



The Abdus Salam  
International Centre for Theoretical Physics



SMR.1663- 12

## *SUMMER SCHOOL ON PARTICLE PHYSICS*

*13 - 24 June 2005*

### Neutrino Physics - Part 2

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# Lecture 2: Neutrino Mixing

John Beacom, The Ohio State University



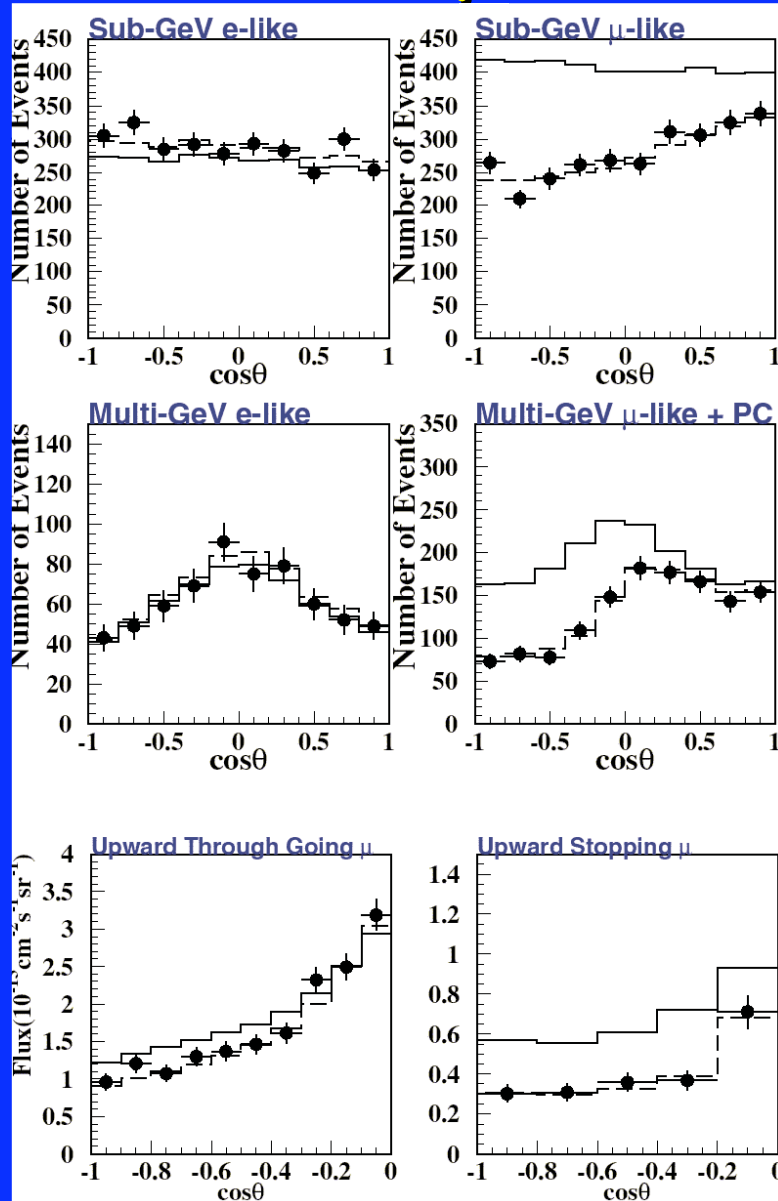
# Elevator Pitch

- Neutrino mixing in vacuum is almost certain
- Neutrino mixing in matter is somewhat certain
- Smoking gun signatures are still needed
- As are precision mixing parameter measurements
- These are keys to understanding neutrino mass

# Neutrino Vacuum Mixing

(On the blackboard)

# Atmospheric Neutrino Results



Super – Kamiokande :

$$\sin^2 2\theta \simeq 1$$

$$\delta m^2 \simeq 3 \times 10^{-3} \text{ eV}^2$$

$\nu_\mu \rightarrow \nu_\tau$  preferred

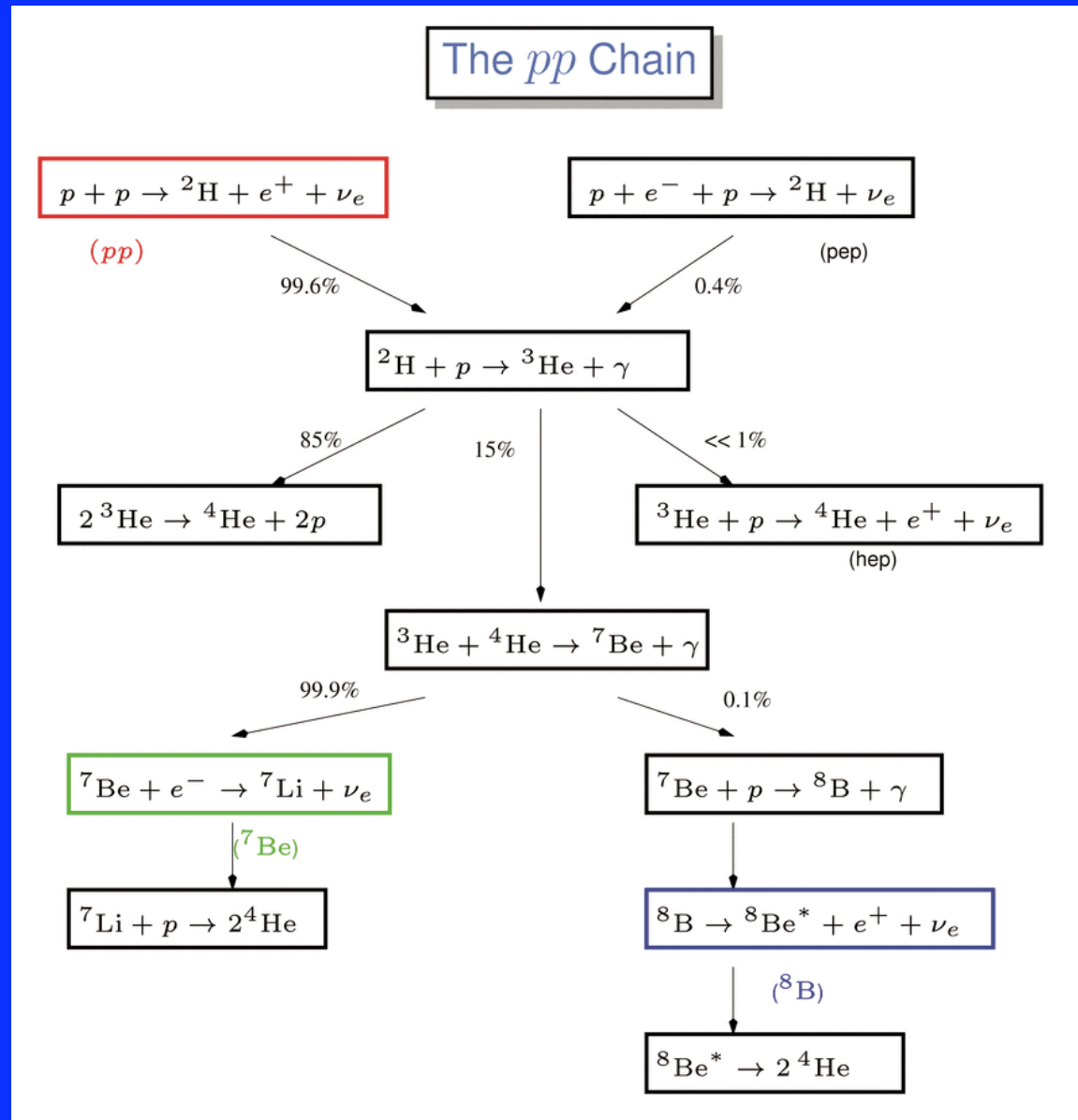
Soudan-2, MACRO, ...

K2K, MINOS, off-axis, JHF, ...

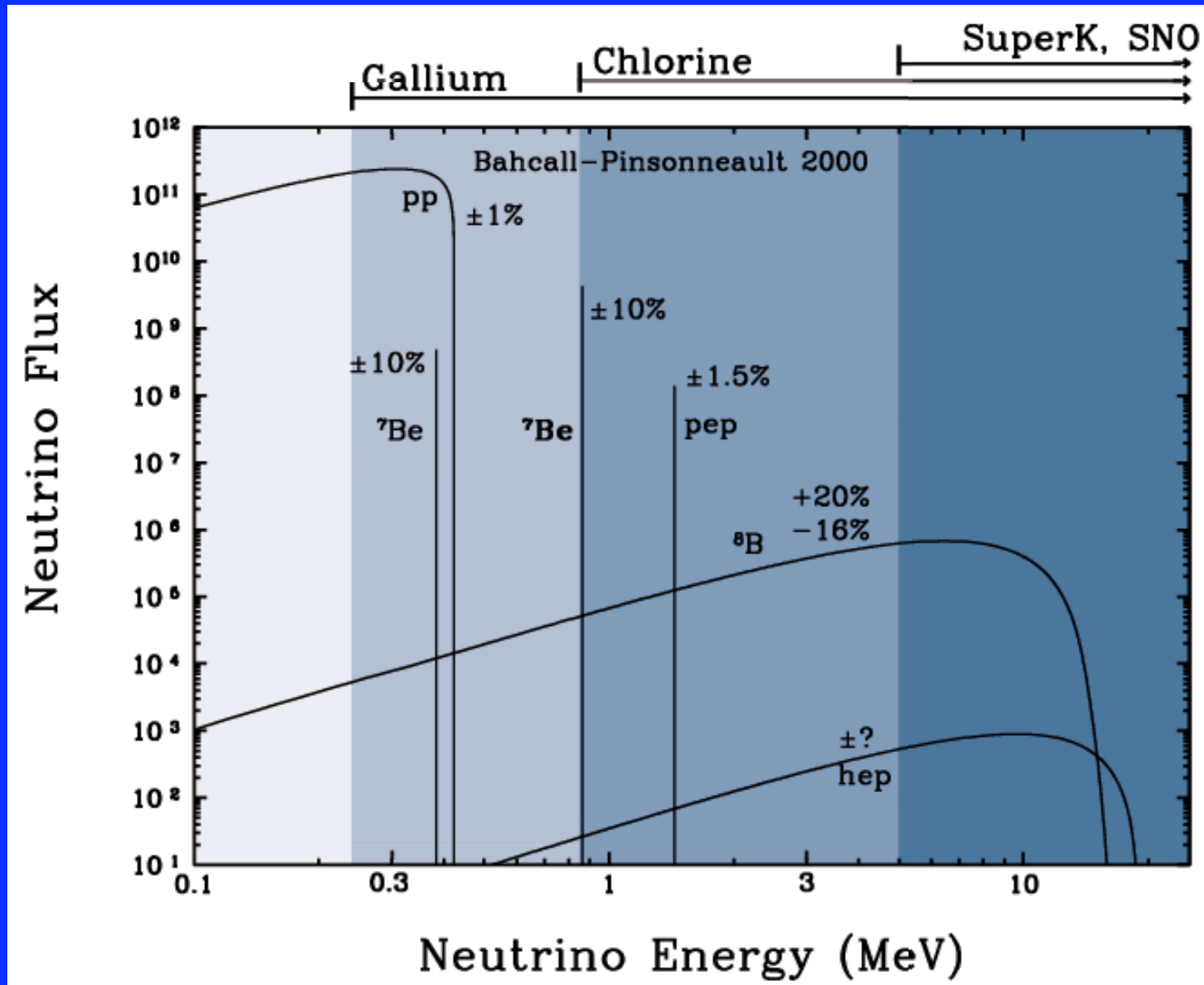
HK, UNO, ...

