



**The Abdus Salam  
International Centre for Theoretical Physics**



**1953-2**

**International Workshop on the Frontiers of Modern Plasma Physics**

*14 - 25 July 2008*

**Non-Newtonian mechanics and manipulation of oscillation centers**

I. Dodin

*Princeton University, Plasma Physics Laboratory, U.S.A.*



# Non-Newtonian mechanics and manipulation of oscillation centers

ILYA Y. DODIN AND NATHANIEL J. FISCH

*Princeton University*

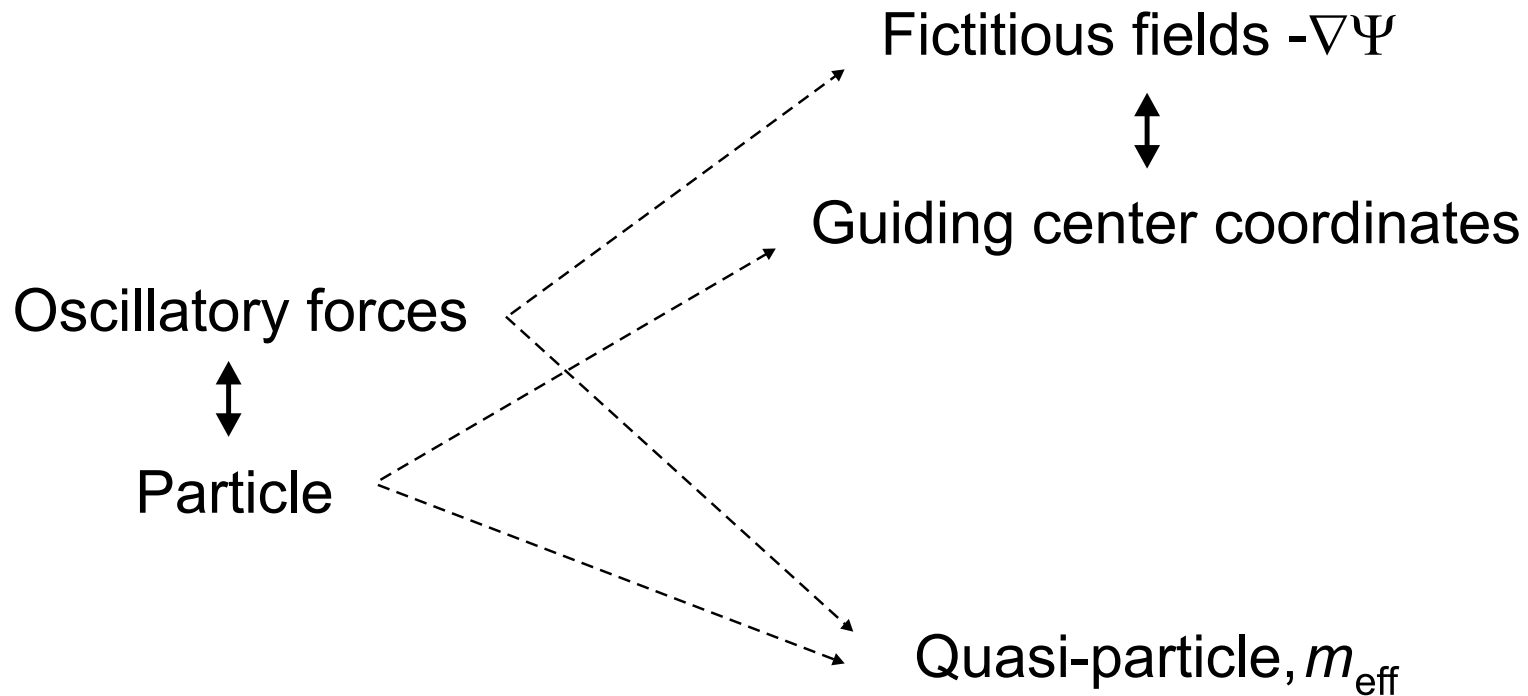
INTERNATIONAL WORKSHOP ON THE FRONTIERS OF  
MODERN PLASMA PHYSICS

ICTP, Trieste, Italy

July 14-25, 2008



# Guding Centers as Quasi-particles



*Fully relativistic object-oriented approach:  
all ponderomotive, diamagnetic etc. effects  
are embedded in the effective mass*



# Examples

$$\mathcal{H} = \sqrt{m_{\text{eff}}^2 c^4 + \bar{p}^2 c^2}, \quad m_{\text{eff}} = c^{-2} (\mathbf{J} \cdot \boldsymbol{\Omega} - \langle L \rangle)_{\langle \bar{\mathbf{v}} \rangle = 0}$$

- Particle in a magnetic field:

$$\mathcal{H} = \sqrt{m^2 c^4 + 2\mu B m c^2 + p_{\parallel}^2 c^2} \approx$$

$$\approx p_{\parallel}^2 / 2m + \mu B$$

$$m_{\text{eff}} = m \sqrt{1 + 2\mu B / m c^2}$$

e.g., Grebogi and Littlejohn, *Phys. Fluids*, 1984; Boozer, *Phys. Plasmas*, 1996

