

The Abdus Salam International Centre for Theoretical Physics



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Introduction to Optofluidics

1 - 5 June 2009

Optical control of electroosmotic fluid flow

P. Ormos Hungarian Academy of Sciences Hungary Introduction to optofluidics

Trieste, 2009

Optical control of electro-osmotic flow

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Concept for optical control of fluidics

- Aim: microfluidics system with total optical control
- Full photonic drive of microfluidic functions (pumping, steering, etc.) may not be optimal (not enough power)

Control flow by light, but drive electrically

Electro-osmosis controlled by light

Other approaches to eletrooptical fluid manipulation electrowetting



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Sensors and Actuators A 104 (2003) 222-228



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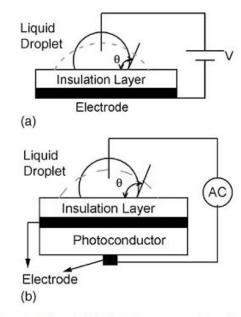
Light actuation of liquid by optoelectrowetting

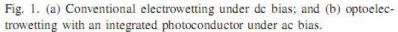
Pei Yu Chiou^{a,*}, Hyejin Moon^b, Hiroshi Toshiyoshi^c, Chang-Jin Kim^b, Ming C. Wu^a

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Other approaches to eletrooptical fluid manipulation

electrowetting

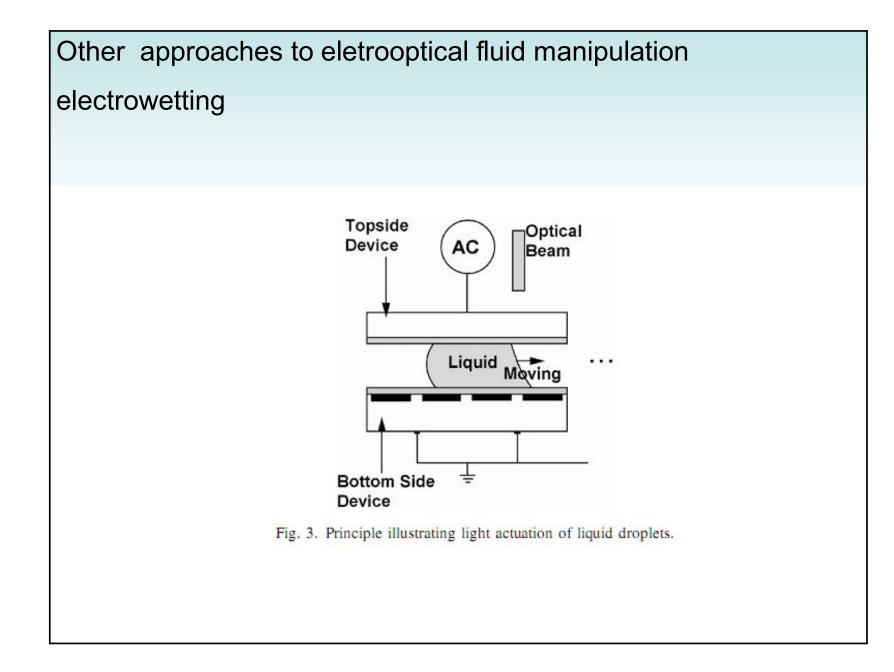


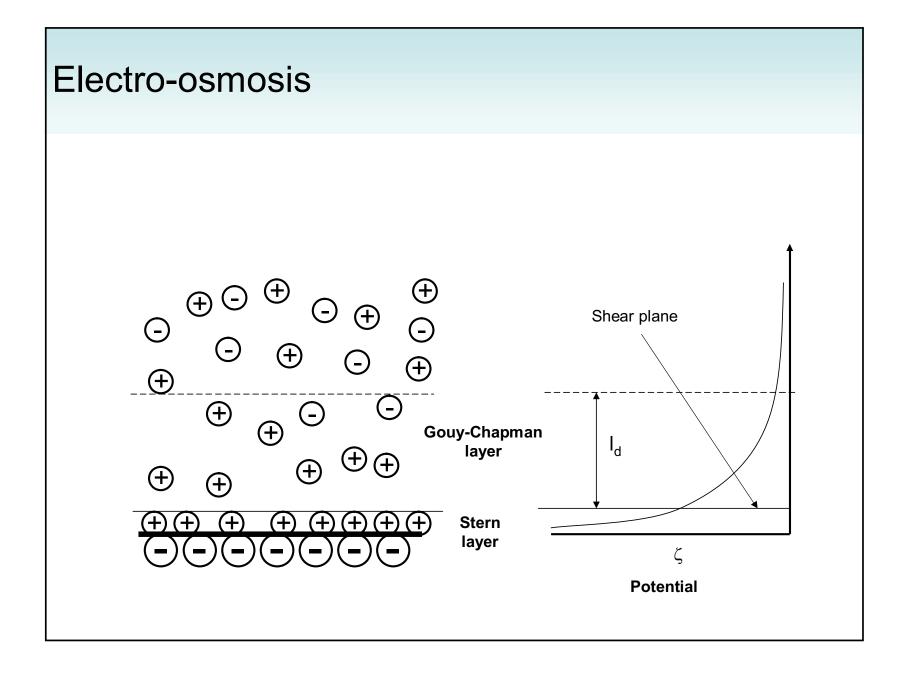


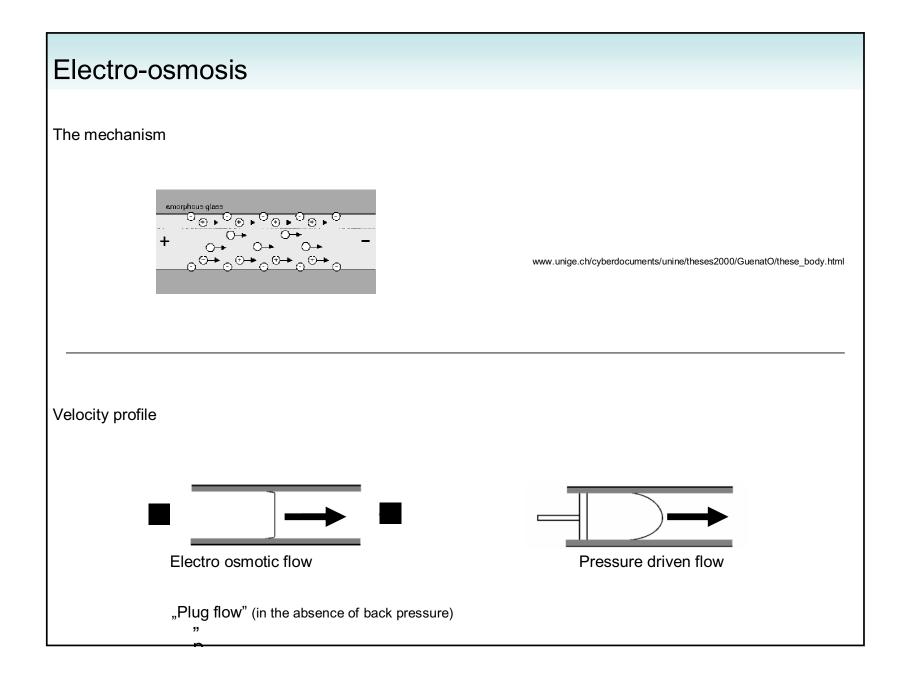
Other approaches to eletrooptical fluid manipulation electrowetting