# Solar Neutrinos

# Solar Nuclear Reactions



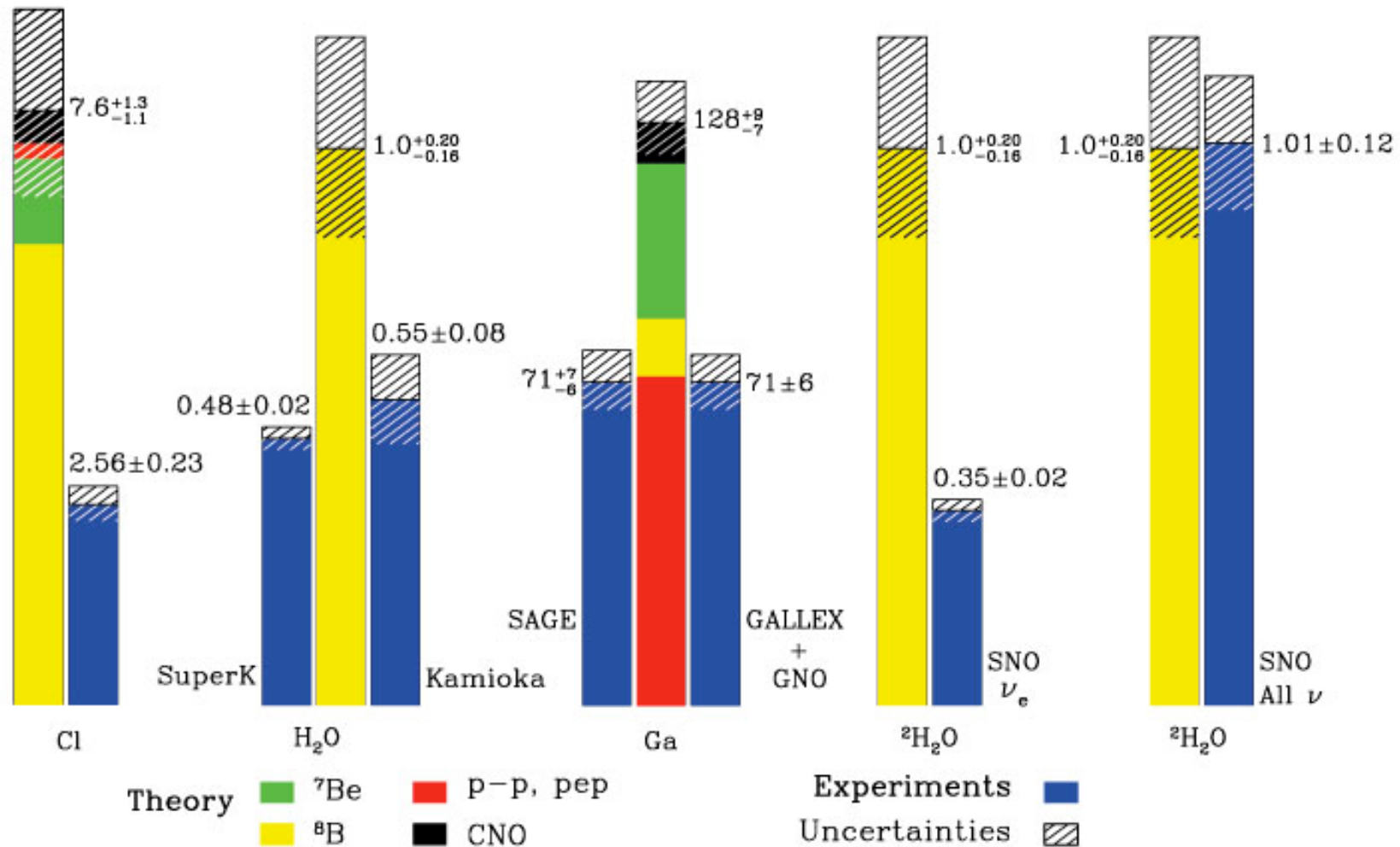
# Solar Neutrino Fluxes



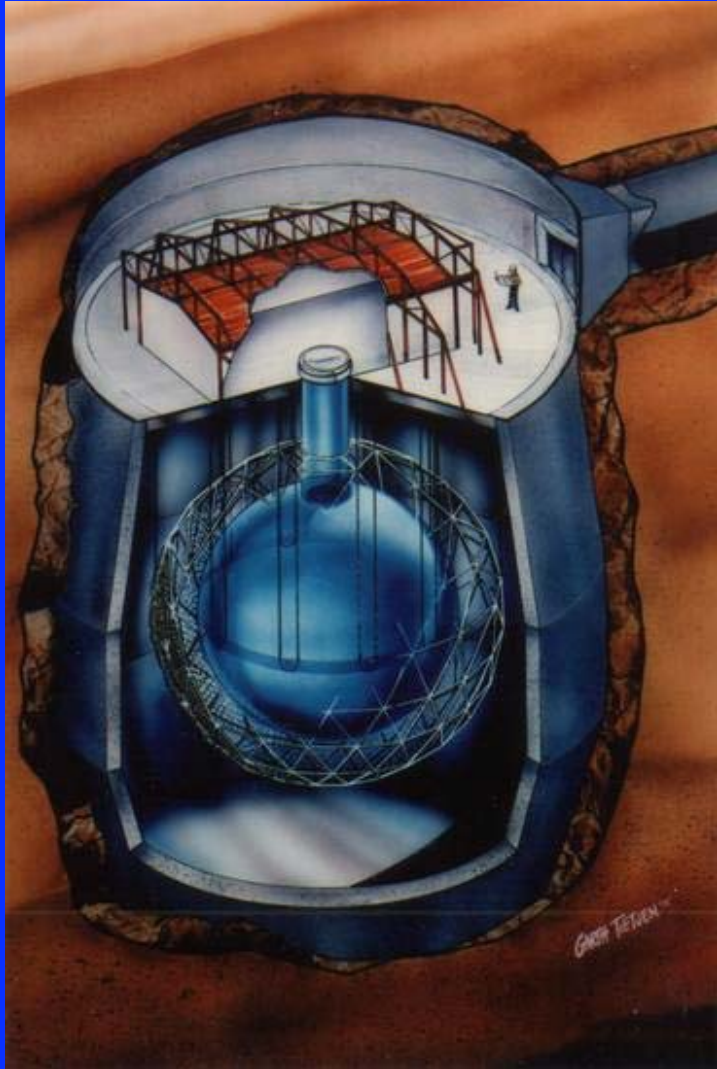


# The Solar Neutrino Problem

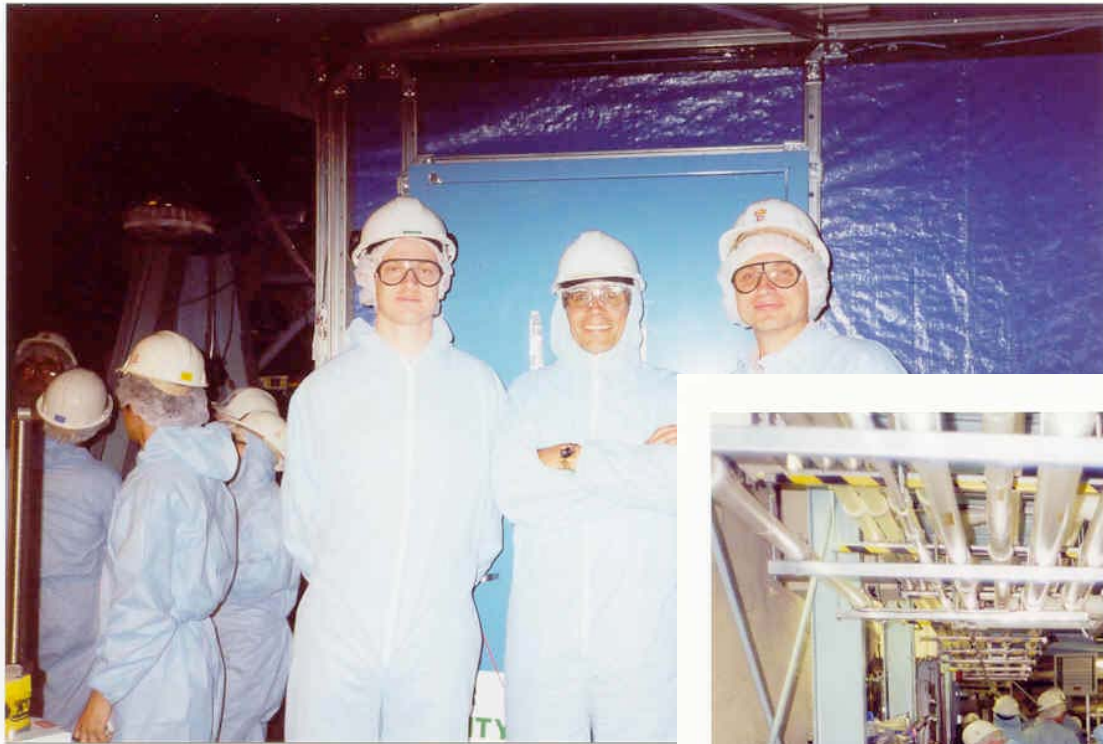
Total Rates: Standard Model vs. Experiment  
Bahcall-Pinsonneault 2000



