



**The Abdus Salam
International Centre for Theoretical Physics**



1855-21

**School and Workshop on Highly Frustrated Magnets and Strongly
Correlated Systems: From Non-Perturbative Approaches to
Experiments**

30 July - 17 August, 2007

Strongly Correlated Electrons on Frustrated Lattices

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Strongly Correlated Electrons on Frustrated Lattices



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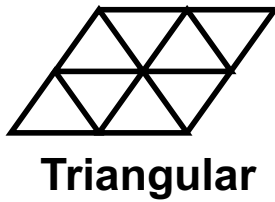
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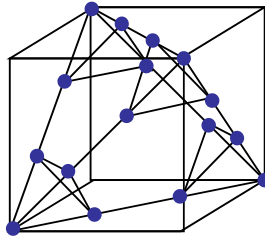
OUTLINE

- Motivation
- [A] Kagomé Lattice Hubbard model
 - Mott transition and charge fluctuations
 - 1D-type spin correlation
 - quasiparticle
- [B] Anisotropic Triangular Lattice Hubbard Model
 - ** effects of reduction of frustration ***
 - slope of Mott transition line - entropy effects
 - heavy quasiparticle formation and metal-insulator transition
- Summary

Correlated electron systems with geometrical frustration



Pyrochlore



Many interesting systems:

- Superconductivity $\text{Na}_x\text{CoO}_2 \cdot y\text{H}_2\text{O}$, AOs_2O_6
- Heavy Fermion LiV_2O_4 , etc
- Quantum spin liquid $\kappa\text{-(ET)}_2\text{Cu}_2(\text{CN})_3$

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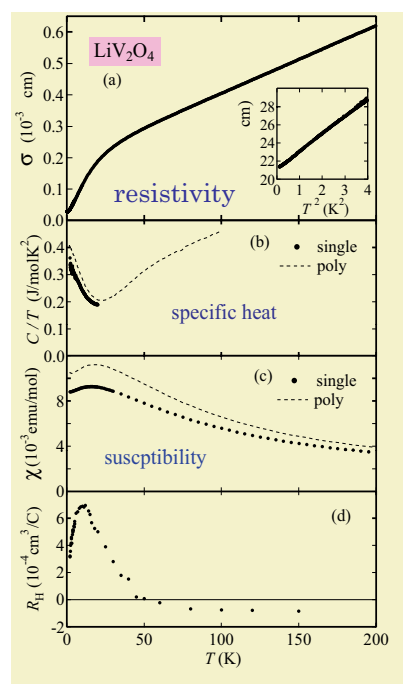
(Ex.) LiV_2O_4 – Heavy Fermion behavior

[Ref: Kondo et al, Urano et al.]

V ions \Rightarrow pyrochlore lattice

- NO f-electrons ($\text{V}^{3.5+} \cdot (3d)^{1.5}$)
- Specific heat $\gamma \sim 200 [\text{mJ}/\text{K}^2 \text{ V-mol}]$
- Wilson ratio $R_W \sim 1.7$

Large low-energy entropy characteristic to frustrated lattice participate to heavy quasiparticle formation?



Large reduction of energy scale is due to GF?

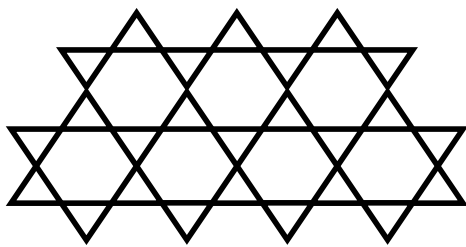
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PART A

Mott Transition in Kagomé Lattice Hubbard Model

[Ohashi, Kawakami, and Tsunetsugu, Phys. Rev. Lett. **97** ('06) 066401]

Kagomé Lattice Hubbard Model



$$H = -t \sum_{\langle i,j \rangle, \sigma} c_{i\sigma}^+ c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

- 2D analog of pyrochlore lattice
- Effective model of $\text{Na}_x\text{CoO}_2 \cdot y\text{H}_2\text{O}$
 - Koshibae & Maekawa PRL 91, 257003 (2003)**
 - Bulut, Koshibae, & Maekawa PRL 95, 37001 (2005)**
- Spin systems on Kagomé lattice
 - ⇒ unusual properties (gapped triplet, gapless singlet excitations)
 - [today's morning sessions]

Relation btw charge fluctuations and spin correlations
Relation btw spin correlations and quasiparticle coherence
[fix density at half filling $n=1$]

Method

Metal-insulator transition in Kagomé lattice Hubbard model at half filling



- strong correlation
- geometrical frustration
- short-range quantum fluctuations \Rightarrow DMFT

Cluster extension of DMFT:
Cellular dynamical mean field theory (CDMFT)

Kotliar, et al. PRL 87, (2001)

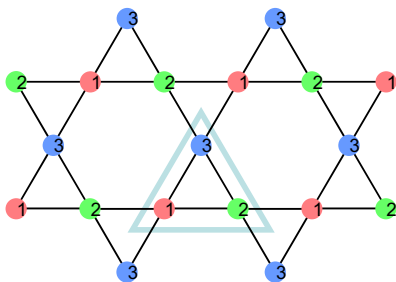
Lichtenstein & Katsnelson, PRB 62, (2000)

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Cellular-DMFT

Kagomé lattice Hubbard model

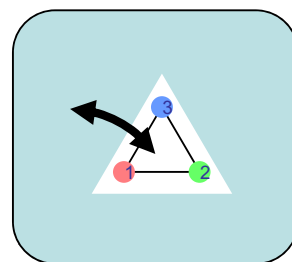


$$H = -t \sum_{\langle i,j \rangle, \sigma} c_{i\sigma}^+ c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Self-energy: 3x3 matrix $\Sigma_{ij}(\omega)$

- Spatially extended correlation
- Geometrical frustration

Effective cluster model

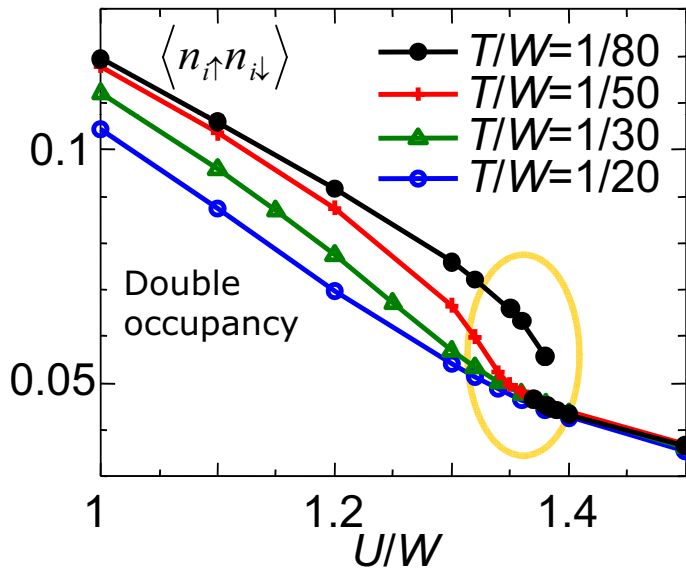


QMC
Hirsch & Fye, (1986)

