

Universal aspects of low-temperature transport from GHD

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June 12, 2020

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Universal aspects of low-temperature transport from GHD

- B. Bertini, L. Piroli, and P. Calabrese,
Universal Broadening of the Light Cone in Low-Temperature Transport,
Phys. Rev. Lett. **120**, 176801 (2018)
- B. Bertini and L. Piroli,
Low-temperature transport in out-of-equilibrium XXZ chains,
J. Stat. Mech. 033104 (2018)
- M. Mestyán, B. Bertini, L. Piroli, and P. Calabrese,
Spin-charge separation effects in the low-temperature transport of one-dimensional Fermi gases,
Phys. Rev. B **99**, 014305 (2019)

Motivations

- GHD introduced in
 - B. Bertini, M. Collura, J. De Nardis, M. Fagotti,
Transport in Out-of-Equilibrium XXZ Chains: Exact Profiles of Charges and Currents,
Phys. Rev. Lett. **117**, 207201 (2016)
 - O. A. Castro-Alvaredo, B. Doyon, T. Yoshimura,
Emergent Hydrodynamics in Integrable Quantum Systems Out of Equilibrium,
Phys. Rev. X **6**, 041065 (2016)
- Original motivations:
 - study of quantum quenches in inhomogeneous settings

Quantum quenches

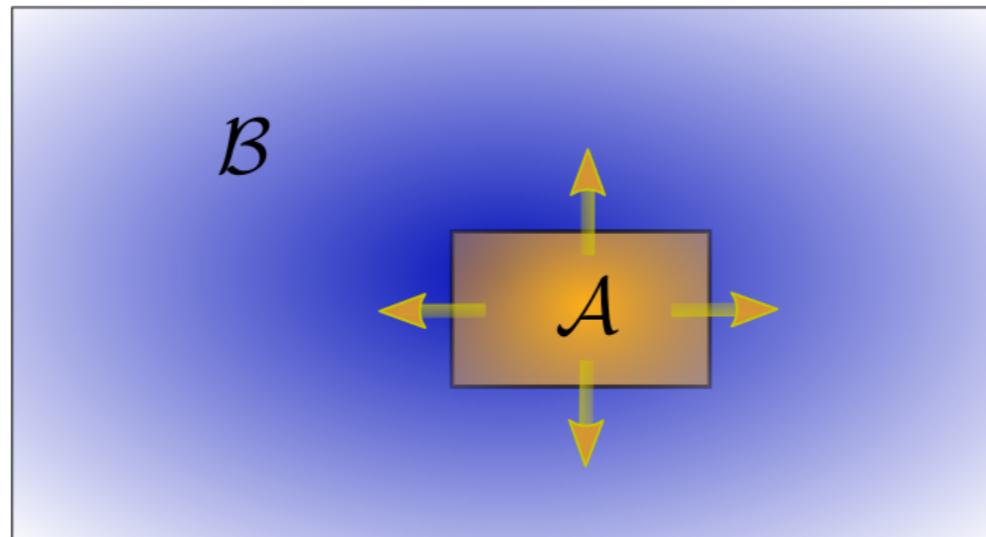
- Well-defined initial state $|\Psi_0\rangle$, e.g. ground-state of

$$H(g) = \sum_j h_j(g)$$

- At time $t = 0$, evolve with different Hamiltonian, e.g. $g \rightarrow g'$
- Compute **local** correlations $\langle \mathcal{O}_{j,\dots,j+q}(t) \rangle$
- Ideal **homogeneous** systems
- Lots of no's:
 - no disorder
 - no trapping potential
 - no external driving

Quantum quenches

- Obvious motivations from **experimental** and **foundational** point of view
- Physical intuition: equilibration occurs **locally**



- **Thermalization** expected for generic systems
- Different behavior in the presence of **conservation laws**

M. Rigol, V. Dunjko, M. Olshanii, Nature **452**, 854 (2008)

M. Rigol, V. Dunjko, V. Yurovsky, M. Olshanii, Phys. Rev. Lett. **98**, 050405 (2007)

- **Integrable systems:** extensive number of local conservation laws

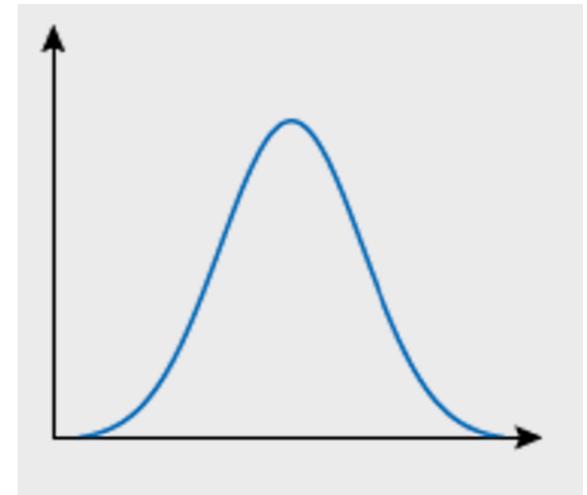
$$H = \sum_j h_j \quad \boxed{\quad} \Rightarrow [H, Q^{(k)}] = 0$$

$$Q^{(k)} = \sum_j q_j^{(k)} \quad \xleftarrow{\text{finite support}}$$

- **Exactly solvable:** spectrum can be computed via **Bethe Ansatz**
- Description in terms of **stable quasi-particles**

- Finite systems: $E = \sum_{j=1}^N \varepsilon(\lambda_j)$ ← “quasi-momenta”
- TD limit: each eigenstate associated with distribution function of quasi-momenta $\rightarrow \rho(\lambda)$

- Thermal states $\rightarrow \rho_T(k) \propto \frac{1}{1 + e^{\varepsilon(k)/K_B T}}$
- "Dressed" dispersion relation



- Quench problem conceptually **solved** for integrable systems:

- at large times, system described by a **GGE**

$$\hat{\rho}_{\text{GGE}} \propto e^{\sum_k \beta_k Q_k}$$

associated with (non-thermal) distribution functions $\rho(\lambda)$

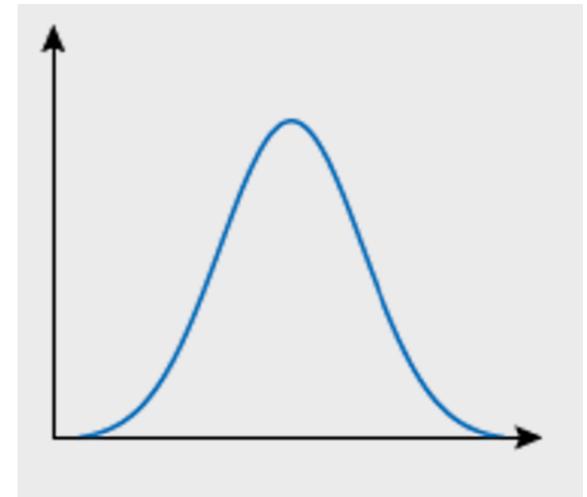
- Recent developments allow us to **extract correlations** from $\rho(\lambda)$
- Several approaches developed to **compute** $\rho(\lambda)$ in practice