The voltage dependence of the contact angle, $\theta(V),$

$$\cos[\theta(V_{\rm A})] = \cos[\theta(0)] + \frac{1}{2} \frac{\varepsilon}{d\gamma_{\rm LV}} V_{\rm A}^2$$







Boltzmann distribution i, c_i:
$$c_i(y) = c_{\infty,i} \exp(-\frac{ye\phi(y)}{kT})$$

 C_i : concentration of ith ion, ϕ : local potential, y: distance from wall,

$$\rho_E = F \sum_{i=1}^N z_i c_i$$

Z: valence of ion, F: Faraday number

The relation of the net charge density and the local potential (Poisson equation):

$$\nabla^2 \phi = \frac{-\rho_{\varepsilon}}{\varepsilon} \tag{1}$$

 ρ : net charge density, ϵ : permittivity

Combine these equations for the plane case:

$$\frac{d^2\phi}{dy^2} = \frac{-F}{\varepsilon} \sum_{i=1}^{N} z_i c_{\infty,i} \exp(-\frac{ze\phi(y)}{kT})$$

With symmetric monovalent electrolyte this becomes:

$$\frac{d^2\phi}{dy^2} = \frac{2Fz_i c_\infty}{\varepsilon} \sinh(\frac{ze\phi(y)}{kT})$$

Poisson-Boltzmann equation

If kT>>ze ϕ (Debye-Hückel limit)

$$\frac{d^2\phi}{dy^2} = \frac{\phi(y)}{\lambda_D^2} \quad \text{where} \quad \lambda_D = (\frac{\varepsilon kT}{2z^2 F^2 c_\infty}) \quad \text{: Debye length}$$

For 10 mM concentration λ ~ few nanometers

Г

Electro osmotic flow

Applied electric field introduces a Lorentz body force: $\rho b = \rho_E E$

The equation of motion for steady flow in the channel: $\nabla p = \mu \nabla^2 u + \rho_E E$

Substitute the Poisson equation (1):
$$\nabla^2 (u - \frac{\varepsilon E}{\mu}\phi) = \frac{\nabla p}{\mu}$$

This equation is linear, so electric field and pressure driven flow can be separated

$$\nabla^{2}(u_{EOF} - \frac{\varepsilon E}{\mu}\phi) = 0$$

Continuity equation: $\frac{\partial u}{\partial x} = 0$
$$\nabla^{2}u_{pressure} = \frac{\nabla p}{\mu}$$

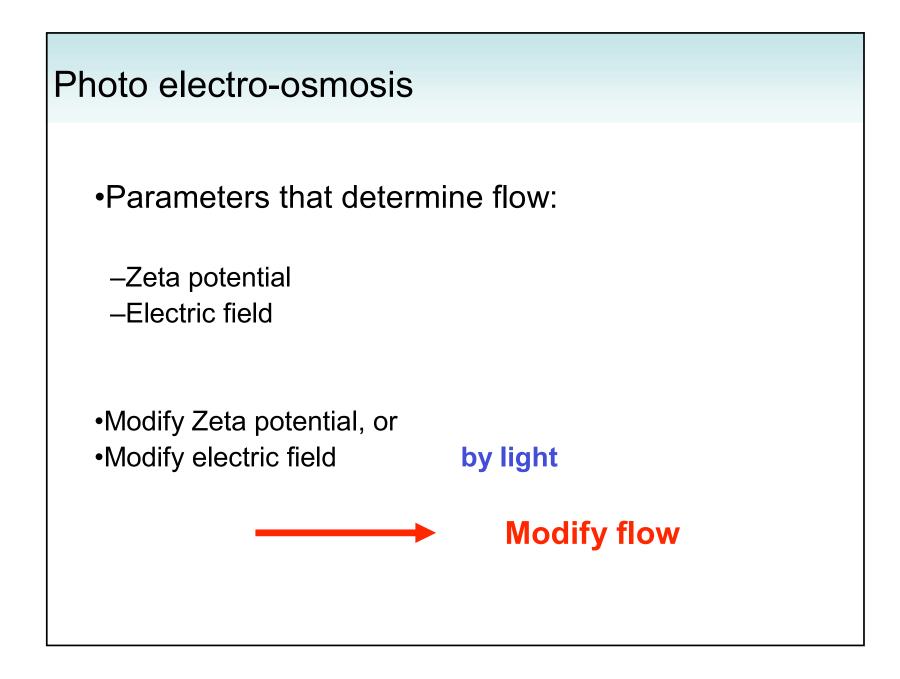
Electro osmotic flow

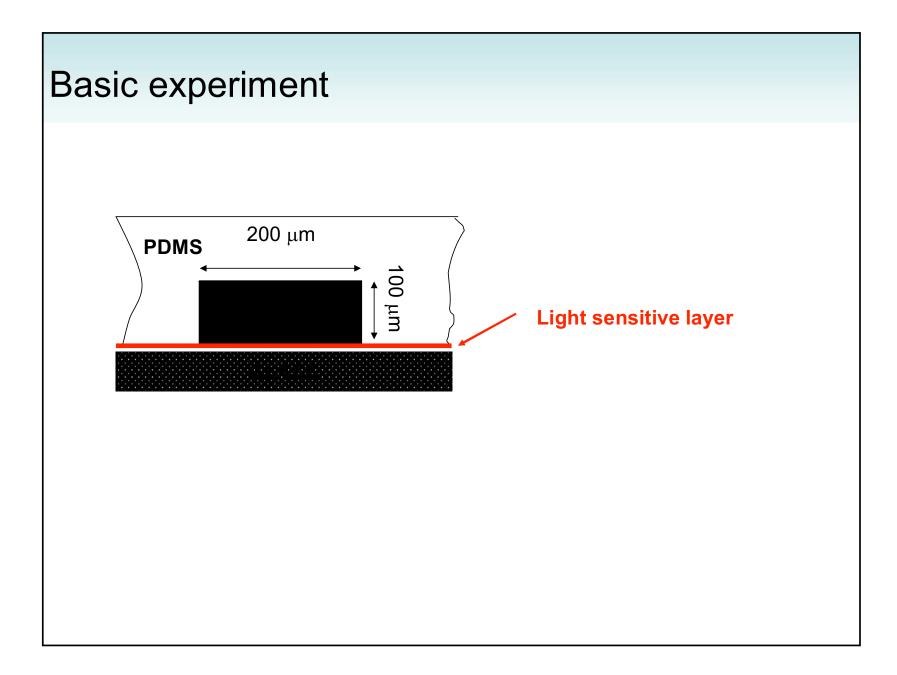
Reasonable case: long, straight channel with uniform electric field

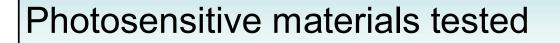
$$u_{EOF} - \frac{\varepsilon E}{\mu}\phi = \frac{-\varepsilon E\zeta}{\mu}$$

or

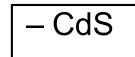
$$u_{EOF}(y,z) = \frac{-\varepsilon E\zeta}{\mu} (1 - \frac{\phi(y,z)}{\zeta})$$



