



Tokyo Tech

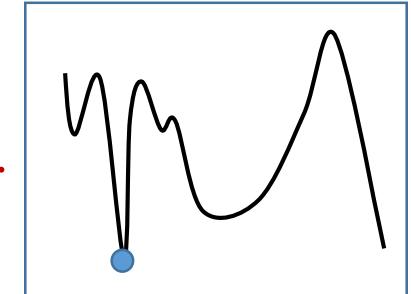
# Quantum simulation by quantum annealing

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# Quantum Annealing

## Original target:

To solve classical **combinatorial optimization problems**.



## Method:

Appropriate control of coefficients of the **transverse-field Ising model**

$$H(t) = - \sum_{i,j} J_{ij}(t) \sigma_i^z \sigma_j^z - \Gamma(t) \sum_i \sigma_i^x$$

## An additional target:

**Quantum simulation** (experiment of quantum systems)

# Recent examples of quantum simulation by quantum annealing

## Static (equilibrium) properties

- Spin glass in 3 dimensions      *Harris et al., Science (2018)*
- Kosterlitz-Thouless transition      *King et al., Nature (2018)*       To be reviewed
- Spin ice      *King et al., arXiv (2020)*
- Shastry-Sutherland model      *Kairys et al., arXiv (2020)*       To be reviewed
- Field theory      *Abel et al., arXiv (2020)*
- $Z_2$  lattice gauge theory      *Zhou et al., arXiv (2020)*
- Griffiths-McCoy singularity      *Nishimura et al., Phys. Rev. A (2020)*       To be explained

## Dynamical (non-equilibrium) properties

- Kibble-Zurek mechanism      *Gardas et al., Sci. Rep. (2018)*      *1d*
- Generalized Kibble-Zurek      *Weinberg et al., Phys. Rev. Lett. (2020)*      *2d*
- *Bando et al., Phys. Rev. Res. (2020)*      *1d*      
- To be explained

# Review (1)

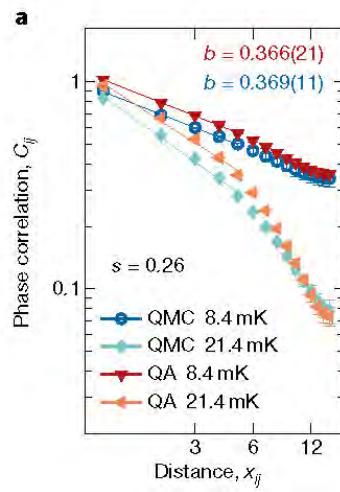
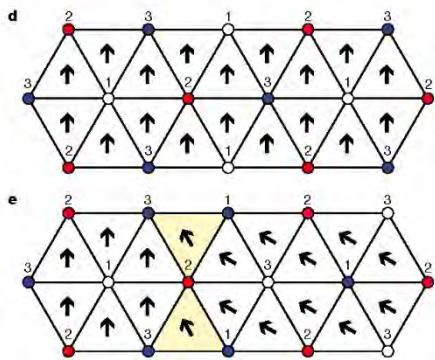
## Quantum simulation of the Kostertlitz-Thouless transition

