



**The Abdus Salam
International Centre for Theoretical Physics**



2141-24

**Joint ICTP-IAEA Workshop on Nuclear Reaction Data for Advanced
Reactor Technologies**

3 - 14 May 2010

Introduction to Nuclear Model Code TALYS

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France*

Introduction
to the nuclear model code
TALYS

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&
*S. Goriely*³



www.talys.eu

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2 Nuclear Research and Consultancy Group, Petten, The Netherlands

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Nuclear data needed for

Understanding basic reaction mechanism between particles and nuclei
Astrophysical applications (Age of the Galaxy, element abundances ...)
Existing or future nuclear reactor simulations
Medical applications, oil well logging, waste transmutation ...

But

Finite number of experimental data (price, safety or counting rates)
Complete measurements restricted to low energies (< 1 MeV)



**Predictive & Robust Nuclear models
(codes) are essential**

- ➔ **General features of TALYS**
- ➔ **Models implemented in TALYS**
- ➔ **Some TALYS results**
- ➔ **Conclusions and prospects**

General features of TALYS



The TALYS team

Authors

Arjan Koning, NRG Petten

Stéphane Hilaire, CEA-DIF

Marieke Duijvestijn, NRG Petten

Computational & Theoretical support, validation

Several members of CEA-DIF

Stéphane Goriely, ULB

Emmeric Dupont, CEA Cadarache

Jura Kopecky, JUKO Research

Robin Forrest, UKAEA

Current version

TALYS-1.2 at www.talys.eu

- **Date of birth : 1998**
- **Fortran 77**
- **50000 lines (+ 20000 lines of ECIS)**
- **Modern programming**
 - modular (270 subroutines)
 - descriptive variable names & well commented (45% of lines are comments)
 - transparent programming
- **Very extensive input handling and checking**
 - Flexible (from default to multi-parameter adjustment : > 200 keywords)
 - Random input to check stability
 - Drip-line to drip-line calculations
- **> 380 pages manual**
- **Compiled and tested with f77, f90, f95, ... over several OS**
- **Externally driven for :**
 - ENDF formatting
 - Random input to check stability

ALICE – LLNL – 1974 – Blann

(Mc-)GNASH – LANL – 1977 – Young, Arthur & Chadwick

TNG – ORNL – 1980 – Fu

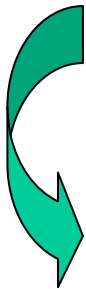
STAPRE – Univ. Vienna – 1980 – Uhl

UNF,MEND – CIAE, Nanking Univ. – 1985 – Cai, Zhang

EXIFON – Univ. Dresden – 1989 – Kalka

EMPIRE – ENEA/IAEA/BNL – 1980 – Herman

TALYS – NRG/CEA – 1998 – Koning, Hilaire & Duijvestijn



**Modern computers (i.e. High speed & Big memory)
already available when TALYS development started**

What TALYS does !

- Simulates a nuclear reaction between a projectile and a target

projectiles : n,p,d,t,³he, ⁴he

target : $3 \leq Z \leq 110$ or $5 \leq A \leq 339$

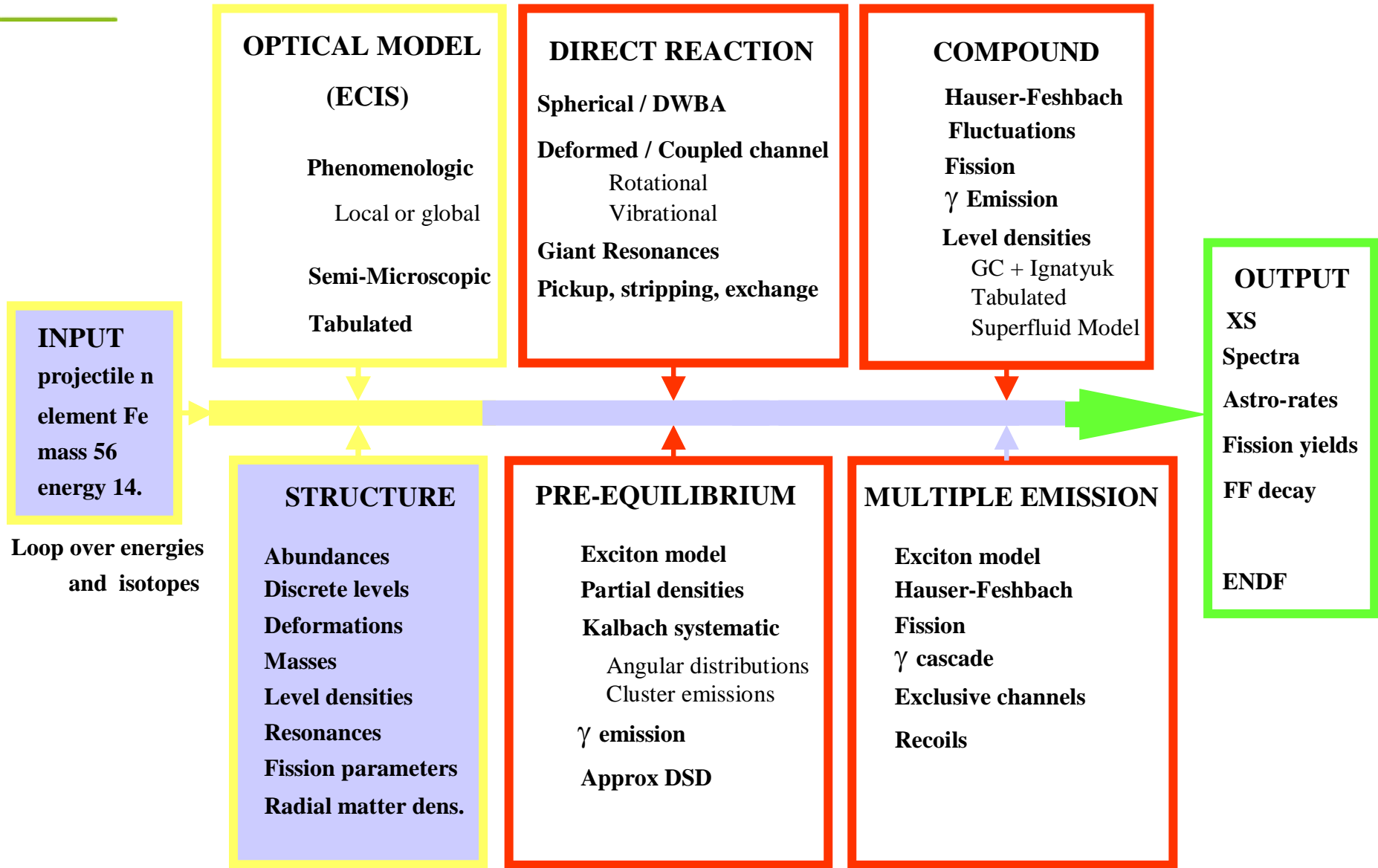
- Projectile energy from 1keV up to 200 MeV

- TALYS mantra : “ Completeness then quality ”



- Optical, pre-equilibrium and statistical model implemented with sets of default parameters
- All opened channels smoothly described
- Possibilities for future improvements anticipated
 - Level densities (stored and interpolated)
 - Parity dependence
- **Still under development (improvement)**

How TALYS works !

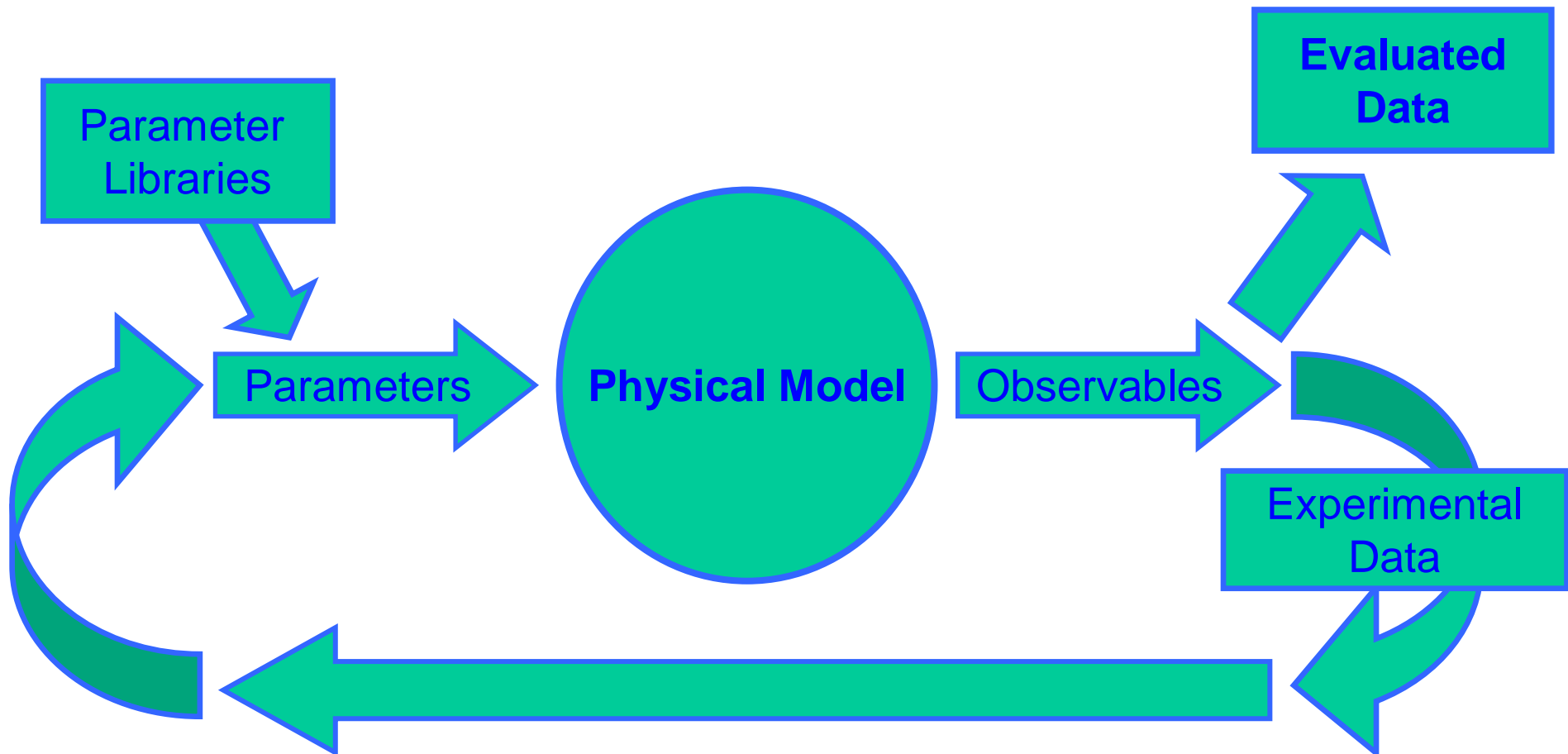


What TALYS yields !

- **Cross sections : total, reaction, elastic (shape & compound), inelastic (per level & total) and all opened channels.**
- **Elastic and inelastic angular distribution**
- **Exclusive reaction channels : xs, spectra & ddx**
- **Exclusive discrete and continuum γ -ray production**
- **Photonuclear reactions & reactions on isomeric targets**
- **Fission cross sections and fission yields**
- **Residuals production and recoils**
- **Total particle production : xs, spectra & ddx**
- **Extrapolation down to thermal energy**
- **Stellar reaction rates**
- **Fission fragment decay**
- **Level density tables**

Nuclear reaction modeling

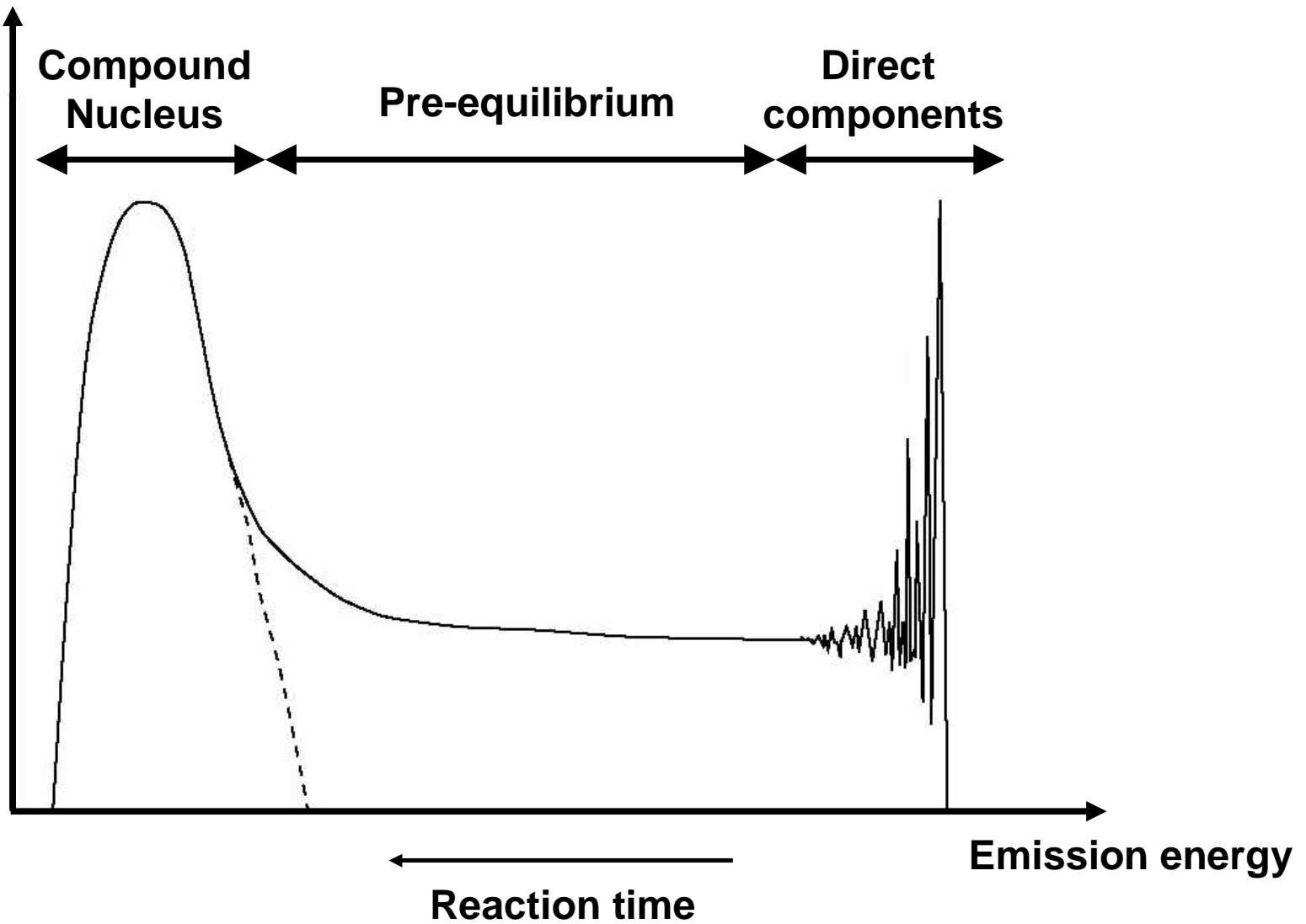
Method which consists in using a physical model (together with sets of parameters) to calculate evaluated data.



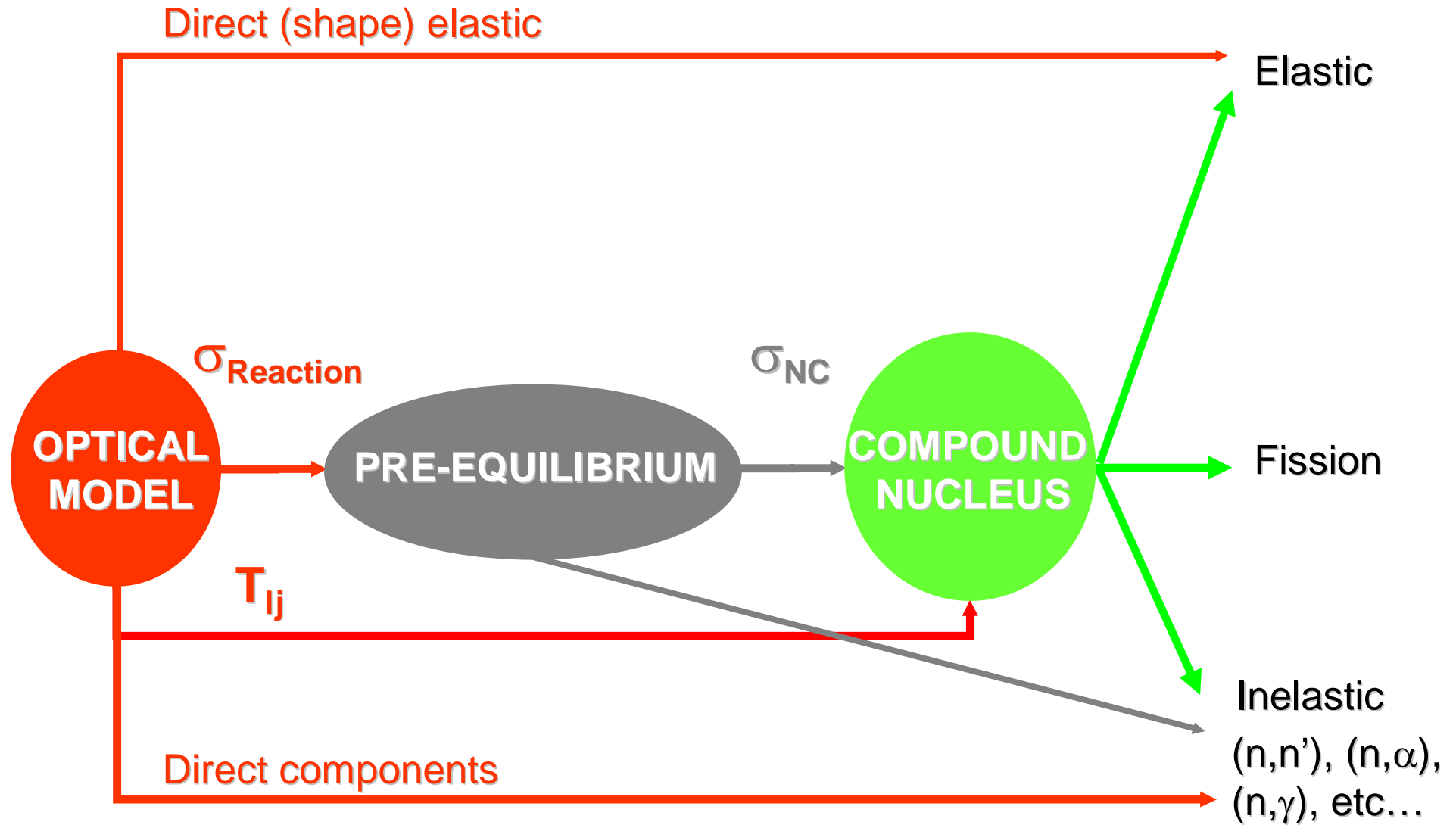
Models implemented in TALYS

Models sequence

$$d^2\sigma / d\Omega dE$$



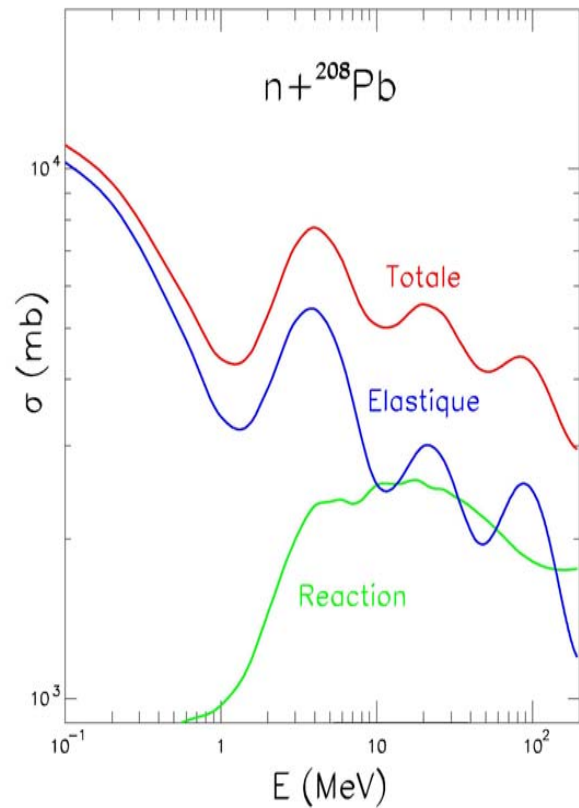
Models sequence



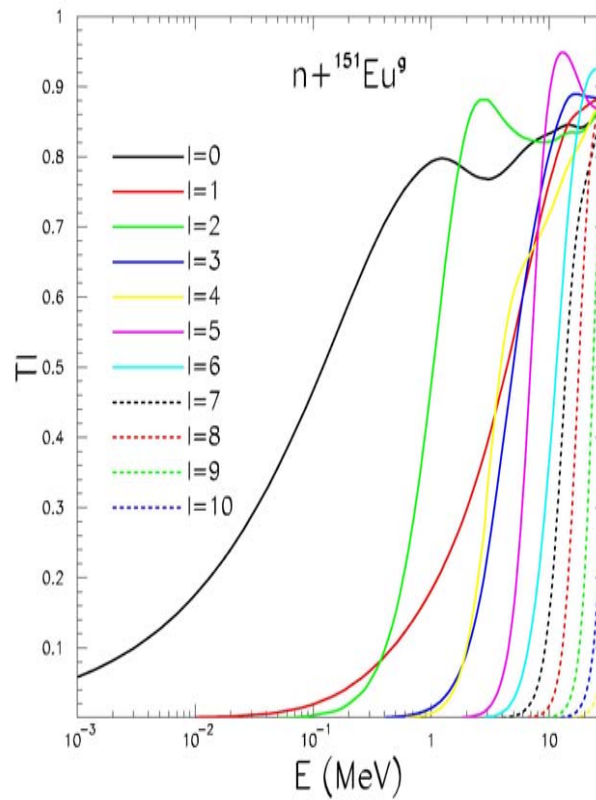
The Optical model

This model yields :

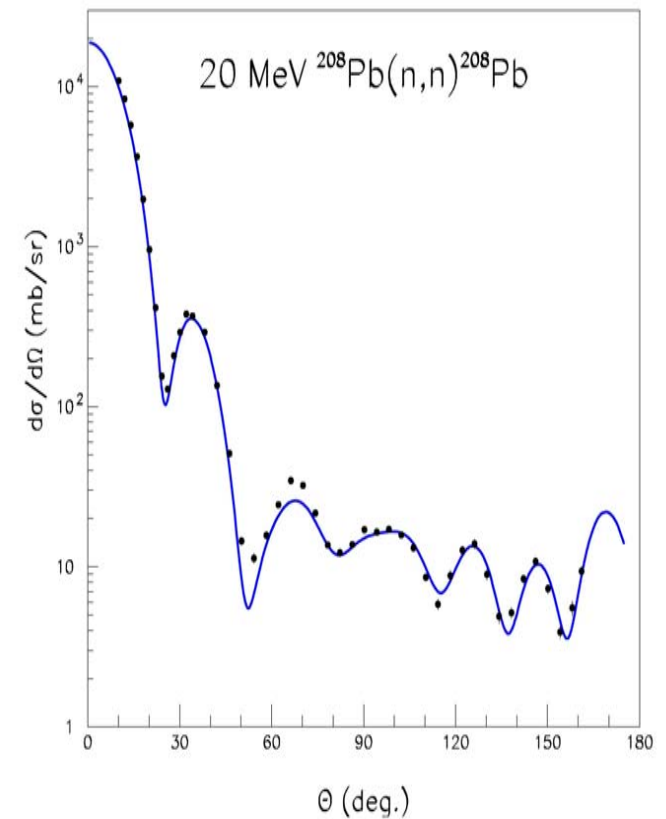
Integrated cross sections



Transmission coefficients



Angular distributions



The Optical model

Direct interaction of a projectile with a target nucleus considered as a whole
 Quantum model → Schrödinger equation

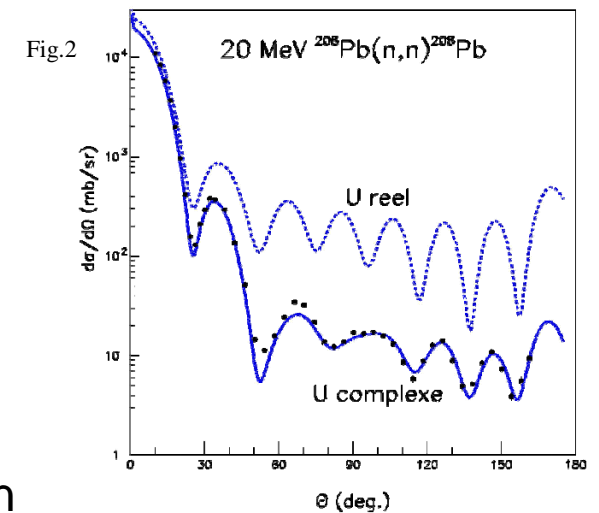
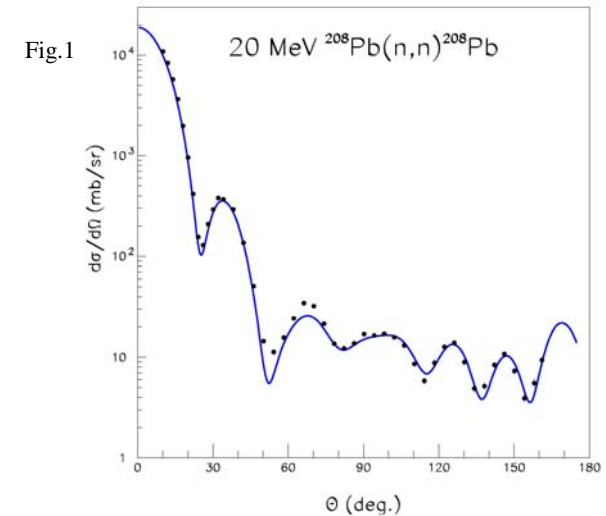
$$\left(-\frac{\hbar^2}{2\mu} \nabla^2 + \mathbf{U} - E \right) \Psi = 0$$

