



Forschungszentrum Jülich
in der Helmholtz-Gemeinschaft

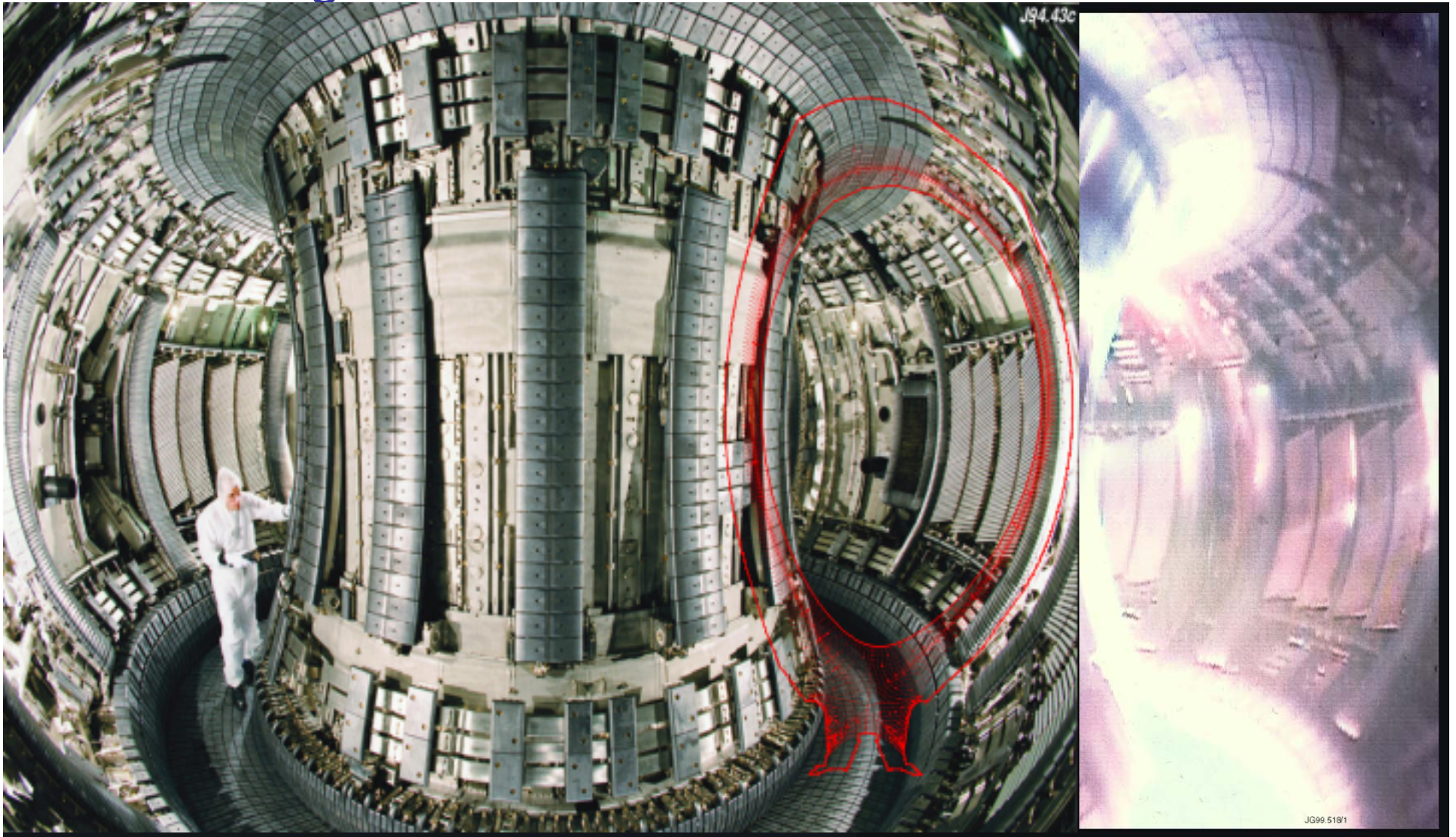


II: The role of hydrogen chemistry in present experiments and in ITER edge plasmas

D. Reiter

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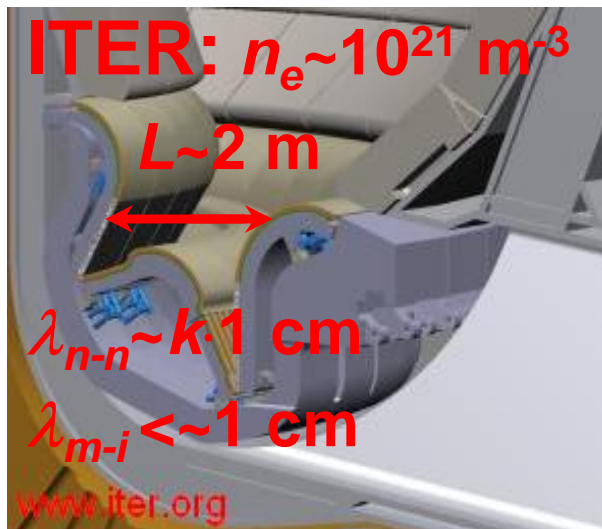
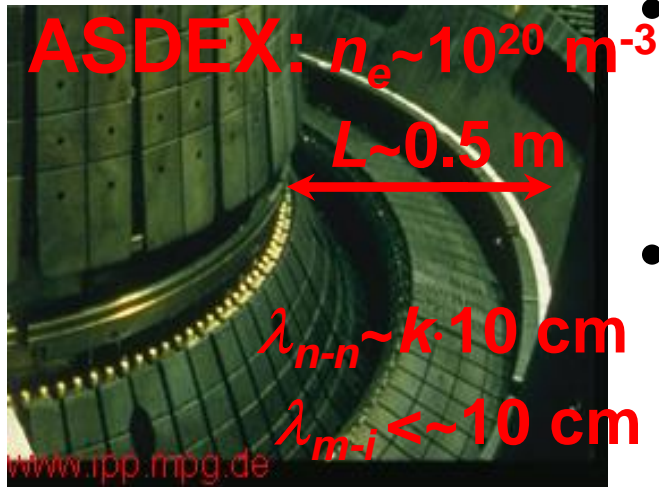
Main focus of coordinated European
edge code model validation is on JET



JET Furnace chamber:

Ø 8.5 m 2.5 m high 3.4 T 7 MA 1 min

Motivation

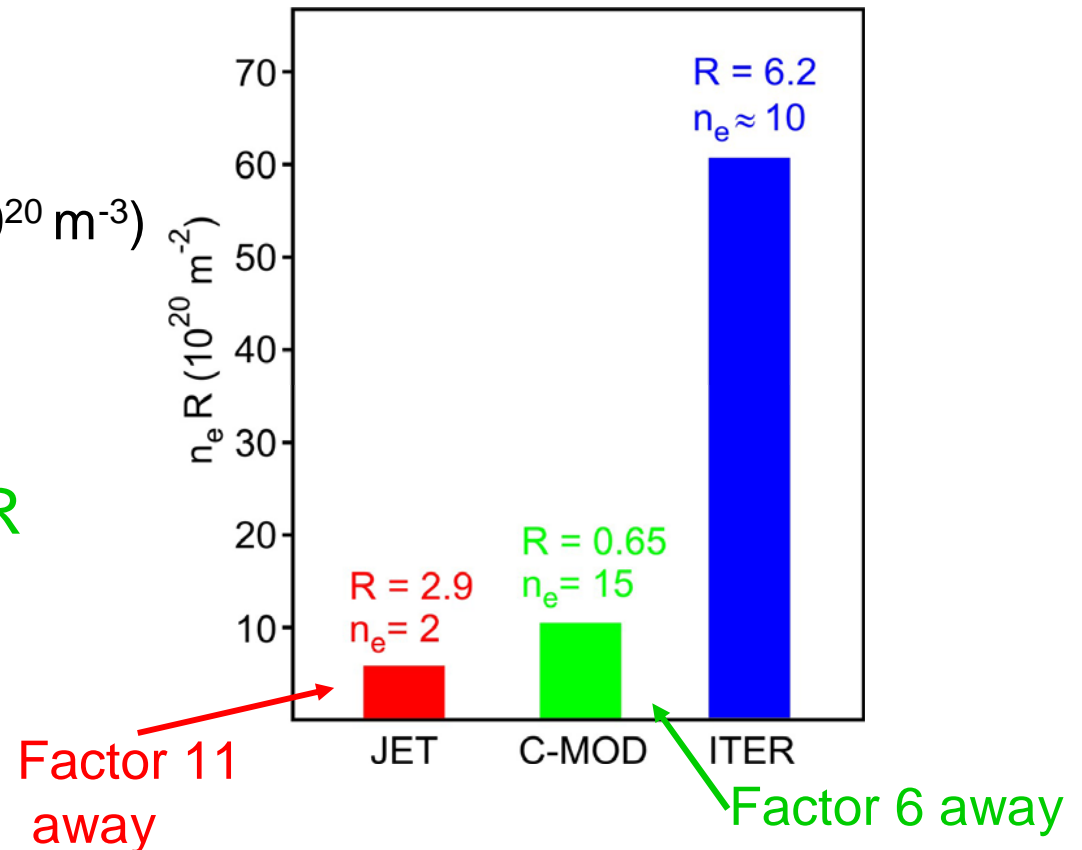


- ITER divertor will have much larger $L \cdot n$ than all existing tokamaks.
- It is necessary to take into account effects which have been safely ignored in divertor modelling for most present devices
- One such effect is *neutral-neutral collisions (viscous effects in gas)*
- Large regions of dense, low temperature plasma requires more accurate simulation of the *molecular reaction kinetics*, including *molecule-ion collisions*

Current hypothesis:
in the “detached state” is the divertor dynamics
and chemistry is controlled by “Collisionality”
(inv. Knudsen number)

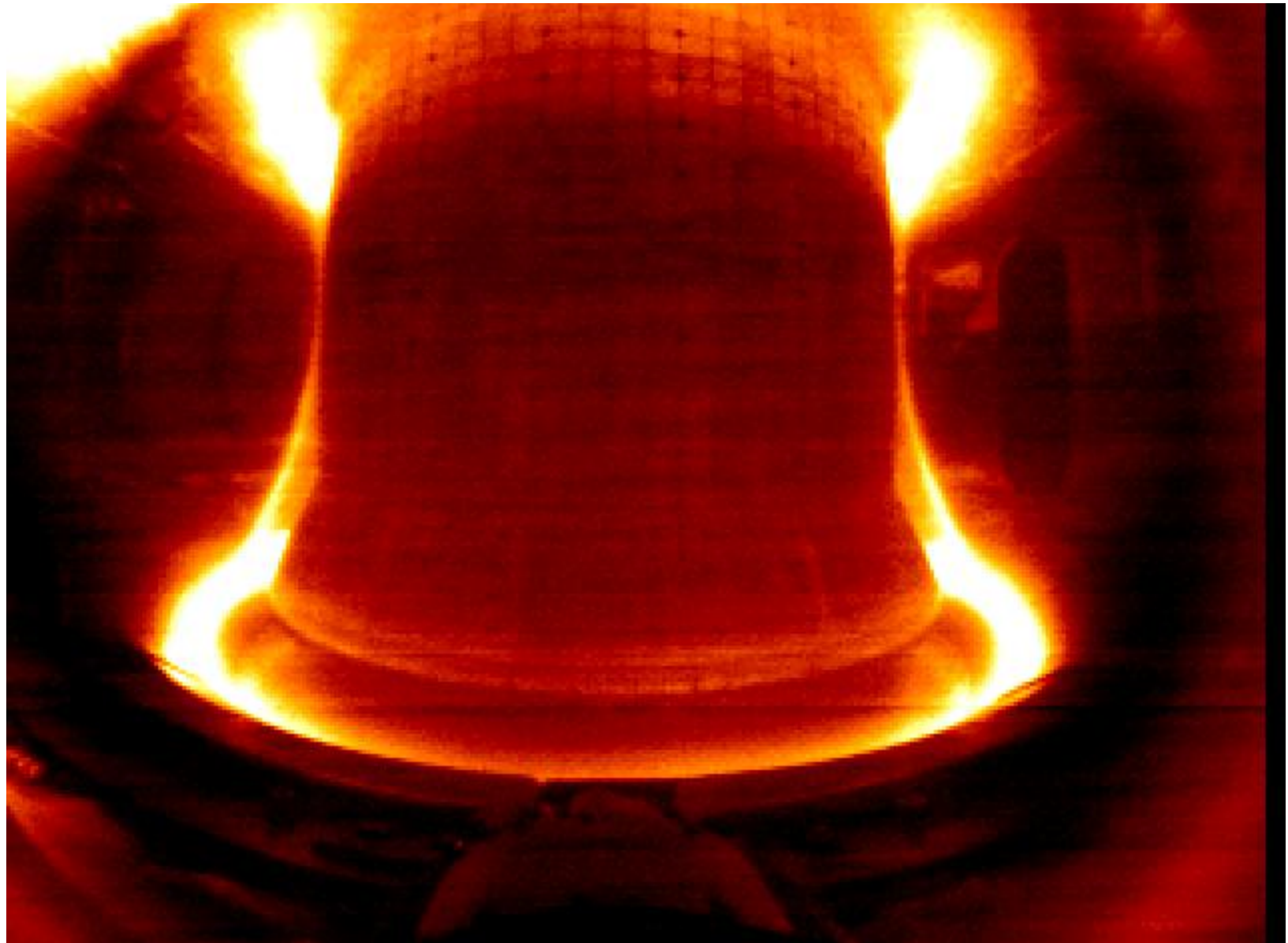
Estimate “Collisionality”: $n_e R$
- n_e -Divertor Plasma density ($\times 10^{20} \text{ m}^{-3}$)
- R - Major Radius (m)

Alcator C-Mod (MIT)
10 times smaller than ITER
similar shape
higher density



Alcator C-Mod (MIT)



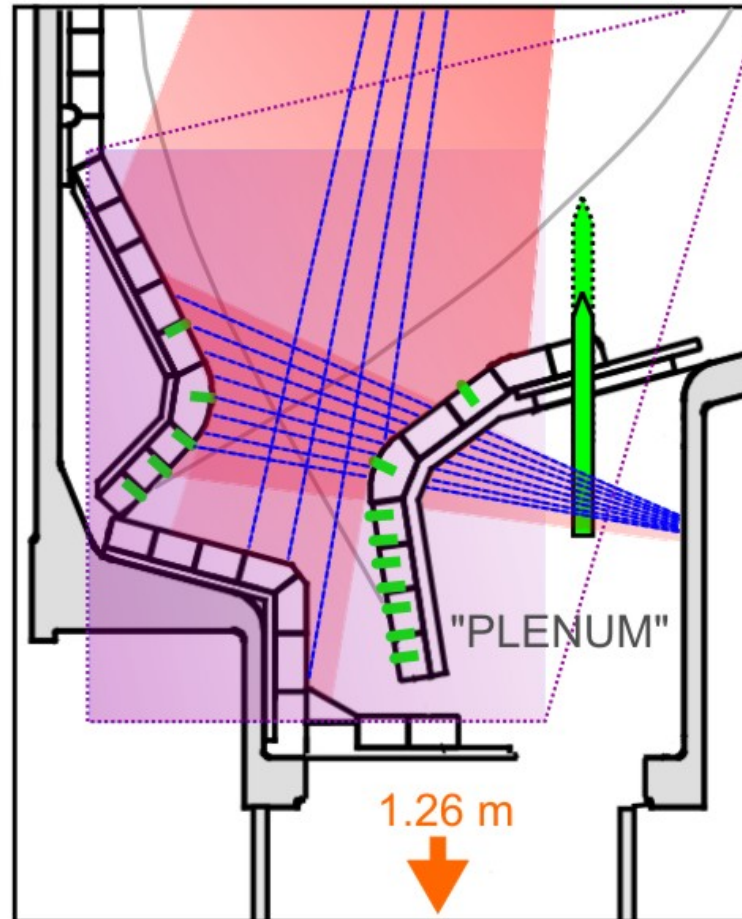


Predictive quality regarding „divertor chemistry“ ???
Separate plasma transport by replacing plasma transport modelling with OSM reconstruction

HIGH
RESOLUTION
DIODE ARRAYS
WITH D_{α} FILTER

TARGET
LANGMUIR
PROBES AND
UPSTREAM
RECIPROCATING
PROBE FOR
 n_e AND T_e

DIVERTOR GAS
PRESSURE
(25 ± 3 mTorr)

