

Collapse Dynamics, Singularity Formation and Regularization

Jens Juul Rasmussen

jens.juul.rasmussen@risoe.dk

Risø National Laboratory,

Department of Optics and Plasma Research, OPL-128,
DK-4000 Roskilde, Denmark

Motivations

- Collapse physics is fascinating
- Collapse related to singularity formation
- Mathematically challenging
- It is real physics
- Important in many disparate fields
- Numerous experimental observations - of its effects

Specifically consider the collapse dynamics governed by NLS:

Nonlinear Schrödinger Equation

Various physical contexts:

- Nonlinear optics - medium with amplitude dependent index of refraction – Self focusing
- Water waves at the free surface
- Plasma waves – Langmuir collapse (*Zakharov Sov. Phys. JETP* **35**, 908 (1972))
- Bose-Einstein Condensate (BEC), often referred as the Gross-Pitaevskii (G-P) equation
- Quantum Electrodynamics EM wave collapse in Radiation background (*Marklund et al PRL* **91**, 163601 (2003))
-

Some Reviews:

.....

E.A. Kuznetsov, A.m. Rubenchik and V.E. Zakharov, *Phys. Rep.* **142**, 103 (1986)

J. Juul Rasmussen and K. Rypdal, *Phys. Scr.* **33**, 481 (1986).

E.A. Kuznetsov, *CHAOS* **6**, 381 (1996).

L. Bergé, *Phys. Rep.* **303**, 259 (1998).

G. Fibich and G.C. Papanicolaou. *SIAM J. Appl. Math.* **60**, 183 (1999).

C. Sulem and P.-L. Sulem, *The Nonlinear Schrödinger Equation: Self-focusing and Wave Collapse* (Springer, Berlin 1999).

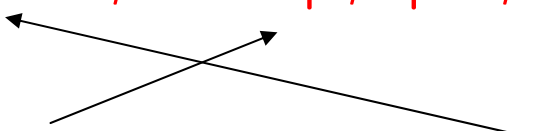
.....

What is collapse – blow-up?

$$\exists t_0 < \infty \quad \|\nabla\psi\| \rightarrow \infty \text{ for } t \rightarrow t_0 \quad \|f\| \equiv (\int |f|^2 d\vec{r})^{1/2}$$

Development of a singularity, infinite amplitude or gradient locally

Nonlinear Schrödinger equation:

$$i \frac{\partial \psi}{\partial t} + \nabla^2 \psi + |\psi|^2 \psi = 0$$


Collapse when nonlinearity dominates over dispersion

Conserved quantities:

Power - Number of particles: $N = \int |\psi|^2 d\vec{r}$,

Hamiltonian: $H = \int (|\nabla\psi|^2 - \frac{1}{2}|\psi|^4) d\vec{r}$

Collapse

Collapse: mean squared width $\langle r^2 \rangle = N^{-1} \int r^2 |\psi|^2 d\vec{r} \rightarrow 0$ for $t \rightarrow t_0$

Virial theorem: $\partial_t^2 \langle r^2 \rangle = \frac{1}{N} 2 (4H - (D - 2) \int |\psi|^4 d\vec{r})$

For dimension $D \geq 2$ collapse ensured for $H < 0$

Collapse implies blow-up, e.i., sufficient condition for blow-up:

$$N \leq \langle r^2 \rangle \int |\nabla \psi|^2 d\vec{r}, \text{ vanishing of } \langle r^2 \rangle \Rightarrow \text{blow-up of } \|\nabla \psi\|$$

Collapse implies concentration of the energy in “one point”

Physics do not “allow” singularities some regularizations take place

NOTE: Blow-up usually appears before collapse!