F. H. L. Essler and M. Fagotti, J. Stat. Mech. 064002 (2016)

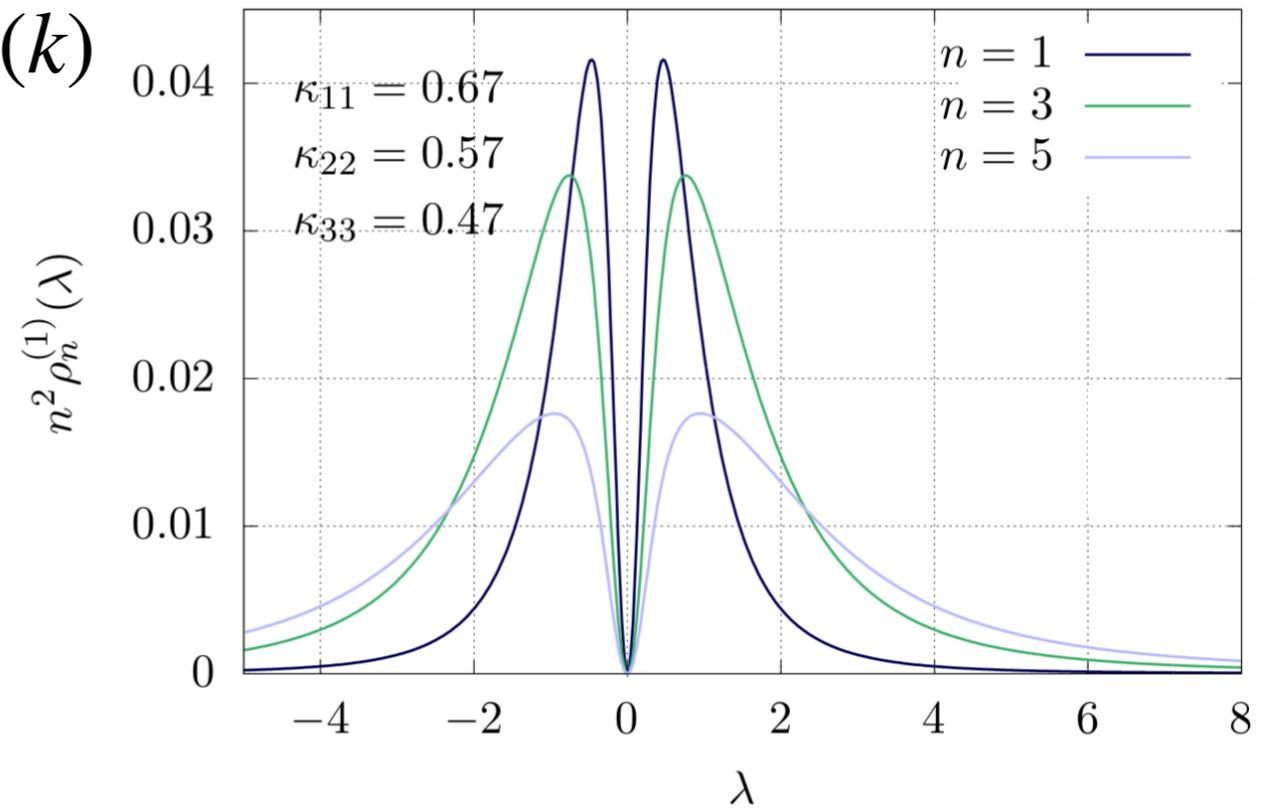
J.-S. Caux, J. Stat. Mech. 064006 (2016)

E. Ilievski, M. Medenjak, T. Prosen, L. Zadnik, J. Stat. Mech. 064008 (2016)

- Thermal states $\rightarrow \rho_T(k) \propto \frac{1}{1 + e^{\varepsilon(k)/K_B T}}$
 "Dressed" dispersion relation



- GGE states $\rightarrow \rho_{GGE}(k) \neq \rho_T(k)$



- Quench problem conceptually solved for integrable systems:

- at large times, system described by a GGE

$$\hat{\rho}_{\text{GGE}} \propto e^{\sum_k \beta_k Q_k}$$

associated with (non-thermal) distribution functions $\rho(\lambda)$

- Recent developments allow us to extract correlations from $\rho(\lambda)$
- Several approaches developed to compute $\rho(\lambda)$ in practice
- What happens in inhomogeneous systems?

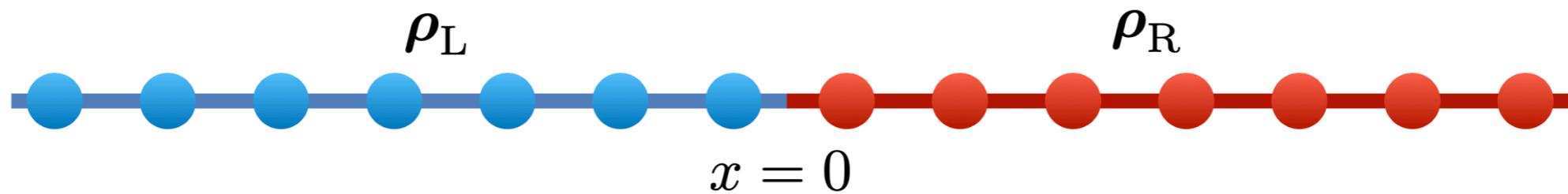
F. H. L. Essler and M. Fagotti, J. Stat. Mech. 064002 (2016)

J.-S. Caux, J. Stat. Mech. 064006 (2016)

E. Ilievski, M. Medenjak, T. Prosen, L. Zadnik, J. Stat. Mech. 064008 (2016)

Inhomogeneous quenches

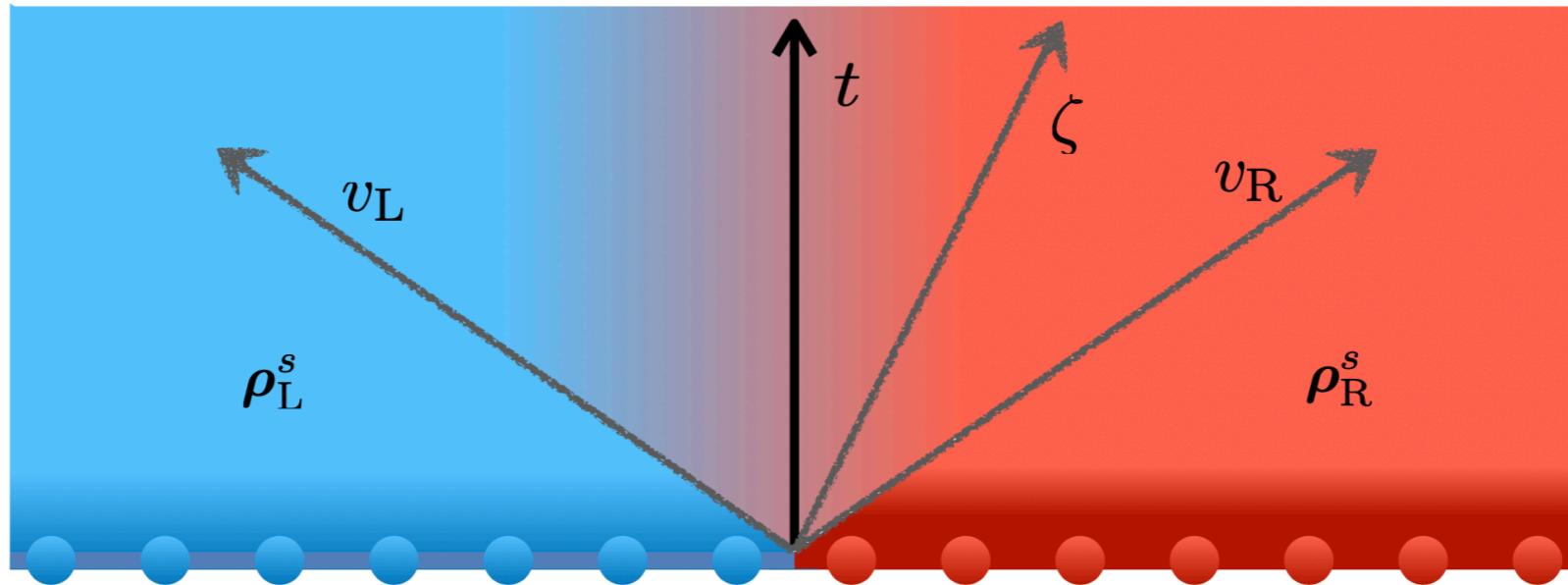
- Simplest setting: **bipartition** protocols



- Ex: $\rho_L = \frac{1}{Z_L} e^{-\beta_L \sum_{x<0} h_x + \mu_L \sum_{x<0} s_x^z}$ $\rho_R = \frac{1}{Z_R} e^{-\beta_R \sum_{x>0} h_x + \mu_R \sum_{x>0} s_x^z}$

