



The Abdus Salam  
International Centre for Theoretical Physics



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**SCHOOL ON QUANTUM PHASE TRANSITIONS  
AND  
NON-EQUILIBRIUM PHENOMENA IN COLD ATOMIC GASES**

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***BCS - BEC Crossover: An introduction***

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# BCS-BEC CROSSOVER: AN INTRODUCTION

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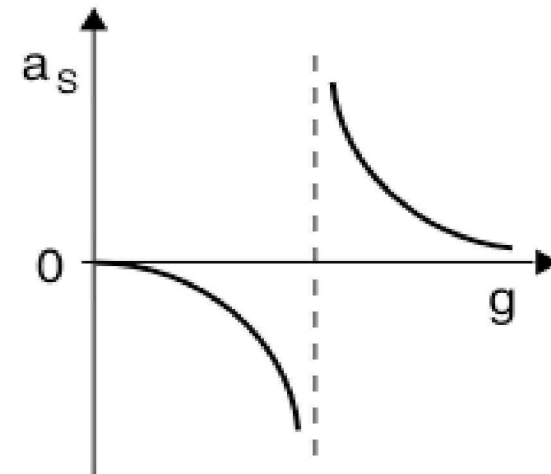
Technische Universität München

problem: Fermions  $\uparrow\downarrow$  with density  $n = k_F^3/3\pi^2$  and

attractive two particle interaction  $V_{\uparrow\downarrow}(\vec{x}) = -\tilde{g} \cdot \delta(\vec{x})$

two-particle bound state :  $g > g_c$

binding energy  $\epsilon_b = \hbar^2/ma_s^2$



motivation: ultracold gases  ${}^6\text{Li}$ ,  ${}^{40}\text{K}$ ,  $T_c \approx 10^{-7} \text{ K}$

idea goes back to H. Stoof et.al. PRL **76**, 10 (1996)

Outline: 1) BCS-theory,  $T = 0$

2) Ginzburg-Landau theory,  $T \approx T_c$

3) Unitarity limit and universality

4) Conclusion and open basic problems

## Superfluidity of Fermions

1911 conventional SC's Hg, Al, ...  $T_c \approx 1 - 23$  K

1960 pairing in nuclei,  $\Delta \approx$  MeV

1972 superfluid  $^3\text{He}$ ,  $T_c = 1$  mK (p-wave)

1975 neutron stars

1986 high- $T_c$  SC's  $\text{La}_2\text{CuO}_4, \dots$   $T_c = 35 - 138$  K

1991 Alkali-doped  $\text{C}_{60}$   $T_c \approx 30$  K

1994 p-wave SC's in  $\text{Sr}_2\text{RuO}_4$   $T_c = 1.5$  K

1998 color SC's  $\langle qq \rangle \neq 0$ ,  $\Delta \approx$  GeV

## 1) BCS-THEORY (1957) (only $T = 0$ )

basic idea: Fermi-sea is unstable towards formation of

Cooper-pairs for arbitrary weak attraction

$$|\psi\rangle_C = \sum_{k > k_F} \frac{1}{2\epsilon_k - E} c_{k\uparrow}^\dagger c_{-k\downarrow}^\dagger |F\rangle$$

exp. small binding  $E = 2\epsilon_F - 2\hbar\omega_c \exp\left(-\frac{\pi}{k_F |a_s|}\right)$

BCS groundstate (variational Ansatz for arbitrary  $g$ )

$$\psi_{\text{BCS},N} = \hat{A} \{ \underbrace{\phi(1\ 2) \phi(3\ 4) \cdots \phi(N-1\ N)} \}$$

ideal Bose gas of identical pairs

pair-state  $\phi(1, 2) = \psi(\vec{x}) \cdot \chi(\vec{r}) \cdot \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$

internal wavefunction  $\chi(\vec{r}) \sim \exp -r/\xi_b \rightarrow$  pair radius  $\xi_b$

