



The Abdus Salam
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Energy Agency



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SUMMER SCHOOL ON PARTICLE PHYSICS

13 - 24 June 2005

Dark Matter, Dark Energy and New Physics beyond the Standard Model

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ICTP SUMMER SCHOOL
ON
PARTICLE PHYSICS
13-24 June 2005

DM, DE

and

NEW PHYSICS

BEYOND THE SM

Antonio Masiero

Univ. and INFN, PADOVA

OBSERVATIONAL HINTS

FOR NEW PHYSICS

BEYOND THE SM

● HIGH ENERGY FLAVOR PHYS. RARE PHYS

A_{FB}^b

$\sim 3\sigma$

$b \rightarrow s q \bar{q}$

average of
s-penguin decays
 $\sim 3.5\sigma$

$(g-2)_\mu$

$\sim 2.6\sigma$

(using
 e^+e^-)

● > PHYSICS: MASSES and MIXINGS

→ LLHH
 $M \rightsquigarrow$ first access to a new mass scale $M > M_W$ (?)

● "CLASHES.. between

PARTICLE PHYSICS

SM

G-W-S SM

COSMOLOGY

SM

BIG BANG SM

GRAND UNIFICATION

PREDICTION OF $\alpha_s(M_Z)$

Given the precision on $\sin^2 \theta_w$

→ use $\sin^2 \theta_w(M_Z)$ and $\alpha_{ew}(M_Z)$ as INPUT
 ← determine $M_x, \alpha_{GUT}, \alpha_s(M_Z)$

Hall, Nomura

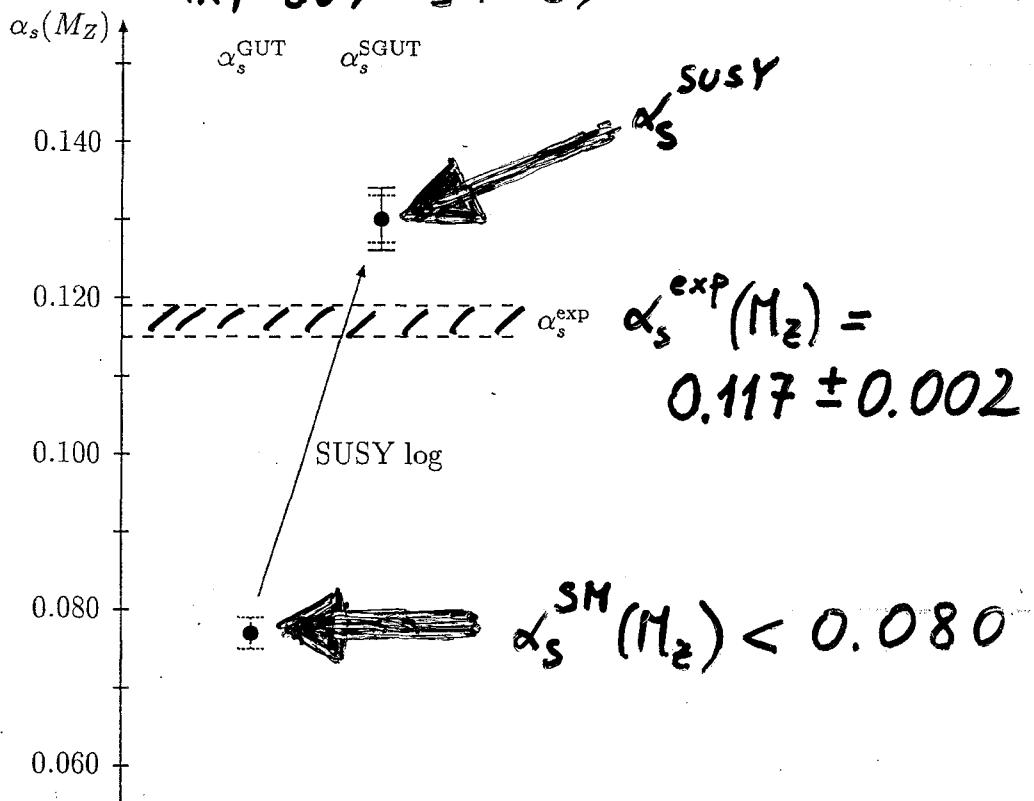


Figure 1: The predictions for $\alpha_s(M_Z)$ in non-supersymmetric grand unification, α_s^{GUT} , and supersymmetric grand unification, α_s^{SGUT} . The solid error bar represents the threshold corrections from the superpartner spectrum. Dotted error bars represent threshold corrections from the unified scale corresponding to a heavy $5 + \bar{5}$ representation with unit logarithmic mass splitting between doublets and triplets.

$$\alpha_s(M_Z) = \frac{\alpha_{ew}(M_Z)}{(1+8x) \sin^2 \theta_w(M_Z) - 3x}$$

$$x \equiv \frac{1}{5} \left(\frac{b_2 - b_3}{b_1 - b_2} \right)$$

MICRO

PARTICLE PHYSICS

GWS STANDARD
MODEL

MACRO

COSMOLOGY

HOT BIG BANG
STANDARD MODEL



HAPPY MARRIAGE

EX. : NUCLEOSYNTHESIS

BUT ALSO

POINTS OF FRICTION

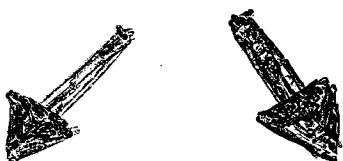
- * * { - COSMIC MATTER-ANTIMATTER ASYMMETRY
- * { - INFLATION
- * ** { - DARK MATTER + DARK ENERGY



"OBSERVATIONAL" EVIDENCE FOR NEW
PHYSICS BEYOND THE (PARTICLE PHYSICS)
STANDARD MODEL

Λ CDM after WMAP

$$\Omega_{\text{MATTER}}^2 h^2 = 0.135 \begin{array}{l} +0.008 \\ -0.009 \end{array}$$



$$\Omega_{\text{BARYON}}^2 h^2 = 0.0224 \pm 0.0009$$

$$\Omega_{\text{CDM}}^2 h^2 = 0.1126 \begin{array}{l} +0.0161 \\ -0.0181 \end{array}$$



consistent with what inferred from earlier observations, but significantly more precise

before : $0.1 < \Omega_{\text{CDM}} h^2 < 0.3$

now (from WMAP) : $0.094 < \Omega_{\text{CDM}} h^2 < 0.129$

Large Scale Structure: the Sloan Digital Sky Survey

New determination of the power spectrum at small

scales from Ly α data

[astro-ph/0405013,07377]

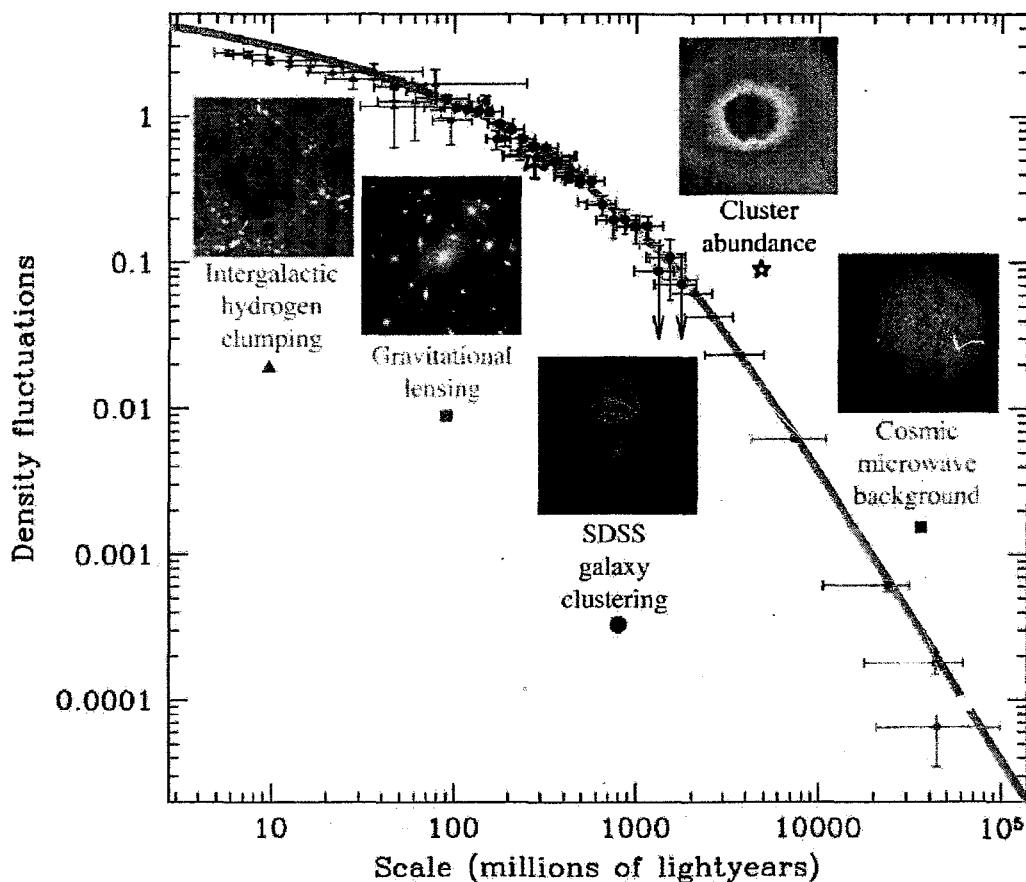


Figure by M. Tegmark

→ better control of systematics

→ longer lever arm in scale



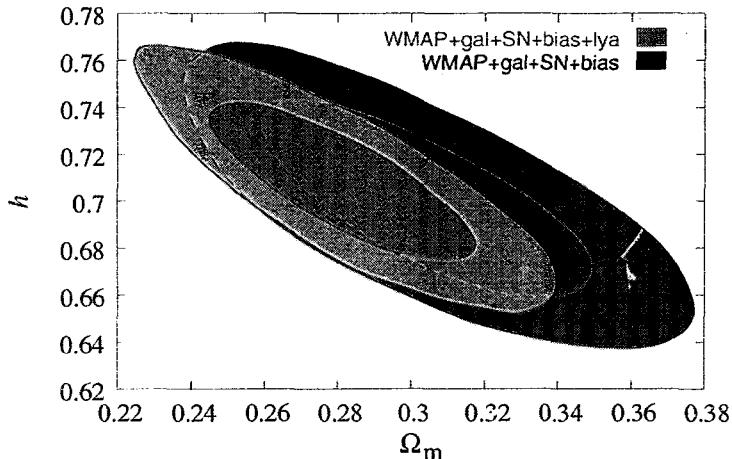
NEW determination of the cosmological parameters
from WMAP data, SDSS galaxy clustering, bias and
Lyman α data, SN Ia data.