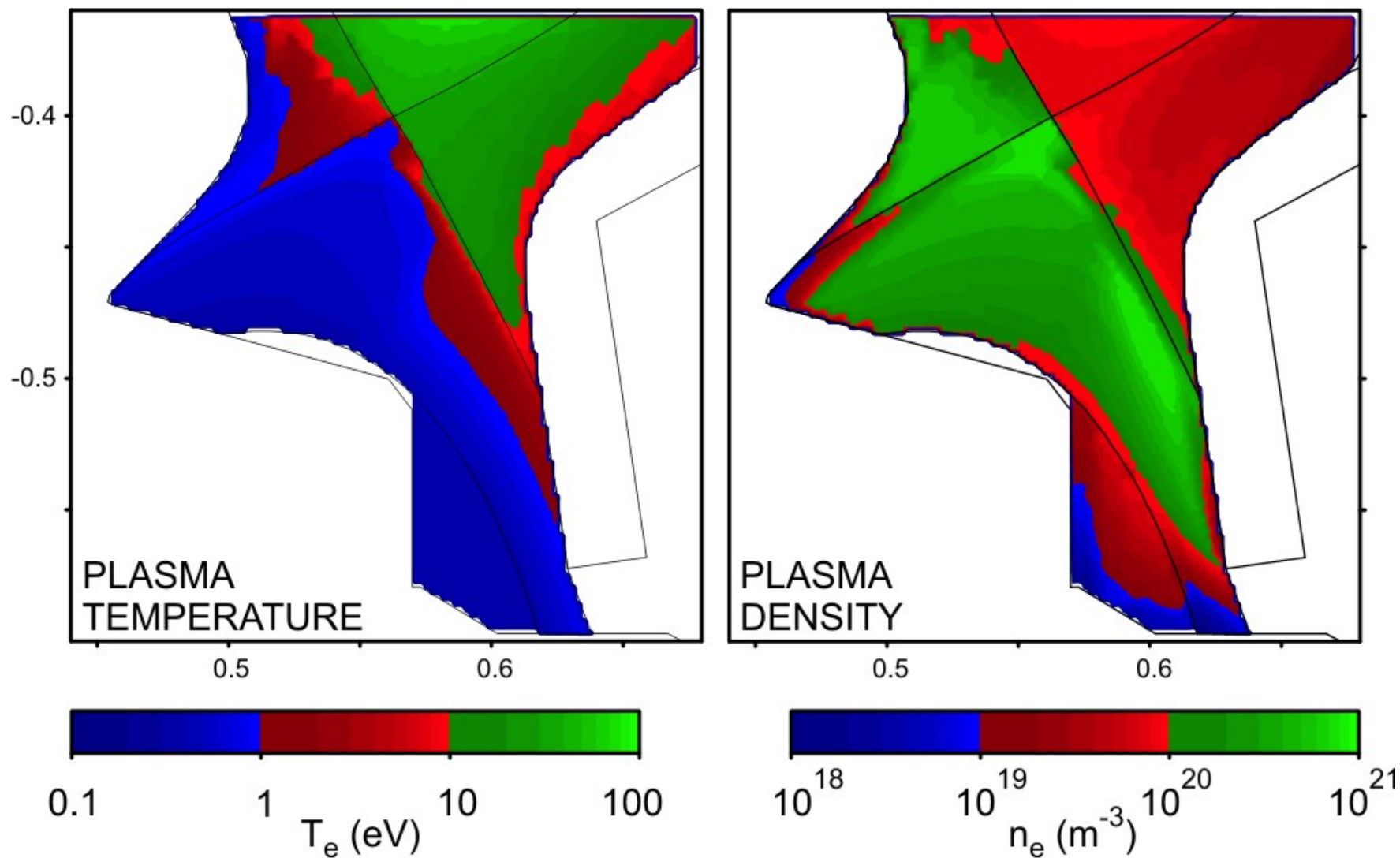


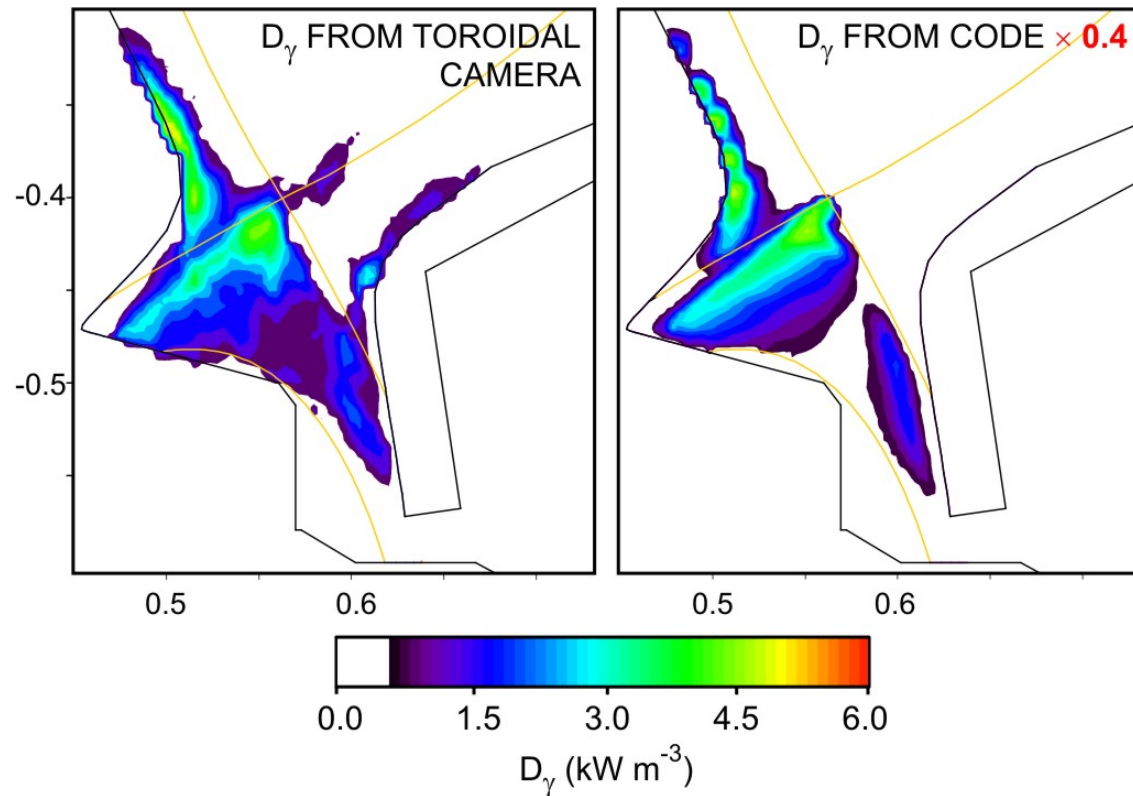
TOROIDALLY
VIEWING CCD
CAMERA WITH
 D_{γ} FILTER

SPECTROMETER
FOR VOLUME
 n_e AND T_e

C. BOSWELL
B. LaBOMBARD
B. LIPSCHULTZ
A. NIEMCZEWSKI
S. PITCHER
J. TERRY

Shot: 990429019, at 950ms,
 $\langle n_e \rangle = 1.5 \cdot 10^{20}$, $I_p = 0.8$ MA, $B_{\text{tor}} = 5.4$ T
OSM reconstruction (Lisgo et al., 2004)

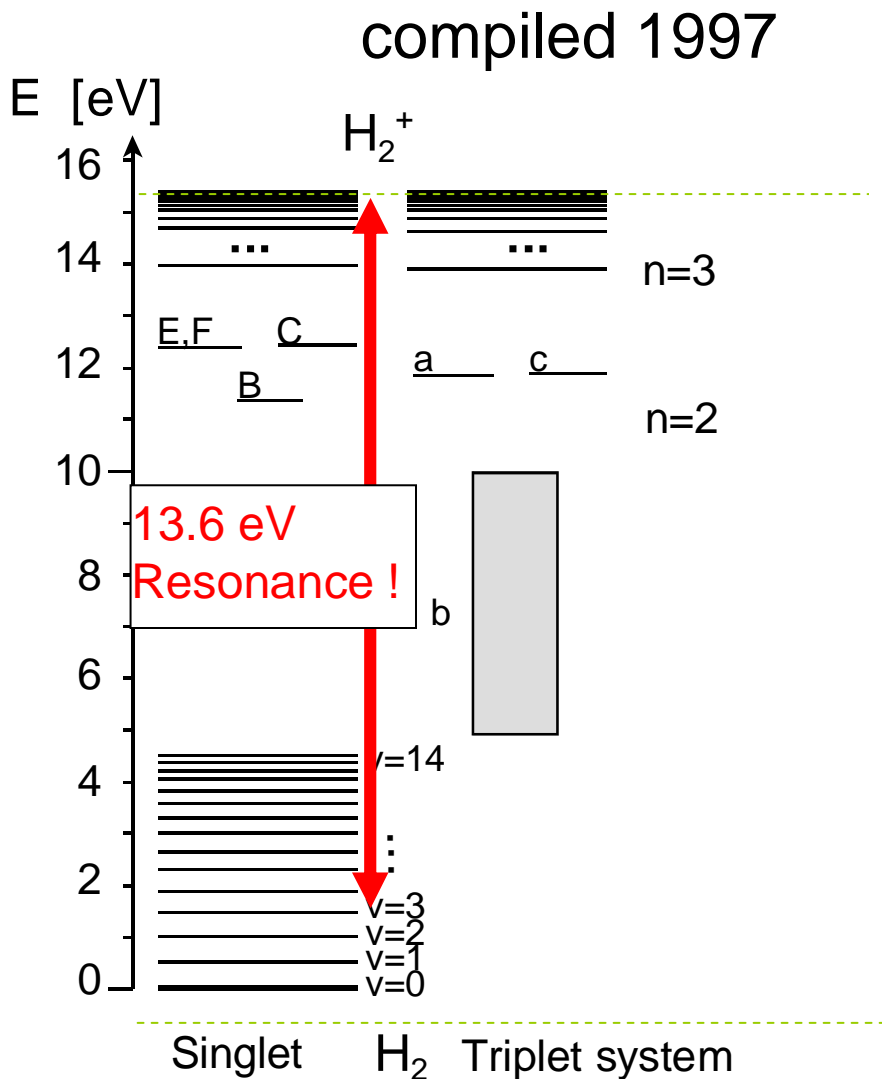




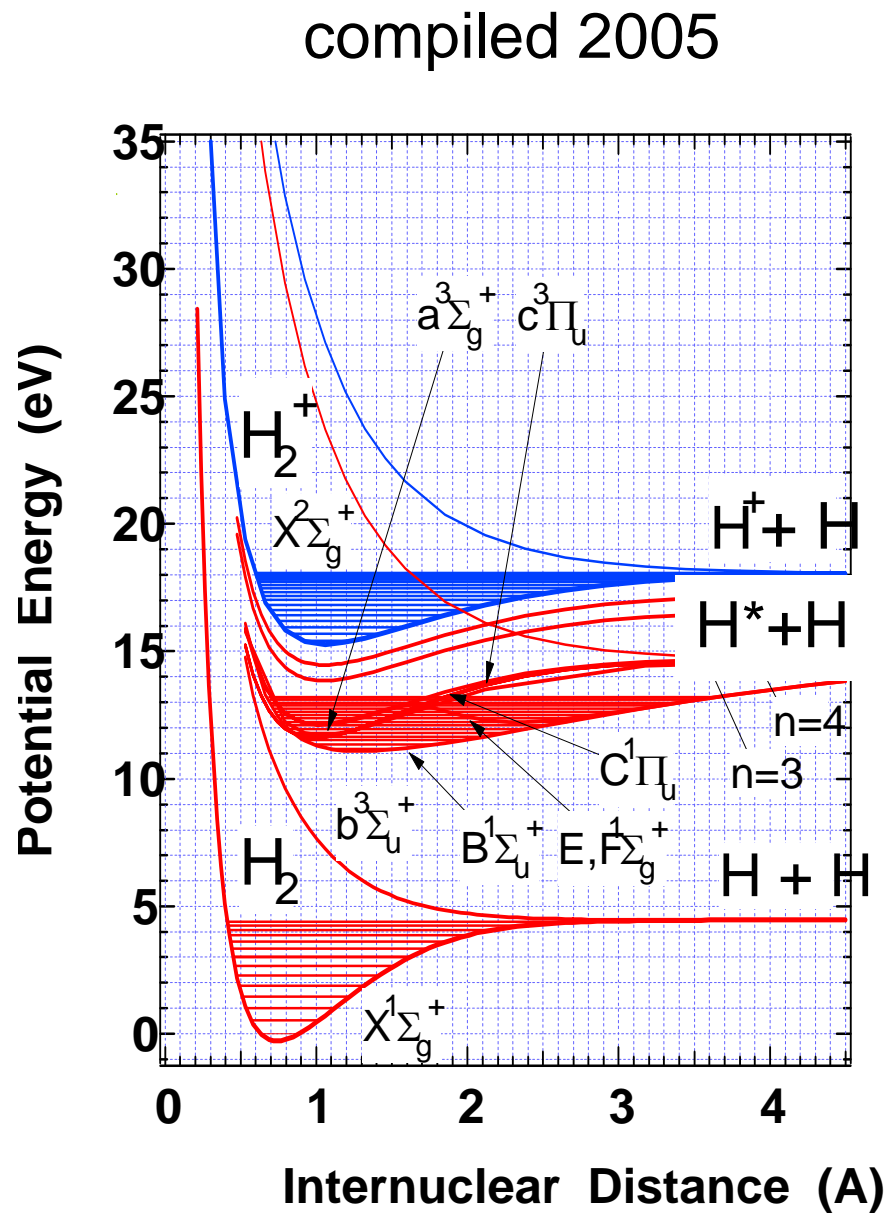
D_γ (from D , D_2 , D^+ , D_2^+):
 Profile matched, but high by factor 2
 Calibration? Atomic Data? Plasma reconstruction?

