

# **4-Nonlinear Advection Equation**

# Nonlinear Advection Equation

- We turn our attention to the scalar conservation law

$$\frac{\partial u}{\partial t} + \frac{\partial f(u)}{\partial x} = 0$$

- Where  $f(u)$  is, in general, a nonlinear function of  $u$ .
- To gain some insights on the role played by nonlinear effects, we start by considering the inviscid Burger's equation:

$$\frac{\partial u}{\partial t} + \frac{\partial}{\partial x} \left( \frac{u^2}{2} \right) = 0$$

# Nonlinear Advection Equation

- We can write Burger's equation also as

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = 0$$

- In this form, Burger's equation resembles the linear advection equation, with the only difference being that the velocity is no longer constant, but it is equal to the solution itself.
- The characteristic curve for this equation is

$$\frac{dx}{dt} = u(x, t) \quad \Longrightarrow \quad \frac{du}{dt} = \frac{\partial u}{\partial t} + \frac{\partial u}{\partial x} \frac{dx}{dt} = 0$$

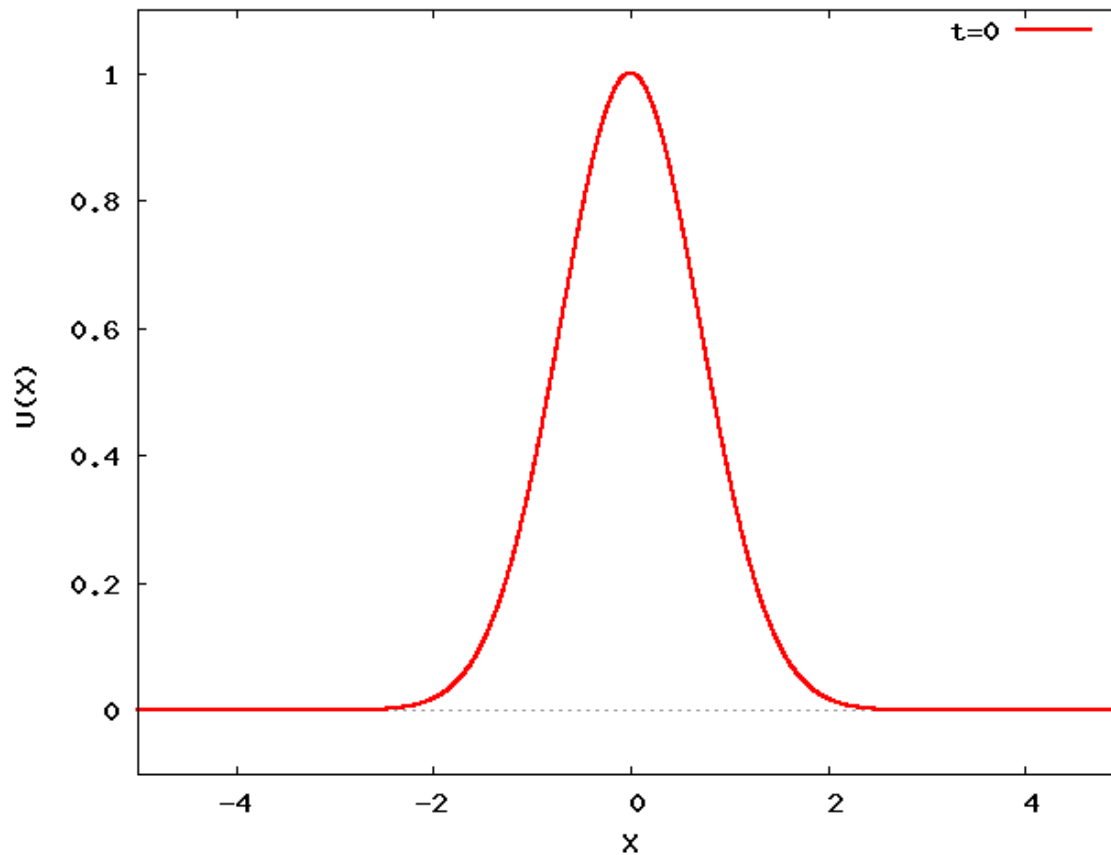
- which tells us that  $u$  is constant along the curve  $dx/dt=u(x,t)$ .
- Along these curves the PDE becomes an ODE.

