

HIGH-DIMENSIONAL QUANTUM PHOTONICS: PUSHING THE LIMITS OF QUANTUM TECHNOLOGIES

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My Academic Journey



MY ACADEMIC JOURNEY



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BBQLAB.ORG BEYOND BORING QUBITS! Heriot-Watt University, Edinburgh









LECTURE 1

- Introduction / Qudits
- Superposition / Entanglement
- Measuring HD Entanglement
- Noise Resistance

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- Superposition / Entanglement
- Measuring HD Entanglement
- Noise Resistance

LECTURE 2

- Unscrambling Entanglement through a Complex Medium
 - Multi-photon HD entanglement
- Free-space Quantum Communication

 \bullet

• Technologies based on the **coherent** control of <u>many</u> quantum systems

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- Form the basis of modern society as we know it

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Transistor

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Transistor



Laser

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LIGO will be getting a quantum upgrade

Quantum 'squeezing' light could lead to daily gravitational wave detections



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Fast and Cheap ££



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Fast and Cheap ££

Low-noise

Efficient transport







Photons are the ideal carriers of <u>quantum information</u>



Classical photonics industry and research!

• Quantum light: <u>sources</u> and *detectors* reaching maturity

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Quantum optical intercor



h a global quantum network

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Quantum optical intercor







- Polarization: 2-level QM system (*qubit*)
 - Quantum optics workhorse



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 - D-level quantum system! (qudit)





 $|\Psi\rangle = |HG_{00}\rangle + |HG_{01}\rangle + |HG_{10}\rangle \dots + |HG_{mn}\rangle$

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- Limits: uncertainty principle + physical system



$$|\Psi\rangle = \frac{1}{\sqrt{D}} [|t_0\rangle + |t_1\rangle + |t_2\rangle \dots + |t_{D-1}\rangle]$$



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 $2^{14} \approx 100^2 \approx 20^3$

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Ecker, MM et al, PRX 9, 041042 (2019)

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 - ★ Challenge: <u>efficient</u> transport, manipulation, and measurement

 $2^{14} \approx 100^2 \approx 20^3$



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$$LG^p_{\ell}(\rho,\phi,z) = \frac{A}{w(z)} \left[\frac{\sqrt{2}\rho}{w(z)}\right]^{|\ell|} L^{\ell}_p \left[\frac{2\rho^2}{w(z)^2}\right] \exp\left[\frac{-\rho^2}{w(z)^2}\right] \exp\left[\frac{-ik\rho^2}{2R(z)^2} - i\varphi(z)\right] \exp[i\ell\phi]$$





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 $E(\rho, \phi, z) = A(\rho, z) \frac{\exp[i\ell\phi]}{\exp[i\ell\phi]}$

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Light beams with azimuthal phase dependence





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- The OAM state space is discrete and unbounded*

 $|\Psi\rangle = |-1\rangle + |0\rangle + |1\rangle + |2\rangle + \dots$



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• Helical phase profile on SLM



Generating Sf

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- Added to linear phase ramp (mod 2π)



GENERATING SF

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- "Forked" diffraction grating
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- Can also generate complex superpositions of OAM modes
- Reverse process used for detection








$$A(x, y) \to \iint |A|^2 \to \text{non-zero}$$

 $B(x, y) \to \iint BA^* \to \text{zero!}$



$$A(x, y) \to \iint |A|^2 \tilde{LG}_{00} \to \text{non-zero}$$
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SUPERPOSITION AND ENTANGLEMENT

- Schrödinger's cat metaphor
 - Cat put in box
 - Radioactive decay triggers poison
 - 50% chance decay takes place



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- Opening box constitutes a *measurement*



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- Schrödinger's cat metaphor
 - Cat put in box
 - Radioactive decay triggers poison
 - 50% chance decay takes place
- Superposition of probability amplitudes
- Opening box constitutes a *measurement*
- Measurement <u>collapses</u> state of cat to dead or alive



$$|\Psi
angle = rac{1}{\sqrt{2}}(|\textit{dead}
angle + |\textit{alive}
angle)$$























• Coherent superposition of two possible states – "Qubit"

 $|\psi\rangle = a |H\rangle + b |V\rangle$



$$|\psi\rangle = a |H\rangle + b |V\rangle$$

 $\rho = |\psi\rangle \langle \psi|$ = $|a|^2 |H\rangle \langle H| + |b|^2 |V\rangle \langle V|$ + $ab^* |H\rangle \langle V| + a^*b |V\rangle \langle H|$



$$|\psi\rangle = a |H\rangle + b |V\rangle$$

$$o = |\psi\rangle \langle \psi|$$

= $|a|^2 |H\rangle \langle H| + |b|^2 |V\rangle \langle V| \rightarrow D$
+ $ab^* |H\rangle \langle V| + a^*b |V\rangle \langle H| \rightarrow O$

→ Off-diagonals



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$$\langle H | \rho | H \rangle \rightarrow |a|^2$$



$$|\psi\rangle = a |H\rangle + b |V\rangle$$

$$\rho = |\psi\rangle \langle \psi|$$

= $|a|^2 |H\rangle \langle H| + |b|^2 |V\rangle \langle V| \rightarrow Diagonals$
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$$\begin{split} \rho &= |\psi\rangle \langle \psi| \\ &= |a|^2 |H\rangle \langle H| + |b|^2 |V\rangle \langle V| & \rightarrow \text{Diagonals} \\ &+ ab^* |H\rangle \langle V| + a^* b |V\rangle \langle H| & \rightarrow \text{Off-diagonals} \end{split}$$

$$\langle H | \rho | H \rangle \rightarrow |a|^2 \quad \langle H | \rho | V \rangle = ab^* \qquad \rho \xrightarrow{\text{PBS}} \langle H | \rho | H \rangle$$



$$|\psi\rangle = a |H\rangle + b |V\rangle$$

$$p = |\psi\rangle \langle \psi|$$

= $|a|^2 |H\rangle \langle H| + |b|^2 |V\rangle \langle V| \rightarrow Diagonals$
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$$\langle H | \rho | H \rangle \rightarrow |a|^{2} \quad \langle H | \rho | V \rangle = ab^{*}$$

$$Re(\langle H | \rho | V \rangle) = \frac{1}{2}(\langle D | \rho | D \rangle - \langle A | \rho | A \rangle)$$

$$Im(\langle H | \rho | V \rangle) = \frac{1}{2}(\langle L | \rho | L \rangle - \langle R | \rho | R \rangle)$$

$$PBS$$

$$\langle H | \rho | H \rangle$$

$$\langle V | \rho | V \rangle$$



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$$HWP PBS \\ \langle D | \rho | D \rangle$$

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$$"Quantum State Tomography"$$



• Conjoined kittens born in Trieste!