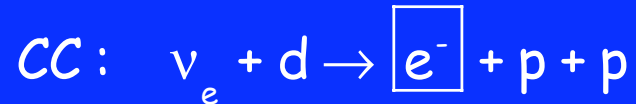
# Sudbury Neutrino Observatory



# Inside SNO

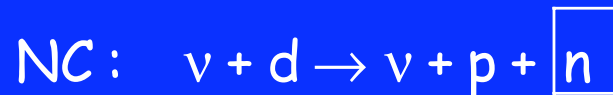


# SNO Detection Reactions



measures  $\boxed{\phi_{\nu_e}}$

where  $E_\nu \simeq E_e + \text{few MeV}$



measures  $\boxed{\phi_{\nu_e} + \phi_{\nu_\mu}}$

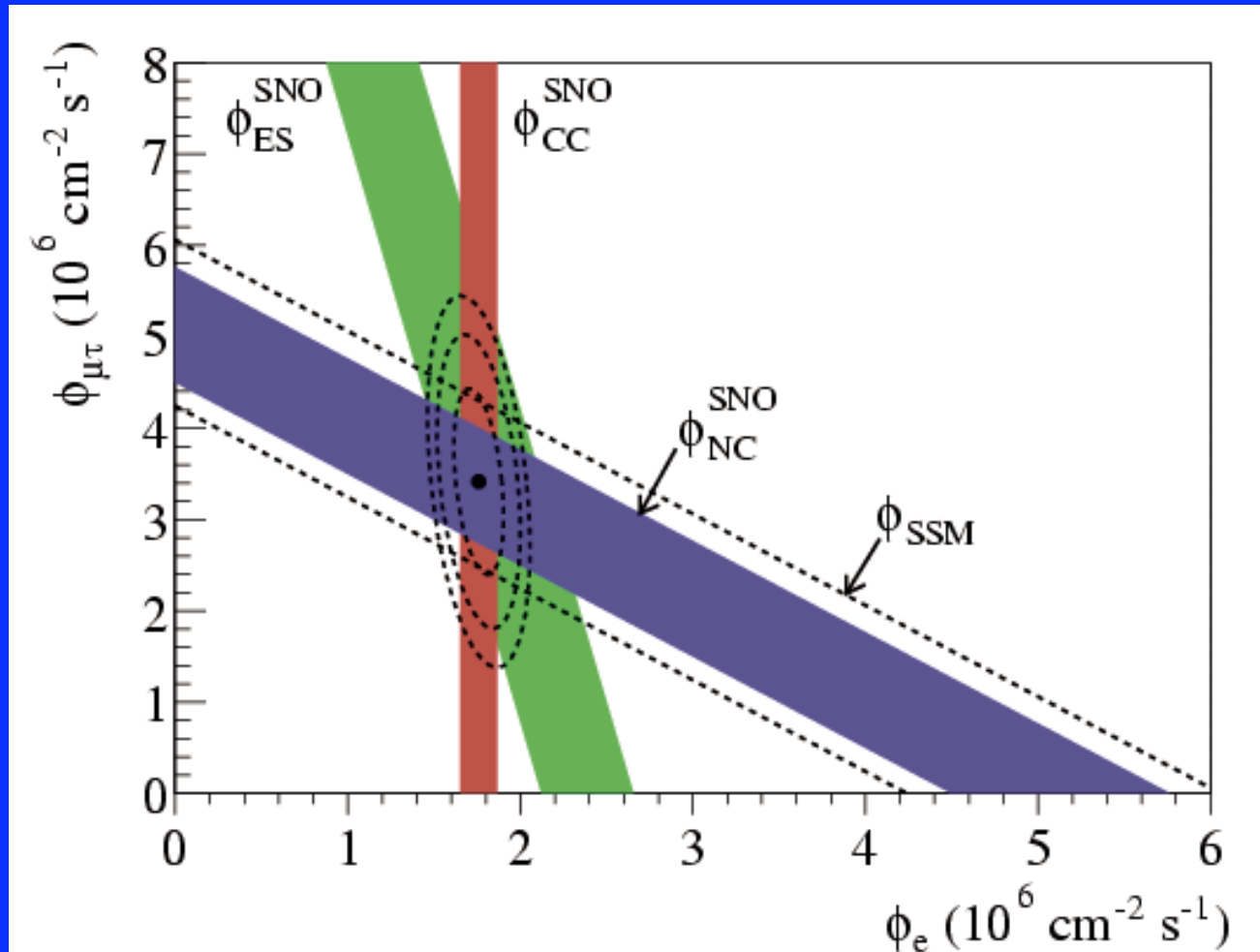
where  $E_\nu$  unknown



measures  $\boxed{\phi_{\nu_e} + \frac{1}{6}\phi_{\nu_\mu}}$

where  $E_\nu \sim 2E_e$

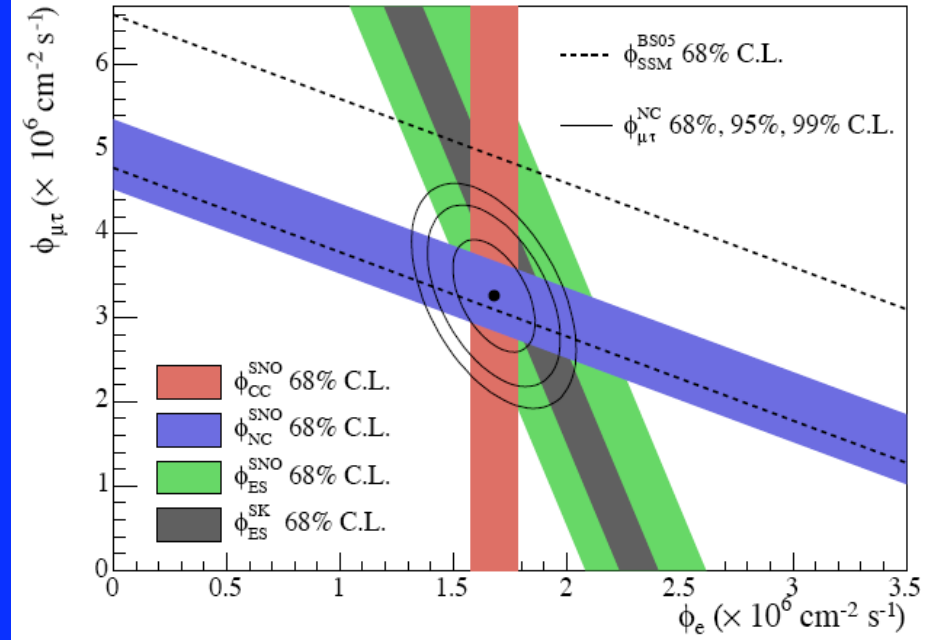
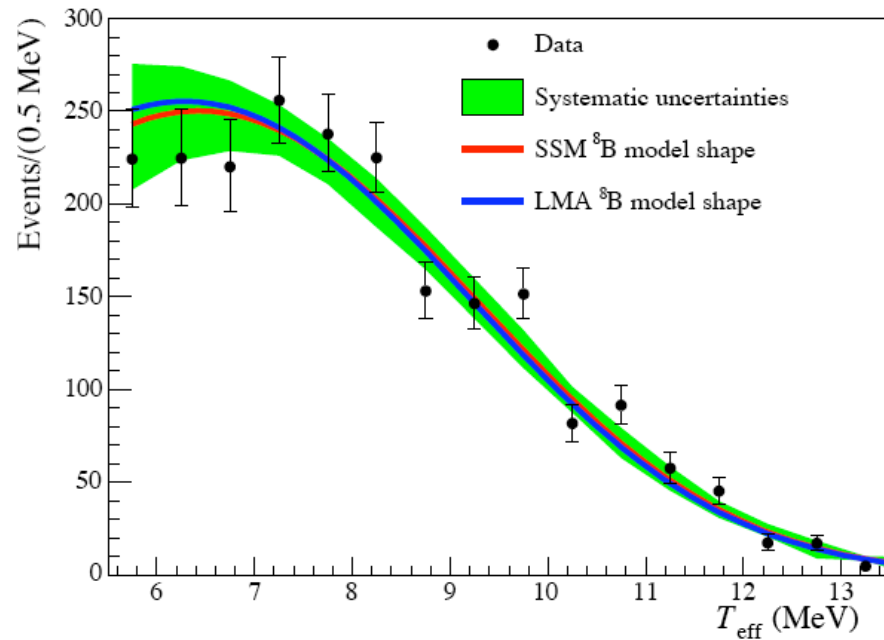
# SNO Flux Measurements