# Nonlinear Advection Equation

- ❑ A quantity that remains constant along a characteristic curve is called a *Riemann invariant*.
- ❑ In this simple case,  $u$  is a Riemann invariant.
- ❑ Considering that  $dx/dt = u(x,t)$  we deduce that characteristic curves are again straight lines: values of  $u$  associated with some fluid element do not change as that element moves.
- ❑ However, since  $u(x,t)$  can change in space, these lines are not necessarily parallel to each other as was the case for the linear advection equation.

# Nonlinear Advection Equation

- Now consider the initial Gaussian profile at  $t=0$ :

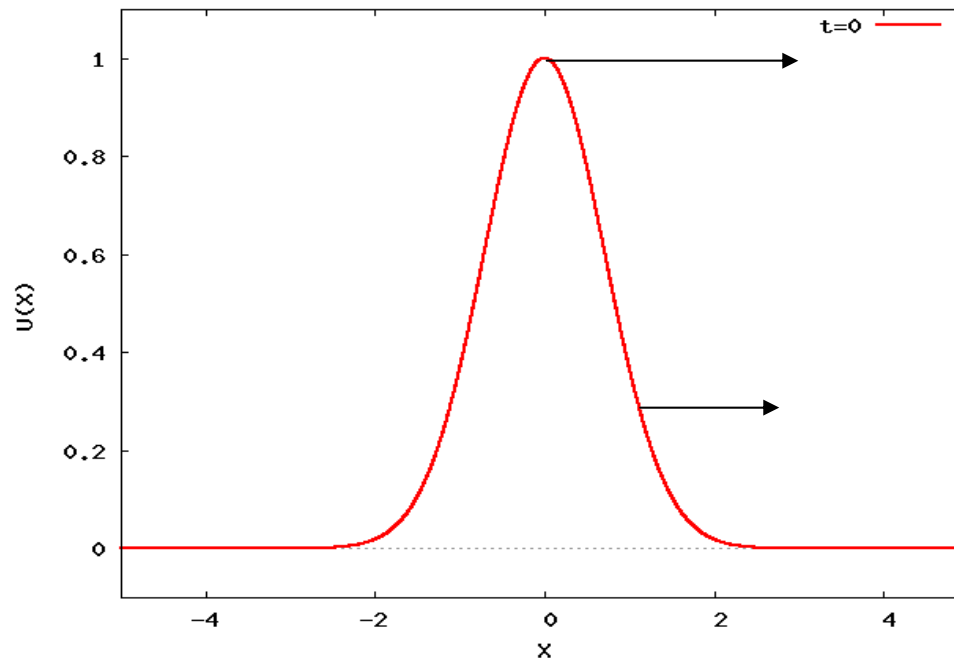


- What's going to happen at  $t > 0$  ?

