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ICTP 40th Anniversary

SMR.1572 - 7

**Workshop on
Novel States and Phase Transitions in Highly Correlated Matter
12 - 23 July 2004**

Electrons on a triangular lattice in Na-doped cobalt oxide

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These are preliminary lecture notes, intended only for distribution to participants

Electrons on a triangular lattice in Na-doped Cobalt Oxide

Yayu Wang, Maw Lin Foo, Lu Li, Nyrissa Rogado,
S. Watauchi, R. J. Cava, N.P.O.
Princeton University

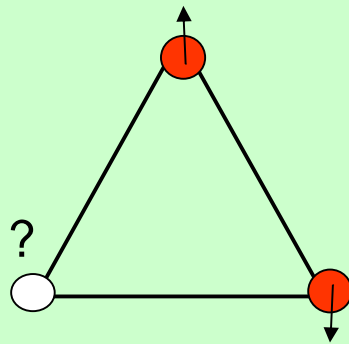
1. Frustration on triangular lattice
2. Large thermopower in Na_xCoO_2
3. Superconductivity
4. ARPES
5. Hall effect
6. Phase diagram

Supported by NSF, ONR

Geometrical Frustration on triangular lattice

$$H = -J \sum_{(i,j)} S_i S_j$$

Antiferromagnetic Ising model



Impossible to have
AF alignment
on all 3 bonds

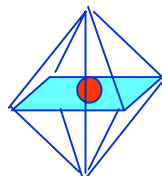
Ground state is disordered and highly degenerate

Resonating valence bond model(s) 1971, 1987

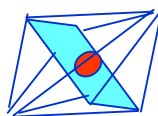
Spin Ice in pyrochlores 1998

Frustrated magnetic states in spinels 1999

building block

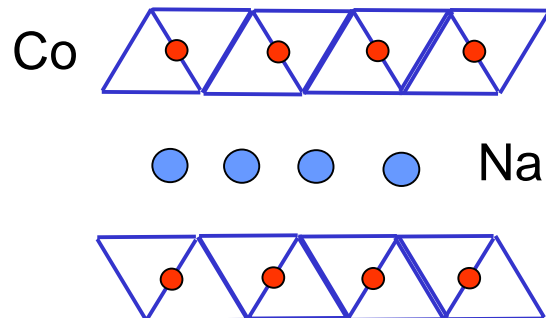


tilt →



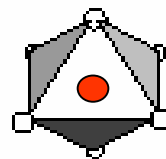
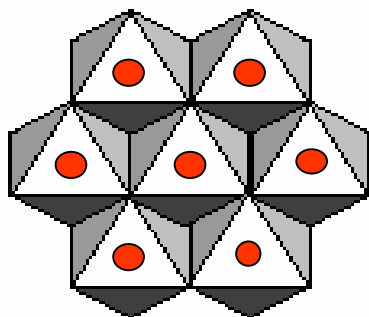
Terasaki, Uchinokura 1997

Octahedra tilted to form a layer



Na ions (dopants) sandwiched btw layers of tilted CoO_2 octahedra

Co ions define a triangular lattice

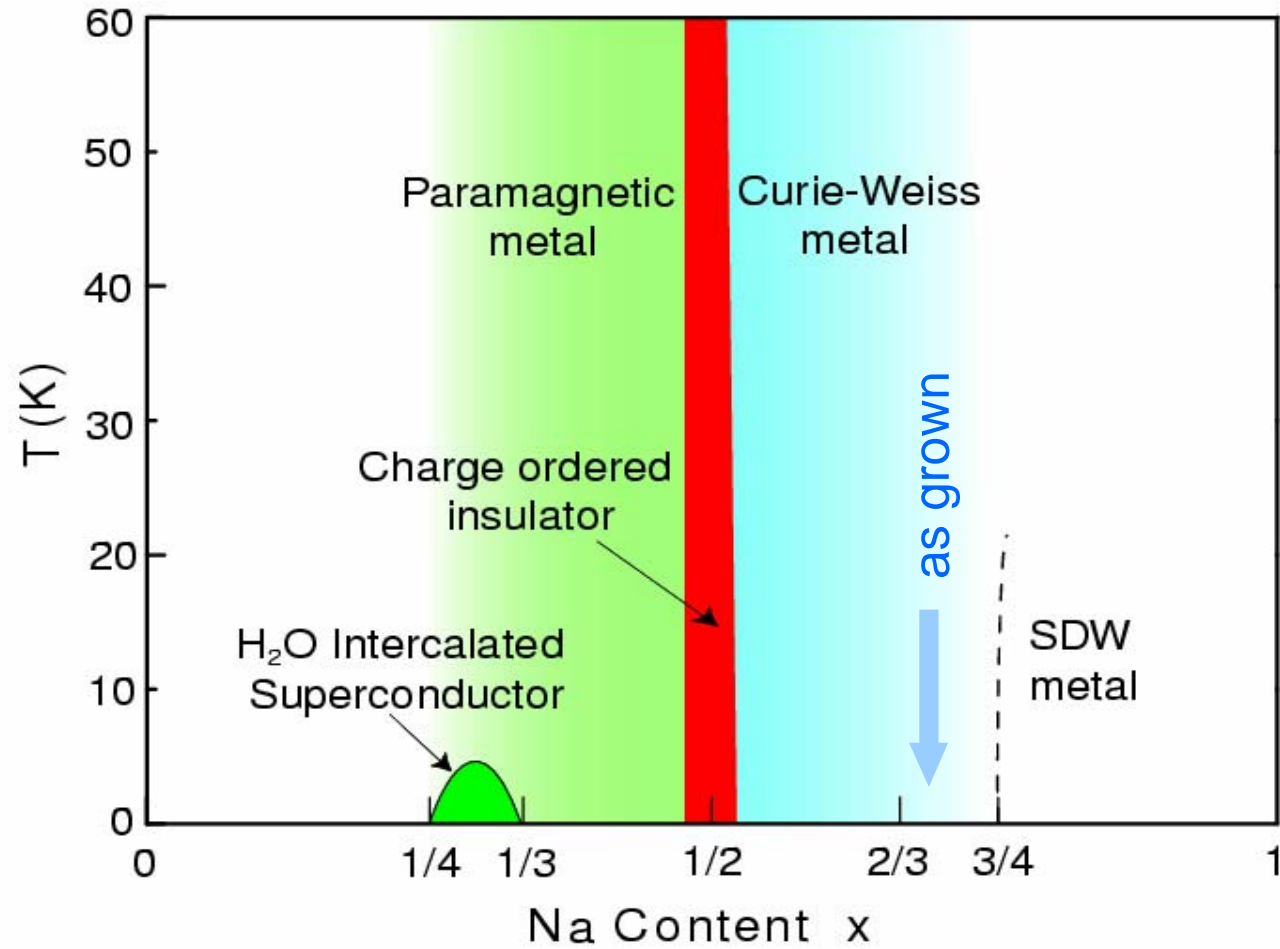
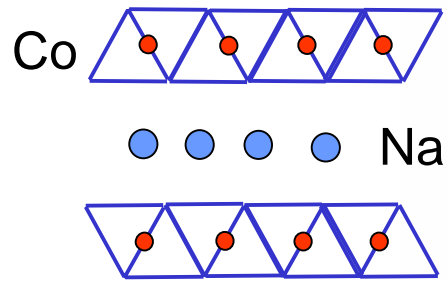


- oxygen
- cobalt



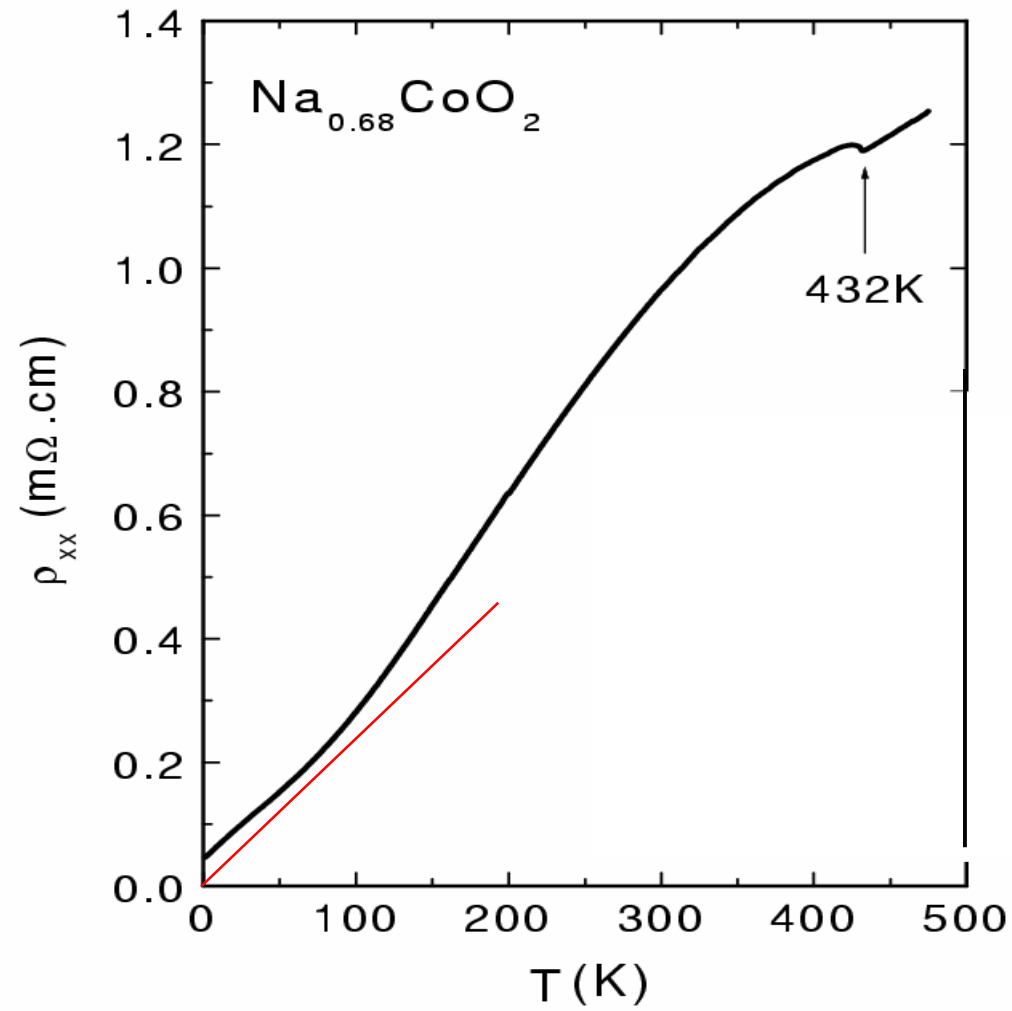
Multiple electronic phases vs. Na content

Foo et al.
PRL '04

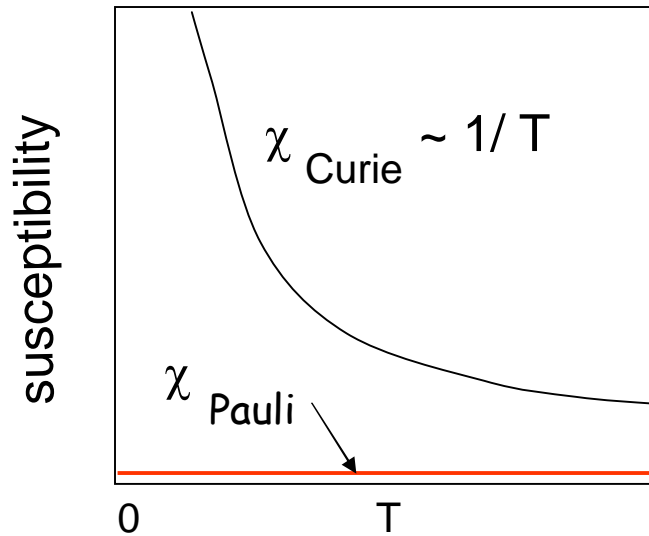


Resistivity of Na_xCoO_2 ($x \sim 0.71$)

Terasaki *et al.*, PRB 1997
Wang *et al.* Nature '03

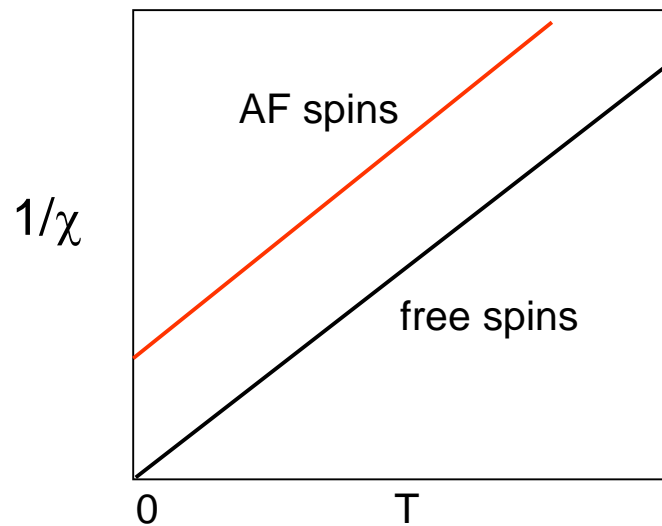
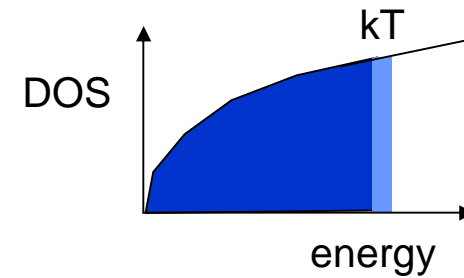


