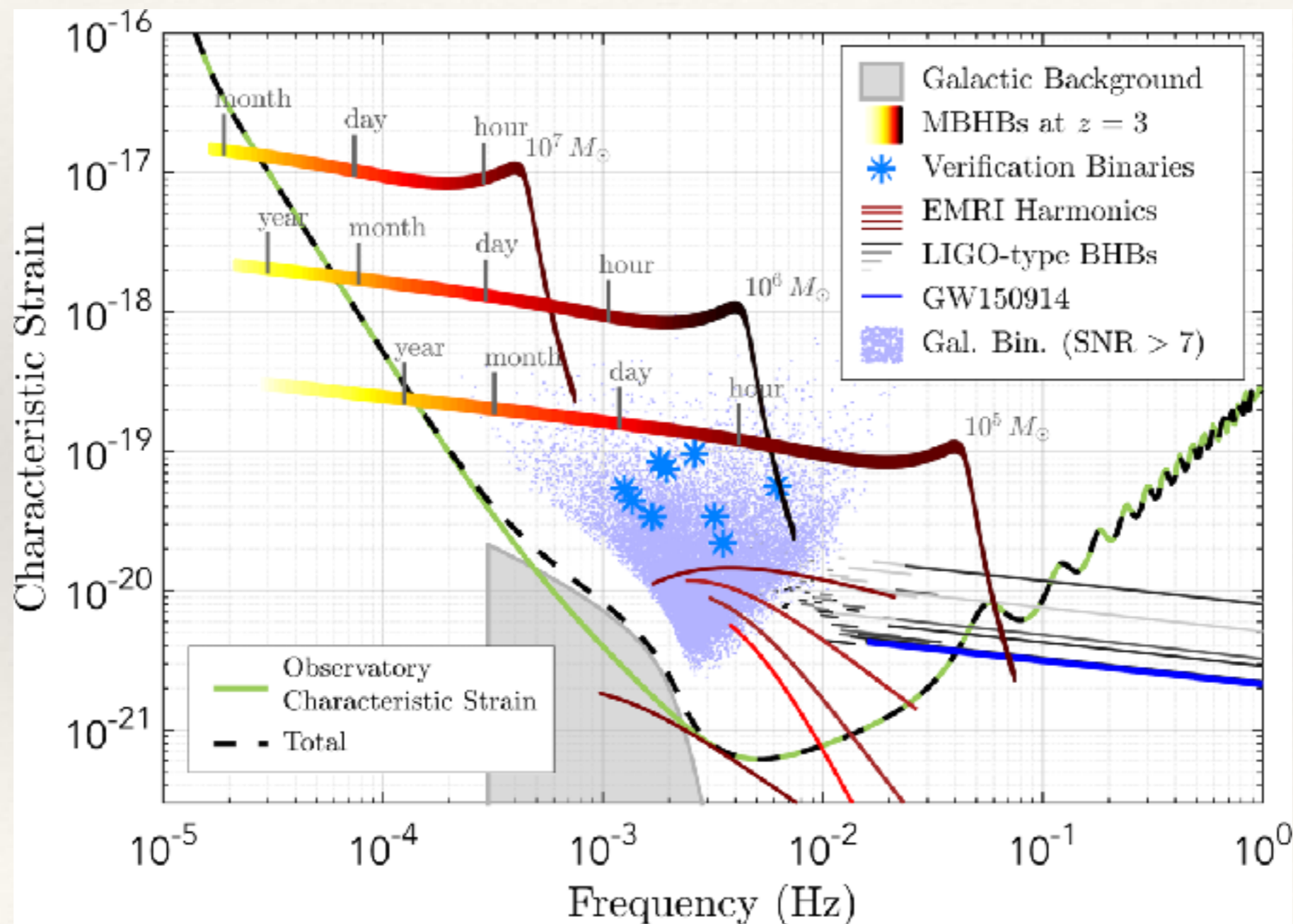
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- MBHB : high uncertainties in the event rate - from few to few hundreds per year
- EMRIs: even more uncertain - from few to few thousands of detectable GW signals per year.
- GW signal from solar mass BBH (LIGO/VIRGO sources). We expect to observe about 10 sources: GW signal first observed in LISA and then 5-10 years later with the ground based detectors.
- Possible detection of the stochastic GW signal from energetic processes in the early Universe.



# LISA data

LISA data analysis is quite a complex task. We organize the LISA data challenge: <https://lisa-ldc.lal.in2p3.fr/home>. We simulate LISA data (noise and GW signals) and anyone can download the data and analyse it.