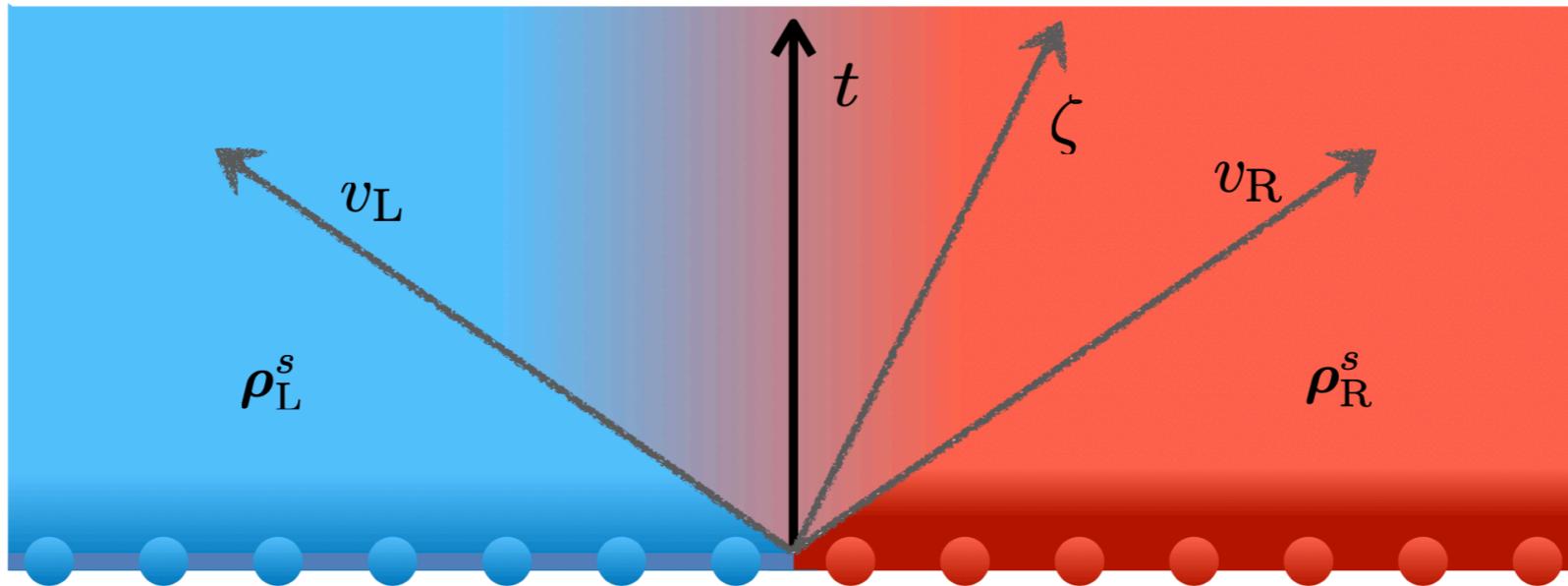
- Let evolve with integrable Hamiltonian at time $t = 0$
- What happens at large times?
- The problem has a long history: previously studied in **free theories** and **conformal systems**

The GHD solution



- Postulate emergence of space-time dependent stationary states $\rho_{x,t}(\lambda)$
- Intuition:
 - information carried by **quasi-particles** moving from the two sides
 - frames at different velocities receive different sets of quasi-particles
- \Rightarrow in the limit $x, t \rightarrow \infty$ stationary states only depend on $\zeta = x/t$

The GHD solution



- States $\rho_{x,t}(\lambda)$ fixed by constraints from **conservation laws**:
 - continuity equations:

$$\partial_x \langle Q \rangle_\zeta + \partial_t \left\langle J_Q \right\rangle_\zeta = 0$$

- expectation values:

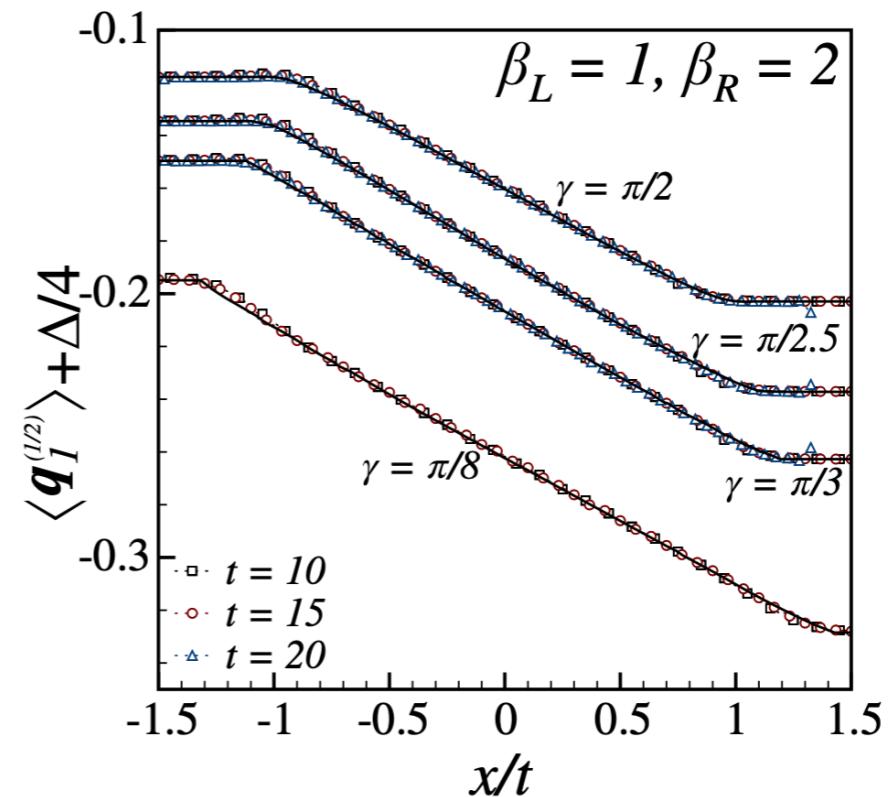
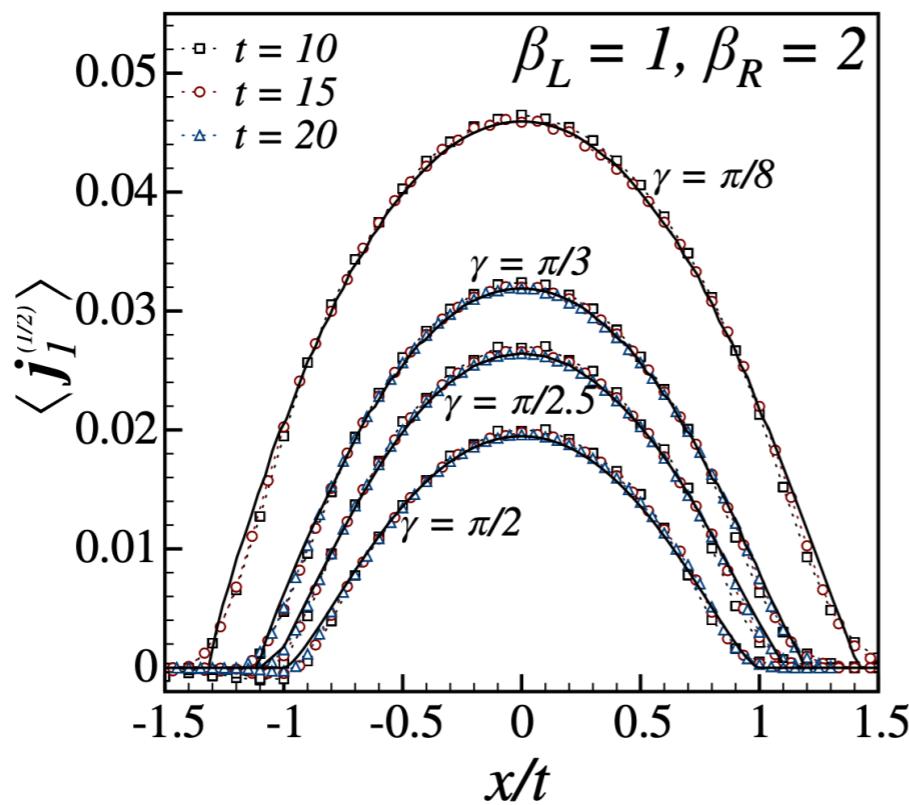
$$\langle Q \rangle = \int dk \rho_\zeta(k) q(k) \quad \left\langle J_Q \right\rangle = \int dk \rho_\zeta(k) q(k) v(k)$$

