



MaX School on Advanced Materials and Molecular Modelling with QUANTUM ESPRESSO

QE-2021: Hands-on session - Day-5 (Density Functional Perturbation Theory – Part 1)

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 Ljubljana, Slovenia





Outline

1. Introduction

- 2. Exercise 1a: Phonons at Gamma in non-polar materials
- 3. Exercise 1b: Phonon dispersion in non-polar materials
- 4. Exercise 2a: Phonons at Gamma in polar materials
- 5. Exercise 2b: Phonon dispersion in polar materials
- 6. Exercise 3: Phonon dispersion of 2D materials (optional)

Outline

1. Introduction

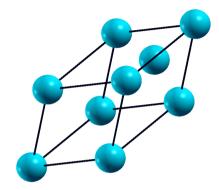
- 2. Exercise 1a: Phonons at Gamma in non-polar materials
- 3. Exercise 1b: Phonon dispersion in non-polar materials
- 4. Exercise 2a: Phonons at Gamma in polar materials
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- 6. Exercise 3: Phonon dispersion of 2D materials (optional)

Basic concepts

Let us consider a unit cell with N_{at} atoms.

 $s = 1...N_{at}$ index of an atom in the unit cell $\alpha = x, y, z$ is the cartesian index

- **R** is the point in the Bravais-lattice, identifying the position of a given unit cell
- $N_{\mathbf{R}}$ is the number of unit cells in the crystal



 $\mathbf{U}_{s\alpha}(\mathbf{R})$ is the α - component of the displacement of the *s*-th atom \mathbf{V} \mathbf{V} $\mathbf{3} \times N_{at}$ variables

Interatomic Force Constants :

$$C_{s\alpha,s'\beta}(\mathbf{R},\mathbf{R}') = C_{s\alpha,s'\beta}(\mathbf{R}-\mathbf{R}') = \frac{\partial^2 E_{tot}}{\partial \mathbf{u}_{s\alpha}(\mathbf{R})\partial \mathbf{u}_{s'\beta}(\mathbf{R}')}$$

Basic concepts

Normal mode frequencies, ω , and eigenvectors, $\mathbf{U}_{s\alpha}$ are determined by the secular equation:

$$\sum_{\boldsymbol{s}',\boldsymbol{\beta}} \tilde{D}_{\boldsymbol{s}\boldsymbol{\alpha},\boldsymbol{s}'\boldsymbol{\beta}}(\mathbf{q}) \, \tilde{\mathbf{u}}_{\boldsymbol{s}'\boldsymbol{\beta}}(\mathbf{q}) = \omega_{\mathbf{q}}^2 \, \tilde{\mathbf{u}}_{\boldsymbol{s}\boldsymbol{\alpha}}(\mathbf{q})$$

where

$$\tilde{D}_{s\alpha,s'\beta}(\mathbf{q}) = \frac{1}{\sqrt{M_s M_{s'}}} \sum_{\mathbf{R},\mathbf{R}'} \frac{\partial^2 E_{tot}}{\partial \mathbf{u}_{s\alpha}(\mathbf{R}) \partial \mathbf{u}_{s'\beta}(\mathbf{R}')} e^{i\mathbf{q}(\mathbf{R}'-\mathbf{R})}$$

is the *dynamical matrix*.

Diagonalization of the dynamical matrix gives phonon modes at **q**.

The phonon code works for a rather wide variety of systems and methods:

- Insulators (also polar insulators, with LO-TO splitting)
- Metals
- Magnetic systems at the scalar relativistic collinear level (LSDA)
- Spin-orbit coupling (fully relativistic approach)
- Electric field calculations: Born effective charges, dielectric tensor

Recent developments:

- Phonons for magnetic systems in the fully relativistic non-collinear approach
- Phonons within the DFT+U approach

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2. Exercise 1a: Phonons at Gamma in non-polar materials

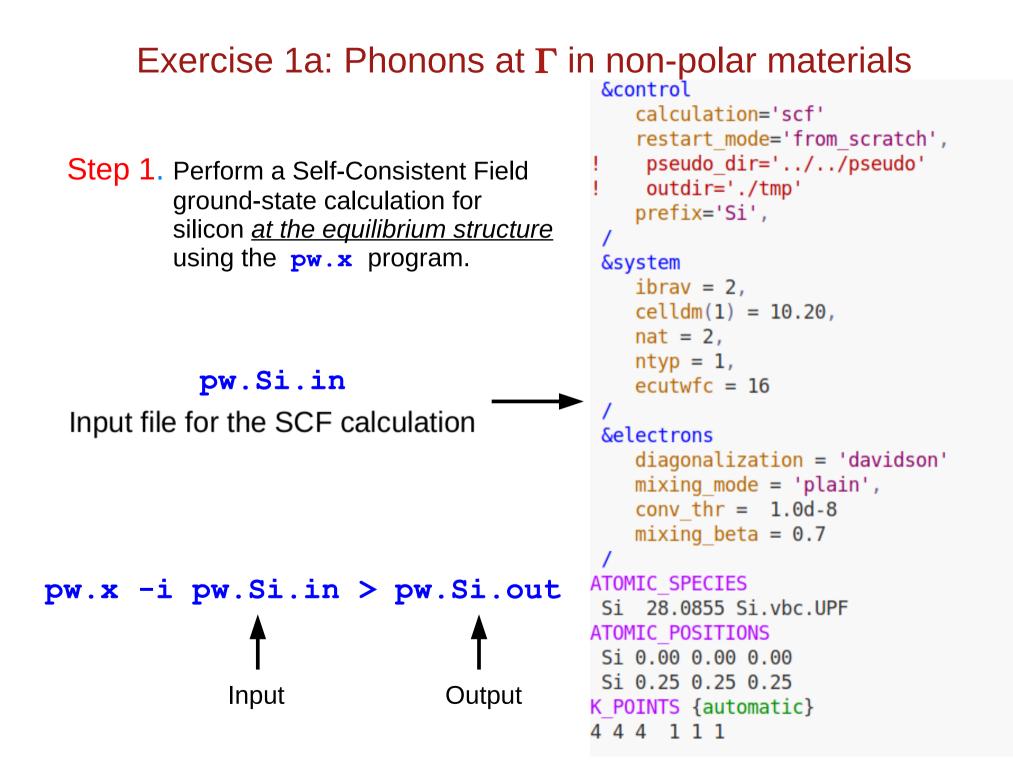
- 3. Exercise 1b: Phonon dispersion in non-polar materials
- 4. Exercise 2a: Phonons at Gamma in polar materials
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Go to the directory with the input files:

cd ~/QE-2021/Day-5/example1a

In this directory you will find:

- $README \cdot md File \ describing \ how \ to \ do \ the \ exercise$
- pw.Si.in Input file for the SCF ground-state calculation
- ph.Si.in Input file for the phonon calculation at Γ
- dynmat.Si.in Input file to impose the acoustic sum rule
- reference Directory with the reference results



Step 2. Perform a phonon calculation at Γ using the **ph.x** program.

ph.Si.in

Input file for the phonon calculation

```
Phonons at Gamma
   &inputph
    prefix = 'Si',
    tr2_ph = 1.0d-14,
    amass(1) = 28.0855,
! outdir = './tmp'
    fildyn = 'Si.dyn',
    /
    0.0 0.0 0.0
```