25000 k-points
 $L_\beta = 12/T$
 10^6 QMC sweeps

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Mott transition



Band width: $W=6t$

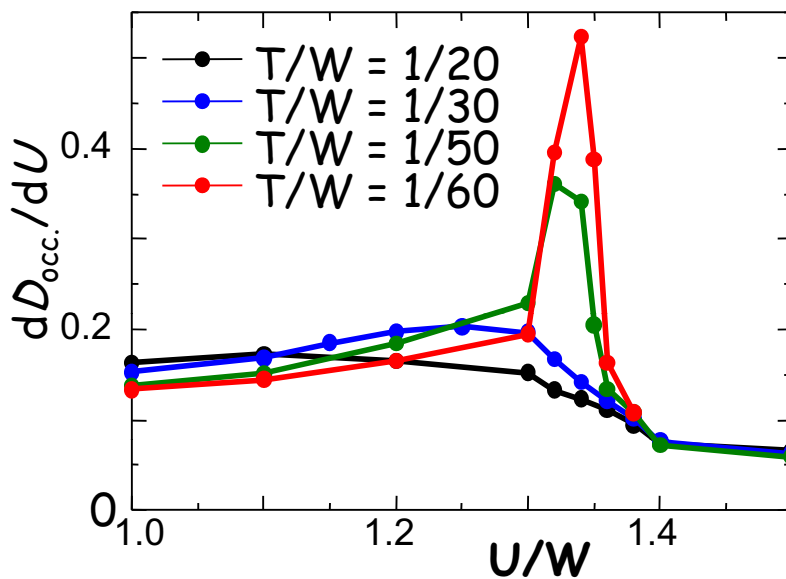
- High temperature $T/W > 1/80$
crossover $U^* \sim 1.35$
- Low temperature $T/W = 1/80$
1st order transition with hysteresis:
 $U_c \sim 1.37$

square lattice:
 $U_c \sim 0.5-1.0$

Larger critical value U_c

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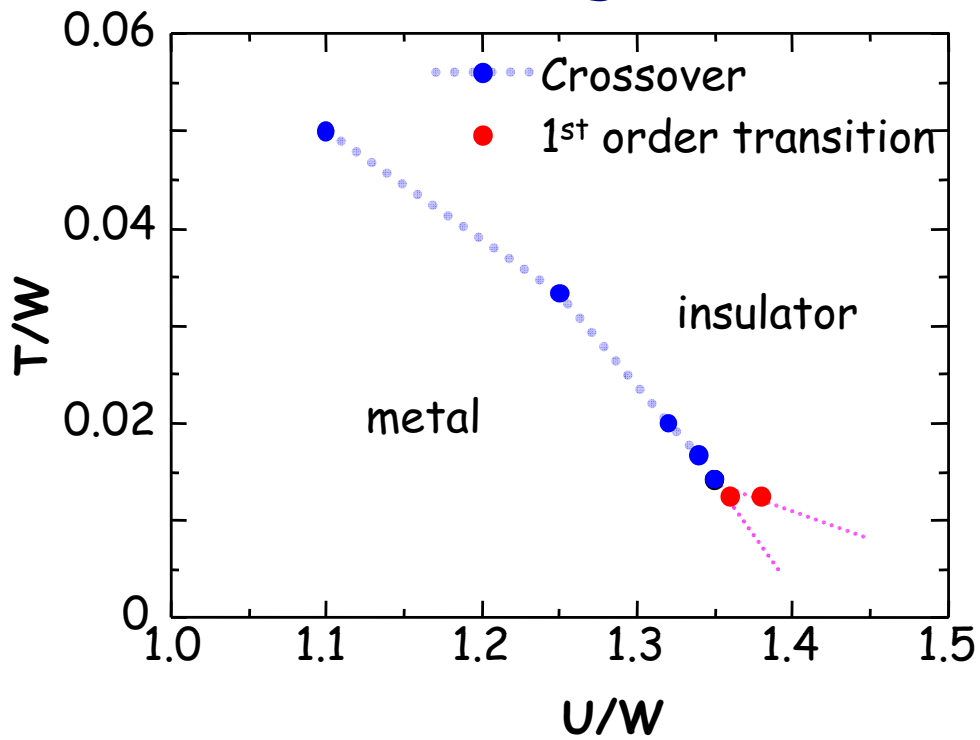
Crossover



Define metal-insulator crossover points $U^*(T)$ by largest change in double occupancy

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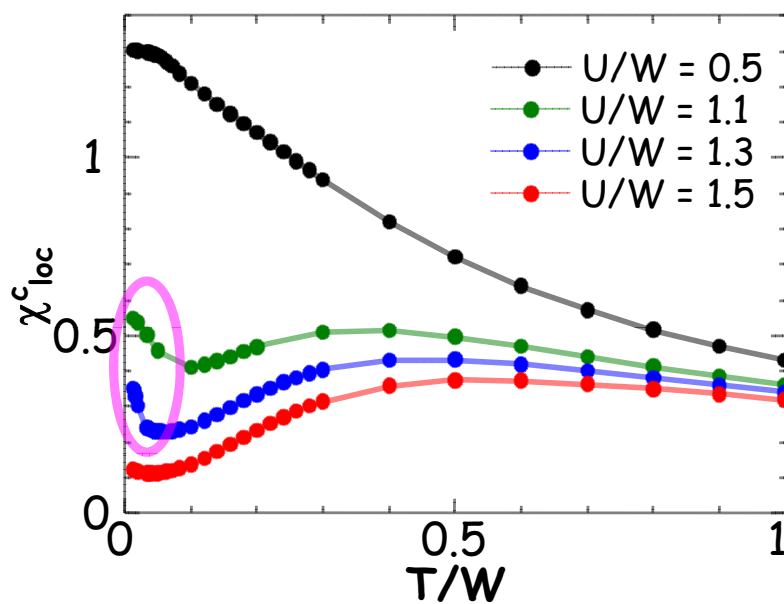
Phase diagram



Mott transition in Kagome Hubbard model

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Local charge susceptibility

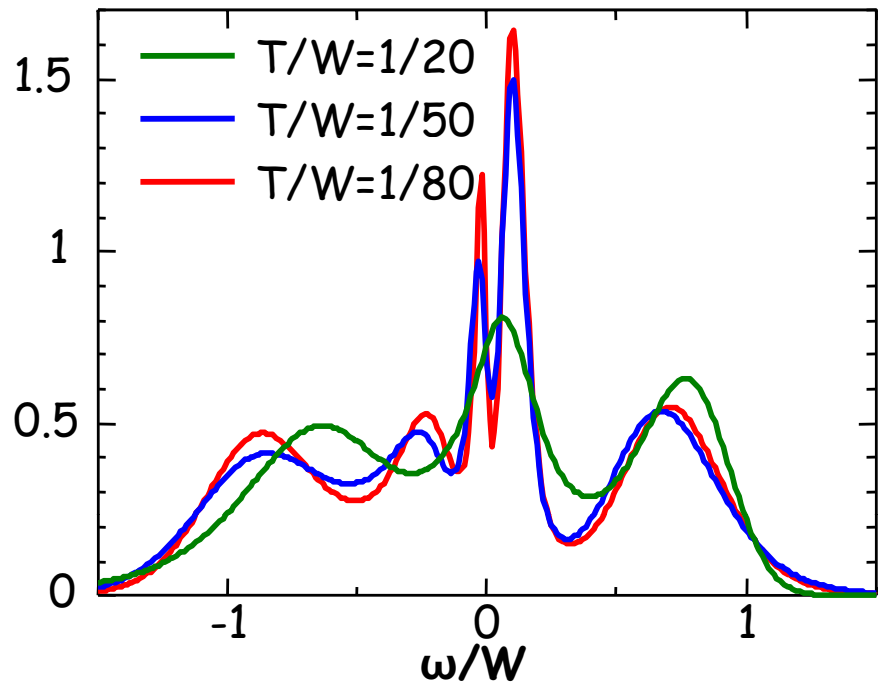


Charge response grows once again in low-T metallic region

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Density of States: $U/W=1.1$