L. covi

New analysis by Seljak *et al* [astro-ph/0407372]

What has changed ???

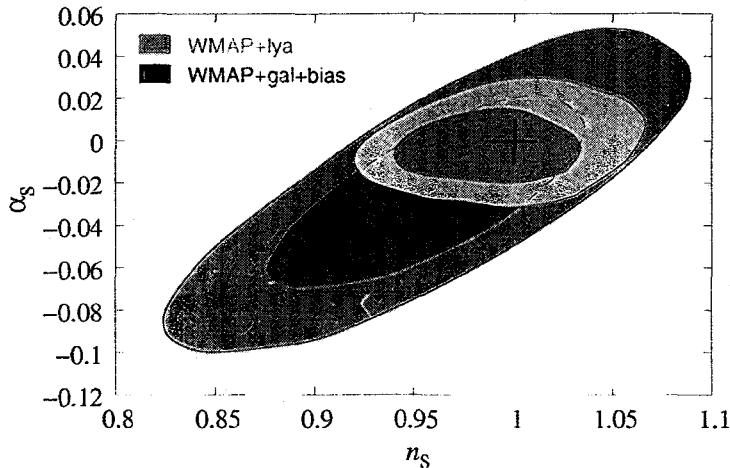
- Improved accuracy on all parameters, e.g.



- More stringent bound on neutrino masses

$$\sum m_\nu \leq 0.66 \text{ eV } (3\nu) \quad \text{or} \quad m_\nu \leq 0.79 \text{ eV } (3 + 1\nu)$$

- Spectral index $n_s = 0.977^{+0.025}_{-0.021}$ and
 $\alpha_s = n'_s = -0.003 \pm 0.010$: NO RUNNING !



Also bound on the tensors as $r \leq 0.45$ at 1σ .

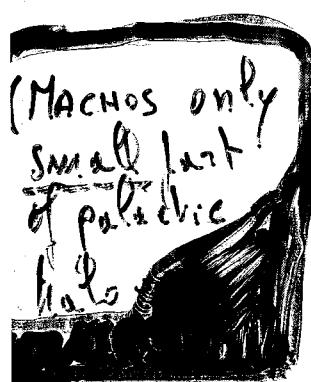
WHAT IS THE DM MADE OF?

Ω_{MATTER} at least 0.2

Ω_Λ cosm. const.

$$\Omega_{\gamma h^2} = 2.48 \times 10^{-5} \quad \text{with } w_y = 0 \quad \Omega_y = \left(\frac{7}{8}\right)^{4/3} \Omega_\gamma$$

SM: relics $\gamma, \nu, \text{ baryons}$ $\Omega_{\text{SM}} < 0.1$



NON-BARYONIC DM

CALLS FOR NEW PHYSICS
BEYOND SM

DM RELIC PARTICLES

THERMAL

ex: $m_\chi \sim 10 \text{ eV}$

$T_{\text{decoupl.}} \sim 10 \text{ GeV}$

HOT

NON-THERMAL

(ex. axions)

STRONG CP PROBLEM

ex: $m_\chi \sim 100 \text{ GeV}$

COLD

$m \ll T_{\text{decoupl.}}$ $T_{\text{decoupl.}} \sim 10 \text{ GeV}$
non relativistic when galaxies
start forming

$m \ll T_{\text{decoupl.}}$

relativistic at

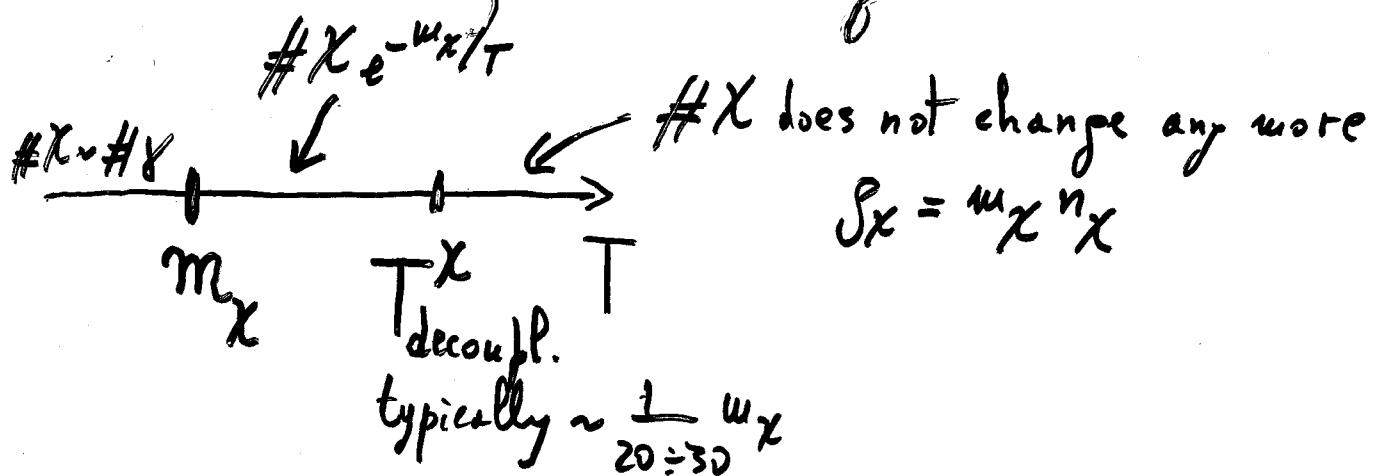
time structure formation (i.e.)

BEST "THERMAL" CDM CANDIDATES



WIMPS

Weakly Interacting Particles



$\Omega_\chi h^2$ depends on particle physics ($\sigma_{\text{annih.}}^\chi$) and
"cosmological" quantities (H, T_0, \dots)

$$\Omega_\chi h^2 \approx \frac{10^{-3}}{\langle \sigma_{\text{annih.}}(\chi\chi) v_\chi \rangle \text{ TeV}^2}$$

$\sim \alpha^2 / M_\chi^2$

from $\sqrt{T^0} M_{\text{Planck}}$

$\Omega_\chi h^2$ in the range $10^{-2} \div 10$ to be cosmologically interesting
(for DM)

$$m_\chi \sim 10^2 \div 10^3 \text{ GeV} \text{ (weak interaction)} \Rightarrow \Omega_\chi h^2 \sim 10^{-2} \div 10 !!$$

NON-BARYONIC

DARK MATTER

HOT \Rightarrow LIGHT NEUTRINOS

COLD \Rightarrow AXIONS

WIMPs



LSP

SUPERHEAVY DM

LIGHTEST K-K

:

BEST WIMP CANDIDATE :

L S P
↓ ↓
LIGHTEST SUSY PARTICLE

SUSY EXTENSION OF THE SM

⇒ MOTIVATION : STABILIZATION OF HIGGS
FOR "LOW MASS AT $\sim M_W$
ENERGY SUSY

SUPPORT : UNIFICATION OF EW + STRONG
GAUGE COUPL.

PROBLEM : PROTON FAST DECAY

REMEDY : IMPOSING (ADDITIONAL)
DISCRETE SYMMETRY (R PARITY)

BONUS : LSP STABLE
↳ WIMP CANDIDATE

LINKING $\Lambda_{\text{NEW PHYSICS}}$ TO THE

ELECTROWEAK BREAKING

$$M_Z = \Lambda_{\text{New Physics}} \times f(\text{phys. param.})$$



MSSM

$$M_Z^2 \sim (90 \text{ GeV})^2 \left(\frac{\langle m_E \rangle}{250 \text{ GeV}} \right)^2 \log \left(\frac{1}{\langle m_E \rangle} \right) + \dots$$

\downarrow
if no fine-tuning with ...

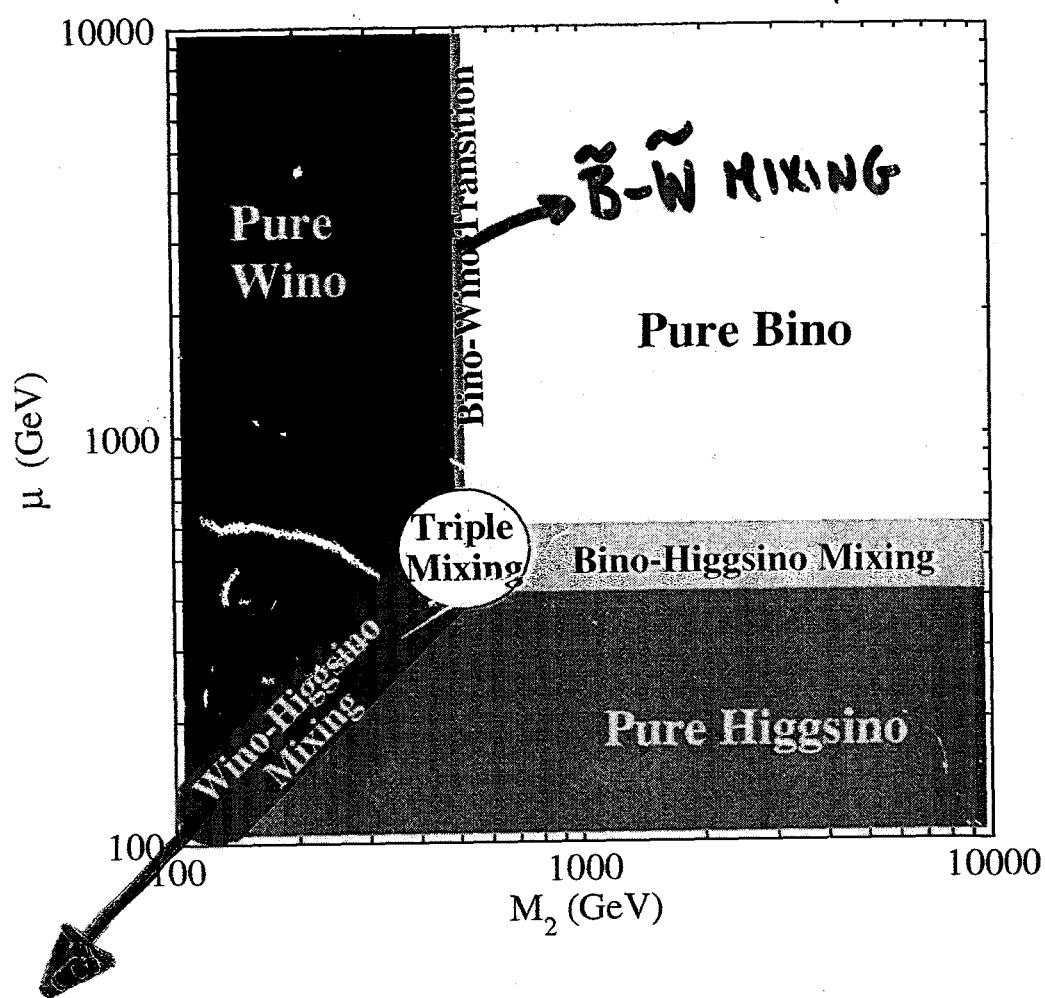
$\Rightarrow \langle m_E \rangle$ upper bounded

but NEED for large $\langle m_E \rangle$ to enhance
tad. corr. to m_{Higgs}

\rightarrow if $m_{\text{Higgs}} \sim 120 \text{ GeV}$ still need
a few percent tuning between
the first term and ... in M_Z^2

