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Novel States and Phase Transitions in Highly Correlated Matter  
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**Anisotropic properties of nodal superconductors  
in a magnetic field**

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These are preliminary lecture notes, intended only for distribution to participants

# **Anisotropic properties of nodal superconductors in a magnetic field**

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*USA*

## Collaborators:

**Peter Hirschfeld (U. Florida)**

Tony Houghton

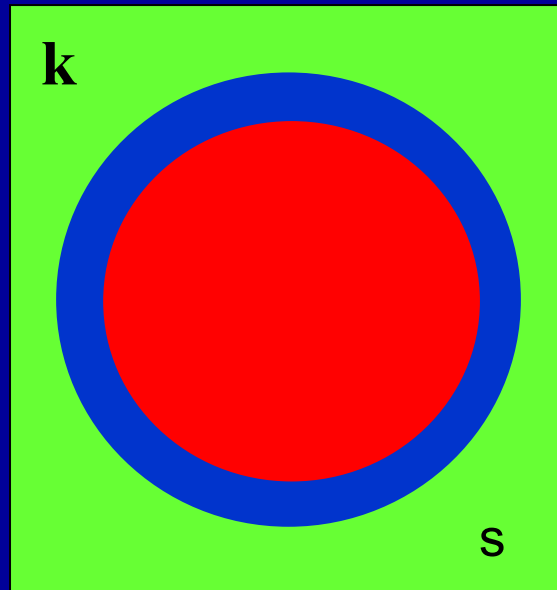
Elisabeth Nicol (U. Guelph)

Jules Carbotte (McMaster)

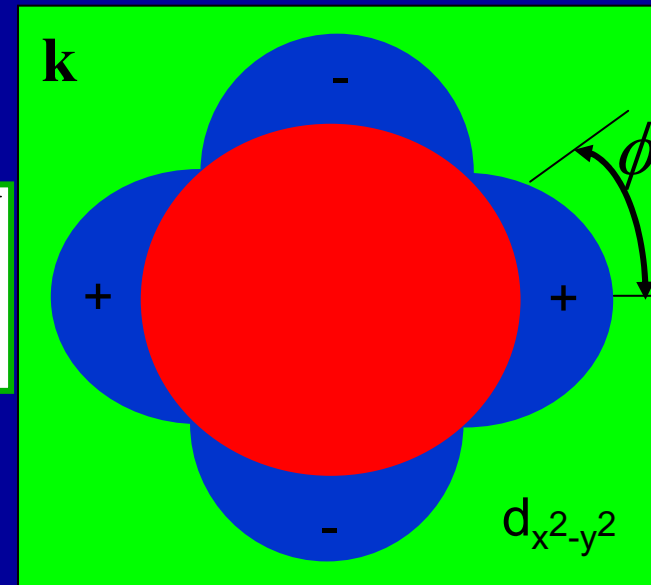
## Many thanks to :

**Yuji Matsuda (U. Tokyo)**

# Anisotropic Superconductors



$$E(\mathbf{k}) = \sqrt{\xi_{\mathbf{k}}^2 + |\Delta(\hat{\mathbf{k}})|^2}$$
$$\xi_{\mathbf{k}} = v_F \cdot (\mathbf{k} - \mathbf{k}_F)$$



**isotropic** gap  $\Delta$

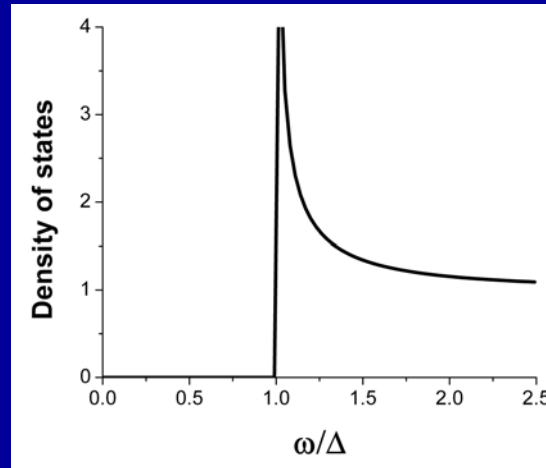
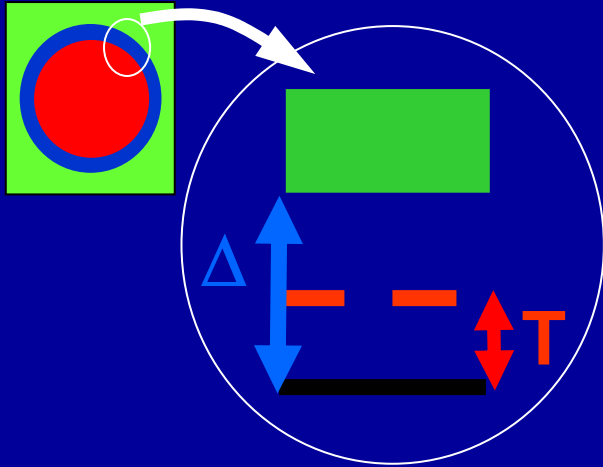
**anisotropic**  $\Delta = \Delta_0 \cos 2\phi$

Anisotropic: heavy fermions, organics, ruthenates, high- $T_c$  ...

**Basic question:** how to determine the symmetry of the superconducting gap in the bulk?

# Low energy excitations

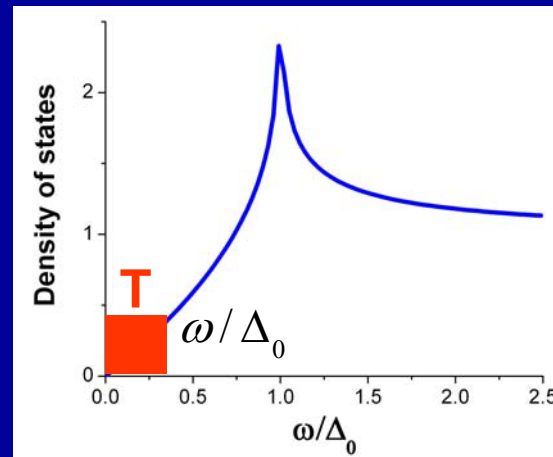
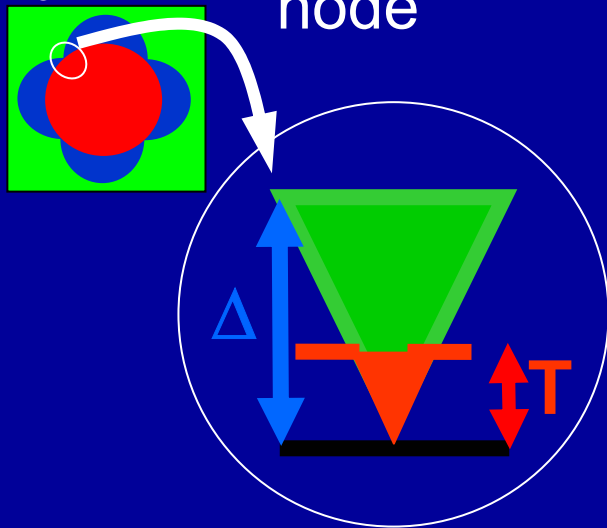
s



No excitations at low T  
 Activated behavior of thermal properties  
 $\exp(-\Delta/T)$

d

node



Density of qp  $\propto T$

Power laws:

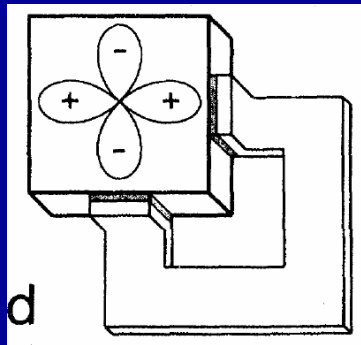
Specific heat  $C(T) \propto T^2$

NMR  $T_1^{-1} \propto T^3$

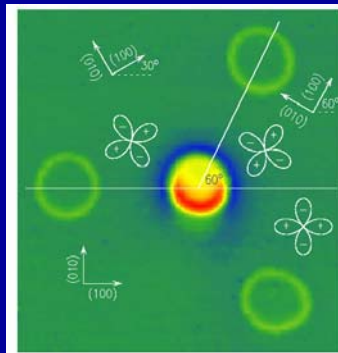
Existence but not position of nodes

# Sign of $\Delta(\mathbf{k})$ : surface probes

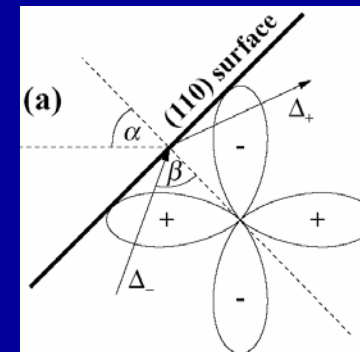
**Cuprates:** change in the phase of the gap  
corner junctions, tricrystal, Andreev bound states



Van Harlingen et al.



Kirtley et al.