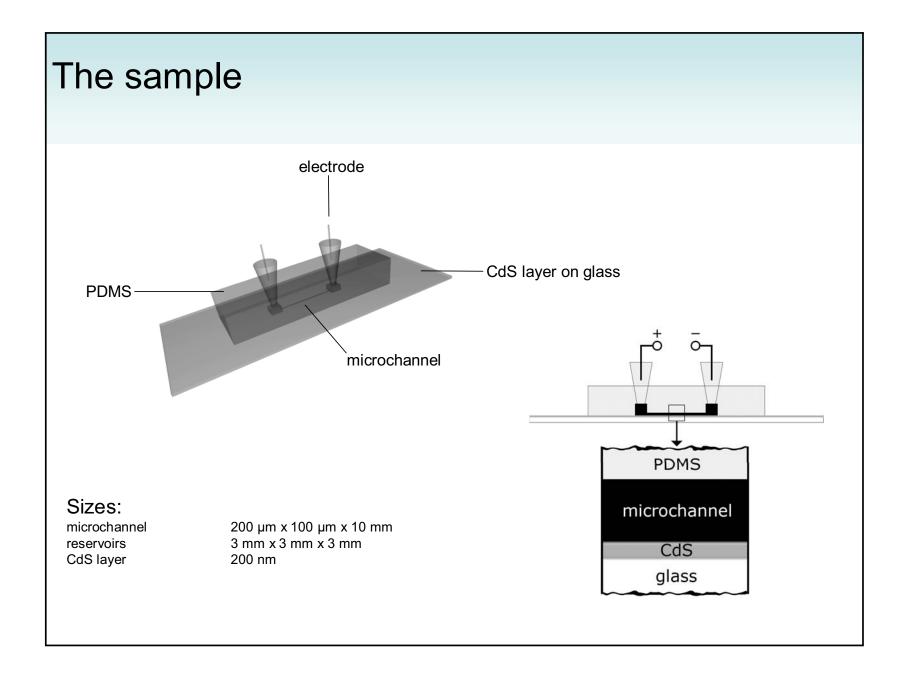




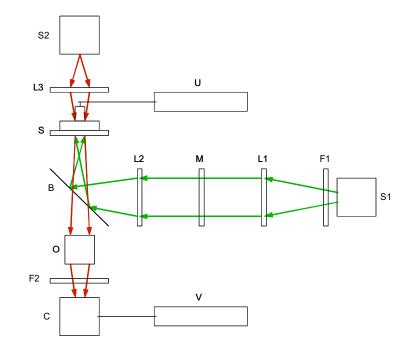
- Change of surface charge
 - $-\operatorname{TiO}_2$
 - Bacteriorhodopsin
 - BSO
- Photoresistor
 - Amorphous Si



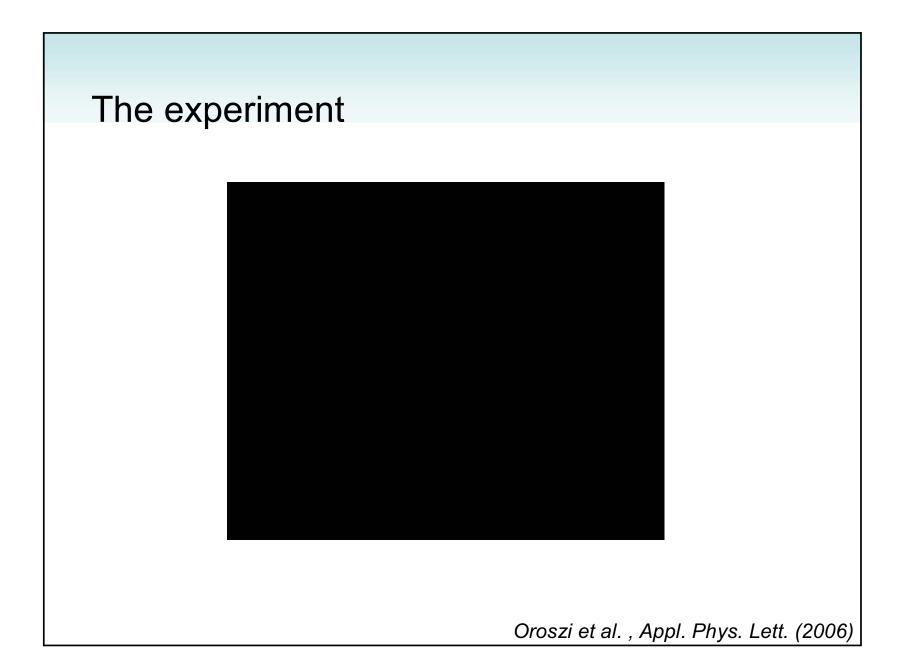
Experiments with CdS are shown



Experimental layout

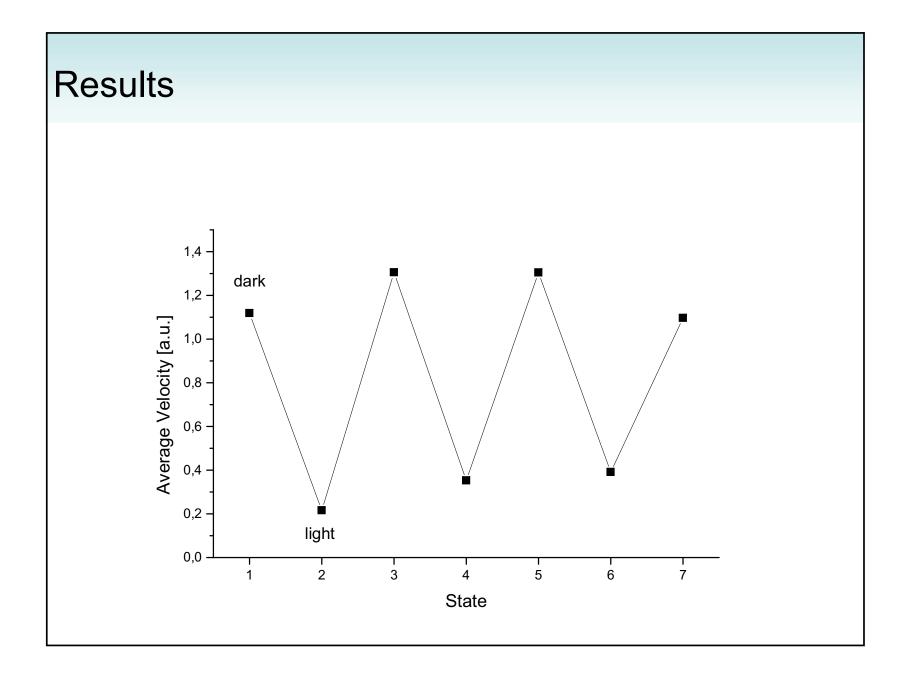


S: microfluidic sample, S1: mercury lamp, S2: LED, F1,F2: optical filters, L1,L2,L3: lenses, M: photomask, B: beam-splitter, O: objective, C: camera, V: video recorder, U: voltage source