COM-coord.  $\vec{x} \rightarrow$  inhomogeneous pair-amplitude  $\psi(\vec{x})$

increase  $g$ : extended  $\rightarrow$  local pairs  $\rightarrow$  smooth crossover

(Keldysh 1965, Eagles 1969, Leggett 1980)

model Hamiltonian: separable interaction between

Fermions of opposite spin (Gorkov)

$$\hat{H} - \mu\hat{N} = \sum_{k\sigma} (\epsilon_k - \mu) c_{k\sigma}^\dagger c_{k\sigma} - \frac{\tilde{g}}{V} \sum_q b_q^\dagger b_q$$

pair-operators

$$b_q^\dagger = \sum_k c_{k+q\downarrow}^\dagger c_{-k\uparrow}^\dagger$$

BCS Hamiltonian  $\hat{H}_{\text{BCS}} = -\frac{\tilde{g}}{V} b_0^\dagger b_0$  (only  $q = 0$  pairs)

exact ground state for fixed  $\mu$  is a coherent state

$$|\psi\rangle_{\text{BCS}} = \prod_k (u_k |00\rangle_k + v_k |11\rangle_k) \sim \exp \sum_k \frac{v_k}{u_k} c_{k\uparrow}^\dagger c_{-k\downarrow}^\dagger |0\rangle$$

at fixed  $N$  (even):  $|\psi_N\rangle_{\text{BCS}} \sim \left( \sum_k \frac{v_k}{u_k} c_{k\uparrow}^\dagger c_{-k\downarrow}^\dagger \right)^{N/2} |0\rangle$

$\chi_k = v_k/u_k =$  Fouriertransf. of internal wavefunction



gap parameter  $\Delta_k$  follows from gap-equation

$$\frac{1}{\bar{g}} = \frac{1}{V} \sum_k \frac{1}{\sqrt{(\epsilon_k - \mu)^2 + \Delta_k^2}}$$

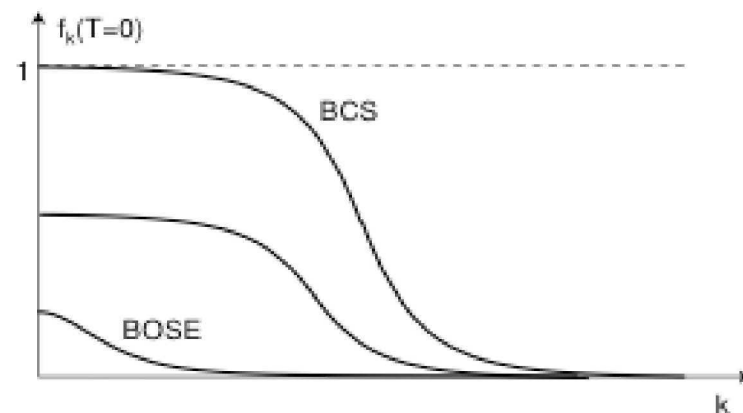
chemical potential  $\mu$  from number equation ( $\mu_{\text{BCS}} = \epsilon_F$ )

$$\langle N \rangle = 2 \sum_k v_k^2 = \sum_k \left( 1 - \frac{\epsilon_k - \mu}{\sqrt{(\epsilon_k - \mu)^2 + \Delta_k^2}} \right)$$

momentum distrib.

$$f_k = v_k^2 \rightarrow \frac{4\pi n \xi_b^3}{1 + (k\xi_b)^2}$$

since  $\mu_{\text{BEC}} \rightarrow -\epsilon_b/2$



exact solution of  $\hat{H}_{\text{BCS}}$  at  $T = 0$  and fixed  $N$

Richardson 1963, Ortiz/Dukelsky cond-mat/0503664

are Cooper pairs Bosons ?

antisymmetrization  $\hat{A}$  reduces the condensate fraction  $N_0/N$  to  $\approx \Delta_0/\epsilon_F \ll 1$  in weak coupling, nevertheless