C.R. Hu, L. Greene,  
L. Alff

**Not always easy/possible**

**Bulk measurements** bargain: **shape** of the gap in the bulk,

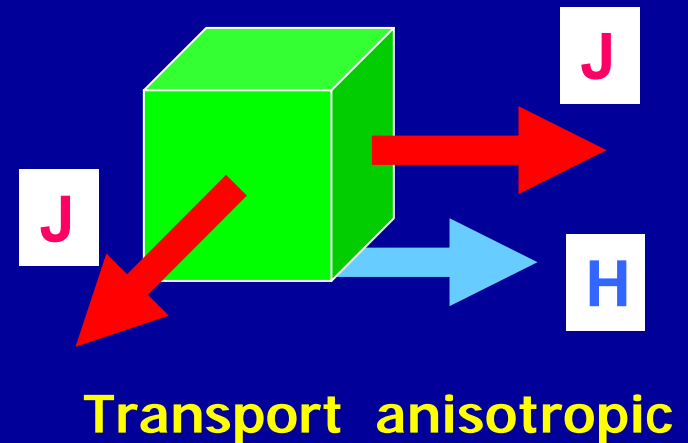
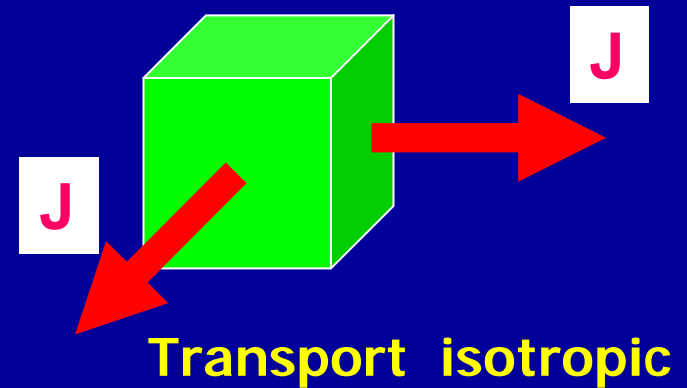
**Bulk:** thermodynamics and transport  
(specific heat, NMR, thermal conductivity)



$$|\Delta(\hat{\mathbf{k}})|$$

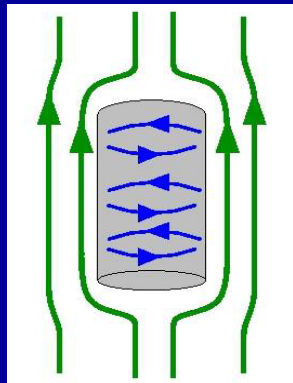
# Why magnetic field?

- Need **directional probe** that provides anisotropy of transport/thermodynamics.
- **Magnetic field?**
- But also need a directional probe that **couples to nodal quasiparticles**
- **How does H do that?**



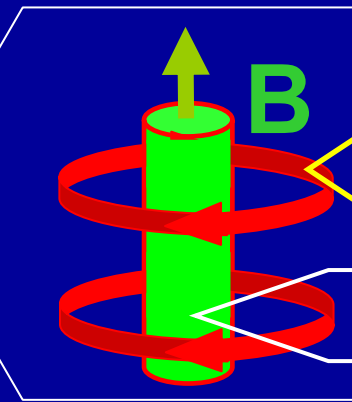
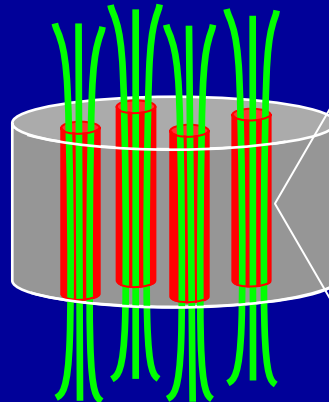
# Magnetic field

- Magnetic field induces vortices (type-II superconductors)



Energy cost  
 $H^2/8\pi.$

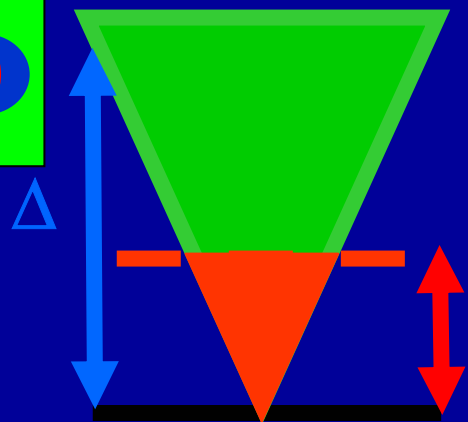
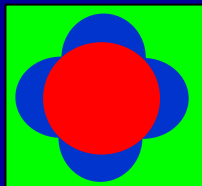
When  
 $H > H_{c1}$



$$j = 2en_s v_s$$

core:  $\Delta = 0$

- In a superconductor with nodes field excites quasiparticles



$$v_s \cdot k$$

Doppler shift  $E'(k, r) = E(k) + v_s(r) \cdot k$

Contribution of Doppler shifted quasiparticles **outside of vortex cores** exceeds that of the bound states in the cores



# “Volovik” Effect

Magnetic field probes mostly nodal quasiparticles.

**Caveat 1.** True for low fields  $H \ll H_{c2}$

**Caveat 2.** True for **line** and **quadratic point** nodes.

For linear point nodes cores and nodes contribute almost equally.

**Caveat 3. Semiclassical approximation:**  $v_s$  nearly uniform  
no Aharonov-Bohm phases (*cf Franz and Tesanovic*)

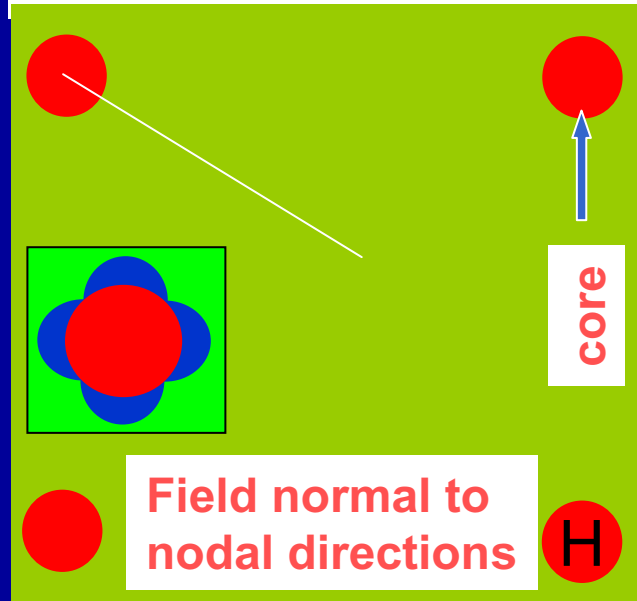
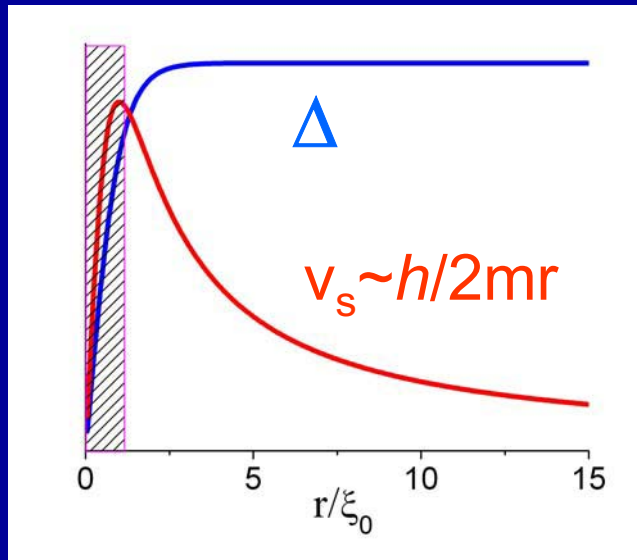
**Caveat 4. Infinite lifetime:** no scattering on vortices.

Good when the mean free path  $l \leq 2R$

Intervortex  
distance

$$R \approx \sqrt{\Phi_0 / H} \approx 450 \text{ \AA} / \sqrt{H, \text{ Tesla}}$$

# Density of states



Characteristic supervelocity

$$v_s \approx \hbar / 2mR$$

Characteristic Doppler energy

$$E_H = v_s k_F \approx v_F / R \propto \sqrt{H}$$

Competes with  $T$ .