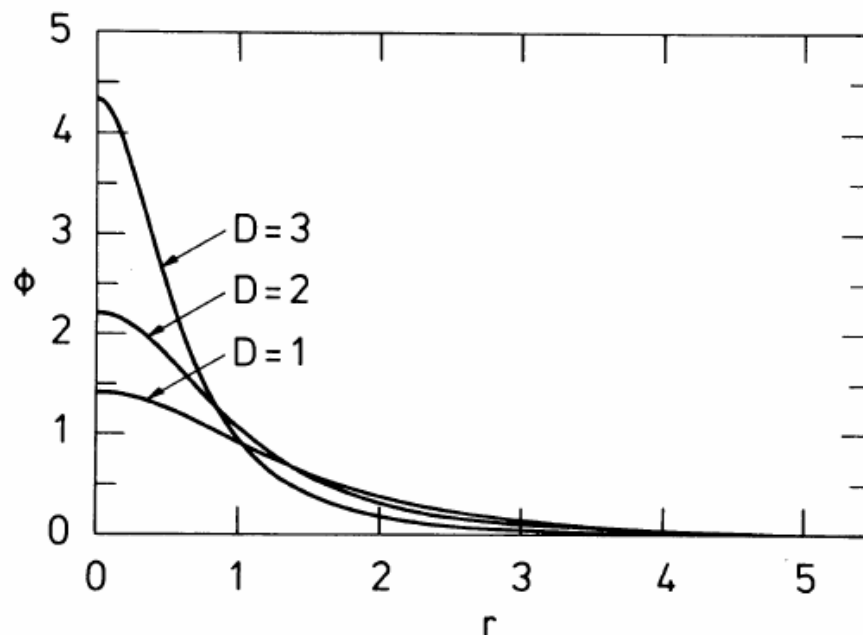
Collapse criteria

$D = 1$ NO collapse,
stable soliton solution (IST)

$D = 2$: $N \geq N_s$ necessary
and sufficient condition
 $N_s = 11.68$

$D = 3$: $H \leq 0$
sufficient condition
sharper criterion related to H_s

Kuznetsov *et al.* *Physica D* **87**, 273 (1995)

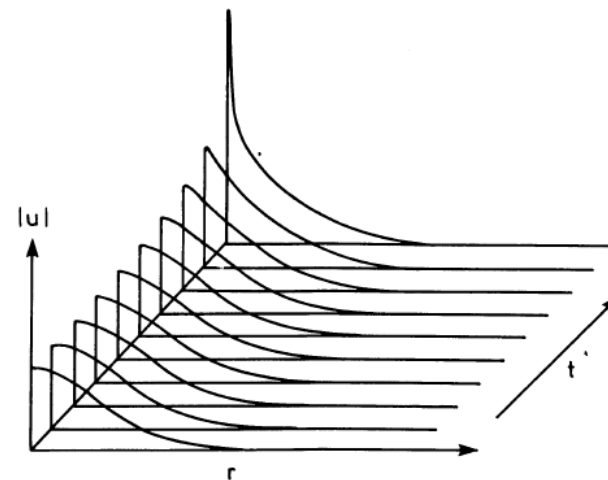
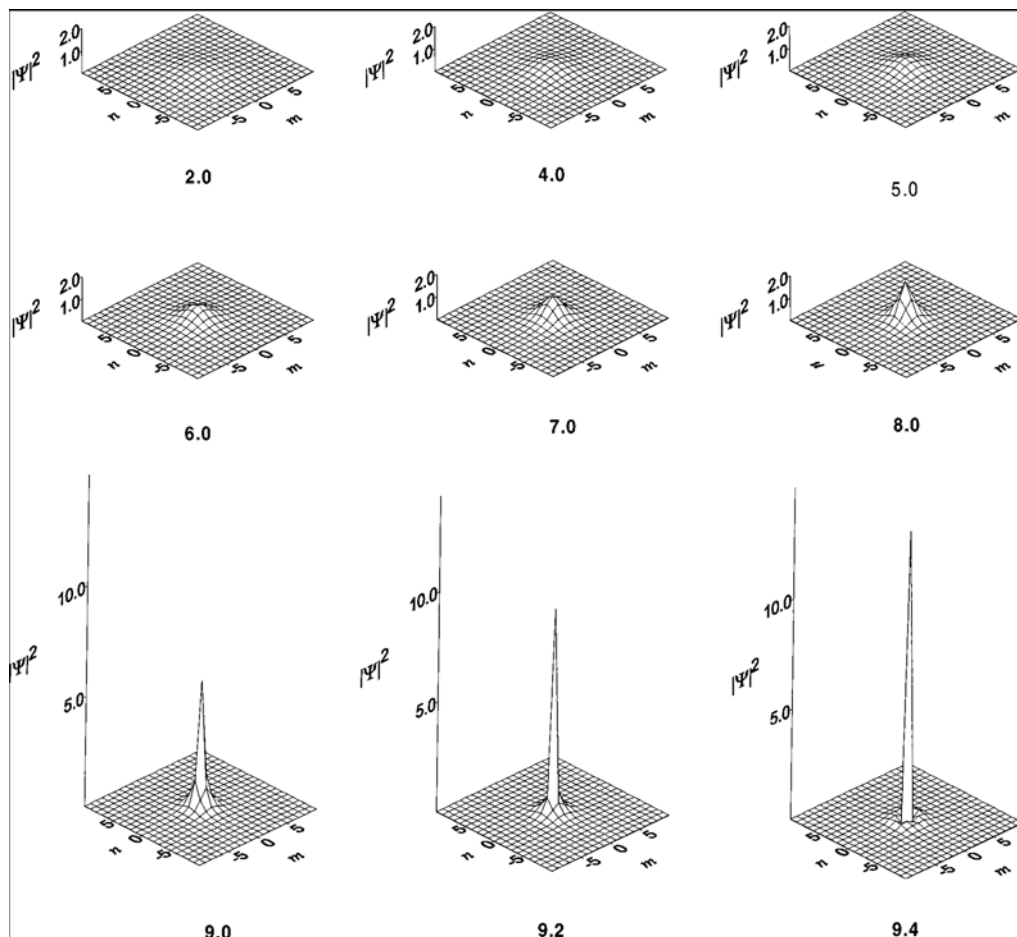


