



The Abdus Salam  
International Centre for Theoretical Physics



**SMR.1763- 18**

**SCHOOL and CONFERENCE  
on  
COMPLEX SYSTEMS  
and  
NONEXTENSIVE STATISTICAL MECHANICS**

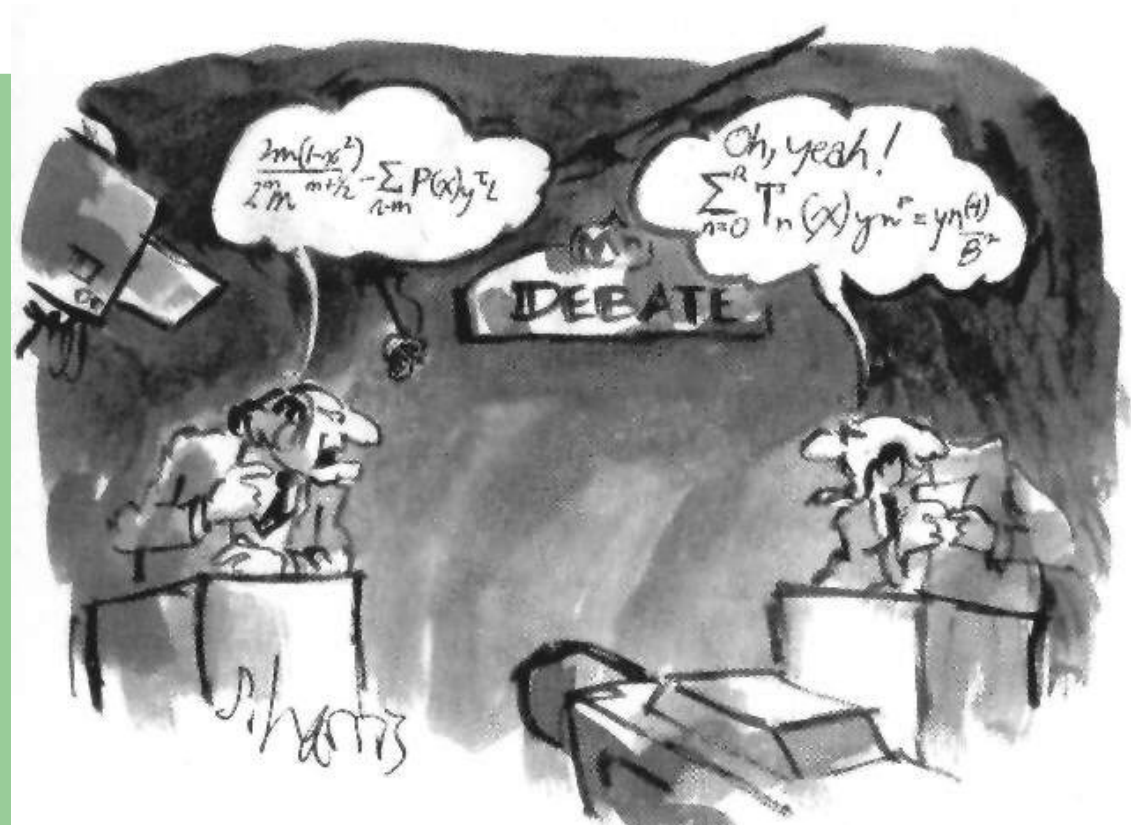
*31 July - 8 August 2006*

**Debate on Statistical Mechanics**

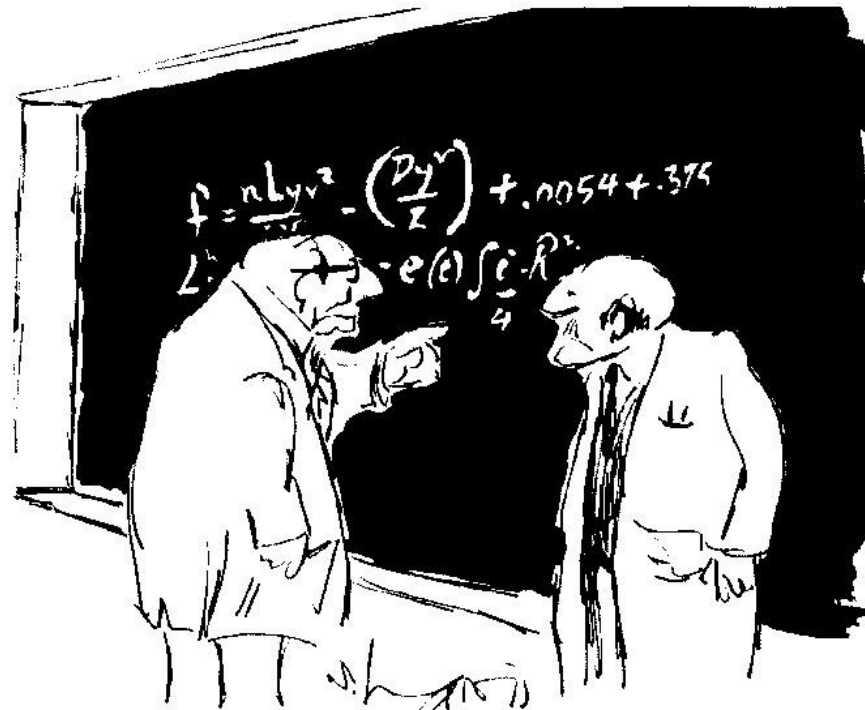
A. Rapisarda

Dipartimento di Fisica e Astronomia  
and INFN sezione di Catania,  
Università di Catania Italy  
[www.ct.infn.it/rapis](http://www.ct.infn.it/rapis)

# Debate on $S_q$ *Statistical Mechanics*



# When does a theory apply?



*"Does this apply always, sometimes, or never?"*

# beyond Boltzmann Gibbs statistical mechanics

Why do we feel the need to go beyond Boltzmann Gibbs standard statistical mechanics, a very successful theory with so many applications?

# There are several reasons

- Many systems relax very fast to equilibrium ... but many others do not ...and they live in long-standing quasi-stationary states which cannot be explained by BG statistics

Moreover BG statistical mechanics has also other limitations

In fact it is based on:

- Ergodicity and equal a priori probability for cells in phase space
- Short range interactions
- Thermodynamical limit
- Small Gaussian fluctuations

# Boltzmann Gibbs statistics does not work when one has

- Long range interacting systems
- Complex systems
- Systems at the edge of chaos
- Fractal phase space
- Ergodicity breaking and vanishing Lyapunov exponents
- Vanishing Lyapunov exponents
- SOC systems
- Turbulence
- Biological, geophysical and economical systems
- etc.