NEUTRALINO COMPOSITION

A.M., PROFUMO, ULLIO



$\tilde{W}-\tilde{H}$ MIXING

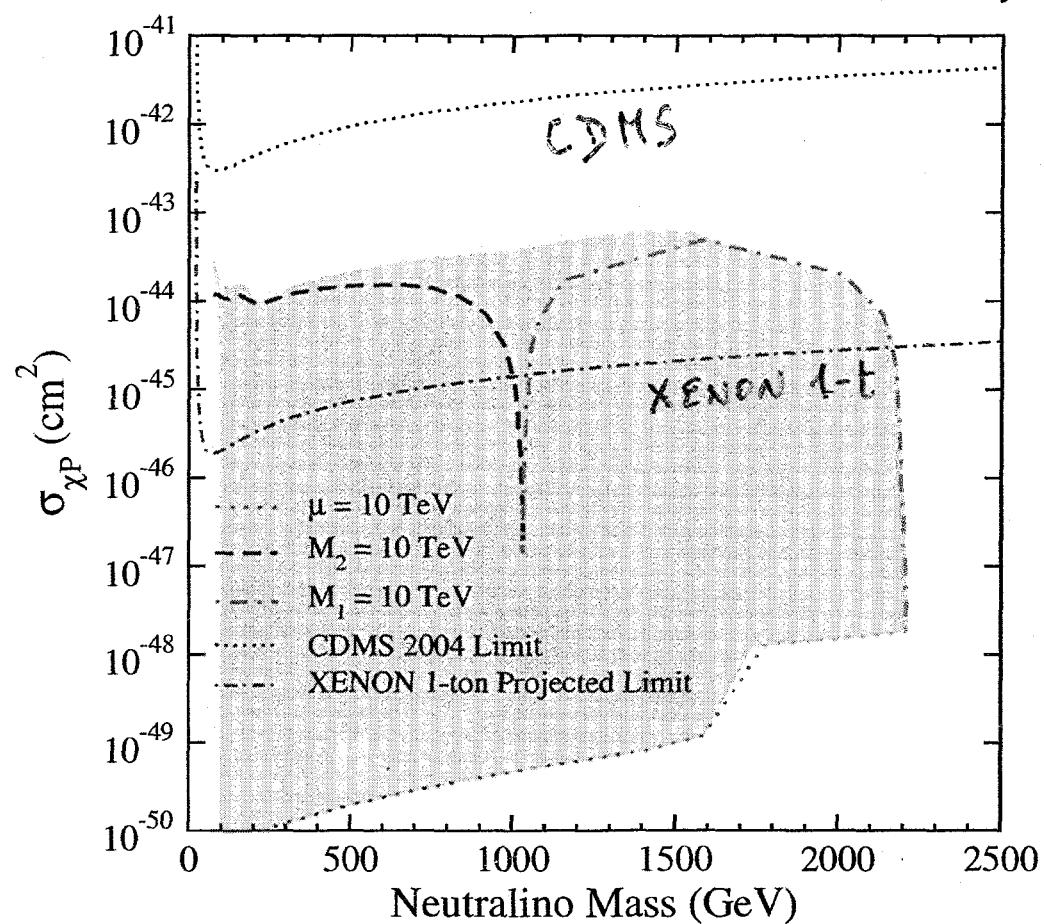
$M_1 = 500 \text{ GeV}$

$$\begin{aligned} \tan \beta &= 50 \\ m_h &= 115 \text{ GeV} \end{aligned}$$

NO GUT RELATION ASSUMED

M_1, M_2, μ ARE CONSIDERED AS
INDEPENDENT PARAMETERS

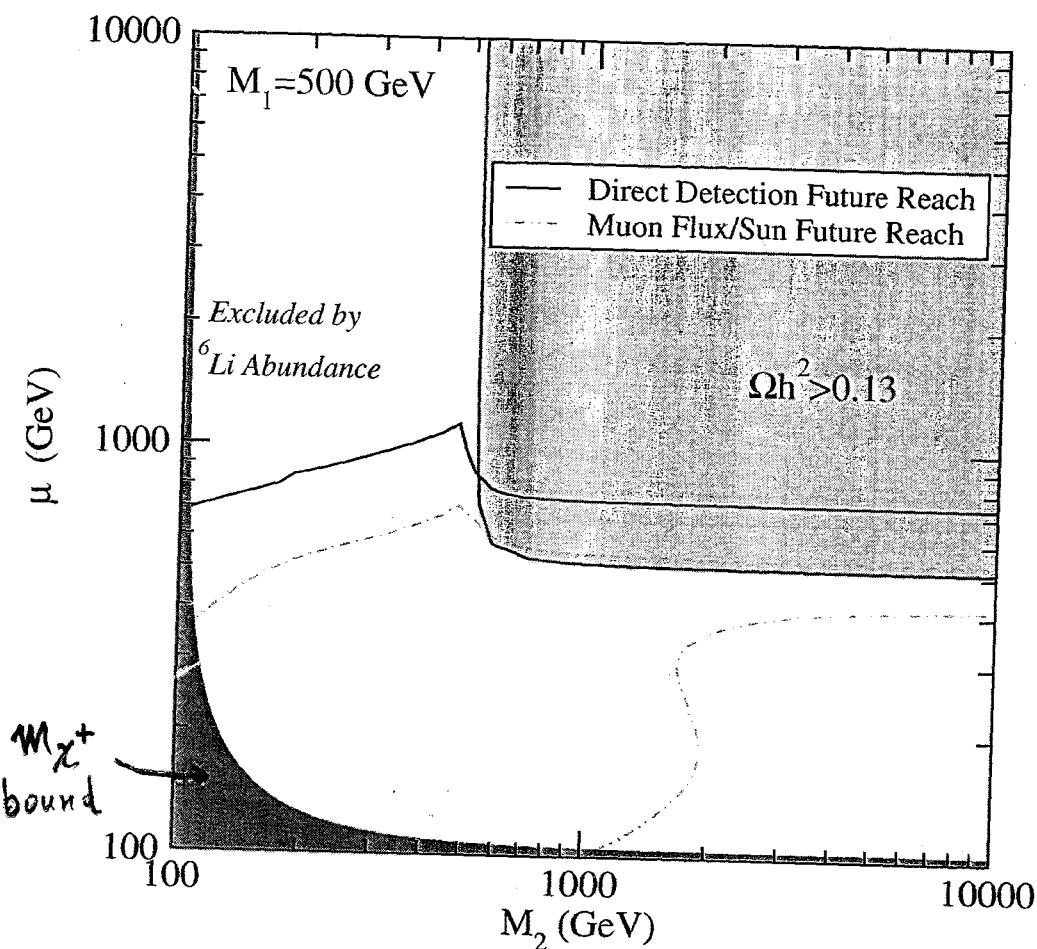
SCANNING FOR M_1, M_2, μ up
to $10\overline{\text{TeV}}$!