Density of states

$$N(\omega, H) / N_0 \approx \omega / \Delta_0 \quad E_H \ll \omega$$

$$N(\omega, H) / N_0 \approx E_H / \Delta_0 \quad E_H \gg \omega$$

Specific heat for  $E_H \gg T$

$$C(T, H) \propto T \sqrt{H} \quad \text{G.Volovik 1993}$$

More formally

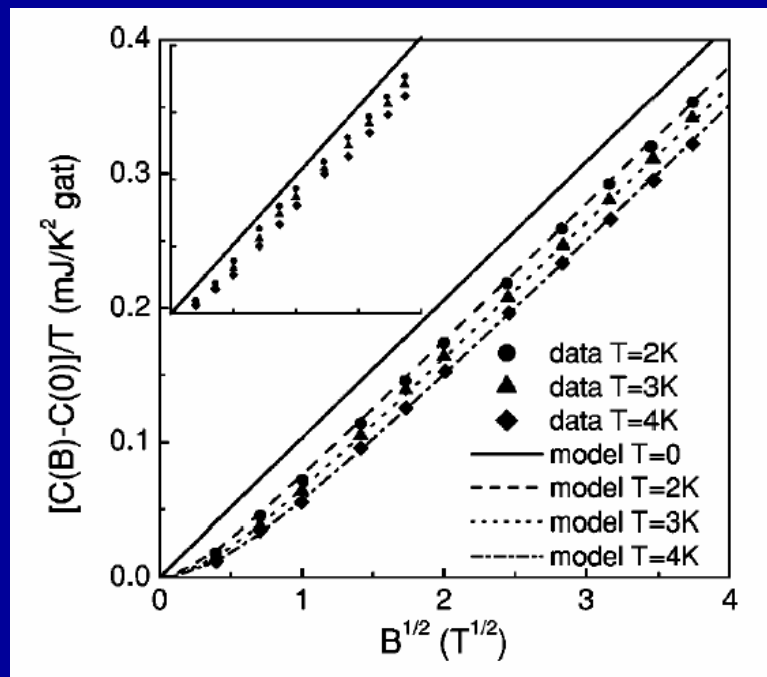
$$G^{-1}(k, \omega_n) = i\omega_n - \xi_k \tau_3 - \Delta(k)(i\tau_2)$$

In field

$$G(k, i\omega_n; r) = G(k, i\omega_n + \mathbf{v}_s(r) \cdot \mathbf{k})$$

Take a realistic  $\mathbf{v}_s(r)$ , compute **local**  $N(r)$ , average

$$N(\omega, H) = A^{-1} \int d^2r N(\omega, r)$$



## Quantitative agreement in YBCO

Expt.:

K. A. Moler et al 1994, B. Revaz et al. 1998

D. Wright et al. 1999, Y. Wang et al. 2001

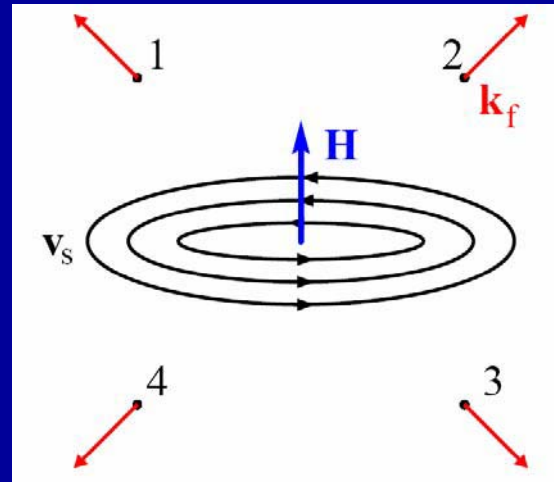
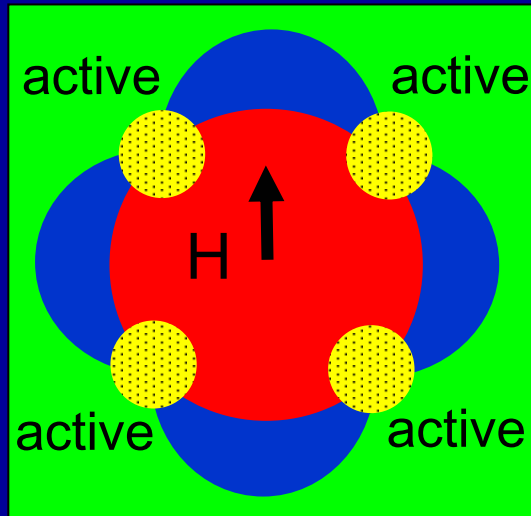
Theory:

Kübert, Hirschfeld 1998,

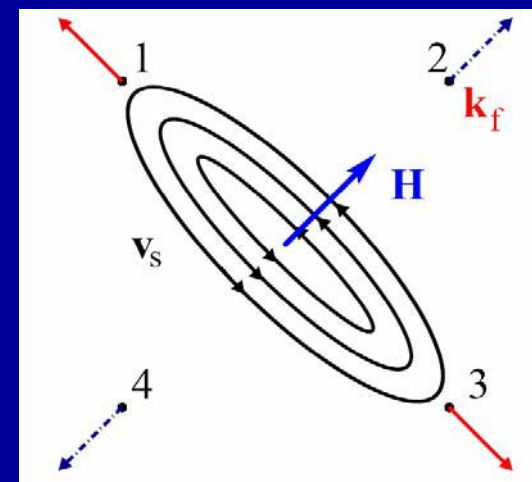
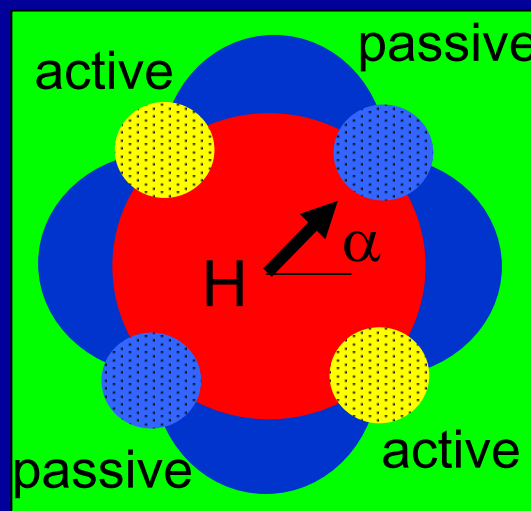
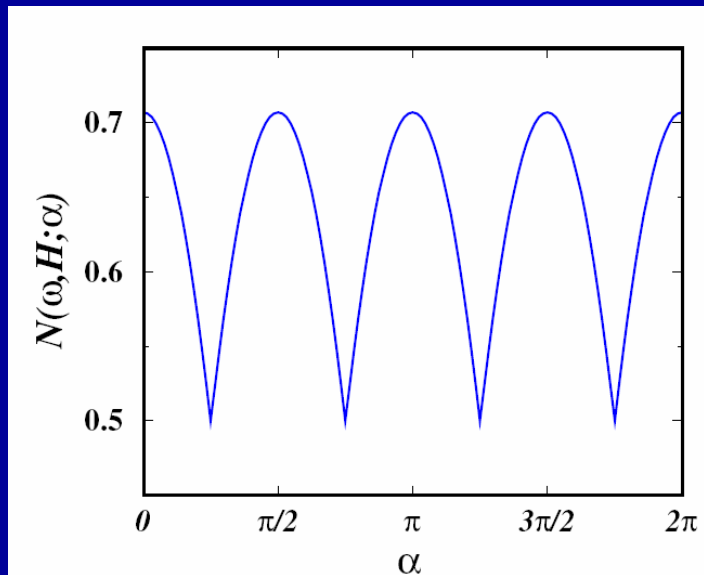
Vekhter et al. 1998-2001

# Anisotropic specific heat

Supervelocity  $\mathbf{v}_s(\mathbf{r}) \perp \mathbf{H}$  and Doppler shift is  $\mathbf{v}_s(\mathbf{r}) \cdot \mathbf{k}_{\text{nodal}}$



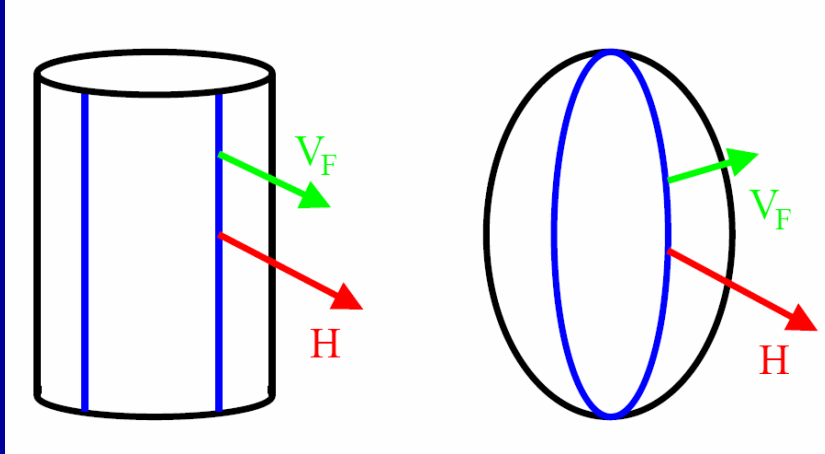
Minima in DOS for  $\mathbf{H} \parallel$  nodes



