



The Abdus Salam
International Centre for Theoretical Physics



SMR 1760 - 10

**COLLEGE ON
PHYSICS OF NANO-DEVICES**

10 - 21 July 2006

Quantum optics using quantum dots

Presented by:

Atac Imamoglu

ETH-Zürich, Switzerland

Quantum optics using quantum dots

A. Imamoglu

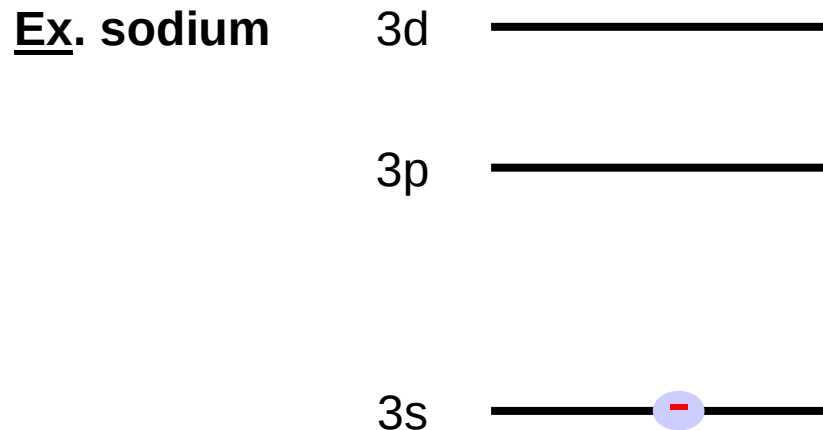
*Quantum Photonics Group, Department of Physics
ETH-Zürich*

Outline

- 1) Brief overview of quantum dots – a.k.a. artificial atoms
- 2) Photon correlation measurements
- 3) Single photon sources

Atoms

- The allowed internal energy states (describing the relative motion of electrons with respect to a nucleus) form a discrete anharmonic set.
- Each discrete atomic orbital has a fine and hyperfine structure arising from the coupling of electronic spin to orbital angular momentum and electronic spin to nuclear spin, respectively.
- The eigenstates describing the center-of-mass motion of a free atom form a continuum; internal & external degrees of freedom could be coupled via light forces $\Rightarrow$ optical trapping.

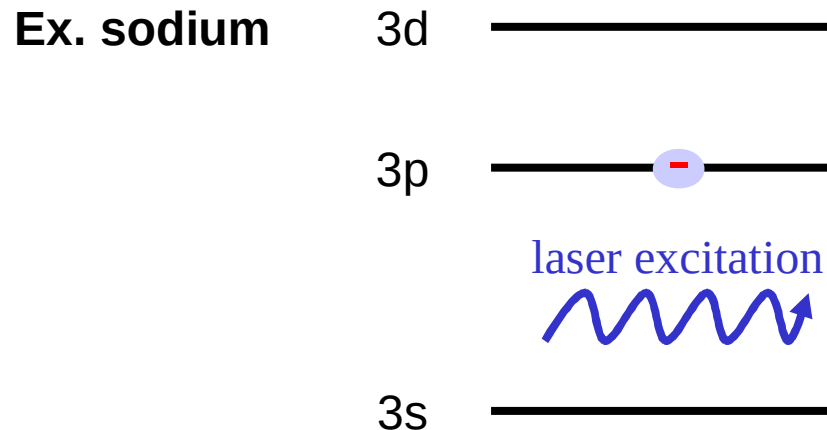


Each atomic orbital consists of a number of discrete states:

Ex.: 3s $\Rightarrow$ 8 distinct eigenstates due to coupled electron (2) and nuclear (4) spin states.

Atoms

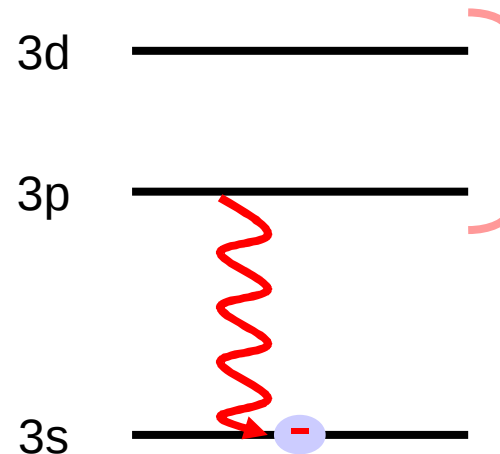
- The allowed internal energy states (describing the relative motion of electrons with respect to a nucleus) form a discrete anharmonic set.
- Each discrete atomic orbital has a fine and hyperfine structure arising from the coupling of electronic spin to orbital angular momentum and electronic spin to nuclear spin, respectively.
- The eigenstates describing the center-of-mass motion of a free atom form a continuum; internal & external degrees of freedom could be coupled via light forces $\Rightarrow$ optical trapping.



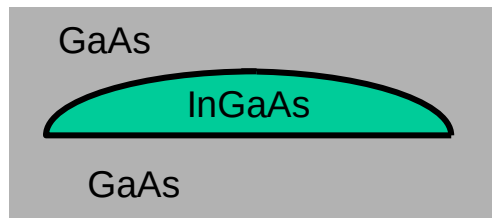
Atoms

- The allowed internal energy states (describing the relative motion of electrons with respect to a nucleus) form a discrete anharmonic set.
- Each discrete atomic orbital has a fine and hyperfine structure arising from the coupling of electronic spin to orbital angular momentum and electronic spin to nuclear spin, respectively.
- The eigenstates describing the center-of-mass motion of a free atom form a continuum; internal & external degrees of freedom could be coupled via light forces $\Rightarrow$ optical trapping.

