



The Abdus Salam  
International Centre for Theoretical Physics



SMR 1773 - 11

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SCHOOL ON PHYSICS AT LHC: "EXPECTING LHC"  
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***Higgs bosons searches at LHC  
Part II***

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*These are preliminary lecture notes, intended only for distribution to participants.*

# The Higgs at the LHC

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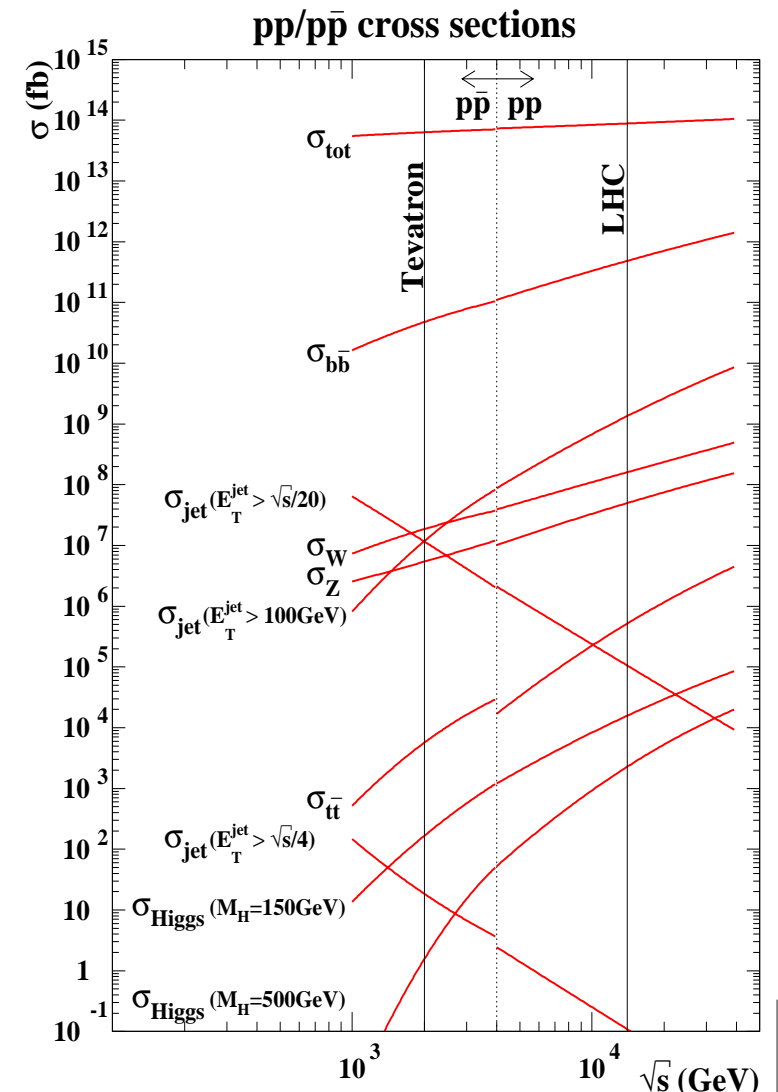
- **The Higgs boson in the Standard Model**
  - **The SM Higgs at the LHC**
    1. **Generalities on Higgs production in pp**
    2. **Higgs production and detection channels**
    3. **Measurement of Higgs properties**
      - **The Higgs boson in SUSY theories**
      - **The SUSY Higgs bosons at the LHC**

# 1. Generalities on Higgs production in pp

**LHC: pp collider**

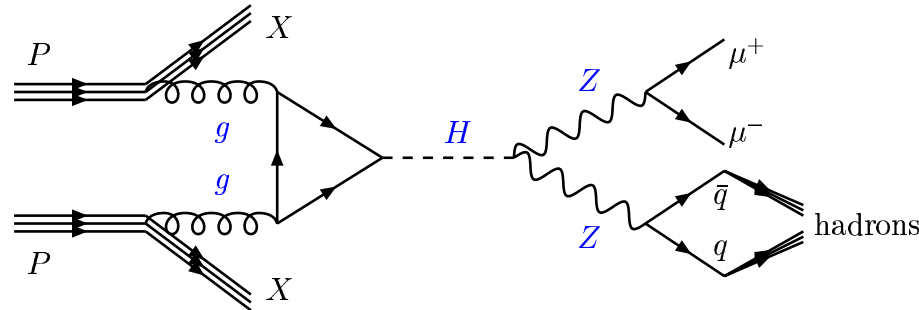
$\sqrt{s}=7+7=14$  TeV  $\Rightarrow \sqrt{s}_{\text{eff}} \sim \sqrt{s}/3 \sim 5$  TeV  
 $\mathcal{L} \sim 10 \text{ fb}^{-1}$  first years and  $100 \text{ fb}^{-1}$  later

- Huge cross sections for QCD processes.
- Small cross sections for EW Higgs signal.  
 $S/B \gtrsim 10^{10} \Rightarrow$  a needle in a haystack!
- Need some strong selection criteria:  
 Trigger: get rid of uninteresting events...  
 Select clean channels:  $H \rightarrow \gamma\gamma, VV \rightarrow \ell\ell$   
 Use different kinematic features for Higgs  
 Combine different decay/production channels  
 Have a precise knowledge of S and B rates.
- Gigantic experimental (+theoretical) efforts!



# 1. Generalities: example of process in pp

An example of process at pp colliders to see how things work.



$$N_{\text{ev}} = \mathcal{L} \times P(g/p) \times \hat{\sigma}(gg \rightarrow H) \times B(H \rightarrow ZZ) \times B(Z \rightarrow \mu\mu) \times BR(Z \rightarrow qq)$$

For a large number of events, all these numbers should be large!

Two ingredients: hard process ( $\sigma$ , B) and soft process (PDF, hadr).

Factorization theorem! Here discuss production process (BR done).

The partonic cross section of the subprocess,  $gg \rightarrow H$ , is:

$$\hat{\sigma}(gg \rightarrow H) = \int \frac{1}{2\hat{s}} \times \frac{1}{2.8} \times \frac{1}{2.8} |\mathcal{M}_{Hgg}|^2 \frac{d^3p_H}{(2\pi)^3 2E_H} (2\pi^4) \delta^4(\mathbf{q} - \mathbf{p}_H)$$

Flux factor, color/spin average, matrix element squared, phase space.

Convolute with gluon densities to obtain total hadronic cross section

$$\sigma = \int_0^1 dx_1 \int_0^1 dx_2 \frac{\pi^2 M_H}{8\hat{s}} \Gamma(H \rightarrow gg) g(x_1) g(x_2) \delta(\hat{s} - M_H^2)$$

# 1. Generalities: example of process in pp

The calculation of  $\sigma_{\text{born}}$  is not enough in general at pp colliders: need to include higher order radiative corrections which introduce terms of order  $\alpha_s^n \log^m(Q/M_H)$  where Q is either large or small...

- Since  $\alpha_s$  is large, these corrections are in general very important.
- Choose a (natural scale) which absorbs/resums the large logs.

Since we truncate pert. series: only NLO/NNLO corrections available.

- The (hopefully) not known HO corrections induce a theoretical error.
- The scale variation is a (naive) measure of the HO: must be small.

Also, precise knowledge of  $\sigma$  is not enough: need to calculate some kinematical distributions (e.g.  $p_T, \eta, \frac{d\sigma}{dM}$ ) to distinguish S from B.

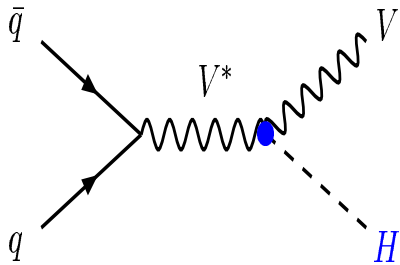
In fact, one has to do this for both the signal and background (unless directly measurable from data): the important quantity is  $\sigma = \frac{N_S}{\sqrt{N_{\text{bjg}}}}$   
 $\Rightarrow$  a lot of theoretical work is needed!

But most complicated thing is to actually see the signal for  $S/B \ll 1!$

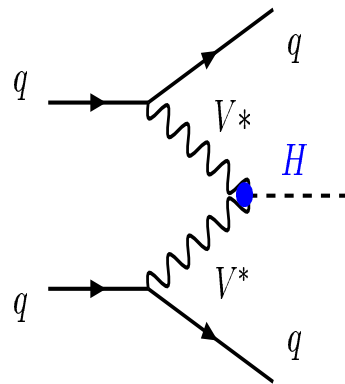
# 1. Generalities: Higgs production in pp