- Final result:

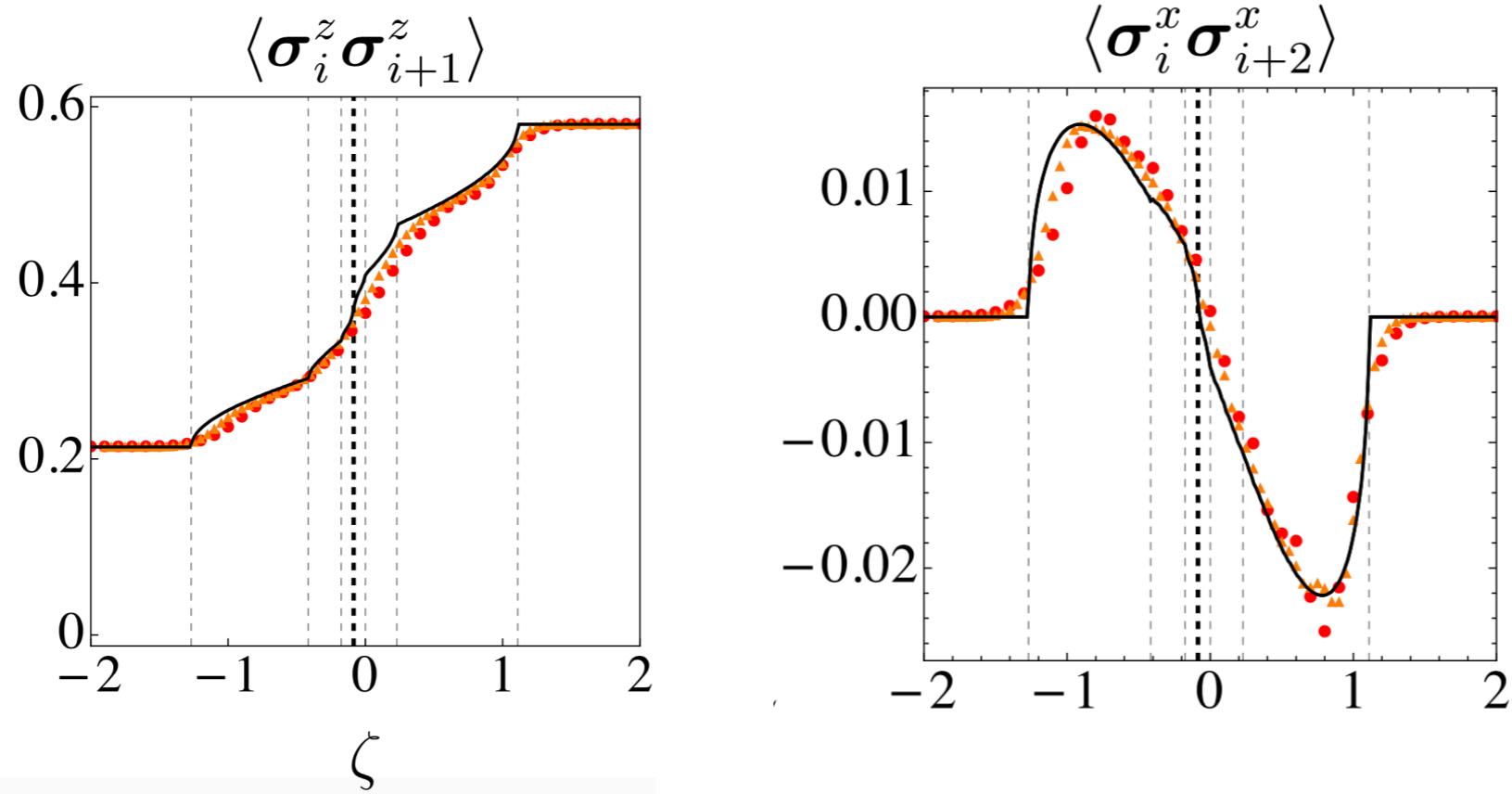
$$\partial_t \rho_\zeta(\lambda) + \partial_x \left(v_\zeta(\lambda) \rho_\zeta(\lambda) \right) = 0$$

$$\lim_{\zeta \rightarrow \pm\infty} \rho_\zeta(\lambda) = \rho_{L,R}(\lambda)$$

- Equations in general solved **numerically**
- Predictions are claimed to be **exact** for $t \rightarrow \infty$



- In principle, all local observables can be computed



- Early literature: “phenomenology” of transport profiles:
 - characteristic features
 - universal aspects
 - solvable limits

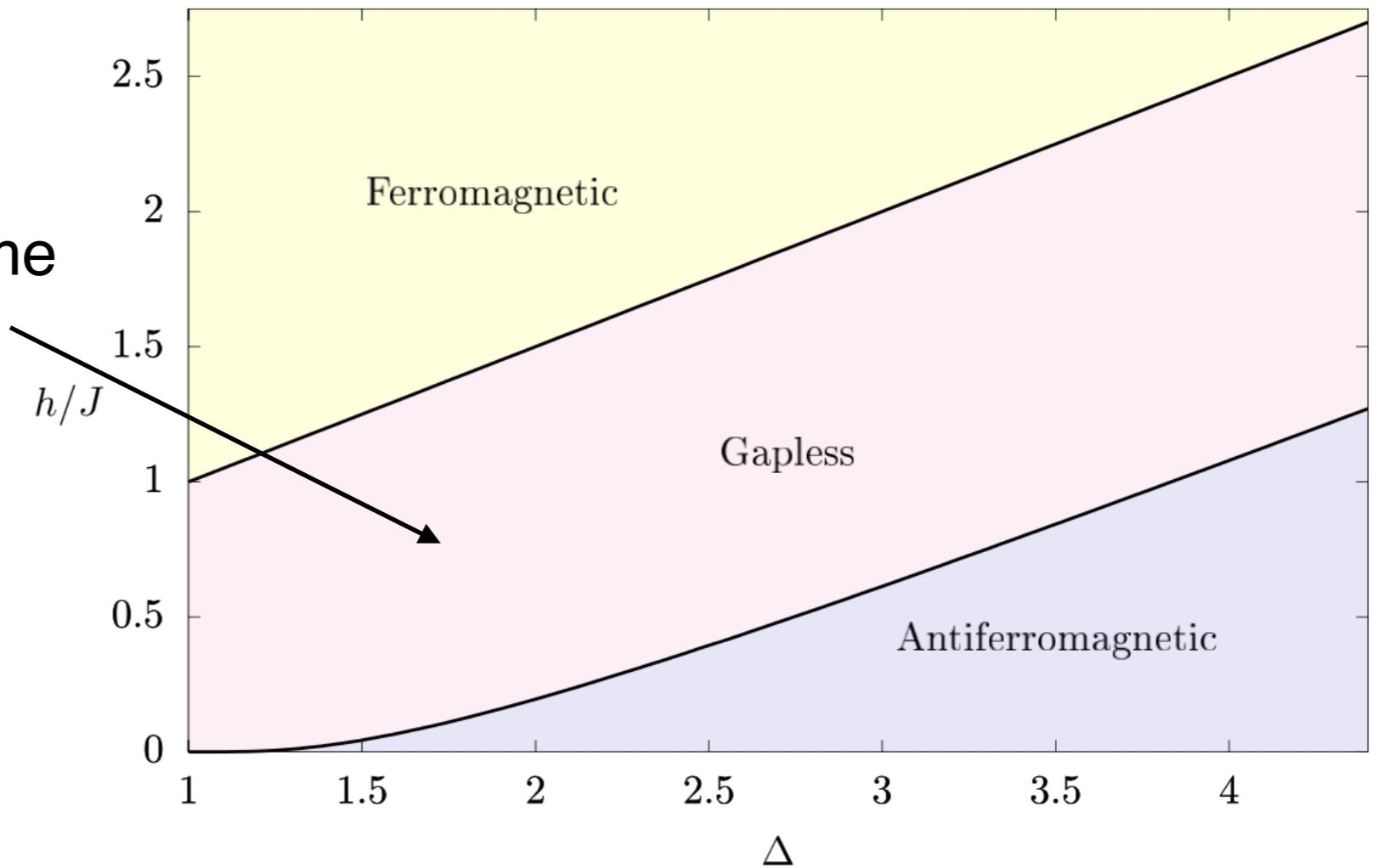
- This talk:
 - focus on **low-temperature** regime of bipartition protocols
- Main points:
 - **test** GHD against previous **CFT** predictions
 - discovery of new **universal** in the Luttinger-Liquid class
 - **spin-charge separation** effects in multi-species integrable models