$$E_H \gg T$$

I. Vekhter et al. 1999-2002,  
see also T. Nakai et al. 2004

# 3D vs 2D



**Anisotropy:**

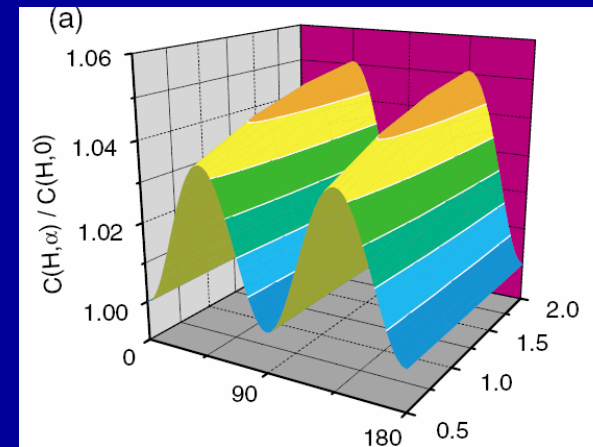
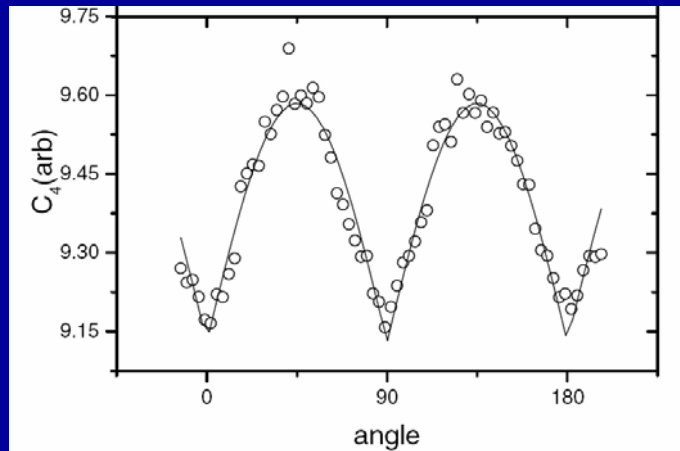
**Amplitude is smaller in 3D** than in **2D**: nodal lines are only partially inactive

**Anisotropy amplitude depends on the shape of the Fermi surface:** (not the salient features)

- what areas have  $\mathbf{v}_F$  parallel to  $\mathbf{H}$ ;
- how close these areas are to the nodes.

I. Vekhter et al `99, K. Maki and H. Won, 2001; K. Maki and P. Thalmeier 2003  
S. Graser, T. Dahm, and N. Schopohl, 2003, ....

# Experiment: borocarbides



$\text{YNi}_2\text{B}_2\text{C}$  and  $\text{LuNi}_2\text{B}_2\text{C}$  : T.Park et al, PRL 2003, 2004

**Nodes or deep minima?**

$$E_H \approx \Delta \sqrt{H / H_{c2}}$$

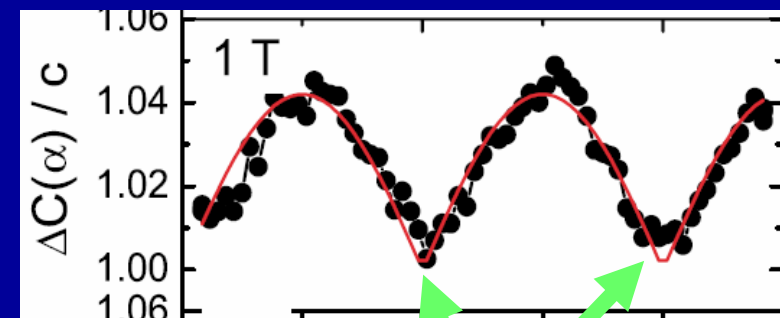
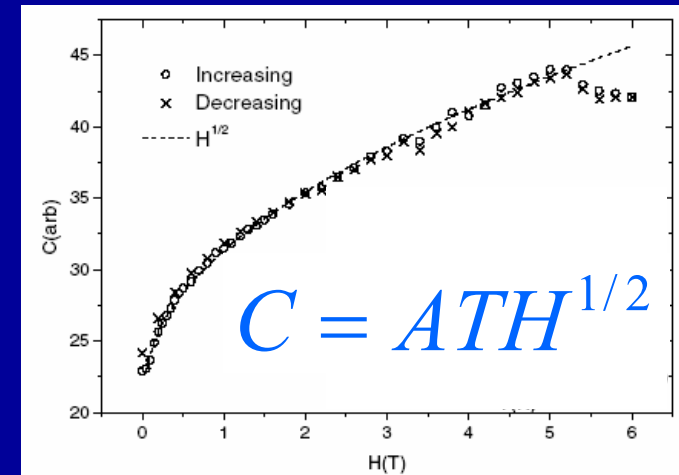
suggests

$$\Delta_{\min} \leq 0.1 \Delta_{\max}$$

- Not a phase sensitive experiment: only anisotropy of  $|\Delta(\hat{\mathbf{k}})|$
- Upper limit:  $E_H$  is a moderately high energy scale
- Combine with other measurements (low-T NMR, penetration depth, etc.) to improve and decide on true nodes.

# Specific heat: summary

- **Anisotropic** superconductor with a known Fermi surface in a field far below  $H_{c2}$ .
- **Measure:**
  - Field dependence of **C/T**;
  - Dependence on the angle



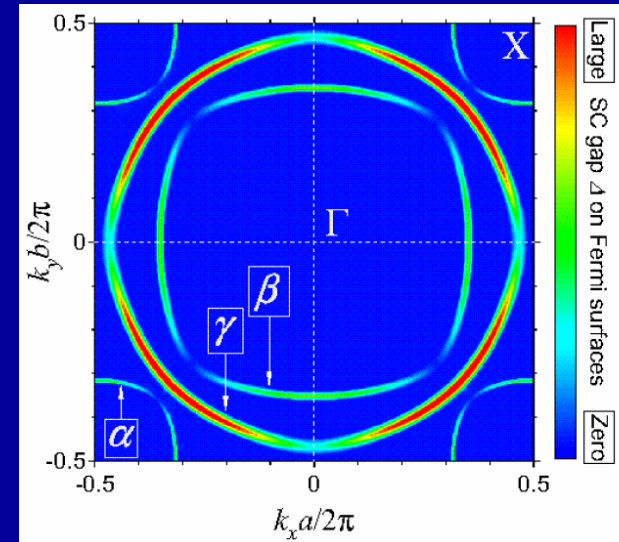
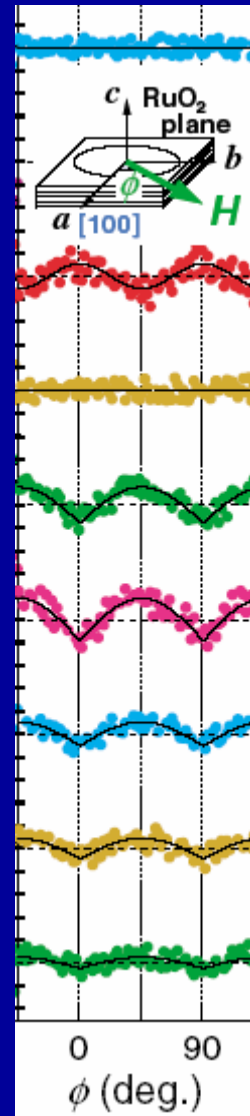
T.Park et al. 04

**When C/T is at a minimum:  $H \parallel v_F$  at 'nodes'.**

# Dirty details

$H_{c2}$

1. **Multiband SC:** interpretation is more difficult.
2. **Don't go too close to  $H_{c2}$ :** it may be anisotropic.



