



**NANO  
LUND**  
AT THE FOREFRONT  
OF NANOSCIENCE

# Realization of an efficient quantum-dot heat engine

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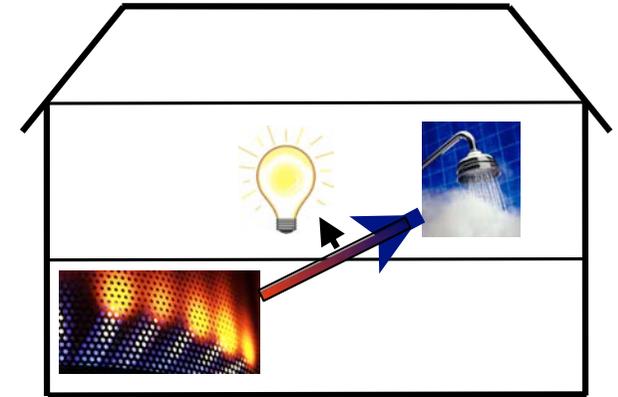
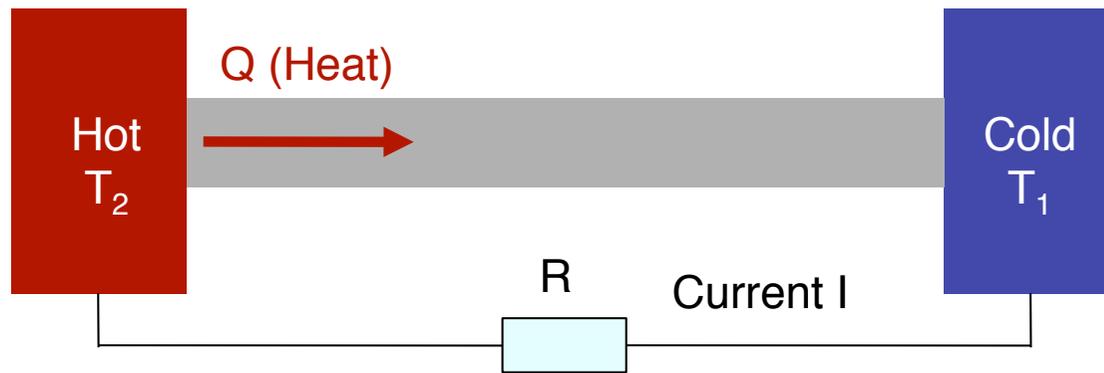
**HEINER LINKE**

NanoLund and Solid State Physics, Lund University, Sweden

Energy Conversion in the Quantum Regime, 27 August 2019, ICTP, Trieste



# Thermoelectrics



- Low parasitic heat conduction by electrons ( $\kappa_{el}$ ) and phonons ( $\kappa_{ph}$ ).
- High Seebeck coefficient  $S = \Delta V / \Delta T$
- Little Joule heating (high conductivity  $\sigma$ )

Figure of merit:

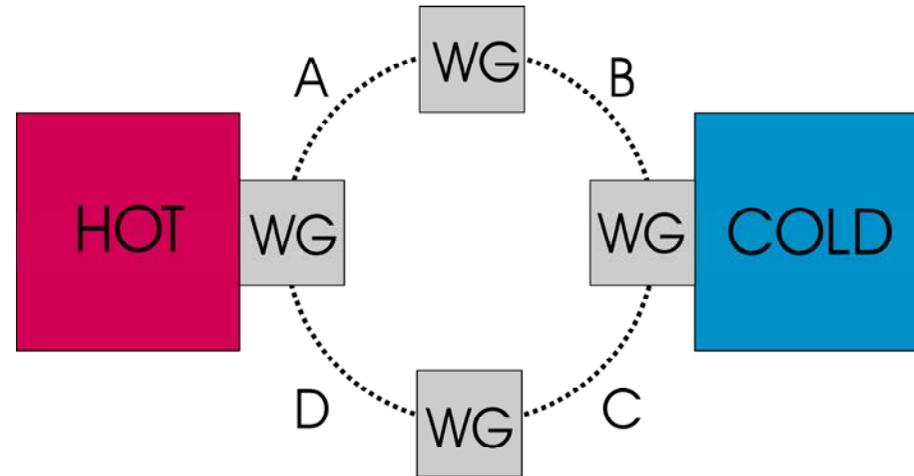
$$Z = \frac{S^2 \sigma}{K_e + K_{ph}} \text{Power factor}$$

# Fundamental efficiency limit of thermoelectrics?

## Classic, cyclic Carnot engine:

Working gas (WG) in contact with only one heat reservoir at a time.

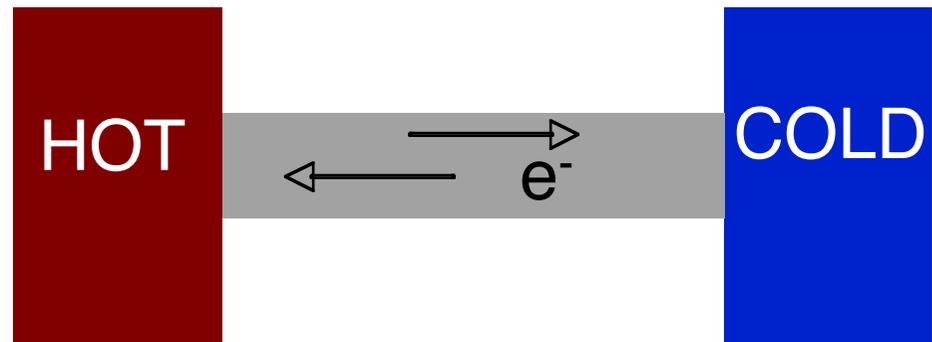
$$\eta_c = 1 - \frac{T_C}{T_H}$$



## Thermoelectric:

In contact with both reservoirs at all times.

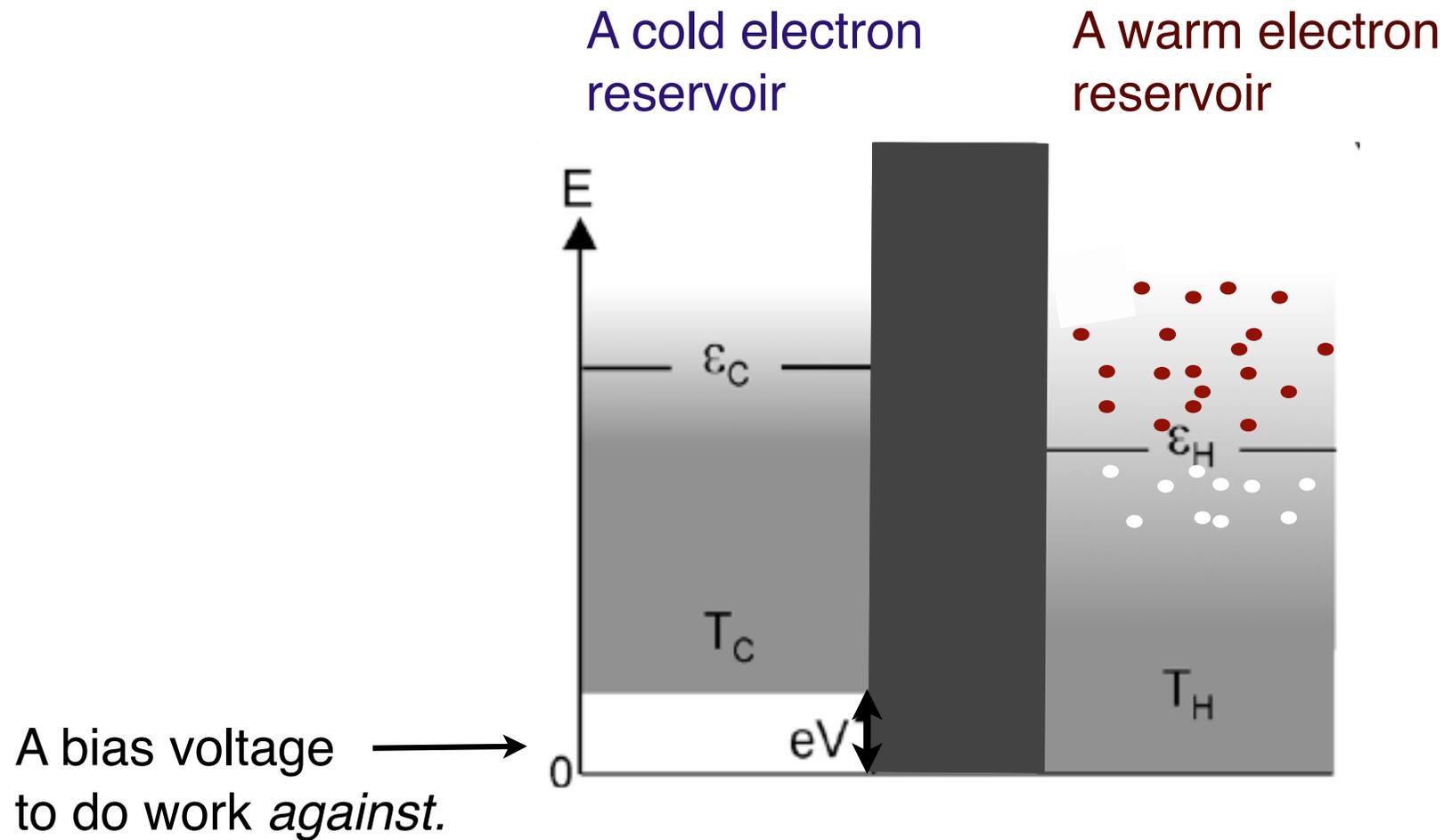
$$Z = \frac{S^2 \sigma}{K_e + K_{ph}}$$



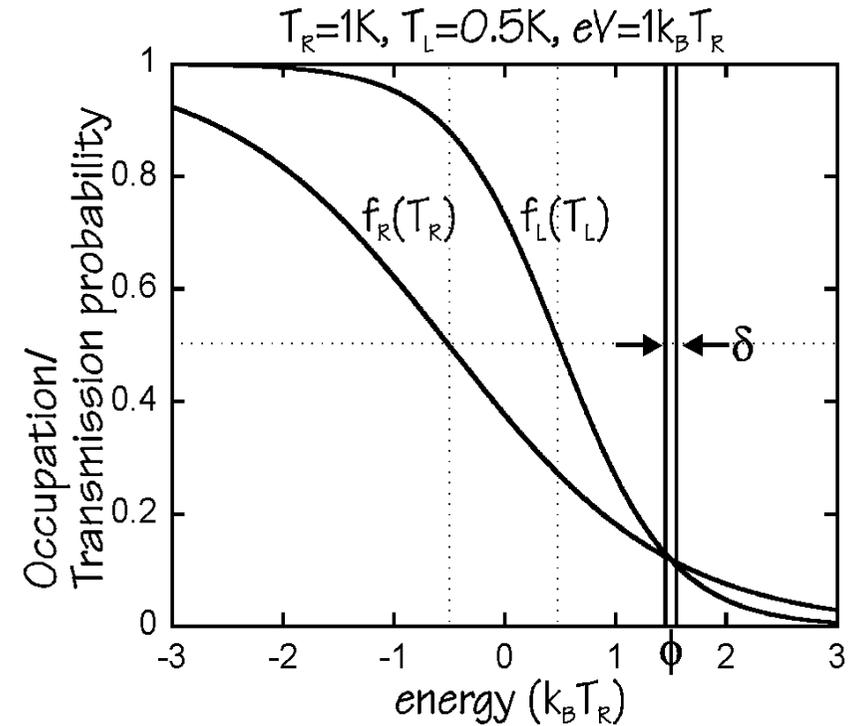
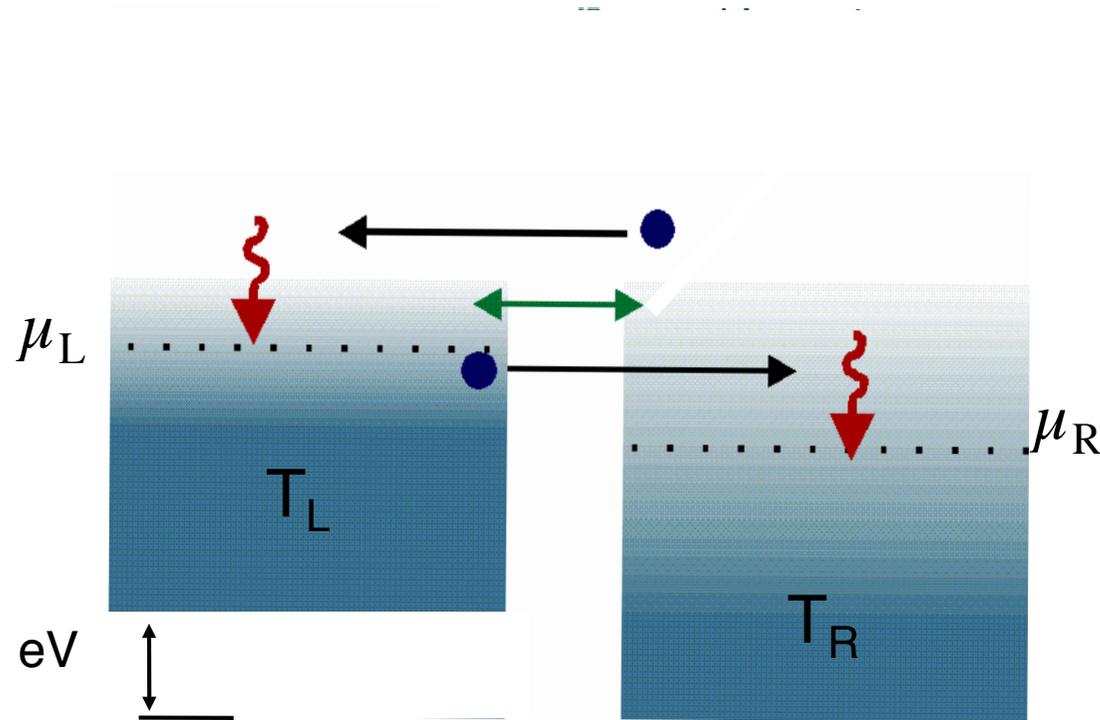
# Outline

- Energy-filtering and energy-specific equilibrium in thermoelectrics
- Realizing “the best thermoelectric”: quantum-dot heat engines
- Experiments: QD heat engine with  $> 70\%$  of Carnot efficiency at finite power
- Single-molecule “quantum dots”
- Application of energy filtering to hot-carrier solar cells

# Fundamental elements of thermoelectrics



# Reversible electron transfer



Transfer of one electron  
at energy  $\varepsilon$  from L to R:

$$\Delta S = \frac{-(\varepsilon - \mu_L)}{T_L} + \frac{(\varepsilon - \mu_R)}{T_R}$$