Evolution of heavy quasiparticles at low temperatures



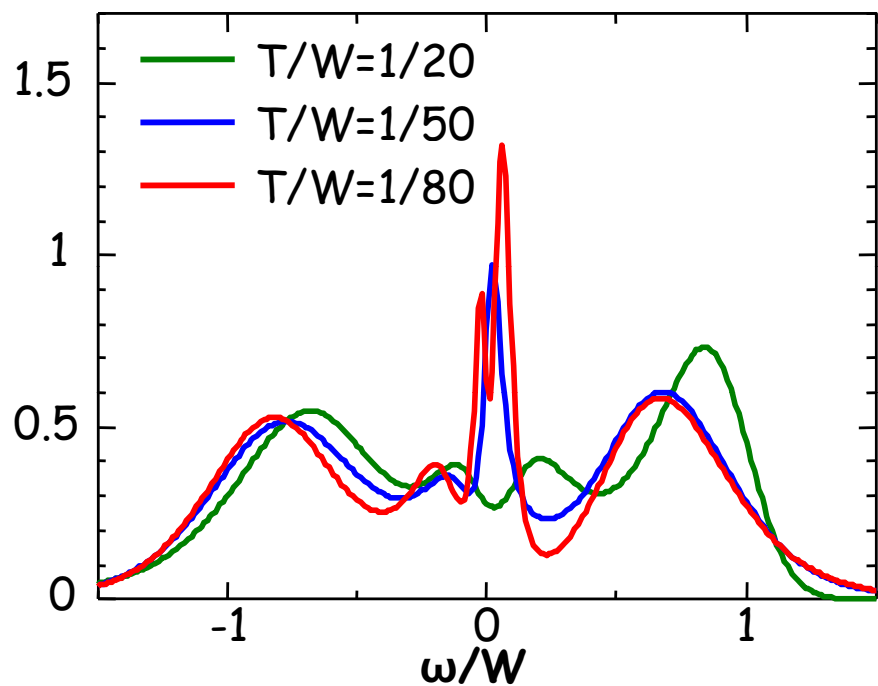
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Density of States: $U/W=1.3$

Precursor of insulating behavior

+

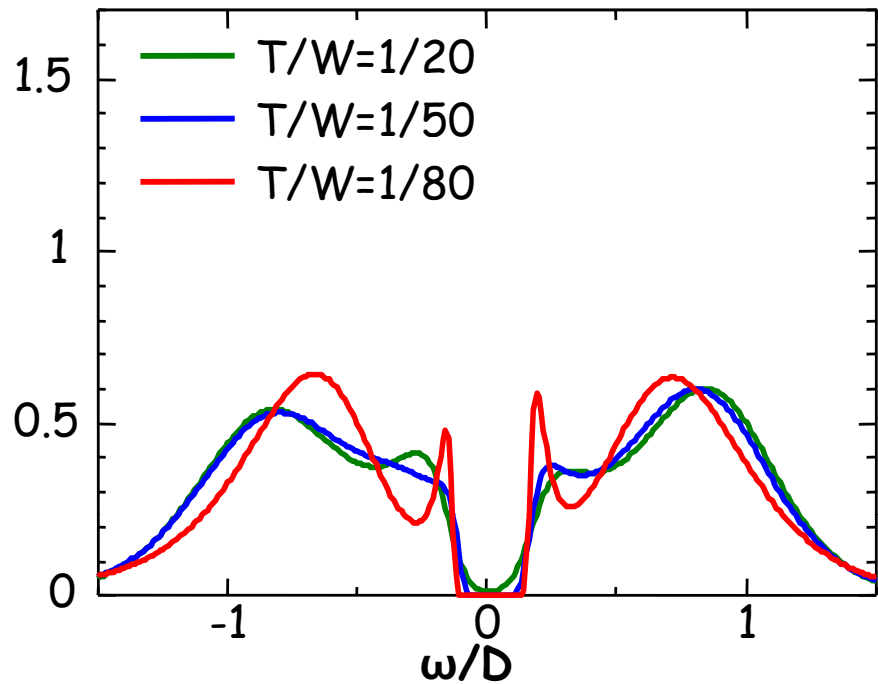
Evolution of heavy quasiparticles at low temperatures



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Density of States: $U/W=1.5$

Clear formation
of Mott gap



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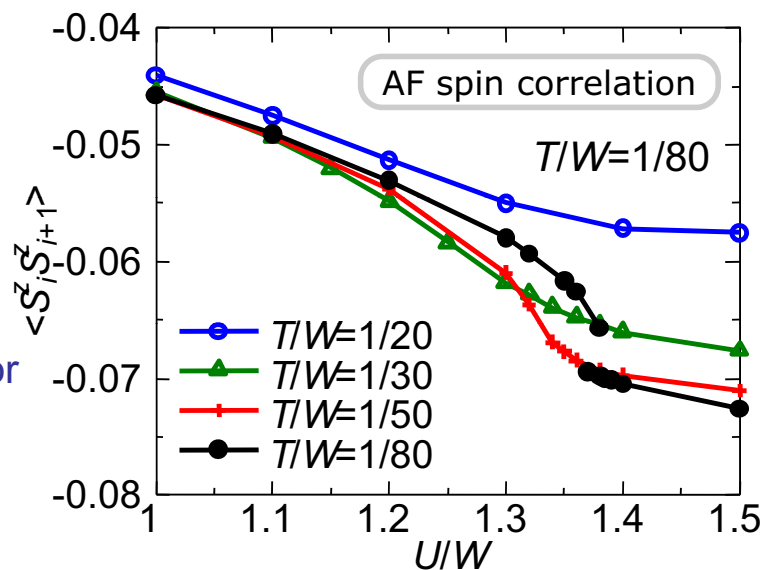
Spin correlation function

Nearest-neighbor spin correlation $\langle S_i^z S_{i+1}^z \rangle$

Insulating phase:
AF correlation \Rightarrow
monotonic
enhancement

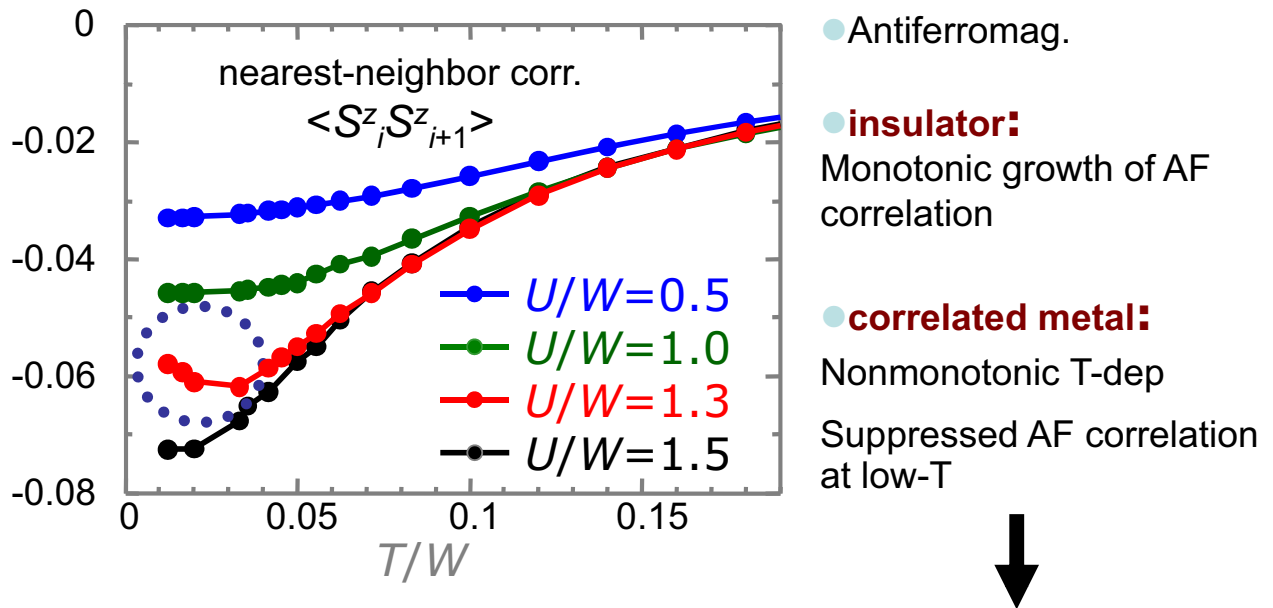
Metallic phase:
Nonmonotonic behavior

Recover of
itinerancy



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Temperature Dependence of Spin Correlations