A. D. King et al., *Nature* **560**, 456 (2018)

# Kosterlitz-Thouless transition

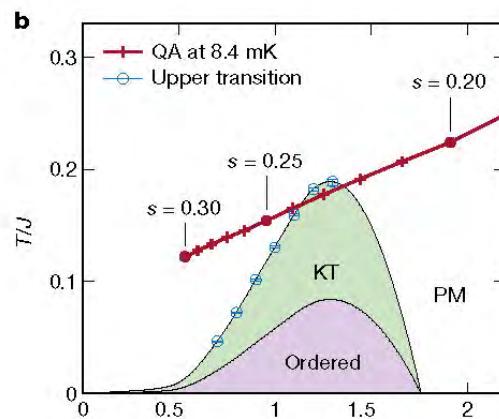
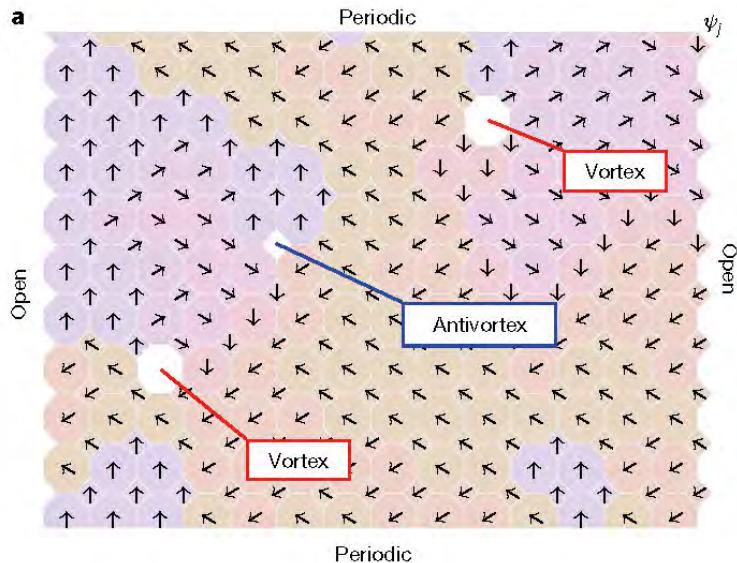
A. D. King et al., *Nature* 560, 456 (2018)

Representation of  $Z_6$  model by the frustrated Ising model embedded on the Chimera graph of D-Wave



Power decay of correlation, a hallmark of the KT transition

Observation of vortex-antivortex pairs in some parameter range



# Review (2)

## Quantum simulation of the Shastry-Sutherland model

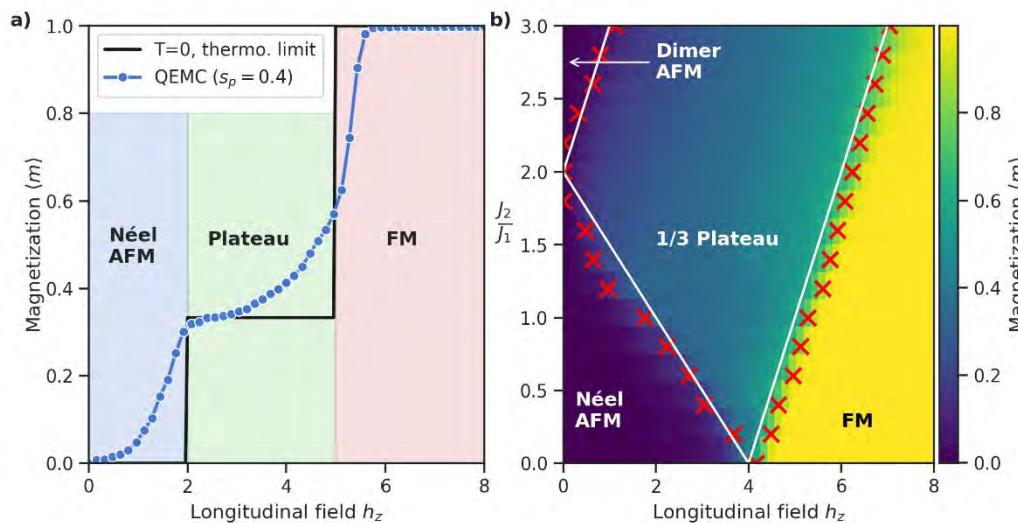
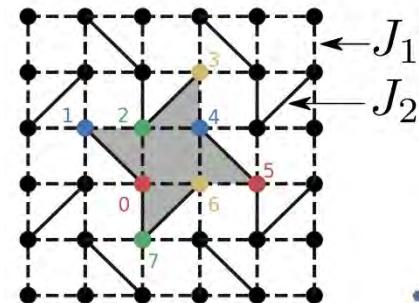
P. Kairys, A. D. King, I. Ozfidan, K. Boothby, J. Raymond, A. Banerjee, T. S. Humble,  
*arXiv:2003.01019*

# Shastry-Sutherland model

P. Kairys et al., arXiv:2003.01019

Ground-state phase diagram of a frustrated classical Ising model on the square lattice

$$H = J_1 \sum_{\langle i,j \rangle} \sigma_{(i)}^z \sigma_{(j)}^z + J_2 \sum_{\langle\langle i,j \rangle\rangle} \sigma_{(i)}^z \sigma_{(j)}^z + h_z \sum_i \sigma_{(i)}^z$$



Magnetization plateau observed.  
Phase diagram confirmed.