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- "Magical" connection: if one cat is found to be <u>alive</u>, the other must be **alive**!





Trieste

$$|\Psi\rangle = |alive\rangle_1 |alive\rangle_2 + |dead\rangle_1 |dead\rangle_2$$

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 - Separated by skilled Italian veterinarians at birth
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- Cats are in a *superposition* of being dead and alive
- "Magical" connection: if one cat is found to be <u>alive</u>, the other must be **alive**!
- Not a causal connection





Trieste

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• Magical connection exists no matter when (how) you open the box!



$$|\Psi\rangle = |alive\rangle_1 |alive\rangle_2 + |dead\rangle_1 |dead\rangle_2$$



 $|\Psi\rangle = |zombie\rangle_1 |zombie\rangle_2 + |vampire\rangle_1 |vampire\rangle_2$

• Two photons can be entangled in their polarisation

$$|\Psi\rangle = \frac{1}{\sqrt{2}}[|H\rangle_1|H\rangle_2 + |V\rangle_1|V\rangle_2]$$

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- Measure in the *HV basis*, polarisations will be **perfectly correlated**

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- Created using type II nonlinear crystals

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e₂

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- Measurements in the **DA basis** also show perfect correlations!
- Created using type II nonlinear crystals
- One blue photon is absorbed
- Two *red* photons are emitted





$$|\Psi\rangle = \frac{1}{\sqrt{2}}[|H\rangle_1|H\rangle_2 + |V\rangle_1|V\rangle_2]$$

$$\begin{aligned} |H\rangle &= |D\rangle + |A\rangle \\ |V\rangle &= |D\rangle - |A\rangle \end{aligned}$$

$$|\Psi\rangle = \frac{1}{\sqrt{2}}[|A\rangle_1|A\rangle_2 + |D\rangle_1|D\rangle_2]$$

ENTANGLEMENT: APPLICATIONS

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Measurement-based Quantum Computation Quantum Error Correction

> Raussendorf & Briegel, Phys. Rev. Lett. 86, 5188 (2001) Lidar & Brun, ed., "Quantum Error Correction," Cambridge University Press (2013)

ENTANGLEMENT: APPLICATIONS



Measurement-based Quantum Computation Quantum Error Correction



Device-Independent Quantum Cryptography

Raussendorf & Briegel, Phys. Rev. Lett. 86, 5188 (2001) Lidar & Brun, ed., "Quantum Error Correction," Cambridge University Press (2013) Acín, Brunner, Gisin, Massar, Pironio, Scarani, Phys. Rev. Lett. 98, 230501 (2007)

• Orbital angular momentum entanglement

$$|\Psi\rangle = |-1,1\rangle + |0,0\rangle + |1,-1\rangle + |2,-2\rangle + |3,-3\rangle + \dots$$

Mair, Vaziri, Weihs, Zeilinger, Nature 412, 313 (2001)



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• Time-bin entanglement

$$\Psi \rangle = |t_0, t_0\rangle + |t_1, t_1\rangle + |t_2, t_2\rangle + |t_3, t_3\rangle + \dots$$

Brendel, Gisin, Tittel, Zbinden, Phys. Rev. Lett. 82, 2594 (1999)

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MEASURING HIGH-DIMENSIONAL ENTANGLEMENT

MEASURE OF ENTANGLEMENT DIMENSIONALITY

SCHMIDT RANK, k

Minimum numbers of levels needed to represent a state and its correlations in any local basis

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Minimum numbers of levels needed to represent a state and its correlations in any local basis



$$\rho_{1d} = \frac{1}{3} (|00\rangle \langle 00| + |11\rangle \langle 11| + |22\rangle \langle 22|)$$

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$$\left|\psi_{2d}^{01}\right\rangle = \frac{1}{\sqrt{2}}(\left|00\right\rangle + \left|11\right\rangle), \ \left|\psi_{2d}^{02}\right\rangle = \frac{1}{\sqrt{2}}(\left|00\right\rangle + \left|22\right\rangle), \ \left|\psi_{2d}^{12}\right\rangle = \frac{1}{\sqrt{2}}(\left|11\right\rangle + \left|22\right\rangle)$$

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$$\rho_{2d} = \frac{1}{3} \left(\left| \psi_{2d}^{01} \right\rangle \left\langle \psi_{2d}^{01} \right| + \left| \psi_{2d}^{02} \right\rangle \left\langle \psi_{2d}^{02} \right| + \left| \psi_{2d}^{03} \right\rangle \left\langle \psi_{2d}^{12} \right| \right) \qquad k = 2$$

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$$|\psi_{3d}\rangle = \frac{1}{\sqrt{3}}(|00\rangle + |11\rangle + |22\rangle)$$
ENTANGLEMENT VS CLASSICAL CORRELATIONS

$$\rho_{1d} = \frac{1}{3} (|00\rangle \langle 00| + |11\rangle \langle 11| + |22\rangle \langle 22|)$$

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Krenn, MM, Erhard & Zeilinger, Phil Trans R Soc A 375 (2017)

PHYSICAL REVIEW A 84, 062101 (2011)

Tomography of the quantum state of photons entangled in high dimensions

Megan Agnew,¹ Jonathan Leach,¹ Melanie McLaren,² F. Stef Roux,² and Robert W. Boyd^{1,3} ¹Department of Physics, University of Ottawa, 150 Louis Pasteur, Ottawa, Ontario, K1N 6N5 Canada ²CSIR National Laser Centre, Pretoria 0001, South Africa ³Institute of Optics, University of Rochester, Rochester, New York 14627, USA (Received 28 September 2011; published 2 December 2011)



★ Full state tomography → $k=8 \times 8$, scales as $d^2(d+1)^2$, time-consuming, unreliable