Complex potential:

$$\mathbf{U} = \mathbf{V} + i\mathbf{W}$$

Refraction

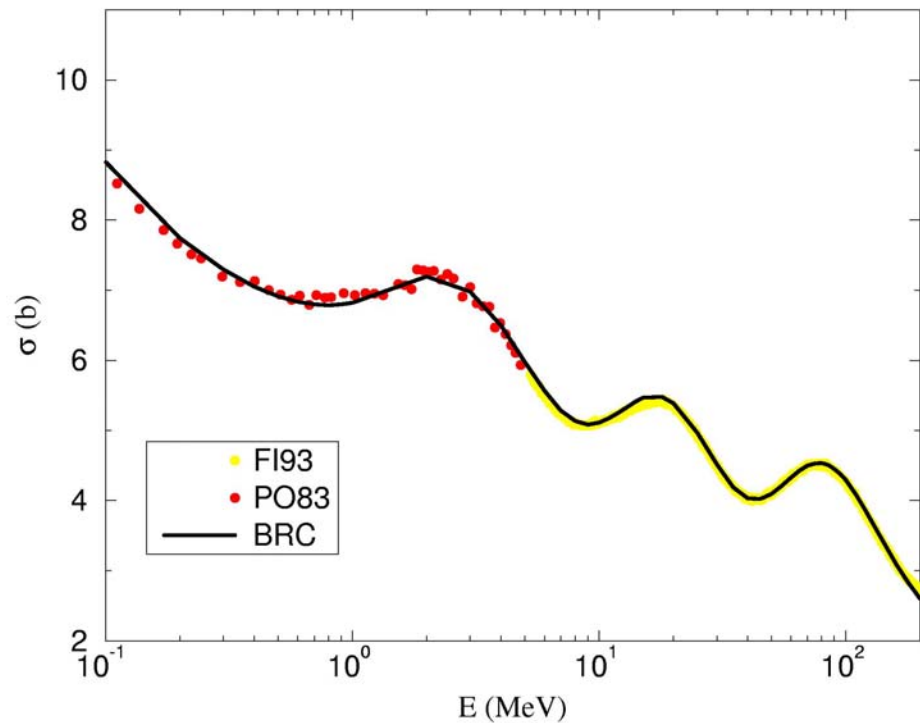
Absorption



Approaches implemented in TALYS

Phenomenologic

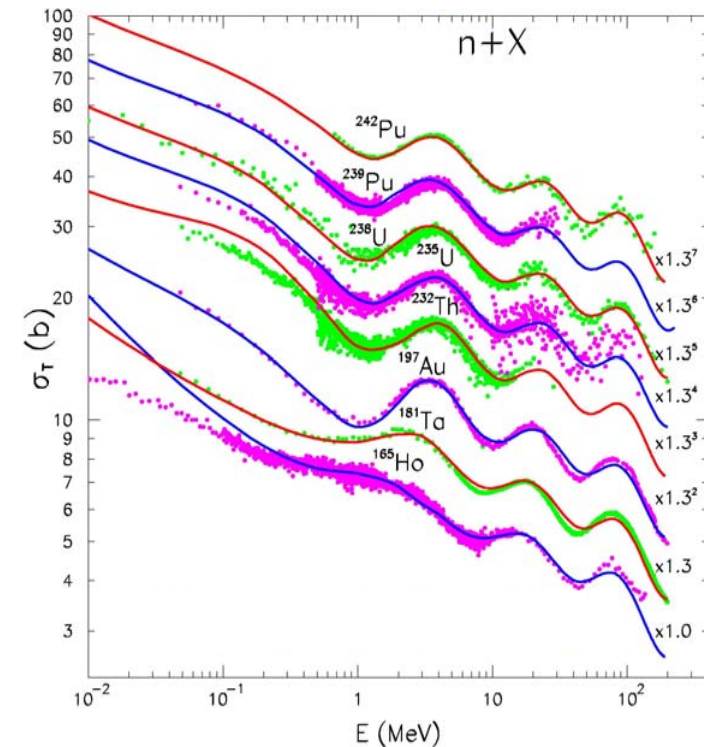
Adjusted parameters
Weak predictive power
Very precise ($\approx 1\%$)
Important work



Semi-microscopic

No adjustable parameters
Usable without exp. data
Less precise ($\approx 5-10\%$)
Quasi-automated

Total cross sections



Phenomenologic :

- Koning-Delaroche for non-fissile nuclei
- Soukhovitsky for fissile nuclei
- Other implementations easy (e.g. α)
- Tabulation possible

Semi-microscopic

- JLM approach based on matter densities
 - \Rightarrow any type of matter density can be used
(Skyrme and Gogny already available)

\Rightarrow OMP calculations essentially performed with ECIS

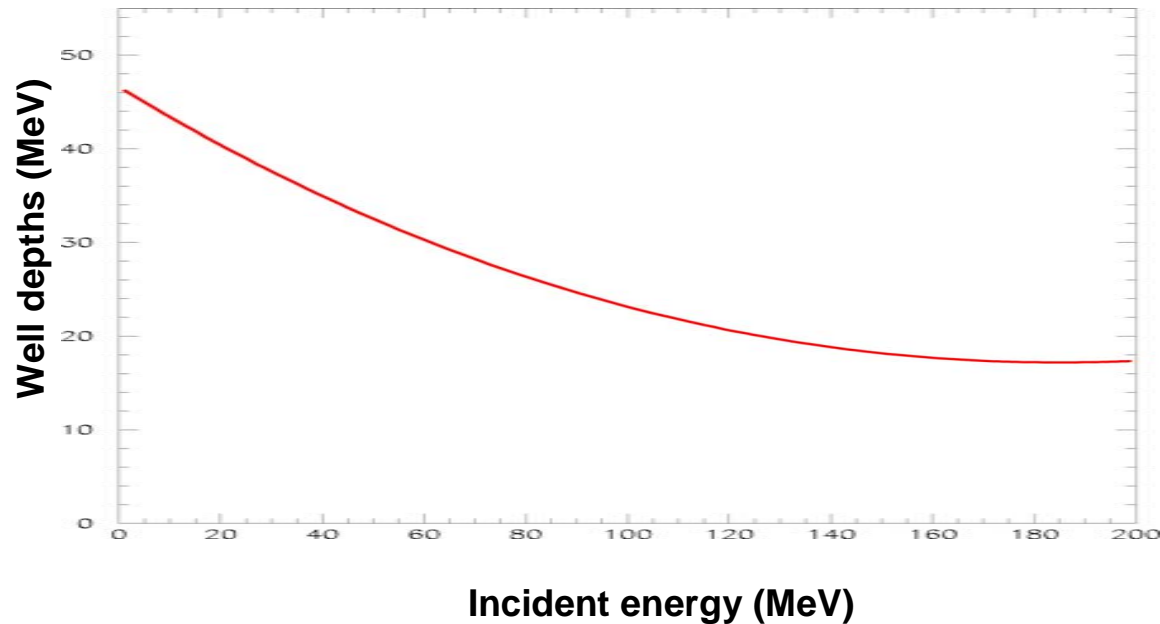


Phenomenological OMP

$$U(r,E) = V(E,r) + i W(E,r)$$

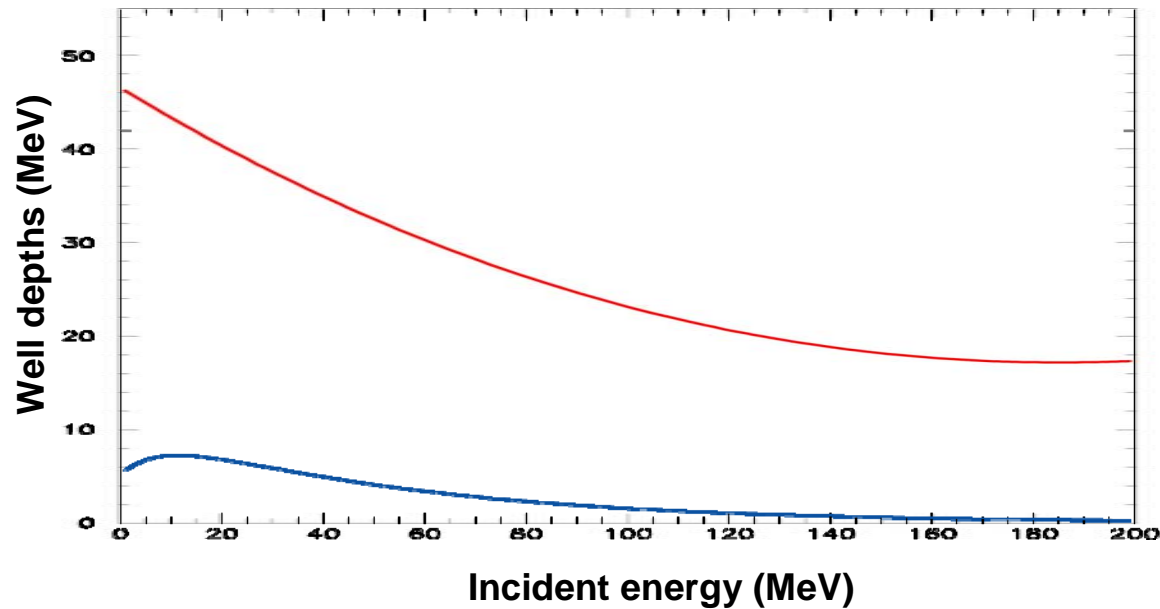
Phenomenological OMP

$$U(r,E) = \begin{bmatrix} V_V(E) f(r, R_V, a_V) + V_S(E) g(r, R_S, a_S) \\ + i \begin{bmatrix} W_V(E) f(r, R_V, a_V) + W_S(E) g(r, R_S, a_S) \end{bmatrix} \end{bmatrix}$$



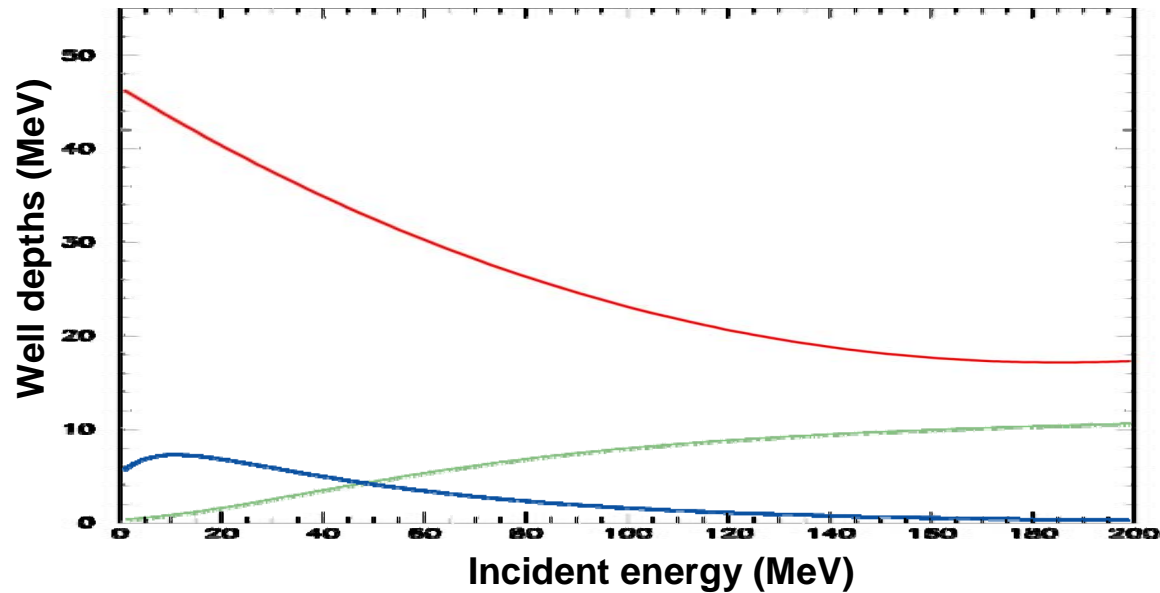
Phenomenological OMP

$$U(r,E) = \begin{bmatrix} V_V(E) f(r, R_V, a_V) + V_S(E) g(r, R_S, a_S) \\ + i \left[W_V(E) f(r, R_V, a_V) + W_S(E) g(r, R_S, a_S) \right] \end{bmatrix}$$



Phenomenological OMP

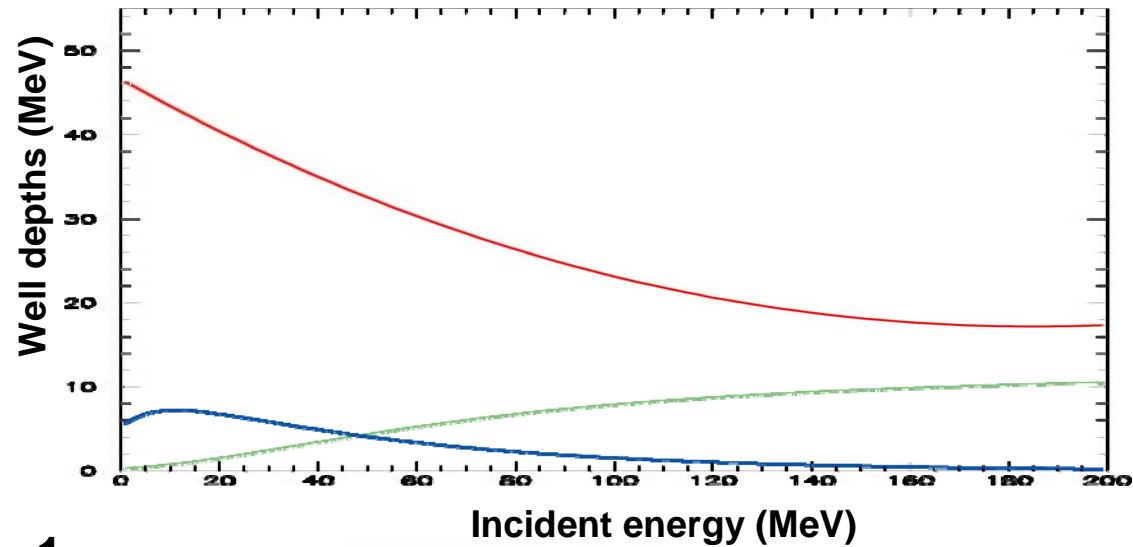
$$U(r,E) = \begin{bmatrix} V_V(E) f(r, R_V, a_V) + V_S(E) g(r, R_S, a_S) \\ + i \begin{bmatrix} W_V(E) f(r, R_V, a_V) + W_S(E) g(r, R_S, a_S) \end{bmatrix} \end{bmatrix}$$



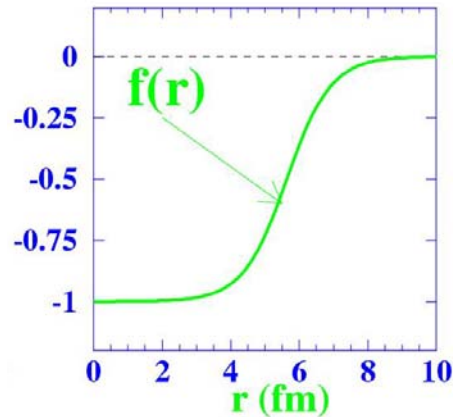


Phenomenological OMP

$$U(r,E) = \left[V_V(E) f(r, R_V, a_V) + V_S(E) g(r, R_S, a_S) \right] + i \left[W_V(E) f(r, R_V, a_V) + W_S(E) g(r, R_S, a_S) \right]$$



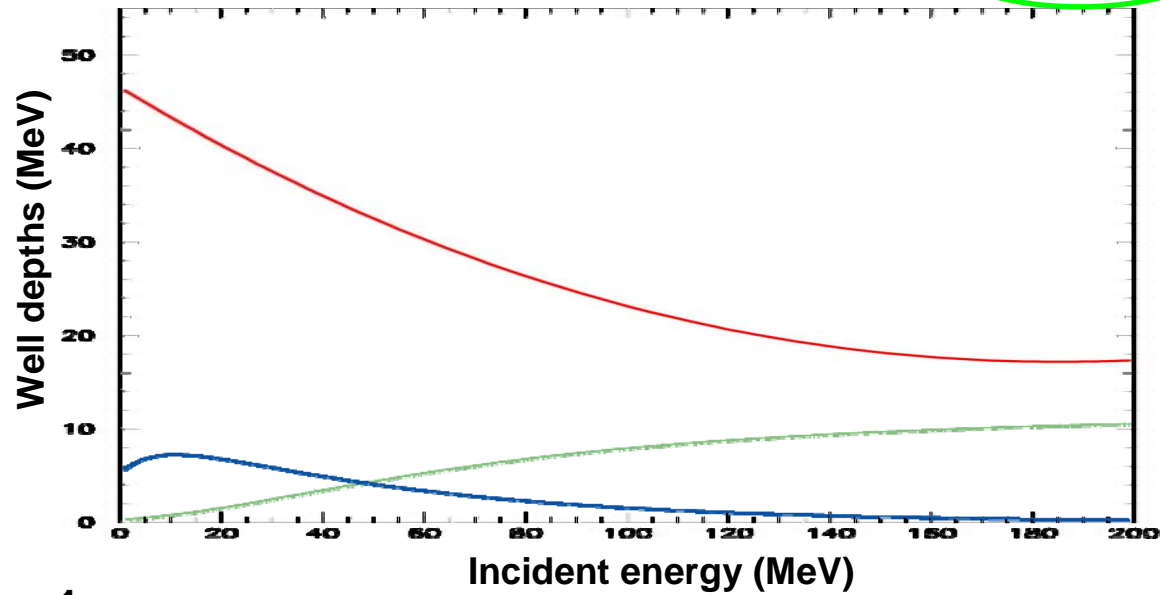
$$f(r,R,a) = \frac{-1}{1 + \exp((r-R)/a)}$$



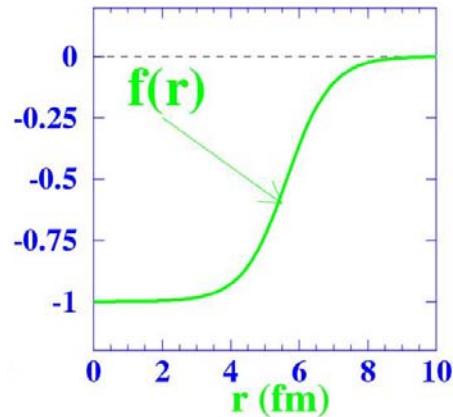


Phenomenological OMP

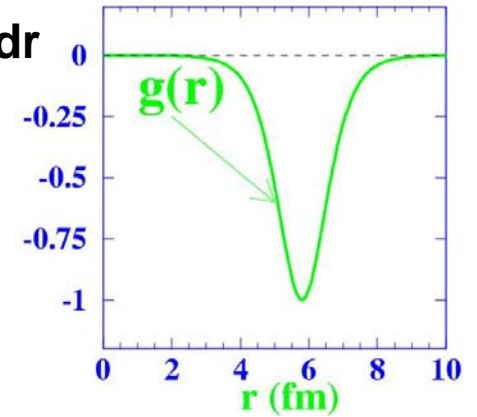
$$U(r,E) = \begin{bmatrix} V_V(E) f(r, R_V, a_V) + V_S(E) g(r, R_S, a_S) \\ + i \begin{bmatrix} W_V(E) f(r, R_V, a_V) + W_S(E) g(r, R_S, a_S) \end{bmatrix} \end{bmatrix}$$



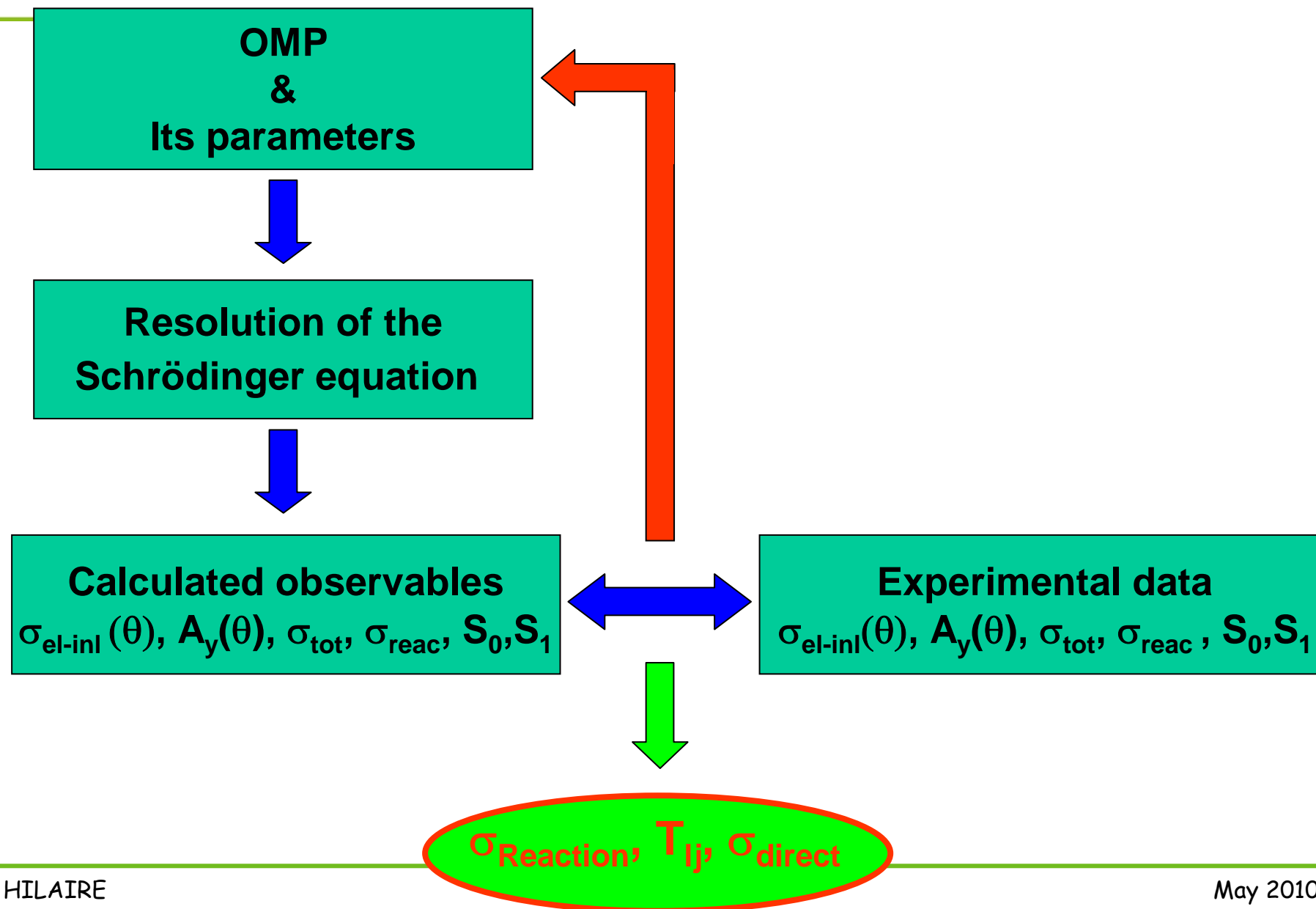
$$f(r,R,a) = \frac{-1}{1 + \exp((r-R)/a)}$$



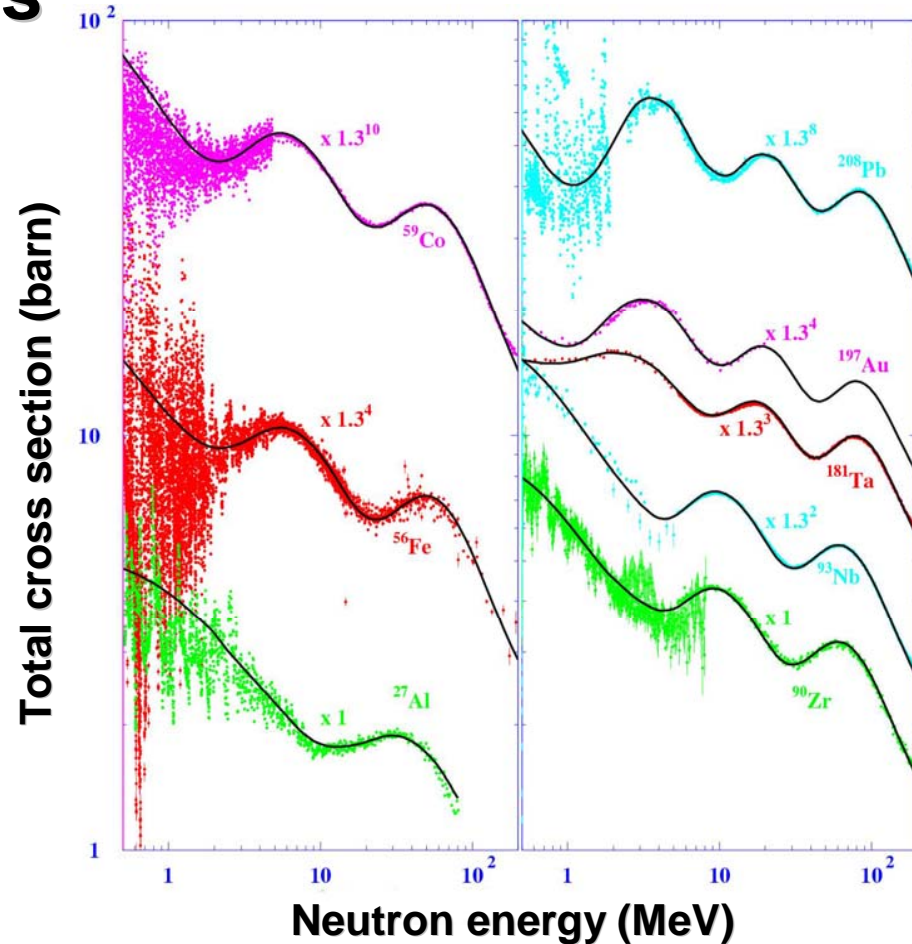
$$g(r,R,a) = -df/dr$$



Phenomenological OMP

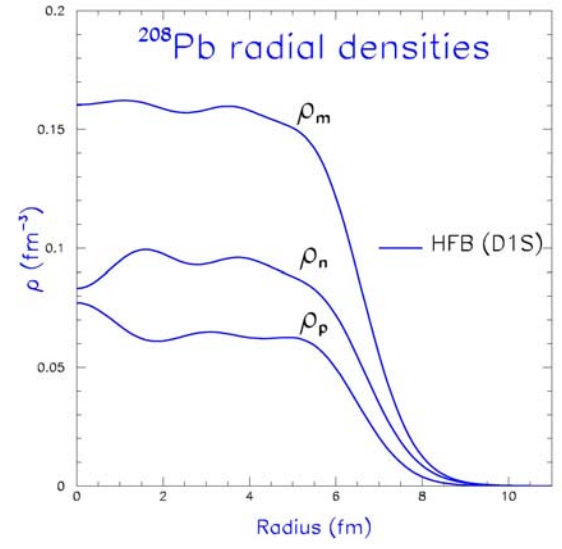
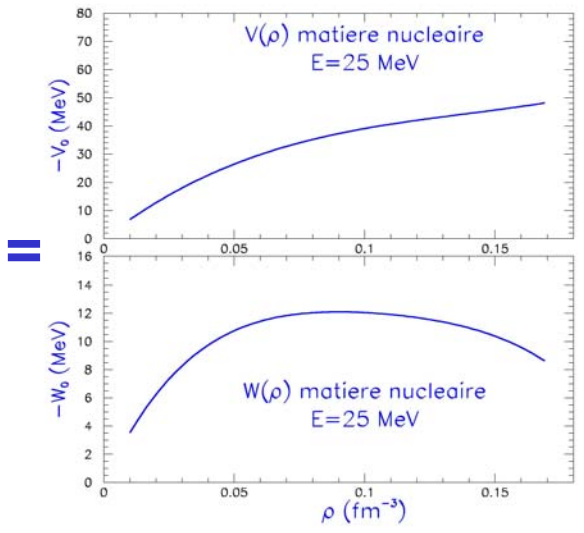
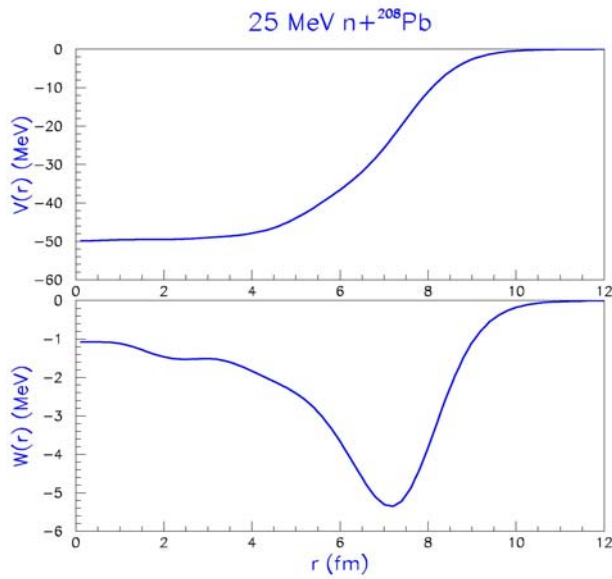