- ← The same prefix as in the SCF calculation
- Threshold for self-consistency
- Atomic mass
- Directory for temporary files
- File containing the dynamical matrix
- Coordinates of the **q** point in units of 2*pi/a in the Cartesian reference system

ph.x -i ph.Si.in > ph.Si.out

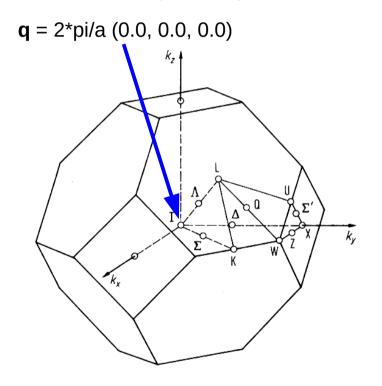
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Input file for the phonon calculation

```
Phonons at Gamma
   &inputph
    prefix = 'Si',
    tr2_ph = 1.0d-14,
    amass(1) = 28.0855,
! outdir = './tmp'
    fildyn = 'Si.dyn',
   /
   0.0 0.0 0.0
```

We consider only the Γ point:



Brillouin Zone

Dynamical matrix file Si.dyn :

7	***************************************						
	2.333248 [cm-1]	:	49 [THz] =	0.06994	1) =	freq (
Acoustic	311 0.000000)	-0.69	0.000000	0.104520	0.000000	0.006270	(
	311 0.000000)	-0.69	0.000000	0.104520	0.000000	0.006270	(
modes	2.333248 [cm-1]	:	49 [THz] =	0.06994	2) =	freq (
	545 0.000000)	0.04	0.000000	0.286121	0.000000	0.644809	(
	545 0.000000)	0.04	0.000000	0.286121	0.000000	0.644809	(
	2.333248 [cm-1]	:	49 [THz] =	0.06994	3) =	freq (
	775 0.000000)	-0.09	0.000000	-0.638130	0.000000	0.290142	(
	775 0.000000)	-0.09	0.00000	-0.638130	0.000000	0.290142	(
7	6.173789 [cm-1]	5)1 [THz] =	15.47450	4) =	freq (
	325 0.000000)	-0.01	0.000000	-0.463023	0.000000	-0.534305	(
	325 0.000000)	0.01	0.000000	0.463023	0.000000	0.534305	(
Optical	6.173789 [cm-1]	5)1 [THz] =	15.47450	5) =	freq ([
-	460 0.000000)	0.64	0.000000	0.217991	0.000000	-0.202506	(
modes	460 0.000000)	-0.64	0.000000	-0.217991	0.000000	0.202506	(
	6.173789 [cm-1]	5)1 [THz] =	15.47450	6) =	freq (
	322 0.000000)	0.29	0.000000	-0.487944	0.000000	0.416545	(
	322 0.000000)	-0.29	0.000000	0.487944	0.000000	-0.416545	(
	******	*****	********	******	********	********	**

Acoustic sum rule at Γ

Because of the numerical inaccuracies the interatomic force constants do not strictly satisfy the acoustic sum rule (ASR).

ASR comes directly from the continous translational invariance of the crystal. If we translate the whole solid by a uniform displacement, the forces acting on the atoms must be zero.

As a consequence:

=or each
$$\alpha, \beta$$
 and i : $\sum_{\mathbf{L}, j} C_{\alpha i, \beta j}(\mathbf{R}_{\mathbf{L}}) = 0$

As a consequence, the frequencies of the acoustic modes must be zero.

Acoustic sum rule at Γ

Because of the numerical inaccuracies the interatomic force constants do not strictly satisfy the acoustic sum rule (ASR), hence the acoustic frequencies might be slightly different from zero.

However, the ASR can be imposed using the **dynmat.x** program.

The input file is **dynmat.Si.in**:

&input
 fildyn = 'Si.dyn',
 asr = 'simple'
/

File containing the dynamical matrix

 A way to impose the acoustic sum rule (simple, crystal, one-dim, zero-dim)

dynmat.x -i dynmat.Si.in > dynmat.Si.out

The program dynmat.x produces the file dynmat.out which contains the new acoustic frequencies, which are *exactly* equal to zero.

diagonalizing the dynamical matrix ...

q = 0.00	000 0	.0000 0.0	0000			
**********	*******	***********	*******	******	*********	***
freq (1) =	(0.000000)[TI	Hz] =	0.000000 [cm	-1]	
(0.00000	0.000000	- 0.707 107	0.000000	0.000000	0.000000)
(0.000000	0.000000	-0.707107	0.000000	0.000000	0.000000)
freq (2) =	(0.000000)[TI	Hz] =	0.000000 [cm	-1]	
(-0.707107	0.000000	0.000000	0.000000	0.000000	0.000000)
(-0.707107	0.000000	0.000000	0.000000	0.000000	0.000000)
freq (3) =	(0.000000)[TI	Hz] =	0.000000 [cm	-1]	
(0.000000	0.000000	0.000000	0.000000	-0.707107	0.000000)
(0.000000	0.000000	0.000000	0.000000	-0.707107	0.000000)
freq (4) =	15.474329 [TI	Hz]= 5	16.168042 [cm	-1]	
(0.000000	0.000000	0.707107	0.000000	0.000000	0.000000)
(0.000000	0.000000	-0.707107	0.000000	0.000000	0.000000)
freq (5) =	15.474329 [TI	Hz]= 5	16.168042 [cm	-1]	
(0.000000	0.000000	0.000000	0.000000	-0.707107	0.000000)
(0.000000	0.000000	0.000000	0.000000	0.707107	0.000000)
freq (6) =	15.474329 [TI	Hz]= 5	16.168042 [cm	-1]	
(-0.707107	0.000000	0.000000	0.000000	0.000000	0.000000)
(0.707107	0.000000	0.000000	0.000000	0.000000	0.000000)
**********	*******	**********	*******	************	*********	***

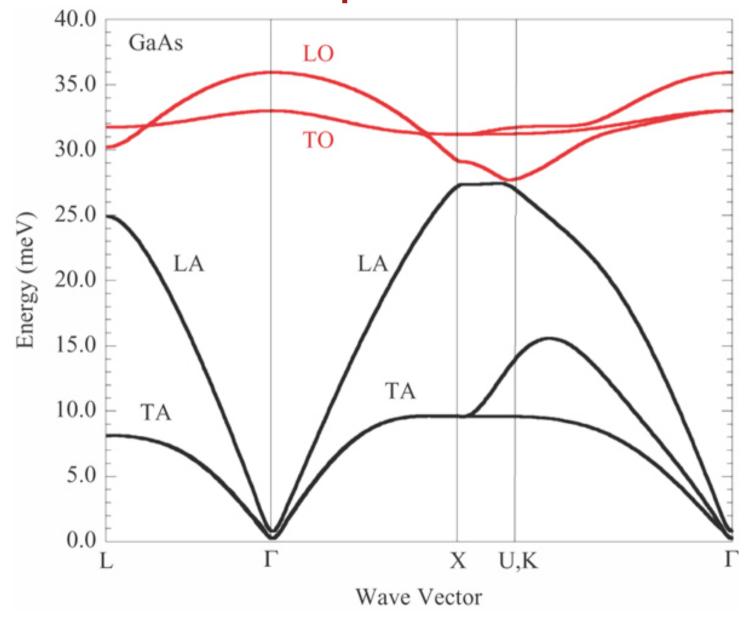
Outline

1. Introduction

2. Exercise 1a: Phonons at Gamma in non-polar materials

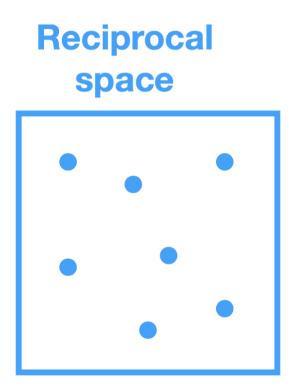
3. Exercise 1b: Phonon dispersion in non-polar materials