# Nonlinear Advection Equation

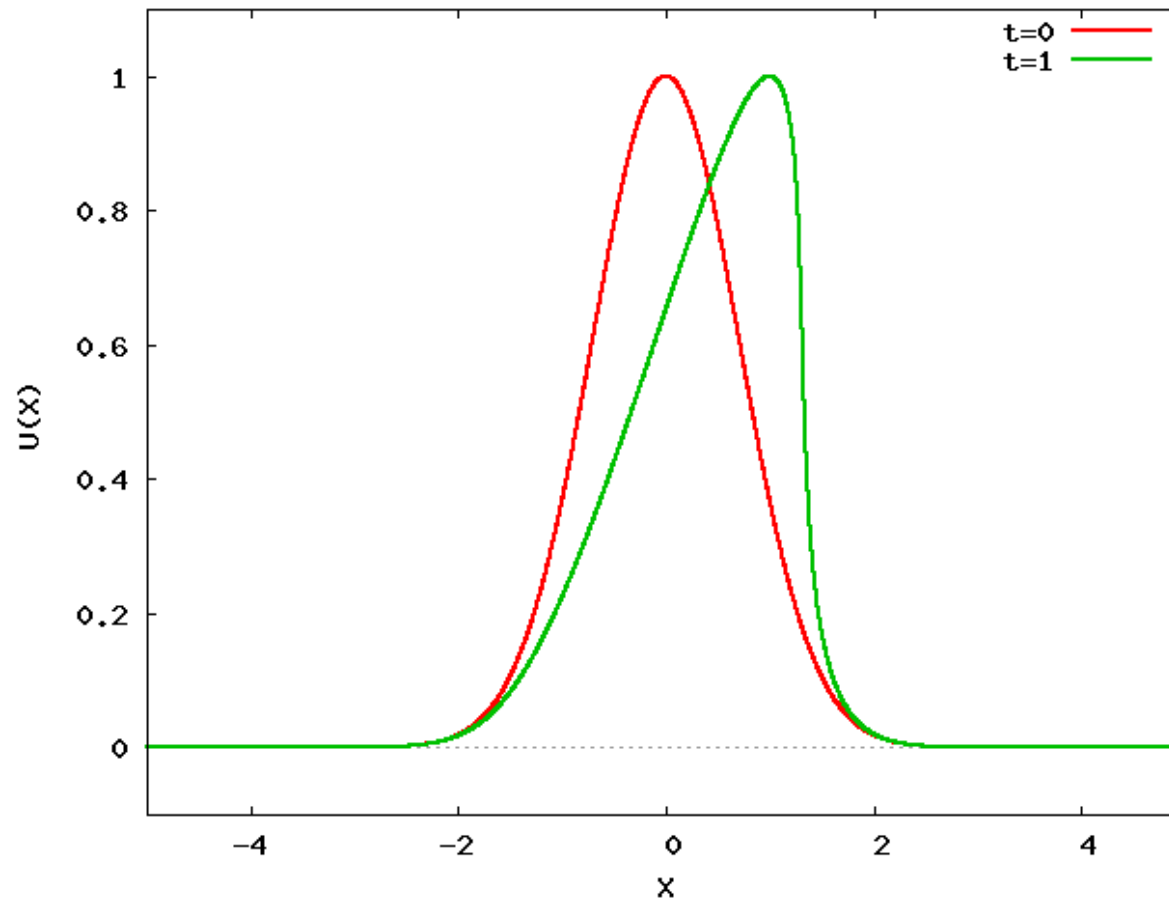
□ From 
$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = 0$$

one can predict that, higher values of  $u$  will propagate faster than lower values: this leads to a wave steepening, since upstream values will advance faster than downstream values.



# Nonlinear Advection Equation

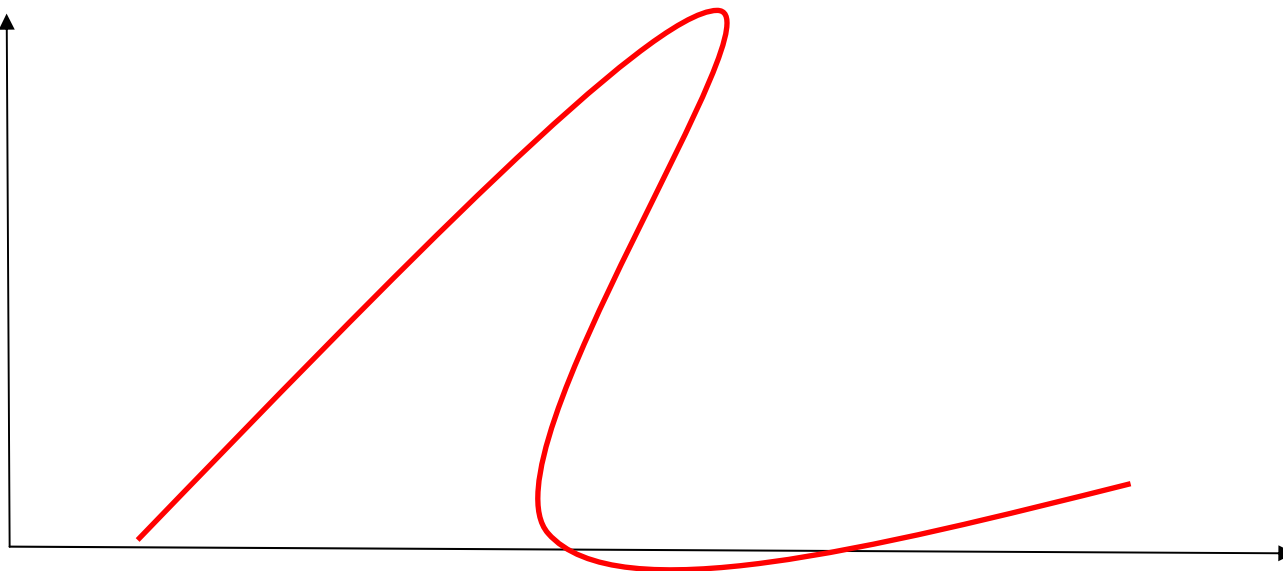
□ Indeed, at  $t=1$  the wave profile will look like:



□ the wave steepens...

