



**The Abdus Salam
International Centre for Theoretical Physics**



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**Joint ICTP-IAEA Advanced Workshop on Multi-Scale Modelling for
Characterization and Basic Understanding of Radiation Damage
Mechanisms in Materials**

12 - 23 April 2010

Radiation damage summary

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**Joint ICTP-IAEA Advanced Workshop on Multi-Scale Modeling
for Characterization and Basic Understanding of Radiation
Damage Mechanisms in Materials, April 12-23, 2010**



The experimental validation of radiation damage modeling: an “historical” perspective

III. Fe, steels and other bcc metals

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Introduction



- The magnetic behavior of Fe has in general not been taken into account in radiation damage calculations performed up to the late 90's
- As a result, these calculations ignored early experimental results that indicated that the minimum energy configuration of interstitial defects was different from that resultant from the calculations

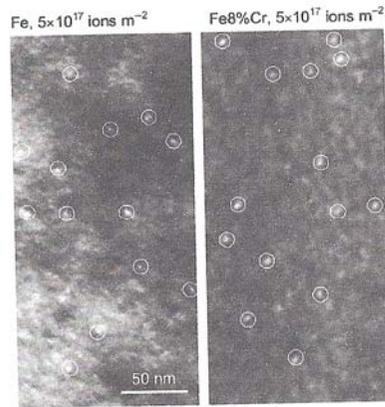
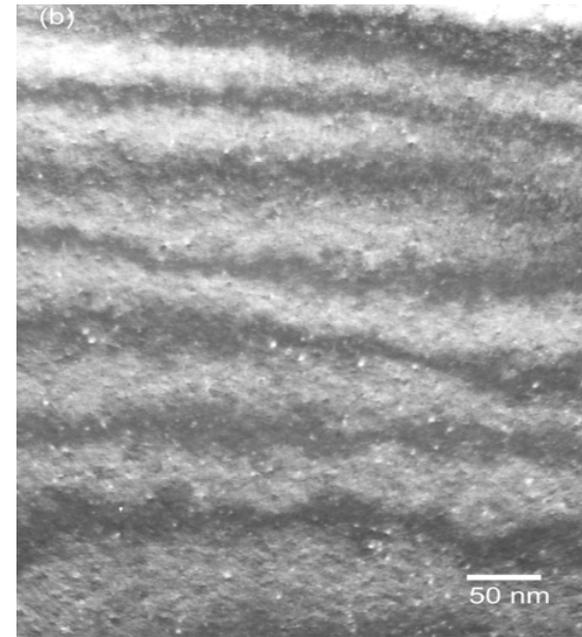


Figure 1. Small dislocation loops in (110) UHP Fe and (113) UHP Fe-8% Cr irradiated with 150 keV Fe⁺ ions to a dose of $5 \times 10^{17} \text{ m}^{-2}$ ions at RT. The micrographs were taken under weak beam conditions, ($g, \sim 3g$) using the reflection $g = 1\bar{1}0$.

Defect microstructures in Fe



Luppo et al (2000)

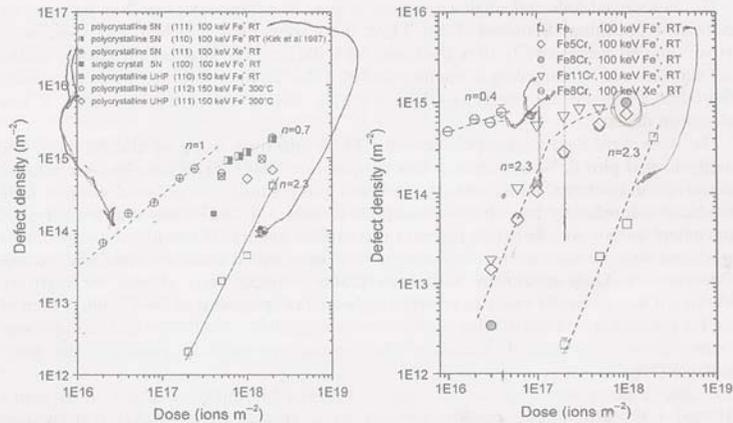
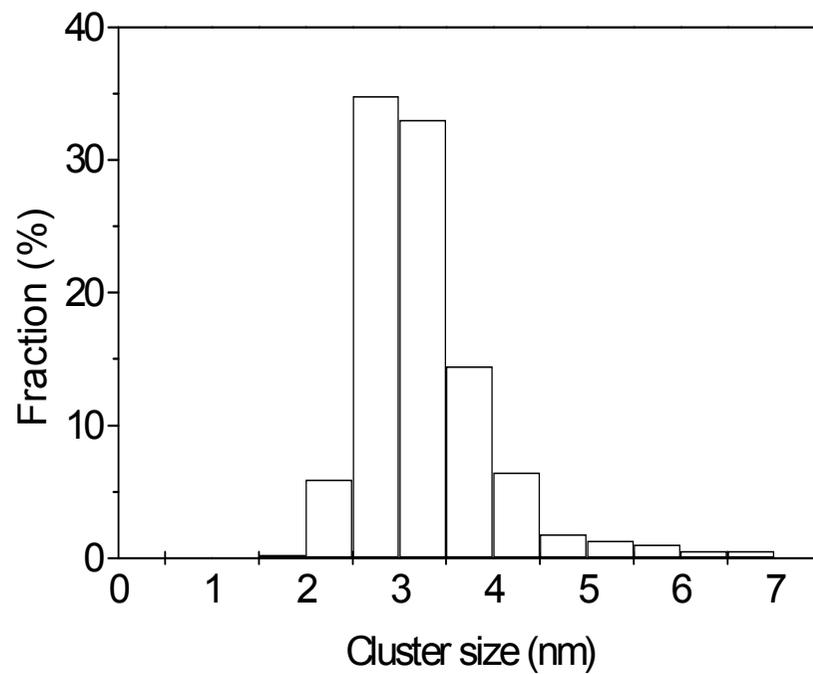


Figure 3. (a) Defect areal densities versus dose for pure Fe. For polycrystalline specimens the approximate orientations of the grains used in the experiments are indicated; the true grain orientations were usually several degrees off the pole. Some representative error bars are shown. Values n of the slopes are shown. (b) Loop areal densities as a function of dose in Fe-Cr alloys irradiated at RT with 100 keV ions in foils oriented close to (111).

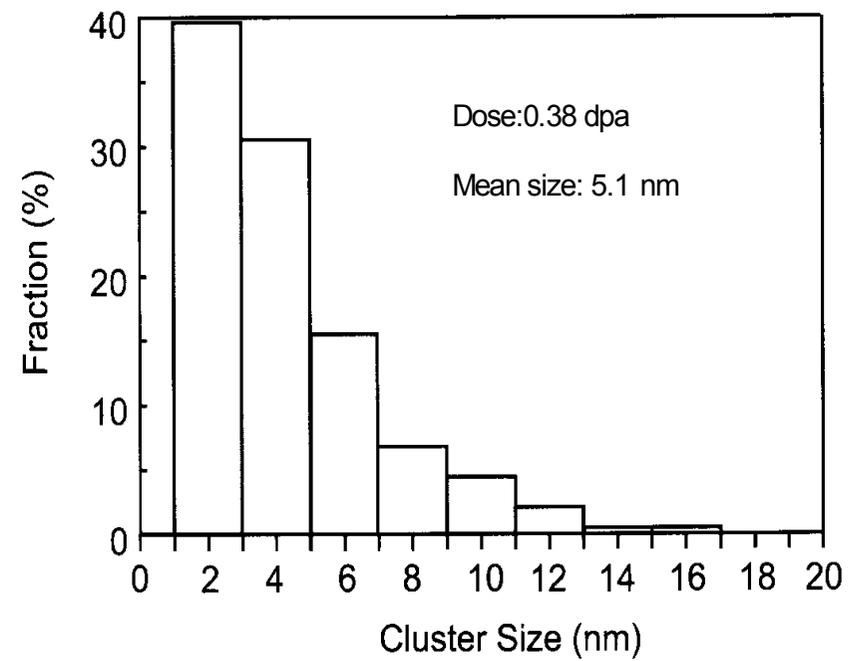
Yao et al (2008)