**Sr<sub>2</sub>RuO<sub>4</sub>**

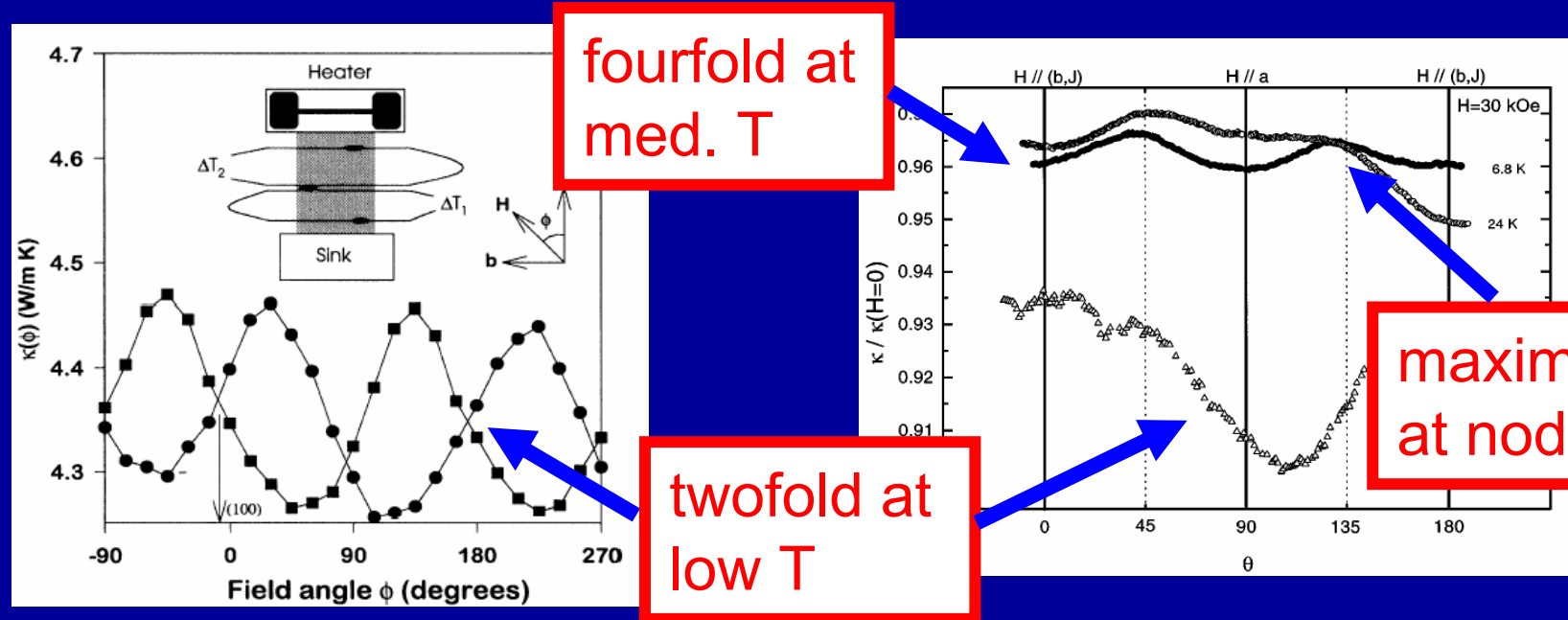
K. Deguchi et al. 2004



# Thermal conductivity

Entropy transport: **only unpaired qps contribute**

Cuprates: experiment predates theory



F. Yu, M. Salamon et al. 1995;

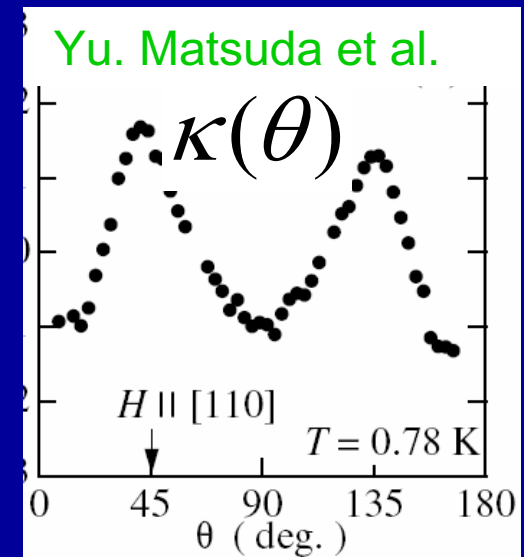
H. Aubin, K. Behnia et al. 1997

**Not at all what is expected from the density of states**

# Transport: a challenge

- Depends on **density of states and lifetime**
- Applied magnetic field
  - **enhances** the local density of states;
  - **modifies** scattering;  
Kübert and Hirschfeld, Vekhter and Hirschfeld
  - **introduces vortex scattering**  
Yu et al., Aubin et al.

$$\kappa \propto TN(0) v_F^2 \tau$$



Low DOS



poor transport for  $H \parallel$  nodes

Reduced scattering

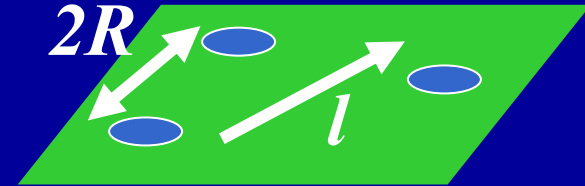


good transport for  $H \parallel$  nodes

**Minima or maxima correspond to nodes?**

# Semiclassical analysis: questions

1. Is there a well-defined **local thermal conductivity**  $\kappa(r)$  ?



• yes, if  $l \leq 2R$

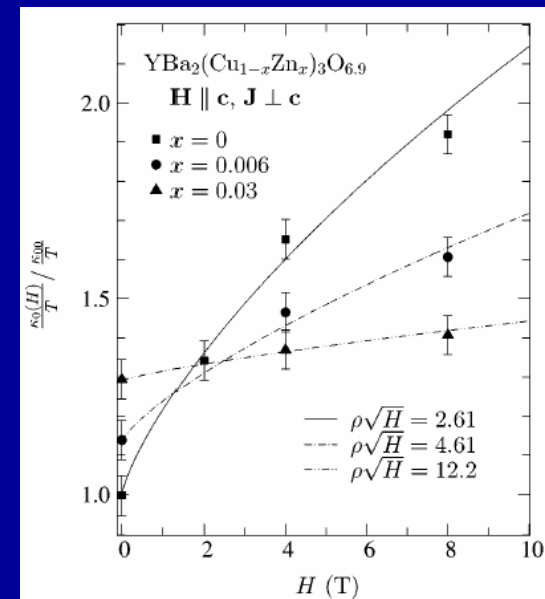
P. Hirschfeld, P. Hirschfeld and I. Vekhter

• possibly otherwise if one takes  $\tau^{-1} = \int d^2r \tau^{-1}(r)$  K. Maki et al.

2. **No vortex scattering** in the model

Fit to data with field normal to the 2D planes in YBCO at low T and H.

May well describe situation when vortex scattering is unimportant



M. Chiao et al.

# What is measured?

$\kappa(r)$ :

$$\frac{\kappa(T, r) / T}{\kappa_n / T_c} = \frac{3}{2\pi^2} \int_0^\infty \frac{d\omega}{T} \left( \frac{\omega}{T} \right)^2 \operatorname{sech}^2 \frac{\omega}{2T} K(\omega)$$

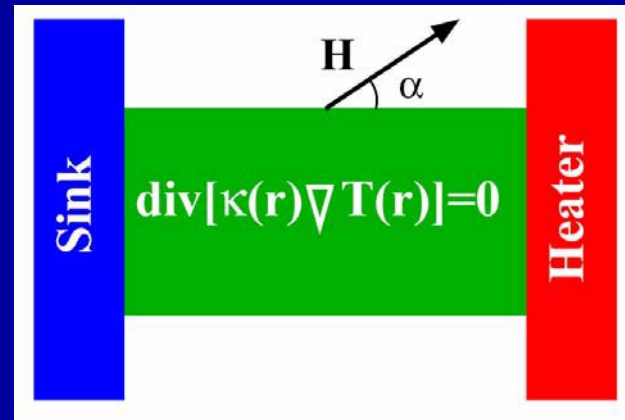
$$K(\omega) = \frac{\Gamma}{\tilde{\omega}' \tilde{\omega}''} \operatorname{Re} \left\langle \frac{(\tilde{\omega}^2 + |\tilde{\omega}|^2 - 2|\Delta_k|^2) k_x^2}{\sqrt{\tilde{\omega}^2 - \Delta_k^2}} \right\rangle_{FS}$$

**Input:** form of the gap,  
impurity scattering,  
Doppler shift.



Mimics gap symmetry:  
4-fold for d-wave etc.

$\kappa$ :



**Input:** local conductivity  
direction of net current.

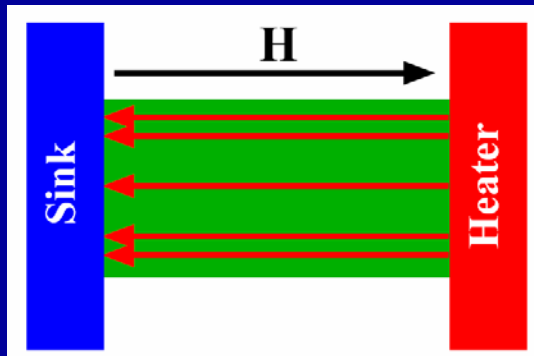


**Does not** mimic gap symmetry  
More complicated dependence

$$\kappa \neq \int d^2 r \kappa(r)$$

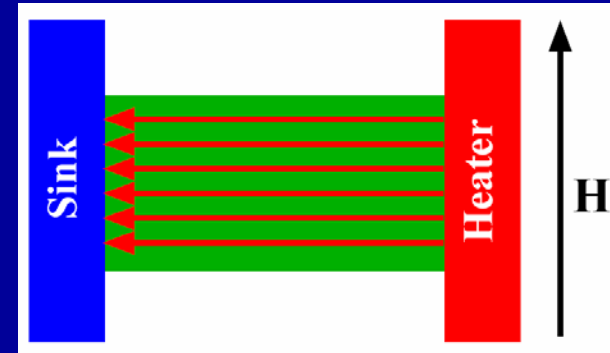
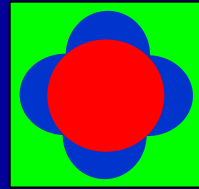
# Twofold angle-dependence