# Tsallis generalized statistics: a possible proposal

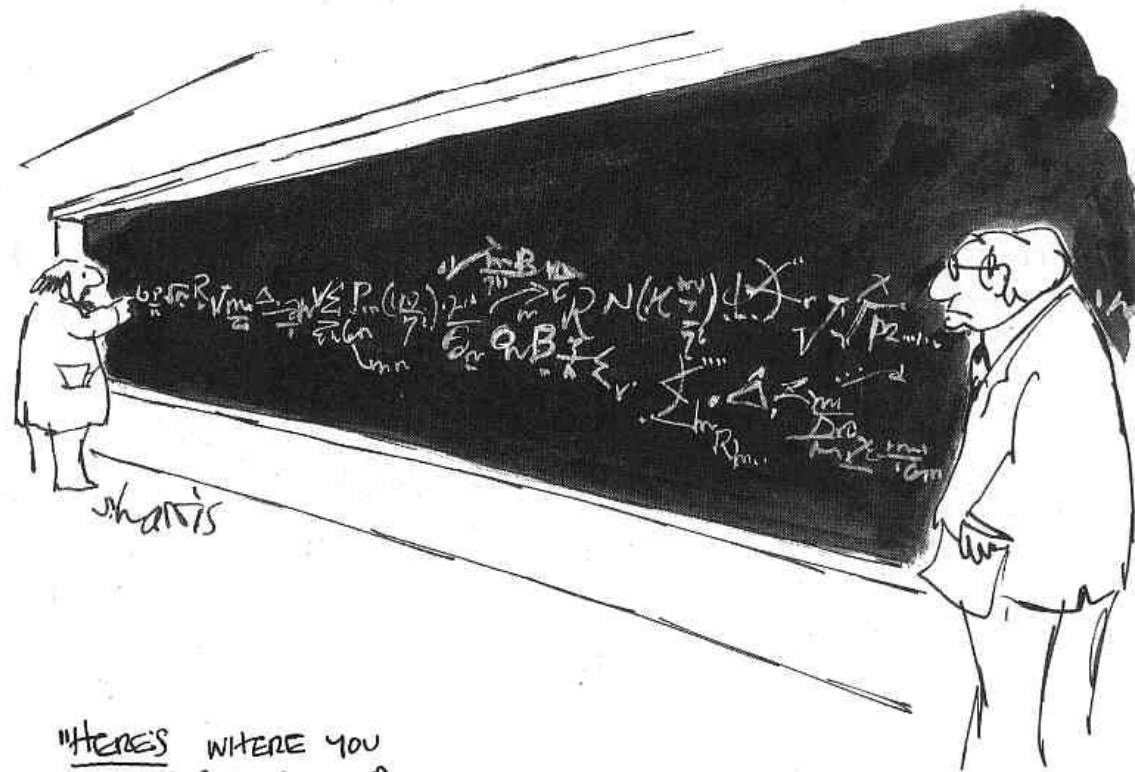
$S_q$  entropy

$$S_q = \frac{\sum_i^N p_i^q - 1}{1 - q} = \sum_i^N p_i \ln_q \frac{1}{p_i}$$

Maximized by the generalized weight

$$e_q(x) = \left[ 1 + (1 - q)x \right]^{1/(1-q)}$$

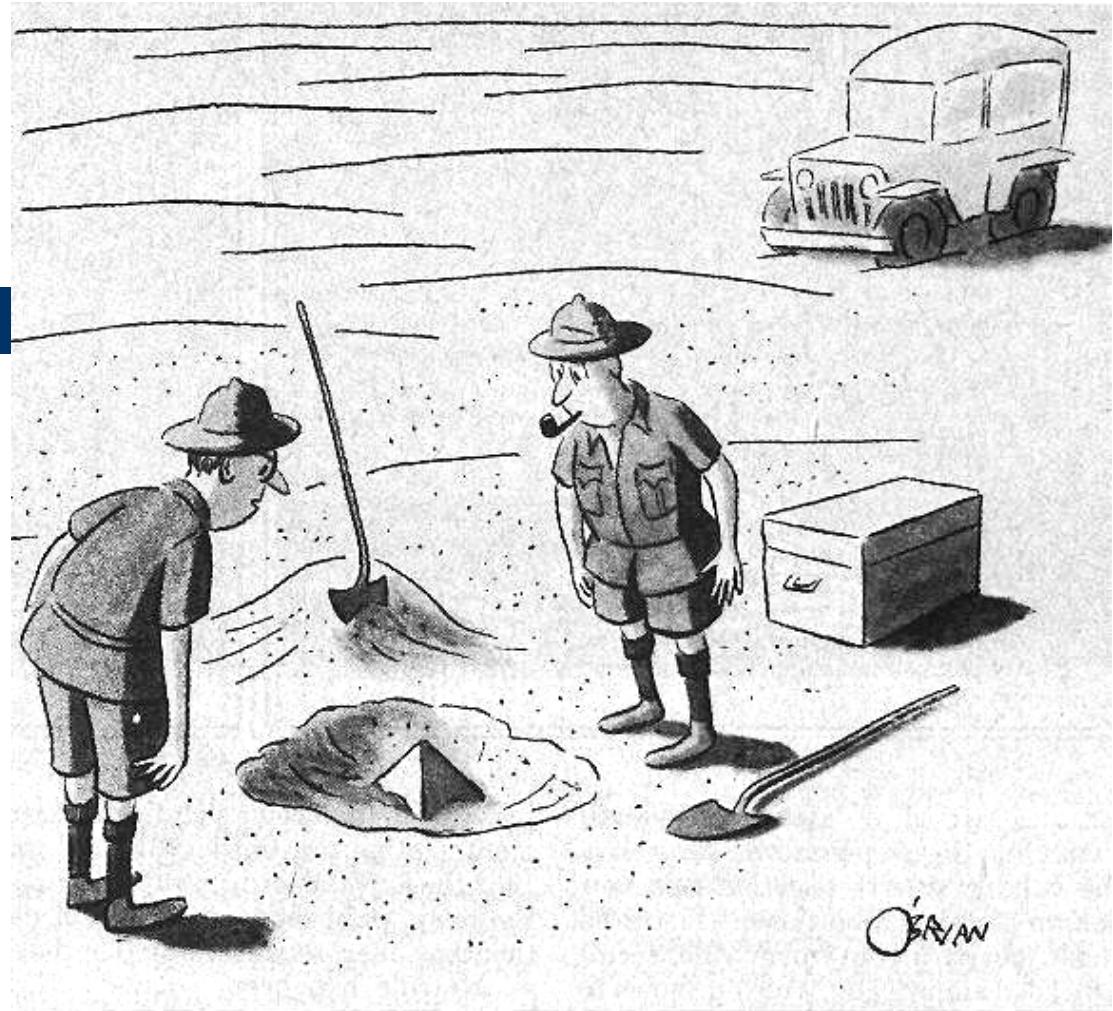
# Is this the situation....?



"HERE'S WHERE YOU  
MADE YOUR MISTAKE."^



Or this...?



