



International Atomic Energy Agency

SMR.1663-15

SUMMER SCHOOL ON PARTICLE PHYSICS

13 - 24 June 2005

QCD Phase Transitions and Heavy ion Collisions - Part 1

R. VENUGOPALAN Department of Physics Brookhaven National Laboratory Building 510A P.O. Box 5000 Upton, NY 11973-5000 U.S.A.

xQCD: QCD at high energies, temperatures & baryon densities

Raju Venugopalan Brookhaven National Laboratory

ICTP, Trieste, June 20th -24th, 2005

Outline of 4 lectures

Lecture I : General Introduction, QCD at high energies.

Lecture II: QCD at high energies-continued.

Heavy Ion Collisions: creating hot & dense matter in the laboratory.



Lecture III: Exploring the phase diagram of QCD

Lecture IV: Heavy Ion Collisions at RHIC.

What has RHIC wrought? What awaits at LHC?



Many body / statistical QCD at

High Temperatures,

High Baryon Densities,

and yes,

at high energies...

Outstanding phenomenological questions in heavy ion collisions

Can we learn about the universal properties of hadronic wavefunctions at high energies ? (Color Glass Condensate)

How does the CGC melt when heated in high energy heavy ion collisions? Does the matter thermalize to form a Quark Gluon Plasma (QGP) ?

Is the QGP formed at RHIC? What can we learn about its properties?

Outline of lecture I

QCD preliminaries

What does a hadron look like at high energies?

Quantum Chromodynamics (QCD) in the news...







The 2004 Nobel prize in physics



For the discovery of asymptotic freedom in QCDthe theory of the strong interaction...



Fundamental particles and their interactions

QUARKS S=1/2		LEPTONS S=1/2		GAUGE BOSONS S=1	
Q = -2/3	Q = -1/3	Q = -1	Q = 0	quanta	
uuu	d d d	e	ν _e	9 ₁ 9 ₈	
m=(1-4) 10 ⁻³	m=(5-8) 10 ⁻³	m=5.11 10 ⁻⁴	m<3 10 ⁻⁹	m< a few 10 ⁻³	
ссс	S S S	μ	ν_{μ}	γ	
m=1.0-1.4	m=0.08-0.15	m=0.10566	m<1.9 10 ⁻⁴	m<2 10 ⁻²⁵	
† † †	bb b	τ	ν_{τ}	W ±, Z ⁰	
m=174.3±5.1	m=4.0-4.5	m=1.7770	m<18.2 10 ⁻³	m_W =80.432 ±0.39,	
				m _Z =91.1876 ±0.0021	

All masses in GeV units

	Interaction	exchanged boson	relative strength	example
QCD →	Strong	Gluon (g)	1	
▼	Electromagnet.	Photon (y)	$\frac{1}{137}$	°u
QED	Weak	w +, w -, z •	10^{-14}	$\gtrsim \sim \sim$
	Gravitation	Graviton (G)?	10 ⁻⁴⁰	$^{u}_{u} > \stackrel{G}{\sim} \stackrel{e}{<}_{e}$



In QCD, one has "anti-screening" because colored gluons can interact with each other...



Photons do not self-interact in QED

The coupling between color charges gets weaker at high energies- or short distances...



Potential between static quark-anti-quark pair grows linearly at large distances provides intuitive picture of confinement

DEEPLY INELASTIC SCATTERING

THE SIMPLEST WAY TO STUDY QCD AT SMALL DISTANCES...

R. Hofstader



Kinematic Invariants:

Center of mass energy square s Momentum resolution square Q^2 $x_{\rm Bj}\,$ -an invariant to be defined shortly

$$x_{\rm Bj} \approx \frac{Q^2}{s}$$

DIS inclusive cross-section:



Rutherford cross-section

Nobel to Friedman, Kendall, Taylor



Bj-scaling - apparent scale invariance of structure functions...

Feynman and Bjorken explanation of scaling puzzle...

Parton model



 $x_{\rm Bj} = {\rm Fraction \ of \ hadron \ momentum \ carried \ by \ a \ parton...}$

Puzzle resolved in QCD



QCD \neq Parton model Logarithmic scaling violations

$$F_2(x,Q^2) = \sum_{\substack{q=u,c,t \\ d,s,b}} e_q^2 \left(x \, q(x,Q^2) + x \, \bar{q}(x,Q^2) \right)$$

The Hadron at high energies: I



QCD-logarithmic corrections



"X"-QCD--RG evolution

$$\int_0^1 \frac{dx}{x} \left(xq(x) - x\bar{q}(x) \right) = 3$$

of valence quarks

$$\int_0^1 \frac{dx}{x} (xq(x) + x\bar{q}(x)) \to \infty$$
quarks...

Renormalization group (RG) evolution:



DGLAP evolution: Linear RG in Q^2

Dokshitzer-Gribov-Lipatov-Altarelli-Parisi



Sum large # of Bremsstrahlung logs -# of gluons grows rapidly at small x...



But... the phase space density decreases -the proton becomes more dilute

Much of the discussion in pQCD has focused on the Bjorken limit:

$$Q^2 \to \infty; s \to \infty; x_{\rm Bj} \approx \frac{Q^2}{s} = \text{fixed}$$

Asymptotic freedom, the Operator Product Expansion (OPE) & Factorization Theorems: machinery of precision physics in QCD...

STRUCTURE OF HIGHER ORDER CONTRIBUTIONS IN DIS



- Coefficient functions C computed to NNLO for many processes, e.g., gg -> H Harlander, Kilgore; Ravindran, Van Neerven, Smith; ...
- Splitting functions -P computed to 3-loops recently! Moch, Vermaseren, Vogt

The other interesting limit-is the Regge limit of QCD:

$$x_{\rm Bj} \to 0; s \to \infty; Q^2 (>> \Lambda_{\rm QCD}^2) = \text{fixed}$$

Physics of strong fields in QCD, multi-particle productionnovel universal properties of theory in this limit ?

BFKL evolution: Linear RG in x

Balitsky-Fadin-Kuraev-Lipatov



of gluons grows even more rapidly...





Proton is a dense many body system at high energies

Mechanism for parton saturation:

Competition between "attractive" bremsstrahlung and "repulsive" recombination effects.

Maximal phase space density =>

$$\frac{1}{2(N_c^2 - 1)} \frac{x G(x, Q^2)}{\pi R^2 Q^2} = \frac{1}{\alpha_S(Q^2)}$$

Saturated for

$$Q = Q_s(x) >> \Lambda_{\rm QCD} \approx 0.2 \,\,{\rm GeV}$$

* Higher twists (power suppressed-in Q^2) are important when: $Q^2 pprox Q_s^2(x) >> \Lambda_{ m QCD}^2$

* Leading twist "shadowing" of these contributions can extend up to at small x. $Q^2 >> Q_s^2(x)$

Need a new organizing principlebeyond the OPE- at small x.

Golec-Biernat & Wusthoff's model



Parameters:

 $Q_0 = 1 \,\text{GeV}; \lambda = 0.3; x_0 = 3 \cdot 10^{-4}$

Geometrical scaling at HERA



Road map of the strong interactions



(Momentum Resolution)

Total Cross-sections

Donnachie & Landshoff





Novel regime of QCD evolution at high energies



Extra Slides

Hard diffractive processes



Diffractive Surprises at HERA



Approximate 10% of events are hard diffractive events!

Comparison with Data

FS model with/without saturation and **IIM CGC model** hep-ph /0411337.