- ≈ 20 adjusted parameters
- Very precise (1%)
- Relatively weak predictive power far away from stability



Semi-microscopic OMP

Optical potential = **Effective Interaction** \otimes **Radial densities**



$$U(r,E) = \frac{U(\rho(r'),E)}{\rho(r')} \otimes \rho(r)$$

Depends on the nucleus

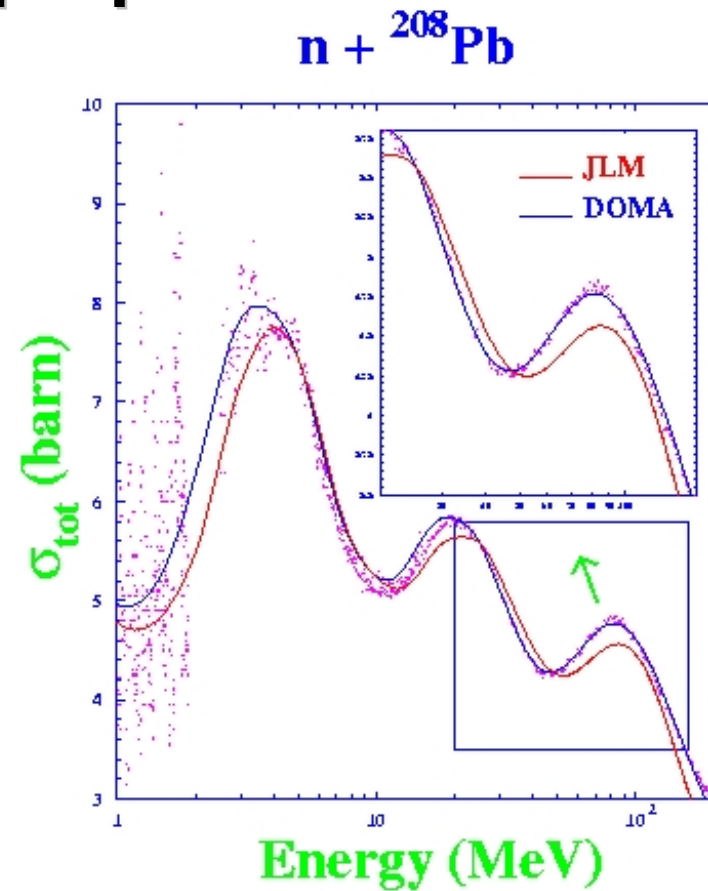
Independent of the nucleus

Depends on the nucleus

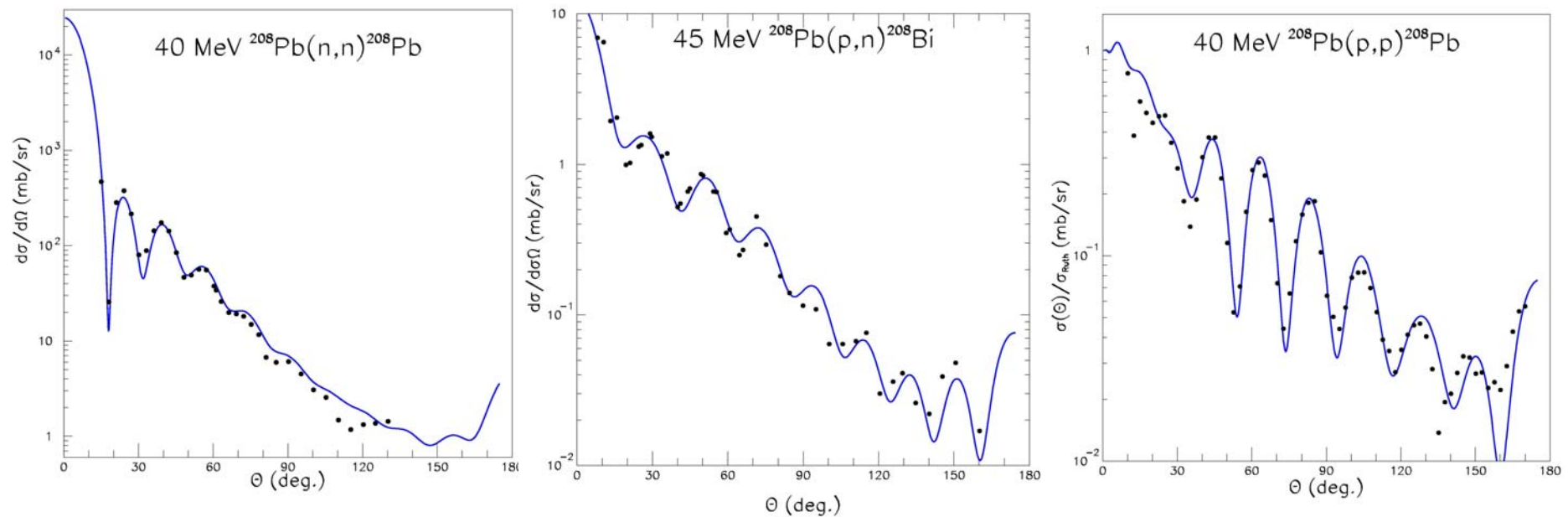
- No adjustable parameters
- Based on nuclear structure properties

⇒ usable for any nucleus

- Less precise than the phenomenological approach

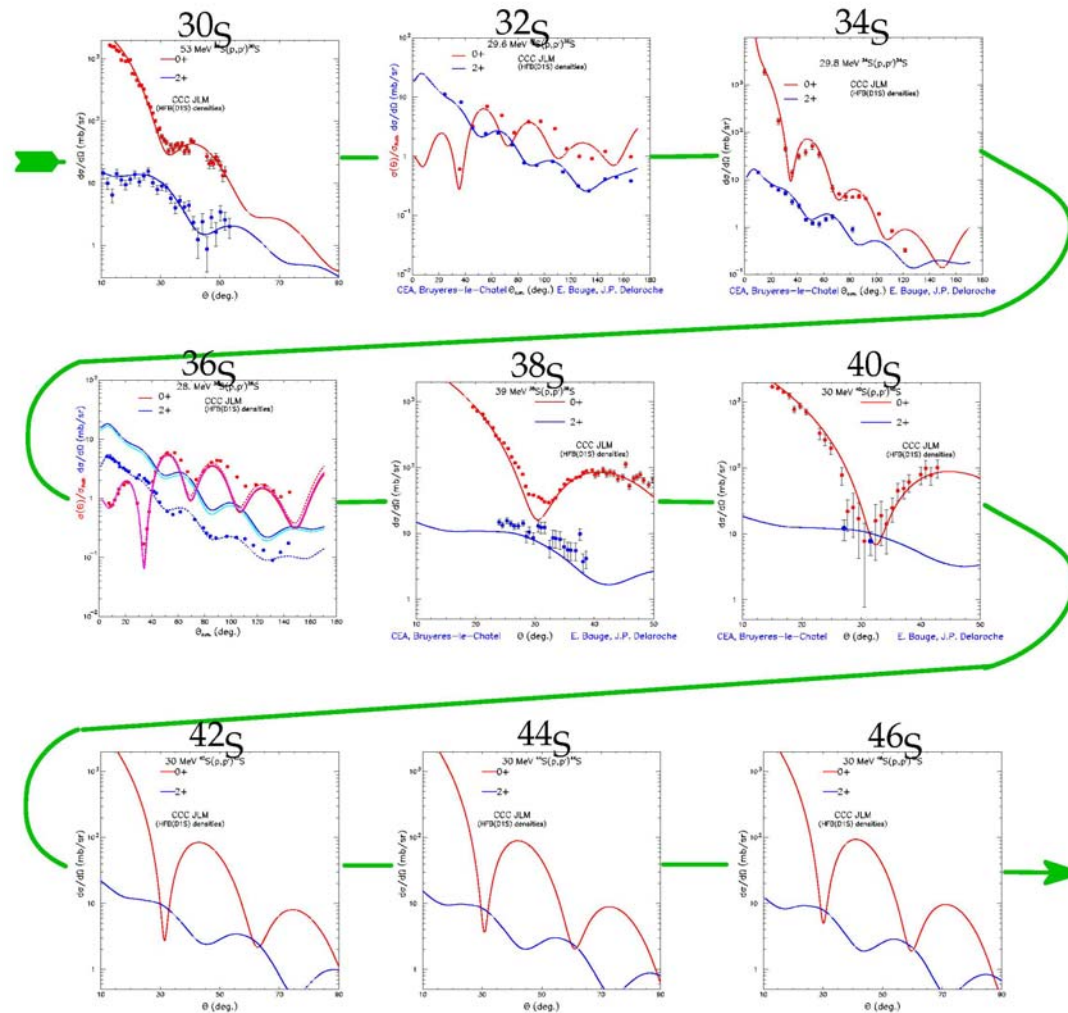


Unique description of elastic scattering (n,n), (p,p) et (p,n)

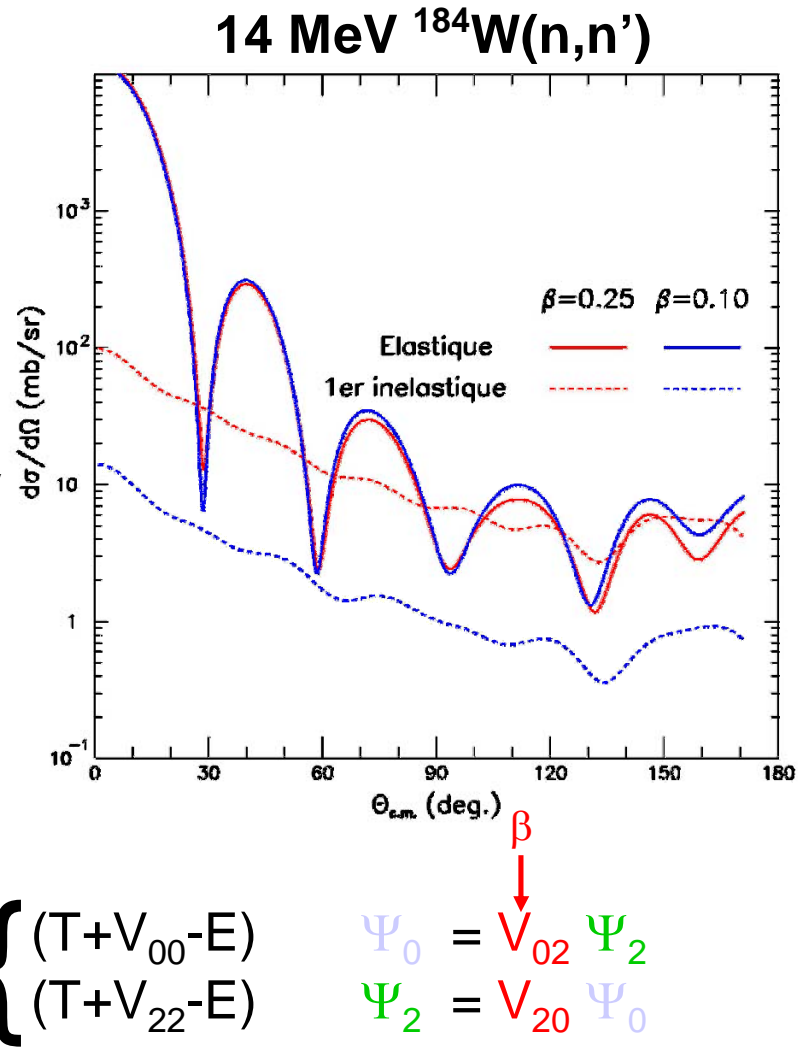
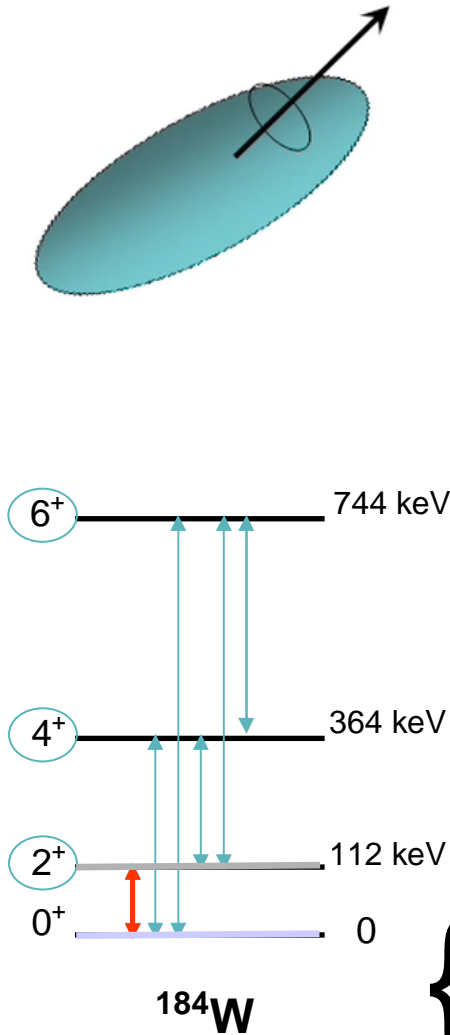
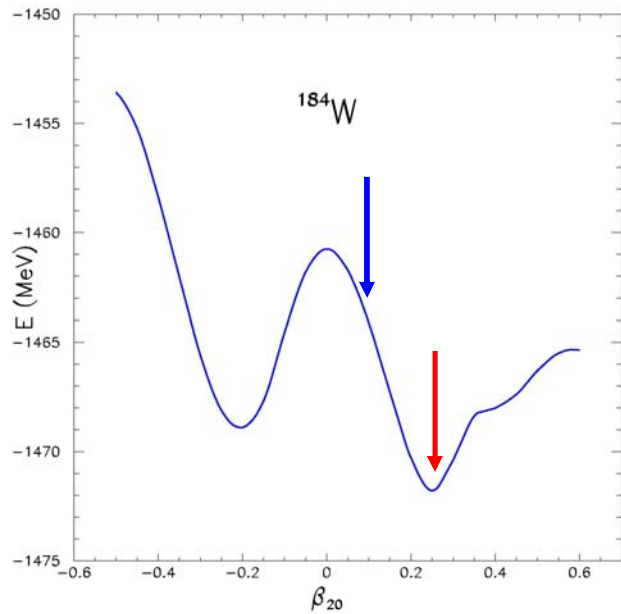


Semi-microscopic OMP

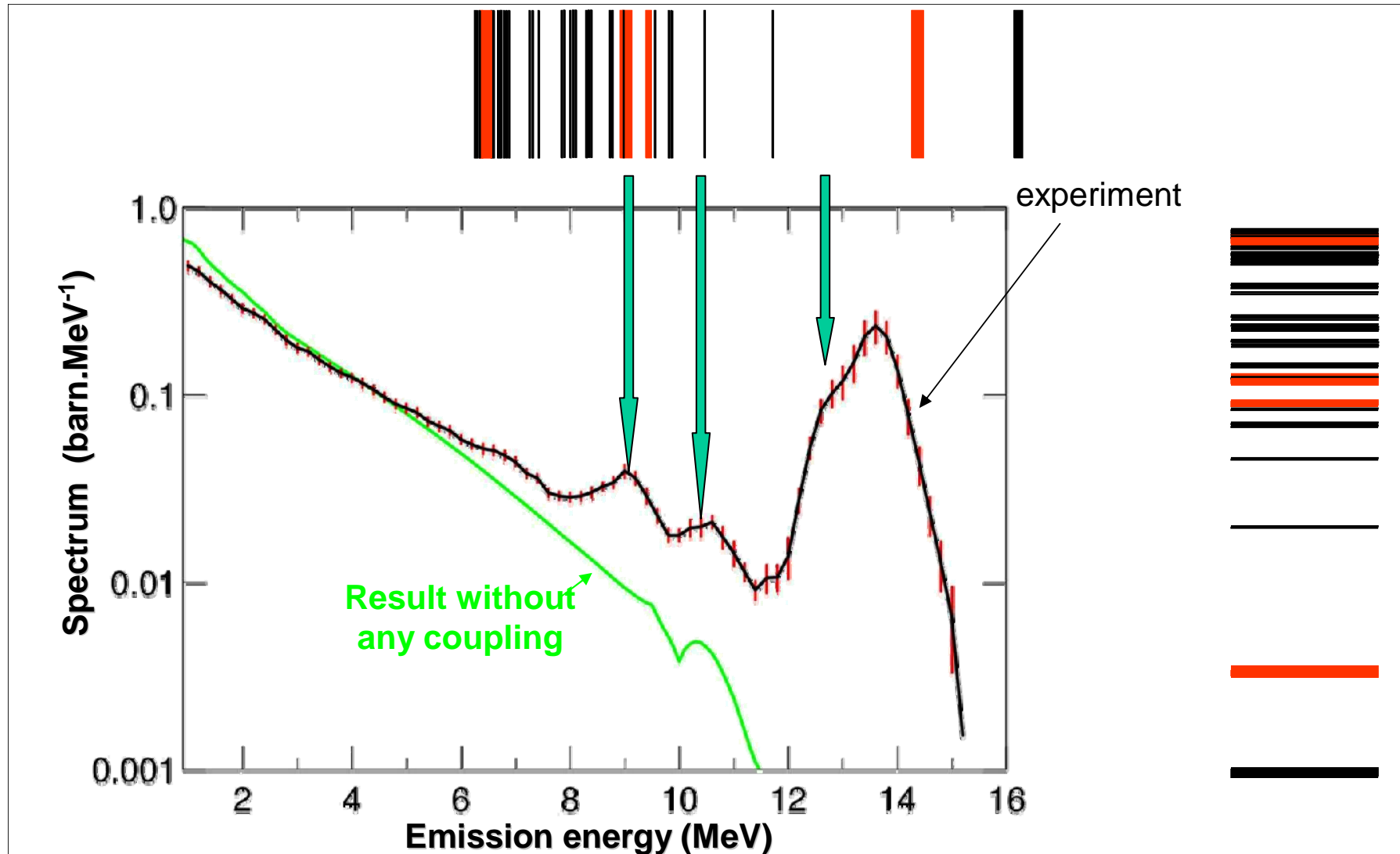
Enables to perform predictions for very exotic nuclei for which
There exist no experimental data



Coupled channels in OMP



Coupled channels in OMP





Coupled channels in TALYS

deformation file z092

```

92 237 5 R D
0 R 0 1.75000 0.65000
1 R 0
2 R 0
3 R 0
5 R 0
92 238 23 R D
0 R 0 1.54606 0.44508
1 R 0
2 R 0
3 R 0
4 R 0
5 V 1 3 0 0.90000
6 V 1
7 R 0
8 V 1
9 V 2 4 0 0.20000
10 V 3 3 1 0.10000
11 V 3
12 V 2
13 V 1
14 V 4 2 0 0.10000
15 V 3
16 V 4
17 V 2
21 V 5 2 2 0.10000
22 R 0
23 V 5
25 V 4
31 V 5

```

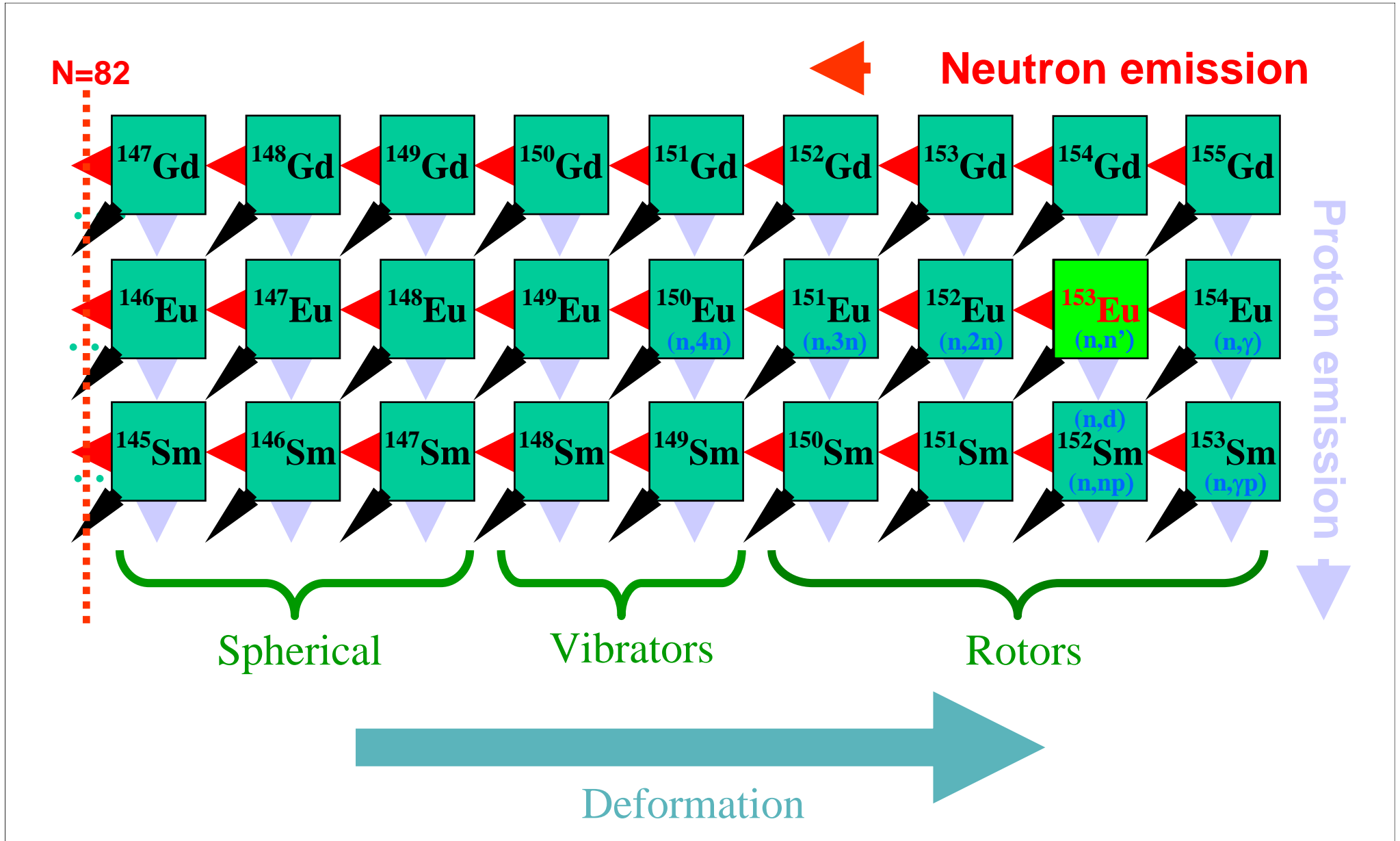
level file z092

```

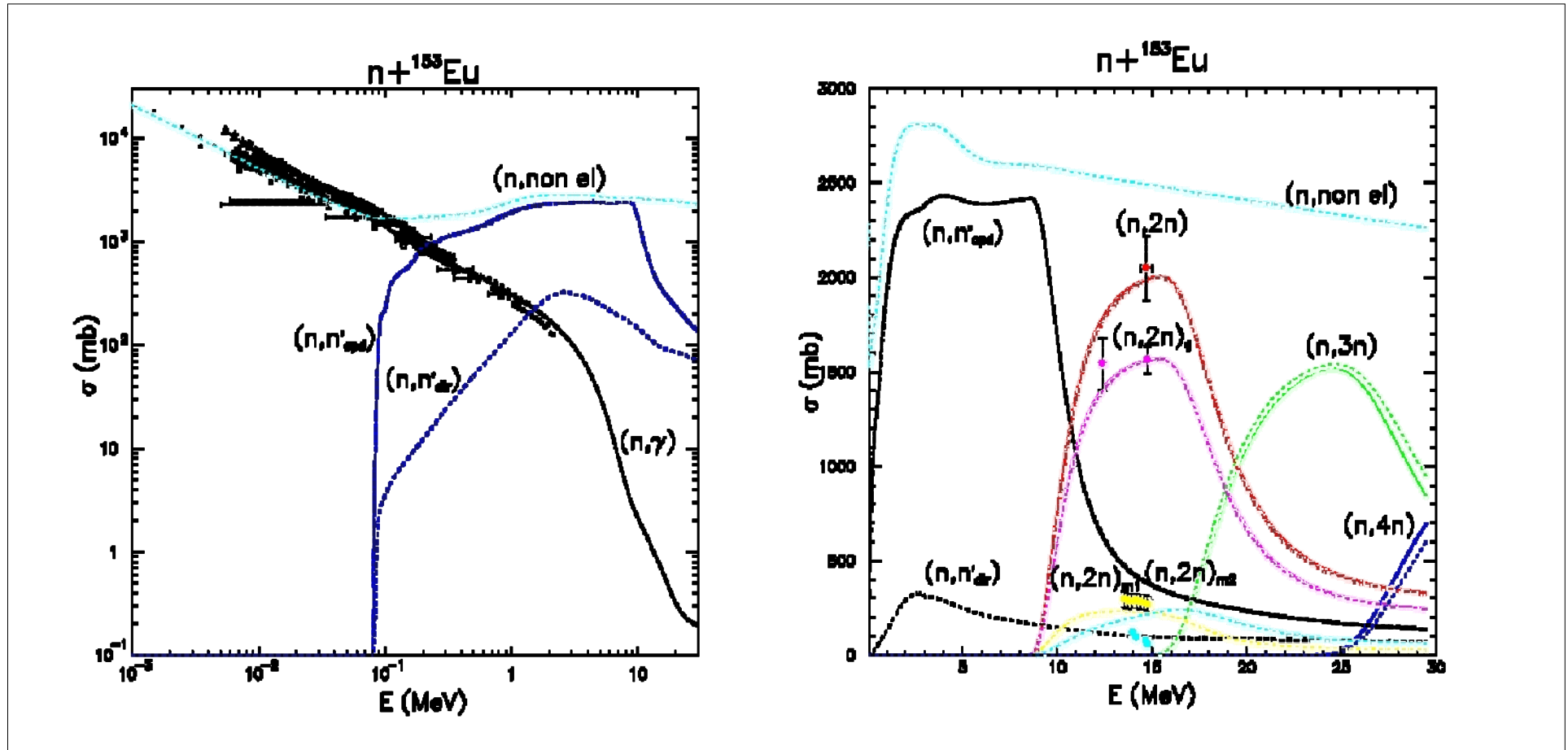
92 238 501 152
0 0.000000 0.0 1 0 1.410E+17 238U 0+
1 0.044916 2.0 1 1 2.060E-10 2+
2 0.148380 4.0 1 1 6.090E+02 4+
3 0.307180 6.0 1 1 1.160E+01 6+
4 0.518100 8.0 1 1 1.870E+00 8+
5 0.680110 1.0 -1 2 2.300E-11 1-
6 0.731930 3.0 -1 3 3.500E-14 3-
7 0.775900 10.0 1 1 9.000E-12 10+
8 0.826640 5.0 -1 2 3.123E+02 5-
9 0.927210 0.0 1 1 0.450225 1.000E-02 0+
10 0.930550 1.0 -1 3 0.549775 7.045E-03 (1-)
11 0.950120 2.0 -1 3 0.335445 1.195E-02 2-
12 0.966130 2.0 1 5 0.664555 7.213E-03 2+
0 0.000000 3.123E+02
1 0.558700 2.000E-02
0 0.441300 2.000E-02
2 0.450225 1.000E-02
1 0.549775 7.045E-03
4 1.000000 3.130E-01
3 0.335445 1.195E-02
2 0.664555 7.213E-03
1 1.000000 1.341E-02
5 0.157218 3.413E-01
1 0.673221 4.650E-03
0 0.169560 4.260E-03
6 0.419039 5.570E-01
5 0.251909 2.685E-01
1 0.329052 4.470E-03
6 0.064358 6.890E-02
5 0.035757 4.380E-02
2 0.418755 1.660E-02
1 0.367863 2.300E-01
0 0.312267 1.200E-02

