



**The Abdus Salam  
International Centre for Theoretical Physics**



**2052-63**

**Summer College on Plasma Physics**

*10 - 28 August 2009*

**The Sachs-Wolfe Effect and Photon-Graviton interactions**

Jose Tito Mendonça  
*IPFN  
Portugal*



INSTITUTO  
SUPERIOR  
TÉCNICO

# The Sachs-Wolfe Effect and Photon-Graviton interactions

J. Tito Mendonça

IPFN and CFIF, Instituto Superior Técnico  
CfFP, Rutherford Appleton Laboratory

Collaborators:

R. Bingham (RAL), C. Wang (Un. Aberdeen), L. Drury (Dublin),  
G. Brodin, M. Marklund (U. Umea), P. Shukla (U. Bochum)



INSTITUTO  
SUPERIOR  
TÉCNICO

## Motivations

- The metric space-time is a physical field (it can exchange energy with other fields);
- Microwave photons of CMB are already used to probe the early universe;
- Detection of G waves is a major challenge;
- Could photons be used as probes of G waves and G fields?



INSTITUTO  
SUPERIOR  
TÉCNICO

# Outline

- **Photon ray equations;**
- **Refractive index of a gravitational field;**
- **Gravitational red shift, a new approach;**
- **The Sachs-Wolf effect;**
- **Photons in a gravitational wave;**
- **Gamma-ray bursts and G waves;**
- **Nonlinear photon-graviton coupling**



INSTITUTO  
SUPERIOR  
TÉCNICO

## Laser experiments with an ionization front

[Dias et al. PRL (1997)]

Laser pulse  
65 fs, 2.5 mJ @ 620 nm



<

> Counter-Propagation

Co-Propagation

<

>

Simultaneous measurement:  
 $\beta=0.942$ ,  $n=4.26 \times 10^{19} \text{ cm}^{-3}$

Can we find similar effects with gravitational fields?



## Photon ray equations

In a gravitational field we can always use geometric optics (scaling)

Electric field:

$$E(x^\nu) = A \exp i\psi(x^\nu)$$

Wavevector components

$$k_i = \partial\psi / \partial x^i \quad x^i = (ct, \vec{r})$$

Photon canonical equations

$$\frac{dx^\alpha}{dt} = \frac{\partial\omega}{\partial k_\alpha}$$

$$\frac{dk_\alpha}{dt} = -\frac{\partial\omega}{\partial x^\alpha}$$

$$\omega = -ck_0$$

( $\alpha=1,2,3$ )



## Variational principle

$$\delta\psi = \delta \int k_j dx^j = \delta \int (k_\alpha dx^\alpha - \omega dt) = 0$$

## Photon frequency shift

$$\frac{d\omega}{dt} = \frac{\partial\omega}{\partial t}$$

$$\omega = \omega(k_\alpha, x^\alpha, t)$$

Not to be confused with the photon proper frequency

$$\omega' = -\frac{\partial\psi}{\partial\tau} = -\frac{\partial\psi}{\partial x^0} \frac{\partial x^0}{\partial\tau} = \frac{\omega}{\sqrt{g_{00}}}$$

Proper time  $\tau = \sqrt{g_{00}} x^0 c = \sqrt{g_{00}} t.$



## Dispersion relation in a plasma

$$k_j k^j = g^{ij} k_i k_j = \frac{\omega_p^2}{c^2}$$

### Photon effective (rest) mass

$$\omega = \omega(k_\alpha, x^i)$$

$$\omega = \omega_1 + \sqrt{\nu^2 + \omega_1^2 - \omega_2^2}$$

$$\omega_1 = Y^{0\alpha} k_\alpha c \quad , \quad \omega_2 = \sqrt{Y^{\alpha\beta} k_\alpha k_\beta} c$$

$$m_{eff} = \hbar \omega_p / c.$$

$$Y^{ij} = g^{ij} / g^{00}$$

$$\nu^2 = \omega_p^2 / g^{00}$$





## G field as an optical medium

$$\omega = \frac{kc}{n(k_\alpha, x^\alpha, t)}$$

## Refractive index of the gravitational field

$$n(k_\alpha, x_\alpha, t) = \frac{k}{Y^{0\alpha}k_\alpha} \left[ 1 + \nu^2 + \sqrt{1 - \frac{Y^{\alpha\beta}k_\alpha k_\beta}{(Y^{0\alpha}k_\alpha)^2}} \right]^{-1}$$