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PUBLISHED ONLINE: 8 MAY 2011 | DOI: 10.1038/NPHYS199

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nature physics

Experimental high-dimensional two-photon entanglement and violations of generalized Bell inequalities

Adetunmise C. Dada¹*, Jonathan Leach², Gerald S. Buller¹, Miles J. Padgett² and Erika Andersson¹

★ Bell-type (CGLMP) inequality \rightarrow k=11 x 11, strict criterion, assumptions on state



Generation and confirmation of a (100×100) -dimensional entangled quantum system

Mario Krenn^{a,b,1}, Marcus Huber^{c,d,e}, Robert Fickler^{a,b}, Radek Lapkiewicz^{a,b}, Sven Ramelow^{a,b,2}, and Anton Zeilinger^{a,b,1}

^aInstitute for Quantum Optics and Quantum Information, Austrian Academy of Sciences, A-1090 Vienna, Austria; ^bVienna Center for Quantum Science and Technology, Faculty of Physics, University of Vienna, A-1090 Vienna, Austria; ^cDepartment of Mathematics, University of Bristol, Bristol BS8 1TW, United Kingdom; ^dInstitut de Ciencies Fotoniques, E-08860 Castelldefels (Barcelona), Spain; and ^eFisica Teorica: Informacio i Fenomens Quantics, Departament de Fisica, Universitat Autonoma de Barcelona, E-08193 Bellaterra (Barcelona), Spain

Contributed by Anton Zeilinger, February 24, 2014 (sent for review December 15, 2013)



★ Subspace Visibility Witness → $k=103 \times 103$, $d^2(d+1)$, assumed conservation of OAM



PUBLISHED ONLINE: 29 FEBRUARY 2016 | DOI: 10.1038/NPHOTON.2016.12

nature photonics

Multi-photon entanglement in high dimensions

Mehul Malik^{1,2*}, Manuel Erhard^{1,2}, Marcus Huber^{3,4,5}, Mario Krenn^{1,2}, Robert Fickler^{1,2†} and Anton Zeilinger^{1,2}

★ 2D Fidelity Witness → $k=3 \times 3 \times 2$, scales as $d^2(d+1)$



Generation and confirmation of a (100×100) -dimensional entangled quantum system

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Multi-photon entanglement in high dimensions

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★ 2D Fidelity Witness → $k=3 \times 3 \times 2$, scales as $d^2(d+1)$

Assumption-free, efficient entanglement witness? \rightarrow QKD, Qinfo



$$|D\rangle = \frac{1}{\sqrt{2}}[|H\rangle + |V\rangle] \qquad |A\rangle = \frac{1}{\sqrt{2}}[|H\rangle - |V\rangle]$$

$$|L\rangle = \frac{1}{\sqrt{2}}[|H\rangle + i|V\rangle] \qquad |R\rangle = \frac{1}{\sqrt{2}}[|H\rangle - i|V\rangle]$$

N = 3





N = d+1

$$|D\rangle = \frac{1}{\sqrt{2}}[|H\rangle + |V\rangle] \qquad |A\rangle = \frac{1}{\sqrt{2}}[|H\rangle - |V\rangle]$$

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$$|L\rangle = \frac{1}{\sqrt{2}}[|H\rangle + i|V\rangle] \qquad |R\rangle = \frac{1}{\sqrt{2}}[|H\rangle - i|V\rangle]$$

$$|j\rangle_k = \frac{1}{\sqrt{d}} \sum_{m=0}^{d-1} \omega^{jm+km^2} |m\rangle , \quad \omega = e^{2\pi i/d}$$



Bavaresco, Herrera, Klöckl, Pivoluska, Erker, Friis, MM, Huber, Nat. Phys. 14, 1032 (2018)



Measure two-photon counts in Schmidt (OAM) basis to obtain diagonal density matrix elements:

$$\langle mn | \rho | mn \rangle = \frac{N_{mn}}{\sum_{i,j} N_{ij}}$$

Bavaresco, Herrera, Klöckl, Pivoluska, Erker, Friis, MM, Huber, Nat. Phys. 14, 1032 (2018)



















$$F(\rho, \Phi) = \operatorname{Tr}(|\Phi \rangle \langle \Phi | \rho)$$





$$F(\rho, \Phi) = \operatorname{Tr}(|\Phi X \Phi| \rho)$$

For a state of rank $k \le d$, fidelity to target state is upper-bounded by:

$$F(\rho, \Phi) \le B_k(\Phi)$$
$$B_k(\Phi) := \sum_{m=0}^{k-1} \lambda_m^2$$



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Eg: If state fidelity is greater than

$$F(\rho, \Phi) > B_3(\Phi) := \lambda_0^2 + \lambda_1^2 + \lambda_2^2$$

Fickler et al, Nat. Commun. 5:4502 (2014)

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Entanglement dimensionality = **4**

Fickler et al, Nat. Commun. 5:4502 (2014)

$$F(\rho, \Phi) = \operatorname{Tr} (|\Phi \rangle \langle \Phi | \rho)$$
$$= \sum_{m,n=0}^{d-1} \lambda_m \lambda_n \langle mm | \rho | nn \rangle$$

$$F(\rho, \Phi) = \operatorname{Tr} (|\Phi| \langle \Phi| \rho)$$

= $\sum_{m,n=0}^{d-1} \lambda_m \lambda_n \langle mm| \rho | nn \rangle$
= $\sum_m \lambda_m^2 \langle mm| \rho | mm \rangle + \sum_{m,n \neq m} \lambda_m \lambda_n \langle mm| \rho | nn \rangle$

$$=F_1(\rho,\Phi)+F_2(\rho,\Phi)$$

$$F(\rho, \Phi) = \operatorname{Tr} (|\Phi \rangle \langle \Phi | \rho)$$

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$$= F_1(\rho, \Phi) + F_2(\rho, \Phi)$$
DIAGONAL ELEMENTS
$$f(\rho, \Phi) = F_1(\rho, \Phi) + F_2(\rho, \Phi)$$







$$\lambda_{m} = \sqrt{\frac{\langle mm | \rho | mm \rangle}{\sum_{n} \langle nn | \rho | nn \rangle}}$$





$$\lambda_{m} = \sqrt{\frac{\langle mm | \rho | mm \rangle}{\sum_{n} \langle nn | \rho | nn \rangle}}$$

$$\omega = e^{2\pi i/d}$$

Construct tilted basis:

$$\left|\tilde{j}\right\rangle = \frac{1}{\sqrt{\sum_{n}\lambda_{n}}} \sum_{m=0}^{d-1} \omega^{jm} \sqrt{\lambda_{m}} \left|m\right\rangle$$





$$\lambda_{m} = \sqrt{\frac{\langle mm | \rho | mm \rangle}{\sum_{n} \langle nn | \rho | nn \rangle}}$$