## Production mechanisms

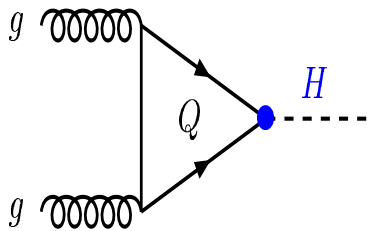
### Higgs-strahlung



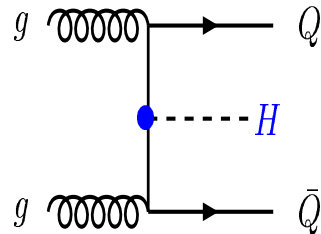
### Vector boson fusion



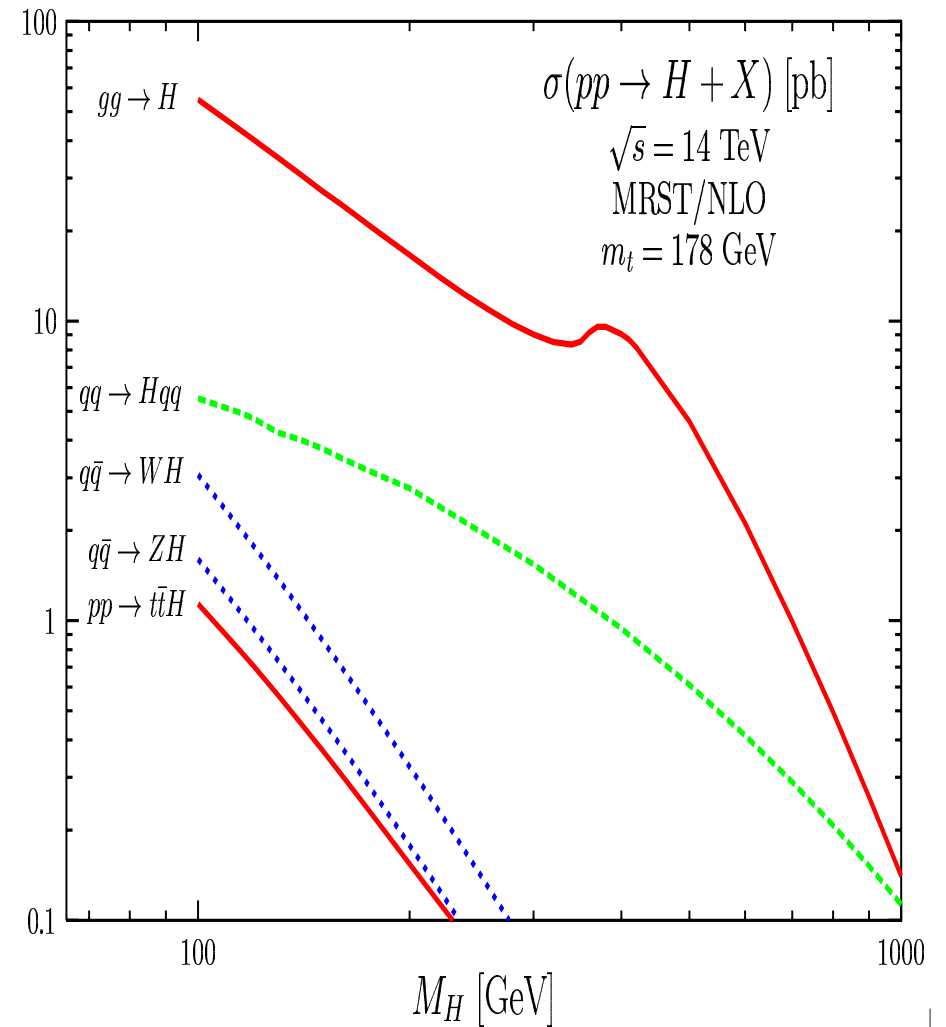
### gluon-gluon fusion



### in associated with $Q\bar{Q}$



## Cross sections at the LHC

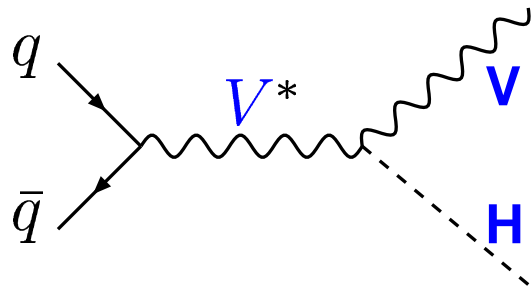


There are also subleading processes,  $gg \rightarrow HH$ , etc...

## 2. SM Higgs production

Let us look at all the main Higgs production channels at the LHC:

The associated HV production:



$$\hat{\sigma}_{\text{LO}}(q\bar{q} \rightarrow VH) = \frac{G_\mu^2 M_V^4}{288\pi \hat{s}} \times (\hat{v}_q^2 + \hat{a}_q^2) \lambda^{1/2} \frac{\lambda + 12M_V^2/\hat{s}}{(1 - M_V^2/\hat{s})^2}$$

Similar to  $e^+e^- \rightarrow HZ$  process used for Higgs searches at LEP2.

Cross section  $\propto \hat{s}^{-1}$  sizable only for low  $M_H \lesssim 200$  GeV values.

Cross section for  $W^\pm H$  approximately 2 times larger than  $ZH$ .

In fact, simply Drell–Yan production of virtual boson with  $q^2 \neq M_V^2$

$$\hat{\sigma}(q\bar{q} \rightarrow HV) = \hat{\sigma}(q\bar{q} \rightarrow V^*) \times \frac{d\Gamma}{dq^2}(V^* \rightarrow HV)$$

$\Rightarrow$  radiative corrections are mainly those of the known DY process

(at 2-loop, need to consider also  $gg \rightarrow HZ$  through box which is  $\neq$ ).

## 2. SM Higgs production: associated HV

### Radiative corrections needed:

- for precise determination of  $\sigma$
- stability against scale variation

### HO also needed to fix scales:

- renormalization  $\mu_R$  for  $\alpha_s$
- factorization  $\mu_F$  for matching.

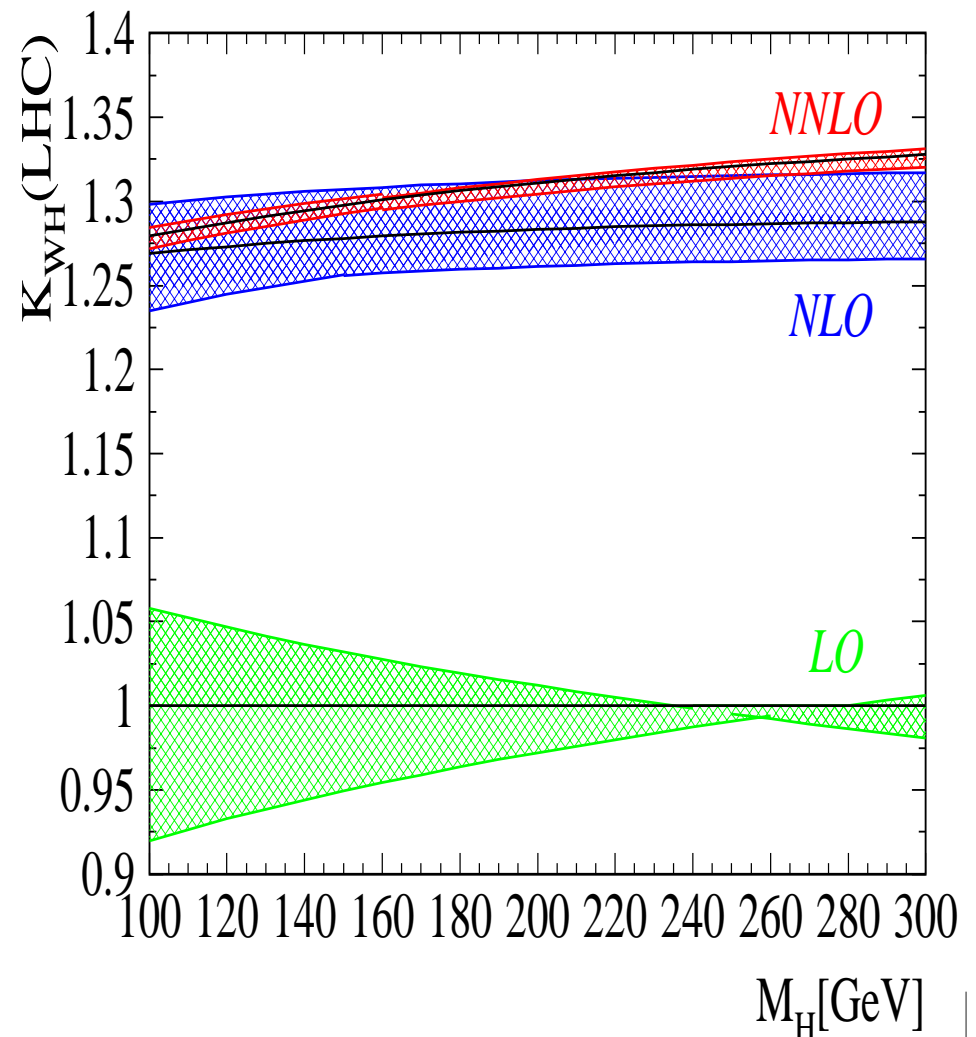
### RC parameterized by K-factor:

$$\mathbf{K} = \frac{\sigma_{\text{HO}}(\text{pp} \rightarrow \text{H} + \text{X})}{\sigma_{\text{LO}}(\text{pp} \rightarrow \text{H} + \text{X})}$$

Can also define K-factor at LO.

QCD RC known up to NNLO.

EW RC known at  $\mathcal{O}(\alpha)$ : small.





## 2. SM Higgs production: associated HV

Up-to-now, it only plays a marginal role at the LHC (small rates etc...).

Interesting final states are:  $WH \rightarrow \gamma\gamma l, b\bar{b}l, 3l$  and  $ZH \rightarrow q\bar{q}\nu\nu$ .

Analyses by ATLAS+CMS:  $5\sigma$  discovery possible with  $\mathcal{L} \gtrsim 100$  fb.

But very clean channel when normalized to  $pp \rightarrow Z$ . Measurements!

**However:**

**WH channel is the most important at Tevatron:**

$M_H \lesssim 130$  GeV:  $H \rightarrow b\bar{b}$ :

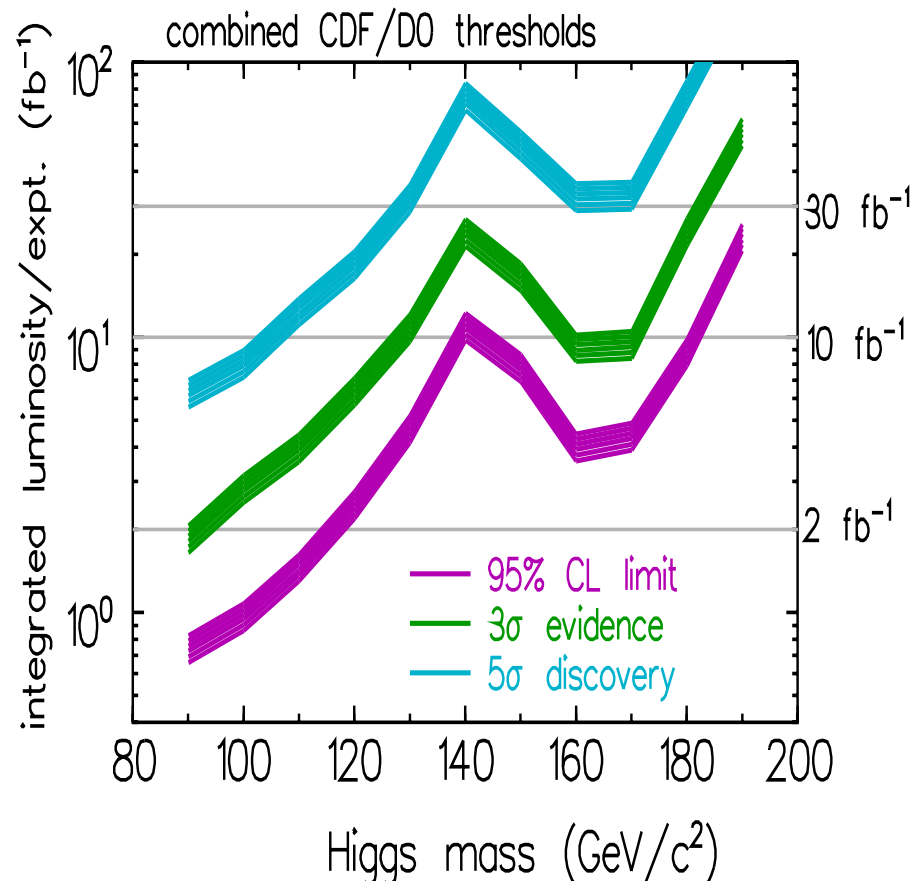
$\Rightarrow l\nu b\bar{b}, \nu\bar{\nu} b\bar{b}, l^+l^-b\bar{b}$

$M_H \gtrsim 130$  GeV:  $H \rightarrow WW^*$

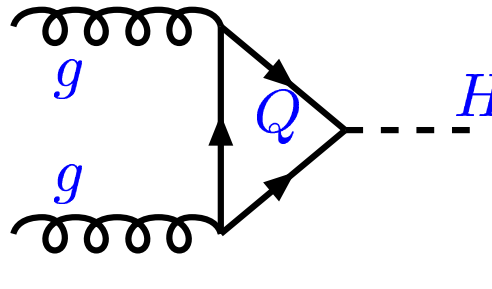
$\Rightarrow l^\pm l^\pm jj, 3l^\pm$

**Possible discovery!!**

**(Report Tevatron HWG)**



## 2. SM Higgs production: gg fusion



$$\hat{\sigma}_{\text{LO}}(\text{gg} \rightarrow \text{H}) = \frac{\pi^2}{8M_{\text{H}}} \Gamma_{\text{LO}}(\text{H} \rightarrow \text{gg}) \delta(\hat{s} - M_{\text{H}}^2)$$

$$\sigma_0^{\text{H}} = \frac{G_{\mu} \alpha_s^2(\mu_{\text{R}}^2)}{288\sqrt{2}\pi} \left| \frac{3}{4} \sum_{\text{q}} A_{1/2}^{\text{H}}(\tau_{\text{Q}}) \right|^2$$

Related to the Higgs decay width into gluons discussed previously.