INSTITUTO  
SUPERIOR  
TÉCNICO

## Robertson metric

$$ds^2 = (1 + 2\Phi)c^2 dt^2 - (1 - 2\Phi)a^2(t) dr^2$$

**Perturbed Newtonian potential**

$$\Phi \equiv \Phi(r, t)$$

**Scale function (universe expansion)**

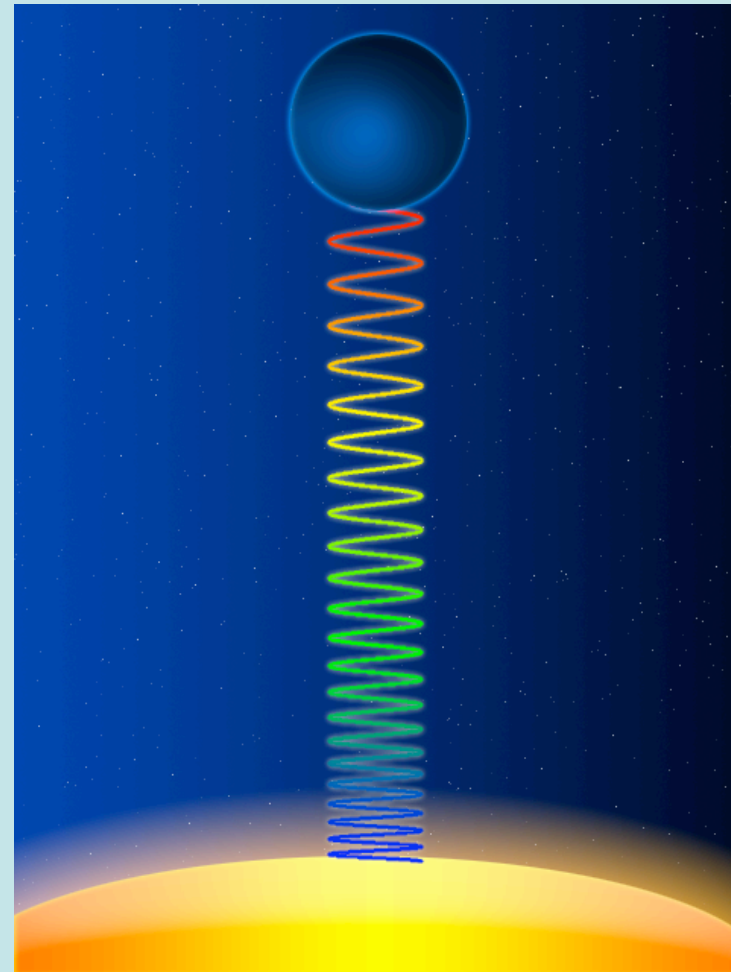
$$a(t)$$

## Gravitational red-shift

**Constant frequency,  $\omega$**

**But proper frequency changes!**

$$\frac{dk_\alpha}{dt} = -\frac{\partial\omega}{\partial x^\alpha} = \frac{kc}{n^2} \frac{\partial n}{\partial x^\alpha}$$





## Gravitational red-shift (cont.)

Initial proper frequency (emitted at  $R$ )

$$\omega_e'^2 = \frac{\omega^2}{g_{00}(R)} = \omega_p^2(R) + \frac{k_e^2 c^2}{1 - 2\Phi(R)}$$

$\omega$  remains constant

Final proper frequency

$$\omega_f'^2 = \omega^2 = k_f^2 c^2 = (1 + 2\Phi(R)) \left[ \omega_p^2(R) + \frac{k_e^2 c^2}{1 - 2\Phi(R)} \right]$$