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- B. Bertini and L. Piroli, J. Stat. Mech. 033104 (2018)
 - B. Bertini, L. Piroli, and P. Calabrese, Phys. Rev. Lett. **120**, 176801 (2018)
 - M. Mestyán, B. Bertini, L. Piroli, and P. Calabrese, Phys. Rev. B **99** 014305 (2019)

- Consider **XXZ** Heisenberg chain

$$H = \frac{J}{4} \sum_{j=-L/2}^{L/2-1} \left[\sigma_j^x \sigma_{j+1}^x + \sigma_j^y \sigma_{j+1}^y + \Delta (\sigma_j^z \sigma_{j+1}^z - 1) \right] - h \sum_{j=-L/2}^{L/2-1} (\sigma_j^z - 1)$$

Focus on gapless regime



- Focus on **XXZ** Heisenberg chain

$$H = \frac{J}{4} \sum_{j=-L/2}^{L/2-1} \left[\sigma_j^x \sigma_{j+1}^x + \sigma_j^y \sigma_{j+1}^y + \Delta (\sigma_j^z \sigma_{j+1}^z - 1) \right] - h \sum_{j=-L/2}^{L/2-1} (\sigma_j^z - 1)$$

- **TBA** equations

$$T_{L/R} \log \eta_j(\lambda) = e_j(\lambda) + T_{L/R} \sum_k \left[T_{jk} * \log (1 + \eta_k^{-1}) \right](\lambda)$$

- **Bethe** equations

$$\rho_{j,\zeta}(\lambda) + \rho_{j,\zeta}^h(\lambda) = a_j(\lambda) - \left[\sum_k T_{jk} * \rho_{k,\zeta} \right](\lambda)$$

$$\vartheta_j(\lambda) = \rho_j \left(\rho_j + \rho_j^h \right)^{-1}$$

- **GHD** solution

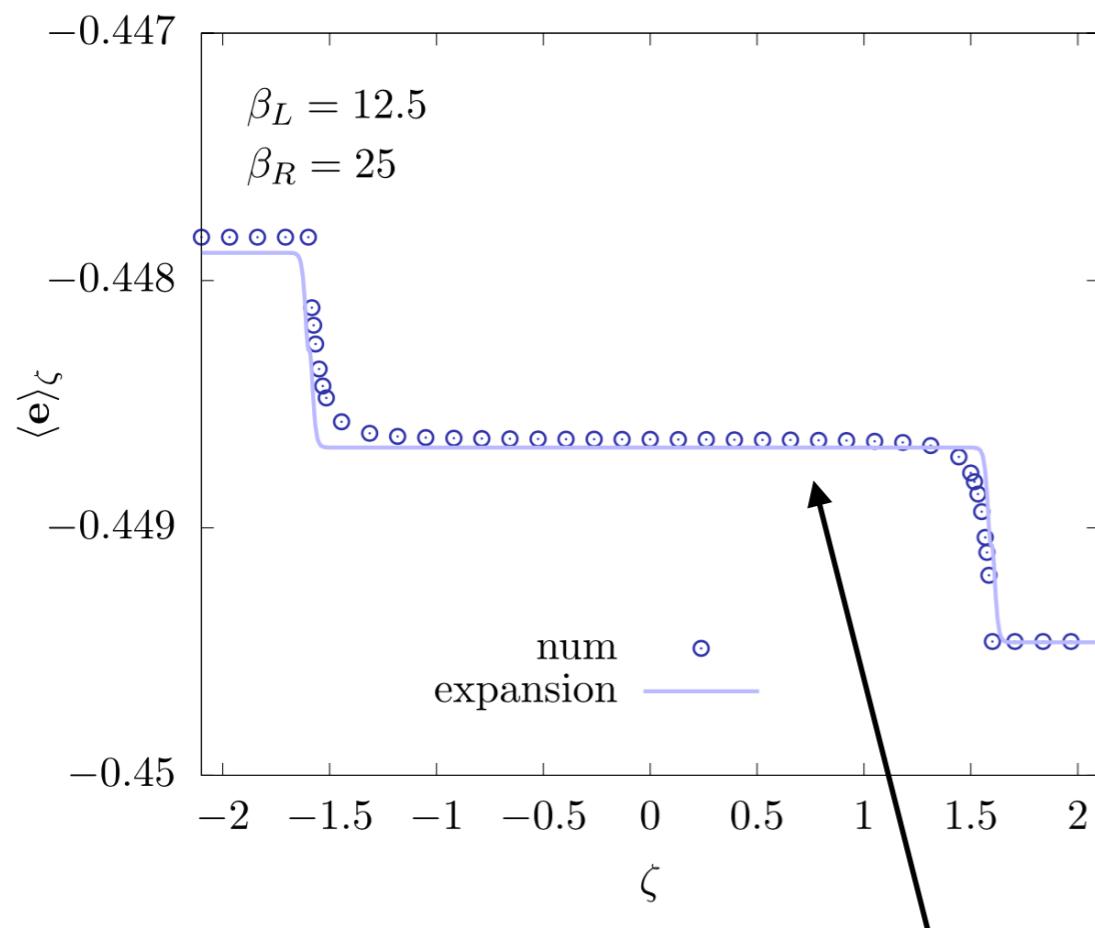
$$\vartheta_{j,\zeta}(\lambda) = \vartheta_j^R(\lambda) H \left(\zeta - v_{j,\zeta}(\lambda) \right) + \vartheta_j^L(\lambda) H \left(v_{j,\zeta}(\lambda) - \zeta \right)$$

- **Velocities**

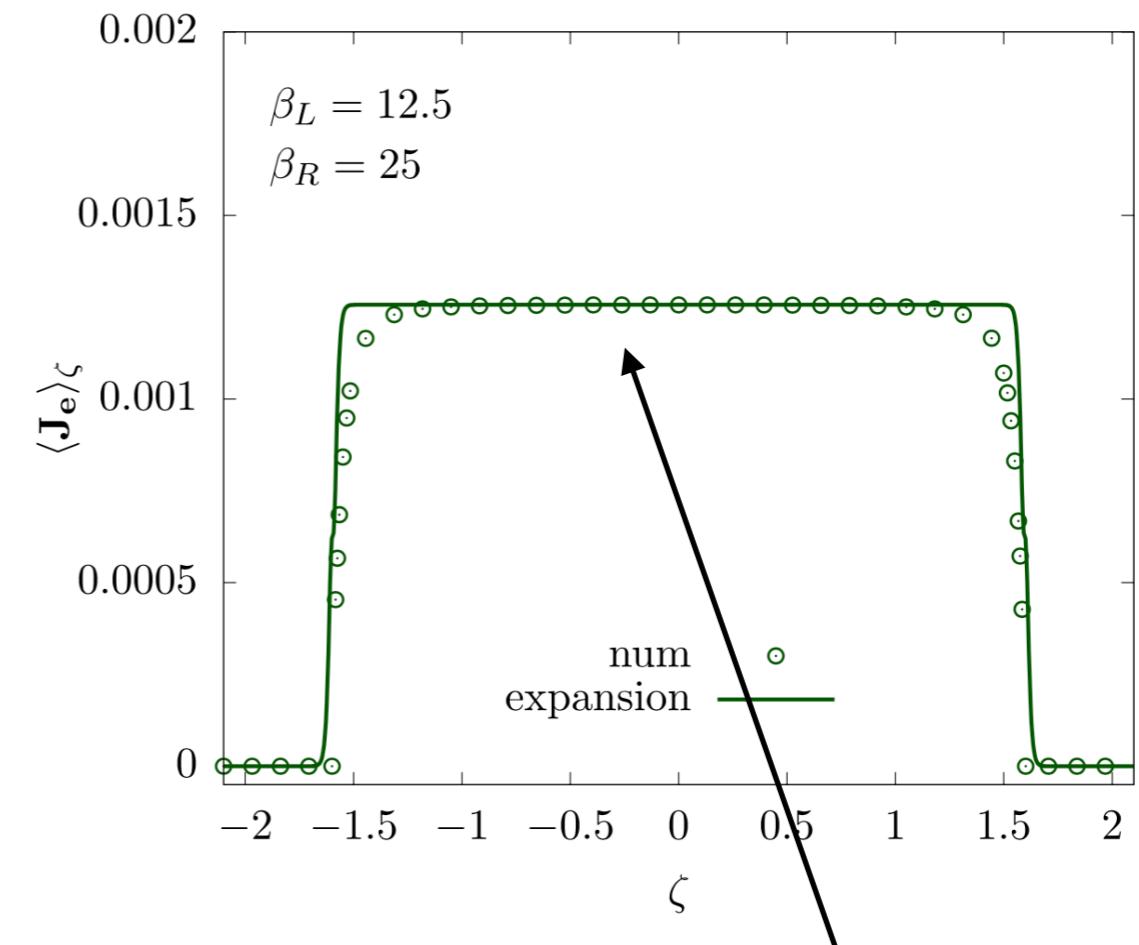
$$v_j(\lambda) \rho_j^t(\lambda) = \frac{a'_j(\lambda)}{2\pi} - \left[\sum_k T_{jk} * v_k \rho_k \right](\lambda)$$