Cluster size distribution in irradiated Fe

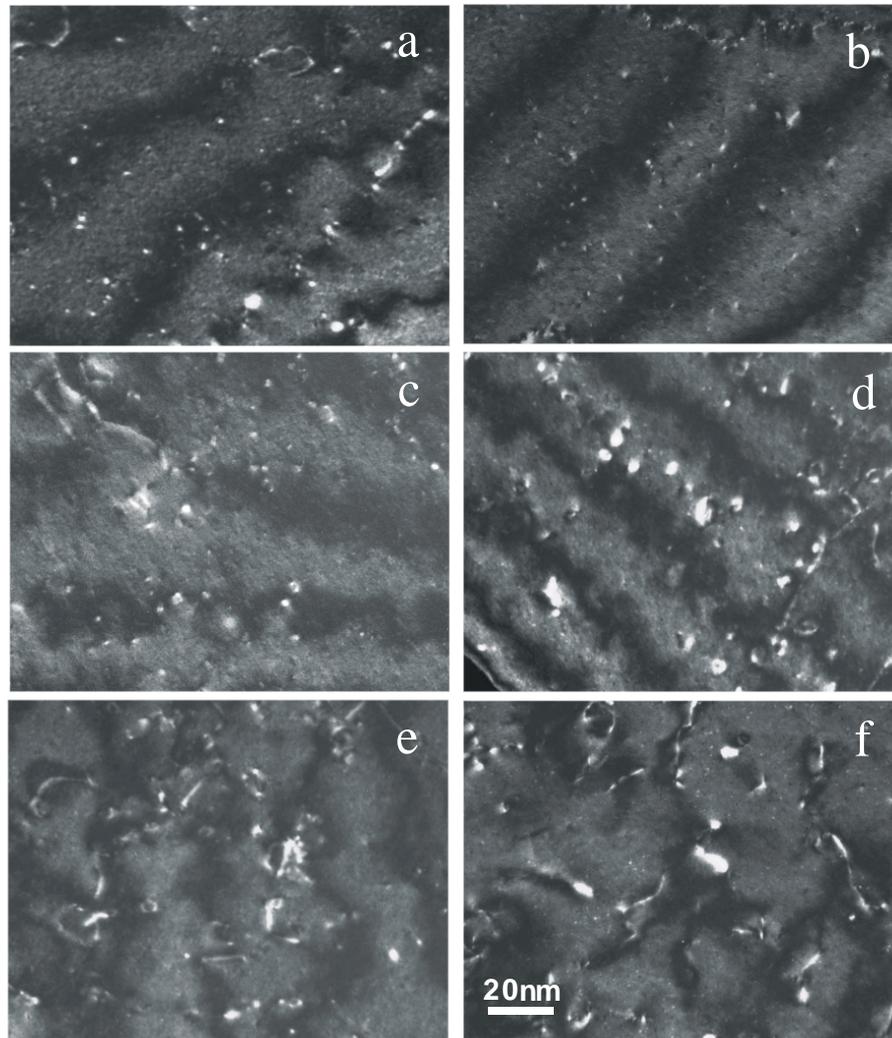


p⁺ (0.33 dpa)



n (0.38 dpa)

Postirradiation defect microstructure



F82H steel irradiated in SINQ

(a) 10 dpa; 140⁰ C

(b) 10.2 dpa; 175⁰ C

(c) 10.2 dpa; 210⁰ C

(d) 10.3 dpa; 255⁰ C

(e) 9.9 dpa; 295⁰ C

Damage accumulation in α -Fe



- Prevailing view (Bacon and Osetsky, 2003) based on modeling using empirical potentials:
 - a) SIA's: $\langle 111 \rangle$ and $\langle 110 \rangle$ configurations have comparable energies
 - b) The same potentials predict that clusters of two or more SIAs form a set of parallel dumbbells moving rapidly in the $\langle 111 \rangle$ direction with activation energies ~ 0.01 eV.
 - c) But **measured** migration values are much higher: 0.3 eV for SIAs and 0.42 eV for small clusters

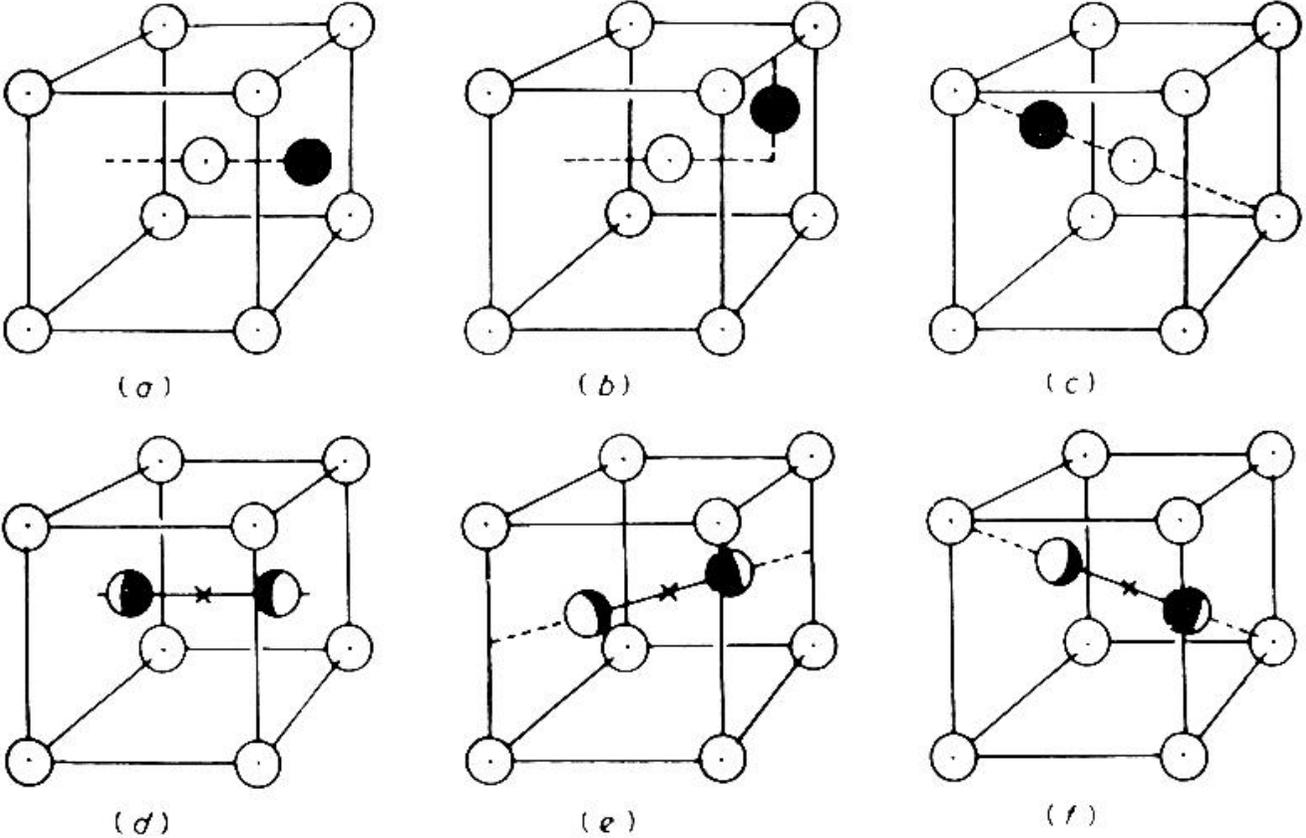
Damage accumulation in α -Fe



In fact experimental results indicate that

- $\langle 110 \rangle$ dumbbell configuration established from a number of experimental methods (Schultz (1982), Takaki (1983), Ehrhart (1986)) with a mobility of $E_{i_m} = 0.27$ eV.
- Very early neutron irradiation experiments performed in not very pure Fe at He temperatures: no annealing peak of RR observed at low T
- Post-mortem TEM observations show only $\langle 111 \rangle$ and $\langle 100 \rangle$ loops

Interstitial configurations in bcc



Self-interstitial configurations in a bcc lattice (Johnson 1973): (a) octahedral, (b) tetrahedral, (c) crowdion, (d) (100)-split (dumbbell), (e) (110)-split, (f) (111)-split.

Magnetism is important in α -Fe and controls its singular behavior amongst bcc metals



Basic parameters of bcc transition metals and formations energies of defects, calculated using DFT. Italics: experimental values.

From D. Nguyen-Manh *et al* (2006)

	V	Nb	Ta	Cr	Mo	W	Fe
a (Å)	3.04	3.32	3.31	2.85	3.17	3.18	2.88
	<i>3.03</i>	<i>3.30</i>	<i>3.30</i>	<i>2.88</i>	<i>3.15</i>	<i>3.16</i>	<i>2.87</i>
B (Mbar)	1.71	1.73	1.99	1.92	2.68	3.05	1.80
	<i>1.62</i>	<i>1.70</i>	<i>2.00</i>	<i>1.90</i>	<i>2.72</i>	<i>3.23</i>	<i>1.68</i>
H^v (eV)	2.51	2.99	3.14	2.64	2.96	3.56	2.07
	<i>2.1–2.2</i>	<i>2.6–3.1</i>	<i>2.8–3.1</i>	<i>2.0–2.4</i>	<i>2.6–3.2</i>	<i>3.5–4.1</i>	<i>1.8–2.2</i>
H_m^v (eV)	0.62	0.91	1.48	0.91	1.28	1.78	0.67
	<i>0.5–1.2</i>	<i>0.6–1.0</i>	<i>0.7–1.9</i>	<i>0.95</i>	<i>1.3–1.6</i>	<i>1.7–2.0</i>	<i>0.55</i>
$\langle 111 \rangle$ dumbbell	3.367	5.253	5.832	5.685	7.417	9.548	4.45
	<i>3.14</i>	<i>4.795</i>	<i>7.157</i>		<i>7.34</i>	<i>8.919</i>	
$\langle 111 \rangle$ crowdion	3.371	5.254	5.836	5.660	7.419	9.551	4.45
	<i>3.15</i>	<i>4.857</i>	<i>7.158</i>		<i>7.34</i>	<i>8.893</i>	
$\langle 110 \rangle$	3.652	5.597	6.382	5.674	7.581	9.844	3.75
	<i>3.48</i>	<i>4.482</i>	<i>6.847</i>		<i>7.51</i>	<i>9.641</i>	
tetrahedral	3.835	5.758	6.771	6.189	8.401	11.05	4.26
	<i>3.69</i>				<i>8.20</i>		
$\langle 100 \rangle$	3.918	5.949	7.003	6.643	9.004	11.49	4.75
	<i>3.57</i>	<i>4.821</i>	<i>8.068</i>		<i>8.77</i>	<i>9.815</i>	
octahedral	3.964	6.060	7.095	6.723	9.067	11.68	4.94
	<i>3.62</i>				<i>8.86</i>		

Disagreement resolved by *ab initio* calculations



The disagreement was solved by DFT calculations (Fu, Willaime and Ordejon, Phys Rev Let 2004) in which the energies of the different interstitial configurations were obtained for Fe, showing that the $\langle 110 \rangle$ dumbbell configuration has an energy lower than the $\langle 111 \rangle$ by 0.7 eV.