Results very sensitive eg. to T_e profile

H₂ molecule, status in present divertor code



More complete models available, still need to be integrated



Critical for Particle
Throughput (convection):

Neutral Plenum Pressure

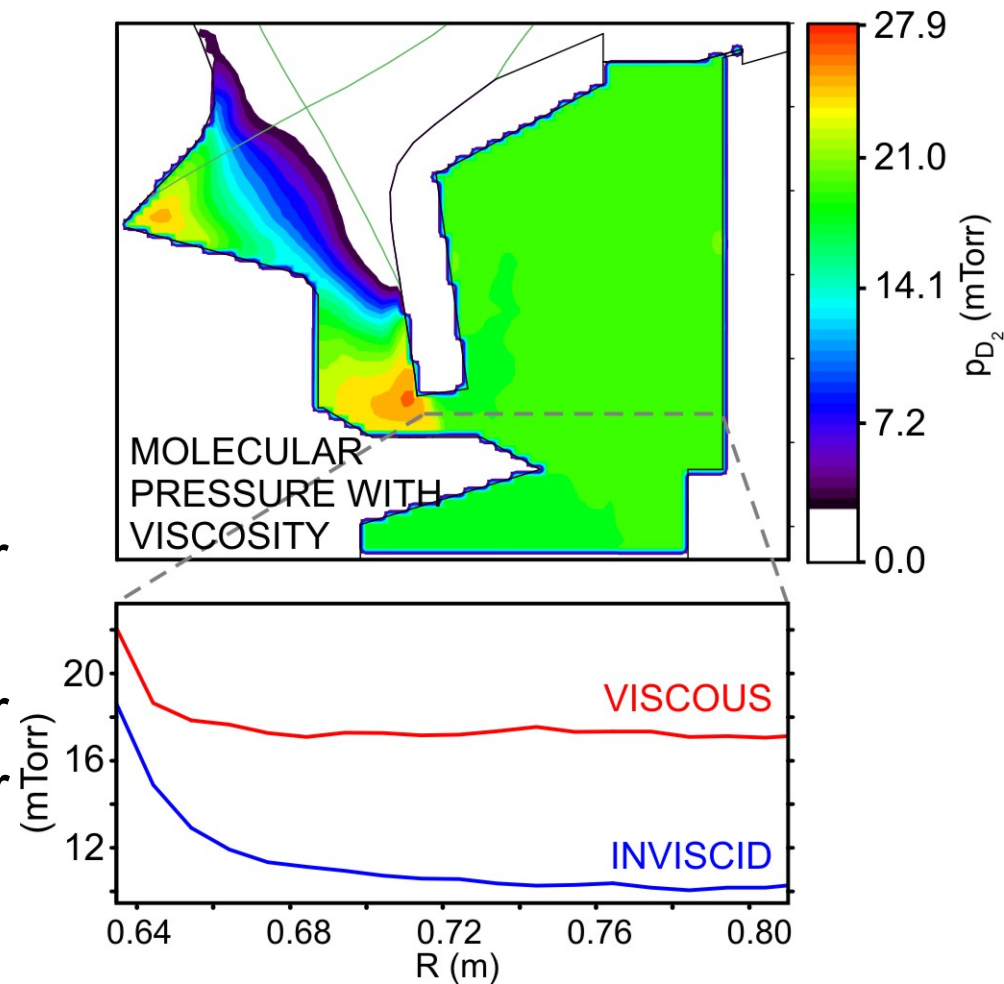
Experiment: 25 mTorr

Calc 2D (2000) 3 mTorr

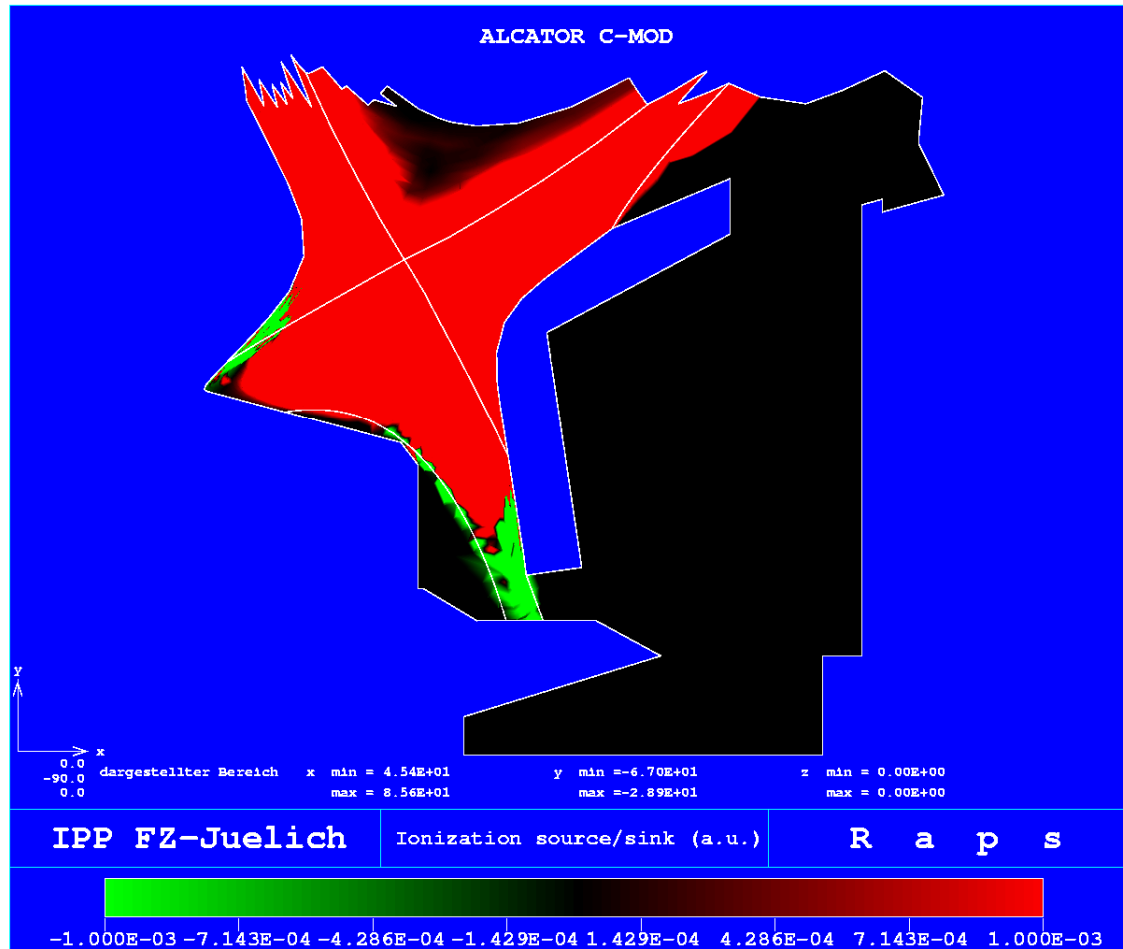
Calc 2D (2003) 27 mTorr

(better A&M data,
better transport model,
better plasma data)

Good match to experiment in 2003,
with full 2D kinetic detailed recycling model



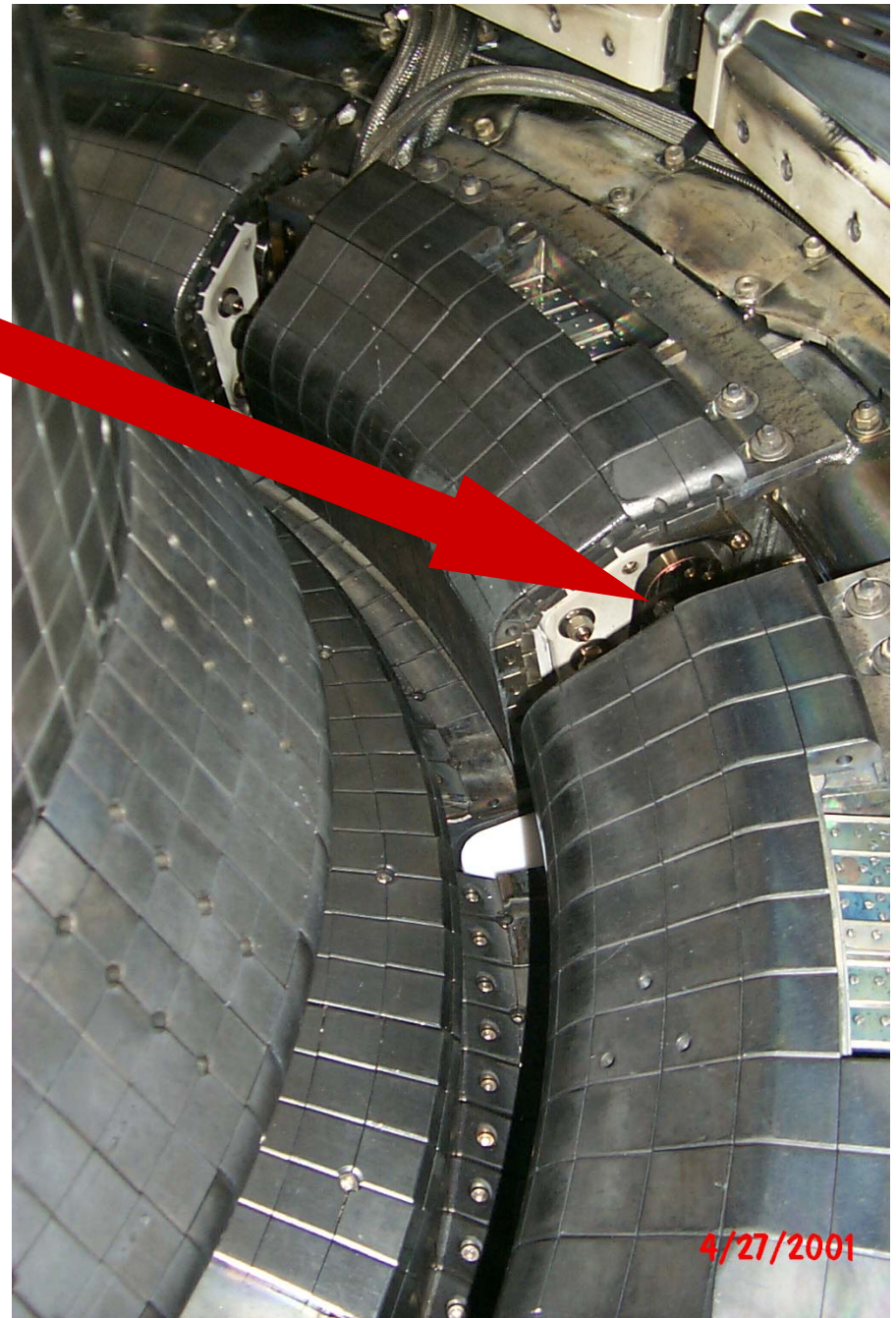
Ionizing area (red) , recombining area (green)



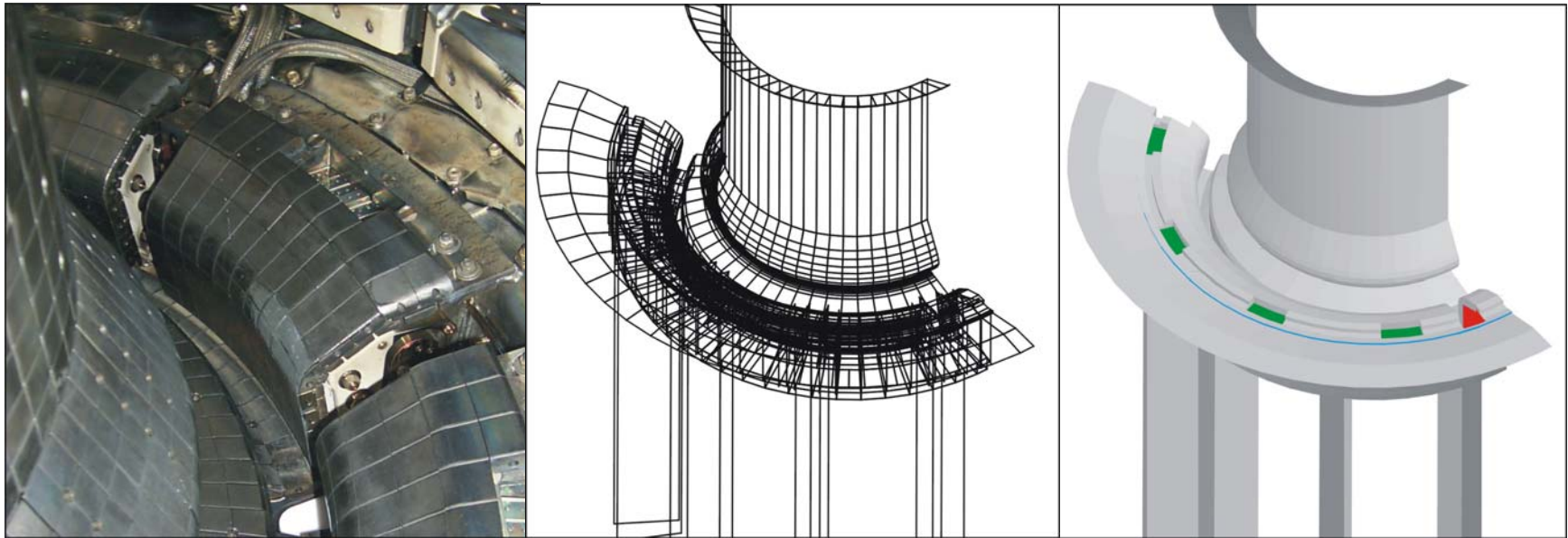
A quite cold edge plasma is sustained by recycling.
It “detaches” from the divertor targets (plasma pressure drop)

Additional
leakage pathways:

2D → 3D
(see later)



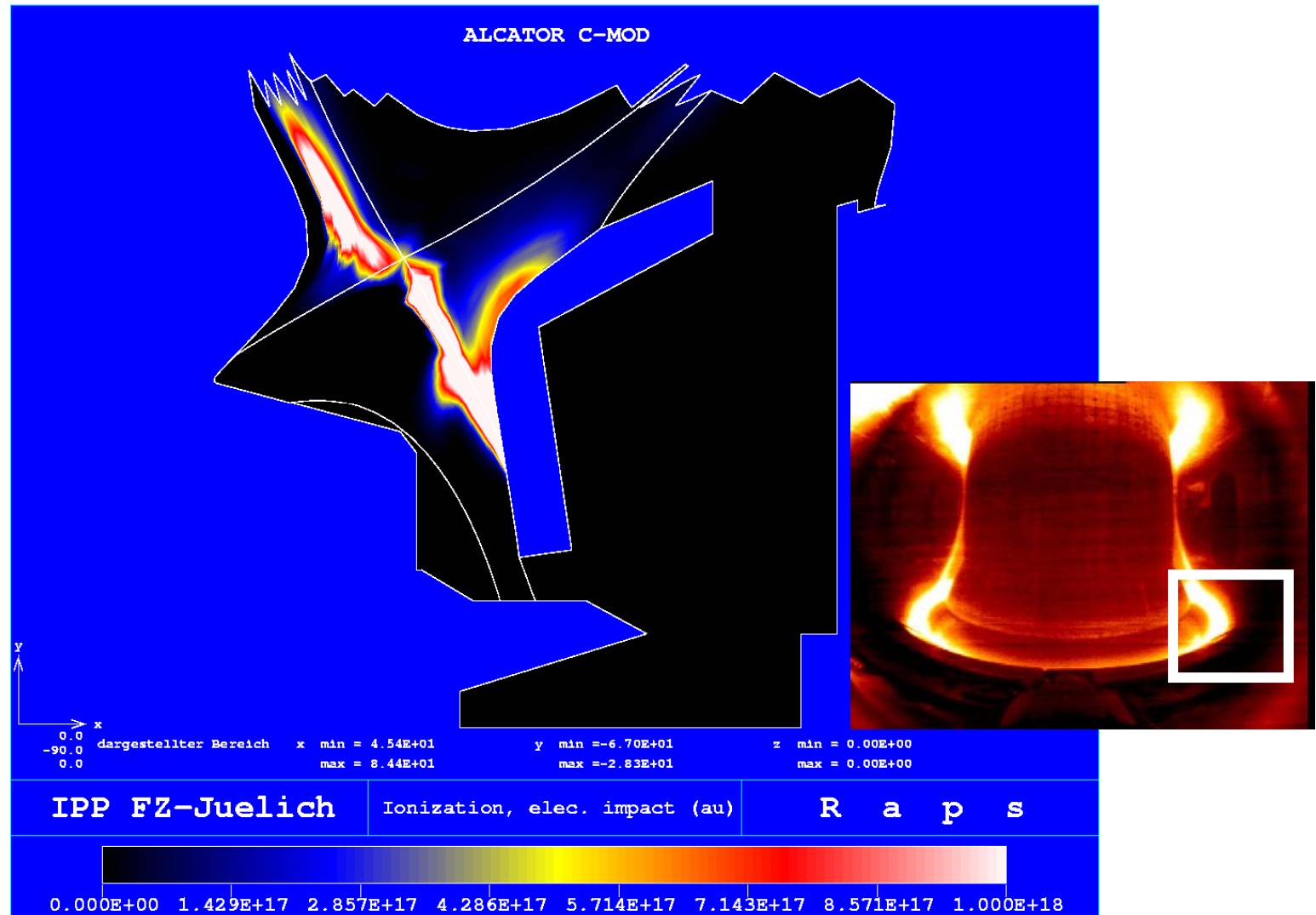
3D Neutral Gas, A&M and PSI Modelling



3D divertor structures (toroidal gap and gussets, bypass and poloidal gap)