- 4. Exercise 2a: Phonons at Gamma in polar materials
- 5. Exercise 2b: Phonon dispersion in polar materials
- 6. Exercise 3: Phonon dispersion of 2D materials (optional)

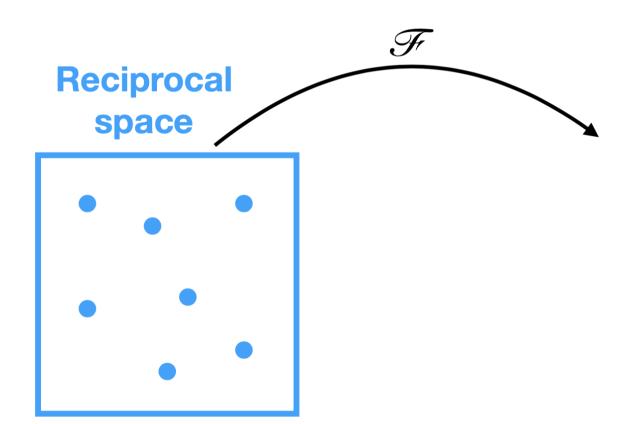




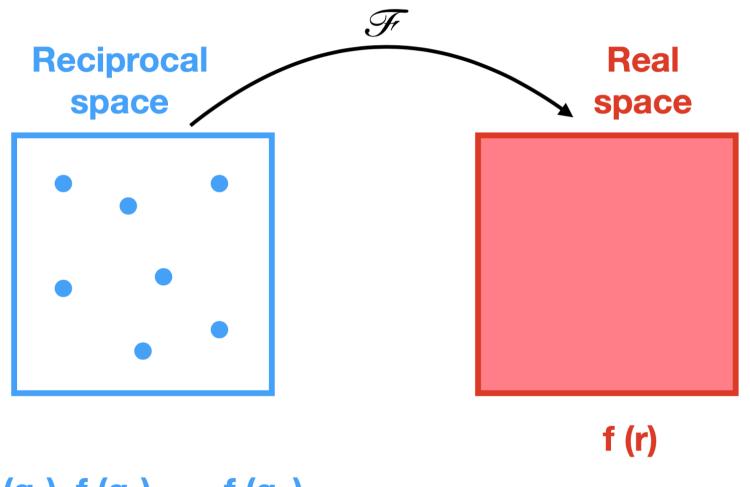




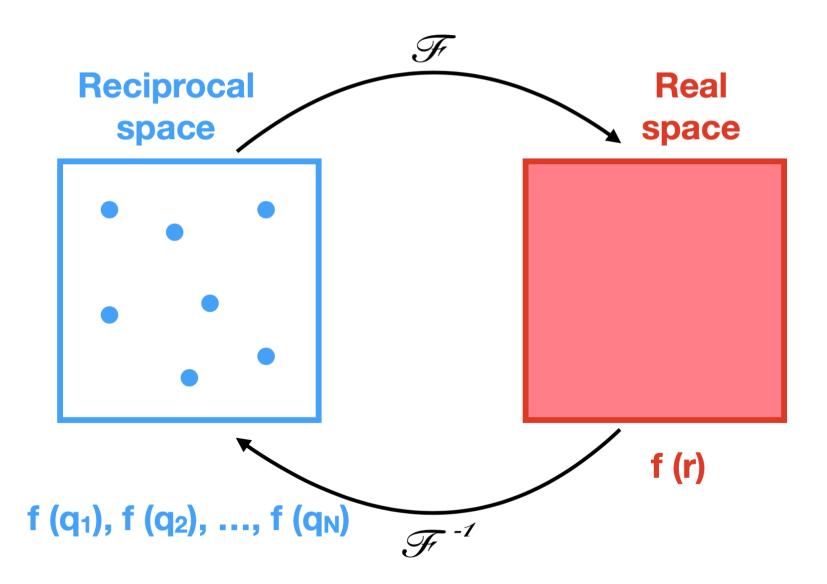
f (q₁), f (q₂), ..., f (q_N)

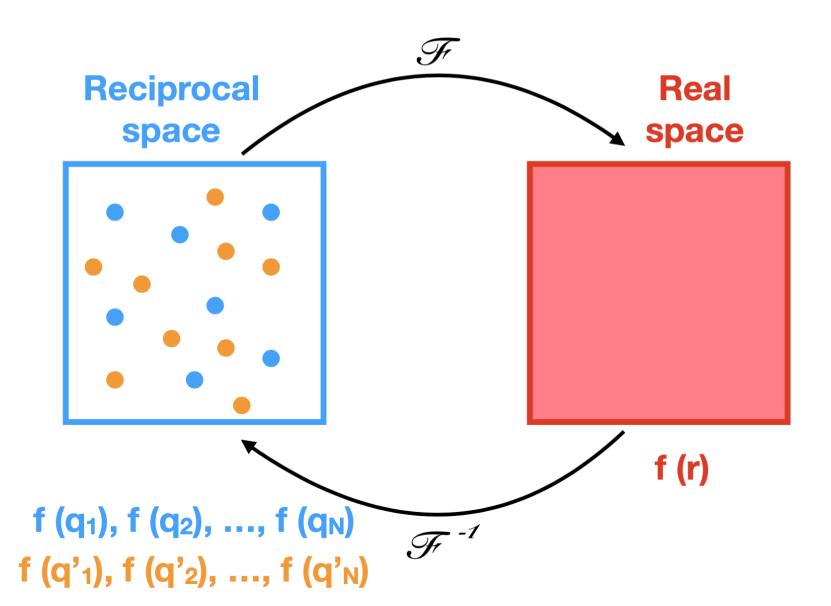


f (q₁), f (q₂), ..., f (q_N)



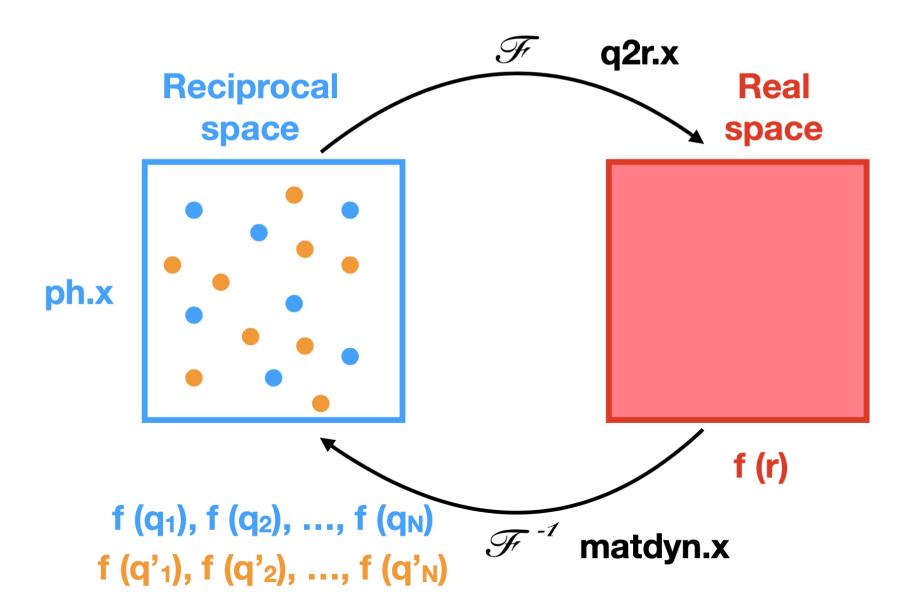
f (q₁), f (q₂), ..., f (q_N)