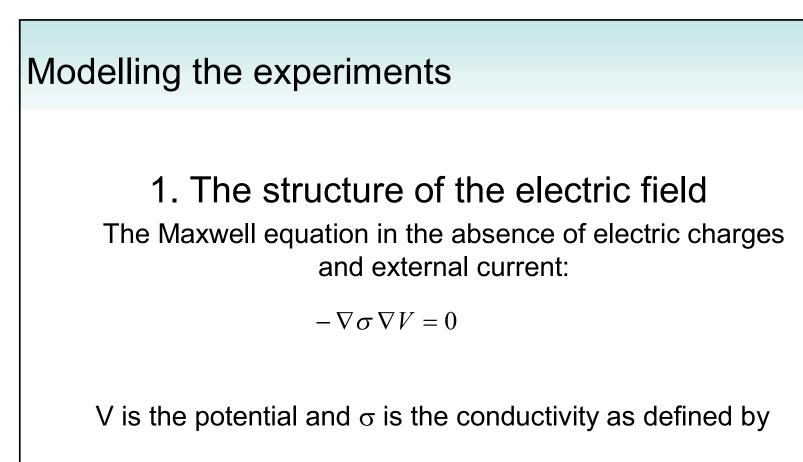


Characterize fluid flow





- Simulation is done with FEMLAB (Comsol)
- Fluid is treated with the Navier-Stokes equation
- The electric field is calculated by electrostatics of conductive materials
- The effect of light is the decrease of the ohmic resistance upon illumination of one surface
- Electro-osmotic driving is treated with the Helmholz-Smoluchowski relation



$$\vec{J}=\sigma\vec{E}=-\sigma\,\nabla V$$

1. The structure of the electric field

The channel has three insulator and one conductive walls.

The conductivity values used were determined from experiments (for the conductive wall in dark and light)

2. Modelling fluid flow

The Navier-Stokes equations for our stationary case:

 $-\eta \nabla^2 \, \vec{v} + \rho \, \vec{v} \nabla \, \vec{v} + \nabla \, p = \vec{F}$ $\nabla \vec{v} = 0$

where v is the velocity field of the fluid, ρ is the density, η is the dynamic viscosity, p is the pressure, and F is a volume force field

3. Modelling osmotic flow

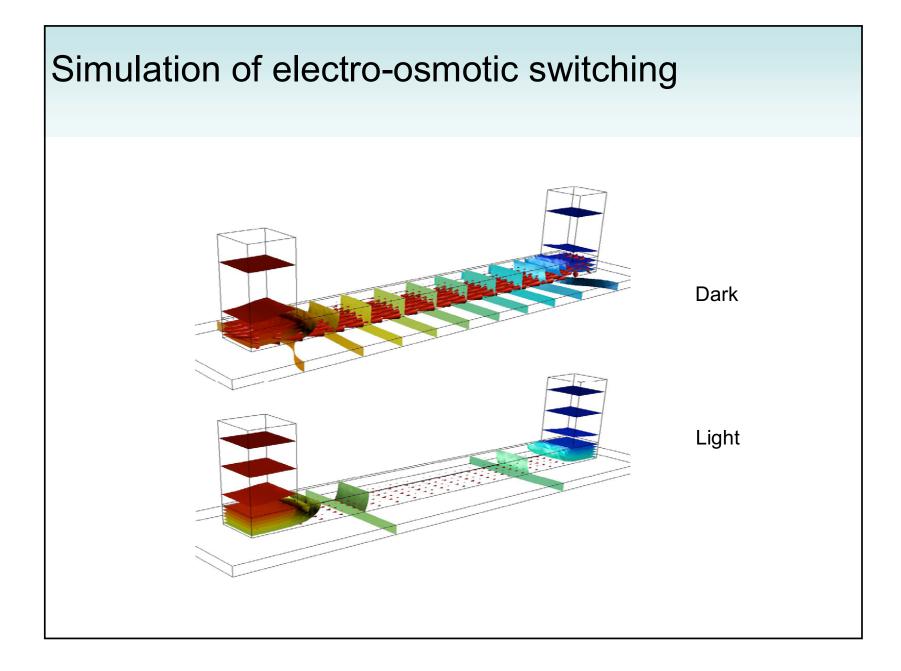
At the wall the velocity components perpendicular to the wall are set to zero, the two in-plane components are set according to the Helmholtz-Smoluchowski equation:

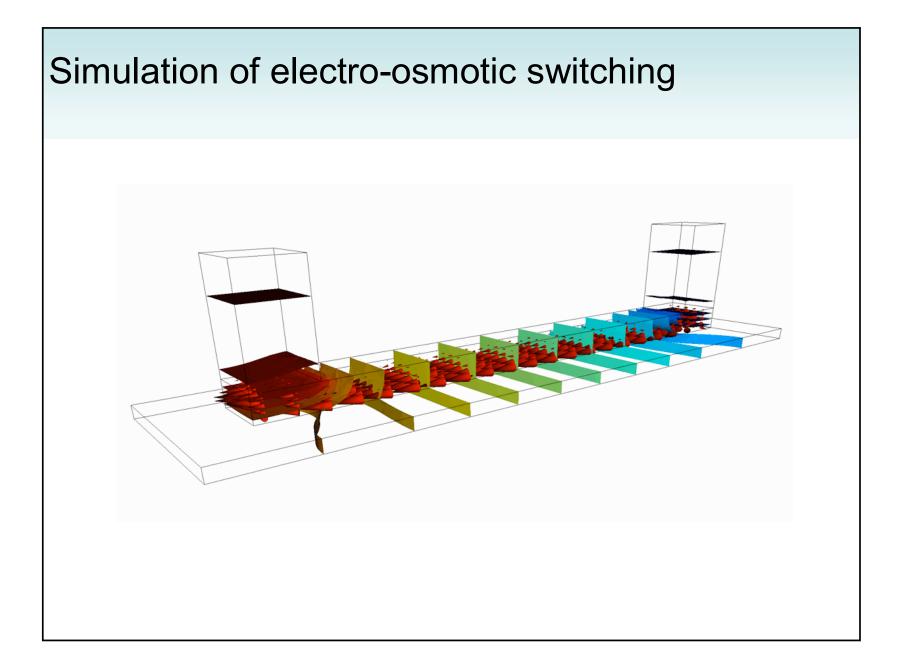
$$\overline{v} = \frac{\varepsilon_w \zeta_0}{\eta} \overline{E}$$

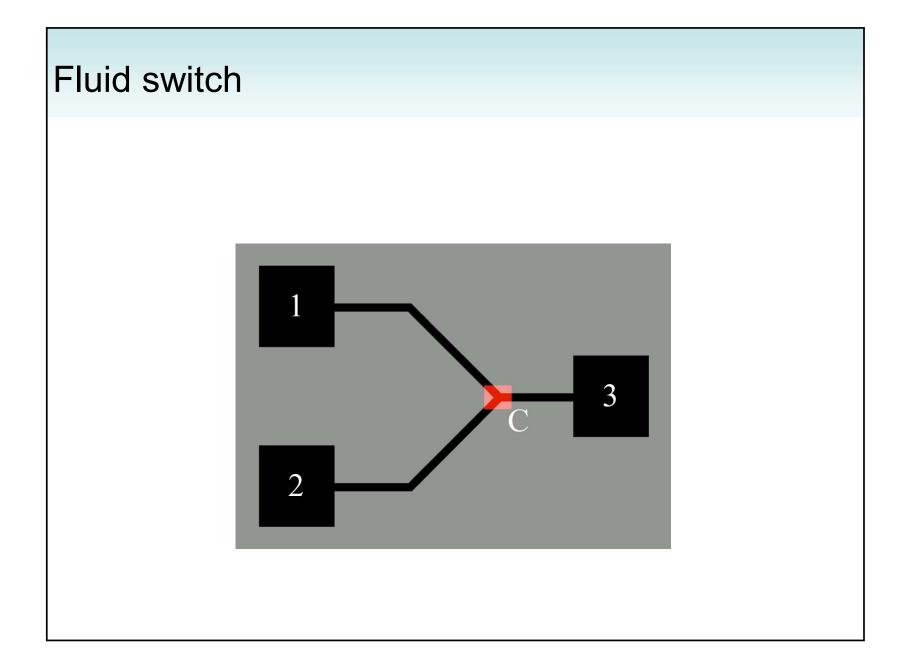
where E is the electric field, ζ is the zeta-potential at the wall and ϵ is the fluid's absolute permittivity

