

The Abdus Salam International Centre for Theoretical Physics



() International At Energy Agency

SMR 1760 - 11

COLLEGE ON

PHYSICS OF NANO-DEVICES

10 - 21 July 2006

Coupling of Single Quantum Dots to Single Nanocavity Modes

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Coupling of Single Quantum Dots to Single Nanocavity Modes

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Outline

- 1) Brief overview of cavity quantum electrodynamics (QED)
- 2) Single quantum dot (QD) as a two-level emitter
- 3) Nano-cavities based on photonic crystal defects
- 4) Cavity-QED using a single QD embedded in a nano-cavity
- 5) Outlook

Cavity Quantum Electrodynamics (cavity-QED)

• Single two-level (anharmonic) emitter coupled to a single cavity mode is described by the Jaynes-Cummings (JC) Hamiltonian

$$H_{JC} = \hbar \omega_{eg} \sigma_{ee} + \hbar \omega_{c} a_{c}^{\dagger} a_{c} + \hbar g_{c} (\sigma_{eg} a_{c} + a_{c}^{\dagger} \sigma_{ge}) \qquad \text{An exactly} \\ \text{solvable model}$$

• ω_{eg} : emitter frequency • ω_{c} : cavity frequency • g_{c} : cavity-emitter coupling strength In all optical realizations: $g_{c} << \omega_{c} \approx \omega_{eg}$ 10^{8} 10^{15} 10^{15}



For electric-dipole coupling:

$$\mathbf{g}_{\mathbf{c}} = \left(\frac{\hbar \,\omega}{2\varepsilon_0 \varepsilon V}\right)^{1/2} \,\mathbf{q} < \mathbf{e} | \, \mathbf{\epsilon}.\mathbf{r} | \mathbf{g} >$$

Eigenstates of the JC Hamiltonian

$$\int 2\sqrt{2} g_{c} \qquad |-;2\rangle = |g; n_{c}=2\rangle - |e; n_{c}=1\rangle$$

$$|-;2\rangle = |g; n_{c}=2\rangle + |e; n_{c}=1\rangle$$

$$\begin{array}{c}
 1 > = |g; n_c=1 > - |e; n_c=0 > \\
 |+;1 > = |g; n_c=1 > + |e; n_c=0 > \\
 \omega_c = \omega_{eg} \\
 |0 > = |g; n_c=0 >
\end{array}$$

- The eigenstates of the coupled system are entangled emitter-cavity states
- The spectrum is anharmonic: the nonlinearity of the two-level emitter ensures that the coupled system is also anharmonic.

Eigenstates of the JC Hamiltonian



- The eigenstates of the coupled system are entangled emitter-cavity states
- The spectrum is anharmonic: the nonlinearity of the two-level emitter ensures that the coupled system is also anharmonic.
- Ex: A laser (ω_L) that is resonant with the |0>-|+;1> transition will be offresonant with all other transitions; the emitter-cavity molecule may act as a two-level system.
- \Rightarrow Single photon blockade!

Dissipative cavity-emitter coupling: an open quantum system

• Single two-level (anharmonic) emitter coupled to a single cavity mode is described by the non-Hermitian JC Hamiltonian

$$H_{JC} = \hbar \omega_{eg} \sigma_{ee} + \hbar \omega_c a_c^+ a_c + \hbar g_c (\sigma_{eg} a_c + a_c^+ \sigma_{ge})$$
$$-i\frac{\hbar}{2}\Gamma_{sp} \sigma_{ee} - i\frac{\hbar}{2}\kappa_c a_c^+ a_c$$

+ noise terms describing quantum jumps associated with spontaneous emission and cavity decay processes.



A simple application of cavity-QED: Purcell effect



There are two channels for radiative decay in the presence of a cavity:

- i) Spontaneous emission rate due to vacuum modes surrounding the emitter remain mostly unchanged: Γ_{sp}
- ii) The interaction with the cavity mode leads to an irreversible decay of emitter excitation; using Fermi's golden rule (valid when $\kappa_c > g_c$) gives

$$\Gamma_{Purcell} = \frac{g_c^2}{\kappa_c} \implies \text{Purcell effect: } \Gamma_{\text{Purcell}} > \Gamma_{\text{sp}}$$

⇒ Purcell regime enables many elementary tasks in quantum information processing, such as intra-conversion of spin and polarization qubits and generation of spin entanglement.

Signatures of the strong coupling regime

- <u>Frequency domain</u>: anti-crossing of the cavity and emitter like resonances as the cavity (emitter) is tuned on to resonance with the emitter (cavity).
- <u>Time-domain</u>: vacuum Rabi oscillations in the timedependence of the spontaneous emission intensity, upon initial excitation of state |e>.
- <u>Photon blockade</u>: proves that strong coupling is achieved using a single anharmonic emitter coupled to the cavity.



Cross-sectional TEM

• Each QD has a different transition energy.