Important concept: Fourier interpolation F **Reciprocal** Real space space ph.x f (r) f (q₁), f (q₂), ..., f (q_N) \mathcal{F}^{-1} f (q'₁), f (q'₂), ..., f (q'_N)

Important concept: Fourier interpolation F q2r.x **Reciprocal** Real space space ph.x f (r) f (q₁), f (q₂), ..., f (q_N) \mathcal{F}^{-1} f (q'₁), f (q'₂), ..., f (q'_N)



Go to the directory with the input files:

cd ~/QE-2021/Day-5/example1b

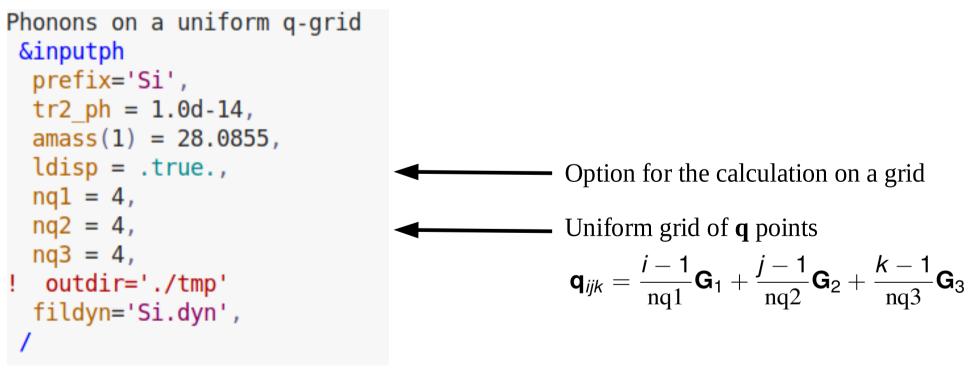
In this directory you will find:

- README . md File describing how to do the exercise
- pw.Si.in Input file for the SCF ground-state calculation
- ph.Si.in Input file for the phonon calculation on a uniform q-grid
- q2r.Si.in Input file for calculation of Interatomic Force Constants
- matdyn.Si.in Input file for Fourier Interpolation for various q points
- plotband.Si.in Input file for plotting a phonon dispersion
- reference Directory containing the reference results

Step 1. Perform a SCF calculation for silicon *at the equilibrium structure* using the **pw.x** program.

pw.x -i pw.Si.in > pw.Si.out

Step 2. Perform a phonon calculation on a uniform grid of **q** points using the **ph.x** program.



ph.x -i ph.Si.in > ph.Si.out

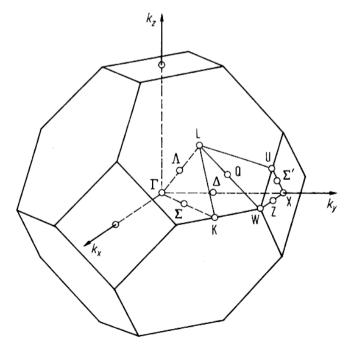
Step 1. Perform a SCF calculation for silicon *at the equilibrium structure* using the **pw.x** program.

pw.x -i pw.Si.in > pw.Si.out

Step 2. Perform a phonon calculation.

```
Phonons on a uniform q-grid
&inputph
  prefix='Si',
  tr2_ph = 1.0d-14,
  amass(1) = 28.0855,
  ldisp = .true.,
  nq1 = 4,
  nq2 = 4,
  nq3 = 4,
! outdir='./tmp'
  fildyn='Si.dyn',
/
```

We sample the Brillouin zone with a uniform grid of 4x4x4 **q** points.



• 4x4x4 = 64 q-points => Use of symmetry => 8 non-equivalent **q** points

The file **Si.dyn0** contains a list of the non-equivalent **q** points (8, in this case).

	4	4	4	(q-grid)	
	8			(number o	non-equivalent q-points)
q1 =	0.0	0000	00000	000000E+00	0.0000000000000E+00 0.000000000000E+00
q2 =	-0.2	25000	00000	000000E+00	0.2500000000000E+00 -0.250000000000E+00
q3 =	0.5	50000	00000	000000E+00	-0.5000000000000E+00 0.500000000000E+00
q4 =	0.0	0000	00000	000000E+00	0.5000000000000E+00 0.000000000000E+00
q5 =	0.7	75000	00000	000000E+00	-0.2500000000000E+00 0.750000000000E+00
q6 =	0.5	50000	00000	000000E+00	0.0000000000000E+00 0.500000000000E+00
q7 =	0.0	0000	00000	000000E+00	-0.1000000000000E+01 0.000000000000E+00
q8 =	-0.5	50000	00000	000000E+00	-0.10000000000000000E+01 0.0000000000000000000000000000000000

 The phonon code ph.x generates a file for every non-equivalent q point (Si.dyn1, Si.dyn2, ..., Si.dyn8), which contain information about dynamical matrices, phonon frequencies and atomic displacements.

Step 3. Calculation of the Interatomic Force Constants (IFC) using the **q2r.x** program.

Fourier transforms of IFC's :

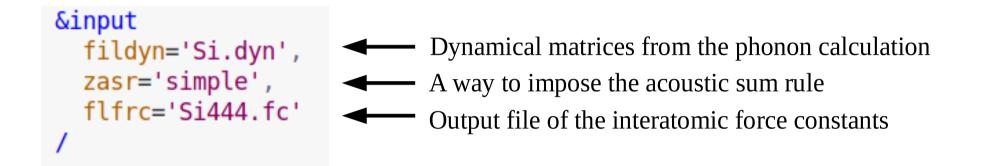
$$\tilde{C}_{\boldsymbol{s}\alpha,\boldsymbol{s}'\beta}(\boldsymbol{\mathsf{q}}_{\boldsymbol{i}\boldsymbol{j}\boldsymbol{k}}) = \frac{\partial^2 E_{tot}}{\partial \tilde{\boldsymbol{\mathsf{u}}}_{\boldsymbol{s}\alpha}^*(\boldsymbol{\mathsf{q}}_{\boldsymbol{i}\boldsymbol{j}\boldsymbol{k}})\partial \tilde{\boldsymbol{\mathsf{u}}}_{\boldsymbol{s}'\beta}(\boldsymbol{\mathsf{q}}_{\boldsymbol{i}\boldsymbol{j}\boldsymbol{k}})}$$

 $\alpha,\beta~$ are Cartesian components, and $\textit{\textbf{S}},\textit{\textbf{S}}'~$ are atomic indices.

$$C_{s\alpha,s'\beta}(\mathbf{R}_{lmn}) = \frac{1}{N_{q}} \sum_{i,j,k} \tilde{C}_{s\alpha,s'\beta}(\mathbf{q}_{ijk}) e^{i\mathbf{q}_{ijk}\cdot\mathbf{R}_{lmn}} \rightarrow C_{s\alpha,s'\beta}(\mathbf{R}_{lmn})$$
Fourier transforms of IFC's on a grid of **q** points nq1 x nq2 x nq3 and in real space

in reciprocal space

```
Input file q2r.Si.in:
```

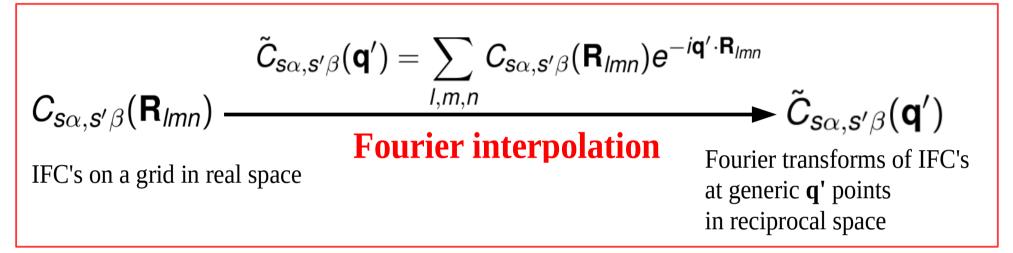


<u>The denser the grid of **q** points, the larger the vectors **R** for which the Interatomic Force Constants are calculated!</u>