Observed frequency shift

$$\Delta\omega' \simeq \omega_e' \Phi(R)$$

for  $|\Phi(R)| \ll 1$ ,

Apparent temperature shift

$$\frac{\Delta T}{T} = \Phi(R)$$

**(Simple) Sachs-Wolfe effect**

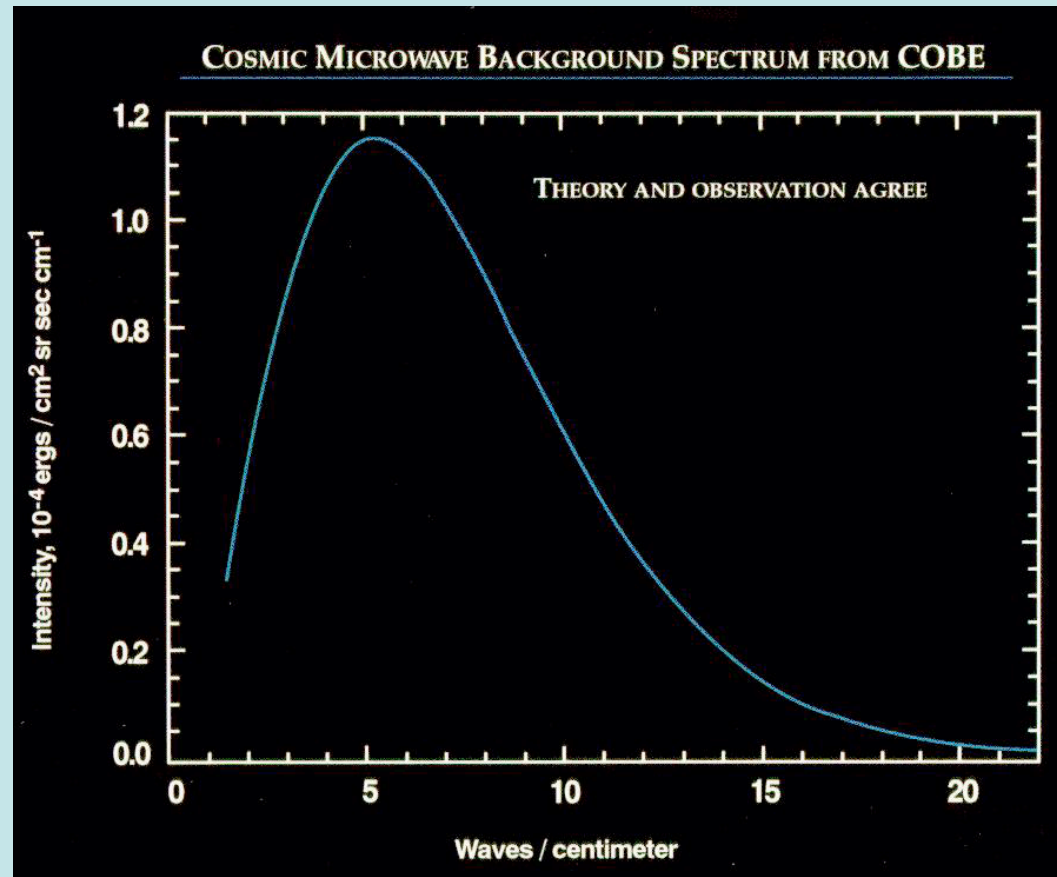


INSTITUTO  
SUPERIOR  
TÉCNICO

# Cosmic Microwave Background

- First Predicted by G. Gamow, R. Alpher and R. Herman (1948)
- First detected by A. Penzias and R. Wilson (1965)

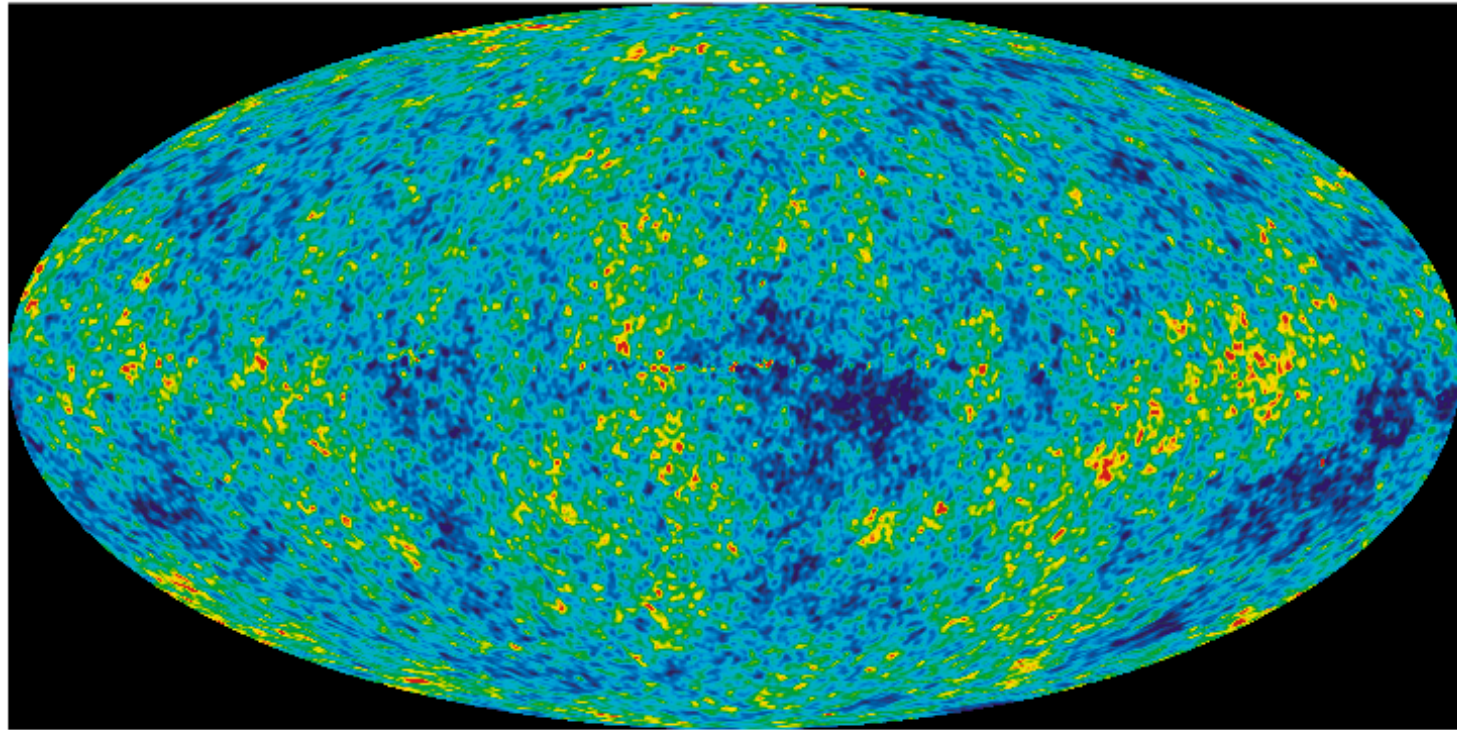
$T = 2.725 \text{ K}$





INSTITUTO  
SUPERIOR  
TÉCNICO

## Cosmic Temperature map



WMAP\_2008



## Cosmological red-shift (locally perturbed)

### Contribution from the scale factor $a(t)$

Continuity equation

$$\dot{\rho} + 3H(\rho + p) = 0$$

$$H = \dot{a}/a$$

Equation of state

$$p = w\rho.$$

Matter dominated universe,  $w = 0$   
Radiation dominated,  $w = 1/3$

Apparent temperature shift ( $aT = \text{const.}$ )

$$\frac{\Delta T}{T} = -\frac{\Delta a}{a} = -\frac{2\Phi(R)}{3(1+w)}$$

**Accumulated SW effect**

$$\frac{\Delta T}{T} = \frac{1+3w}{3(1+w)} \Phi(R)$$

Matter dominated universe  
(well known result)

$$\Delta T/T = \Phi(R)/3.$$



## Integrated Sachs-Wolfe

$$\Phi(r, t)$$

$$\frac{dk}{dt} = -\frac{\partial\omega}{\partial r} = -\frac{2k^2c^2}{\omega(1-2\Phi)^2} \frac{\partial\Phi}{\partial r} - \frac{(1+2\Phi)}{2\omega} \frac{\partial\omega_p^2}{\partial r}$$