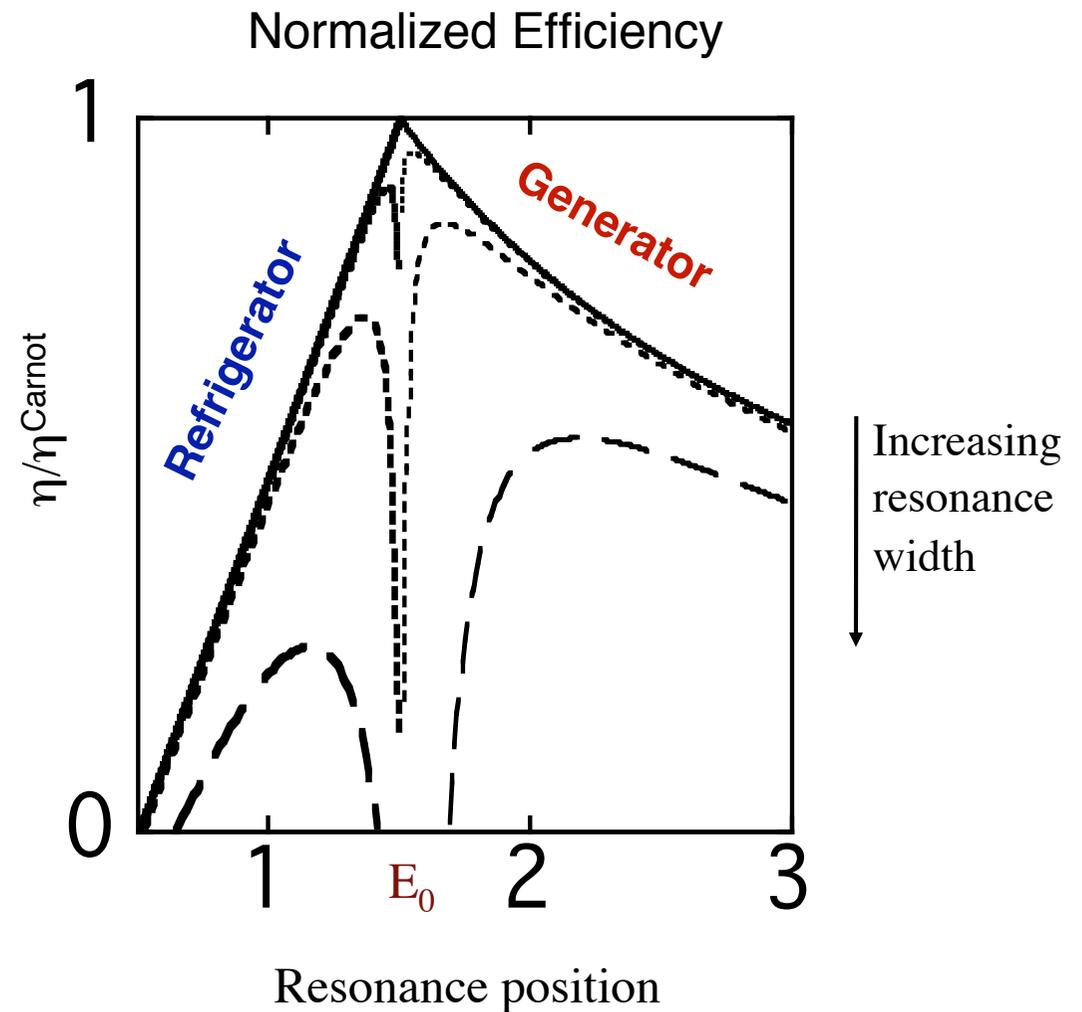
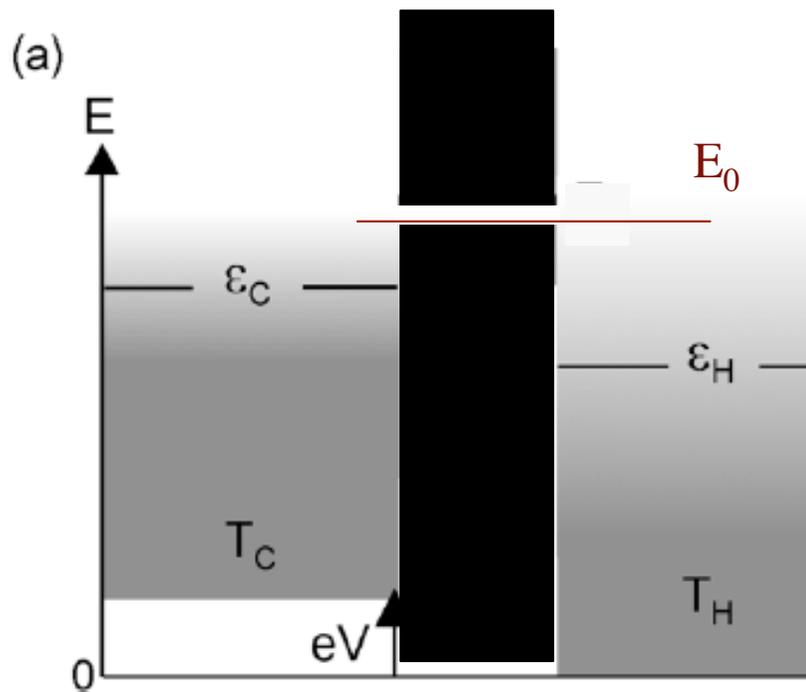
$$\Delta S = 0 \quad \text{for} \quad \varepsilon = \left( \frac{\mu_L T_R + \mu_R T_L}{T_R - T_L} \right)^{\ddagger}$$

“Energy-specific equilibrium”

T. E. Humphrey and H. Linke, PRL **89**, 116801 (2002)

T. E. Humphrey, H Linke, PRL **94**, 096601 (2005)

# Power generation or Refrigeration with tuneable efficiency and power



# The “best thermoelectric”



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## The best thermoelectric

*Proc. Natl. Acad. Sci. USA*  
Vol. 93, pp. 7436–7439, July 1996  
Applied Physical Sciences

G. D. MAHAN\*† AND J. O. SOFO‡

\*Department of Physics and Astronomy, The University of Tennessee, Knoxville, TN 37996-1200; †Solid State Division, Oak Ridge National Laboratory, P.O. Box 2008, Oak Ridge, TN 37831-6030; and ‡Instituto Balseiro, Centro Atomico Bariloche, (8400) Bariloche, Argentina

*Contributed by G. D. Mahan, May 20, 1996*

**ABSTRACT** What electronic structure provides the largest figure of merit for thermoelectric materials? To answer that question, we write the electrical conductivity, thermopower, and thermal conductivity as integrals of a single function, the transport distribution. Then we derive the mathematical function for the transport distribution, which gives the largest figure of merit. A delta-shaped transport distribution is found to maximize the thermoelectric properties. This result indicates that a narrow distribution of the energy of the electrons participating in the transport process is needed for maximum thermoelectric efficiency. Some possible realizations of this idea are discussed.

$$Z = \frac{S^2 \sigma}{K_e + K_{ph}}$$

Figure of merit

# Efficiency at maximum power: Curzon-Ahlborn efficiency



$$\eta_c = 1 - T_c / T_h$$

Carnot efficiency requires reversible operation, which is equivalent to zero power output.

$$\eta_{CA} = 1 - \sqrt{T_c / T_h}$$

Curzon-Ahlborn efficiency describes the efficiency of an ideal Carnot engine operated at maximum power (neglecting dissipation in reservoirs)

$$\eta_{CA} = 1 - \sqrt{T_c / T_h} = \eta_c / 2 + \eta_c^2 / 8 + \dots$$

II Novikov, *J Nuclear Energy* **7**, 125 (1954).

F. Curzon and B. Ahlborn, *Am. J. Phys.* **43**, 22 1975.

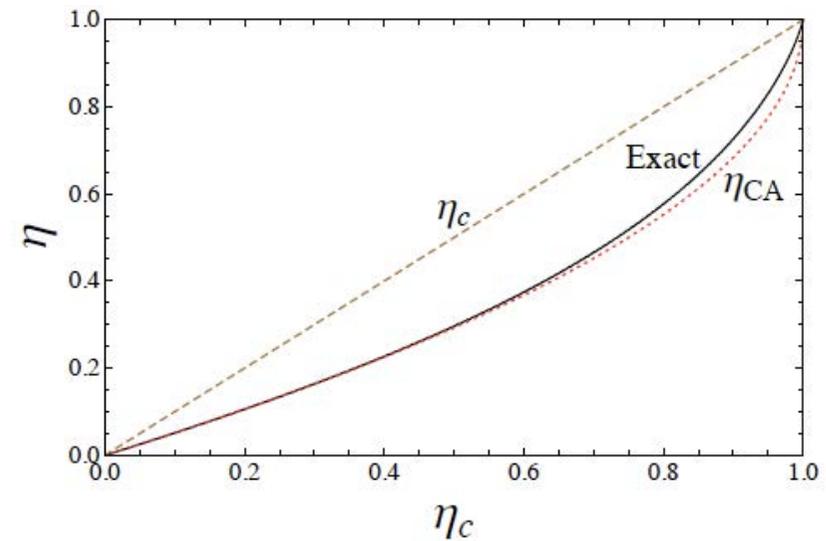
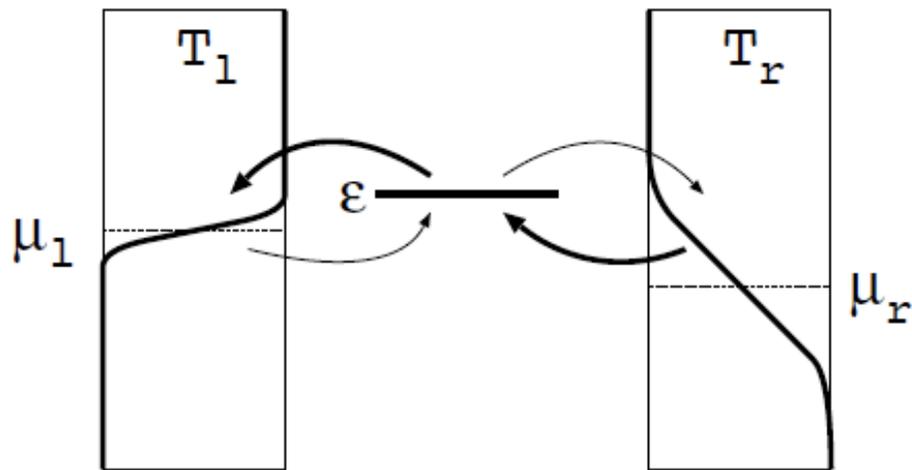
# Thermoelectric efficiency at maximum power in a quantum dot

Massimiliano Esposito\* and Katja Lindenberg

*Department of Chemistry and Biochemistry and Institute for Nonlinear Science,  
University of California, San Diego, La Jolla, CA 92093-0340, USA*

Christian Van den Broeck

*Hasselt University, B-3590 Diepenbeek, Belgium*

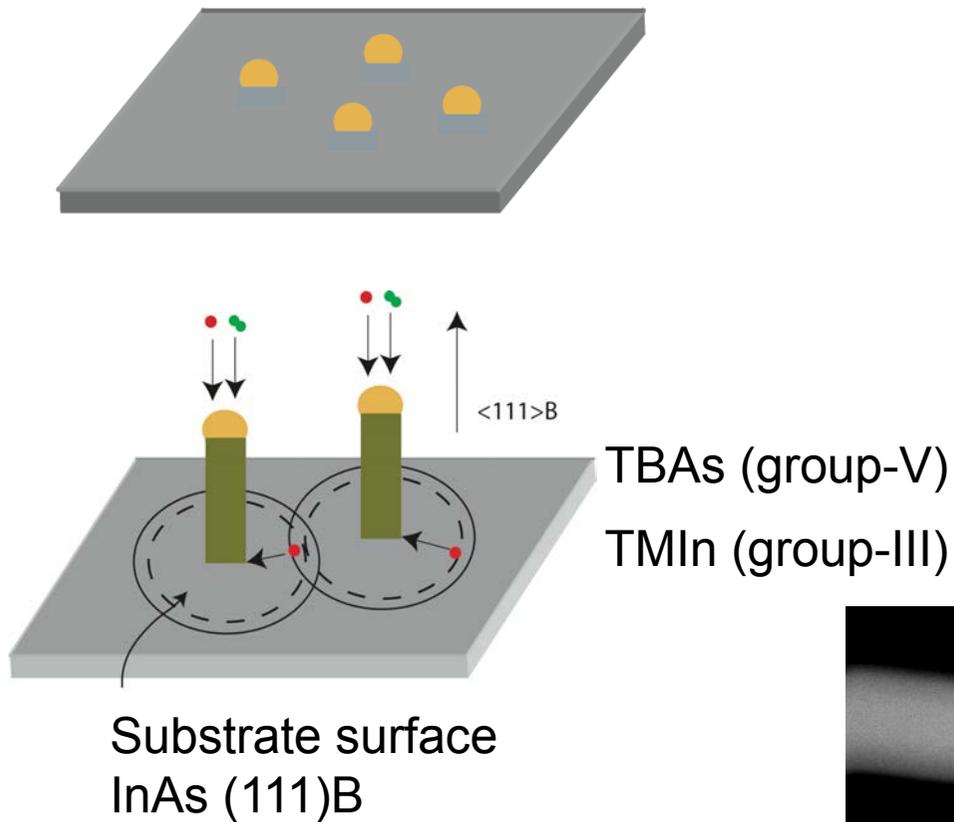


$$\eta = \frac{\eta_c}{2} + \frac{\eta_c^2}{8} + \frac{[7 + \operatorname{csch}^2(a_0/2)]}{96} \eta_c^3 + \mathcal{O}(\eta_c^4)$$

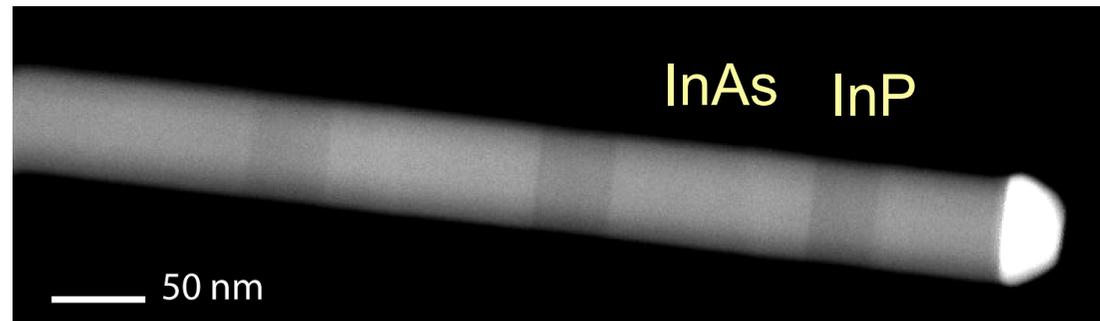
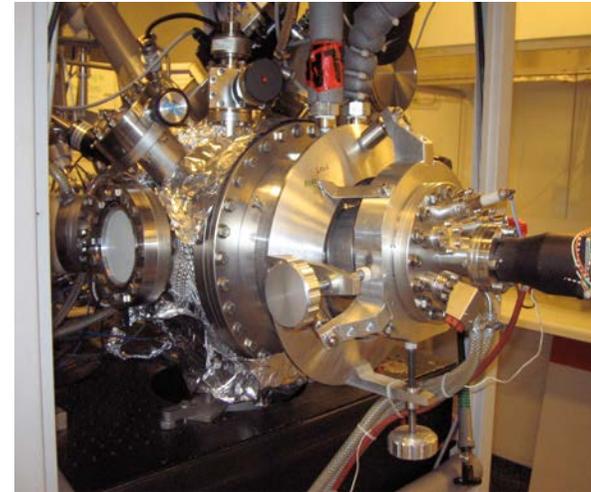
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# Epitaxially grown nanowires, e.g. InAs/InP



