



SMR.1766 - 12

**Miniworkshop on
New States of Stable and Unstable Quantum Matter
(14 - 25 August 2006)**

Fermi Surfaces that Wobble and Pop

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These are preliminary lecture notes, intended only for distribution to participants



CCLRC **ISIS** 

Fermi surfaces that wobble and pop

Pomeranchuk and Topological Fermi Surface Instabilities from Central Interactions

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Credits

This work:
(mostly...)

[cond-mat/0601103](https://arxiv.org/abs/cond-mat/0601103)
(to appear in *PRB*)

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A.J. Schofield



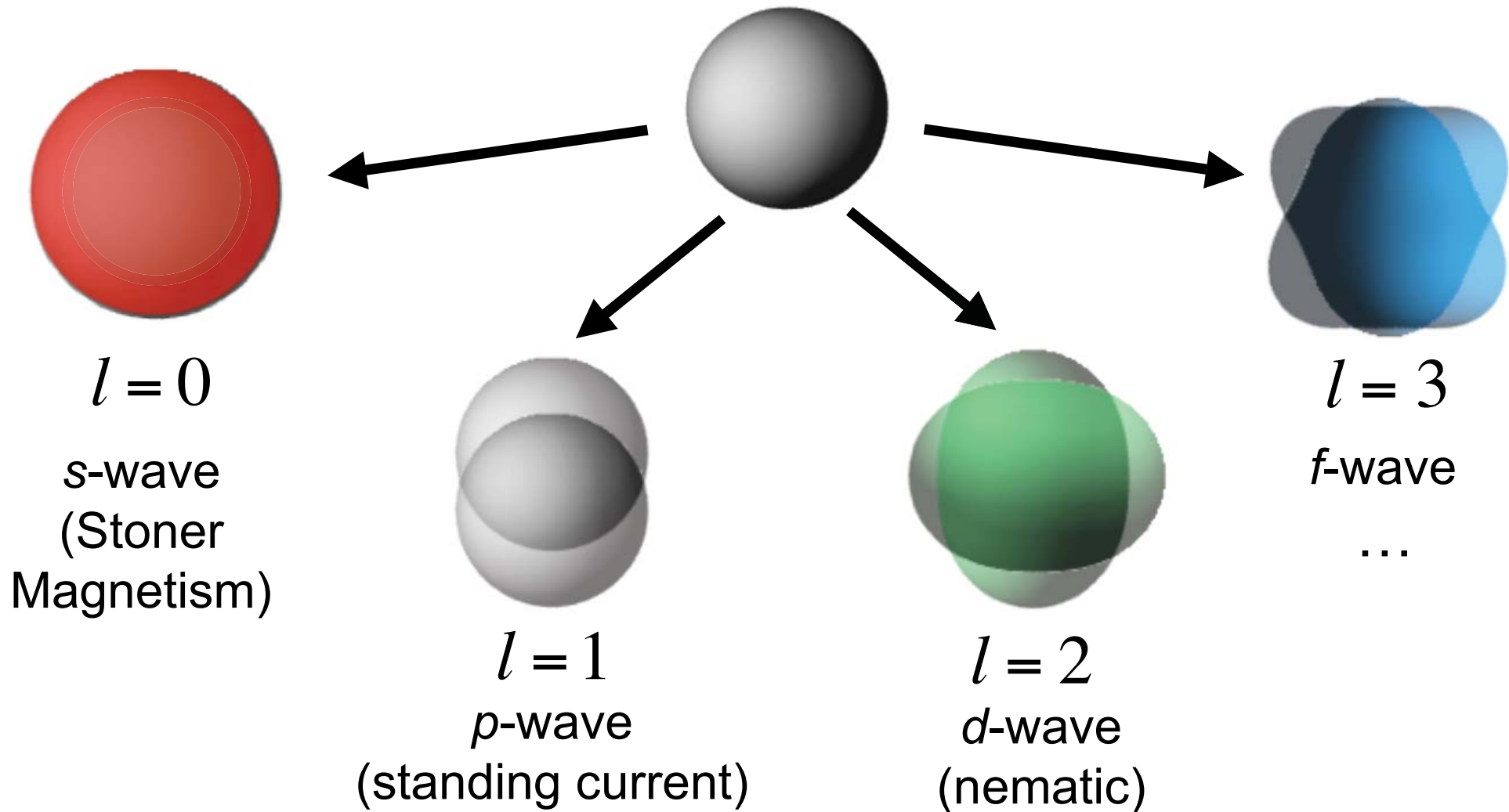
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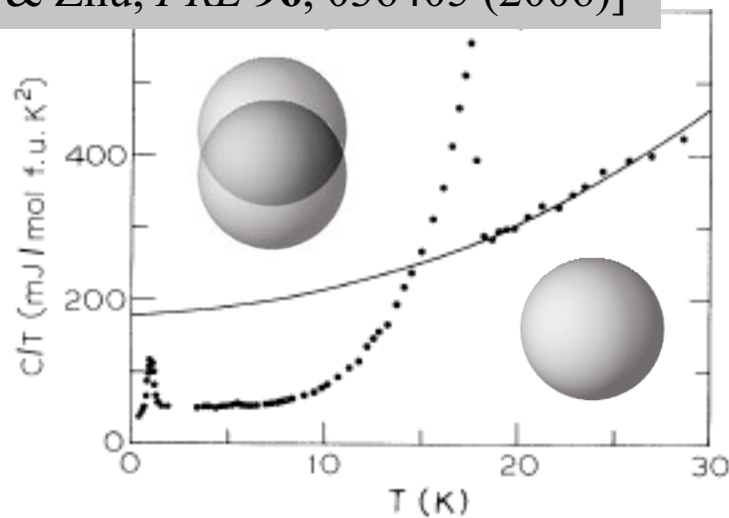
The Pomeranchuk Instability



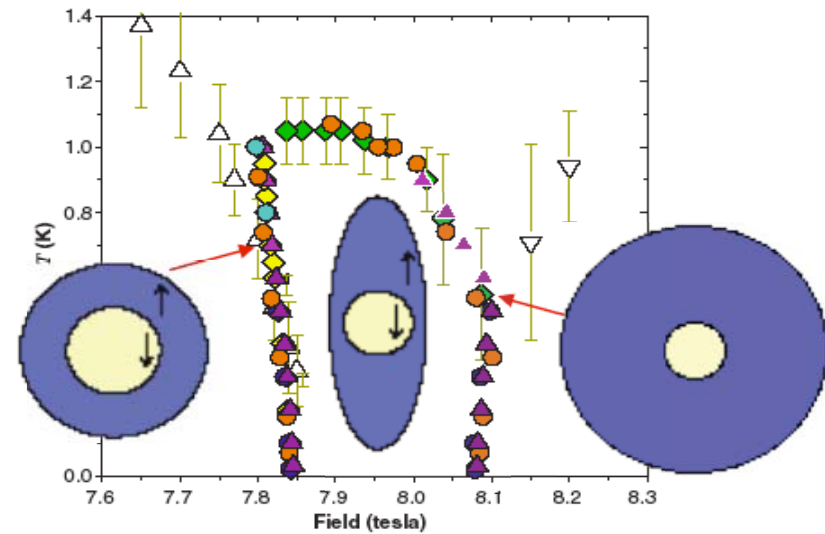
Unconventional Phase Transitions



[Varma & Zhu, *PRL* **96**, 036405 (2006)]