- In SM: only top quark loop relevant, b-loop contribution  $\lesssim 5\%$ .
- For  $m_{\text{Q}} \rightarrow \infty$ ,  $\tau_{\text{Q}} \sim 0 \Rightarrow A_{1/2} = \frac{4}{3} = \text{constant}$  and  $\hat{\sigma}$  finite.
- Approximation  $m_{\text{Q}} \rightarrow \infty$  valid for  $M_{\text{H}} \lesssim 2m_{\text{t}} = 350 \text{ GeV}$ .

Gluon luminosities large at high energy+strong QCD and Htt couplings

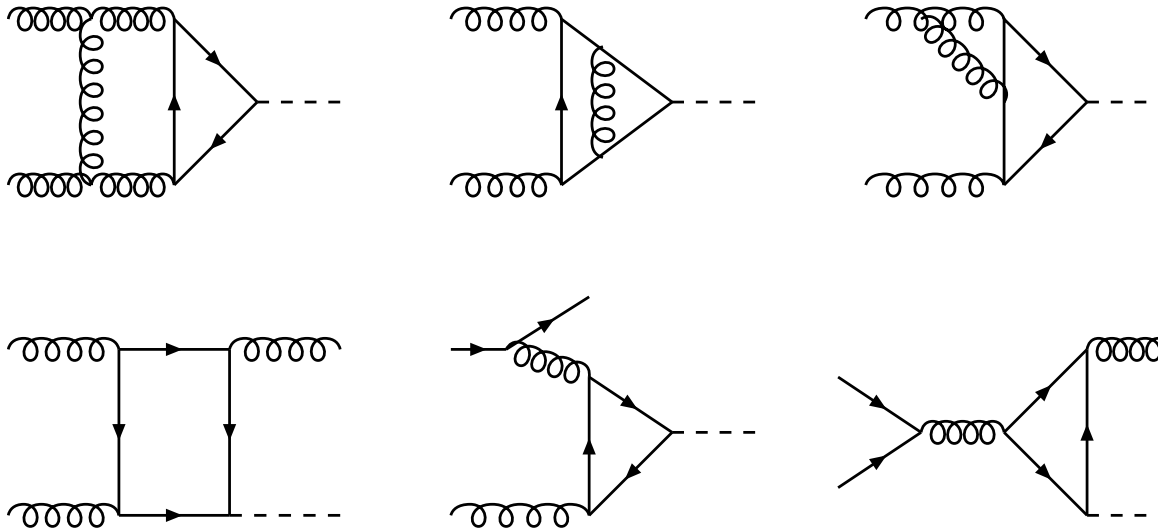
*gg* → *H* is the leading production process at the LHC.

- Very large QCD RC: the two- and three-loops have to be included.
- Also the Higgs  $P_{\text{T}}$  is zero at LO, must be generated at NLO.

## 2. SM Higgs production: gg fusion

### QCD radiative corrections to $gg \rightarrow H$ : NLO case

Typical diagrams for virtual and real QCD corrections to  $gg \rightarrow H$  at NLO:



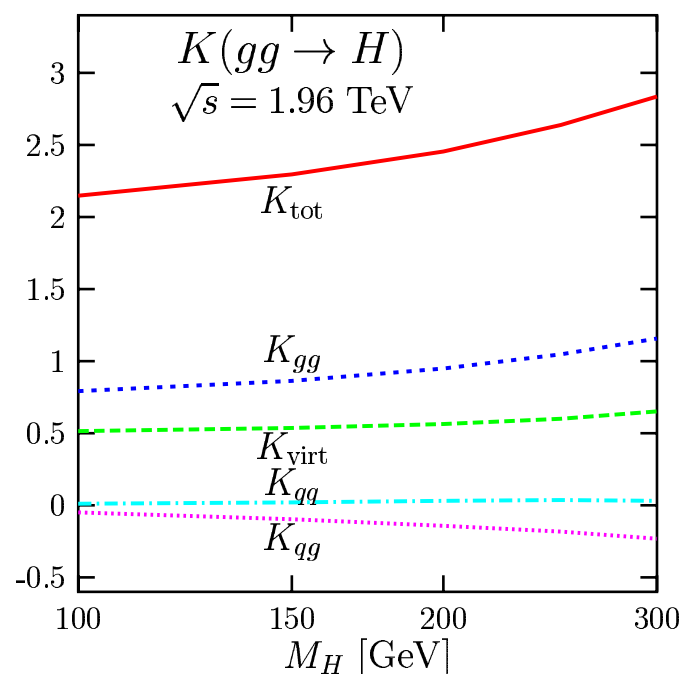
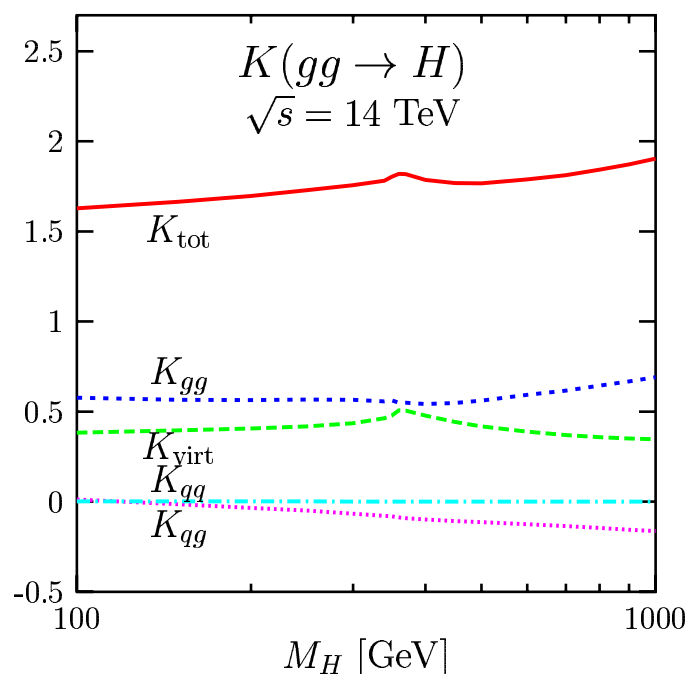
- Regularization of UV divergences from virtual and IR+collinear divergences from real corrections in dimensional regularization.
- UV divergences cancelled by corresponding counterterms.
- IR divergences cancel in sum of virtual+real corrections.
- Collinear singularities are left: absorbed in PDF renormalization.

## 2. SM Higgs production: gg fusion

- Corrections known exactly, i.e. for finite  $m_t$  and  $M_H$ , at NLO:
  - quark mass effects are important for  $M_H \gtrsim 2m_t$ .
  - $m_t \rightarrow \infty$  is still a good approximation for masses below 300 GeV.
  - corrections are large, increase cross section by a factor 1.6–1.9.

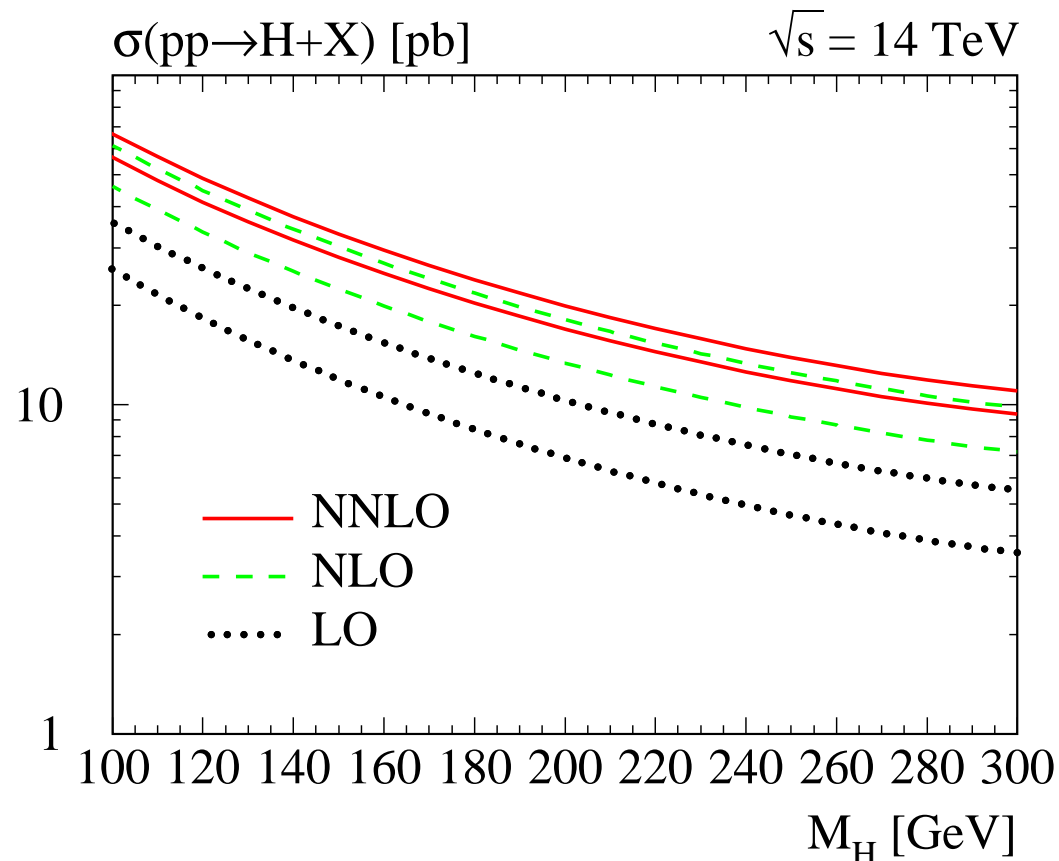
Note 1: NLO corrections to  $P_T, \eta$  distributions are also known.

Note 2: NLO EW corrections are also available, they are rather small.