*"This could be the discovery of the century. Depending, of course, on how far down it goes."*

## No a priori reasons to reject $S_q$ statistics

- It preserves the Legendre structure of thermodynamics
- Most important theorems have been generalized for  $S_q$
- It conserves simultaneously concavity, robustness, etc.
- It is extensive for special global correlations
- It reduces to BG formalism for  $q \rightarrow 1$
- The index  $q$  characterizes the underlying dynamics and, in several cases, it has already been calculated a priori
  
- The point is then:  
*“To what extent is  $S_q$  formalism supported by experiments (real data or numerical ones) or is it useful for practical applications”?*

# Where $S_q$ formalism has been successfully applied\*

\* just a few examples

|  |   |
|--|---|
| <b>Probability theory</b>                          | <b>Rigorous mathematical results on the generalization of the Central limit theorem</b><br>(Umarov, Tsallis Gell-mann, Mojano)  |
| <b>Maps at the edge of chaos</b>                   | <b>Rigorous mathematical results using renormalization group techniques</b><br>(Robledo et al)  |
| <b>Cold atoms in a dissipative optical lattice</b> | <b>Very good experimental confirmation of Lutz prediction within <math>S_q</math> statistics</b><br>(Renzoni et al Prl 2006)  |
| <b>Turbulence</b>                                  | <b>Pdf of several turbulent fluids, solar wind reproduced with a very good approximation. Very good description of defect turbulence</b><br>(Beck Cohen Swinney Daniels Bodenshatz et al) |
| <b>Long-range Hamiltonians</b>                     | <b>Good reproduction of relaxation and anomalous diffusion for QSS</b><br>(Rapisarda, Anteneodo, Mojano, Pluchino, Tamarit )  |
| <b>Biological systems</b>                          | <b>Very good description of anomalous diffusion and decay of correlation for Hydra viridissima.</b><br>(Upadhyaya et al)  |
| <b>Econophysics</b>                                | <b>Very good description of the volatility smile</b><br>(Borland)   |
| <b>Computational physics</b>                       | <b>Important applications for Optimization algorithms</b><br>(Straub, Andriociai)   |
| <b>Signal and image processing</b>                 | <b>Very performant techniques for image reconstruction</b><br>(de Albuquerque , Ben Hanza et al)  |

# q-product

$$x \otimes_q y \equiv \left[ x^{1-q} + y^{1-q} - 1 \right]^{\frac{1}{1-q}}$$

*Properties :*

i)  $x \otimes_1 y = x y$

ii)  $\ln_q(x \otimes_q y) = \ln_q x + \ln_q y$

[whereas  $\ln_q(x y) = \ln_q x + \ln_q y + (1-q)(\ln_q x)(\ln_q y)$ ]

[L. Nivanen, A. Le Mehaute and Q.A. Wang, Rep. Math. Phys. 52, 437 (2003);  
E.P. Borges, Physica A 340, 95 (2004)]

## q - GENERALIZED CENTRAL LIMIT THEOREM: (mathematical proof)

S. Umarov, C.T. and S. Steinberg [cond-mat/0603593]

q-Fourier transform:

$$F_q[f](\xi) \equiv \int_{-\infty}^{\infty} e_q^{ix\xi} \otimes_q f(x) dx = \int_{-\infty}^{\infty} e_q^{\frac{ix\xi}{[f(x)]^{1-q}}} f(x) dx \quad (\text{nonlinear!})$$

q-correlation:

Two random variables  $X$  [with density  $f_X(x)$ ] and  $Y$  [with density  $f_Y(y)$ ] are said  $q$ -correlated if

$$F_q[X+Y](\xi) = F_q[X](\xi) \otimes_q F_q[Y](\xi),$$

i.e., if

$$\int_{-\infty}^{\infty} dz e_q^{iz\xi} \otimes_q f_{X+Y}(z) = \left[ \int_{-\infty}^{\infty} dx e_q^{ix\xi} \otimes_q f_X(x) \right] \otimes_q \left[ \int_{-\infty}^{\infty} dy e_q^{iy\xi} \otimes_q f_Y(y) \right],$$

$$\text{with } f_{X+Y}(z) = \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy h(x, y) \delta(x + y - z) = \int_{-\infty}^{\infty} dx h(x, z - x) = \int_{-\infty}^{\infty} dy h(z - y, y)$$

where  $h(x, y)$  is the joint density.

$$\left( \begin{array}{lll} q\text{-correlation means} & \text{independence} & \text{if } q=1, \text{ i.e., } h(x, y) = f_X(x) f_Y(y) \\ & \text{global correlation} & \text{if } q \neq 1, \text{ hence } h(x, y) \neq f_X(x) f_Y(y) \end{array} \right)$$

**CENTRAL LIMIT THEOREMS:  $N^{1/[\alpha(2-q)]}$  - SCALED ATTRACTOR  $\mathbb{F}(x)$  WHEN SUMMING  $N \rightarrow \infty$   
 $q$  - CORRELATED IDENTICAL RANDOM VARIABLES WITH SYMMETRIC DISTRIBUTION  $f(x)$**

|   | $q = 1$ [independent]   | $q \neq 1$ (i.e., $Q \equiv 2q - 1 \neq 1$ ) [globally correlated]  |
|---|---|---|
| $\sigma_Q < \infty$<br>$(\alpha = 2)$               | $\mathbb{F}(x) = \text{Gaussian } G(x)$ ,<br>with same $\sigma_1$ of $f(x)$<br><br>Classic CLT  | $\mathbb{F}(x) = G_{\frac{3q-1}{q+1}}(x) \equiv \frac{3q-1}{q+1}$ - Gaussian,<br>with same $\sigma_Q \left[ \equiv \int dx x^2 [f(x)]^Q / \int dx [f(x)]^Q \right]$ of $f(x)$<br>$G_{\frac{3q-1}{q+1}}(x) \begin{cases} \approx G(x) & \text{if }  x  \ll x_c(q, 2) \\ \sim f(x) \sim C_q /  x ^{(q+1)/(q-1)} & \text{if }  x  \gg x_c(q, 2) \end{cases}$<br>with $\lim_{q \rightarrow 1} x_c(q, 2) = \infty$<br>S. Umarov, C. T. and S. Steinberg (2006) [cond-mat/0603593]  |
| $\sigma_Q \rightarrow \infty$<br>$(0 < \alpha < 2)$ | $\mathbb{F}(x) = \text{Levy distribution } L_\alpha(x)$ ,<br>with same $ x  \rightarrow \infty$<br>asymptotic behavior<br>$L_\alpha(x) \begin{cases} \approx G(x) & \text{if }  x  \ll x_c(1, \alpha) \\ \sim f(x) \sim C_\alpha /  x ^{1+\alpha} & \text{if }  x  \gg x_c(1, \alpha) \end{cases}$<br>with $\lim_{\alpha \rightarrow 2} x_c(1, \alpha) = \infty$<br><br>Levy-Gnedenko CLT | $\mathbb{F}(x) = L_{\frac{2\alpha q - \alpha + 3}{\alpha + 1}, \alpha}$ stable distribution,<br>with $L_{\frac{2\alpha q - \alpha + 3}{\alpha + 1}, \alpha} \sim f(x) \sim C_{q, \alpha}^{(L)} /  x ^{(1+\alpha)/(1+\alpha q - \alpha)}$<br>or<br>$\mathbb{F}(x) = L_{\frac{\alpha q + q - 1}{\alpha + q - 1}, \alpha}$ stable distribution,<br>with $L_{\frac{\alpha q + q - 1}{\alpha + q - 1}, \alpha} \sim f(x) \sim C_{q, \alpha}^{(*)} /  x ^{2(\alpha + q - 1)/\alpha(q - 1)}$<br>S. Umarov, C. T., M. Gell-Mann<br>and S. Steinberg (2006)<br>[cond-mat/0606038] and [cond-mat/0606040] |

