

SMR.1566 - 13

**Introductory School on**  
**RECENT DEVELOPMENTS**  
**IN SUPERSYMMETRIC GAUGE THEORIES**

**14 - 25 June 2004**

**S-DUALITY IN  $N=2$  SUPERSYMMETRIC**  
**GAUGE THEORIES**  
**(Part II)**

**Luis Alvarez-Gaumé**  
**Theoretical Physics Division, CERN**

# Lecture 4

We want to determine the quantum geometry of mod. space

$$\begin{aligned} \mathcal{I} &= \frac{\partial \mathcal{A}}{\partial a} = \frac{i}{2\pi} (2N_c - N_f) \ln \frac{a}{\Lambda} \\ \mathcal{F} &= \frac{i}{4\pi} (N_c - N_f/2) a^2 \ln \frac{a^2}{\Lambda^2} \end{aligned} \quad \left. \vphantom{\begin{aligned} \mathcal{I} \\ \mathcal{F} \end{aligned}} \right\}$$

$$K = \text{Im } A_D \bar{A}, \quad \hbar = a n + a_D n m$$

CMS  $\begin{pmatrix} a_D(u) \\ a(u) \end{pmatrix} \quad u \rightarrow \infty \text{ we find monodromy}$

$$\pi_1(\mathcal{M}) \rightarrow G$$

$$\begin{cases} \ln u \rightarrow \ln u + 2\pi i \\ \ln a \rightarrow \ln a + i\pi \end{cases} \quad \left\{ \begin{array}{l} a_D = \frac{i}{\pi} \left( a \ln \frac{u}{\Lambda^2} + a \right) \\ a \end{array} \right.$$

$$\begin{pmatrix} a_D \\ a \end{pmatrix} \rightarrow \begin{pmatrix} -1 & 2 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} a_D \\ a \end{pmatrix}$$

$$T = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}, \quad S = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad P = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$M_\infty = PT^{-2}$$

$\text{Im } \tau$  is harmonic, we need more singularities to avoid  $\text{Im } \tau < 0$ , and also we need to maintain positivity.

$$\begin{aligned} ds^2 &= \text{Im} \frac{da_D}{du} \frac{d\bar{a}}{d\bar{u}} du d\bar{u} \\ &= -\frac{i}{2} \epsilon_{\alpha\beta} \frac{da^\alpha}{du} \frac{d\bar{a}^\beta}{d\bar{u}} du d\bar{u} \end{aligned}$$

$$M = \text{ISL}(2, \mathbb{R}) \quad \text{if } N_f = 0 \quad \text{SL}(2, \mathbb{R})$$

$$Z \text{ invariant } M \in \text{SL}_2(\mathbb{Z}) \cong T$$

Hence the trans. functions are in  $\cong T$ .

$$\text{Im} \int \tau (F + i\bar{F})^2 + \int A_D \wedge dF \quad (2)$$

↓

$$\tau \rightarrow -1/\tau.$$

$$\text{Im} \int A_D(q) \bar{A}(q) = \text{Im} \int A_D \bar{A}$$

$$\begin{pmatrix} A_D(q) \\ A(q) \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \begin{pmatrix} A_D \\ A \end{pmatrix} \Rightarrow \tau(q) = \frac{\alpha\tau + \beta}{\gamma\tau + \delta}$$

$$\begin{pmatrix} q_m & q_e \\ r & s \end{pmatrix} \in \text{SL}_2(\mathbb{Z}) \quad q_m, q_e \text{ are coprime}$$

Imagine  $a_D(q) \approx 0$  at some  $u_0$ . Then the

$$\tau_D \approx -\frac{i}{\pi} \log a_D(q)$$

$$a(q) = - \int \tau_D(q) da_D(q) = \frac{i}{\pi} a_D(q) \log a_D(q) + a_0.$$

$$\left. \begin{array}{l} a_D(q) \rightarrow a_D(q) \\ a(q) \rightarrow a(q) - 2a_D(q) \end{array} \right\} \begin{pmatrix} 1 & 0 \\ -2 & 1 \end{pmatrix}$$

monopole. Then conj. back

$$M(q_m, q_e) = \begin{pmatrix} 1 + 2q_m q_e & 2q_e^2 \\ -2q_m^2 & 1 - 2q_m q_e \end{pmatrix}$$

$$M(m, n) = M_\infty M^{-1}(m', n') \quad | \quad (q_m, q_e) \cdot M = (q_m, q_e)$$

$$m = \pm 1, \quad m' = \pm 1, \quad n = \pm(n' + m')$$

conjugate with  $M_\infty$ .