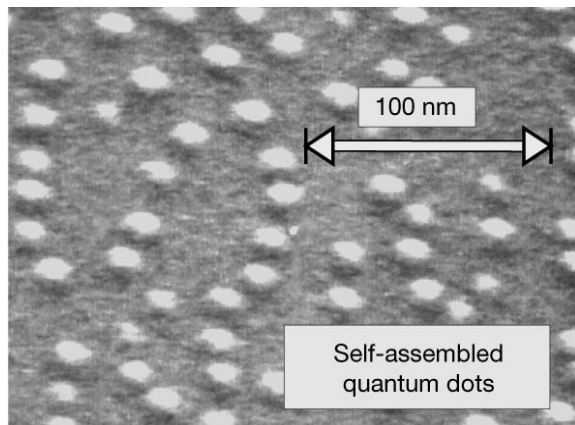
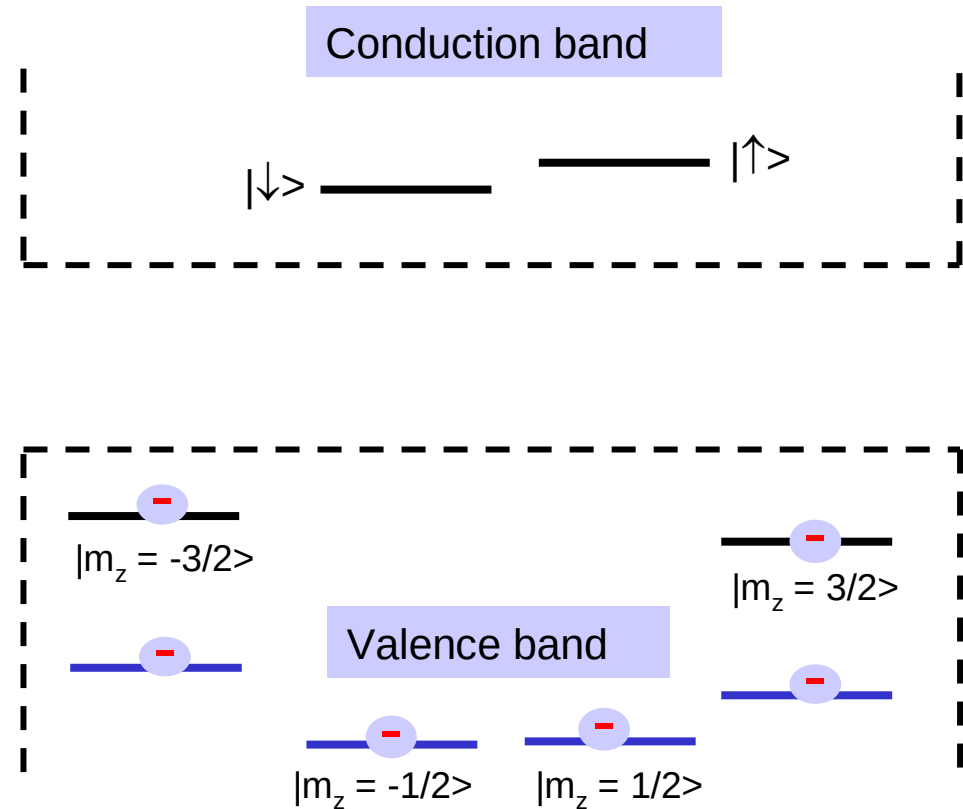
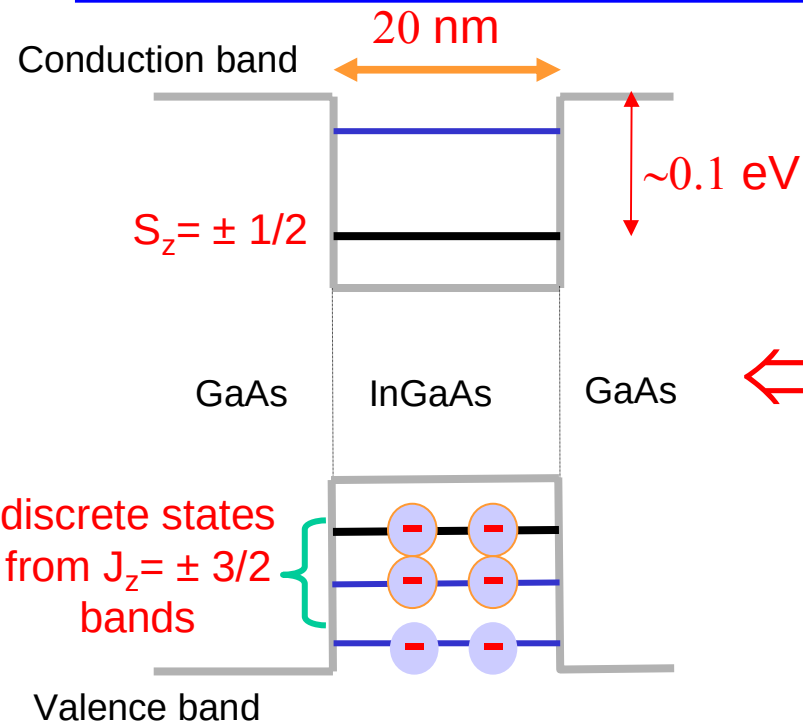
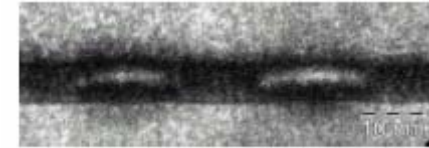
Ex. sodium



All excited states are broadened in energy due to spontaneous emission: the absorption/emission lineshape is a Lorentzian.



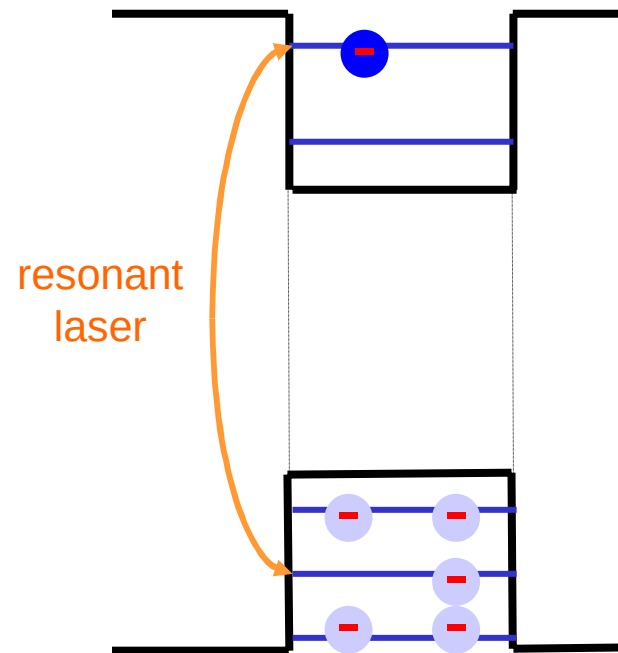
Quantum dots (QD)



- QDs have discrete states for electrons & holes.
- QD location is fixed in a macroscopic crystal by growth.
- There are $\sim 10^5$ atoms in each QD.