# Logistic map at the edge of chaos

Entropy growth is linear only for  $q^* = 0.2425\dots$

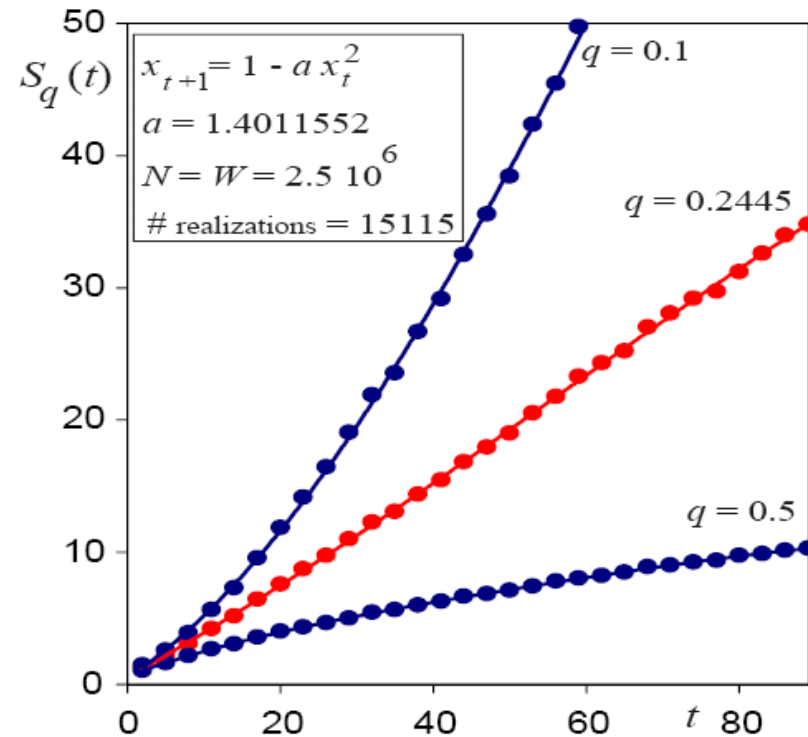
$$\frac{1}{1 - q^*} = \frac{\ln \alpha_F}{\ln 2}$$

Generalized Pesin-like theorem

$$K_q = \lambda_q$$

$$\lambda_q = \frac{1}{1 - q}$$

(weak chaos, i.e., zero Lyapunov exponent)



C. T. , A.R. Plastino and W.-M. Zheng, Chaos, Solitons & Fractals 8, 885 (1997)  
 M.L. Lyra and C. T. , Phys. Rev. Lett. 80, 53 (1998)  
 V. Latora, M. Baranger, A. Rapisarda and C. T. , Phys. Lett. A 273, 97 (2000)  
 E.P. Borges, C. T. , G.F.J. Ananos and P.M.C. Oliveira, Phys. Rev. Lett. 89, 254103 (2002)  
 F. Baldovin and A. Robledo, Phys. Rev. E 66, R045104 (2002) and 69, R045202 (2004)  
 G.F.J. Ananos and C. T. , Phys. Rev. Lett. 93, 020601 (2004)  
 E. Mayoral and A. Robledo, Phys. Rev. E 72, 026209 (2005), and references therein

Feigenbaum constant has been measured in many real experiments

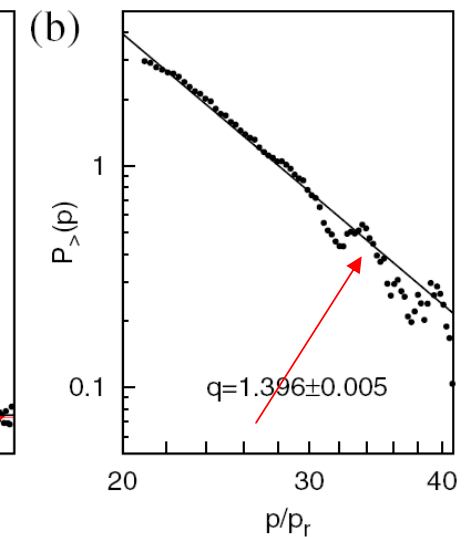
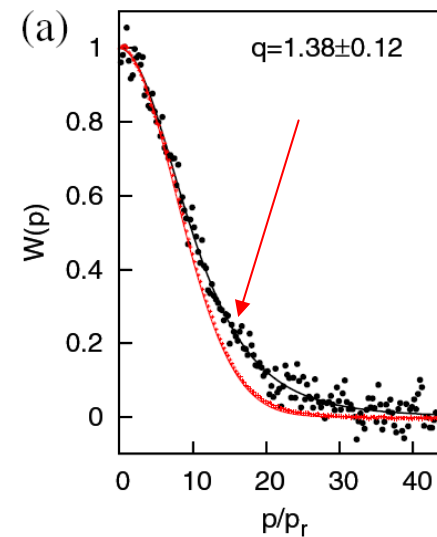
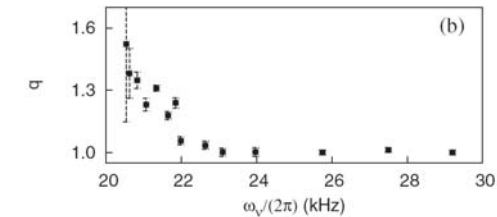
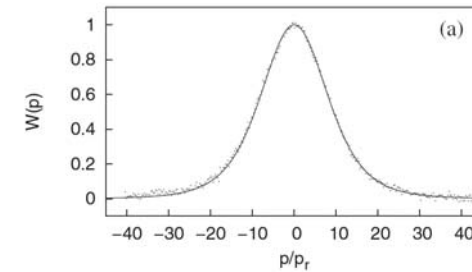
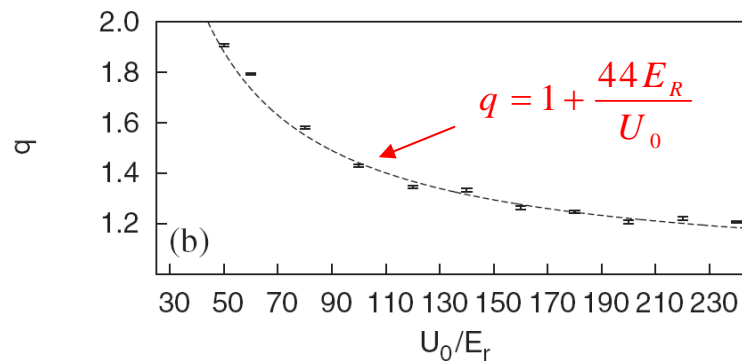
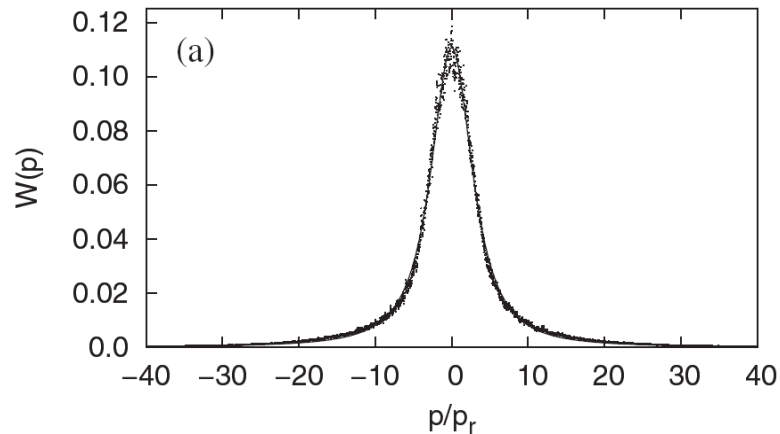
$$\alpha_F = 2.5029078\dots$$