# Our contribution (1)

## Quantum simulation of the generalized Kibble-Zurek mechanism

Y. Bando, Y. Susa, H. Oshiyama, N. Shibata, M. Ohzeki, F. Gómez-Ruiz, D. A. Lidar,  
S. Suzuki, A. del Campo, and H. Nishimori, *Phys. Rev. Res.* **2**, 033369 (2020)

cf. Talk by Y. Bando in Session 3 for additional information

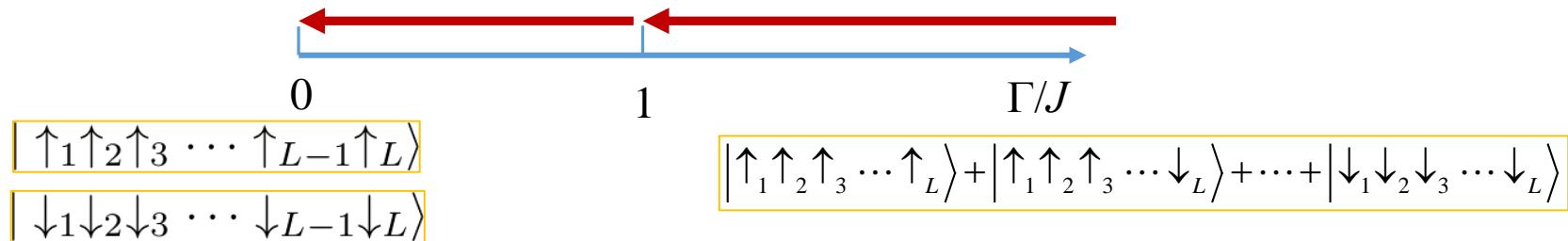
Supported by IARPA QEO / DARPA QAFS

# Kibble-Zurek mechanism

*Kibble (1976), Zurek (1985)*

Quantum phase transition of the 1d transverse-field Ising model

$$H = -J \sum_{i=1}^L \sigma_i^z \sigma_{i+1}^z - \Gamma \sum_{i=1}^L \sigma_i^x \quad (J > 0)$$



Defect (kink)



When we change  $\Gamma/J$  at a finite speed, the system goes out of equilibrium (diabatic). Defects are created in the final state at  $\Gamma/J=0$ .

Problem: Set a quantitative measure of diabaticity.  
How many defects ( $n$ ) are created as a function of the annealing time ( $t_a$ )?

$n$  as a function of  $t_a$

# Prediction of the Kibble-Zurek theory

$$\frac{n}{L} = \rho \propto (t_a)^{-\frac{d\nu}{1+\nu z}}$$

$d$ : system dimension

$n$ : average of the number of kinks

$L$ : chain length

$t_a$ : elapsed time = 1 / (speed of parameter change)

$$\xi \sim (\Gamma - \Gamma_c)^{-\nu}$$

$$\tau \sim (\Gamma - \Gamma_c)^{-\nu z}$$

## 1d ferromagnetic transverse-field Ising model

$$\rho \propto t_a^{-0.5}$$

Isolated system:  $\nu = z = 1$

$$\rho \propto t_a^{-0.28}$$

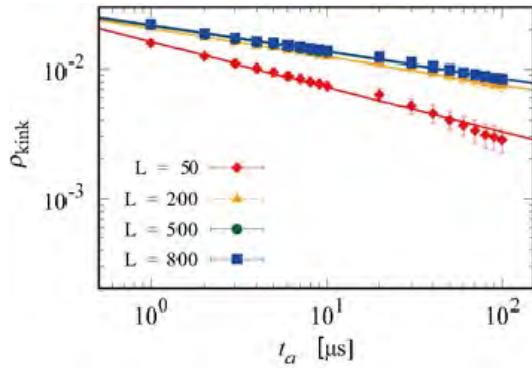
Under bosonic environment:  $\nu = 0.63$ ,  $z = 1.98$  (*Werner et al., 2005*)