A.M., PROFUMO, ULLIO

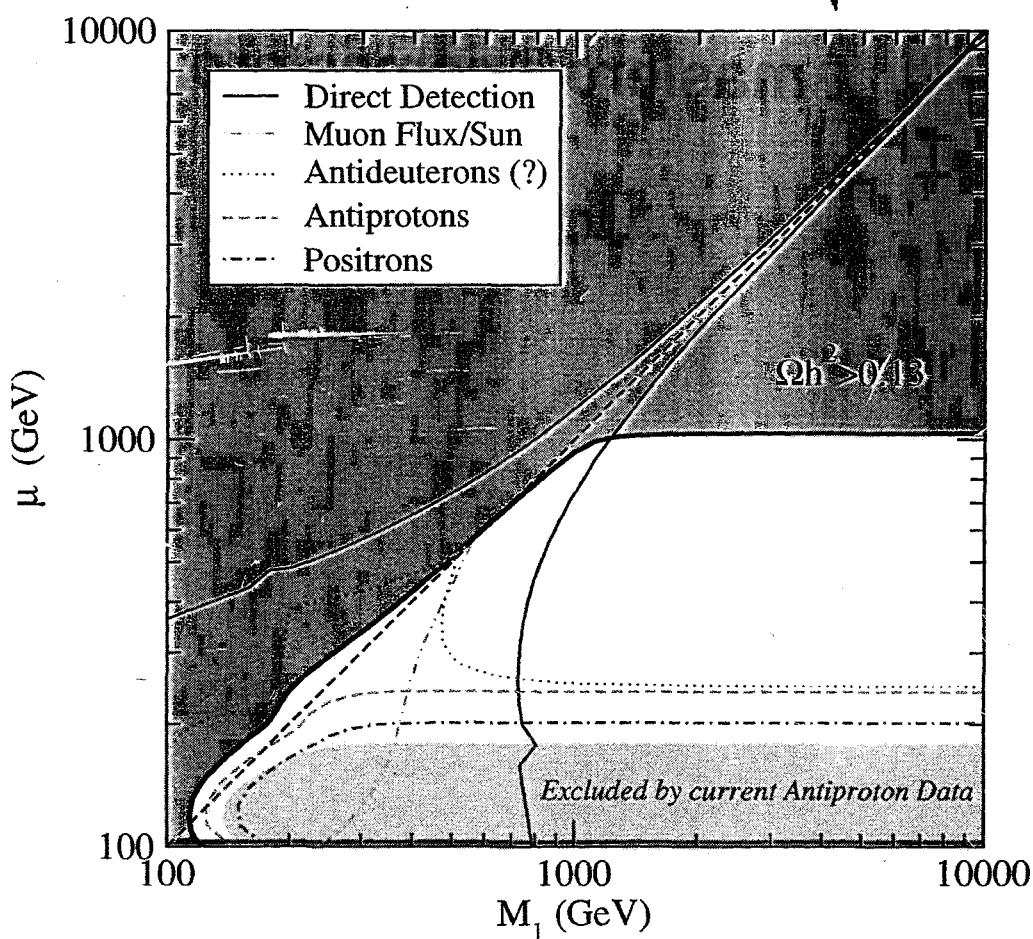
SUSY EXTENSION WITH
HEAVY S-FERMIONS, H, A
NO M_1, M_2, M_3 RELATIONS

A.H., PROFUMO, ULLIO



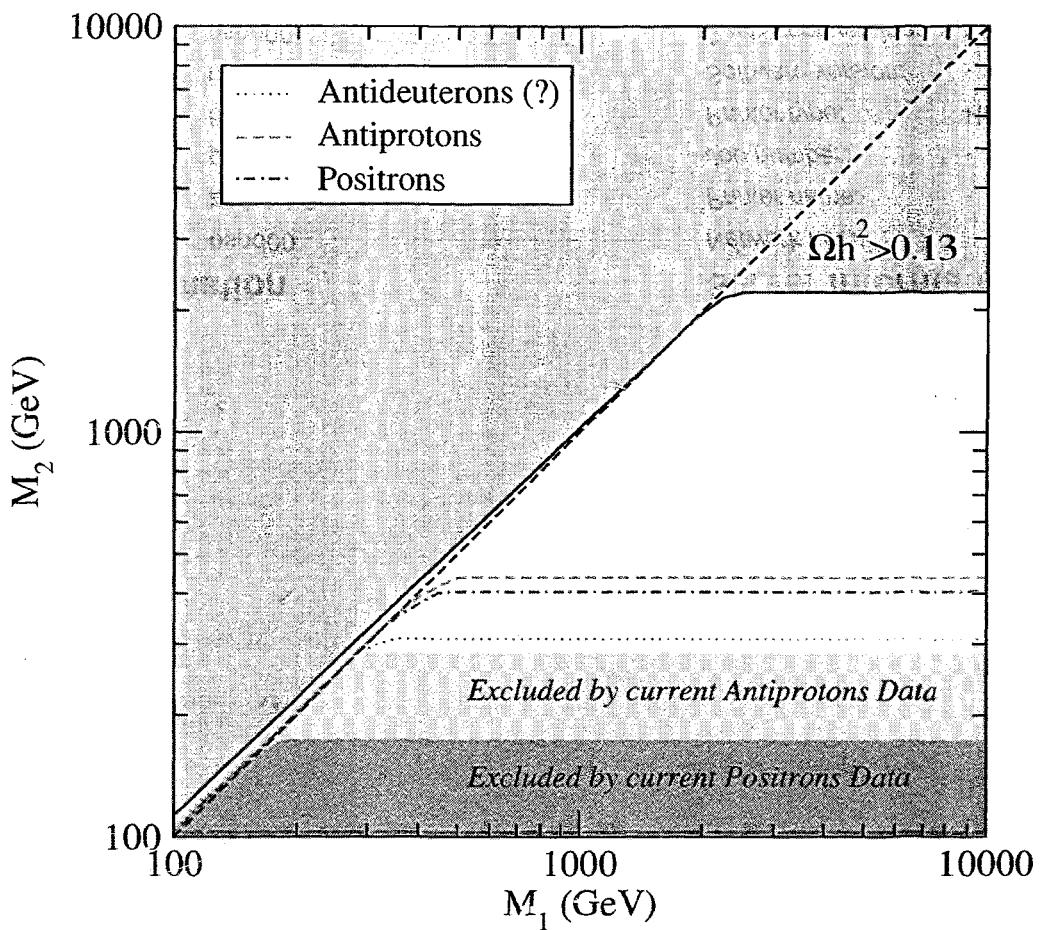
THE REACH IN FUTURE FACILITIES
IN THE DIRECT DETECTION —
AND μ -flux from the Sun ——
CHANNELS ON THE (M_2, μ) PLANE
AT $M_1 = 500$ GeV

A. H., Profumo, Ullio



FUTURE EXCLUSION LIMITS ON
THE (M_1, μ) PLANE AT $M_2 = 10 \text{ TeV}$
Cored Burkert profile assumed

A.M., PROFUMO, ULLIO



FUTURE EXCLUSION LIMITS
ON THE (M_1, M_2) plane at ~~10 TeV~~
 $\mu = 10 \text{ TeV}$
CORED BURKERT PROFILE IS ASSUMED

THE DE PUZZLE

MISINTERPRETATION OF THE DATA ?

1) SNe data 2) CMB $\Omega_T = 1$ 3) measures of f_{DN}

two variables : Ω_{DN}, Ω_{DE} 3 indep. meas.

\Rightarrow cosmic concordance for an ACCELERATING UNIVERSE

"The lengths to which it seems necessary to go in order to avoid concluding that the Universe is accelerating is a strong argument in favor of the concordance model." S. CARROLL

BREAKDOWN OF GENERAL RELATIVITY ?

\rightarrow possibility that gravitation might deviate from conventional GR on scales corresponding to the radius of the entire Universe

Ex. : gravity can be FOUR-DIMENSIONAL below a certain (very large) length scale , but HIGHER-DIMEN
At larger distances \rightarrow universe acceleration at late times

Dvali, Gabadadze, Porrati
Arhanti-Hamed, Dimopoulos, Dvali, Gabadadze
Deffayet, Dvali, Gabadadze
Dvali, Gruzinov, Calderaiga
Lue, Starkman

or 4-dim. modification of GR :

$$\text{ex. } S = \int d^4x \sqrt{|g|} R \rightarrow S = \int d^4x \sqrt{|g|} \left(R - \frac{\kappa^4}{R} \right)$$

Carroll, Duvvuri, Trodden, Turner

"The difficulty in finding a simple extension of GR that does away with the cosmological constant provides yet more support for the standard scenario (ΛCDM)

CARROLL

● ACCELERATION ONLY IN OUR "LOCAL"
RIPPLE ? Kolb, Matarrese, Notari, Riotto

IF : UNIV. IS HOMOGENEOUS, ISOTROPIC AND ACCELERATING

IF : GENERAL RELATIVITY HOLDS



NEED FOR SOME SORT OF DARK ENERGY SOURCE



UNCLUSTERISED → SMOOTHLY-DISTRIBUTED, PERSISTENT
ENERGY DENSITY DOMINATING
THE UNIVERSE ENERGY DENSITY

DARK ENERGY : COSMOLOGICAL
CONSTANT

OR

DYNAMICAL
DARK ENERGY ?

VACUUM ENERGY

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = -8\pi G_N T_{\mu\nu} + \lambda g_{\mu\nu}$$

$[\lambda] = L^{-2}$

if the vacuum energy is non-zero :

$$\langle T_{\mu\nu} \rangle = -\langle \rho \rangle g_{\mu\nu}$$

$$\Rightarrow R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = -8\pi G_N T_{\mu\nu} + \underbrace{8\pi G_N \langle \rho \rangle g_{\mu\nu}}_{\lambda_{\text{eff}} = \frac{\Lambda^4}{M_P^2}}$$

$$\frac{1}{|\lambda_{\text{eff}}|^{1/2}} = \frac{M_P}{\Lambda} \gtrsim \frac{1}{H_0} = 10^{60} l_P \Rightarrow \Lambda \lesssim 10^{-30} M_P$$

$l_P = \sqrt{8\pi G_N} \sim 10^{-32} \text{ cm}$

$$\rightarrow M_P = \sqrt{\frac{1}{8\pi G_N}} \sim 10^{18} \text{ GeV}$$

$$H^2 = \frac{\dot{a}^2}{a^2} = \frac{8\pi G_N}{3} S_{\text{TOT}} \Rightarrow \dot{a}^2 \propto a^2 \rho$$

acceleration (\dot{a} increasing) in an expanding universe \Rightarrow if ρ falls off more slowly than a^{-2}

$$\rho_{\text{MATTER}} \propto a^{-3}, \quad \rho_{\text{RADIATION}} \propto a^{-4}$$

COSM. CONST. - CONST. VACUUM ENERGY

SMOOTHLY-DISTRIBUTED SOURCES OF DARK ENERGY

VARYING SLOWLY WITH TIME

POSSIBILITY OF FINDING A
DYNAMICAL SOLUTION TO

THE COINCIDENCE PROBLEM

\rightarrow "DE TRACKING SOME MATTER
COMPONENT + RECENT TAKEOVER
BY DE (for a wide range of param. of the theory)

HOW TO MAKE VACUUM ENERGY DYNAMICAL



SIMPLEST CASE : EVOLVING SCALAR FIELD
which has not reached its state of minimum energy

Ex. the energy of the true vacuum is zero, but not all fields have evolved to their state of minimum energy \Rightarrow field classically unstable rolling towards its lowest energy state

$$S = \frac{1}{2} \dot{\phi}^2 + V(\phi) \quad P = \frac{1}{2} \dot{\phi}^2 - V(\phi)$$

eq. of motion: $\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$

$$\omega = S/P = \left(\frac{1}{2} \dot{\phi}^2 + V(\phi) \right) / \left(\frac{1}{2} \dot{\phi}^2 - V(\phi) \right)$$

\downarrow
can take any value from +1 to -1

can vary with time

Bronstein (1933) "decaying cosmological constant"
Fried et al. '32; Debye '33; Ritter & Rebbi '58;

Frieman et al.; Turner, White '97;
R. Caldwell, Dave, Steinhardt '98

↳ QUINTESSENCE scalar field

Candidates: pseudo-Goldstone bosons, axions,
scalar fields with a potential

ex: $V(\phi) = e^{1/\phi}$ [decreasing to zero for infinite values
of the field]

$V(\phi) = \frac{1}{\phi^n}$ such a behaviour occurs naturally in
models of dynamical SUSY breaking

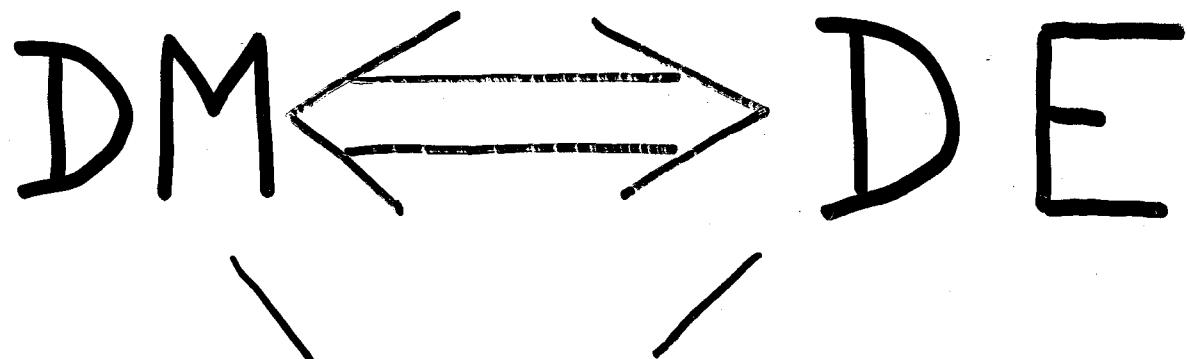
BINETRUY;

A.M., Pietroni, Rosati

⇒ possibility of "tracking" behaviour that makes
the current energy density largely independent
of the initial conditions Zlatev, Wang, Steinhardt

but: no solution to the coincidence problem

⇒ when the scalar field begins to dominate
is still set by tuning parameters of the theory



DO THEY "KNOW" EACH OTHER?