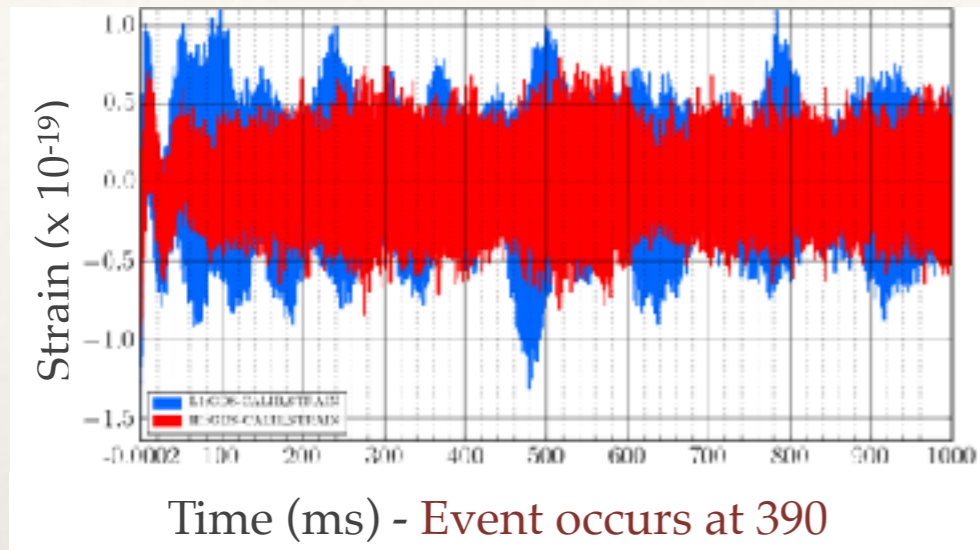




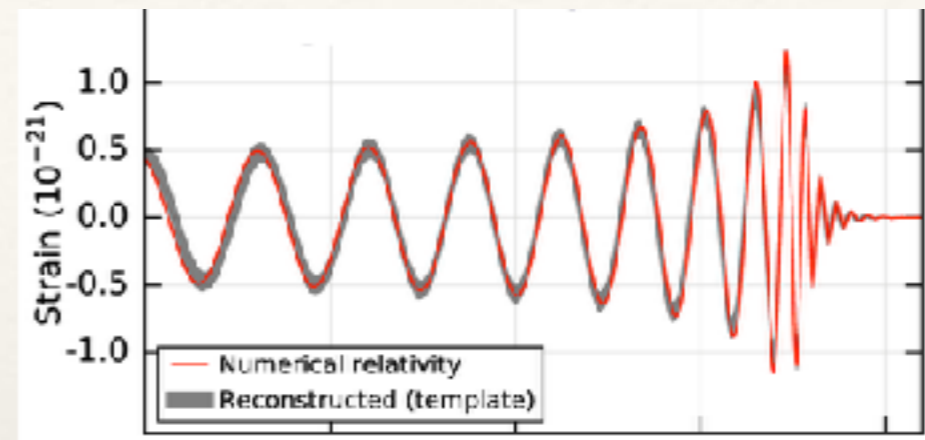
# Data analysis: Matched filtering

GW150914

Raw data



GW signal from merging BHs (we search for)



*Matched filtering*: is used when we are searching for a signal of known form in the noisy data. The basis: we correlate the data with expected signal and search for a maximum of correlation.

wavform / template we search for.

$$\rho \sim 4\Re \int_0^\infty \frac{\tilde{d}(f)\tilde{h}^*(f)}{S(f)} df$$

data

noise power spectral density

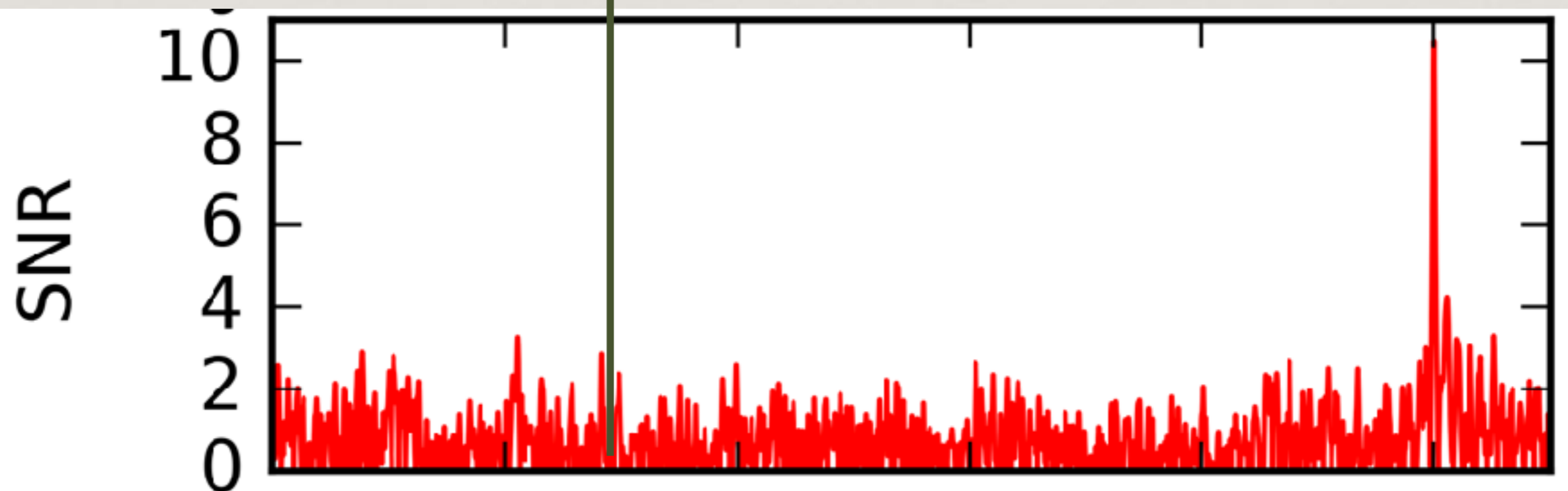
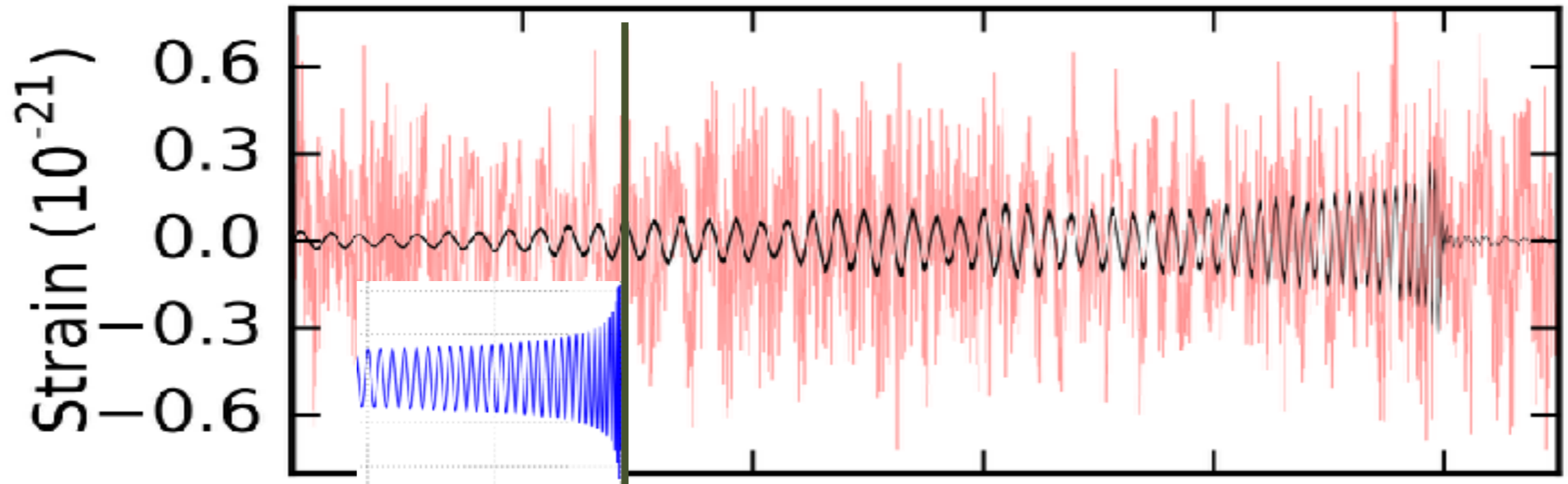
Correlation in frequency domain,  
weighted by detector sensitivity