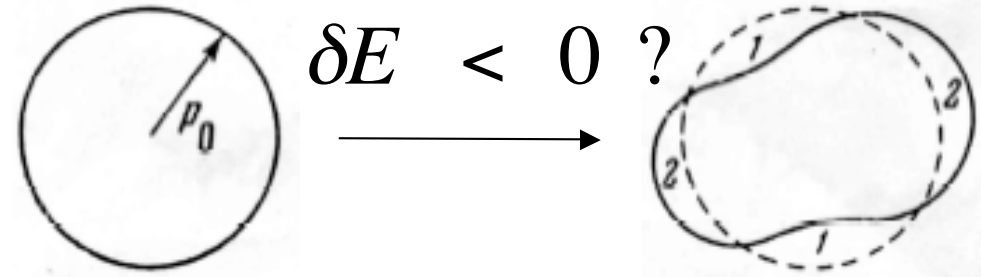
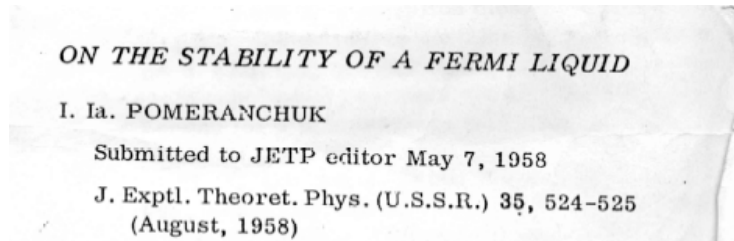


[Palstra et al.,
PRL **55**, 2727 (1985)]



[S.A.Griger et al.,
Science **306**, 1154 (2004)]

Conditions for a Pomeranchuk instability



Pomeranchuk instability condition:

$$1 + \frac{1}{2l+1} F_l < 0$$

$l = 0, 1, 2, 3 \dots$ symmetry of deformation

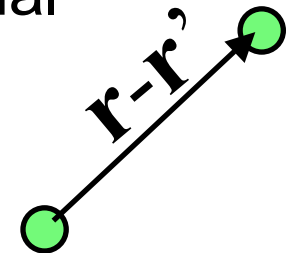
$F_0, F_1, F_2, F_3 \dots$ Landau parameters

Conditions for a Pomeranchuk instability

Microscopic model:

$$H = \underbrace{\int d^3\mathbf{r} \sum_{\sigma} \hat{c}_{\mathbf{r},\sigma}^+ \left[\frac{-\hbar^2 \nabla^2}{2m} - \mu \right] \hat{c}_{\mathbf{r},\sigma}}_{\text{free fermions}} + \underbrace{\frac{1}{2} \sum_{\sigma, \sigma'} \int d^3\mathbf{r} \int d^3\mathbf{r}' \hat{c}_{\mathbf{r},\sigma}^+ \hat{c}_{\mathbf{r}',\sigma'}^+ V(|\mathbf{r} - \mathbf{r}'|) \hat{c}_{\mathbf{r}',\sigma'} \hat{c}_{\mathbf{r},\sigma}}_{\text{central potential}}$$

Trial ground state: $|\Psi\rangle = \prod_{\varepsilon_{\sigma}(\mathbf{k}) < 0} \hat{c}_{\mathbf{k},\sigma}^+ |0\rangle$



The dispersion relation $\varepsilon_{\sigma}(\mathbf{k})$ is determined **variationally**.

Conditions for a Pomeranchuk instability

$\varepsilon_\sigma(\mathbf{k})$ becomes anisotropic under the following instability condition:

$$\int d^3\mathbf{r} V(|\mathbf{r}|) \left[j_l(k_F |\mathbf{r}|)^2 - j_1(k_F |\mathbf{r}|)^2 \right] > \frac{2\pi^2 \hbar^2}{mk_F}$$

(Note: The term $j_1(k_F |\mathbf{r}|)^2$ is highlighted with a red bracket and labeled "= 0 for l=1")

(Note: The term $j_l(k_F |\mathbf{r}|)^2$ is highlighted with a green bracket and labeled "for $k_F r \ll \pi$ { ≈ 1 for $l=0$; ≈ 0 for $l=1,2,3,\dots$ ")

(Note: The index $l=0,1,2,3,\dots$ is shown with a green arrow pointing to the $l=0$ term and a red 'X' over the $l=1$ term)

Pomeranchuk instability with $l=1$ **never takes place.**

$l=2,3,\dots$ requires range of interaction $r_0 \sim \pi/k_F$ or larger.

A simple example

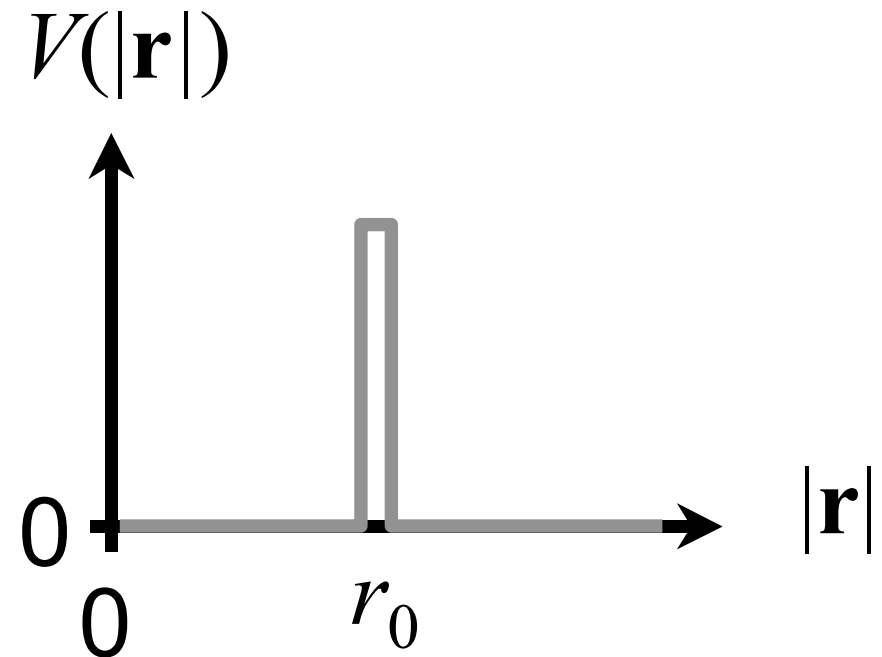
We can solve the theory analytically for **repulsion at a finite distance** r_0 :