## 2. SM Higgs production: gg fusion

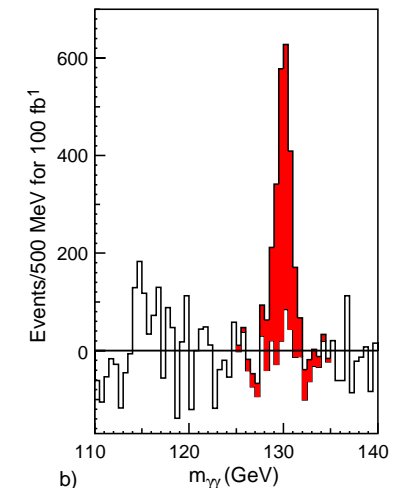
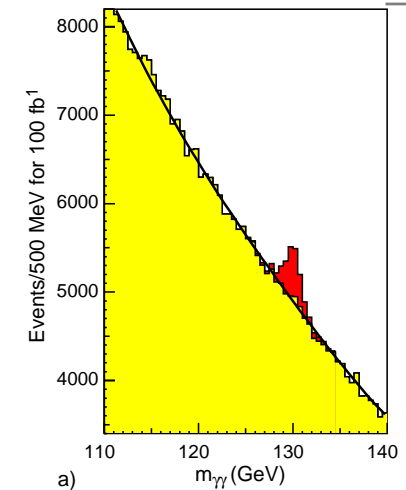
- Corrections have been calculated in  $m_t \rightarrow \infty$  limit at NNLO.
  - moderate increase of cross section by 30% (good behavior of PT!).
  - large stabilization with renormalization and factorization scales.
  - soft-gluon resummation performed up to NNLL:  $\sim 5\%$  effects.



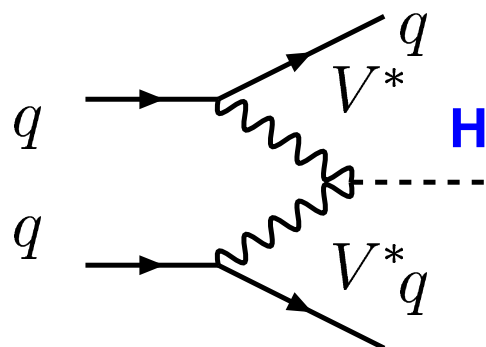
## 2. SM Higgs production: gg fusion

### Relevant detection signals

- $H \rightarrow b\bar{b}, \tau^+\tau^-, t\bar{t}$ : hopeless.
- $H \rightarrow \gamma\gamma$  for  $M_H \lesssim 150$  GeV:
  - large  $\sigma$  and small BR: many events left.
  - huge irreducible bkg from jets:  $10^6$  rejection.
  - large physics bkg from  $q\bar{q}/gg \rightarrow \gamma\gamma + X$ .
  - measure  $d\sigma/dM_{\gamma\gamma}$  on both sides of peak.
  - $S/B = 1/30$  for  $M_{\gamma\gamma} \sim 2$  GeV (good  $\gamma\gamma$  res.).
- $H \rightarrow WW \rightarrow \ell\nu\nu$  for  $M_H \sim 130\text{--}200$  GeV:
  - large  $\sigma \times \text{BR}$  in this range but no  $M_H^{\text{recons}}$
  - large bkg from  $WW/tt$  but use spin-correlations!
- $H \rightarrow ZZ \rightarrow 4\ell^\pm$  for  $M_H \gtrsim 180\text{--}500$  GeV:
  - gold plated mode, clean and small/measurable ZZ bkg.
- $H \rightarrow ZZ \rightarrow \ell\ell jj, \ell\nu\nu, WW \rightarrow \ell\nu jj$  for  $M_H = 0.5\text{--}1$  TeV.



## 2. SM Higgs production: WW fusion



$$\hat{\sigma}_{\text{LO}} = \frac{16\pi^2}{M_{\text{H}}^3} \Gamma(\text{H} \rightarrow \text{V}_{\text{L}} \text{V}_{\text{L}}) \frac{d\mathcal{L}}{d\tau} |_{\text{V}_{\text{L}} \text{V}_{\text{L}}/qq}$$

$$\frac{d\mathcal{L}}{d\tau} |_{\text{V}_{\text{L}} \text{V}_{\text{L}}/qq} \sim \frac{\alpha}{4\pi^3} (\mathbf{v}_{\text{q}}^2 + \mathbf{a}_{\text{q}}^2)^2 \log\left(\frac{\hat{s}}{M_{\text{H}}^2}\right)$$

Three-body final state: analytical expression rather complicated...

Simple form in LVBA:  $\sigma$  related to  $\Gamma(\text{H} \rightarrow \text{V}\text{V})$  and  $\frac{d\mathcal{L}}{d\tau} |_{\text{V}_{\text{L}} \text{V}_{\text{L}}/qq}$

Not too bad approximation at  $\sqrt{\hat{s}} \gg M_{\text{H}}$ : a factor 2 accurate.

Large cross section: in particular for small  $M_{\text{H}}$  and large c.m. energy:

$\Rightarrow$  most important process at the LHC after  $gg \rightarrow \text{H}$ .

QCD radiative corrections small: order 10% (also for distributions).

In fact: at LO in/out quarks are in color singlets and at NLO: no gluons are exchanged between first/second incoming (outgoing) quarks:

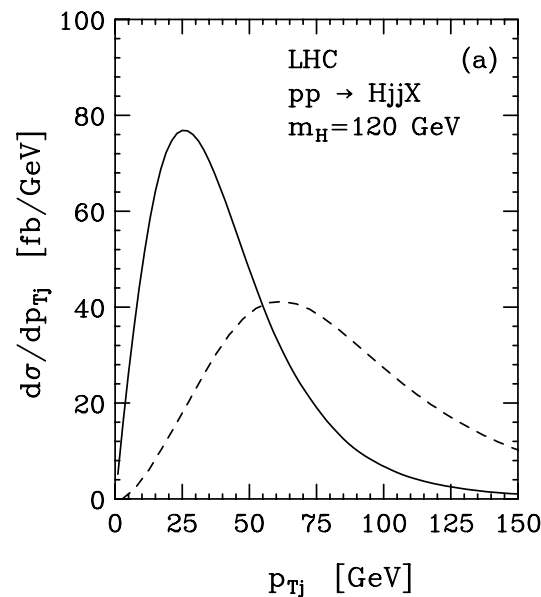
QCD corrections only consist of known corrections to the PDFs!

## 2. SM Higgs production: WW fusion

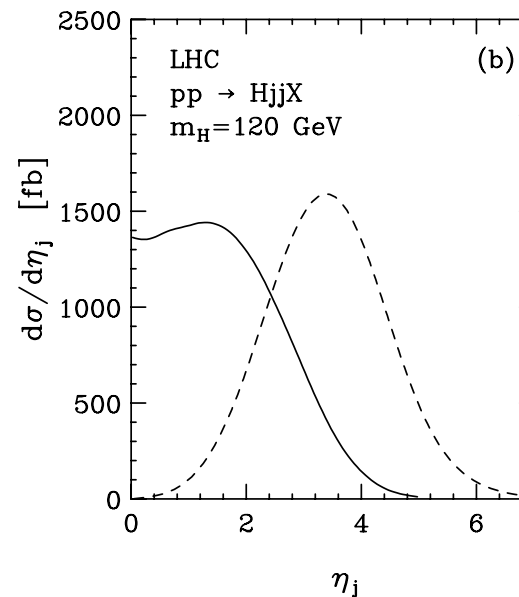
**Kinematics of the process:** a very specific kinematics indeed....

- Forward jet tagging: the two final jets are very forward peaked.
- They have large energies of  $\mathcal{O}(1 \text{ TeV})$  and sizeable  $P_T$  of  $\mathcal{O}(M_V)$ .
- Central jet vetoing: Higgs decay products are central and isotropic.
- Small hadronic activity in the central region no QCD (trigger upon).

Allow to suppress the background to the level of H signal:  $S/B \sim 1$ .



———— lowest/central jet



----- highest/central jet



## 2. SM Higgs production: WW fusion

### Relevant detection signals

- $H \rightarrow \tau^+ \tau^-$  for  $M_H \lesssim 150$  GeV:

first to be established: needs  $\mathcal{L} \sim 30 \text{fb}^{-1}$

$M_{\tau^+ \tau^-}^{\text{recons.}}$  against WW/tt/Zjj bkg.

$\tau$  polarization usefull against  $Z \rightarrow \tau^+ \tau^-$

- $H \rightarrow \gamma\gamma$  for  $M_H \lesssim 150$  GeV:

very clean with small/measurable bkg

rare/needs  $\mathcal{L}$ +combine with other channels

- $H \rightarrow WW \rightarrow \ell\nu\nu$

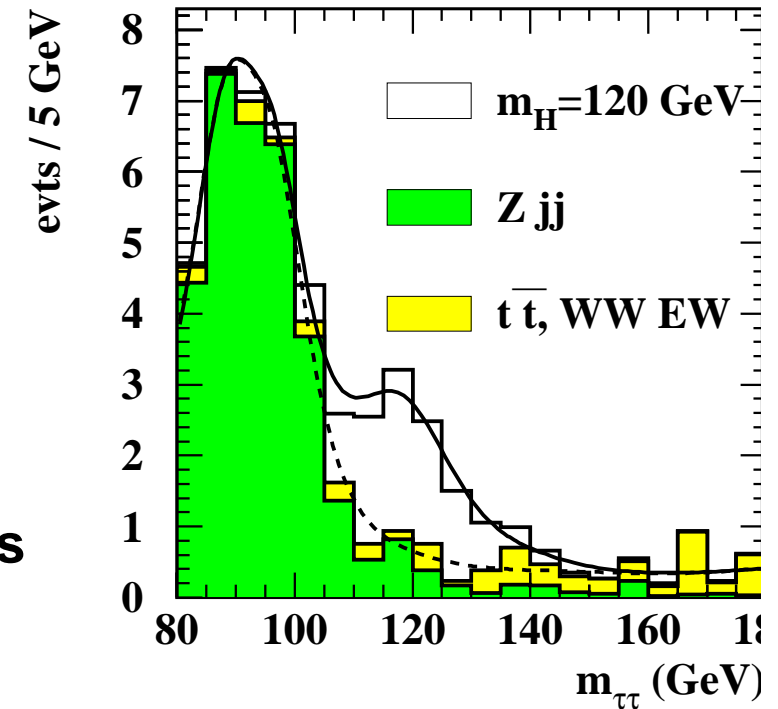
very difficult as you need to know background.

but feasible at low  $M_H$  and efficient at high  $M_H$ .