CBE



Conduction band edge

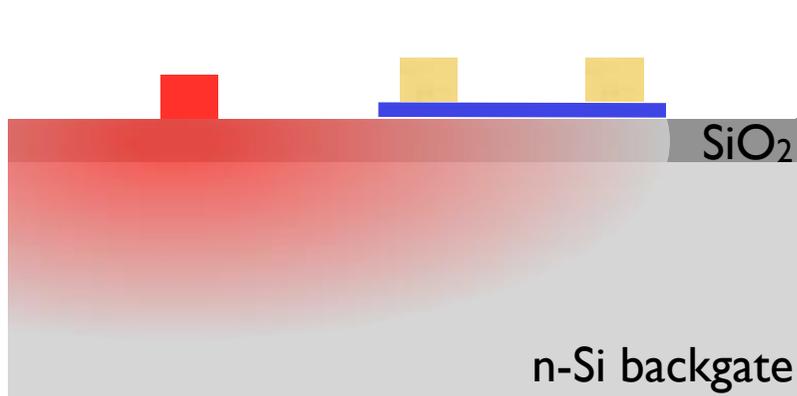
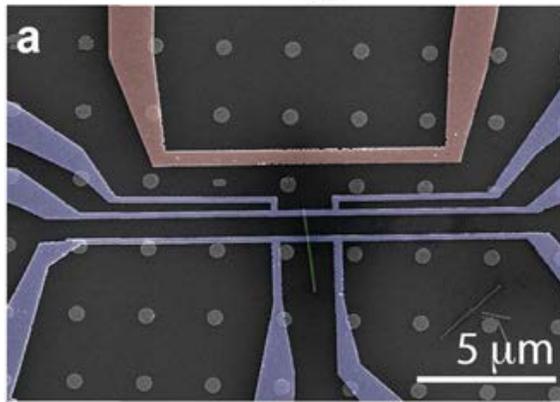


# Applying a temperature gradient along a nanowire



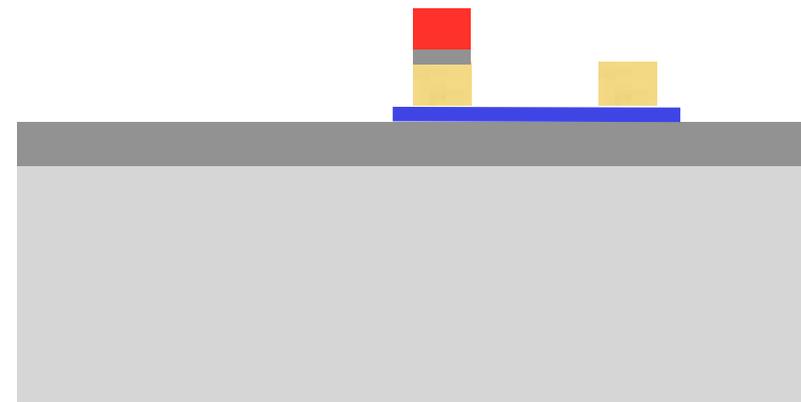
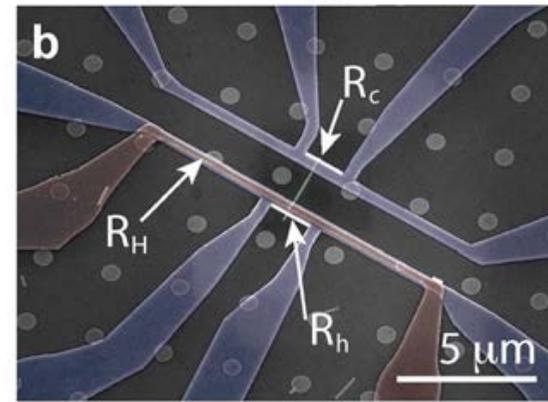
LUND  
UNIVERSITY

## Traditional side-heater



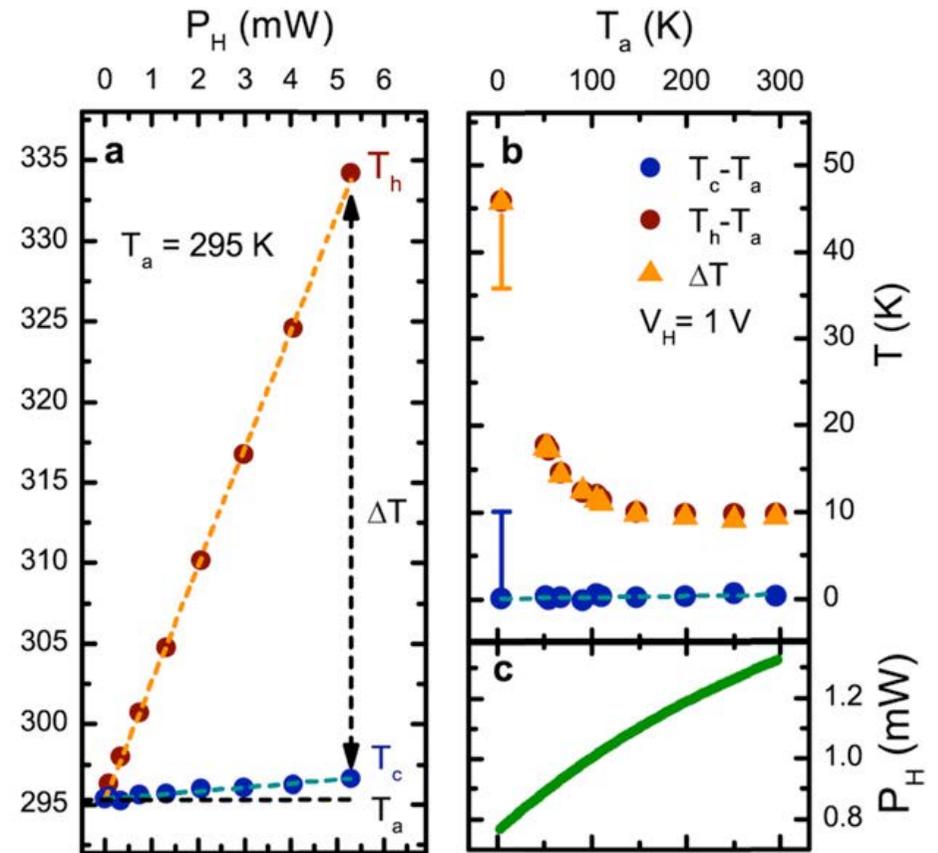
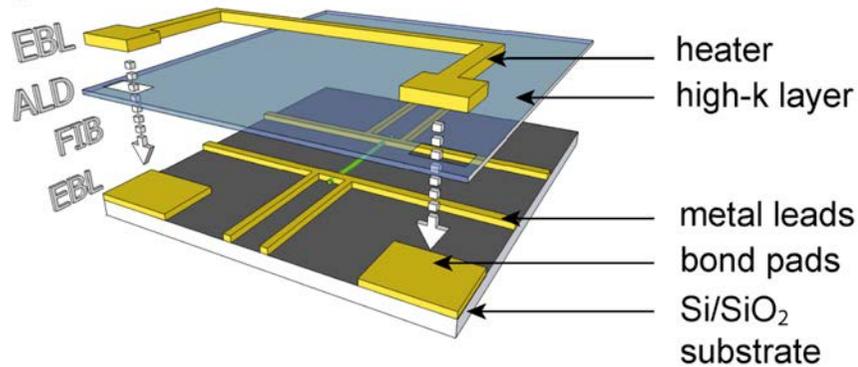
Substantial global device heating limits use for low-temperature experiments

## Local top-heater



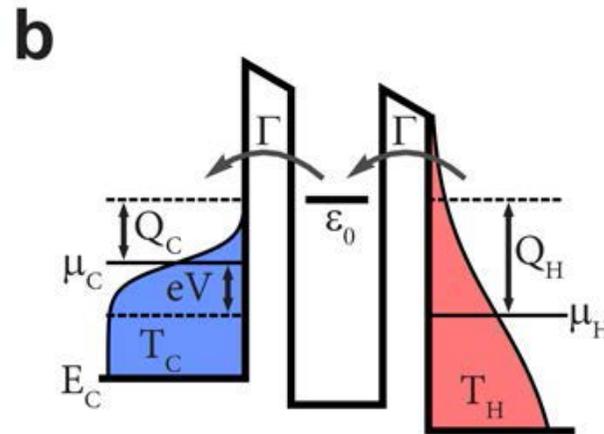
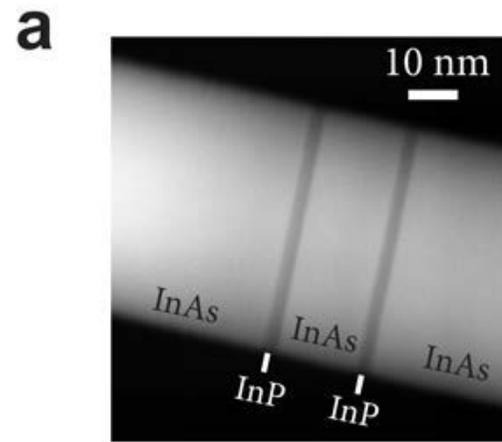
Local, direct heating of the warm contact without electrical interference.

# Top-heaters to enable high $\Delta T$ with minimal heating

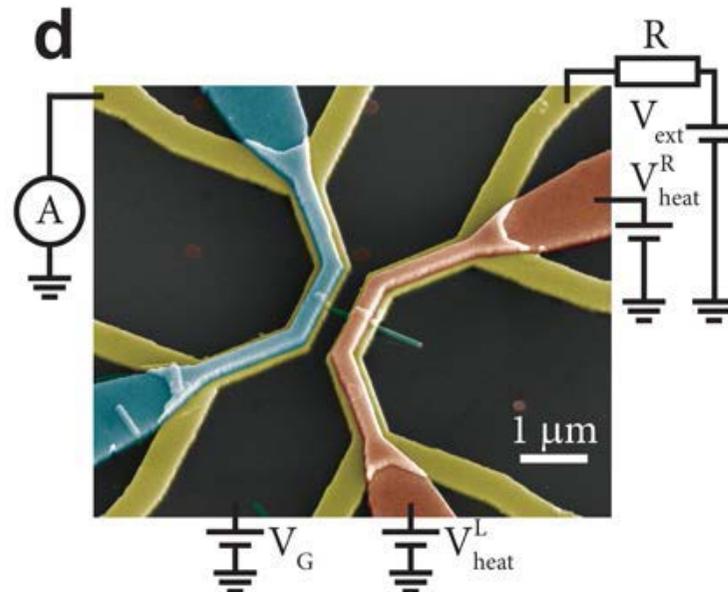
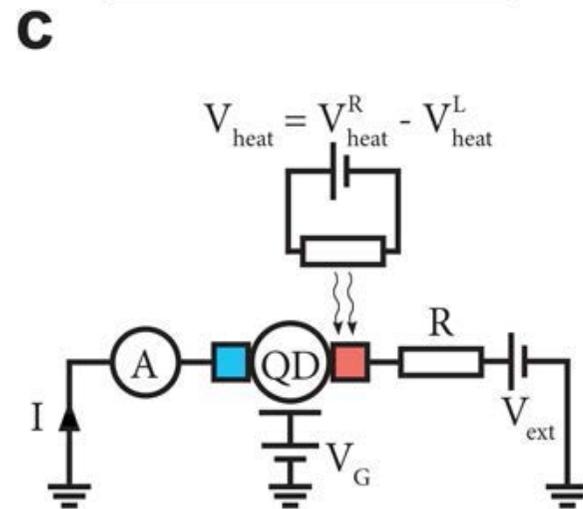


J. Gluschke et al  
Nanotechnology **25**, 385704 (2014)

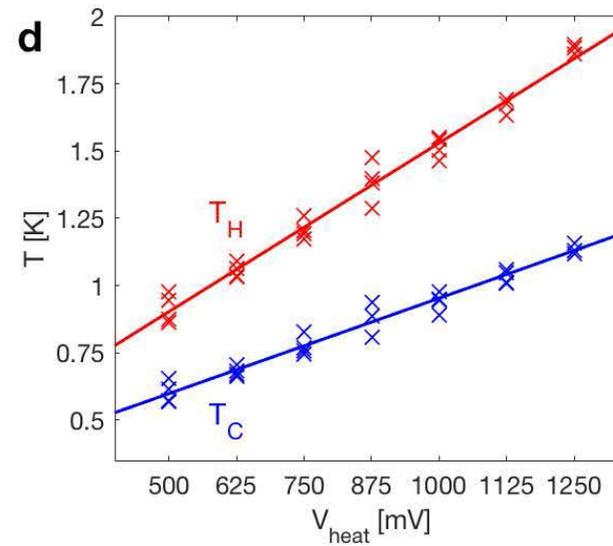
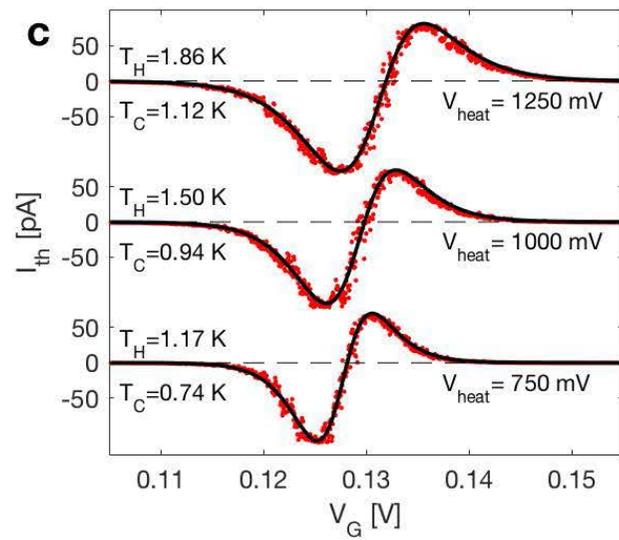
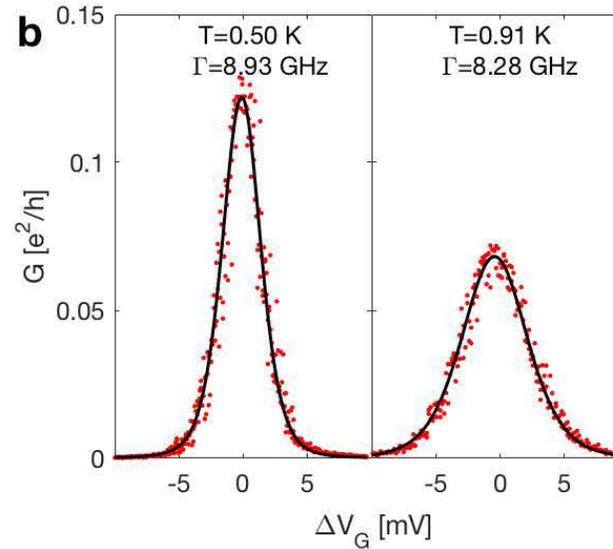
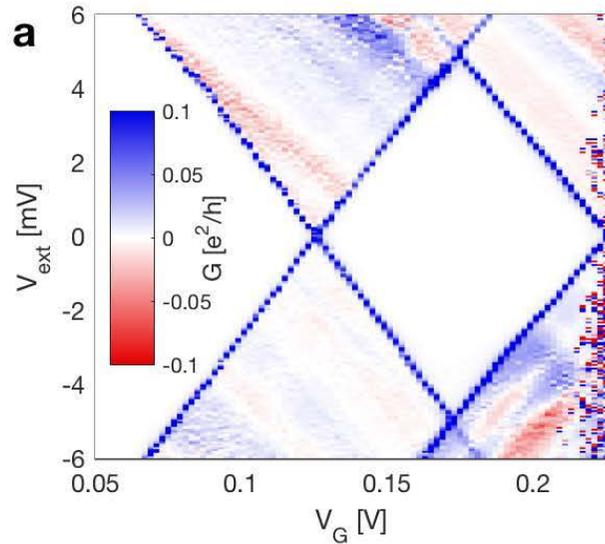
# Quantum-dot heat engine: device