Susceptibility of insulators vs metals



Susceptibility
 $\chi = dM/dH$

In metals, χ small and indept of T

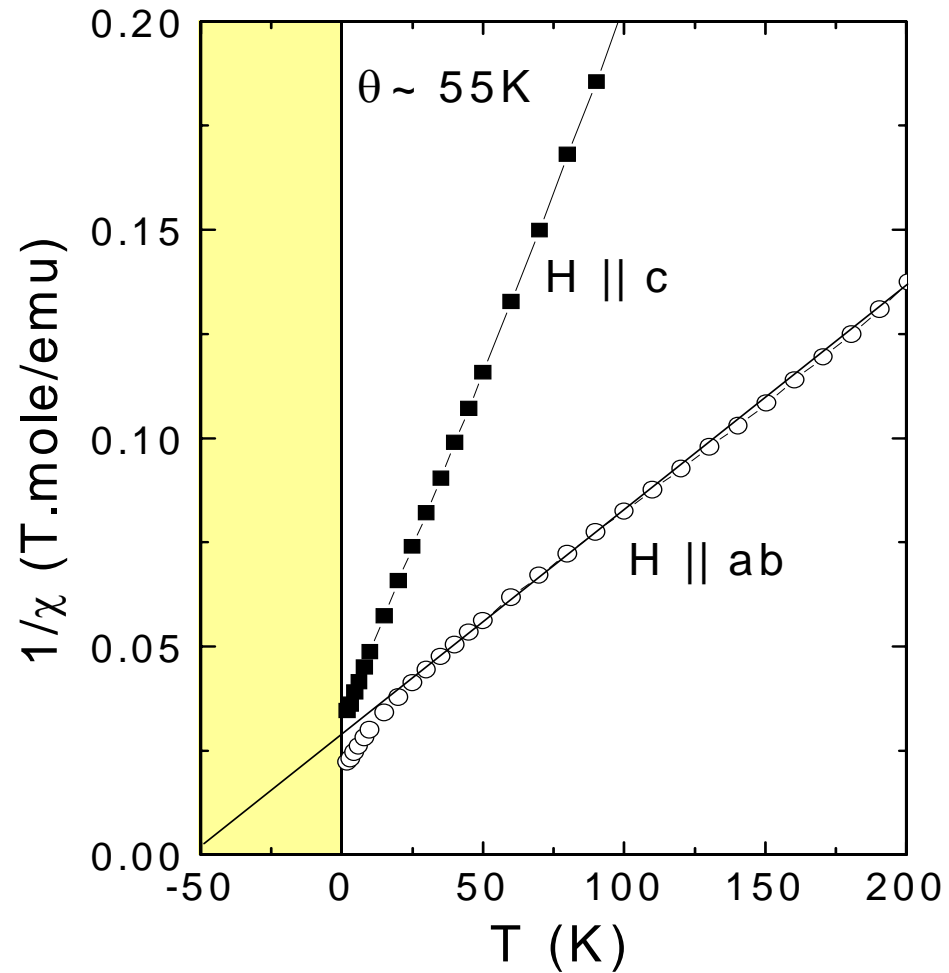


In Antiferromagnets

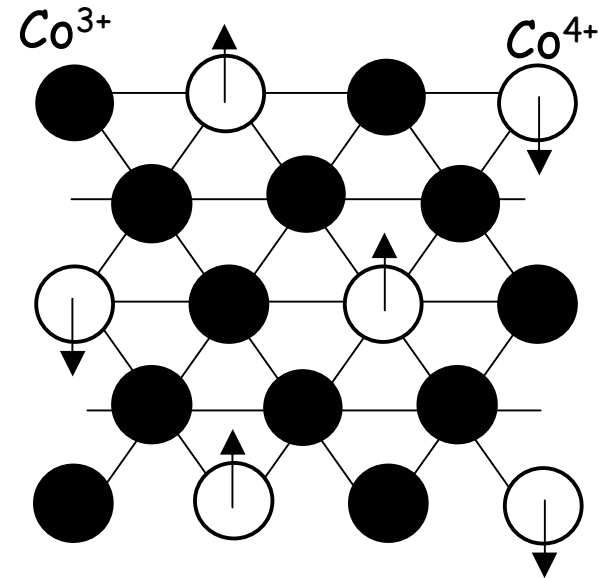
$$\chi = C/(T + \theta)$$

$$\theta = T_N (\text{Neel temp})$$

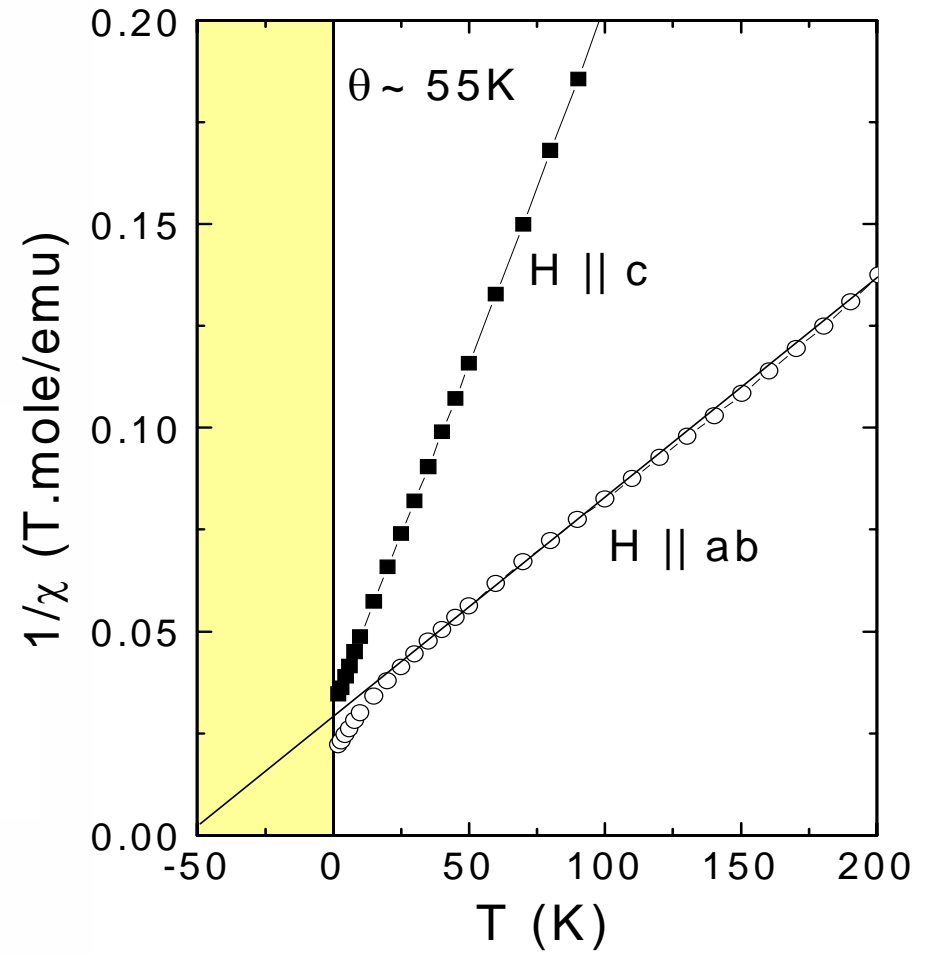
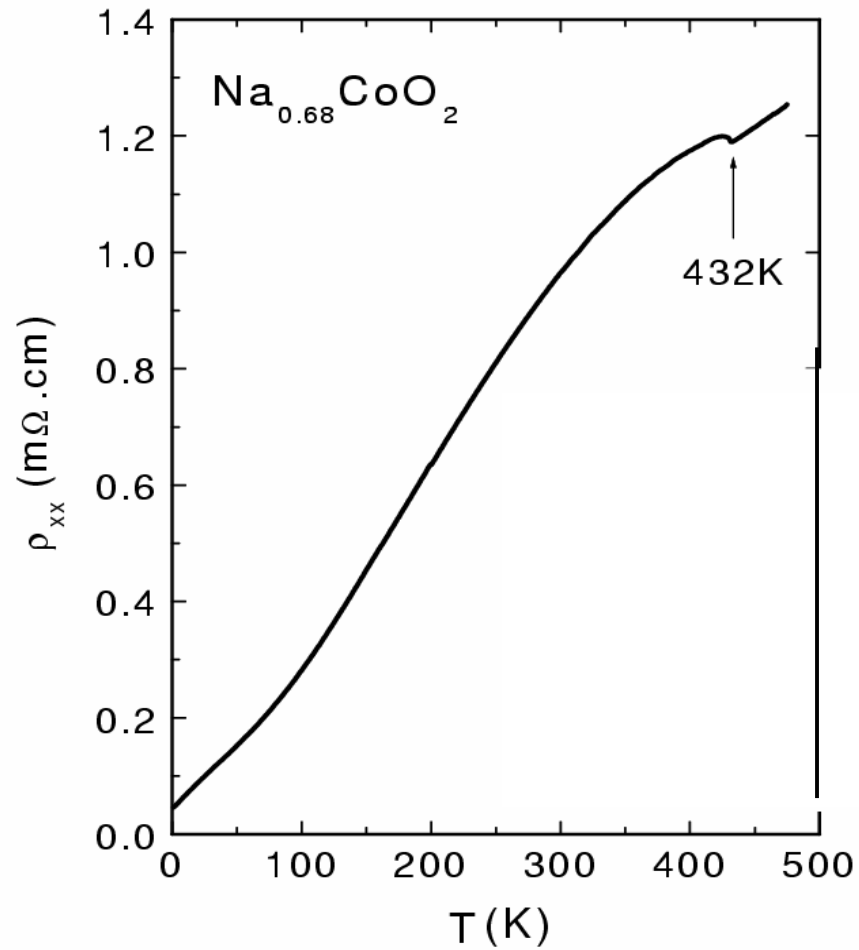
Susceptibility χ has Curie-Weiss form



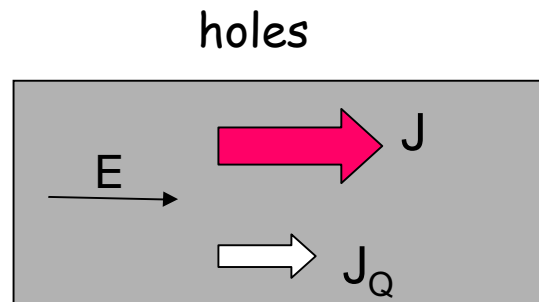
- AF Neel temperature
 $T_N \sim 60-100$ K
- Magnitude of χ implies
Co⁴⁺ ions spin $S = \frac{1}{2}$
Co³⁺ is diamagnetic ($S = 0$),



Metallic resistivity but antiferromagnetic in spin response
(Curie-Weiss Metal)



Thermopower and Peltier coef.



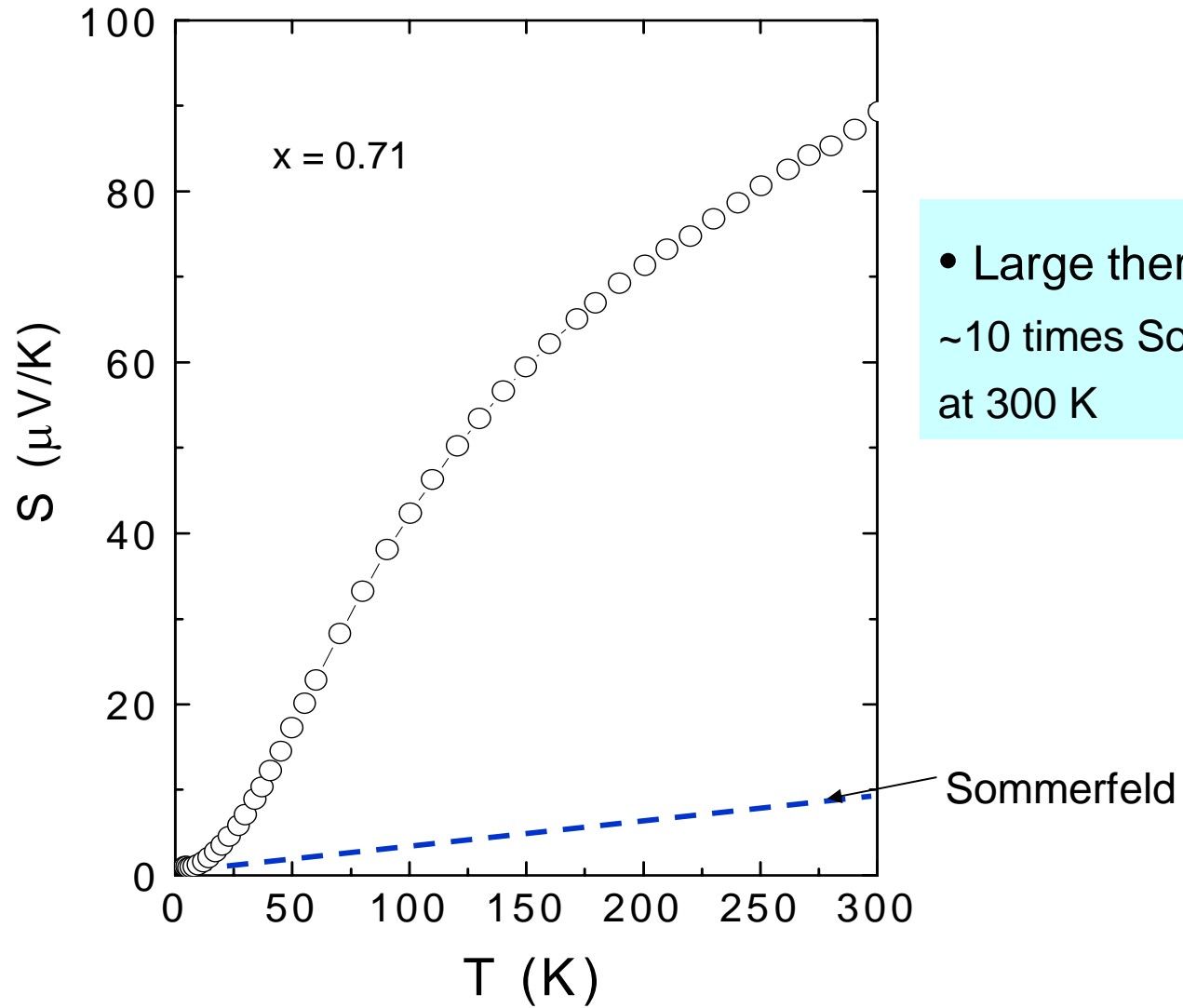
Heat current density J_Q
accompanies charge current density J

Ratio of currents $J_Q/J = \Pi$ (Peltier coef)

$$S = \Pi / T = J_Q / JT$$

Large thermopower S of Na_xCoO_2

Terasaki *et al* Phys. Rev. B (1997)



- Large thermopower
~10 times Sommerfeld value
at 300 K

Thermopower

Classical gas

$$J = nev$$

$$J_Q = n k_B T v$$

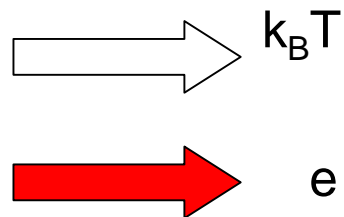
Peltier coef.

$$\Pi = J_Q/J = k_B T/e$$

Seebeck coef.

$$S = \Pi / T = k_B/e$$

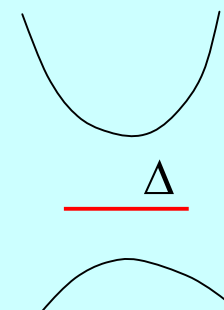
Natural unit of S
 $k_B/e = 86 \mu\text{V/K}$