- $H \rightarrow ZZ \rightarrow \ell\nu\nu, \ell\ell jj$ : have large bkg

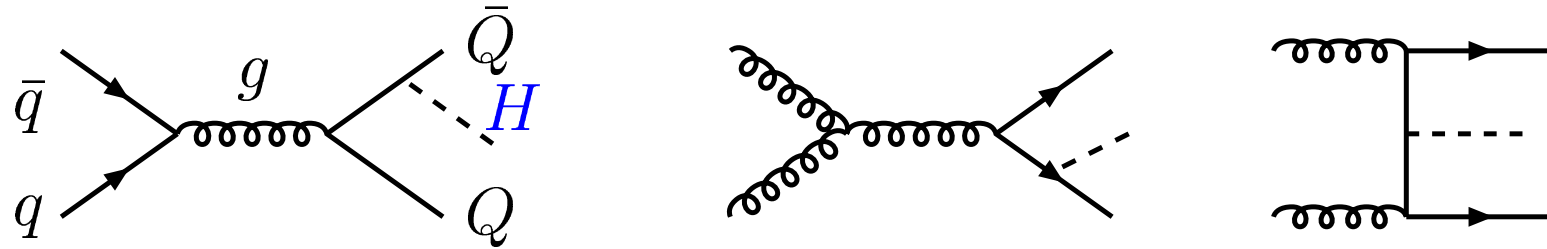
need high  $\mathcal{L}$ , usefull at high masses in combination.

- $H \rightarrow b\bar{b}, t\bar{t}$  very difficult and  $H \rightarrow \mu^+ \mu^-$  needs high  $\mathcal{L}$ .

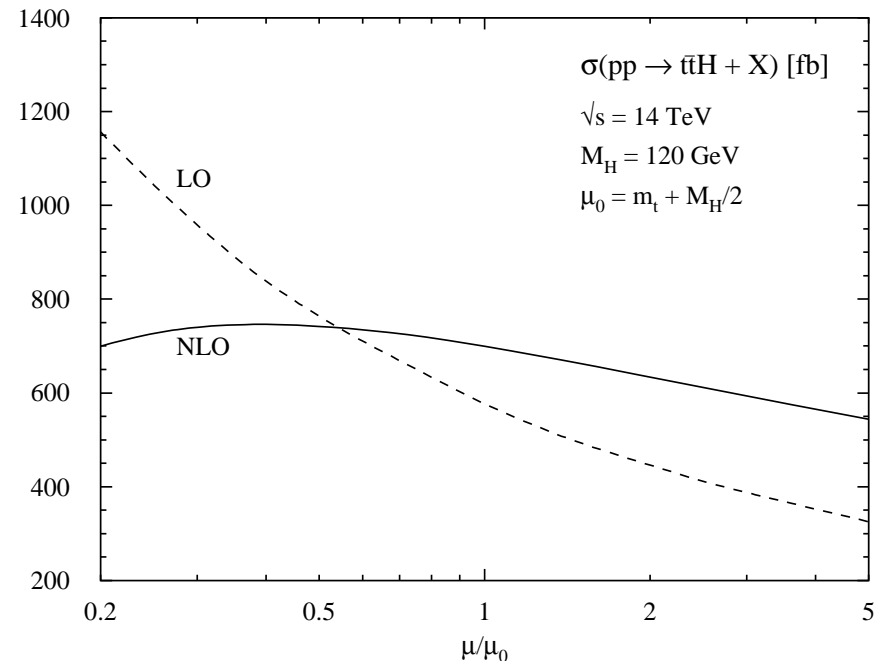
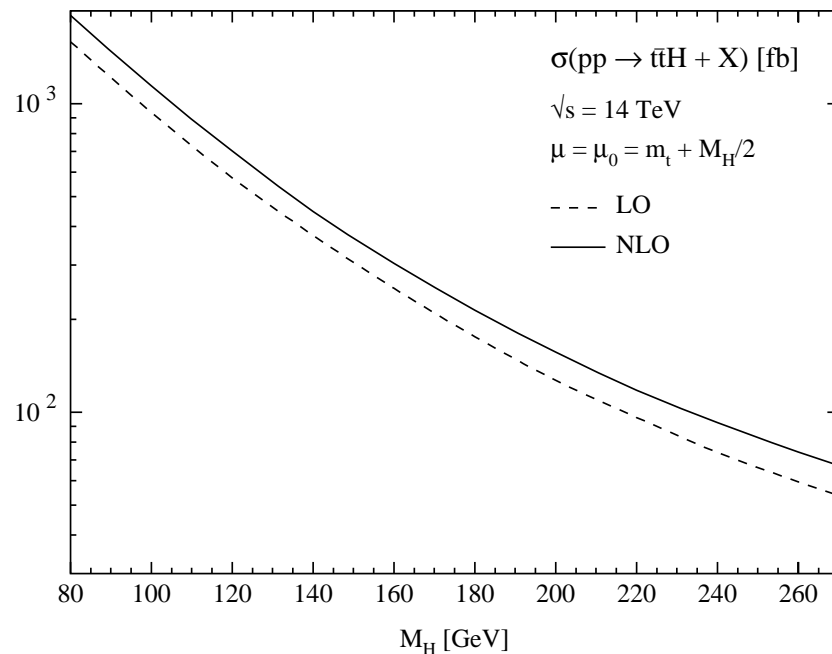


## 2. SM Higgs production: Htt production

Most complicated process for Higgs production in pp: many channels:



NLO corrections recently calculated (Zerwas et al., Dawson et al.):  
small K-factors ( $\sim 1.2$ ) but strong reduction of scale variation!



## 2. SM Higgs production: Htt production

Small corrections to kinematical distributions (e.g:  $p_T^{\text{top}}$ ,  $P_T^{\text{H}}$ ), etc...

- Rather tiny uncertainties from higher orders, PDFs.
- Other possible processes involving heavy quarks work only in BSM:
  - Single top+Higgs production:  $pp \rightarrow tH + X$ .
  - Associated production with bottom quarks:  $pp \rightarrow bbH$ .

Interesting signals at the LHC for this process are:

- $pp \rightarrow Htt \rightarrow \gamma\gamma l^\pm$ : clean but rather small rates.
- $pp \rightarrow Htt \rightarrow b\bar{b}l^\pm$ : needs efficient b tagging.
- $pp \rightarrow Htt \rightarrow l^\mp l^\pm \nu\nu$ : large bckgs from ttWjj, etc...

Possibility for a 3–5 signal at  $M_H \lesssim 140$  GeV with high luminosity.

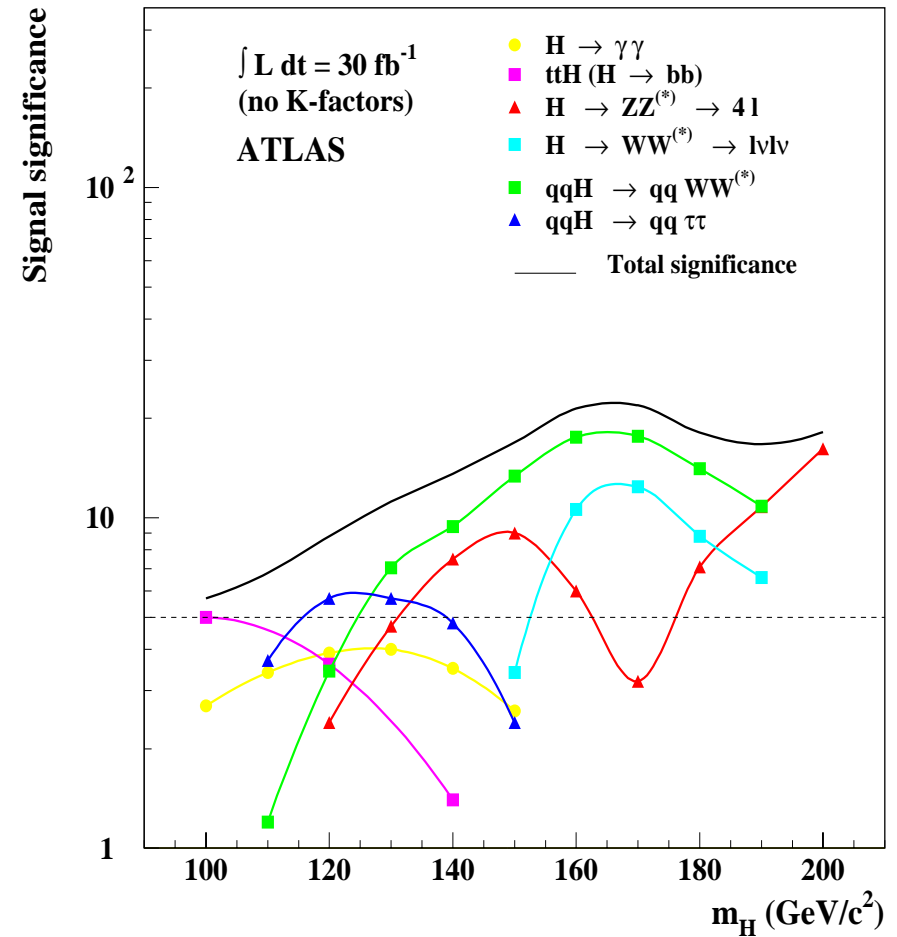
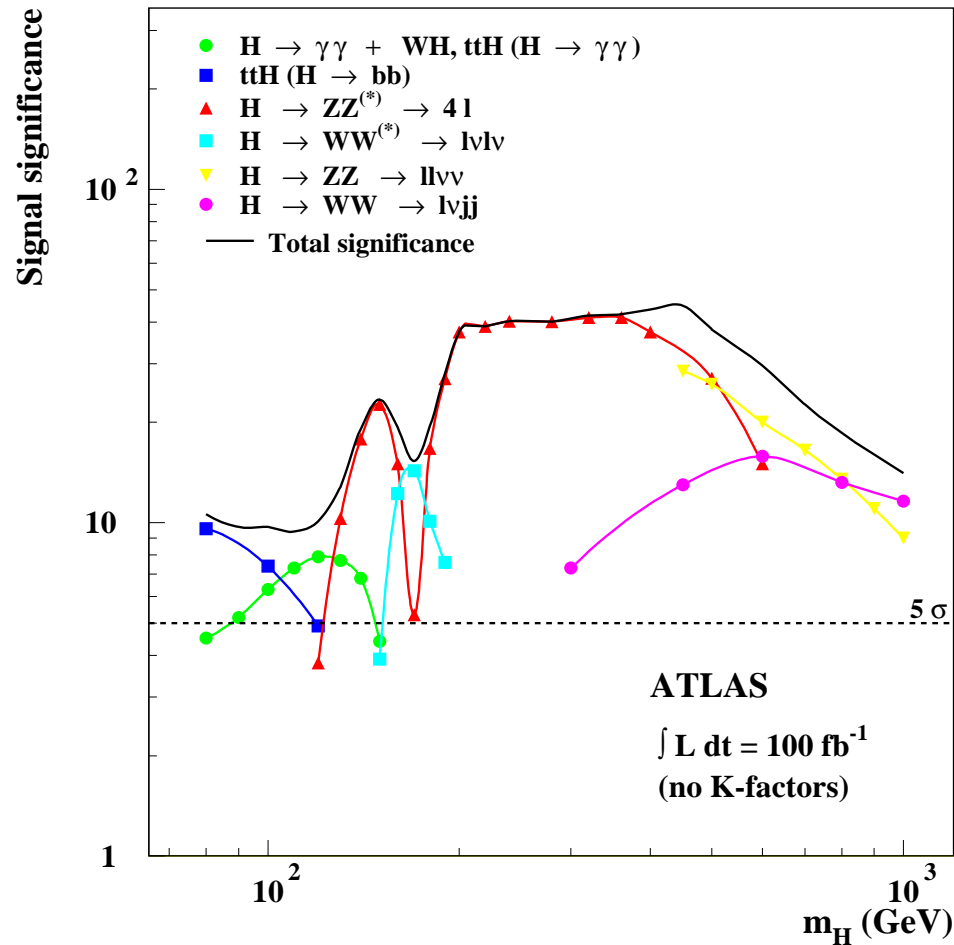
Needs to be combined with similar channels and topologies (eg:

$pp \rightarrow WH \rightarrow l\gamma\gamma, lb\bar{b}$ ) to increase total signal significance.

