# Beyond SiO<sub>2</sub>: New tetrahedral and octahedral structures in IV-VI compounds

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## Outline

- 1. Brief overview of  $SiO_2$  and  $CO_2$
- 2. CS<sub>2</sub> experiment, our theoretical results
- 3. SiS<sub>2</sub> experiment, our theoretical and experimental results
- 4. Conclusions

# IV - VI AB<sub>2</sub> compounds



Difference

CO<sub>2</sub>, CS<sub>2</sub> - molecular
SiO<sub>2</sub>, SiS<sub>2</sub> - non molecular

# Paradigm: SiO<sub>2</sub> (silica)

- important component of Earth's crust (10 %)
- very rich polymorphism
- highly important amorphous phase
- important in many practical applications
- artificially created on Si surface in microelectronics
- most common form is  $\alpha$ -quartz







# SiO<sub>2</sub> polymorphism - tetrahedral 1

Form	Crystal symmetry Pearson symbol, group No.	ρ g/cm <sup>3</sup>	Notes	Structure
α-quartz	rhombohedral (trigonal) hP9, P3 <sub>1</sub> 21 No.152 <sup>[33]</sup>	2.648	Helical chains making individual single crystals optically active; $\alpha$ -quartz converts to $\beta$ -quartz at 846 K	
β-quartz	hexagonal hP18, P6 <sub>2</sub> 22, No. 180 <sup>[34]</sup>	2.533	closely related to $\alpha$ -quartz (with an Si-O-Si angle of 155°) and optically active; $\beta$ -quartz converts to $\beta$ -tridymite at 1140 K	
a-tridymite	orthorhombic oS24, C222 <sub>1</sub> , No.20 <sup>[35]</sup>	2.265	metastable form under normal pressure	
β-tridymite	hexagonal hP12, P6 <sub>3</sub> /mmc, No. 194 <sup>[35]</sup>		closely related to $\alpha$ -tridymite; $\beta$ -tridymite converts to $\beta$ -cristobalite at 2010 K	
a-cristobalite	tetragonal tP12, P4 <sub>1</sub> 2 <sub>1</sub> 2, No. 92 <sup>[36]</sup>	2.334	metastable form under normal pressure	
coesite	monoclinic mS48, C2/c, No.15 <sup>[43]</sup>	2.911	$\rm Si_4O_8$ and $\rm Si_8O_{16}$ rings; 900 K and 3–3.5 GPa	

# SiO<sub>2</sub> polymorphism - tetrahedral 2

m	orthorhombic ol12, Ibam, No.72 <sup>[45]</sup>			like SiS <sub>2</sub> consisting of edge sharing chains, melts at ~1700 K	
	oganite	monoclinic mS46, C2/c, No.15 <sup>[42</sup>	2]	$Si_4O_8$ and $Si_6O_{12}$ rings	
ke	eatite	tetragonal tP36, P4 <sub>1</sub> 2 <sub>1</sub> 2, No. 92 <sup>[</sup>	[41] 3.01	Si <sub>5</sub> O <sub>10</sub> , Si <sub>4</sub> O <sub>14</sub> , Si <sub>8</sub> O <sub>16</sub> rings; synthesised from glassy silica and alkali at 600–900 K and 40–400 MPa	
me	elanophlogite	cubic (cP*, P4 <sub>2</sub> 32, No.208) <sup>[10]</sup> or tetrago (P4 <sub>2</sub> /nbc) <sup>[39]</sup>	nal 2.04	$Si_5O_{10}$ , $Si_6O_{12}$ rings; mineral always found with hydrocarbons in interstitial spaces-a clathrasil <sup>[40]</sup>	
fai	ujasite	cubic cF576, Fd3m, No.227	7 <sup>[38]</sup> 1.92	sodalite cages connected by hexagonal prisms; 12-membered ring pore opening; faujasite structure. <sup>[14]</sup>	
β-	·cristobalite	cubic cF104, Fd3m, No.227	7[37]	closely related to $\alpha$ -cristobalite; melts at 1978 K	

fibrous

W-silica<sup>[4]</sup>

# SiO<sub>2</sub> polymorphism - octahedral

stishovite	Tetragonal tP6, P4 <sub>2</sub> /mnm, No.136 <sup>[44]</sup>	4.287	One of the densest (together with seifertite) polymorphs of silica; rutile-like with 6-fold coordinated Si; 7.5–8.5 GPa	
seifertite	orthorhombic oP, Pbcn <sup>[46]</sup>	4.294	One of the densest (together with stishovite) polymorphs of silica; is produced at pressures above 40 GPa. <sup>[47]</sup>	