- Particle in a vacuum laser pulse:

$$\mathbf{A} = (m c^2 / e) \mathbf{a}(\mathbf{r}, t) \propto e^{i\xi}$$

$$\xi = \omega t - \mathbf{k} \cdot \mathbf{r}, \quad \mathbf{k} \approx \hat{\mathbf{z}} \omega / c$$

$$m_{\text{eff}} = m \sqrt{1 + \langle a^2 \rangle_{\xi}}$$

e.g., Kibble, *Phys. Rev.*, 1966; Bauer et al., *Phys. Rev. Lett.*, 1995; Mora and Jr., *Phys. Plasmas*, 1997

- $m_{\text{eff}}$  can also be a negative or even nonanalytic function of the particle average velocity and the oscillation energy; hence non-Newtonian dynamics.



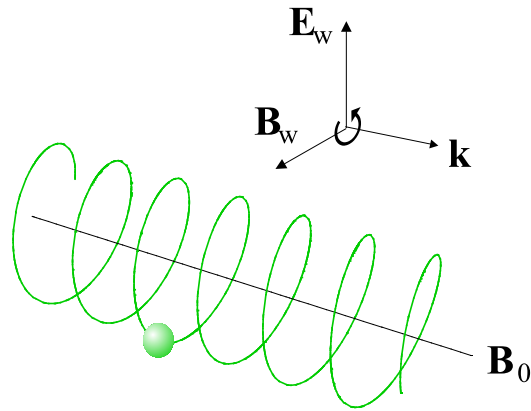
# Outline

---

- Example: laser field over a dc magnetic field
- Particles with anisotropic dispersion  $\bar{p}(\bar{v})$ :
  - Negative-mass plasma maser
  - Hamiltonian ratchet and current drive
  - Frequency doubling
- Selective barriers:
  - Nonrelativistic hybrid potentials
  - One-way walls and cooling effects
- Quantum-like effects



# Example: Laser Field + DC Magnetic field

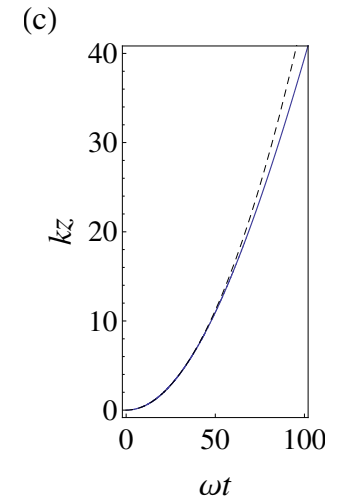
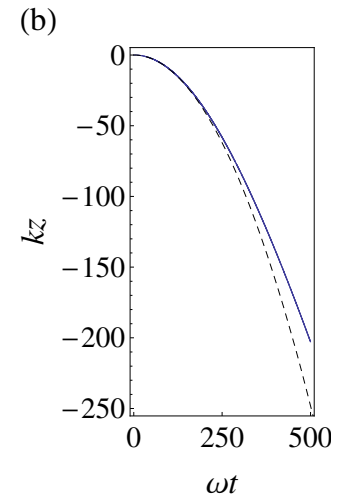
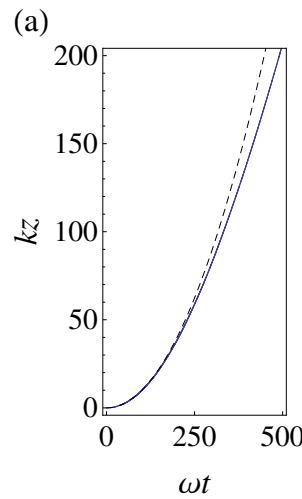
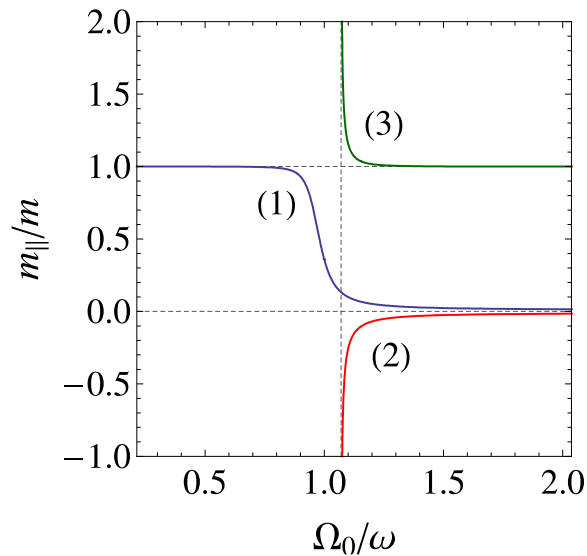


$$m_{\parallel}(\mathbf{E}_w, \mathbf{B}_0) \frac{d^2 \bar{z}}{dt^2} = F_{\parallel}$$

$$m_{\parallel} = \frac{\partial}{\partial \bar{v}} \left( \bar{\gamma} m_{\text{eff}} \bar{v} - \frac{c^2}{\bar{\gamma}} \frac{\partial m_{\text{eff}}}{\partial \bar{v}} \right)$$