Artis Svilans



# Quantum-dot heat engine: characterisation

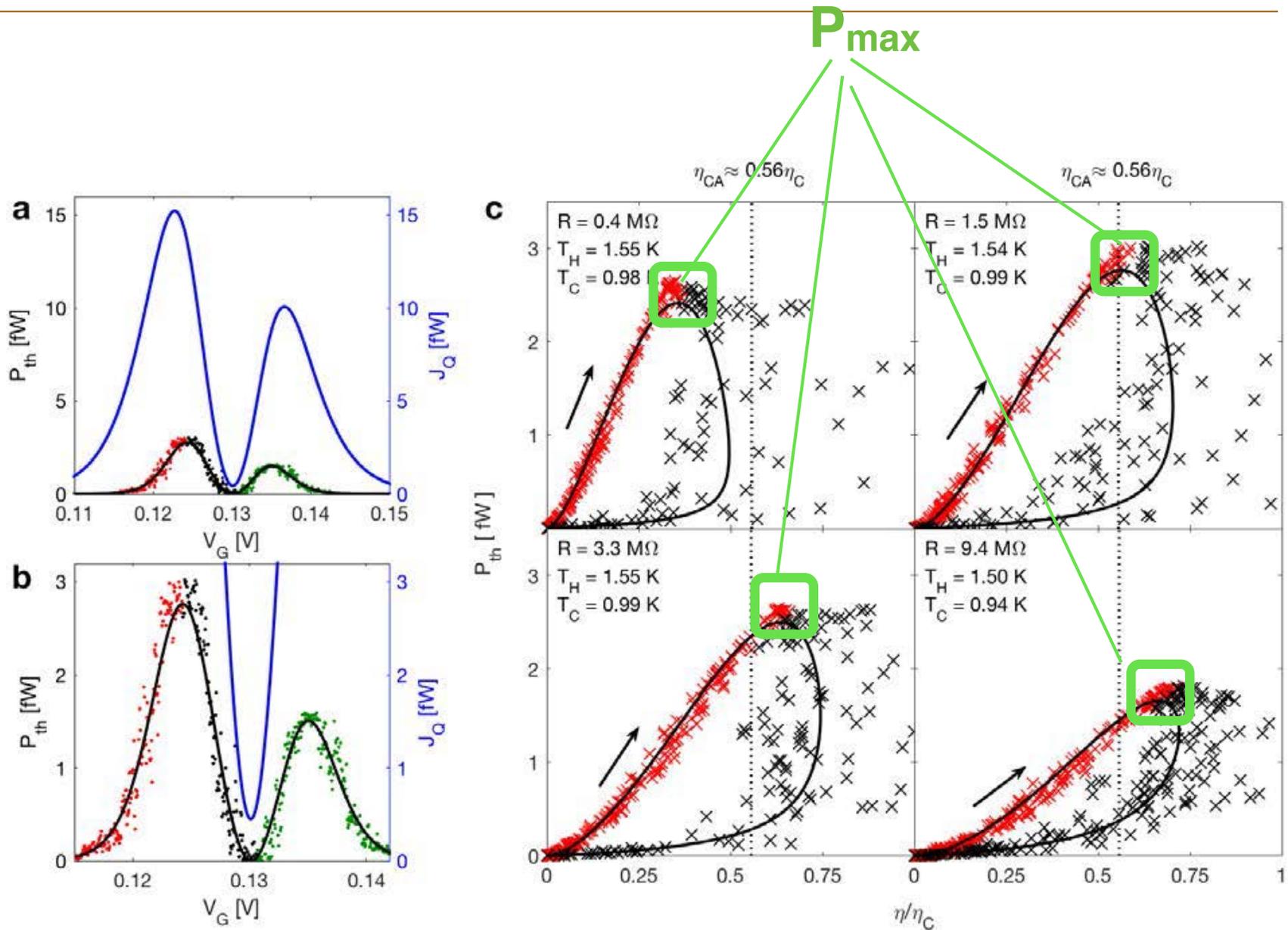


Artis Svilans

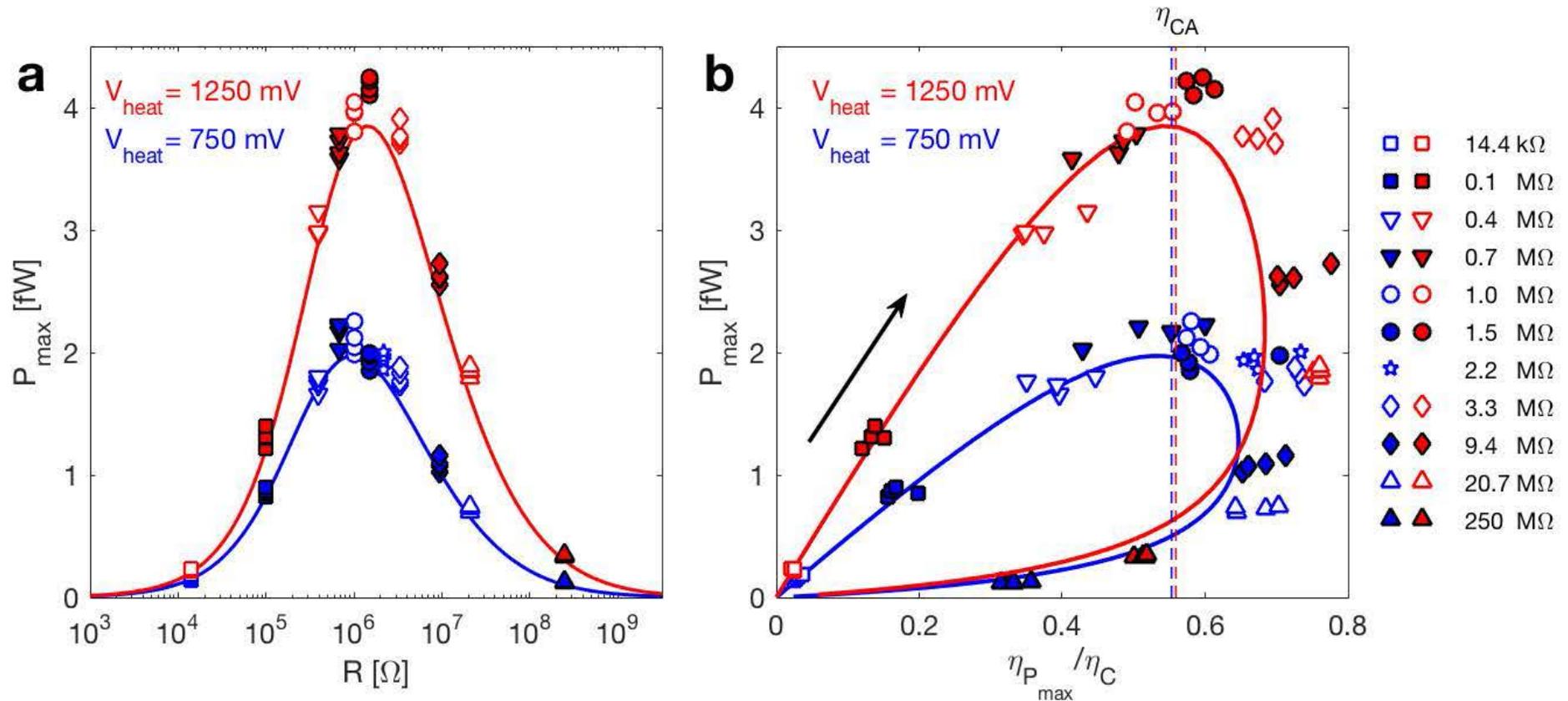


Martin Josefsson

# Quantum-dot heat engine: performance



# Quantum-dot heat engine: 70% of Carnot efficiency demonstrated



Quantum-dot heat engine achieves Curzon-Ahlborn efficiency at maximum power and about 70 % of Carnot efficiency with finite power output

# Outline

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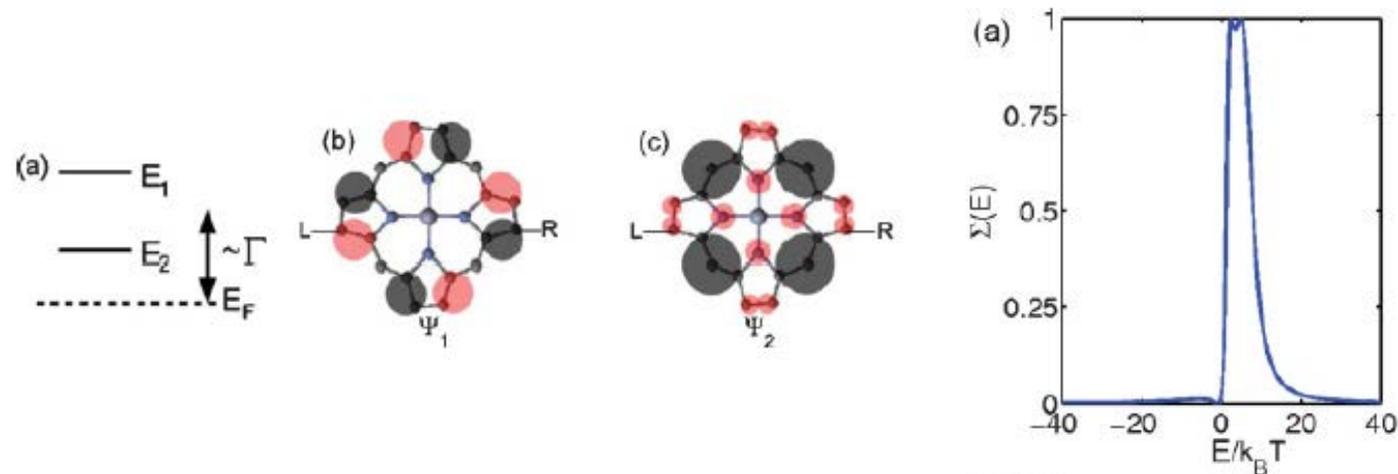
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# Higher maximum power by using interference

PHYSICAL REVIEW B 84, 113415 (2011)

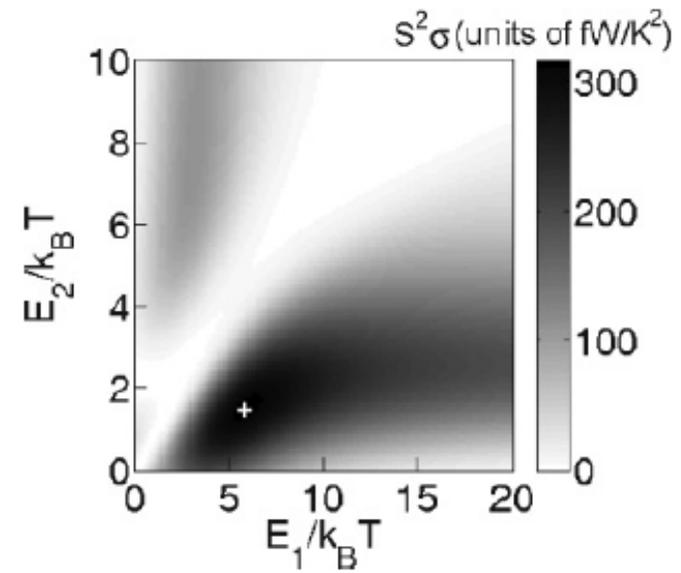
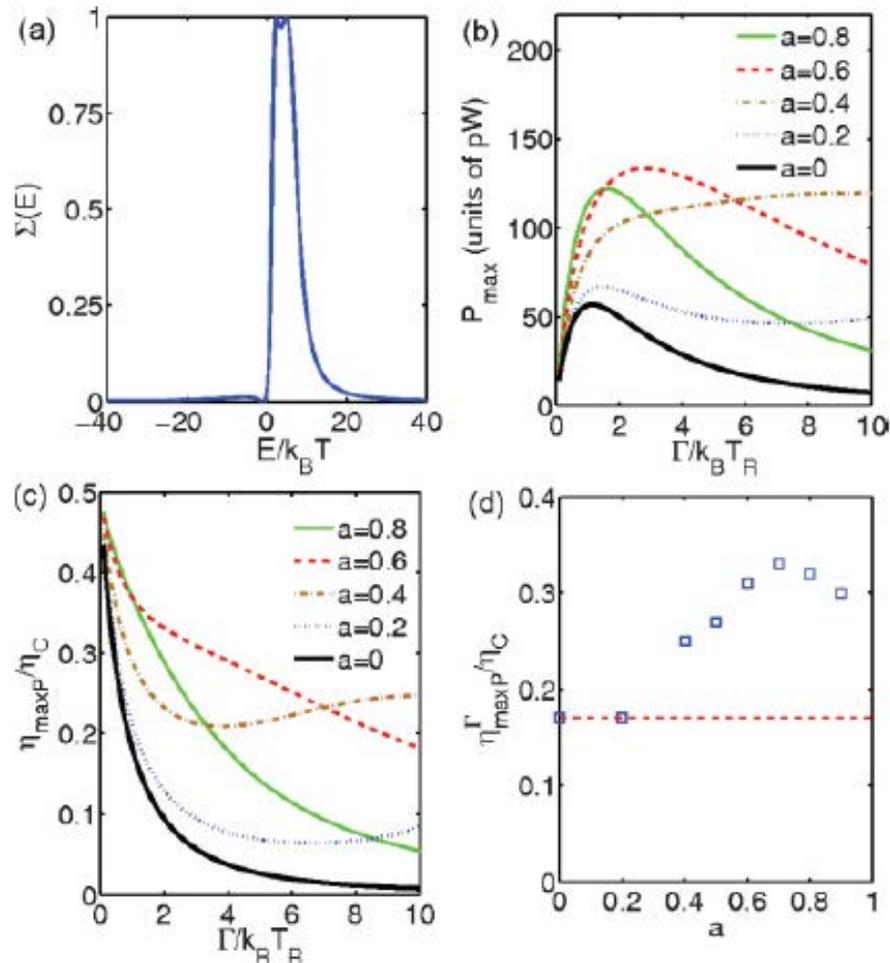
## Increasing thermoelectric performance using coherent transport