```

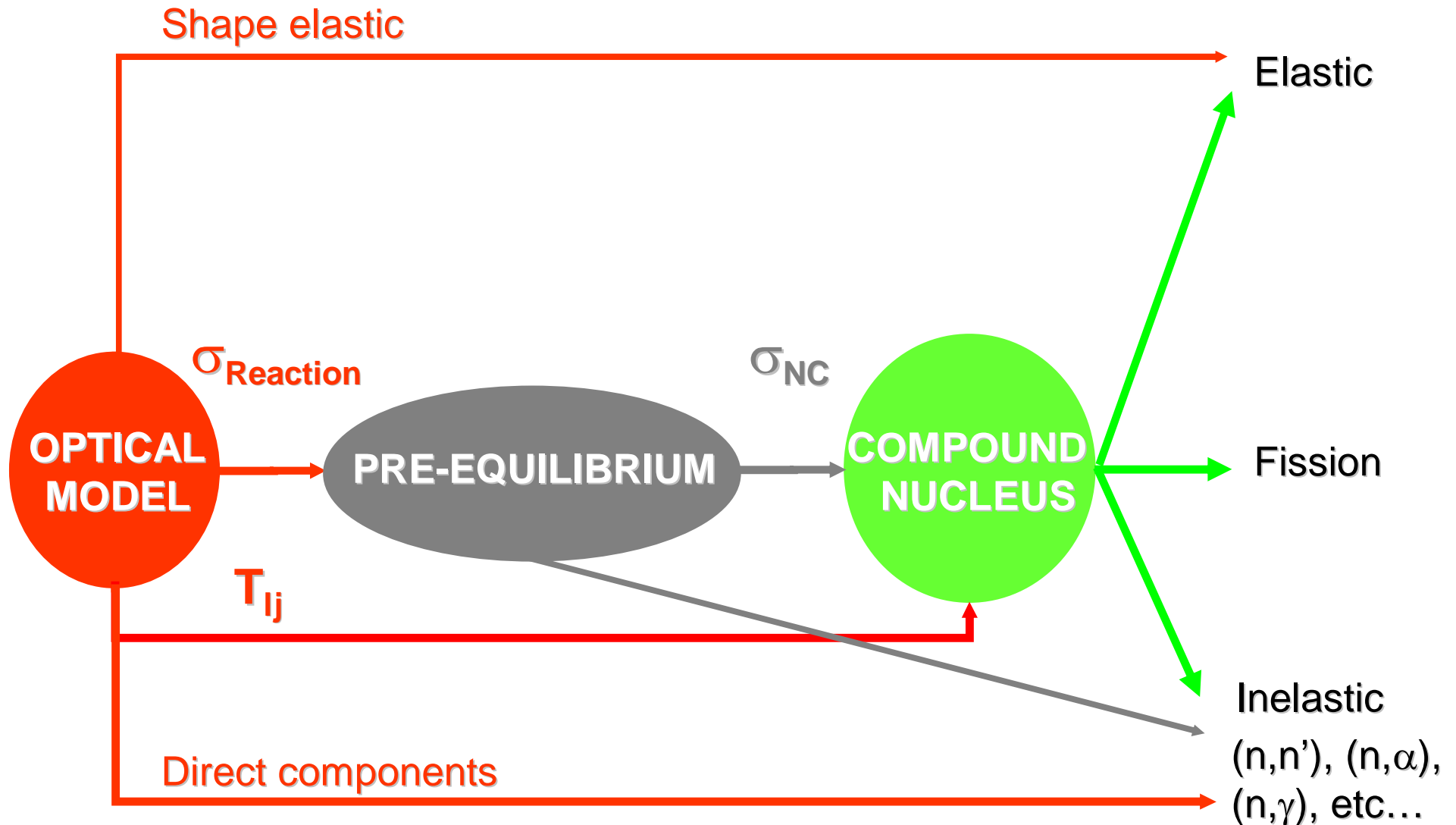
Decay-dependent OMPs in TALYS



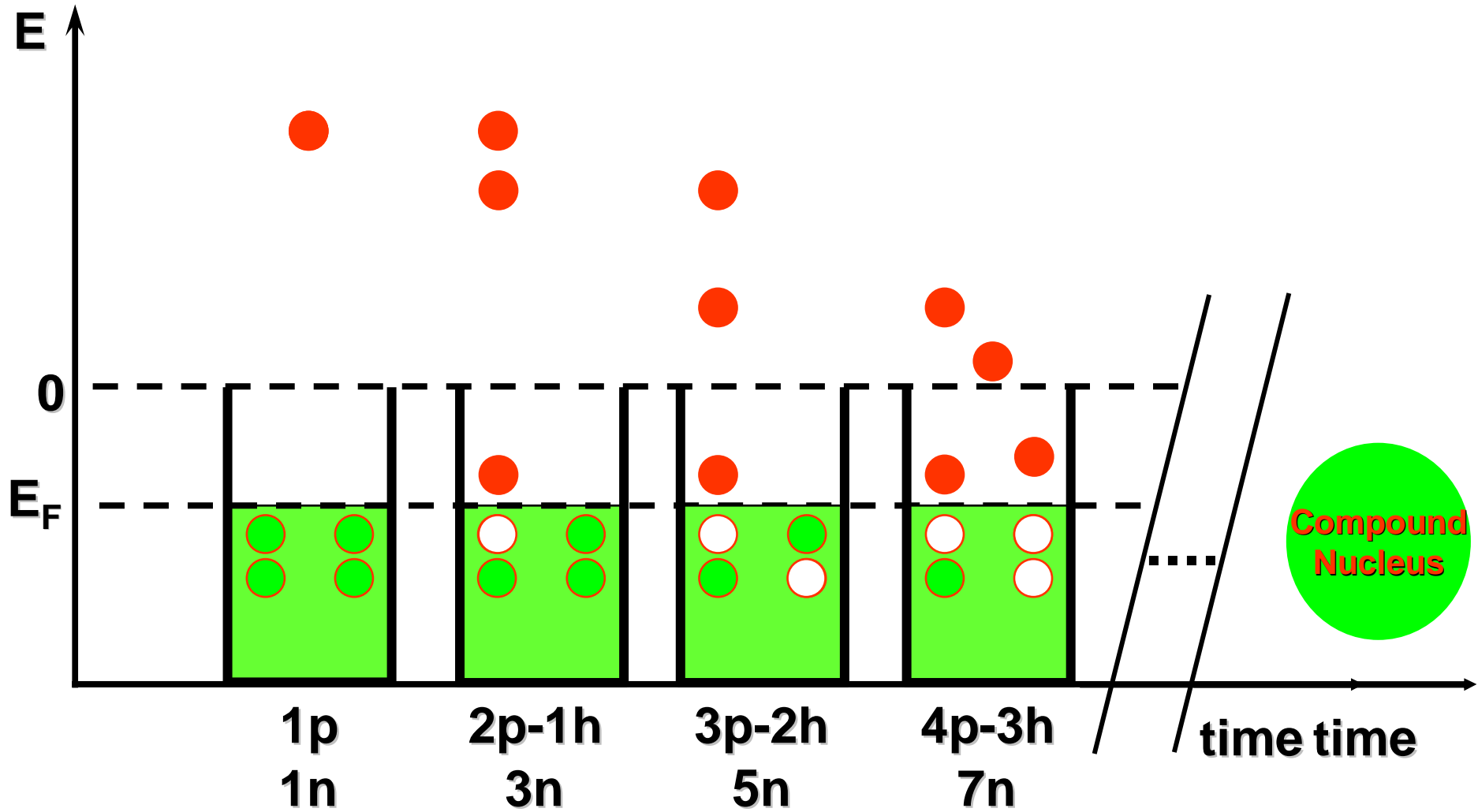
Decay-dependent OMPs in TALYS



Models sequence



Pre-equilibrium model(s)



Pre-equilibrium exciton model

$P(n, E, t)$ = **Probability** to find for a given time t the composite system with an energy E and an **excitons number** n .

$\lambda_{a, b}(E)$ = Transition rate from an initial state a towards a state b for a given energy E .

Evolution equation

$$\frac{dP(n, E, t)}{dt} = \overset{\text{Apparition}}{P(n-2, E, t) \lambda_{n-2, n}(E)} + \overset{\text{Disparition}}{P(n+2, E, t) \lambda_{n+2, n}(E)} - P(n, E, t) \left[\lambda_{n, n+2}(E) + \lambda_{n, n-2}(E) + \lambda_{n, \text{emiss}}(E) \right]$$

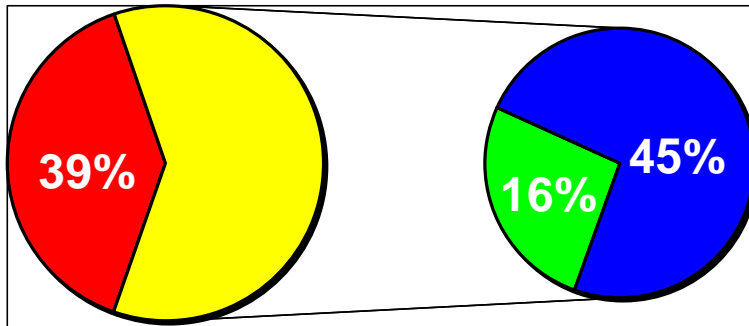
Emission cross section in channel c

$$\sigma_c(E, \varepsilon_c) d\varepsilon_c = \sigma_R \int_0^{t_{\text{eq}}} \sum_{n, \Delta n=2} P(n, E, t) \lambda_{n, c}(E) dt d\varepsilon_c$$

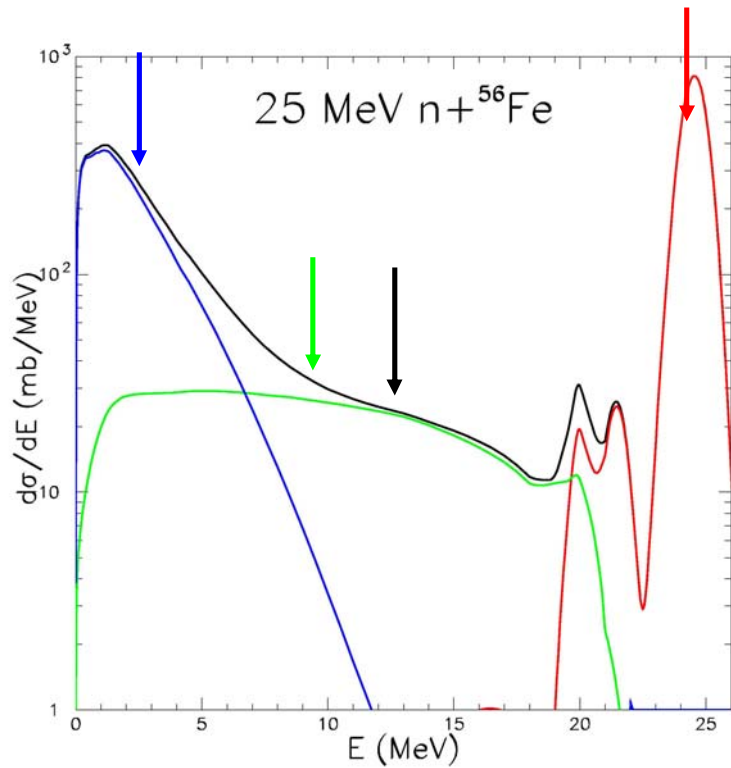
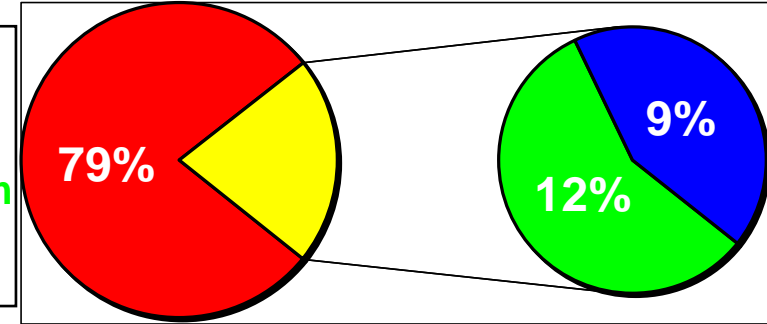
Pre-equilibrium model

Cross section

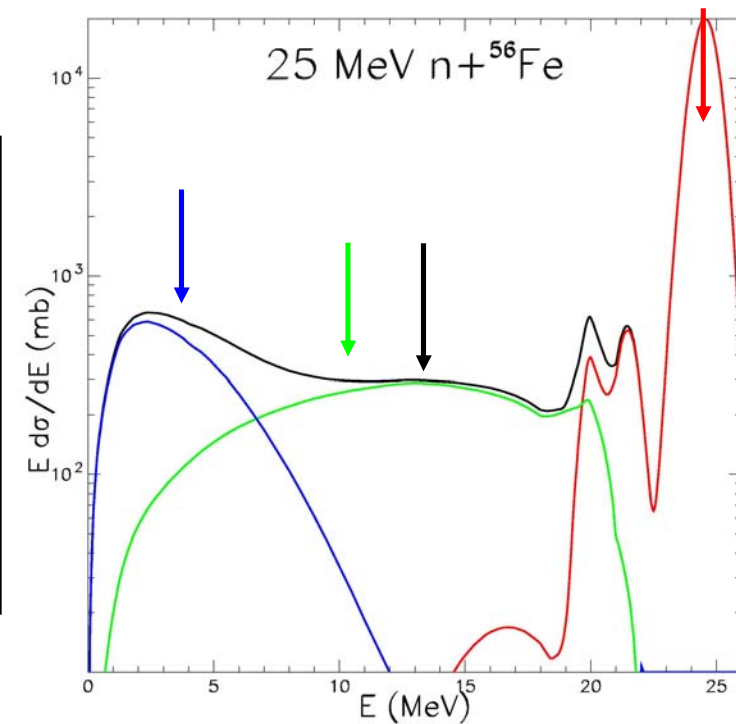
Outgoing energy

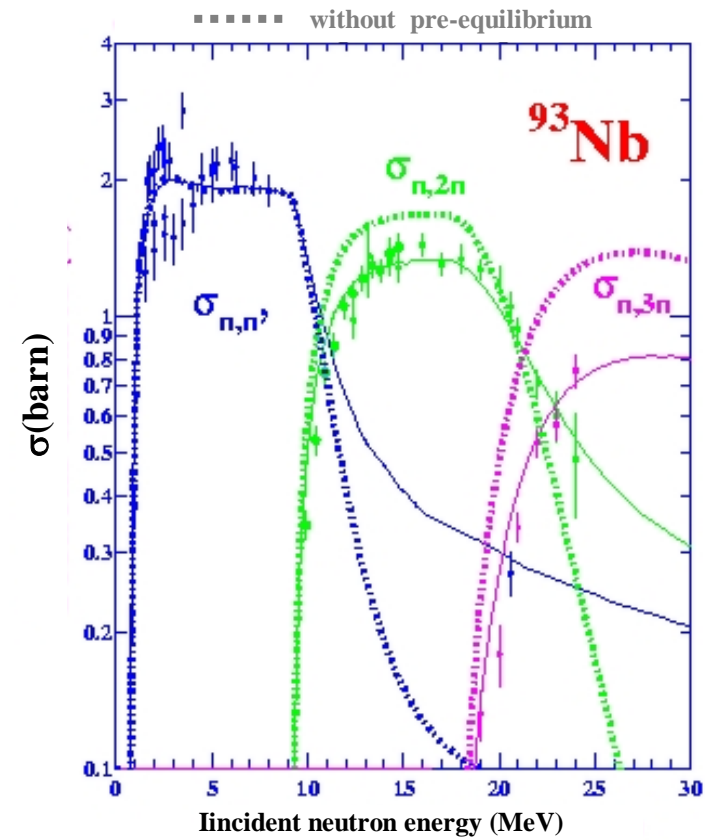
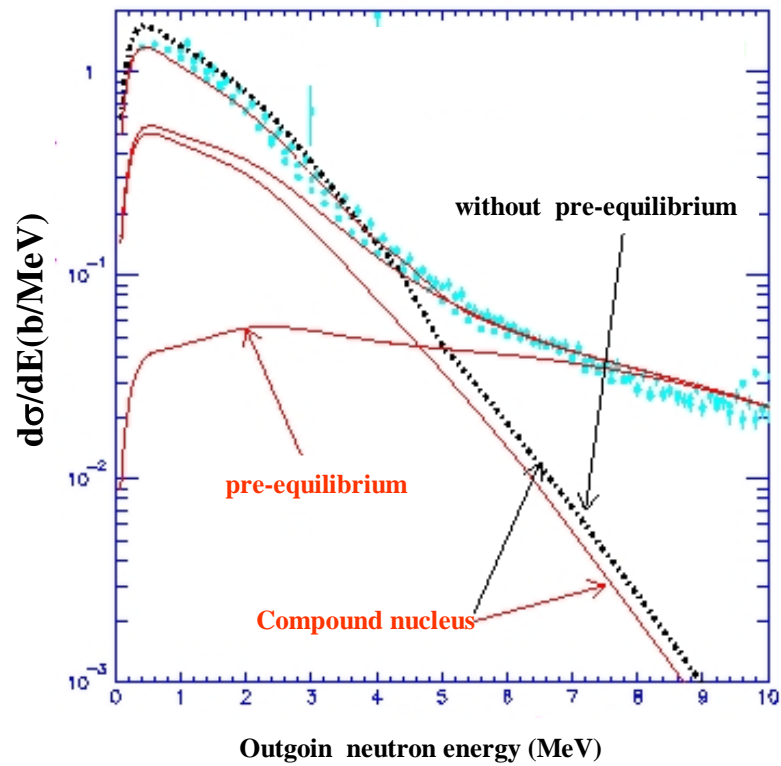


Total
Direct
Pre-equilibrium
Statistical



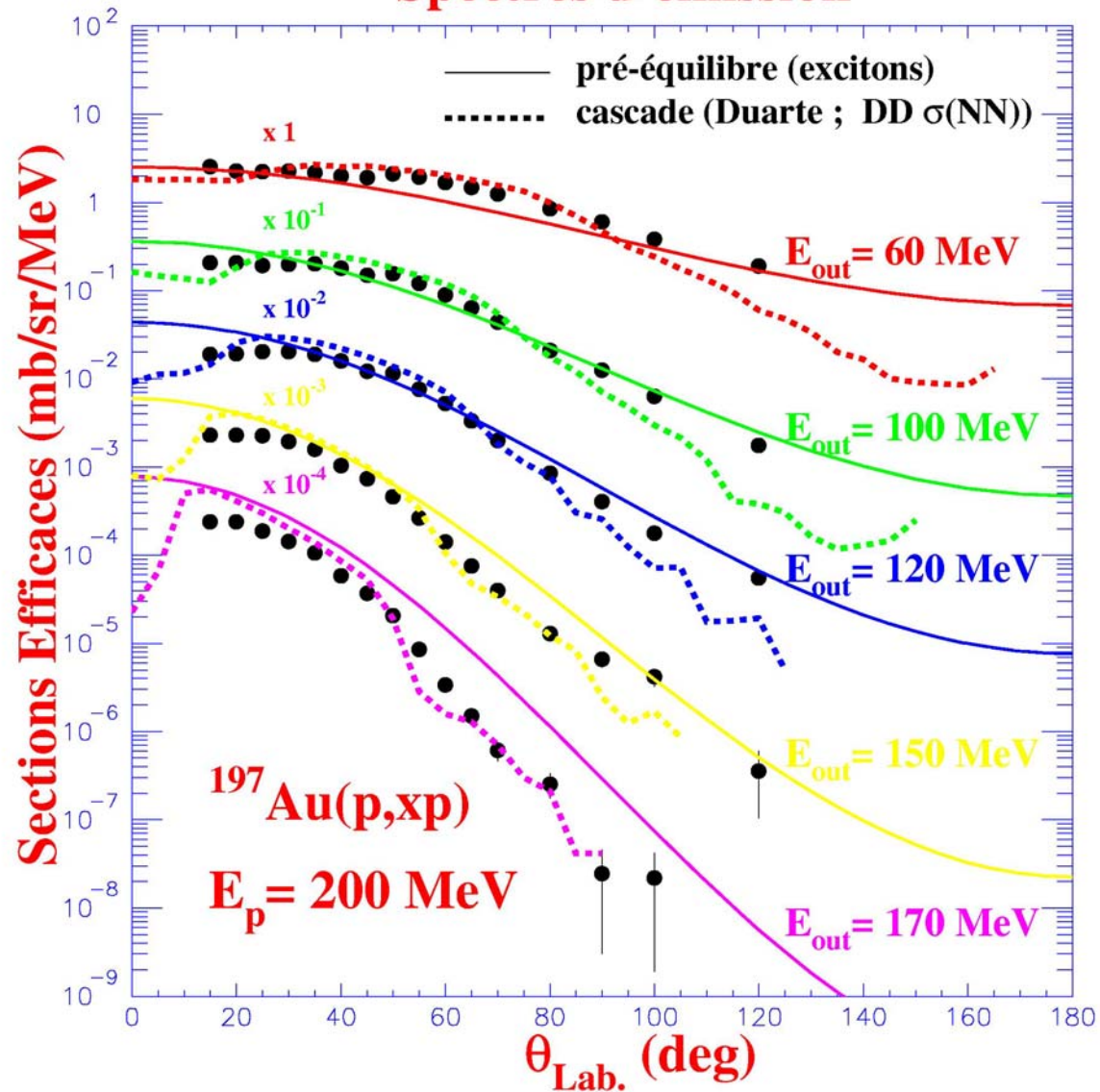
$\langle E_{\text{Tot}} \rangle = 12.1$
 $\langle E_{\text{Dir}} \rangle = 24.3$
 $\langle E_{\text{PE}} \rangle = 9.32$
 $\langle E_{\text{Sta}} \rangle = 2.5$
 (MeV)



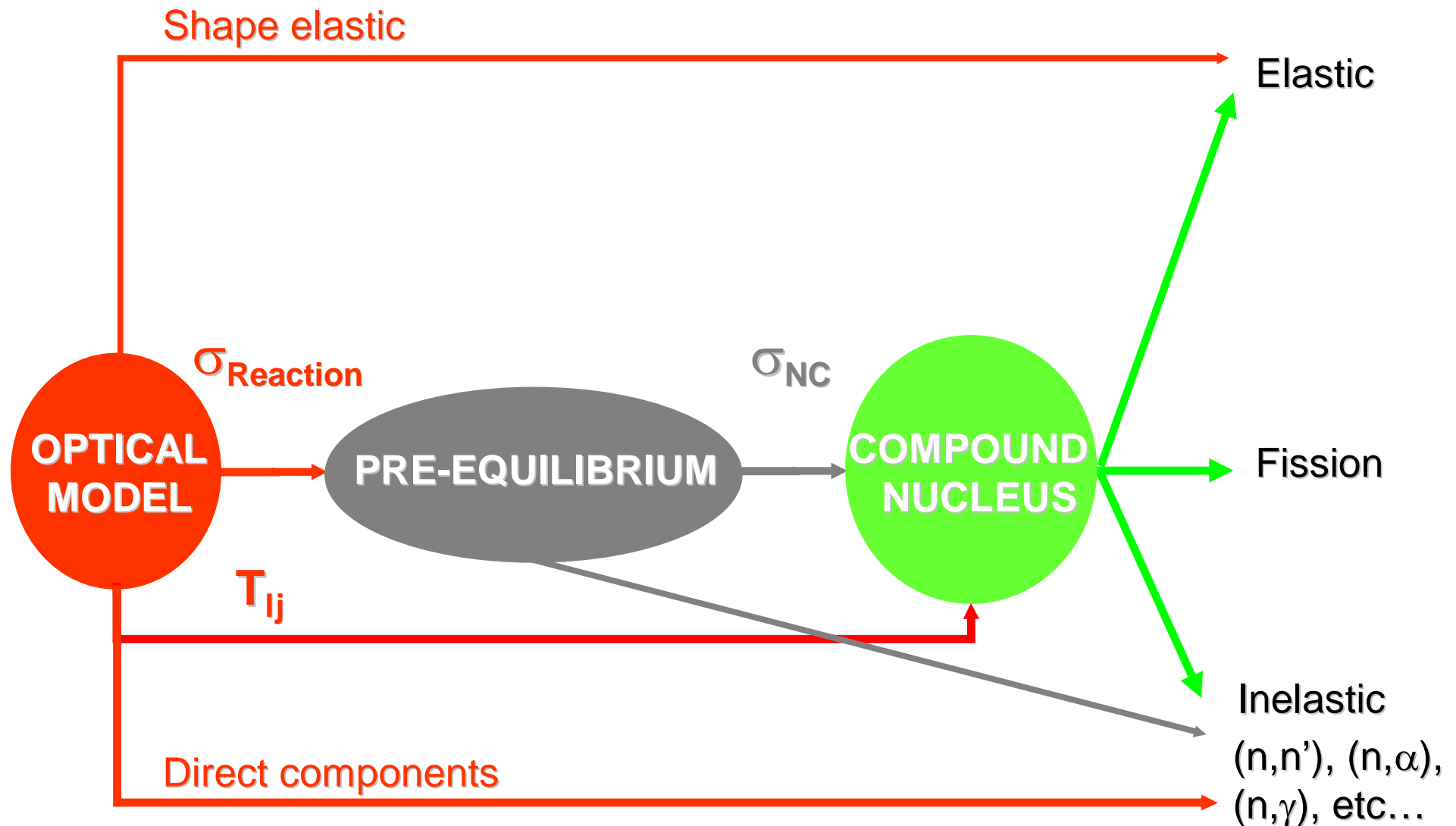
14 MeV neutron + ^{93}Nb 

Link with high energy cascade

Spectres d'émission




Models sequence



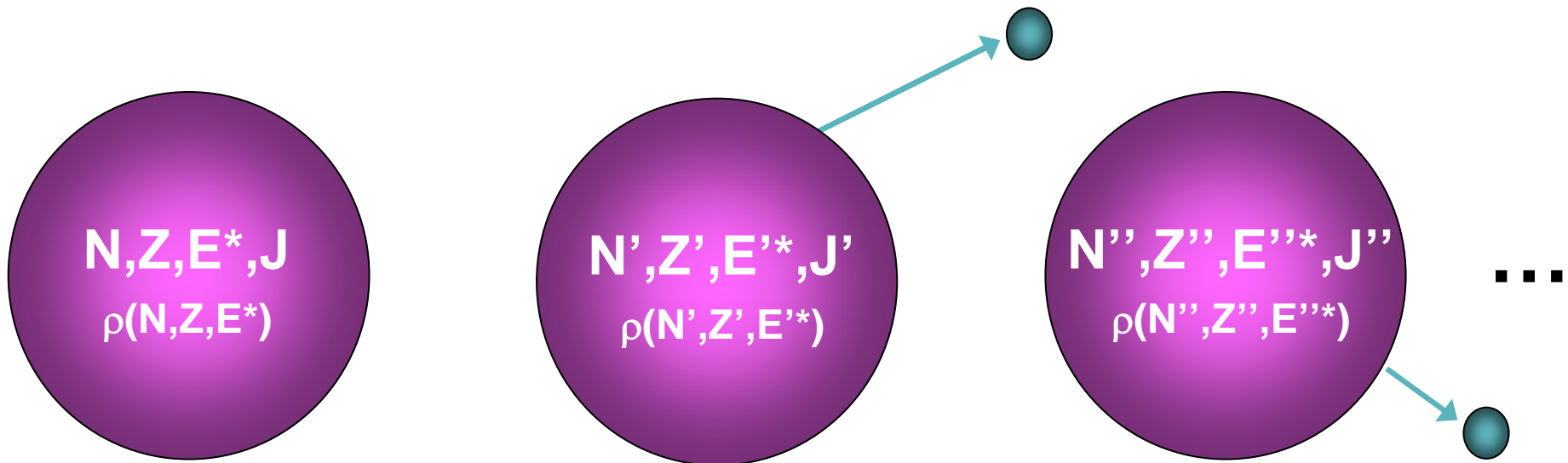
- ➔ **Generalities and definitions**
- ➔ **Model ingredients**
- ➔ **Fission**
- ➔ **Level densities**