$$V(|\mathbf{r}|) = g\delta(|\mathbf{r}| - r_0)$$

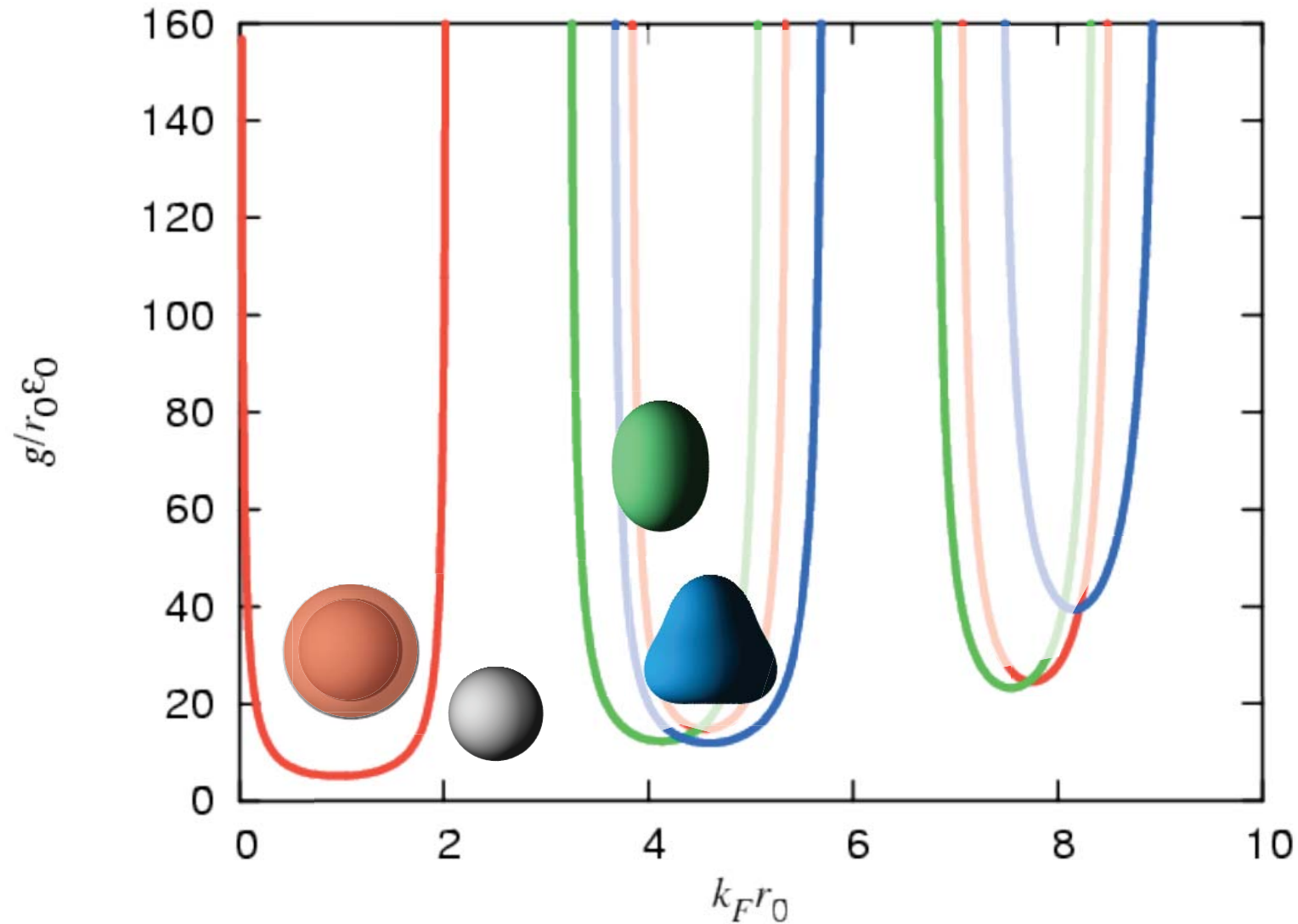
Characteristic length: r_0

Characteristic energies: g/r_0

$$\hbar^2/2mr_0^2 \equiv \varepsilon_0$$



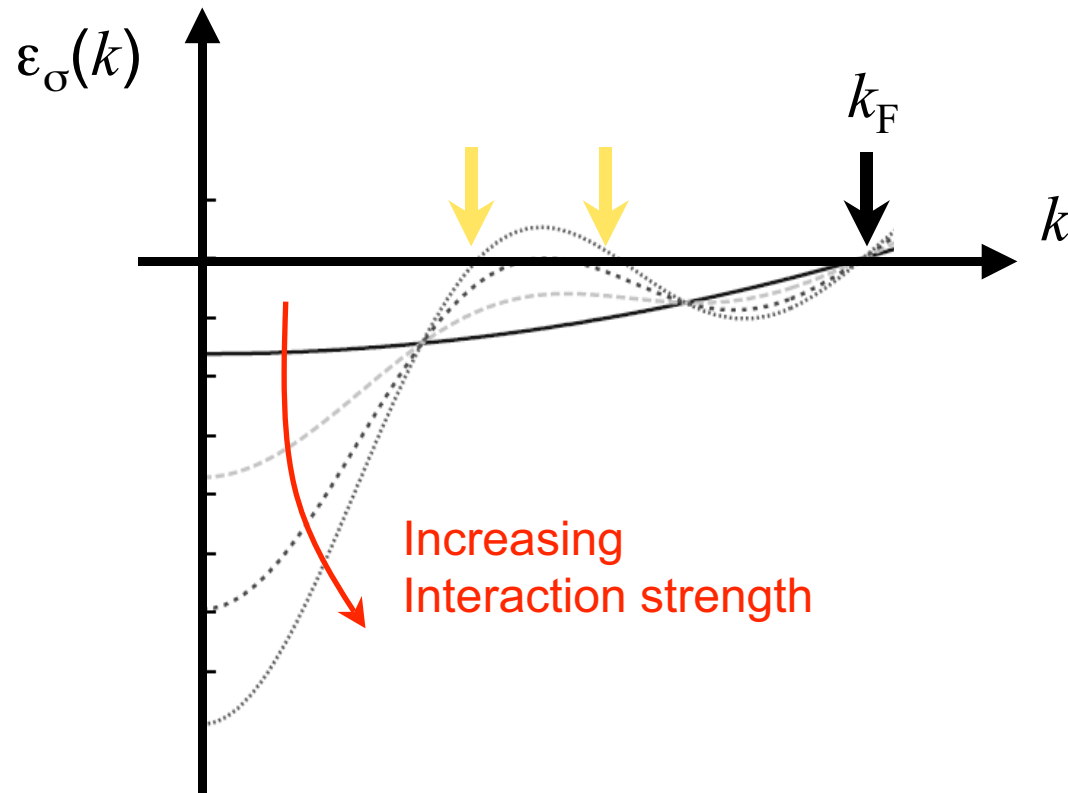
A simple example



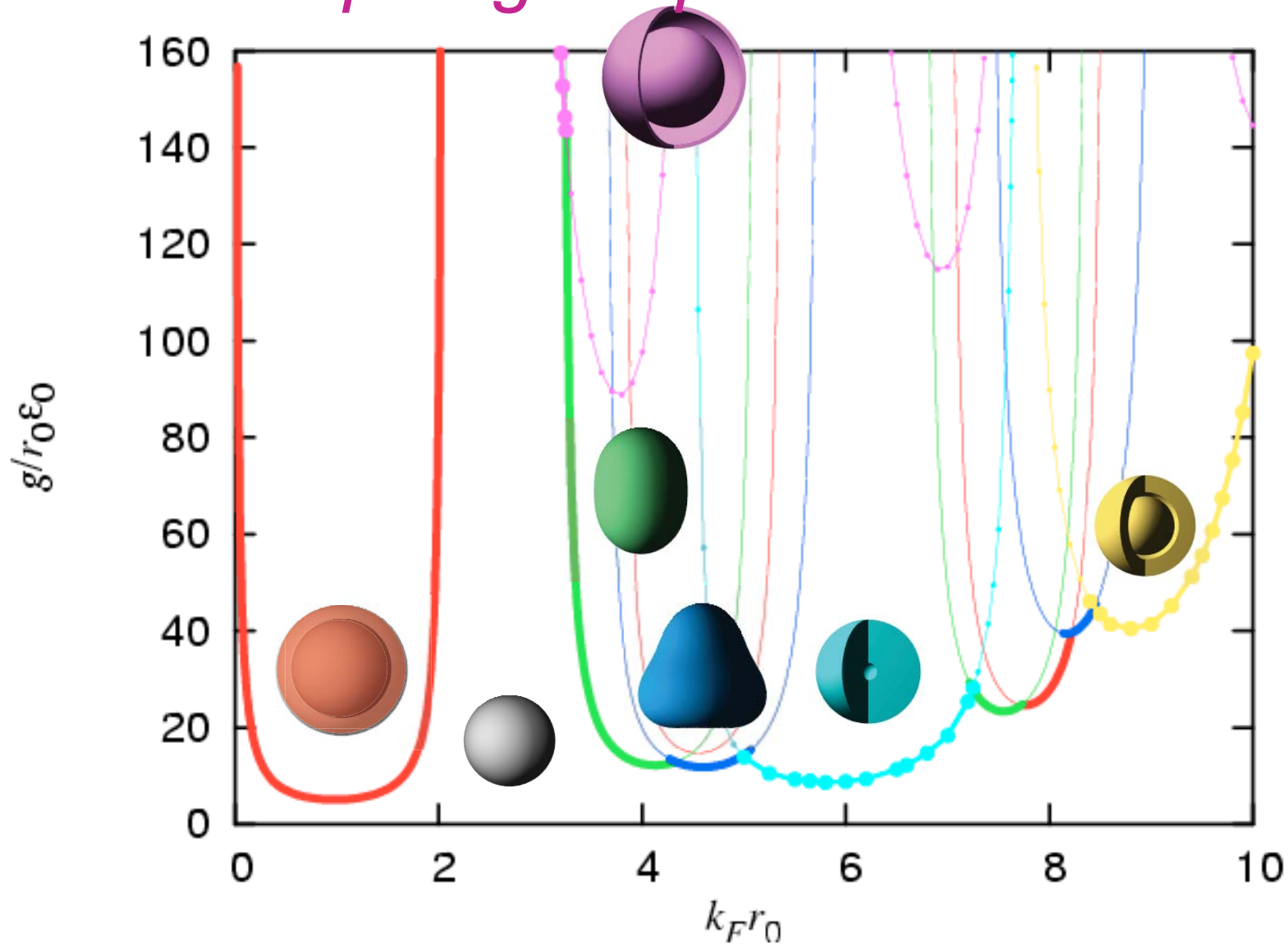
Topological phase transitions