Wikipedia

- further increasing pressure in Mbar range silica acquires pyrite-like structure
- no layered structures

#### Most important crystalline phases of silica





Figure 2. High *P-T* phase diagram of SiO<sub>2</sub> determined from static compression experiments. The solid lines are equilibrium phase boundaries consistent with the available data (see text). Data points for the high-pressure melting curves are shown: closed circles, Jackson (1976); open circles, Kanzaki (1990), Zhang et al. (1993); squares, Shen and Lazor (1994); triangles, shock-compression data from Kondo and Ahrens (1983) and Schmitt and Ahrens (1989). The dashed line is the stishovite melting line suggested by Zhang et al. (1993). The dotted lines are metastable extensions of quartz and coesite melting calculated from thermochemical data (Zhang et al., 1993). The stishovite I-II transition (Cohen, 1992; Kingma, 1994), and the (metastable) quartz I-II (Kingma et al., 1993b) and coesite I-II (Hemley, 1987) transitions, all observed 300 K, are shown. The approximate pressures of extensive room-temperature static pressure amorphization are shown; progressive amorphization associated with microfracture begins at lower pressures (horizontal arrow) (Kingma et al., 1993a). The ambient-pressure amorphization of coesite and stishovite is indicated by the vertical arrow (e.g. Richet, 1989). *Inset*: Higher pressure range determined from shock-wave experiments, adapted from Schmitt and Ahrens (1989) and Lyzenga et al. (1983).

from "High-pressure behaviour of silica", Hemley, Prewitt, Kingma (1994)

# Amorphous silica (glass)

material used by mankind since many years
 Syria, Mesopotamia, Egypt (2500 BC)



polyamorphism - LDA, HDA silica

Carbon dioxide

Carbon dioxide



- molecular, at ambient conditions gas
- carbon dioxide counts among most important materials on Earth and in the Solar system
- solid CO<sub>2</sub> has a number of molecular phases I, II, III, IV
- well-known molecular phase is dry ice (phase I)
- upon compression above 20 GPa dry ice transforms to another molecular phase III (*Cmca*)
- at high pressure double bonds in CO<sub>2</sub> molecules are destabilized and polymeric phases with single bonds are created (similar to those found in SiO<sub>2</sub>)

# Dry ice





#### sublimates at -78.5 °C

Wikipedia

### Polymerization of CO<sub>2</sub> at high pressure

- SiO<sub>2</sub> at low pressure forms tetrahedral covalent structures
- CO<sub>2</sub> at low pressure forms molecular crystals
- molecular phases of CO<sub>2</sub> under pressure transform to tetrahedral polymeric ones, similar to those of SiO<sub>2</sub>
- many open questions about structures and transformation paths remain



Phase diagram of CO<sub>2</sub>, Kume et al. (2007)

#### Phase III $\rightarrow \alpha$ -cristobalite-type phase metadynamics, P = 800 kbar, T = 100 K

Jian Sun, Dennis D. Klug, Roman Martoňák, Javier Antonio Montoya, Mal-Soon Lee, Sandro Scandolo and Erio Tosatti, PNAS **106**, 6077–6081 (2009)

#### view along z

#### view along x



## Structure of polymeric CO<sub>2</sub>

PRL 108, 125701 (2012)

PHYSICAL REVIEW LETTERS

week ending 23 MARCH 2012

#### Structure of Polymeric Carbon Dioxide CO<sub>2</sub>-V

 Frédéric Datchi,<sup>1</sup> Bidyut Mallick,<sup>1</sup> Ashkan Salamat,<sup>2</sup> and Sandra Ninet<sup>1</sup>
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 <sup>2</sup>European Radiation Synchrotron Facility, F-38043 Grenoble Cedex, France (Received 2 November 2011; published 19 March 2012)

The structure of polymeric carbon dioxide (CO<sub>2</sub>-V) has been solved using synchrotron x-ray powder diffraction, and its evolution followed from 8 to 65 GPa. We compare the experimental results obtained for a 100% CO<sub>2</sub> sample and a 1 mol % CO<sub>2</sub>/He sample. The latter allows us to produce the polymer in a pure form and study its compressibility under hydrostatic conditions. The high quality of the x-ray data enables us to solve the structure directly from experiments. The latter is isomorphic to the  $\beta$ -cristobalite phase of SiO<sub>2</sub> with the space group  $I\bar{4}2d$ . Carbon and oxygen atoms are arranged in CO<sub>4</sub> tetrahedral units linked by oxygen atoms at the corners. The bulk modulus determined under hydrostatic conditions,  $B_0 = 136(10)$  GPa, is much smaller than previously reported. The comparison of our experimental findings with theoretical calculations performed in the present and previous studies shows that density functional theory very well describes polymeric CO<sub>2</sub>.