O. Karlström,<sup>1,\*</sup> H. Linke,<sup>1</sup> G. Karlström,<sup>2</sup> and A. Wacker<sup>1</sup>



$$\Sigma(E) = \Gamma^2 \left| \frac{1}{E - E_1 + i\Gamma} - \frac{a^2}{E - E_2 + ia^2\Gamma} \right|^2$$

# Higher maximum power by using interference



$$\Sigma(E) = \Gamma^2 \left| \frac{1}{E - E_1 + i\Gamma} - \frac{a^2}{E - E_2 + ia^2\Gamma} \right|^2$$

# First step: interference effects yield measurable difference in thermopower



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**NANO** LETTERS

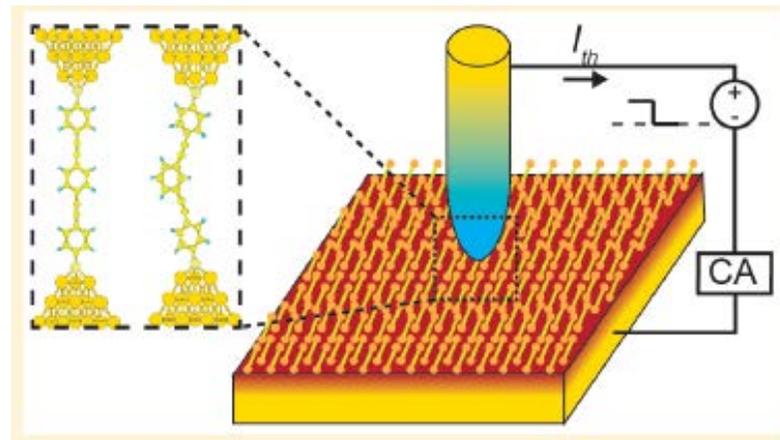
Letter

[pubs.acs.org/NanoLett](https://pubs.acs.org/NanoLett)

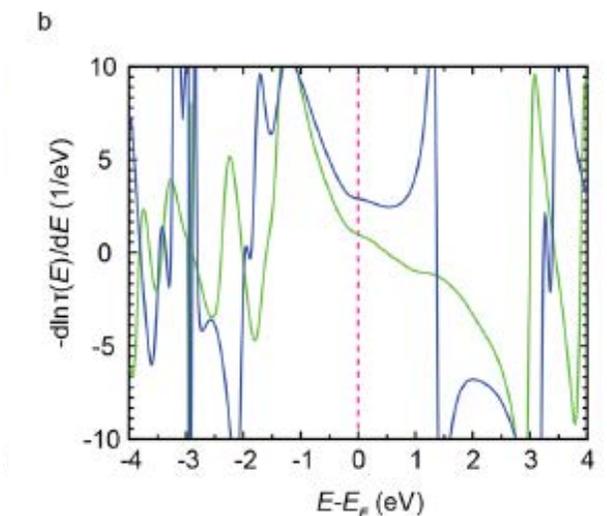
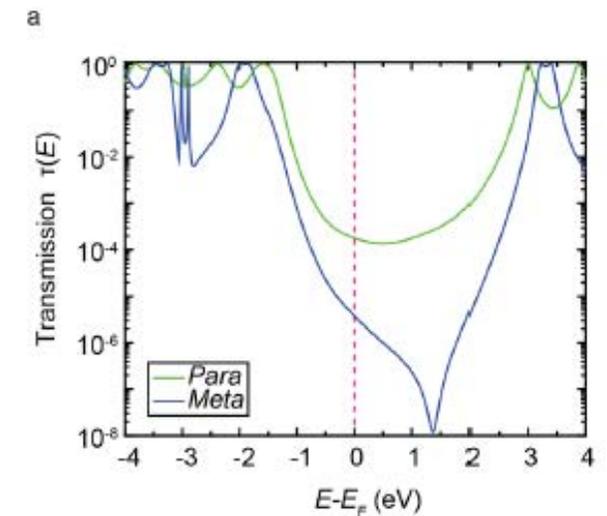
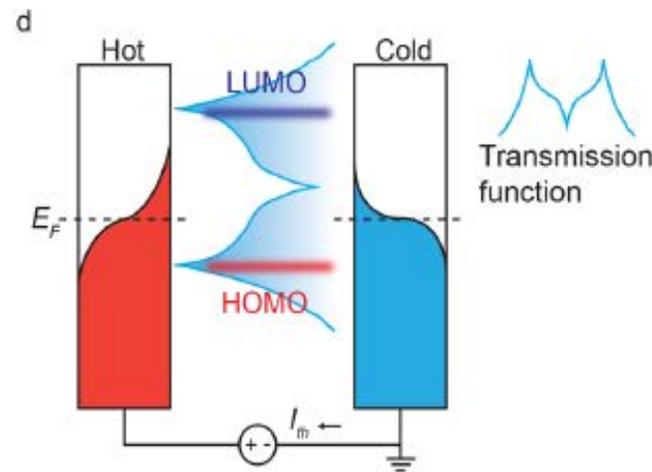
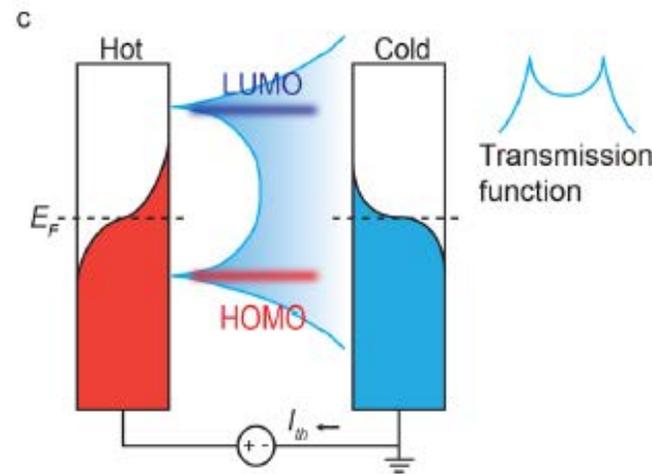
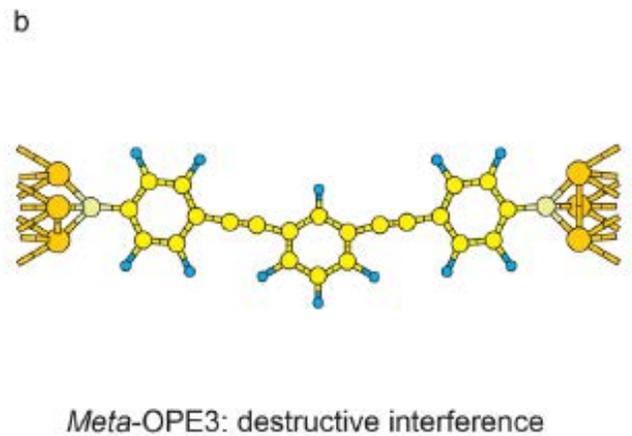
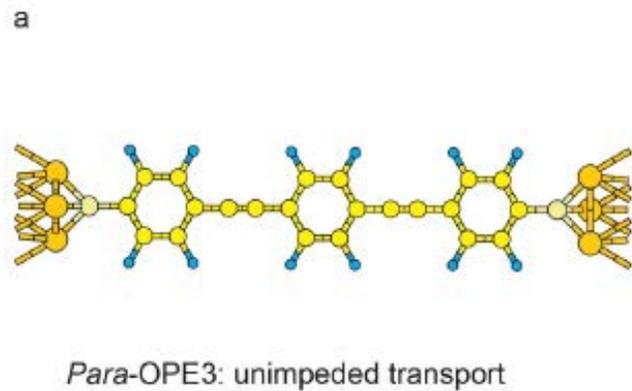
## 1 Influence of Quantum Interference on the Thermoelectric Properties 2 of Molecular Junctions

3 Ruijiao Miao,<sup>†</sup> Hailiang Xu,<sup>‡,§</sup> Maxim Skripnik,<sup>||,⊥</sup> Longji Cui,<sup>†,⊕</sup> Kun Wang,<sup>†</sup> Kim G. L. Pedersen,<sup>#,▽</sup>  
4 Martin Leijnse,<sup>‡,⊖</sup> Fabian Pauly,<sup>||,⊥,⊕</sup> Kenneth Wärnmark,<sup>‡,§,⊖</sup> Edgar Meyhofer,<sup>\*,†,⊖</sup>  
5 Pramod Reddy,<sup>\*,†,◆,⊖</sup> and Heiner Linke<sup>\*,‡,⊖</sup>

August 2018



# First step: interference effects yield measurable difference in thermopower

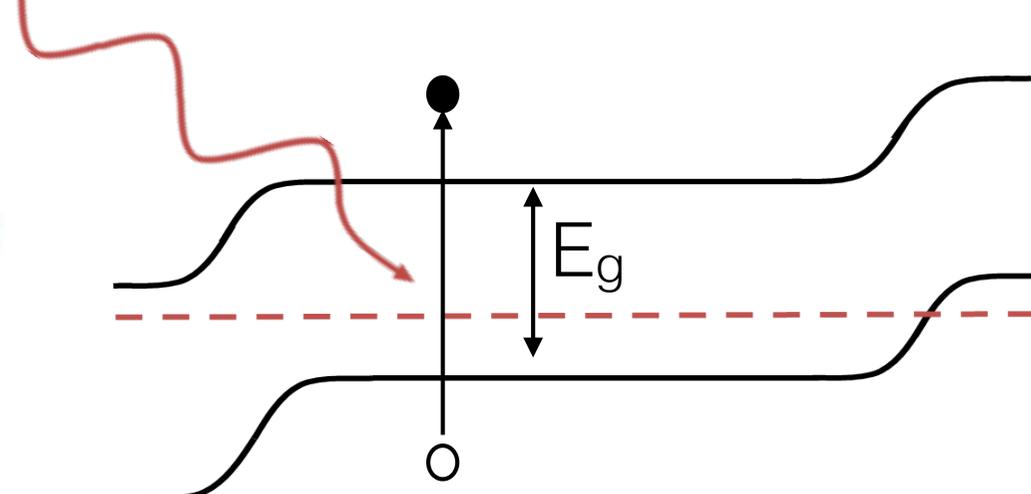
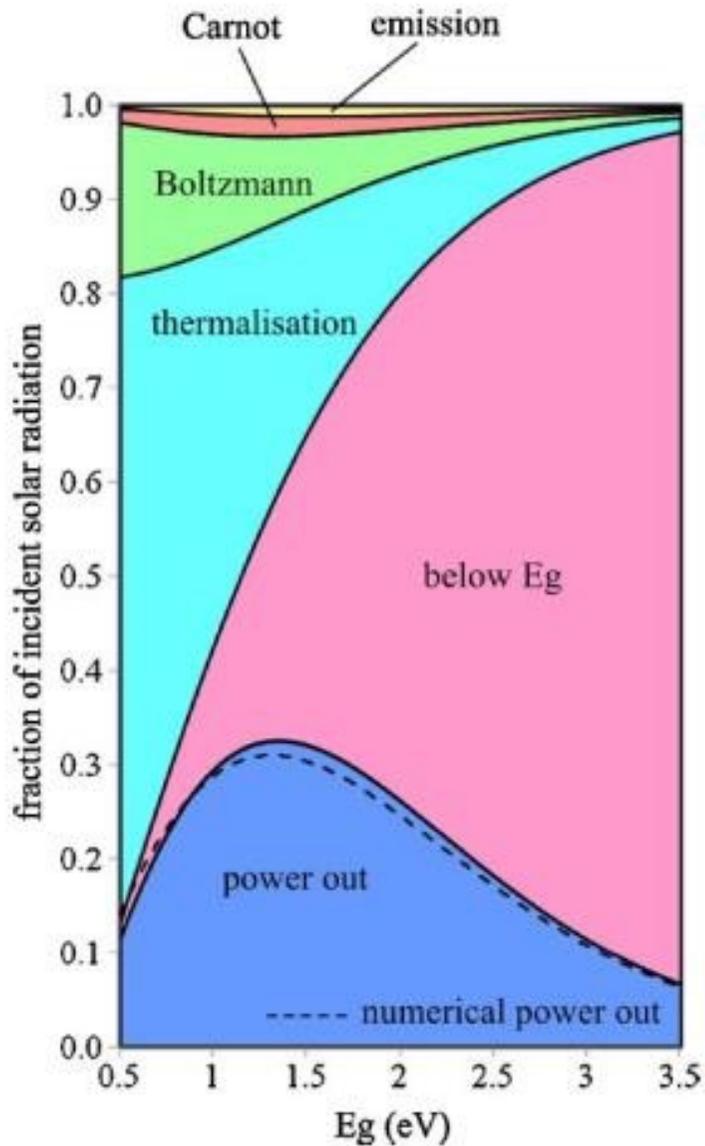