Ground-state solutions of NLS

$D=2$: “Townes profile”

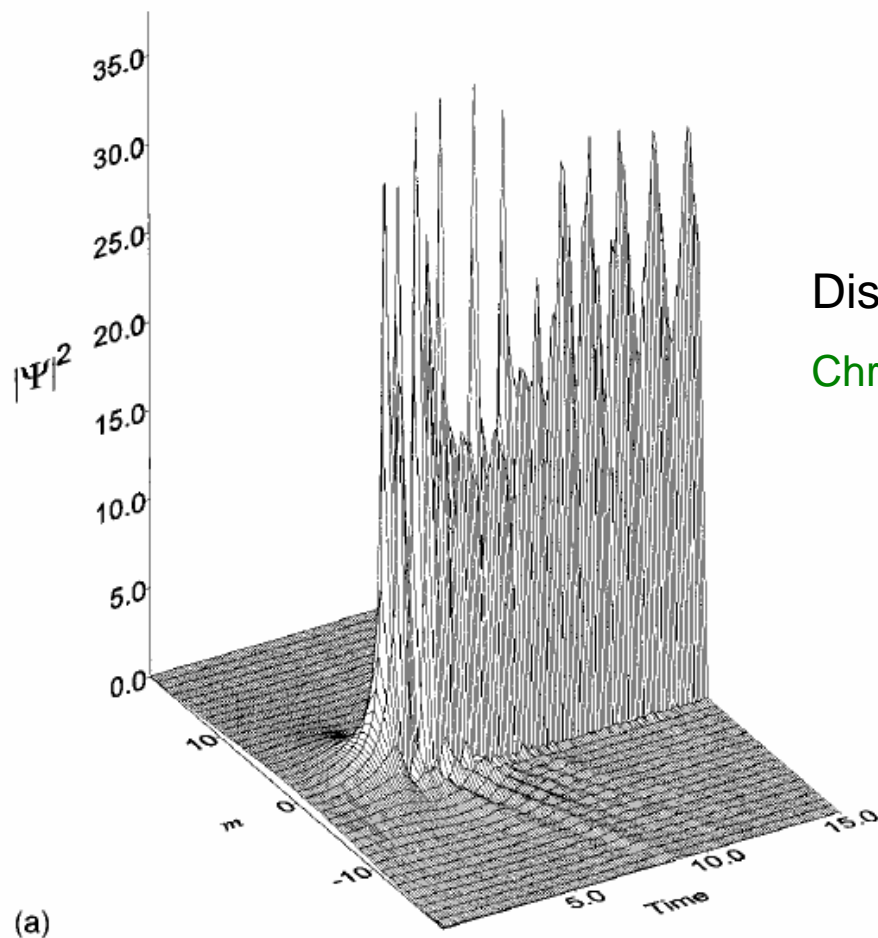
(Chiao *et al.*, *Phys. Rev. Lett.* **13**,479 (1964))

Collapse/blow-up in 2D NLS



NO global self-similar description
NO few parameter trial function
governs this evolution, e.i.,
standard variational approach
with Gaussian trial function
cannot describe the full evolution

Singularity regularization



Discrete 2D NLS

Christiansen *et al.* Phys. Rev. B **57**, 11303 (1998)

Filamentation of fs laser pulse in air

Femtosecond LIDAR



Collapse types

Weak collapse : The power/number of particles , N_{spike} tends to 0 in the singular spike

Typical for $D = 3$ (here strong collapse is unstable)

Strong collapse : The power/number of particles , N_{spike} tends to a finite value in the singular spike

For $D = 2$ always strong collapse : $N_{\text{spike}} = N_s$

Zakharov and Kuznetsov, *Sov. Phys. JETP* **64**,773 (1986)

Singularity Regularization

Passive : “missing physics” in the model :

Dissipations

Extra effects, extra dimension...