● DIRECT INTERACTION ϕ (quintessence)
WITH DM



DANGER:

ϕ very LIGHT

$$m_\phi \sim H_0^{-1} \sim 10^{-33} \text{ eV}$$

\Rightarrow threat of violation of the equivalence principle
constancy of the fundamental "constants", ...

● INFLUENCE of ϕ ON THE NATURE
AND THE ABUNDANCE OF CDM

\rightarrow modifications of the standard
picture of WIMPS FREEZE-OUT

CDM CANDIDATES

KINATION

DOMINATION BY THE KINETIC ENERGY
 OF THE QUINTESSENCE FIELD ϕ AT
 EPOCHS BEFORE BBN, IN PARTICULAR
 AT THE TIME WIMPS $\xrightarrow{\text{FREEZE OUT}}$
 JORCE; JOYCE, PROKOPEC; FERREIRA, JORCE \rightarrow SALATI

$$\rho_\phi = T_0^\circ = \frac{\dot{\phi}^2}{2} + V(\phi)$$

\hookrightarrow assumption: for some time $\dot{\phi}^2/2$
 dominates

$$\rho_\phi = \dot{\phi}/2 \propto \bar{a}^6 \quad \rho_{\text{rad}} \propto T^4 \sim \bar{a}^{-4} \quad \hookrightarrow \text{when WIMPS decouple}$$

$$\gamma_\phi = \frac{\rho_\phi}{\rho_\gamma} \quad \text{def: } \gamma_\phi^\circ = \frac{\rho_\phi^\circ}{\rho_\gamma^\circ} = \left. \frac{\rho_\phi}{\rho_\gamma} \right|_{T_{BBN}}$$

\Rightarrow the expansion rate H of the Universe increases
 by a factor $\sqrt{\gamma_\phi} T/T_{BBN}$ with respect to
 conventional radiation dominated cosmology

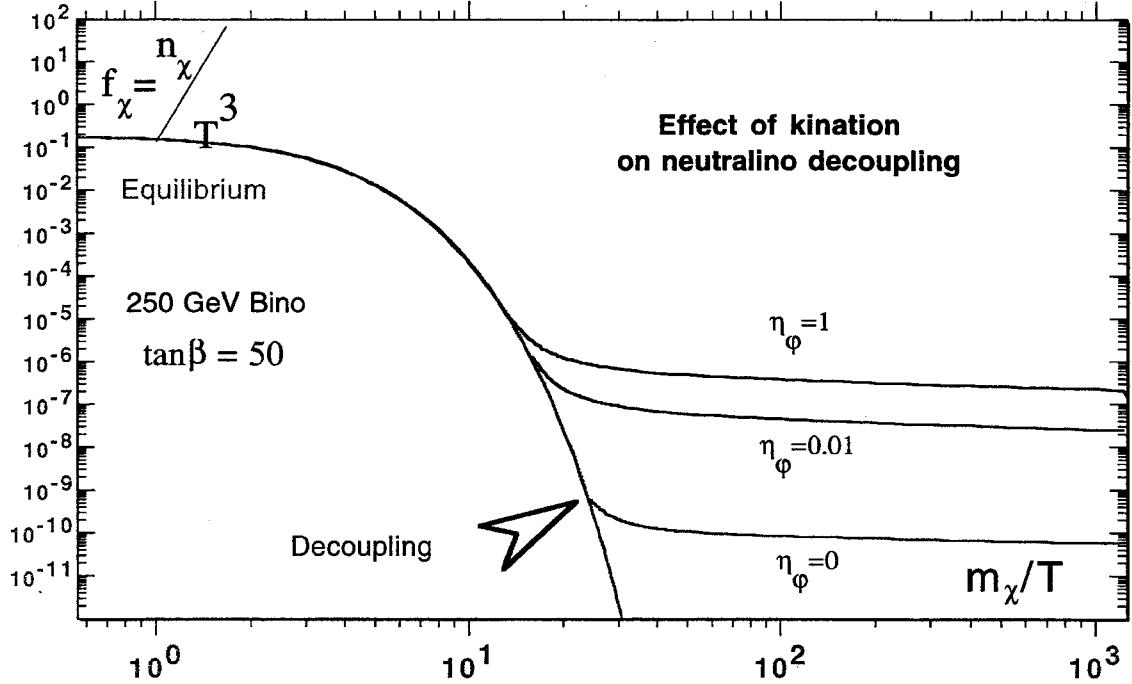


FIG. 1. Neutralino codensity as a function of the mass-to-temperature ratio $y = m_\chi/T$ for three different values of the kination parameter. For $\eta_\phi = 0$, we recover the standard radiation dominated cosmology whereas for $\eta_\phi = 0.01$ and $\eta_\phi = 1$, the expansion rate H is significantly increased. This leads to an earlier decoupling and to a much larger asymptotic value for the neutralino codensity.

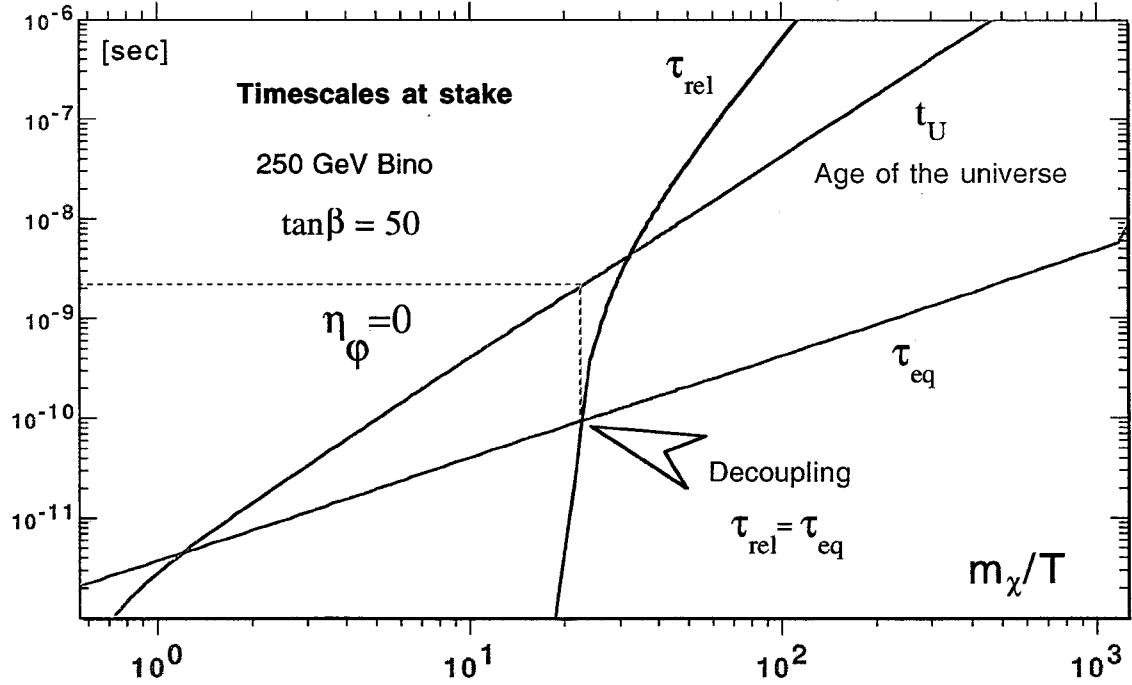
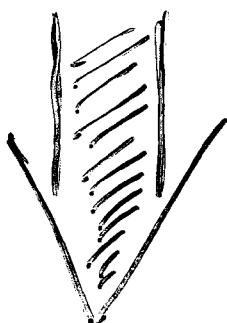


FIG. 2. The age of the universe t_U as well as the typical time scales τ_{rel} and τ_{eq} are featured as a function of the mass-to-temperature ratio $y = m_\chi/T$. The freeze-out occurs at $y_F = 22.7$ when τ_{rel} overcomes τ_{eq} . The standard radiation dominated cosmology is assumed here with $\eta_\phi = 0$ so that t_U evolves like y^2 .