### Low plasma density limit

$$k(t) = k(0) \exp \left[ -2c \int_{r(0)}^{r(t)} \frac{\partial\Phi(r(t'), t')}{\partial r} dt' \right]$$

$$\omega_p^2 \ll 4k^2c^2|\Phi|$$

### Low frequency shift limit

$$k(t) = k(0) \left( 1 - 2 \int \frac{\partial\Phi}{\partial r} ds \right)$$

$$\frac{\Delta T}{T(0)} = -2 \int_{r(0)}^{r(t)} \frac{\partial\Phi}{\partial r} ds$$

Mendonça, Bingham and Wang, CQG (2008)



## Moving perturbations

$$\Phi(r, t) = \Phi(r - ut)$$

### Dynamical invariant

$$\Omega = \omega(\eta) - uk, \text{ with } \eta = r - ut$$

### Generalized Sachs-Wolfe

$$\frac{\Delta T}{T} \simeq \frac{\Phi_0}{(1 - \beta)} \left[ 1 + \beta \left( 1 - \frac{\omega_{p0}^2}{2\omega_T^2} \right) \right] + \frac{\beta\omega_{p0}^2}{2\omega_T^2(1 - \beta)}$$

### Purely vacuum perturbations

$$\frac{\Delta T}{T} \simeq \Phi_0 \frac{(1 + \beta)}{(1 - \beta)}$$

$$\omega_{p0} \simeq 0$$

**It is not a Doppler correction!**





## Photons in a plane G wave

### Perturbed flat space time

$$g^{ij} = \eta^{ij} + h^{ij}$$

Minkowski metric tensor

$$\eta^{\alpha\beta} = -\delta^{\alpha\beta}, \eta^{00} = 1$$

### Weak gravitational wave

$$a = A \sin(k_0 x^0 + k_1 x^1 + \phi) = A \sin(qx - \Omega t + \phi)$$

$$h_{22} = -h_{33} = a$$

### Photon dispersion relation

$$\omega = kc \left\{ 1 + \frac{a}{2} \left[ \left( \frac{k_y}{k} \right)^2 - \left( \frac{k_z}{k} \right)^2 \right] \right\}$$