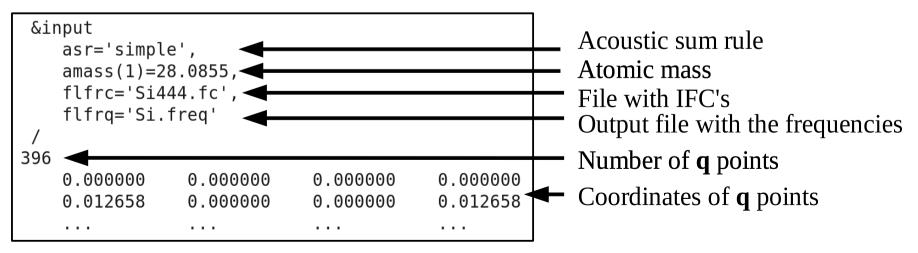
To perform the calculation:

q2r.x -i q2r.Si.in > q2r.Si.out

Step 4. Calculate phonons at generic **q'** points using IFC by means of the code **matdyn.x**



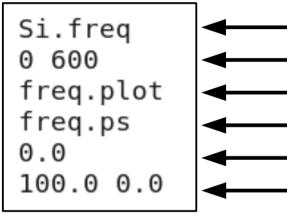
Input file **matdyn**.Si.in:



matdyn.x -i matdyn.Si.in > matdyn.Si.out

Step 5. Plot the phonon dispersion using the plotband.x program and gnuplot.

Input file plotband.Si.in :



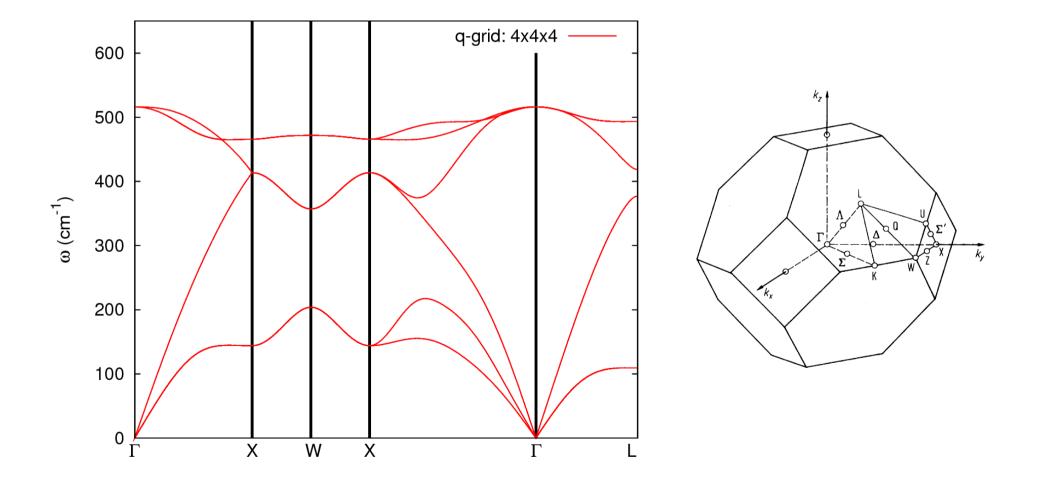
Input file with the frequencies at various q' points
Range of frequencies for a visualization
Output file with frequencies which will be used for plot
Plot of the dispersion (we will produce another one)
Fermi level (needed only for band structure plot)
Freq. step and reference freq. on the plot freq.ps

plotband.x < plotband.Si.in > plotband.Si.out

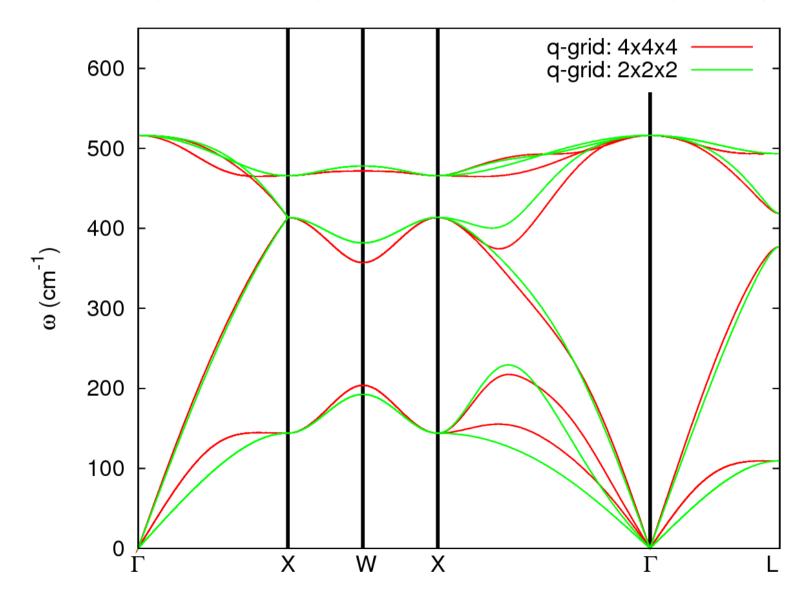
Use **gnuplot** and the file **plot_dispersion.gp** in order to plot the phonon dispersion of silicon (look at the file **experimental_data.dat** for the experimental reference).

You will get a postscript file **phonon_dispersion.eps** which you can visualize.

Phonon dispersion of silicon along some high-symmetry directions in the Brillouin zone (file phonon_dispersion.eps):

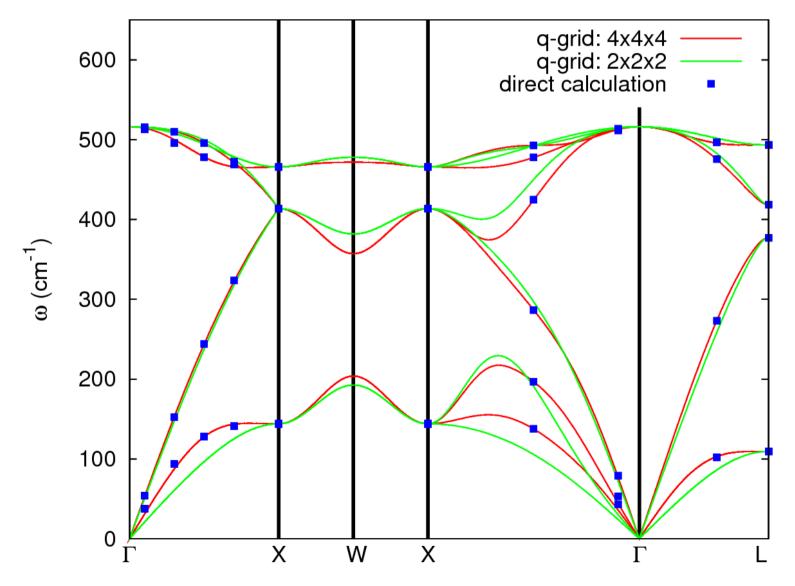


How to determine whether the quality of the Fourier interpolation is satisfactory? -> Compare with the direct calculation (no interpolation)!