Higher thermopower  $S$  predicted in presence of destructive interference

# Outline

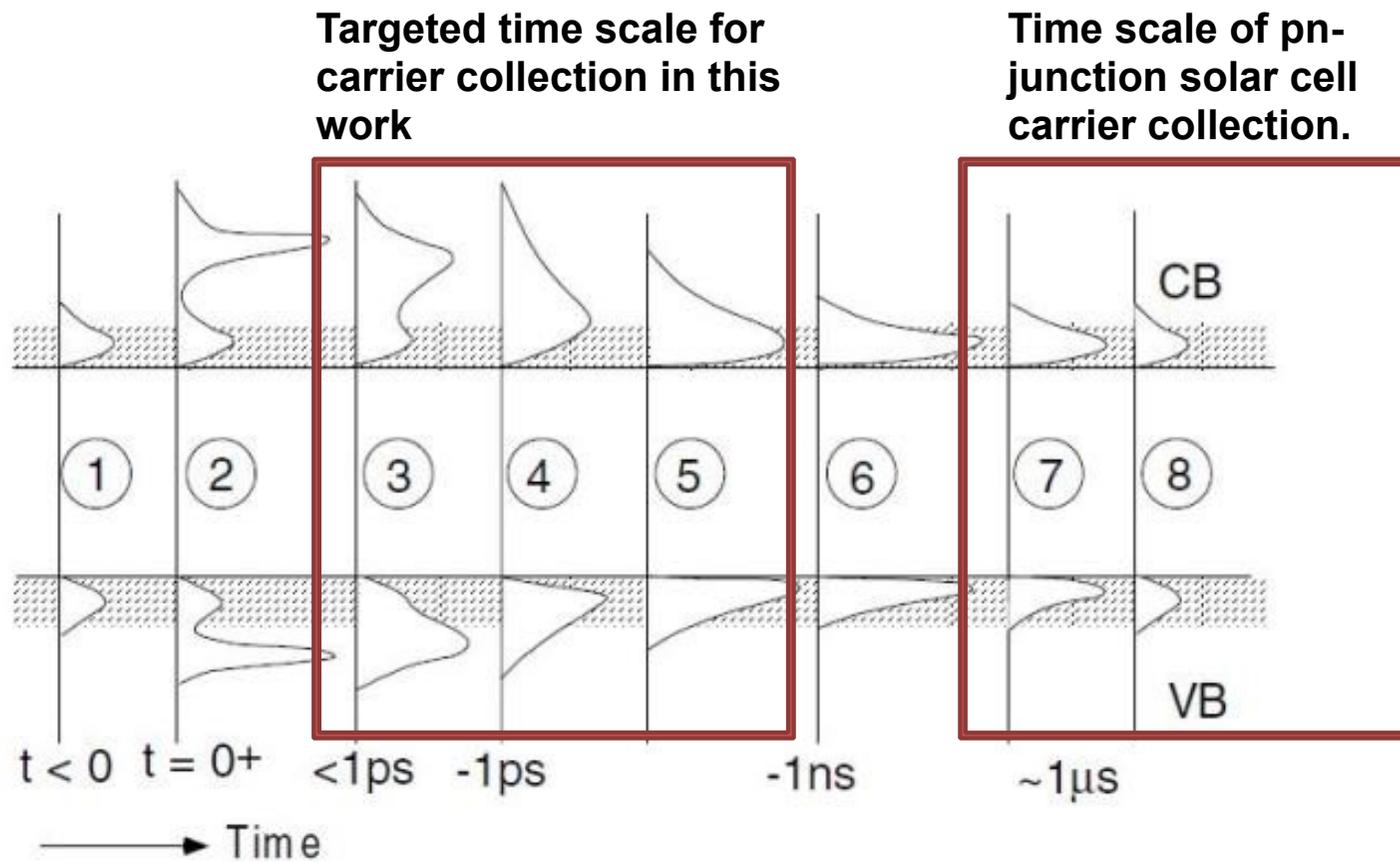
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- Application of energy filtering to hot-carrier solar cells



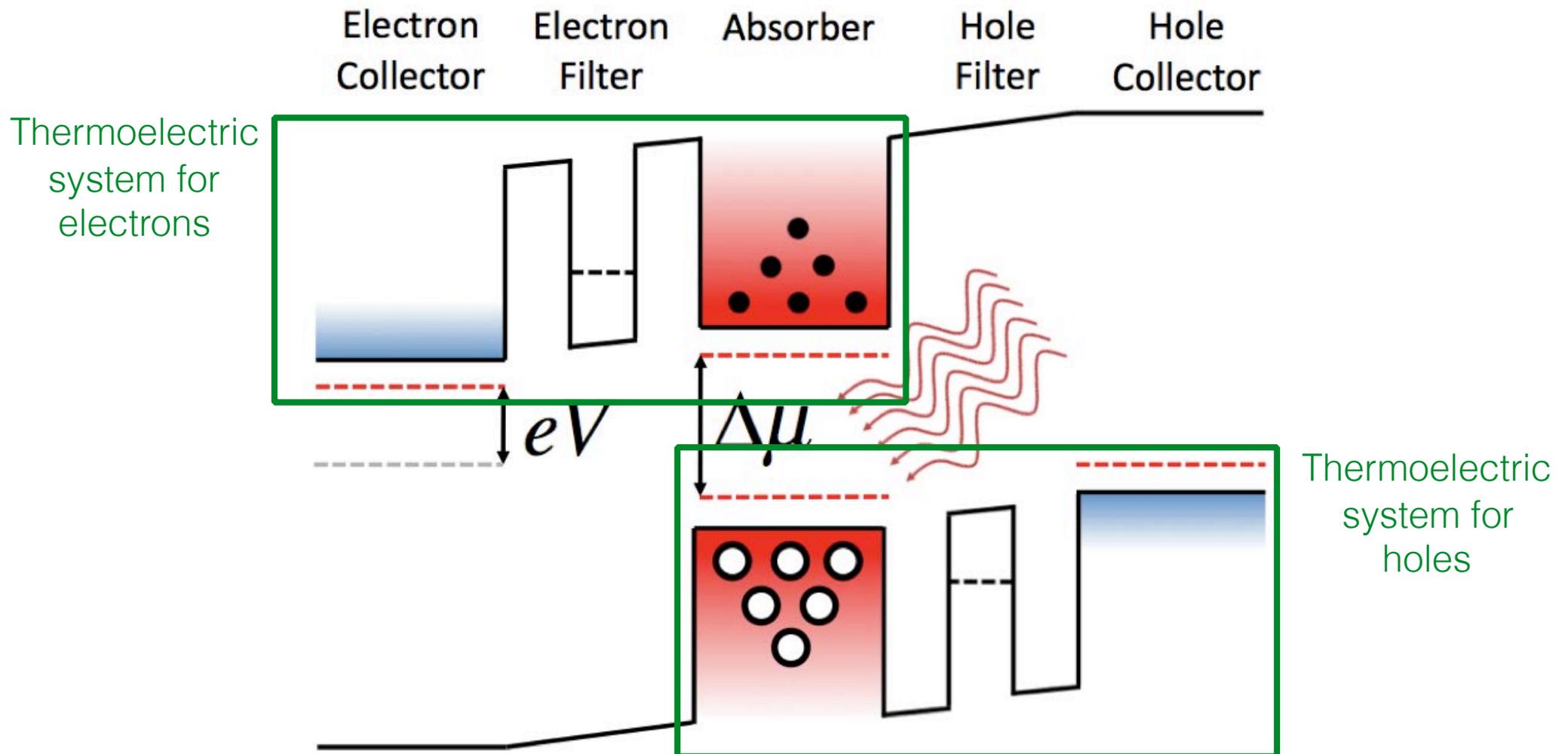
- Carrier cooling decreases the energy each carrier can provide to an external circuit.
- pn-junction solar cells of small bandgap materials are rarely made due to the magnitude of thermalisation losses.

**Figure 6.** Intrinsic loss processes and hence, power out are shown to be dependent on  $E_g$ . All incident radiation is accounted for, illustrating why intrinsic loss mechanisms lead to fundamental limiting efficiency.

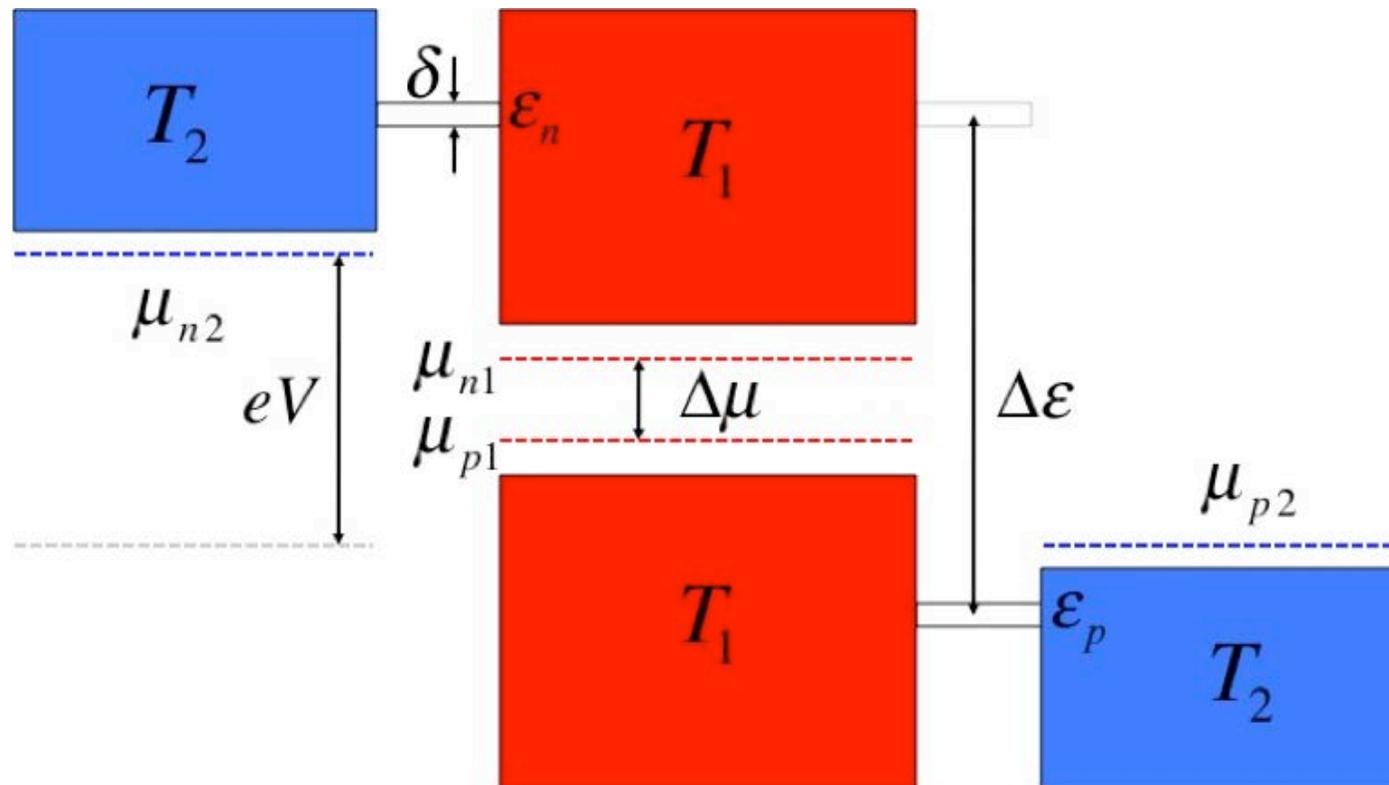


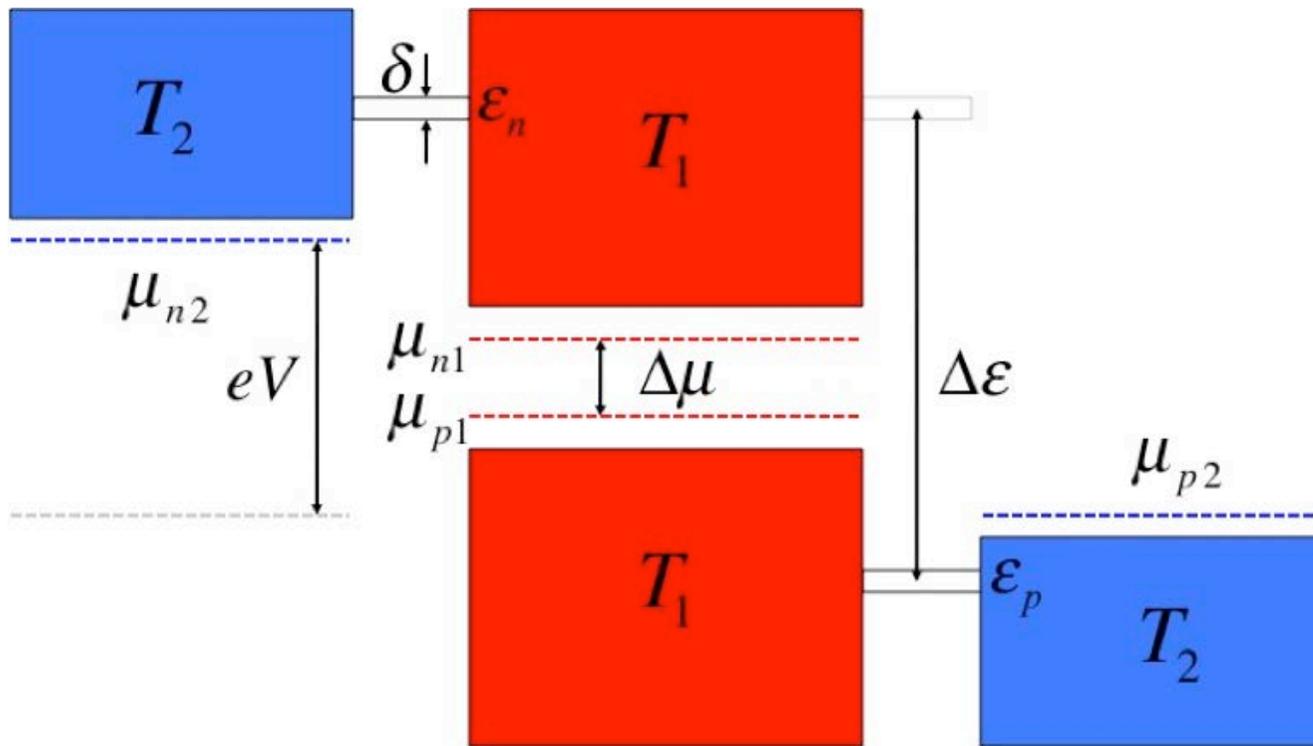
**Fig. 6.2:** Time evolution of electron and hole distributions in a semiconductor subject to a short, high intensity, monochromatic pulse of light from a laser: (1) Thermal equilibrium before pulse; (2) “coherent” stage straight after pulse; (3) carrier scattering; (4) thermalisation of “hot carriers”; (5) carrier cooling; (6) lattice thermalised carriers; (7) recombination of carriers; (8) return to thermal equilibrium.