This is a consequence of **magnetism** in Fe. For all other bcc metals in the Periodic Table, it is the $\langle 111 \rangle$ dumbbell configuration which has the lower energy

The effects of magnetism in Fe and its alloys has led to the need to develop **magnetic potentials** (Dudarev and Derlet (2005)) or semi-empirical potentials where the effects of magnetism ([110] dumbbell as minimum energy configuration) are included.

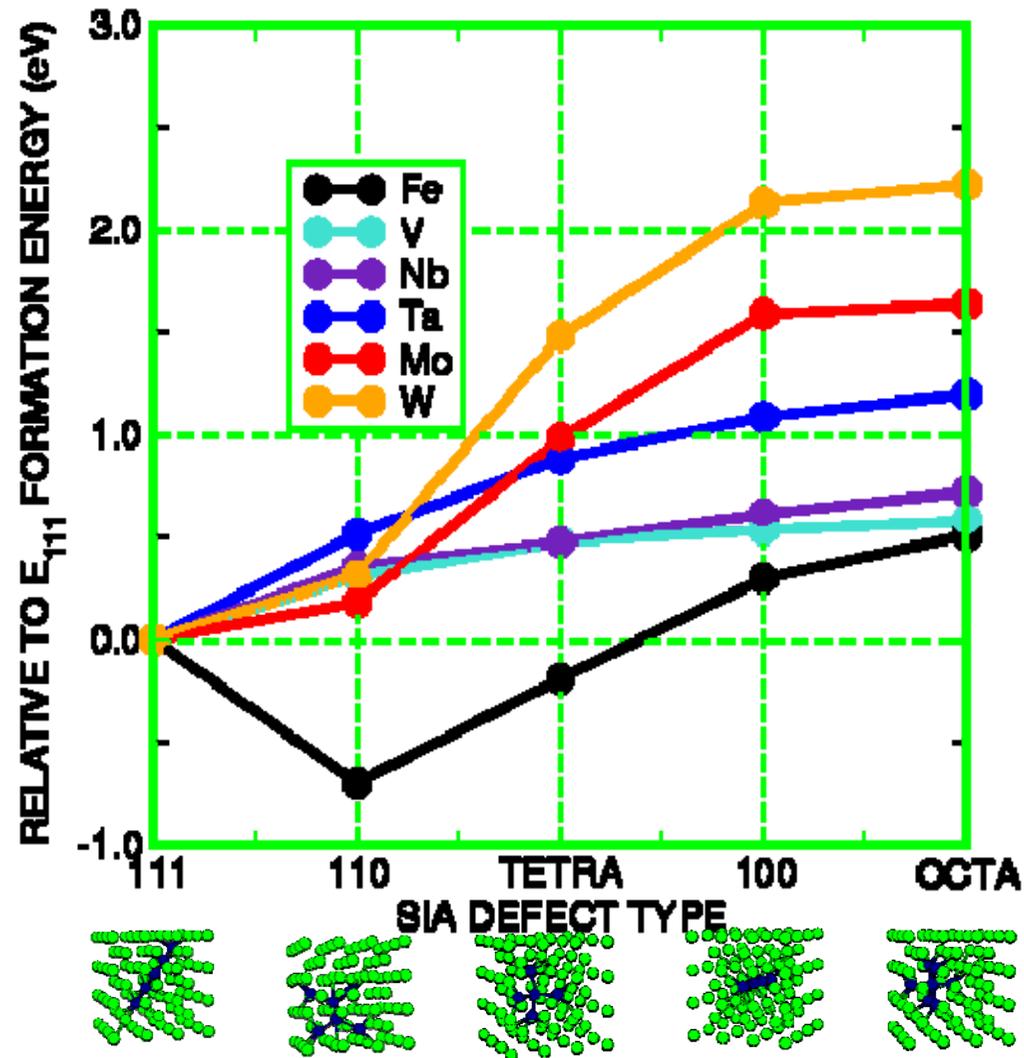
Practically all the microscopic processes associate with the irradiated microstructure evolution are **thermally activated**. The rate at which a given event takes place is given by the Arrhenius law

$$\nu = \nu_0 \exp\left(-\frac{Q}{k_B T}\right)$$

Because it is in an exponential Q needs to be defined with good precision. Present day DFT calculations can reach the accuracies needed, of the order of 0.1 eV. The migration of the SIA in Fe takes place by a translation-rotation mechanism with a migration energy of 0.34 eV

DFT modelling of SIA defects

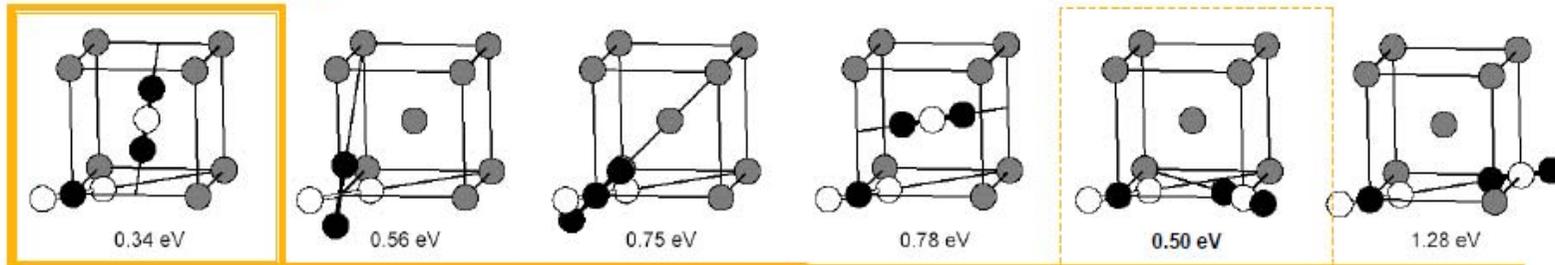
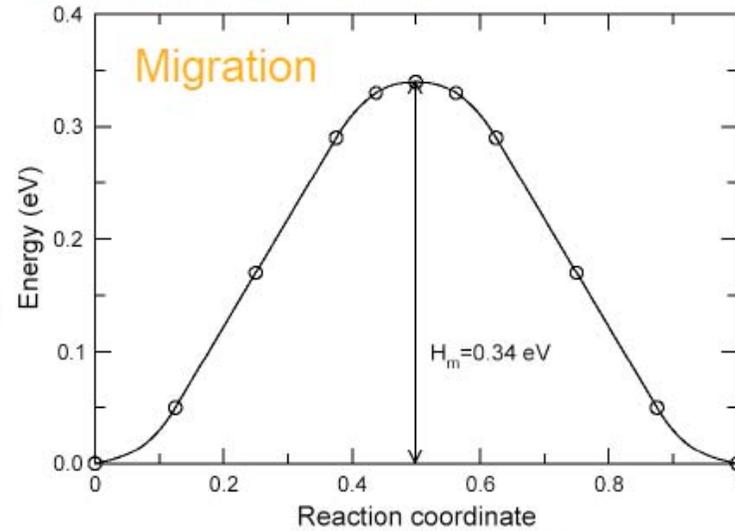
D. Nguyen Manh, S. L. Dudarev and A. Horsfield (2004)



Self-interstitials in BBC Fe

Chu Chun Fu & F. Willaime 2003-2004

Migration mechanisms
of $\langle 110 \rangle$ dumbbells
and diffusion path



Translation / rotation 1st neighbour jump \rightarrow 3D Migration
Excellent agreement with experiment : $H_m = 0.3$ eV

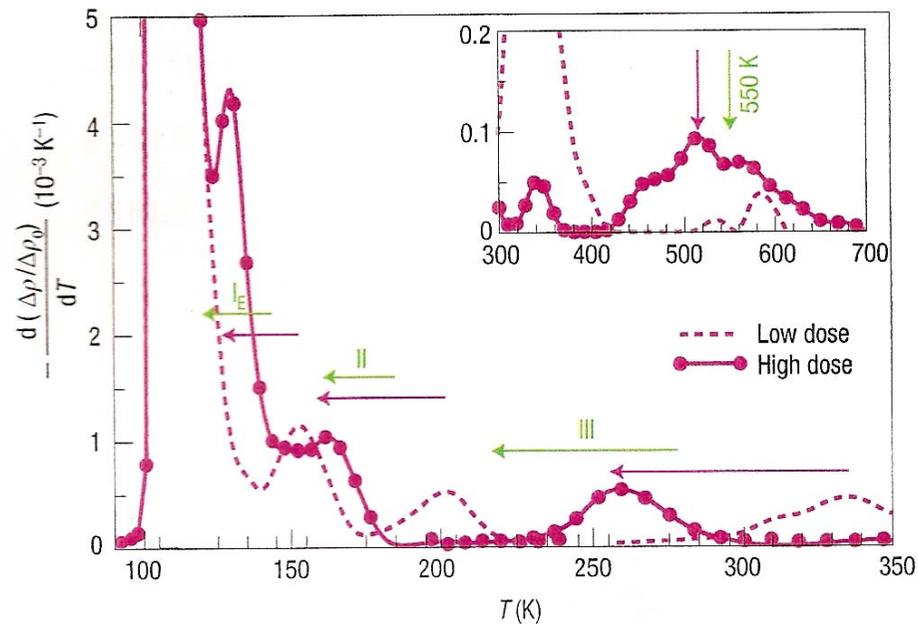


Figure 5 Effect of irradiation dose. Derivative with respect to temperature of the resistivity recovery in the case of high irradiation dose (magenta solid line). The low-dose simulation is from Fig. 4 (dashed line). The beginning and the end of the horizontal arrows indicate the stage temperatures at low and high irradiation doses (2×10^{-6} and 200×10^{-6} dpa respectively): green arrows for experiments and magenta arrows for the simulation. Stage (IV) is shown in the inset.