#### β-cristobalite-like I-42d



Carbon disulfide



 metastable compound - enthalpy of formation 88.7 kJ/mol

 $CS_2$ 

- at ambient conditions molecular liquid
- below 161 K freezes to Cmca molecular crystal





# CS<sub>2</sub> at high pressure

#### Nobel prize 1946 - Percy Williams Bridgman

- Bridgman (1941) compressed Cmca CS<sub>2</sub> to 4.5 GPa at 175 C
- transformation to black polymer observed
- Whalley (1960) proposed a polymeric structure based on a group





• exact crystal structure is not known

### CS<sub>2</sub> at high pressure

PHYSICAL REVIEW B 84, 144104 (2011)

#### Insulator-metal transition of highly compressed carbon disulfide

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We present integrated spectral, structural, resistance, and theoretical evidences for simple molecular CS<sub>2</sub> transformations to an insulating black polymer with threefold carbon atoms at 9 GPa, then to a semiconducting polymer above 30 GPa, and finally to a metallic solid above 50 GPa. The metallic phase is a highly disordered three-dimensional network structure with fourfold carbon atoms at the carbon-sulfur distance of ~1.70 Å. Based on first-principles calculations, we present two plausible structures for the metallic phase:  $\alpha$ -chalcopyrite and tridymite, both of which exhibit metallic ground states and disordered diffraction features similar to that measured. We also present the phase and chemical transformation diagram for carbon disulfide, showing a large stability field of the metallic phase to 100 GPa and 800 K.



FIG. 1. (Color online) Microphotographs of carbon disulfide under high pressure showing its transformation from (a) transparent fluid to (b) and (c) molecular solid (*Cmca*) at ~1 GPa, to (d), (e), and (g) black polymer above 10 GPa ((-S-(C=S)-)<sub>p</sub> or *CS3*) and eventually to (f) and (h) a highly reflecting extended solid above 48 GPa (*CS4*) at ambient temperature. The rightmost image (h) illustrates the metallic reflectivity of CS<sub>2</sub> samples above 55 GPa similar to those of Platinum (Pt) metal probes in a four-probe configuration for resistance measurements.

#### CS<sub>2</sub> at high pressure



# **Proposed interpretation**

- based on analogy with CO<sub>2</sub>, tridymite and  $\beta$ -cristobalite were proposed as candidates for tetrahedral CS<sub>2</sub>
- comparison to experiment elusive because of disorder
- does the analogy with CO<sub>2</sub> really work well for CS<sub>2</sub>?
- why not considering also analogy with e.g. SiS<sub>2</sub>?

# Our goal

 address the problem with state-of-the-art ab initio crystal structure search techniques

## Crystal structure prediction ab initio genetic algorithms

- we work at T=0 and optimize enthalpy H = E + PV
- we used the USPEX software (Oganov, Glass 2006)
- ab initio calculations and structural relaxations performed by VASP
- plain PBE functional used for structural search
- enthalpies calculated by the optB86b-vdW scheme of Klimeš et al. (2010, 2011) based on the vdW functional of Dion et al. (2004)
- phonon, Raman and IR calculations performed by Quantum Espresso using LDA functional

# Problem with decomposition in CS<sub>2</sub>

intrinsic metastability towards decomposition



Solution - constraint on C-C and S-S bond lengths

# **Results of EA search**

- search performed at p= 0, 26, 38, 75, 120, and 170 GPa
- at p=0 we reproduced the molecular *Cmca* phase
- at higher p we found various tetrahedral structures,  $\alpha$  and  $\beta$  cristobalite
- new tetrahedral layered structure with space group  $P2_1/c$
- also various octahedral structures with high enthalpy



# **Results of EA search**



### Best tetrahedral structure P2<sub>1</sub>/c (HP1)



"shahabite" - layered with edge-sharing octahedra

# 3rd Pauling's rule (1929)

The sharing of edges and particularly faces by two anion polyhedra decreases the stability of an ionic structure.