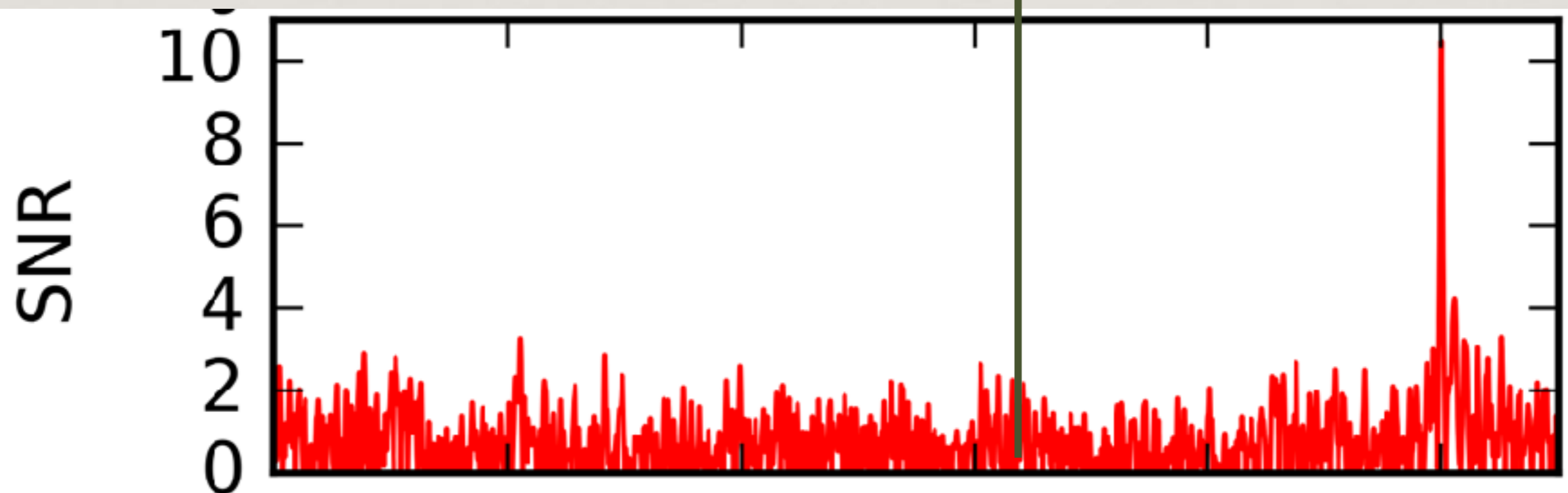
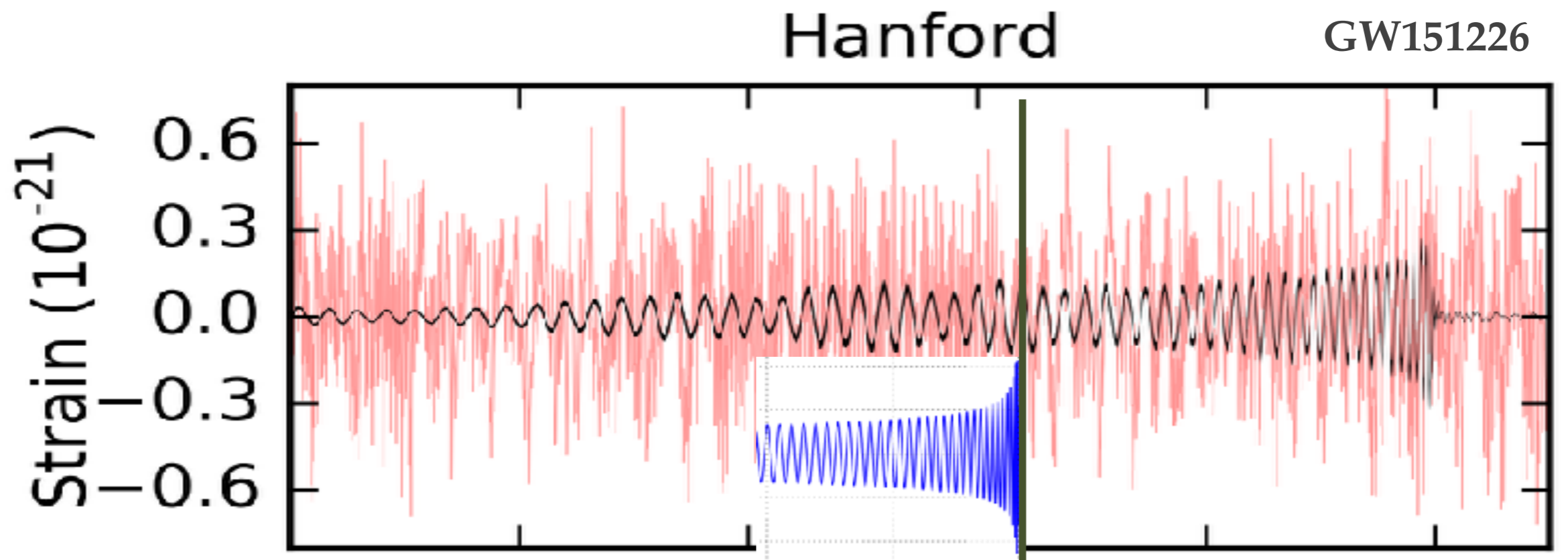
# Matched filtering