Characteristic for frustrated systems near MIT

- recovery of coherence
- relax frustration

Dynamical Susceptibility near Mott Transition

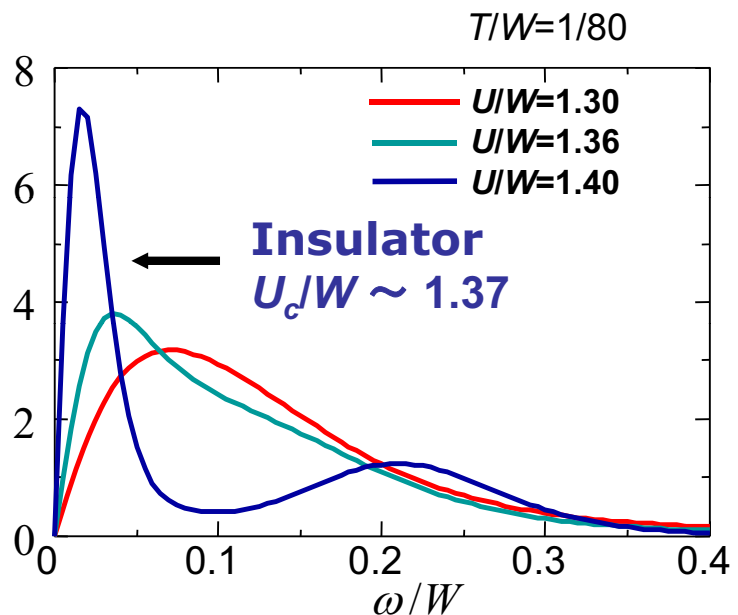
Imaginary part of local susceptibility

$$\chi_{loc}(\omega) = -i \int \langle [S_i^z(t), S_i^z(0)] \rangle e^{-it\omega} dt$$

metal \Rightarrow insulator

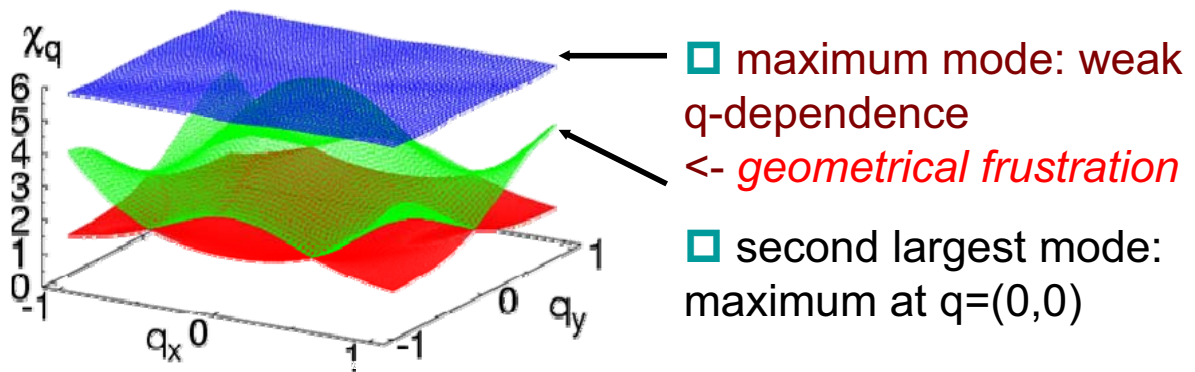
Double peak

Metallic phase:
Renormalized single peak



Magnetic instabilities

Static susceptibility $\chi_m(\mathbf{q})$: 3 magnetic modes



Consistent with previous studies

▣ FLEX approximation

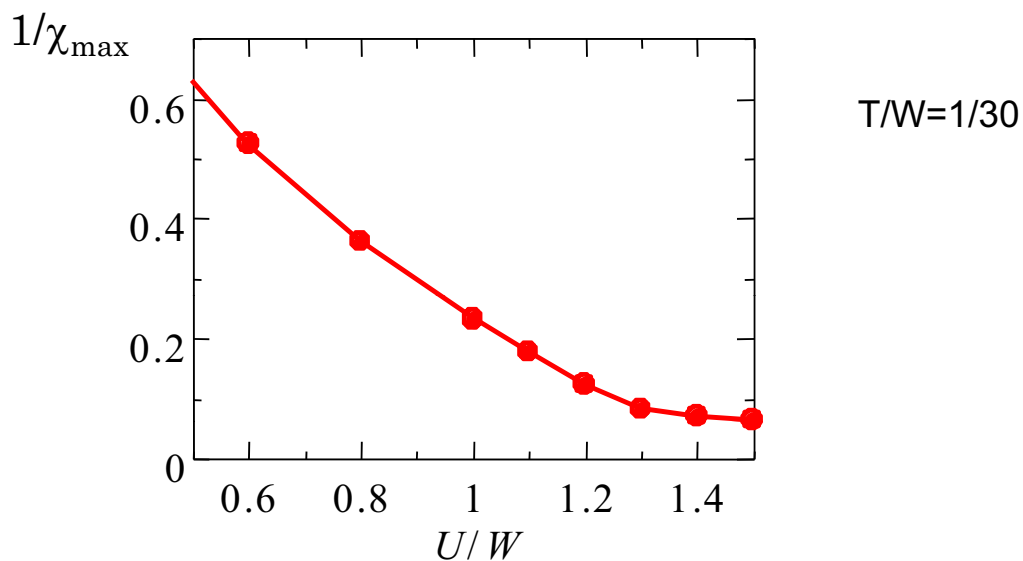
Imai, et al., PRB 68, 195103 (2003)

▣ QMC simulation

Bulut, et al. PRL 95, 37001 (2005)