$$H = -J \sum_{i=1}^{N_x} \sigma_i^z \sigma_{i+1}^z - \Delta \sum_{i=1}^{N_x} \sigma_i^x + \sum_{i,k} \{ C_k (a_{i,k}^\dagger + a_{i,k}) \sigma_i^z + \omega_{i,k} a_{i,k}^\dagger a_{i,k} \},$$

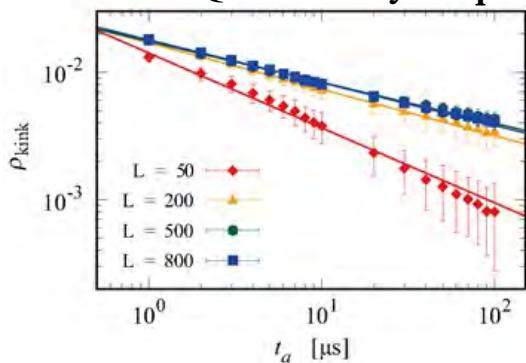
# Result

## Average number of kinks $n$

DW2000Q@NASA : exponent 0.20



DW2000Q@Burnaby: exponent 0.34



Data for  $L=50$  are excluded because they don't satisfy the condition for KZ theory to be valid,  $n>1$ .

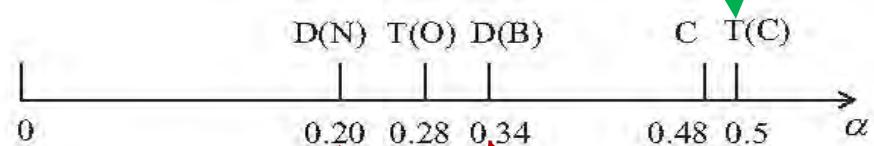
Exponent for the average  $n$

Under bosonic environment

Theory: 0.28

Isolated system

Theory: 0.5



Experiment 0.20(NASA), 0.34(Burnaby)

$$\begin{aligned} H = & -J \sum_{i=1}^{N_x} \sigma_i^z \sigma_{i+1}^z - \Delta \sum_{i=1}^{N_x} \sigma_i^x \\ & + \sum_{i,k} \{C_k(a_{i,k}^\dagger + a_{i,k})\sigma_i^z + \omega_{i,k} a_{i,k}^\dagger a_{i,k}\}, \end{aligned}$$

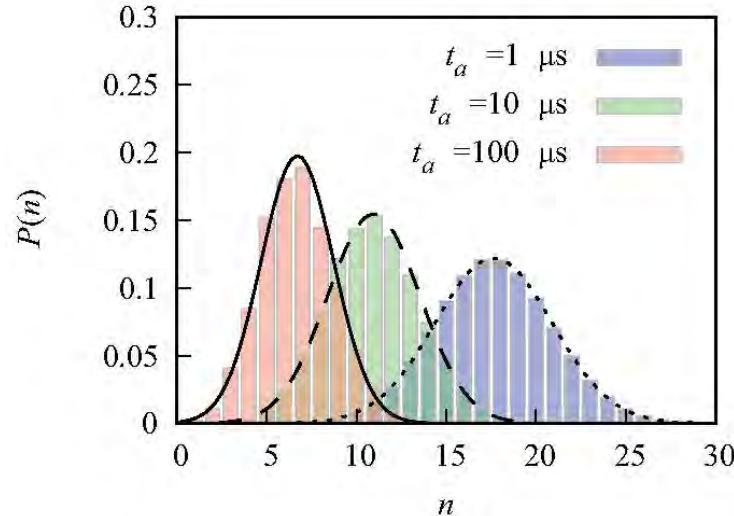
Deviations would come from non-universal effects.