Hanford

GW151226

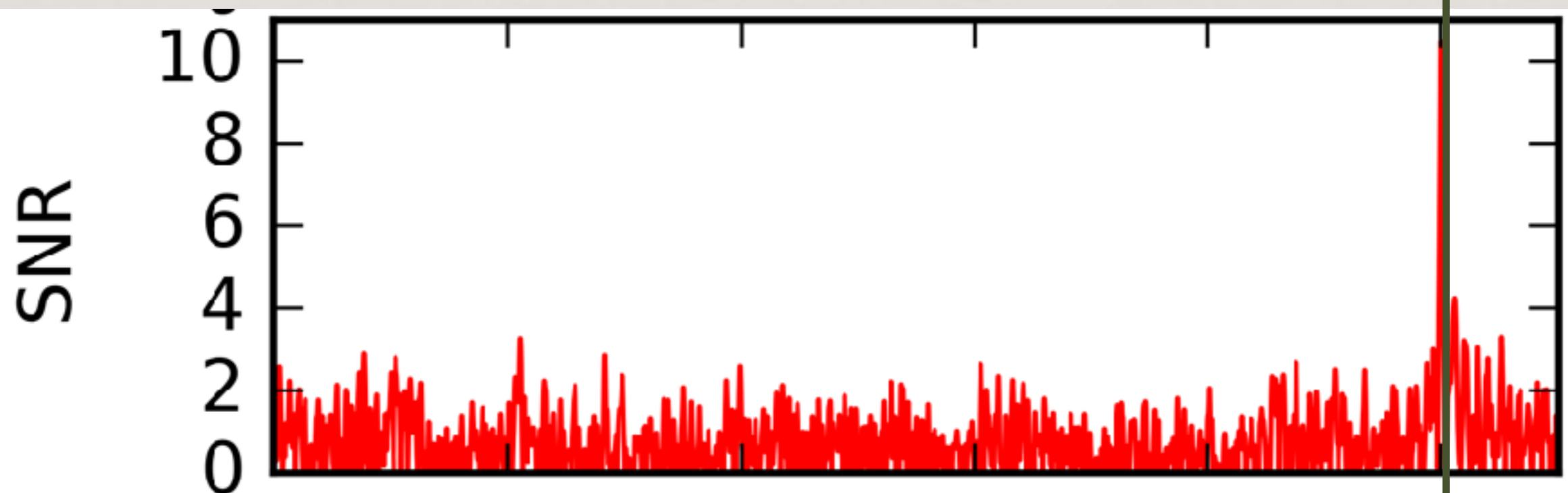
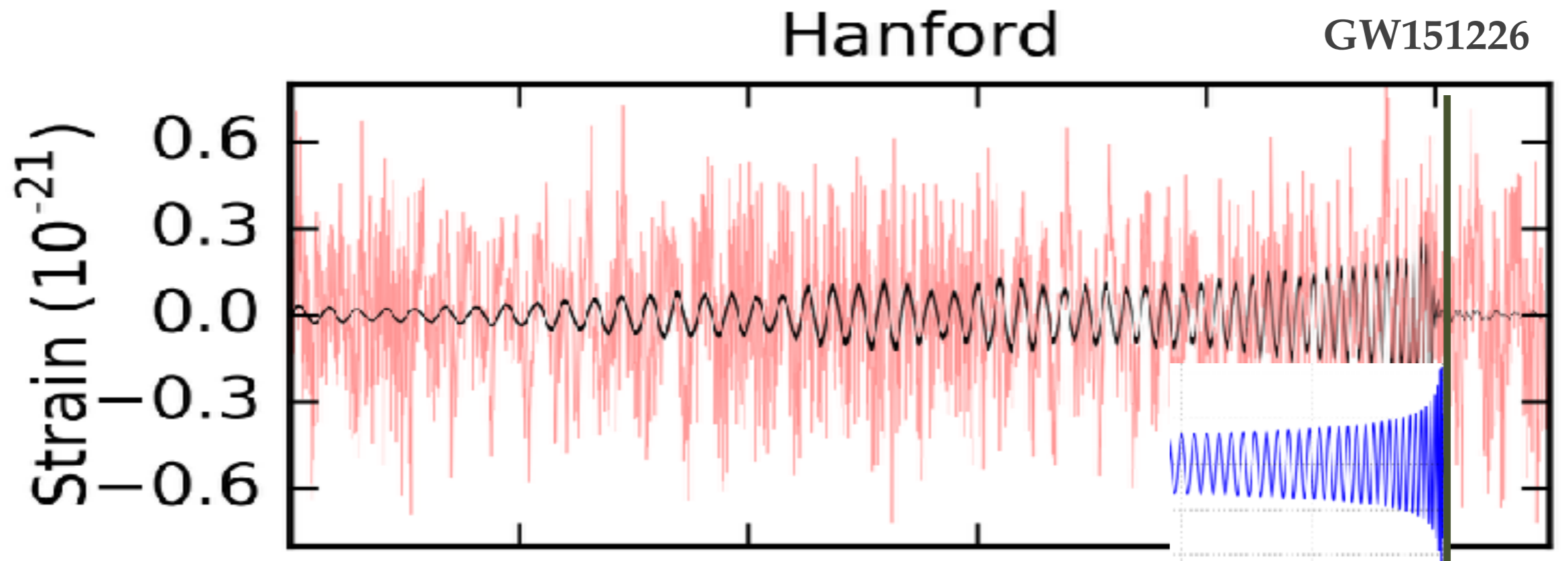