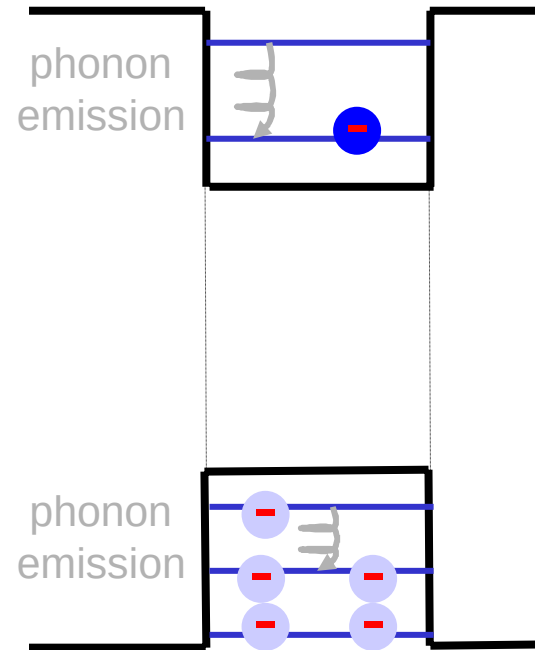
Light generation by a quantum dot

a) Optical excitation of an electron



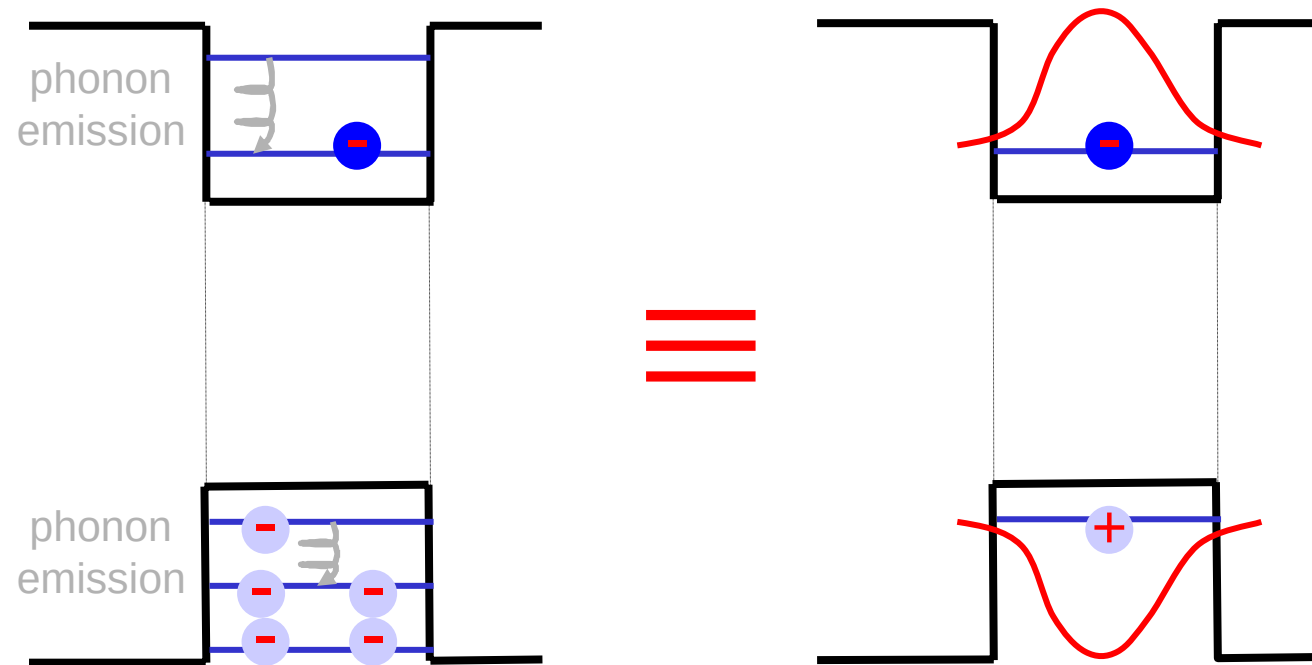
Light generation by a quantum dot

b) Relaxation by phonon emission (~ 10 ps)



Light generation by a quantum dot

b) Relaxation by phonon emission (~ 10 ps)

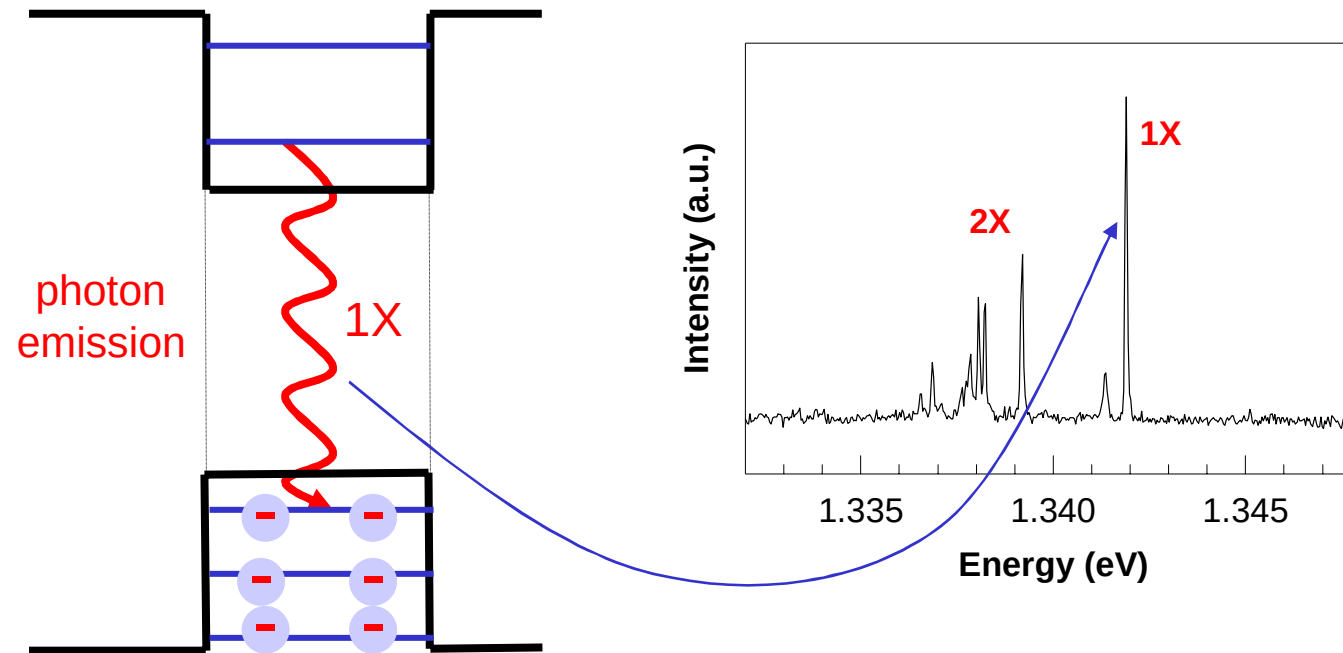


electron-hole pair = exciton

Light generation by a quantum dot

c) Emission of a photon (~ 1 ns)

Spectrum of emitted photons

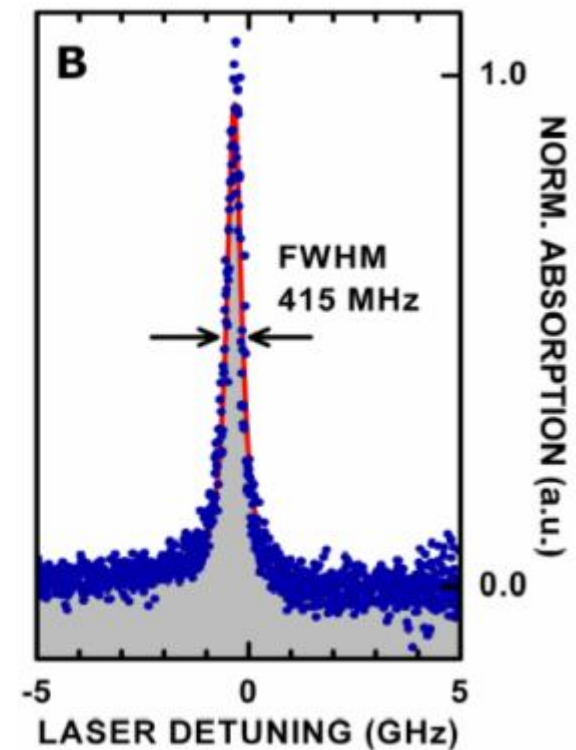
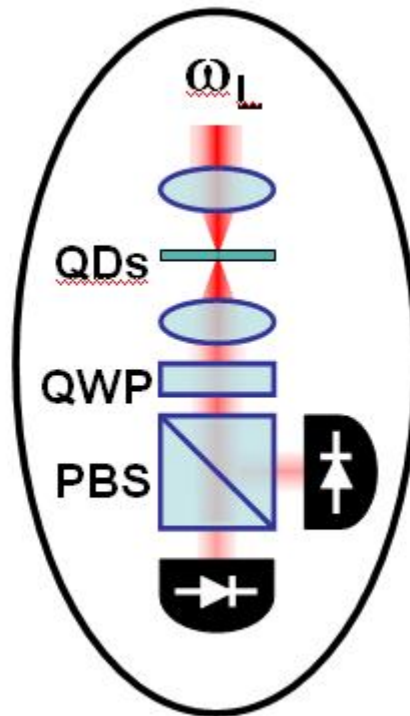
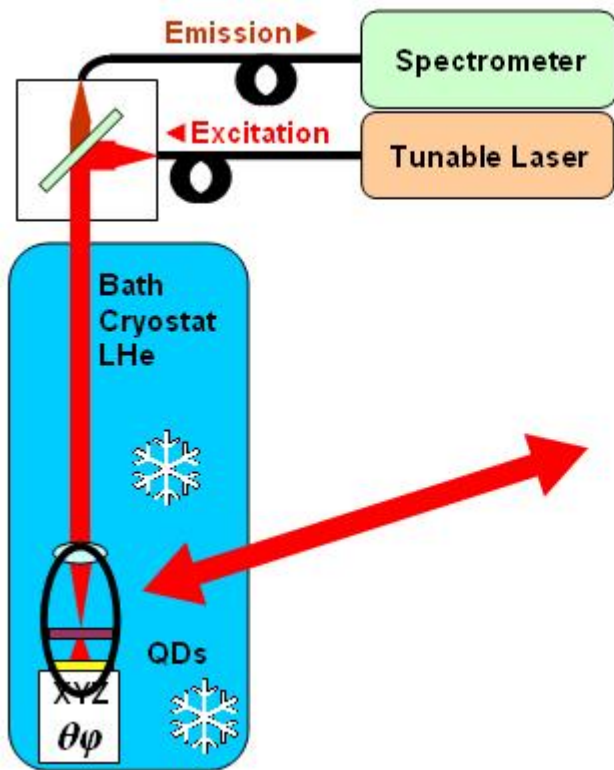