# Result

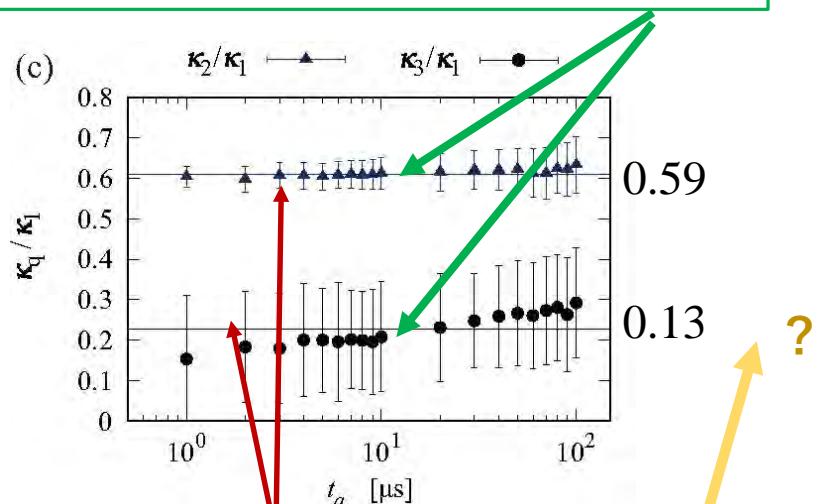
## Generalized KZ theory on the distribution of $n$

*Adolfo del Campo (2018)*

Distribution of the number of kinks

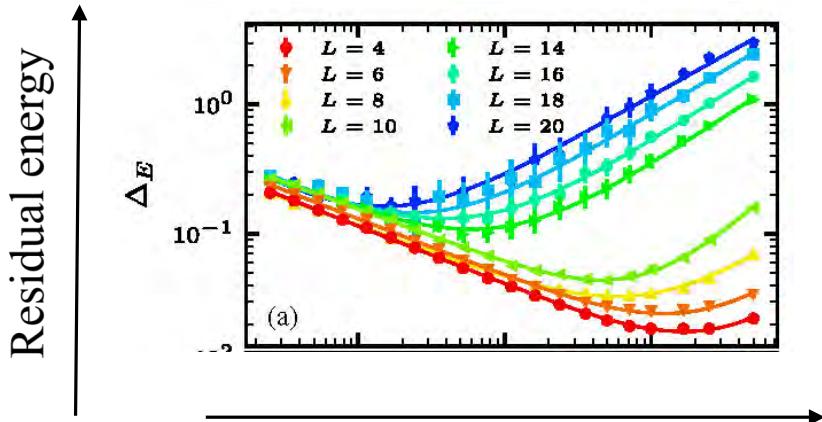


Theoretical prediction for isolated system



- Our quantum simulation (experiment) predicts that the theoretical prediction for isolated system should remain valid even under bosonic environment.
- Prompting theoretical development for confirmation.
- This is (probably) the first time that quantum simulation has gone ahead of theory.

# Related study 2d square lattice



Speed of change  $v = (\text{annealing time } t_a)^{-1}$   
(the other way round from our notation)

Necessary:

- Theory for 2d system under bosonic environment.
- To understand why 1d has no minimum but 2d has.