Even without symmetry breaking, interactions may change $\varepsilon_{\sigma}(\mathbf{k})$ qualitatively...

...even leading to changes of Fermi surface **topology**



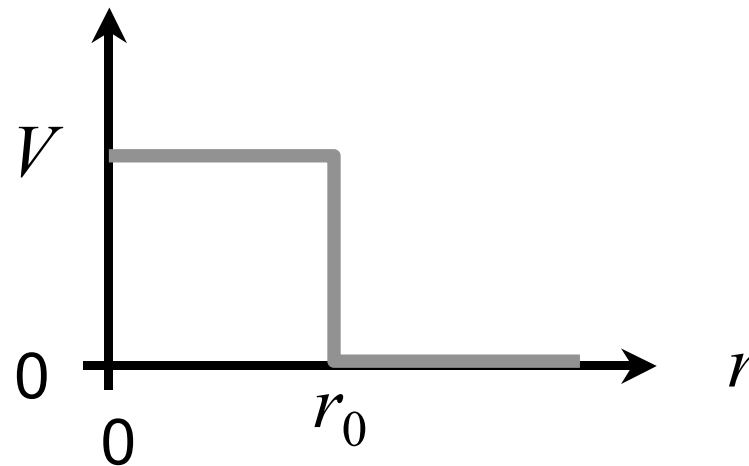
Topological phase transitions



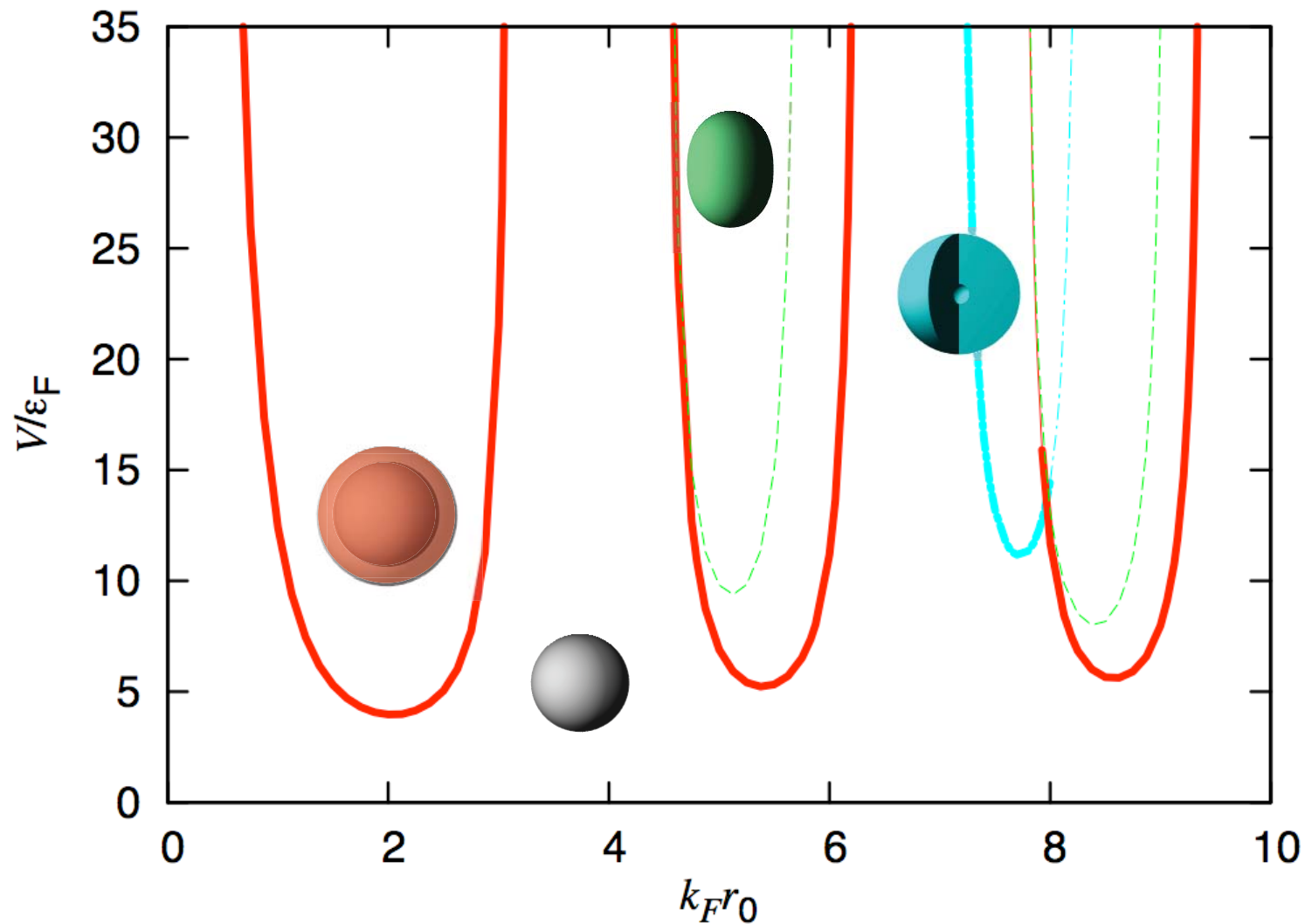
Another example: “hard” core potential

“Top hat” potential of radius r_0 :