→ strong toroidal variations in the divertor neutral pressure

Ionization by electron impact on neutral gas

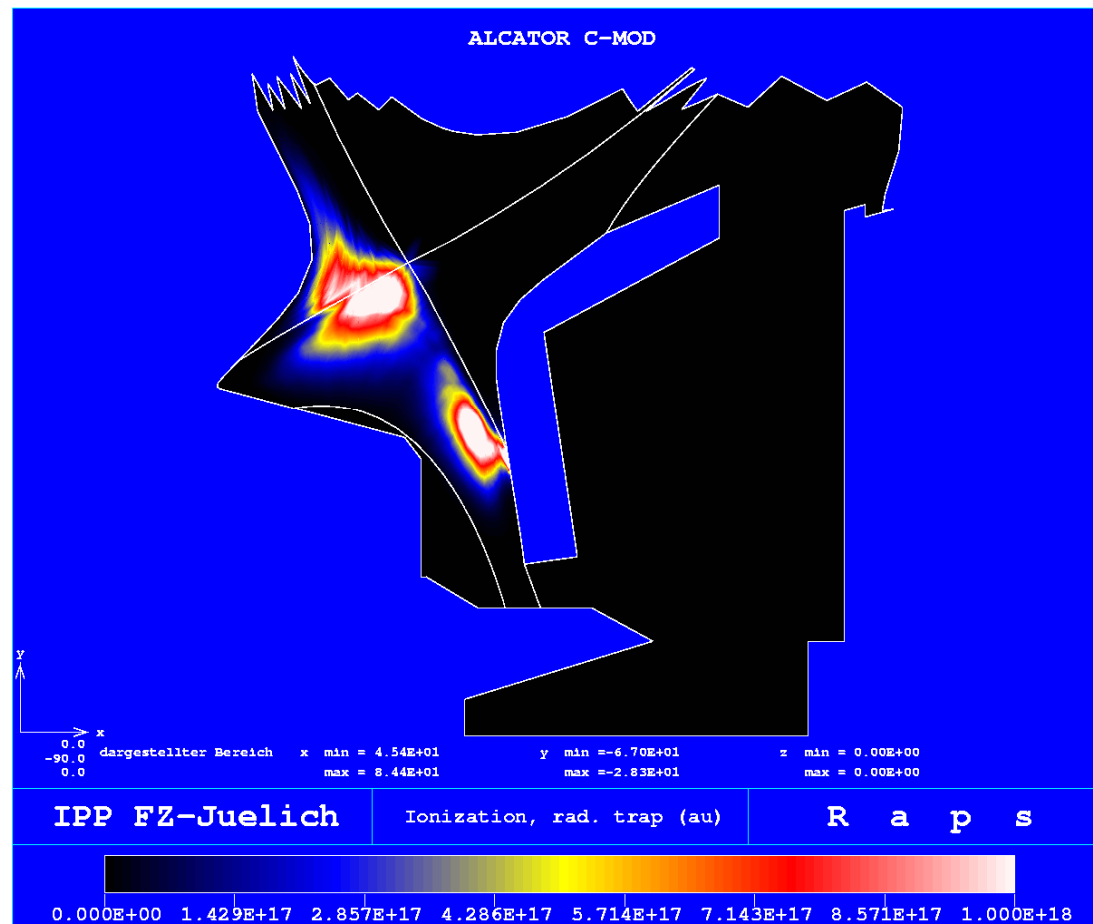


Radiation transfer: opacity of Ly-lines

(though completely elementary, has long remained unnoticed in divertor modelling)

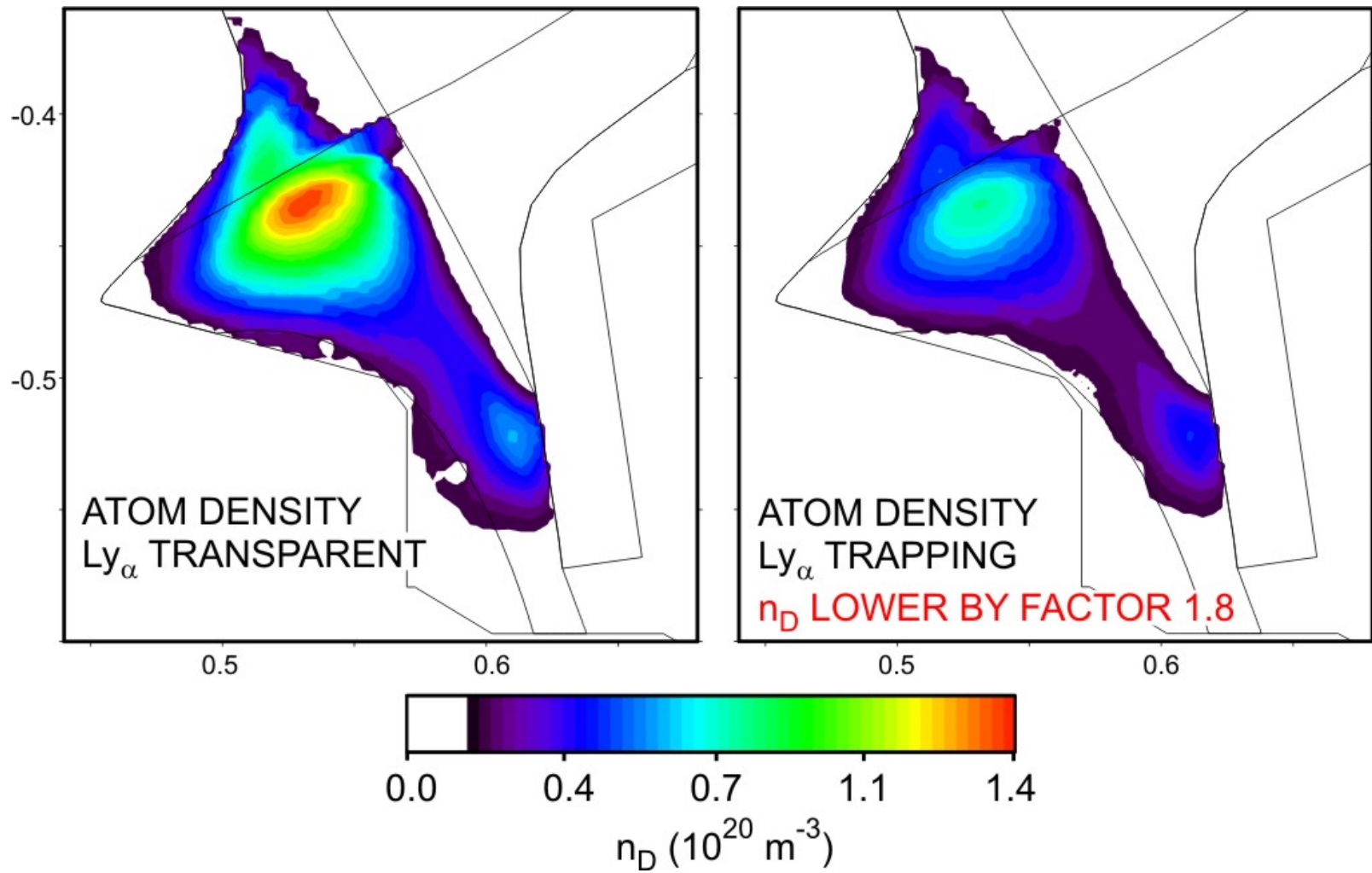


(additional path for ionization in dense, low T_e divertors)



Calc. 2003

Calc. 2004

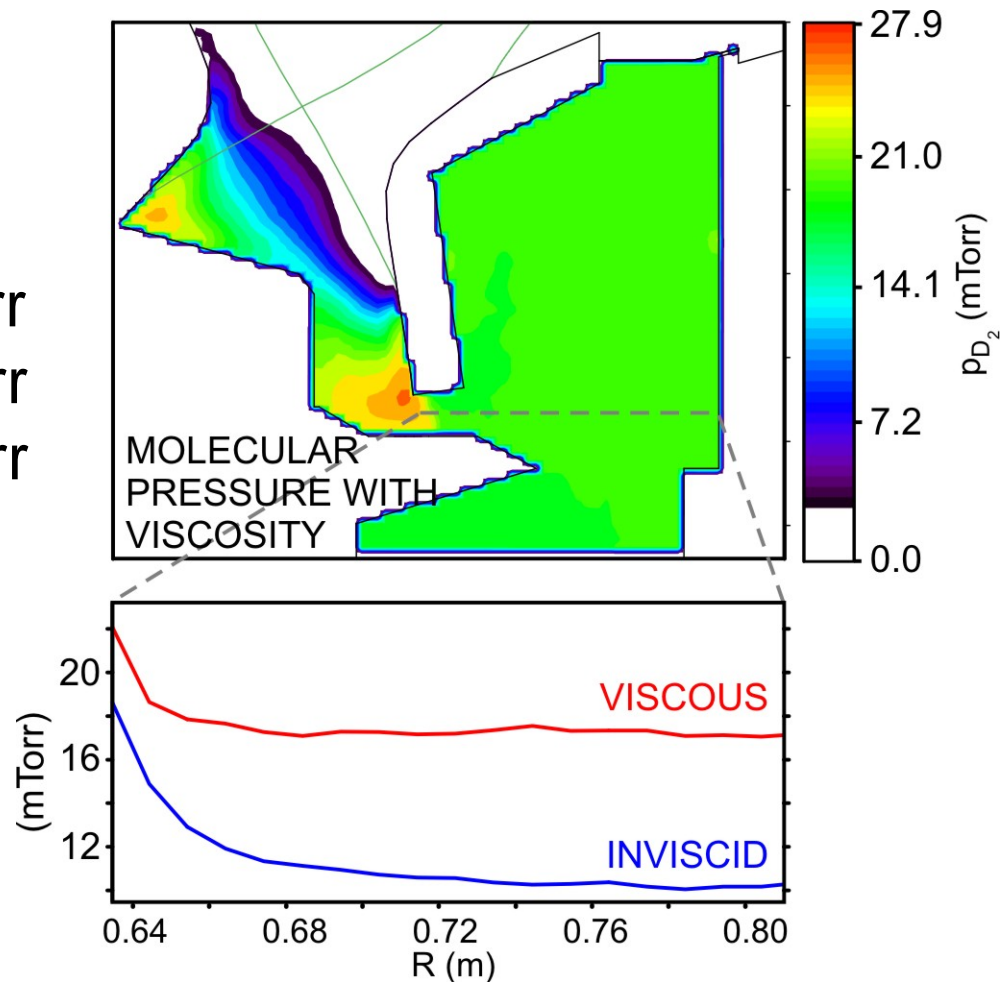


Neutral Pressure

Exp: 25 mTorr
Calc 2D (2000) 3 mTorr
Calc 2D (2003) 27 mTorr
(better A&M data,
better Plasma data)

However

Ly-opacity: 17 mTorr
3D: 11 mTorr



Model validation in the presence of many free parameters:

include ALL edge physics that we are sure must be operative even while our capability to confirm these directly remains limited

Example 2: Hydrogenic ionisation-recombination balance

The full database must contain a large number (~ several 100) of individual processes

Fortunately: very different timescales are involved (numerically stiff problem):
⇒ an underlying reduced model exists.

This is often referred to as:

Collisional radiative models (Astrophysics)

Multistep ionisation and recombination models (Plasma physics)

Condensed case approximations (Semiconductor physics)

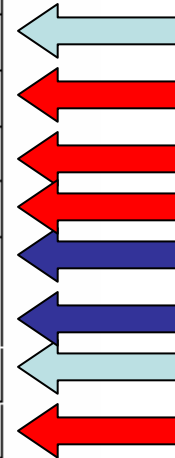
Lumped species concepts (Atmospheric research)

Intrinsic low dimensional manifolds (combustion flames)

Atomic and Molecular Database for EIRENE:
 see: www.eirene.de/htmlseiten/amjuel/

ionisation – recombination – transport balance: atoms

multi	direct	Path	Comment	File
gscr		$e + H(n=1) \rightarrow H^+$	via H^*	fort.10 (sum)
gscr2		$e + H(n=2) \rightarrow H^+$	via H^*	fort.101 (sum)
gscr21		$e + H(n=2) \rightarrow H(n=1)$	via H^*	fort.102 (sum)
gscr12		$e + H(n=1) \rightarrow H(n=2)$	via H^*	fort.103 (sum)
alcr-bel alcr	bel	$e + H^+ \rightarrow H(n=1)$	rad.rec (2.1.8)	fort.11.1
		$e + H^+ \rightarrow H(n=1)$	via H^*	fort.11.2
		$e + H^+ \rightarrow H(n=1)$		fort.11.3 (sum)
alcr2		$e + H^+ \rightarrow H(n=2)$		fort.111 (sum)



conventional multistep model

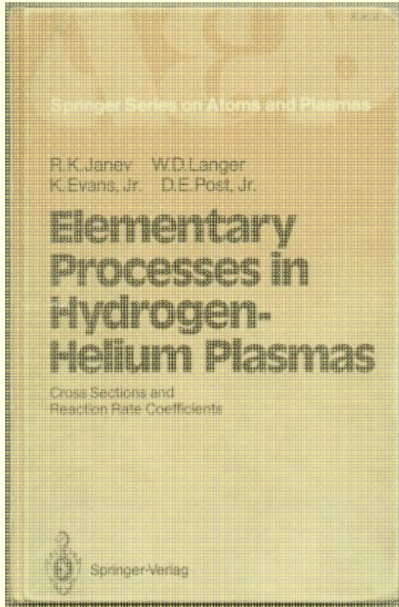


Extensions for Lyman- α radiation transfer (photoexcitation)



Extension for (future) photoionisation simulations

ionisation – recombination - transport : molecules



+

Sawada/
Fujimoto
CR-Model

multi	direct	Path	Comment	File
gscrh2	bq1	$e + H_2 \rightarrow H_2^+$ $e + H_2 \rightarrow H_2^+$	direct (2.2.9) via H_2^*	fort.12 (sum)
	bq7	$p + H_2 \rightarrow H_2^+ + H$	from $H_2(v=0)$	
	bq8	$p + H_2 \rightarrow H_2^+ + H$	from $H_2(v \geq 1)$	
	bq6	$e + H_2 \rightarrow H + H^-$	from $H_2(v=0)$	
	bq5	$e + H_2 \rightarrow H + H^-$	from $H_2(v \geq 1)$	
gs2c	bq2	$e + H_2 \rightarrow p + H(n=1)$ $e + H_2 \rightarrow p + H(n=1)$ $e + H_2 \rightarrow p + H(n=1)$ $e + H_2 \rightarrow p + H(n=1)$ $e + H_2 \rightarrow p + H(n=1)$	direkt (2.2.10) via $H_2 \rightarrow H + H^*$ via $H_2^* \rightarrow p + H$ via $H_2^* \rightarrow H + H^*$	fort.14 (sum)
missing	—	$e + H_2 \rightarrow p + H(n=2)$	above H_2^+	—
missing	h1s (bq4?)	$e + H_2 \rightarrow H + H(n=1)$ $e + H_2 \rightarrow H + H(n=1)$ $e + H_2 \rightarrow H + H(n=1)$ $e + H_2 \rightarrow H + H(n=1)$ $e + H_2 \rightarrow H + H(n=1)$	direkt (2.2.5) via $H_2 \rightarrow H + H^*$ via $H_2^* \rightarrow H + H$ via $H_2^* \rightarrow H + H^*$	fort.13 (sum)
total	h2s	$e + H_2 \rightarrow H + H(n=2)$ $e + H_2 \rightarrow H + H(n=2)$ $e + H_2 \rightarrow H + H(n=2)$ $e + H_2 \rightarrow H + H(n=2)$ $e + H_2 \rightarrow H + H(n=2)$	direkt via $H_2^* \rightarrow H + H(n=2)$ via $H_2 \rightarrow H + H^*$ via $H_2^* \rightarrow H + H^*$	fort.131 (sum)
missing	—	$e + H_2^+ \rightarrow p + p$	direkt (2.2.11)	fort.16 (sum)
gs2e	gstq1	$e + H_2^+ \rightarrow p + p$	via $p + H^*$	fort.16 (sum)
missing	bq3	$e + H_2^+ \rightarrow p + H(n=1)$ $e + H_2^+ \rightarrow p + H(n=1)$ $e + H_2^+ \rightarrow p + H(n=1)$ $e + H_2^+ \rightarrow p + H(n=1)$	direkt (2.2.12) via $H + H^*$ via $p + H^*$	fort.15.1 fort.15.2 fort.15.3
total	h3s	$e + H_2^+ \rightarrow p + H(n=2)$ $e + H_2^+ \rightarrow p + H(n=2)$	direkt via $p + H^*$	fort.151 (sum)
missing	—	$e + H_2^+ \rightarrow H + H(n=1)$ $e + H_2^+ \rightarrow H + H(n=1)$	direkt: not possible via $H + H^*$ (2.2.14)	fort.17
missing	h4s	$e + H_2^+ \rightarrow H + H(n=2)$ $e + H_2^+ \rightarrow H + H(n=2)$	direkt via $H + H^*$ (2.2.14)	fort.171
missing	gstq2mm	$e + H_2^+ \rightarrow H + H(n=2)$	via $H + H^*$ (2.2.14)	fort.171
missing	bq9	$e + H^- \rightarrow H + 2e$	direkt, (7.1.1)	
missing	gstq3	$p + H^- \rightarrow p + H$ $p + H^- \rightarrow p + H$	direkt (7.2.1) via $H + H^*$	
missing	gstq3m	$p + H^- \rightarrow H + H(n=1)$ $p + H^- \rightarrow H + H(n=1)$	direkt via $H + H^*$	
missing	h5s	$p + H^- \rightarrow H + H(n=2)$ $p + H^- \rightarrow H + H(n=2)$	direkt (7.2.2) via $H + H^*$	
missing	gstq3mm	$p + H^- \rightarrow H + H(n=2)$	via $H + H^*$	