Quasi-2D system: analytic solution possible when



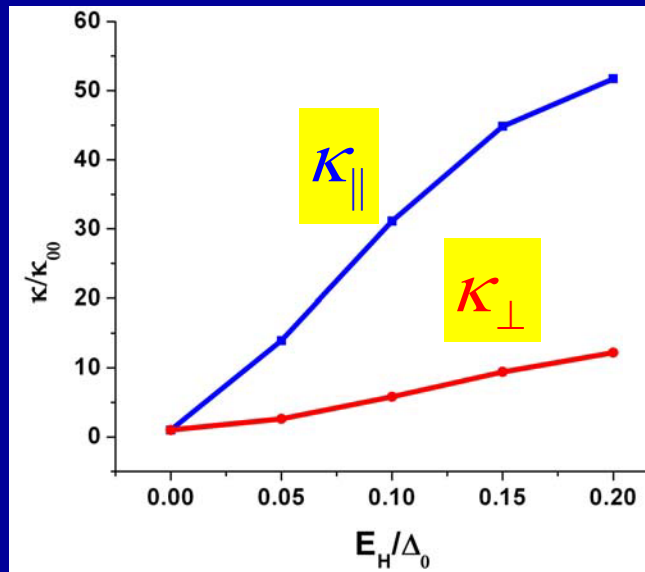
$$\kappa_{\parallel} = \int dy \kappa(y)$$

series

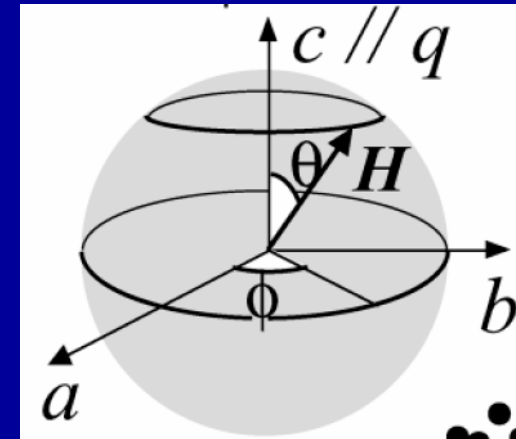
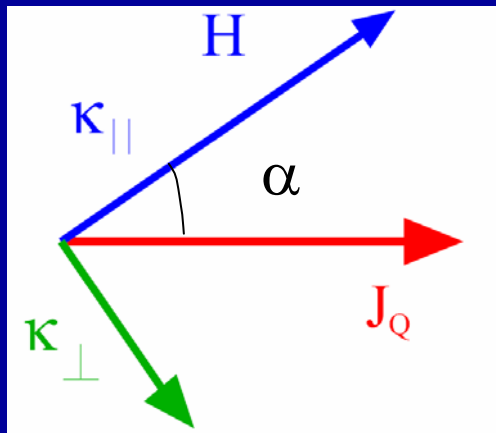


$$\kappa_{\perp}^{-1} = \int dx \kappa^{-1}(x)$$

parallel



# 3D vs 2D



## 2D:

Need to rotate  $H$  wrt  $J$ ,

**Measured  $\kappa$**  is some convolution of

**2-fold** (vortex scattering) and

**nodal patterns** (4-fold)

## 3D:

**conical rotation**

**Directly nodal patterns**

For all angles  $\phi$  convoluted in the same way with vortex scattering.

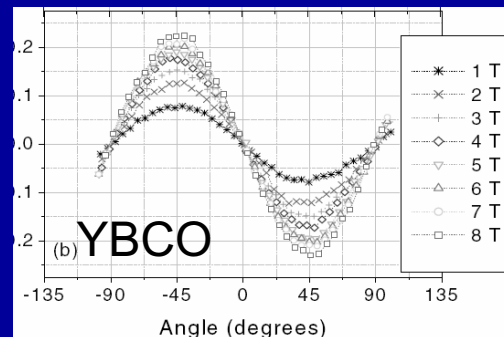
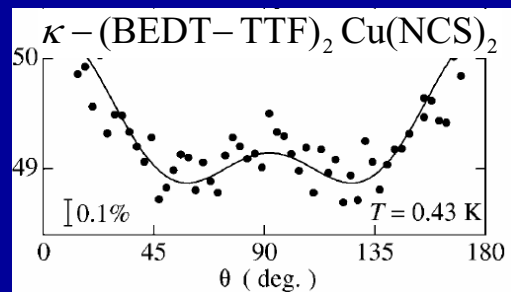
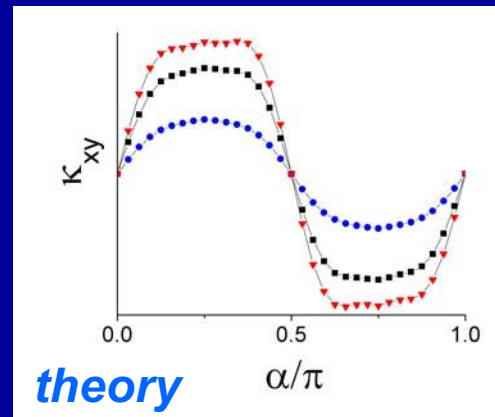
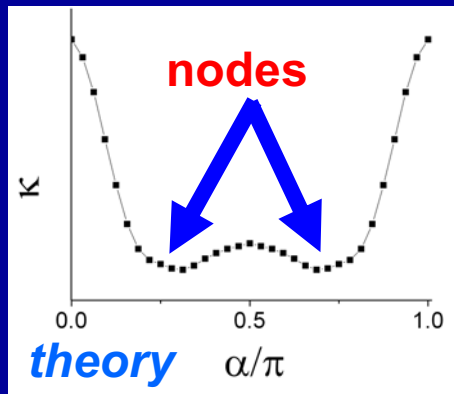
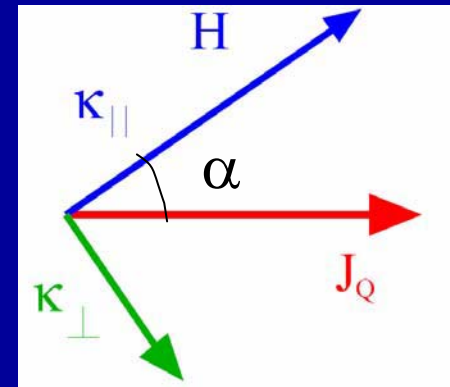
Yu. Matsuda et al. 01-04

# Effective medium approach

- Treat  $\kappa_{\parallel}$  and  $\kappa_{\perp}$  as **principal axes**
- **Steady state technique: fixed  $J_Q$**

$$\kappa_{xx} = \frac{\kappa_{\parallel} \kappa_{\perp}}{\kappa_{\perp} \cos^2 \alpha + \kappa_{\parallel} \sin^2 \alpha}$$

$$\kappa_{xy} = \frac{\kappa_{\parallel} \kappa_{\perp} (\kappa_{\parallel} - \kappa_{\perp}) \sin 2\alpha}{(\kappa_{\perp} \cos^2 \alpha + \kappa_{\parallel} \sin^2 \alpha)^2}$$



Minima almost always remain at the nodes (tentative).

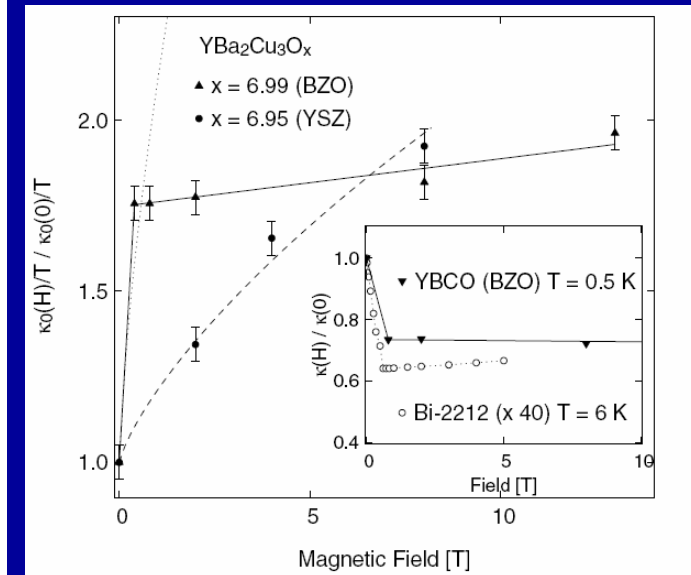
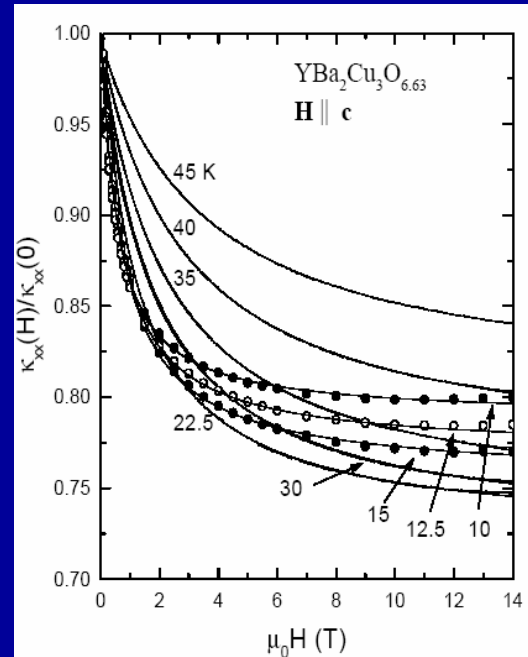
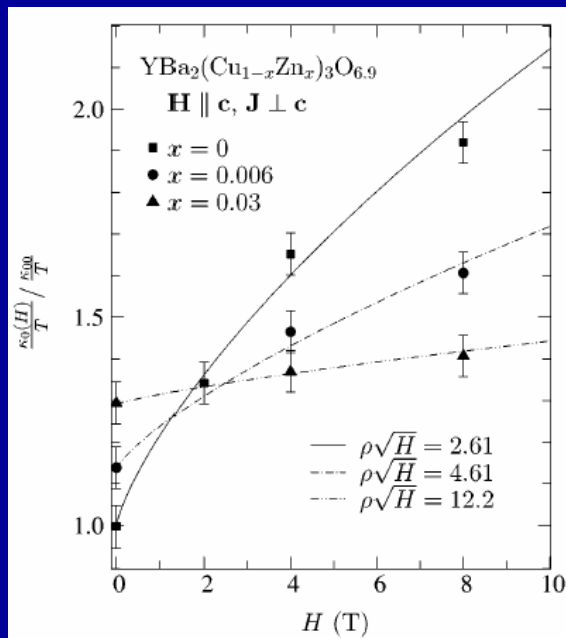
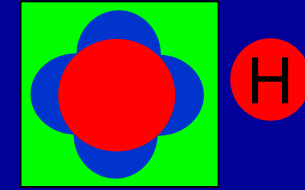
Thermal Hall of the same order as 4-fold part of  $\kappa_{xy}$

K.Izawa et al.