$$V(|\mathbf{r}|) = V \Theta(r_0 - |\mathbf{r}|)$$



Another example: "hard" core potential

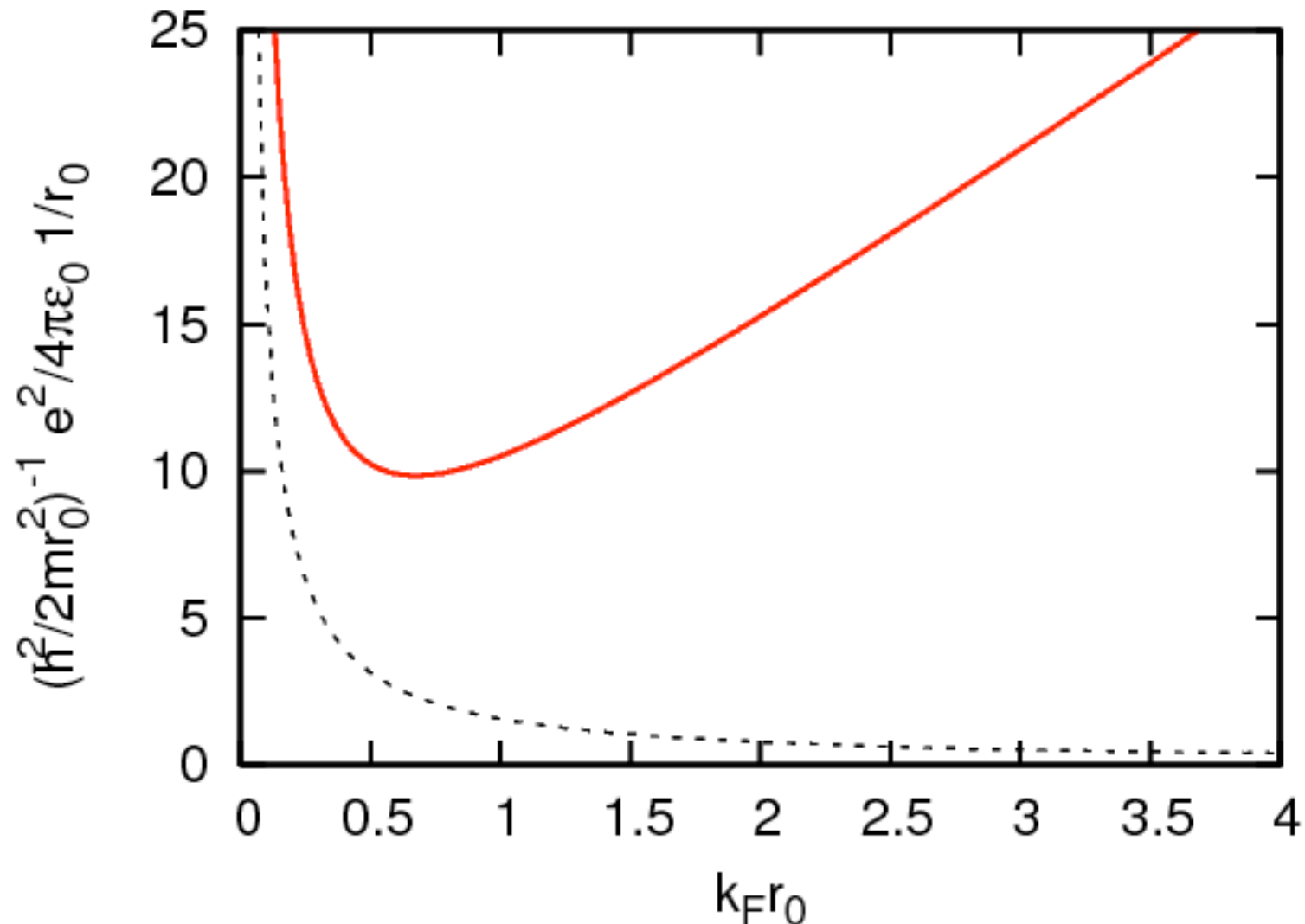


How “sharp” does the feature have to be?