Nonlocal interactions ([Bang et al *Phys Rev E* **66**, 046619 \(2002\)](#))

Active : external forcing / control.

Potential landscape, optical lattice

Self-focusing of light in nonlinear materials

$$i\partial_z\psi + \nabla_{\perp}^2\psi - \beta\partial_t^2\psi + |\psi|^2\psi = \delta\rho\psi + \mathcal{R}(\psi)$$

$$\partial_t\delta\rho = \sigma_K\rho_0|\psi|^{2K} + \sigma|\psi|^2\delta\rho - \frac{1}{\tau}\delta\rho$$

β group velocity dispersion:

normal $\beta > 0$, anomalous $\beta < 0$

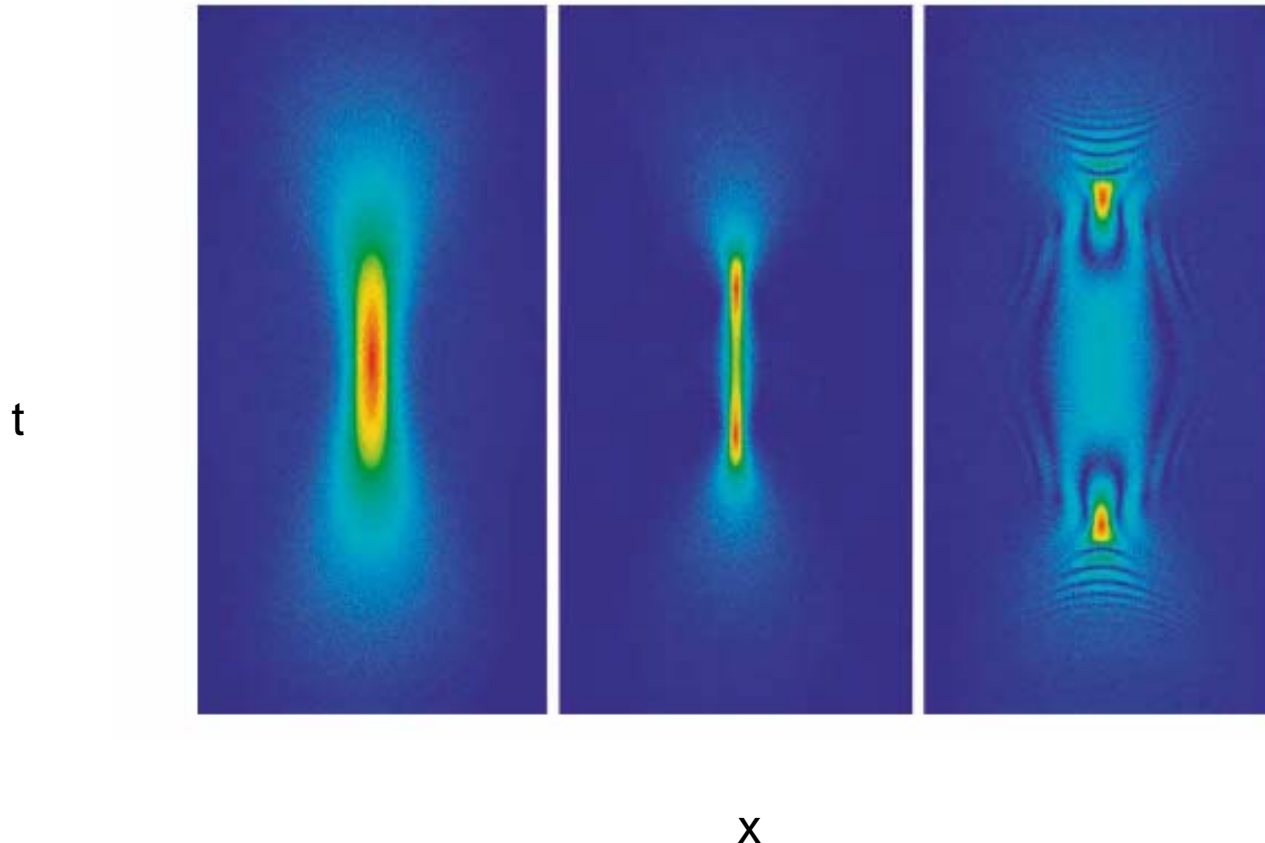
$\delta\rho$ plasma density, focused beam burns a channel

$\mathcal{R}(\psi)$ Raman scattering etc.