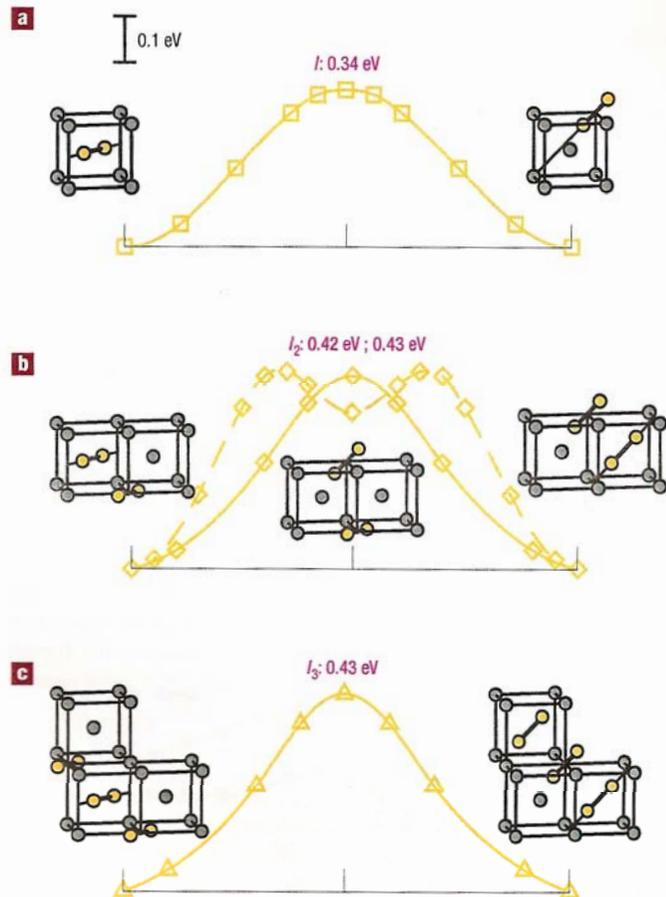


Figure 2 Migration of interstitial-type defects. Energy variation along the most favourable migration pathways for I_1 , I_2 and I_3 , determined by *ab initio* calculations. The corresponding migration energies and schematic representations of the migration mechanisms starting from the defect ground-states are shown: the orange and grey spheres indicate self-interstitial and lattice atoms respectively. The energy scale is the same for all graphs.

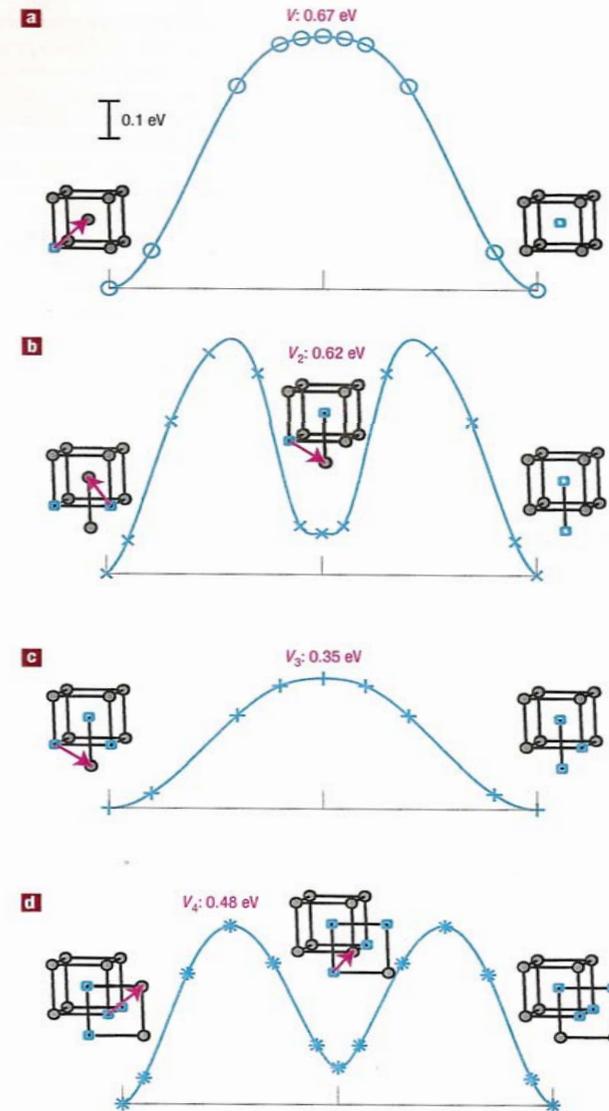
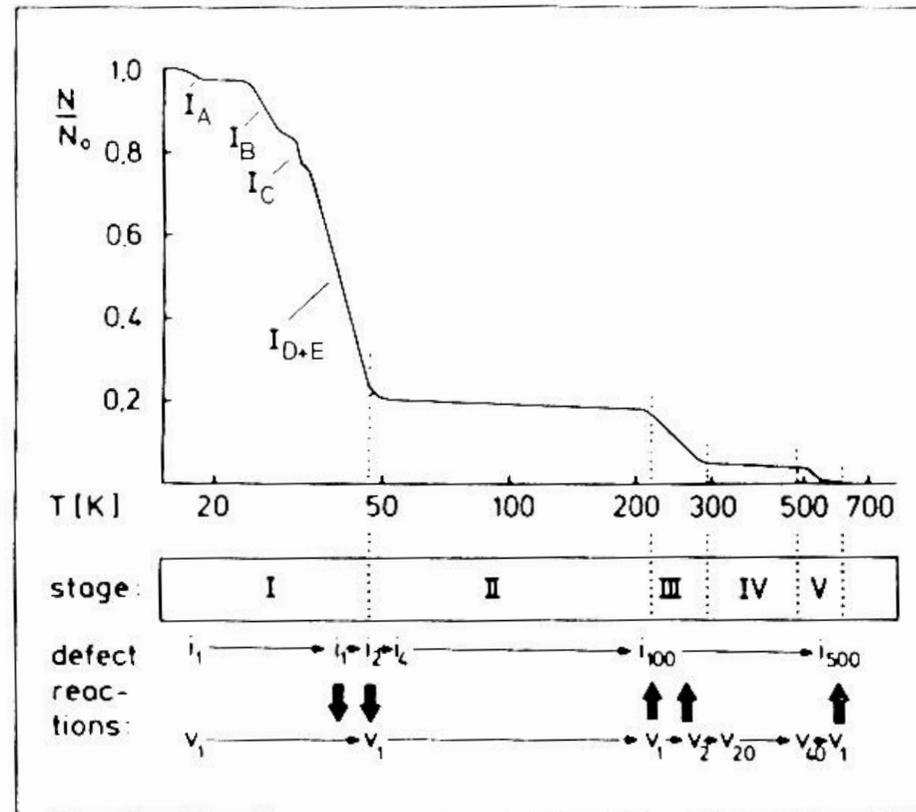


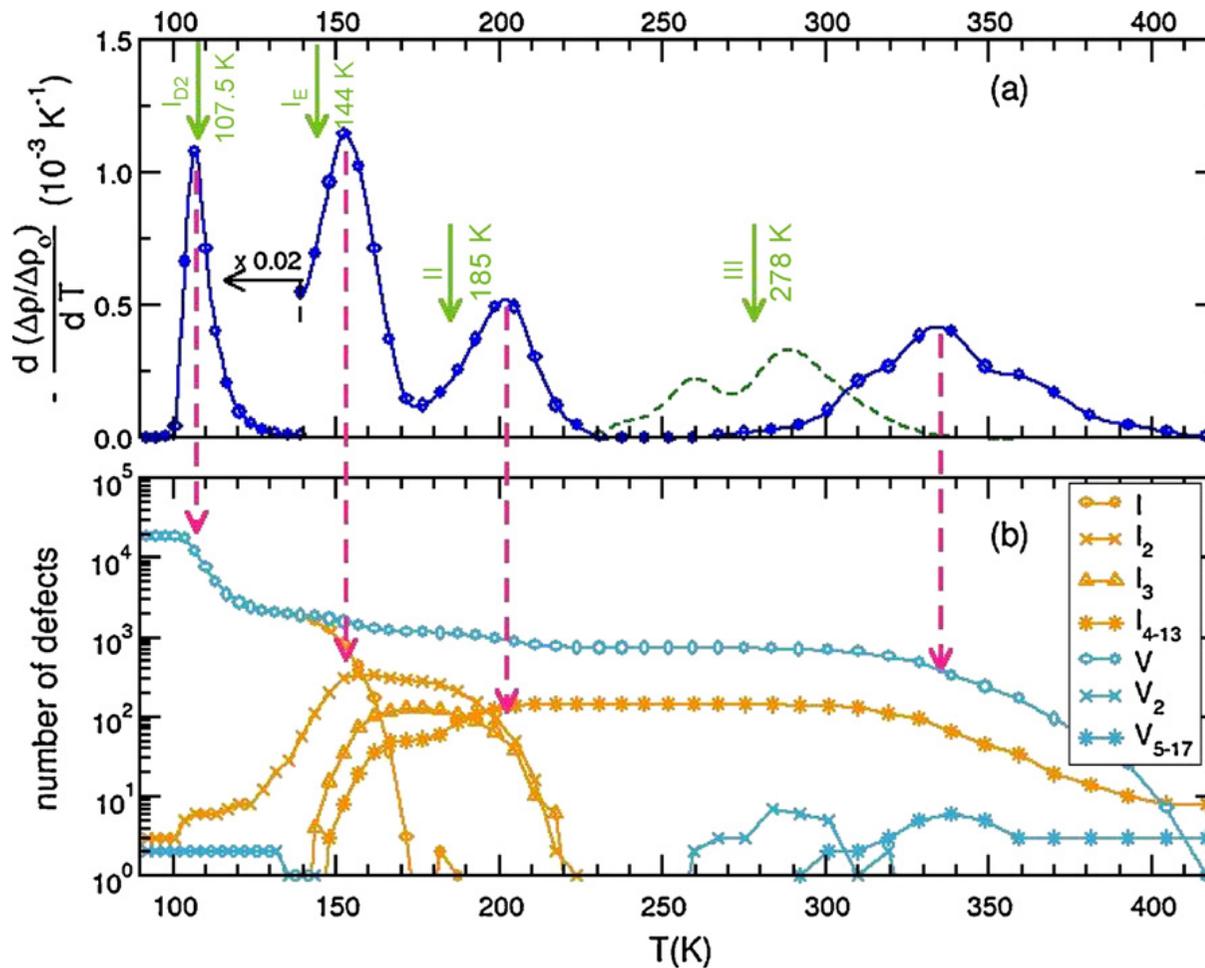
Figure 3 Migration of vacancy-type defects. Energy variation along the most favourable migration pathways for V_1 , V_2 , V_3 and V_4 . The vacant sites are represented by blue cubes, and their migration jumps by magenta arrows. The energy scales are the same as in Fig. 2 for comparison.