*Dodin and Fisch, Phys. Rev. E, 2008*

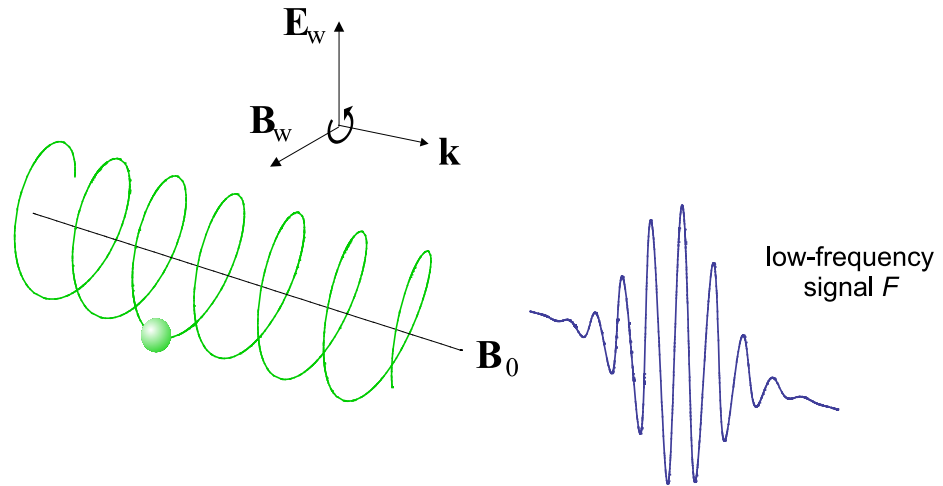
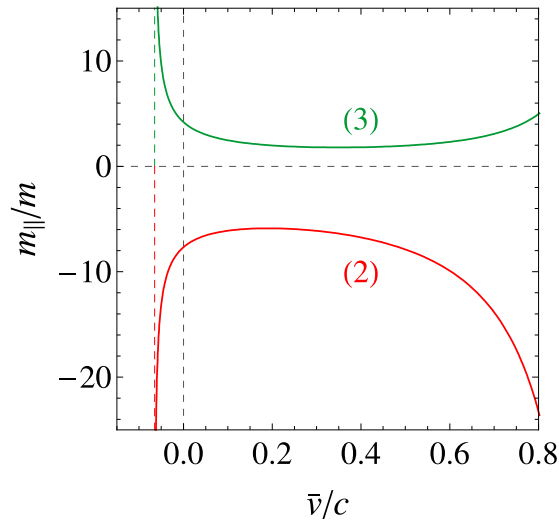
- Three different masses:



(a)  $m_{\parallel 1} > 0$ ,      (b)  $m_{\parallel 2} < 0$ ,      (c)  $m_{\parallel 3} > 0$ .  
 Numerical (solid) vs  $z(t) = F_{\parallel} t^2 / 2m_{\parallel}$  (dashed).



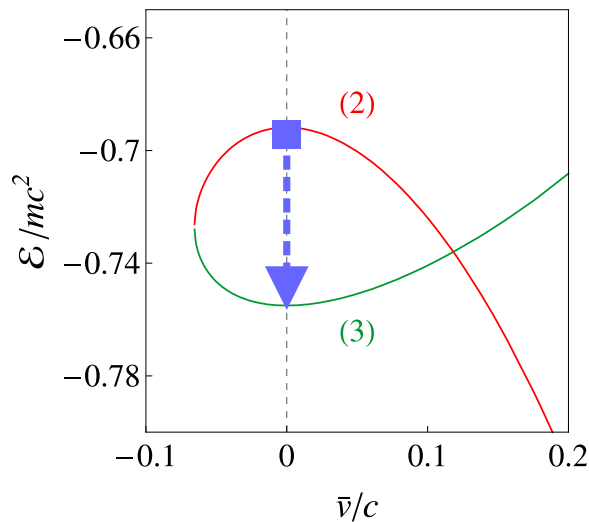
# Anisotropic Dispersion. Negative-Mass Amplifier



- Particles cannot absorb energy at branch 2; instead, their own energy goes to the pulse.

$$\Delta\mathcal{E} = (m_{\text{eff},2} - m_{\text{eff},3}) c^2 > 0$$

- Lasing mechanism using  $m_{\parallel} < 0$

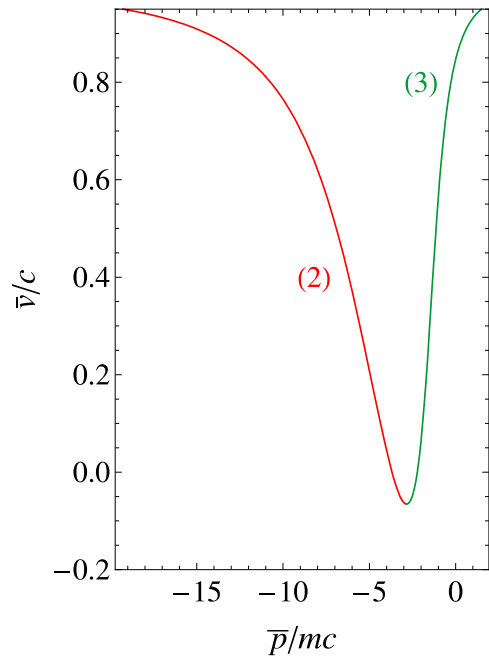


*Krömer, Phys. Rev., 1958; Andronov et al., Pis'ma Zh. Eksp. Teor. Fiz., 1984*



# Hamiltonian Ratchet. Frequency doubling

- Dispersion law:

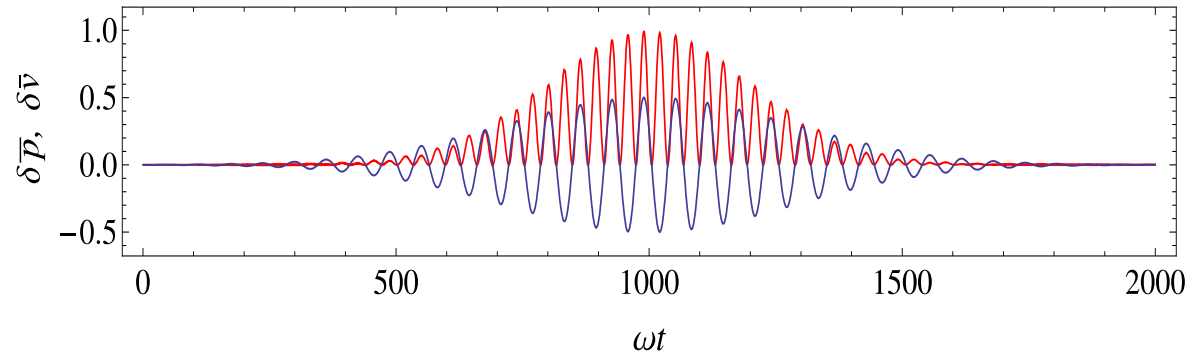
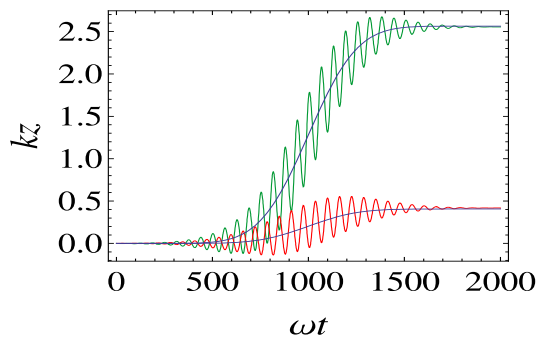


$$\bar{p} = \bar{\gamma} m_{\text{eff}} \bar{v} - (c^2/\bar{\gamma}) \partial_{\bar{v}} m_{\text{eff}}$$

$$\dot{\bar{p}} = \mathcal{F} \sin \omega t \quad \Rightarrow \quad \bar{p} = -(\mathcal{F}/\omega) \cos \omega t$$

$$\bar{v} \approx \bar{v}_0 + \Delta \bar{p} \partial_{\bar{p}} \bar{v} + (\Delta \bar{p}^2/2) \partial_{\bar{p}\bar{p}}^2 \bar{v}$$

$$\bar{v} \approx \bar{v}_0 - \underbrace{\frac{\mathcal{F}^2}{4\omega^2 m_{\parallel}} \frac{\partial m_{\parallel}^{-1}}{\partial \bar{v}}}_{\bar{v}_{\text{drift}}} + \underbrace{v_1 \cos \omega t + v_2 \cos 2\omega t}_{\bar{v}_{\sim}}$$







# Nonrelativistic Effective Potential

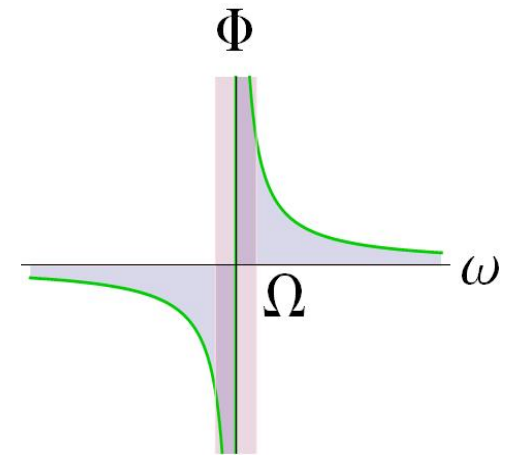
- Nonrelativistic energies:
  - For small  $\delta m$ ,  $\Psi$  can yet be a significant effective potential

$$\mathcal{H} \approx \frac{\bar{p}^2}{2m} + \Psi, \quad \Psi = \delta m c^2$$

$$\delta m \equiv m_{\text{eff}} - m \ll m$$

- Natural oscillators:
  - Atoms and molecules – quantum oscillators
  - Clusters (ion core + electron cloud),  $\Omega = \omega_p / \sqrt{3}$
  - Charged particles in a magnetic field,  $\Omega = eB / mc$

$$\Psi \approx \Phi + J\Omega, \quad \Phi \approx \frac{\kappa_1 |E|^2}{\omega - \Omega}$$