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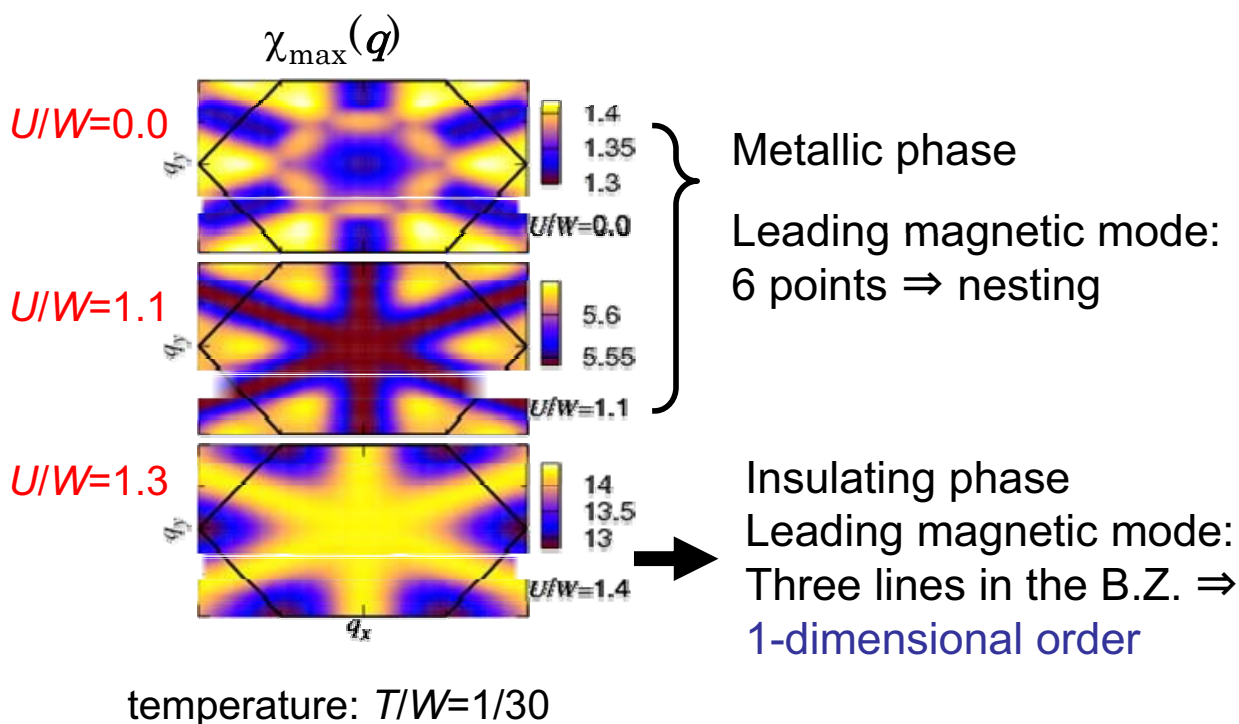
Suppressed Magnetic Instability



Mott transition : $U_c/W \sim 1.35$

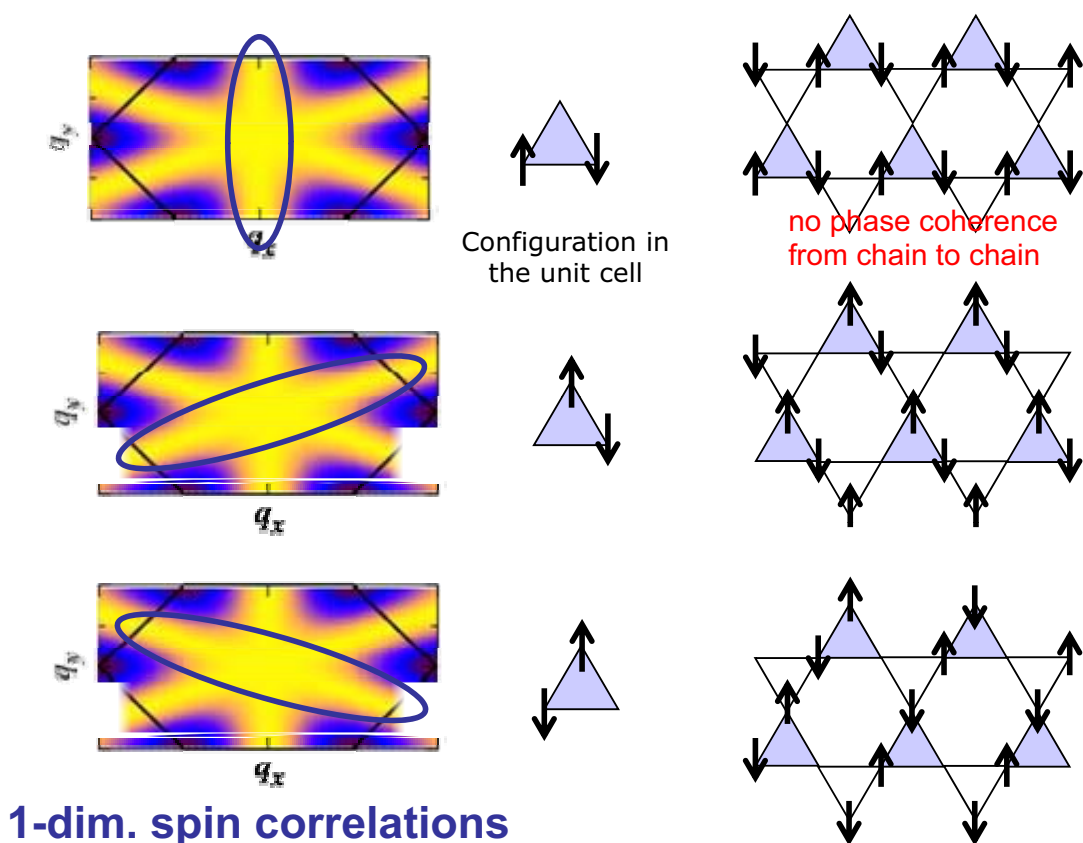
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Wavevector dependence of dominant mode



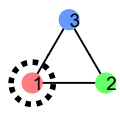
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Spin correlations in the real space



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Quasiparticle Renormalization Factor



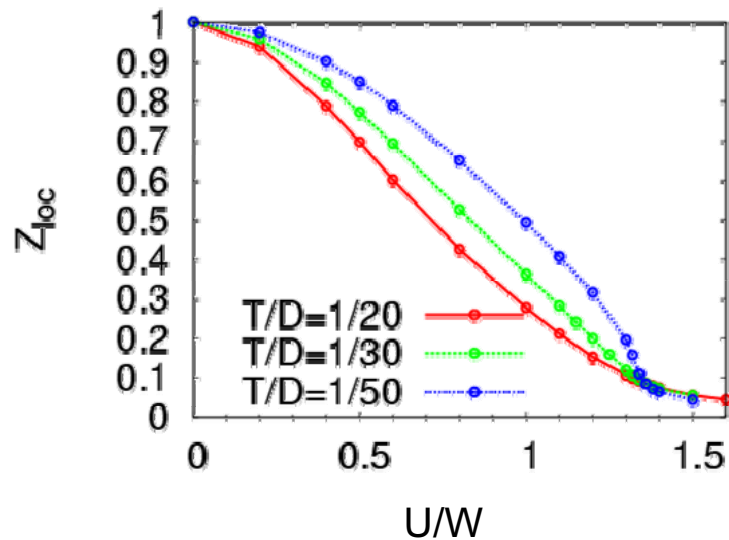
self energy

$$\Sigma_{11}(i\omega_n)$$



Renormalization factor Z_{loc}

Mass enhancement
 $\sim 10-20$
near Mott transition



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PART B

Mott Transition in Anisotropic Triangular-Lattice Hubbard Model

- Phase Boundary Topology
- Heavy Quasiparticles

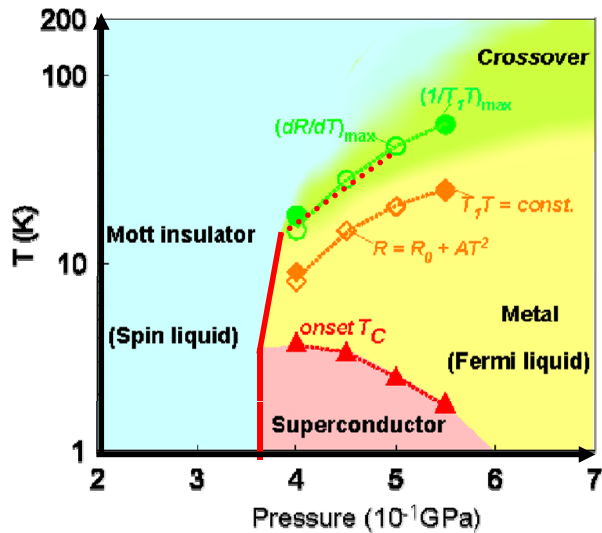
[Ohashi, Momoi, Tsunetsugu, and Kawakami, in preparation]