⇒ QDs typically exhibit several sharp emission lines

⇒ At low pump power, a single (exciton) line dominates the spectrum; the width of this line is resolution limited at $\sim 30 \mu\text{eV} \ll kT \sim 300 \mu\text{eV}$.

How narrow are quantum dot exciton lines?

- The least perturbing measurement that can be done is a resonant absorption measurement



$\Rightarrow$ Observed absorption linewidth is $1.5 \mu\text{eV} \sim 1.5 \Gamma_{\text{spont}}$

Photon correlation measurements

- Intensity (photon) correlation function:

→ gives the likelihood of a second photon detection event at time $t+\tau$, given an initial one at time t ($\tau=0$).

$$g^{(2)}(\tau) = \frac{\langle : I(t) I(t + \tau) : \rangle}{\langle I(t) \rangle^2}$$

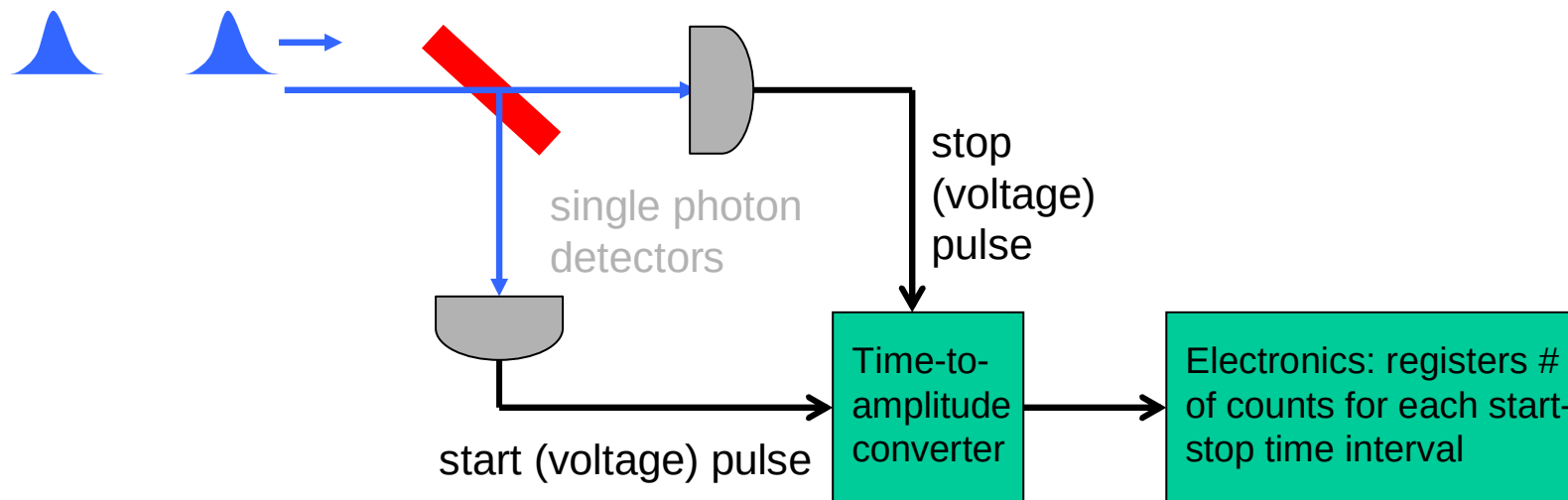
Photon correlation measurements

- Intensity (photon) correlation function:

→ gives the likelihood of a second photon detection event at time $t+\tau$, given an initial one at time t ($\tau=0$).

$$g^{(2)}(\tau) = \frac{\langle : I(t) I(t + \tau) : \rangle}{\langle I(t) \rangle^2}$$

- Experimental set-up for photon correlation [$g^{(2)}(\tau)$] measurement:



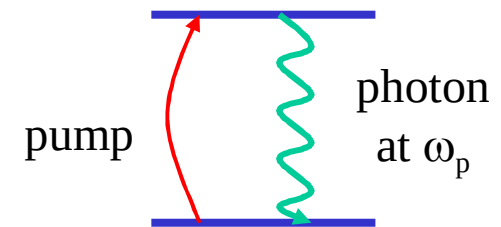
Photon antibunching

- Intensity correlation ($g^{(2)}(\tau)$) of light generated by a single two-level (anharmonic) emitter.

- Assume that at $\tau=0$ a photon is detected:
 - We know that the system is necessarily in the ground state $|g\rangle$
 - Emission of another photon at $\tau=0+\varepsilon$ is impossible.

⇒ Photon antibunching: $g^{(2)}(0) = 0$.

- $g^{(2)}(\tau)$ recovers its steady-state value in a timescale given by the spontaneous emission time.



nonclassical light

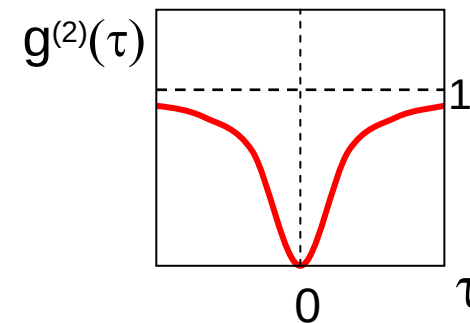
- If there are two or more 2-level emitters, detection of a photon at $\tau=0$ can not ensure that the system is in the ground state ($g^{(2)}(0) > 0.5$).

Signature of photon antibunching

- Intensity (photon) correlation function:

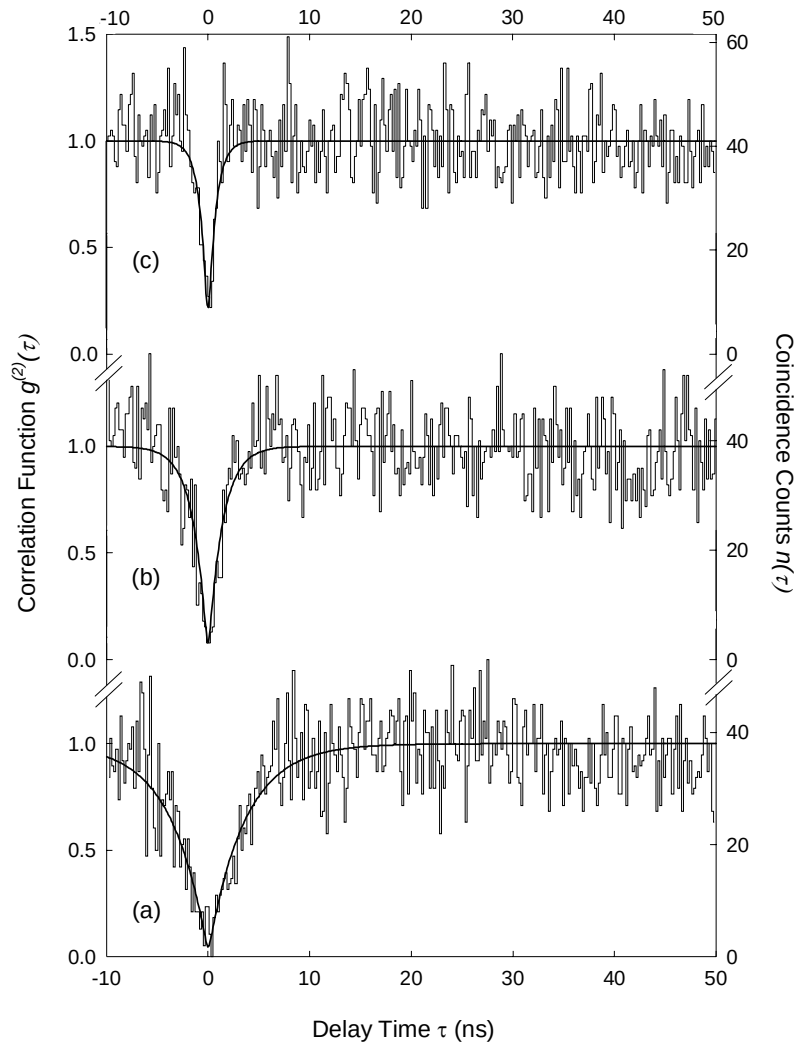
$$g^{(2)}(\tau) = \frac{\langle : I(t) I(t + \tau) : \rangle}{\langle I(t) \rangle^2}$$

- Single quantum emitter (i.e. an atom) driven by a cw laser field exhibits photon antibunching.



Photon antibunching from a Single Quantum Dot

(continuous-wave excitation)