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Wootters and Fields, Ann. of Phys. 191, 363 (1989)



Measure two-photon counts in tilted basis to obtain:

$$\left< \tilde{j}\tilde{j}^* \middle| \rho \left| \tilde{j}\tilde{j}^* \right> = \frac{\tilde{N}_{jj}}{\sum_{i,j}\tilde{N}_{ij}} c_{\lambda} \right.$$



Use measurements in two bases to calculate fidelity bound:

$$F(\rho, \Phi) \geq \sum_{m} \lambda_{m}^{2} \langle mm | \rho | mm \rangle + \frac{\left(\sum_{m} \lambda_{m}\right)^{2}}{d} \sum_{j=0}^{d-1} \left\langle \tilde{j}\tilde{j}^{*} | \rho | \tilde{j}\tilde{j}^{*} \right\rangle$$

$$-\sum_{m,n=0}^{d-1} \lambda_{m} \lambda_{n} \langle mn | \rho | mn \rangle - \sum_{\substack{m\neq m', m\neq n \\ n\neq n', n'\neq m'}} \tilde{\gamma}_{mm'nn'} \sqrt{\langle m'n' | \rho | m'n' \rangle \langle mn | \rho | mn \rangle}$$

$$\tilde{\gamma}_{mm'nn'} = \begin{cases} 0 \text{ if } (m-m'-n-n') \mod d \neq 0 \\ \sqrt{\lambda_{m}\lambda_{m'}\lambda_{n}\lambda_{n'}} \text{ otherwise.} \end{cases}$$



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=
$$\left\{ \sqrt{\lambda_m \lambda_{m'} \lambda_n \lambda_{n'}} \text{ otherwise.} \right\}$$



Calculate entanglement dimensionality from target state bound:

$$F(\rho, \Phi) \le B_k(\Phi) := \sum_{m=0}^{k-1} \lambda_m^2$$

 \mathbf{O}

EXPERIMENTAL RESULTS







d	$d_{_{\mathrm{ent}}}$	$\widetilde{\boldsymbol{F}}(ho, \Phi^{+})$	$\widetilde{\pmb{F}}(ho,\Phi)$
3	3	91.5 <u>+</u> 0.4%	92.5 <u>+</u> 0.4%
5	5	89.9 <u>+</u> 0.4%	90.0±0.5%
7	6	84.2±0.5%	86.9±0.6%
11	9	74.8±0.4%	76.2 <u>+</u> 0.6%

Bavaresco, Herrera, Klöckl, Pivoluska, Erker, Friis, MM, Huber, Nat. Phys. 14, 1032 (2018)




Bavaresco, Herrera, Klöckl, Pivoluska, Erker, Friis, MM, Huber, Nat. Phys. 14, 1032 (2018)







Highest entanglement dimension certified without any assumptions on the state!

Bavaresco, Herrera, Klöckl, Pivoluska, Erker, Friis, MM, Huber, Nat. Phys. 14, 1032 (2018)



ADAPTIVE WITNESS: COMPARISON 1-outcome *d* = 11 TOMOGRAPHY $d^{2}(d+1)^{2}$ 17,424 $(d+1)^2$ 144 **2D WITNESS** 12 $d^{2}(d+1)$ 1452 d+1 **ADAPTIVE** 2 $2d^2$ 242 2 **WITNESS**













































Fidelity > 94.1%, <u>18-dim</u> entanglement in a 19-dim space



TWO APPLICATIONS OF HIGH-DIM ENTANGLEMENT

NOISE-RESISTANT ENTANGLEMENT DISTRIBUTION



REALISTIC QUANTUM NETWORKS

Imperfect Source



Noisy Channel 💊



Imperfect Detector



2D-Entanglement + Noise = \bigotimes



$$|\Psi\rangle = \frac{1}{\sqrt{2}} [|H\rangle_1 |H\rangle_2 + |V\rangle_1 |V\rangle_2]$$

$$F(\rho,\Psi) \geq 0.92$$

2D-Entanglement + Noise = \bigotimes



H₁

 V_1



 A_2

D₂

 D_1

 A_1

 V_2

 H_2

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$$F(\rho,\Psi) \geq 0.92$$

$$F(\rho,\Psi) \geq 0.65$$

$$F(\rho, \Psi) \ge 0.35$$

PATHWAYS TO NOISE RESISTANCE



S. Ecker, F. Bouchard, et al, arXiv:1904.01552 (2019)

PATHWAYS TO NOISE RESISTANCE



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PATHWAYS TO NOISE RESISTANCE



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$$\left|\Phi^{+}\right\rangle = \frac{1}{\sqrt{d}} \sum_{m=0}^{d-1} \left|mm\right\rangle$$

$$\rho_{\text{noisy}} = p \left| \Phi^+ X \Phi^+ \right| + \frac{1-p}{d^2} \mathbb{1}$$

$$\left|\Phi^{+}\right\rangle = \frac{1}{\sqrt{d}} \sum_{m=0}^{d-1} \left|mm\right\rangle$$

$$\rho_{\text{noisy}} = p \left| \Phi^+ \chi \Phi^+ \right| + \frac{1-p}{d^2} \mathbb{1}$$
White Noise

Detector Dark Counts Background Light Loss-induced Accidentals

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White Noise

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Bavaresco, Herrera, Klöckl, Pivoluska, Erker, Friis, MM, Huber, Nat. Phys. 14, 1032 (2018)

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Detector Dark Counts Background Light Loss-induced Accidentals

Eg. - For a **7-dim** state with noise level p = 0.43,

We can certify **no** entanglement with 1 MUB, but d = 4 with 7 MUBs!





