## Nearly parallel photon propagation

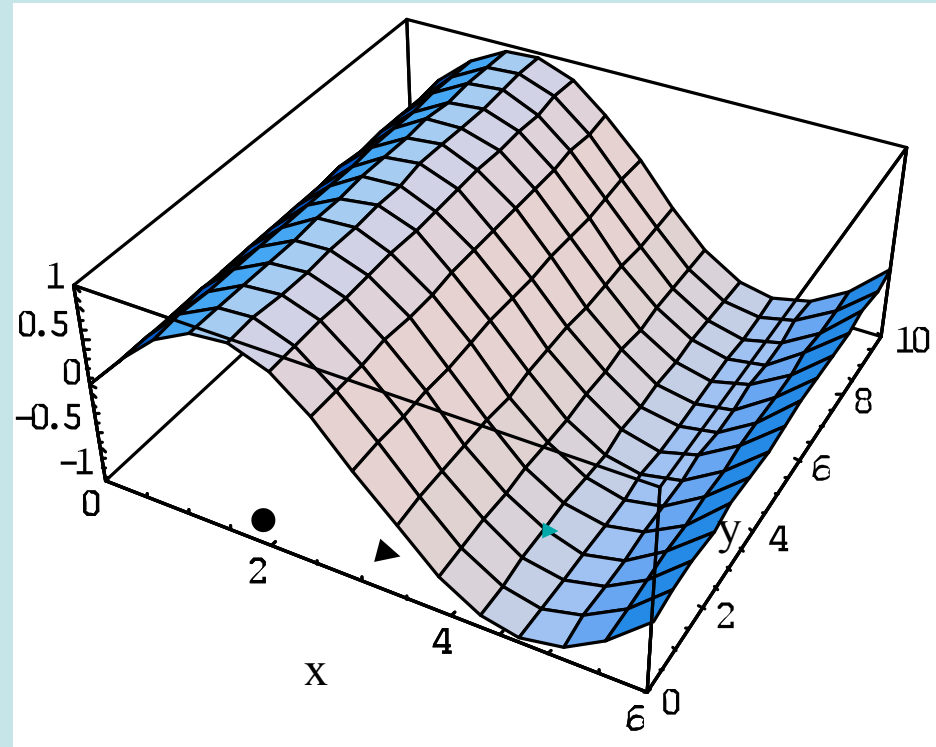
### Perpendicular photon motion

$$y(t) = ck_y \int^t \frac{1 + a(t')}{k(t')} dt'$$

### Parallel photon motion

$$\frac{dx}{dt} = c \frac{k_x}{k}$$

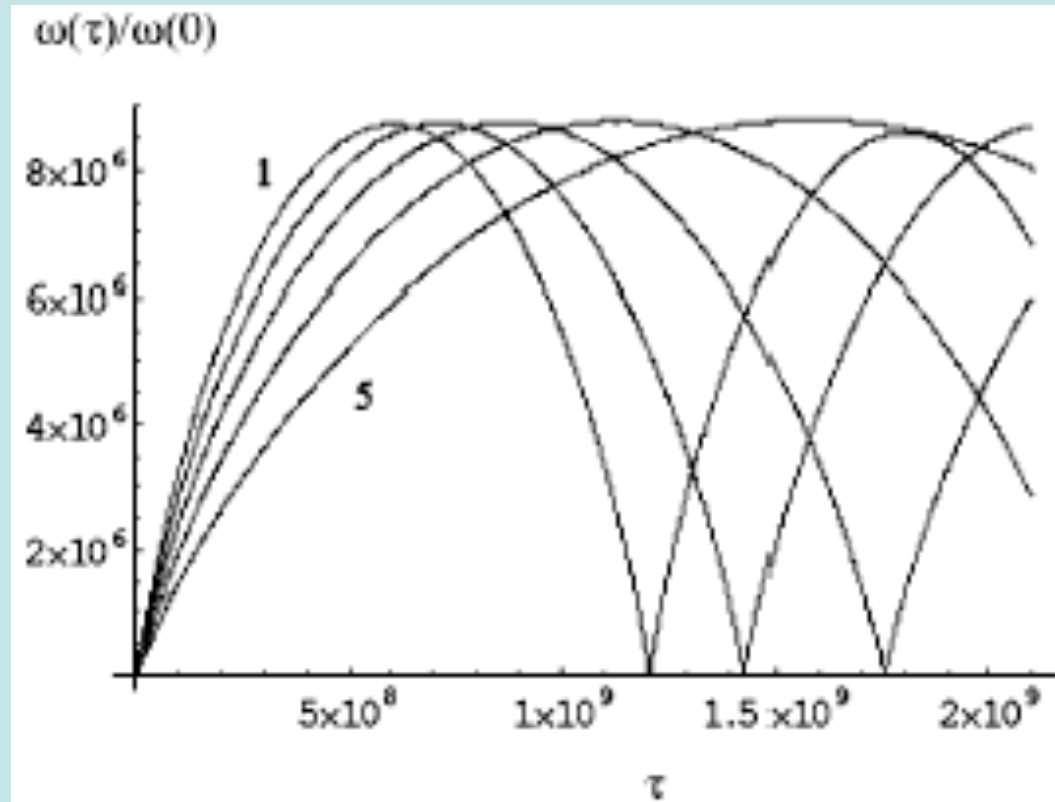
$$\frac{dk_x}{dt} = -\frac{1}{2} \left( \frac{k_y}{k} \right)^2 \frac{\partial a}{\partial t}$$





# Typical photon trajectories

$$\frac{d\omega}{dt} = c \frac{dk_x}{dt}$$



✚ Relevance of this numerical example:

Gravitational wave with frequency:  $\Omega = 10^4 \text{ s}^{-1}$

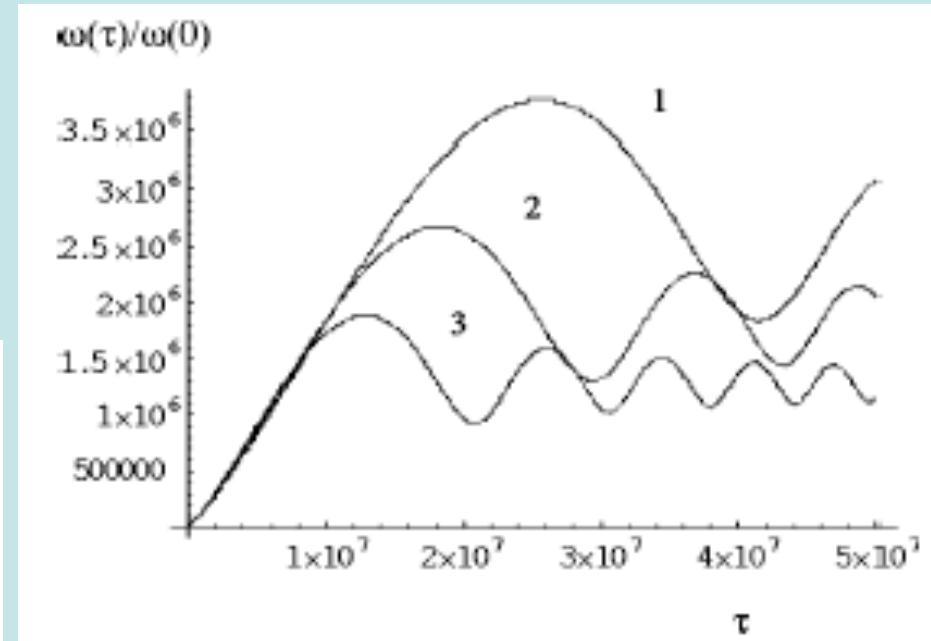
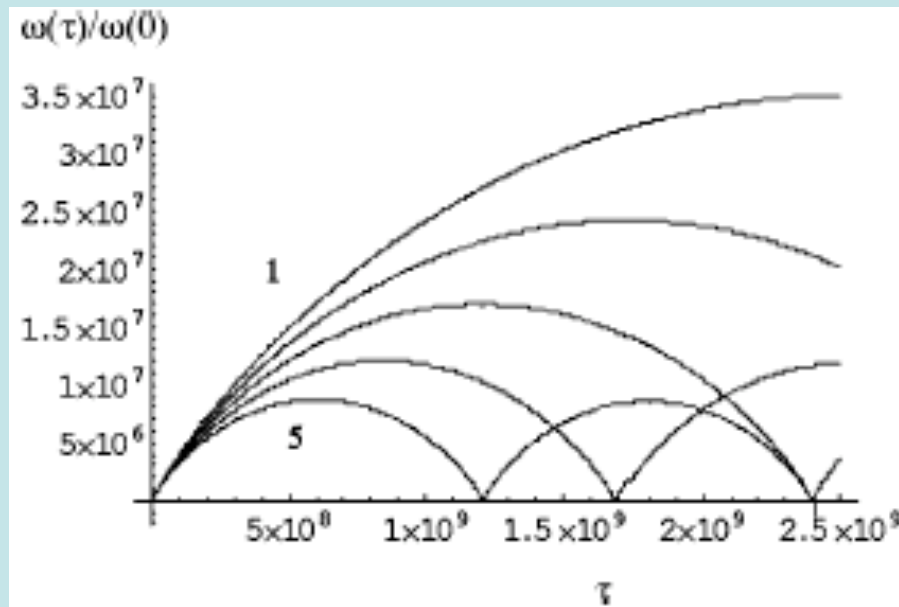
$$\tau = A c t / 2$$