$$(q_m, q_e) \begin{pmatrix} -1 & 2 \\ 0 & -1 \end{pmatrix} = (-q_m, -q_e + 2q_m)$$

$$(1, 0), \quad (1, -1)$$

$$u = \frac{1}{2} a^2$$

$$u = \infty \quad \left\{ \begin{array}{l} a \approx \sqrt{2u} \\ a_D \approx \frac{i\sqrt{2u}}{\pi} \ln u \end{array} \right.$$

$$u = 1 \quad \left\{ \begin{array}{l} a_D \approx c_0(u-1) \\ a \approx a_0 + \frac{i}{\pi} a_D \ln a_D \end{array} \right.$$

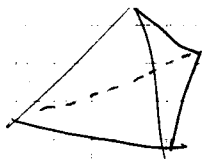
$$u = -1 \quad \left\{ \begin{array}{l} \text{same with } a_D \rightarrow a - a_D \end{array} \right.$$

What about  $m=0$  gluons?

Mathematics.

cpt. Riemann surface.

$$\chi = 2 - 2g \quad g = \text{genus or} \\ \# \text{ handles.}$$



$$V = 4 \quad F = 4 \quad E = 6 \\ \chi = 4 - 6 + 4 = 2$$



$g=0$



$g=1$



$g=2$

covering

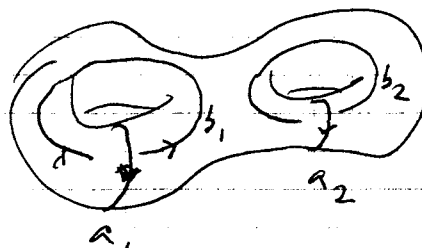
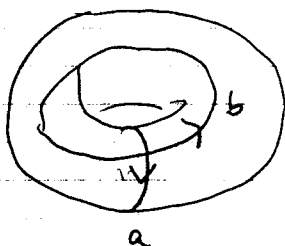
$S^2$

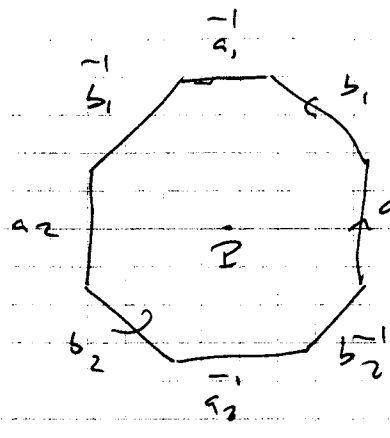
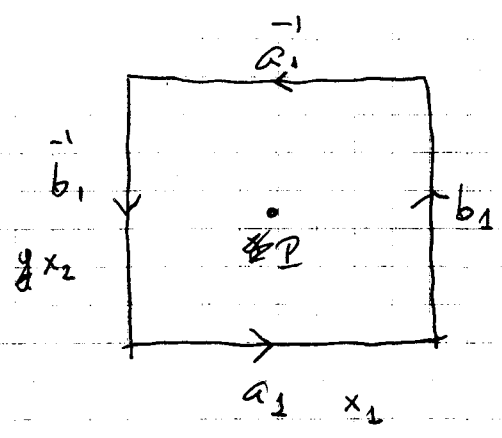
$\mathbb{C}$

$\mathbb{H}$

uniformization Thms.

$g \geq 1$  not simply connected.  $H^1 \neq 0, \pi^1 \neq 0.$





$$\prod_i a_i b_i \bar{a}_i \bar{b}_i = 1$$

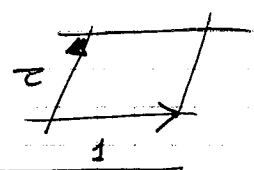
Think of them as cycles or as loops  $S^1 \rightarrow \Sigma$   
 $H^1$   $\pi_1$

Here you can break

Here not.