Semiconductor

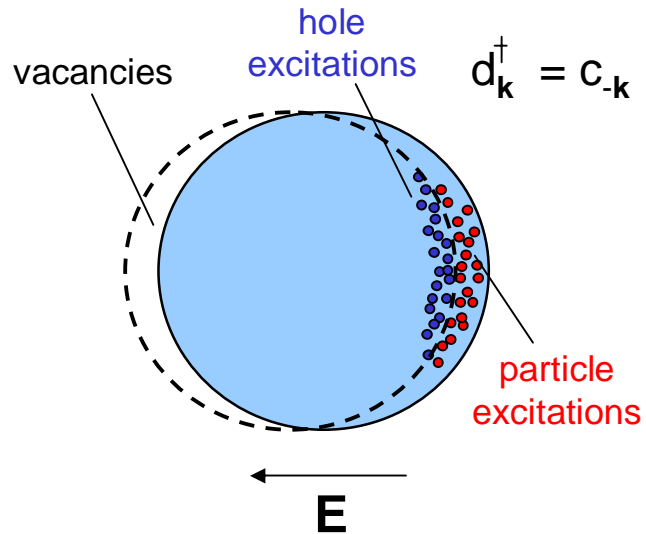
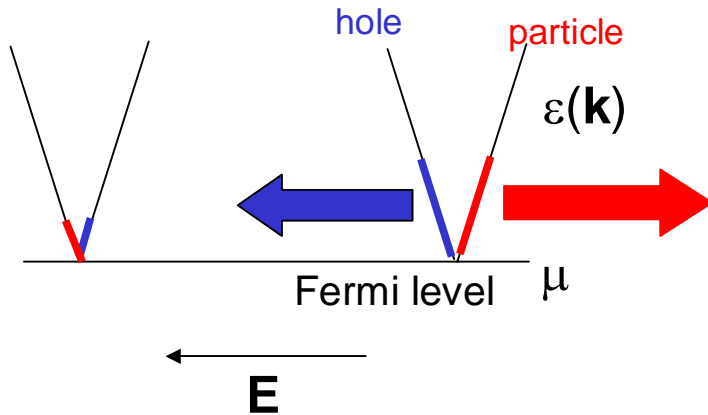
$$J_Q = n v \Delta$$

$$S = (k_B/e)(\Delta/k_B T)$$



Thermopower of conventional metals

“Excitation picture”



Fermi Gas in **E** field

Charge currents **add**
mass currents **cancel**

Heat currents **cancel**

$S = J_Q/JT$ strongly suppressed

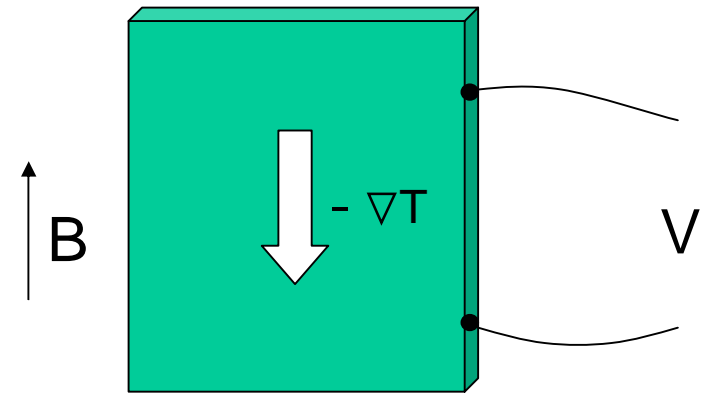
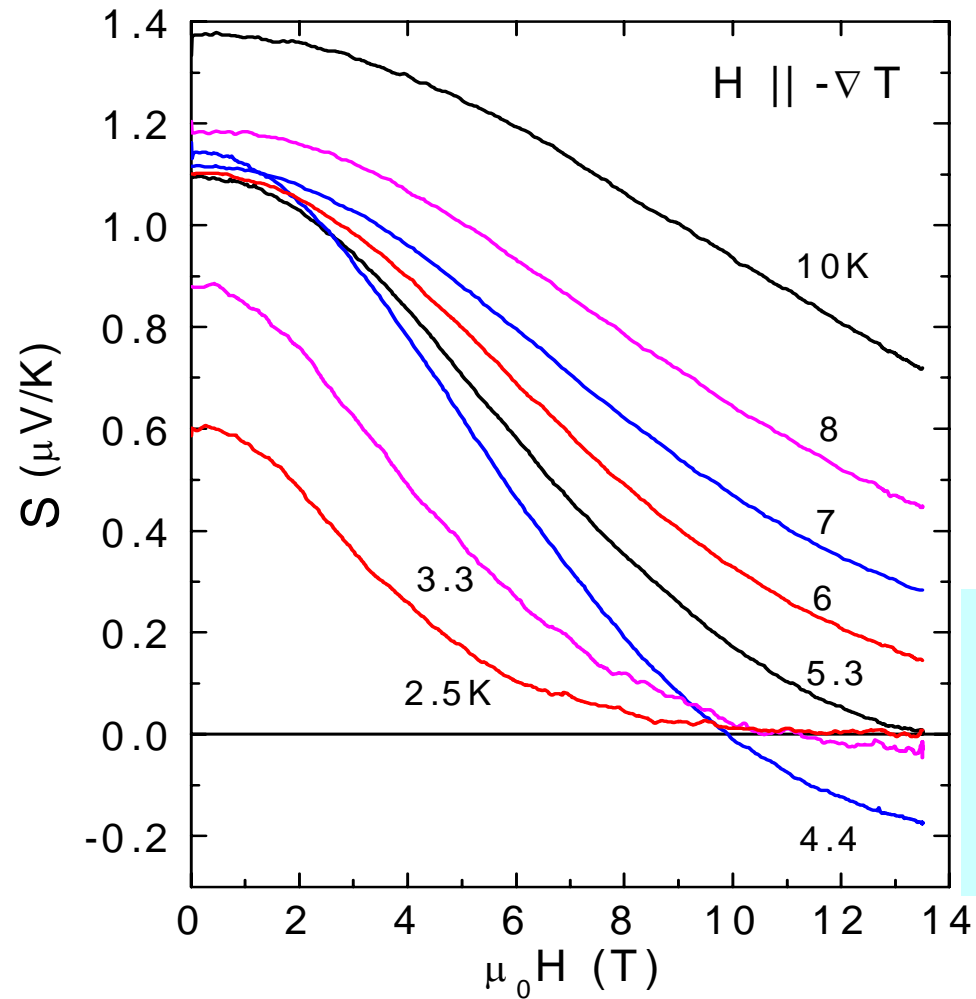
$$S \sim (k_B/e) (T/T_F) \quad T_F \sim 50,000 \text{ K}$$

$$\sim 86 \times 10^{-2} \mu\text{V/K}$$

S virtually **indepndnt** of H

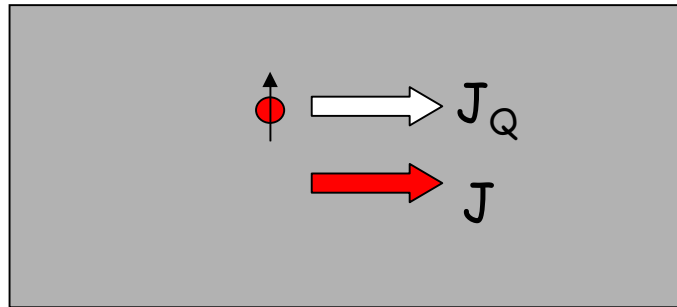
Field dependence of S in Na_xCoO_2

Wang et al. Nature '03



- In-plane field $H \parallel -\nabla T$
- Strong field suppression of Thermopower

Spin contribution to thermopower (Chaikin Beni, 1976)



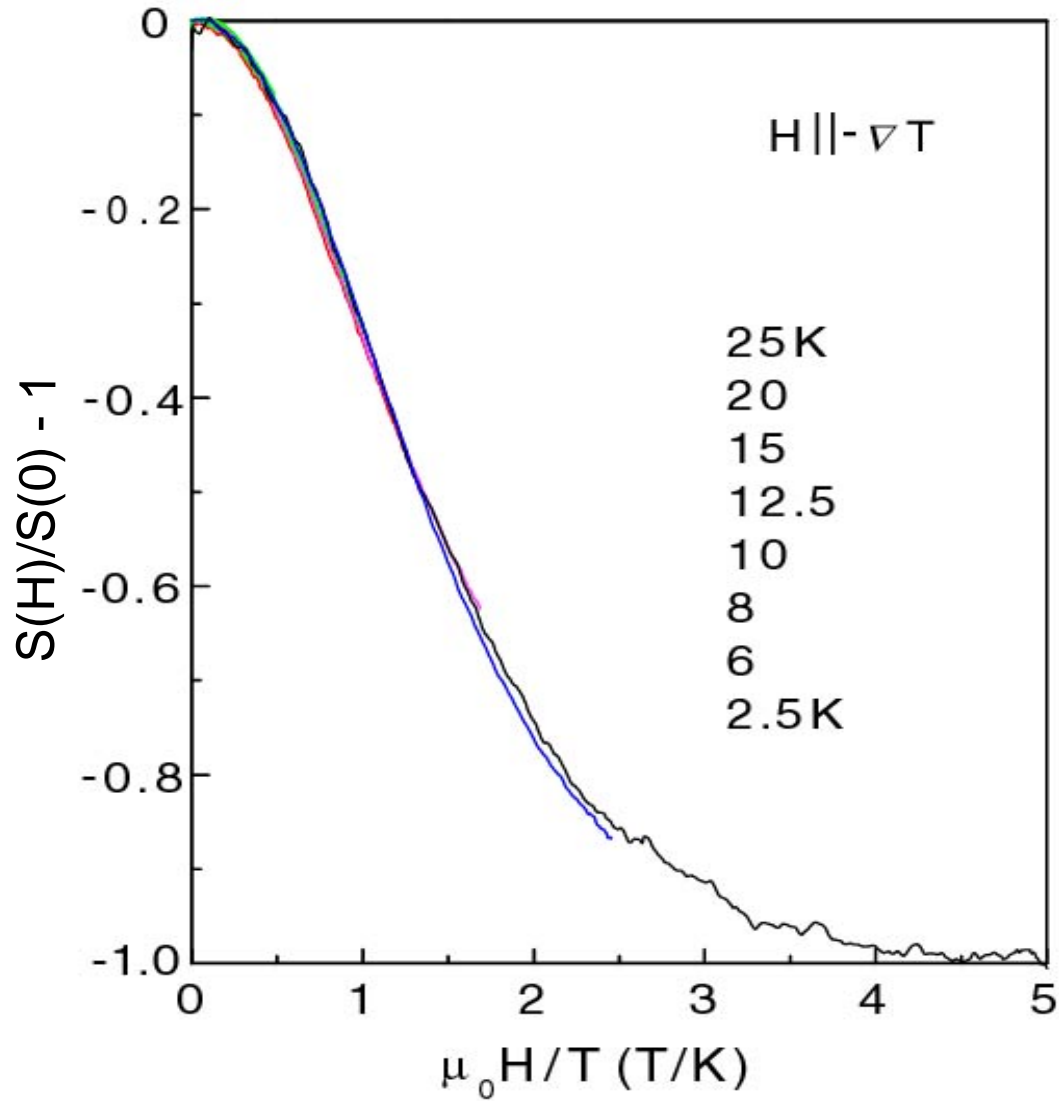
$$J = nev$$

$$\text{Spin entropy per carrier} = k_B \log 2$$

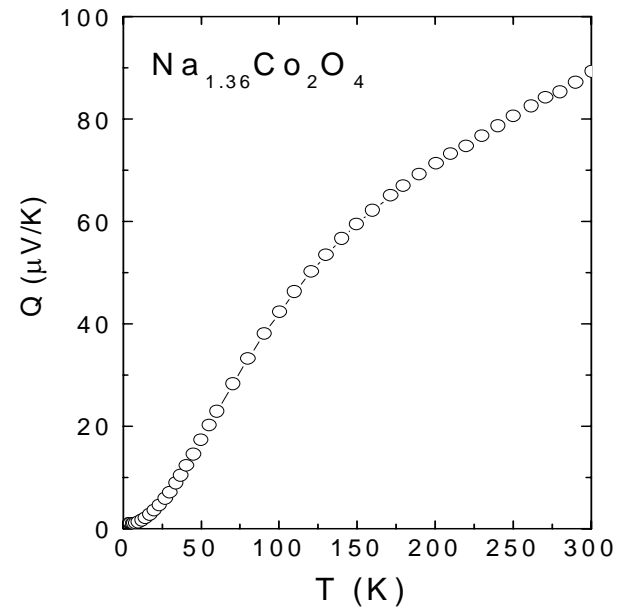
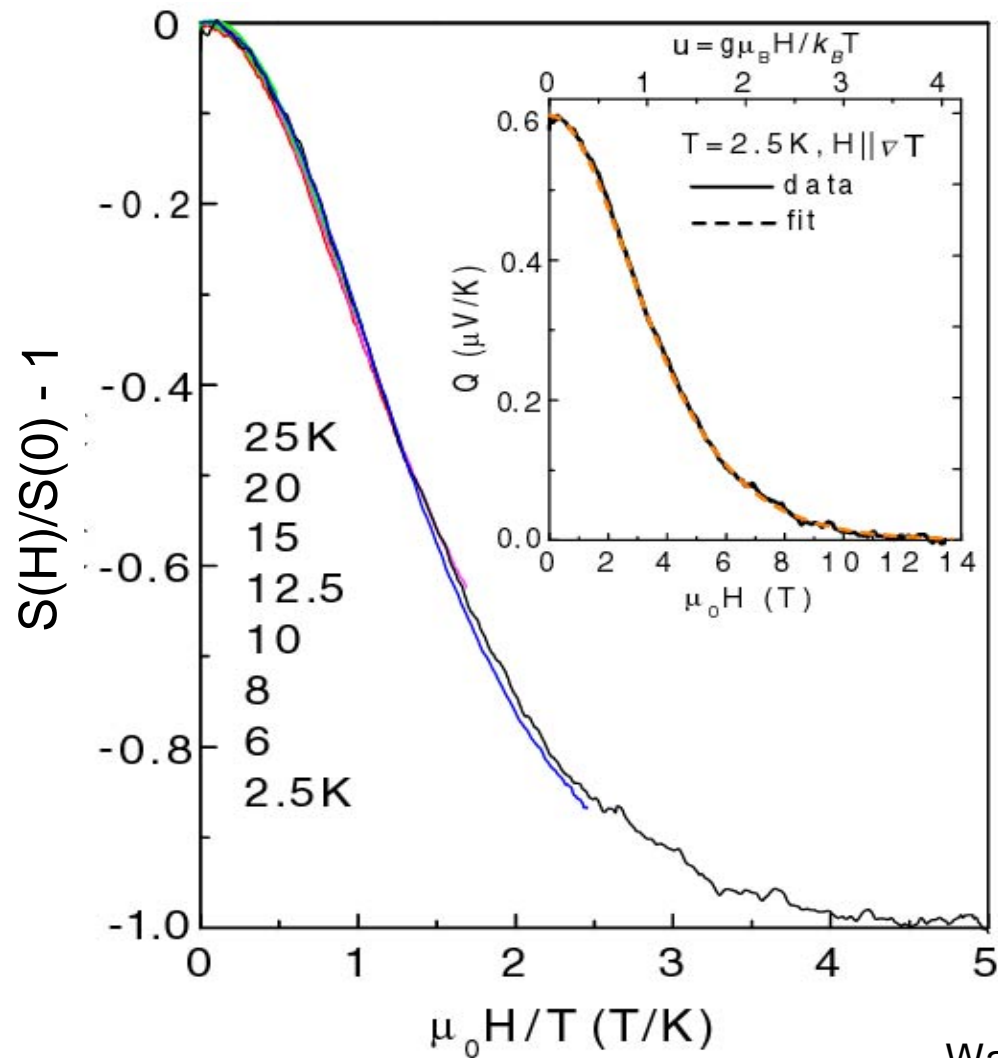
$$J_Q = nv k_B T \log 2$$