Quantum dots as two-level emitters in cavity-QED

- Quantum dots do not have random thermal motion and they can easily be integrated in nano-scale cavities.
- Just like atoms, quantum dots have spontaneous emission broadened emission lines.
- Exciton transition oscillator (dipole) strength ranges from 10 to 100 (depending on the QD type) due to a collective enhancement effect.
- QDs nucleate in random locations during molecular beam epitaxy growth;
- The size and hence the emission energy of each QD is different; the standard deviation in emission energy is ~50 meV.
- ⇒ Obtaining spatial and spectral overlap between the QD and the cavity electric field is a major challenge.

Light confinement in photonic crystals

- A photonic crystal (PC) has a periodic modulation in dielectric constant and modifies the dispersion relation of electromagnetic fields to yield allowed energy bands for light propagation. For certain crystal structures there are band-gaps: light with frequency falling within these band-gaps cannot propagate in PCs.
- A defect that disrupts the periodicity in a PC confines light fields (at certain frequencies) around that defect, allowing for optical confinement on length scales ~ wavelength.
- A two dimensional (2D) PC membrane/slab defect can confine light in all three dimensions: the confinement in the third dimension is achieved thanks to total internal reflection between the high index membrane and the surrounding low-index (air) region.

<u>Note</u>: Total internal reflection and mirrors based on periodic modulation of index-of-refraction are key ingredients in all solid-state nano-cavities.

Photonic Crystal Slabs





Our slabs: 126 nm-thick GaAs a=260 nm, r/a=0.29

Photonic band gap ~ 990nm -760nm

Optimized High-Q, small V_{eff} cavities



Vertical confinement by TIR, planar confinement by PBG!

Photonic crystal fabrication





 $Q \approx 2 \times 10^4$, $V_{eff} \approx 0.7 \ (\lambda/n)^3$



Single nano-cavity in a 2D GaAs photonic crystal membrane

(t = 180 nm, a = 300 nm, r = 111 nm)

S1 defect cavity

180 nm thick GaAs membrane (with a high density of embedded QDs) obtained by etching away the Al containing layer underneath.



• Recent results (Noda *et al.*) demonstrated Q > 1,000,000 for $V_{cav} = (\lambda/n)^3 \sim 0.1 \ \mu m^3$. Predicted $g_c^2 > 10,000 \ \Gamma_{sp} \ \kappa_c$ (had that cavity contained a single QD)!

•Two major obstacles for doing QD cavity-QED:

a) Spatial overlap of a single QD and the cavity mode

b) Spectral overlap of QD and cavity resonances

Active positioning a single QD at the anti-node of a defect cavity (A. Badolato & K. Hennessy)



Tuning a cavity-mode into near-resonance with a QD line by digital etching



Cavity-QD resonance by temperature tuning



⇒ Resonant enhancement of X- luminescence intensity by a factor of 700: cavity-QD coupling $g_c \sim 80 \mu eV$ and Purcell factor $F_P \sim 40$.

Direct measurement of the Purcell effect



<u>Deterministic cavity-QED</u>: we were able to observe resonance enhancement of luminescence in all 4 structures we studied.

How well does digital etching work?



The magnitude of cavity blue-shift at each etching step depends on the cavity:

 $S1 \Rightarrow 3 \text{ nm/step}$

 $L3 \Rightarrow 2 \text{ nm/step}$

Higher-order optical excitations in a quantum dot



Polarization entangled photons from QDs



- In the absence of anisotropic exchange splitting (ΔE_{exch}) of the exciton line, cascaded emission from a QD biexciton generates polarization entangled photons at ω_{XX} and ω_X .
- When $\Delta E_{exch} > \Gamma_{spon}$, the polarization can be identified by measuring energy; the output state has <u>classical</u> polarization correlations.
 - By coupling the QD to a nano-cavity, we can ensure that the Purcell enhanced decay rate of the excitonic states well exceed ΔE_{exch} ; the output state will be entangled if the Purcell factor is > 20.

The nano-cavity must support two (nearly) degenerate modes with orthogonal polarization (Milburn et al.)

Photonic crystal nano-cavity supporting degenerate modes with orthogonal polarizations



SEM image of a calzone cavity



Polarization analysis of the two nearly-degenerate modes

- •The observed mode-splitting ranges from 2 nm to \sim 0 nm.
- The Q values range from 5000 to 10000.

Power of nanotechnology: fine-tuning of cavity modes







By spatially selective oxidation of the PC membrane surface using AFM, we can change the relative energy spacing of the two cavity modes with a resolution < 100 μ eV.

Future prospects and challenges

- Active positioning and digital etching should allow for essentially all cavity-QED tasks requiring a single emitter coupled to a single cavity mode.
- Using stacked QDs in a cavity, it is possible to solve the "spatial overlap" problem for 2 or 3 QDs in a single defect cavity. However, size non-uniformity of QDs would still limit the realisability of cavity-mediated interactions.
- Other challenges:

i) Increasing the Q value of GaAs PC cavities below 1 µm;
ii) <u>Confining light in sub-wavelength length scales</u>.

• Alternative systems for cavity-QED: microwave transitions, magnetic-dipole transitions, collective nuclear spin excitations...

- A. Badolato and K. Hennessy
- M. Winger, S. Gulde and D. Gerace, M. Atature, J. Dreiser (ETH Zurich)
- E. Hu and P. Petroff (UCSB)