Mott transition in κ -type organic materials

STRONG frustration

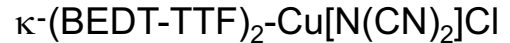


Y. Kurosaki et al., PRL 95, 177001 (2005)

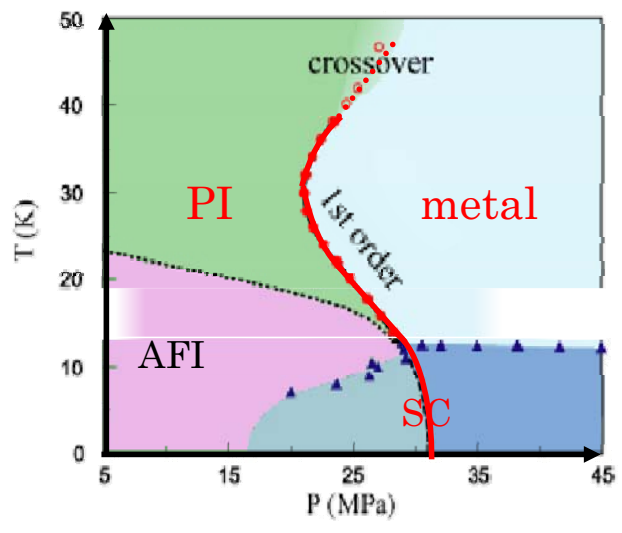


[talk by Kanoda, Aug. 8]

INTERMED. frustration



F. Kagawa et al., PRB 69, 064511 (2004)

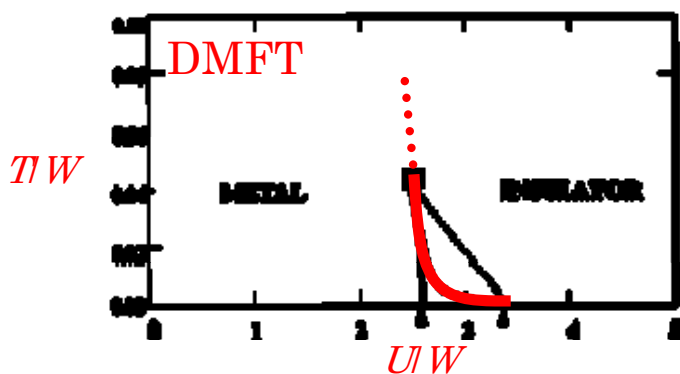


Reentrant !!

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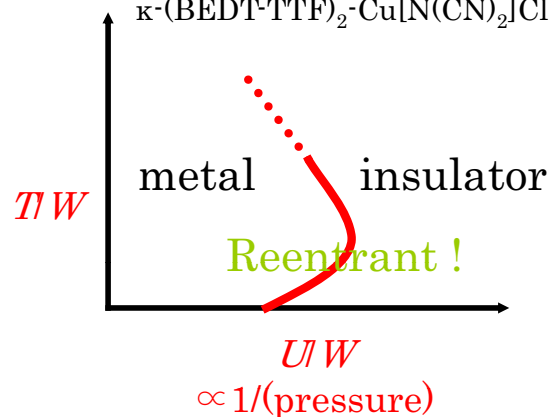
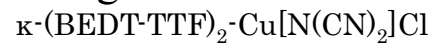
Mott transition line in Phase Diagram

$d=\infty$ Hubbard model



Georges et al., RMP, 68, 13 (1996)

organic conductor



$\propto 1/(\text{pressure})$

Anisotropic
Hubbard model

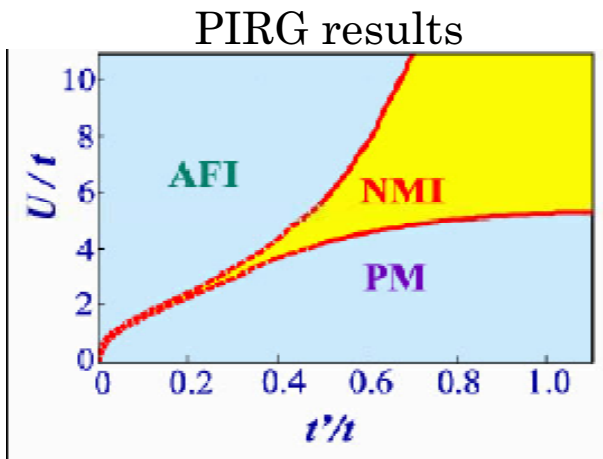


Cellular-DMFT

effects of 1-site approx. or frustration?

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Previous studies



Ground state

- PIRG

Kashima & Imada, JPSJ 70, 3052 (2001).

- Cellular-DMFT + ED

Kyung & Tremblay, PRL 97, 046402 (2006).

- Variational Monte Carlo

Watanabe et al., JPSJ 75, 074707 (2006).

- Exact Diagonalization

Koretsune, Motome, & Furusaki, cond-mat/0703111.

Finite temperature

- Cellular-DMFT + QMC

Parcollet et al., PRL 92, 226402 (2004)

- Correlator Projection method

Onoda & Imada, PRB 67, 161102 (2003).

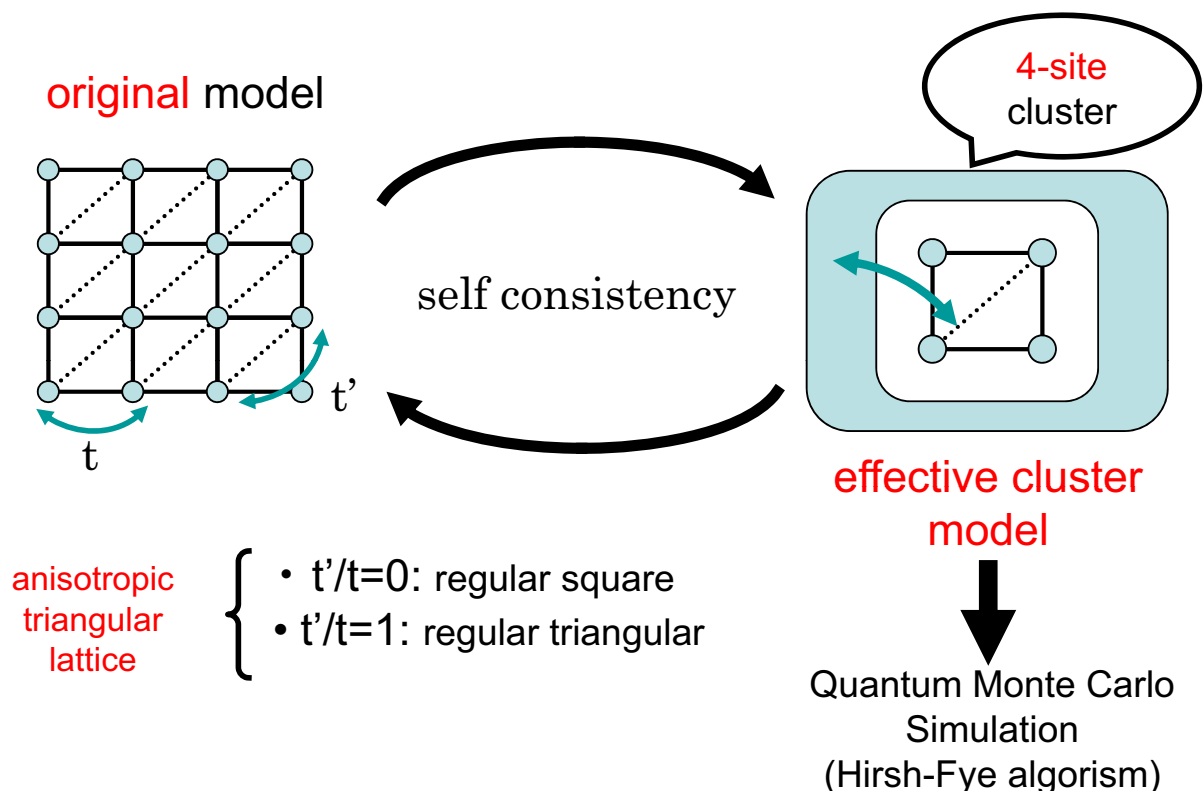
Finite-T Mott transition

- RG & MF

Onoda & Nagaosa, JPSJ 72, 2445 (2003).