# Matched filtering



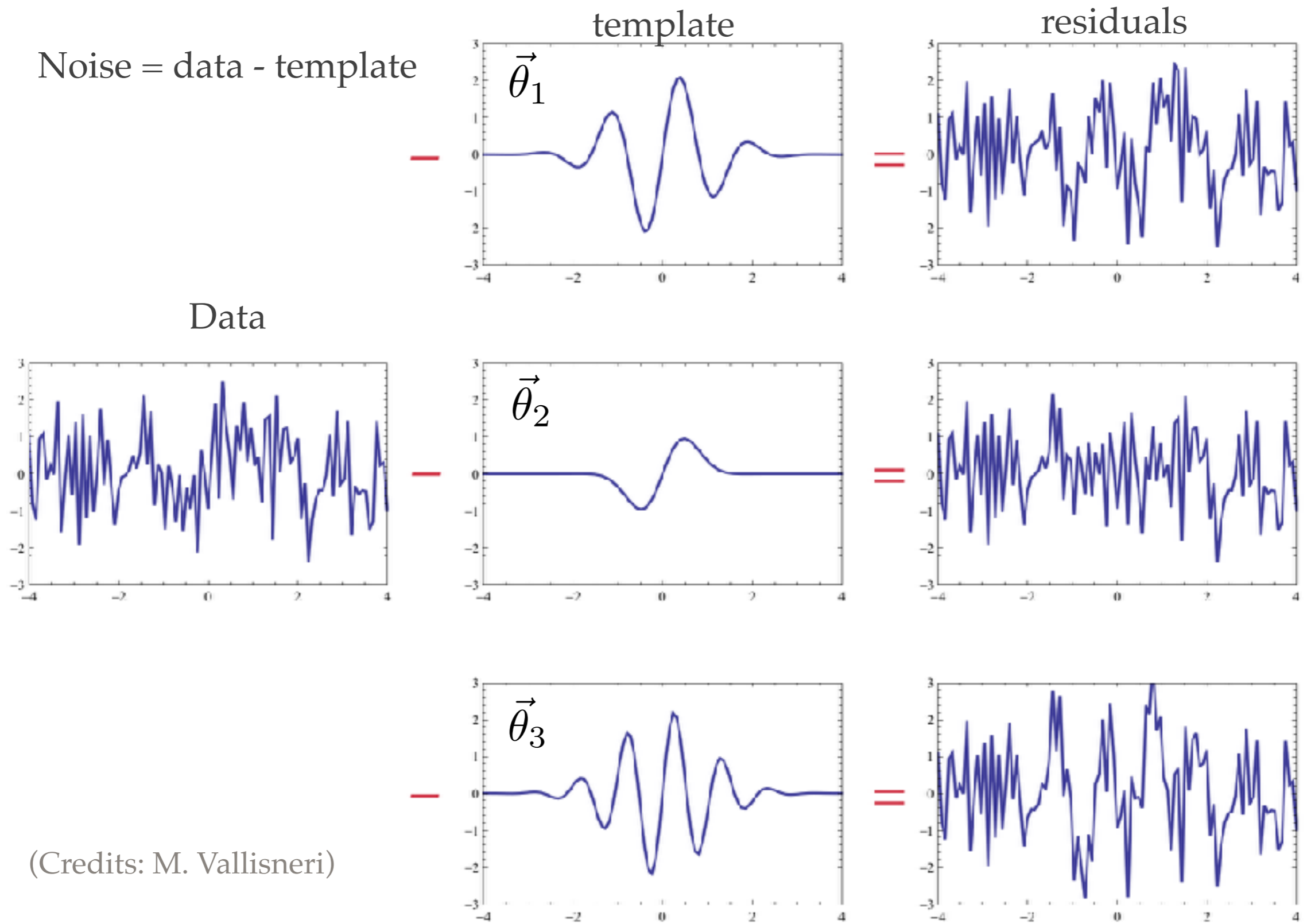


# Matched filtering



# Matched filtering and parameter estimation

Noise = data - template



(Credits: M. Vallisneri)