Gapless regime

- Three-step profiles for **energy density** and currents:



$$\langle e \rangle_{\text{ness}} = \frac{\pi}{12} (T_l^2 + T_r^2)$$



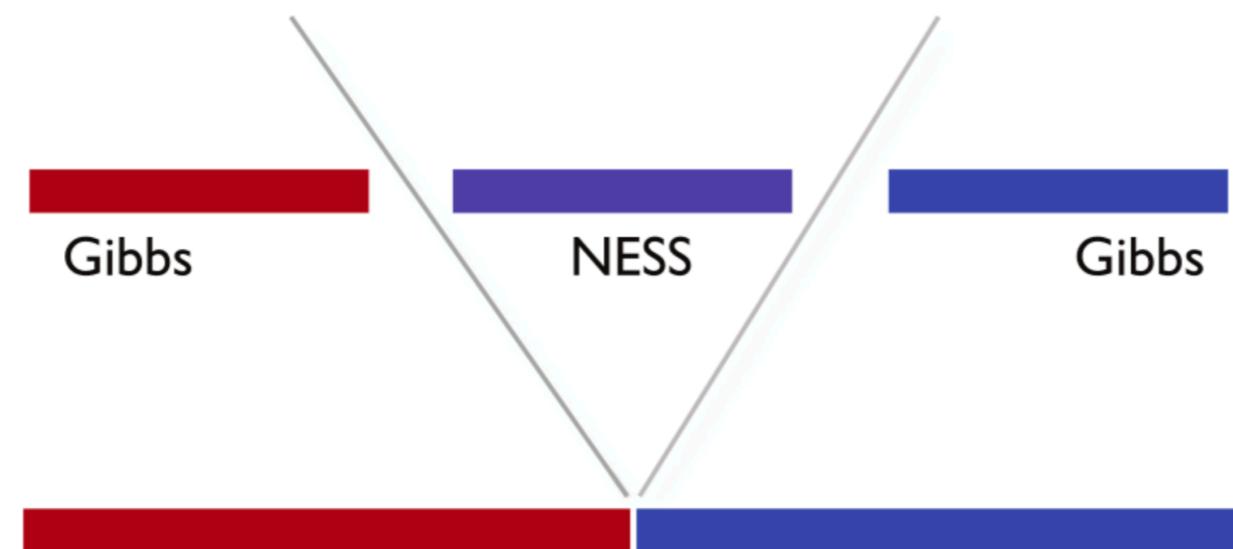
$$\langle j \rangle_{\text{ness}} = \frac{\pi}{12} (T_l^2 - T_r^2)$$

Gapless regime

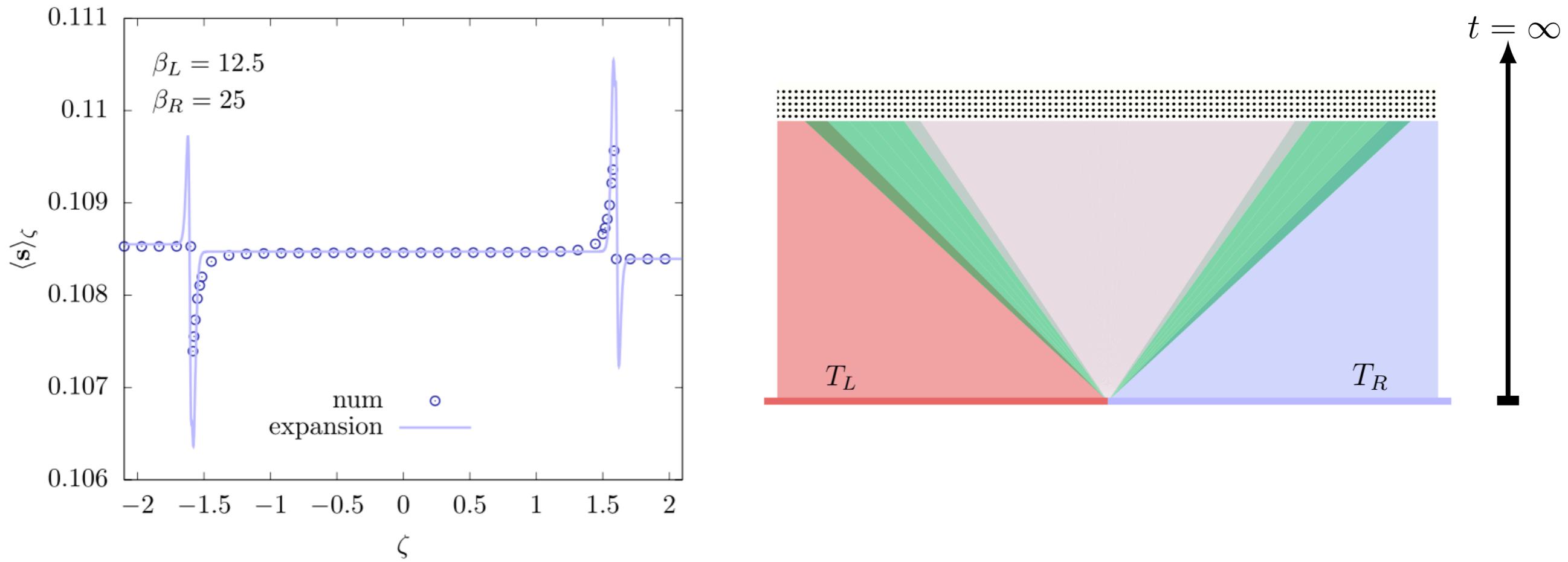
- We recover analytical results by **CFT**

$$\langle e \rangle_{\text{ness}} = \frac{c\pi}{12} (T_l^2 + T_r^2) \quad \langle j \rangle_{\text{ness}} = \frac{c\pi}{12} (T_l^2 - T_r^2)$$