Weinberg *et al.*,  
*PRL 124, 090502(2020)*

$$f(v) = a_L v^\alpha + b_L v^{-\beta}$$

Kibble-Zurek

Experiment

$$\alpha = 0.74 \pm 0.02$$

Theory for isolated system

Bath

$$\alpha = 0.77$$

$$\mathcal{V}_{\text{noise}}(t) = \lambda \sum_i \eta_i(t) \sigma_i^x,$$

Confirmed qualitatively for 1d numerically

# Our contribution (2)

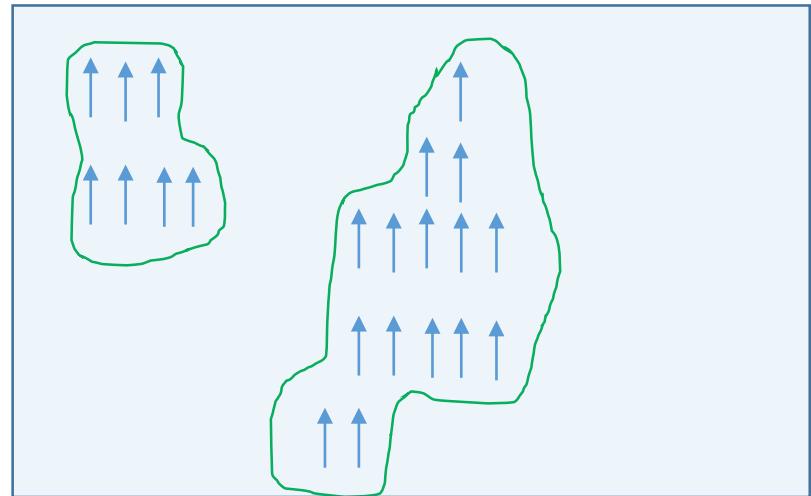
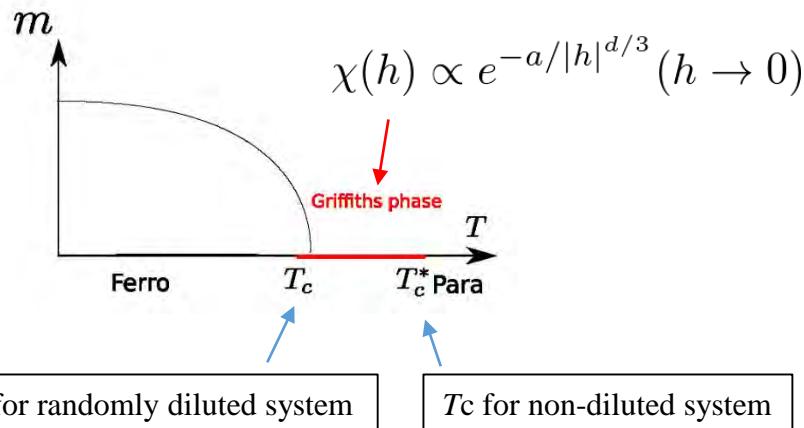
## Quantum simulation of the Griffiths-McCoy singularity

Kohji Nishimura, H. Nishimori, Helmut G. Katzgraber,  
*Phys. Rev. A (to be published)*, arXiv:2006.16219

Supported by IARPA QEO / DARPA QAFS

# Griffiths-McCoy singularity

**Problem:** Does the Griffiths-McCoy singularity exist on the Chimera graph?



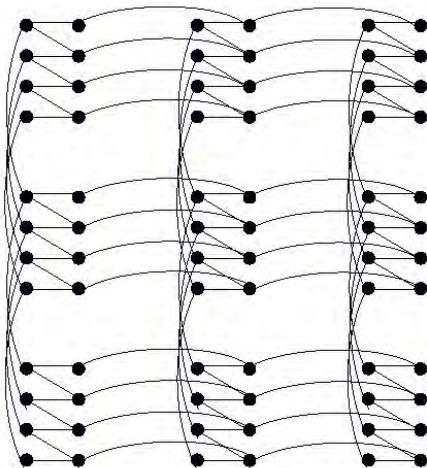
## Ising ferromagnet on a diluted lattice

- Large (but very rare) clusters respond very strongly to an external field even in the paramagnetic phase for  $T_c(\text{diluted}) < T < T_c(\text{non-diluted})$ .  $\rightarrow$  (Weak) singularity in  $\chi$  at  $h=0$ .
- This is **enhanced in low-dimensional quantum spin systems**, resulting in **divergence of non-linear susceptibility**.

# GM singularity on diluted Chimera graph

- Numerical and experimental studies show evidence of the Griffiths-McCoy singularity in low-dimensional quantum magnets.
- But studies of the 2d randomly-diluted ferromagnet have been rare. Mostly spin glass.
- There has been no study for the Chimera graph.
- There has been no study by quantum simulation on quantum device.

**Regularly-diluted** Chimera graph with smaller connectivity  
 → Stronger Griffiths-McCoy singularity if randomness is introduced.