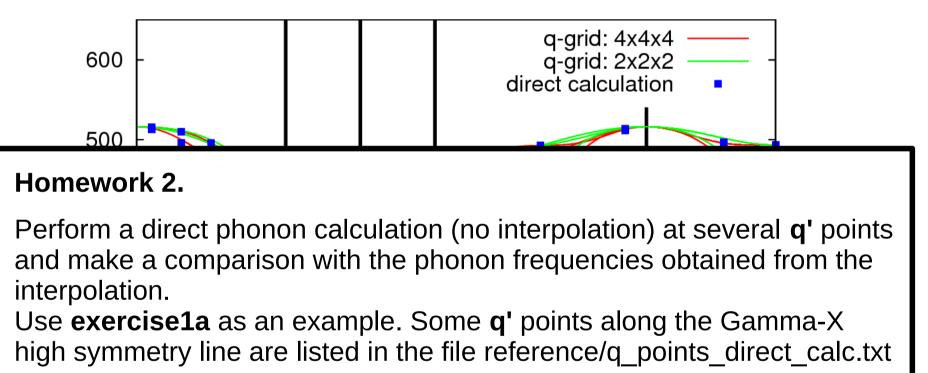


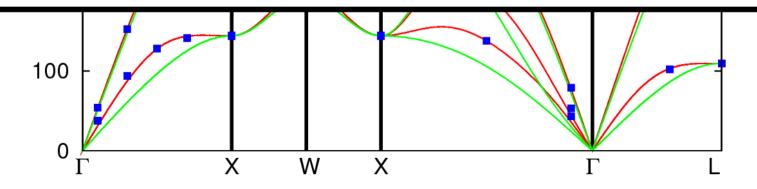
How to determine whether the quality of the Fourier interpolation is satisfactory? - Compare with the direct calculation (no interpolation)! q-grid: 4x4x4 600 q-grid: 2x2x2 500 Homework 1. Perform a phonon dispersion calculation for several **q**-points grids (eg. 2x2x2, 4x4x4, 6x6x6) and compare the dispersions. Do they converge? 100 0 Х Х W

Comparison of the phonon dispersion computed using the Fourier interpolation with the direct calculation at several \mathbf{q} points. The Fourier interpolation for Si with the \mathbf{q} -grid 4x4x4 is satisfactory.

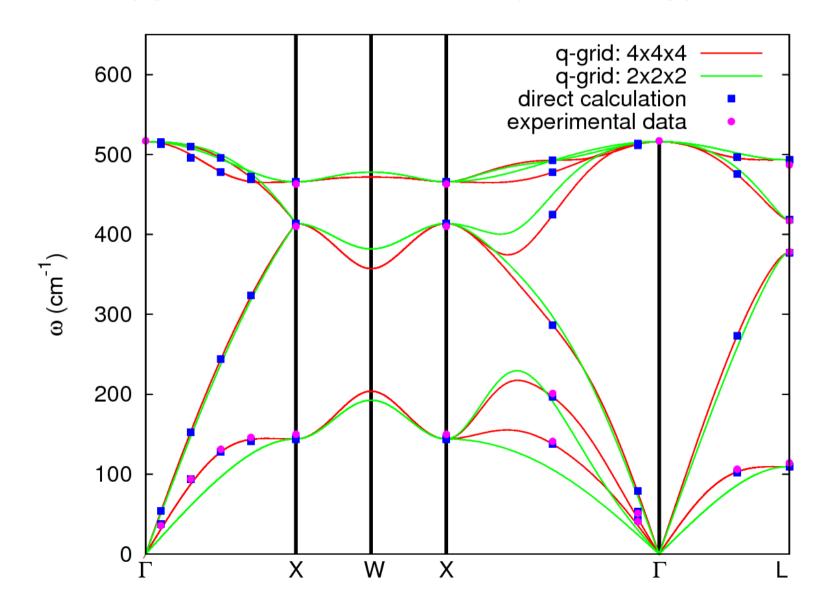


Comparison of the phonon dispersion computed using the Fourier interpolation with the direct calculation at several \mathbf{q} points. The Fourier interpolation for Si with the \mathbf{q} -grid 4x4x4 is satisfactory.





The agreement of *ab initio* calculation of the phonon dispersion with the experimental data is very good when we use the Fourier interpolation on a \mathbf{q} -grid 4x4x4.



Phonon modes visualizer

The phonon modes are not always easy to visualize, especially if we are not at Γ . An online phonon visualizer is very helpful in this regard.

https://www.materialscloud.org/work/tools/interactivephonon

Drag to rotate, scroll to zoom Phonon band structure (select phonon) 600 Settings **Repetitions:** 500 ^ update 3 3 3 400 Camera: x y z (cm-1) Cell: 🔽 on Frequency (005 Amplitude: 0,65 0 200 Vectors: 🔽 on 100 Speed: pause

0

Using the phonon visualiser

Phonon dispersion: pw.Si.in & matdyn.modes

The Fourier interpolation works well if the Interatomic Force Constants (IFC's) are known on a sufficiently large supercell, i.e. on a large enough mesh of **q** points in the phonon calculation.

<u>There are cases when the IFC's are long range and the Fourier interpolation</u> <u>does not work properly:</u>

- When there are Kohn anomalies in metals. In this case the dynamical matrices are not a smooth function of **q** and the IFC's are long range.
- In polar insulators, where the atomic displacements generate long range electrostatic interactions and the dynamical matrix is not analytical for q→0. However, this case can be addressed by calculating the Born effective charges and the dielectric tensor of the material.

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- 3. Exercise 1b: Phonon dispersion in non-polar materials

4. Exercise 2a: Phonons at Gamma in polar materials

- 5. Exercise 2b: Phonon dispersion in polar materials
- 6. Exercise 3: Phonon dispersion of 2D materials (optional)

Polar materials in the q = 0 limit: a macroscopic electric field appears as a consequence of the long-range character of the Coulomb interaction (incompatible with Periodic Boundary Conditions).

A non-analytic term must be added to Interatomic Force Constants at $\mathbf{q} = \mathbf{0}$:

$$ilde{C}_{slpha,s'eta}(\mathbf{q}) = ilde{C}_{slpha,s'eta}^{\mathrm{analytic}}(\mathbf{q}) + rac{4\pi}{\Omega} rac{(\mathbf{q} \cdot \mathbf{Z}_s^{\star})_{lpha}(\mathbf{q} \cdot \mathbf{Z}_{s'}^{\star})_{eta}}{\mathbf{q} \cdot \epsilon_{\infty} \cdot \mathbf{q}}$$

Effective charges Z_s^{\star} are related to polarization **P** induced by a lattice distortion:

$$Z^{\star}_{\boldsymbol{s},\alpha\beta} = \Omega \frac{\partial \mathbf{P}_{\alpha}}{\partial \boldsymbol{u}_{\boldsymbol{s}\beta}}$$

Dielectric tensor $\epsilon_{\infty}^{\alpha\beta}$ is related to polarization **P** induced by an electric field **E** :

$$\epsilon_{\infty}^{\alpha\beta} = \delta_{\alpha\beta} + 4\pi \left. \frac{\partial \mathbf{P}_{\alpha}}{\partial \mathbf{E}_{\beta}} \right|_{u_{s}(q=0)=0}$$

All of the above can be calculated from (mixed) second order derivatives of the total energy.

Go to the directory with the input files:

cd ~/QE-2021/Day-5/example2a

Step 1. Perform a Self-Consistent Field ground-state calculation for the polar semiconductor AIAs.