R. Ocaña and P. Esquinazi

I. Vekhter,  
 P.Hirschfeld,  
 unpublished

# Vortex scattering: lessons from cuprates



**Low T,H:** density of states effect dominant. Semiclassical theory.

**High T:** DOS from **T**, **H** - vortex scattering

**Ultrapure sample.** Low T is also dominated by vortex scattering

M. Chiao et al. 1998

K. Krishana et al. 1997

R. Hill et al. 2004



# Theory for $H \parallel c$

**Input:**

vortex **lattice**,

account for supervelocity **to all orders**,

**average Green's function** over vortex unit cell

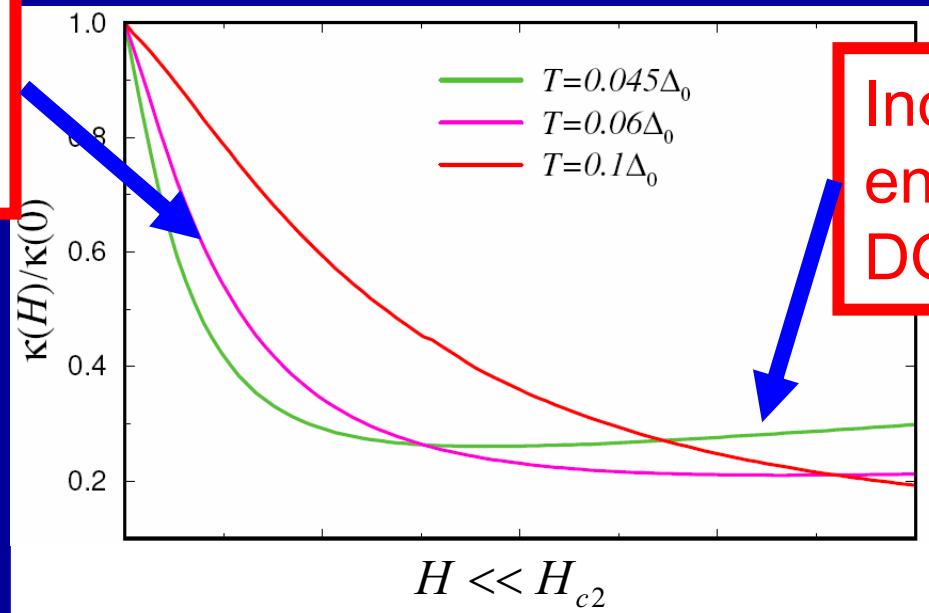
I. Vekhter and A. Houghton'99 based on U. Brandt, W. Pesch, L. Tewordt '68,

$$\tau^{-1}(\omega, \hat{\mathbf{k}}) = \tau_{imp}^{-1}(\omega) + \Delta^2(\hat{\mathbf{k}}) F(\omega / \sqrt{H})$$

Away from node: more scattering

Nodal qp feel impurities

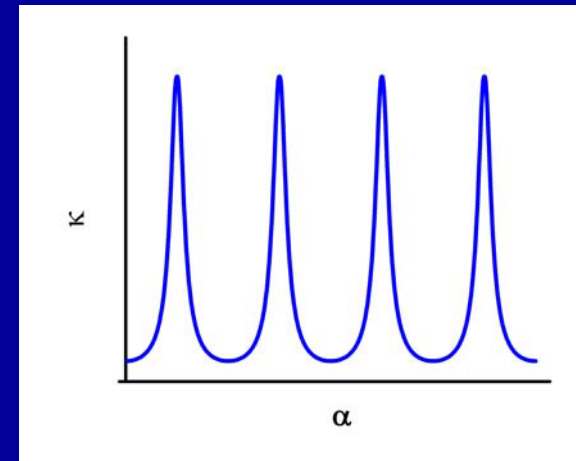
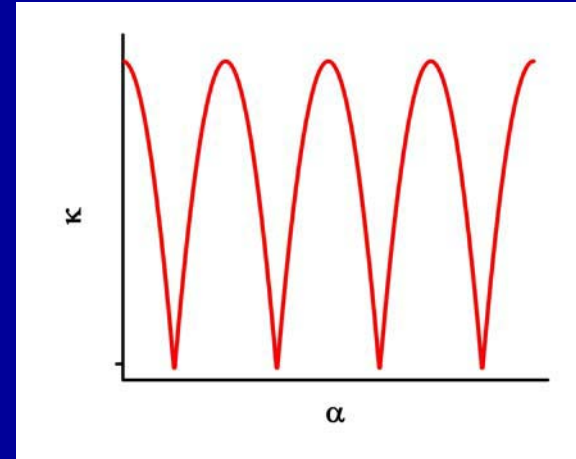
Decrease due to scattering on vortices



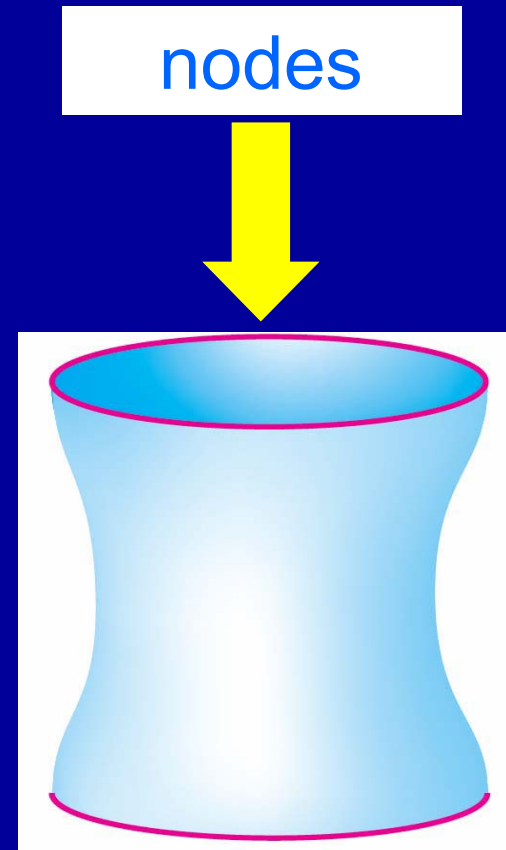
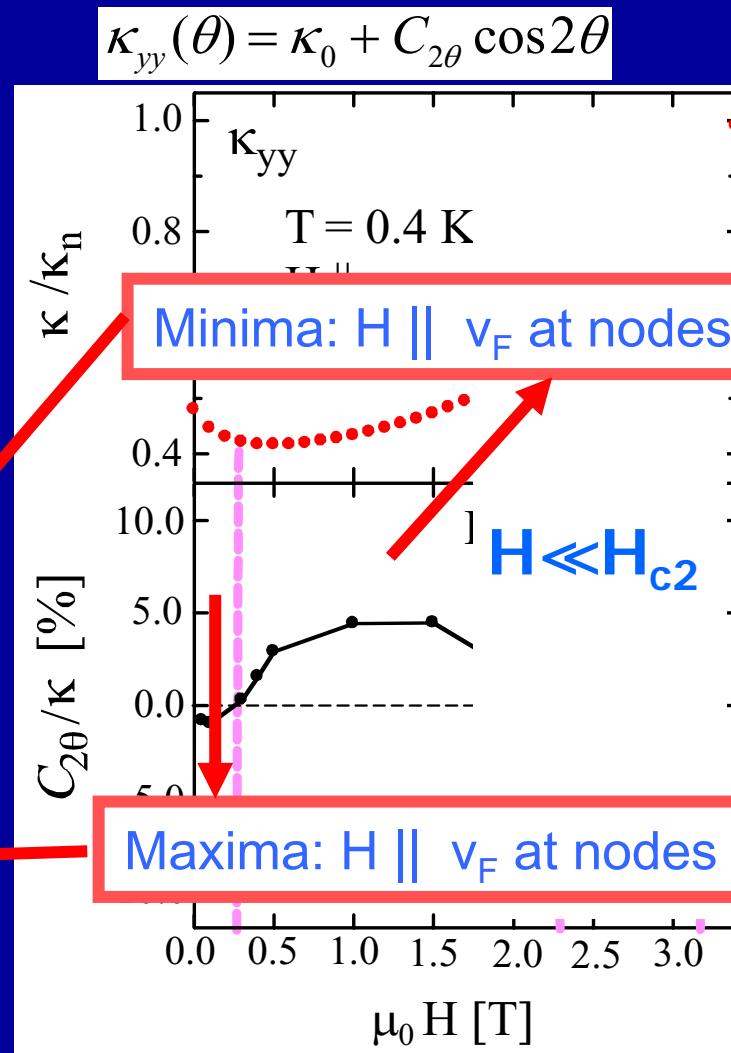
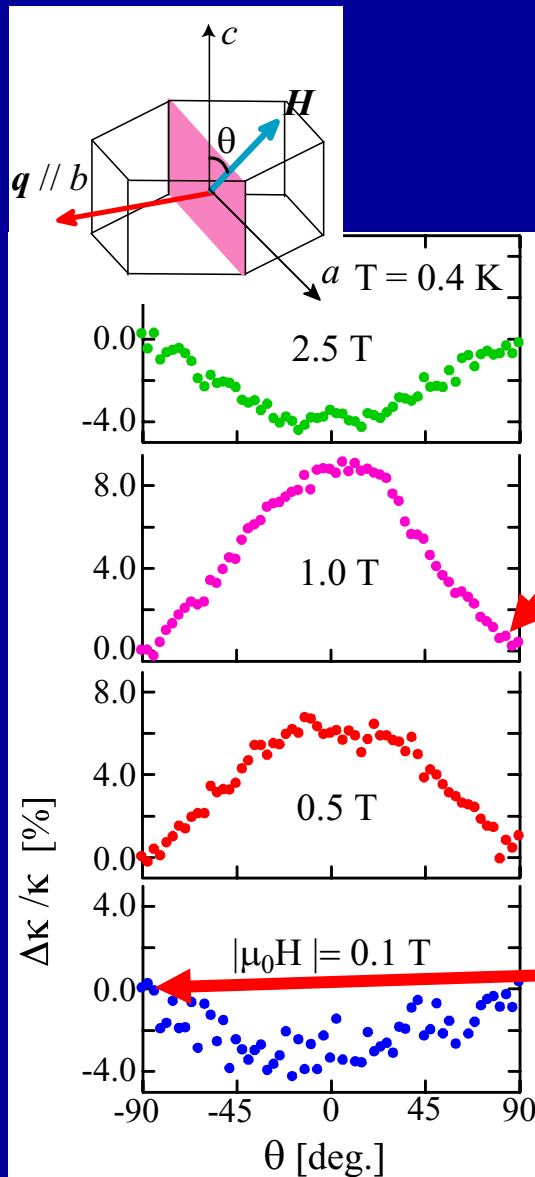
Increase due to enhancement of DOS