But process very important for measurement of Htt Yukawa coupling!

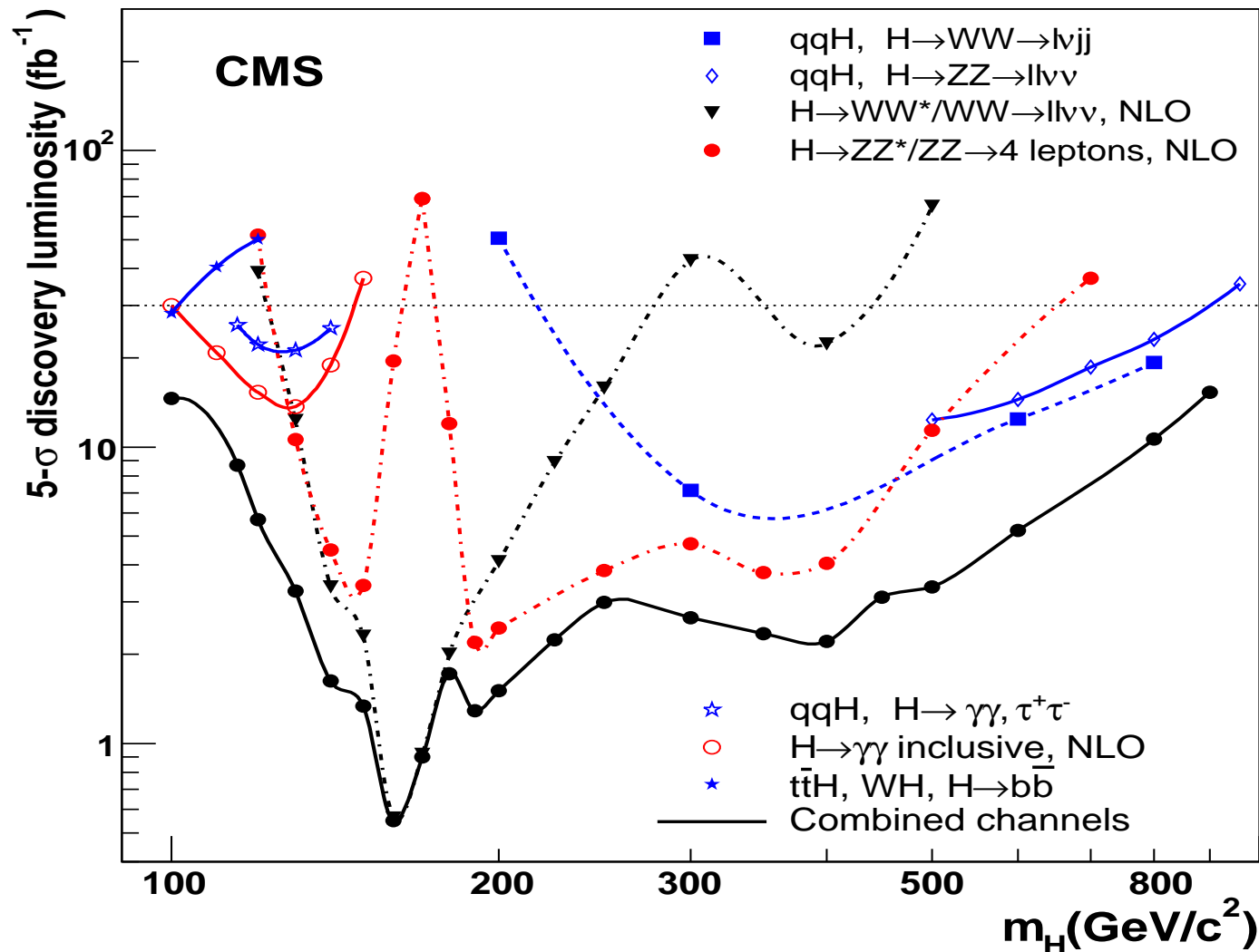
## 2. SM Higgs production: summary

All in all, when you do the hard experimental work, you will get:



## 2. SM Higgs production: summary

Another way too summarize the expectations: in terms of luminosity



# 3. Measurement of Higgs properties

So in 2-3 years from now we will find the Higgs (and maybe nothing else): we celebrate, shake hands, drink champagne, take care of our bets, ... and should we declare Particle Physics closed and go home or fishing? No! We need to check that it is indeed responsible of spontaneous EWSB

**Measure its fundamental properties in the most precise way:**

- its mass and total decay width,
- its spin–parity quantum numbers and check  $J^{PC} = 0^{++}$ ,
- its couplings to fermions and gauge bosons and check that they are indeed proportional to the particle masses (fundamental prediction!),
- its self–couplings to reconstruct the potential  $V_H$  that makes EWSB.

**A very ambitious and challenging program!**

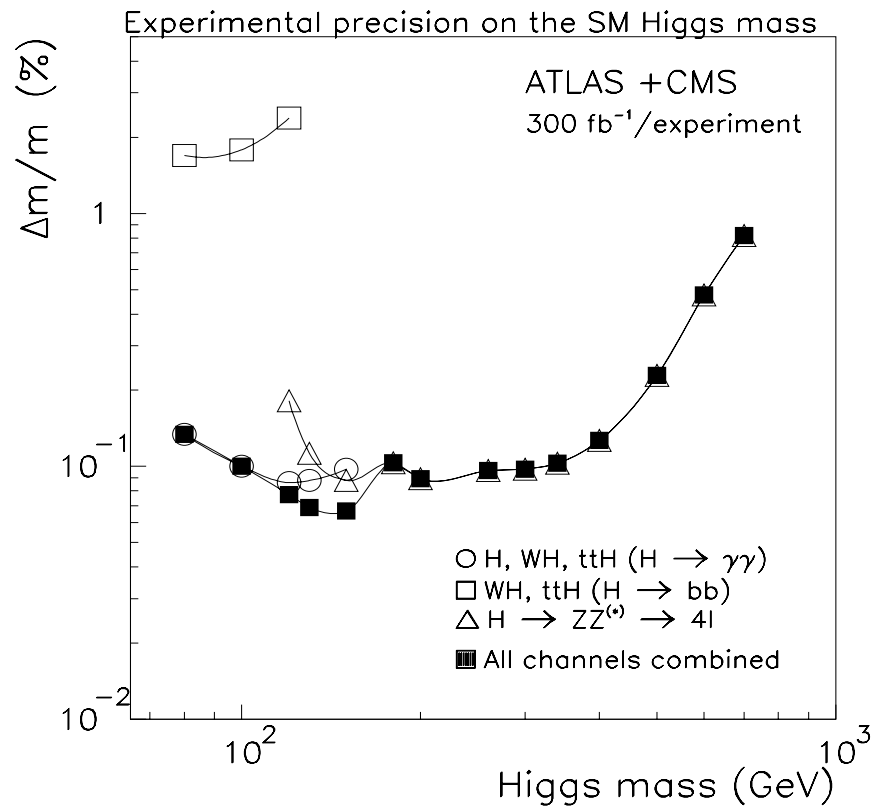
which is even more difficult to achieve than the Higgs discovery itself...

### 3. Higgs properties: mass and width

#### Higgs boson mass from:

- $H \rightarrow \gamma\gamma$  for  $M_H \lesssim 130$  GeV
- $H \rightarrow ZZ \rightarrow 4\ell^\pm$  beyond

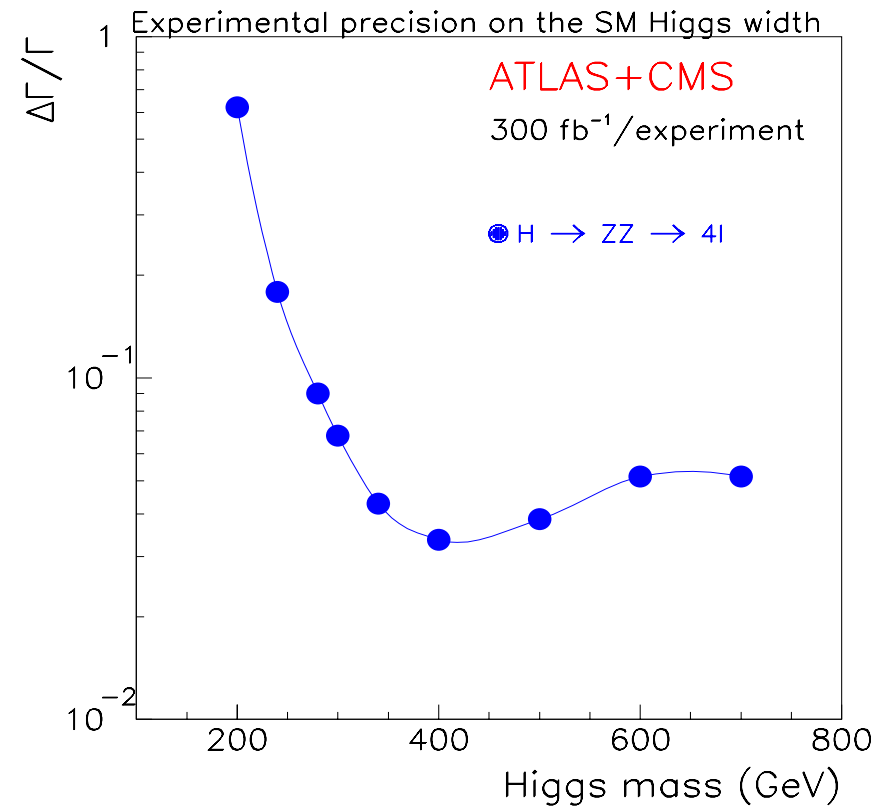
Final  $\Delta M_H/M_H \sim 0.1\%$  to  $1\%$ .



#### Higgs boson total width:

- Too small for  $M_H \lesssim 2M_Z$
- $H \rightarrow ZZ \rightarrow 4\ell^\pm$  beyond

Final  $\Delta\Gamma_H/\Gamma_H \sim$  a few %



However: for large  $M_H$  effects from large width are important!