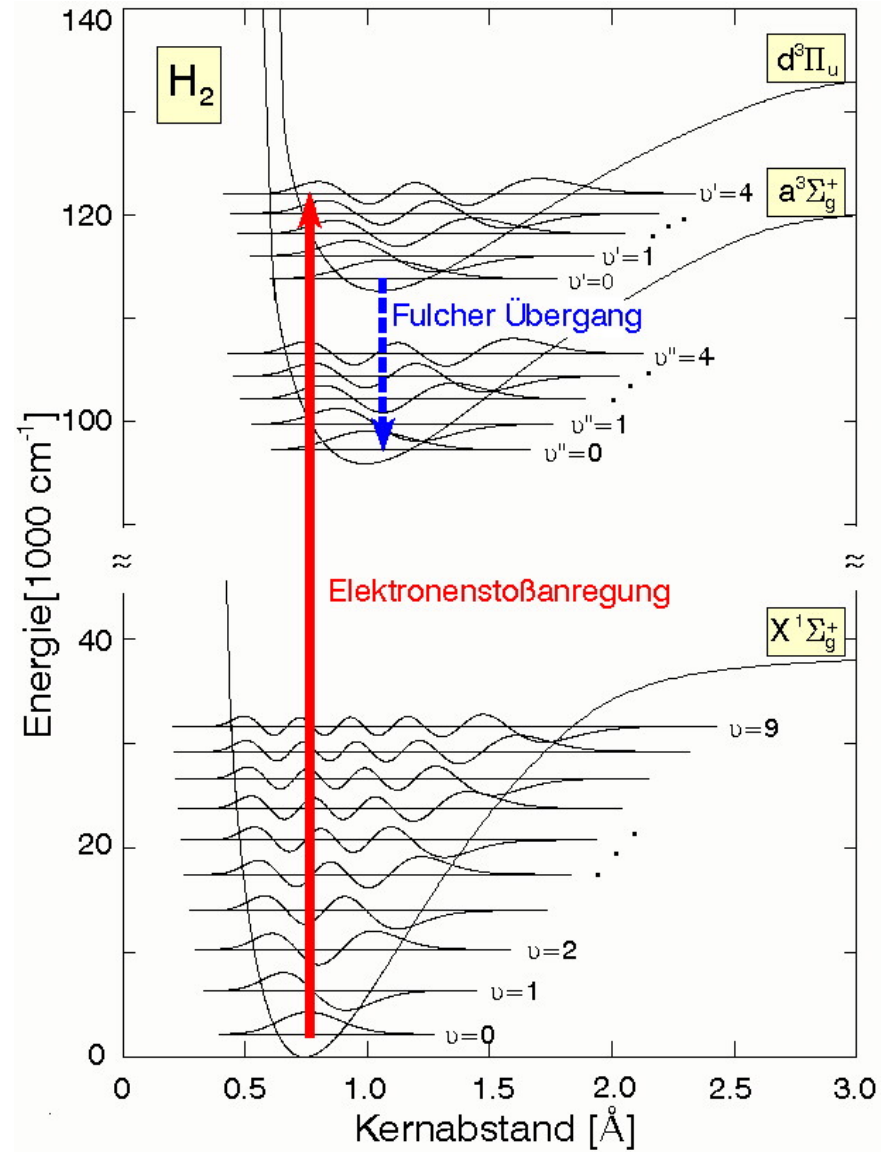
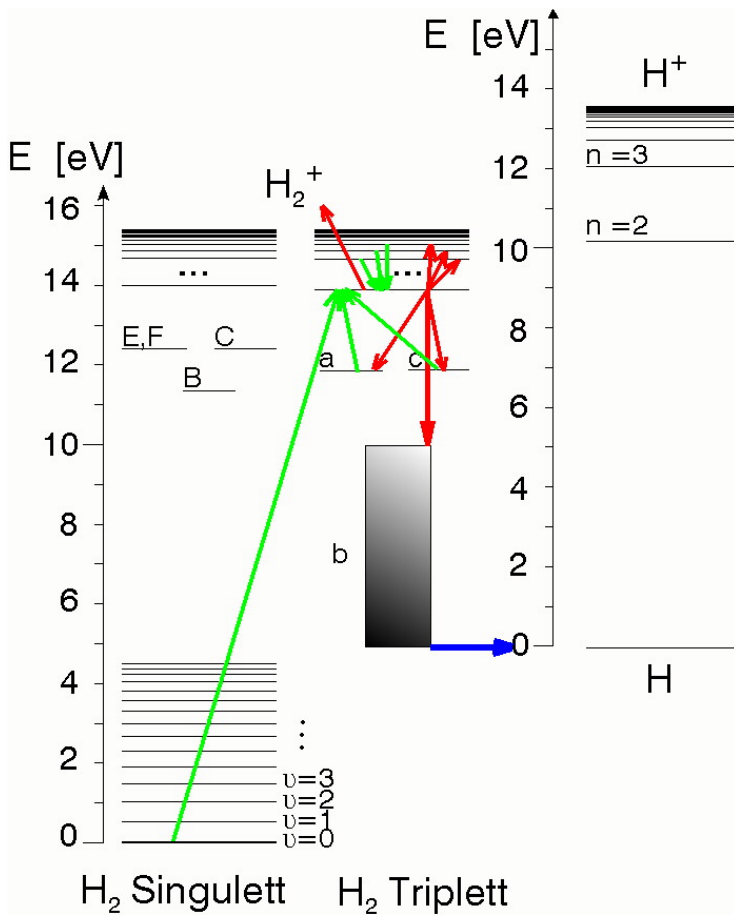
(MAR₁)

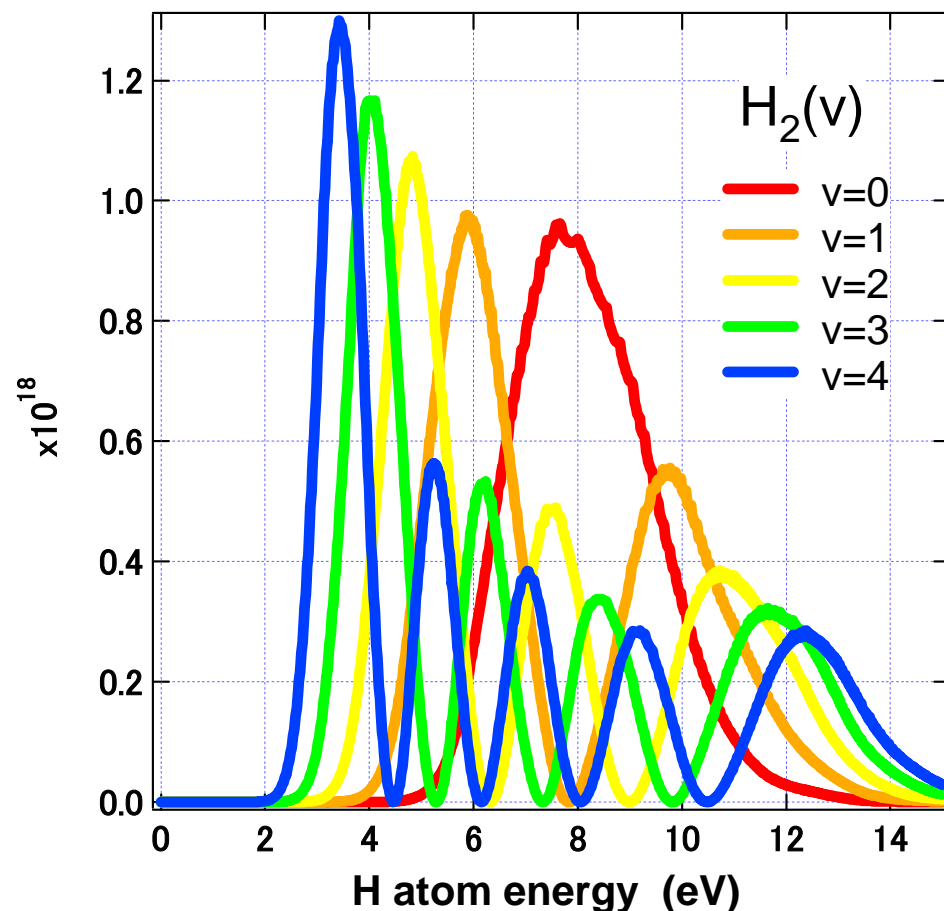
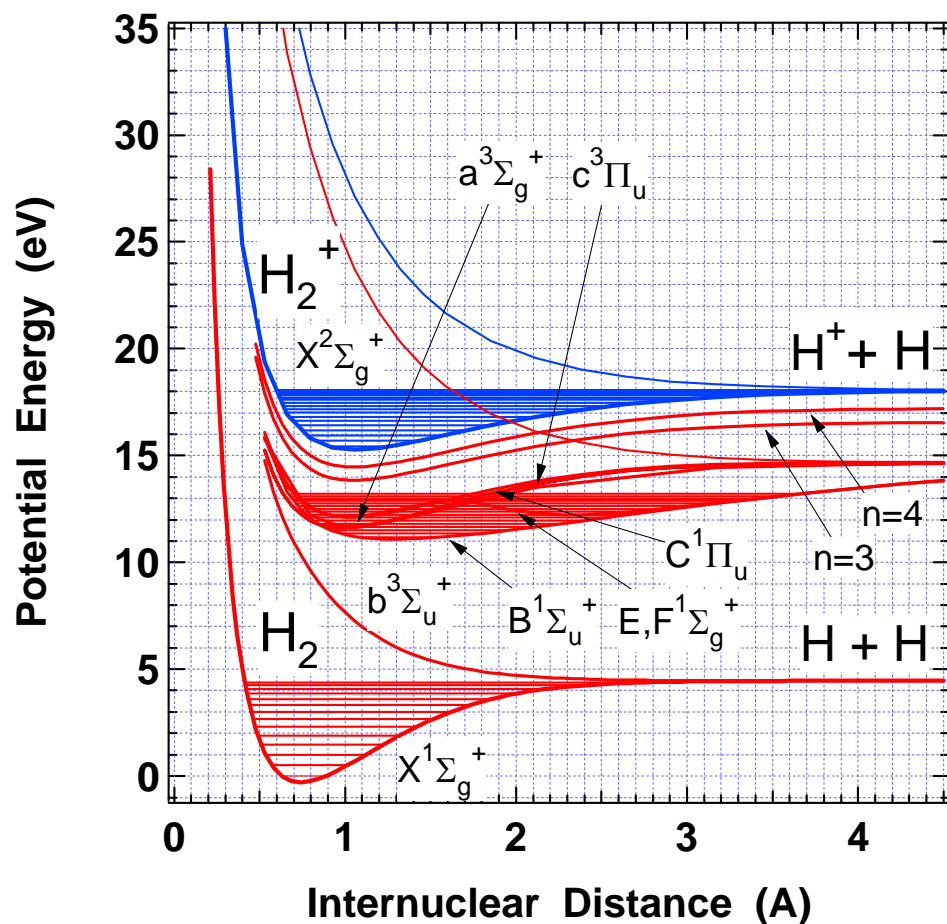
(MAD₂)

(MAR₂)

Molecule spectroscopy (late nineties): identification of important role of molecule reaction kinetics on overall Divertor dynamics

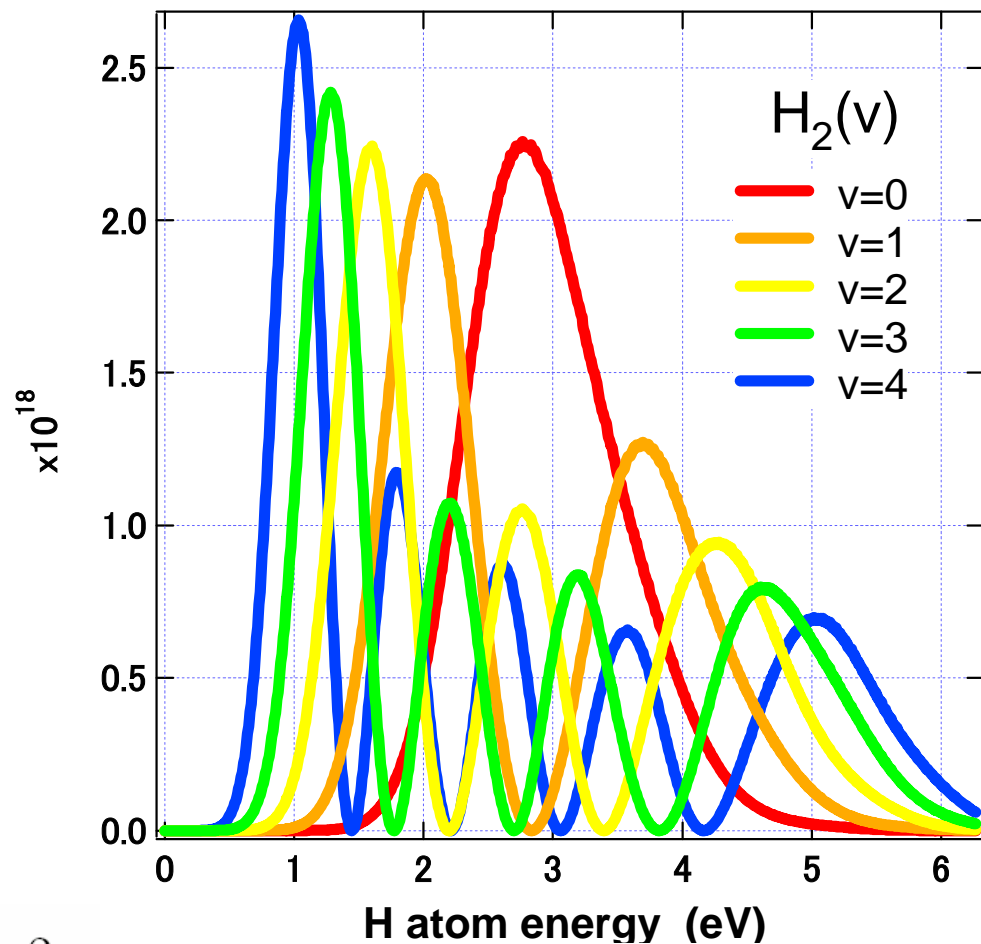
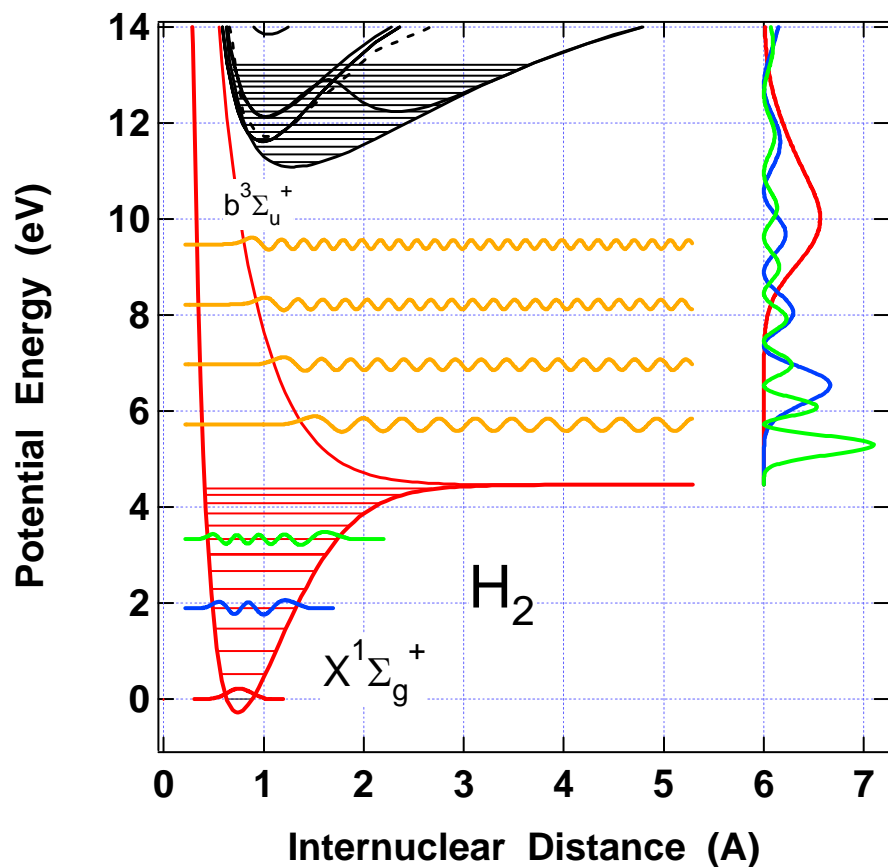
U. Fantz et al.