Compound Nucleus model

After direct and pre-equilibrium emission

$$\sigma_{\text{reaction}} = \sigma_{\text{dir}} + \sigma_{\text{pre-eq}} + \sigma_{\text{NC}}$$


N_0	$N_0 - dN_D$	$N_0 - dN_D - dN_{PE} = E$
Z_0	$Z_0 - dZ_D$	$Z_0 - dZ_D - dZ_{PE} = Z$
E^*_0	$E^*_0 - dE^*_D$	$E^*_0 - dE^*_D - dE^*_{PE} = E^*$
J_0	$J_0 - dJ_D$	$J_0 - dJ_D - dJ_{PE} = J$



Compound nucleus hypothesis

- Continuum of excited levels
- Independence between incoming channel **a** and outgoing channel **b**

$$\sigma_{ab} = \sigma_a^{(CN)} P_b$$

$$\sigma_a^{(CN)} = \frac{\pi}{k_a^2} T_a$$

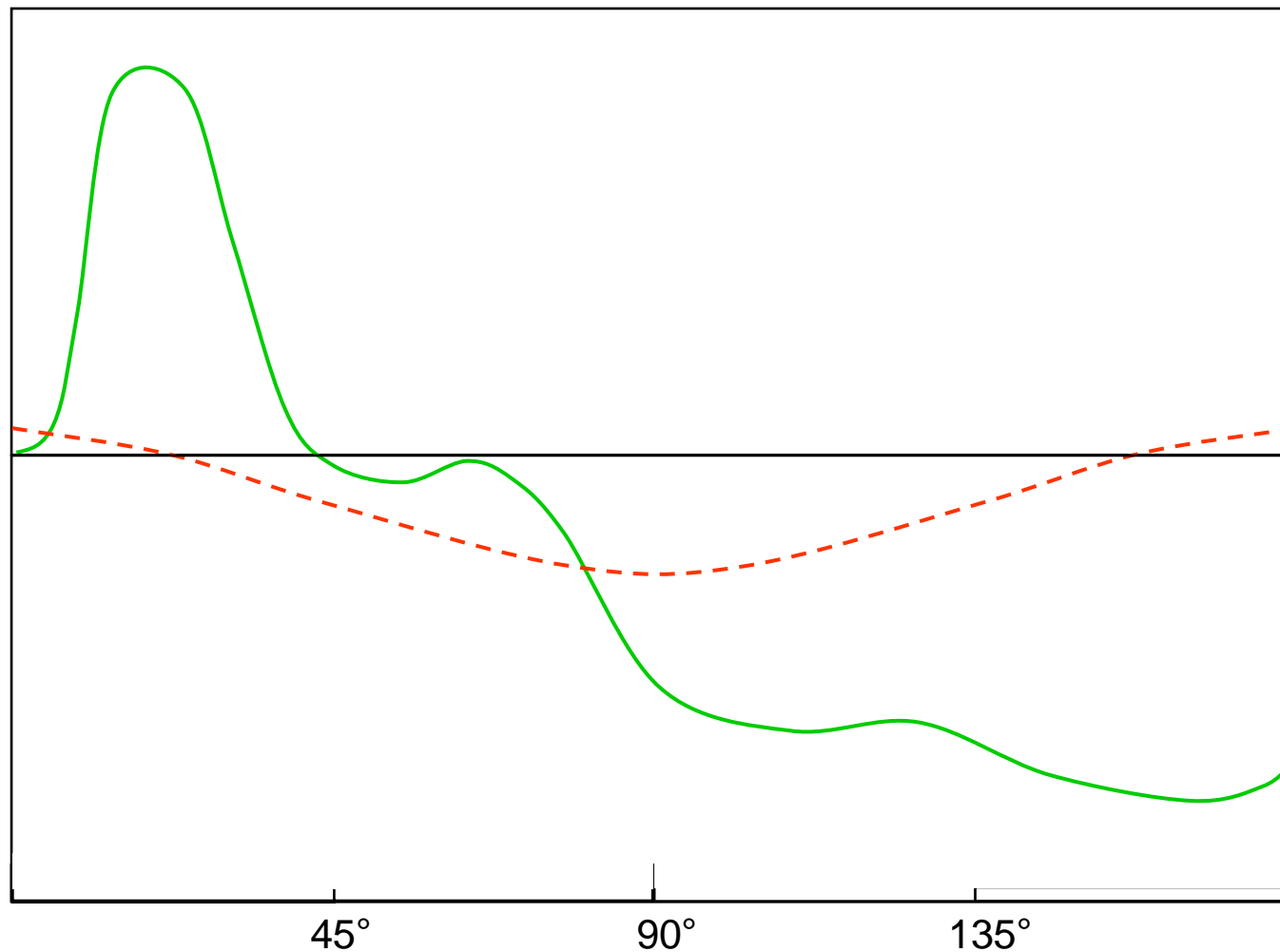
$$P_b = \frac{T_b}{\sum_c T_c}$$

⇒ Hauser- Feshbach formula

$$\sigma_{ab} = \frac{\pi}{k_a^2} \frac{T_a T_b}{\sum_c T_c}$$

Compound Nucleus model

Compound angular distribution & **direct** angular distributions



Channel Definition



Incident channel $a = (\vec{l}_a, \vec{j}_a = \vec{l}_a + \vec{s}_a, \vec{J}_A, \pi_A, E_A, E_a)$

Conservation equations

- Total energy : $E_a + E_A = E_{\text{CN}} = E_b + E_B$
- Total momentum : $\vec{p}_a + \vec{p}_A = \vec{p}_{\text{CN}} = \vec{p}_b + \vec{p}_B$
- Total angular momentum : $\vec{l}_a + \vec{s}_a + \vec{J}_A = \vec{J}_{\text{CN}} = \vec{l}_b + \vec{s}_b + \vec{J}_B$
- Total parity : $\pi_A (-1)^{l_a} = \pi_{\text{CN}} = \pi_B (-1)^{l_b}$

In realistic calculations, all possible quantum number combinations have to be considered

$$\sigma_{ab} = \frac{\pi}{k_a^2} \sum_{J=|I_A - s_a|}^{I_A + s_a + l_a^{max}} \sum_{\pi = \pm} \frac{(2J+1)}{(2I_A+1)(2s_a+1)}$$

$$\sum_{j_a=|J-I_A|}^{J+I_A} \sum_{l_a=|j_a-s_a|}^{j_a+s_a} \sum_{j_b=|J-I_B|}^{J+I_B} \sum_{l_b=|j_b-s_b|}^{j_b+s_b}$$

$$\delta_{\pi}(a) \delta_{\pi}(b) \frac{T_{a, l_a, j_a}^{J\pi} T_{b, l_b, j_b}^{J\pi}}{\sum_c T_{c, l_c, j_c}^{J\pi}} W_{a, l_a, j_a, b, l_b, j_b}^{J\pi}$$

Breit-Wigner resonance integrated and averaged over an energy width
Corresponding to the incident beam dispersion

$$\langle \sigma_{ab} \rangle = \frac{\pi}{k_a^2} \frac{2\pi}{D} \left\langle \frac{\Gamma_a \Gamma_b}{\Gamma_{tot}} \right\rangle$$

Or $T_\alpha \approx \frac{2\pi \langle \Gamma_\alpha \rangle}{D}$

$$\langle \sigma_{ab} \rangle = \frac{\pi}{k_a^2} \frac{T_a T_b}{\sum_c T_c} \mathbf{W}_{ab}$$

with $\mathbf{W}_{ab} = \frac{\left\langle \frac{\Gamma_a \Gamma_b}{\Gamma_{tot}} \right\rangle}{\frac{\langle \Gamma_a \rangle \langle \Gamma_b \rangle}{\langle \Gamma_{tot} \rangle}}$

- Tepel method

Simplified iterative method

- Moldauer method

Simple integral

- GOE triple integral

« **exact** » result

Elastic enhancement with respect to the other channels

The GOE triple integral

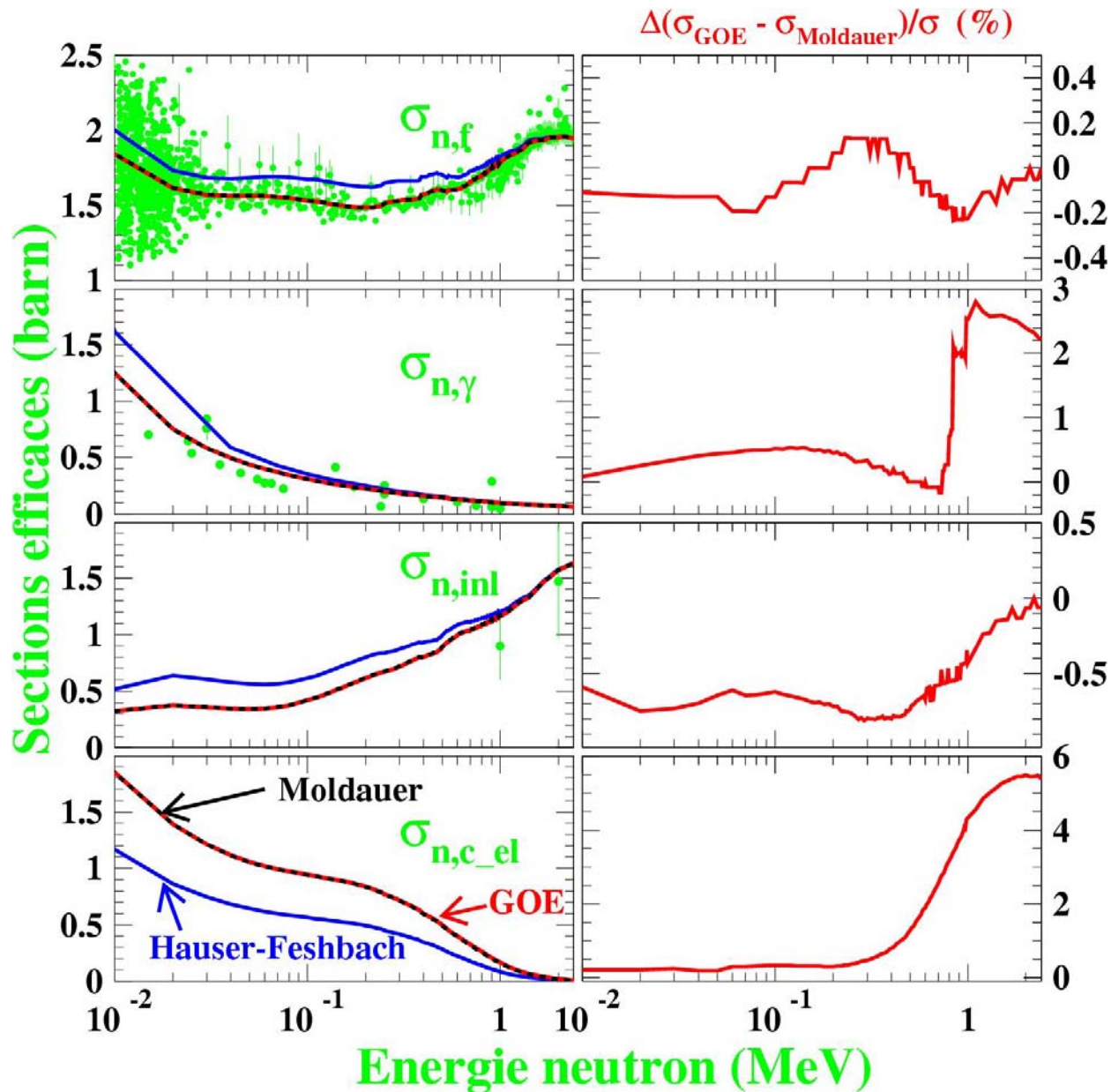
$$W_{a,l_a,j_a,b,l_b,j_b} = \int_0^{+\infty} d\lambda_1 \int_0^{+\infty} d\lambda_2 \int_0^1 d\lambda \frac{\lambda(1-\lambda)|\lambda_1 - \lambda_2|}{\sqrt{\lambda_1(1+\lambda_1)\lambda_2(1+\lambda_2)(\lambda+\lambda_1)^2(\lambda+\lambda_2)^2}}$$

$$\prod_c \frac{(1 - \lambda T_{c,l_c,j_c}^J)}{\sqrt{(1 + \lambda_1 T_{c,l_c,j_c}^J)(1 + \lambda_2 T_{c,l_c,j_c}^J)}} \left\{ \delta_{ab}(1 - T_{a,l_a,j_a}^J) \right.$$

$$\left[\frac{\lambda_1}{1 + \lambda_1 T_{a,l_a,j_a}^J} + \frac{\lambda_2}{1 + \lambda_2 T_{a,l_a,j_a}^J} + \frac{2\lambda}{1 - \lambda T_{a,l_a,j_a}^J} \right]^2 + (1 + \delta_{ab})$$

$$\left[\frac{\lambda_1(1 + \lambda_1)}{(1 + \lambda_1 T_{a,l_a,j_a}^J)(1 + \lambda_1 T_{b,l_b,j_b}^J)} + \frac{\lambda_2(1 + \lambda_2)}{(1 + \lambda_2 T_{a,l_a,j_a}^J)(1 + \lambda_2 T_{b,l_b,j_b}^J)} \right.$$

$$\left. + \frac{2\lambda(1 - \lambda)}{(1 - \lambda T_{a,l_a,j_a}^J)(1 - \lambda T_{b,l_b,j_b}^J)} \right] \left. \right\}$$



$$\sigma_{\text{NC}} = \sum_{\mathbf{b}} \sigma_{\mathbf{a}\mathbf{b}} \quad \text{où } \mathbf{b} = \gamma, \text{ n, p, d, t, } \dots, \text{ fission}$$

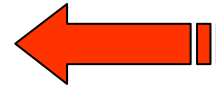
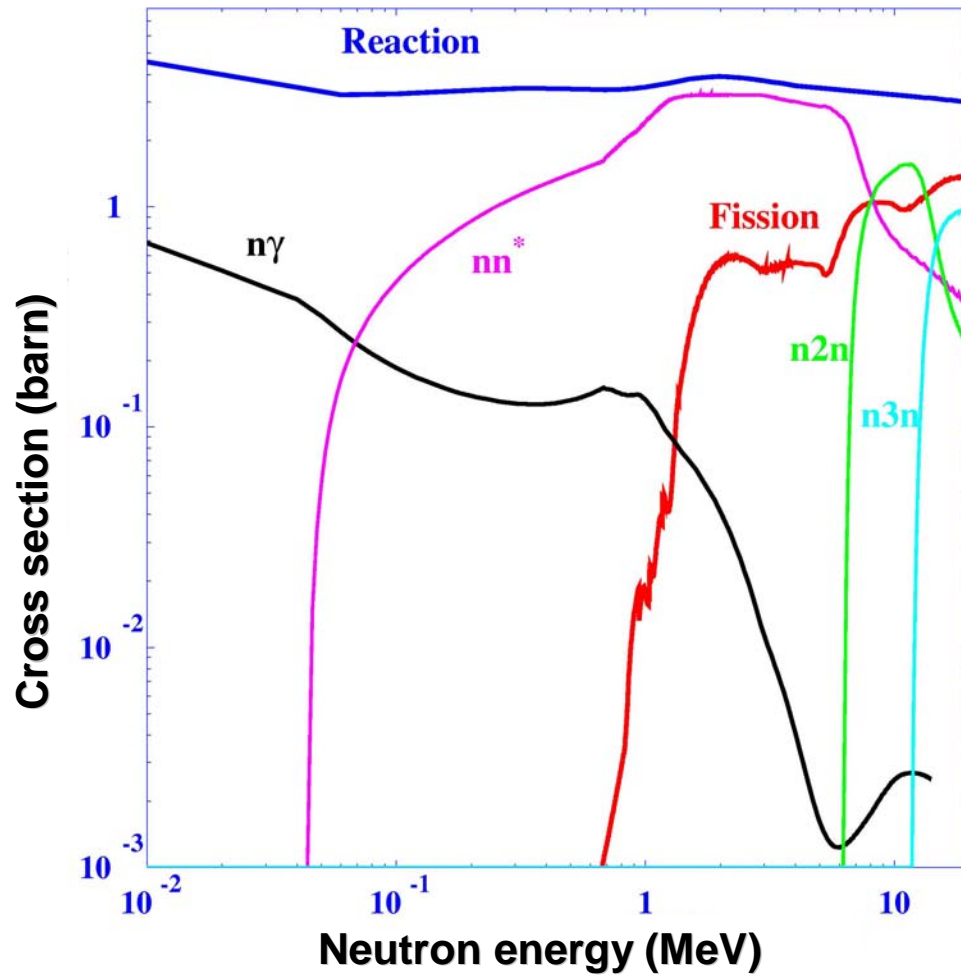
$$\sigma_{\mathbf{a}\mathbf{b}} = \frac{\pi}{k_{\mathbf{a}}^2} \sum_{\mathbf{J}, \pi} \sum_{\alpha, \beta} \frac{(2\mathbf{J}+1)}{(2s+1)(2\mathbf{I}+1)} T_{\mathbf{l}\mathbf{j}}^{\mathbf{J}\pi}(\alpha) \frac{\langle T_{\mathbf{b}}^{\mathbf{J}\pi}(\beta) \rangle}{\sum_{\delta} \langle T_{\mathbf{d}}^{\mathbf{J}\pi}(\delta) \rangle} W_{\alpha\beta}$$

with $\mathbf{J} = \mathbf{l}_{\alpha} + \mathbf{s}_{\alpha} + \mathbf{I}_{\text{A}} = \mathbf{j}_{\alpha} + \mathbf{I}_{\text{A}}$ et $\pi = (-1)^{\mathbf{l}_{\alpha}} \pi_{\text{A}}$

and $\langle T_{\mathbf{b}}^{\mathbf{J}\pi}(\beta) \rangle$ = transmission coefficient for outgoing channel β associated with the outgoing particle \mathbf{b}

Compound Nucleus Model

$n + {}^{238}\text{U}$

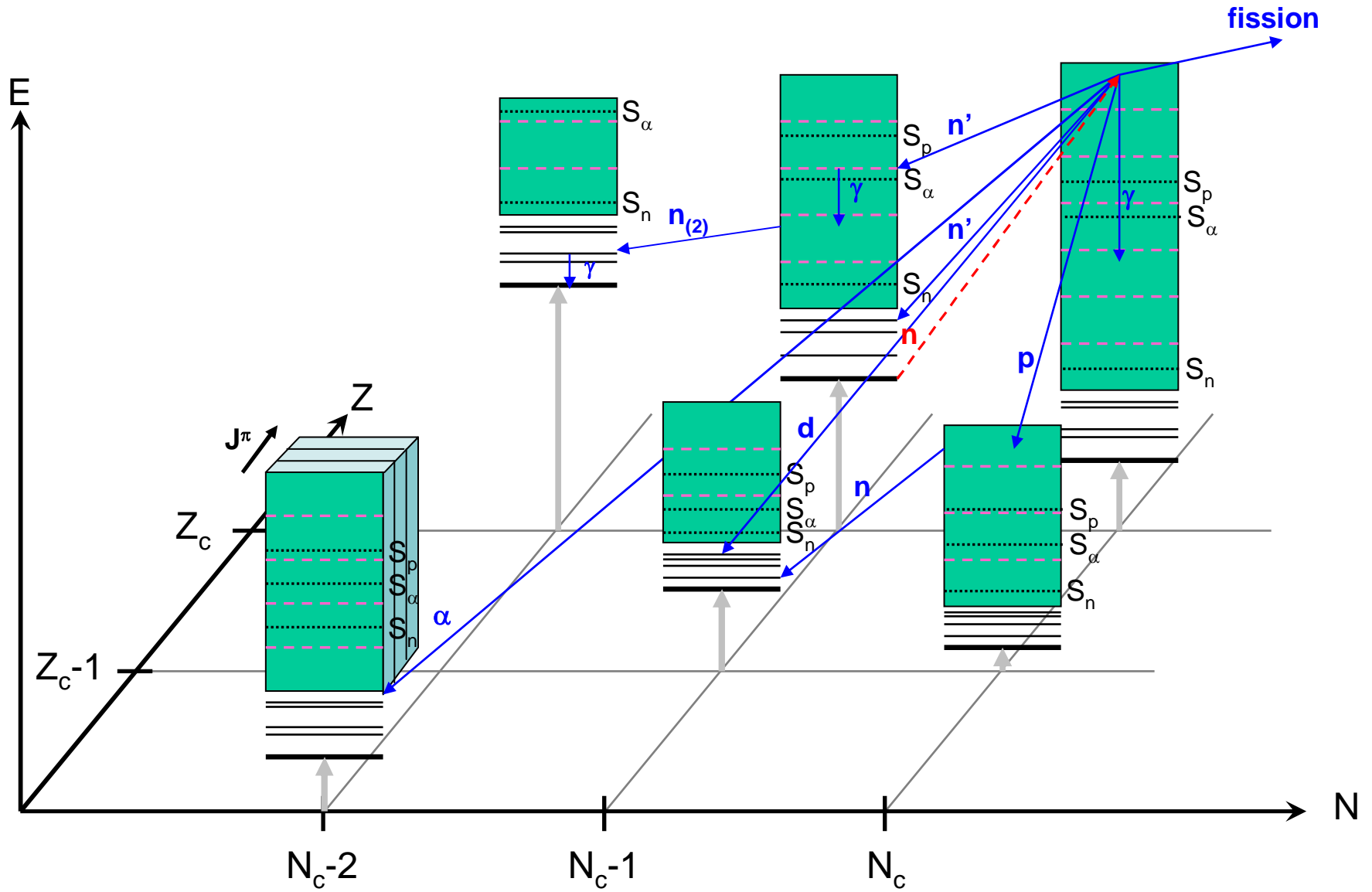


Optical model
 +
Statistical model
 +
Pre-equilibrium model

$$\sigma_R = \sigma_d + \sigma_{PE} + \sigma_{CN}$$

$$= \sigma_{nn'} + \sigma_{nf} + \sigma_{n\gamma} + \dots$$

Multiple Hauser-Feshbach



Possible decays

- Emission to a discrete level with **energy E_d**

$$\langle T_b(\beta) \rangle = T_{ij}^{J\pi}(\beta) \quad \text{given by the O.M.P.}$$

- Emission in the level continuum

$$\langle T_b(\beta) \rangle = \int_E^{E+\Delta E} T_{ij}^{J\pi}(\beta) \rho(E, J, \pi) dE$$

$\rho(E, J, \pi)$ **density of residual nucleus' levels** (J, π) with excitation energy E

- Emission of photons, fission

Specific treatment

Gamma decay

$$T^{k\lambda}(\varepsilon_\gamma) = 2\pi \int_E^{E+\Delta E} \Gamma^{k\lambda}(\varepsilon_\gamma) \rho(E) dE$$

$$= 2\pi f(\mathbf{k}, \lambda, \varepsilon_\gamma) \varepsilon_\gamma^{2\lambda+1}$$

k : transition type EM (E ou M)

λ : transition multipolarity

ε_γ : outgoing gamma energy

$f(\mathbf{k}, \lambda, \varepsilon_\gamma)$: gamma strength function (several models)

Decay selection rules from a level $J_i^{\pi_i}$ to a level $J_f^{\pi_f}$:

Pour $E\lambda$: $\pi_f = (-1)^\lambda \pi_i$

Pour $M\lambda$: $\pi_f = (-1)^{\lambda+1} \pi_i$ $|J_i - \lambda| \leq J_f \leq J_i + \lambda$