Annealing of electrical resistivity



Annealing stages and defect reactions in pure copper after electron irradiation; i_1 and v_1 denote single interstitials and vacancies, respectively, i_2 and v_2 di-interstitials and divacancies (Schilling et al. 1975.)

The time evolution of the microstructure



Recovery of electrical resistivity and defect population during isochronal annealing after electron irradiation:

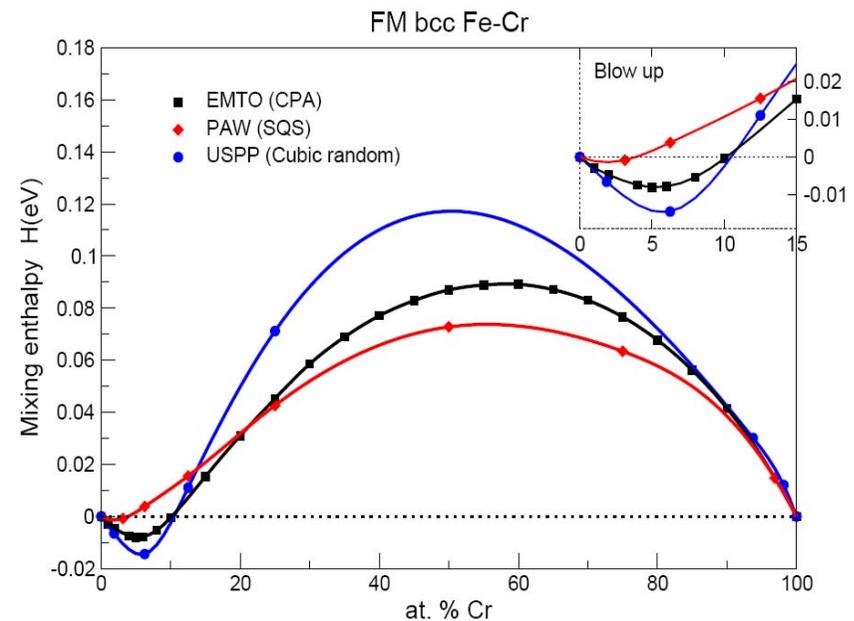
- temperature derivative of the simulated resistivity recovery stages. Green arrows indicate the positions of experimental peaks.
- Associated defect population evolution for 16×10^9 atoms. The magenta dashed lines indicate the position of the peak in the simulation.

The Fe-Cr system



- FeCr is the base alloy of the 9-12% Cr ferritic/martensitic steels that are the main candidates for structural components of fusion and Gen IV fission reactors

In Fe Cr alloys, antiferromagnetic Cr is added to a matrix of ferromagnetic Fe. This defines the thermodynamic behavior of the alloy: there is a tendency to short range order for low Cr and to clustering for higher Cr contents, resulting in an anomaly of the heat of formation

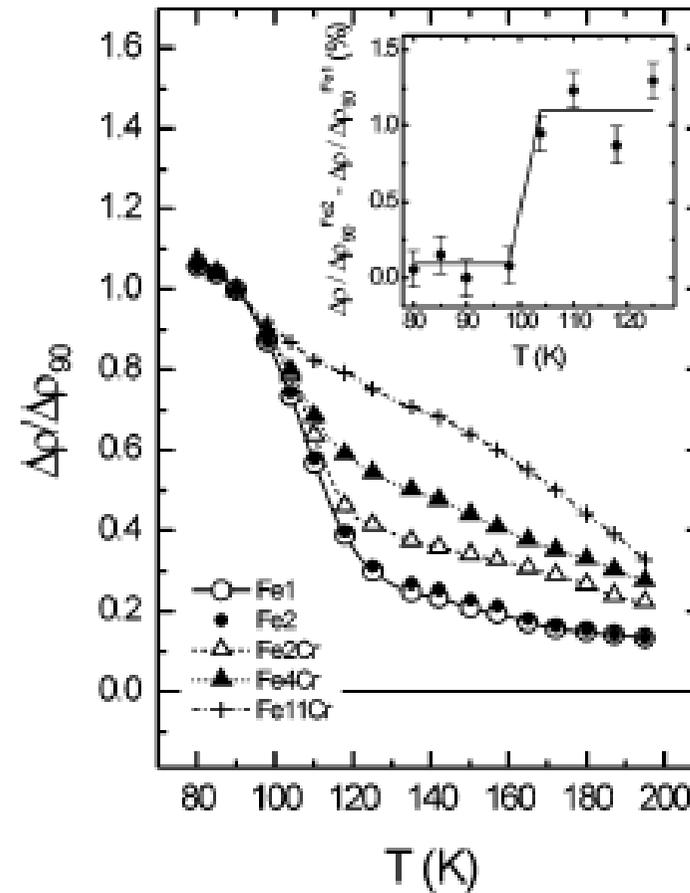


The case of Fe-Cr alloys

A number of measurements have been performed for dilute alloys (Maury (1987), Benkaddour (1987) Nikolaev (1997, 1999), Abe (1999))

The features associated with the annealing of RR of concentrated alloys in stage I can be explained by the formation, migration and trapping of mixed Fe-Cr dumbbells at a second Cr atom

Irradiation that produces small cascades stimulates the formation of mobile di-interstitials probably formed at the cascade quenching



Defect accumulation in FeCr

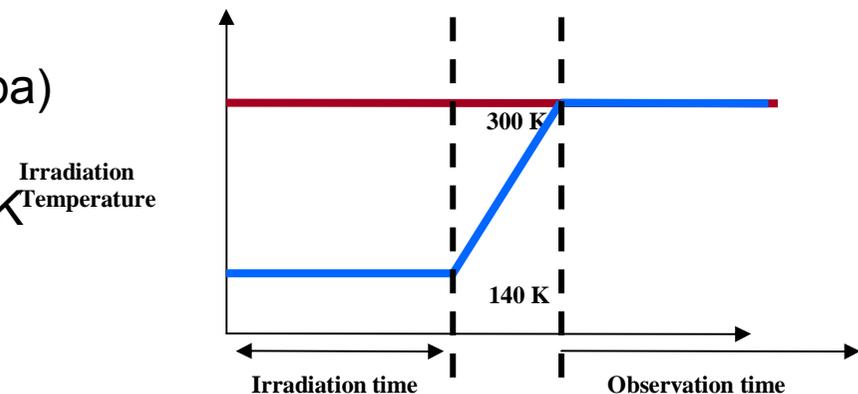


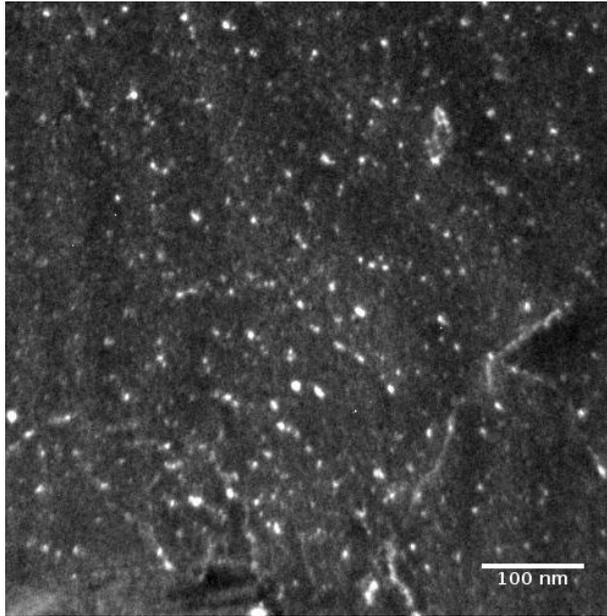
- The mobile [111] cluster that will result from the [110] clustering have a high mobility and are expected to rapidly move out from the (thin) sample observed in TEM. But as indicated in the talks of Malerba and Samaras the clusters are pinned by the presence of Cr. It is therefore interesting to see their behaviour after irradiation of alloys with increasing Cr concentration.
- Irradiation experiments have been performed in a series of FeCr alloys with increasing Cr content