**BCS-superfluidity is BEC of pairs !**

**BCS-LIMIT:** weak coupling if  $k_F |a_s| < 0.5$  (in practice)

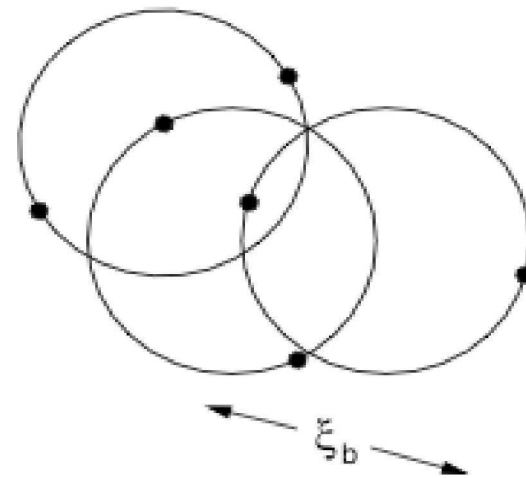
pairs form and condense at

$$T_c \approx 0.3 T_F \exp -\frac{\pi}{2k_F |a_s|} \ll T_F \quad 2\Delta_0 = 3.52 k_B T_c$$

superfluidity is destroyed by pair- breaking

pair size  $\xi_b = \xi_0 = \frac{\hbar v_F}{\pi \Delta_0}$  is large  $k_F \xi_b \gg 1$

$$k_F \xi_0 \approx \begin{cases} 10^3 & \text{supercond.} \\ 10^2 & {}^3\text{He} \\ 10^1 & \text{high-}T_c \end{cases}$$



## BCS-transition in a harmonic trap

local density approx.  $\xi_b \ll R_{TF}$  ( $\Delta_0 \gg \hbar\omega$ ) valid

if  $N \gg N^* \approx \exp 3\pi/2k_F|a_s| = 10^5$  at  $k_F|a_s| = 0.4$

$N \ll N^*$ : pairs formed within a single shell  $\Delta_0 \ll \hbar\omega$

as in atomic nuclei (Bruun, Heiselberg, Mottelson 02)

**BEC limit:** strong coupling  $|\epsilon_b| \gg \epsilon_F \rightarrow k_F a_s \ll 1$

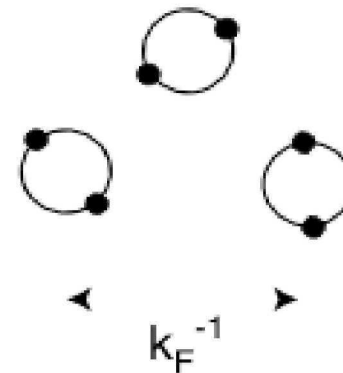
point bosons  $\xi_b \approx a_s \rightarrow 0$  form far above  $T_c$

gap parameter  $\Delta_0 \rightarrow \sqrt{2|\epsilon_b|\epsilon_F} \gg \epsilon_F$  is irrelevant

pair size

$$k_F \xi_b \ll 1$$

$\xi_b \neq$  coherence length  $\xi_0$



condensation temperature  $T_c = 0.218 T_F \sim n^{2/3}$

from ideal Bose gas with  $n_B = n/2$  and  $m_B = 2m$

BEC is destroyed by excitations to finite momenta  $\rightarrow$

BCS-theory does not describe crossover at finite  $T$  !

still,  $T_c$  increases monotonically (?) with coupling:

Nozieres/Schmitt-Rink JLTP **59**, 195 (1985)

## 2) Ginzburg–Landau theory

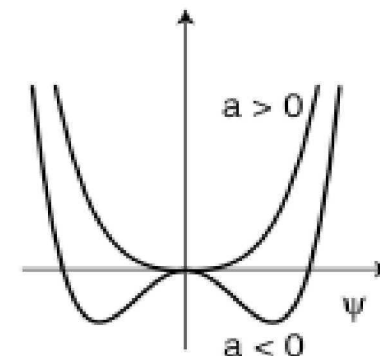
Drechsler/Zw. Ann. Physik **1**, 15, 1992,  $d = 2$