Calculation of H atom velocity

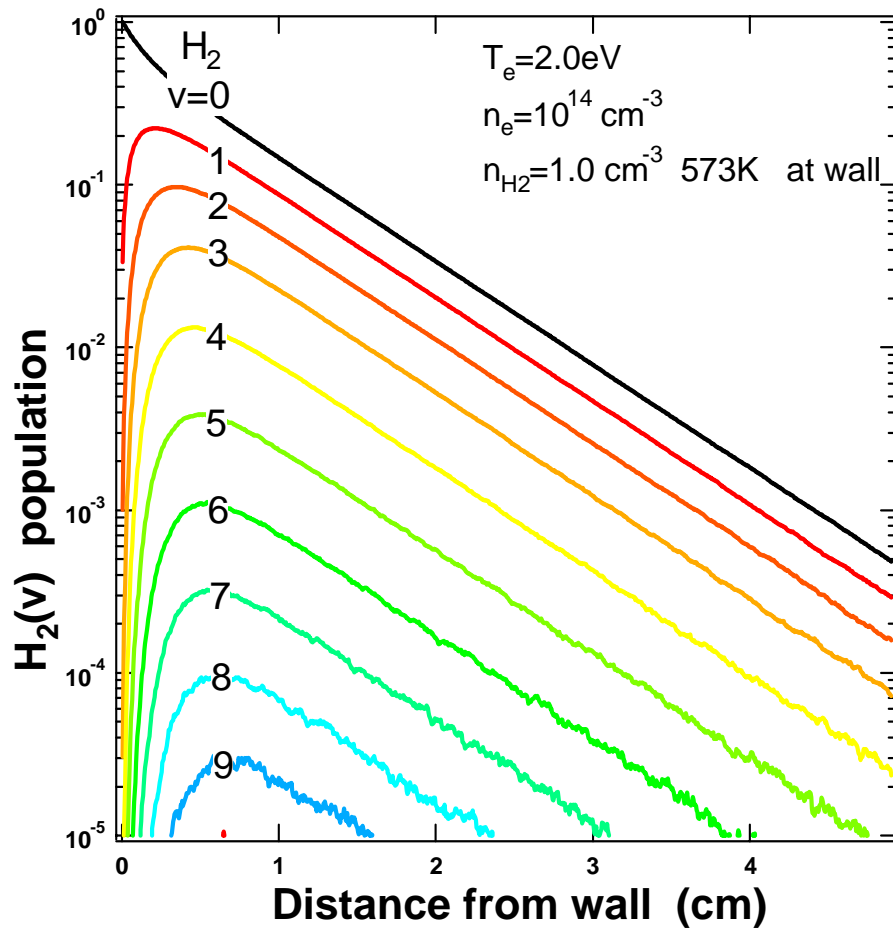
$H_2(X) \rightarrow H_2(b) \rightarrow H + H$



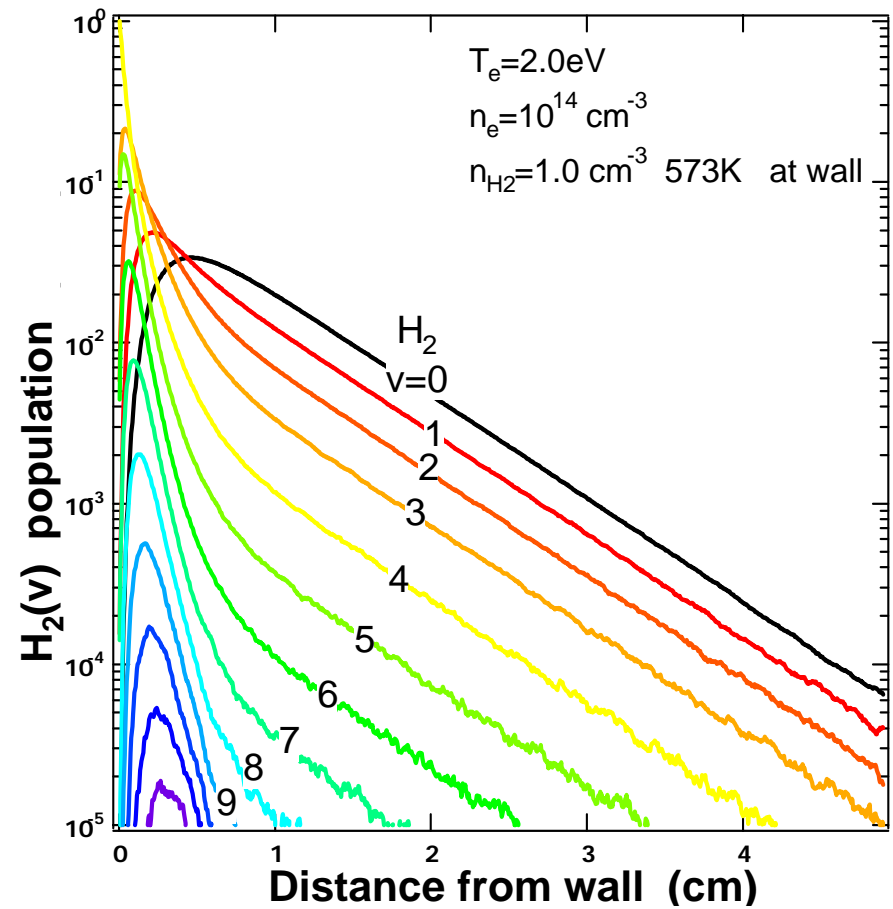
$$q_{\kappa v} = \left| \int \chi_{\kappa}^b(R) \chi_v^X(R) dR \right|^2$$

Sensitivity to surface produced vibr. excitation

1-D Monte-Carlo Model of Neutral-Particle Transport (K. Sawada)

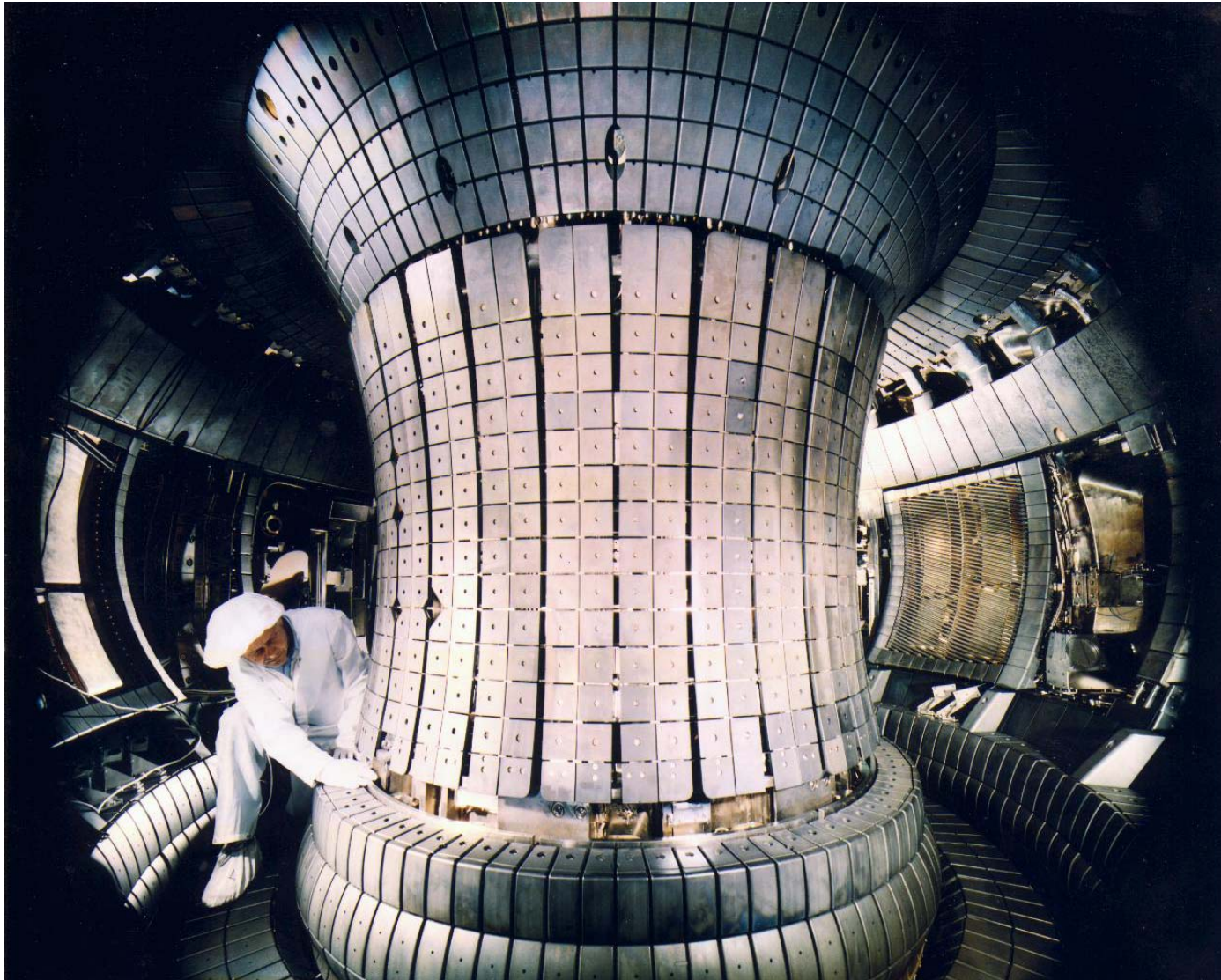


All H_2 from wall: $v=0$



All H_2 from wall: $v=4$

ASDEX-UPGRADE (IPP Garching)





Trapping of neutral particles in the Divertor: high recycling and detachment regime

• COUPLING EIRENE TO BRAAMS CODE: ASDEX UPGRADE SINGLE-NULL

SCALING FACTORS

FACT-X= 1.500E+02

FACT-Y= 1.500E+02

ORIGIN

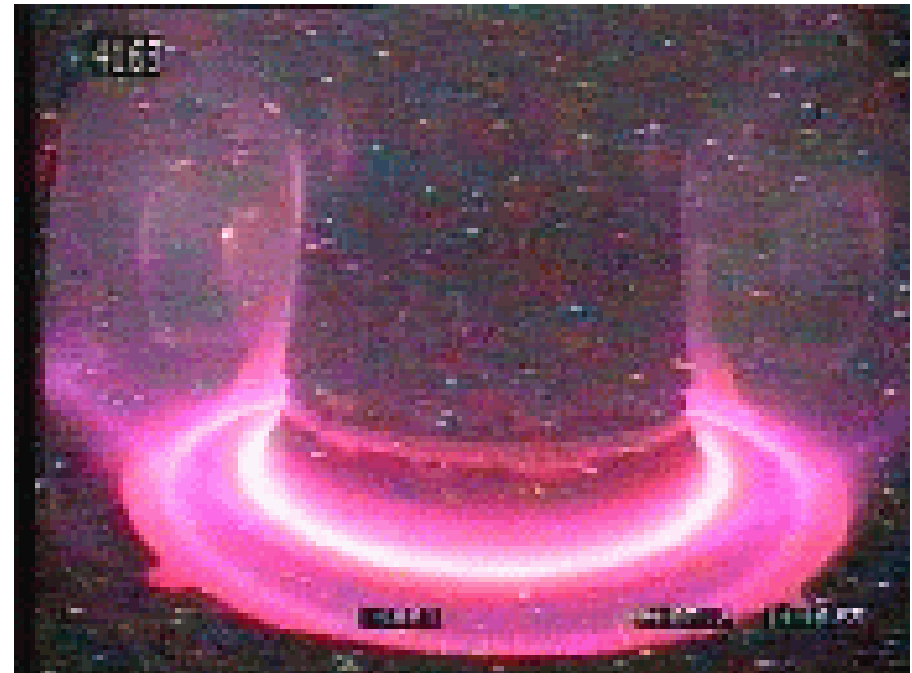
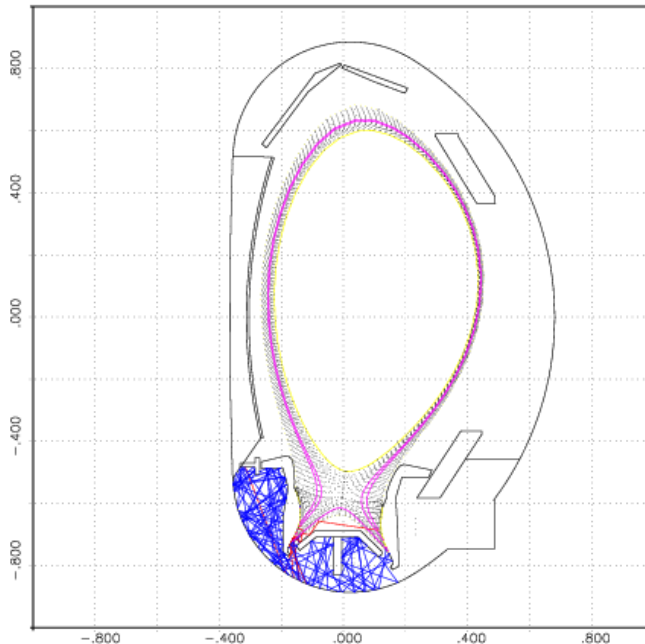
CH2X0= 1.500E+02

CH2Y0= 0.000E+00

PLOTTED AT

Z = 0.000E+00

- + LOCATE(1)
- ★ ELECTR. IMPACT(2)
- ELASTIC COLL.(3)
- × CHARGE EXCHANGE(4)
- FOKKER PLANCK(5)
- ▲ SURFACE(6)
- ◀ SPLITTING(7)
- ▼ RUSSIAN ROULETTE(8)
- ▶ PERIODICITY(9)
- ▶ RESTART-A. SPLT.(10)
- ▶ SAVE:COND. EXP.(11)
- ▶ RESTART:COND EXP(12)
- △ TIME LIMIT(13)
- ◇ GENERATION LIMIT(14)
- ★ FLUID LIMIT(15)
- ERROR DETECTED