Experimental methodology



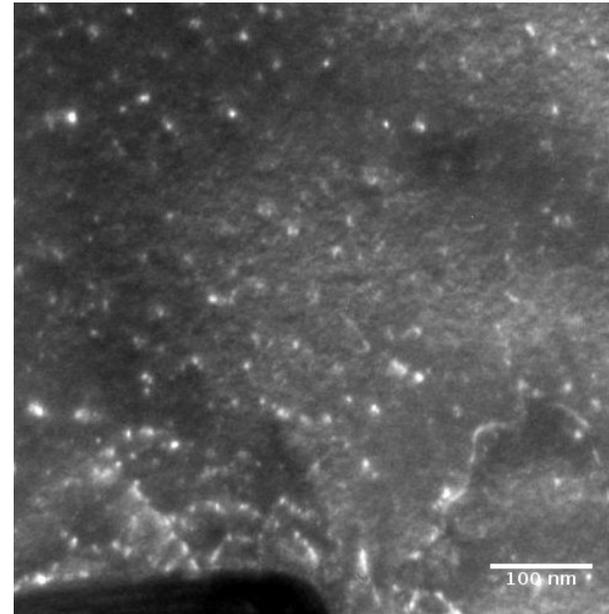
- Materials:
 - Alloys prepared from 99,99 Fe and Cr by arc melting in a pure Ar atmosphere. Nominal concentrations: 1, **2**, 3, **5**, 7, 10, **12** y 15% at Cr.
 - Effect of Cr content
- Ion irradiation conditions
 - Irradiations performed in the AIM facility of the Forschungszentrum Dresden-Rossendorf (FZD)*:
 - Fe⁺ 150 keV
 - Dose: 2×10^{18} ions/m² (~0.8 dpa)
 - Flux: 2×10^{15} ions/m² s
 - Temperatures: 140K and 300K
- TEM observation at RT
 - JEOL 200 keV



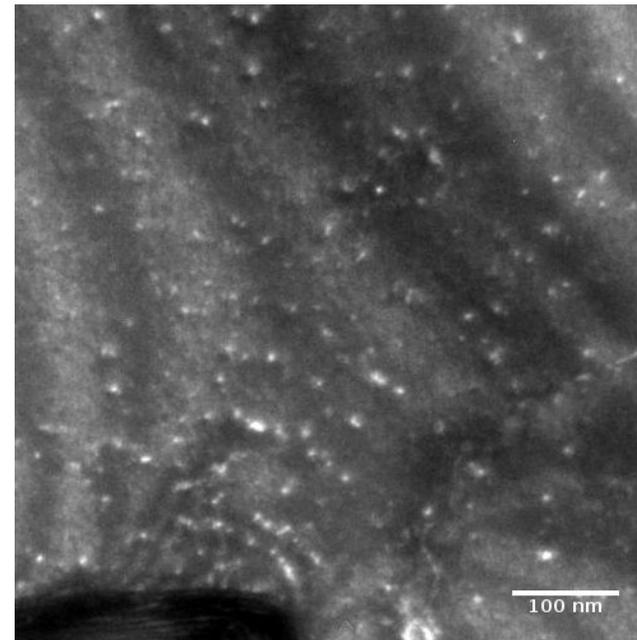


Fe12Cr 300K
g=200
1/3 of $\langle 100 \rangle$ are visible
All the $\langle 111 \rangle$ are visible

Fe12Cr 300K g=10-1
2/3 of $\langle 100 \rangle$ are visible
1/2 of $\langle 111 \rangle$ are visible

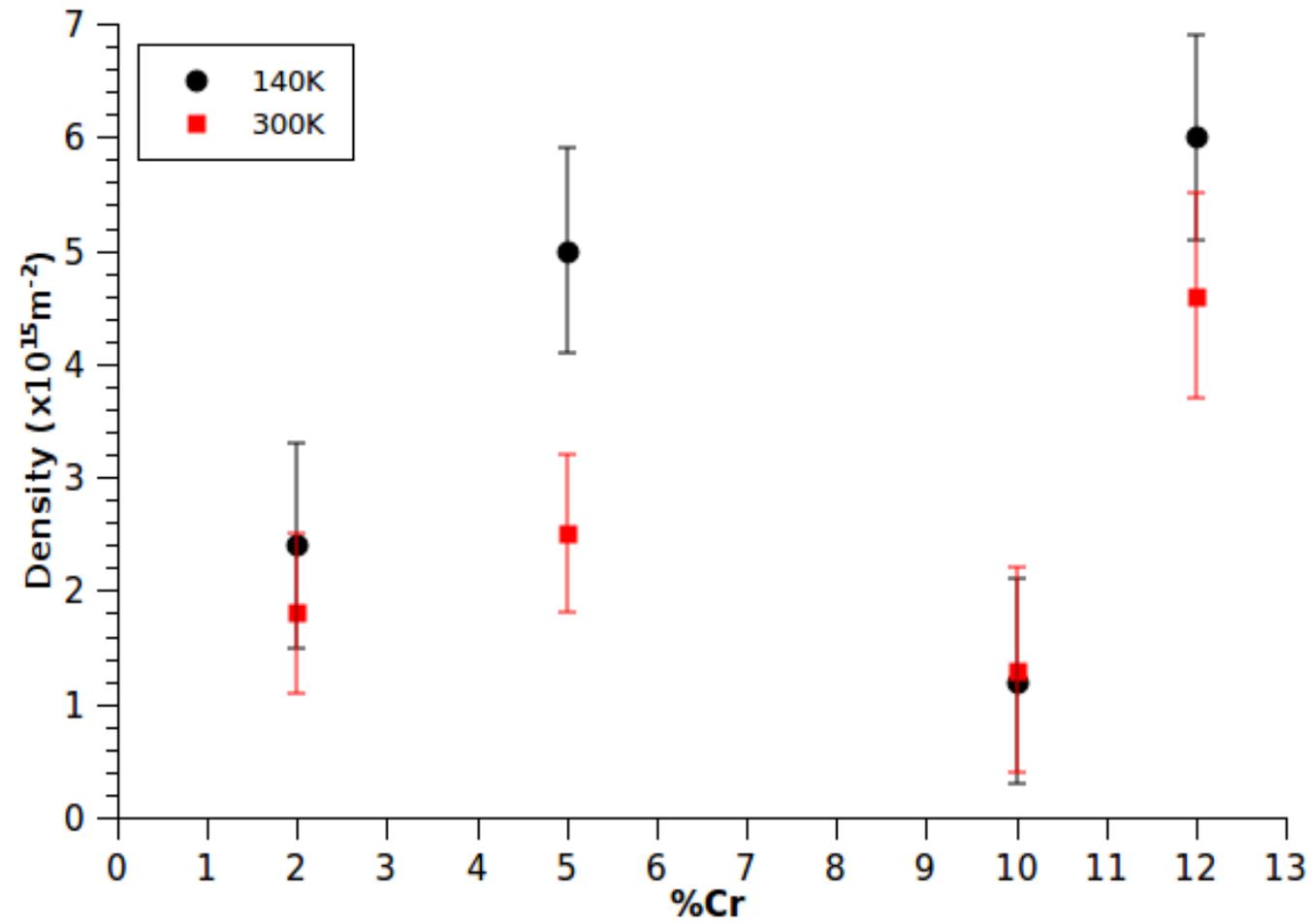


Fe12Cr 300K g=0-11
2/3 of $\langle 100 \rangle$ are visible
1/2 of $\langle 111 \rangle$ are visible