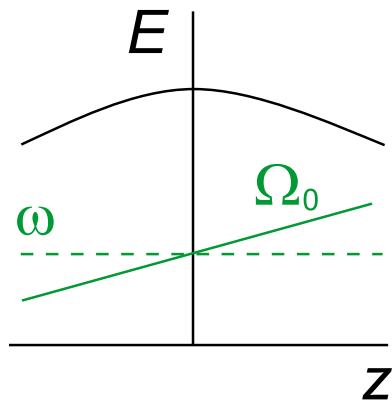


*Dodin and Fisch, Phys. Lett. A, 2006*

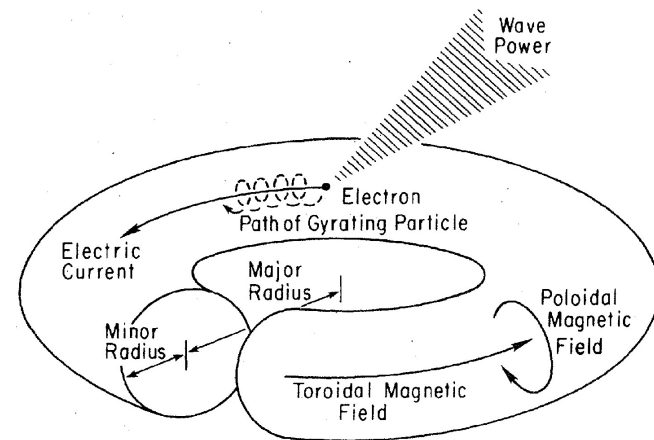
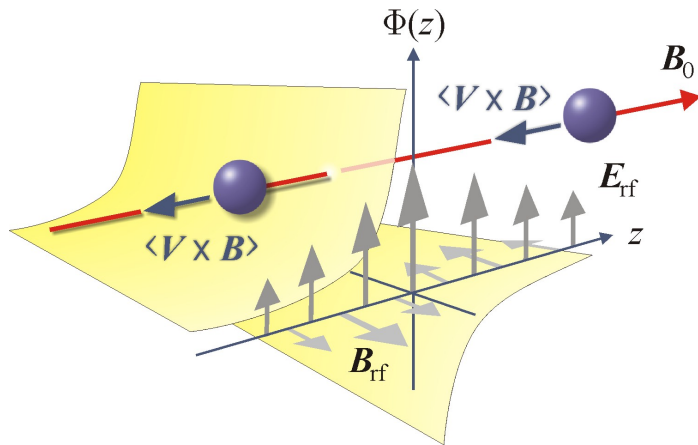
$\kappa_1 > 0$ ;  $J$  is the action,  $\Omega$  is the natural oscillation frequency, e.g.,  $J\Omega = \mu B$



# Singular Asymmetric Ponderomotive Potential



- Asymmetric potential  $\Phi \propto \frac{|E|^2}{\omega - \Omega(z)}$ :
  - Repulsive at  $\Omega(z) < \omega$
  - Attractive at  $\Omega(z) > \omega$
- “Maxwell demon” effect
- Possible current drive mechanism



Fisch et al., *Phys. Rev. Lett.*, 2003; Dodin et al., *Phys. Plasmas*, 2004; Suvorov and Tokman, *Fiz. Plazmy*, 1988



# One-way Walls for Nonlinear Oscillators. Translational Cooling

- Energy-dependent ponderomotive potentials:

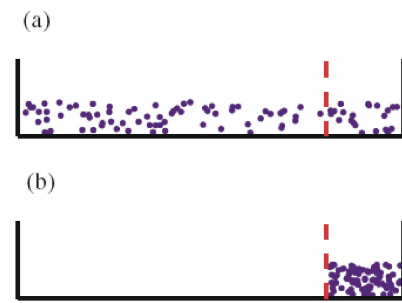
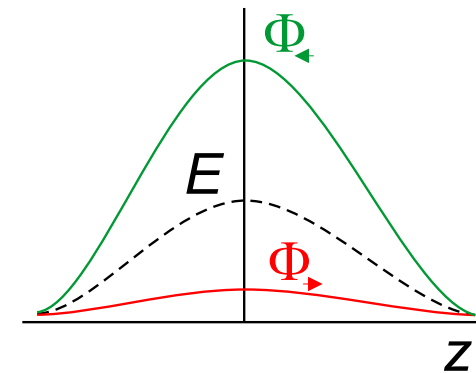
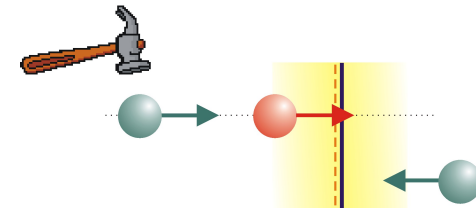
- Nonlinear frequency, at  $\omega \approx \Omega$ :

$$\Psi = \frac{\kappa_1 |E|^2}{\omega - \Omega(J)}$$

- Coupled oscillations, at  $\omega \approx \Omega_1 - \Omega_2$ :

$$\Psi = \kappa_2 |E|^2 \frac{J_1 - J_2}{\Omega_1 - \Omega_2 - \omega}$$

- Possible to manipulate particles with fixed  $\Omega$
- Translational cooling is possible