- comparison: in CO<sub>2</sub> the *P21/c* structure is higher in enthalpy with respect to  $\beta$ -cristobalite by 0.4 eV/molecule
- Bader charge analysis
- partial charges on atoms C (+2), O (-1)
- in CS<sub>2</sub> it is quite different C (-0.55), S (+0.27)
- not only smaller charge, but opposite polarity
- electronegativities: C 2.55, S 2.58
- edge-sharing more plausible in less ionic CS<sub>2</sub>

### **Diffraction patterns**



#### Phonons and spectroscopy



FIG. 5. (Color online) Calculated phonon spectrum of the  $P2_1/c$  layered structure of CS<sub>2</sub> at 50 GPa.



FIG. 6. (Color online) Predicted Raman spectra of two competing CS<sub>2</sub> structures at 50 GPa, compared with measurements at 50 GPa, 297 K.<sup>7</sup> The high frequency secondary peak near 800 cm<sup>-1</sup> is only present in layered  $P_{21}/c$  and absent in  $\beta$ -cristobalite. Also the low frequency spectrum is better reproduced by HP1 than by  $\beta$ -cristobalite.





-FIG, 7. (Color online) Calculated IR spectra of different CS<sub>2</sub> structureSm@tote the stiffer frequencies of layered HP1 compared with β-cristobalite. Blue arrows indicate the experimental IR peak positions of molecular CS<sub>2</sub>.<sup>28,29</sup>

## Metallization



- in PBE metallization at 30 GPa
- in B3LYP at 50 GPa
- closes indirect gap
- unlike tetrahedral CO<sub>2</sub>, CS<sub>2</sub> metalizes very easily with pressure



#### PHYSICAL REVIEW B 91, 224108 (2015)

#### **High-pressure layered structure of carbon disulfide**

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Solid CS<sub>2</sub> is superficially similar to CO<sub>2</sub>, with the same *Cmca* molecular crystal structure at low pressures, which has suggested similar phases also at high pressures. We carried out an extensive first-principles evolutionary search in order to identify the zero-temperature lowest-enthalpy structures of CS<sub>2</sub> for increasing pressure up to 200 GPa. Surprisingly, the molecular *Cmca* phase does not evolve into  $\beta$ -cristobalite as in CO<sub>2</sub> but transforms instead into phases HP2 and HP1, both recently described in high-pressure SiS<sub>2</sub>. HP1 in particular, with a wide stability range, is a layered  $P2_1/c$  structure characterized by pairs of edge-sharing tetrahedra and is theoretically more robust than all other CS<sub>2</sub> phases discussed so far. Its predicted Raman spectrum and pair correlation function agree with experiment better than those of  $\beta$ -cristobalite, and further differences are predicted between their respective IR spectra. The band gap of HP1-CS<sub>2</sub> is calculated to close under pressure, yielding an insulator-metal transition near 50 GPa, in agreement with experimental observations. However, the metallic density of states remains modest above this pressure, suggesting a different origin for the reported superconductivity.

# SiS<sub>2</sub> at ambient pressure



#### 1D chains of edge-sharing tetrahedra

# SiS<sub>2</sub> at high pressure

#### **Inorganic Chemistry**

Article

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#### Two High-Pressure Phases of SiS<sub>2</sub> as Missing Links between the Extremes of Only Edge-Sharing and Only Corner-Sharing Tetrahedra

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Supporting Information

**ABSTRACT:** The ambient pressure phase of silicon disulfide (NP-SiS<sub>2</sub>), published in 1935, is orthorhombic and contains chains of distorted, edge-sharing SiS<sub>4</sub> tetrahedra. The first high pressure phase, HP3-SiS<sub>2</sub>, published in 1965 and quenchable to ambient conditions, is tetragonal and contains distorted corner-sharing SiS<sub>4</sub> tetrahedra. Here, we report on the crystal structures of two monoclinic phases, HP1-SiS<sub>2</sub> and HP2-SiS<sub>2</sub>, which can be considered as missing links between the orthorhombic and the tetragonal phase. Both monoclinic phases contain edge- as well as corner-sharing SiS<sub>4</sub> tetrahedra. With increasing pressure, the volume contraction  $(-\Delta V/V)$ 