Differential forms  $dx_1, dx_2$

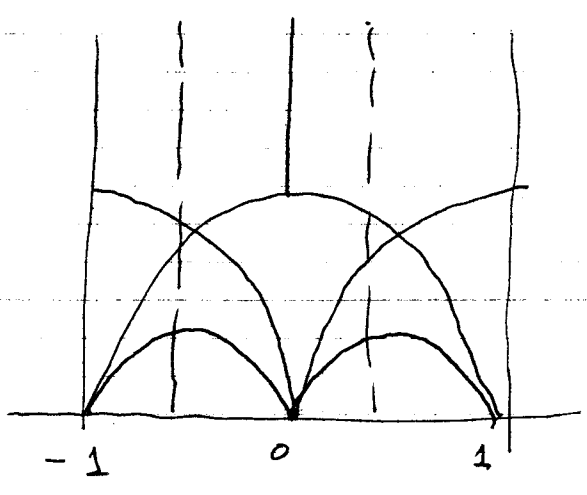
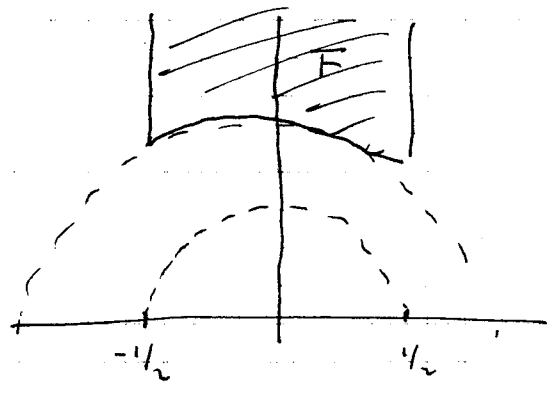
$$T^2 = \mathbb{C} / \{ \mathbb{Z}\tau + \mathbb{Z} \}$$



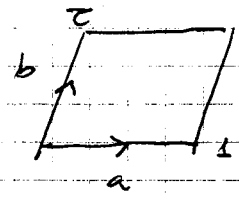
The set of tori is given by  $\boxed{\text{Im } \tau > 0}$   
 and  $SL(2, \mathbb{Z})$  acts as usual

$$\Gamma / \Gamma(2) = SL(2, \mathbb{Z}_2)$$

- $\tau \rightarrow \tau + 1$
- $\tau \rightarrow -1/\tau$
- $\tau \rightarrow 1/2 + \tau$
- $\tau \rightarrow \tau/2 + 1$
- $\tau \rightarrow (\tau+1)/2$



$$dz = dx_1 + i dx_2$$



$$\oint_{\gamma_i} dx_i = 1$$

$$\oint_a dz = 1$$

$$\oint_b dz = i$$

In general  $z$  complex coordinate

$$\omega_i(z) = \oint f_i(z) dz \quad \text{globally defined}$$

Normalize.

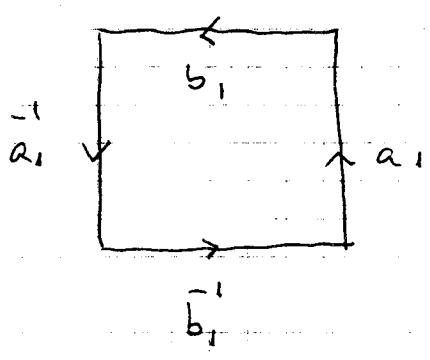
$$\oint_{a_i} \omega_j = \delta_{ij}$$

$$\oint_{b_i} \omega_j = z_{ij}$$

$$dz \wedge dz = 0$$

1).  $z_{ij} = z_{ji}$

2).  $\text{Im } z_{ij} > 0$



$$0 = \int \omega_i \wedge \omega_j$$

$$= \int_{\Sigma} d\varphi_i \wedge \omega_j$$

$$= \int_{\partial \Sigma_c} \varphi_i \omega_j = \int_{a_1} + \int_{b_1} + \int_{a_1^{-1}} + \int_{b_1^{-1}}$$

$$\int_{a_1} \varphi_i \omega_j + \int_{a_1^{-1}} \varphi_i \omega_j = \int_{a_1} \varphi_i \omega_j - \int_{a_1 + b_1} \varphi_i \omega_j$$

$$= \int_{a_1} \varphi_i \omega_j - \int_{a_1} \varphi_i (p + b_1) \omega_j =$$

$$= - \int_a (\varphi_i(p + b_1) - \varphi_i(p)) \omega_j$$

$$= - \oint_b \omega_i \oint_a \omega_j$$

6

$$\int_{\Sigma} \omega_1 \wedge \omega_2 = \sum_i \left( \int_{a_i} \omega_1 \int_{b_i} \omega_2 - \int_{b_i} \omega_1 \int_{a_i} \omega_2 \right)$$

$$\int_{\Sigma} \omega_i \wedge \omega_j = \sum_k \int_{a_k} \omega_i \int_{b_k} \omega_j - \int_{a_k} \omega_j \int_{b_k} \omega_i$$

$$= \sum_k (\delta_{ik} \tau_{kj} - \delta_{jk} \tau_{ki}) = \tau_{ij} - \tau_{ji} = 0$$