$$\phi_{\nu_e} = \frac{1}{3}$$

$$\phi_{\nu_e} + \phi_{\nu_{\mu}} = 1$$

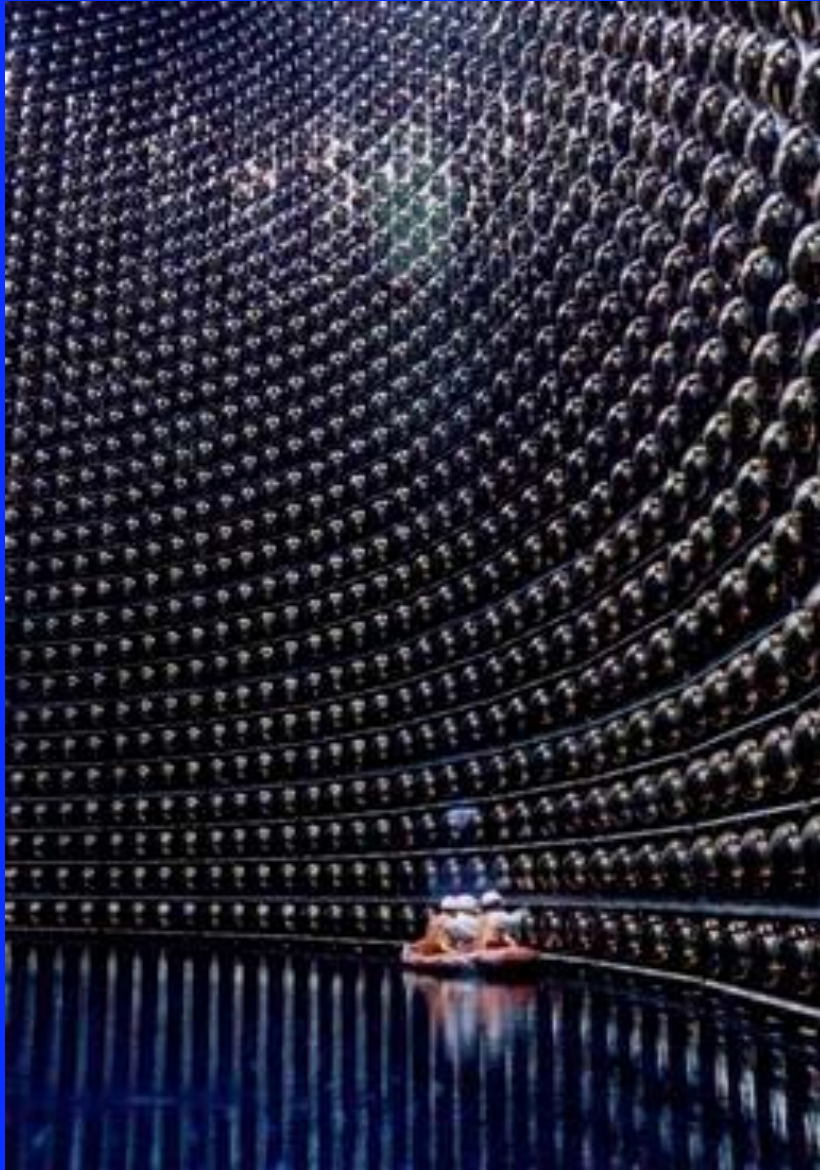
# Latest SNO Results



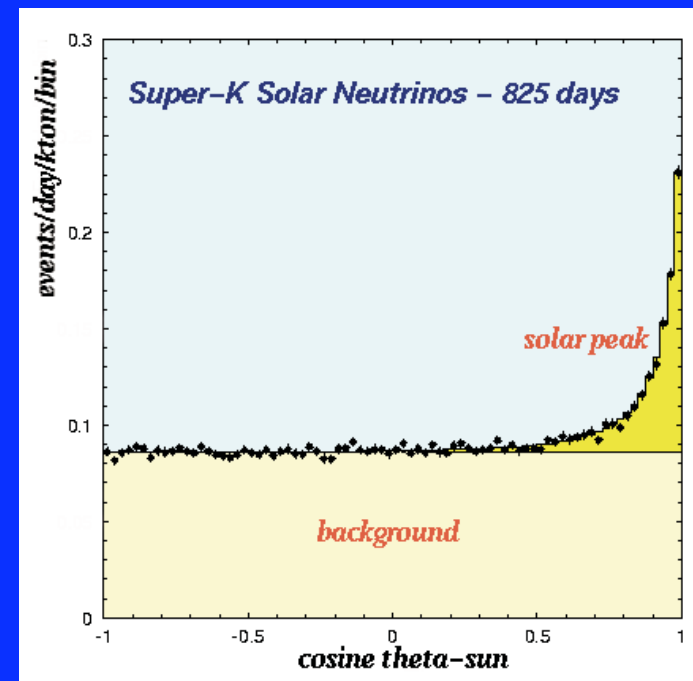
Aharmim et al, nucl-ex/0502021

Still NO evidence of:  
spectral distortion or day-night effect!

# Super-Kamiokande



$$\nu + e^{-} \rightarrow \nu + e^{-}$$



# Super-Kamiokande Results



$$\text{Ratio} = \frac{\text{measured } e^- \text{ spectrum}}{\text{expected } e^- \text{ spectrum}}$$

$$\phi_{\nu_e} + \frac{1}{6} \phi_{\nu_\mu} = 0.34 + \frac{1}{6} 0.66 = 0.45$$

Perfect Agreement!



# MS & W Papers

## 1) RESONANT AMPLIFICATION OF NEUTRINO OSCILLATIONS IN MATTER AND SOLAR NEUTRINO SPECTROSCOPY.

By S.P. Mikheev, A.Yu. Smirnov (Moscow, INR),. 1986.

Published in **Nuovo Cim.C9:17-26,1986**

TOPCITE = 1000+

[References](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [BibTeX](#) | [Keywords](#) | Cited [1225 times](#)

## 2) RESONANCE AMPLIFICATION OF OSCILLATIONS IN MATTER AND SPECTROSCOPY OF SOLAR NEUTRINOS.

By S.P. Mikheev, A.Yu. Smirnov (Moscow, INR),. 1986.

Published in **Sov.J.Nucl.Phys.42:913-917,1985, Yad.Fiz.42:1441-1448,1985**

TOPCITE = 1000+

[References](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [BibTeX](#) | [Keywords](#) | Cited [1838 times](#) | [More Info](#)

## 3) NEUTRINO OSCILLATIONS AND STELLAR COLLAPSE.

By L. Wolfenstein (Carnegie Mellon U.),. COO-3066-133, Aug 1979. 6pp.

Published in **Phys.Rev.D20:2634-2635,1979**

TOPCITE = 500+

[References](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [BibTeX](#) | [Keywords](#) | Cited [755 times](#)

[ADS Abstract Service](#)

[Phys. Rev. D Server](#)

[Scanned Version](#) (KEK Library)

## 4) NEUTRINO OSCILLATIONS IN MATTER.

By L. Wolfenstein (Carnegie Mellon U.),. COO-3066-102, Nov 1977. 19pp.

Published in **Phys.Rev.D17:2369,1978**

TOPCITE = 1000+

[References](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [BibTeX](#) | [Keywords](#) | Cited [2503 times](#) | [More Info](#)

[ADS Abstract Service](#)

[Phys. Rev. D Server](#)

# MSW Effect

$$H_{\text{flavor}} = \frac{\delta m^2}{4E} \begin{bmatrix} \zeta(t) - \cos 2\theta_v & \sin 2\theta_v \\ \sin 2\theta_v & -\zeta(t) + \cos 2\theta_v \end{bmatrix}$$

$$\zeta(t) = \frac{2\sqrt{2}G_F N_e(t)}{\delta m^2/E}$$

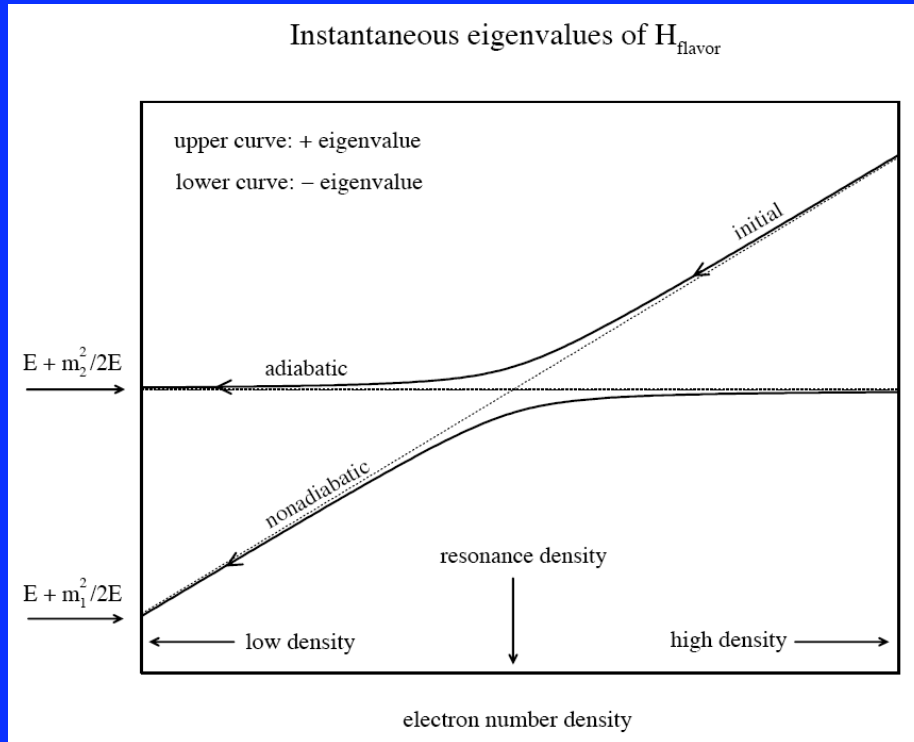
$$\begin{bmatrix} \Psi_1(t) \\ \Psi_2(t) \end{bmatrix} = \begin{bmatrix} \cos \theta(t) & -\sin \theta(t) \\ \sin \theta(t) & \cos \theta(t) \end{bmatrix} \begin{bmatrix} \Psi_e(t) \\ \Psi_\mu(t) \end{bmatrix}$$

$$E \pm \frac{m_1^2 + m_2^2 + \delta m^2 \zeta(t)}{4E} \mp \frac{\delta m^2}{4E} \sqrt{\sin^2 2\theta_v + (\zeta(t) - \cos 2\theta_v)^2}$$

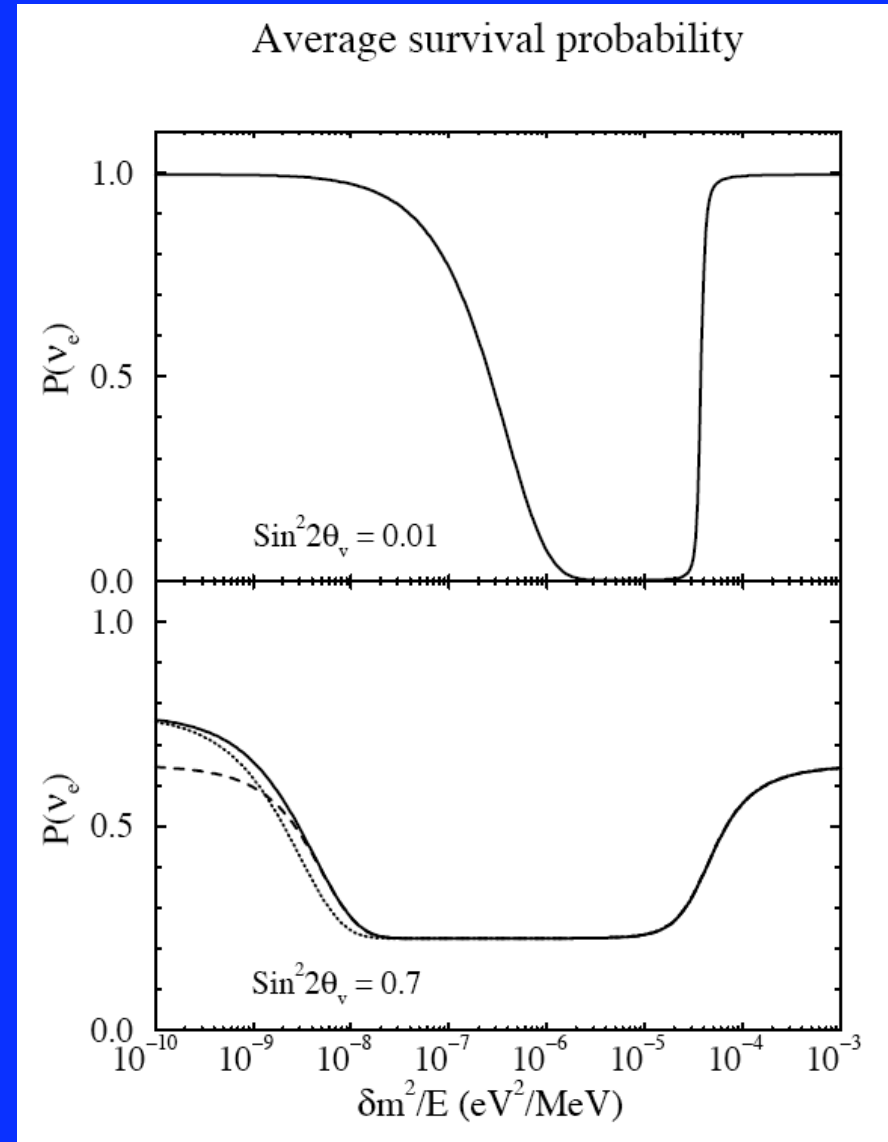
Parke,  
PRL 1986:

$$P(\nu_e \rightarrow \nu_e) = \frac{1}{2} [1 + (1 - 2P_{\text{hop}}) \cos 2\theta_i \cos 2\theta_v]$$