Plasma defocusing, self-guiding

Skupin and Bergé *Physica D* **220**, 14 (2006)

Normal GVD



Focusing in the transverse plane (x,y), defocusing in time

Germaschewski *et al.*, *Physica D* **151**, 175 (2001)

Dissipation

$$i\partial_t\psi + \nabla^2\psi + |\psi|^2\psi + i\eta|\psi|^4\psi = i\gamma\psi,$$

Nonlinear 3-body recombination and linear damping, both can regulate the singularity. *Passot et al. Physica D* **203**, 167 (2005); *LeMesurier Physica D* **138**, 334 (2000)

nonlinear dissipation acts for $|\psi|^2 \approx \eta|\psi|^4$

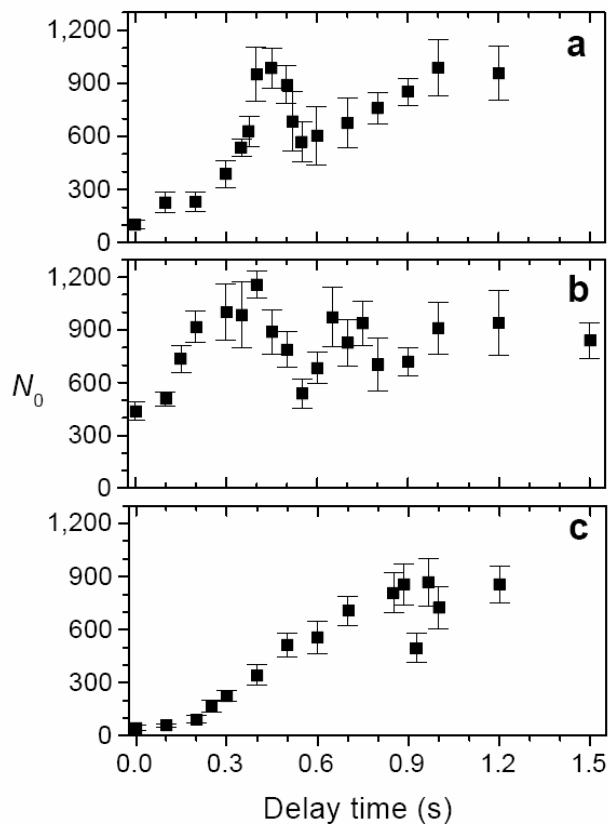
Strong collapse (D=2) power dissipated in blow-up events quantized $\approx N_s$ independent of η !

Weak collapse (D=3) power dissipated in blow-up events $\rightarrow 0$ for $\eta \rightarrow 0$, i.e. dependent on η

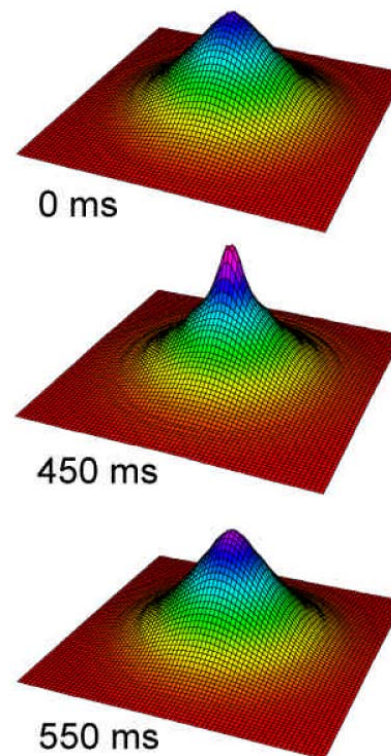
Collapse in Bose-Einstein condensates

^7Li Attractive interaction

J.M. Gerton et al *Nature* **408**, 692 (2000)



