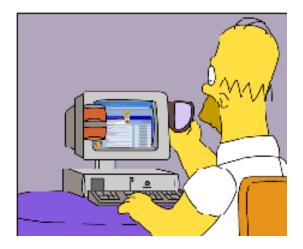
<u>6-Higher order methods,</u> <u>multidimensional integration</u> <u>and the PLUTO Code</u>



"Don't worry brain, the computer will do our thinking now"

-- Homer Simpson



<u>High Order Methods:</u> <u>The PLUTO code</u>

Multi-physics, multi-algorithm modular code for the solution of compressible, highly supersonic flows, based on the RSA:

Newtonian Hydrodynamics

Relativistic Hydrodynamics

Ideal/resistive Magneto Hydrodynamics (MHD)

Relativistic MHD

oriented towards the treatment of astrophysical flows in presence of discontinuities;

□ For more information: <u>*http://plutocode.to.astro.it*</u>

The PLUTO code

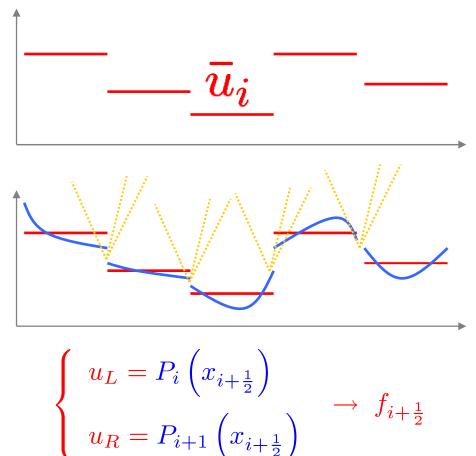
- Pluto achieves second or better accuracy by following a strategy commonly adopted in shock-capturing schemes: the *Reconstruct*-Solve-Average (RSA).
- Reconstruct: piecewise polynomial interpolation is provided inside each cell to increase the spatial order of integration;
- Solve: A Riemann problem is solved between left and right edge interpolated values;
- Average: the final update is done in a conservative way by using either dimensionally split or unsplit schemes.

<u>Reconstruct-Solve-Average</u>

- Start from zone averages, break the problem into 3 pieces:
- 1 Piecewise polynomial reconstruction

 $u_i(x) = P_i(x)$, for $x_{i-\frac{1}{2}} < x < x_{i+\frac{1}{2}}$

- 2 Solve Riemann problem between left and right states
- 3 Form new averages (evolve):



$$\bar{u}_i^{n+1} = \bar{u}_i^n - \frac{\Delta t}{\Delta x} \left[f_{i+\frac{1}{2}} - f_{i-\frac{1}{2}} \right]$$

Must be consistent with data representation

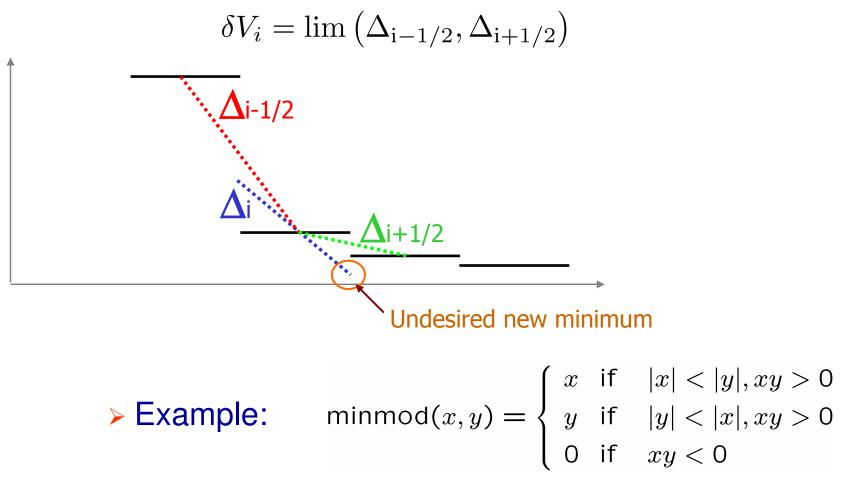
$$\frac{1}{\Delta x_i} \int_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} P_i(x) dx = \overline{u}_i$$

Satisfy monotonicity constraints:

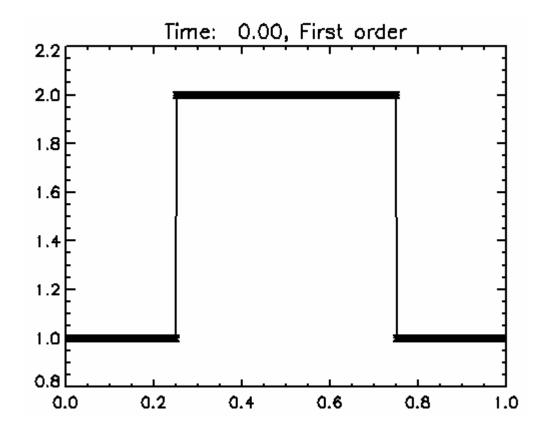
$$\min(P_i(x)) \ge \min\left(\bar{u}_{i-1}, \bar{u}_i, \bar{u}_{i+1}\right)$$
$$\max(P_i(x)) \le \max\left(\bar{u}_{i-1}, \bar{u}_i, \bar{u}_{i+1}\right)$$