$$S = J_Q/JT = (k_B/e) \log 2 \sim 60 \mu\text{V/K}$$

Not signif. in conv. metals



$S(H,T)$ curve is a function of H/T only

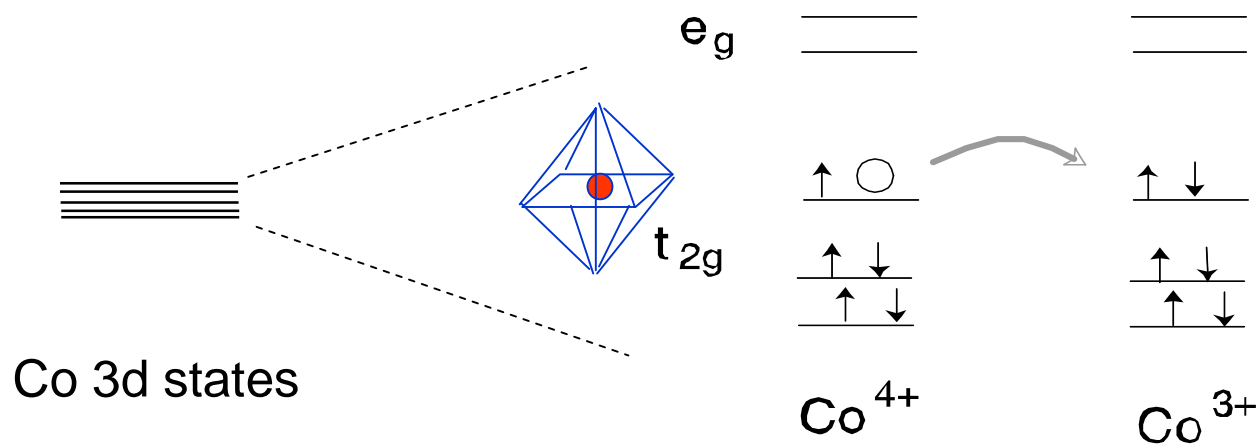
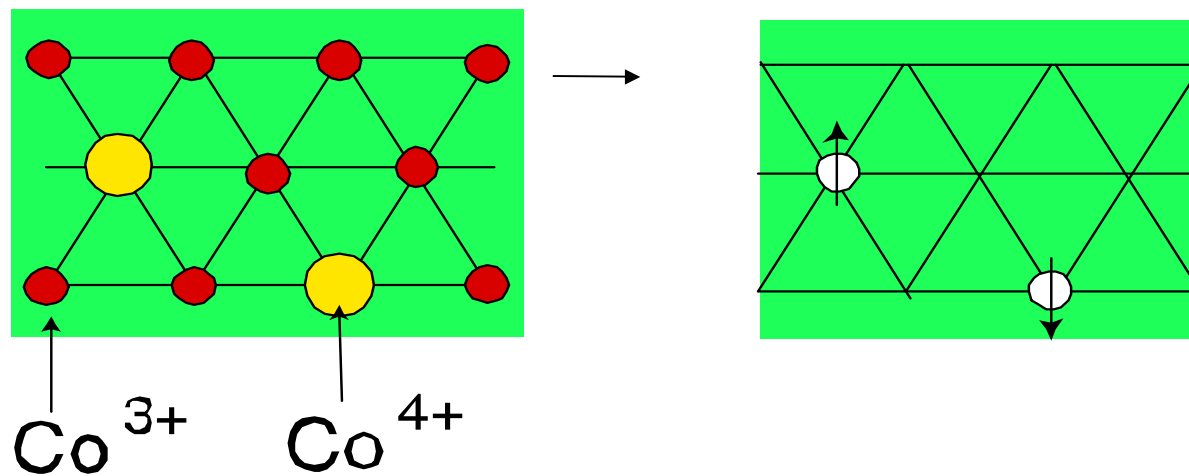


Conclusion:

1. Spin entropy is the source for enhanced thermopower
2. Key for new thermo-electric materials -- Spin

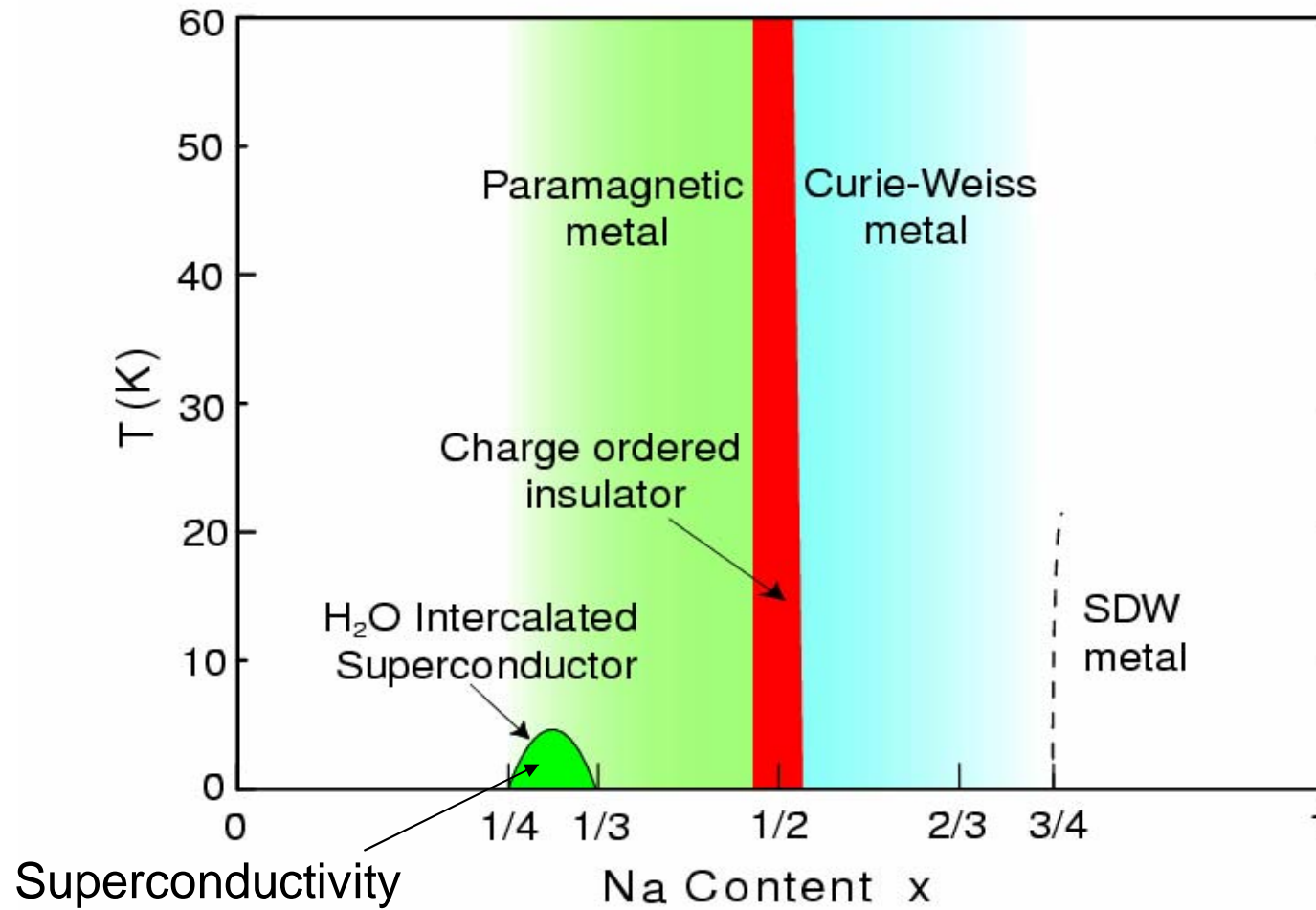
Wang et al. Nature '03

In Na_xCoO_2 , hole density $n_h = 1-x$



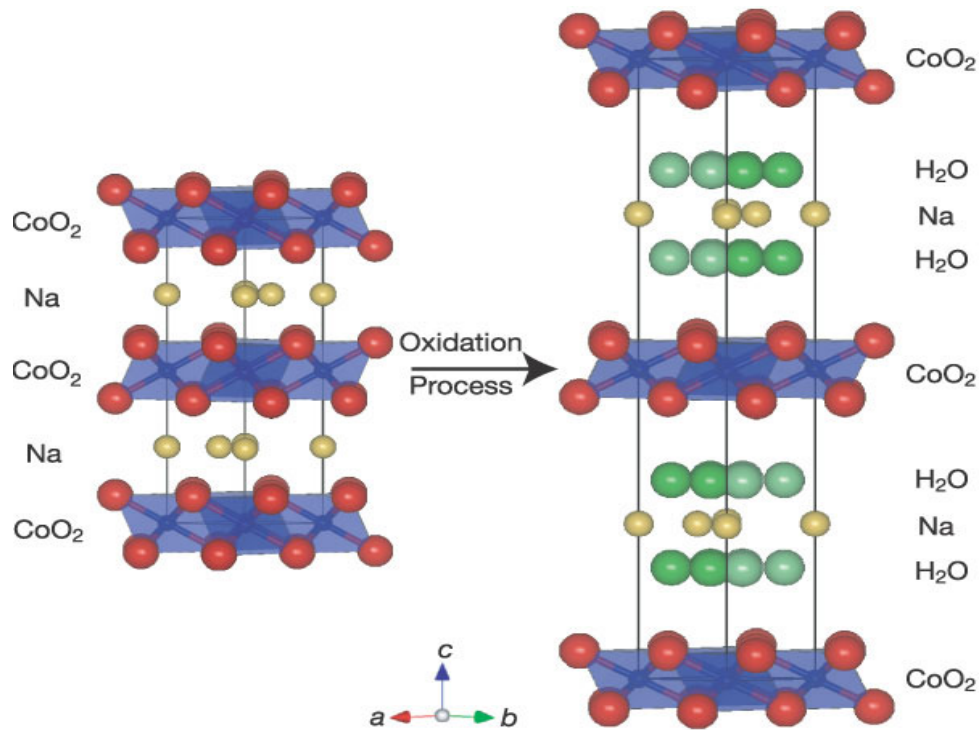
Na_xCoO_2
Multiple electronic phases vs. Na content

Foo et al.
PRL '04



Water intercalated superconductor

Takada et al., Nature (2003).



- pairing symmetry:
 s , p or d -wave?
- Why is water essential?
- What is pairing mechanism:
 e - ph or e - e or magnetic?

$\text{Na}_x\text{CoO}_2 \cdot y \text{H}_2\text{O}$, $x \sim 0.35$, $y \sim 1.30$

Superconductor with $T_c \sim 4.5\text{K}$

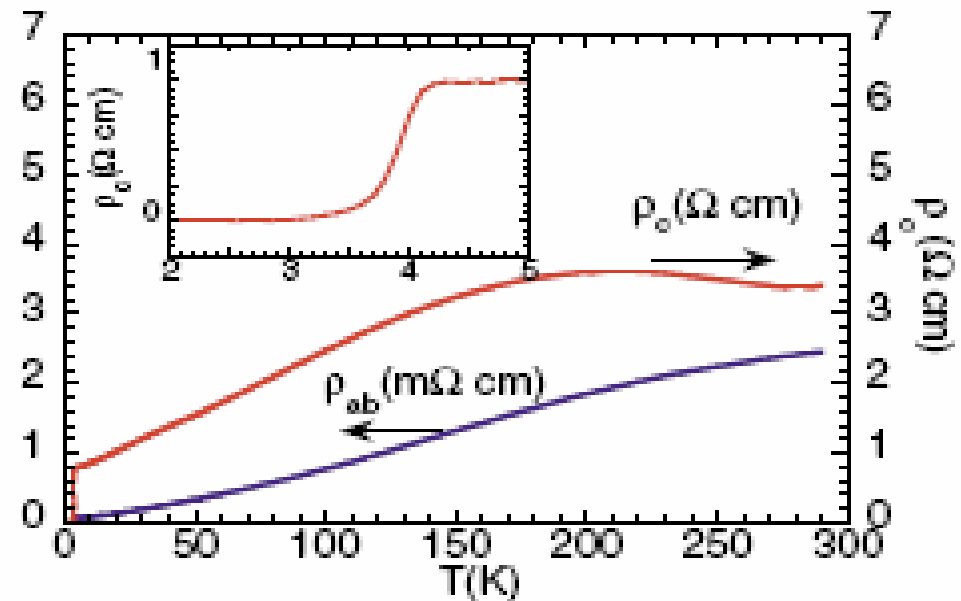
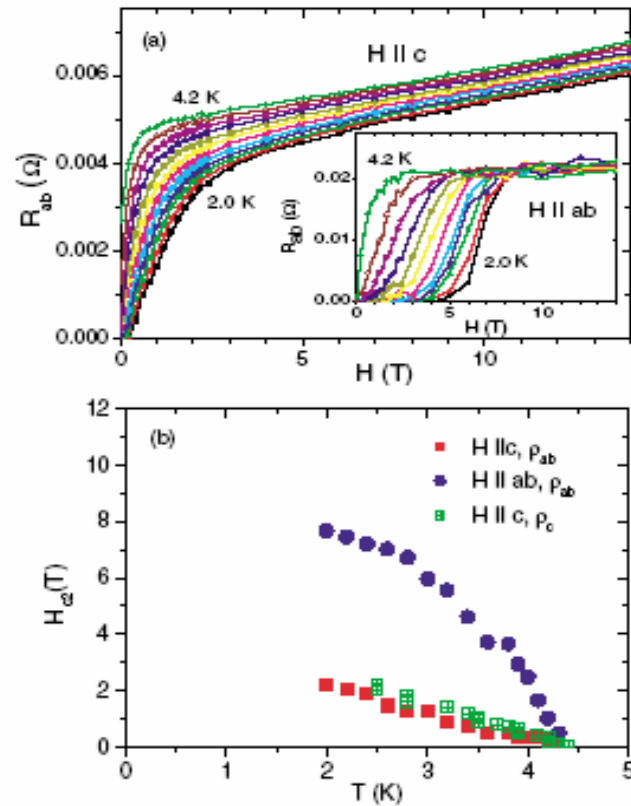
Thermodynamic and Transport Measurements of Superconducting $\text{Na}_{0.3}\text{CoO}_2 \cdot 1.3\text{H}_2\text{O}$ Single Crystals Prepared by Electrochemical Deintercalation