### 3. Higgs properties: $J^{PC}$ numbers

- **Higgs spin:**

Higgs can be observed in  $H \rightarrow \gamma\gamma$  decays: rules out  $J=1$  and fixes  $C=+$ ,

- argument not generalizable to  $H \leftrightarrow gg$  since no  $g/q$  distinction,

- other particle spin-assignments might be possible  $J=2$  (radion), etc.

- **Higgs parity:**

- Higgs can be observed in  $H \rightarrow ZZ \rightarrow 4\ell^\pm$  rules out CP-odd state.

- Higgs spin-correlations in  $gg \rightarrow H \rightarrow WW^*$  also useful here...

But we need to check that H is pure CP-even with no CP-odd mixture:

- it becomes then a challenging high-precision measurement,

- can be done roughly by looking at correlations in  $H \rightarrow ZZ, WW$ .

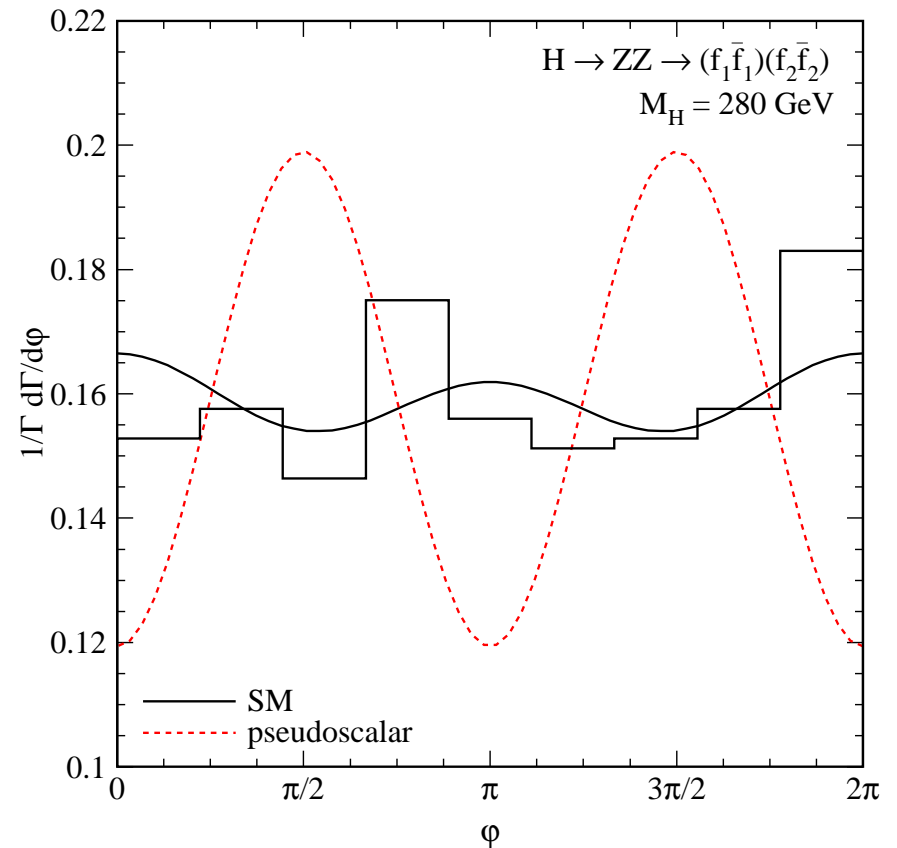
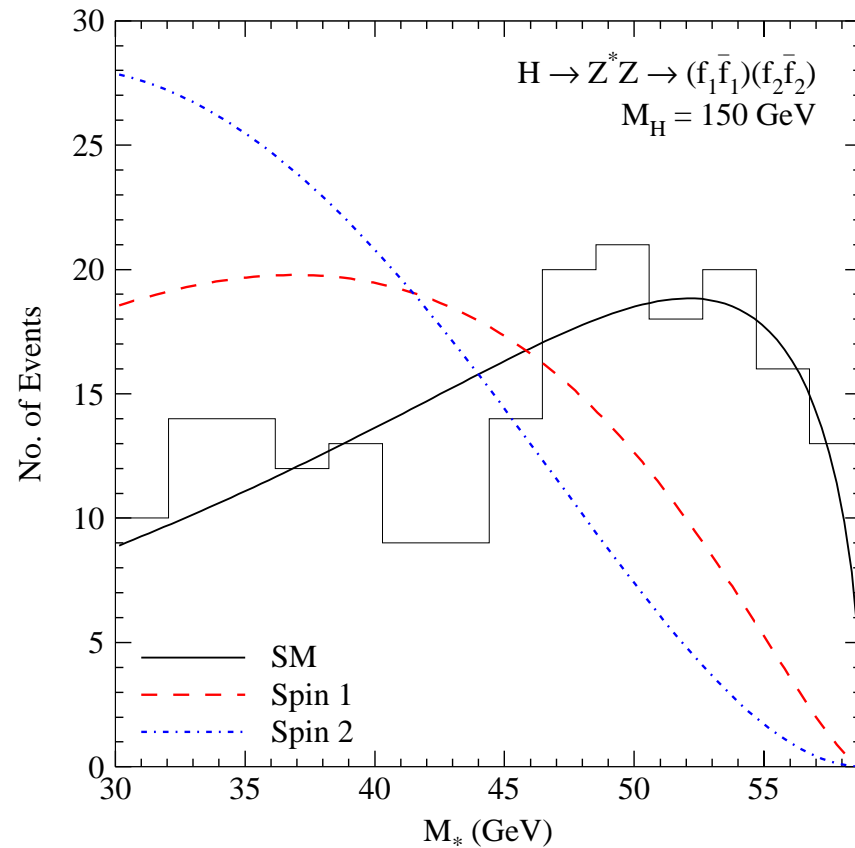
**Drawback:** If H is mostly CP-even, rates for  $A \rightarrow VV$  too small...

More convincing to look at more democratic Higgs-fermion couplings.

Possible channels:  $H \rightarrow \tau^+\tau^-$  or  $pp \rightarrow t\bar{t}H$ : very challenging!!



### 3. Higgs properties: $J^{PC}$ numbers



**left: threshold behavior of  $d\Gamma(H \rightarrow ZZ^*)/dM_*$  distribution for  $J=0,1,2$**

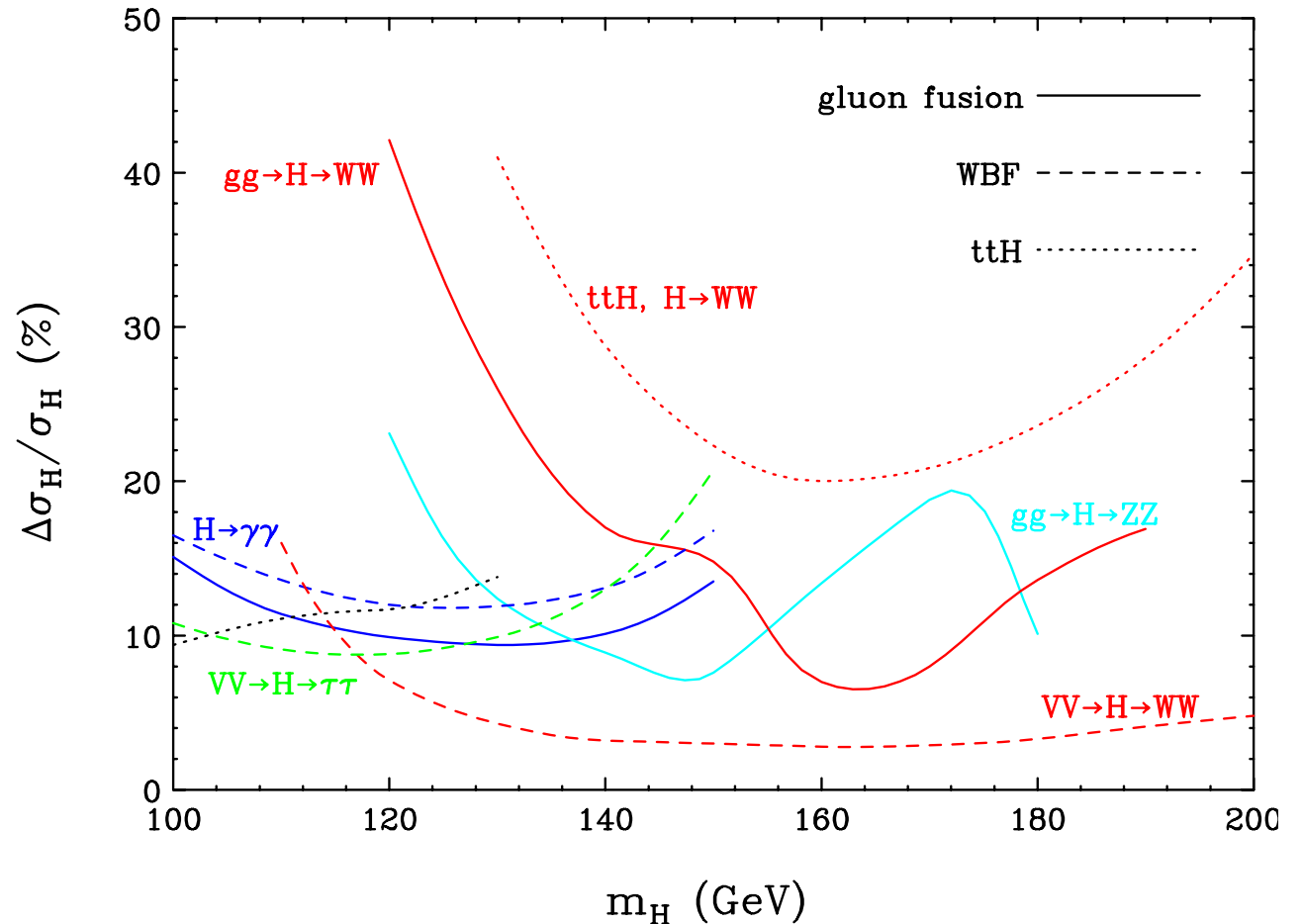
**right: azimuthal distributions  $d\Gamma(H \rightarrow ZZ)/d\phi$  for SM and CP-odd A**

**ATLAS simulation including bkg with  $\int \mathcal{L} dt = 300 \text{ fb}^{-1}$  at the LHC**

### 3. Higgs properties: Higgs couplings

Higgs couplings can be determined by looking at various Higgs production and decay channels and measuring  $N_{ev} = \sigma \times BR$ .

LHC with  $\mathcal{L} = 300\text{fb}^{-1}$ .  
Only statistical errors.  
D. Zeppenfeld et al.



### 3. Higgs properties: Higgs couplings

The errors are in general rather large unfortunately:

- experimental errors: statistics, systematics, parton luminosity,...
- theoretical errors: PDFs, HO+scale variation, model dependence...