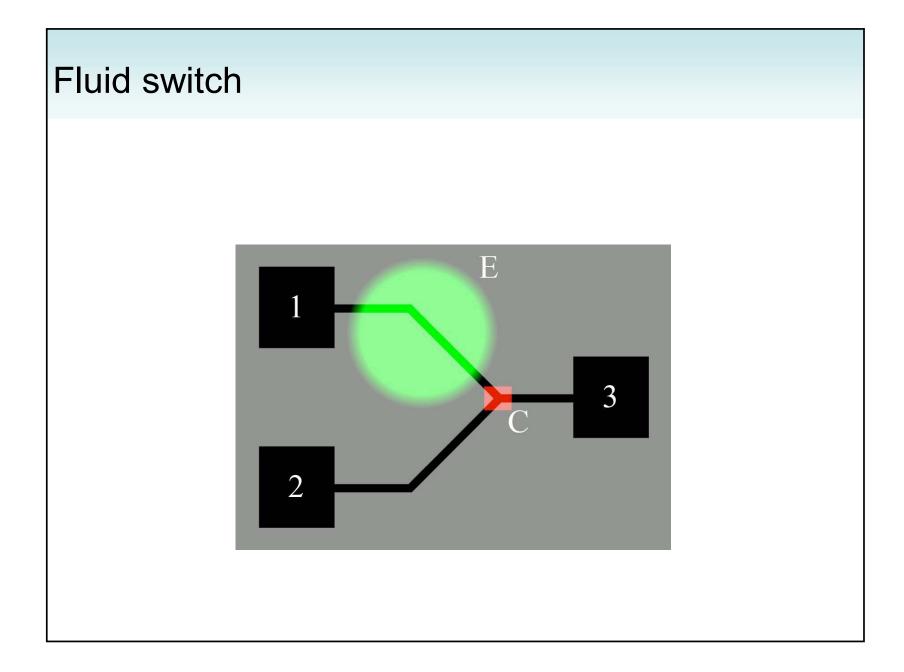
Realistic values were given to the Zeta potential:

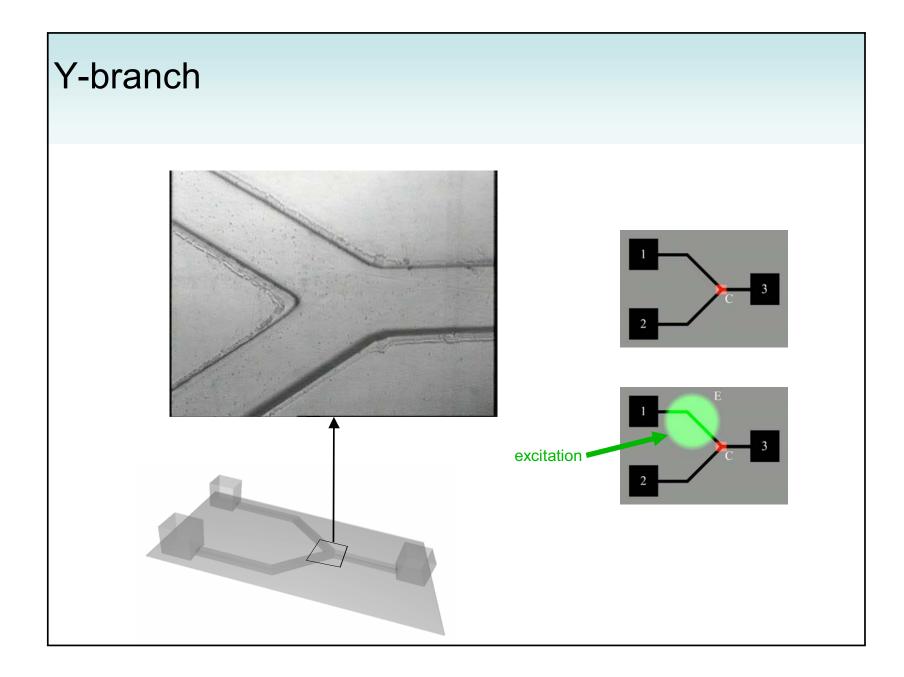
PDMS: -100 mV, CdS: -5 mV

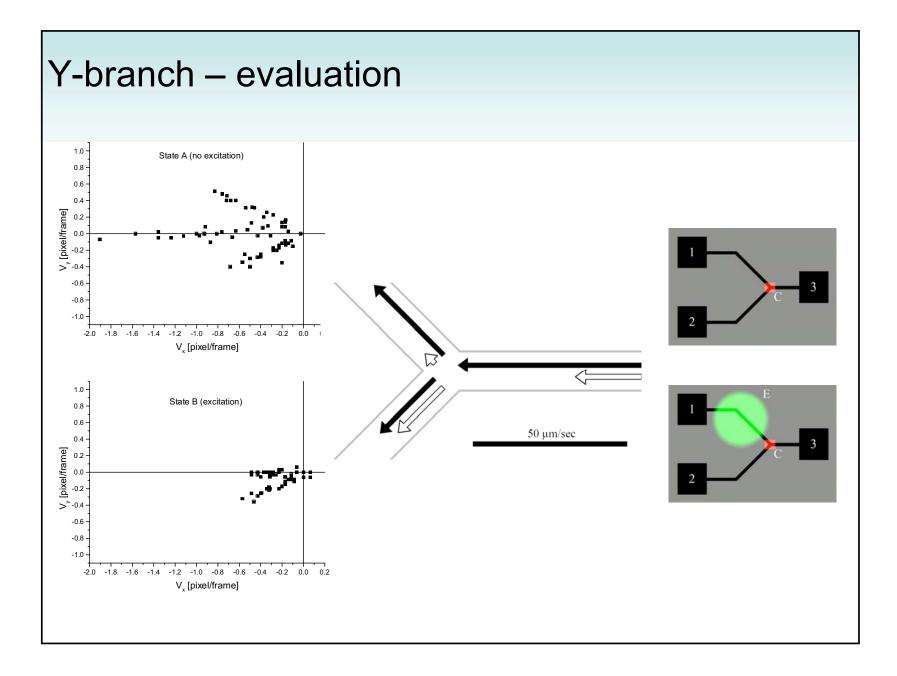


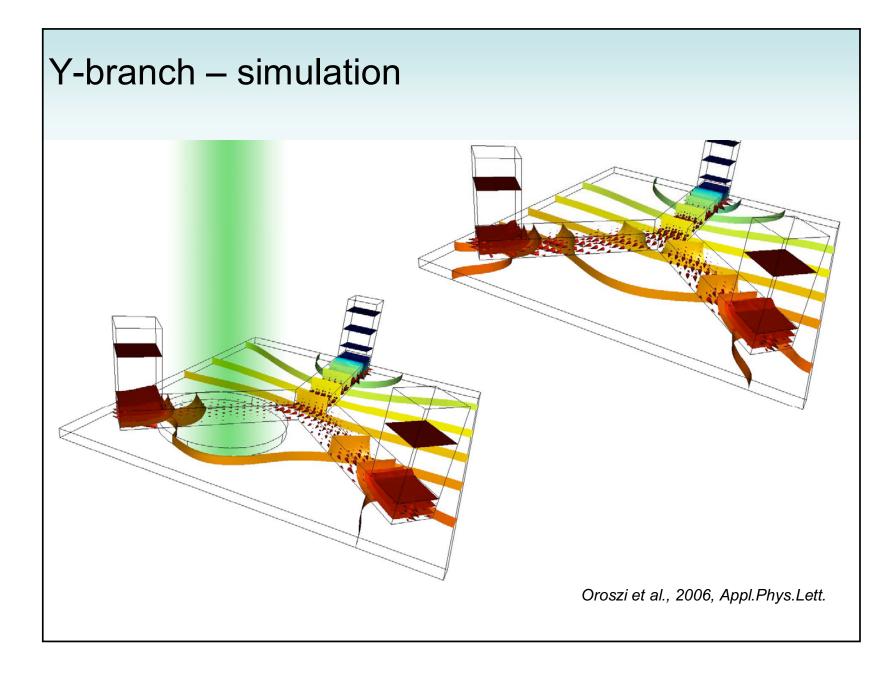












Manipulation of flow pattern in an electro-osmotic system

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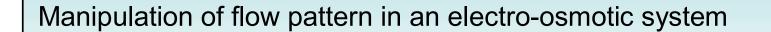
PHYSICAL REVIEW LETTERS

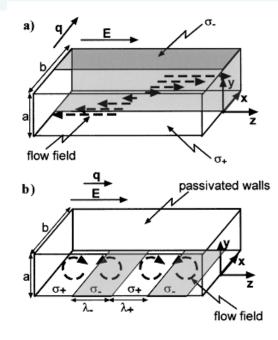
10 April 2000

Patterning Electro-osmotic Flow with Patterned Surface Charge

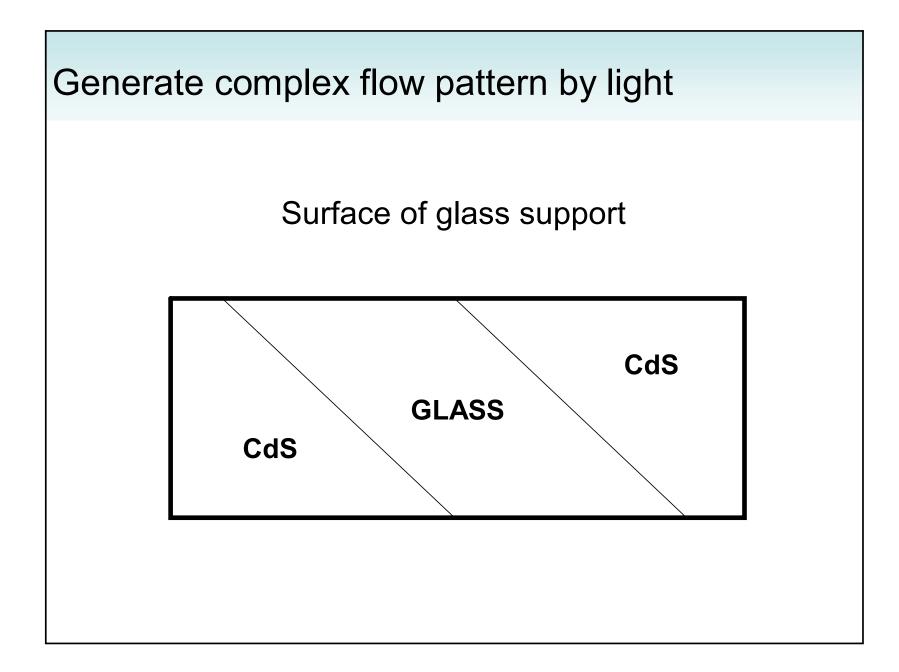
Abraham D. Stroock, Marcus Weck, Daniel T. Chiu, Wilhelm T. S. Huck, Paul J. A. Kenis, Rustem F. Ismagilov, and George M. Whitesides* Department of Chemistry and Chemical Biology, Harvard University, Cambridge, Massachusetts 02138 (Received 2 November 1999)

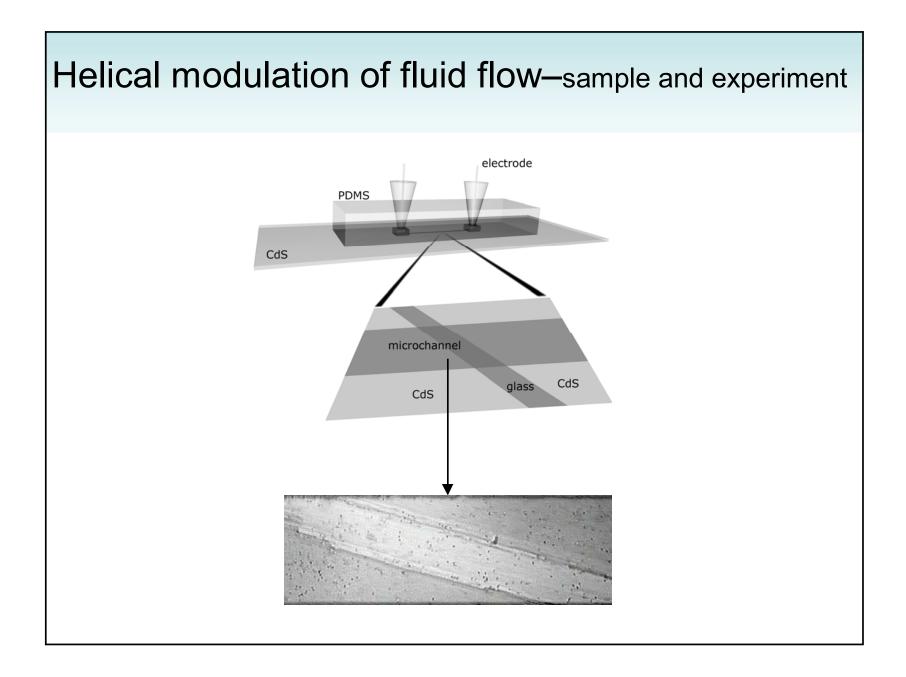
This Letter reports the measurement of electro-osmotic flows (EOF) in microchannels with surface charge patterned on the 200 μ m scale. We have investigated two classes of patterns: (1) Those in which the surface charge varies along a direction perpendicular to the electric field used to drive the EOF; this type of pattern generates multidirectional flow along the direction of the field. (2) Those in which the surface charge pattern varies parallel to the field; this pattern generates recirculating cellular flow, and thus causes motion both parallel and perpendicular to the external field. Measurements of both of these flows agree well with theory in the limit of thin double layers and low surface potential.