F. C. Chou,¹ J. H. Cho,^{2,*} P. A. Lee,^{1,2} E. T. Abel,^{1,2} K. Matan,^{1,2} and Y. S. Lee^{1,2}

¹Center for Materials Science and Engineering, Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA

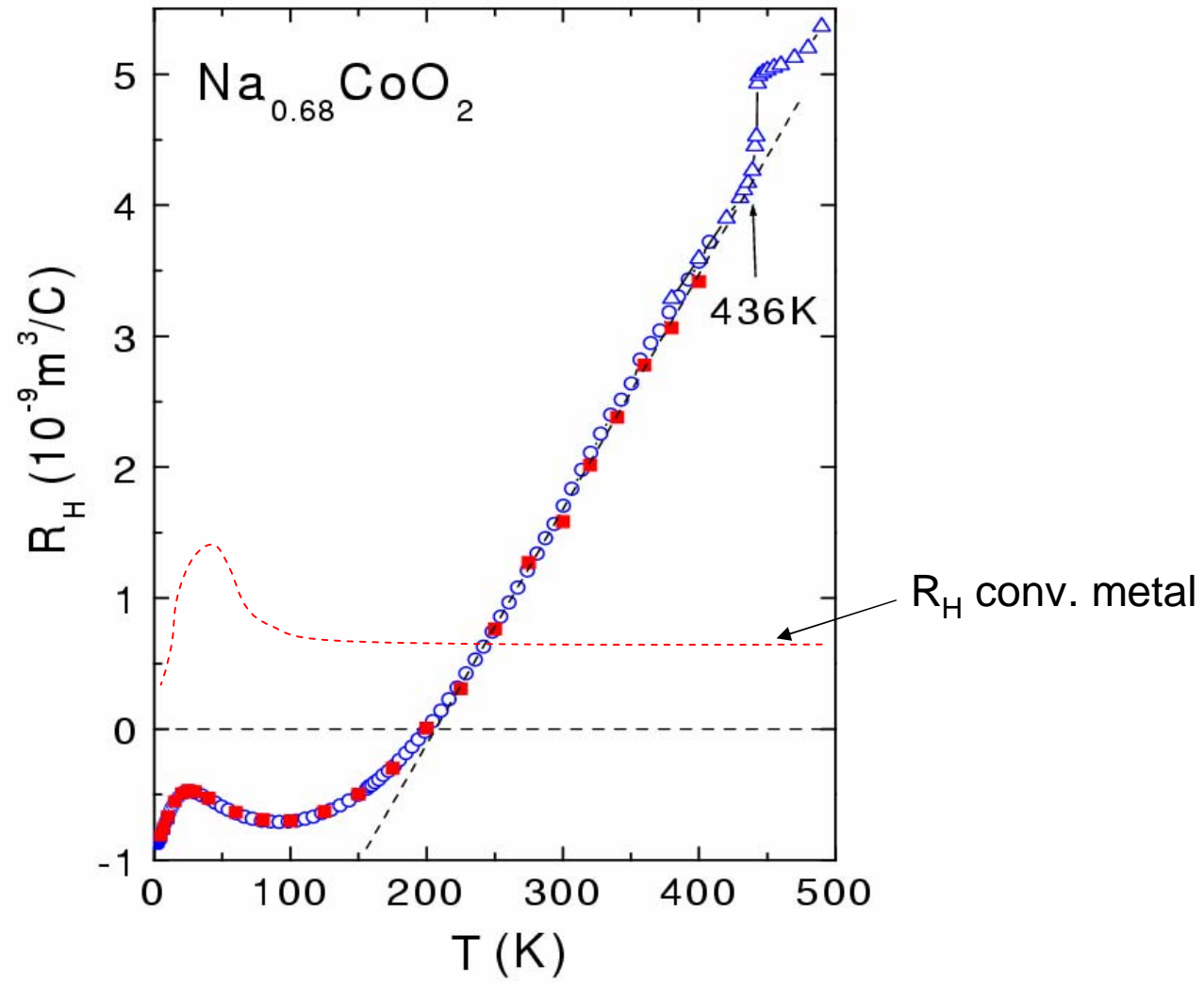
²Department of Physics, Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA

(Received 27 June 2003; published 15 April 2004)



T-linear Hall coefficient

Yayu Wang, 03

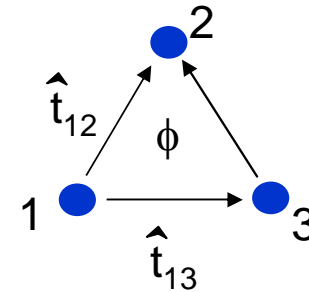


Why is R_H T -linear?

Hopping Hall current in triangular lattice (Holstein, '61)

$$\sigma_H \sim \hat{t}_{12} \hat{t}_{23} \hat{t}_{31} \sim i t^3 \exp(i\alpha)$$

Peierls phase $\alpha = 2\pi \phi/\phi_0$



High-frequency R_H^* in tJ model (B.S. Shastry '93, '03)

$$\sigma_H \sim i(\beta t)^3 \exp(i\alpha) \quad \beta t \ll 1 \quad (\beta = 1/T)$$

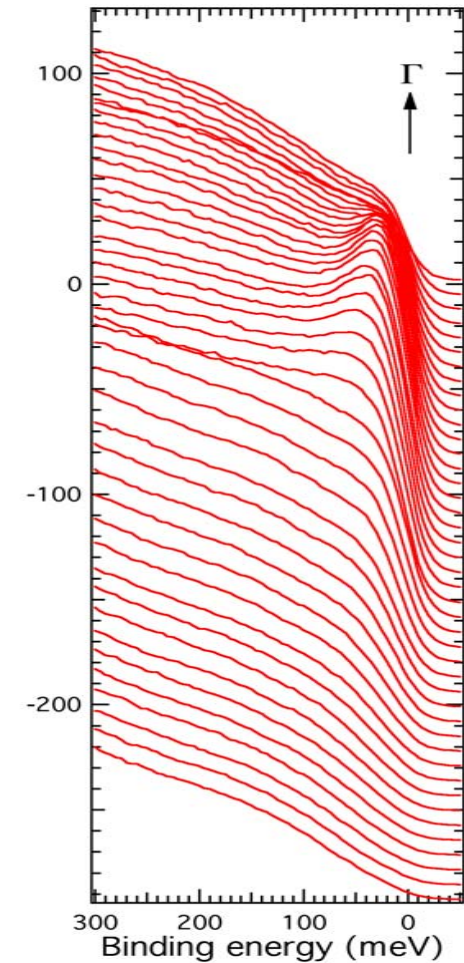
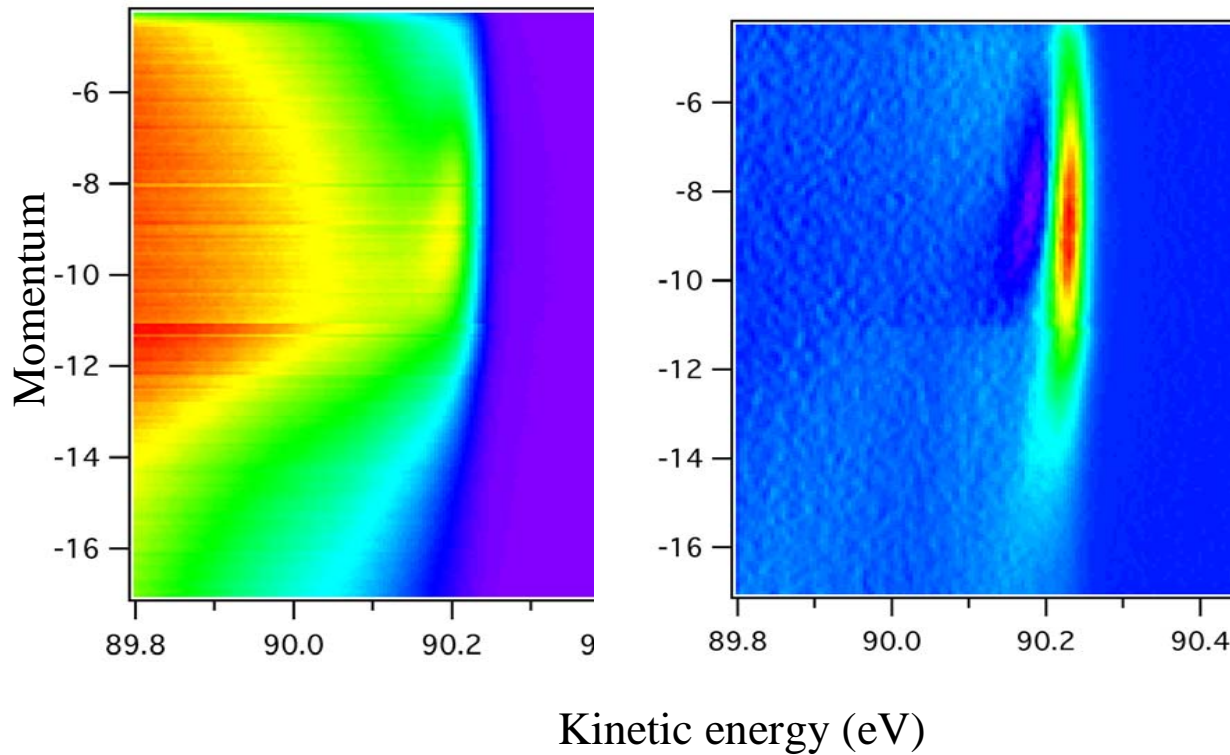
$$\sigma \sim (\beta t)^2$$

$$R_H^* \sim \sigma_H / H \sigma^2 \sim (\beta t)^{-1} \quad T\text{-linear}$$

Single-particle hopping :

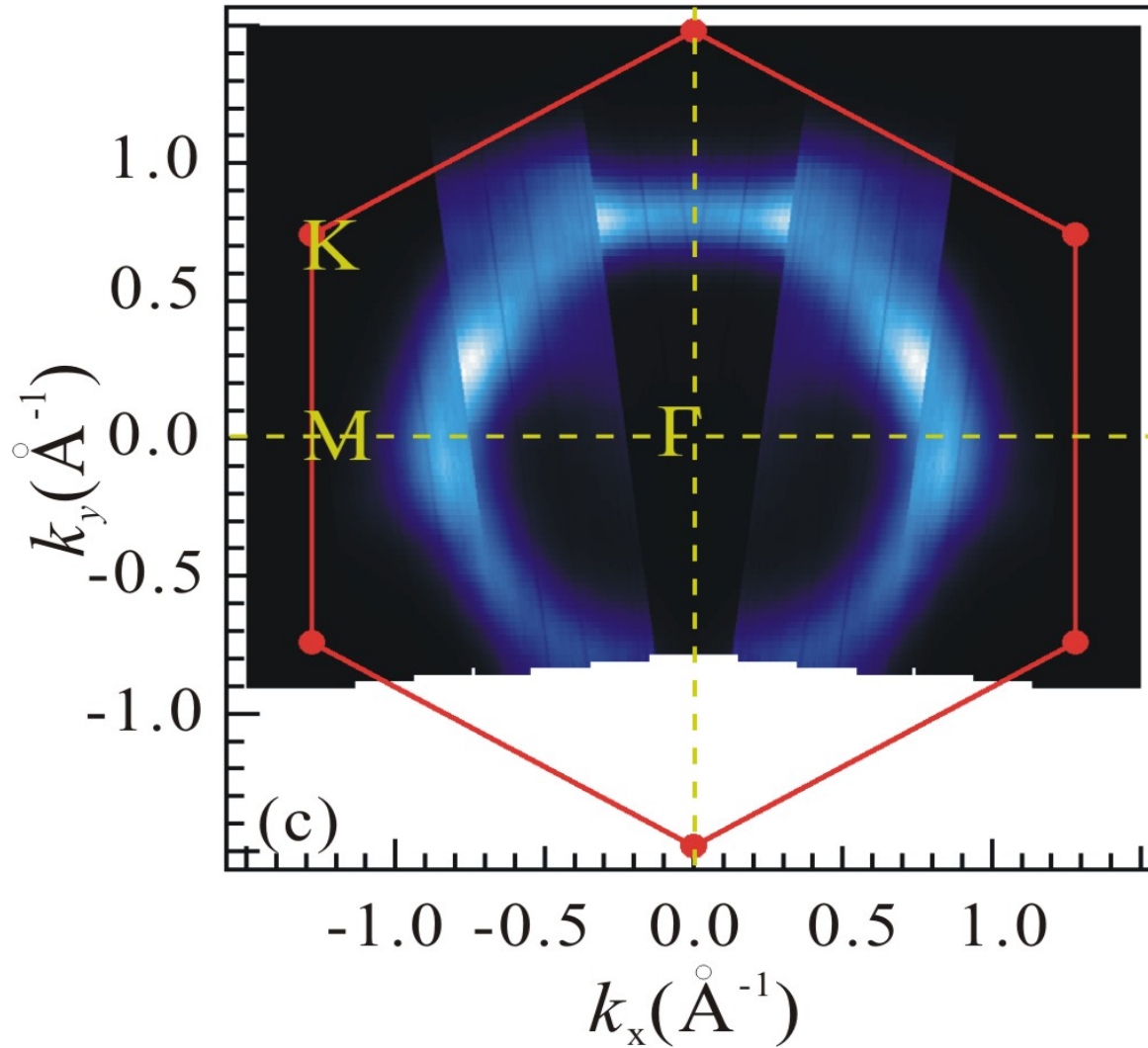
Small bandwidth \rightarrow Low degeneracy T

$t < 0$ and $|t| \sim 10$ meV (bandwidth < 100 meV)



Fermi Surface of $\text{Na}_{0.71}\text{CoO}_2$ measured by ARPES

Hasan et al.