Step 2. Perform a phonon calculation at Gamma for AIAs.

```
Phonons at Gamma
&inputph
prefix='AlAs',
tr2_ph = 1.0d-14,
amass(1) = 26.98,
amass(2) = 74.92,
! outdir='./tmp'
fildyn='AlAs.dyn',
epsil = .true.
/
0.0 0.0 0.0
```

If .true. will calculate and store the dielectric tensor and effective charges

In the file **ph.AlAs.out** you will find information about the dielectric tensor and effective charges:

Dielectric constant in cartesian axis

	(13.743494475 (-0.000000000 (0.000000000		-0.0000 13.7434 0.0000	0.00000000 0.000000000 13.743494475		
	Effect	ive charges	(d	Force / dE)	in car	rtesian axis
	atom	1 Al				
Ex	(1.88274		0.00000		0.00000)
Ey	(0.00000		1.88274	-	0.00000)
Ez	(0.00000		-0.00000		1.88274)
	atom	2 As				
Ex	(-3.23377		-0.00000		0.00000)
Ey	(-0.00000		-3.23377		0.00000)
Ez	(-0.00000		0.00000	-	3.23377)

Diagonalizing the dynamical matrix

0.00000000 0.00000000 0.00000000) a = (

freq (1) =	0.147836	[THz] =	4.931268	[cm-1]
freq (2) =	0.147836	[THz] =	4.931268	[cm-1]
freq (3) =	0.147836	[THz] =	4.931268	[cm-1]
freq (4) =	11.258207	[THz] =	375.533365	[cm-1]
freq (5) =	11.258207	THz] =	375.533365	[cm-1]
freq (6) =	11.258207	[THz] =	375.533365	[cm-1]
***	k ak	*****	******	*****	****

No LO-TO splitting

Step 3. Impose Acoustic Sum Rule and <u>add the non-analytic LO-TO splitting</u> using the dynmat.x program.

Input file dynmat.AlAs.in:

```
&input
    fildyn = 'AlAs.dyn',
    asr='simple',
    amass(1)=26.98,
    amass(2)=74.92
    q(1) = 1.0,
    q(2) = 0.0,
    q(3) = 0.0
/
```

Direction in the Brillouin zone along which we want to compute the LO-TO splitting

Output file dynmat.out :

	freq (1) =	-0.000000 [TH	z] =			
(0.538473	0.000000	-0.458309	0.000000			
(0.538473	0.000000	-0.458309	0.000000			
	freq (2) =	-0.000000 [TH	z] =			
(0.000000	0.000000	0.000000	0.000000			
(0.000000	0.000000	0.000000	0.000000			
	freq (3) =	0.000000 [TH	z] =			
(-0.458309	0.000000	-0.538473	0.000000			
(-0.458309	0.000000	0.538 473	0.000000			
	freq (4) =	11.257454 [TH	z] =			
(0.000000	0.000000	0.940852	0.000000			
(-0.000000	0.000000	-0.338817	0.000000			
	freq (5) =	11.257454 TH	z] =			
(0.000000	0.000000	0.00000	0.000000			
(0.000000	0.000000	0.000000	0.000000			
	freq (6) =	12.307908 TH	z] =			
(0.940852	0.000000	-0.000000	0.000000			
(-0.338817	0.000000	0.00000	0.000000			

LO-TO splitting

dynmat.x < dynmat.AlAs.in > dynmat.AlAs.out

Outline

1. Introduction

- 2. Exercise 1a: Phonons at Gamma in non-polar materials
- 3. Exercise 1b: Phonon dispersion in non-polar materials
- 4. Exercise 2a: Phonons at Gamma in polar materials
- 5. Exercise 2b: Phonon dispersion in polar materials
- 6. **Exercise 3:** Phonon dispersion of 2D materials (optional)

Go to the directory with the input files:

cd ~/QE-2021/Day-5/example2b

• Calculate phonon dispersion in AIAs following the same steps as in exercise 1b.

• Where necessary insert the missing information in the input files.

Step 1. Perform a SCF ground-state calculation for AIAs using **pw.x**

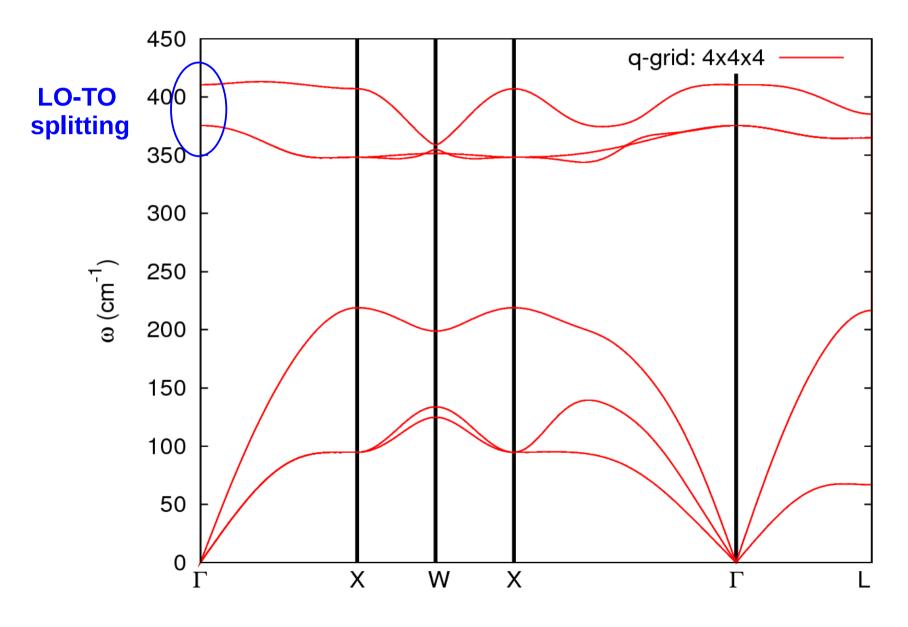
- Step 2. Perform a phonon calculation on a 4x4x4 **q**-grid using **ph.x** (dielectric tensor and effective charges will be calculated)
- Step 3. Perform Fourier Transforms (FT) of $\tilde{C}_{s\alpha,s'\beta}(\mathbf{q})$ in order to get Interatomic Force Constants in real space $C_{s\alpha,s'\beta}(\mathbf{R})$ using $\mathbf{q}^{2\mathbf{r}} \cdot \mathbf{x}$.

A term having the same behaviour for $\mathbf{q} \rightarrow \mathbf{0}$ as the non-analytic term is subtracted from $\tilde{C}_{s\alpha,s'\beta}(\mathbf{q})$ before the FT and re-added to $C_{s\alpha,s'\beta}(\mathbf{R})$, so that no problem related to non-analytic behaviour and related long-rangeness arises in the FT.

Step 4. Calculate phonons at generic **q'** points using Interatomic Force Constants (including the non-analytic term) using the code **matdyn**.**x**

Step 5. Plot the phonon dispersion of AlAs using plotband.x and gnuplot.