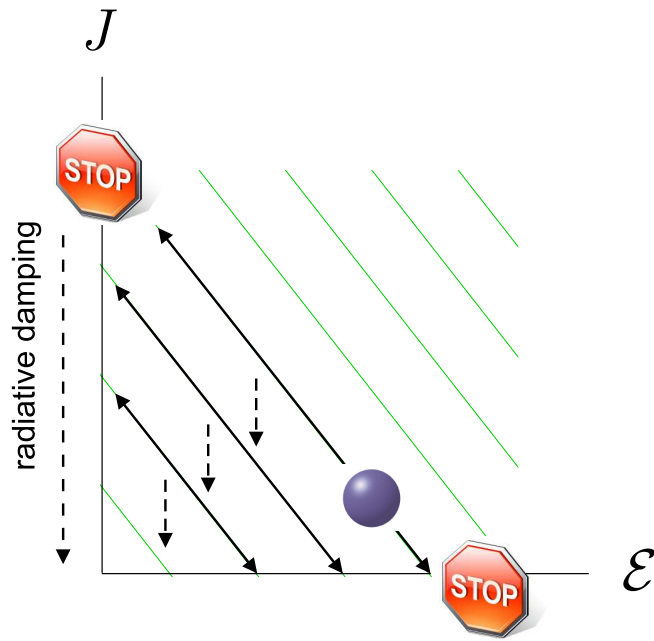


*Dodin and Fisch, Phys. Lett. A, 2006; Raizen et al., Phys. Rev. Lett., 2005; Ruschhaupt and Muga, Phys. Rev. A, 2004; Price et al., Phys.*

*Rev. Lett., 2008; Thorn et al., Phys. Rev. Lett., 2008*



# Sisyphus-like Translational Cooling



- Translational motion is coupled with internal oscillations:

$$\mathcal{E} + (\Omega - \omega)J = \text{const}, \quad \Omega > \omega$$

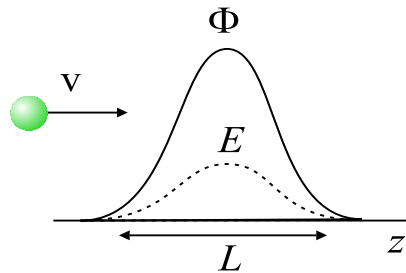
- Radiative damping reduces the free oscillation energy; thus the oscillation action  $J$  is decreased.
- “Dressed” particle energy  $\mathcal{E} \equiv mv^2/2 + \Phi$  is reduced accordingly.

Assumptions: no collisions, negligible photon recoil

*Dodin and Fisch, Phys. Lett. A, 2006; Dodin and Fisch, Phys. Plasmas, 2007*



# Quantum Analogy



$$E_{\sim}(z, t) = E(z) \cos \omega t$$

$$\mathbf{d} = \hat{\alpha} \mathbf{E}, \quad \text{if } \epsilon \ll 1$$

$$\epsilon = \lambda_{\text{eff}}/L, \quad \lambda_{\text{eff}} = \bar{v}T$$

$\lambda_{\text{eff}}$  is the average displacement  
on the oscillation period  $T$

$$\bar{z} = z - \epsilon^0 z_{\sim}^{(0)}(z, t) - \epsilon^1 z_{\sim}^{(1)}(z, t) + \dots$$

$$m\ddot{\bar{z}} = -\Phi'(\bar{z})$$

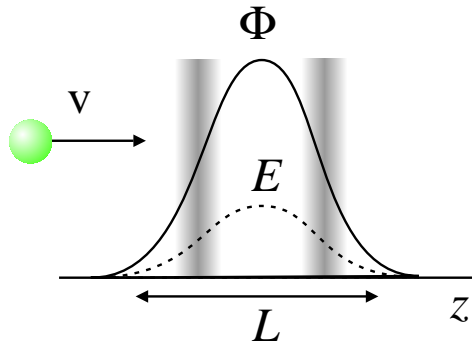
$$\bar{z} = z - \epsilon^0 z_{\sim}^{(0)}(z, t) + \delta\bar{z}, \quad \delta\bar{z} \sim \lambda_{\text{eff}}$$

<i>Adiabatic interaction</i>	$\lambda_{\text{eff}} \ll L$	<i>"Quasi-classical" interaction</i>
<i>Nonadiabatic interaction</i>	$\lambda_{\text{eff}} \gtrsim L$	<i>Quantum-like interaction</i>
<i>Location uncertainty</i>	$\lambda_{\text{eff}} = \bar{v}T$	<i>Effective de Broglie wavelength</i>
<i>Guiding center</i>	$\leftrightarrow$	<i>Quantum object</i>