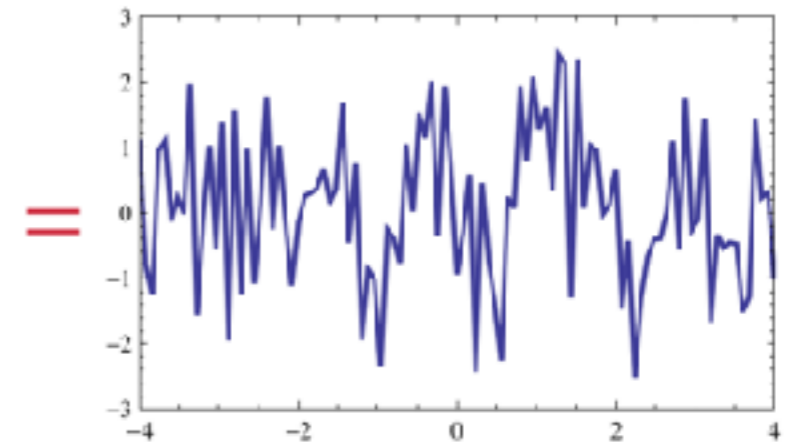
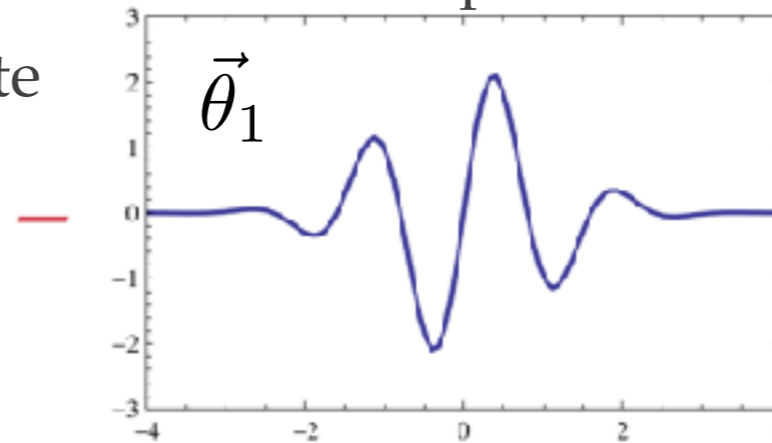


# Matched filtering and parameter estimation

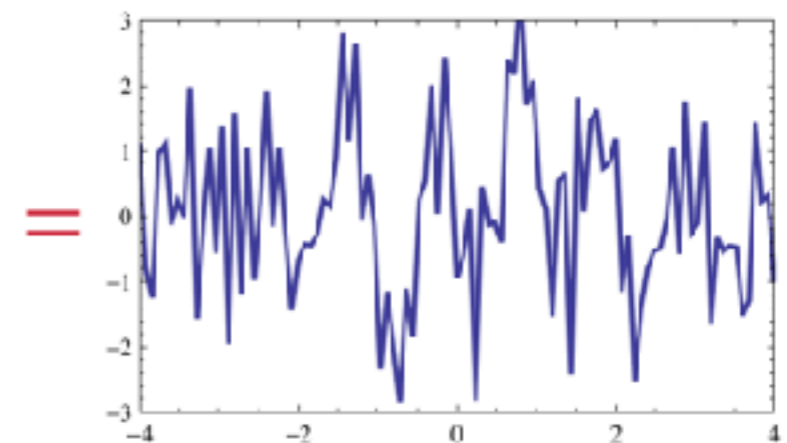
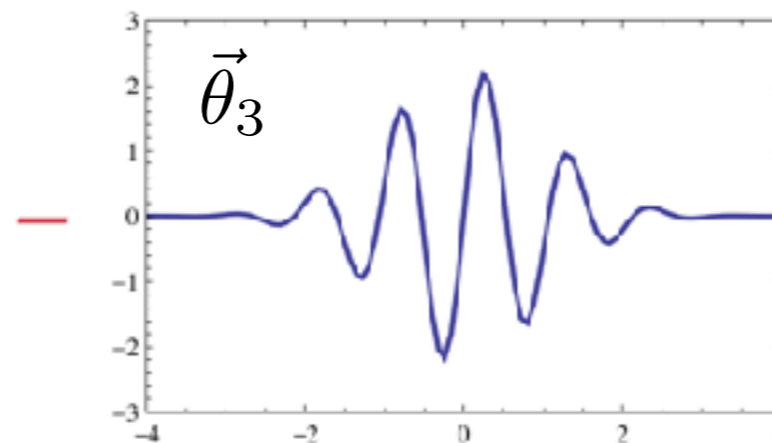
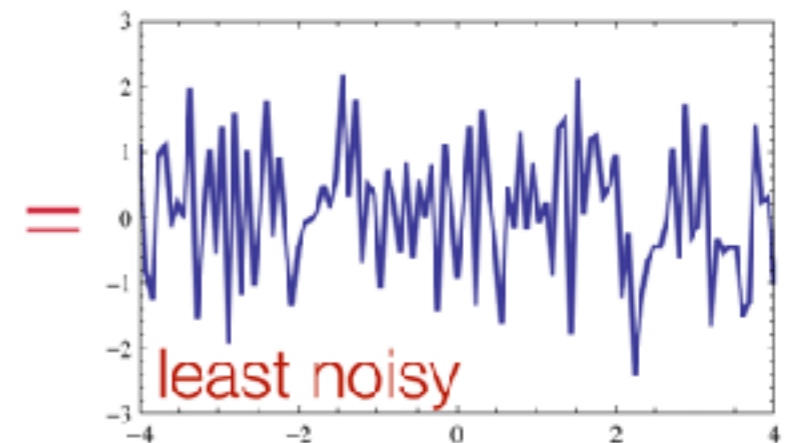
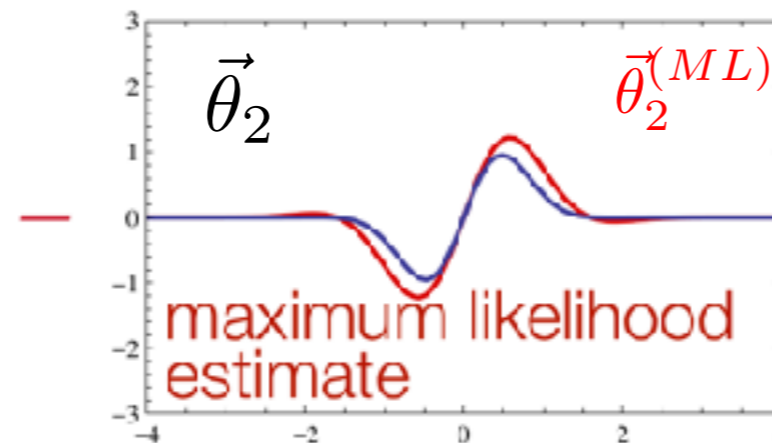
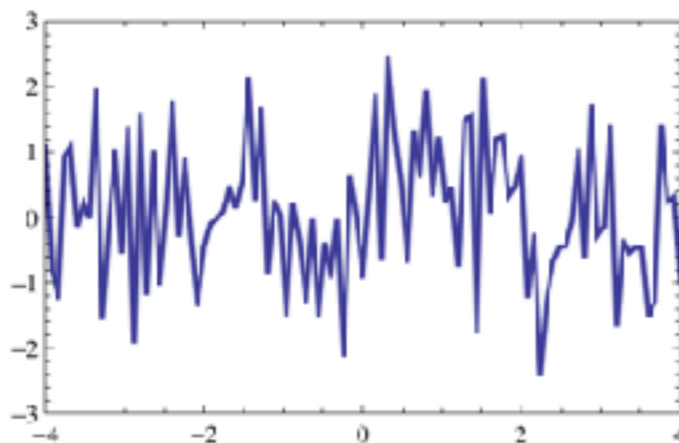
Noise = data - template