- *¬* no new extrema allowed (Total Variation Diminishing (TVD) schemes)
- \rightarrow Oscillation free solution

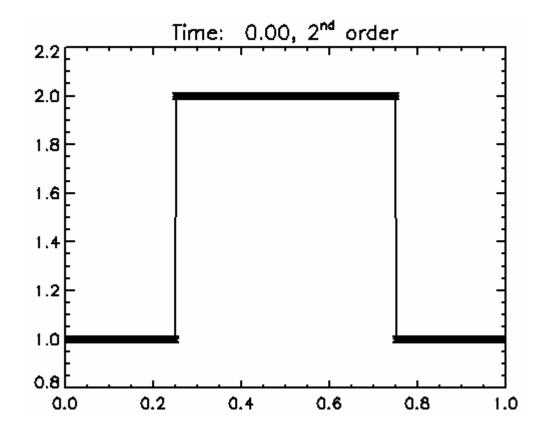
□ For 2nd-order interpolant, we use $V(x) = V_i + \frac{\delta V}{\Delta x}(x - x_i)$ □ Use slope limiters to avoid introducing new extrema:



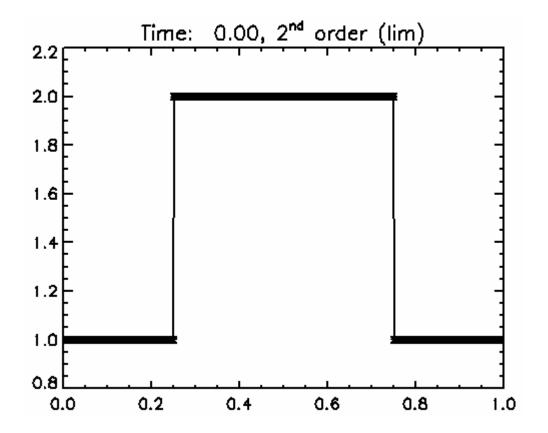
For example, consider the advection of a square pulse. For a first-order spatial reconstruction will evolve like:



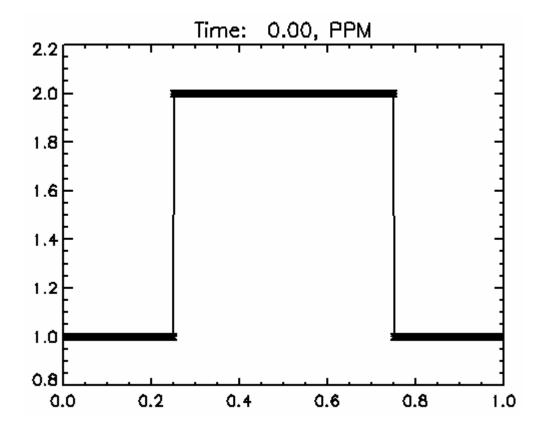
with a second-order unlimited reconstruction:



u with 2nd order, limited reconstruction:



□ with limited parabolic reconstruction:



Multi Dimensional Integration

- Integration in more than one dimensions can be achieved using two distinct approaches:
 - 1. Dimensionally Split schemes: solve the PDE as a sequence of 1-D sub-problems.

$$\mathbf{q}^* = \mathbf{q}^n - \Delta t \mathcal{L}_x(\mathbf{q}^n) \qquad \mathbf{q}^{n+1} = \mathbf{q}^* - \Delta t \mathcal{L}_y(\mathbf{q}^*)$$

2. Dimensionally Unsplit schemes: solve the full problem:

$$\mathbf{q}^{n+1} = \mathbf{q}^n - \Delta t \mathcal{L}_x(\mathbf{q}^n) - \Delta t \mathcal{L}_y(\mathbf{q}^n)$$

High order Integration in time

A simple and effective way to achieve 2nd or 3rd order accuracy in time is to treat the PDE in semi-discrete form:

$$\int \left(\frac{\partial \mathbf{q}}{\partial t} + \nabla \cdot \mathbf{F}\right) dV = 0 \quad \Longrightarrow \quad \frac{d\bar{\mathbf{q}}}{dt} = -\oint \mathbf{F} \cdot d\mathbf{S}$$

and treating the right hand side as an ODE, e.g., using predictor/corrector methods,

$$\begin{cases} \bar{\mathbf{q}}^* = \bar{\mathbf{q}}^n - \Delta t \left(\oint \mathbf{F}^n \cdot d\mathbf{S} \right) \\ \bar{\mathbf{q}}^{n+1} = \frac{1}{2} \left(\bar{\mathbf{q}}^n + \bar{\mathbf{q}}^* - \Delta t \oint \mathbf{F}^* \cdot d\mathbf{S} \right) \end{cases}$$

Practice Session: Shock-Tube integration with high order method

 \Box file \rightarrow pluto1d

Practice Session: 2D Jet Simulation

□ file \rightarrow pluto2d