Next do the same with:

$$\frac{i}{2} \int \omega \wedge \bar{\omega}$$

$$\omega \wedge \bar{\omega} = \underbrace{f \bar{f}}_{(dx+idy) \wedge (dx-idy)} dZ \wedge d\bar{Z} \\ = 2i dx \wedge dy$$

$$\frac{i}{2} \int_{\Sigma} \omega \wedge \bar{\omega}$$

$$\omega = \omega_i \mathbb{P}_i$$

$$= \frac{i}{2} \sum_{i,j} \underbrace{\tau_{ij}}_{\text{Im } \tau_{ij}} \int_{\Sigma} \omega_i \wedge \bar{\omega}_j \geq 0$$

$$\text{Im } \tau_{ij}$$

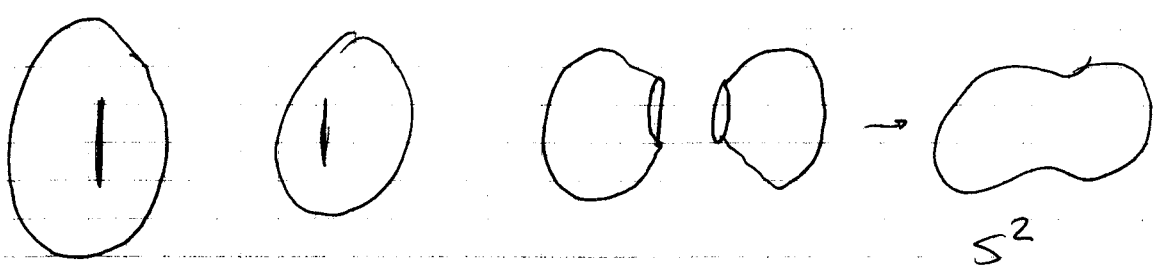
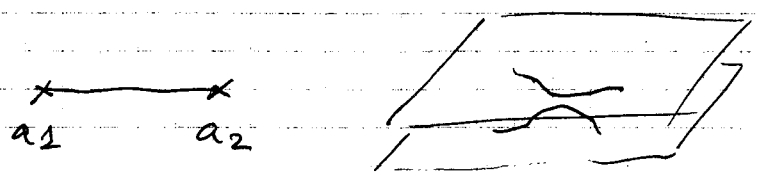
QED

$$\tau_{ij} = \tau_{ji} \\ \text{Im } \tau > 0$$

Every RS is a branch covering over the sphere.

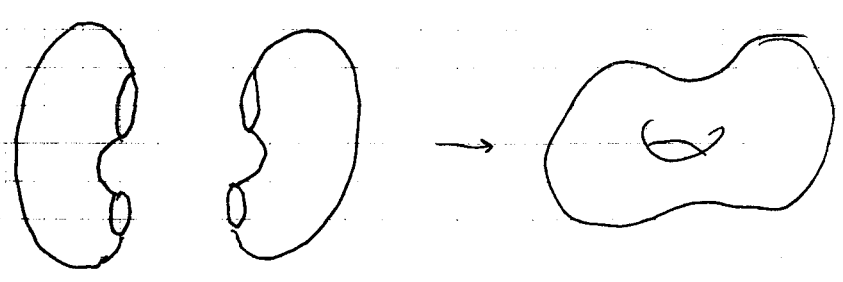
$$M \rightarrow \mathbb{P}^1 = S^2 = \mathbb{C} \cup \{\infty\}$$

$$w = \sqrt{(z-a_1)(z-a_2)} \quad w^2 = (z-a_1)(z-a_2)$$



$$w = \sqrt{(z-a_1)(z-a_2)(z-a_3)}$$

4 branch points  $a_1, a_2, a_3, \infty$ , again two sheets.