# Neutrino Matter Mixing

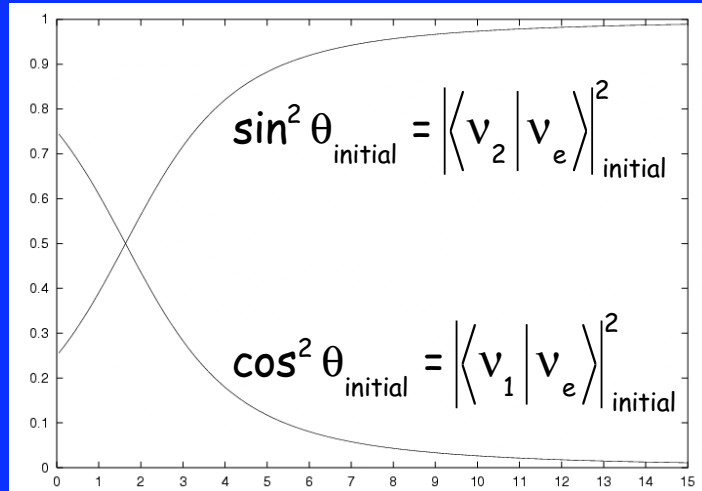


(graphics from Beacom and Balantekin, hep-th/9709117)



# Solar Neutrinos Don't Oscillate

1

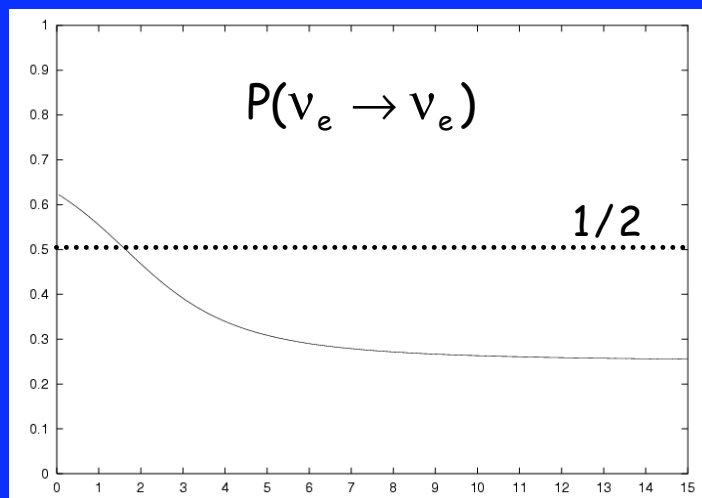


Matter effects  $\sim \frac{G_F N_e}{\delta m^2 / E}$

Propagation is adiabatic

0

1



$\phi_{\nu_e} \sim \left| \langle \nu_e | \nu_2 \rangle \right|_{\text{vacuum}}^2 = \sin^2 \theta_{\text{vacuum}}$

$\phi_{\nu_\mu} \sim \left| \langle \nu_\mu | \nu_2 \rangle \right|_{\text{vacuum}}^2 = \cos^2 \theta_{\text{vacuum}}$

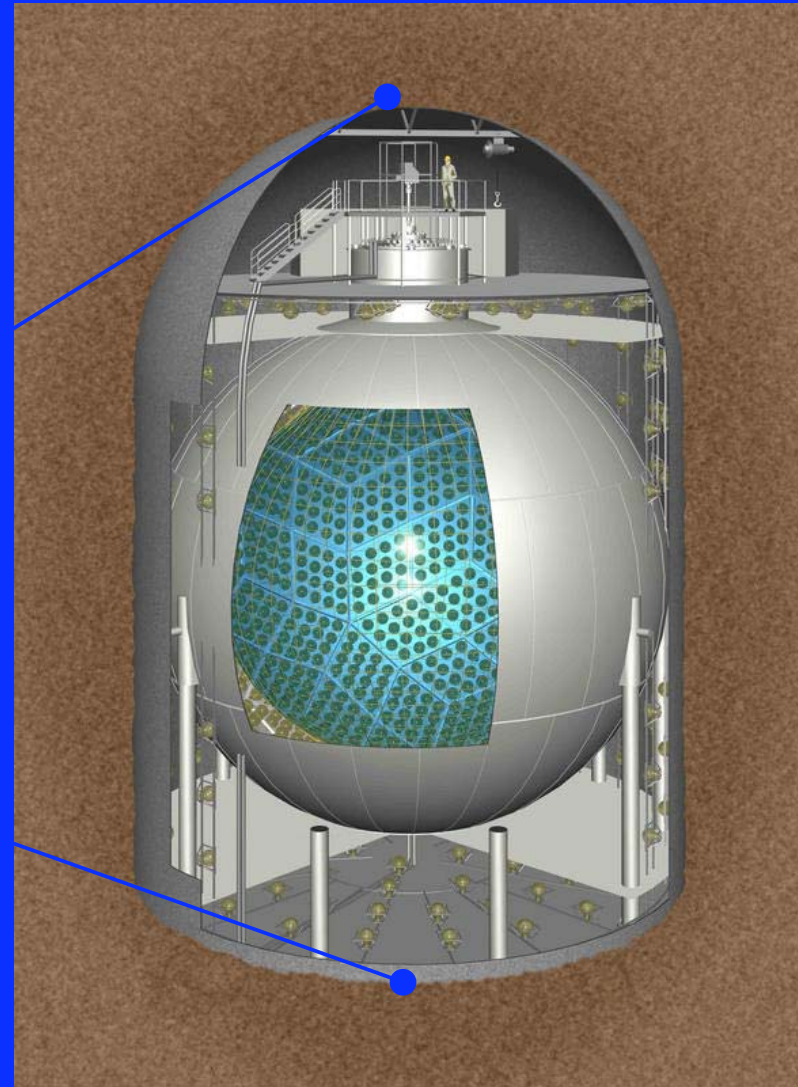
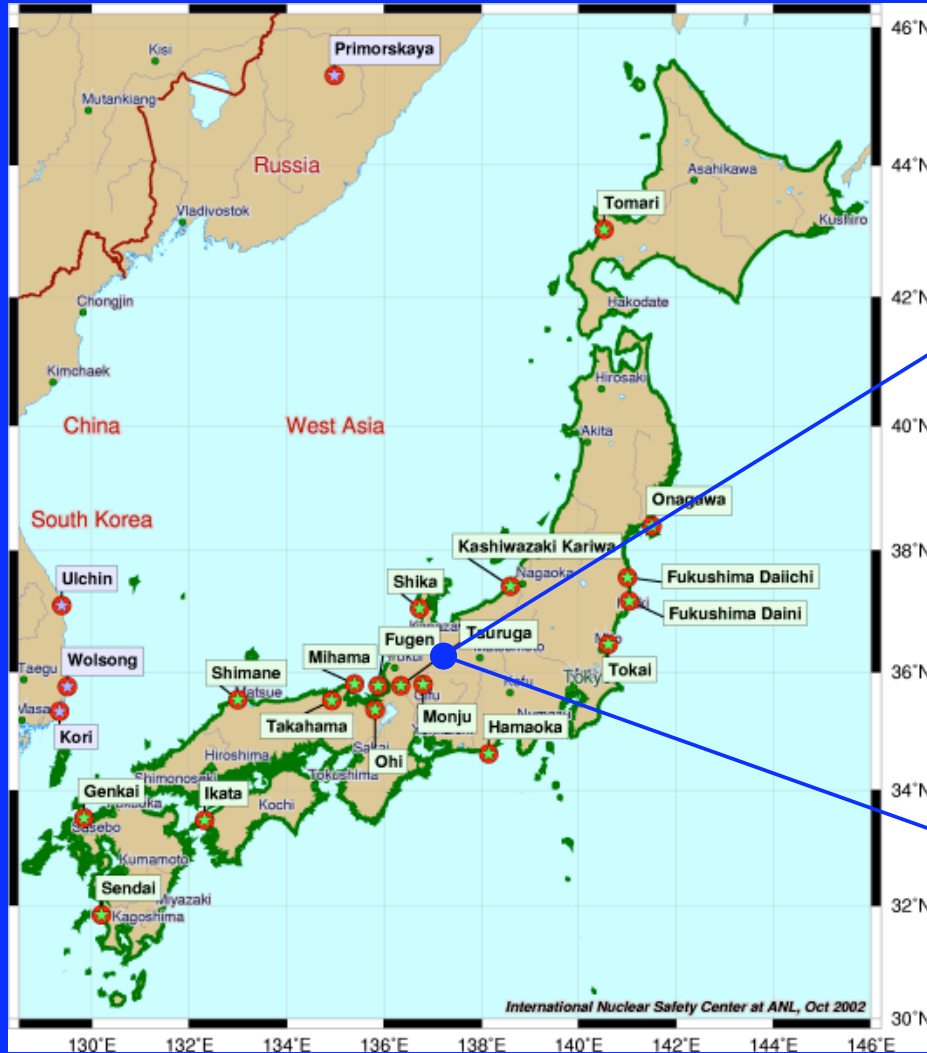
(above about 5 MeV)

0

E in [0,15 MeV]

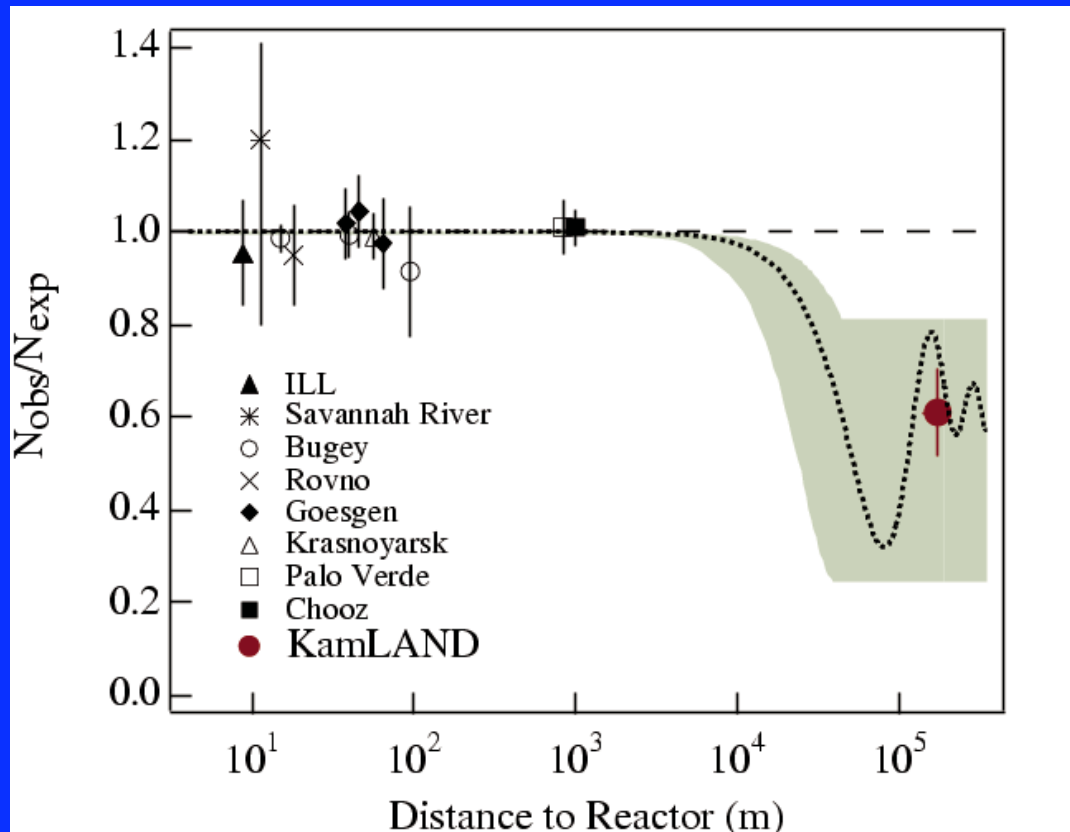
# Reactor Neutrinos

# KamLAND



# Vacuum Oscillation Probability

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu) = \sin^2(2\theta) \sin^2\left(1.27 \frac{\delta m^2 L}{E}\right)$$