# Basic idea of a hot-carrier cell: photothermoelectrics



# Can a hot-carrier photovoltaic system be run reversibly?





$$\begin{aligned} \Delta S &= \Delta S_{n1,ext} + \Delta S_{n2,inj} + \Delta S_{p1,ext} + \Delta S_{p2,inj} \\ &= \frac{\Delta \epsilon - eV}{T_2} - \frac{\Delta \epsilon - \Delta \mu}{T_1}. \end{aligned}$$

$$\Delta S = 0 \text{ when } \frac{\Delta \epsilon - \Delta \mu}{T_1} = \frac{\Delta \epsilon - eV}{T_2}$$

**(equivalent to energy-specific equilibrium across both junctions)**

# Open-circuit voltage



Steven Limpert

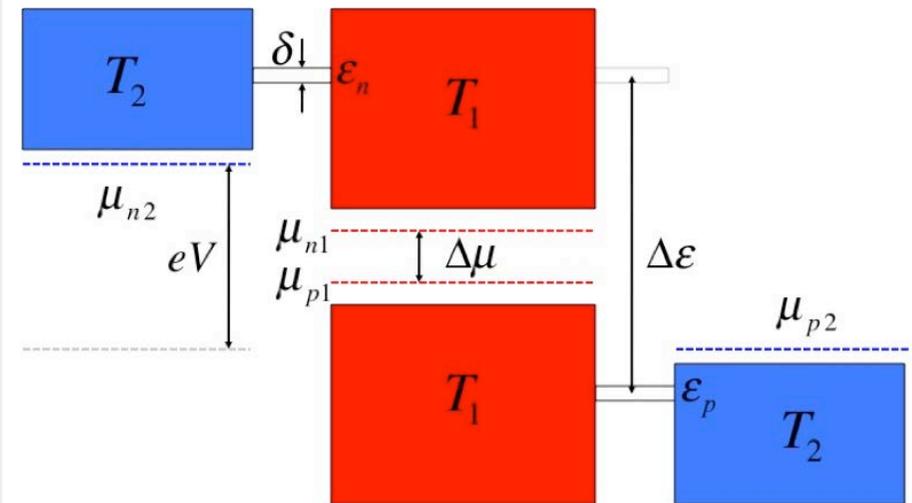
$$\begin{aligned}\Delta S &= \Delta S_{n1,ext} + \Delta S_{n2,inj} + \Delta S_{p1,ext} + \Delta S_{p2,inj} \\ &= \frac{\Delta\varepsilon - eV}{T_2} - \frac{\Delta\varepsilon - \Delta\mu}{T_1}.\end{aligned}$$

$$eV = \Delta\varepsilon \left(1 - \frac{T_2}{T_1}\right) + \Delta\mu \frac{T_2}{T_1} - T_2 \Delta S.$$

$$= Q_1 \eta_{Carnot} + \Delta\mu - T_2 \Delta S,$$

Heat engine      Solar cell

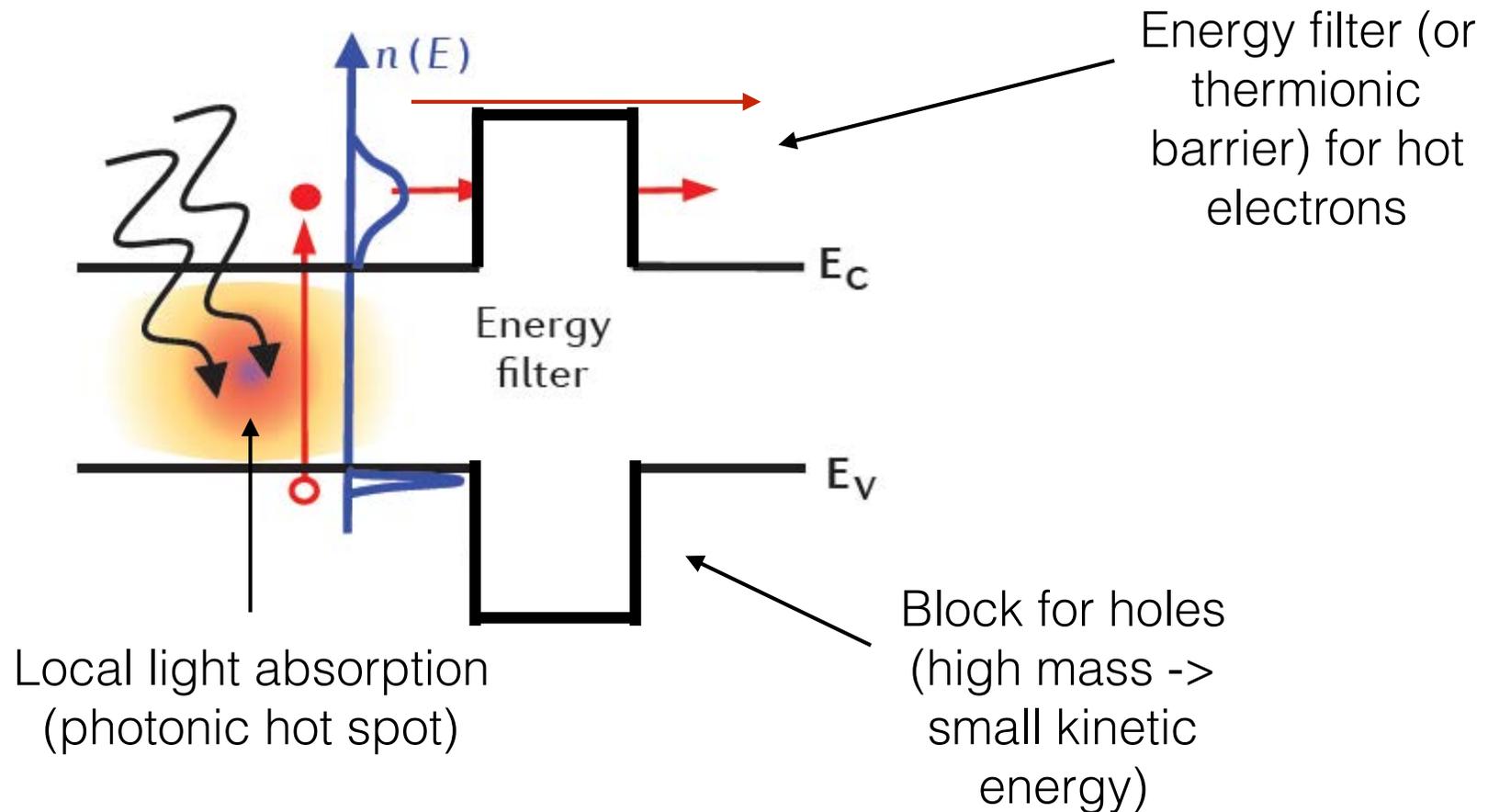
Explicit term describing the  
reduction of voltage due to  
irreversibility



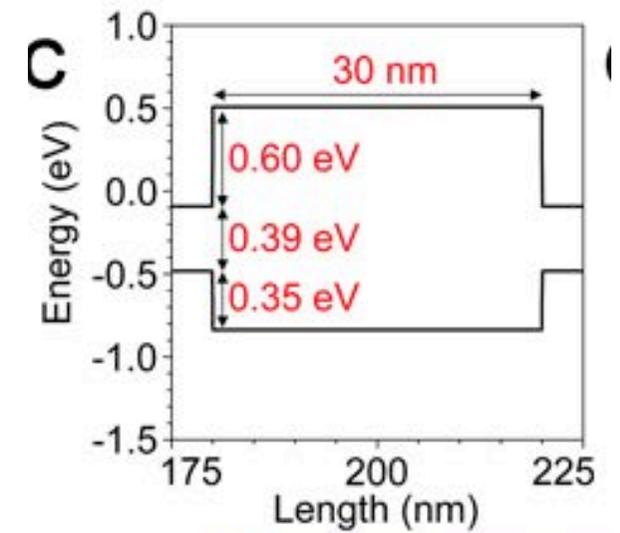
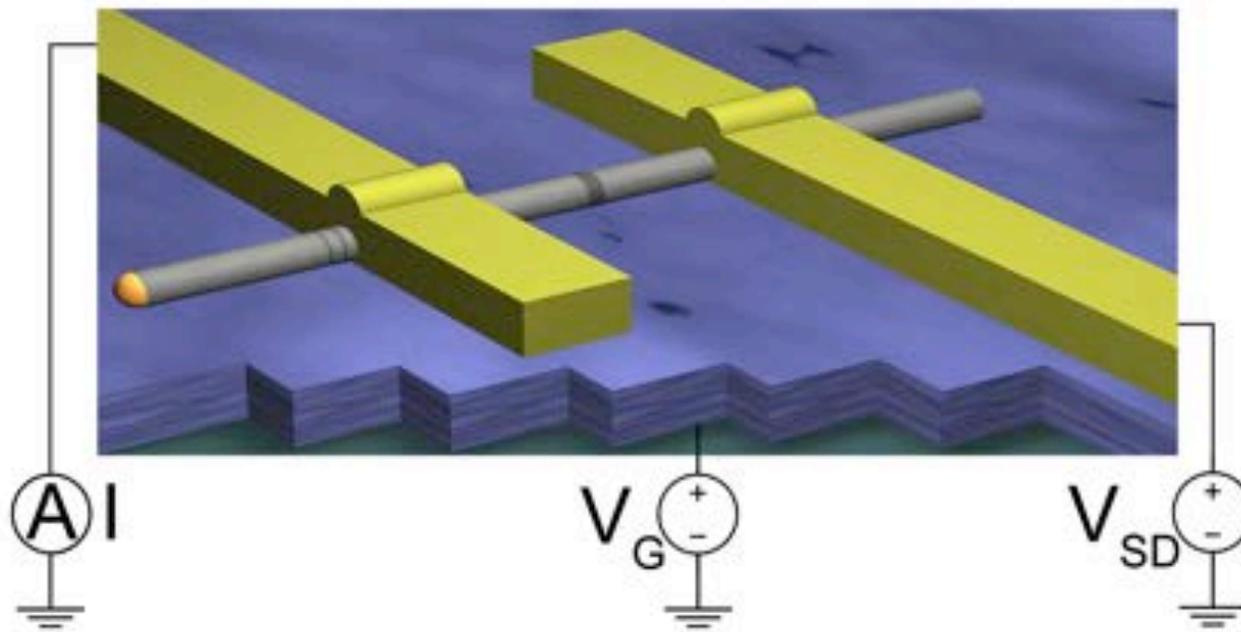
$$Q_1 = Q_{n1} + Q_{p1} = (\varepsilon_n - \mu_{n1}) + (\mu_{p1} - \varepsilon_p) = \Delta\varepsilon - \Delta\mu.$$

# Basic idea for hot-carrier experiments

Heterostructure nanowire with small band gap and high electron-hole mass asymmetry (e.g. InAs/InP)