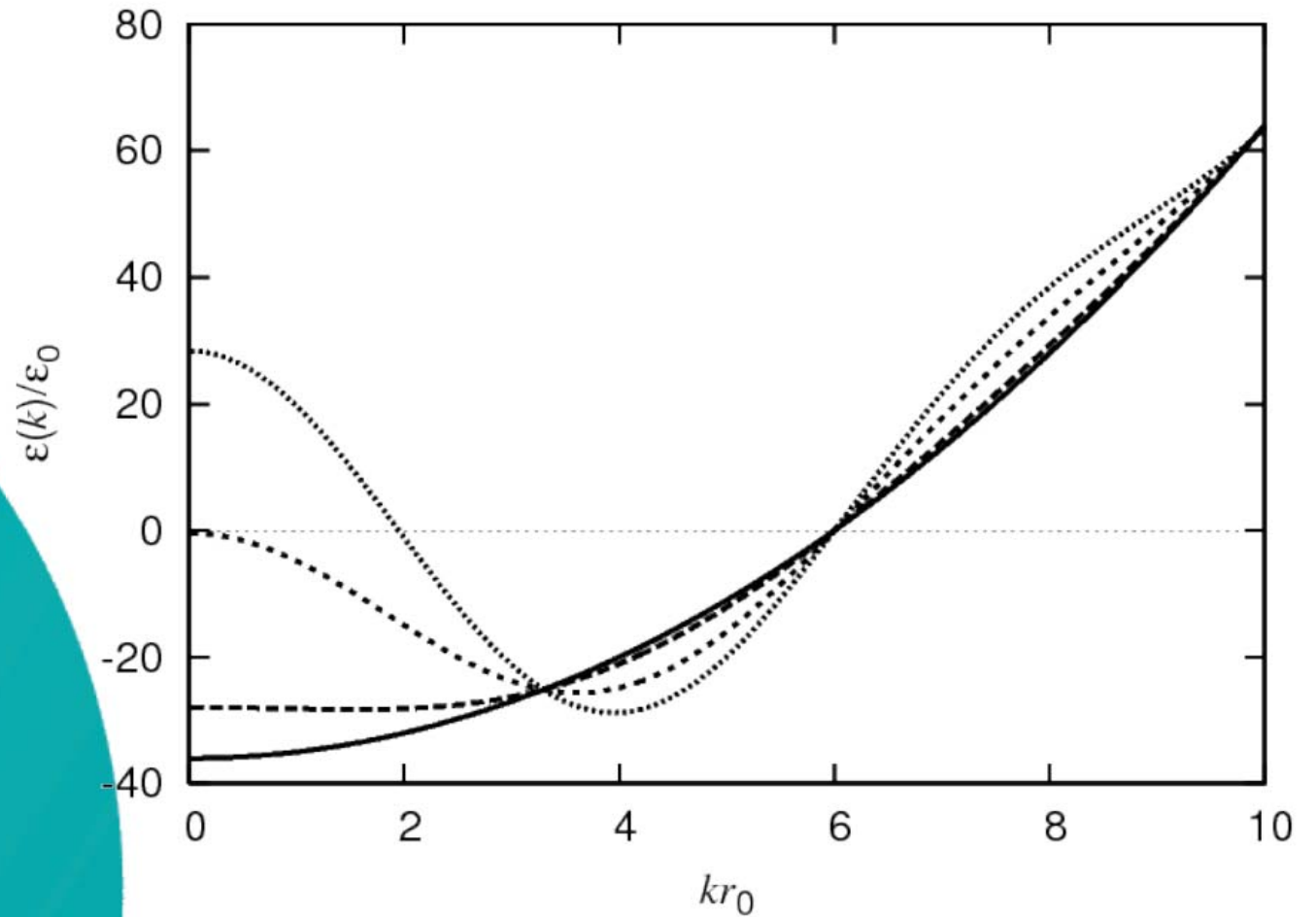
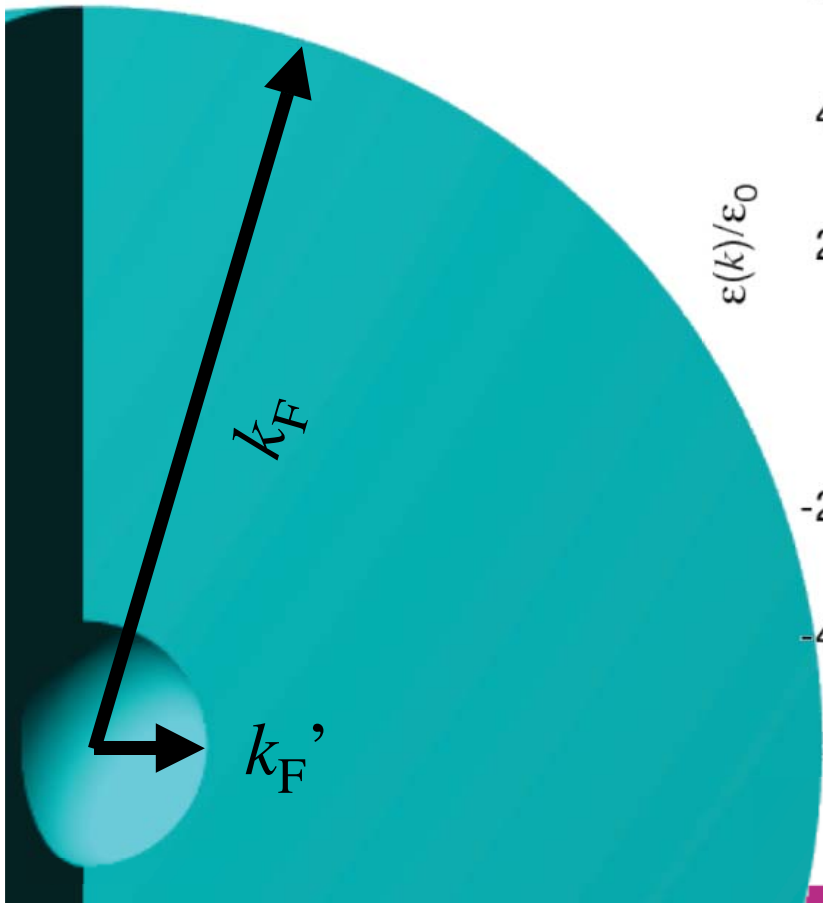
Screened Coulomb
Interaction:

$$V(|\mathbf{r}|) = \frac{e^2}{4\pi\epsilon_0} \frac{e^{-|\mathbf{r}|/r_0}}{|\mathbf{r}|}$$