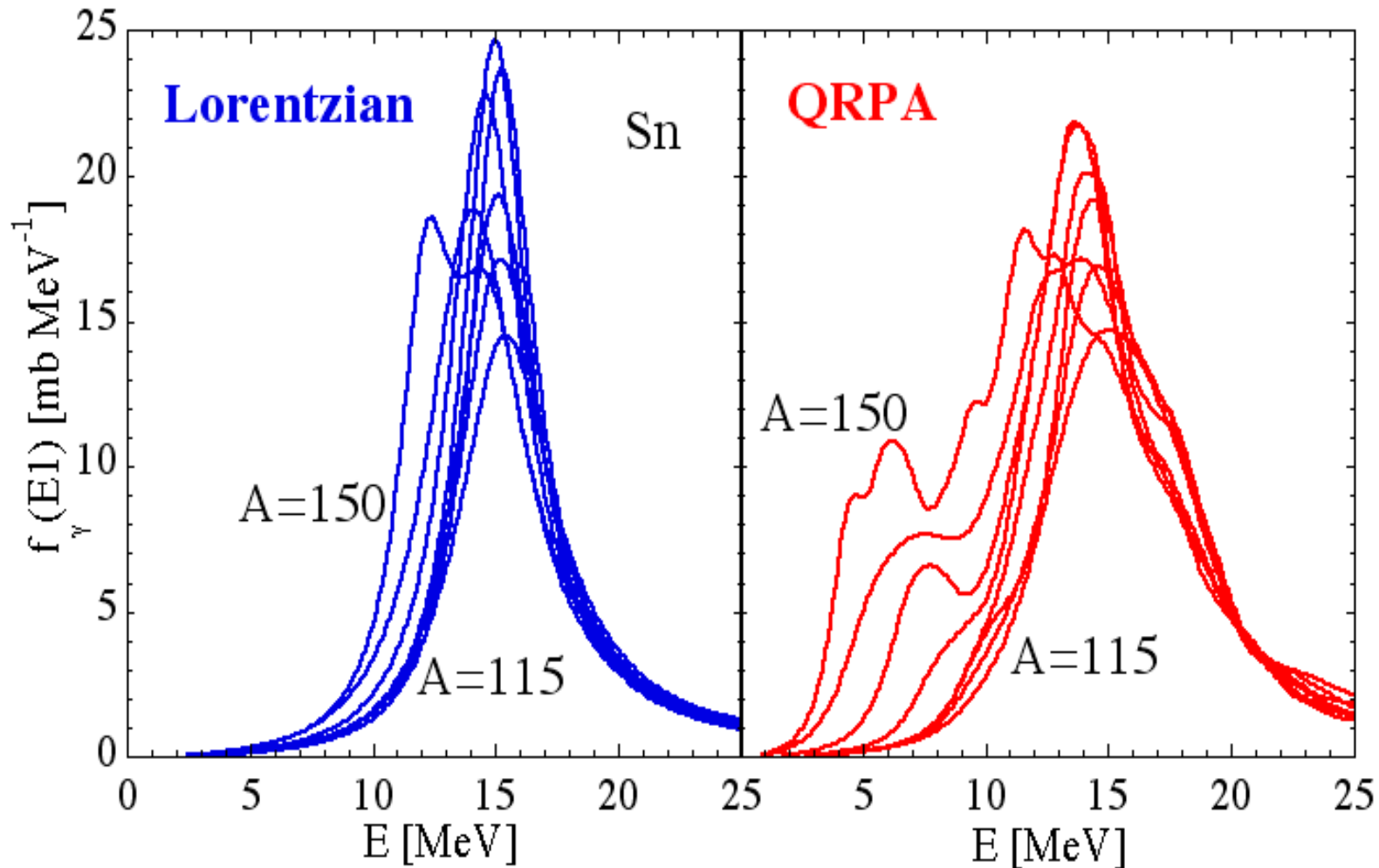
Renormalisation technique for thermal neutrons

$$\langle T_\gamma \rangle = \sum_{J_i, \pi_i} \sum_{k\lambda} \sum_{J_f, \pi_f} \int_0^{B_n} T^{k\lambda}(\varepsilon) \rho(B_n - \varepsilon, J_f, \pi_f) S(\lambda, J_i, \pi_i, J_f, \pi_f) d\varepsilon = 2\pi \langle \Gamma_\gamma \rangle \rho(B_n)$$

$$\langle T_\gamma \rangle = \mathbf{C} \sum_{J_i, \pi_i} \sum_{k\lambda} \sum_{J_f, \pi_f} \int_0^{B_n} T^{k\lambda}(\varepsilon) \rho(B_n - \varepsilon, J_f, \pi_f) S(\lambda, J_i, \pi_i, J_f, \pi_f) d\varepsilon = 2\pi \langle \Gamma_\gamma \rangle \frac{1}{D_0}$$

experiment

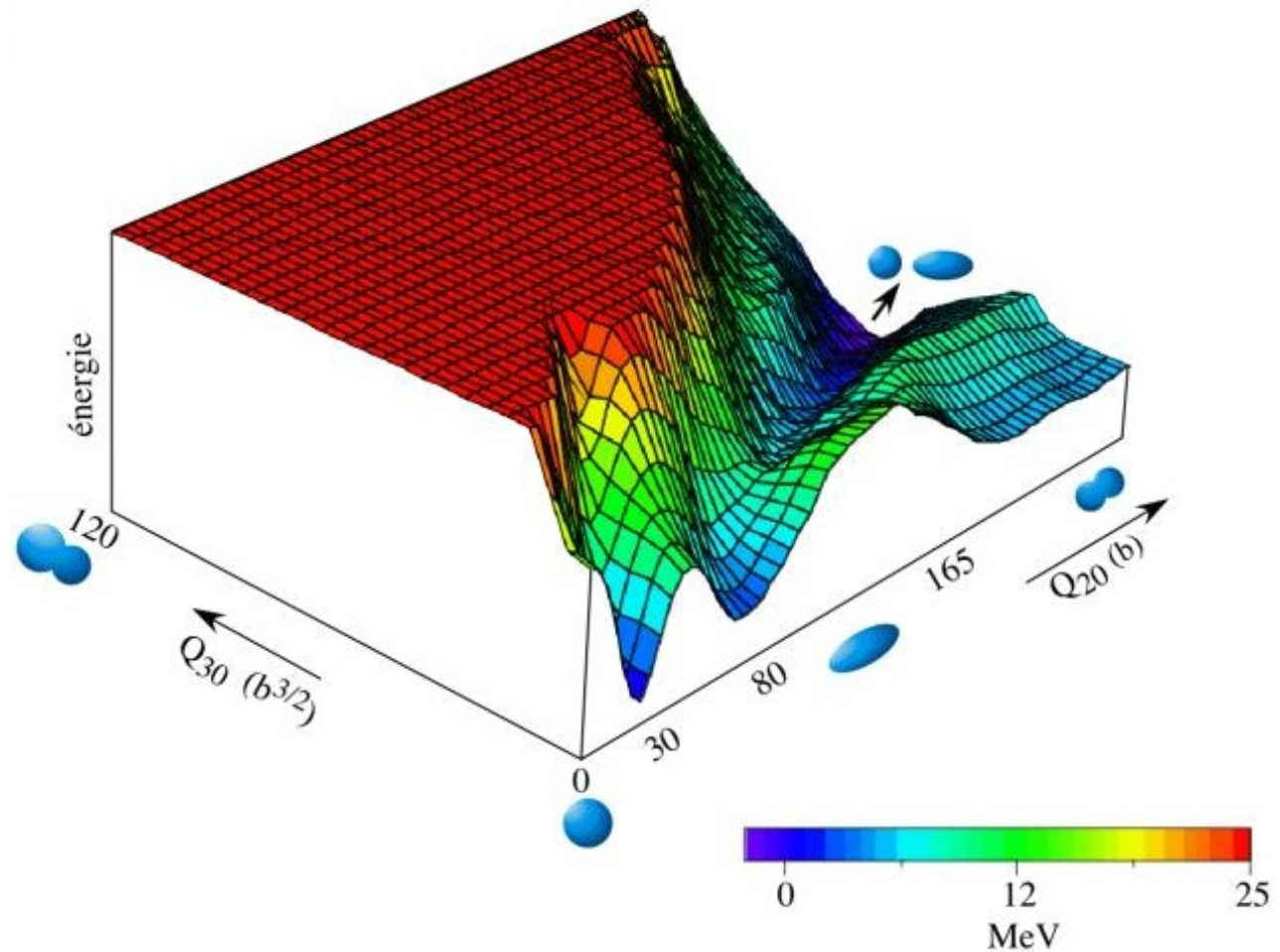
Gamma strength functions options



See S. Goriely & E. Khan, *NPA* 706 (2002) 217.

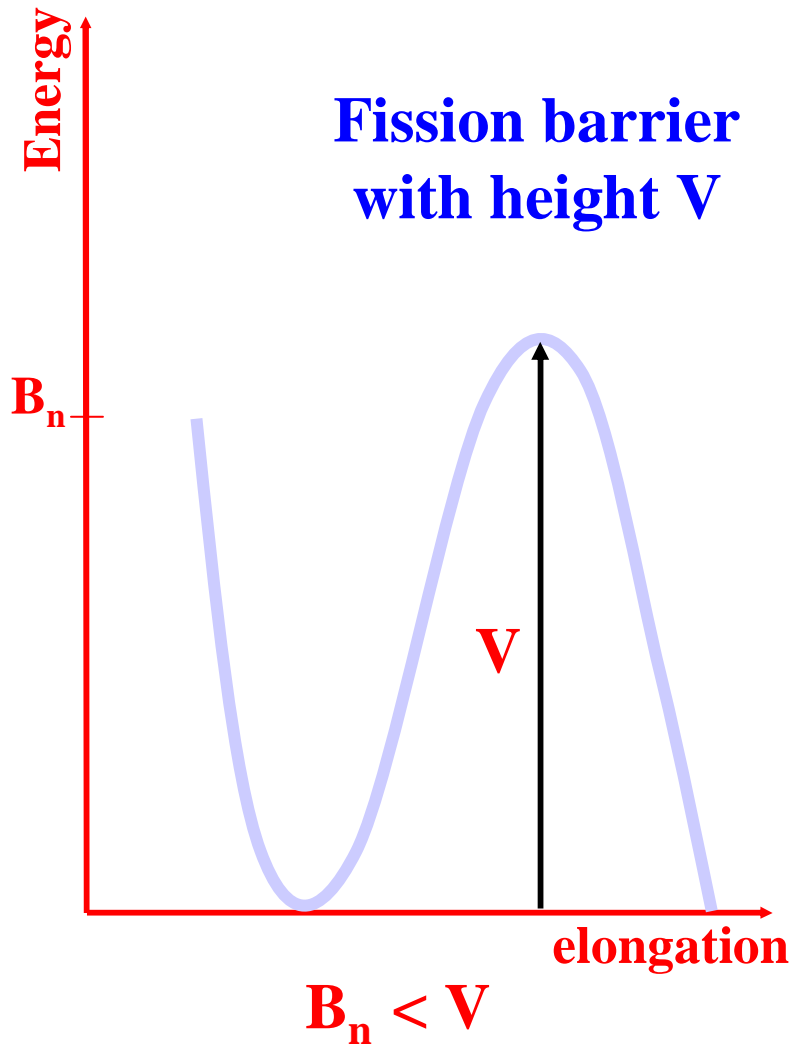
S. Goriely et al., *NPA* 739 (2004) 331.

The fission process

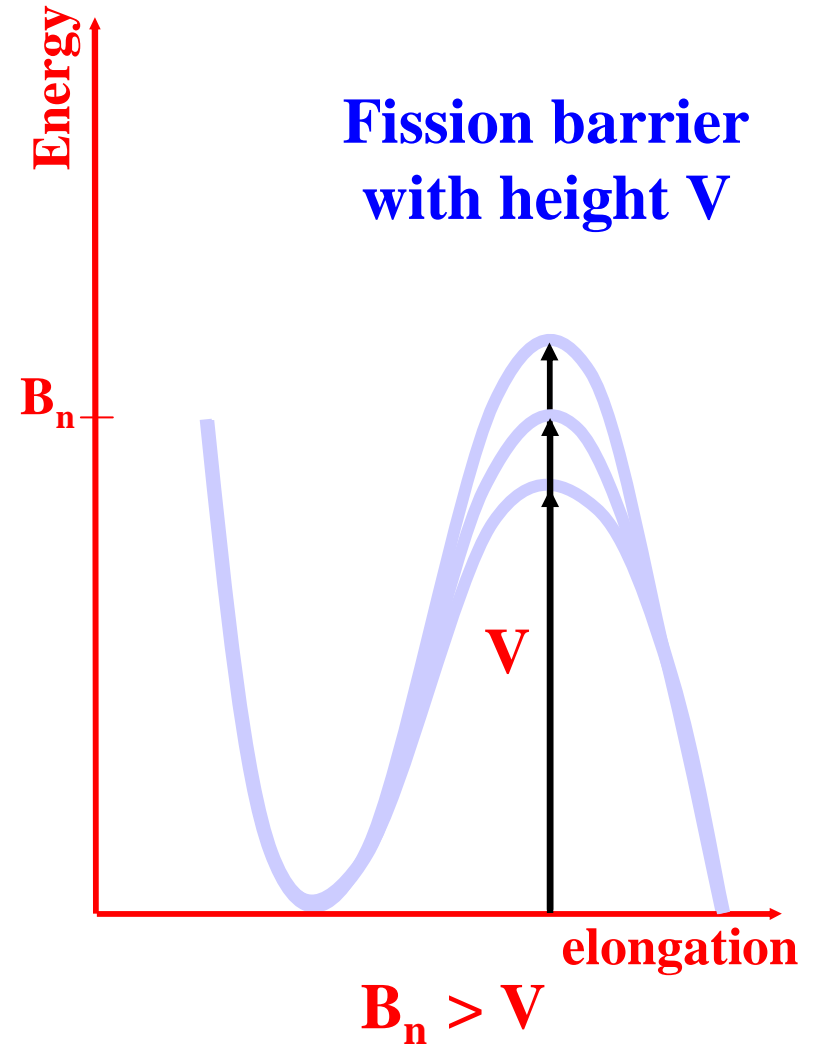
Surface ^{238}U 



Fissile/Fertile

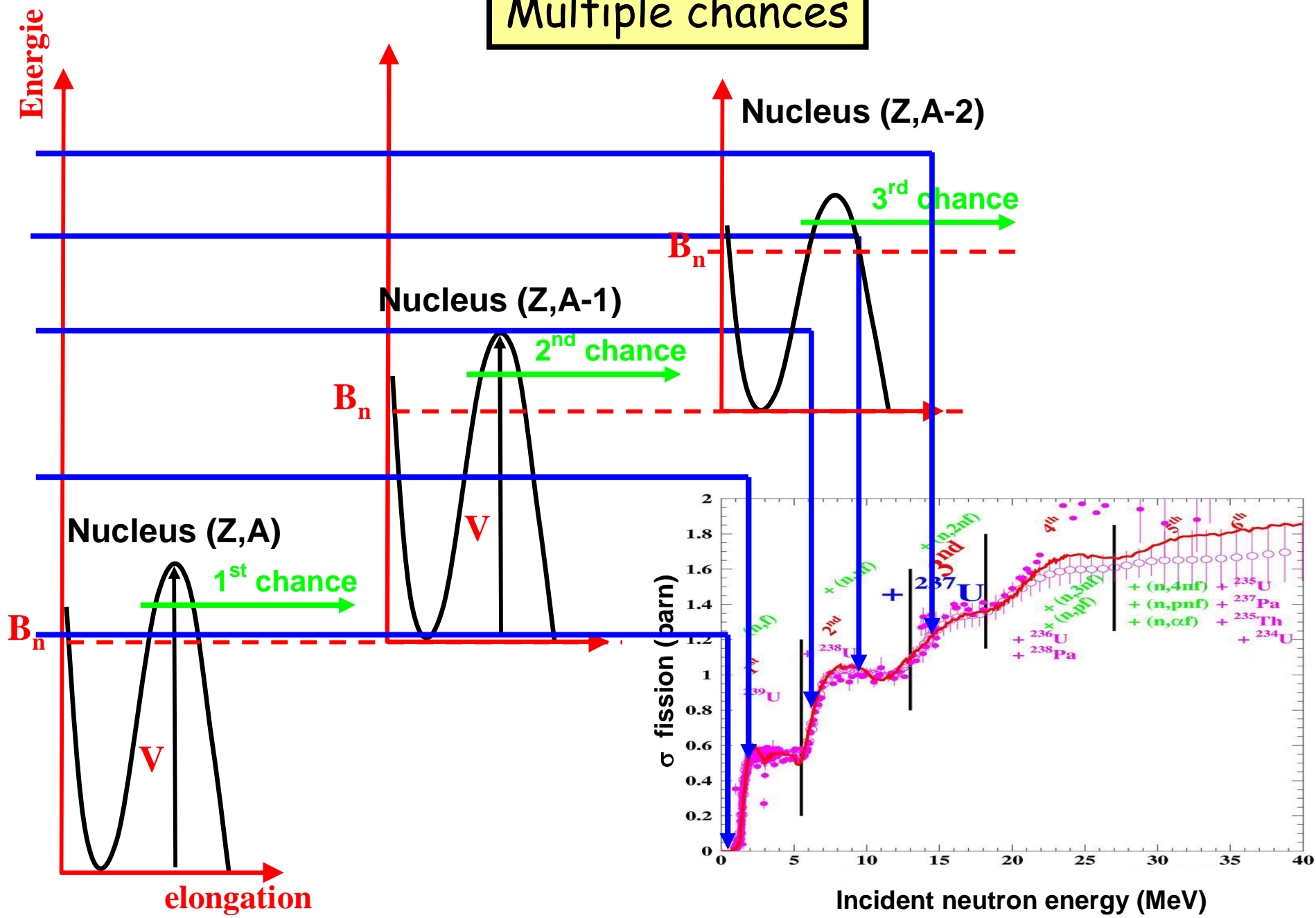


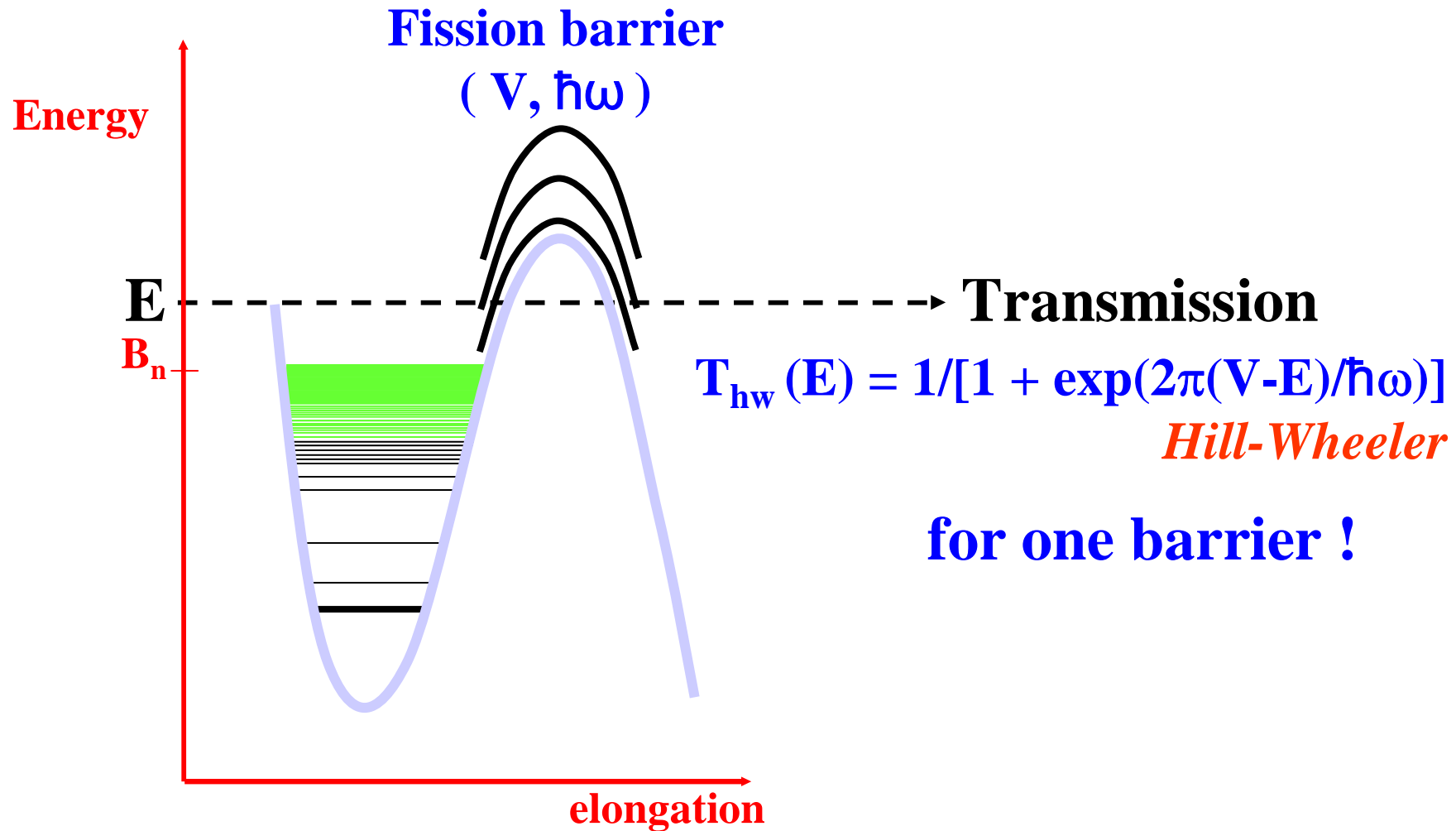
Fertile target (^{238}U)



Fissile target (^{235}U)

Multiple chances





+ transition state on top of the barrier !

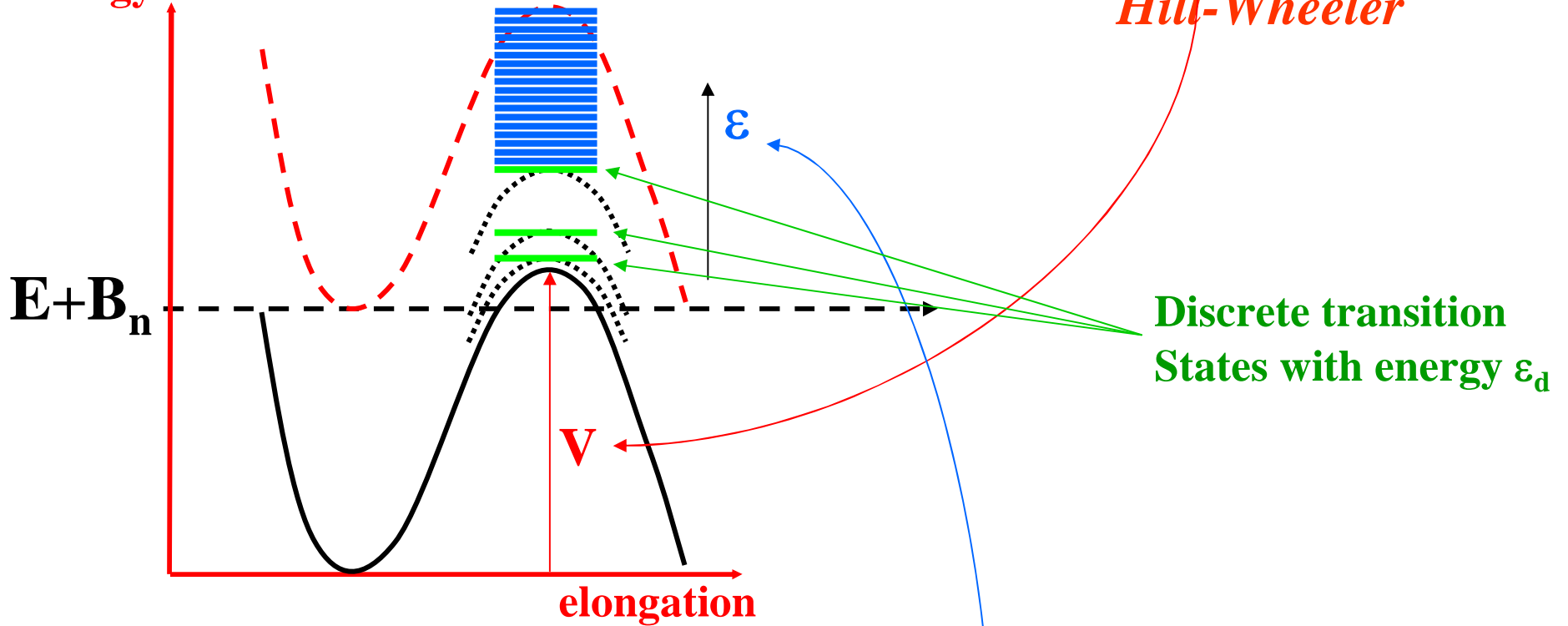
Fission transmission coefficients



Energy

$$T_{hw}(E) = 1/[1 + \exp(2\pi(\mathbf{V}-E)/\hbar\omega)]$$

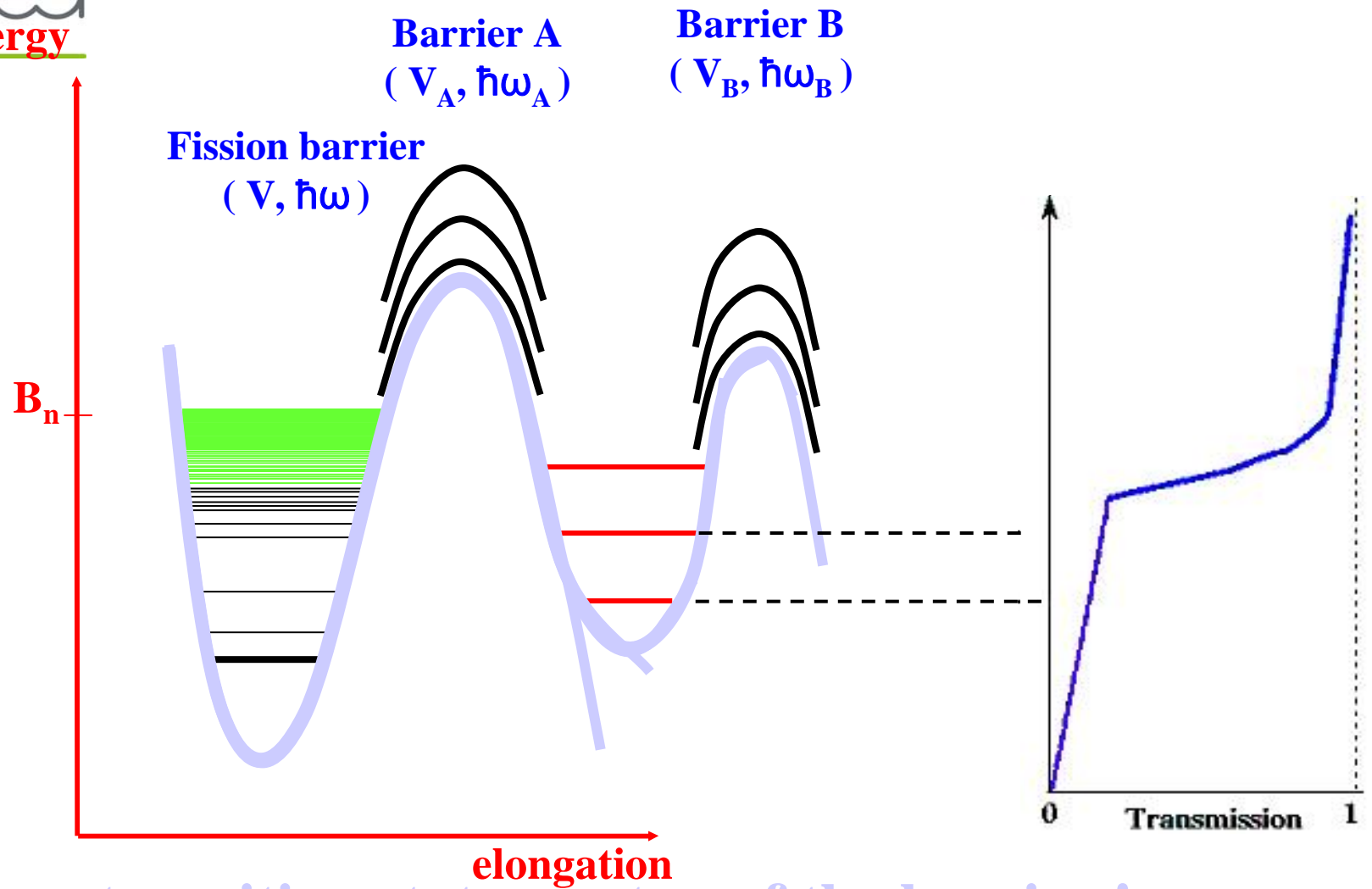
Hill-Wheeler



Discrete transition States with energy ϵ_d

$$T_f(E, J, \pi) = \sum_{\substack{\text{discrets} \\ J, \pi}} T_{hw}(E - \epsilon_d) + \int_{E_s}^{E+B_n} \rho(\epsilon, J, \pi) T_{hw}(E - \epsilon) d\epsilon$$

Multiple humped barriers



+ transition states on top of the barrier !

+ transition states on top of each barrier !

+ class II states in the intermediate well !