and the density, compared to the orthorhombic NP-phase, increase from only edge-sharing tetrahedra to only corner-sharing tetrahedra. The lattice and the positional parameters of NP-SiS<sub>2</sub>, HP1-SiS<sub>2</sub>, HP2-SiS<sub>2</sub>, and HP3-SiS<sub>2</sub> were derived in good agreement with the experimental data from group–subgroup relationships with the CaF<sub>2</sub> structure as aristotype. In addition, the Raman spectra of SiS<sub>2</sub> show that the most intense bands of the new phases HP1-SiS<sub>2</sub> and HP2-SiS<sub>2</sub> (408 and 404 cm<sup>-1</sup>, respectively) lie between those of NP-SiS<sub>2</sub> (434 cm<sup>-1</sup>) and HP3-SiS<sub>2</sub> (324 cm<sup>-1</sup>). Density functional theory (DFT) calculations confirm these observations.

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# SiS<sub>2</sub> at high pressure (Evers et al.)



surprise: HP1 is identical to shahabite!

# SiS<sub>2</sub> at high pressure (Evers et al.)



What happens in SiS<sub>2</sub> at even higher pressure? A denser tetrahedral phase or an octahedral one?

#### Evolutionary structure search for SiS<sub>2</sub>

- Xtalopt code and VASP used
- 6 and 12 atoms in unit cell
- p = 10, 30, 60, 100 GPa
- more than 1000 structures at each pressure
- 3 low-enthalpy layered octahedral polytypes were found, *P-3m1*, *P6<sub>3</sub>mc*, *R-3m*
- the most stable one is *P-3m1*, Cdl2-type structure, one layer per unit cell
- becomes thermodynamically stable at 4 GPa, densification by 19 %
- the same structure is found at p=0 in SnS<sub>2</sub>, SiTe<sub>2</sub>

#### Layered octahedral structures of SiS<sub>2</sub>













### Thermodynamical stability of phases



- vdW calculated within the Tkatchenko-Scheffler scheme
- HP3 appears metastable, could be a kinetic effect
- at p=0 NP and HP1 appear nearly degenerate

# Band structure of the P-3m1 phase



indirect gap semiconductor, metallizes at 40 GPa

### Pressure evolution of e-DOS at E<sub>F</sub>



remains poor metal even at high pressure

#### Octahedral monolayer SiS<sub>2</sub> band structure



1 eV indirect band gap semiconductor

# SCIENTIFIC REPORTS

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#### **Creating new layered structures at** high pressures: SiS<sub>2</sub>

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Received: 27 September 2016 Accepted: 03 November 2016 Published: 25 November 2016 Old and novel layered structures are attracting increasing attention for their physical, electronic, and frictional properties. SiS<sub>2</sub>, isoelectronic to SiO<sub>2</sub>, CO<sub>2</sub> and CS<sub>2</sub>, is a material whose phases known experimentally up to 6GPa exhibit 1D chain-like, 2D layered and 3D tetrahedral structures. We present highly predictive ab initio calculations combined with evolutionary structure search and molecular dynamics simulations of the structural and electronic evolution of SiS<sub>2</sub> up to 100 GPa. A highly stable CdI2-type layered structure, which is octahedrally coordinated with space group  $P_{\overline{3}}m_1$  surprisingly appears between 4 and up to at least 100 GPa. The tetrahedral-octahedral switch is naturally expected upon compression, unlike the layered character realized here by edge-sharing SiS<sub>6</sub> octahedral units connecting within but not among sheets. The predicted phase is semiconducting with an indirect band gap of about 2 eV at 10 GPa, decreasing under pressure until metallization around 40 GPa. The robustness of the layered phase suggests possible recovery at ambient pressure, where calculated phonon spectra indicate dynamical stability. Even a single monolayer is found to be dynamically stable in isolation, suggesting that it could possibly be sheared or exfoliated from bulk P3m1-SiS<sub>2</sub>.