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Cellular-DMFT for Anisotropic Triangular Lattice Model



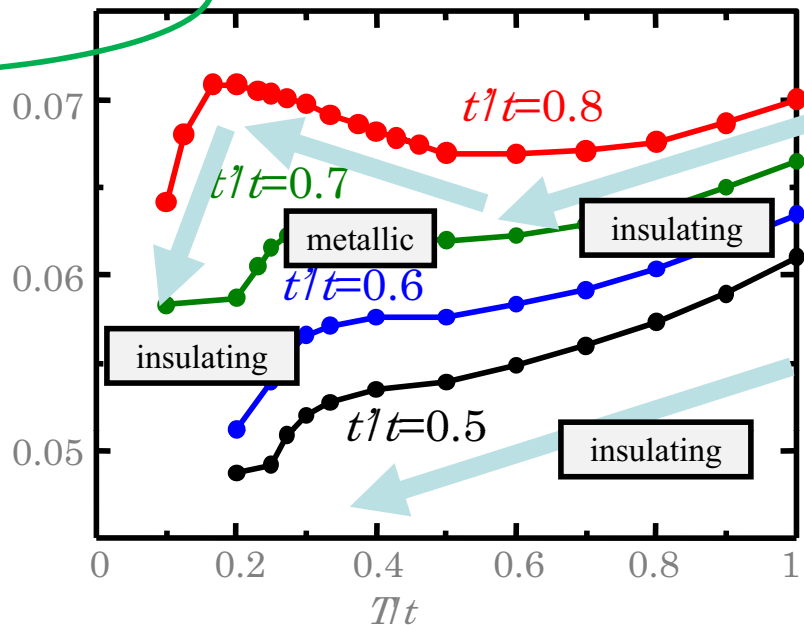
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T-Dependence of Double Occupancy

Insulator-Metal-Insulator Transition?

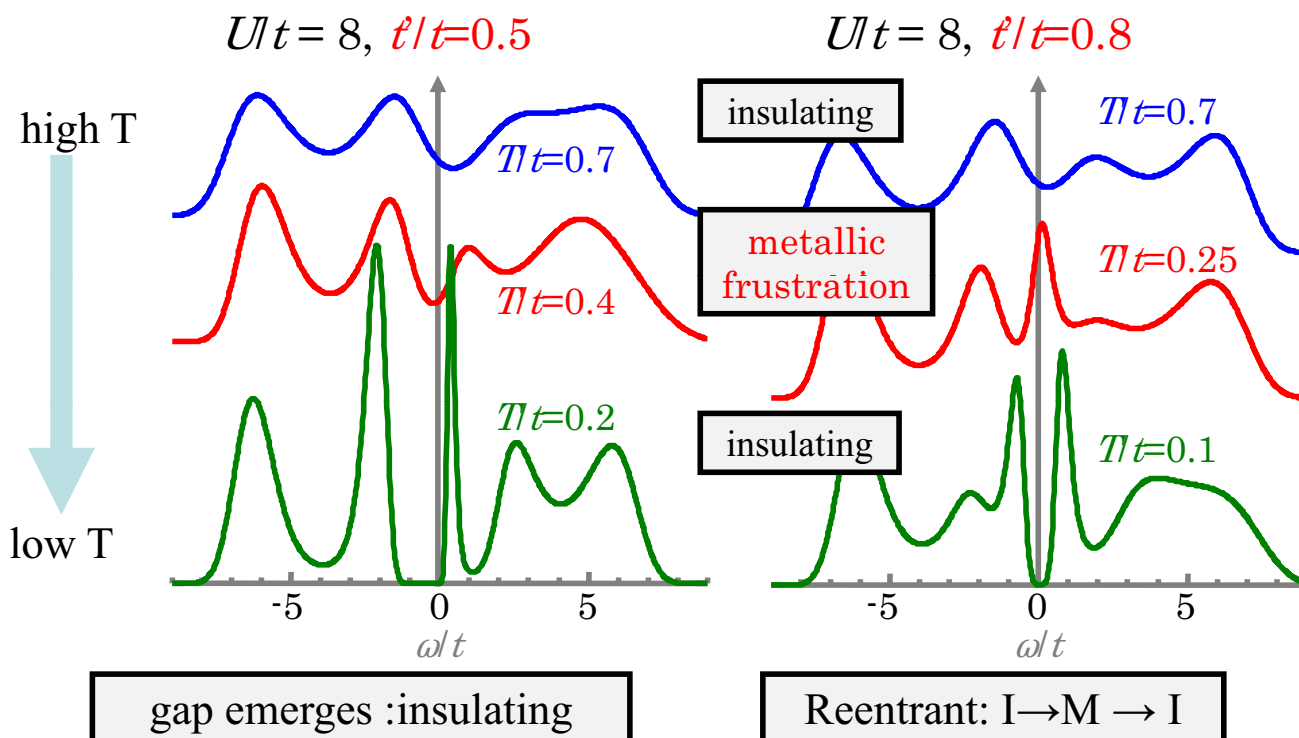
- large t' -strong GF: nonmonotonic T-dep.
- small t' -weak GF: almost monotonic insulating

Repulsion $U/t = 8$



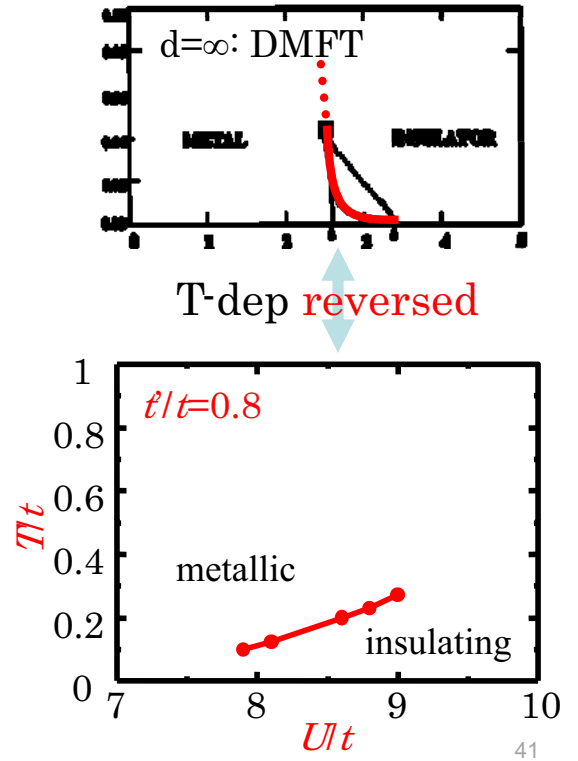
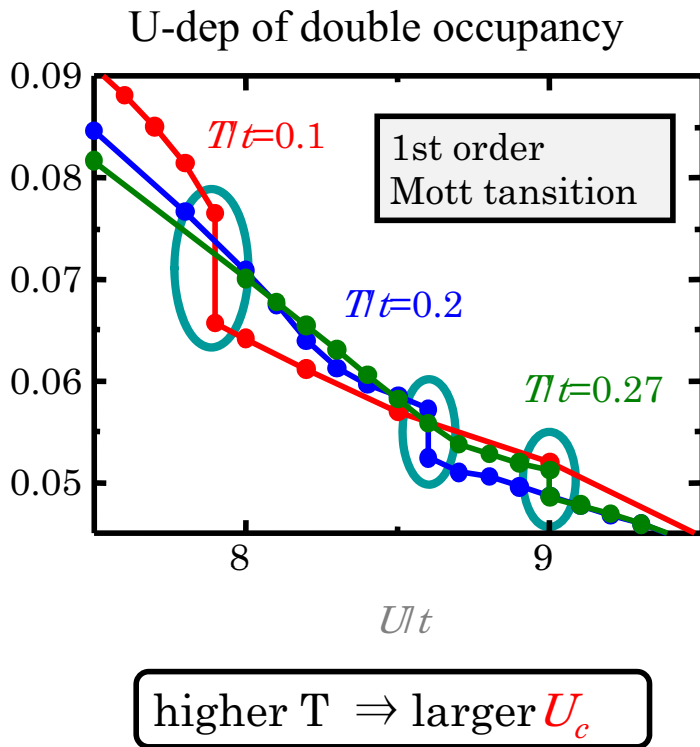
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Electron Spectral Function