Number of particles

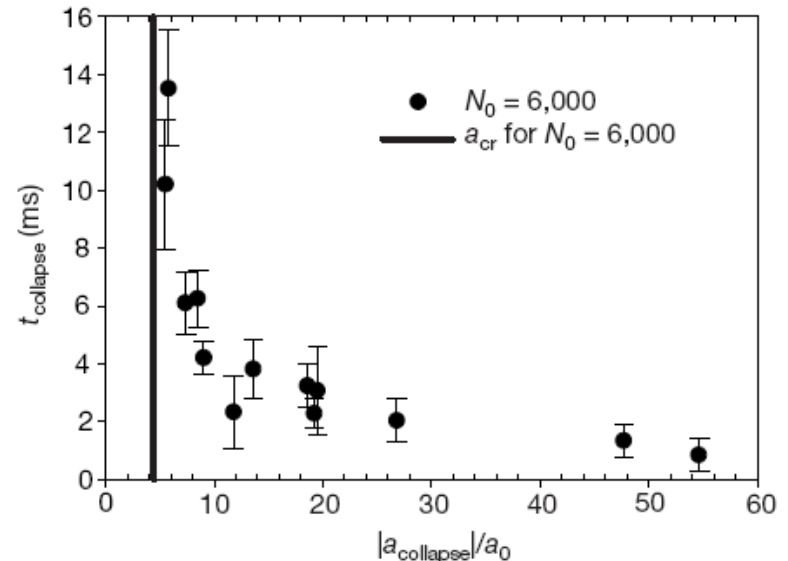
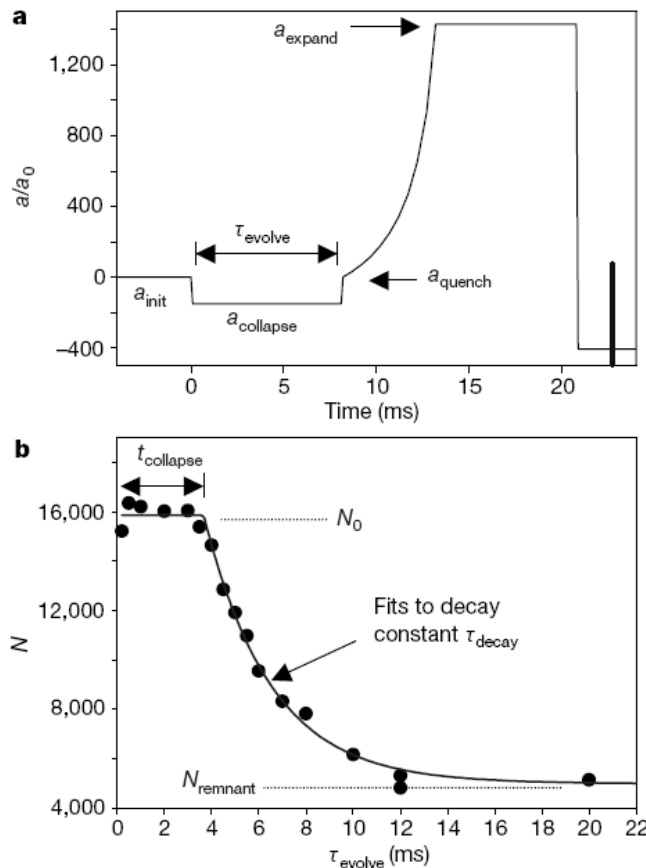


Phase-contrast images, “wave-function”

Collapse in Bose Einstein Condensate: Controlled Collapse

^{85}Ru : Feshbach resonance **shift from repulsive to attractive interaction**

E.A. Donley et al *Nature* **408**, 692 (2000)



Collapse time versus nonlinearity strength

Number of particles

BEC: Mean field theory: Gross-Pitaevskii (G.-P.) equation

$$i\partial_t\psi = -\nabla^2\psi + r^2\psi + a|\psi|^2\psi - i\eta|\psi|^4\psi + i\gamma\psi,$$

Parabolic confinement potential

Where $r = \tilde{r}/l_0$, $t = \tilde{t}\omega/2$, $\psi = \tilde{\psi}(8\pi l_0^2|a_0|)^{1/2}$

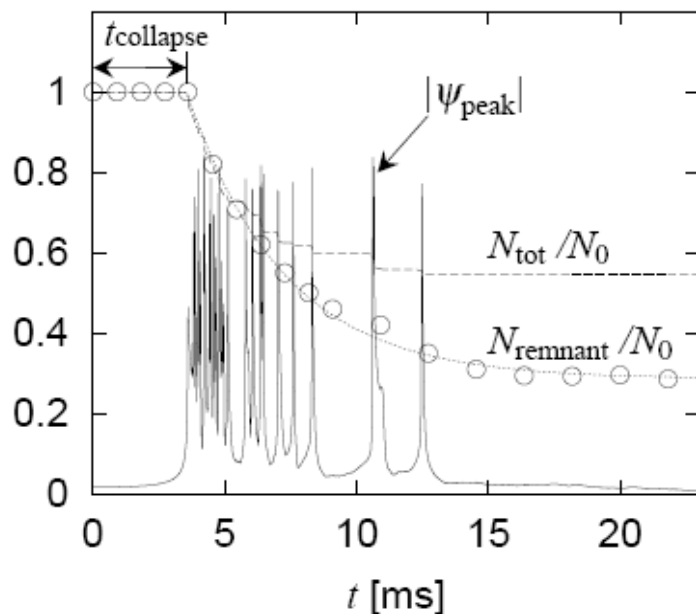
$a = \pm 1$ for positive/negative scattering length, η is 3-body recombination coefficient and γ is source coefficient