$\sin^2 \theta_{\text{solar}}$  is large

$\sin^2 \theta_{13}$  is small

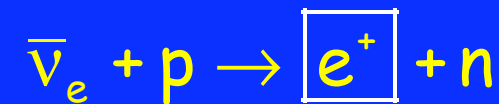
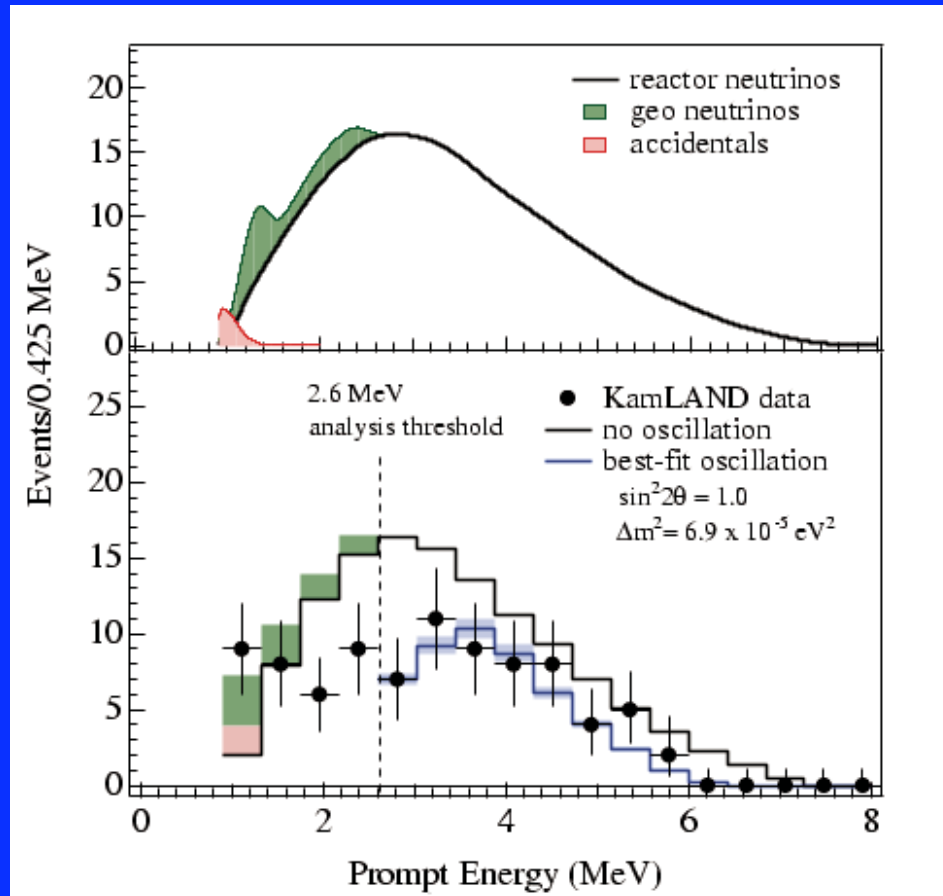
$$\delta m^2 \simeq 5 \times 10^{-5} \text{ eV}^2$$

$$E \simeq 4 \text{ MeV}$$

$\Rightarrow$  Detector at  $\simeq 10^5 \text{ m}$

Eguchi et al., PRL 90, 021802 (2003)

# KamLAND Neutrino Detection

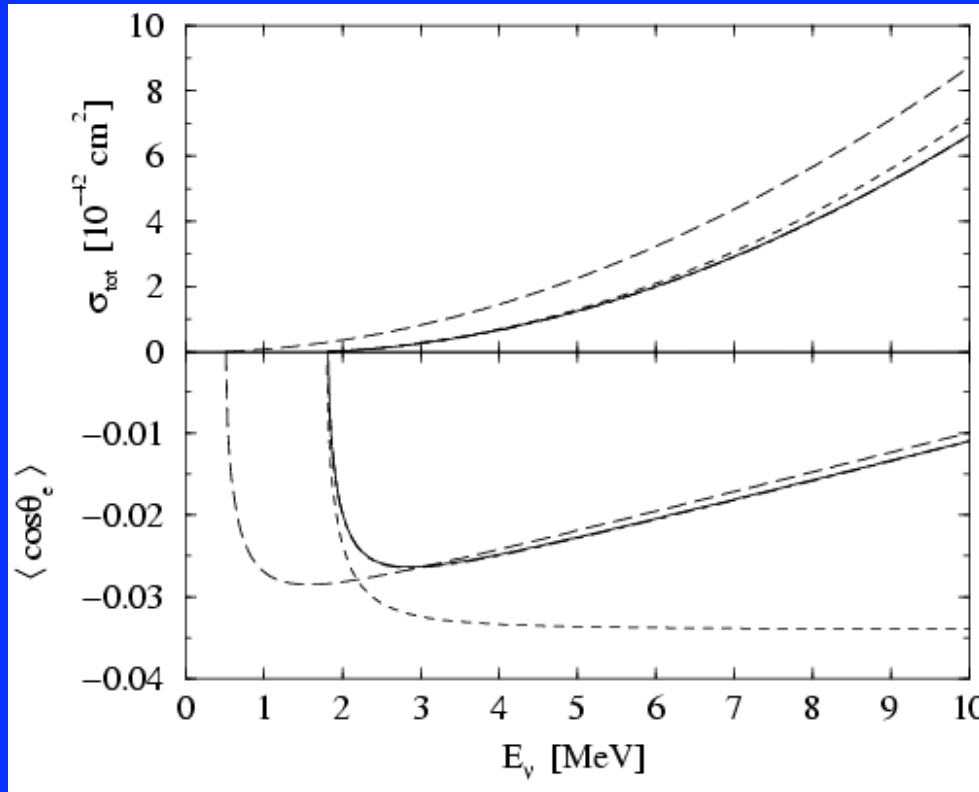


Detection rate  
 $\sim 1$  event per day!

Eguchi et al., PRL 90, 021802 (2003)



# Inverse Beta Cross Section



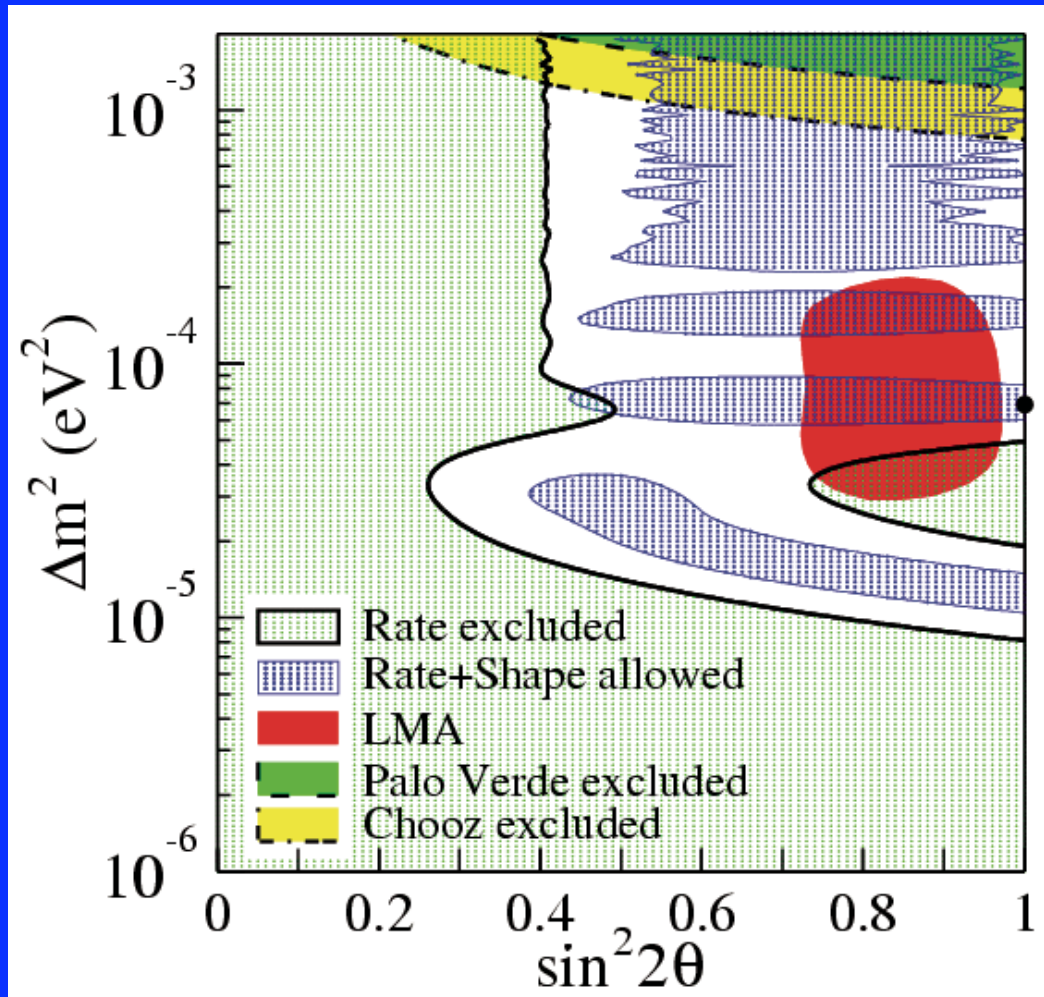
$$\sigma^{(0)} = \frac{2\pi^2 / m_e^5}{f^R \tau_n} E_e^{(0)} p_e^{(0)}$$

$$E_e^{(0)} = E_\nu - (M_n - M_p)$$

Corrections of order  $1/M_p$   
are very important

Vogel and Beacom, PRD 60, 053003 (1999)