In general  $f: M \rightarrow \mathbb{P}^1$   
 $P \rightarrow f(P)$

choose local coord.  $z, S, S = z^n$

$n =$  ramification at  $P$

$b_f(P) \equiv n-1$ . The # of  $P \in M$  ~~with~~  $b_f(P) > 1$  is finite.

Use Euler

$$f: M \rightarrow N \quad (E, F, V)$$

make sure that  $S = \{f(P), P \in M, b_f(P) > 0\}$  is in  $V$ .

Then  $f^{-1}(E, F, V)$  gives a triangulation with  $nF, nE$ , but for the vertices we get  $nV - B$   $B = \sum_{P \in M} b_f(P)$



Hence:

(8)

$$2-2g = n\tilde{F} - nE + nV - B = \\ = n(2-2\gamma) - B$$

For the sphere then  $\gamma = 0$ .

$$2-2g = 2n - B \Rightarrow 2+B-2n = 2g$$

For double covering  $n=$

$$2-2g = n(2-2\gamma) - B$$

$$\gamma = 0 \quad 1-g = n - B/2$$

$$n = 2 \quad -g-1 = -B/2 \quad B = 2g+2$$

$$y^3 = z(z+1)(z-1). \quad \text{no branch at } \infty.$$

$$n = 3, \quad B = 2 \cdot 3 = 6 \quad 1-g = 3 - 6/2 = 0 \\ g = 1$$

$\Sigma$  is hyperelliptic if  $\exists z$  of degree 2 i.e. two first order poles.

A hyperelliptic curve has a f'n of degree  $2g+2$ ,  
 i.e. it has two first order poles

$$z \cong M \rightarrow \mathbb{P}^1 \quad (z) = \frac{Q_3 Q_4}{Q_1 Q_2}$$

now  $g=0$

$$1-g = 2 - B/2 \quad | \quad B = 2g+2$$

let  $P_1, \dots, P_{2g+2}$  be the branched points.  
 $z(P_i) \neq \infty$

$$w = \sqrt{\prod_{j=1}^{2g+2} (z - z(P_j))}$$

$z - z(P_i) = S_i^2$ , hence the divisor of  $w =$

$$(w) = \frac{P_1 \dots P_{2g+2}}{Q_1^{g+1} Q_2^{g+1}}$$

Abelian diff.

$$z^j dz / w \quad j = 0, 1, \dots, g-1$$

$$(z) = \frac{Q_3 Q_4}{Q_1 Q_2} \quad (dz) = \frac{P_1 \dots P_{2g+2}}{Q_1^2 Q_2^2}$$

$$\left( \frac{z^j dz}{w} \right) = \frac{Q_1^{g+1} Q_2^{g+1}}{P_1 \dots P_{2g+2}} \cdot \frac{P_1 \dots P_{2g+2}}{Q_1^2 Q_2^2} \cdot \frac{Q_3^j Q_4^j}{Q_1^j Q_2^j}$$

$$= \frac{Q_1^{g-j-1} Q_2^{g-j-1}}{Q_3^j Q_4^j}$$

$j = 0, 1, \dots, g-1$  holomorphic

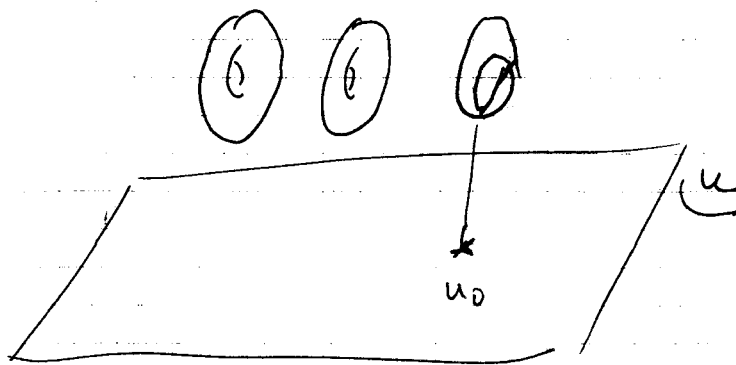
If  $j = g$  however we have:

$$\begin{matrix}
 -1 & -1 & j & j \\
 Q_1 & Q_2 & Q_3 & Q_4
 \end{matrix}$$

on the torus  $x dx/w$  has a double pole at infinity.