# Nonlinear Advection Equation

- If wait more, we should get something like this:

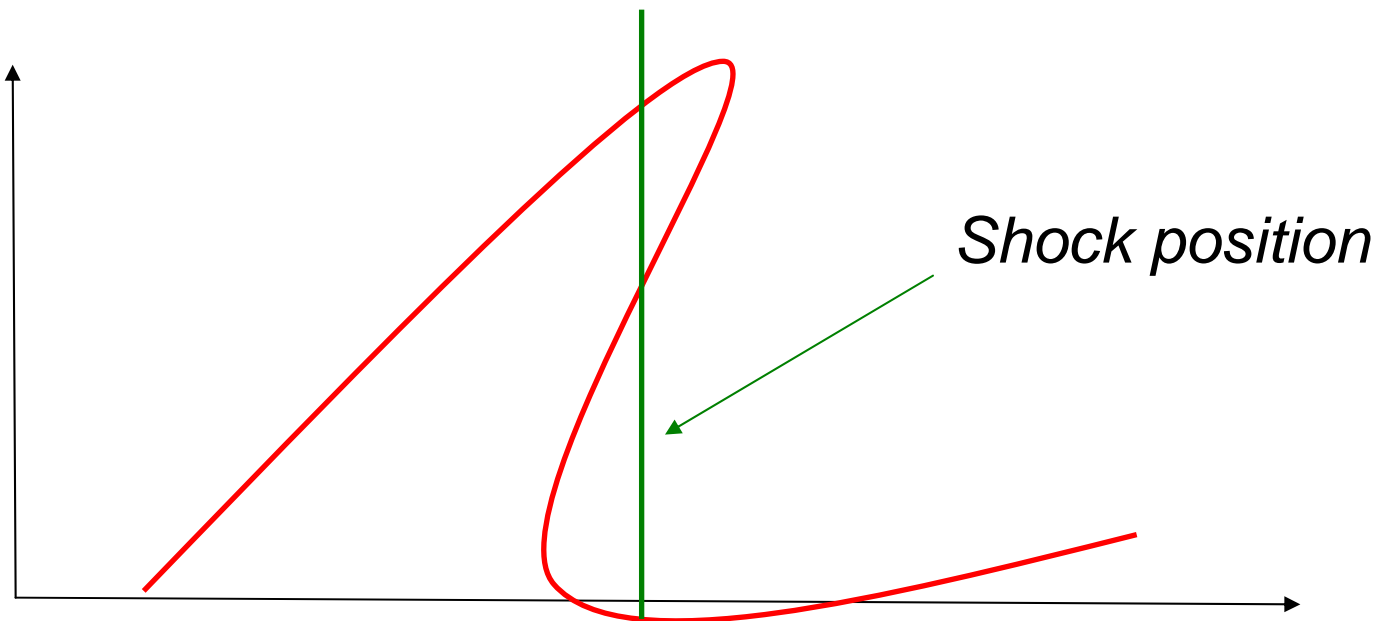


- A multivalued functions ??!?! → Clearly Unphysical !!