105 W/cm²

$g^2(0)=0.1$

$\tau = 750$ ps

55 W/cm²

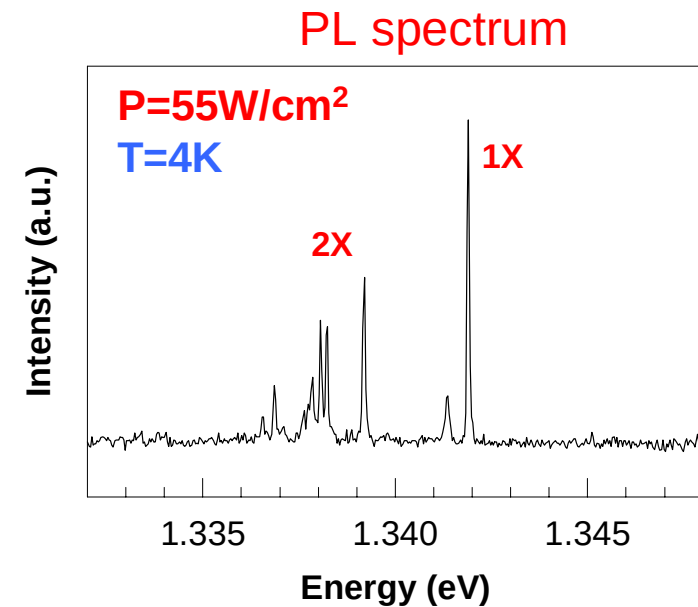
$g^2(0)=0.0$

$\tau = 1.4$ ns

15 W/cm²

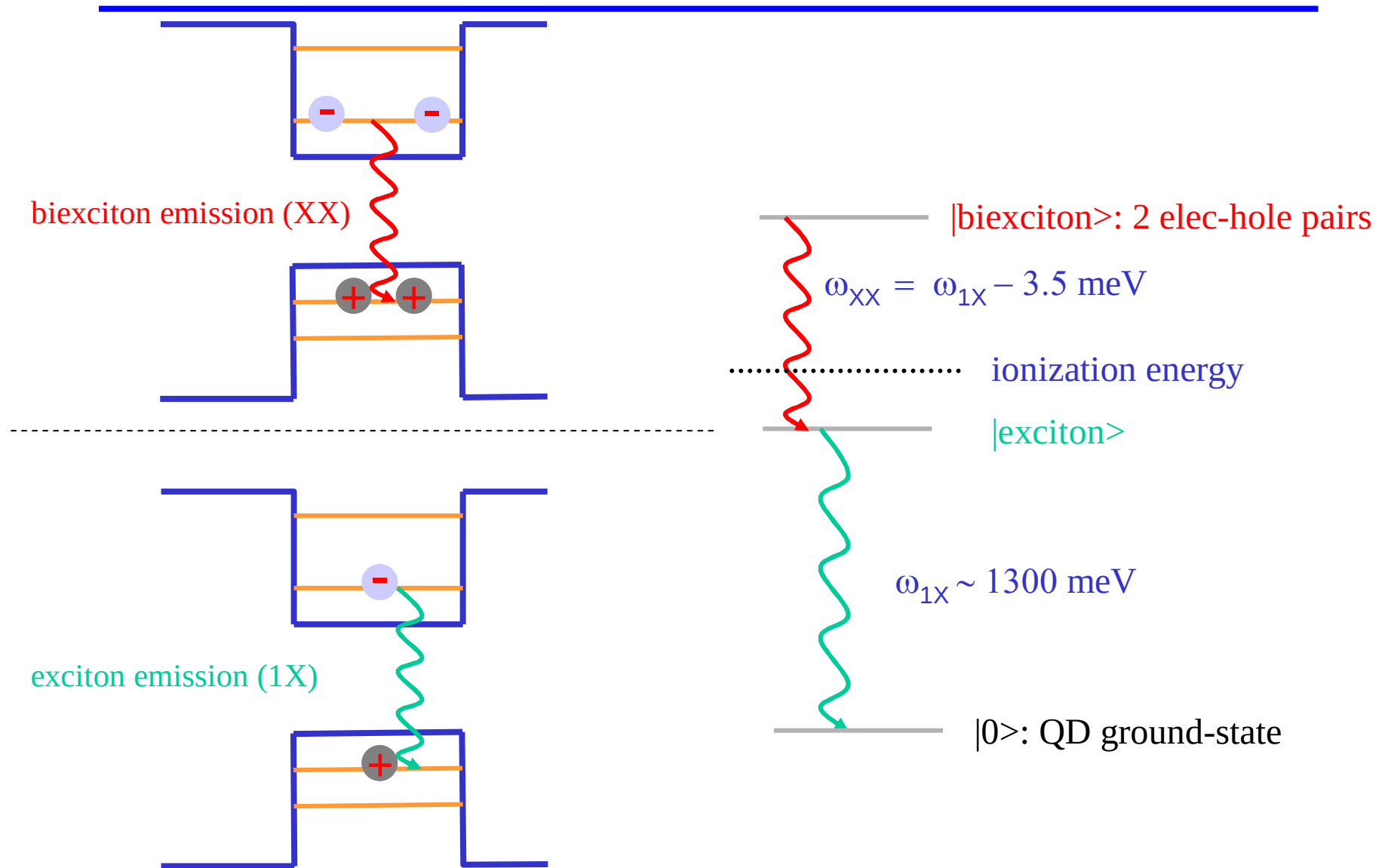
$g^2(0)=0.0$

$\tau = 3.6$ ns



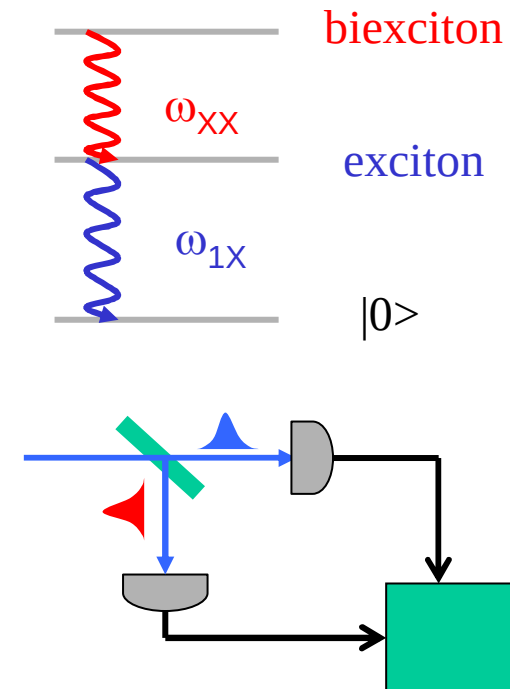
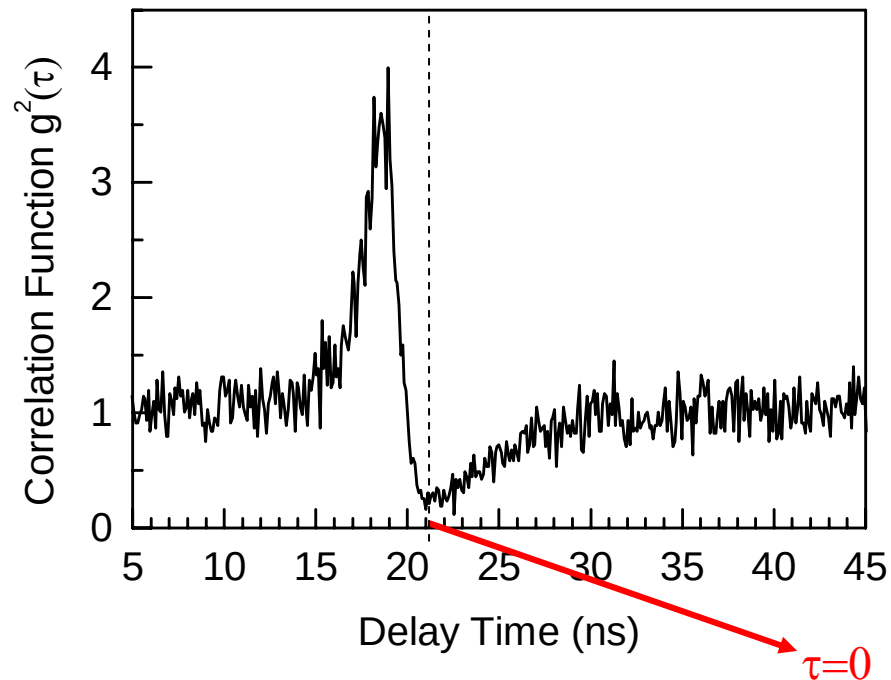
⇒ Proof of atom-like behaviour

Higher-order optical excitations in a quantum dot



Exciton/biexciton (X1/XX) cross-correlation

Pump power *well below* saturation level



⇒ At low excitation regime (average number of excitons < 1):

When a biexciton is observed, the QD projected is onto X1 state; as a result observation of an X1 photon becomes more likely than it is on average ⇒ **bunching**

Why quantum dots?

- Quantum dots do not have random thermal motion and they can easily be integrated in nano-scale cavities.
 - Just like atoms, quantum dots have (nearly) spontaneous emission broadened emission lines.
 - Exciton transition oscillator (dipole) strength ranges from 10 to 100 (depending on the QD type) due to a collective enhancement effect.
-
- A single confined electron interacts weakly with $\sim 10^5$ nuclei; the nature of hyperfine interactions is completely different.
 - novel non-Markovian decoherence mechanism for electron spin
 - single electron controls an isolated mesoscopic (nuclear) spin ensemble.
 - Interactions between electrons and lattice vibrations:
 - may enable optical manipulation of nano-mechanical systems.
 - Kondo-type many-body effects between QD electron spin and the spins of a nearby electron reservoir (n-type region).

Single Photon Sources

A regulated sequence of optical pulses that contain one-and-only-one photon

⇒ Why are we interested in single-photon sources?

i) **Quantum communication:**

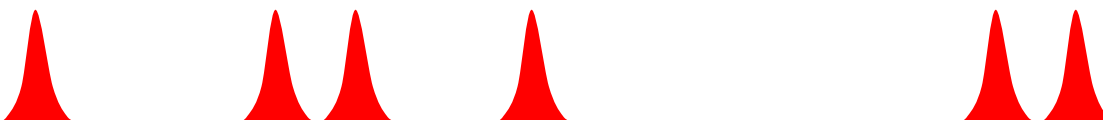
- key distribution in quantum cryptography
- quantum teleportation: transmission of quantum information without transmitting the physical qubit (which remains in an unknown state)

i) **Quantum computation:** single photons + linear optical elements + photo-detection
can eliminate the need for photon-photon interactions.

Typical light fields as a stream of photons

- thermal light (light bulb)
- 
- The diagram shows a horizontal timeline with an arrow pointing to the right labeled 'time'. There are four blue peaks of varying heights and widths. The first three peaks are relatively narrow and tall, while the fourth peak is significantly wider and shorter, indicating a higher probability of multiple photons arriving together (bunching).

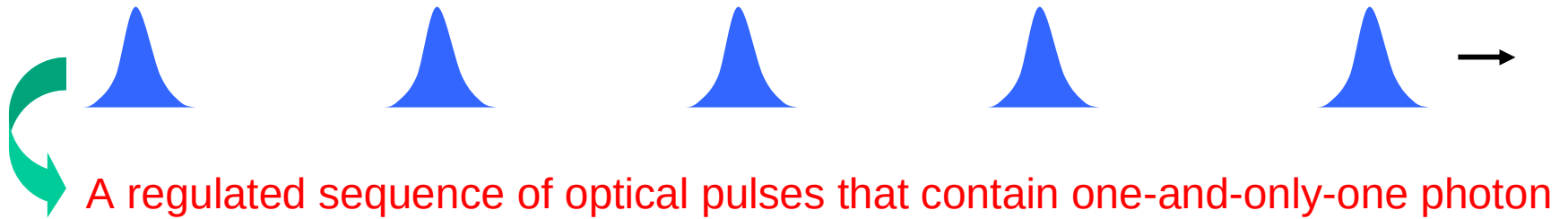
⇒ Photons are bunched together giving rise to large (super Poissonian) fluctuations in detected energy per unit time.

- coherent light (laser)
- 
- The diagram shows a horizontal timeline with an arrow pointing to the right labeled 'time'. There are four red peaks. The first, third, and fourth peaks are single, narrow, and tall. The second peak consists of two narrow peaks very close together, indicating a higher probability of single photons arriving independently (Poissonian distribution).

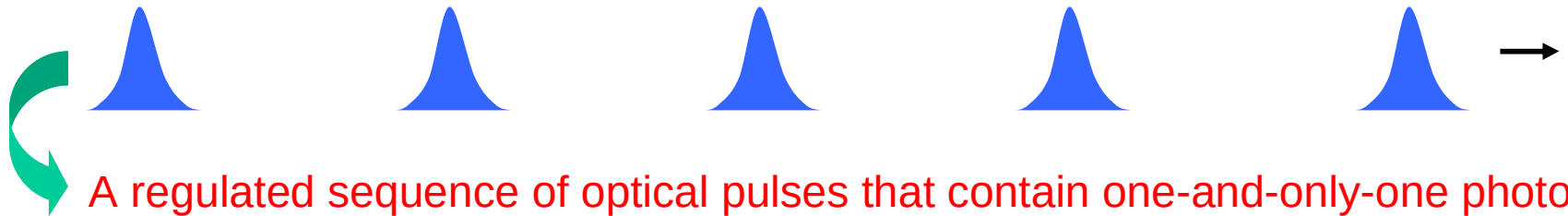
⇒ Photons are completely uncorrelated, giving rise to Poissonian photon detection events (shot noise in photo-current) – i.e. $\Delta n^2 = \langle n \rangle$.