# KamLAND Oscillation Results



Very different technique

No solar physics

No MSW effects

$\bar{\nu}_e$ , not  $\nu_e$

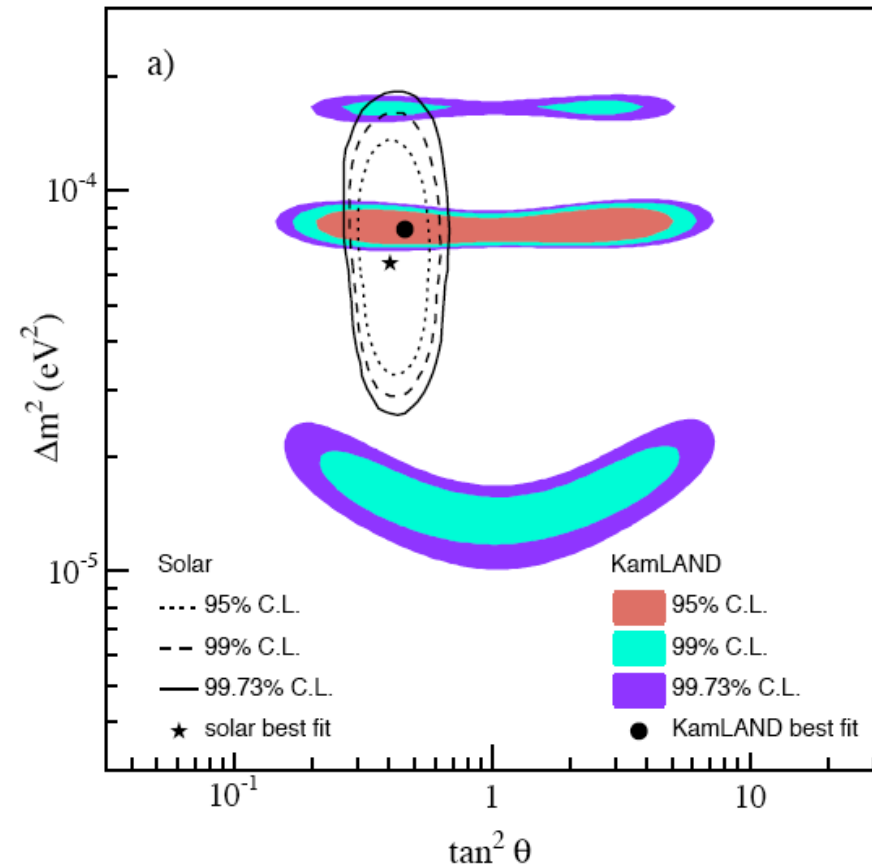
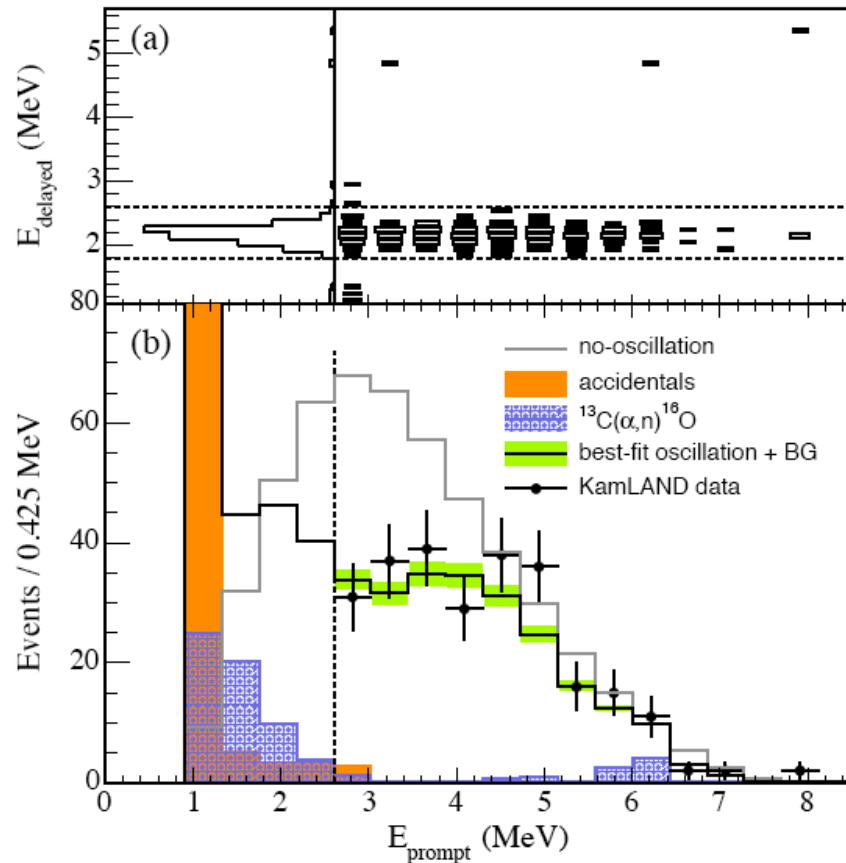
No neutrino decay

No magnetic fields

No matter FCNC

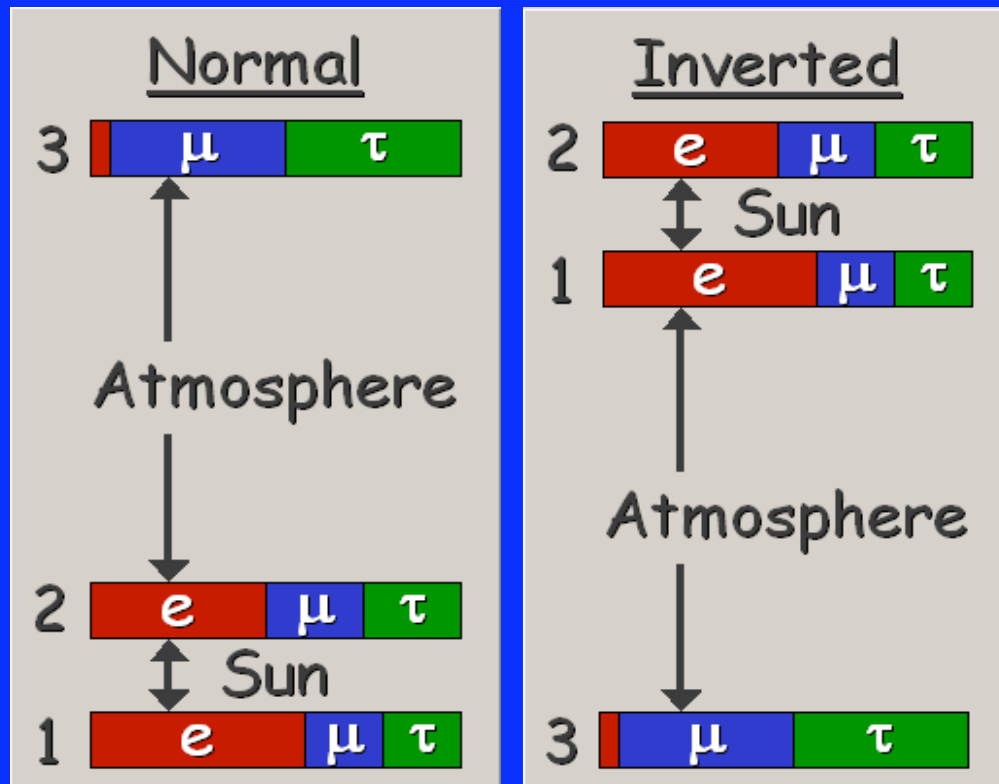
Eguchi et al., PRL 90, 021802 (2003)

# New KamLAND Data



Araki et al., PRL 94, 081801 (2005)

# Neutrino Mixing



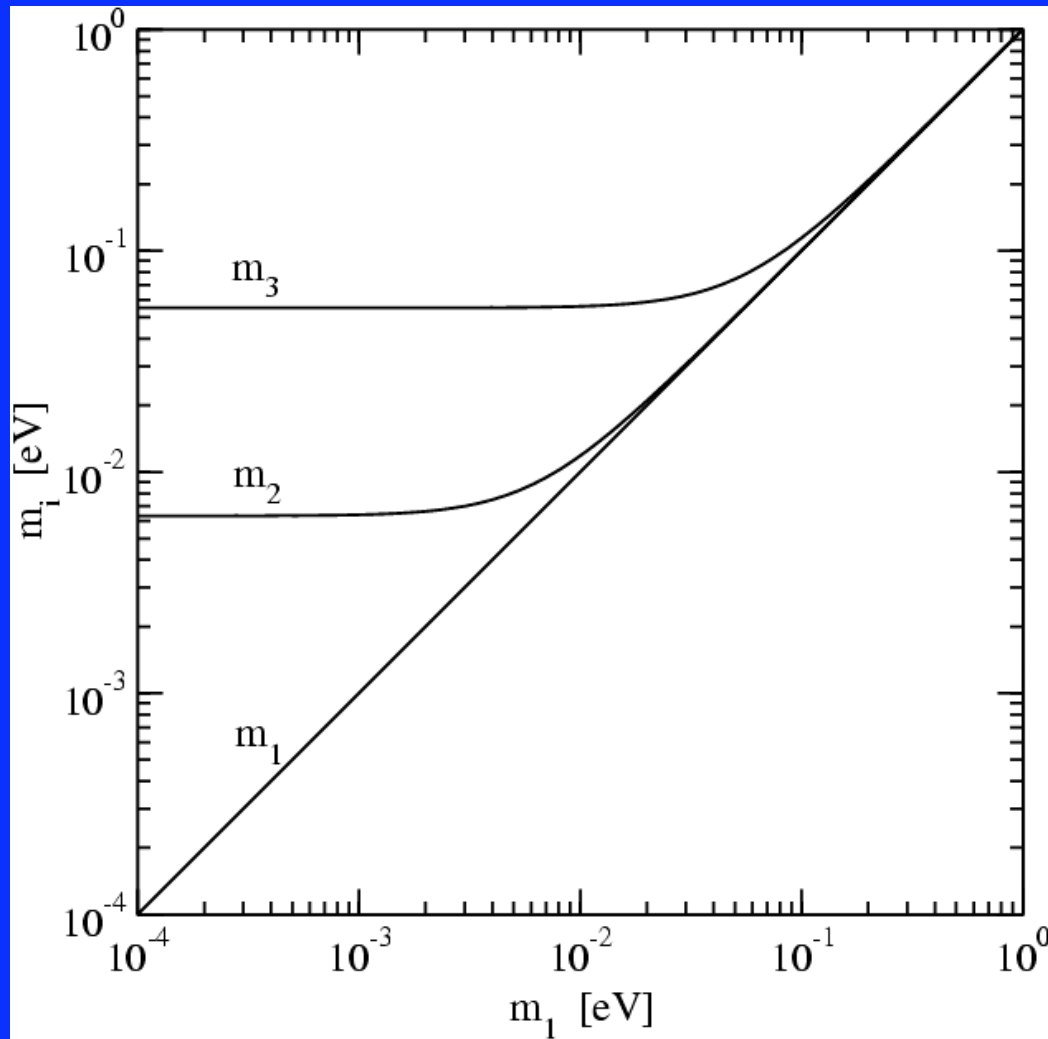
(graphic from Georg Raffelt)