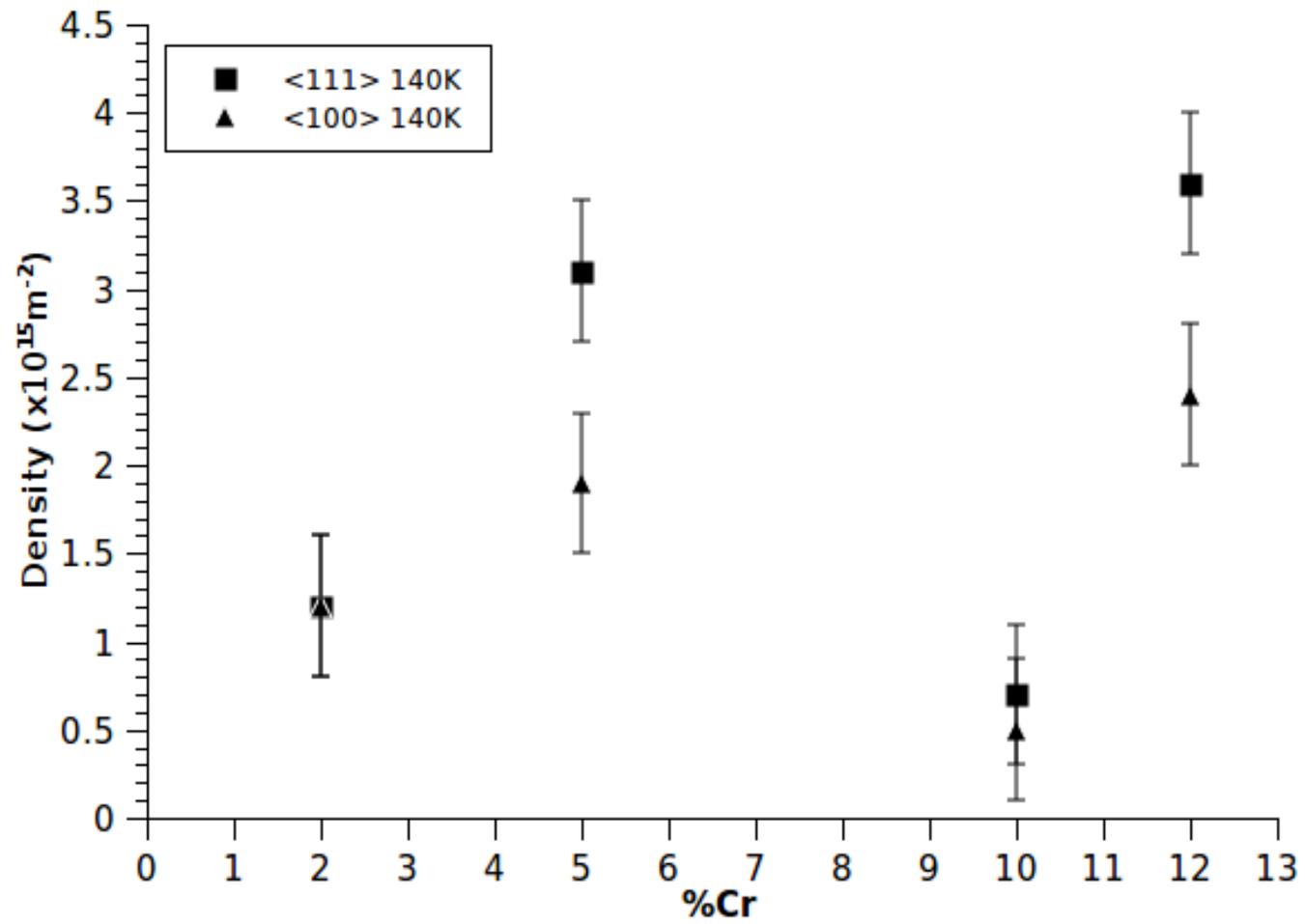


density of loops



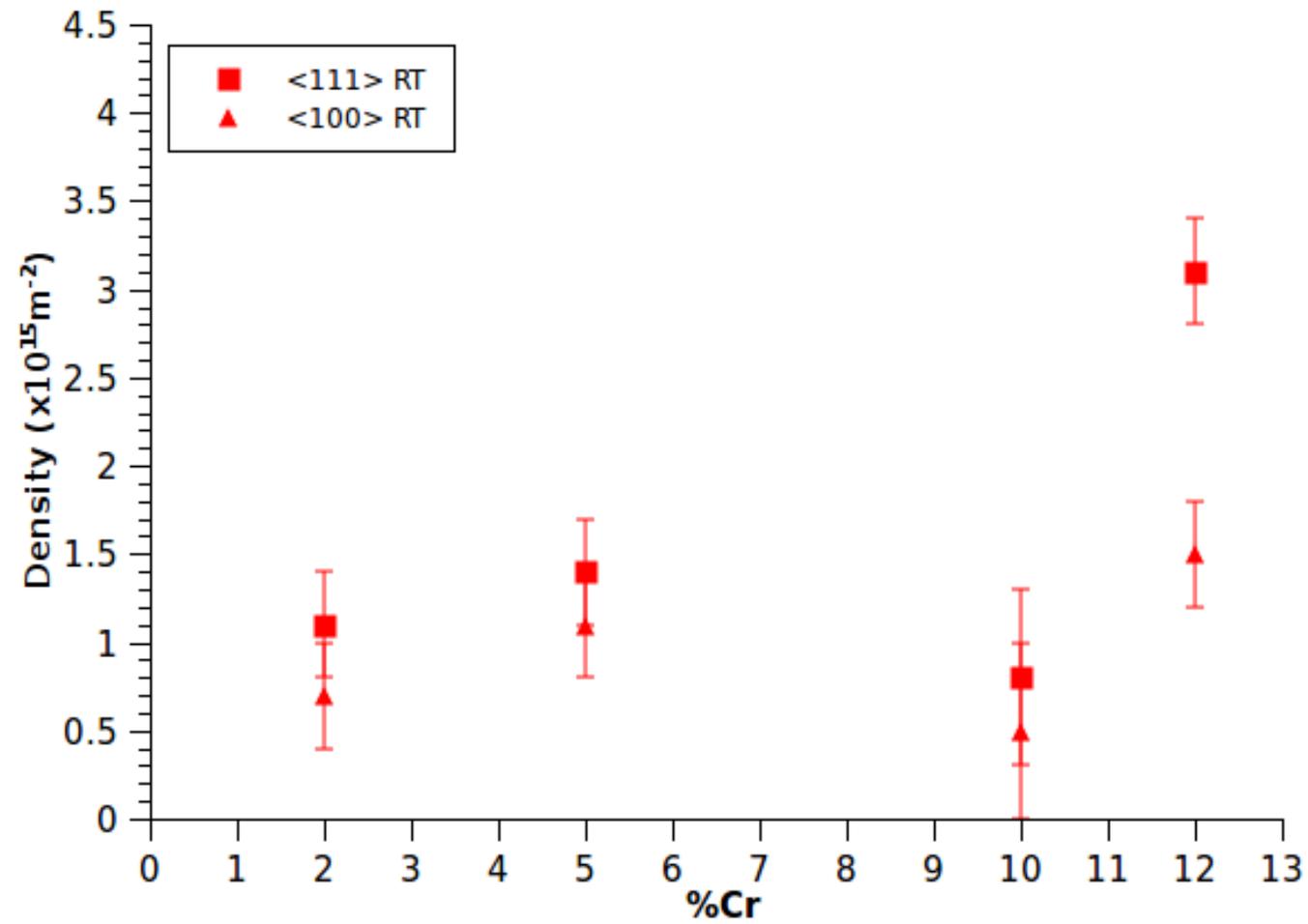


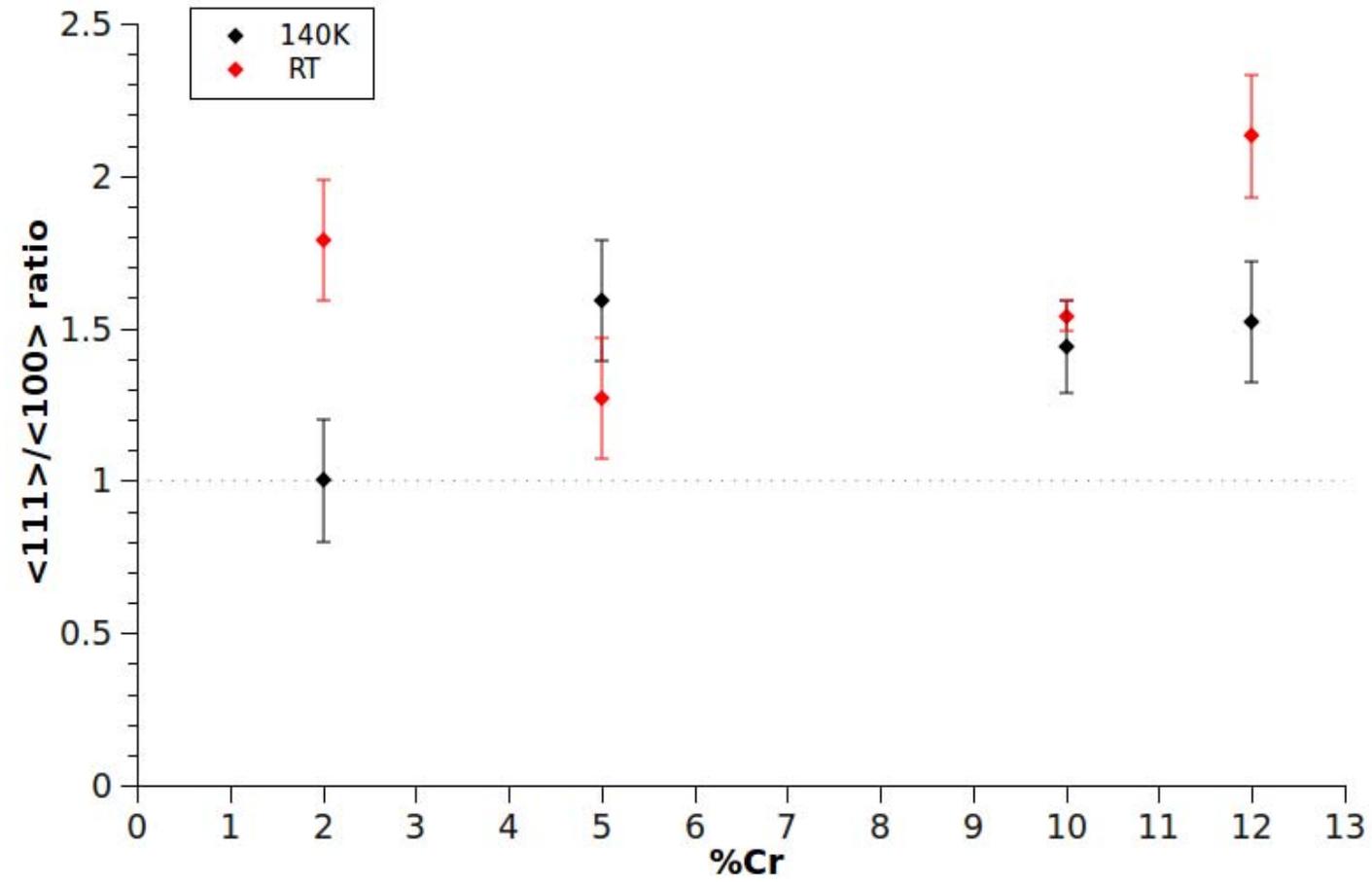
$T_{irr} = 140K$





$T_{irr} = RT$

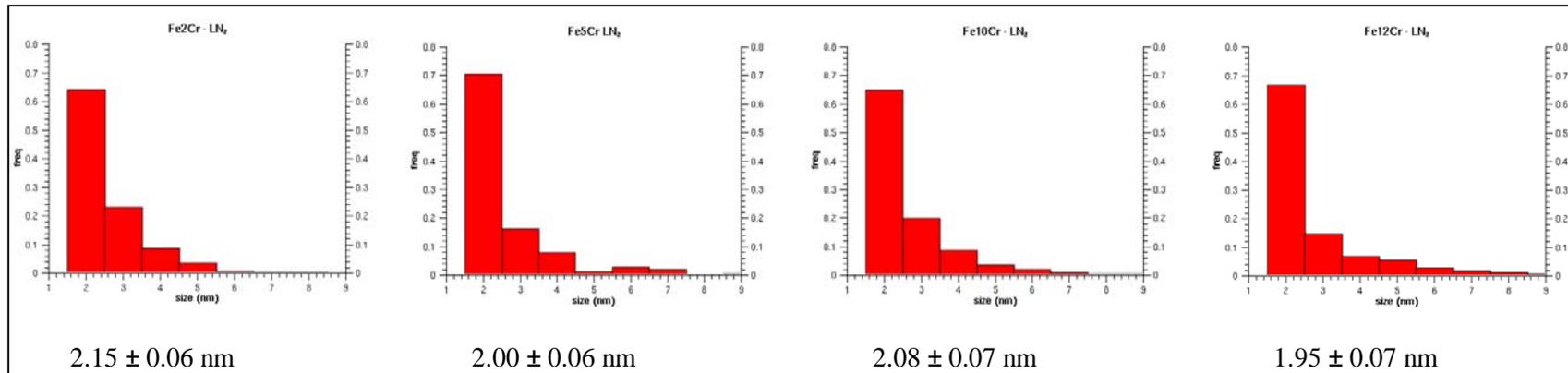




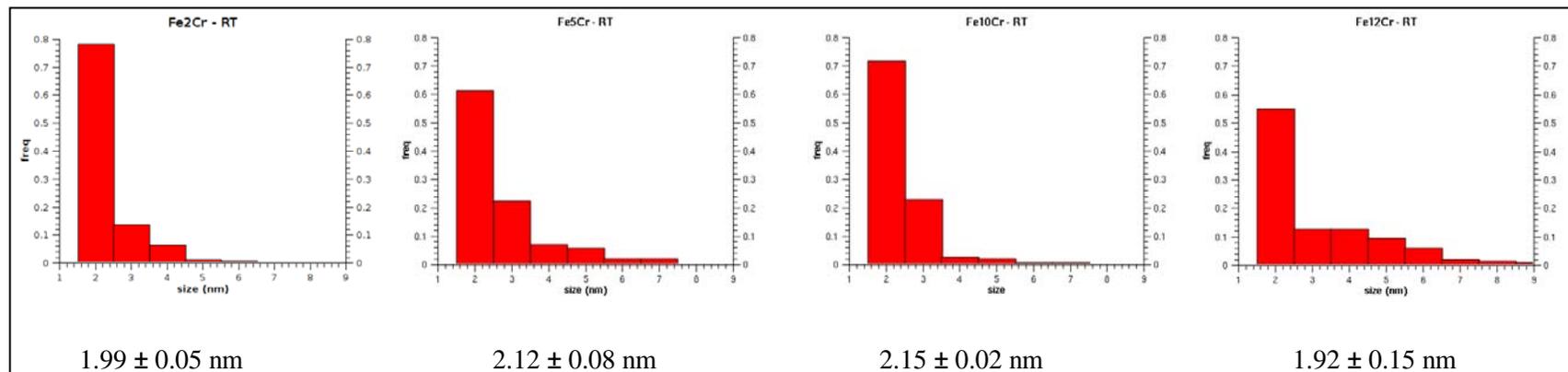
Effect on size distribution



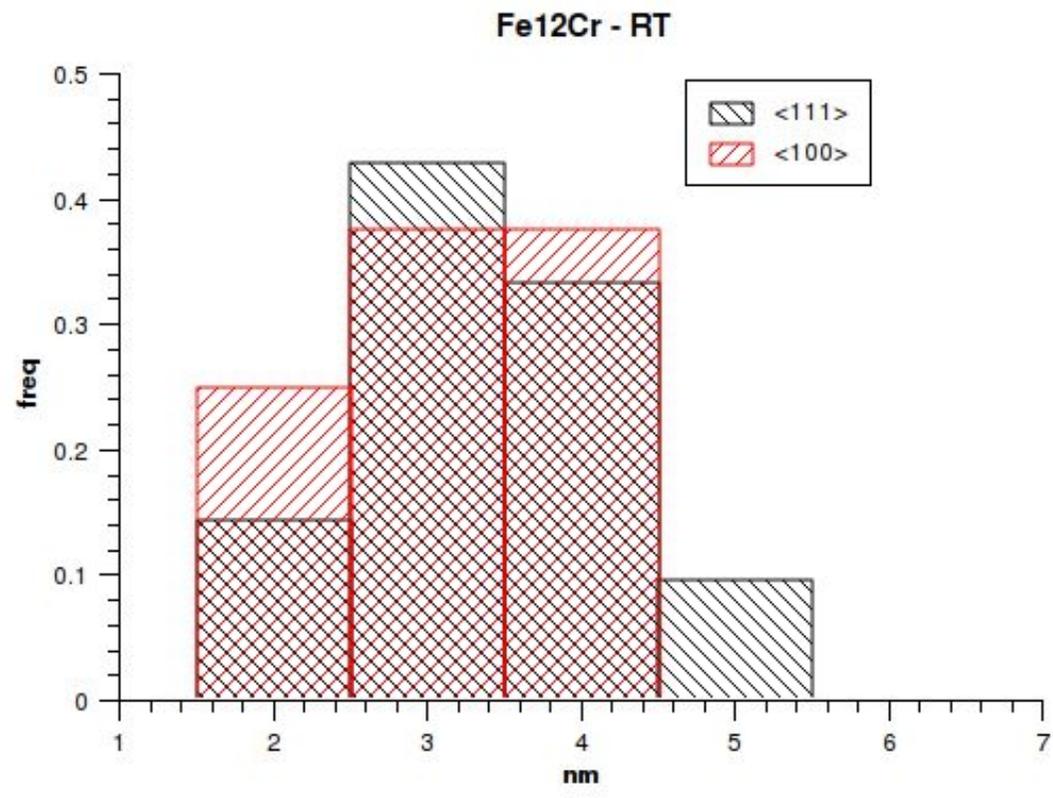
T irr: 140 K



T irr: 300 K



- No effect on average size is observed for different Cr content or irradiation temperature
- Maximum size seems to increase for higher Cr content





The question that arises from these results is the presence of sessile $[100]$ loops. They had been observed previously in Fe and steels