K-DIM ENTANGLEMENT IN D-DIM HILBERT SPACE



K-DIM ENTANGLEMENT IN D-DIM HILBERT SPACE



$$Q_{\text{opt}} = 3k + 2\sqrt{2}\sqrt{(k-2)(k-1)} - 4$$

$$d_{\rm opt} = \sqrt{2}\sqrt{k^2 - 3k + 2} + k - 1$$

K-DIM ENTANGLEMENT IN D-DIM HILBERT SPACE





32 x 32 pixel SPAD camera

$$Q_{\text{opt}} = 3k + 2\sqrt{2}\sqrt{(k-2)(k-1)} - 4$$

$$d_{\rm opt} = \sqrt{2} \sqrt{k^2 - 3k + 2} + k - 1$$
K-DIM ENTANGLEMENT IN D-DIM HILBERT SPACE





32 x 32 pixel SPAD camera 49 x 49 pixels \rightarrow 343 times lower SNR!

$$Q_{\text{opt}} = 3k + 2\sqrt{2}\sqrt{(k-2)(k-1)} - 4$$

$$d_{\rm opt} = \sqrt{2} \sqrt{k^2 - 3k + 2} + k - 1$$

***** 2 > * * * * 가 가 가 차 * * Y: 🥵 🐝 🐝 🎇 🎼 🎲 🔆 🎋 🏂 🤧 🔆 泠 🧌 🐝 🔆 🔆 🔆 🎋 * * * * * **\$** *** *** ** 兆 🎋 🐇

LECTURE 1

- Introduction / Qudits
- Superposition / Entanglement
- Measuring HD Entanglement
- Noise Resistance

LECTURE 2

- Unscrambling Entanglement through a Complex Medium
 - Multi-photon HD entanglement
- Free-space Quantum Communication

 \bullet

UNSCRAMBLING ENTANGLEMENT THROUGH A COMPLEX MEDIUM

REALISTIC QUANTUM NETWORKS

Imperfect Source



Noisy Channel 💊



Imperfect Detector



REALISTIC QUANTUM NETWORKS

Imperfect Source







Imperfect Detector



COMPLEX SCATTERING MEDIA

COMPLEX SCATTERING MEDIA









CLASSICAL COMMS: SDM

nature photonics

ARTICLES

PUBLISHED ONLINE: 13 JULY 2015 | DOI: 10.1038/NPHOTON.2015.112

Seeing through chaos in multimode fibres

Martin Plöschner¹, Tomáš Tyc² and Tomáš Čižmár^{1*}

¹School of Engineering, Physics and Mathematics, College of Art, Science & Engineering, University of Dundee, Nethergate, Dundee DD1 4HN, UK. ²Department of Theoretical Physics and Astrophysics, Masaryk University, Kotlarska 2, 61137 Brno, Czech Republic. *e-mail: t.cizmar@dundee.ac.uk





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Spatial Light Modulator (SLM)





Spatial Light Modulator (SLM)







Spatial Light Modulator (SLM)







Spatial Light Modulator (SLM)





$$\left| \Phi^{+} \right\rangle = \frac{1}{\sqrt{d}} \sum_{i=1}^{d} \left| ii \right\rangle$$

● Statements about **state** ⇔ Statements about <u>channel</u>



$$\left| \Phi^{+} \right\rangle = \frac{1}{\sqrt{d}} \sum_{i=1}^{d} \left| ii \right\rangle$$

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$$\left| \Phi^{+} \right\rangle = \frac{1}{\sqrt{d}} \sum_{i=1}^{d} \left| ii \right\rangle$$

Statements about state ⇔ Statements about <u>channel</u>



- Statements about **state** ⇔ Statements about <u>*channel*</u>
- Choi-Jamiolkowski isomorphism Φ^+ $\gamma = \frac{1}{\sqrt{2}} \sum |ii\rangle$ $\rho = (\mathbb{I} \otimes \mathbf{T}) \left| \Phi^+ \right\rangle \left\langle \Phi^+ \right|$ Φ^+

M.D. Choi, Lin. Alg. Appl. 10, 285 (1975) A. Jamiołkowski, Rep. Math. Phys. 3, 275 (1972)

T MEASUREMENT: QUANTUM





T MEASUREMENT: QUANTUM





T MEASUREMENT: QUANTUM





$$\begin{split} |\Phi\rangle_{\rm T} = (\mathbb{I} \otimes {\rm T}) \left| \Phi^+ \right\rangle \\ \propto \sum t_{ji} \left| ij \right\rangle \end{split}$$





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$$\left|\Phi\right\rangle_{\mathrm{T}} = \left(\mathbb{I}\otimes\mathrm{T}^{-1}\mathrm{T}\right)\left|\Phi^{+}\right\rangle = \left|\Phi^{+}\right\rangle$$





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$$(A \otimes B) \left| \Phi^+ \right\rangle = (B^T A \otimes I) \left| \Phi^+ \right\rangle = (I \otimes B A^T) \left| \Phi^+ \right\rangle$$





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INITIAL STATE



Valencia, Goel, McCutcheon, Defienne, Malik, arXiv:1910.04490 (2019)

INITIAL STATE



Valencia, Goel, McCutcheon, Defienne, Malik, arXiv:1910.04490 (2019)

CHOI STATE: AFTER FIBRE

a)

$$|\Phi\rangle_{T} = \sum_{ij} t_{ji} |ij\rangle$$
b)

$$\int \Phi_{ij} = \sum_{ij} t_{ji} |ij\rangle$$



-5Se

CHOI STATE: AFTER FIBRE

$$|\Phi\rangle_{T} = \sum_{ij} t_{ji} |ij\rangle$$

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COMPLEX T MEASUREMENT: PHASE-STEPPING



COMPLEX T MEASUREMENT: PHASE-STEPPING



COMPLEX T MEASUREMENT: PHASE-STEPPING



Reference Measurement



Reference Measurement



TRANSMISSION MATRIX

$$\mathbf{T} = (\mathbf{S}\mathbf{E}^{-1})^{\dagger}$$



"Scrambling" Entanglement



"Scrambling" Entanglement



MUTUALLY UNBIASED TS



Pixel Basis

MUTUALLY UNBIASED TS





d+1 = 8 MUBs



MUTUALLY UNBIASED TS



Pixel Basis

d+1 = 8 MUBs



$$\mathbf{T}_{\mathbf{M}i} = \mathbf{M}_i^* \, \mathbf{T} \, \mathbf{M}_i^{\mathrm{T}}$$

Valencia, Goel, McCutcheon, Defienne, Malik, arXiv:1910.04490 (2019)













5

3

 $| ilde{g}_v^5
angle$

1



Valencia, Goel, McCutcheon, Defienne, Malik, arXiv:1910.04490 (2019)

0.02

7



Valencia, Goel, McCutcheon, Defienne, Malik, arXiv:1910.04490 (2019)



6-dimensional entanglement!