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} U_{\alpha j} \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

$$U \approx \begin{bmatrix} c_\odot & s_\odot & s_{13}e^{-i\delta} \\ -s_\odot/\sqrt{2} & c_\odot/\sqrt{2} & 1/\sqrt{2} \\ s_\odot/\sqrt{2} & -c_\odot/\sqrt{2} & 1/\sqrt{2} \end{bmatrix}$$

$$\theta_{\text{atm}} \approx 45^\circ, \quad \theta_{\text{solar}} \approx 35^\circ, \quad \theta_{13} \leq 10^\circ$$

# Neutrino Mass Splittings



Normal Hierarchy

$$m_1 = m_1$$

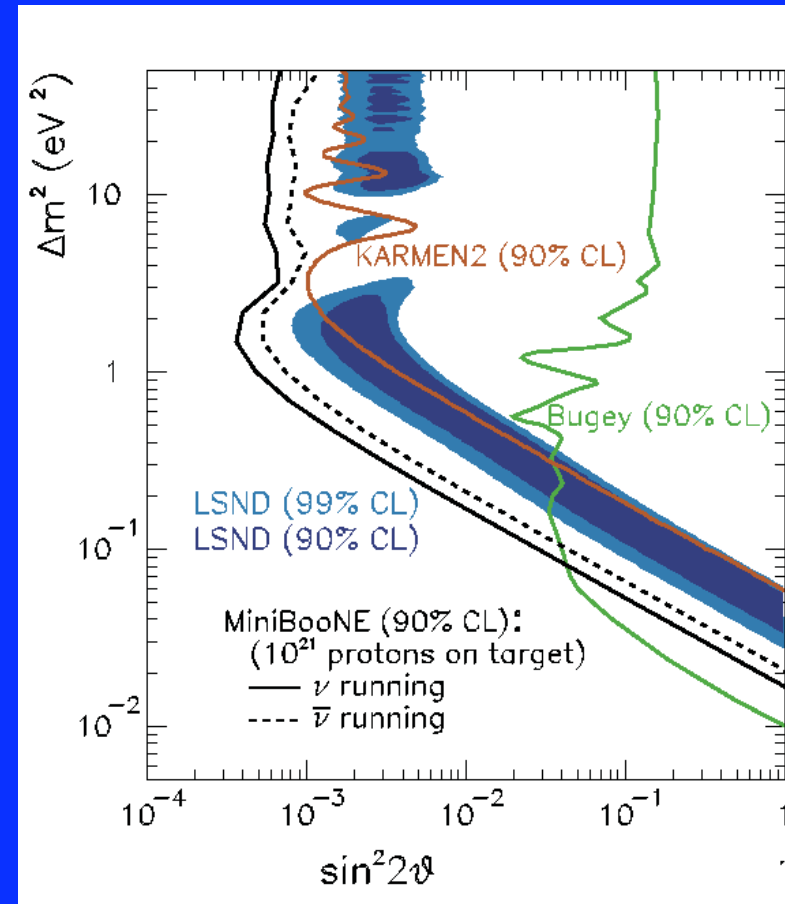
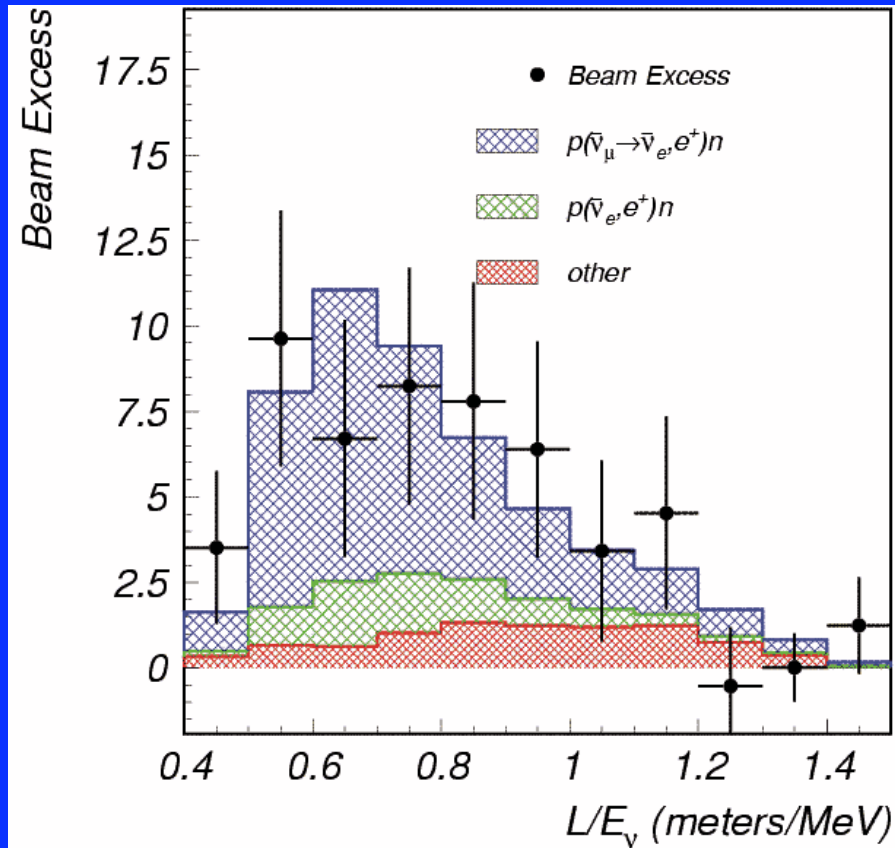
$$m_2 = \sqrt{m_1^2 + \delta m_{\text{solar}}^2}$$

$$m_3 = \sqrt{m_1^2 + \delta m_{\text{solar}}^2 + \delta m_{\text{atm}}^2}$$

$$\frac{m_3}{m_2} \leq \frac{\sqrt{\delta m_{\text{atm}}^2}}{\sqrt{\delta m_{\text{solar}}^2}} \leq 10$$

Beacom and Bell, PRD 65, 113009 (2002)

# What About LSND?



LSND:

$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  appearance signal

MiniBooNE:

$\nu_\mu \rightarrow \nu_e$  appearance search

# Conclusions

- The atmospheric neutrino problem was solved by neutrino oscillations
- And the solar neutrino problem was solved by neutrino oscillations
- Ok, well, just almost sure...
- Essential to complete the mixing parameter measurements and to find the final answers on the above claims (and LSND)!

# Further Reading

- See yesterday's list
- Smirnov, *all of his papers*



- Mass eigenstates:

$$i \frac{\partial}{\partial t} \begin{bmatrix} \psi_1(t) \\ \psi_2(t) \end{bmatrix} = H_{\text{mass}} \begin{bmatrix} \psi_1(t) \\ \psi_2(t) \end{bmatrix}$$

$$H_{\text{mass}} = \begin{bmatrix} E_1 & 0 \\ 0 & E_2 \end{bmatrix}$$

~~Mass~~

$$|v(t)\rangle = \psi_1(t) |v_1\rangle + \psi_2(t) |v_2\rangle$$

fixed, orthonormal kets

$$\begin{aligned} \text{Prob}(m_1) &= |\langle v_1 | v(t) \rangle|^2 \\ &= |\psi_1(t)|^2 = |e^{-iE_1 t}|^2 |\psi_1(0)|^2 \\ &= |\psi_1(0)|^2 \end{aligned}$$

Born in M.E., stay in M.E.

No off-diagonal couplings.

$$E_1 \simeq E + \frac{M_1^2 + M_2^2}{4E} - \frac{\delta m^2}{4E}$$

$$E_2 \simeq E + \frac{M_1^2 + M_2^2}{4E} + \frac{\delta m^2}{4E}$$

$$\delta m^2 = m_2^2 - m_1^2$$

$E \simeq p$ , expand at lowest order  
around common momentum

$$H_{\text{mass}} \simeq \left( E + \frac{M_1^2 + M_2^2}{4E} \right) \mathbb{1} + \frac{\delta m^2}{4E} \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\psi_1(t) \sim e^{+i\delta m^2 t / 4E} \psi_1(0)$$

$$\psi_2(t) \sim e^{-i\delta m^2 t / 4E} \psi_2(0)$$

$$t \sim L/c \sim L$$

Easy, dull.

Neutrino masses basically irrelevant.

• Flavor basis:

Other bases related by unitary transformation, which is just a rotation in 2d.

Rotate amplitudes, with fixed kets.

Useful if states are born and detected in another basis.

~~$$\begin{bmatrix} \psi_e(t) \\ \psi_\mu(t) \end{bmatrix} = U \begin{bmatrix} \psi_1(t) \\ \psi_2(t) \end{bmatrix}$$~~

$$\begin{bmatrix} \psi_1(t) \\ \psi_2(t) \end{bmatrix} = R(\theta) \begin{bmatrix} \psi_e(t) \\ \psi_\mu(t) \end{bmatrix} = \begin{bmatrix} +\cos\theta & -\sin\theta \\ +\sin\theta & -\cos\theta \end{bmatrix} \begin{bmatrix} \psi_e(t) \\ \psi_\mu(t) \end{bmatrix}$$

Physical definition of this basis.

Steps to

$$i \frac{\partial}{\partial t} \begin{bmatrix} \psi_e(t) \\ \psi_\mu(t) \end{bmatrix} = \underbrace{A^\dagger H_{\text{mass}} A}_{H_{\text{fl}}} \begin{bmatrix} \psi_e(t) \\ \psi_\mu(t) \end{bmatrix}$$

Specifically get

$$H_{\mu} = \frac{\delta m^2}{4E} \begin{bmatrix} -\cos 2\theta & +\sin 2\theta \\ +\sin 2\theta & +\cos 2\theta \end{bmatrix} + ( ) \mathbb{1}$$

$$\psi_e(t) = \cos\theta \psi_1(t) + \sin\theta \psi_2(t)$$

↑                      ↑  
put in solutions

$$\psi_e(0) = 1 \quad \text{means} \quad \psi_1(0) = \cos\theta, \quad \psi_2(0) = \sin\theta$$

projection, propagation, deprojection

Grind out  $|\psi_e(t)|^2$  or  $|\psi_\mu(t)|^2$

$$|\psi_e(t)|^2 = 1 - \sin^2 2\theta \sin^2 \left( \frac{\delta m^2 L}{4E} \right)$$

↑ put in units

- Need both mass and mixing
- Explain limits, allowed/disallowed regions
- Importance of  $\frac{1}{2}$  for  $P_{\nu e}$