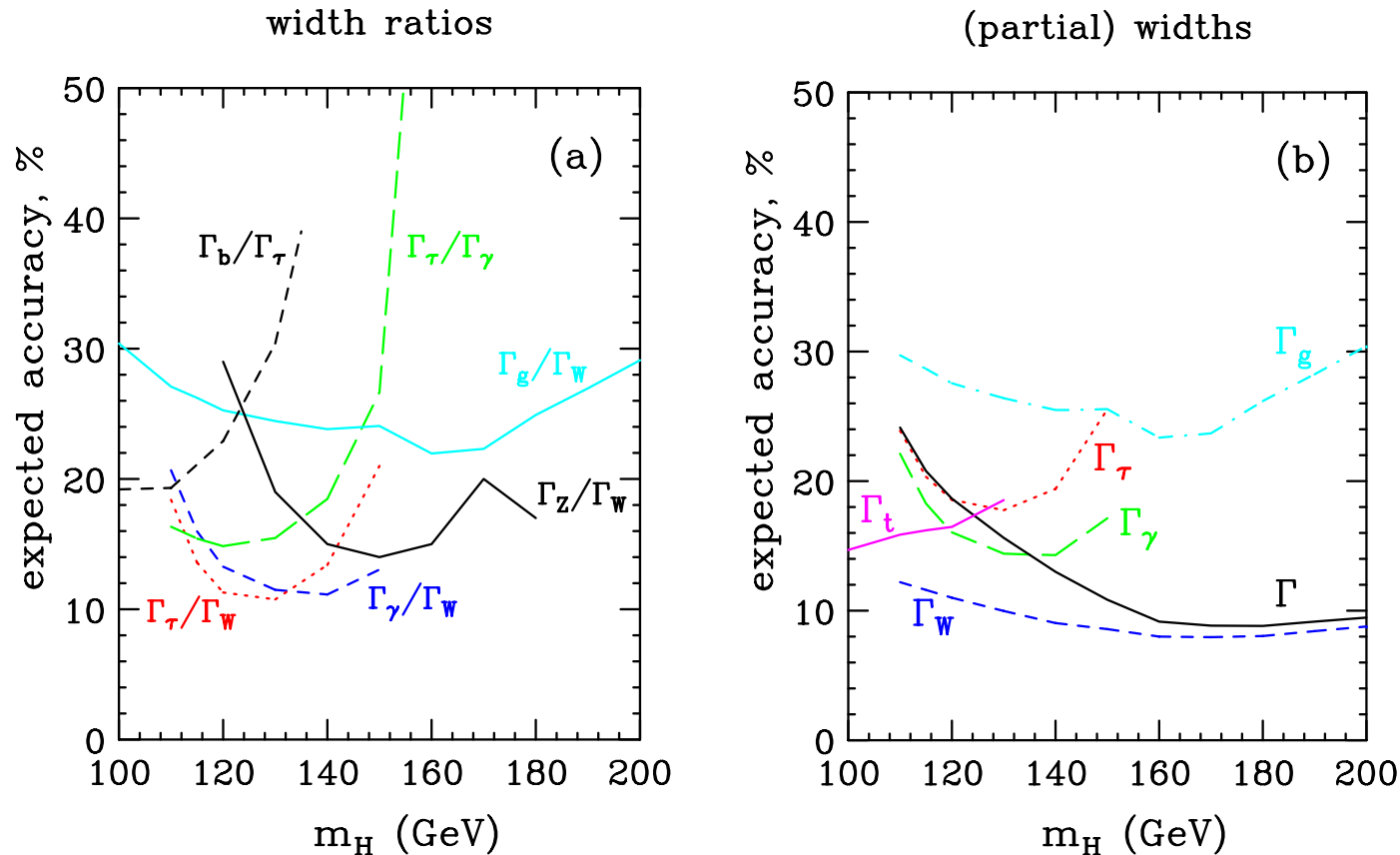
⇒ ratios of  $\sigma \times \text{BR}$ : many errors drop out!

Process	Measurement quantity	Error	Mass range
$\frac{(t\bar{t}H+WH)\rightarrow\gamma\gamma+X}{(t\bar{t}H+WH)\rightarrow b\bar{b}+X}$	$\frac{\text{BR}(H\rightarrow\gamma\gamma)}{\text{BR}(H\rightarrow b\bar{b})}$	$\sim 15\%$	80 – 120 <b>GeV</b>
$\frac{H\rightarrow\gamma\gamma}{H\rightarrow 4\ell^+}$	$\frac{\text{BR}(H\rightarrow\gamma\gamma)}{\text{BR}(H\rightarrow ZZ^*)}$	$\sim 7\%$	120 – 150 <b>GeV</b>
$\frac{t\bar{t}H\rightarrow\gamma\gamma,b\bar{b}}{WH\rightarrow\gamma\gamma,b\bar{b}}$	$(g_{Htt}/g_{HWW})^2$	$\sim 15\%$	80 – 120 <b>GeV</b>
$\frac{H\rightarrow ZZ^*\rightarrow 4\ell^+}{H\rightarrow WW^*\rightarrow 2\ell^\pm 2\nu}$	$(g_{HZZ}/g_{HWW})^2$	$\sim 10\%$	130 – 190 <b>GeV</b>

**Note:** for  $M_H \gtrsim 2M_Z$  only few processes accessible:  $H \rightarrow WW/ZZ$ .  
while  $\sigma(\text{gg} \rightarrow H)$  provides  $g_{Htt}$  but indirectly since loop mediated.

### 3. Higgs properties: Higgs couplings

- Then translate into partial widths  $\propto$  Higgs coupling squared.
- Precision on coupling measurement is:  $\Delta g_{HXX} = \frac{1}{2} \frac{(\Delta^{\text{exp}} \Gamma + \Delta^{\text{th}} \Gamma)}{\Gamma}$
- Some theoretical assumptions (no invisible, SU(2) invariance, some couplings are known, etc..) allow to extract additional couplings....

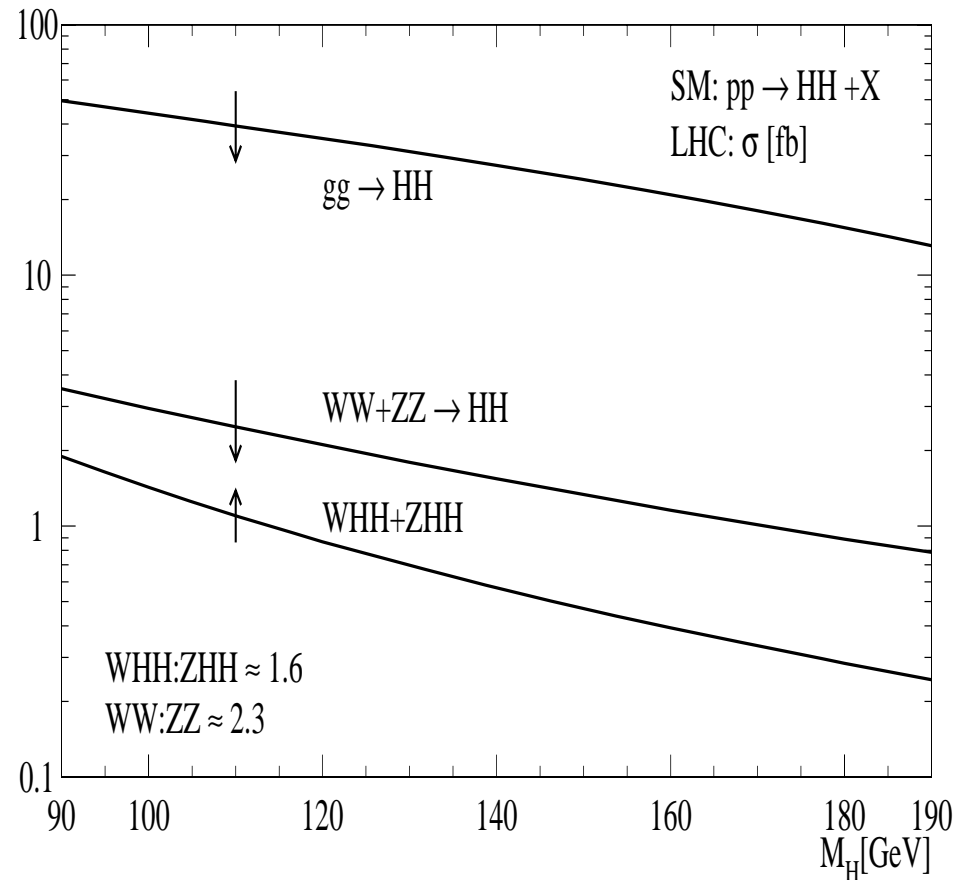
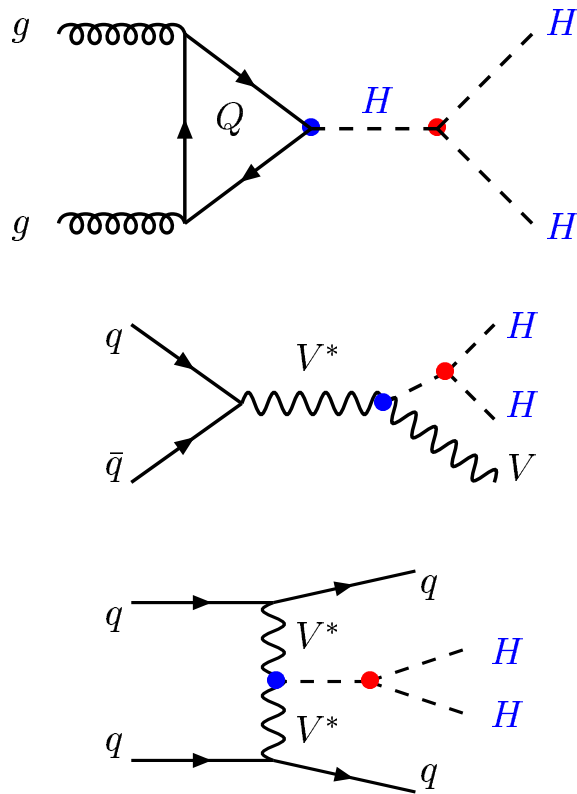


### 3. Higgs properties: Higgs self-couplings

Important couplings to be measured:  $g_{H^3}, g_{H^4} \Rightarrow$  access to  $V_H$ .

- $g_{H^3}$  is accessible in double Higgs production:  $pp \rightarrow HH + X$
- $g_{H^4}$  is hopeless to measure, needs  $pp \rightarrow HHH + X$  with too low rates.

Relevant processes for HH prod:



### 3. Higgs properties: Higgs self-couplings

Cross sections small, except maybe for  $gg \rightarrow HH$  at  $M_H \lesssim 200$  GeV:

- $H \rightarrow \gamma\gamma$  decay too rare,
- $H \rightarrow b\bar{b}$  decay not clean
- $H \rightarrow WW$  at low  $M_H$ ?

Yes, it has been tried:

- parton level analysis...
- look for  $2\ell^\pm, 3\ell^\pm + \nu$ +jets+
- look at IM distributions
- use large luminosity.

Some hope to set limits....

Needs to go to SLHC...

U. Baur et al.