Self-similar analysis of the spike evolution provides the number of atoms lost over a sequence of weak collapses (Bergé and Rasmussen, Phys. Lett. A **304**, 136 (2002))

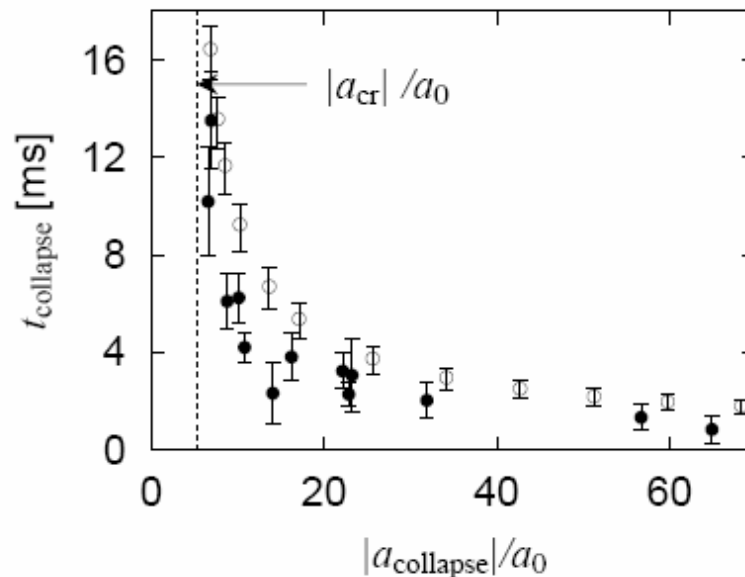
$$N_n = (1 - \Delta N/N_0)^n N_0$$

Numerical solutions of the G-P-equation

M. Ueda, H. Saito, J. Phys. Soc. Jpn. Vol. 72 (2003) Suppl. C pp.127-133



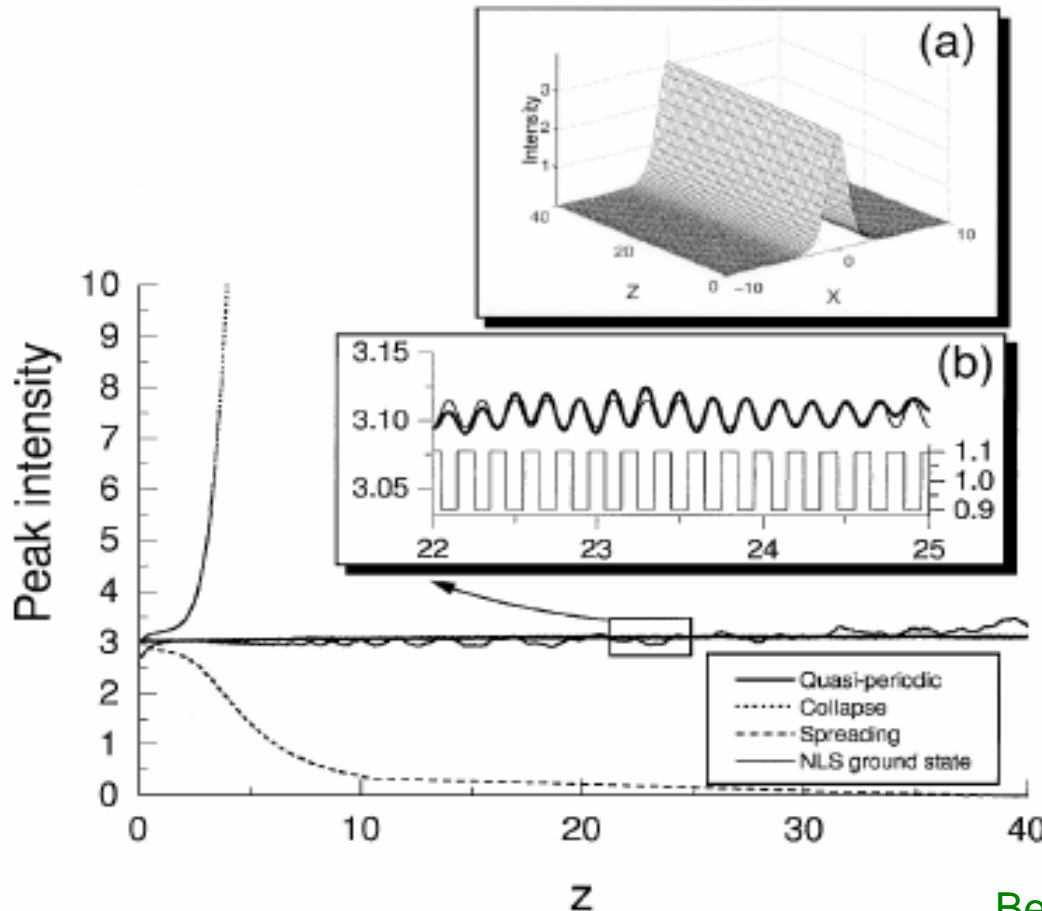
Multiple weak collapses



Collapse time versus scattering “strength”