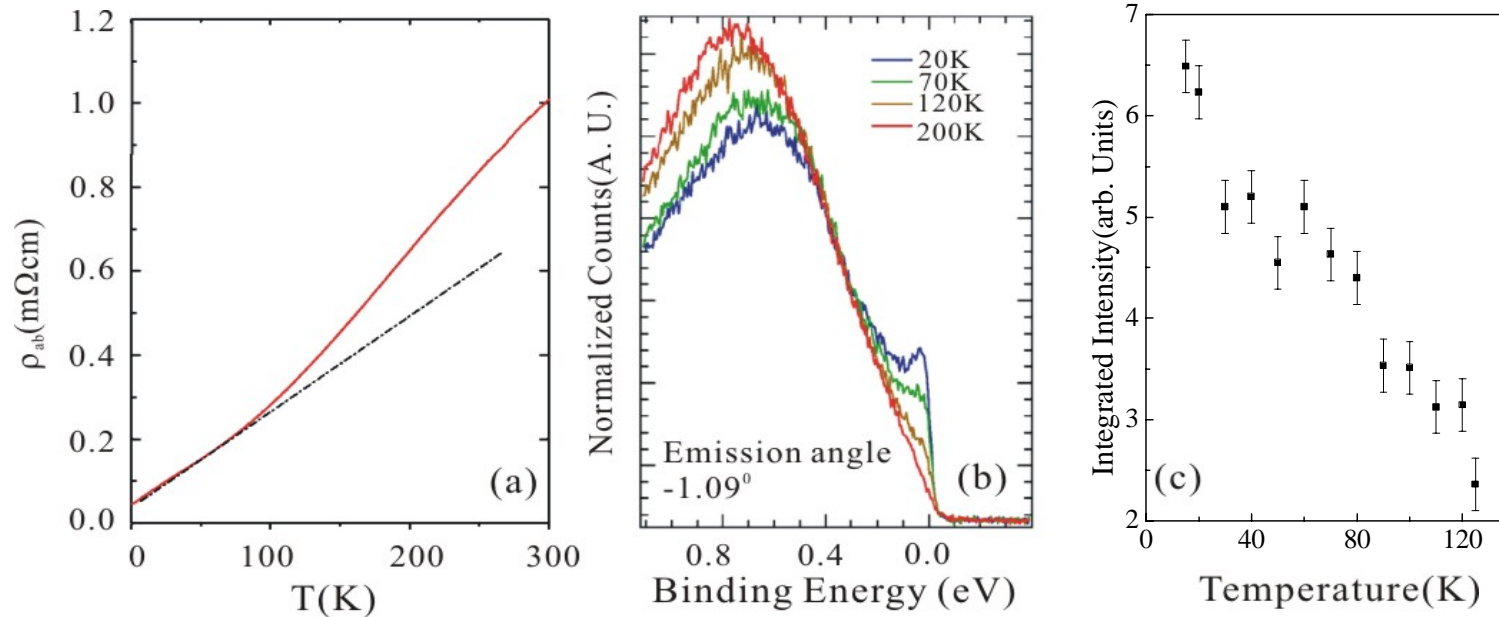


Large hole-like FS

Hopping integral
 $t \sim 10$ meV

Fermi velocity < 0.4 eV. \AA

Behavior of quasi-particles versus temperature



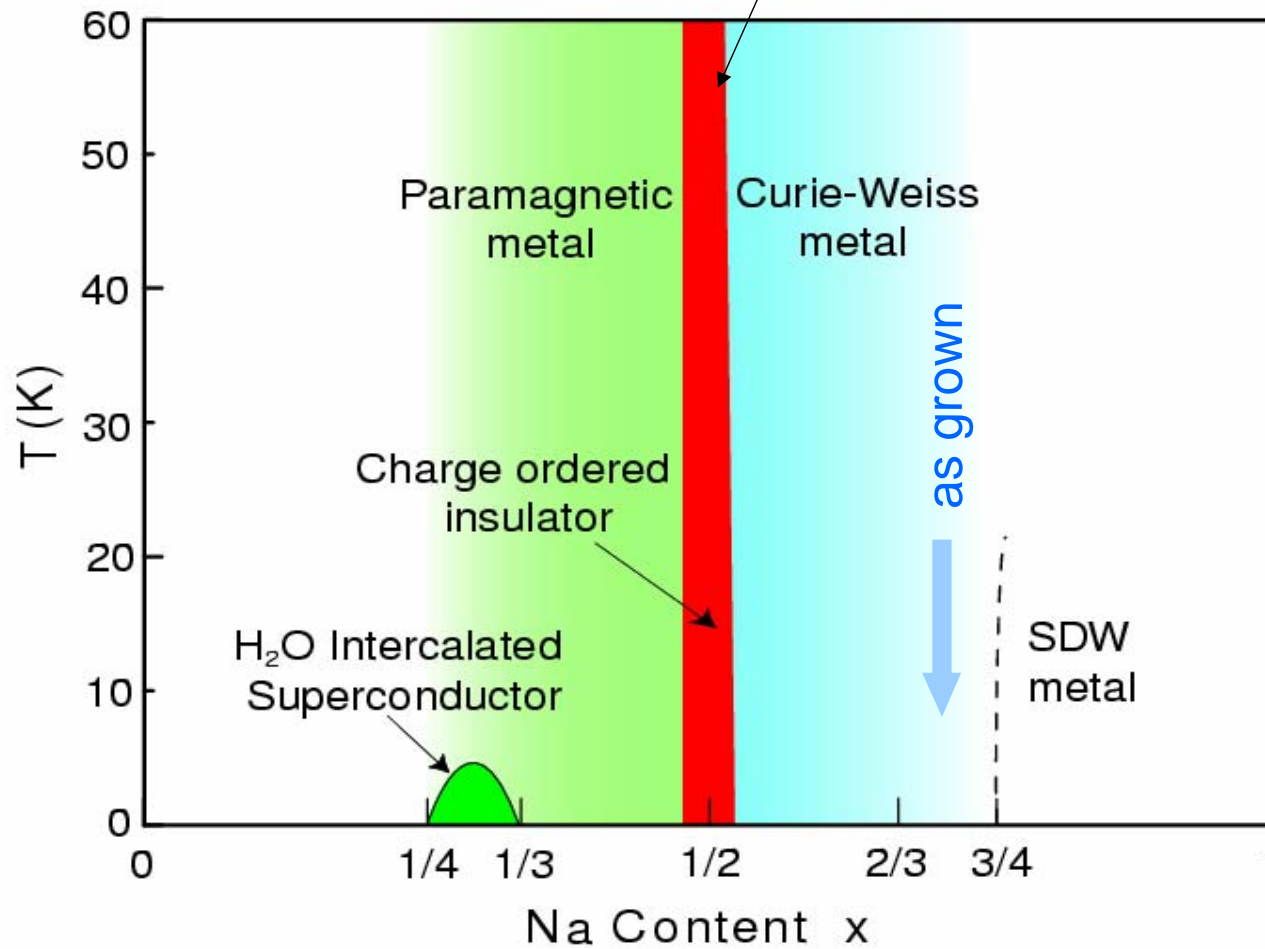
Resistivity is T-linear
below 100K

ARPES Quasiparticles are coherent
only below 150K

Na_xCoO_2
Multiple electronic phases vs. Na content

Insulating state

Foo et al.
PRL '04



Fine-tuning of Na content in Na_xCoO_2 single crystals

Foo et al., condmat/0312174 (2003), PRL '04

- Reduce the Na content by a series of chemical de-intercalation
- $x = 0.75$, as grown crystals of Floating zone or flux method

$x = 0.68$: NaClO_3 in water

$x = 0.50$: I_2 in Acetonitrile

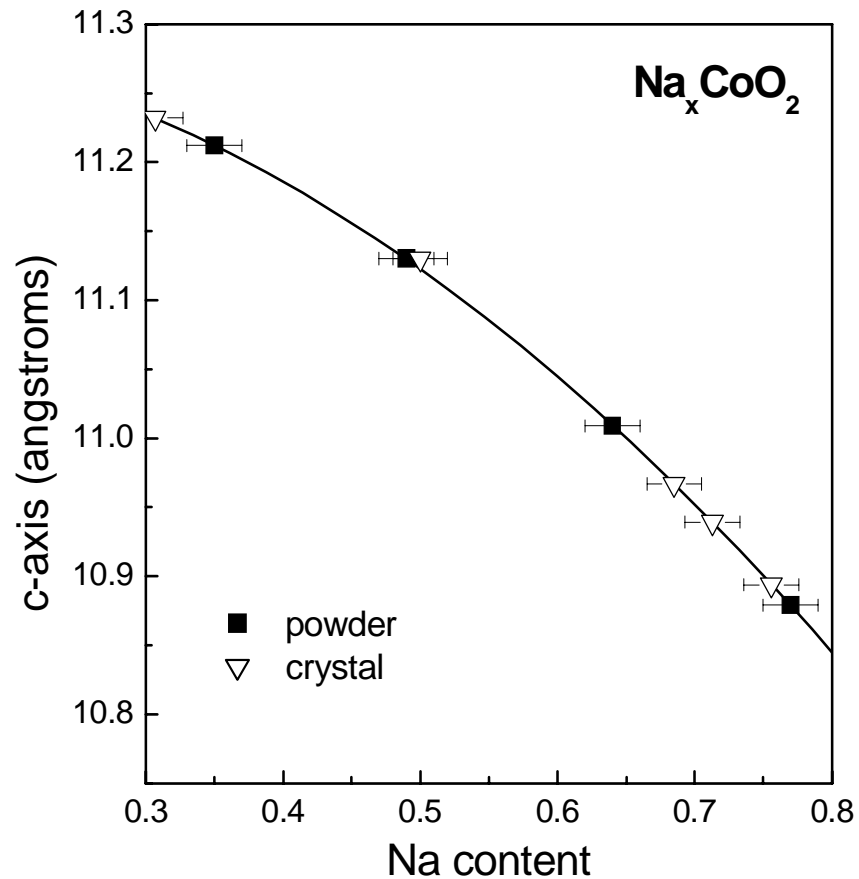
$x = 0.31$: Br_2 in Acetonitrile



Stronger oxidation agent

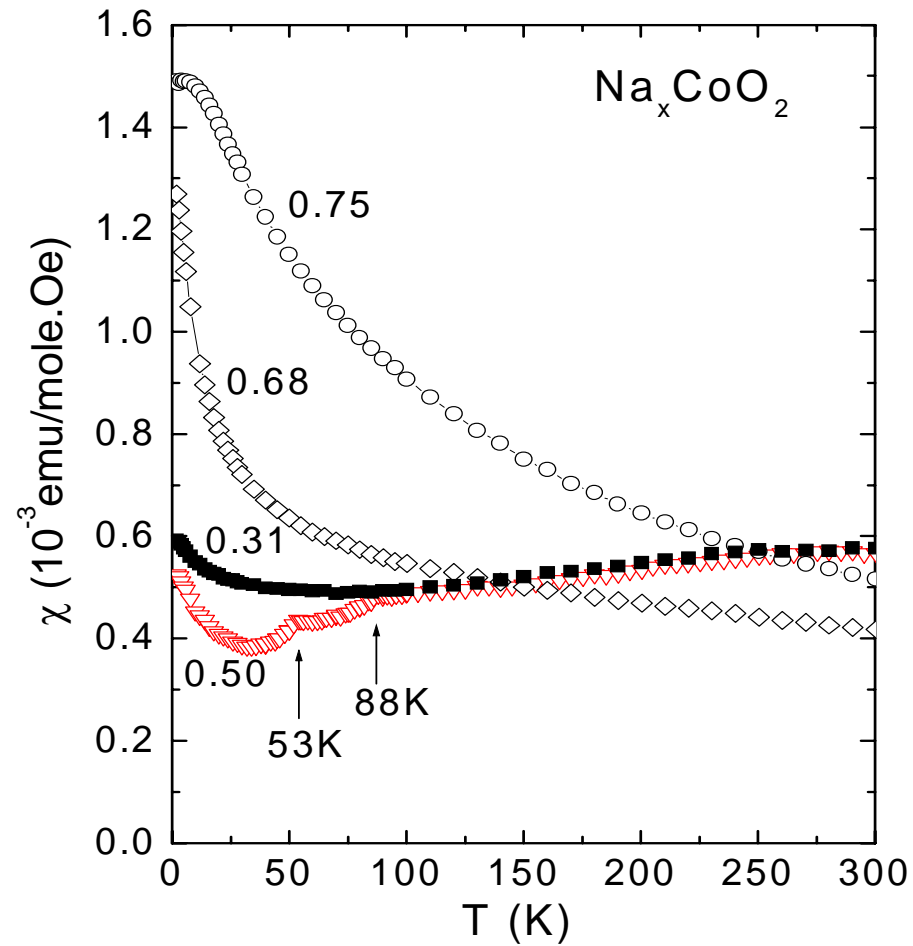
High-quality crystals with Na content $0.31 < x < 0.75$