Multiple humped barriers

Two barriers A and B

$$T_f = \frac{T_A T_B}{T_A + T_B}$$

Three barriers A, B et C

$$T_f = \frac{\frac{T_A T_B}{T_A + T_B} \times T_C}{\frac{T_A T_B}{T_A + T_B} + T_C}$$

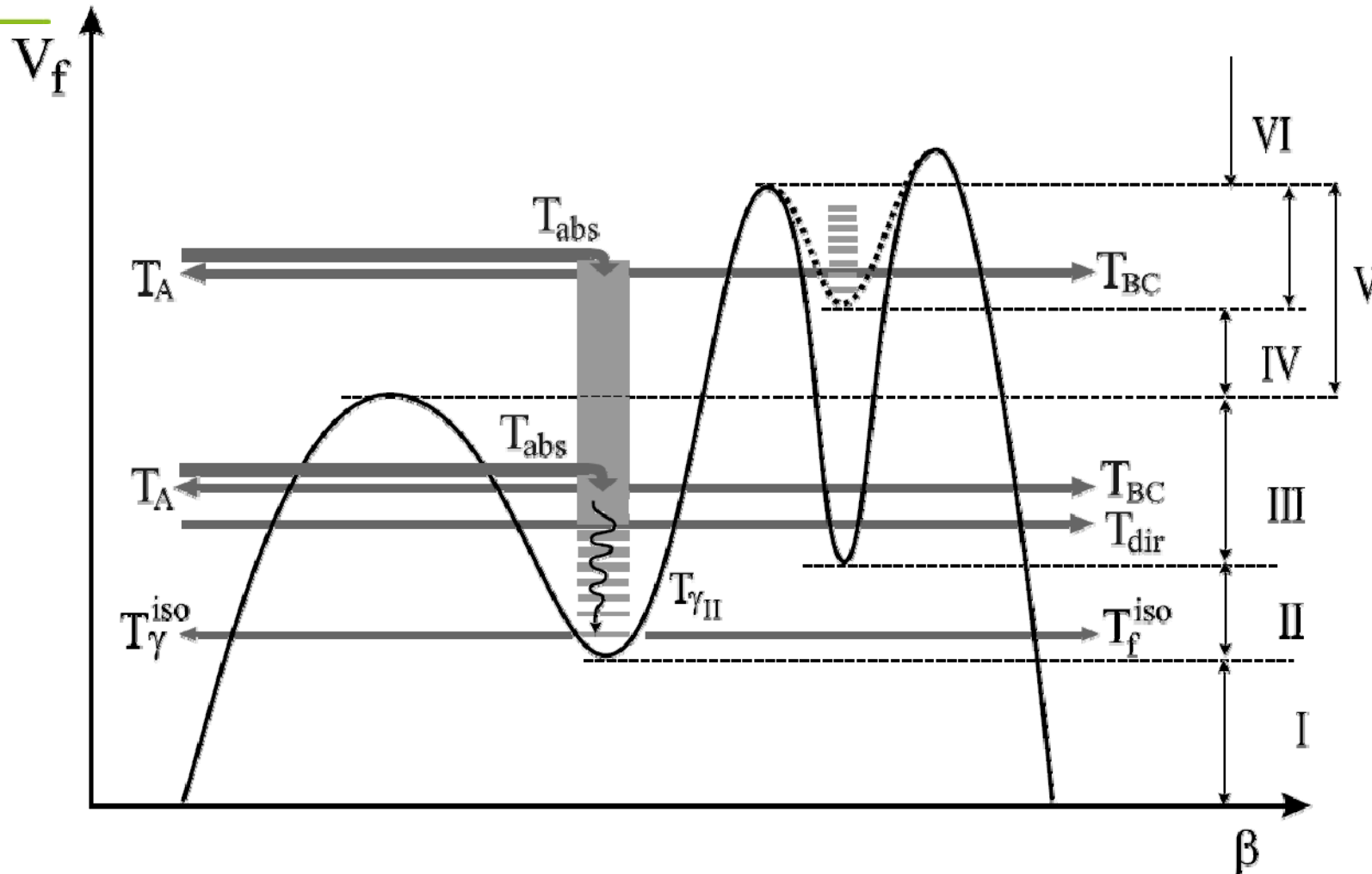
Resonant transmission



$$T_f = \frac{T_A T_B}{T_A + T_B} \times \frac{4}{T_A + T_B}$$

More exact expressions in Sin et al., PRC 74 (2006) 014608

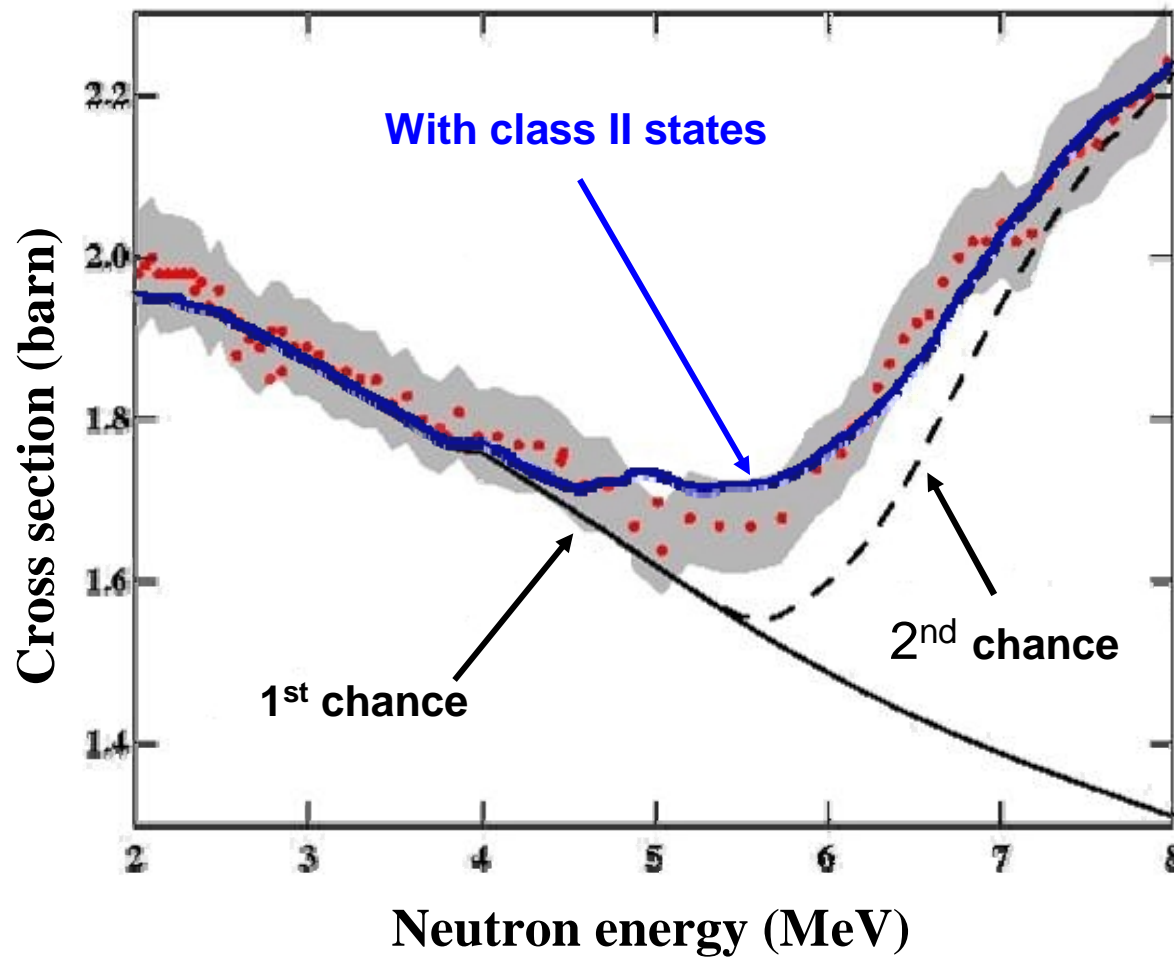
Multiple humped barriers with maximum complexity



See in Sin et al., PRC 74 (2006) 014608

Bjornholm and Lynn, Rev. Mod. Phys. 52 (1980) 725.

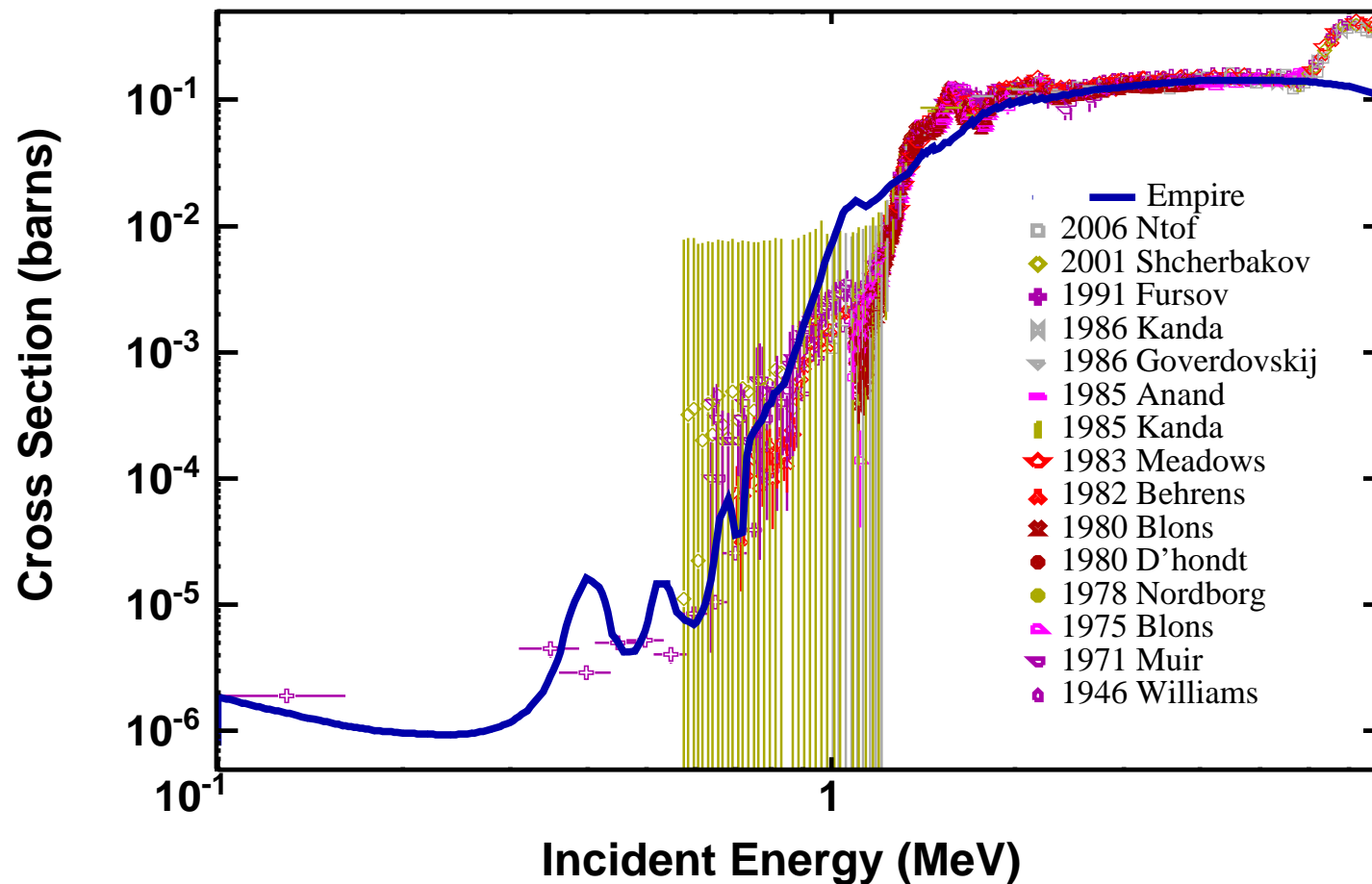
Impact of class II states

 ^{239}Pu (n,f)

Impact of class II states

Case of a fertile nucleus

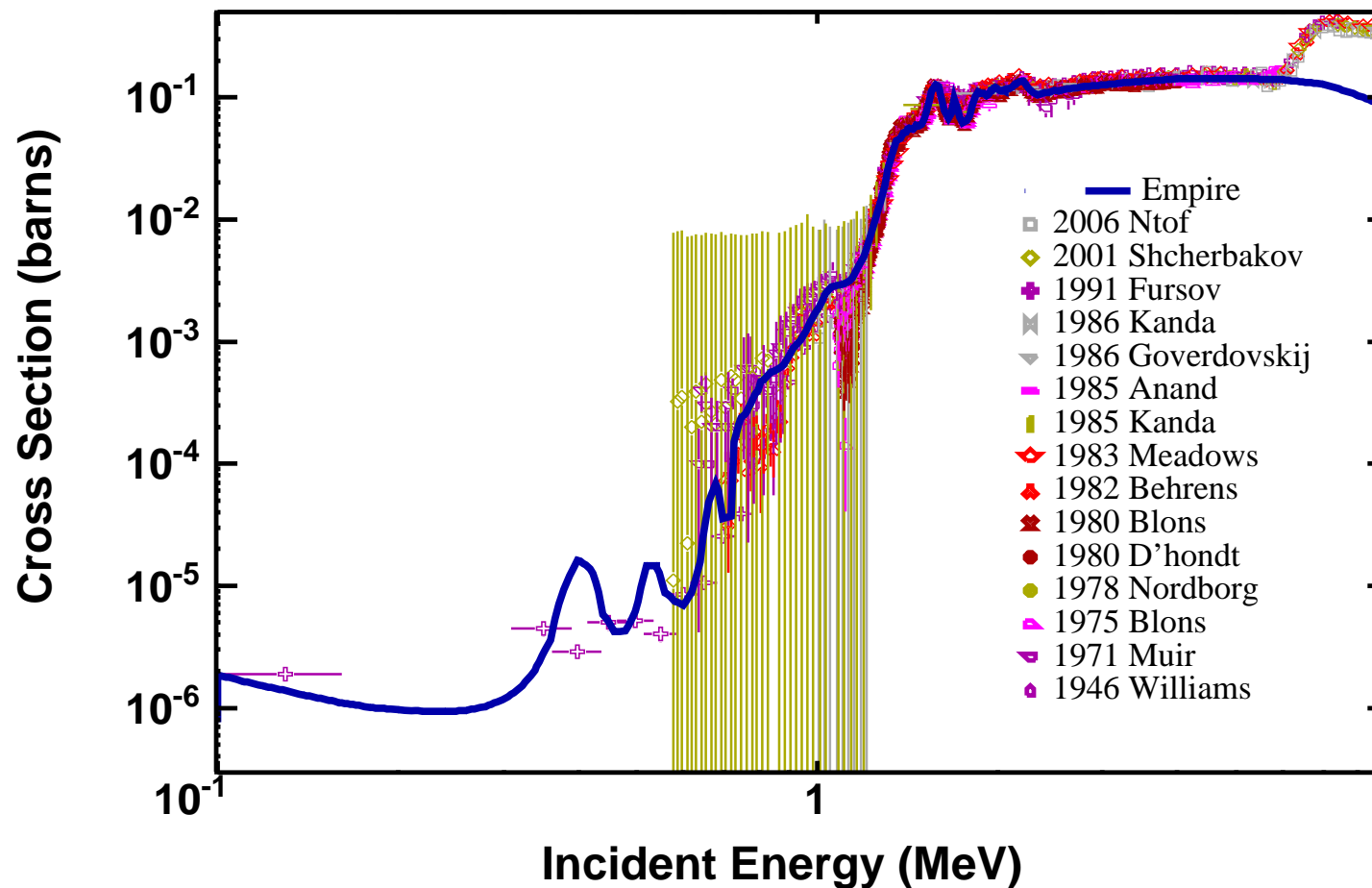
Partially damped class II states. No class III states (fully damped).



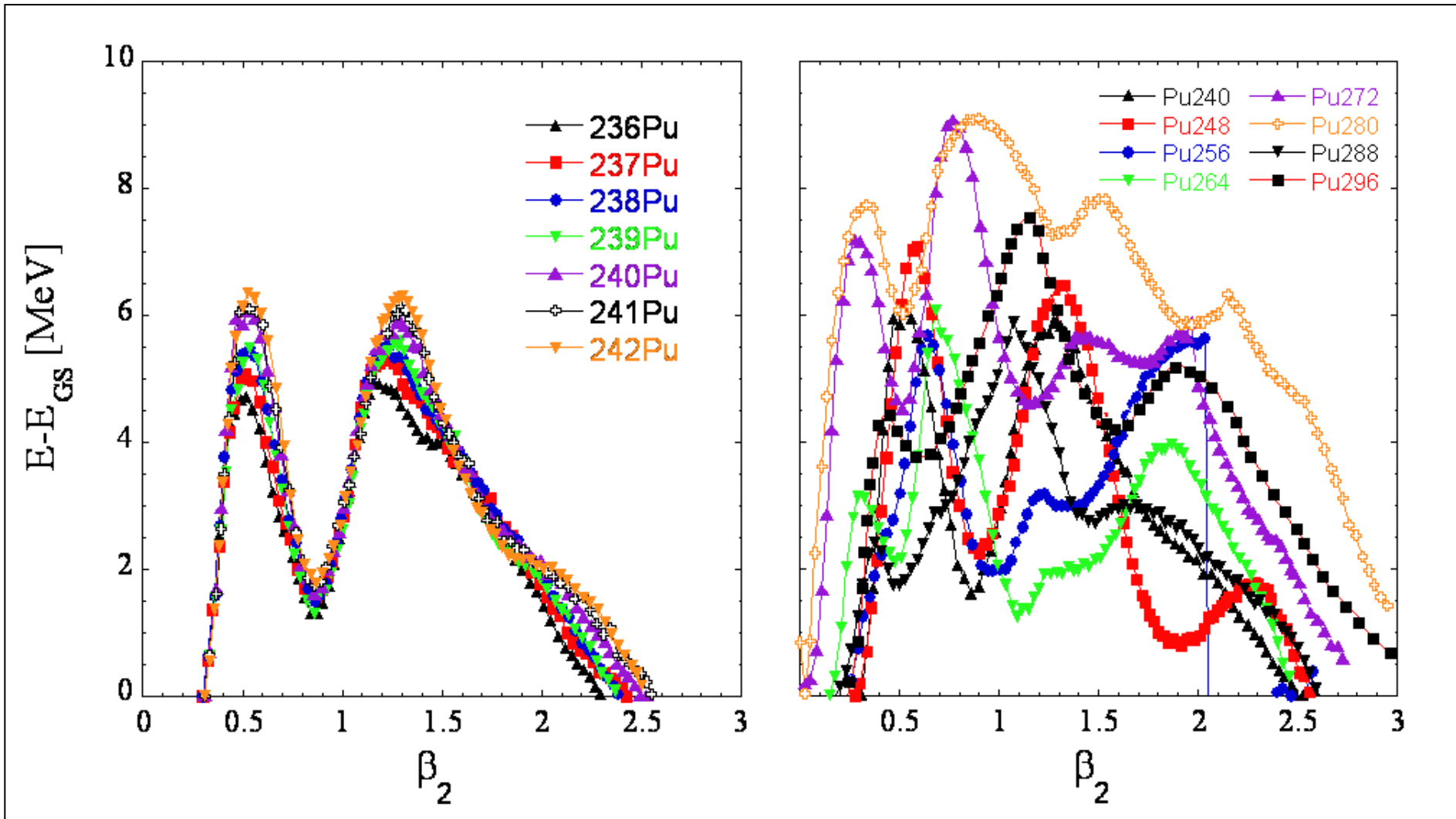
Impact of class II+III states

Case of a fertile nucleus

Class II + III states. Partial damping.

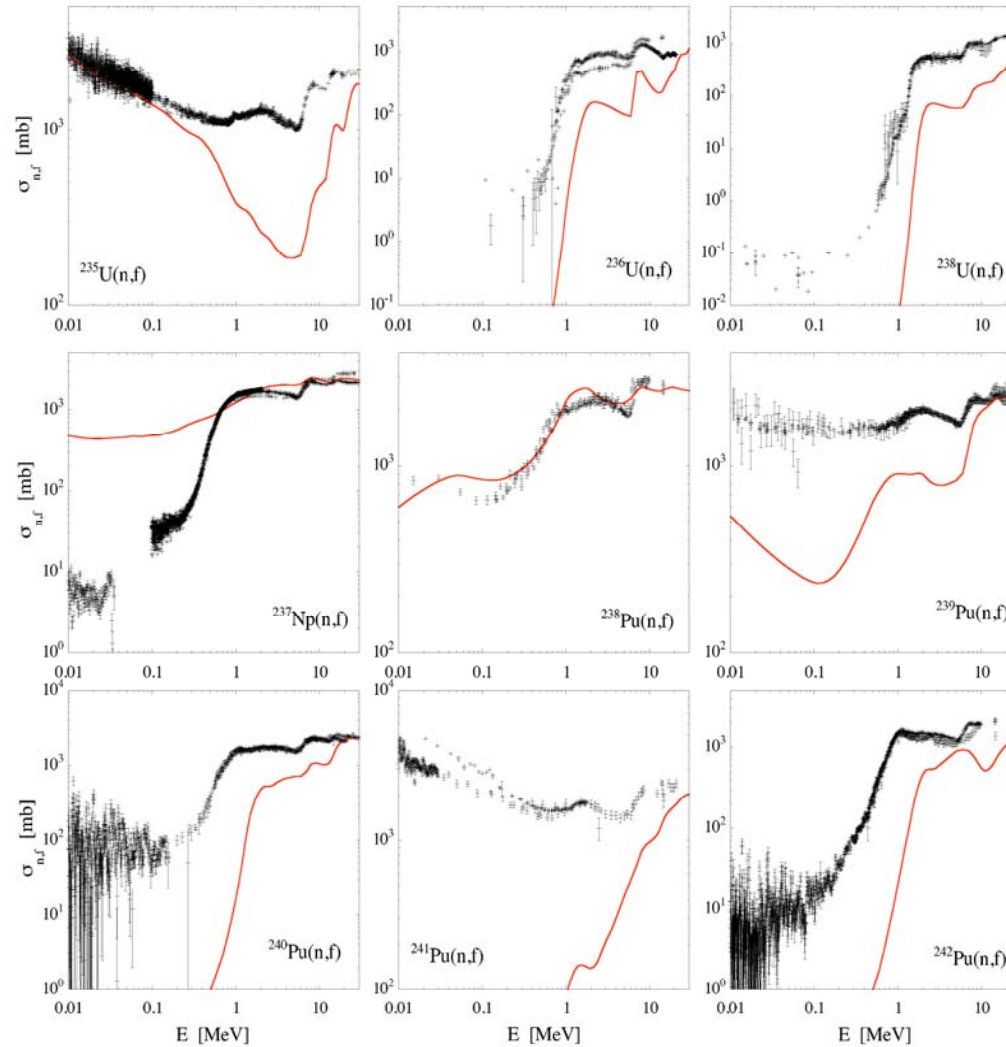


Microscopic fission barrier shapes



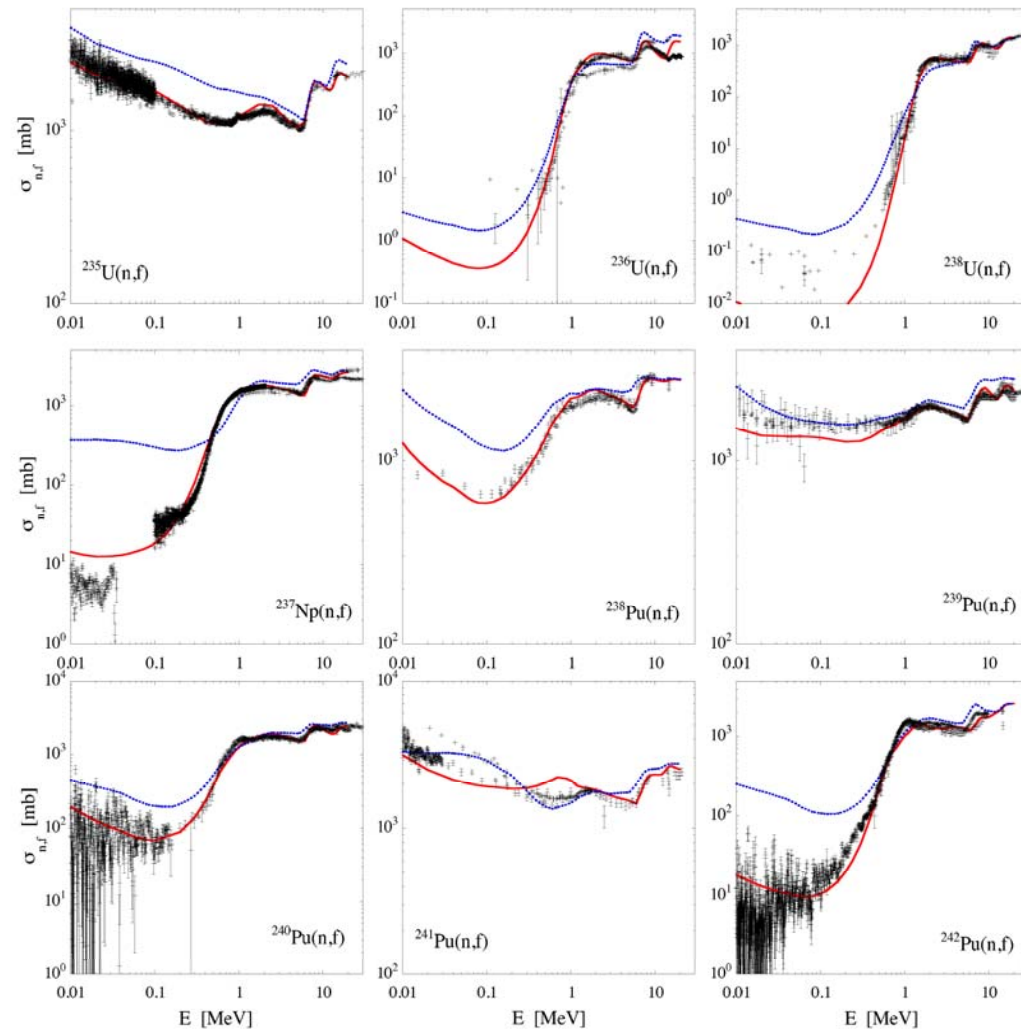
⇒ For exotic nuclei : strong deviations from Hill-Wheeler.

Microscopic fission cross sections



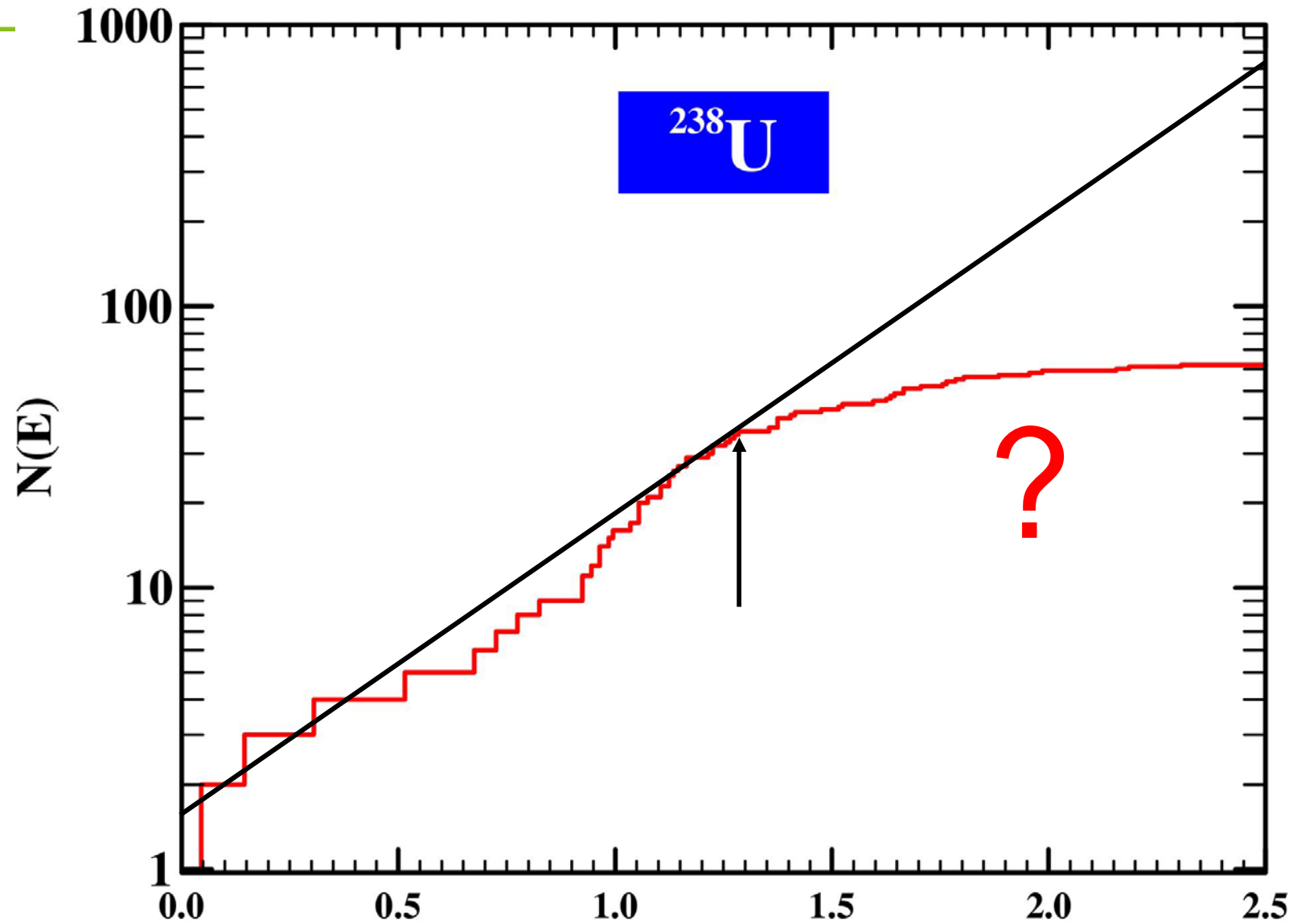
⇒ Default calculations not sufficient for applications.