"HARMLESS" QUINTESSENCE



NO EQUIVALENCE PRINCIPLE
VARYING COUPLINGS ... PROBLEMS



SCALAR-TENSOR (ST)

GRAVITY THEORY

⇒ MATTER HAS A PURELY
METRIC COUPLING
WITH GRAVITY

JORDAN ; FIERZ ;

BRANS - DICKE

DOUBLE ATTRACTOR

IN ST GRAVITY

WITH QUINTESSENCE
BARTOLO-PIETRONI

large class of ST gravities

have an attractor mechanism

towards General Relativity (GR)

\Rightarrow expansion of the Universe

during matter-domination

tends to drive the scalar

fields towards a state where

the theory becomes indistinguishable

from GR

2 attraction mechanisms {
 \rightarrow TRACKER SOLUTION
 \rightarrow accelerating Univ
 \rightarrow Towards GR \rightarrow ultra light scalars safe}

$$S = S_{\text{grav}} + S_{\text{matter}}$$

ST theories of gravity

Damour, Nordtvedt;
 Damour, Polyakov;
 Santiago, Kellis, Wagoner
 Damour, Pichon

$$S_{\text{grav}} = \frac{1}{16\pi} \int d^4x \sqrt{-\tilde{g}} \left[\phi^2 \tilde{R} + 4\omega(\phi) \tilde{g}^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - 2\tilde{V}(\phi) \right]$$

$$S_{\text{matter}} = S_{\text{matter}} [\psi_m, \tilde{g}_{\mu\nu}]$$

ψ_m coupled only to the metric tensor
 not to ϕ

\tilde{R} = Ricci scalar constructed from the physical metric $\tilde{g}_{\mu\nu}$

each ST model identified by $\begin{cases} \omega(\phi) \\ \tilde{V}(\phi) \end{cases}$

Jordan-Fierz-Brans-Dicke $\begin{cases} \omega(\phi) = \omega \text{ const.} \\ \tilde{V}(\phi) = 0 \end{cases}$

$\tilde{T}^{\mu\nu} = \frac{2}{\sqrt{-\tilde{g}}} \frac{\delta S_m}{\delta \tilde{g}_{\mu\nu}}$ is conserved, masses and
 non-gravit. couplings are
 $\tilde{g}_{\mu\nu}, \phi$ "physical" \rightarrow time independent, non-grav.
 Jordan variables physics laws take usual form

but eqs. of motion are cumbersome in Jordan frame as they mix spin-2 and spin-0

$$\text{Einstein frame} \quad \left\{ \begin{array}{l} \tilde{g}_{\mu\nu} = A^2(\varphi) g_{\mu\nu} \\ \phi^2 = A^{-2}(\varphi) G_*^{-1} \\ \tilde{V}(\phi) = A^{-4}(\varphi) V(\varphi) \end{array} \right.$$

$$\alpha(\varphi) \equiv \frac{d}{d\varphi} \log A(\varphi)$$

$$\text{taking } \alpha^2(\varphi) = \frac{1}{2w(\phi)+3}$$

$$S_{\text{grav}}^{\text{"Einstein frame"}} = \frac{1}{16\pi G_*} \int d^4x \sqrt{-g} \left[R - 2g^{\mu\nu} \partial_\mu \varphi \partial_\nu \varphi - 2V(\varphi) \right]$$

but now S_{matter} contains also the scalar field

$$S_{\text{matter}}^{\text{Einstein frame}} \quad [\psi_m, A^2(\varphi) g_{\mu\nu}]$$

\downarrow
 masses and non-grav. coupl. const. are field-dependent; G_* time-independent and field eqs. have a simple form

$$\left\{ \begin{array}{l} R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{T_{\mu\nu}^\phi}{M_*^2} + \frac{T_{\mu\nu}}{M_*^2} \\ \partial_\phi^2 + \frac{1}{M_*^2} \frac{\partial V}{\partial \phi} = - \frac{1}{M_*^2} \frac{\alpha(\phi)}{\sqrt{2}} T \end{array} \right.$$

\Downarrow
 $T = g^{\mu\nu} T_{\mu\nu}$

$T_{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta S_m}{\delta g^{\mu\nu}}$ matter-energy momentum tensor
in the Einstein frame

$$T_{\mu\nu}^\phi = M_*^2 \partial_\mu \phi \partial_\nu \phi - g_{\mu\nu} \left[M_*^2 \frac{g^{\rho\sigma}}{2} \partial_\rho \phi \partial_\sigma \phi - V(\phi) \right]$$

when $\alpha(\phi) \rightarrow 0$ ST \rightarrow GR

\Rightarrow MODIFICATION OF THE HUBBLE

PARAMETER $\tilde{H} = \frac{d}{d\tilde{\tau}} \log \tilde{a}$

$$d\tilde{\tau} = A(\phi) d\tau \quad \tilde{a} = A(\phi) \frac{d\tilde{\tau}}{d\tau} a$$

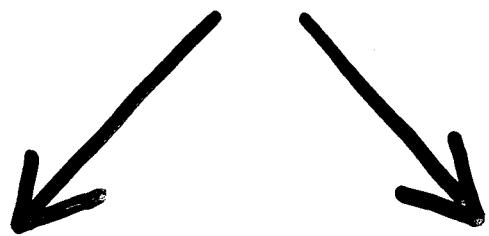
DAHOUR-
PICHON

$$\tilde{H}^2 = \frac{H^2}{A^2(\phi)} \left[1 + \alpha(\phi) \phi' \right]^2$$

$$H = \frac{d}{d\tau} \log a; \quad \tilde{H} = \frac{d}{d\tilde{\tau}} \log \tilde{a}$$

$$\phi' = \frac{d\phi}{d\tau} \quad \tau = \log a$$

IMPLEMENTATION OF THE DOUBLE ATTRACTOR MECHANISM



$$V(\phi) = \frac{\Lambda^4}{\phi^\delta}$$

$\delta > 0 \Rightarrow$ "TRACKER..
SOLUTIONS

(attractors in
field space)

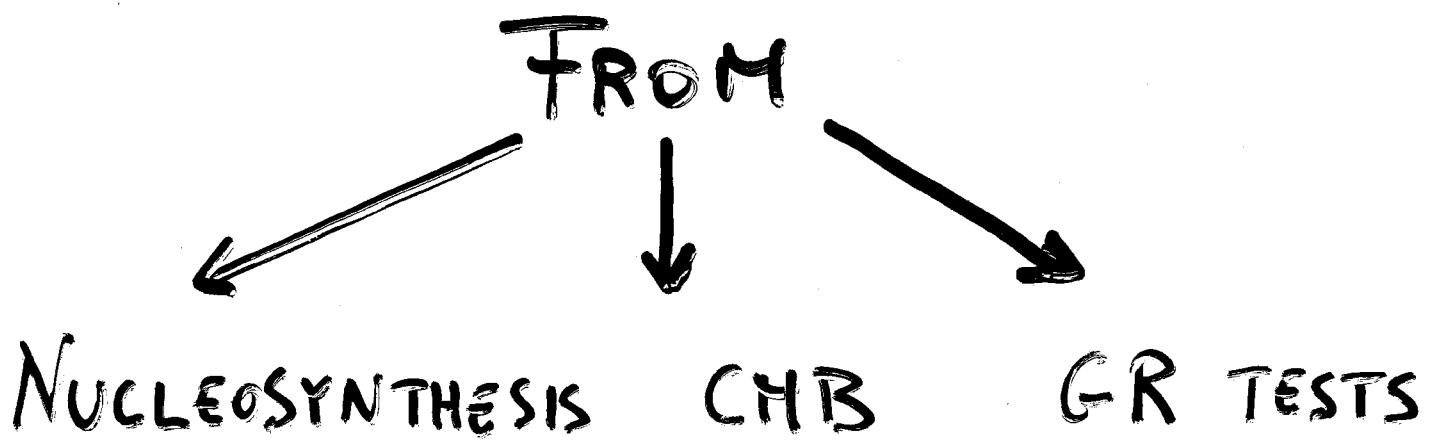
$\alpha(\phi)$ run-away
behavior with
positive slope

$$\text{Ex. : } A(\phi) = 1 + Be^{-\beta\phi}$$

$$\alpha(\phi) = \frac{\beta Be^{-\beta\phi}}{1 + Be^{-\beta\phi}}$$

$$\phi \rightarrow \infty \quad \alpha(\phi) \rightarrow 0 \\ \text{ST} \rightarrow \text{GR}$$

PHENOMENOLOGICAL BOUNDS



① Bound on $A(\phi)$ from BBN

$$\tilde{H}_{BBN}^2 \approx A^2(\phi) \frac{1}{3M_*^2} \tilde{\rho} \Big|_{BBN}$$

$$\left| \frac{\Delta \tilde{H}^2}{\tilde{H}^2} \right|_{BBN} = \left| \frac{\tilde{H} - \tilde{H}_{GR}^2}{\tilde{H}_{GR}^2} \right|_{BBN} = \frac{A^2(\phi_{BBN}) - A^2(\phi_0)}{A^2(\phi_0)}$$

$$\Rightarrow \frac{A(\phi_{BBN})}{A(\phi_0)} < 1.08$$

(using the bound $\Delta N_s < 1$ for BBN)



GENERAL RELATIVITY TESTS

At the post-newtonian level \rightarrow deviations from GR may be parametrized in terms of an effective field-dependent newtonian constant

$$G = G(\phi) \equiv G_* A(\phi)^2 (1 + \alpha^2(\phi))$$

+ 2 dimensionless parameters γ_{PN} β_{PN}

$$\gamma_{PN} - 1 = -2 \frac{\alpha^2}{1 + \alpha^2}; \quad \beta_{PN} - 1 = \frac{k \alpha^2}{(1 + \alpha^2)^2}$$

$$k \equiv \partial \alpha / \partial \phi$$

RECENT SEVERE BOUND ON $\gamma_{PN} - 1$

FROM CASSINI MISSION

$$\gamma_{PN} - 1 = (2.1 \pm 2.3) \times 10^{-5} \quad \text{→}$$

$\nearrow \beta$ large

$\searrow \beta$ very small $\rightarrow \alpha \rightarrow 0 \quad \text{ST} \rightarrow \text{GR}$



BOUNDS FROM THE CMB

POWER SPECTRUM

main change w.r.t. theory of DE based on GR is due to a difference in the expansion rate affecting the angular scale of anisotropies