Randeria et.al. PRL **71**, 3202, 1993,  $d = 3$

idea: condensation is associated with broken

gauge symmetry  $\rightarrow$  complex order parameter  $\psi$

$$F[\psi] = \int d^d x \left[ a |\psi|^2 + \frac{b}{2} |\psi|^4 + c |\vec{\nabla} \psi|^2 \right]$$



Thermodynamics  $Z = \text{Tr} \exp -\beta(\hat{H} - \mu\hat{N})$

linearize  $\sum_q b_q^\dagger b_q$  via Hubbard-Stratonovich  $\rightarrow$

slowly varying order parameter  $\psi(\vec{x}, \tau)$

$$Z = Z_0 \int D\psi(\vec{x}, \tau) \exp(-F[\psi])$$

expansion up to  $\psi^4$  near  $T_c$



quantum GL-functional  $F[\psi] =$

$$\frac{1}{\hbar} \int_0^{\beta\hbar} d\tau \int d^d x \left[ -\mu^* |\psi|^2 + \frac{\hbar^2}{2m^*} |\vec{\nabla}\psi|^2 + \frac{g^*}{2} |\psi|^4 + \hbar \psi^* \partial_\tau \psi \right]$$

GL-theory is coherent state representation of

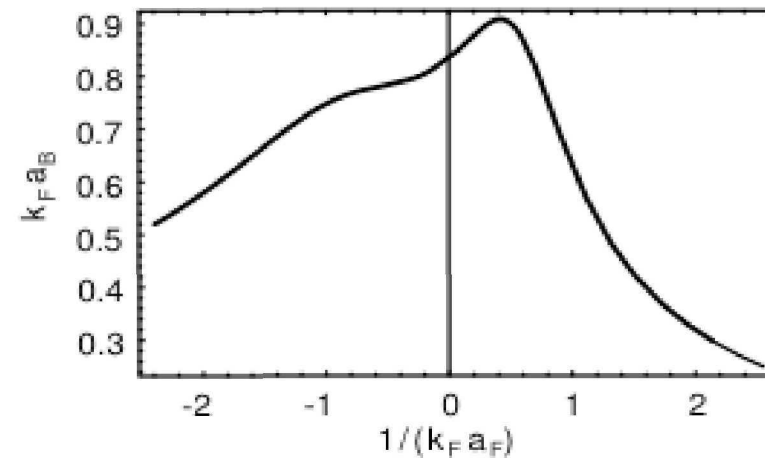
interacting Bosons, smooth crossover from Gorkov's

BCS GL-functional to Gross-Pitaevski for BEC's

$g^* = 4\pi\hbar^2 a_B / m^*$  defines dimer scatt. length  $a_B$

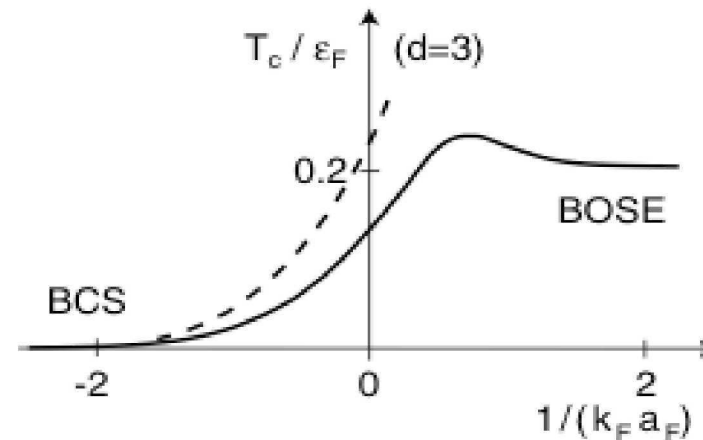
Born approximation

$$a_B = \begin{cases} 6|a_F| & \text{BCS} \\ 2a_F & \text{Bose} \end{cases}$$



Pauli principle gives statistical, repulsive interaction between composite Bosons

critical temperature  
 in Gaussian approx.  
 (Thouless criterion)