Valencia, Goel, McCutcheon, Defienne, Malik, arXiv:1910.04490 (2019)

THREE QUDITS

• 1935: Einstein, Podolsky, Rosen







Einstein, Podolsky, Rosen, Phys. Rev. 47, 777 (1935)

• 1935: Einstein, Podolsky, Rosen





• 1964: John Bell

Physics Vol. 1, No. 3, pp. 195-200, 1964 Physics Publishing Co. Printed in the United States

ON THE EINSTEIN PODOLSKY ROSEN PARADOX*

J. S. BELL[†] Department of Physics, University of Wisconsin, Madison, Wisconsin

(Received 4 November 1964)



• 1935: Einstein, Podolsky, Rosen



$$|\Psi\rangle_{AB} = \frac{1}{\sqrt{2}}[|H\rangle_A|H\rangle_B + |V\rangle_A|V\rangle_B]$$

- 1964: John Bell
 - * <u>Statistical</u> test of QM



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- 1964: John Bell
 - * <u>Statistical</u> test of QM
- 1989: Greenberger, Horne, Zeilinger

GOING BEYOND BELL'S THEOREM

Daniel M. Greenberger¹, Michael A. Horne², and Anton Zeilinger³. ¹City College of the City University of New York, New York, New York ²Stonehill College. North Easton, Massachusetts ³Atominstitut der Oesterreichischen Universitaeten, Wien, Austria



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- Mermin parameter, M = XXX XYY YXY YYX



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$$\xrightarrow{0} Classical 2 \xrightarrow{2} M$$



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- 12 photons: USTC China
 - * Zhong et al, Phys. Rev. Lett.
 121, 250505 (2018)



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- 5 in 20 ions: IQOQI Innsbruck
 - Friis et al, Phys. Rev. X 8, 021012 (2018)

- 12 photons: USTC China
 - Zhong et al, Phys. Rev. Lett.
 121, 250505 (2018)



- Pulsed laser
 - 5 in 20 ions: IQOQI Innsbruck
 - Friis et al, Phys. Rev. X 8, 021012 (2018)

- 20 SQUIDs: ICQI China
 - Song et al, Science 365, 574 (2019)





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HIGH-DIMENSIONAL MULTIPARTITE ENTANGLEMENT?

HIGH-DIMENSIONAL MULTIPARTITE ENTANGLEMENT?

 $|\Psi\rangle_{222} = \frac{1}{\sqrt{2}}[|000\rangle + |111\rangle]$


$|\Psi\rangle_{333} = \frac{1}{\sqrt{3}}[|000\rangle + |111\rangle + |222\rangle]$



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$$|\Psi\rangle_{332} = \frac{1}{\sqrt{3}}[|000\rangle + |111\rangle + |221\rangle]$$





• <u>Asymmetric</u> states: only possible when *N* and *D* are both greater than 2

 $|\Psi\rangle_{333} = \frac{1}{\sqrt{3}}[|000\rangle + |111\rangle + |222\rangle]$



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- <u>Asymmetric</u> states: only possible when *N* and *D* are both greater than 2
- Vast family of multi-particle entangled states!

 $|\Psi\rangle_{333} = \frac{1}{\sqrt{3}}[|000\rangle + |111\rangle + |222\rangle]$





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- <u>Asymmetric</u> states: only possible when *N* and *D* are both greater than 2
- Vast family of multi-particle entangled states!
- Creation of 332 and 333 entanglement

VOLUME 78

21 APRIL 1997

NUMBER 16

Three-Particle Entanglements from Two Entangled Pairs

Anton Zeilinger,¹ Michael A. Horne,² Harald Weinfurter,¹ and Marek Żukowski^{1,3}
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²Stonehill College, North Easton, Massachusetts 02357
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Based on <u>erasing</u> "which source" information

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- Based on <u>erasing</u> "which source" information
- Entangled photons from **two** crystals (two pairs)

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- Based on <u>erasing</u> "which source" information
- Entangled photons from **two** crystals (two pairs)
- *"which crystal"* information erased





• Polarizing beam splitter (PBS): reflects V, transmits H



- Polarizing beam splitter (PBS): reflects V, transmits H
- Mixes polarizations of two input photons









• Coincidences only obtained between same polarizations



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- Coincidences only obtained between same polarizations
- Coherent superposition four photon GHZ state!

$$|\Psi\rangle_{GHZ} = \frac{1}{\sqrt{2}} [|H\rangle_1 |H\rangle_2 |H\rangle_3 |H\rangle_4 + |V\rangle_1 |V\rangle_2 |V\rangle_3 |V\rangle_4]$$



- Coincidences only obtained between same polarizations
- Coherent superposition four photon GHZ state!

• Why polarisation? Tools exist!

• Why polarisation? Tools exist!



• Why polarisation? Tools exist!



• Why polarisation? Tools exist!



• Why polarisation? Tools exist!



- Why polarisation? Tools exist!
- Challenge: develop tools for manipulating photonic OAM





kist!

- Why polaris
- Challenge: p for manipulatin M
- Interferometer with rotated Dove prism in each arm



 $\alpha/2$

A1

kist!

- Why polaris
- Challenge: p for manipulatin M
- Interferometer with rotated Dove prism in each arm
- Introduces l-dependent phase



kist!

- Why polaris
- Challenge: p for manipulatin M
- Interferometer with rotated Dove prism in each arm
- Introduces l-dependent phase
- Sorts OAM into even and odd values



Input

kist!