Riazuelo, Uzan;
Chen, Kamionkowski;
Perrotta, Baccigalupi,
Matarrese

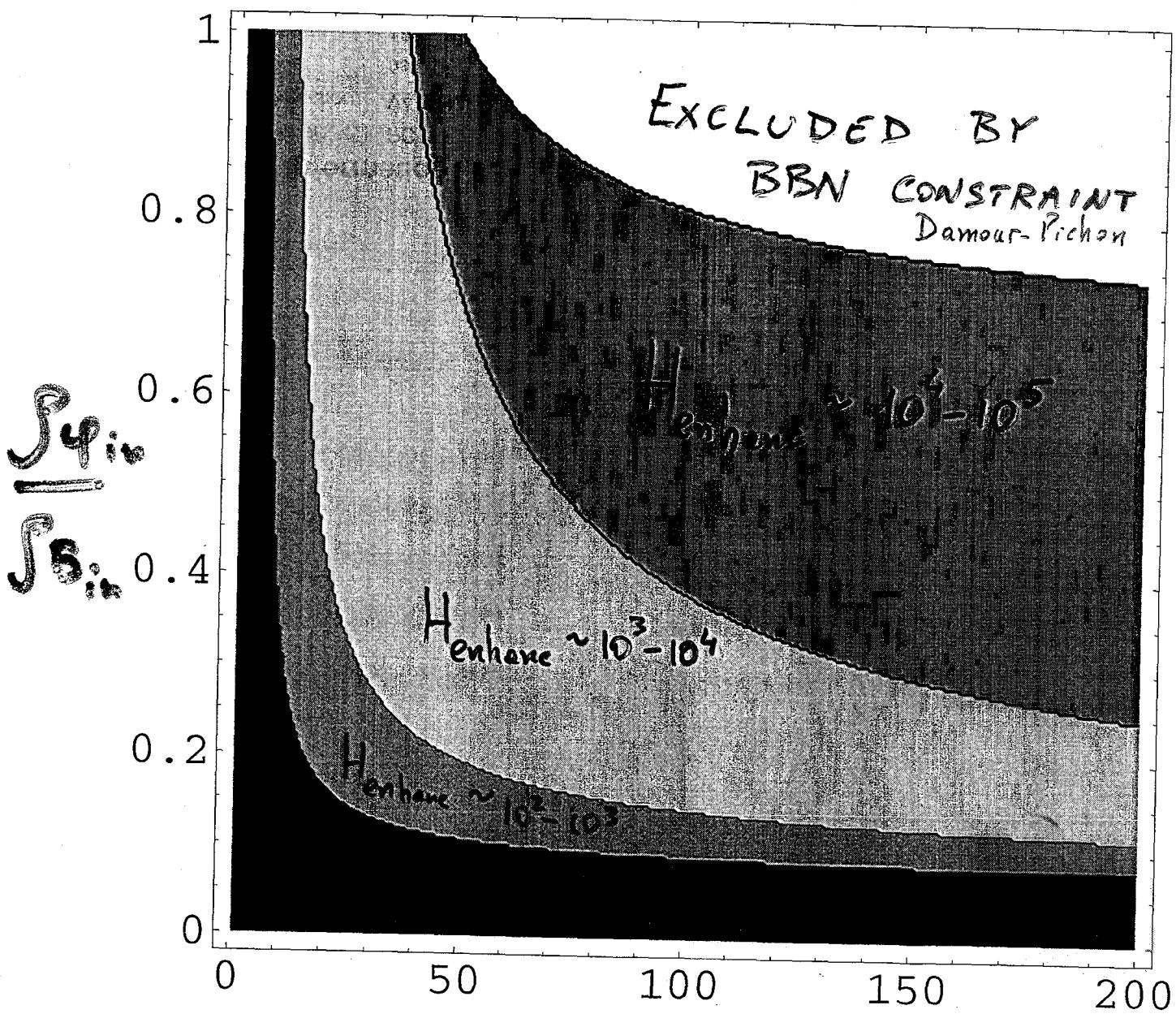
$$\frac{\Delta \ell_{\text{peak}}}{\ell_{\text{peak}}} \approx \frac{4}{3} \frac{A(\tilde{a}_{\text{dec}}) - 1}{A(\tilde{a}_{\text{dec}})}$$

⇒ once the BBN bound on A has been imposed $\rightarrow A(\tilde{a}_{\text{dec}})$ is so close to 1 that shifts in the peak multiples are smaller than the exp. error

CMB spectrum does NOT provide significant bounds on the scenarios we study here

ENHANCEMENT OF H at $T \approx 10$ GeV
 in ST theories of Gravity $\zeta \approx T_{\text{freeze-out}}^x$

$$\beta = 6$$

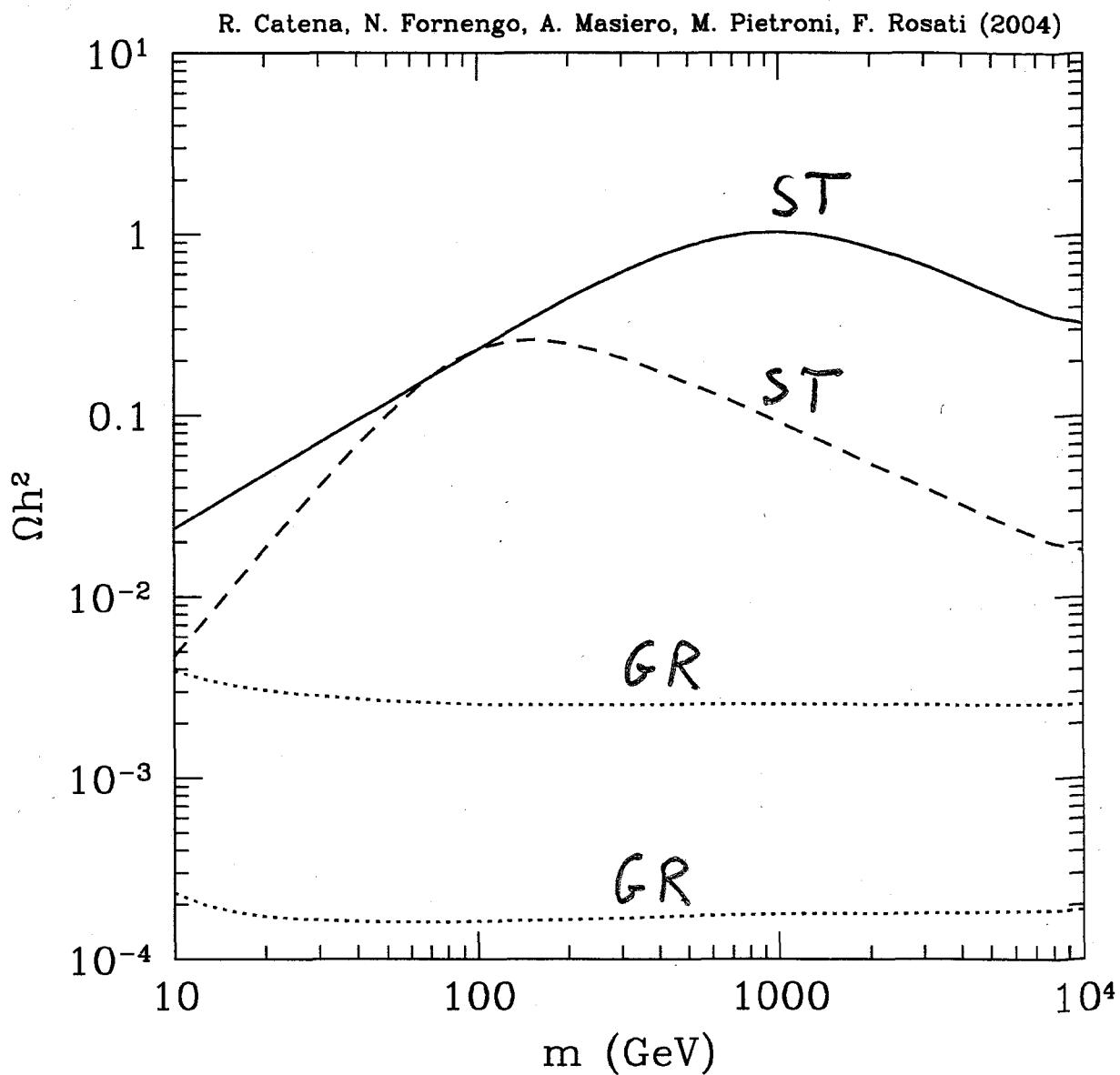


FORNENGO, CATENA, A.M., PIETRONI, ROSATI

EXAMPLE: χ relic abundance in the cases

$$\langle \delta_{\text{ann}} \nu \rangle \equiv a = 10^{-7} \text{ GeV} \quad ---$$

$$\langle \delta_{\text{ann}} \nu \rangle \equiv \frac{b}{\chi} = 10^{-4} \text{ GeV}/\chi \quad ---$$



→ $\Omega_{\text{CDM}}^{\chi} h^2 |$ too small $< 10^{-2}$

jumps up to $\Omega_{\text{CDM}}^{\chi} h^2 |_{\text{ST}} \approx 0.1$ in ST

SHARING THE UNIVERSE ENERGY BUDGET

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} \rho - \frac{k}{a^2}$$

$$\text{FLAT UNIV} \rightarrow k=0 \quad \Omega_i \equiv \rho_i/\rho_c \quad P = P_c = \frac{3H^2}{8\pi G} = 1.88 \cdot 10^{-29} h^2 \text{ g cm}^{-3}$$

$$= (2.73 \pm 0.36) \times 10^{11} \text{ eV}^4$$

5 SHARE HOLDERS :

BARYONS, PHOTONS, NEUTRINOS + DARK MATTER + DARK ENERGY

$$\left(\Omega_x \sim 5 \times 10^{-5} \right)$$

$$\Omega_B \sim 5 \times 10^{-2}$$

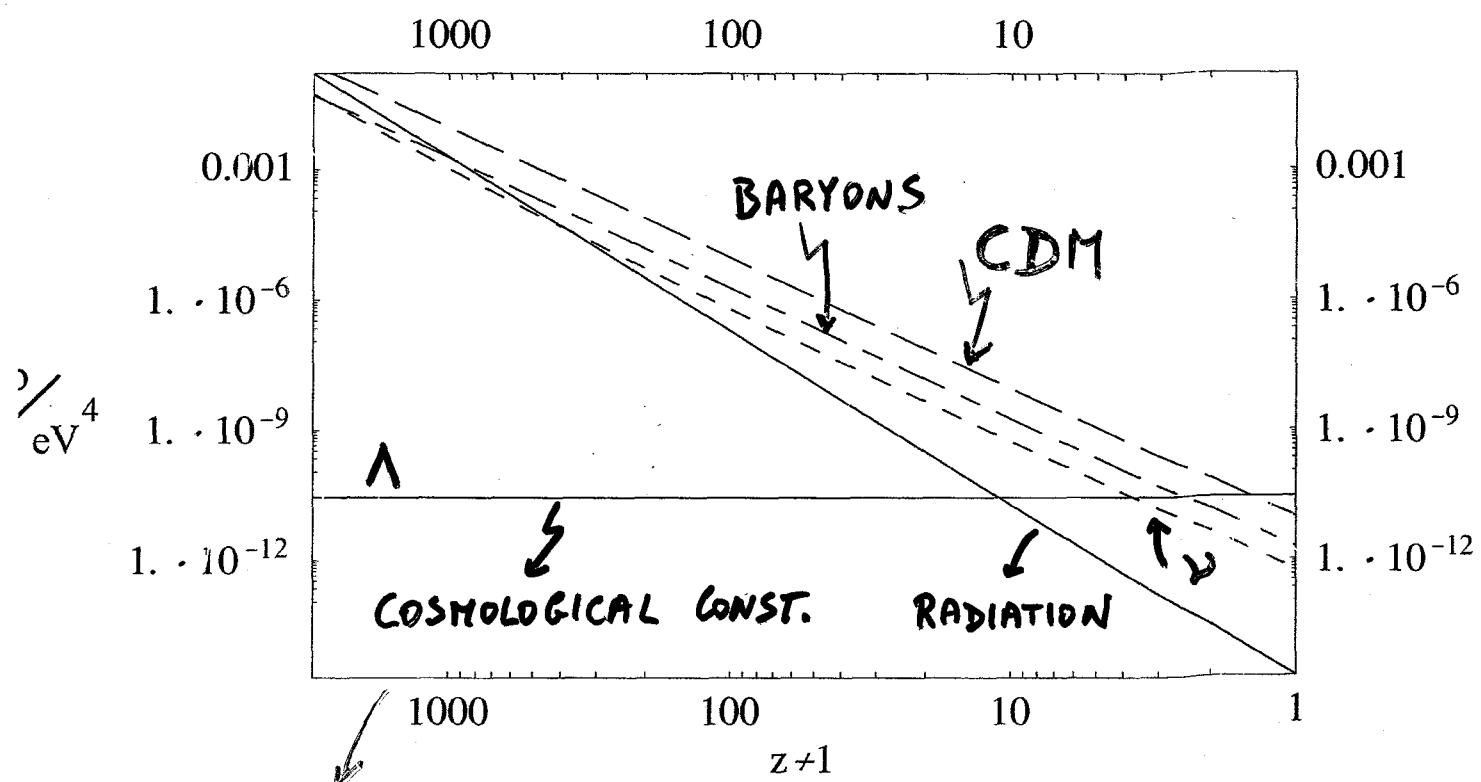
$$\bullet 7 \times 10^{-4} < \Omega_y < 2 \times 10^{-2}$$