The phonon dispersion of AlAs:



Outline

1. Introduction

- 2. Exercise 1a: Phonons at Gamma in non-polar materials
- 3. Exercise 1b: Phonon dispersion in non-polar materials
- 4. Exercise 2a: Phonons at Gamma in polar materials
- 5. Exercise 2b: Phonon dispersion in polar materials
- 6. Exercise 3: Phonon dispersion in 2D materials (optional)

Go to the directory with the input files:

cd ~/QE-2021/Day-5/example3

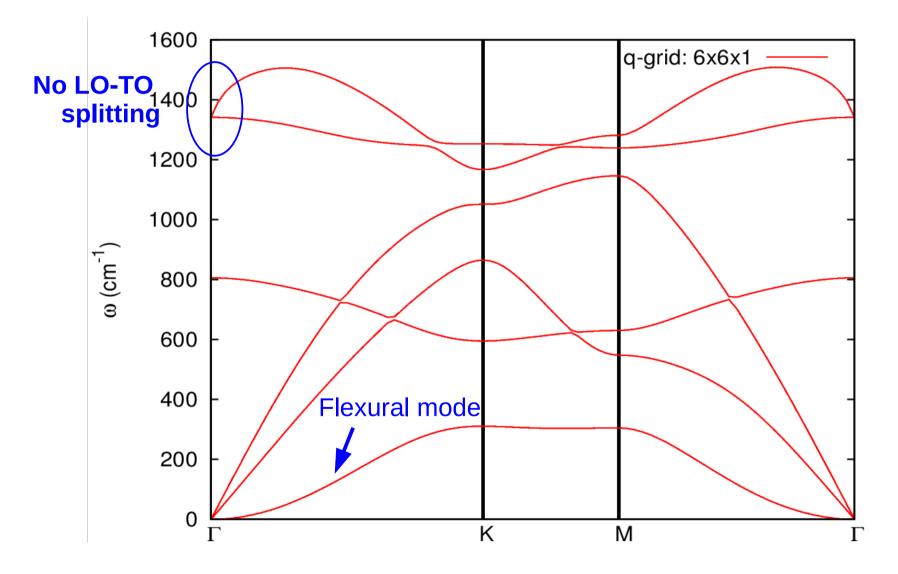
- Calculate phonon dispersion in 2D hexagonal BN following the same steps as in **exercise 1b** and **exercise 2b**.
- Notice that the options assume_isolated='2D' in pw.bn.in and loto_2d=.true. in q2r.bn.in and matdyn.bn.in have been set to properly deal with 2D materials.

Step 1. Perform a SCF ground-state calculation for 2D h-BN using pw.x

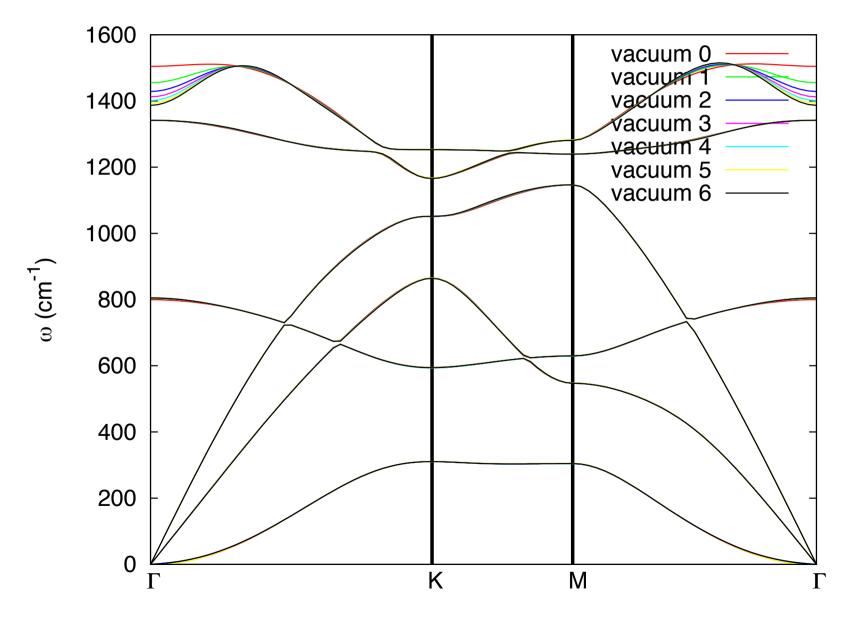
Step 2. Perform a phonon calculation on a 6x6x1 **q**-grid using **ph**.**x**

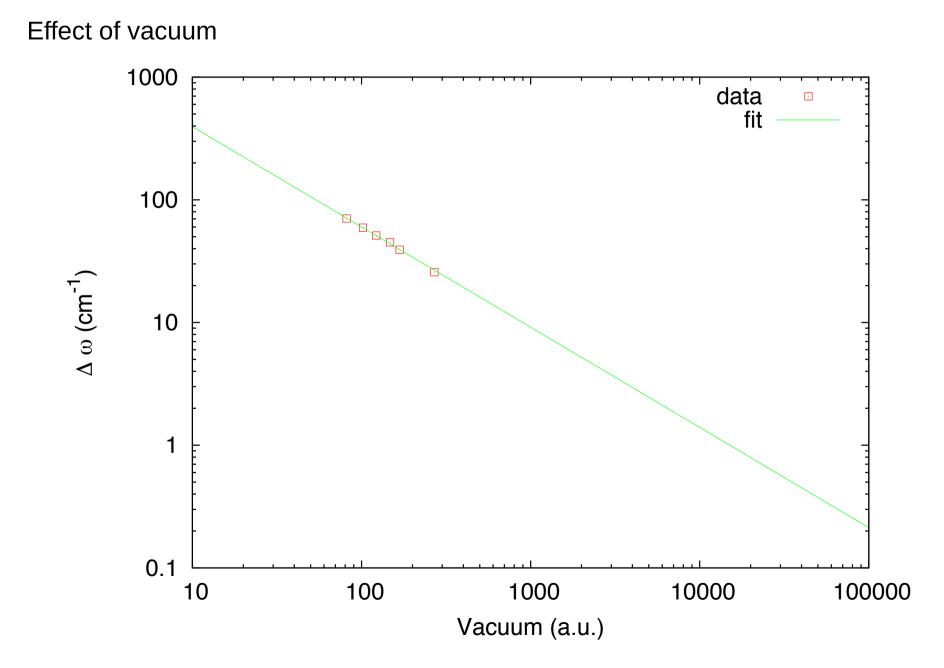
- Step 3. Perform Fourier Transformations (FT) of $\tilde{C}_{s\alpha,s'\beta}(\mathbf{q})$ in order to obtain Interatomic Force Constants in real space $C_{s\alpha,s'\beta}(\mathbf{R})$ using $\mathbf{q}_{2\mathbf{r}}$.
- Step 4. Calculate phonons at generic **q'** points using Interatomic Force Constants using the code **matdyn**.**x**
- Step 5. Plot the phonon dispersion of 2D h-BN using plotband.x and gnuplot.

The phonon dispersion of 2D hexagonal BN:



Effect of vacuum





Bibliography

Books

- 1. N. W. Ashcroft, N. D. Mermin, Solid State Physics (for general discussion)
- **2**. M. Fox, Optical properties of solids (a few words about LO-TO splitting)

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- S. Baroni, P. Giannozzi, and A. Testa, Green's-function approach to linear response in solids, Phys. Rev. Lett. 58, 1861 (1987).
- P. Giannozzi, S. de Gironcoli, P. Pavone, and S. Baroni, Ab initio calculation of phonon dispersions in semiconductors, Phys. Rev. B 43, 7231 (1991).
- **3**. S. Baroni, S. de Gironcoli, A. Dal Corso, and P. Giannozzi, *Phonons and related crystal properties from density-functional perturbation theory*, Rev. Mod. Phys. **73**, 515 (2001).
- 4. T. Sohier, M. Gibertini, M. Calandra, F. Mauri, and N. Marzari, Breakdown of Optical Phonons' splitting in Two-Dimensional Materials, Nano Lett. 2017, 17, 6, 3758-3763.