# Nonlinear Advection Equation

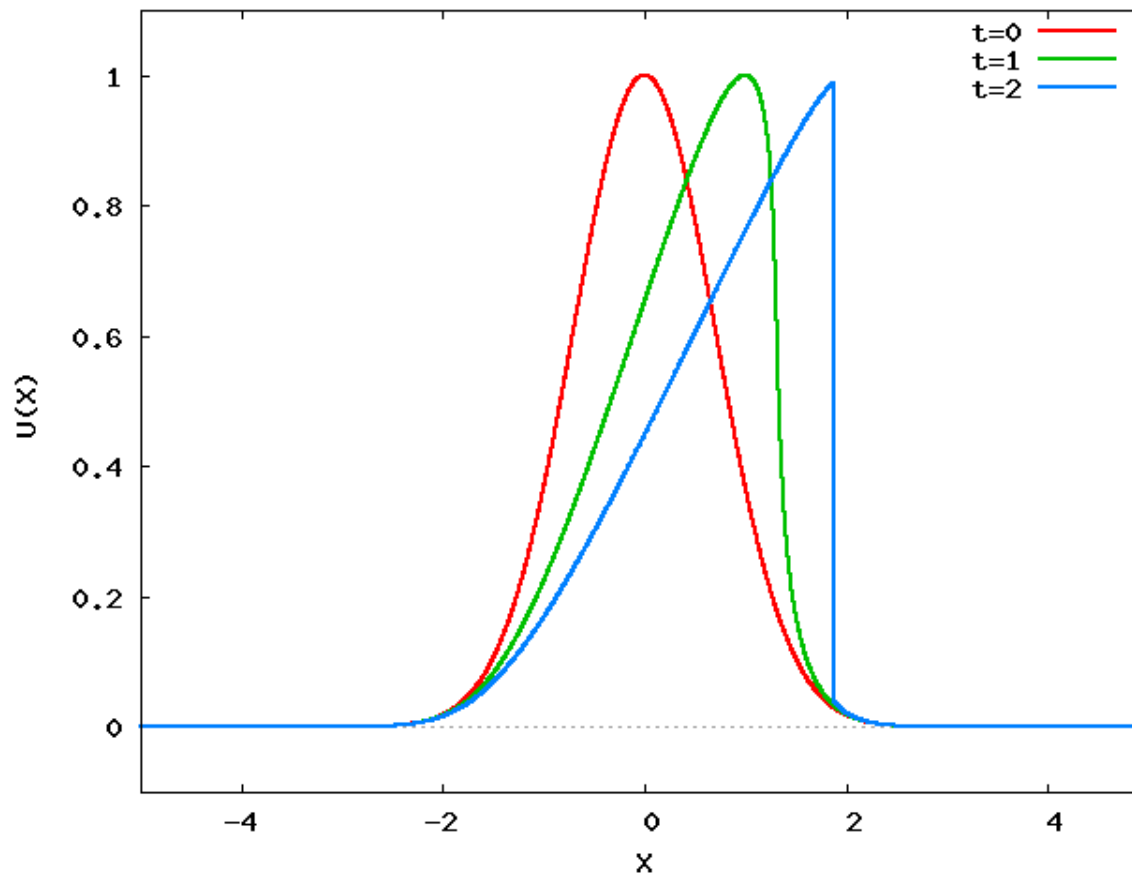
- The physical solution is to place a discontinuity there:  
a *shock wave*.



- Since the solution is no longer smooth, the *differential form* is not valid anymore and we need to consider the *integral form*.

# Nonlinear Advection Equation

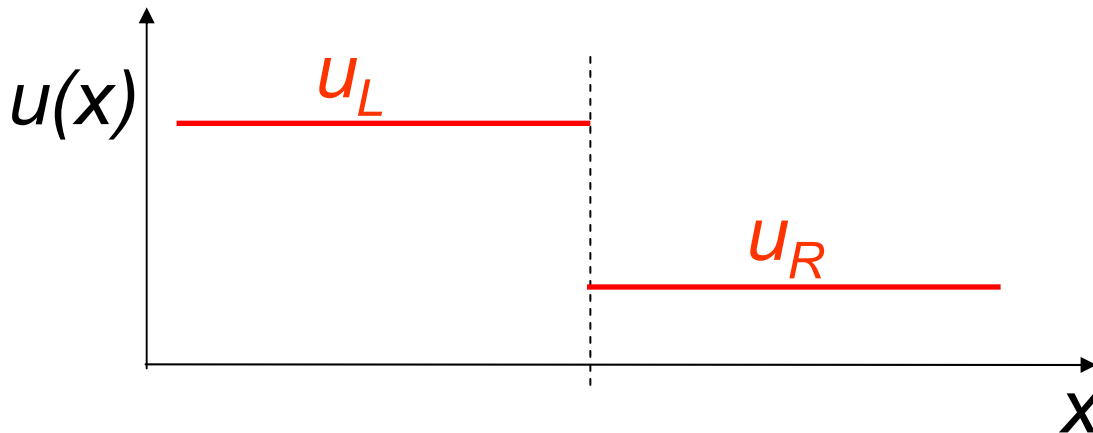
□ This is how the solution should look like:



□ Such solutions to the PDE are called **weak solutions**.