and amplitude  $A = 10^{-4}$



## Decay of G wave (irreversibility)

### Influence of G wave amplitude



### Model for gamma ray bursters?

Mendonça and L.Drury, PRD (2002)



INSTITUTO  
SUPERIOR  
TÉCNICO

# SN explosions and GR bursts: Coupling between G waves and photons (and neutrinos)

Spherical gravitational  
Wave burst



Accelerated escaping photons



[Acceleration of neutrinos could also be relevant to  
Energetic Cosmic rays - not affected by the GZK limit]



## Photons in a spherical G wave

Space-time metric

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$

$\eta_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$  and  $h_{\mu\nu} \rightarrow 0$  at infinity.

Photon equations of motion

$$\frac{dx^\gamma}{dt} = 1 + h^{0\gamma} - \frac{1}{k} h^{\gamma\beta} k_\beta + \frac{1}{2k^3} h^{\alpha\beta} k_\alpha k_\beta k^\gamma$$

$$\frac{dk_\gamma}{dt} = -(\partial_\gamma h^{0\alpha}) k_\alpha + \frac{1}{2} \left[ -(\partial_\gamma h^{00}) + (\partial_\gamma h^{\alpha\beta}) \frac{k_\alpha k_\beta}{k^2} \right]$$

Quadrupole moment

$$Q_{\alpha\beta} = \int \rho(x) x^\alpha x^\beta d^3x$$

$$\bar{h}_{\alpha\beta}(t) = \frac{2}{r} \frac{d^2}{dt^2} Q_{\alpha\beta}(t-r) + O(1/r^2)$$

Wang, Mendonça, Bingham (2009?)



## Photon up-shift and bunching

$$\frac{dr}{dt} = 1 - \frac{1}{2}h^{rr}$$

$$\frac{dk}{dt} = k \frac{\partial}{\partial r} h^{rr}$$

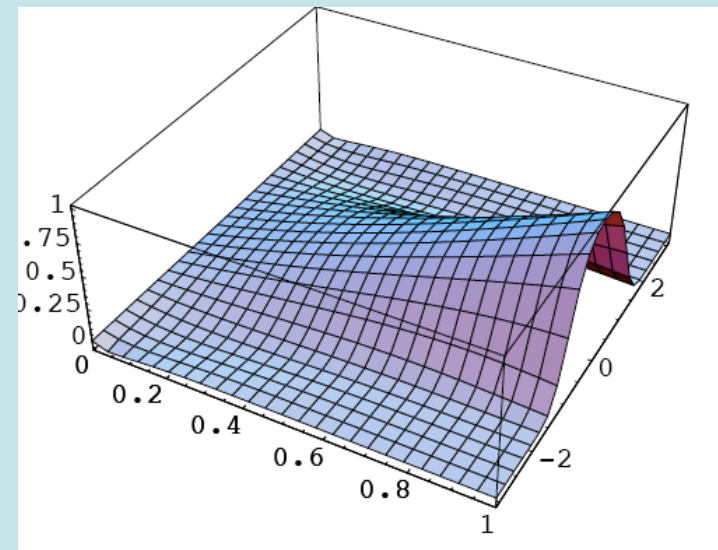
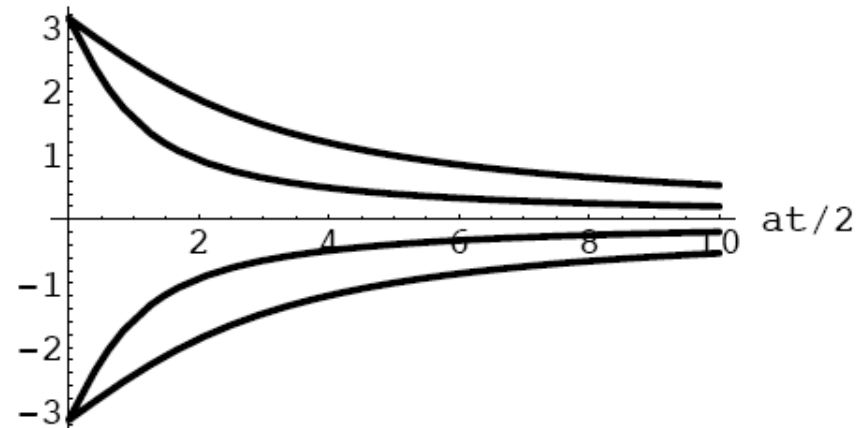
### Photon momentum