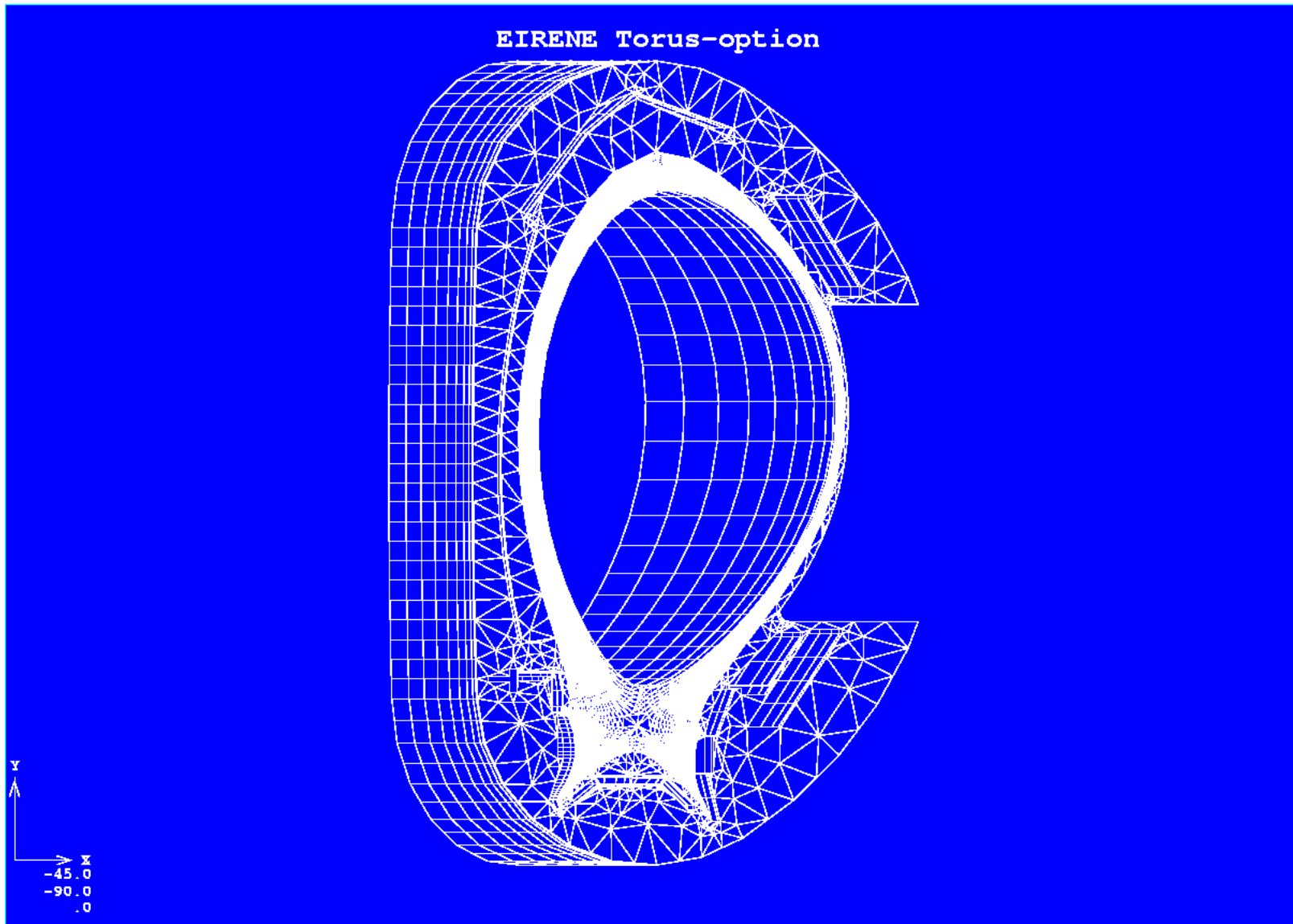
Particle Simulation: PMI, A&M

Visible light from ASDEX-U Divertor





EIRENE Torus-option



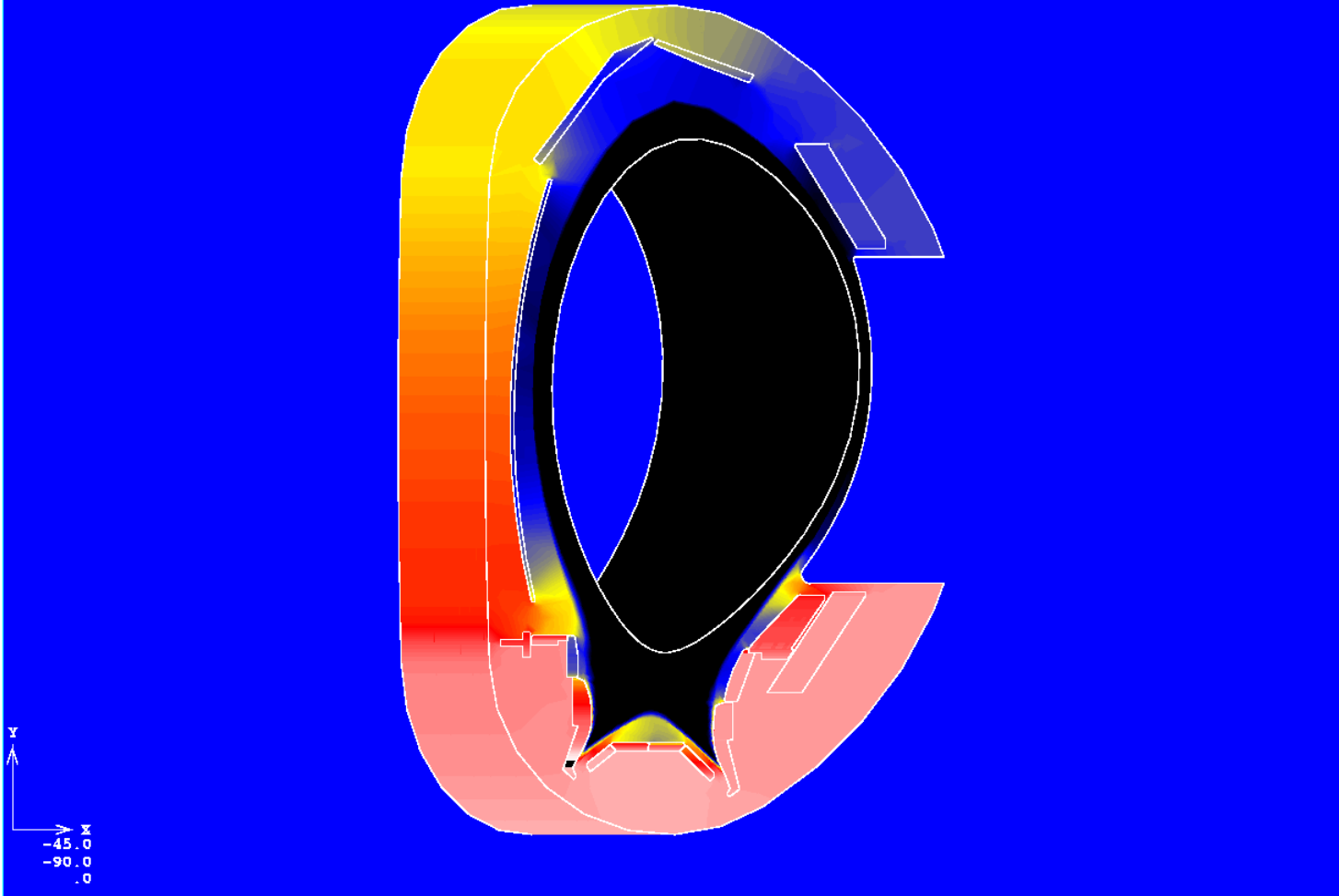
y
-45.000
-90.000
0.000

Fz-Juelich

R A P S



EIRENE Graphics: ASDEX-U (ref. 5c)



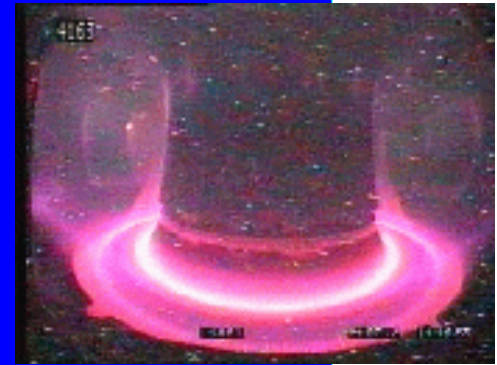
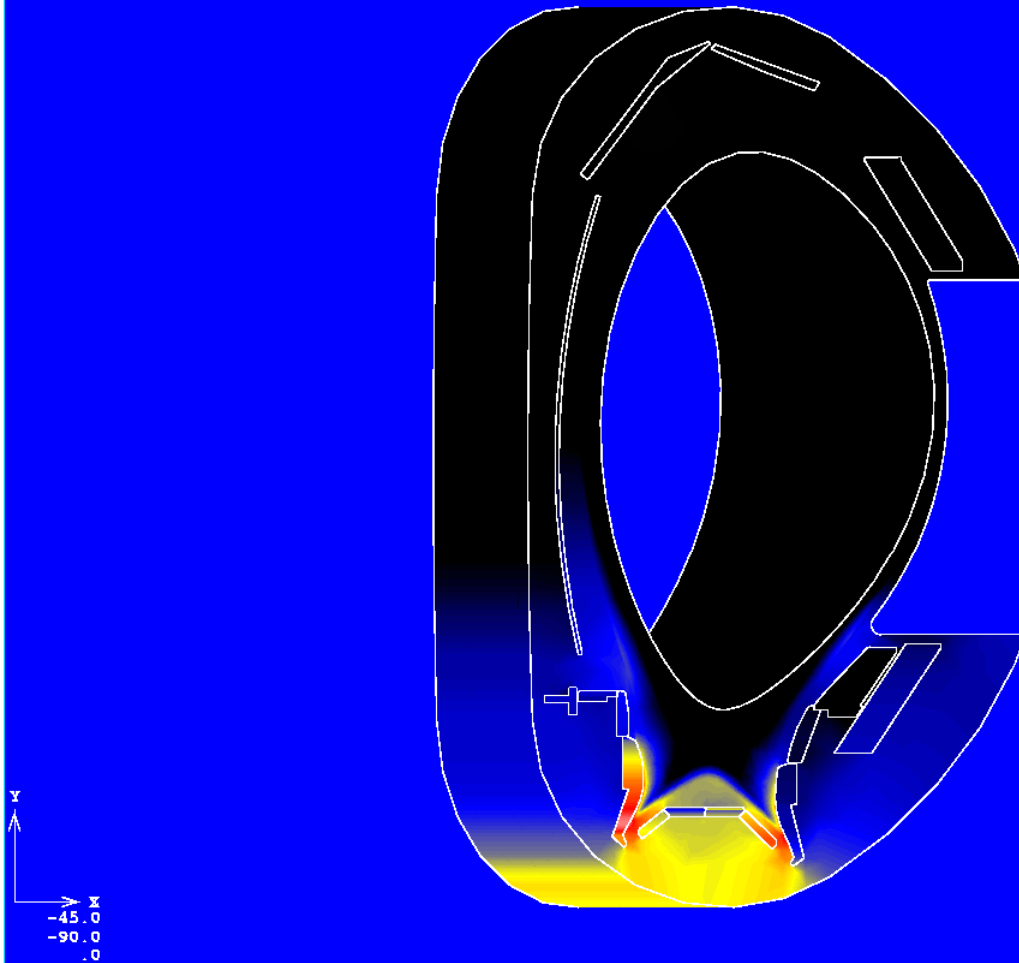
Fz-Juelich

Molecule density (cm**⁻³) log

R A P S

1.000E+01 1.057E+01 1.114E+01 1.171E+01 1.229E+01 1.286E+01 1.343E+01 1.400E+01

EIRENE Graphics: ASDEX-U (ref. 5c)



Fz-Juelich

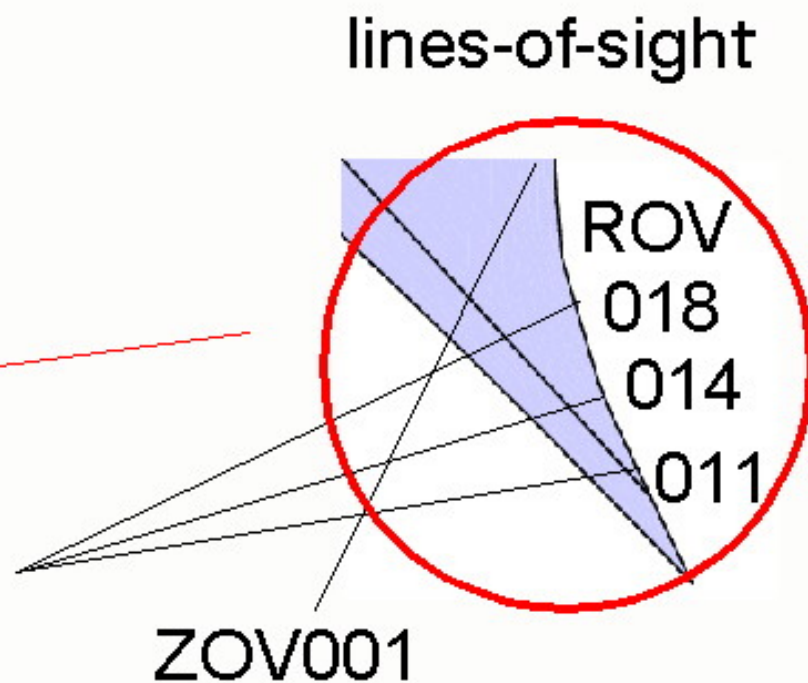
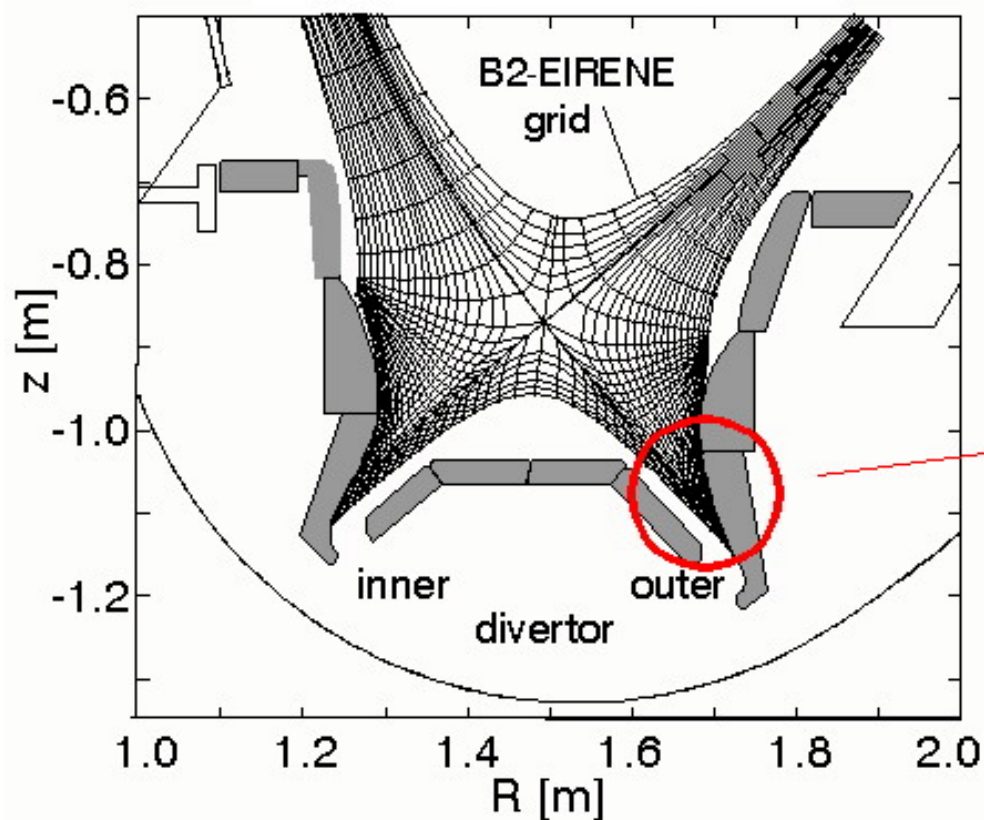
Atom density (cm⁻³) log

R A P S

1.000E+01 1.057E+01 1.114E+01 1.171E+01 1.229E+01 1.286E+01 1.343E+01 1.400E+01

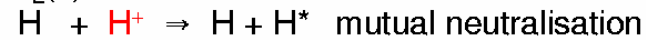
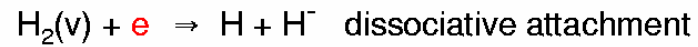
Molecular spectroscopy to verify or disprove model predictions

ASDEX Upgrade Lyra divertor

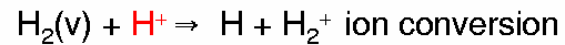


Radiation in divertor plasmas

Molecular Assisted Recombination (MAR)