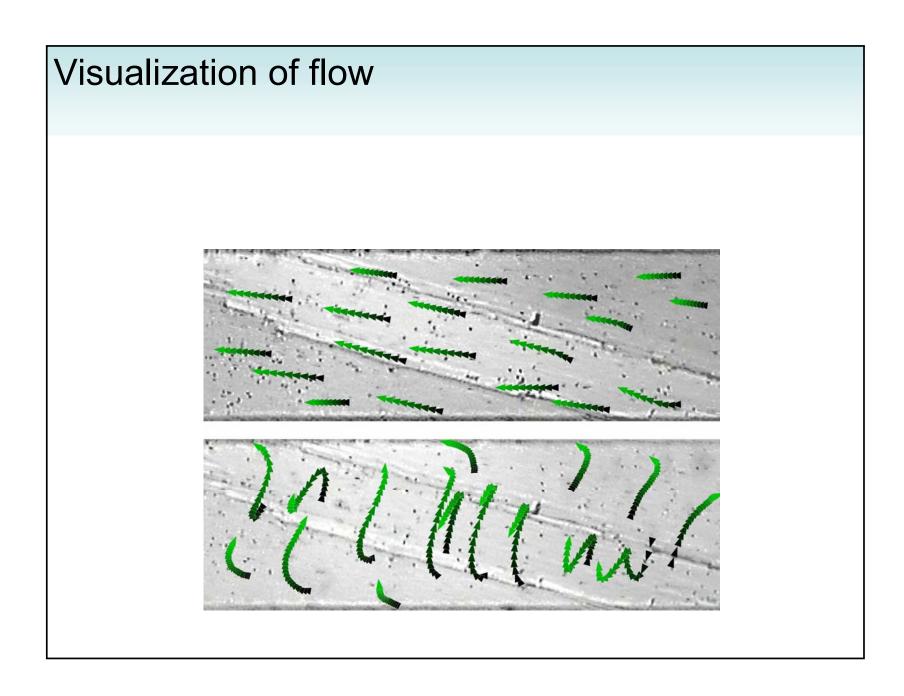


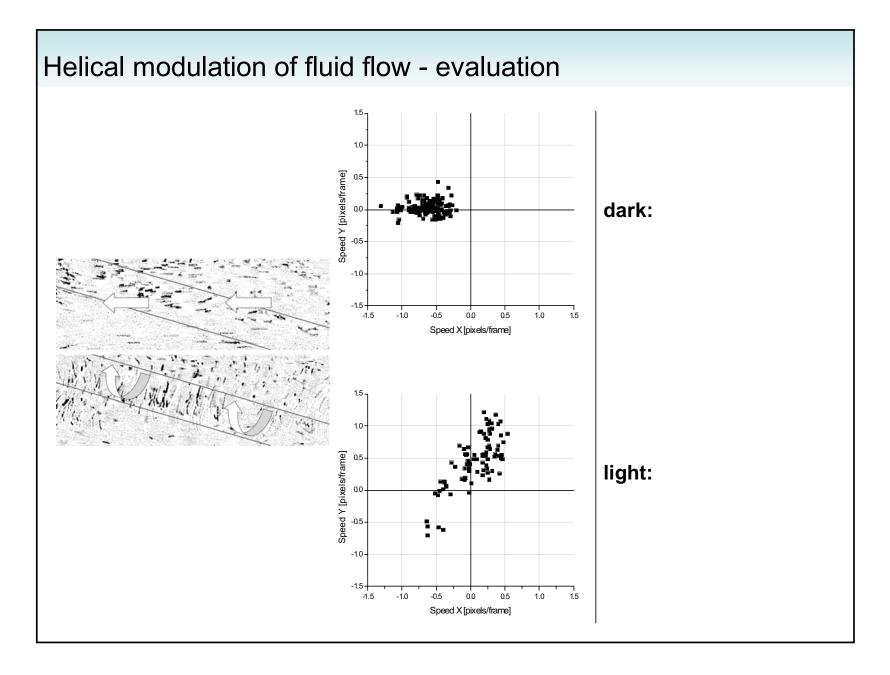


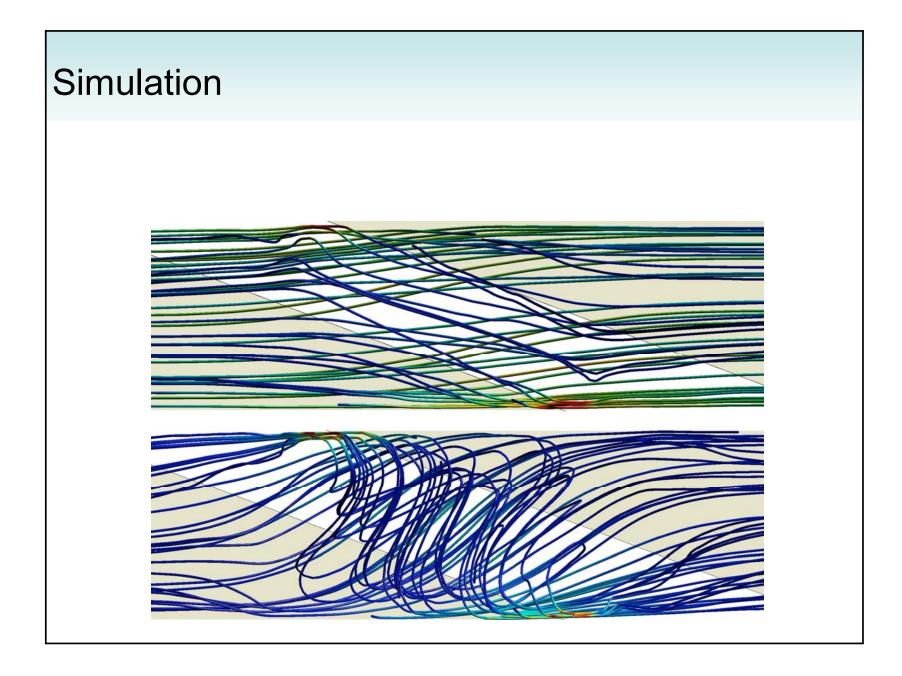
The driving field is patterned with patterned electrodes