template

residuals



Data



(Credits: M. Vallisneri)

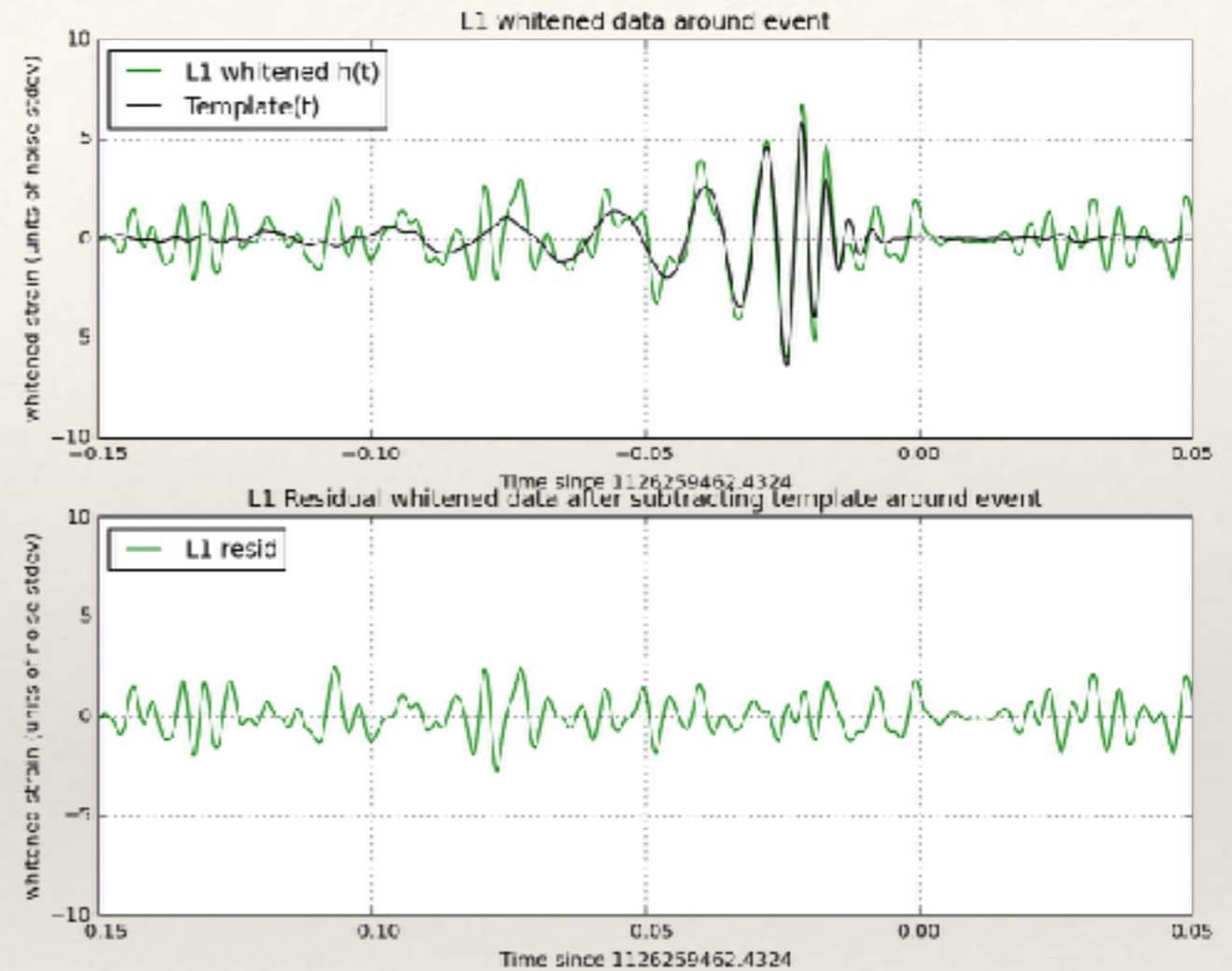
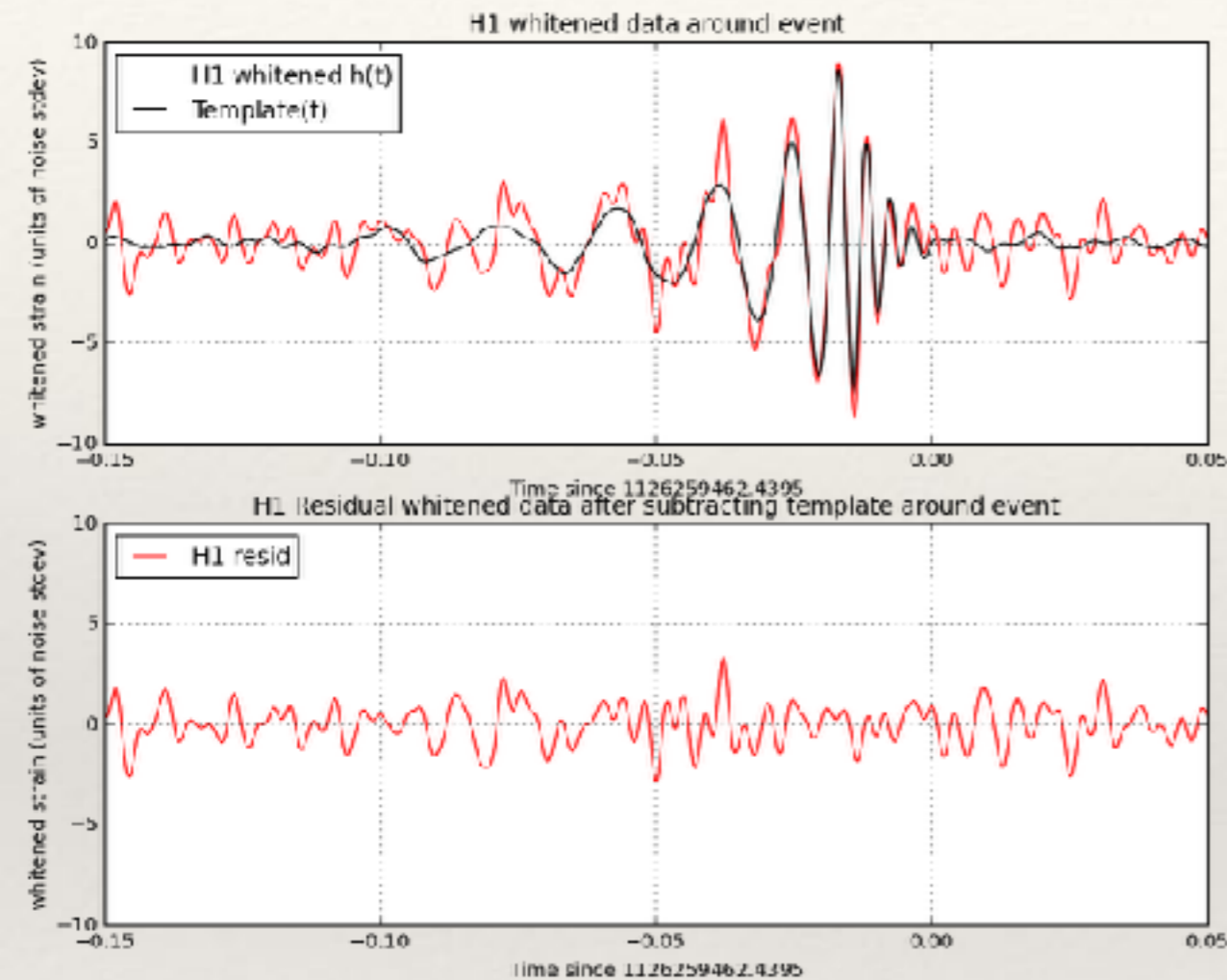




# Matched filtering for GW150914

H1

L1



[LOSC: <https://losc.ligo.org/tutorials/>]



# Likelihood

Let us assume that the data contains the signal.: hypothesis (model)  $H_1$

$$d(t) = n(t) + s(t, \theta_i) \quad \text{signal "s" depends on parameters } \theta_i$$

data = noise + signal

If the template matches the GW signal exactly  $h(t, \theta_i) = s(t, \theta_i) \longrightarrow d(t) - h(t, \theta_i) = n(t)$

$$p(d(t)|H_1, \vec{s}(t, \lambda)) = p(d(t) - s(t, \vec{\lambda})) = p_n$$

- Assume that the noise is Gaussian (but not necessarily white): non white noise has different variance at different frequencies. The the likelihood can be written as

**Likelihood:**  $p(d|H_1, \theta_i) \propto e^{-\frac{1}{2}(d-h(\theta_i)|d-h(\theta_i))}$

The inner product: matched filtering  $(a|b) \equiv 4\Re \int_0^\infty \frac{\tilde{a}(f)\tilde{b}^*(f)}{S_n(f)} df$

- We search for parameters which maximize the likelihood: making the residuals most noise-like — maximum likelihood estimators for parameters  $\hat{\theta}_i$