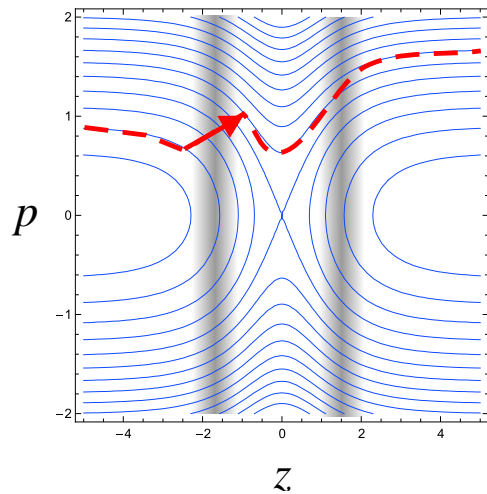
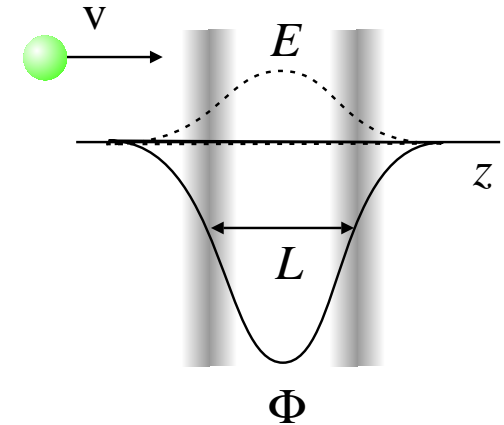
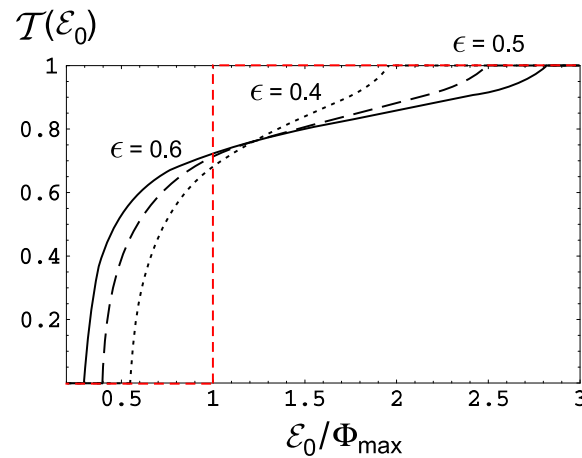
*Dodin and Fisch, Phys. Rev. Lett., 2005*



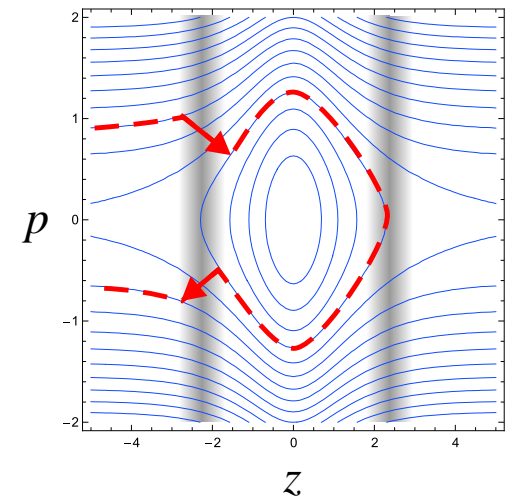
# Quantum-like Tunneling and Reflection



$$\epsilon \sim \bar{v} / \omega L(z), \quad L = E / |E'|$$



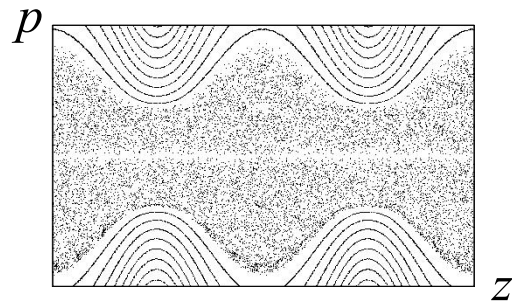
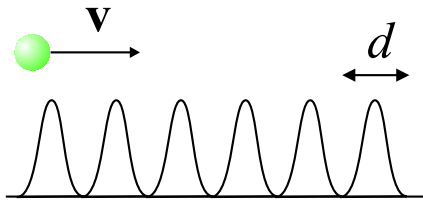
$$\mathcal{T} \approx \frac{1}{2} - \frac{1}{\pi} \arcsin \frac{\Phi_{\max} - \mathcal{E}_0}{\delta \mathcal{E}(\epsilon)}$$