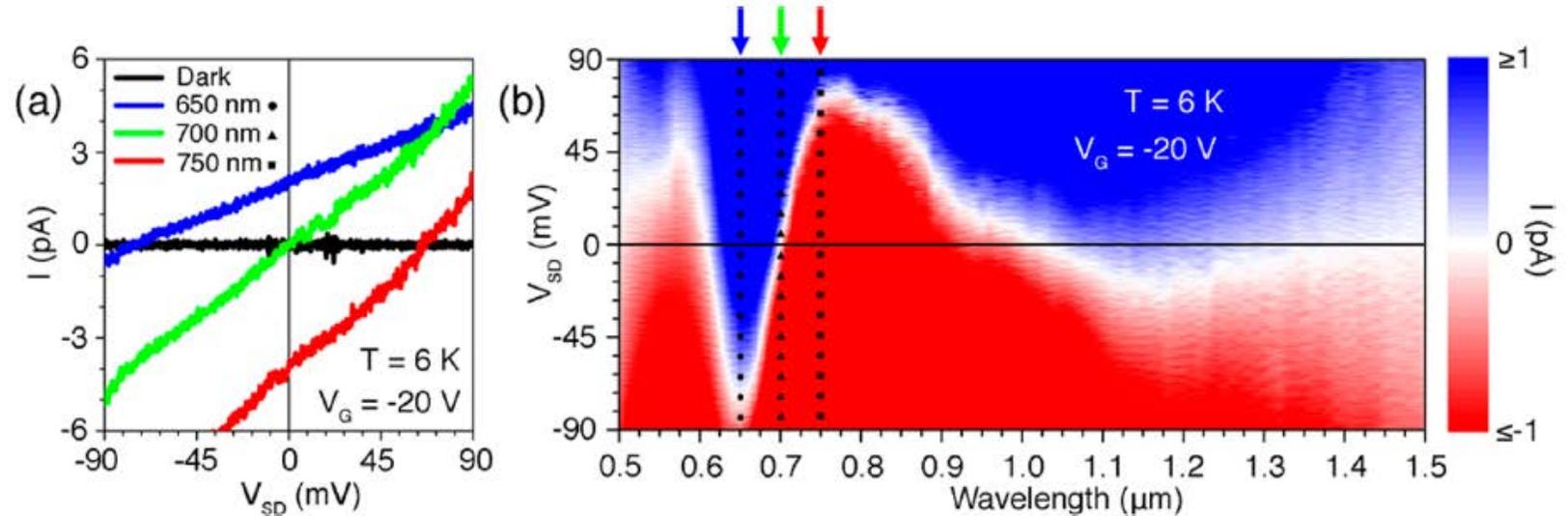
# Device



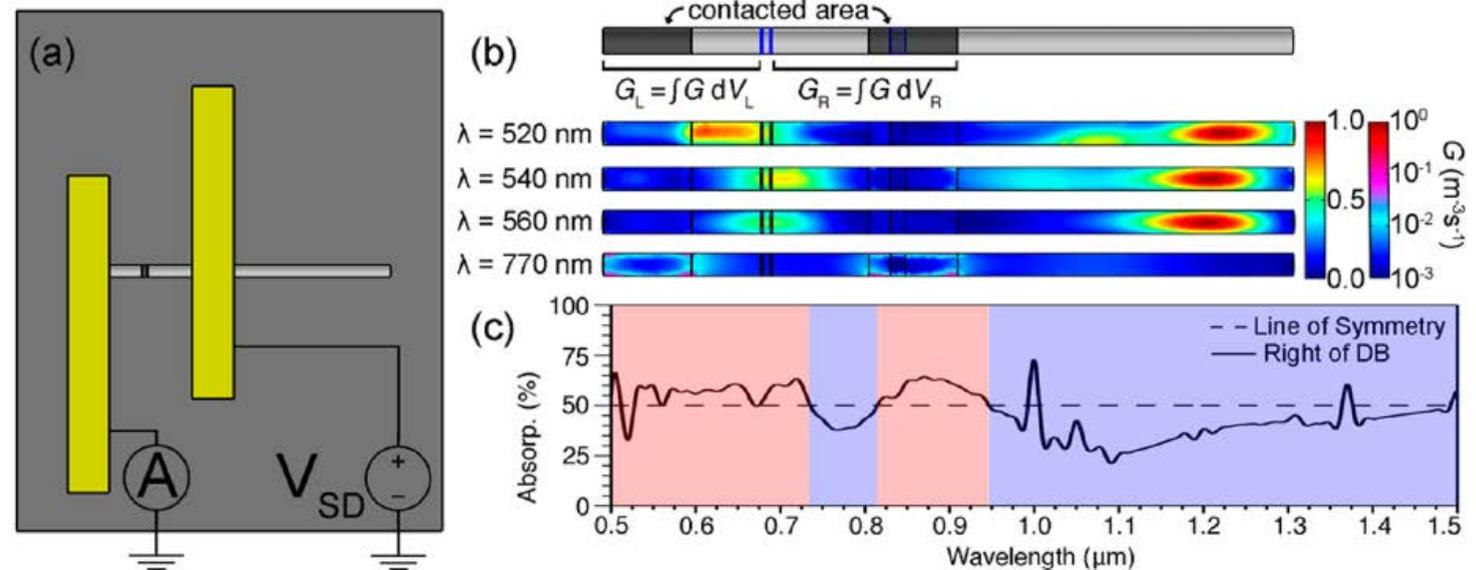
CBE grown InAs/InP/InAs nanowire

# Wavelength-sensitivity (Double-barrier device)

## Experiment

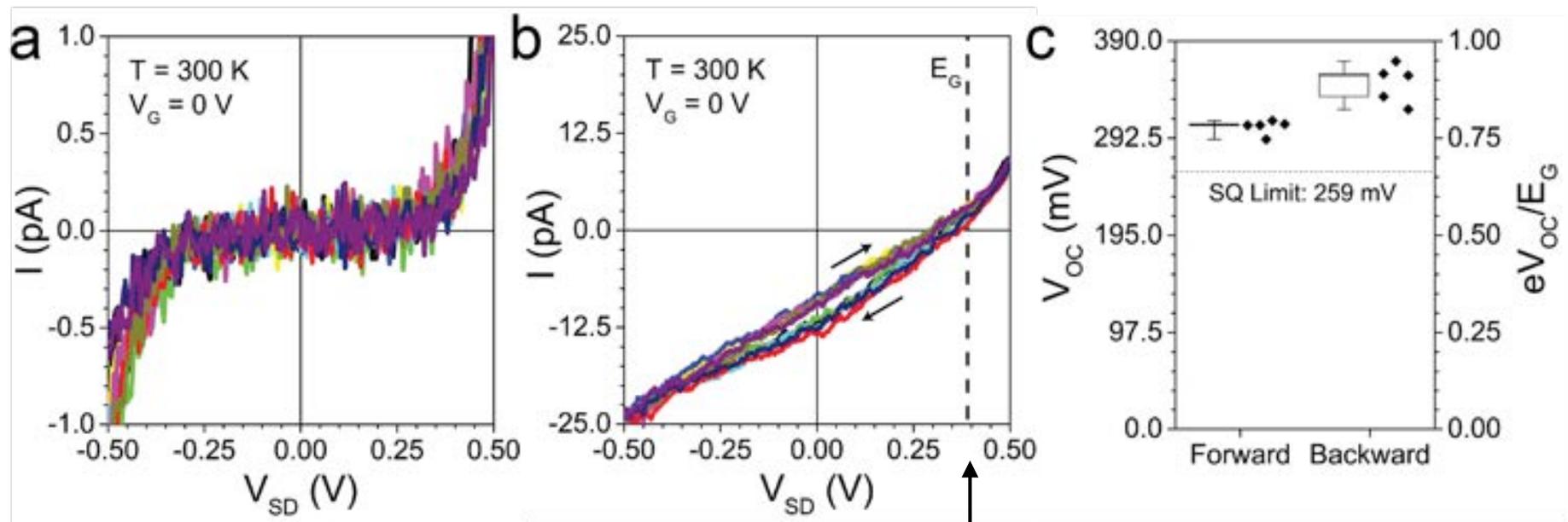


## Model



# Photovoltaic power production (without pn-junction!)

Single-barrier (thermionic) device  
 $V_{oc} > 90\%$  of the bandgap



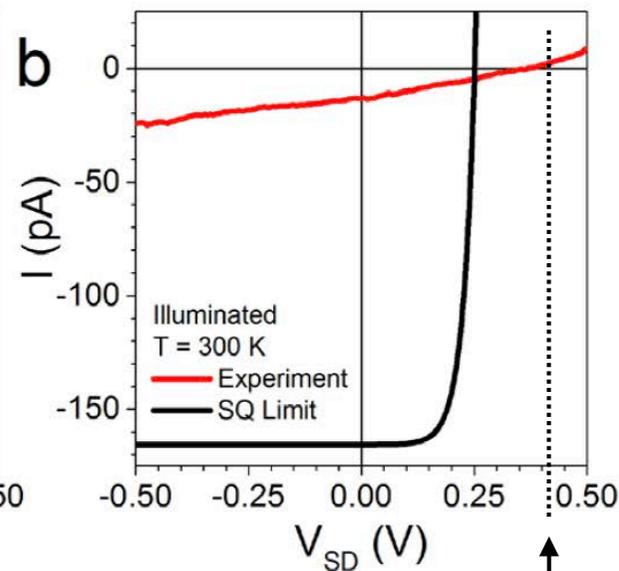
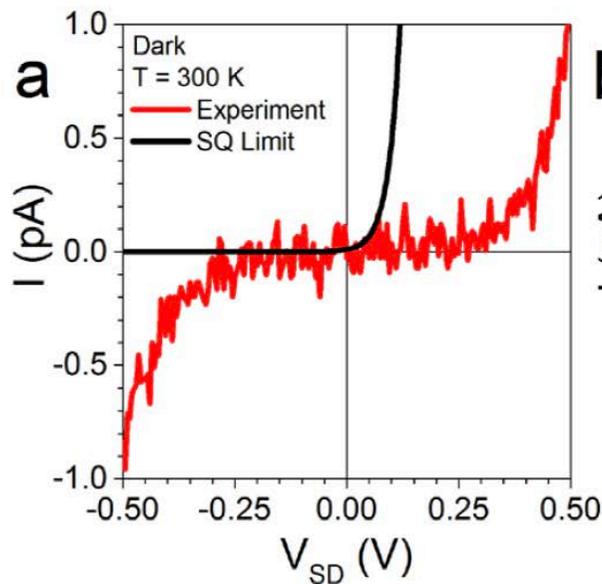
$E_G$  of WZ InAs  
 $\approx 0.39\text{ eV}$

S. Limpert et al, Nano Lett. **17**, 4055 (2017)

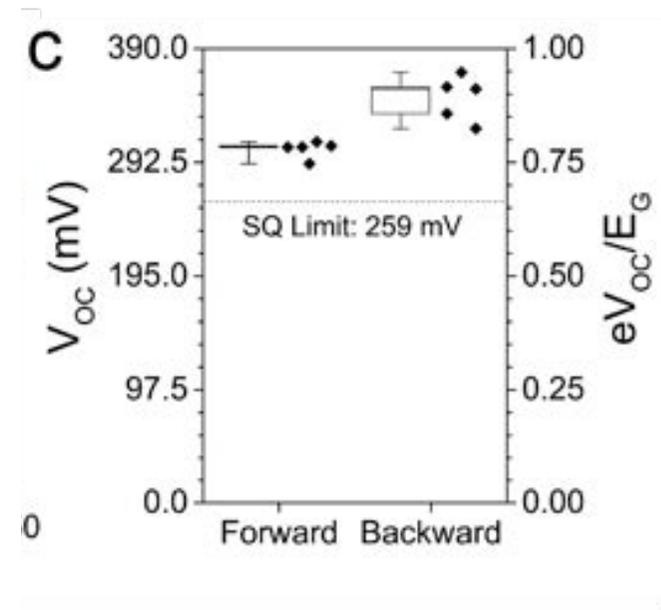
S. Limpert et al., Nanotechnology **28**, 43 (2017)

# Photovoltaic power production (without pn-junction!)

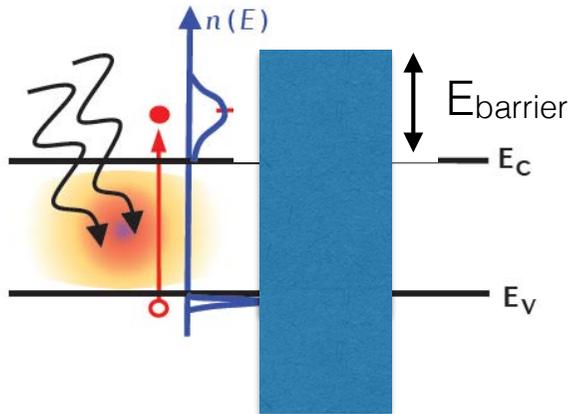
Single-barrier (thermionic) device  
 $V_{oc} > 90\%$  of the bandgap



$E_g$  of WZ InAs  
 $\approx 0.39$  eV



# Thermionic interpretation



Thermionic interpretation:

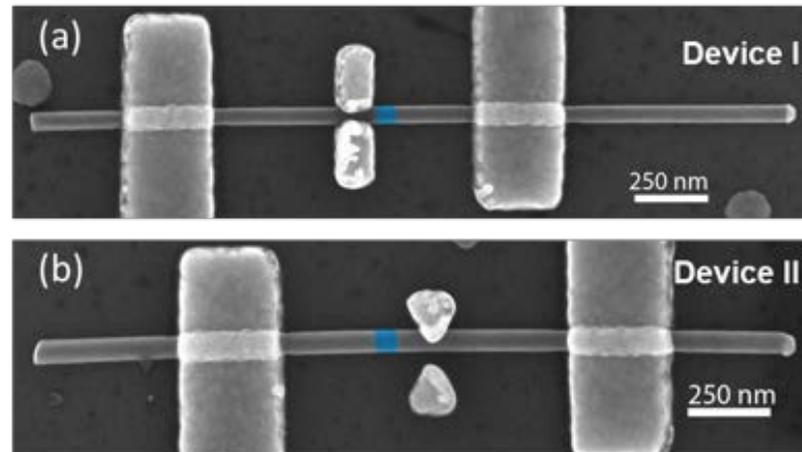
$$V_{oc} = (k/e) (2 + E_{\text{barrier}}/kT) \Delta T_{\text{carrier}}$$

$V_{oc} \approx 0.35$  V is consistent with  $\Delta T_{\text{carrier}} \approx 170$  K

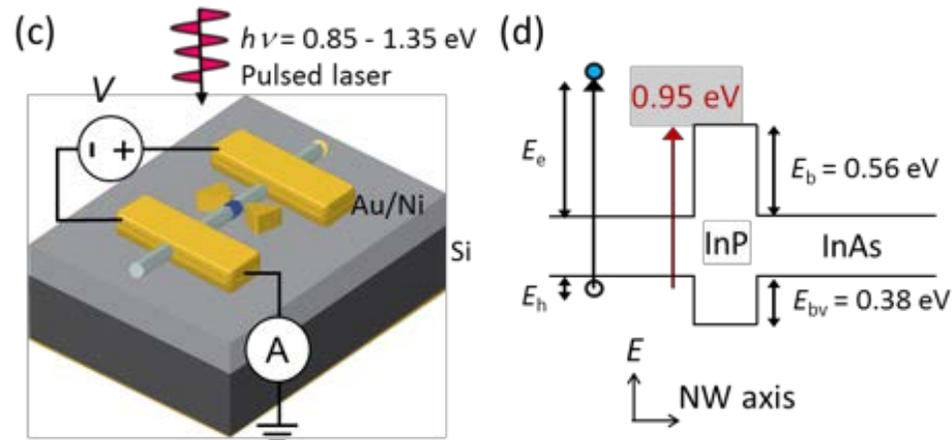
S. Limpert et al, Nano Lett. 17, 4055 (2017)

**Since  $\Delta T$  in this interpretation is the carrier temperature, phonon-mediated heat flow is irrelevant to the efficiency analysis.**

# Controlling the light-absorption hot spot

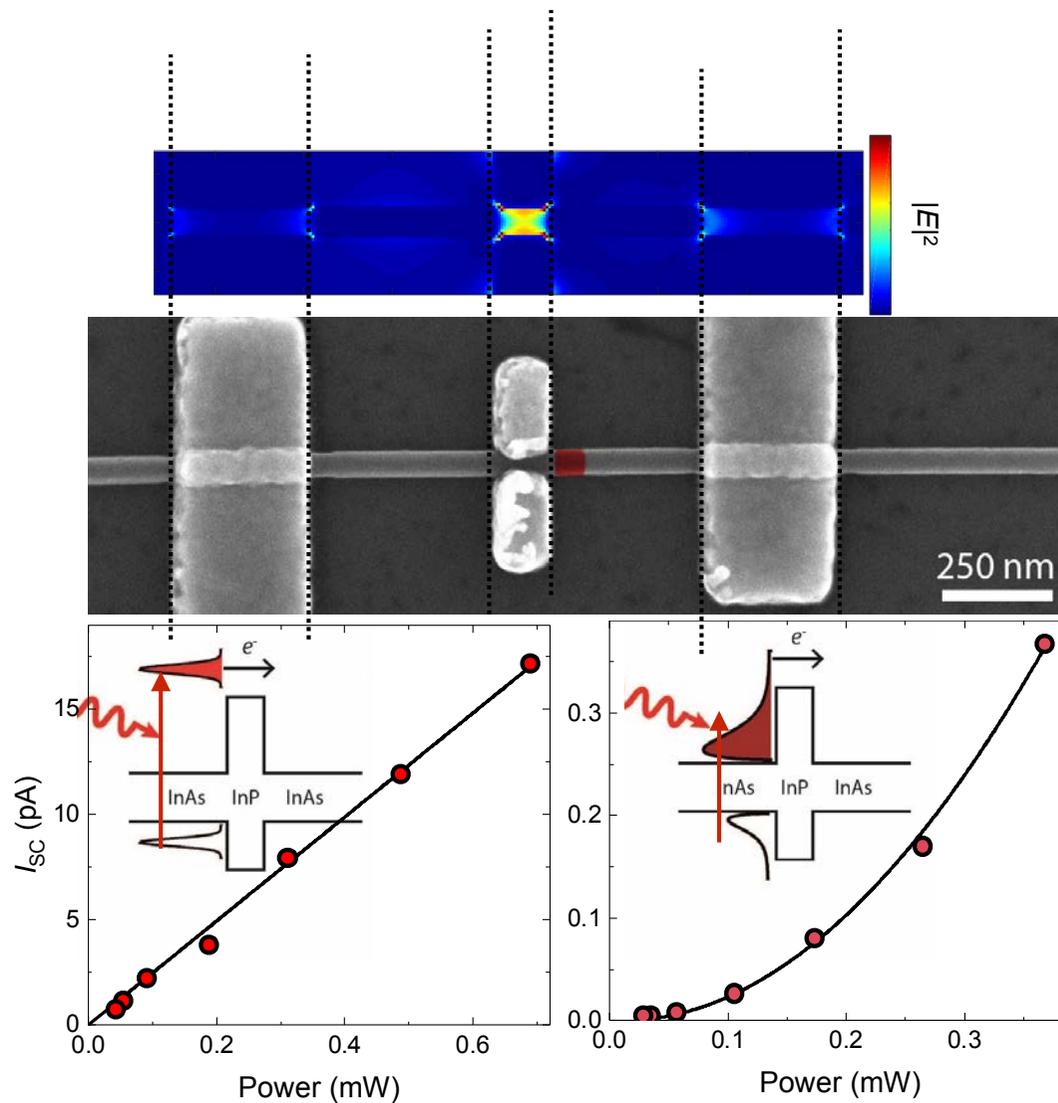


I-Ju Chen



I-Ju Chen et al.  
in preparation

# Evidence of quasi-ballistic extraction of hot carriers



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I-Ju Chen et al.  
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Theory:

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Collaboration:

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# NANO LUND

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