central charge



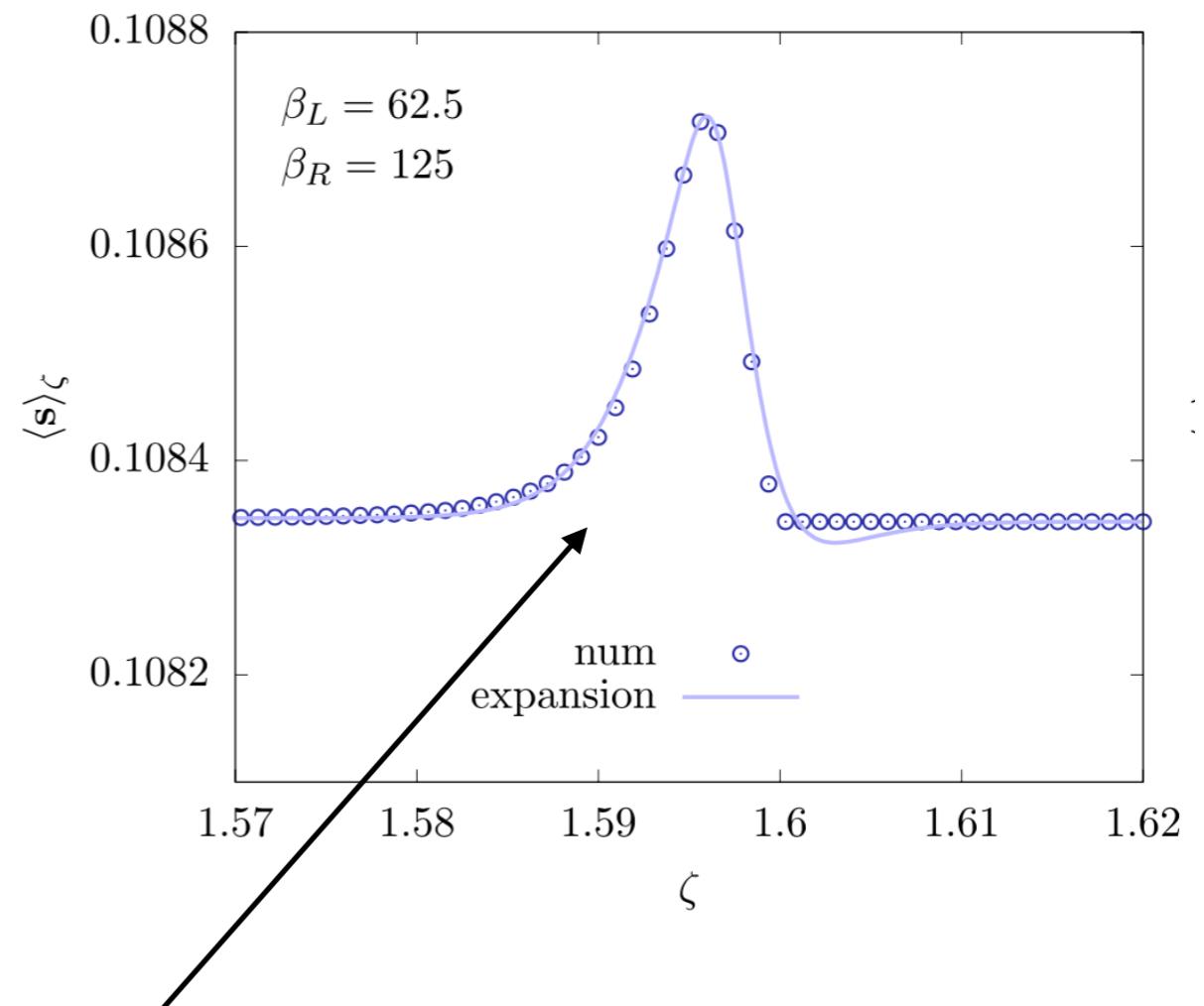
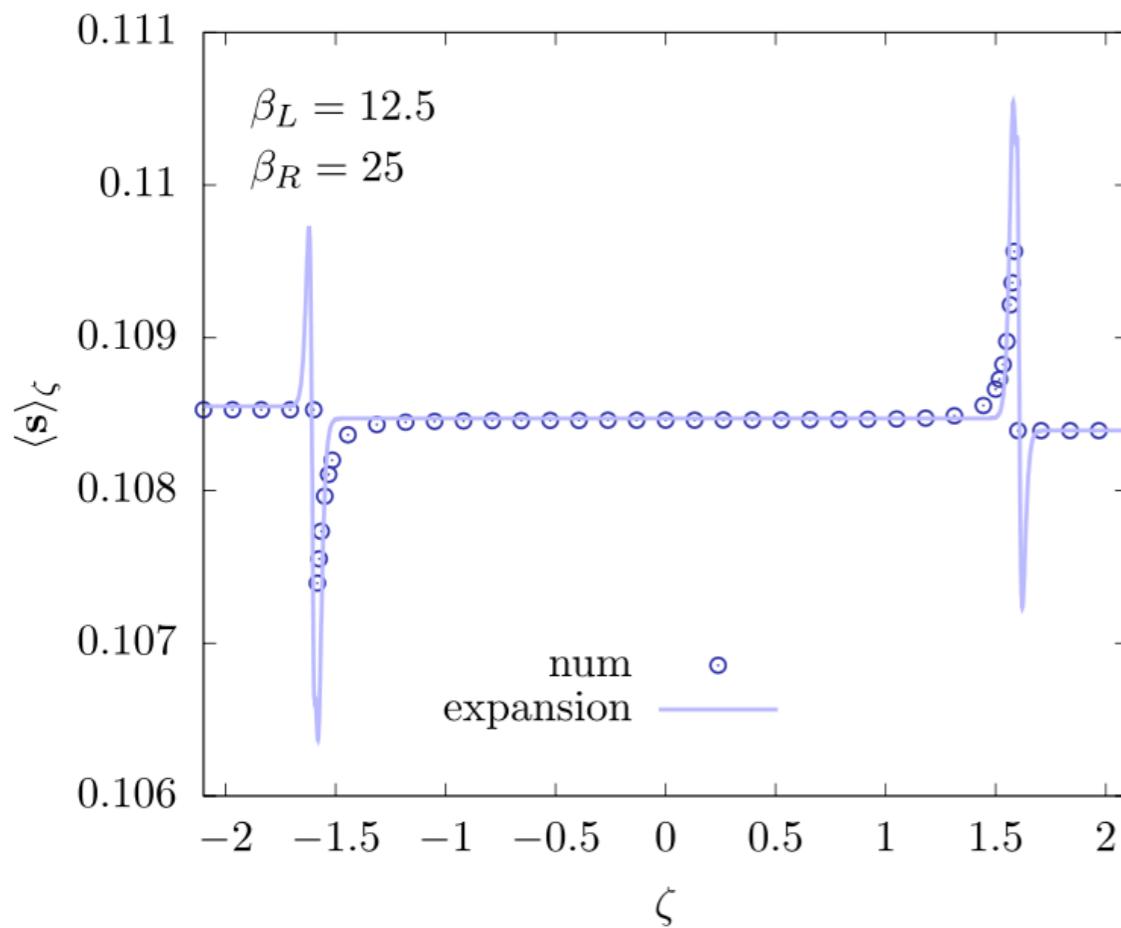
- **Universal broadening of the light-cone for generic observables**



B. Bertini and L. Piroli, J. Stat. Mech. 033104 (2018)

B. Bertini, L. Piroli, and P. Calabrese, Phys. Rev. Lett. **120** 176801 (2018)

- **Universal broadening of the light-cone for generic observables**



$$\mathcal{D}(z) \equiv T_L \log(1 + e^{z/T_L}) - T_R \log(1 + e^{z/T_R})$$

All observables proportional to the **same** function

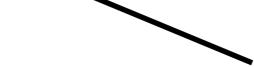
Non-linear Luttinger liquid description

- The broadening of the light cone can be obtained from **universal** non-linear Luttinger liquid description
- Starting from linear Luttinger liquid, dominant irrelevant term introduces a **curvature** in the dispersion of left and right movers