- Why polaris
- Challenge: p for manipulatin M
- Interferometer with rotated Dove prism in each arm
- Introduces l-dependent phase
- Sorts OAM into even and odd values
- Scalable, but not practical



 $\alpha/2$

A1

• Interferometric sorter: <u>one-input</u>, *two-output* device




• Interferometric sorter: <u>one-input</u>, *two-output* device



INPUT

- Interferometric sorter: <u>one-input</u>, *two-output* device
- Can be used as a **two-input**, *two-output* device



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- Interferometric sorter: <u>one-input</u>, *two-output* device
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- Interferometric sorter: <u>one-input</u>, *two-output* device
- Can be used as a **two-input**, *two-output* device
- Works like a polarizing beam splitter: reflects ODD, transmits EVEN modes











Even modes transmitted!



Even modes transmitted!



Odd modes reflected!

Even modes transmitted!





Odd modes reflected!

Even modes transmitted!



Odd modes reflected!

Even modes transmitted!





























$$|\Psi\rangle_{332} = |1\rangle_B |1\rangle_C |-1\rangle_D + |0\rangle_B |0\rangle_C |0\rangle_D + |1\rangle_B |-1\rangle_C |1\rangle_D$$



• Superposition measurements of OAM

$$|\Psi\rangle_{332} = |1\rangle_B |1\rangle_C |-1\rangle_D + |0\rangle_B |0\rangle_C |0\rangle_D$$



- Superposition measurements of OAM
- Look for two-photon interference from independent crystals

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- Look for two-photon interference from independent crystals

$$|\Psi\rangle_{332} = |1\rangle_B |1\rangle_C |-1\rangle_D + |0\rangle_B |0\rangle_C |0\rangle_D$$



- Superposition measurements of OAM
- Look for two-photon interference from independent crystals
- Appearance of interference dip indicates coherent superposition:

$$|\Psi\rangle_{332} = |1\rangle_B |1\rangle_C |-1\rangle_D + |0\rangle_B |0\rangle_C |0\rangle_D + |1\rangle_B |-1\rangle_C |1\rangle_D$$

EXPERIMENTAL SETUP



Experimental Details

Experimental Details








EXPERIMENTAL DETAILS







• Full reconstruction of our (large) state is prohibitive

- Full reconstruction of our (large) state is prohibitive
- Entanglement Witness!



Entanglement Witness (W)

- Full reconstruction of our (large) state is prohibitive
- Entanglement Witness!
- Prove: measured state cannot be decomposed into <u>entangled states</u> of *smaller dimensionality structure* (e.g. - 322, 222, 22)



Best achievable fidelity of a <u>322 state</u> (*σ*) with an *ideal* <u>332 state</u> (Ψ):



Best achievable fidelity of a <u>322 state</u> (*σ*) with an *ideal* <u>332 state</u> (Ψ):

$$F_{max} := \max_{\sigma \in (322)} \operatorname{Tr}(\sigma |\Psi\rangle \langle \Psi|) = 0.6\overline{6}$$



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Measured fidelity of <u>generated state</u> (*ρ_{exp}*) with an *ideal* <u>332 state</u> (Ψ)



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- Measured fidelity of <u>generated state</u> (*ρ_{exp}*) with an *ideal* <u>332 state</u> (Ψ)
- Density matrix elements: 18 diagonal, 3 off-diagonal



• Density matrix \rightarrow 171 elements

- Density matrix \rightarrow 171 elements
- Non-zero: 3 diagonal, 3 off-diagonal

$$\hat{D}_1 = |000\rangle\langle 000| \qquad \hat{A}_1 = |000\rangle\langle 1\bar{1}1|$$

$$\hat{D}_2 = |1\bar{1}1\rangle\langle 1\bar{1}1| \qquad \hat{A}_2 = |1\bar{1}1\rangle\langle 11\bar{1}|$$

$$\hat{D}_3 = |11\bar{1}\rangle\langle 11\bar{1}| \qquad \hat{A}_3 = |11\bar{1}\rangle\langle 000|$$

- Density matrix \rightarrow 171 elements
- Non-zero: 3 diagonal, 3 off-diagonal
- Lab: projective measurements only!

$$\hat{D}_{1} = |000\rangle\langle000| \qquad \hat{A}_{1} = |000\rangle\langle1\bar{1}1|$$
$$\hat{D}_{2} = |1\bar{1}1\rangle\langle1\bar{1}1| \qquad \hat{A}_{2} = |1\bar{1}1\rangle\langle11\bar{1}|$$
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$$\hat{A}_{1} = |000\rangle \langle 1\overline{1}1|$$
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- Density matrix \rightarrow 171 elements
- Non-zero: 3 diagonal, 3 off-diagonal
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- Each off-diagonal element requires **64** projective measurements:

$$|\Psi\rangle_{332} = |1\rangle_B |1\rangle_C |\bar{1}\rangle_D + |0\rangle_B |0\rangle_C |0\rangle_D + |1\rangle_B |\bar{1}\rangle_C |1\rangle_D$$

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Best achievable fidelity of a <u>322 state</u> (*σ*) with an *ideal* <u>332 state</u> (Ψ):



$$|\Psi\rangle_{332} = \frac{1}{\sqrt{3}}[|000\rangle + |111\rangle + |221\rangle]$$

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$$F_{max} := \max_{\sigma \in (322)} \operatorname{Tr}(\sigma |\Psi\rangle \langle \Psi|) = 0.6\overline{6}$$



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• *First* high-dimensional multipartite entangled state!