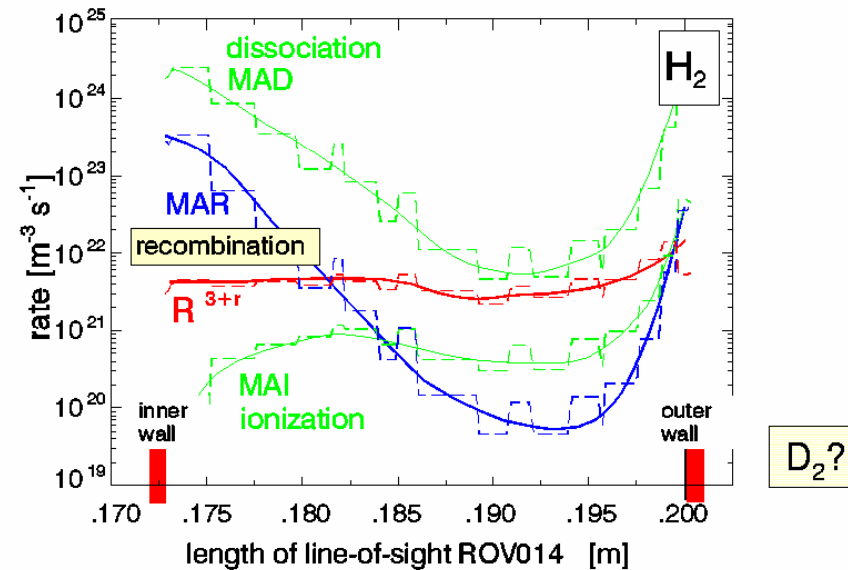
or



← dominant channel



⇒ combination of spectroscopy, CR-model and B2-EIRENE:

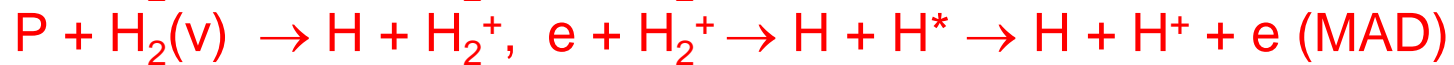
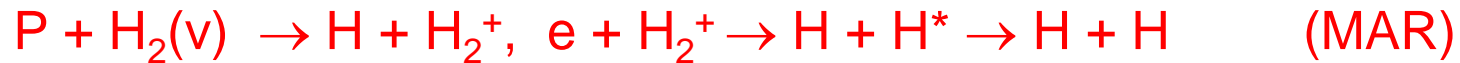


molecules contribute near the surfaces to plasma recombination

- dependence of $\text{H}_2(\text{v})$ on wall material and temperature
- recombination of atoms at surfaces and reflection coefficients



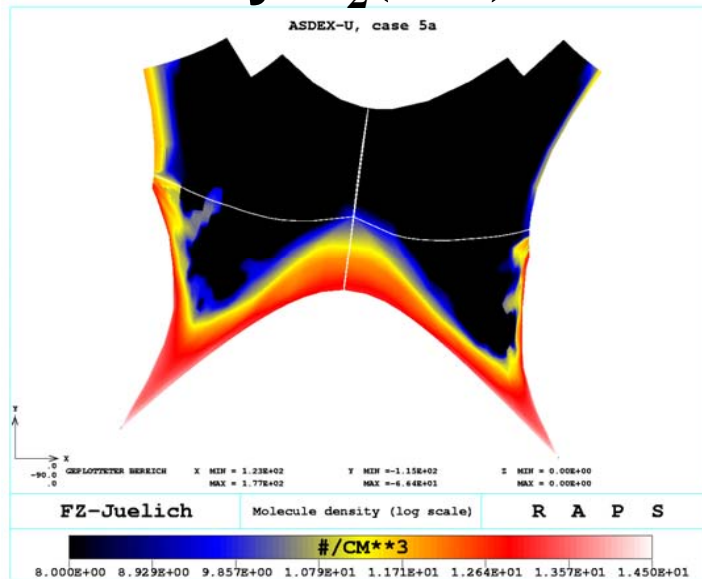
Self sustained neutral cushion? B2-EIRENE simulation: MAR or MAD?



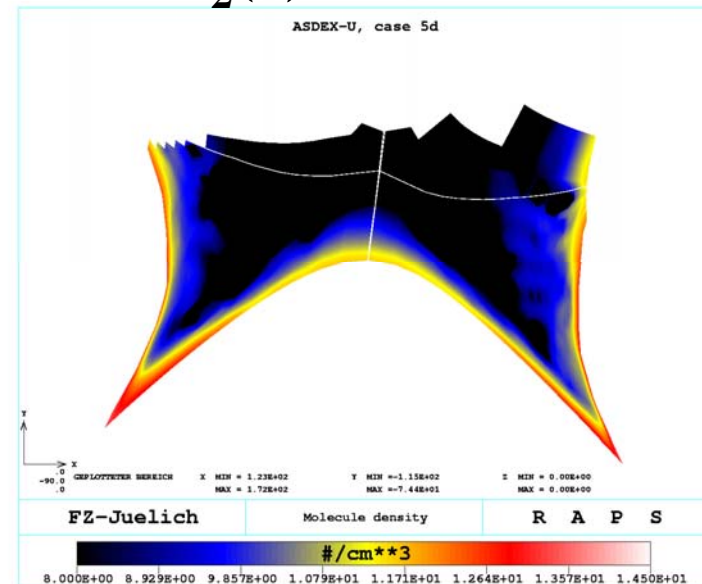
H₂ density field in Divertor:

Neutral gas “cushion” is strongly reduced by H₂ + p ion conversion

only H₂(v=0)



H₂(v) included



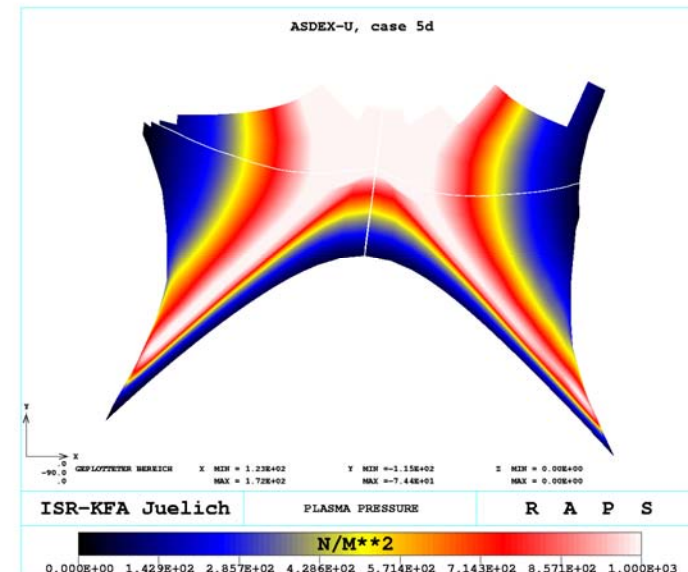
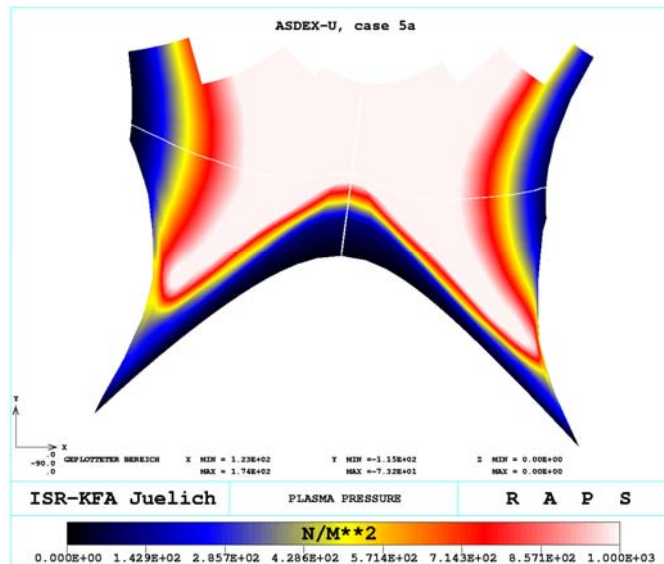


Self sustained neutral cushion? B2-EIRENE simulation: MAR or MAD?

Pressure drop at inner target disappears:
re-attachment due to MAD

$H_2(v=0)$

$H_2(v)$

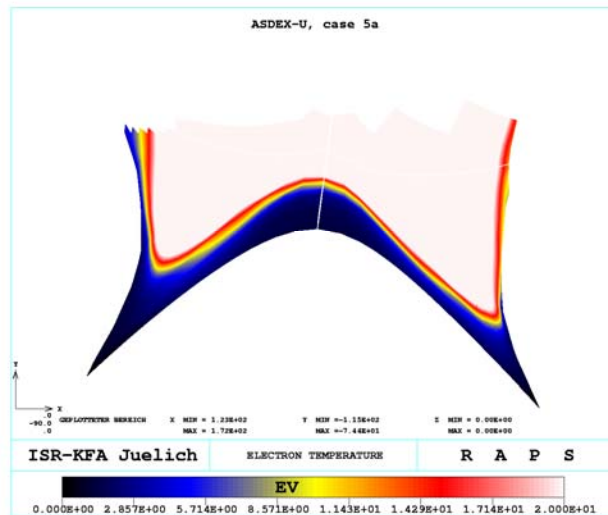




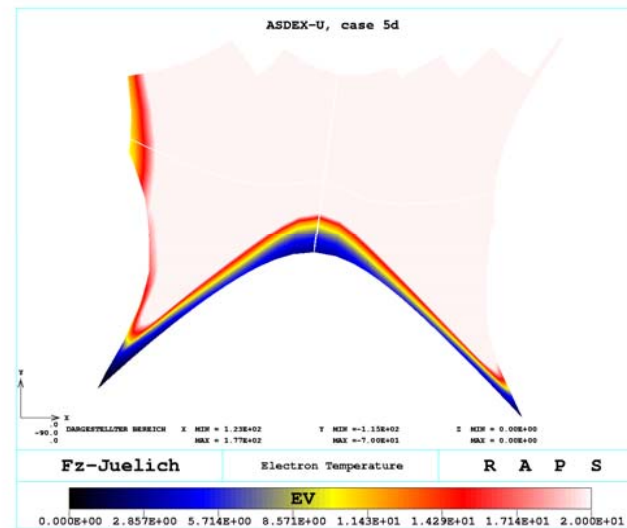
Self sustained neutral cushion? B2-EIRENE simulation: MAR or MAD?

Electron temperature field in Divertor

$H_2(v=0)$



$H_2(v)$

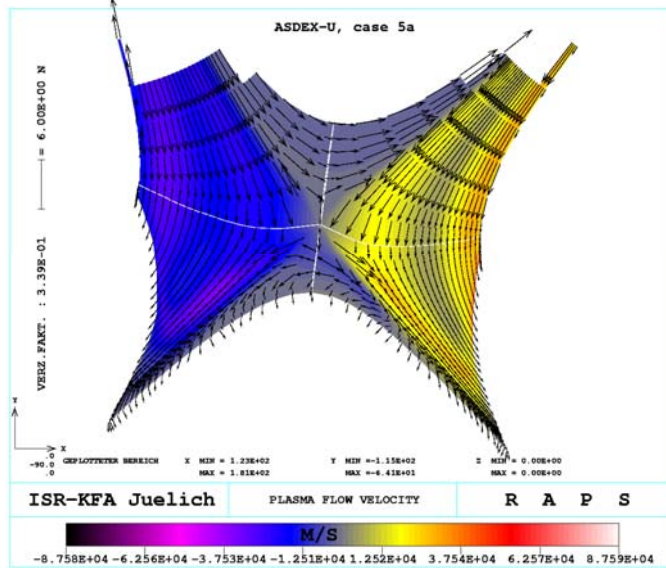




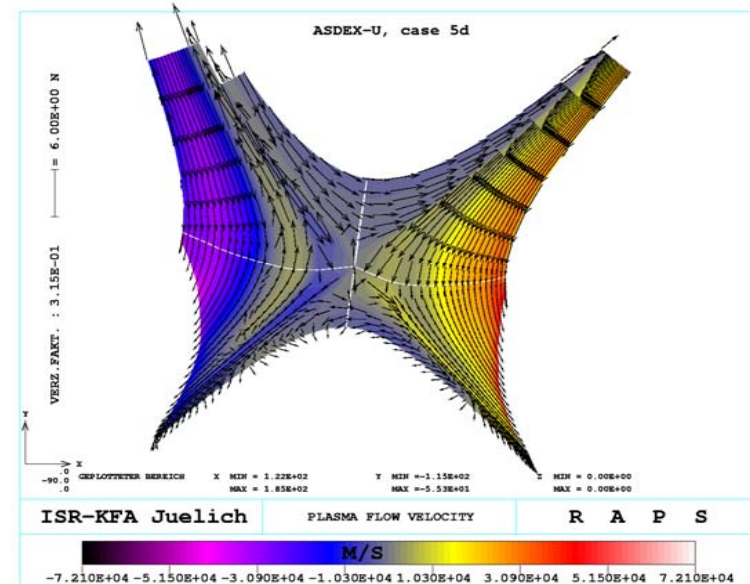
Self sustained neutral cushion? B2-EIRENE simulation: MAR or MAD?

Changed flow pattern: flow reversal (right)
(dominant force (friction) changes sign

$$H_2(v=0)$$



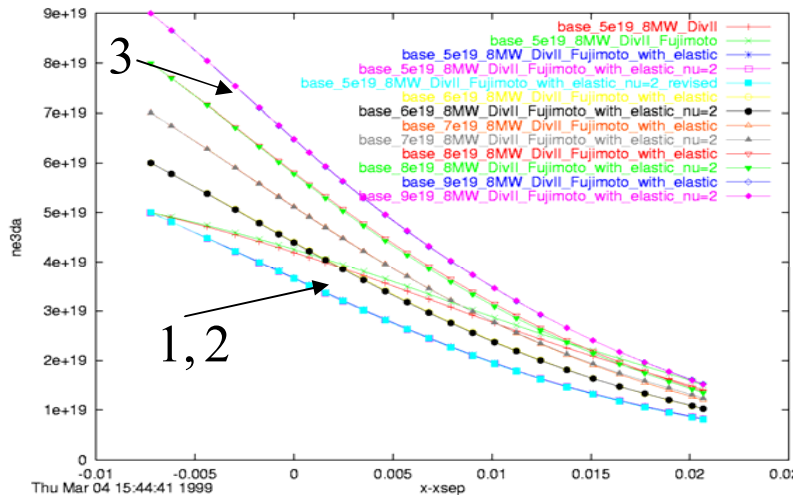
$$H_2(v)$$



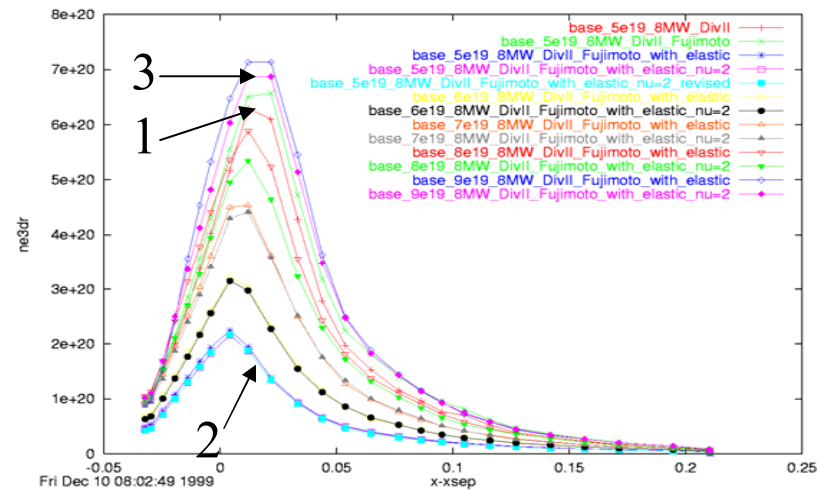
The detached divertor state due to plasma chemistry is a very robust finding from plasma codes. However: Molecule chemistry can change the overall divertor state drastically

The initial detached state (1) (no mol. vibr. chemistry) can be recovered from (2) (with mol.vibr.) by either:

- increasing the upstream density (3) or
- assuming different (anomalous) cross field transport



Outer midplane density profile



Outer target density profile

CONCLUSIONS

Magnetic Confinement Fusion Reactors must operate at reduced target fluxes and temperatures (“detached regime”).

The plasma regime in the Divertor (near target region) is then that of (high electron density) technical gas discharges / fluorescent lamps / RF discharges.

Divertor detachment physics involves a rich complexity of plasma chemistry not otherwise encountered in fusion devices.

- The role of vibrationally excited molecules on overall Divertor dynamics seems accepted.
- The relative role of surface processes (vs. volume processes) to establish this vibrational excitation seems less clear.