Active Regularization

Layered nonlinear medium; nonlinearity management;



2-D soliton

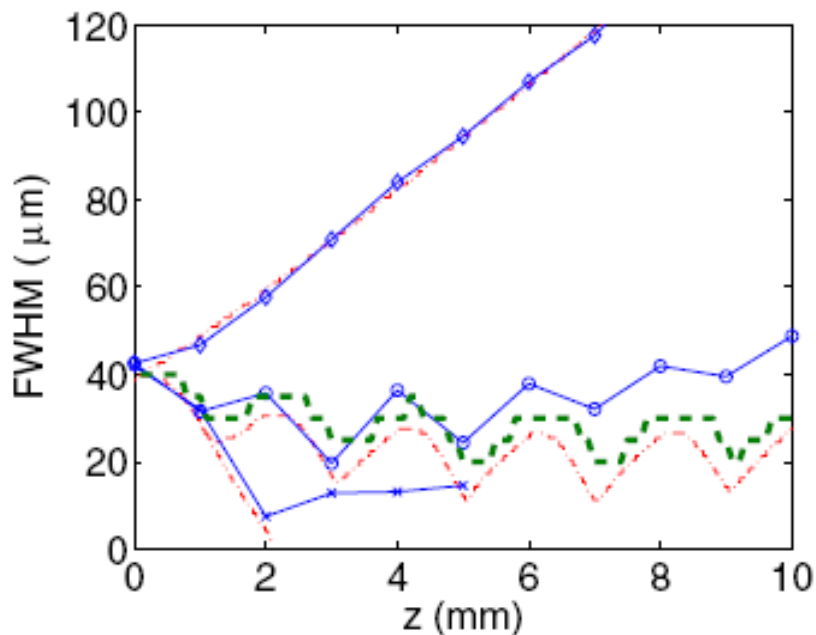
3-D nonlinearity
managements possible??

Bergé *et al. Optics Letters* **25**, 1037 (2000)

Active Regularization

Layered nonlinear medium; nonlinearity management; **experiment!**

Centurion *et al.* *Phys. Rev. Letters* **97**, 033903 (2006)



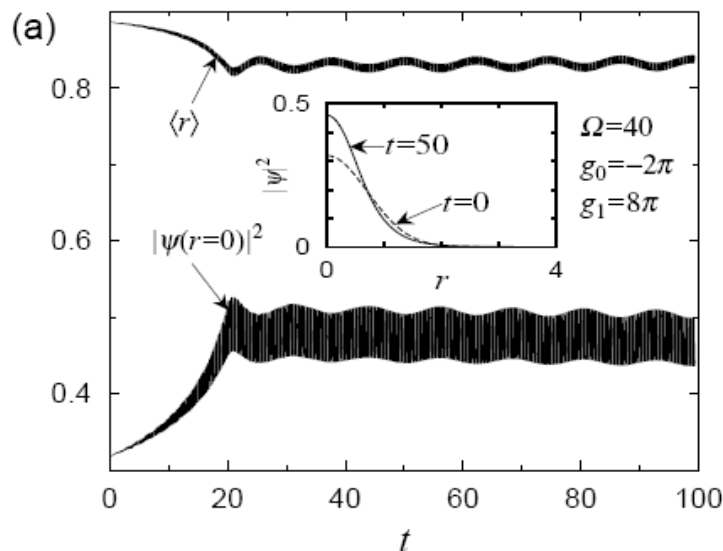
BEC Nonlinearity management

Modulate scattering length a by means of the Feshbach resonance. 2-D possible, 3-D other effects are important

No experimental observations as yet

Numerical simulations 2-D G-P equations with $a(t)$ stable solitons, 3-D?

Saito and Ueda *Phys. Rev. Lett.* **90**, 040403 (2003)



Conclusions

- Collapse dynamics important and fascinating in many branches of physics
- Collapse can be controlled and utilized
- Covered only a limited part of investigations
- Collapse in hydrodynamics : Singularity in Navier-Stokes equations: 2-D no 3-D ??
- Wave breaking
-