⇒ Both thermal and coherent light correlations can be explained using classical Maxwell's equations.

Single Photon Sources



Single Photon Sources



⇒ How can we generate single photon pulses?

- An attenuated laser with average photon number per pulse $\langle n \rangle = 1$ is not a single-photon source, since standard deviation $\Delta n = 1$.
- Photons generated by a single quantum dot driven by a periodic pump (i.e. laser or electrical current) realise a single photon source:

Proposal: A. Imamoglu and Y. Yamamoto, Phys. Rev. Lett. **72**, 210 (1994)

Experiments: P. Michler *et al.*, Science **290**, 2282 (2000)

Z. Yuan *et al.*, Science **295**, 102 (2002)

C. Santori *et al.*, Nature **419**, 584 (2002)

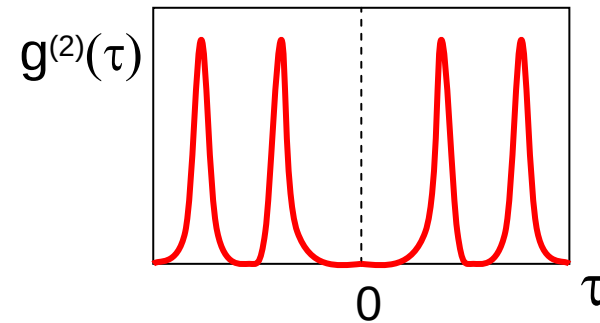
Signature of a single-photon source

- Intensity (photon) correlation function:

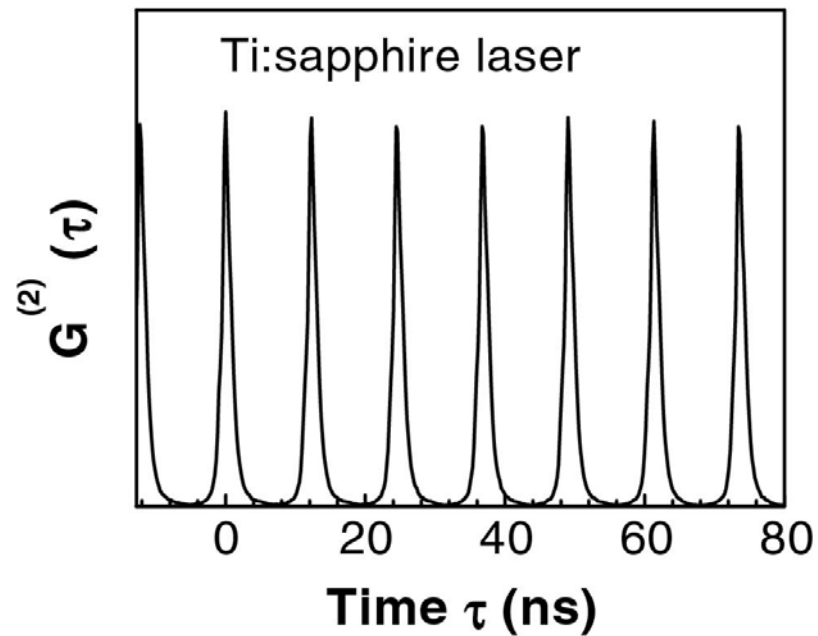
→ gives the likelihood of a second photon detection event at time $t+\tau$, given an initial one at time t ($\tau=0$).

$$g^{(2)}(\tau) = \frac{\langle : I(t)I(t+\tau) : \rangle}{\langle I(t) \rangle^2}$$

- Triggered single photon source: absence of a peak at $\tau=0$ indicates that none of the pulses contain more than 1 photon.

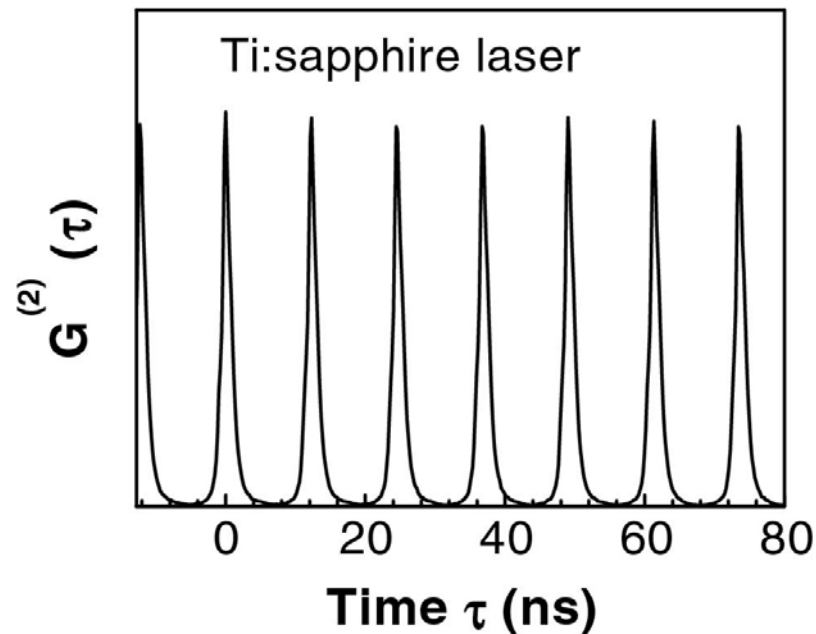


Single QD driven by a pulsed laser

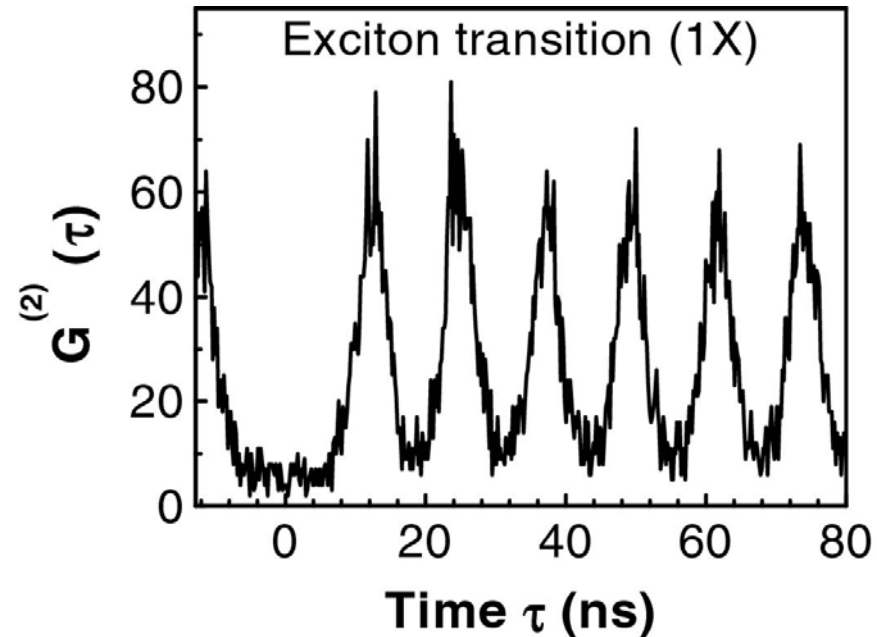


- ➔ all peaks in $G^{(2)}(\tau)$ have the same intensity, including the one at $\tau=0$
- ➔ pulsed coherent light