# Minima vs. maxima: a conjecture

- Electronic  $\kappa$  **increases** with H
  - **Density of states** dominates
  - **Minima** when H  $\parallel$  nodes
  
- Electronic  $\kappa$  **decreases** with H
  - **Scattering** dominates
  - **Maxima** when H  $\parallel$  the nodes

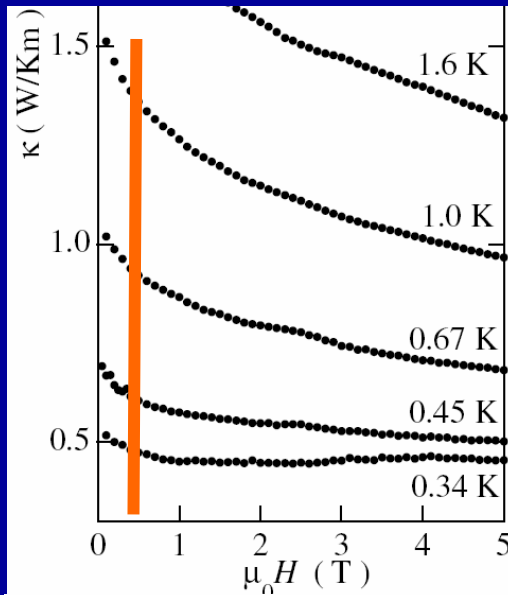


# Example: UPd<sub>2</sub>Al<sub>3</sub>



Yu.Matsuda et al. '04  
Also for YNiBC

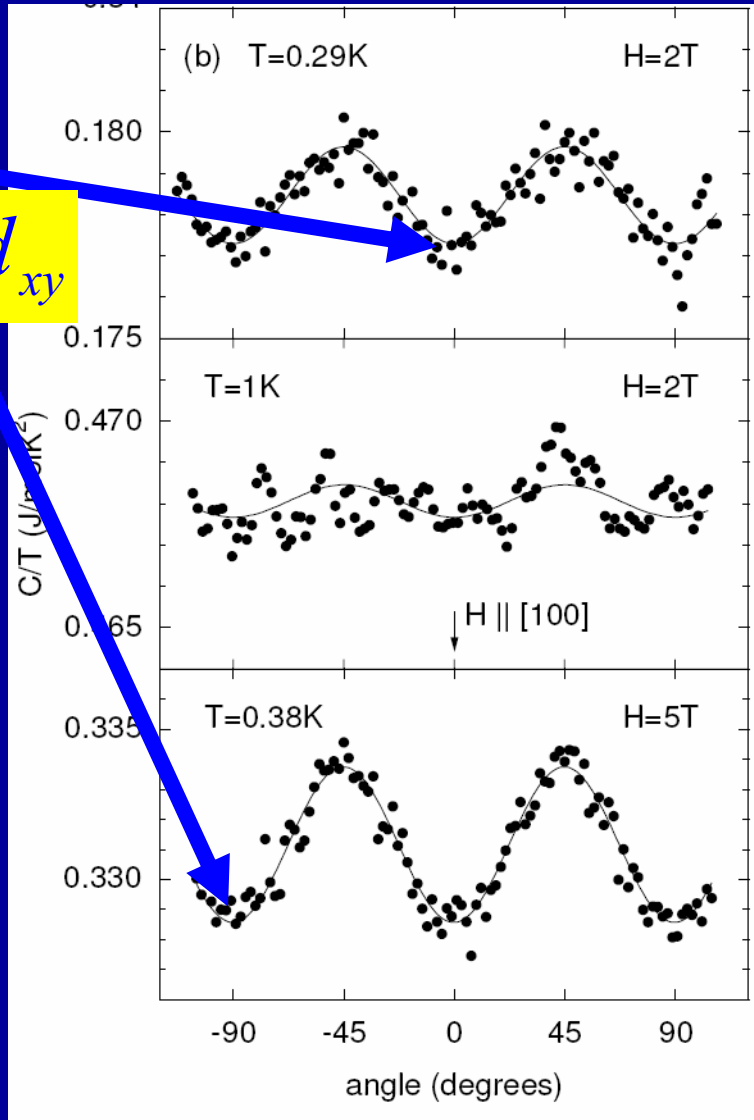
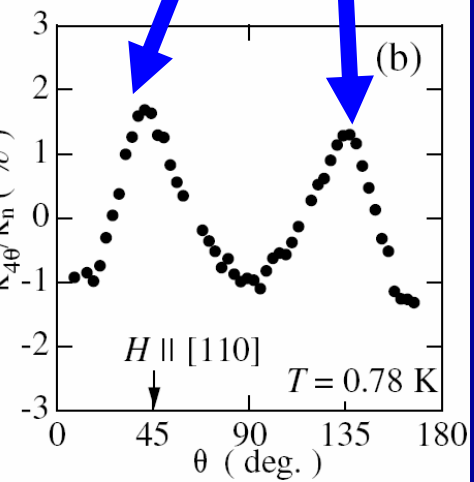
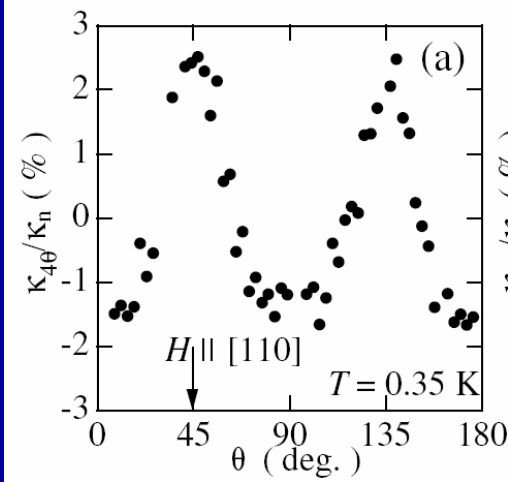
# CeCoIn<sub>5</sub>: a puzzle



nodes

$d_{xy}$

$d_{x^2-y^2}$



K. Izawa et al. '01

H. Aoki et al. '03

# Summary

- **System:** nodal superconductors
- **Foundation:**  
“**Volovik Effect**”: magnetic field probes near-nodal quasiparticles.
- **Rotation** of magnetic field with respect to nodes:
- **Provides:** map of the amplitude of the gap.
- **Specific heat:** direct probe
- **Thermal transport:** vortex scattering?