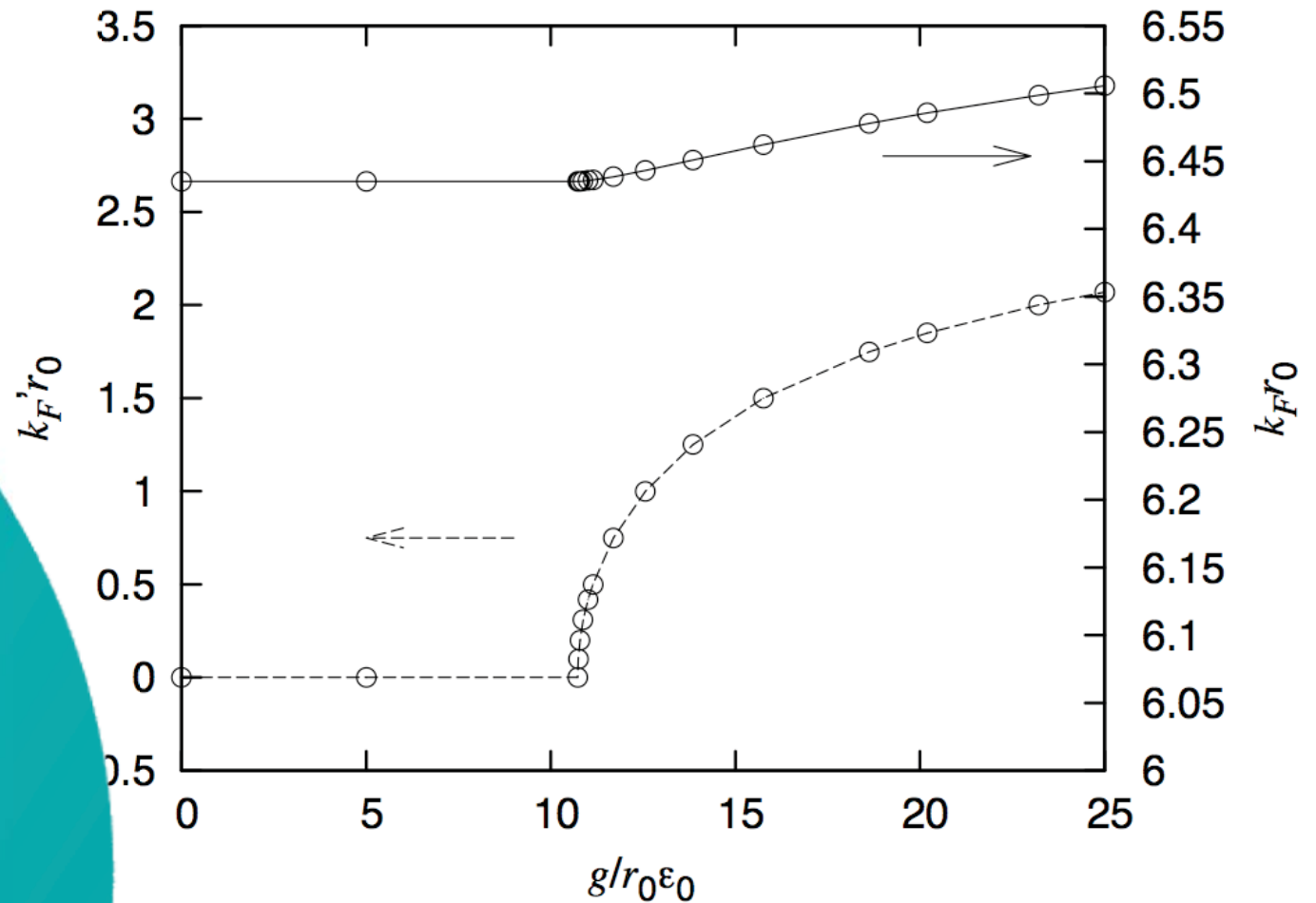
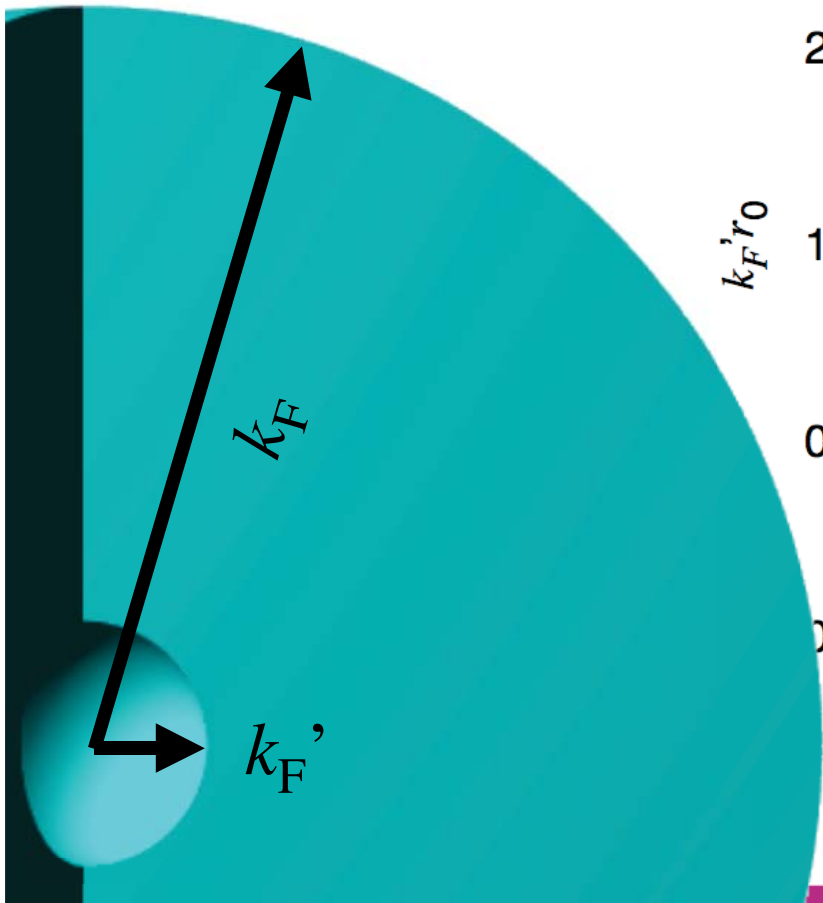
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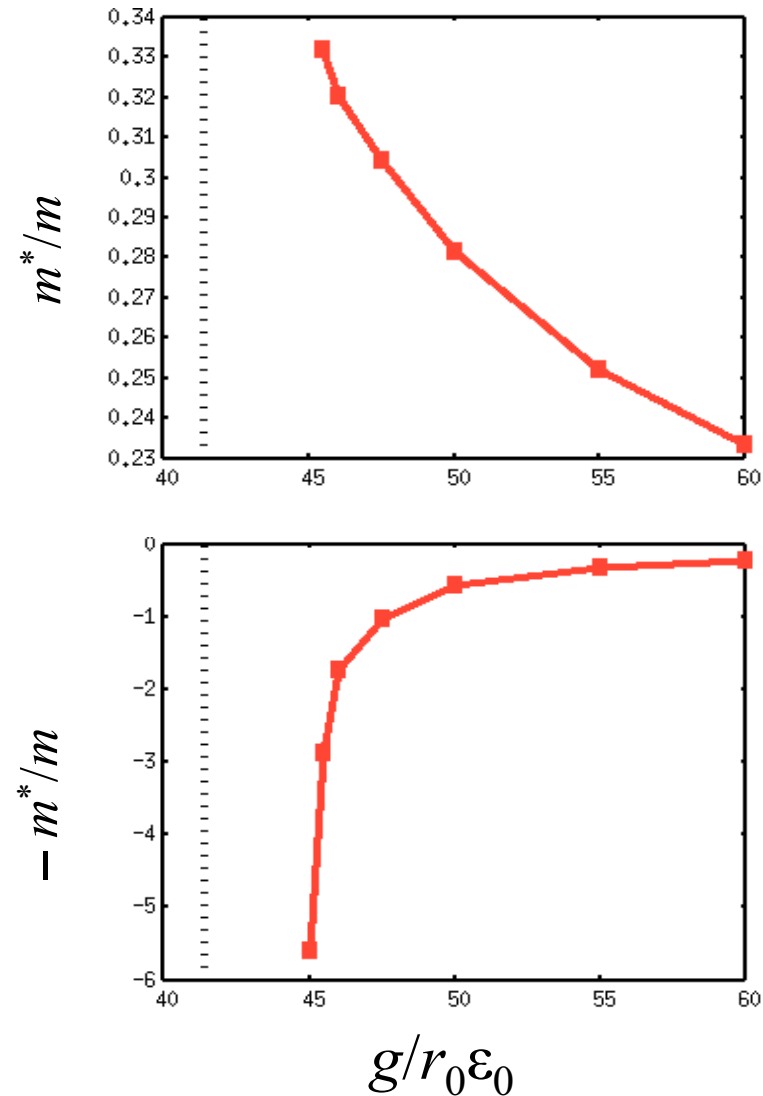
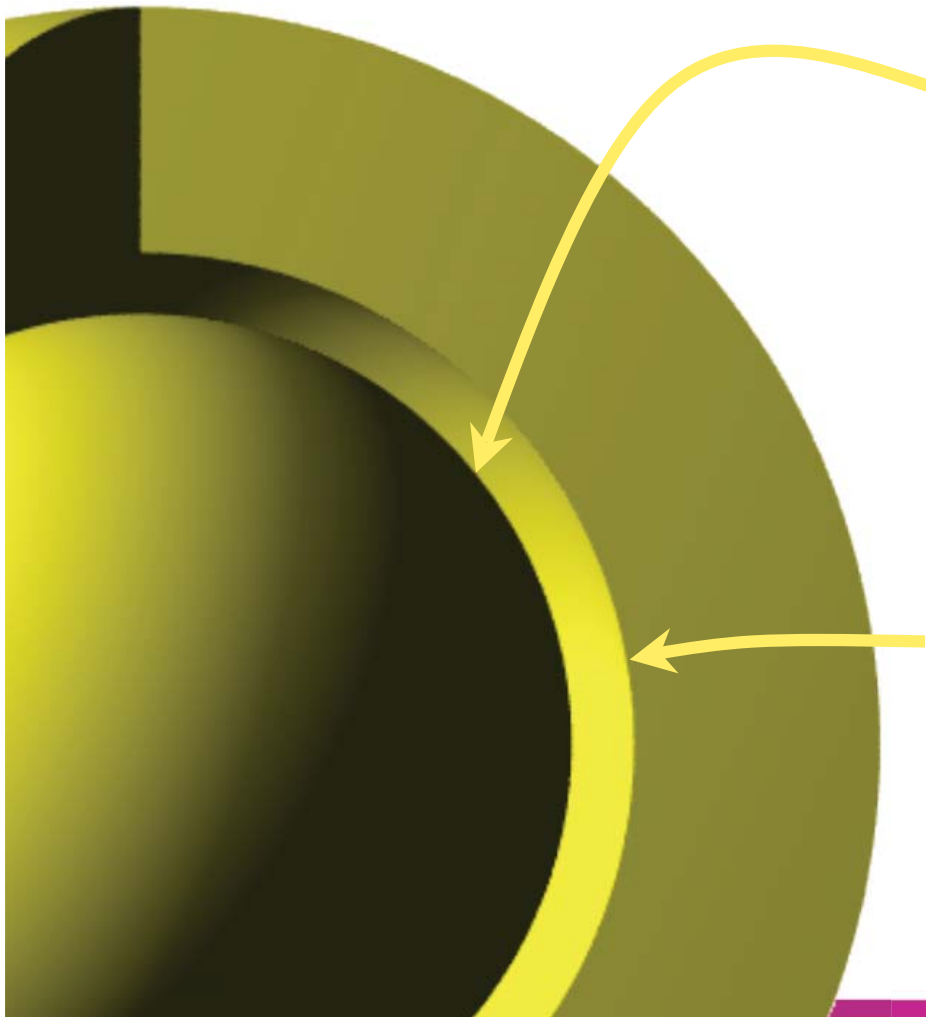
Topological phase transitions