Photon correlation of a single-photon source



- ➔ all peaks in $G^{(2)}(\tau)$ have the same intensity, including the one at $\tau=0$
- ➔ pulsed coherent light

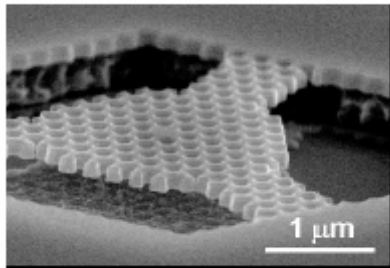


- ➔ the peak at $\tau=0$ disappears.
- ➔ single photon turnstile device with at most one photon per pulse

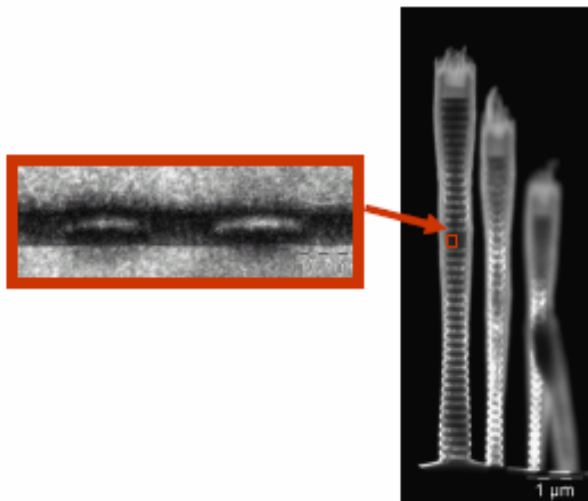
[P. Michler et al., Science **290**, 2282 (2000)]

Single-photon sources: state-of-the-art

(I. Robert-Phillip and co-workers, LPN Paris)

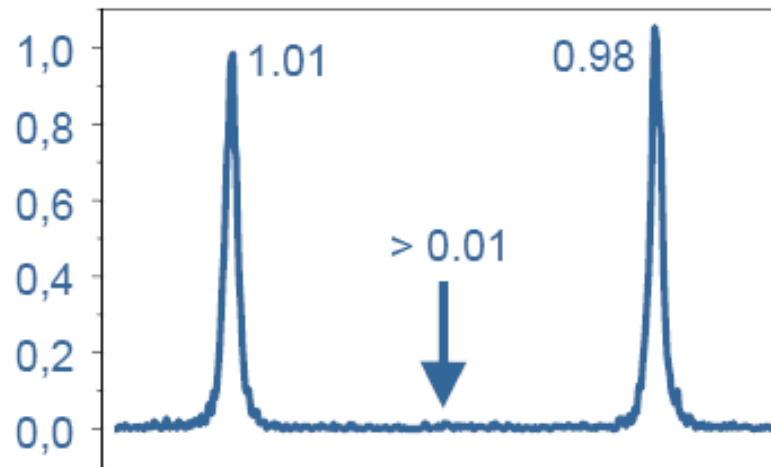


PC cavity
enables large
cavity-QD
coupling



Micropillar cavity allows for well-
collimated output: efficiency ~10%

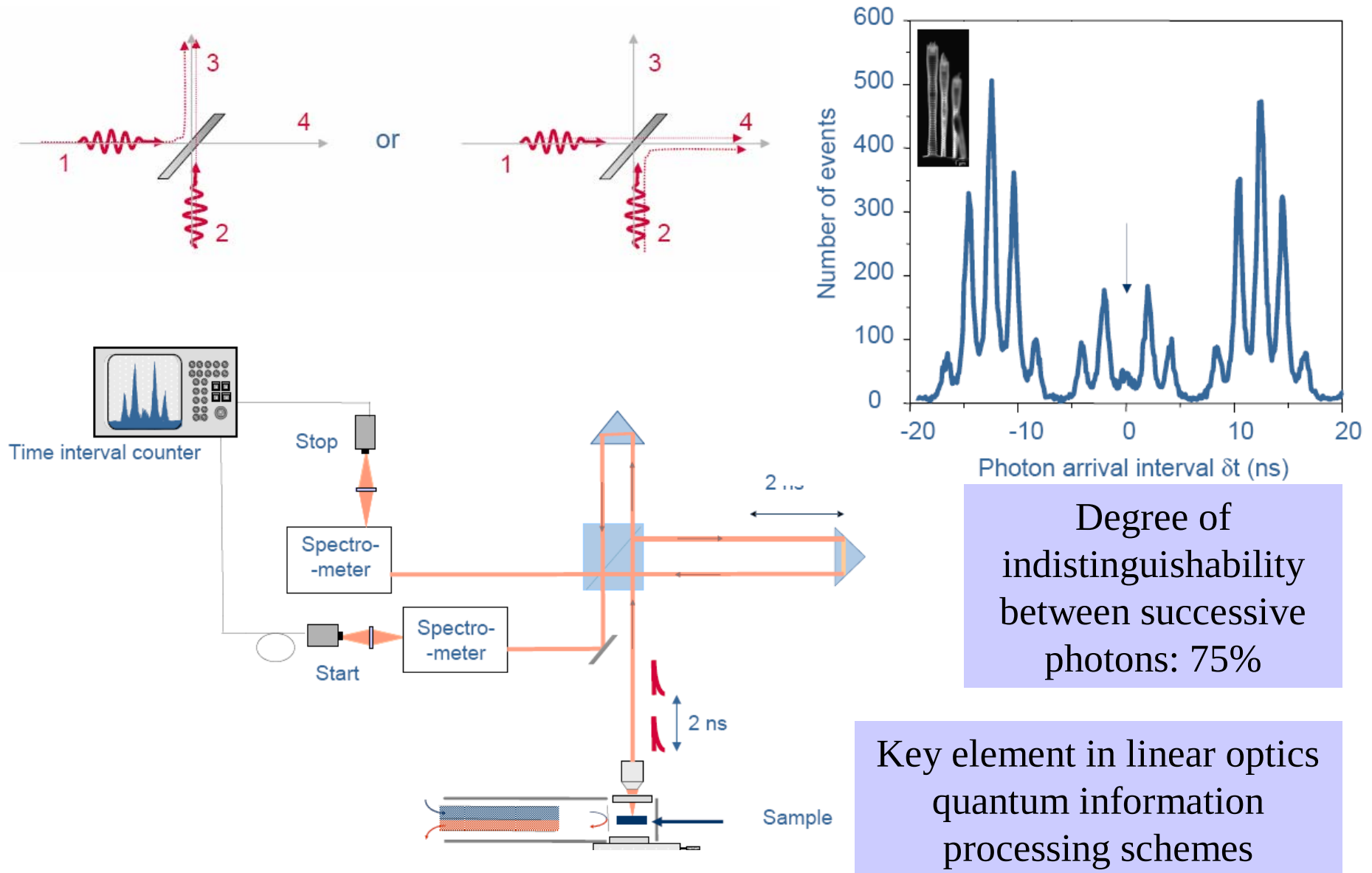
Photon correlation measurements
under pulsed excitation



Probability that any given pulse
contains 2 or more photons is 1%

Indistinguishable photons from a QD

(I. Robert-Phillip and co-workers, LPN Paris)



Future challenges

- Enhancing the efficiency of the single photon sources using cavities
- Generating indistinguishable photons from different quantum dots
- A quantum memory for the single photons