# Cold atoms in a dissipative optical lattice

experimental verification by P. Douglas, S. Bergamini and F. Renzoni, Phys Rev Lett **96**, 110601 (2006)

Theoretical prediction by

E.Lutz PRA **67** (2003) 051402(R)



(Computational verification:  
quantum Monte Carlo simulations)

(Experimental verification)



*The solution of*

$$\frac{\partial p(x,t)}{\partial t} = D \frac{\partial^2 [p(x,t)]^{2-q}}{\partial x^2} \quad [p(x,0) = \delta(0)] \quad (q < 3)$$

*is given by*

$$p(x,t) \propto \left[ 1 + (1-q) x^2 / (\Gamma t)^{2/(3-q)} \right]^{1/(1-q)} \equiv e_q^{-x^2 / (\Gamma t)^{2/(3-q)}} \quad (\Gamma \propto D)$$

*hence*

$x^2$  *scales like*  $t^\gamma$  (e.g.,  $\langle x^2 \rangle \propto t^\gamma$ )

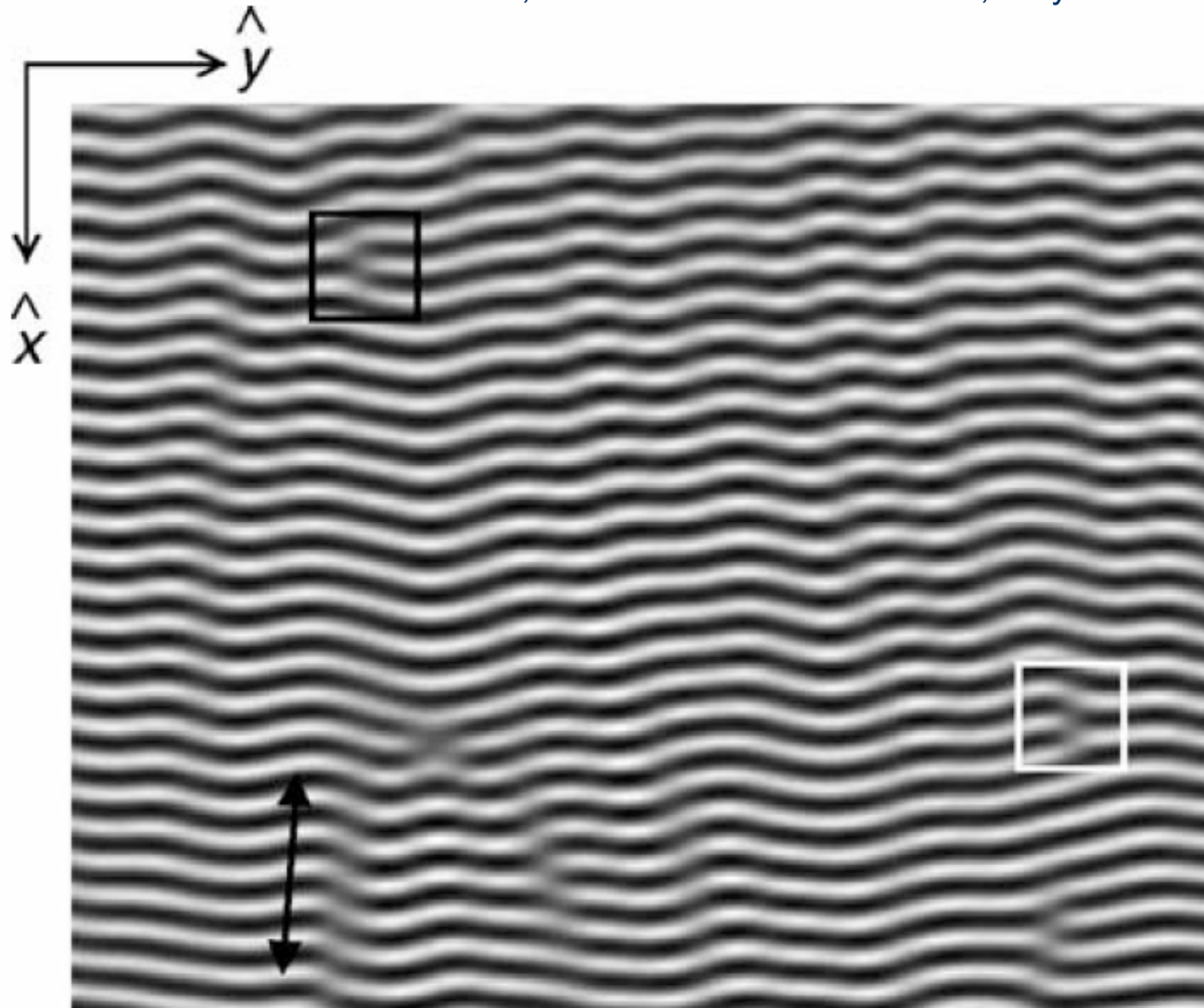
*with*

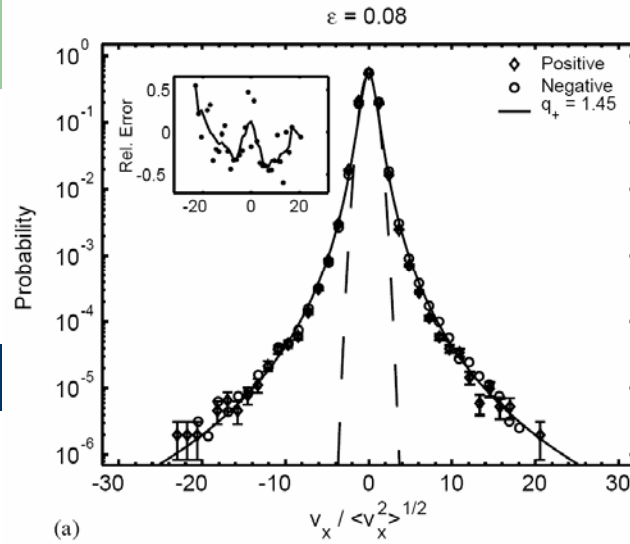
$$\gamma = \frac{2}{3-q}$$

C.Tsallis and D.J. Bukman, Phys Rev E **54**, R2197 (1996)

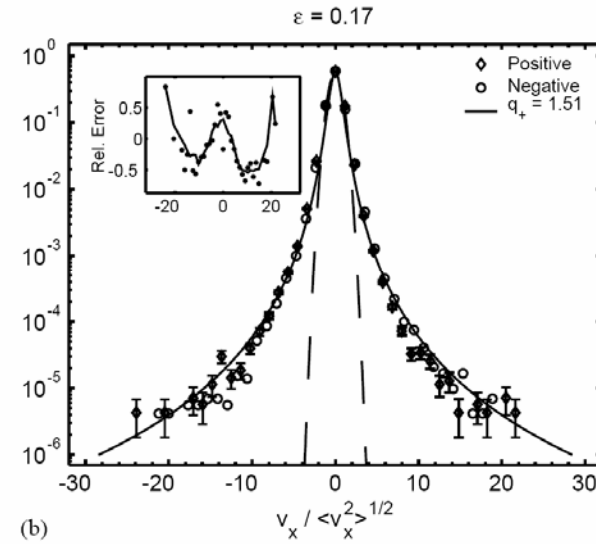
# Defect Turbulence

K.E. Daniels, C. Beck and E. Bodenschatz, Physica D 193, 208 (2004)

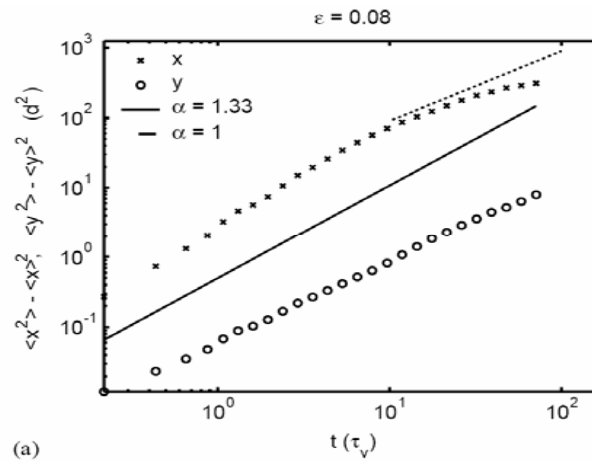




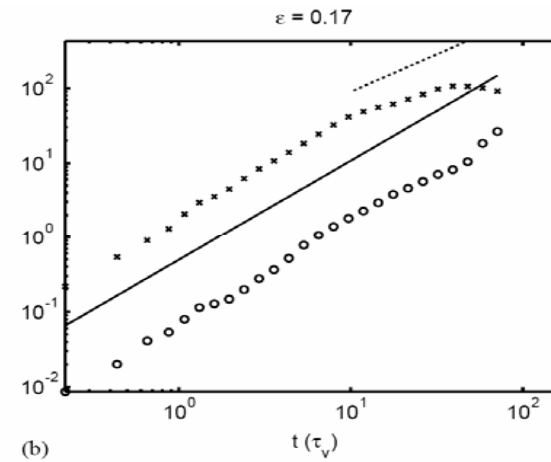
(a)



(b)



(a)



(b)

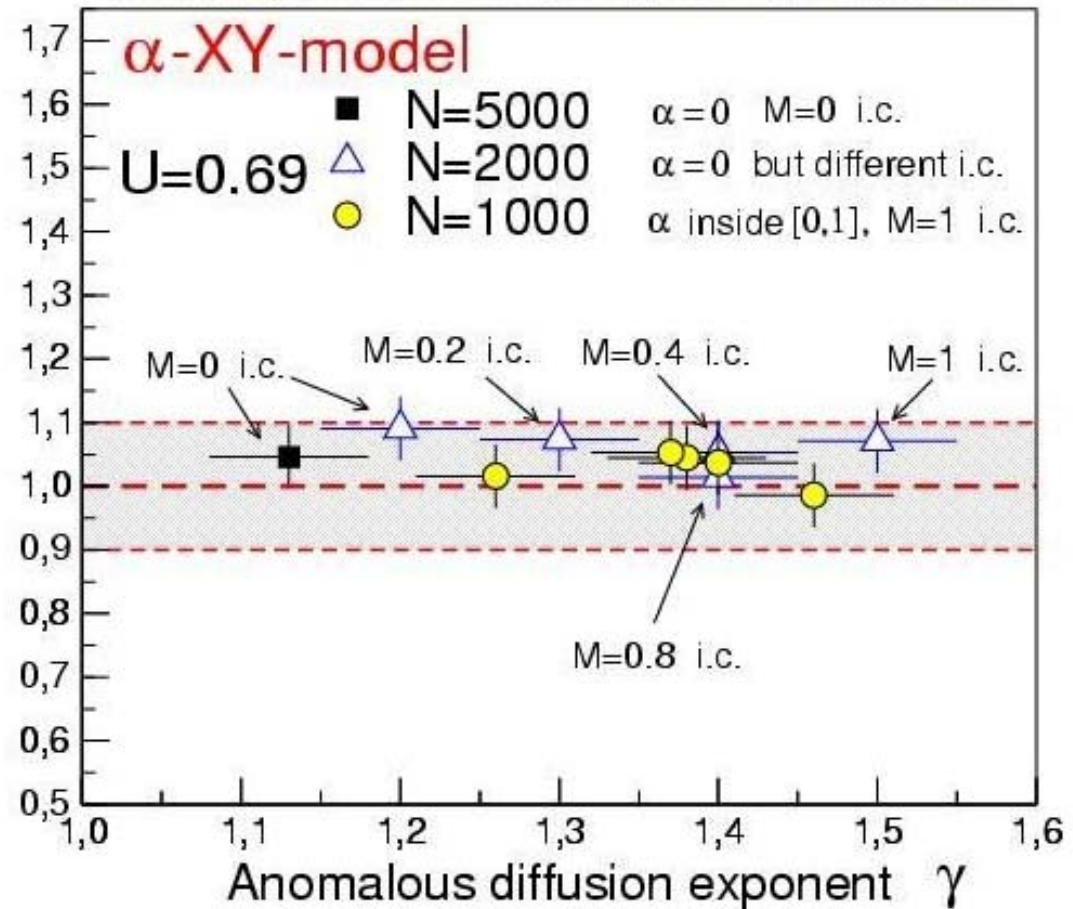
*$q \approx 1.5$  and  $\gamma \approx 4/3$  are consistent with  $\gamma = \frac{2}{3-q}$*