$$\varepsilon(k) = \nu k + \frac{1}{2m_*}k^2 \quad p(k) = k$$

- Simple calculation yields

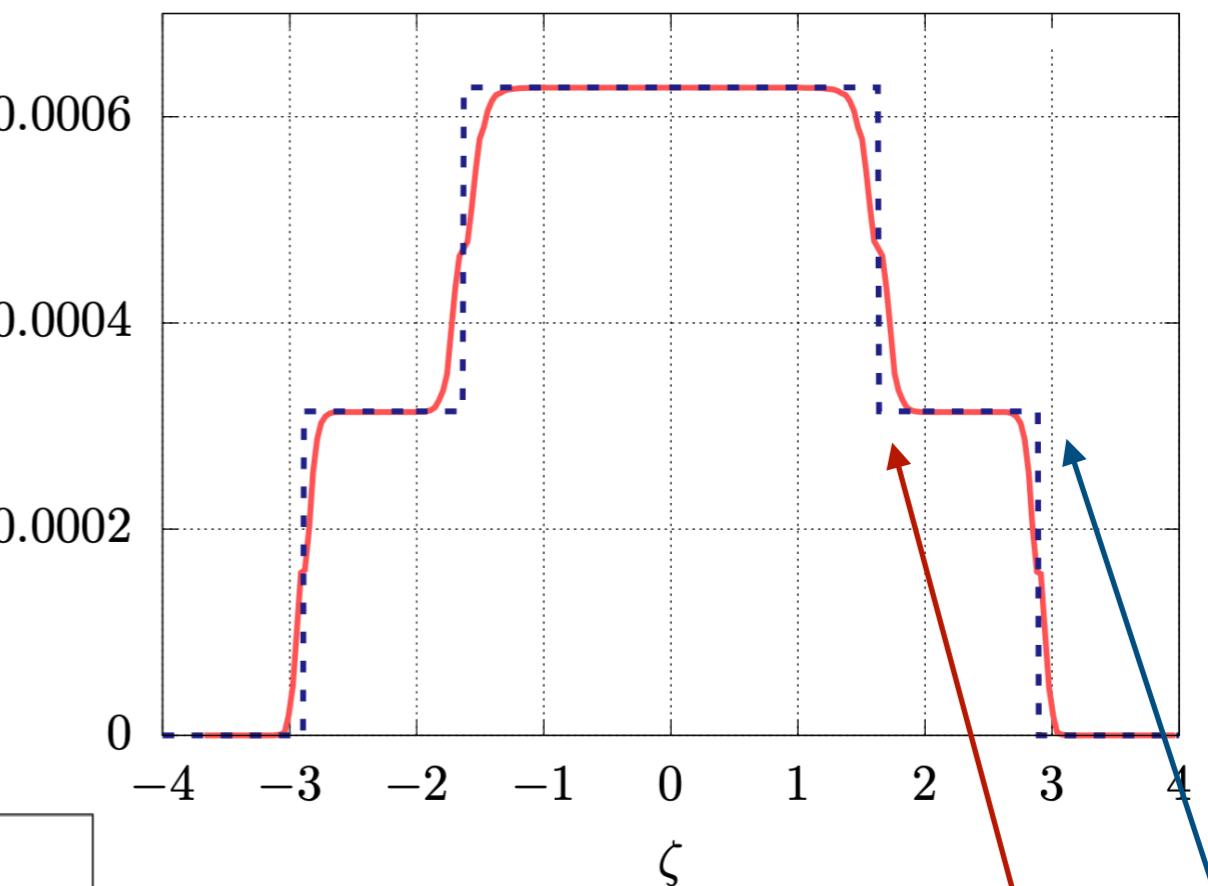
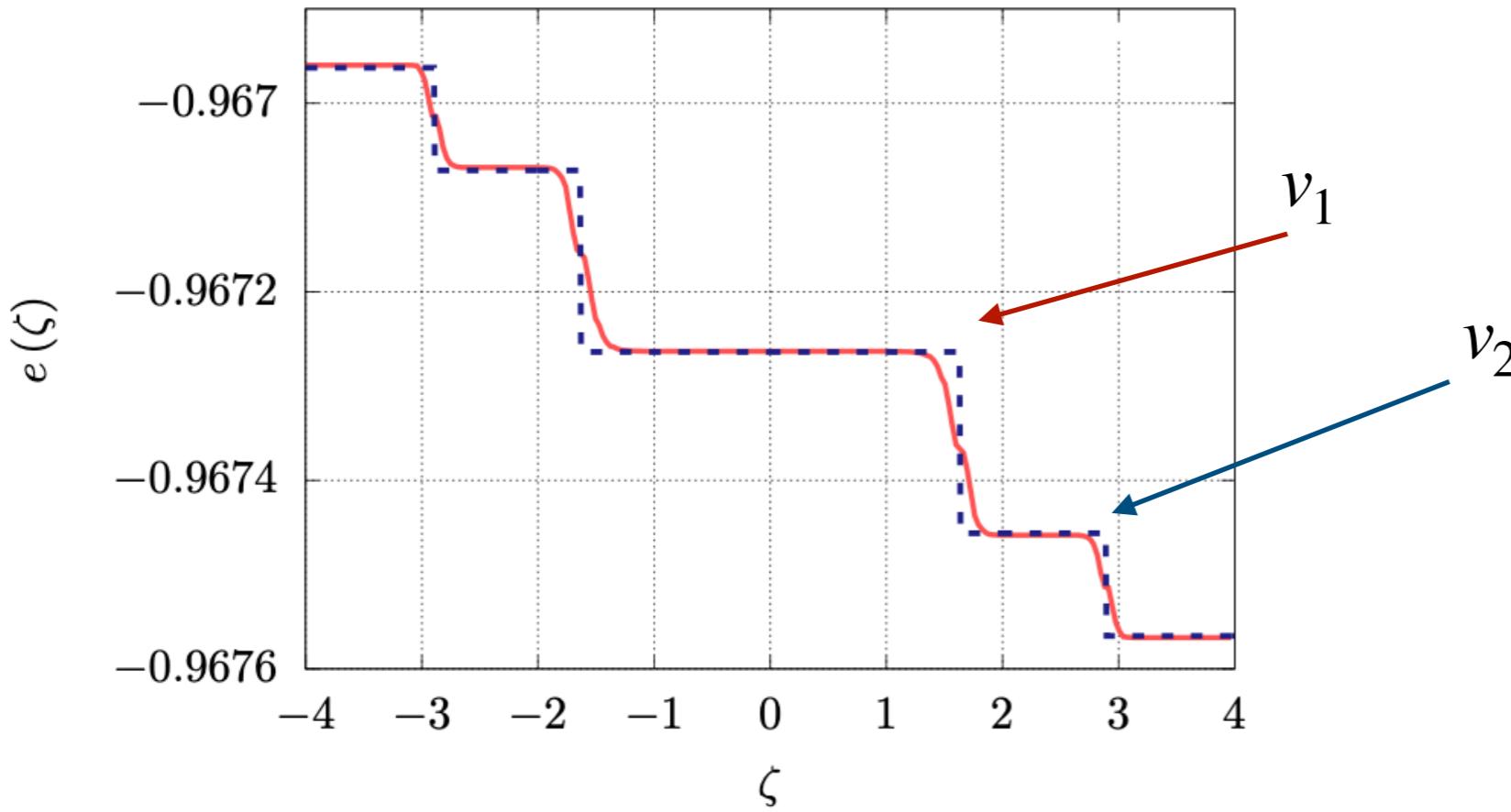
$$\langle \mathbf{q} \rangle_\zeta = \langle \mathbf{q} \rangle_{\text{GS}} + \frac{\alpha_1}{2\pi\nu^2} \mathcal{D} [m_*\nu(\zeta - \nu)] + \frac{\alpha_2}{2\pi\nu^2} \mathcal{D} [m_*\nu(\zeta + \nu)]$$

 **non-perturbative** effect in m_*

- Low-temperature analysis gives **qualitatively** different picture in **gapped** phases
- Particularly interesting features emerging in multi-component “**nested**” integrable systems
- Prototypical example: Yang-Gaudin model of **spinful fermions**

$$\hat{H} = - \int_{-L/2}^{L/2} dx \left[\sum_{\alpha=\pm} \psi_\alpha^\dagger(x) (\partial_x^2 + A + \alpha h) \psi_\alpha(x) \right] + c \int_{-L/2}^{L/2} dx \left[\sum_{\alpha,\beta=\pm} \psi_\alpha^\dagger(x) \psi_\beta^\dagger(x) \psi_\beta(x) \psi_\alpha(x) \right]$$

Emergence of **two distinct velocities**
associated with spin and charge
degrees of freedom



Outlook

- Many works have now extended GHD to more general settings (trapping potentials, inhomogeneities, dephasing..)
- GHD predictions for trap quenches observed in 1D Bose gases!
[M. Schemmer, I. Bouchoule, B. Doyon, J. Dubail, PRL **122** (2019)]
- Do low-temperature features **survive** in the presence of experimentally feasible settings?
- Ex: spin-charge separation effects in trapped multi-component Fermi gases*? [G. Pagano, et al. Nature Phys **10** (2014)]

*Ongoing work with Stefano Scopa, Pasquale Calabrese

Thank you for your attention!