$$k(r) = k_0 \left( \frac{r}{r_0} \right)^{\pm Q}$$

#### Example:

$k_0 = 1$  eV,  $r_0 = 150$  Km,  
 $k = 1$  MeV,  $r = 1$  AU

Position x





INSTITUTO  
SUPERIOR  
TÉCNICO

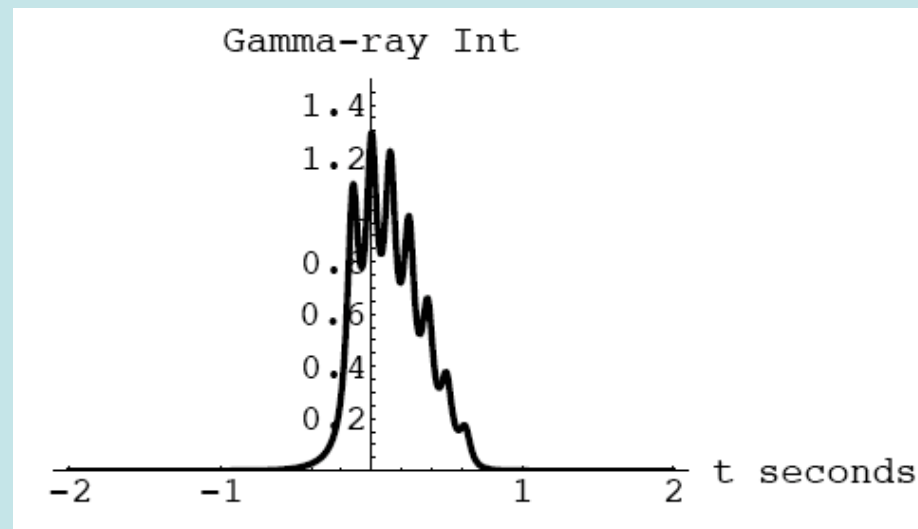
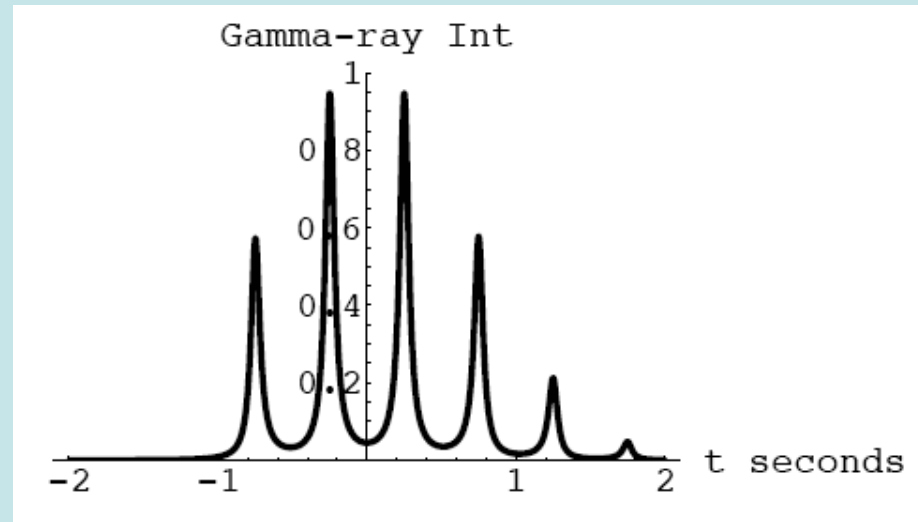
## Expected Gamma-ray signals

Complete and incomplete bunching  
inside Gwave packet

Question:

- Is it gauge dependent?
- Is there any conformal field?

$$G^{ij} = A g^{ij}$$







INSTITUTO  
SUPERIOR  
TÉCNICO

# Observed Gamma-ray bursts

nature

Vol 444 | 21/28 December 2006 | doi:10.1038/nature05376

## LETTERS

### A new $\gamma$ -ray burst classification scheme from GRB 060614

N. Gehrels<sup>1</sup>, J. P. Norris<sup>1</sup>, S. D. Barthelmy<sup>1</sup>, J. Granot<sup>2</sup>, Y. Kaneko<sup>3</sup>, C. Kouveliotou<sup>4</sup>, C. B. Markwardt<sup>1,5</sup>, P. Mészáros<sup>6,7</sup>, E. Nakar<sup>8</sup>, J. A. Nousek<sup>6</sup>, P. T. O'Brien<sup>9</sup>, M. Page<sup>10</sup>, D. M. Palmer<sup>11</sup>, A. M. Parsons<sup>1</sup>, P. W. A. Roming<sup>6</sup>, T. Sakamoto<sup>1,12</sup>, C. L. Sarazin<sup>13</sup>, P. Schady<sup>6,10</sup>, M. Stamatikos<sup>1,12</sup> & S. E. Woosley<sup>14</sup>

Are we already detecting  
G waves?