#### Experiment: octahedral SiS<sub>2</sub>

#### **Inorganic Chemistry**



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### More Than 50 Years after Its Discovery in SiO<sub>2</sub> Octahedral Coordination Has Also Been Established in SiS<sub>2</sub> at High Pressure

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**S** Supporting Information

**ABSTRACT:** SiO<sub>2</sub> exhibits a high-pressure—high-temperature polymorphism, leading to an increase in silicon coordination number and density. However, for the related compound SiS<sub>2</sub> such pressure-induced behavior has not been observed with tetrahedral coordination yet. All four crystal structures of SiS<sub>2</sub> known so far contain silicon with tetrahedral coordination. In the orthorhombic, ambient-pressure phase these tetrahedra share edges and achieve only low space filling and density. Up to 4 GPa and 1473 K, three phases can be quenched as metastable phases from high-pressure high-temperature to ambient conditions. Space occupancy and density are increased first by edge and corner sharing and then by corner sharing alone. The structural situation of SiS<sub>2</sub> up to the current study resembles that of SiO<sub>2</sub> in 1960: Then, in its polymorphs only Si–O<sub>4</sub> tetrahedra were known. But in 1961, a polymorph with rutile structure was discovered: octahedral Si-O<sub>6</sub> coordination was established. Now, 50 years later, we report here on the transition from 4-fold to 6-fold coordination in SiS<sub>2</sub>, the sulfur analogue of silica.



HP4-SiS2 12.7 GPa Si – S: 2.29(1) Å Si – Si: 3.22(1) Å S – S: 3.22(1),3.24(1) Å

#### **XRD** pattern and Raman spectrum



theoretical prediction of the *P-3m1* octahedral phase fully confirmed, Evers et al. (2016)

#### Direct synthesis of SiS2 at high pressure

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#### CrossMark ¢ click for updates

#### Synthesis and Raman spectroscopy of a layered SiS<sub>2</sub> phase at high pressures

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Dichalcogenides are known to exhibit layered solid phases, at ambient and high pressures, where 2D layers of chemically bonded formula units are held together by van der Waals forces. These materials are of great interest for solid-state sciences and technology, along with other 2D systems such as graphene and phosphorene. SiS<sub>2</sub> is an archetypal model system of the most fundamental interest within this ensemble. Recently, high pressure (GPa) phases with Si in octahedral coordination by S have been theoretically predicted and also experimentally found to occur in this compound. At variance with stishovite in SiO<sub>2</sub>, which is a 3D network of SiO<sub>6</sub> octahedra, the phases with octahedral coordination in SiS<sub>2</sub> are 2D layered. Very importantly, this type of semiconducting material was theoretically predicted to exhibit continuous bandgap closing with pressure to a poor metallic state at tens of GPa. We synthesized layered SiS<sub>2</sub> with octahedral coordination in a diamond anvil cell at 7.5-9 GPa, by laser heating together elemental S and Si at 1300-1700 K. Indeed, Raman spectroscopy up to 64.4 GPa is compatible with continuous bandgap closing in this material with the onset of either weak metallicity or of a narrow bandgap semiconductor state with a large density of defect-induced, intra-gap energy levels, at about 57 GPa. Importantly, our investigation adds up to the fundamental knowledge of layered dichalcogenides. *Published by AIP Publishing*. https://doi.org/10.1063/1.5011333



FIG. 1. Raman spectra of the Si/S mixture under pressure, before (elemental Si and S) and after (layered, l, octahedral, o SiS<sub>2</sub>) laser heating. Inset: sample configuration. Blue and grey: diamond anvils and gasket. Yellow and black slabs: S and Si, respectively. Horizontal arrows at the top and at the bottom point to green and near IR laser beams used for Raman spectroscopy and for sample heating (double side), respectively.

#### **Experiment Raman**



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metallization at 57 GPa detected by asymmetry of Raman lines (Fano resonance)

#### Conclusions

- new prediction CS<sub>2</sub> creates at 30 GPa a layered tetrahedral P2<sub>1</sub>/c structure - better candidate for the phase observed by Dias et al.
- the P21/c phase of CS2 is identical to the HP1 phase of SiS2 at 3.5 GPa
- CS<sub>2</sub> is at low pressure similar to CO<sub>2</sub> and at higher pressure to SiS<sub>2</sub>
- challenge to prepare it in more crystalline form low temperature compression?
- new prediction SiS<sub>2</sub> creates at 4 GPa a layered octahedral phase of CdI<sub>2</sub> type (confirmed experimentally)
- $SiS_2$  is similar both to  $CS_2$  and to  $CO_2$
- undergoes an interesting evolution in dimensionality from 1D to 2D to 3D to 2D
- beyond SiO<sub>2</sub> many layered structures are found

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