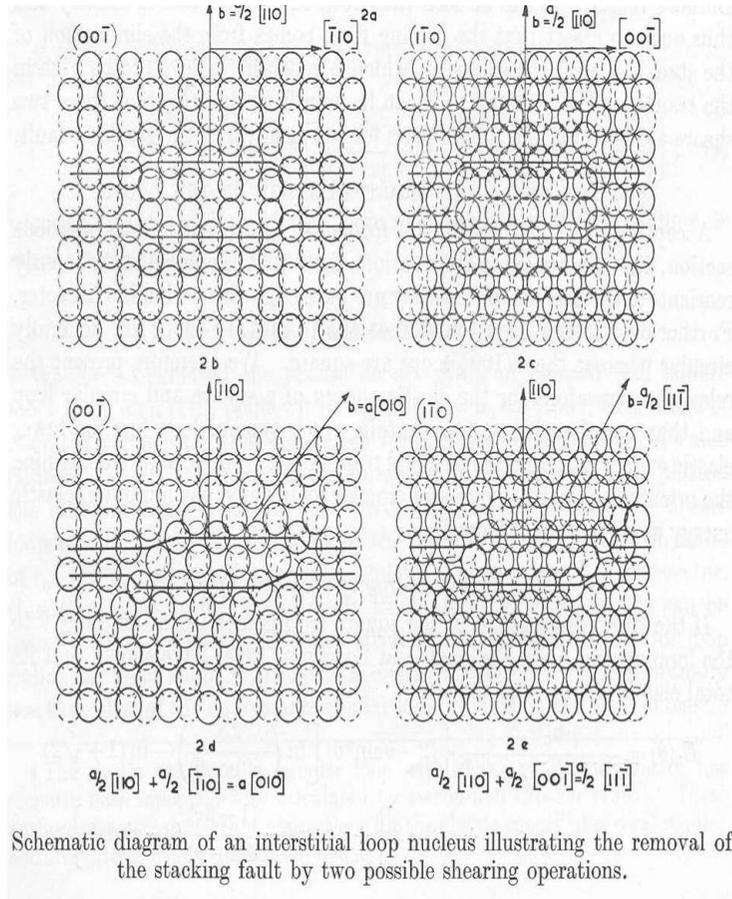
But their behavior in the present set of results is difficult to understand.

Two models have been proposed to explain the formation from $1/2 [111]$ loops, but none of the two would explain the presence of these loops at low irradiation temperatures (140 K)

It is thought at present that the $[100]$ Loops originate from the evolution of the cascade itself

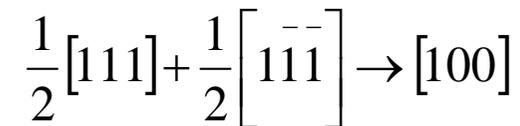
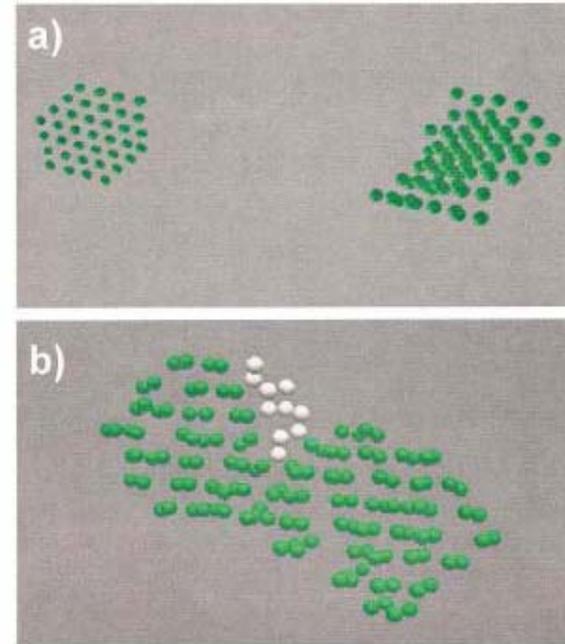
Mechanism of formation of (100) loops

Shear of <110> loops



Eyre & Bullough (1965)

Interaction between <111> loops



Marian et al (2002)