# Nonlinear Advection Equation

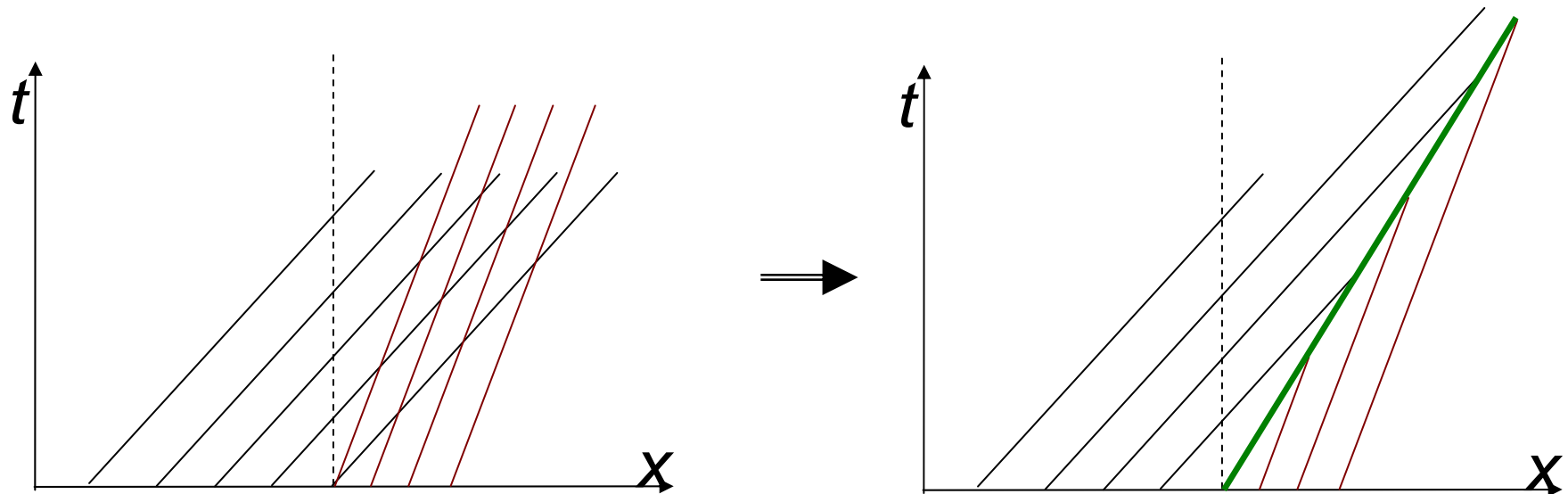
- Let's try to understand what happens by looking at the characteristics.
- Consider two states initially separated by a jump at an interface:



- Here, the characteristic velocities on the left are greater than those on the right.

# Nonlinear Advection Equation

- The characteristic will intersect, creating a *shock*:



- The shock speed is such that  $\lambda(u_L) > S > \lambda(u_R)$ . This is called the *entropy condition*.

# Nonlinear Advection Equation

- The shock speed  $S$  can be found using the Rankine-Hugoniot jump conditions, obtained from the integral form of the equation:

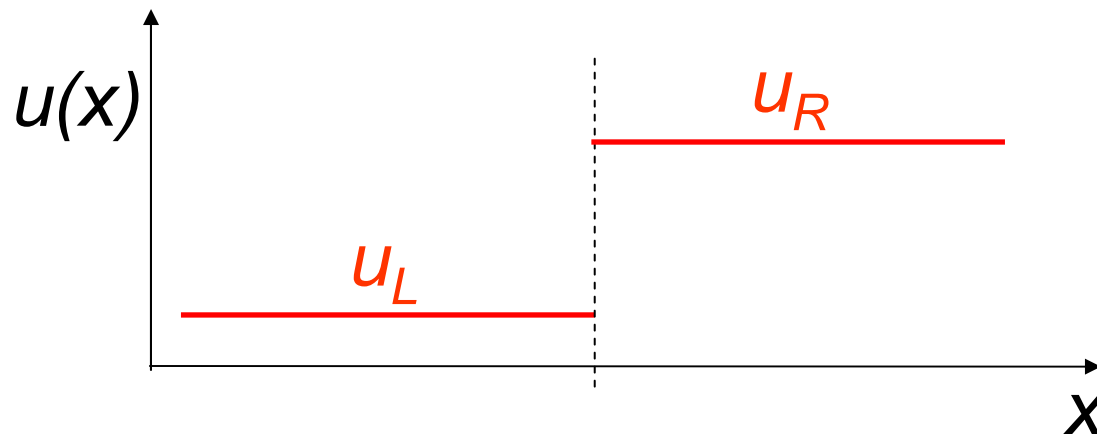
$$f(u_R) - f(u_L) = S(u_R - u_L)$$

- For Burger's equation  $f(u) = u^2/2$  so that one finds the shock speed as

$$S = \frac{u_L + u_R}{2}$$

# Nonlinear Advection Equation

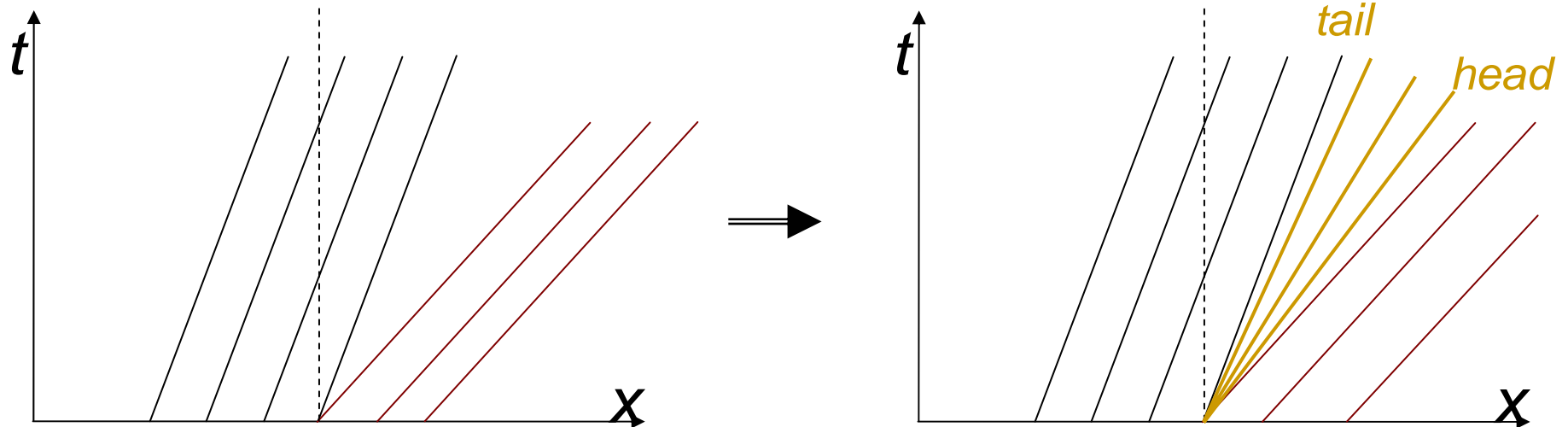
□ Let's consider the opposite situation:



□ Here, the characteristic velocities on the left are smaller than those on the right.

# Nonlinear Advection Equation

□ Now the characteristics will diverge:



□ Putting a shock wave between the two states would be incorrect, since it would violate the entropy condition. Instead, the proper solution is a *rarefaction wave*.

# Nonlinear Advection Equation

- ❑ A rarefaction wave is a nonlinear wave that smoothly connects the left and the right state. It is an expansion wave.
- ❑ The solution between the states can only be self-similar and takes on the range of values between  $u_L$  and  $u_R$
- ❑ The head of the rarefaction moves at the speed  $\lambda(u_R)$ , whereas the tail moves at the speed  $\lambda(u_L)$ .
- ❑ The general condition for a rarefaction wave is  $\lambda(u_L) < \lambda(u_R)$
- ❑ Both rarefactions and shocks are present in the solutions to the Euler equation. Both waves are nonlinear.