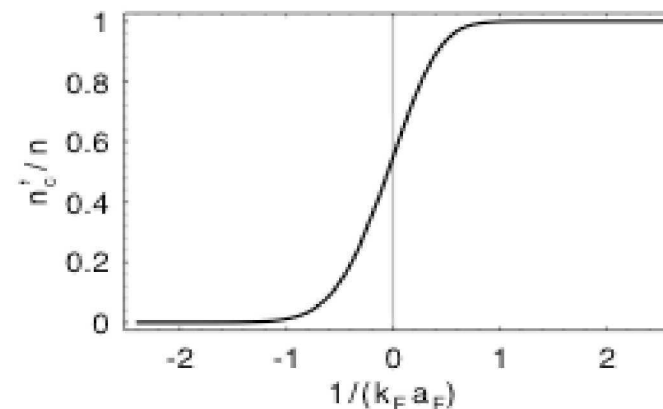
BCS-limit: density int.  $T_c = T_c^{(0)}/2.2$  (Gorkov 1961)

BEC-limit: dilute Bose-gas  $\Delta T_c/T_0 \approx 1.3 n_B^{1/3} a_B > 0!$

exact result  $a_B = 0.6 a_F$  (Petrov/Shlyap./Salomon 2003)

pair formation is just a matter of shifting the equilibrium from  $\bullet + \bullet$  to  $\bullet\bullet$  it coincides with the superfluid transition only in weak coupling

preformed pairs  
at condensation



nature of the transition changes in a narrow range of coupling !

optimally doped high- $T_c$  superconductors

determine  $\xi(0)$  from  $\left. \frac{dH_{c2}}{d \ln T} \right|_{T_c} = \frac{\phi_0}{2\pi\xi^2(0)}$  gives

$$k_F \xi(0) \approx \begin{cases} 5 & \text{YBCO} \\ 8 & \text{BSCCO} \end{cases}$$

interaction strength  $g$  determines

$$\rightarrow \kappa = \frac{\lambda}{\xi} \approx \begin{cases} 99 & \text{YBCO} \\ 87 & \text{BSCCO} \end{cases} \quad \text{as observed}$$

conclusion optimally doped high- $T_c$ 's are far from  
weak coupling but still on the BCS-side

(Stintzing/ Zw. PR **B56**, 9004 (1997))

pseudogap in underdoped high- $T_c$ 's is not related  
to the BCS-BEC crossover (precursor fluctuations)

### 3) Unitarity limit and universality

interacting Fermions with scattering amplitude

$$f(k) = \frac{-1}{a^{-1} + r^*k^2/2 + ik} \rightarrow i/k$$

if  $k_F r^* \ll 1$       broad Feshbach resonance

$\rightarrow k_F \sim n^{1/3}$  and  $\epsilon_F$  are the only scales left       $\rightarrow$

universal thermodynamics  $F(T)/N = \epsilon_F \cdot a(T/T_F)$

(e.g. specific heat, expansion energy, ... Thomas 2004)

and dynamics, e.g. viscosity  $\eta \approx 0.1\hbar n$  (Shuryak 2004)

universal numbers at resonance  $(k_F a)^{-1} = 0$

$\epsilon_{int} = -0.34\epsilon_F$ ,  $\Delta_0 \approx 0.5\epsilon_F$ ,  $T_c \approx 0.1\epsilon_F$ ,  $v_s = \dots$

equation of state at  $T = 0$  from QMC

(Carlson et.al. PRL **91**, Astrakharchik et.al. PRL **93**)



#### 4) Conclusion and open basic problems

- 1) BCS–BEC is a smooth crossover for s-waves, quantitative results exist only asymptotically
- 2) Exact solution for arb. coupling is possible only for  $\hat{H}_{\text{BCS}}$  or in 1D (Fuchs/Recati/Zw. PRL 93)

3) quantitative theory at all  $0 < T < T_c$

4) interacting fermions at the unitarity limit

5) p-wave crossover  $\rightarrow$  quantum phase transition