# Long-range Hamiltonian systems

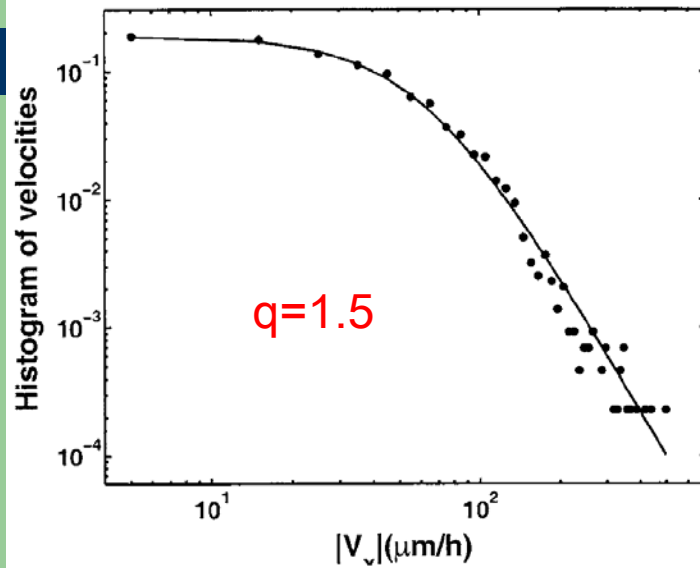
Anomalous diffusion vs q-exponential decay of correlation function

$$\frac{\gamma}{2/(3-q)}$$

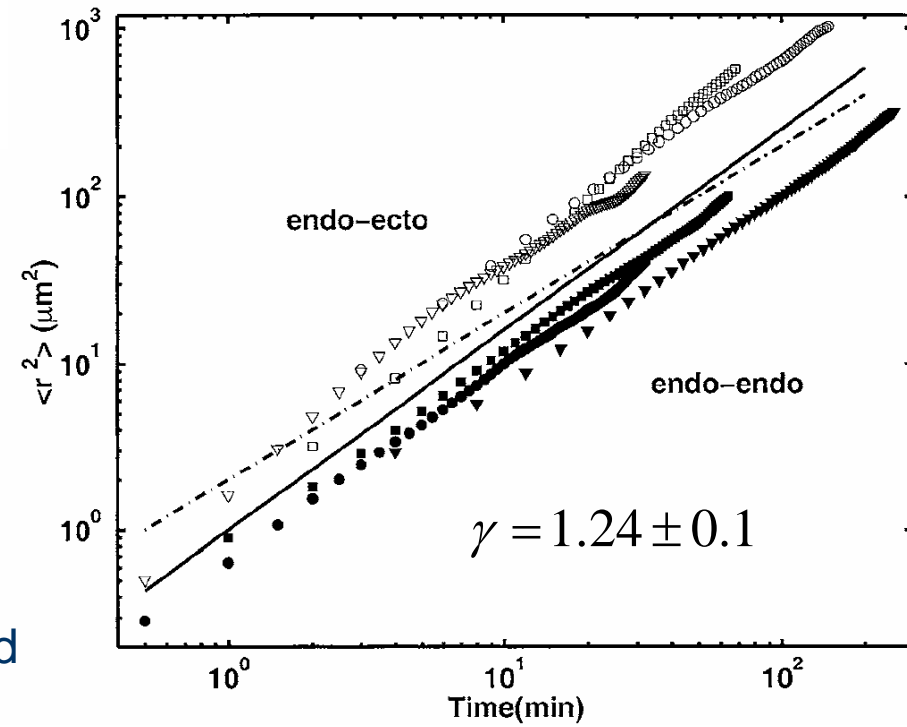
Rapisarda and Pluchino,  
Europhysics News 36 (2005) 202



# Anomalous diffusion for *Hydra viridissima*



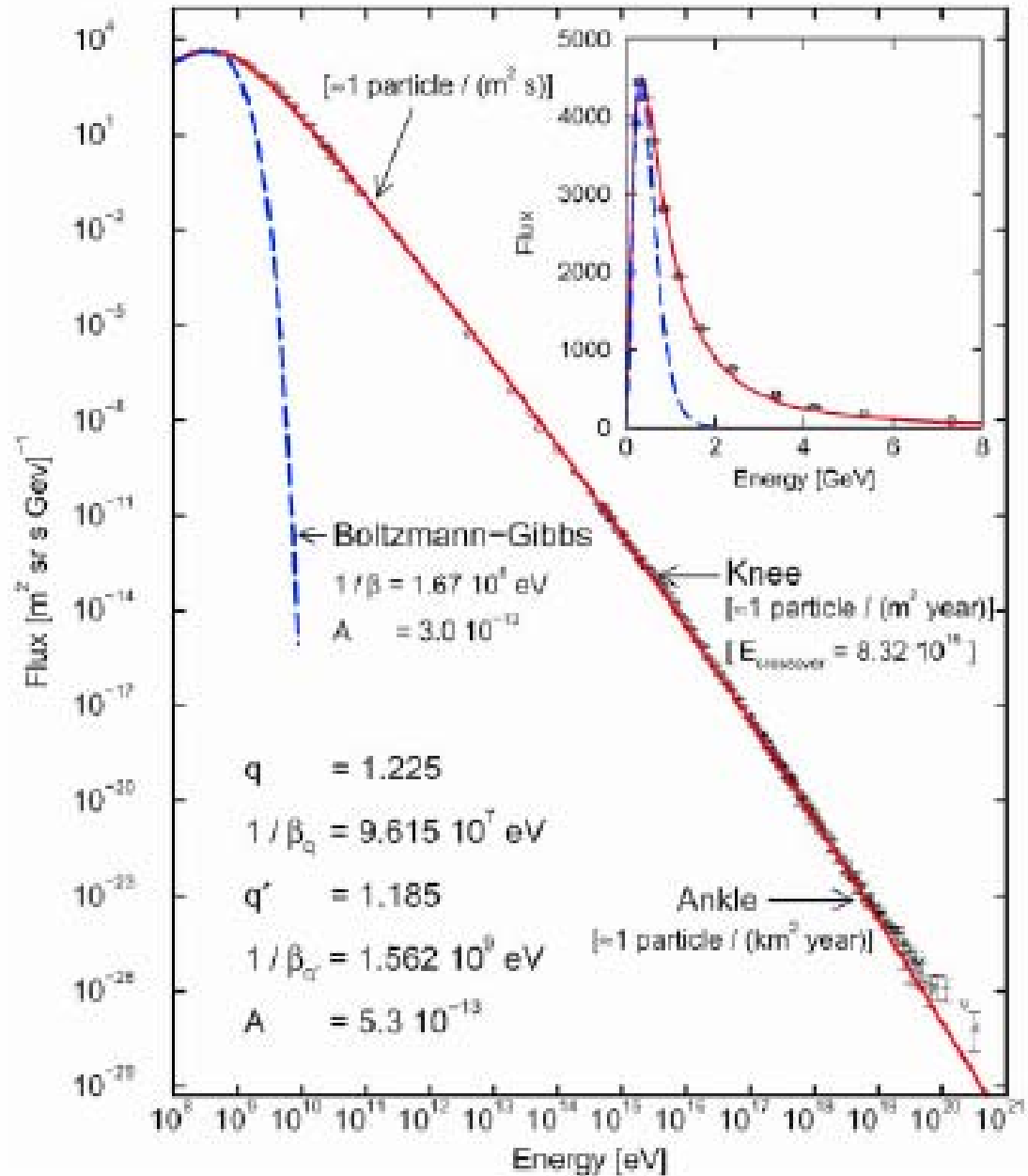
So  $\gamma = \frac{2}{3-q}$  is confirmed