**Randomly diluted ferromagnet**

$$P(J_{ij}) = \frac{1}{6} \sum_{k=0}^5 \delta(J_{ij} + 0.2k),$$

# Analysis and result

Non-linear susceptibility  
(stronger singularity than linear)

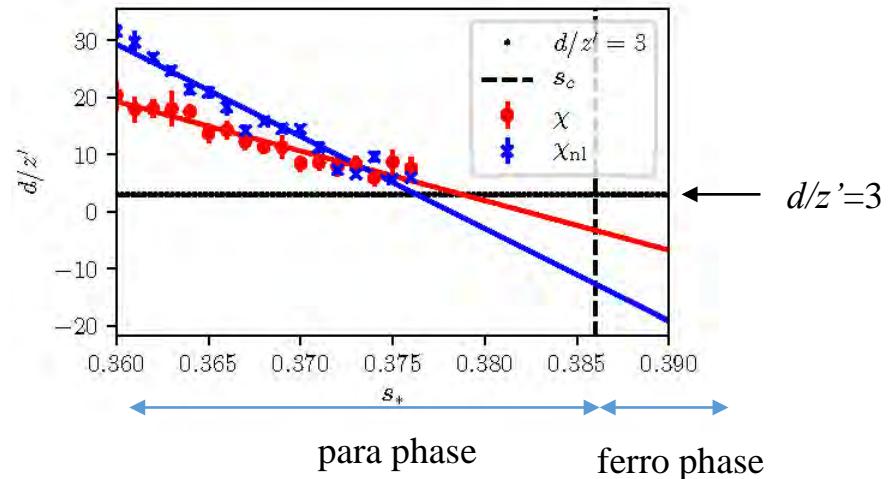
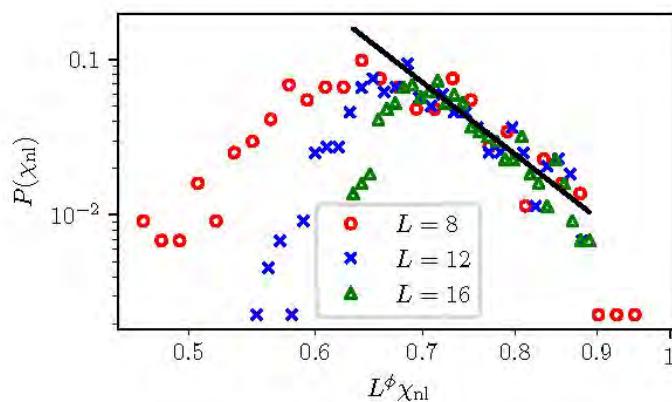
$$\chi_{\text{nl}} = \frac{\partial^3 m}{\partial h^3} \Big|_{h=0} = -(\langle m^4 \rangle - 3\langle m^2 \rangle^2)$$

Distribution over samples and local environments

$$P(\chi_{\text{nl}}) \sim \chi_{\text{nl}}^{-(d/(3z')+1)}$$

Average diverges if  $d/z' < 3$

$$\langle \chi_{\text{nl}} \rangle = \int^{\infty} \chi_{\text{nl}} P(\chi_{\text{nl}}) d\chi_{\text{nl}} \sim \int^{\infty} \chi_{\text{nl}}^{-d/(3z')} d\chi_{\text{nl}}$$



- Quantum simulation has shown: non-linear susceptibility is likely to diverge in the para phase.
- Consistent with the existence of the Griffiths-McCoy singularity on the Chimera graph.
- Backed up by classical simulations (“quantum” Monte Carlo)

# Conclusion

- Non-equilibrium (dynamical) phenomenon of Kibble-Zurek mechanism and its generalization.  
Quantum experiment has gone ahead of theory: A theory for qubits under environment should be developed.
- Equilibrium (static) phenomenon of the Griffiths-McCoy singularity.  
Nonlinear susceptibility has been shown to be likely to diverge even in the paramagnetic phase on the Chimera graph.
- These result motivate further quantum simulations on the existing and future quantum annealers for discoveries/verifications of new/existing physics.