Layered Quantum Cryptography



Layered Quantum Cryptography



Pivoluska, Huber, and MM, PRA 97, 032312 (2018)

LAYERED QUANTUM CRYPTOGRAPHY



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THE QUEST FOR 3D GHZ

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THE QUEST FOR 3D GHZ



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Automated Search for new Quantum Experiments

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Automated Search for new Quantum Experiments

Mario Krenn,^{1,2,*} Mehul Malik,^{1,2} Robert Fickler,^{1,2,†} Radek Lapkiewicz,^{1,2,‡} and Anton Zeilinger^{1,2}



Krenn, MM, Fickler, Lapkiewicz, Zeilinger, Phys. Rev. Lett. 116, 090405 (2016)

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	2,2,2 3,2,	2 4,2,2	5,2,2	6,2,2	7,2,2	8,2,2	9,2,2	10,2,2	11,2,2	12,2,2
	3,3,	2 4,3,2	5,3,2	6,3,2	7,3,2	8,3,2	9,3,2	10,3,2	11,3,2	12,3,2
	3,3,	3 4,3,3	5,3,3	6,3,3	7,3,3	8,3,3	9,3,3	10,3,3	11,3,3	12,3,3
		4,4,2	5,4,2	6,4,2	7,4,2	8,4,2	9,4,2	10,4,2	11,4,2	12,4,2
		4,4,3	5,4,3	6,4,3	7,4,3	8,4,3	9,4,3	10,4,3	11,4,3	12,4,3
		4,4,4	5,4,4	6,4,4	7,4,4	8,4,4	9,4,4	10,4,4	11,4,4	12,4,4
			5,5,2	6,5,2	7,5,2	8,5,2	9,5,2	10,5,2	11,5,2	12,5,2
			5,5,3	6,5,3	7,5,3	8,5,3	9,5,3	10,5,3	11,5,3	12,5,3
			5,5,4	6,5,4	7,5,4	8,5,4	9,5,4	10,5,4	11,5,4	12,5,4
			5,5,5	6,5,5	7,5,5	8,5,5	9,5,5	10,5,5	11,5,5	12,5,5
				6,6,2	7,6,2	8,6,2	9,6,2	10,6,2	11,6,2	12,6,2
				6,0,5	7,0,5	8.6.4	9,0,5	10,6,5	11,0,5	12,0,5
				665	7,0,4	865	9,0,4	10,0,4	11,0,4	12,0,4
				666	766	866	966	10,0,5	11,0,5	12,0,5
xl ¹				0,0,0	7.7.2	8.7.2	9.7.2	10,7,2	11.7.2	12,7,2
					7,7,3	8,7,3	9,7,3	10.7.3	11.7.3	12,7,3
					7,7,4	8,7,4	9,7,4	10,7,4	11,7,4	12,7,4
					7,7,5	8,7,5	9,7,5	10,7,5	11,7,5	12,7,5
\sim					7,7,6	8,7,6	9,7,6	10,7,6	11,7,6	12,7,6
					7,7,7	8,7,7	9,7,7	10,7,7	11,7,7	12,7,7
						8,8,2	9,8,2	10,8,2	11,8,2	12,8,2
						8,8,3	9,8,3	10,8,3	11,8,3	12,8,3
						8,8,4	9,8,4	10,8,4	11,8,4	12,8,4
						8,8,5	9,8,5	10,8,5	11,8,5	12,8,5
						8,8,6	9,8,6	10,8,6	11,8,6	12,8,6
$\sim / \vee \setminus / $						8,8,7	9,8,7	10,8,7	11,8,7	12,8,7
						8,8,8	9,8,8	10,8,8	11,8,8	12,8,8
							9,9,2	10,9,2	11,9,2	12,9,2
							9,9,3	10,9,3	11,9,3	12,9,3
							9,9,4	10,9,4	11,9,4	12,9,4
ハ							9,9,5	10,9,5	11,9,5	12,9,5
5.5.5							9,9,6	10,9,6	11,9,6	12,9,0

Automated Search for new Quantum Experiments



Krenn, MM, Fickler, Lapkiewicz, Zeilinger, Phys. Rev. Lett. 116, 090405 (2016)



Erhard, MM, Krenn and Zeilinger, Nature Photonics 12, 759 (2018)



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Erhard, MM, Krenn and Zeilinger, Nature Photonics 12, 759 (20) (20) (20)





 $\rangle_{A} \rightarrow |2\rangle_{A}$

 $\ell = 0$









 $\ell = 0$









 $\rangle_{A} \rightarrow |2\rangle_{A}$



(|0)_B)





Relative path difference (µm)







 $\rangle_{A} \rightarrow |2\rangle_{A}$



(|0)_B)





Relative path difference (µm)



$$|0,0\rangle_{A,B}$$

1/ 2 ($|0,2\rangle_{A,B}$ - $|-2,0\rangle_{A,B}$

$$10_{A} + 12_{A} = 10_{B}$$

$$|1,-1\rangle_{AB}\otimes |-1,1\rangle_{CD}$$



 $|1,-1\rangle_{AB} \otimes |-1,1\rangle_{CD}$



 $|1,-1\rangle_{AB} \otimes |-1,1\rangle_{CD}$



Verifying 333 Entanglement

$$|\Psi\rangle_{333} = \frac{1}{\sqrt{3}}[|-3, -1, -1\rangle + |-2, 0, 0\rangle - |1, 1, 1\rangle]$$

VERIFYING 333 ENTANGLEMENT



• 27 diagonal, 3 off-diagonal elements measured

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Erhard, MM, Krenn and Zeilinger, Nature Photonics 12, 759 (2018)


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- Certified to be **3D-GHZ entangled** by <u>3 standard deviations</u>



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OAM OUTSIDE THE LAB



Krenn, Fickler, Fink, Handsteiner, MM, Scheidl, Ursin, Zeilinger, New J. Phys. 16, 113028 (2014)





Krenn, Fickler, Fink, Handsteiner, MM, Scheidl, Ursin, Zeilinger, New J. Phys. 16, 113028 (2014)



Krenn, Fickler, Fink, Handsteiner, Ursin, MM, Zeilinger, PNAS 113, 13648 (2016)

a) $\ell = |1\rangle - i |-1\rangle$ b) $\ell = |2\rangle + |-2\rangle$ c) $\ell = |3\rangle - |-3\rangle$ d) $\ell = |3\rangle + |-3\rangle$





Krenn, Fickler, Fink, Handsteiner, Ursin, MM, Zeilinger, PNAS 113, 13648 (2016)

• <u>Manipulation</u>: Efficient + General Local Unitaries

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- <u>Transport</u>: Free-space/Fibre
- <u>Transduction</u>: Quantum States of Matter
- <u>Theory</u>: Certification Strategies, Non-Locality Tests

BBQLAB.ORG BEYOND BORING QUBITS! Heriot-Watt University, Edinburgh