# Nonlinear Advection Equation

- These results can be used to write the general solution to the Riemann problem for the Burger's equation:
  - If  $u_L > u_R$  the solution is a *shock wave*. In this case

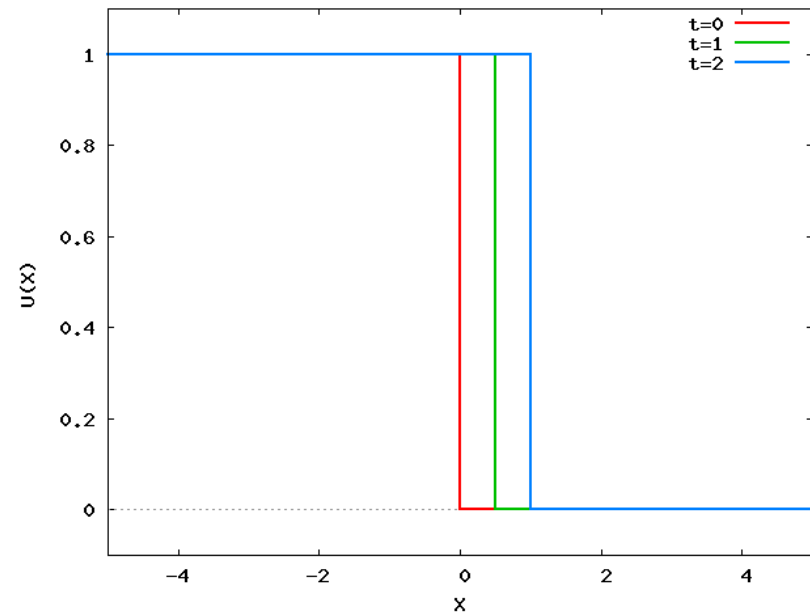
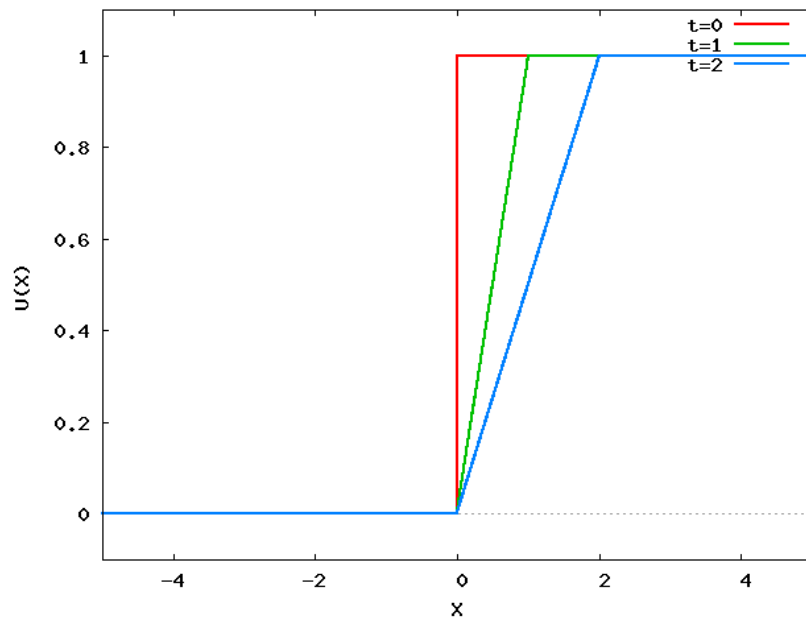
$$u(x, t) = \begin{cases} u_L & \text{if } x - St < 0 \\ u_R & \text{if } x - St > 0 \end{cases}, \quad S = \frac{u_L + u_R}{2}$$

- If  $u_L < u_R$  the solution is a *rarefaction wave*. In this case

$$u(x, t) = \begin{cases} u_L & \text{if } x/t \leq u_L \\ x/t & \text{if } u_L < x/t < u_R \\ u_R & \text{if } x/t > u_R \end{cases}$$

# Nonlinear Advection Equation

□ Solutions look like



for a rarefaction and a shock, respectively.

# Nonlinear Advection Equation

- An implementation is given in `burger.c`.