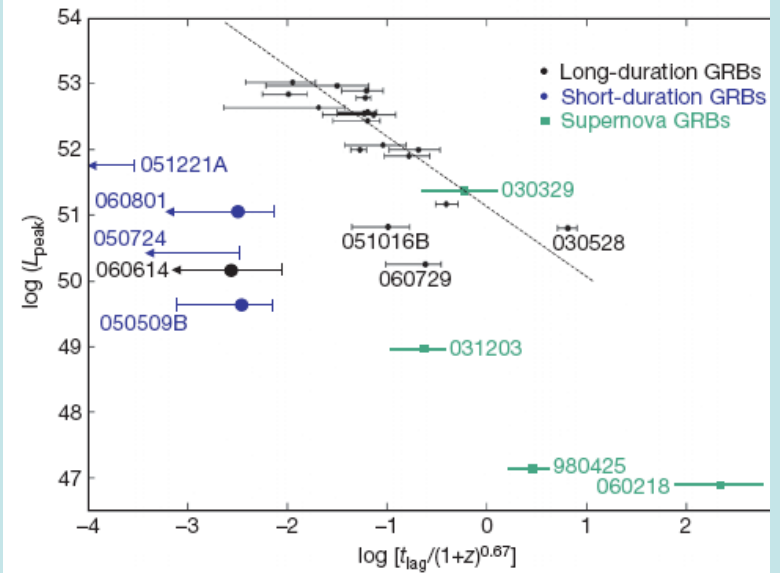


Figure 2 | Spectral lag as a function of peak luminosity showing GRB 060614 in the region of short-duration GRBs. The lags and peak

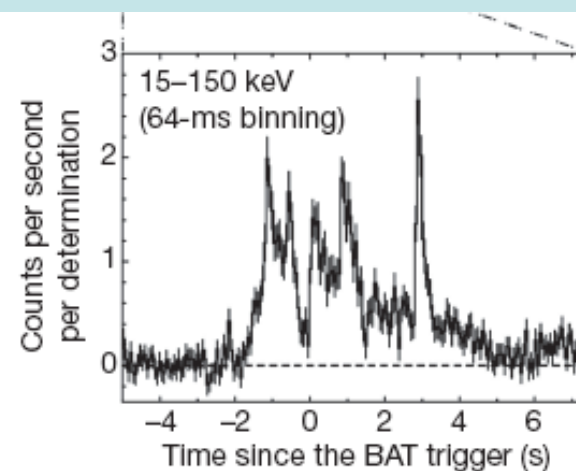


Figure 1 | The light curve of GRB 060614 as observed with the BAT.



# Photon-graviton interactions

## Maxwell's equations in a gravitational field

$$F_{;i}^{ik} = -\frac{4\pi}{c} J^i$$

$$F_{ik} = \partial_i A^k - \partial_k A^i$$

Electromagnetic  
field tensor

In vacuum ( $J^i = 0$ )

$$\partial_k g^{ij} g^{km} F_{jm} = -\frac{F^{ik}}{2g} \partial_k g$$

Einstein equation (in the weak field approximation:  $g_{ik} = \eta_{ik} + h_{ik}$ )

$$\partial^l \partial_l h_{ij} = -16\pi G T_{ij}$$

E.m. energy momentum tensor

$$T_{ij} = \frac{1}{4\pi} \left( -F_{il} F_k^l + \frac{1}{4} g_{ik} F_{lm} F^{lm} \right)$$



# Nonlinear wave coupling

## Wave perturbations

$$A_j = e_j A \exp(ik_l x^l)$$

$$h_{ij} = \varepsilon_{ij} h \exp(iq_l x^l)$$

## Mode coupled equations

$$2ik^l \partial_l A = Wh^* A' \exp(i\phi)$$

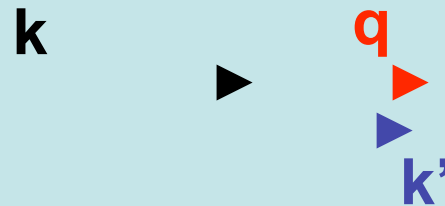
$$2ik'^l \partial_l A' = W' h A \exp(-i\phi)$$

$$2iq^l \partial_l h = W_G AA' \exp(-i\phi)$$

## Phase matching

$$\phi = (k'_l - q_l - k_l) x^l \approx 0$$

**Important question: allowed  
geometric configurations**



**No coupling for co-  
propagation,  $W = 0$**



**Coupling exists for counter-  
propagation,  $W \neq 0$**



INSTITUTO  
SUPERIOR  
TÉCNICO

## Conclusions

- **Photons can be energized by Gravitational fields and waves;**
- **Geometric optics can be efficiently used;**
- **The gravitational field behaves as an optical medium;**
- **A generalized Sachs-Wolfe effect (including dynamical and plasma corrections) can be obtained;**
- **Gravitational waves can energize photons;**
- **Possible explanation for gamma-ray bursts (?);**
- **Gravitational waves are probably being observed;**
- **Nonlinear photon-graviton interactions are possible.**