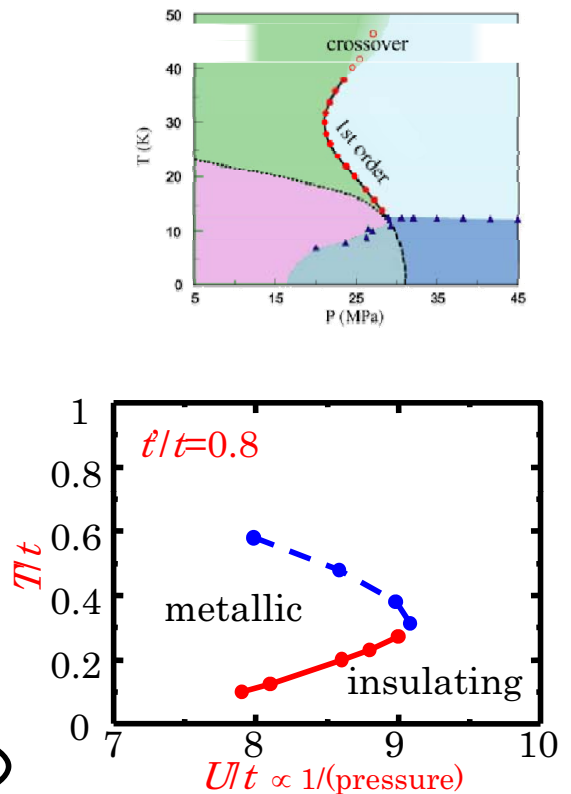
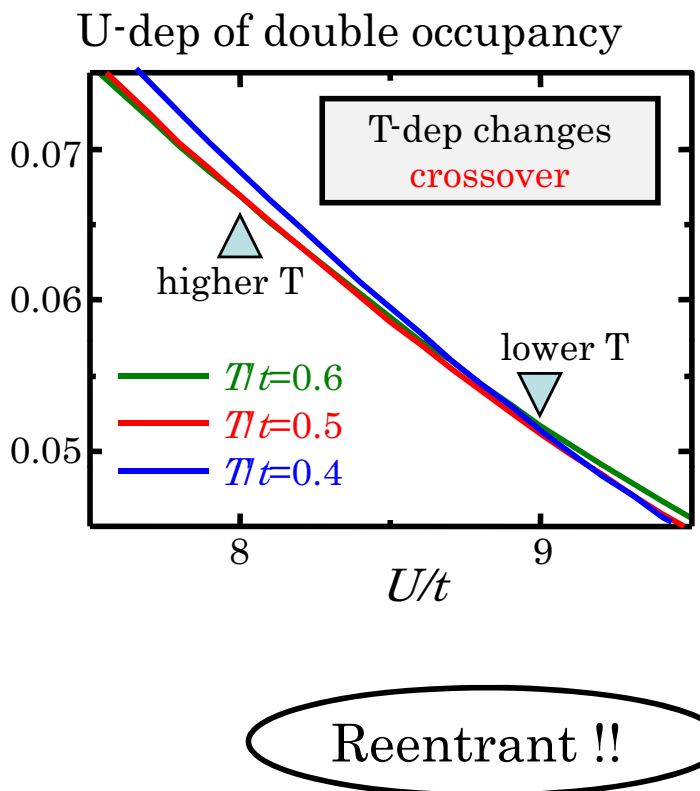


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Mott Transition

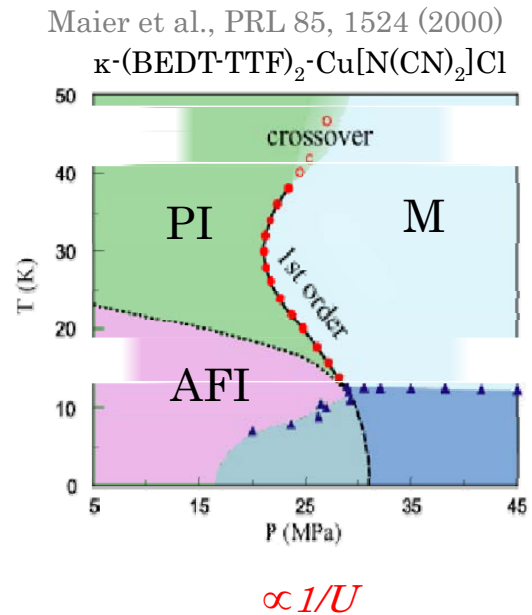
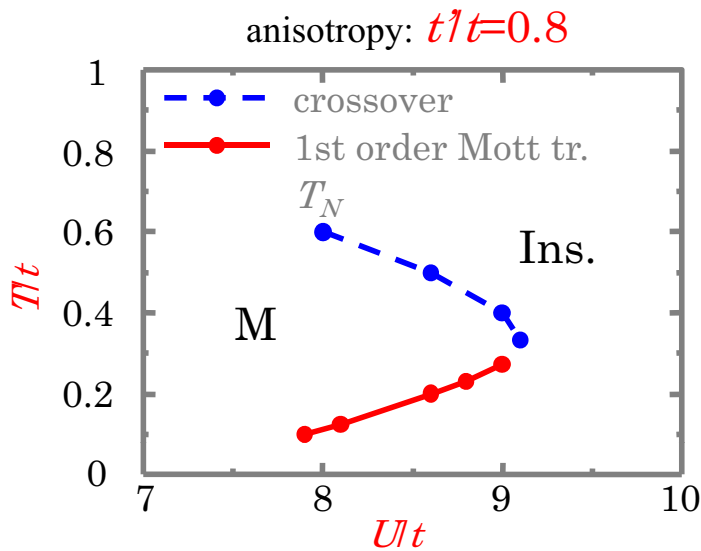


Crossover at a higher temperature



Comparison with Organic Materials

consistent with experiments



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Summary (1)

Kagome lattice Hubbard model

Cellular dynamical mean field theory

- Metal-insulator transition
 - 1st order transition : $U_c/W \sim 1.37$
- Strongly correlated metal
 - Whole bands are renormalized
 - large mass enhancement
 - nonmonotonic temperature dependence of spin correlation functions
- Magnetic instability
 - one-dimensional spin correlations

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Summary (2)

Anisotropic Triangular Lattice Hubbard model (mainly $t'/t=0.8$)

Cellular dynamical mean field theory

- Metal-insulator transition
 - different shape of transition line from unfrustrated systems
 - entropy effects
- Intermediate Correlation Regime:
 - “reentrant” insulator \rightarrow metal \rightarrow insulator transition
 - heavy quasiparticle formation in the intermediate metallic phase
 - gap formation inside heavy qp band
- Magnetic instability
 - transition to paramagnetic insulator phase