Calibration of the Na content vs. c-axis lattice parameter



Calibration procedure

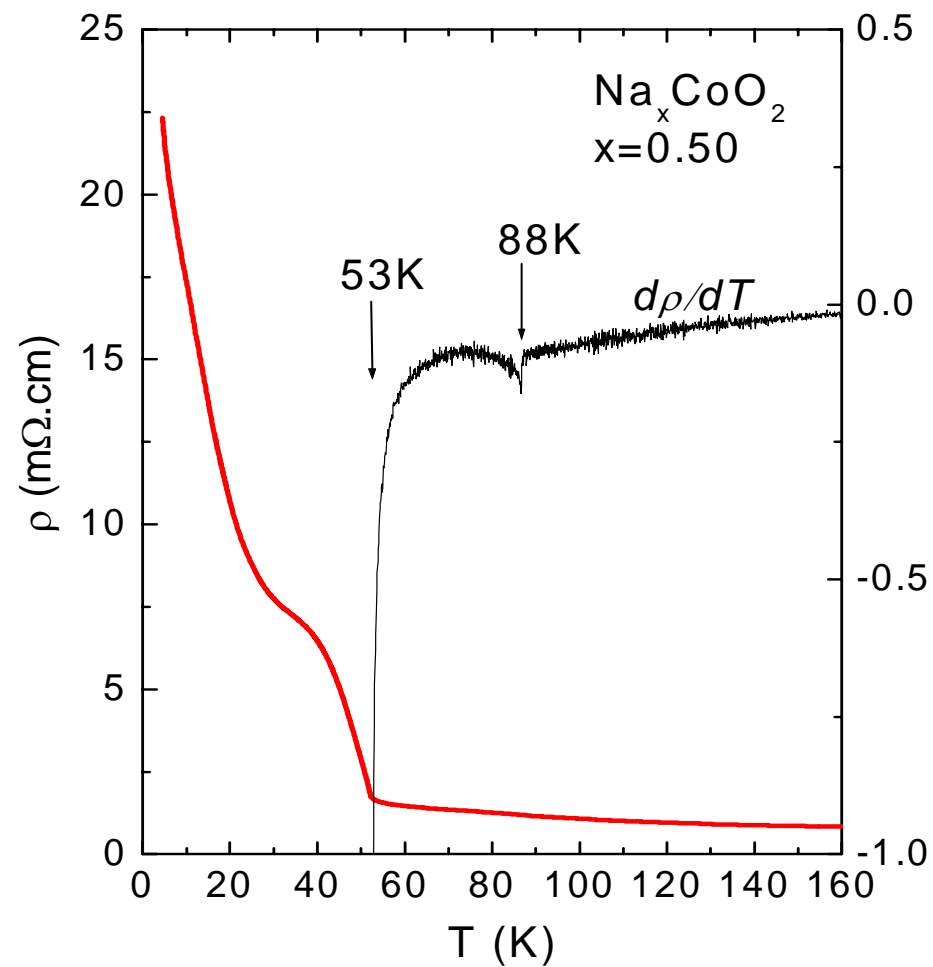
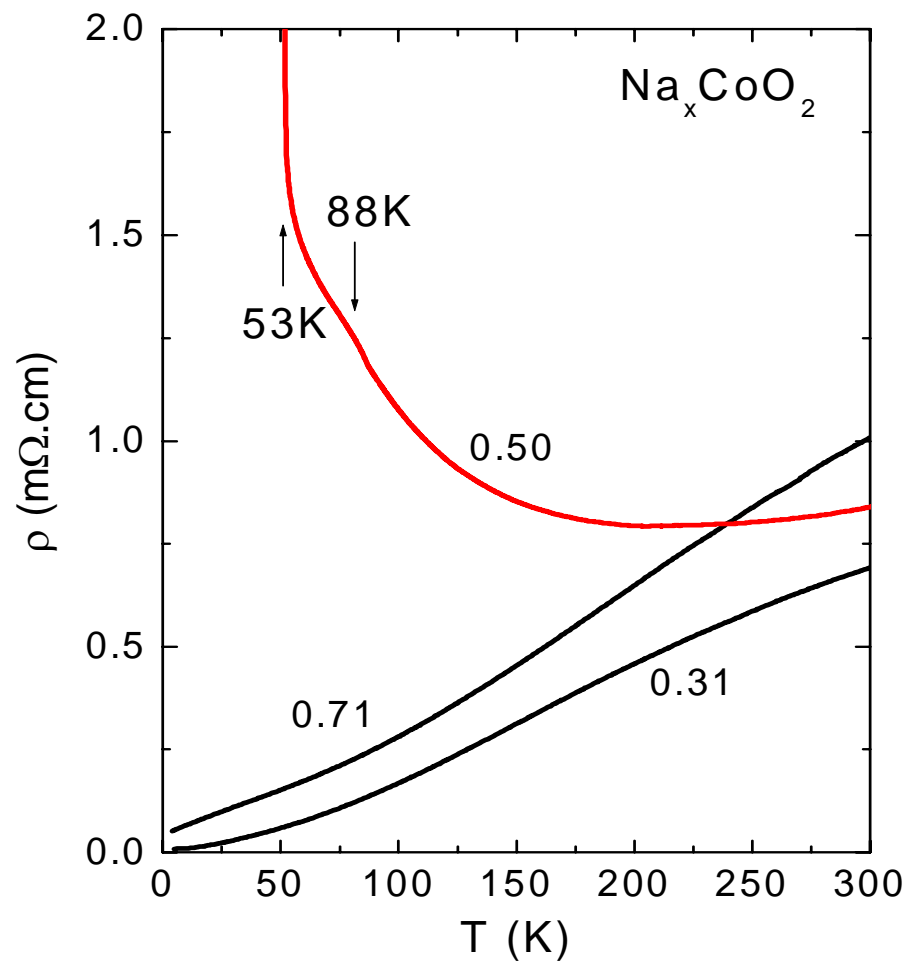
- treat powder and crystals under same conditions
- powder x-ray diffraction to get c-axis lattice constant
- ICP-AES to determine the Na contents of powders
- x vs. c-axis calibration curve
- from the c-axis of crystal, extract the Na content

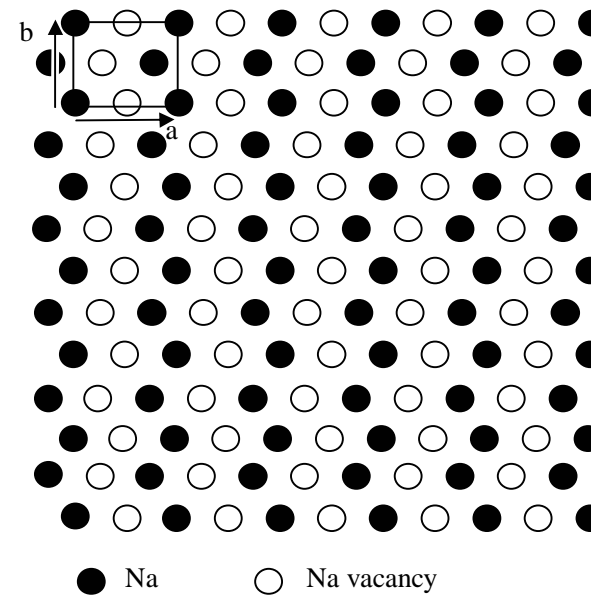
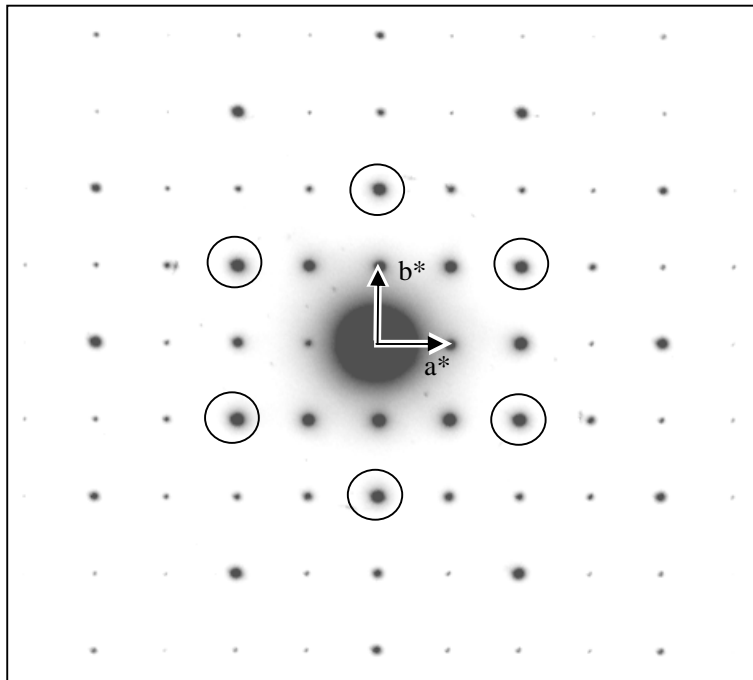


$x = 0.50$ (1/2):

- Two kinks at $T_{c1}=88\text{K}$ and $T_{c2}=53\text{K}$ in χ
- Resistivity shows insulating behavior below $T=53\text{K}$

An unexpected insulator at $x = \frac{1}{2}$

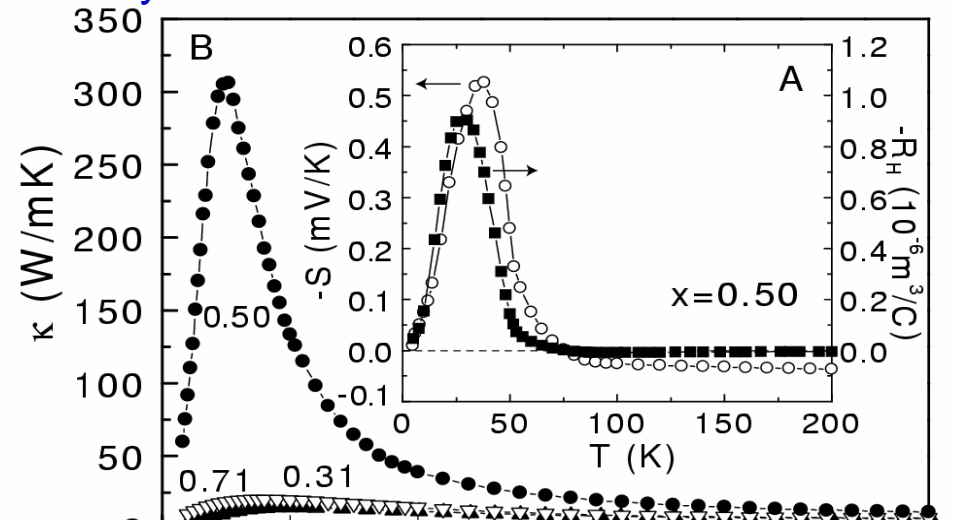




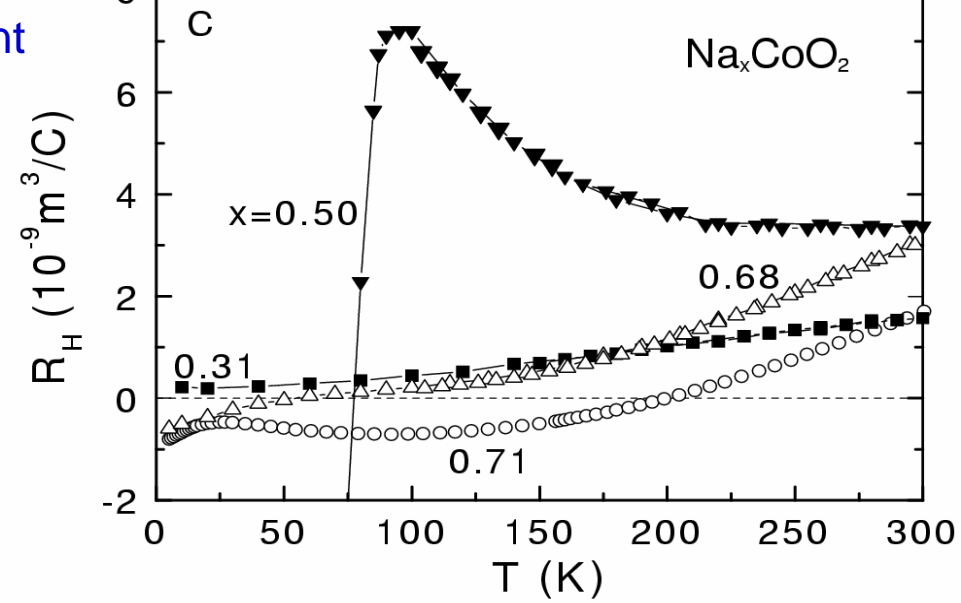
Electron diffraction at 300K shows the superlattice formed by the Na ions, consistent with a zig-zag order

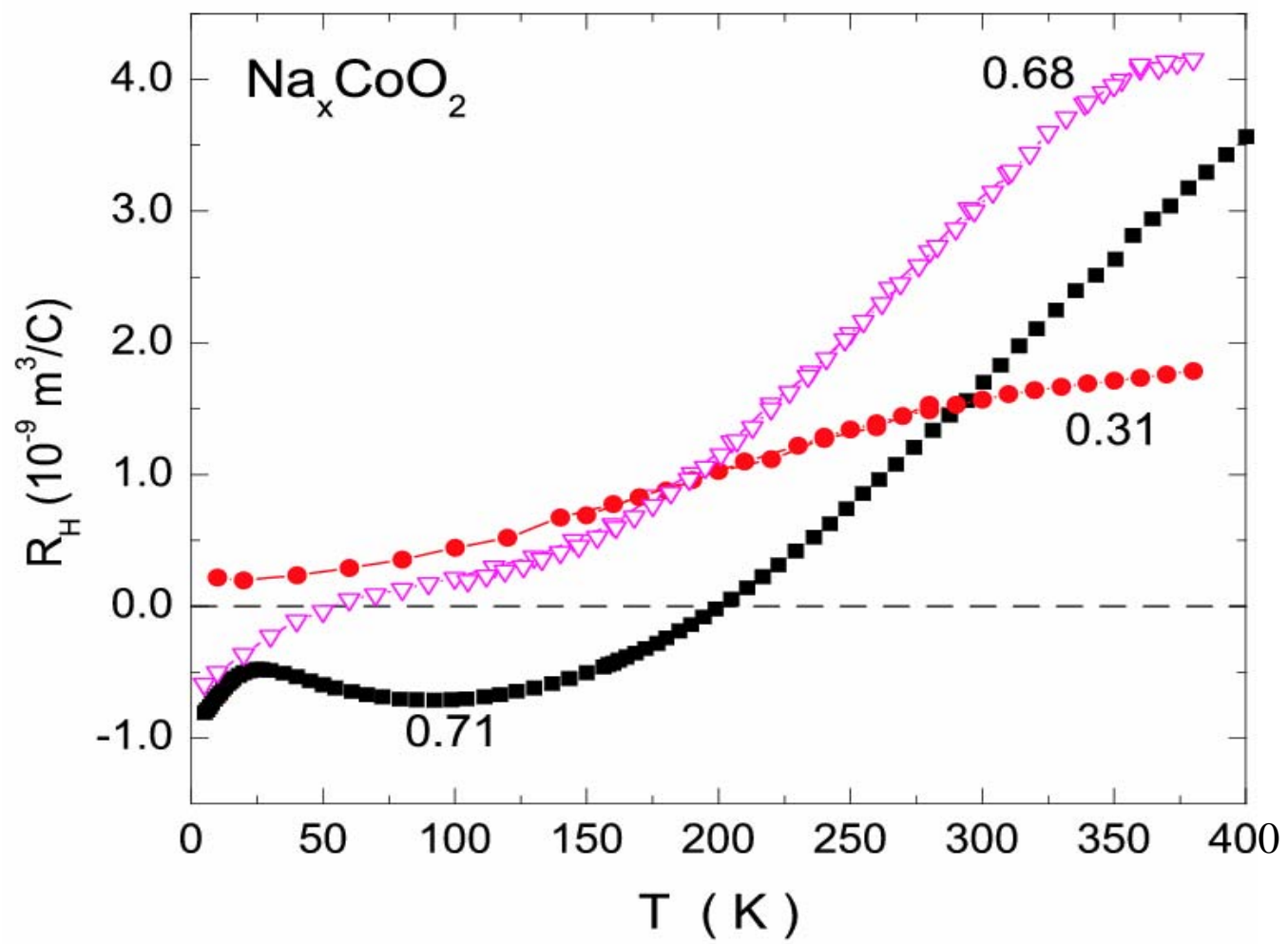
Zendbergen et al., condmat/0403206 (2004)