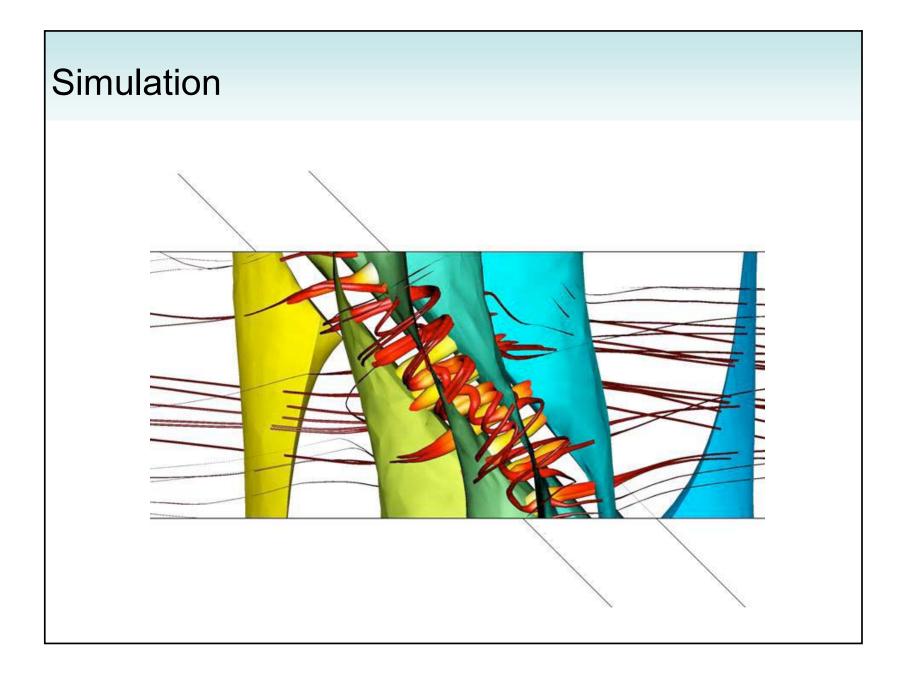


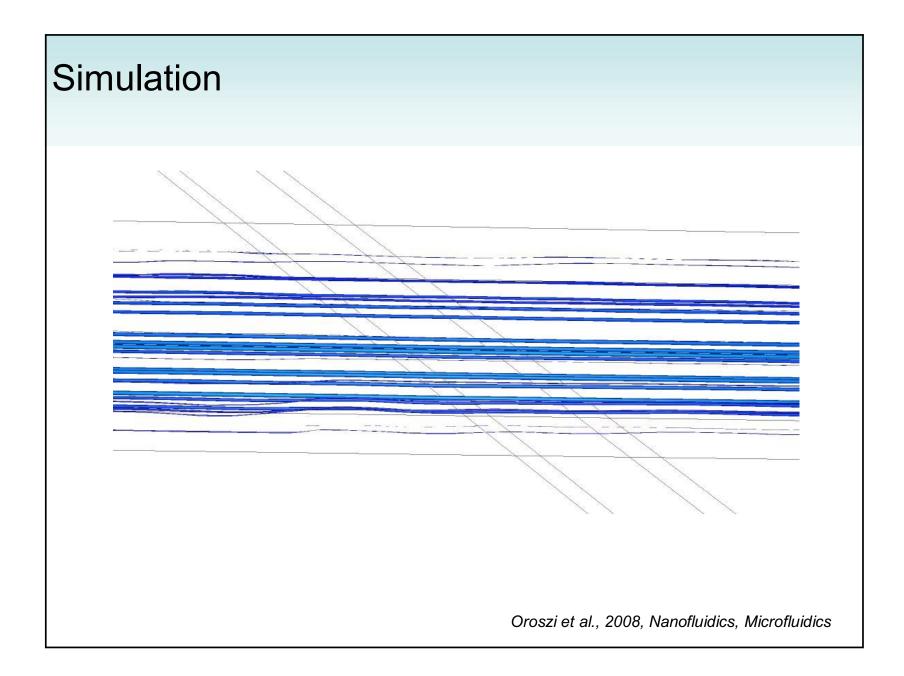




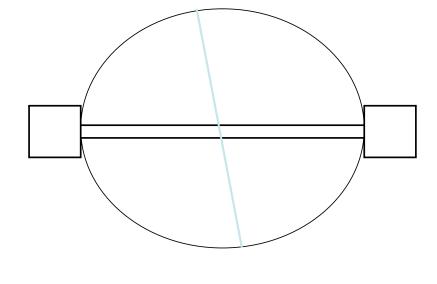




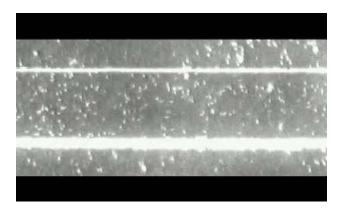


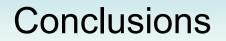


Change flow pattern by projecting image on the photoconductor



Change flow pattern by projecting image on the photoconductor





 Optoelectroosmosis has promise in manipulationg local flows (mixing in reactors, etc.). Other approaches to eletrooptical fluidic manipulation

2. Electro-optical trapping

Electrophoretic assembly of colloidal crystals with optically tunable micropatterns

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Other approaches to eletrooptical fluidic manipulation

2. Electro-optical trapping

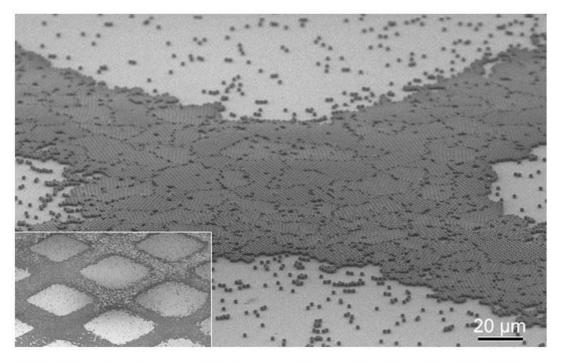
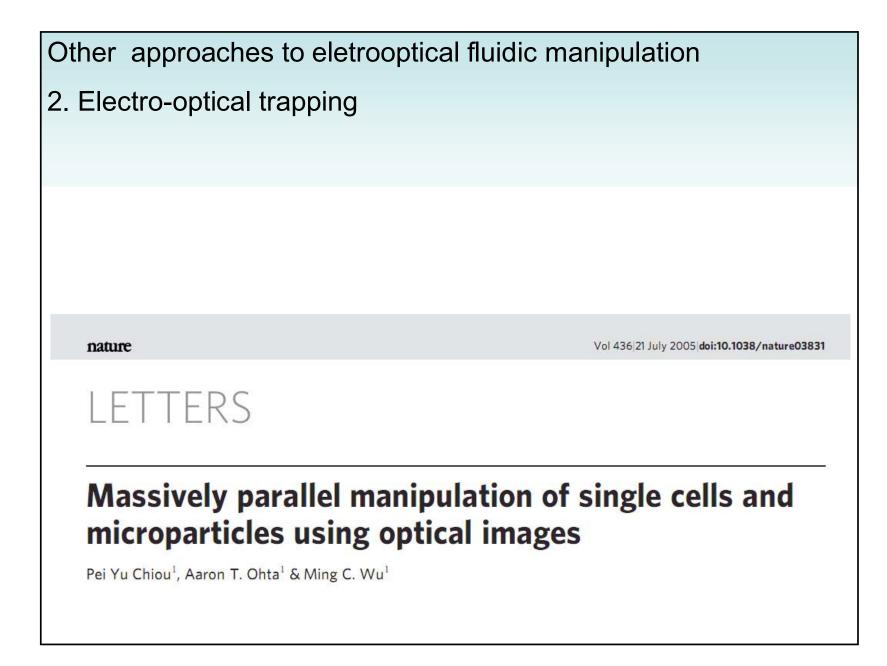


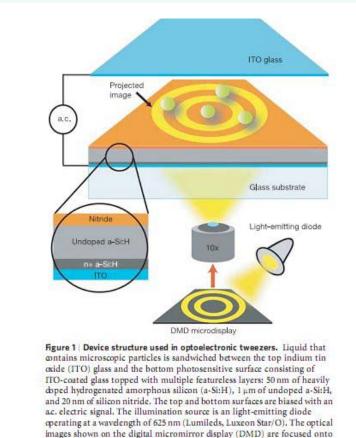
Figure 3 Scanning electron microscope image of a pattern produced by method B. Note that dense-packing of particles results in crystalline (ordered) domains that vary in size $(10-20 \ \mu m)$. The inset ($\sim 600 \ \mu m$ wide) shows the overall appearance of the pattern.

Some areas shown in the inset appear white due to particle charging in the SEM. Charging in these regions may result from the higher particle density which inhibits charge transfer from the uppermost particles to the conducting substrate.



Other approaches to eletrooptical fluidic manipulation

2. Electro-optical trapping



the photosensitive surface and create the non-uniform electric field for DEP

manipulation