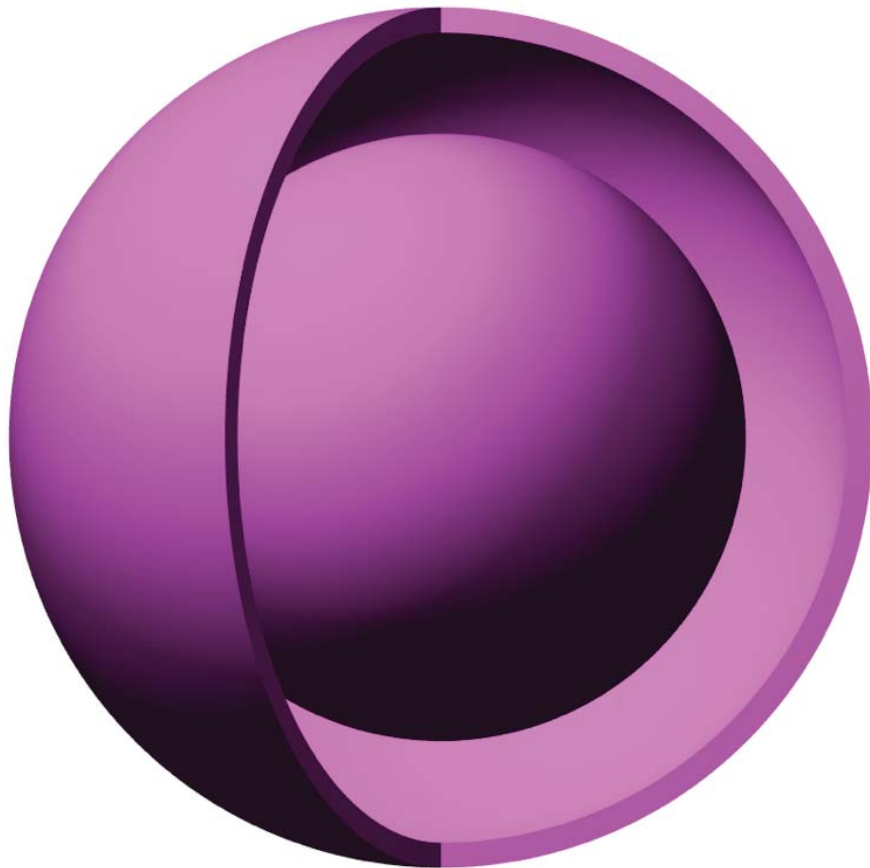
Topological phase transitions



Topological phase transitions



Specific heat signature of topological Fermi surface shape instability



$$C_V = \frac{\pi^2 k_B^2}{3} \times \frac{dn}{d\epsilon} \times T \begin{cases} \mu(T) \approx \epsilon_F \\ \epsilon_\sigma(\mathbf{k}, T) \approx \epsilon_\sigma(\mathbf{k}) \end{cases}$$

$$C_V = C_{V,0} + C_{V,1} + C_{V,2}$$

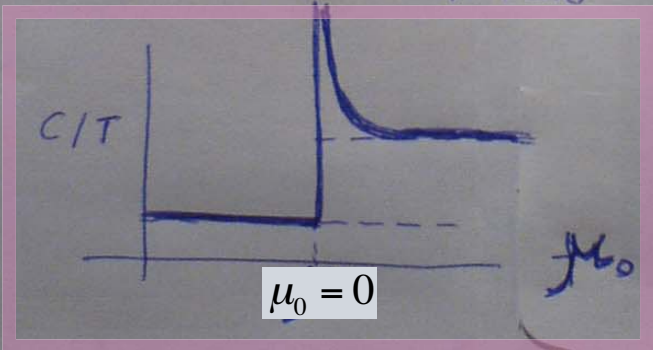
$$\frac{dn}{d\epsilon} = \frac{P_F m^*}{\pi^2 \hbar^3}$$

$$C_V = \frac{\pi^2 k_B^2}{3} \left(\frac{P_F m^*}{\pi^2 \hbar^3} + \frac{P_F' m^{*1}}{\pi^2 \hbar^3} + \frac{P_F'' m^{*11}}{\pi^2 \hbar^3} \right) T$$

$$K_F^{crit} = \frac{K_F' + K_F''}{2}$$

$$\epsilon_{crit}(k) = \frac{\hbar^2}{2m_0} (k - K_F^{crit})^2 \rightarrow \mu_0$$

$$\therefore m^{*1} \approx m^{*11} \approx \frac{1}{2} \frac{1}{\sqrt{\mu_0 / \frac{\hbar^2 P_F^2}{2m_0}}} \text{ as } \mu_0 \rightarrow 0$$



Conclusion and work in progress...

What could be happening ?

Pomeranchuk instability: a subtle symmetry-breaking phase transition, and therefore a candidate for hidden order.

Even subtler: **topological Fermi surface instabilities.**

Why would it happen ?

Key ingredient: **characteristic length** scale $r_0 \gtrsim k_F^{-1}$.

A guide to realisations? (optical lattices, MOSFETS...)

How would we find out ?

Topological may have **unusual thermod. signatures.**

Nesting, STM, coupling to the lattice...

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