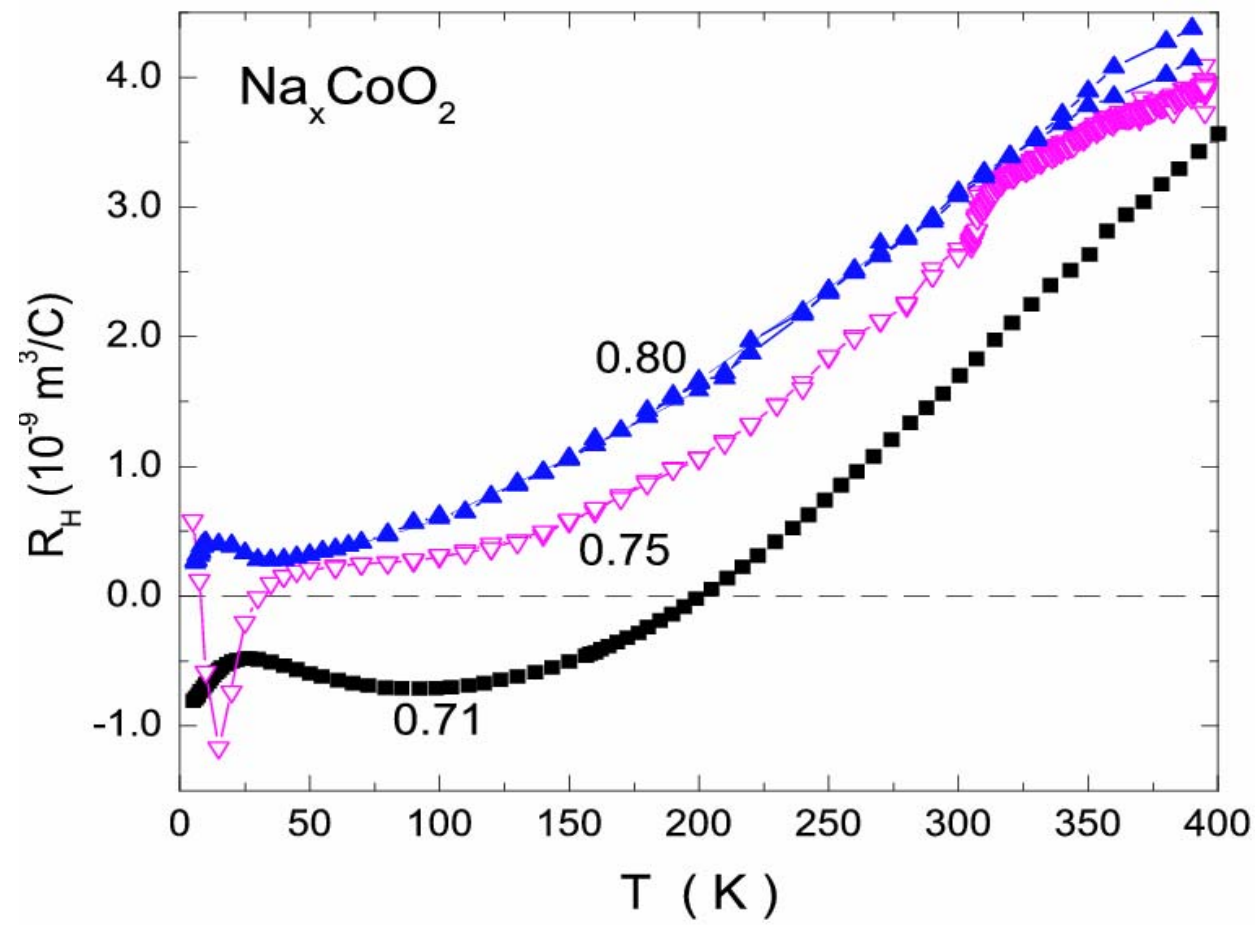
Thermal Conductivity



Hall coefficient



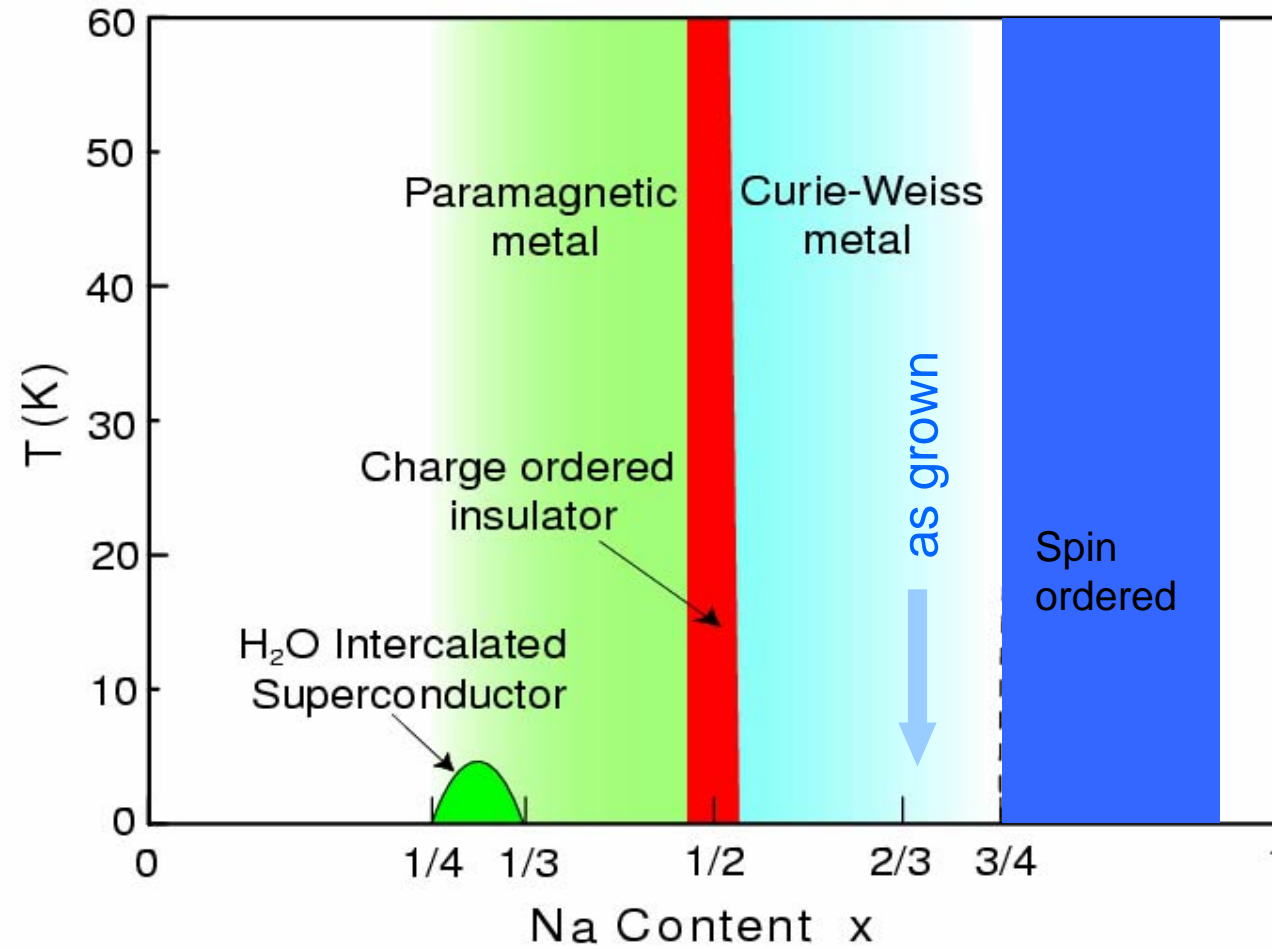




Na_xCoO_2

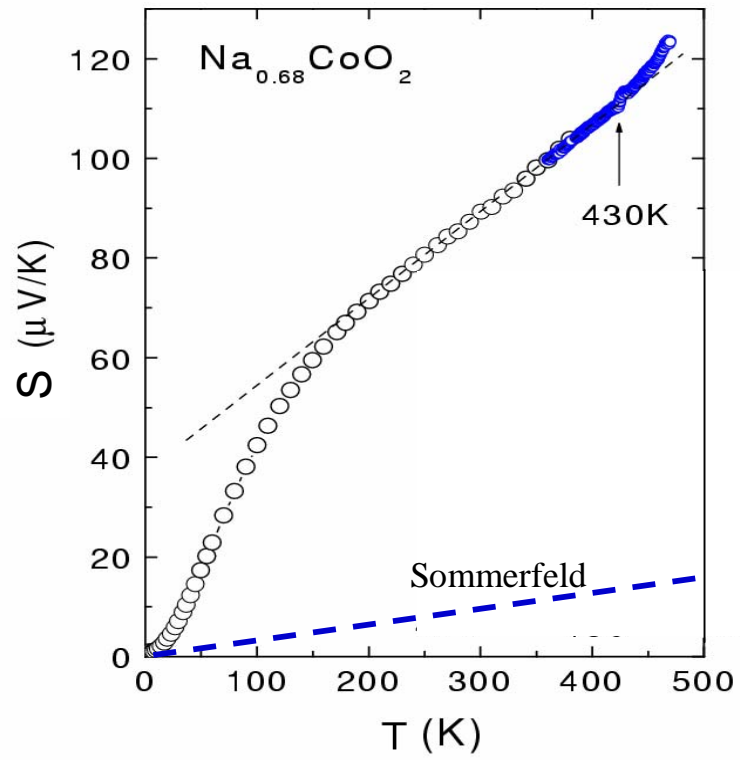
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PRL '04

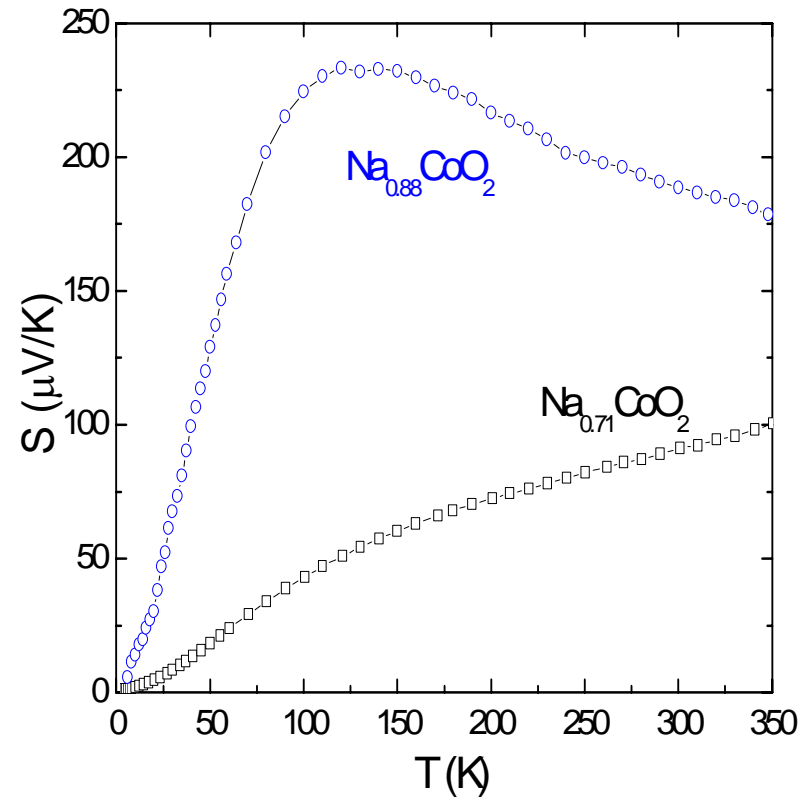


Further enhancement of thermopower

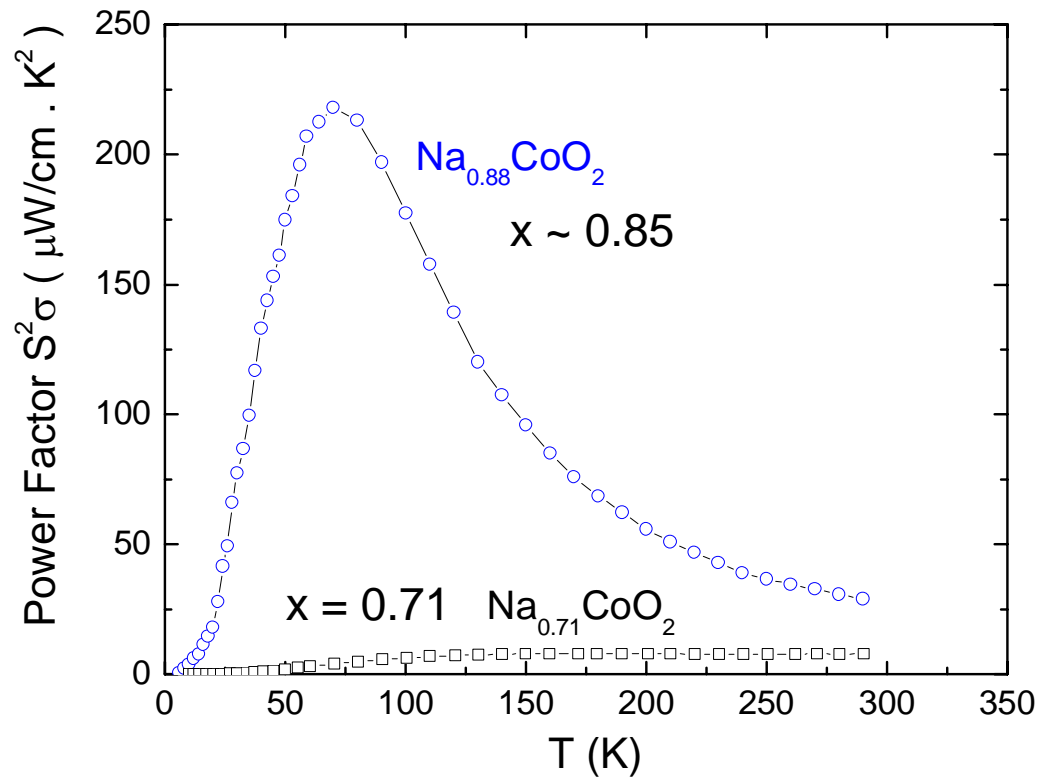
$x = 0.71$



$x = 0.88$



Power factor of Na_xCoO_2 $P = S^2\sigma$



Unusual electronic behavior in Na_xCoO_2

Strongly correlated $s = 1/2$ holes hopping on triangular lattice

- **Paramagnetic Metal ($x \sim 1/3$)**

High conductivity, superconducting with H_2O intercalatn.

- **Charge-ordered Insulator ($x = 1/2$)**

Na ion ordering, hole ordering (stripes?),
giant thermal conductivity

- **Curie-Weiss metal ($x \sim 2/3$)**

Curie-Weiss susceptibility, metallic cond., large thermopower
from spin entropy, T-linear Hall coef.

- **Spin Ordered Phase ($x > 3/4$)**

Even larger thermopower, field-induced metamagnetism