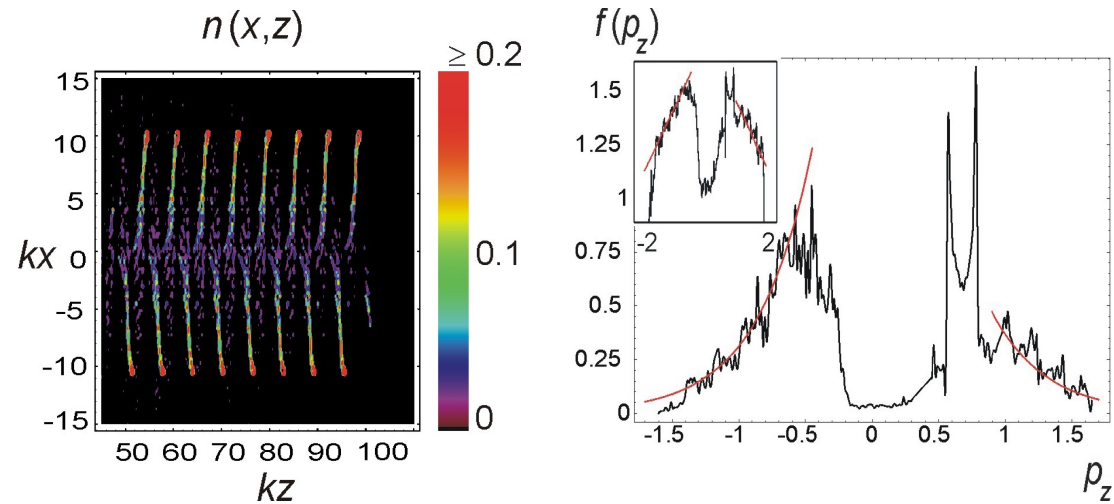
*Dodin and Fisch, Phys. Rev. E, 2006*



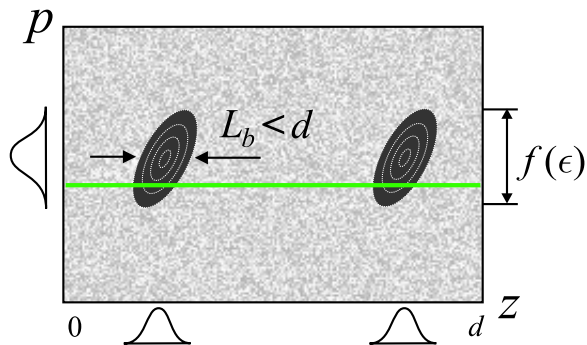
# Particle Beam Slicing



- Beam is sliced to bunches with  $L_b < d$



*Dodin and Fisch, Phys. Rev. Lett., 2007*

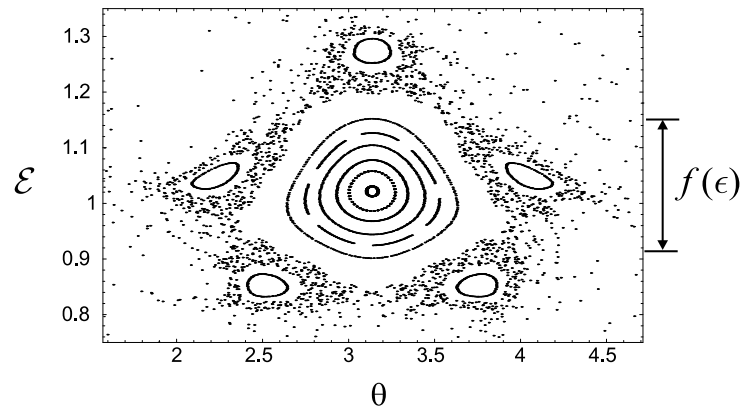
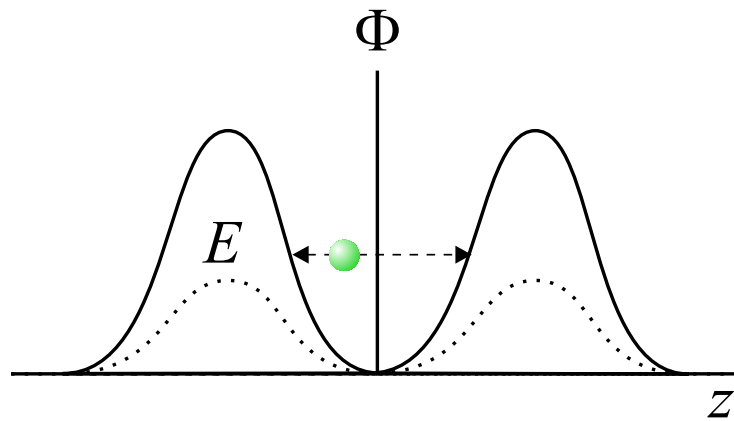


- Analogous technique is used to slice atomic beams using quantum accelerator modes.

*d'Arcy et al., Phys. Rev. A, 2003*



# Quantized Energy States



- Only periodic orbits survive at  $\epsilon \gtrsim 1$
- Guiding center phase:  $\Delta\psi = 2\pi n$
- Bohr-Sommerfeld rule:

$$\oint k dz = 2\pi n$$

$$k = 2\pi/\lambda_{\text{eff}}, \quad \lambda_{\text{eff}} = \bar{v}T$$

*Dodin and Fisch, Phys. Rev. Lett., 2005*





## Summary

---

- Guiding centers can be viewed as non-Newtonian quasi-particles with a rest mass  $m_{\text{eff}}$  different from the particle true mass.
- $m_{\text{eff}}$  depends on the field parameters and can be a nonanalytic function of the particle average velocity and the oscillation energy; hence the applications.
- Adiabatic effects:
  - negative-mass plasma maser,
  - signal rectification,
  - frequency doubling.
- Nonadiabatic and quantum-like effects:
  - one-way walls for current drive and translational cooling,
  - phase-dependent transmission (“tunneling”),
  - particle beam slicing,
  - selective confinement using quantized energy levels.