$$\bullet \quad \Omega_{DM} \sim 0.3$$

$$\Omega_{DE} \sim 0.7$$

COSMOLOGICAL ENERGY DENSITIES AS A FUNCTION OF REDSHIFT IN THE Λ CDM MODEL

FARDON, NELSON, WEINER

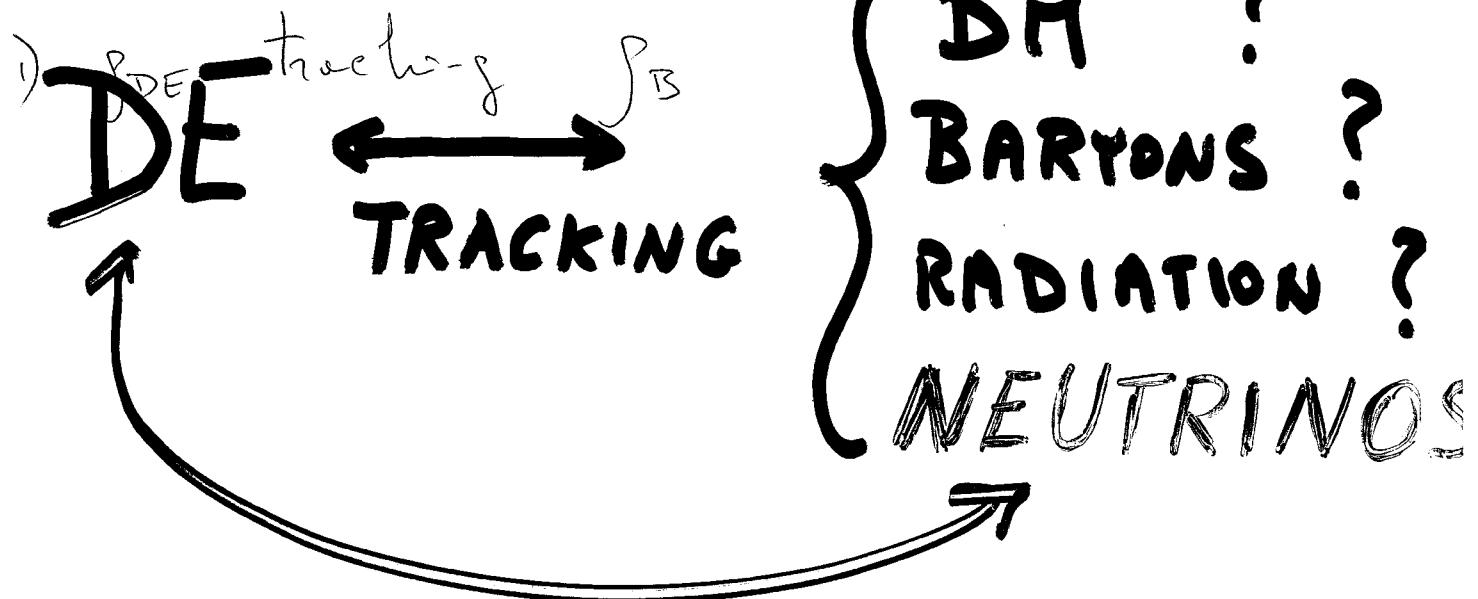


$z \approx 1100$ from CMB \rightarrow DM dominance

$$\frac{\rho_{\Lambda}}{\rho_{DM}} \sim \frac{1}{3 \times 10^8} \quad \text{at recombination}$$

$$\frac{\rho_{\Lambda}}{\rho_{DM}} \sim 2 \div 3 \quad \text{today}$$

Domestic:



$$\rho_{\text{DARK}} = \rho_\nu + \rho_{\text{DE}}$$

FARDON, NELSON, WEINER
FNW PROPOSAL

$$\rho_\nu \sim \rho_\Lambda \text{ within a factor of } 10^3$$

is NOT a coincidence \Rightarrow but a relationship holding over a large portion of the history of the Universe

m_ν ?

$$m_\nu \propto \frac{1}{n_\nu} + \text{effect of } \nu \text{ clustering}$$

when ν becomes non-relativ.

PECCEI $\Rightarrow \nu$'s cluster, i.e. gravity pulls some ν into existing DM halos

$$\text{ex: } m_\nu = 0.6 \text{ eV} \Big|_{\text{at } z=1} \quad z = \frac{T}{T_0} - 1 \quad \overline{T}_{z=1} = 2 T_0$$

✓ ν clustering produces an overdensity of ~ 30 in the local group

$$m_\nu^0 \Big|_{\text{in our vicinity}} \sim 0.6 \times 2^3 \times \frac{1}{30} \sim 0.15 \text{ eV}$$

$$m_\nu^0 \Big|_{\text{outside gravit. bound systems}} \sim 0.6 \times 2^3 \sim 5 \text{ eV}$$

$\hookrightarrow w \approx -0.8$

in the non-relativistic limit $\Rightarrow \rho_\nu = m_\nu n_\nu$

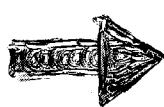
$$V(m_\nu(\phi)) = m_\nu(\phi) n_\nu + V_0(m_\nu(\phi))$$

FNW.

SCALAR POTENTIAL

Kaplov, Nelson THERMAL BACKGROUND NEUTRINOS
Weiner ACT AS A SOURCE DRIVING $m_\nu \downarrow$

Assume $V_0(m_\nu(\phi))$ is MINIMIZED FOR LARGE $m_\nu \uparrow$

 COMPETITION BETWEEN $m_\nu n_\nu$ and V_0
 \rightarrow minimum at an intermediate value of m_ν

Universe expansion \Rightarrow density $\downarrow \Rightarrow$ source term $\downarrow \Rightarrow m_\nu$ is driven to larger values

Assume: $\dot{\rho}_{\text{DARK}}$, i.e. $V(m_\nu)$ is STATIONARY

w.r.t. variations in m_ν :

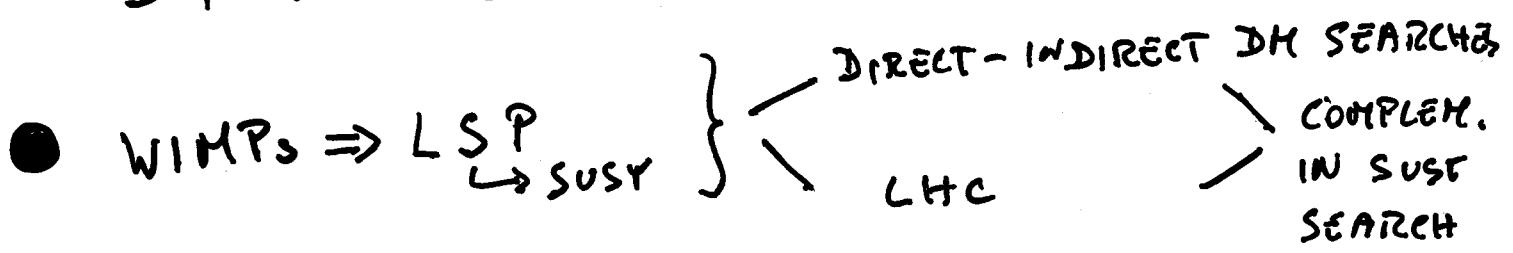
$$V'(m_\nu) = n_\nu + V'_0(m_\nu) = 0$$

\rightarrow conservation of energy: $\dot{\rho} = -3H(\rho + p)$

$$\nu \equiv \frac{\rho_{\text{dark}}}{\rho_{\text{dark}}} \quad w+1 = \frac{\Omega_\nu}{\Omega_\nu + \Omega_\phi} = \frac{m_\nu n_\nu}{m_\nu n_\nu + V_0(m_\nu)} \rightarrow m_\nu \approx \frac{1}{n_\nu}$$

CONCLUSIONS

- STRONG EVIDENCE FOR NON-BARYONIC DM \Rightarrow EVIDENCE FOR NEW PHYSICS BSM



● $\Omega_{DE} \sim \frac{2}{3}$ $\Omega_{DM} \sim \frac{1}{3}$

- WHAT IS DE ? COSM. CONST., DYNAMICAL QUANTITY (ex. SCALAR QUINTESSENCE)

- $\Omega_{DE} \sim 2 \Omega_{DM}$ | TODAY PROBLEM

ARE DE - DM SOURCES "CORRELATED"?
DE - DM INTERACTION?

DE presence can largely influence DM even though no direct DE-DM interaction exists \rightarrow ex. scalar-tensor theories of gravity,

OUTLOOK

THEORY

Superstrings, Brane theories

RARE PHYSICS

F_{CNC}, CP \neq , g-2

NEW PHYSICS

AT THE TEV
SCALE

HIGH ENERGY

PHYSICS

LHC

ILC

ν PHYSICS

ASTRO PARTICLE
PHYSICS

- CMB: WMAP, PLANCK
- γ : GLAST, AGILE, HAWC
- ANTIM: PAMELA, AMS
- ν : ANTARES, NEMO, AMANDA
- DIRECT DM SEARCHES ICECUBE...