A. Upadhyaya, J.-P. Rieu, J.A. Glazier and Y. Sawada  
*Physica A* 293, 549 (2001)

# Cosmic rays

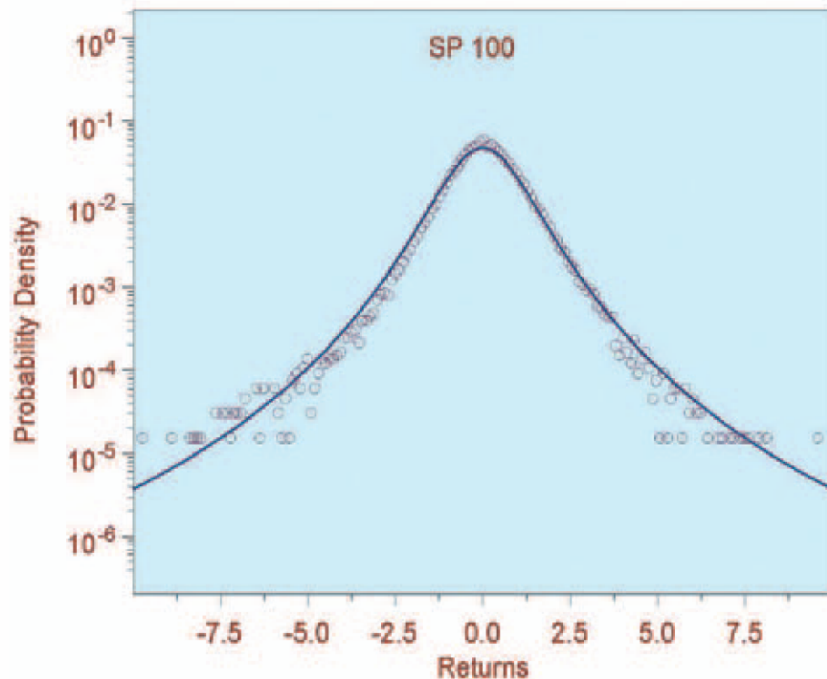
Tsallis Anjos Borges,  
PLA 310,372 (2003)



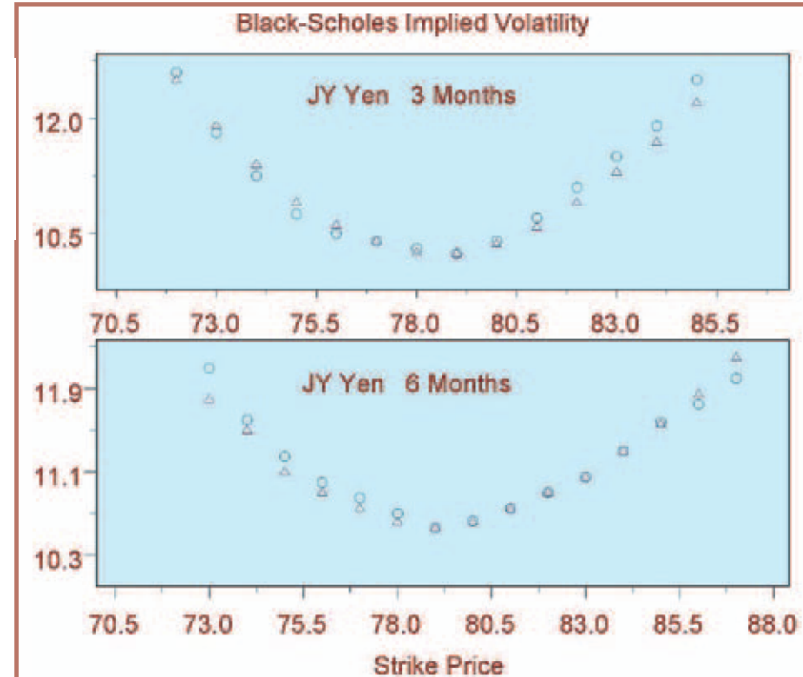
# Econophysics

## q-GENERALIZED BLACK-SCHOLES EQUATION:

L Borland, Phys Rev Lett **89**, 098701 (2002), and Quantitative Finance **2**, 415 (2002)  
L Borland and J-P Bouchaud, cond-mat/0403022 (2004)  
L Borland, Europhys News **36**, 228 (2005)  
See also H Sakaguchi, J Phys Soc Jpn **70**, 3247 (2001)  
C Anteneodo and CT, J Math Phys **44**, 5194 (2003)



▲ **Fig.2:** The empirical distribution of daily returns from the stocks comprising the SP 100 (red) is fit very well by a  $q$ -Gaussian with  $q = 1.4$  (blue).



▲ **Fig.3:** Theoretical implied Black-Scholes volatilities from the  $q = 1.4$  model (triangles) match empirical ones (circles) very well, across all strikes and for different times to expiration.

[REMARK : Student  $t$ -distributions are the particular case of  $q$ -Gaussians when  $q = \frac{n+3}{n+1}$  with  $n$  integer]



Statistical  
mechanics

Probability  
theory

Dynamics



## A few concluding remarks

- $S_q$  statistics is a very useful and elegant theory.
- It has also experimental and numerical support.
- There are still some open problems and further work is certainly needed.
- Future investigations will clarify its limitations and how fundamental it is.