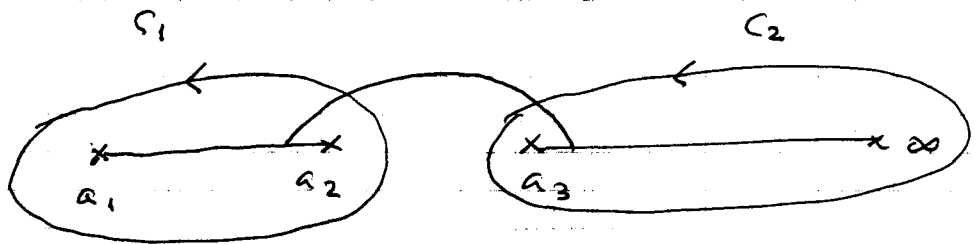
Now we can formulate the solution in detail.

We have seen that  $z_{ij} \Rightarrow \text{Im } z > 0$ . In fact  $z$  is a  $g \times g$  matrix, thus we are looking for a torus. In fact we are looking for a 1-parameter family of tori, parametrized by  $u$ . We also know that at certain points  $z$  get log. singularities. In RS. This can only happen when the surface degenerates, i.e. there are vanishing cycles.



we know how to construct the cycles. In the standard rep. a torus is just

$$w^2 = (z - a_1)(z - a_2)(z - a_3)$$



fix the curve by making sure there is a  $\mathbb{Z}_2$  action:

$$y^2 = (x-1)(x+1)(x-u)$$

more carefully

$$y^2 = (x-\Lambda^2)(x+\Lambda^2)(x-u)$$

We want to find a differential  $\lambda$

$$a(x) \rightarrow \oint_{\gamma} \lambda$$

we want  $\lambda$  to be meromorphic with zero residues so that the pairing is invariant under deformations of  $\gamma$ .

thus we need a  $\lambda$  either hol. or with a 2nd order pole. ~~there are~~ There are:

$$\frac{dx}{y}; \frac{x dx}{y}$$

$$b_i = \oint_{\gamma_i} \lambda_1$$

$$\frac{b_1}{b_2} = z(u)$$

$$\lambda = a_1(u) \lambda_1 + a_2(u) \lambda_2$$

$$a_D = \oint_{\gamma_1} \lambda$$

$$a = \oint_{\gamma_2} \lambda$$

$$\frac{d\lambda}{du} = f(u) \lambda_1 = f(u) \frac{dx}{y}$$

$$\frac{da_D}{du}$$

$$da_D/du = \oint_{\gamma_1} d\lambda/du$$

$$da/du = \oint_{\gamma_2} d\lambda/du$$

we need

$$\frac{d\lambda}{du} = f(u) \lambda_1$$

so that

$$\frac{da_D}{du} = f(u) b_1$$

$$\frac{da}{du} = f(u) b_2$$

$$z = \frac{da_D/du}{da/du} = \frac{b_1}{b_2} = z(u)$$

In fact  $f = -\sqrt{2}/4\pi$  is enough.

$$\lambda_1 = \frac{dx}{\sqrt{(x^2-1)(x-u)}}$$

$$\int du \lambda_1 = -\frac{dx}{\sqrt{x^2-1}} \cdot 2\sqrt{x-u}$$

$$\begin{aligned} \lambda &= -\frac{\sqrt{2}}{4\pi} \dots = \frac{\sqrt{2}}{2\pi} \frac{dx \sqrt{x-u}}{\sqrt{x^2-1}} \\ &= \frac{\sqrt{2}}{2\pi} \frac{dx (x-u)}{y} = \frac{\sqrt{2}}{2\pi} (\lambda_2 - u\lambda_1) \end{aligned}$$

$$\lambda = \frac{\sqrt{2}}{2\pi} (\lambda_2 - u\lambda_1)$$

— x —

Confinement.

$$\sqrt{2} a_D \tilde{M} \tilde{M} + m V(A_D)$$

Use standard Seiberg arguments.

SU(N\_c)

$$\begin{aligned} y^2 &= (x^2 - u)^2 - \Lambda^4 \\ &= (x^2 - u - \Lambda^2)(x^2 - u + \Lambda^2) \\ &= (x - \sqrt{u^2 + \Lambda^2})(x + \sqrt{u^2 + \Lambda^2})(x - \sqrt{u^2 - \Lambda^2}) \\ &\hspace{15em} (x^2 + \sqrt{u^2 - \Lambda^2}) \end{aligned}$$

$$u = \Lambda^2 \quad (x^2 - 2\Lambda^2) x^2$$