Microscopic fission cross sections

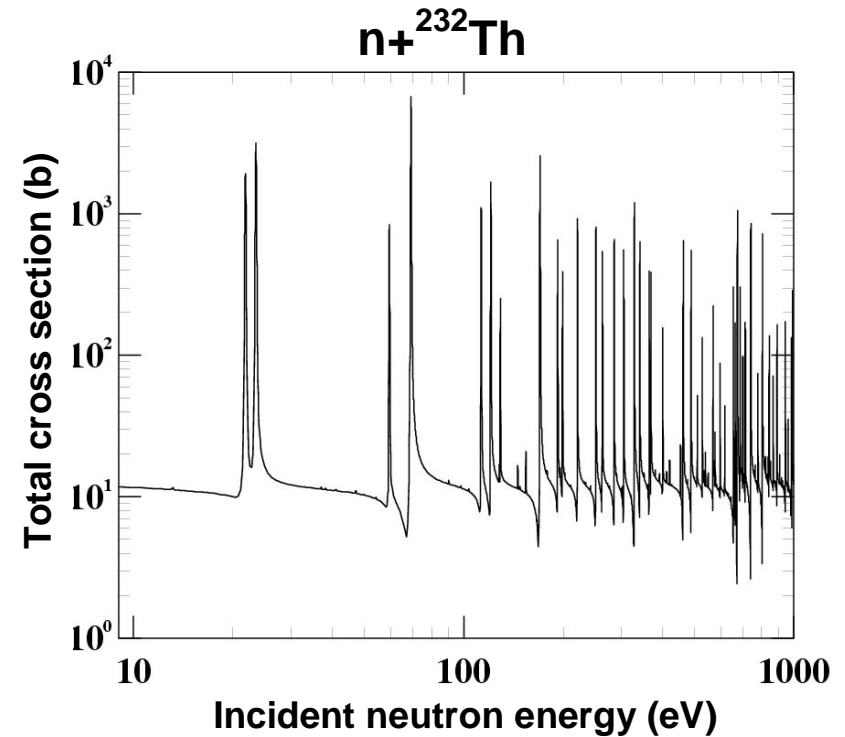
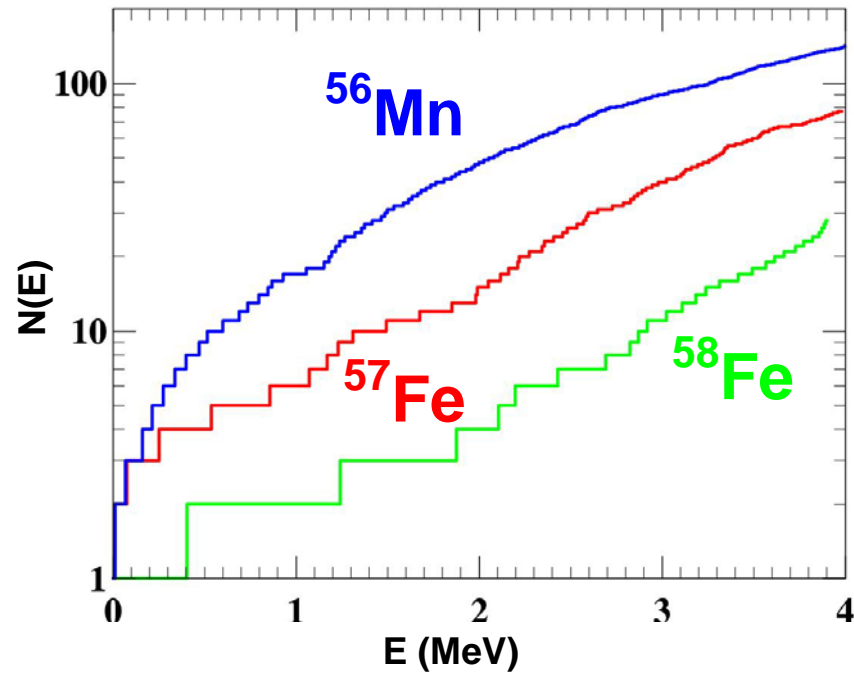


⇒ Not ridiculous after few adjustments.

Level densities : principle



Level densities : qualitative aspects



- Exponential increase of the cumulated number of discrete levels $N(E)$ with energy

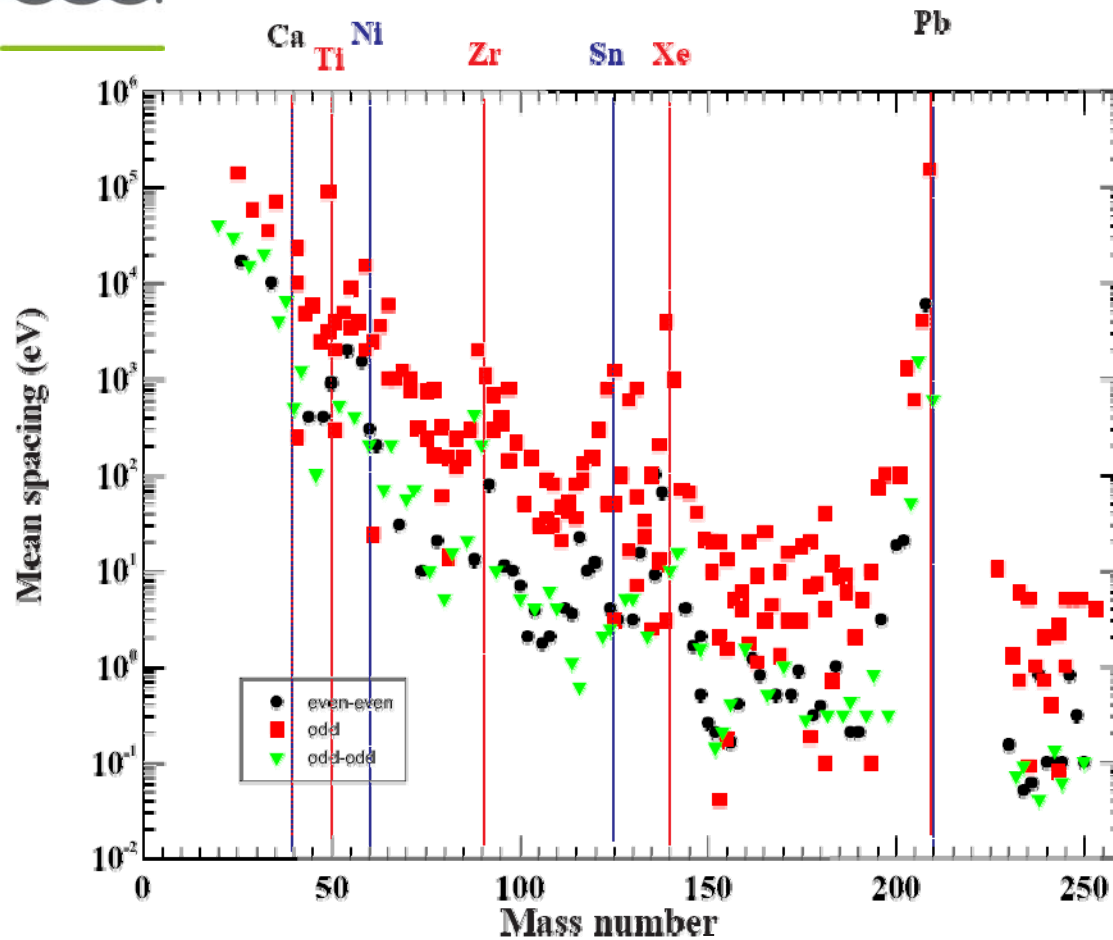
$$\Rightarrow \rho(E) = \frac{dN(E)}{dE} \text{ Increases exponentially}$$

\Rightarrow odd-even effects

- Mean spacings of s-wave neutron resonances at B_n of the order of few eV

$$\Rightarrow \rho(B_n) \text{ of the order of } 10^4 - 10^6 \text{ levels / MeV}$$

Level densities : qualitative aspects



Ilijin et al., NPA 543 (1992) 517.

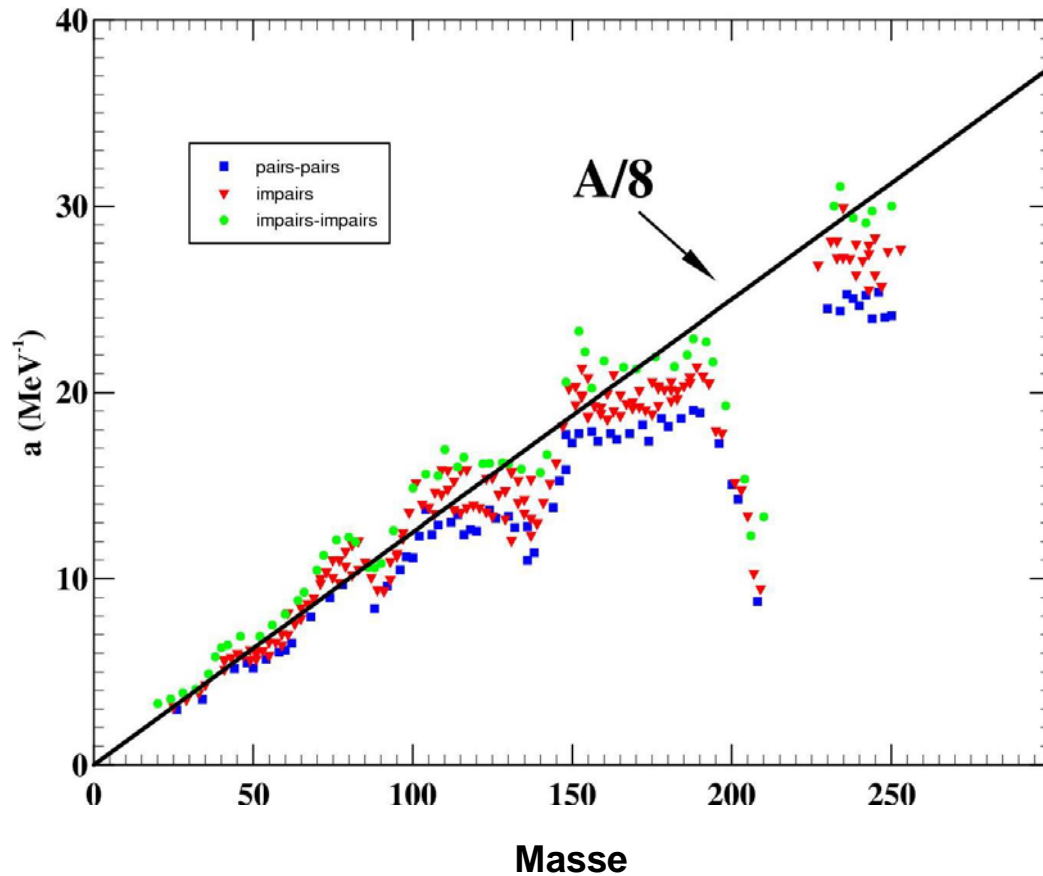
⇒ **Mass dependency**
Odd-even effects
Shell effects

$$\frac{1}{D_0} = \rho(B_n, 1/2, \pi_t) \text{ for an even-even target}$$

$$= \rho(B_n, I_t + 1/2, \pi_t) + \rho(B_n, I_t - 1/2, \pi_t) \text{ otherwise}$$

Level densities : quantitative analysis

$$\rho(\mathbf{U}, \mathbf{J}, \pi) = \frac{1}{2} \frac{\sqrt{\pi}}{12} \frac{\exp(2\sqrt{aU})}{a^{1/4} U^{5/4}} \frac{2J+1}{2\sqrt{2\pi} \sigma^3} \exp\left[-\frac{(J+1/2)^2}{2\sigma^2}\right]$$

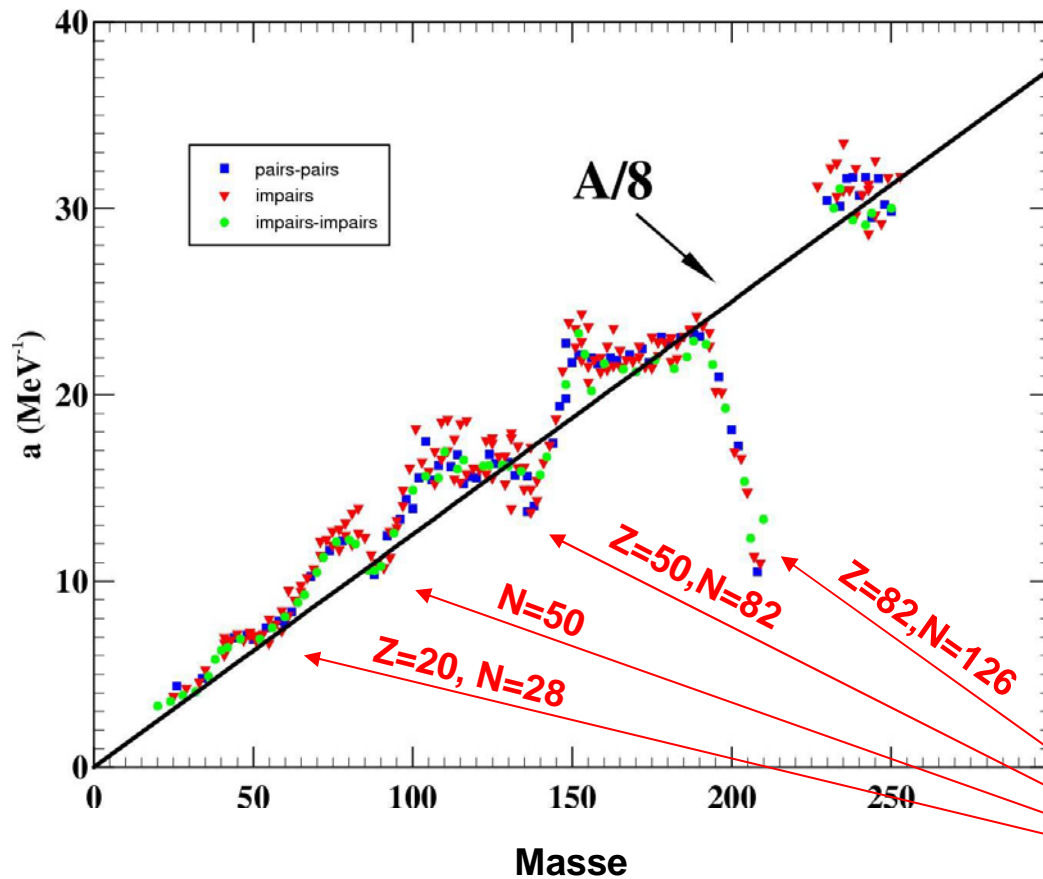


$$\sigma^2 = I_{\text{rig}} \sqrt{\frac{U}{a}}$$

\Rightarrow odd-even effects

Level densities : quantitative analysis

$$\rho(\mathbf{U}, \mathbf{J}, \pi) = \frac{1}{2} \frac{\sqrt{\pi}}{12} \frac{\exp(2\sqrt{aU})}{a^{1/4} U^{5/4}} \frac{2J+1}{2\sqrt{2\pi} \sigma^3} \exp - \left[\frac{(J+1/2)^2}{2\sigma^2} \right]$$



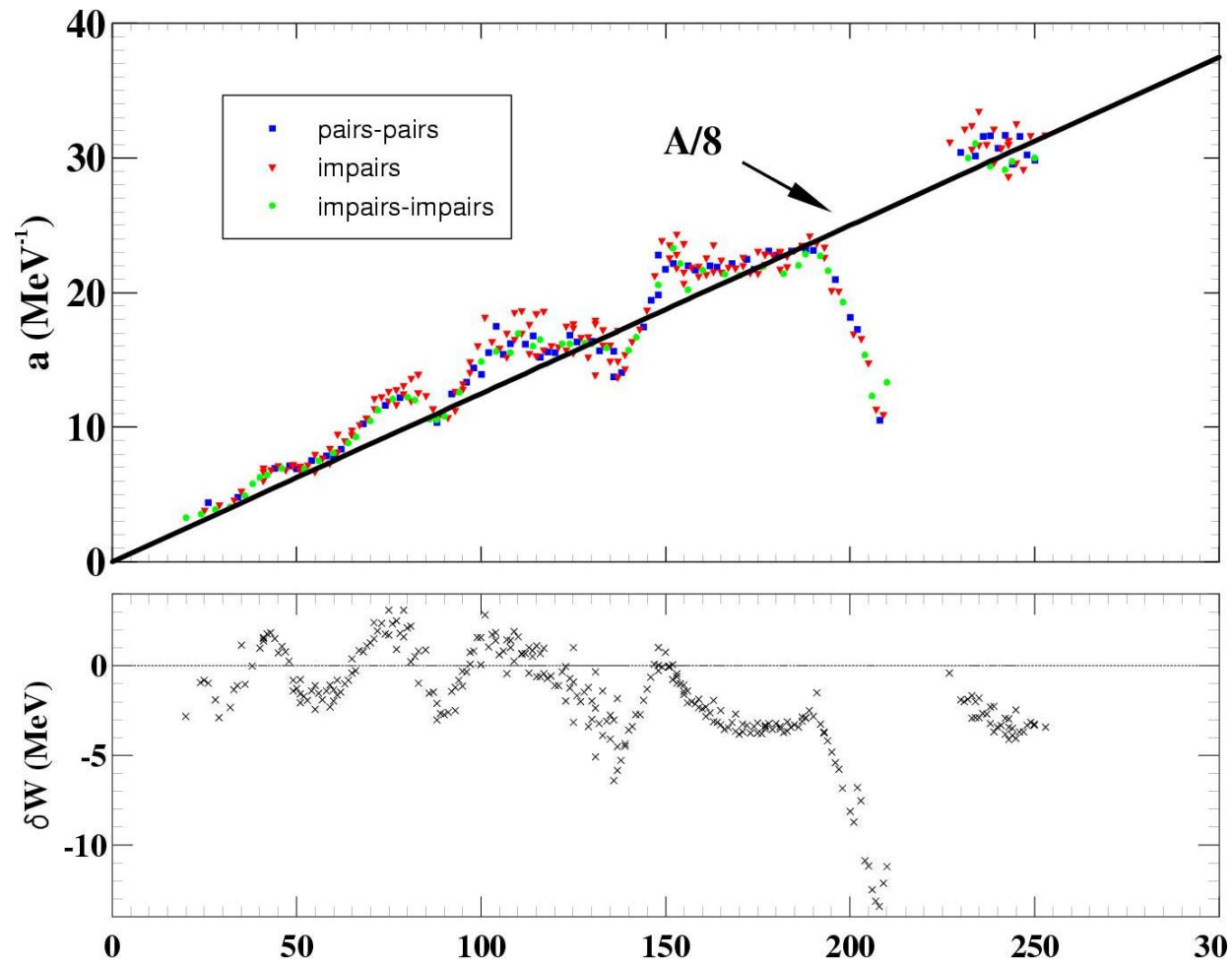
Odd-even effects
accounted for

$$U \rightarrow U^* = U - \Delta$$

$$\Delta = \begin{cases} 0 & \text{odd-odd} \\ 12/\sqrt{A} & \text{odd-even} \\ 24/\sqrt{A} & \text{even-even} \end{cases}$$

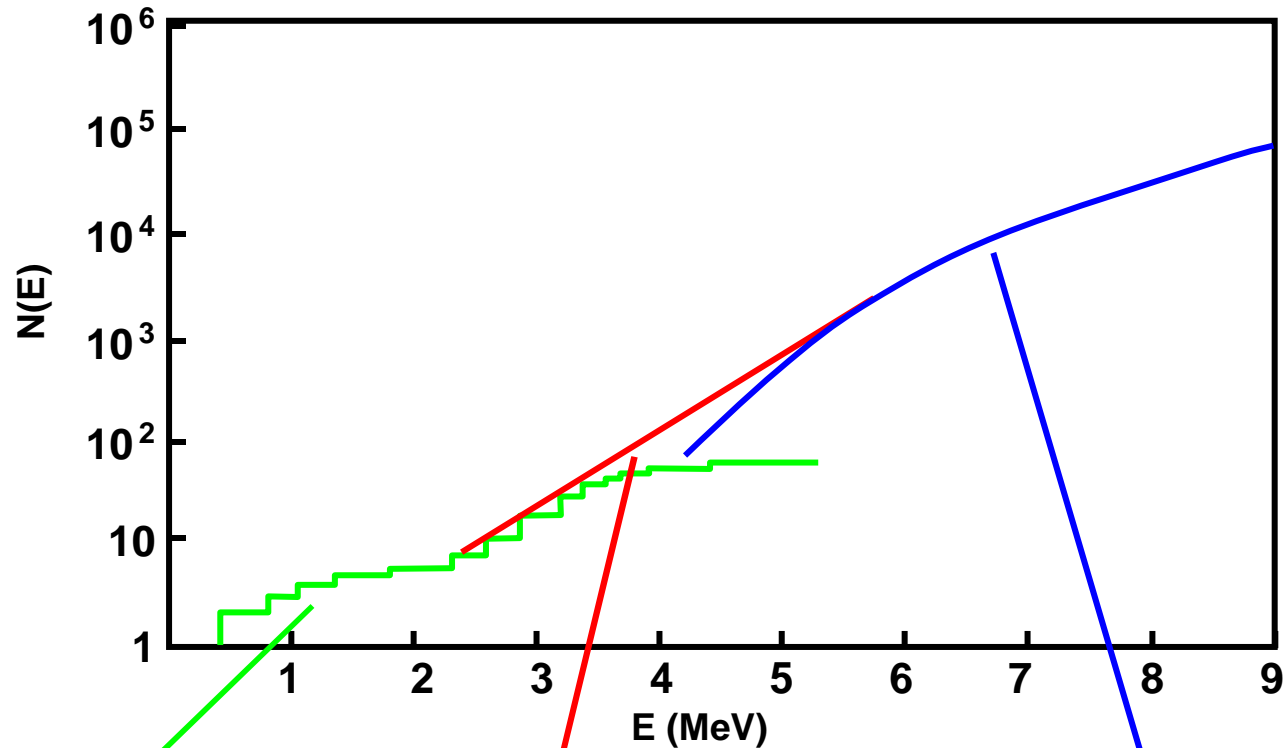
Shell effects

Level densities : Ignatyuk formula



$$a(N, Z, U^*) = \tilde{a}(A) \left[1 + \delta W(N, Z) \frac{1 - \exp(-\gamma U^*)}{U^*} \right]$$

Full description of level densities



Discrete levels
(spectroscopy)

Temperature law

$$N(E) = \exp\left(\frac{E - E_0}{T}\right)$$

Fermi gaz (adjusted at B_n)

$$\rho(E) = \alpha \frac{\exp\left(2\sqrt{aU^*}\right)}{a^{1/4}U^{*5/4}}$$

The combinatorial method

See PRC 78 (2008) 064307 for details

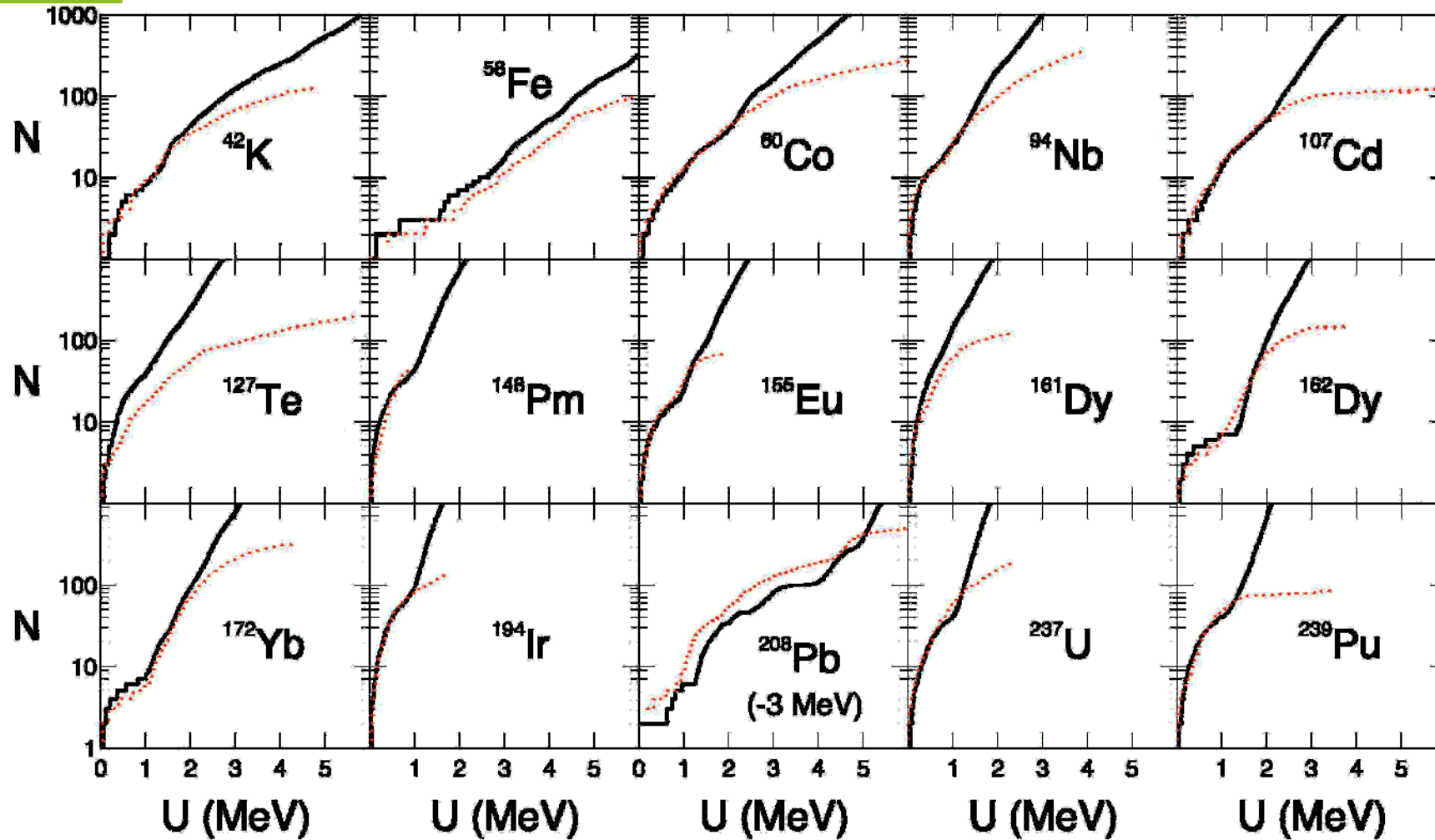
- HFB + effective nucleon-nucleon interaction \Rightarrow single particle level schemes
- Combinatorial calculation \Rightarrow intrinsic p-h and total state densities $\omega_i(U, K, \pi)$
- Collective effects \Rightarrow from **state** to **level** densities $\rho(U, J, \pi)$

~~2006 Approximation : 1) construction of rotational bands
2) multiplication by vibrational enhancement~~

Current treatment : 1) folding of intrinsic and vibrational state densities
2) construction of rotational bands

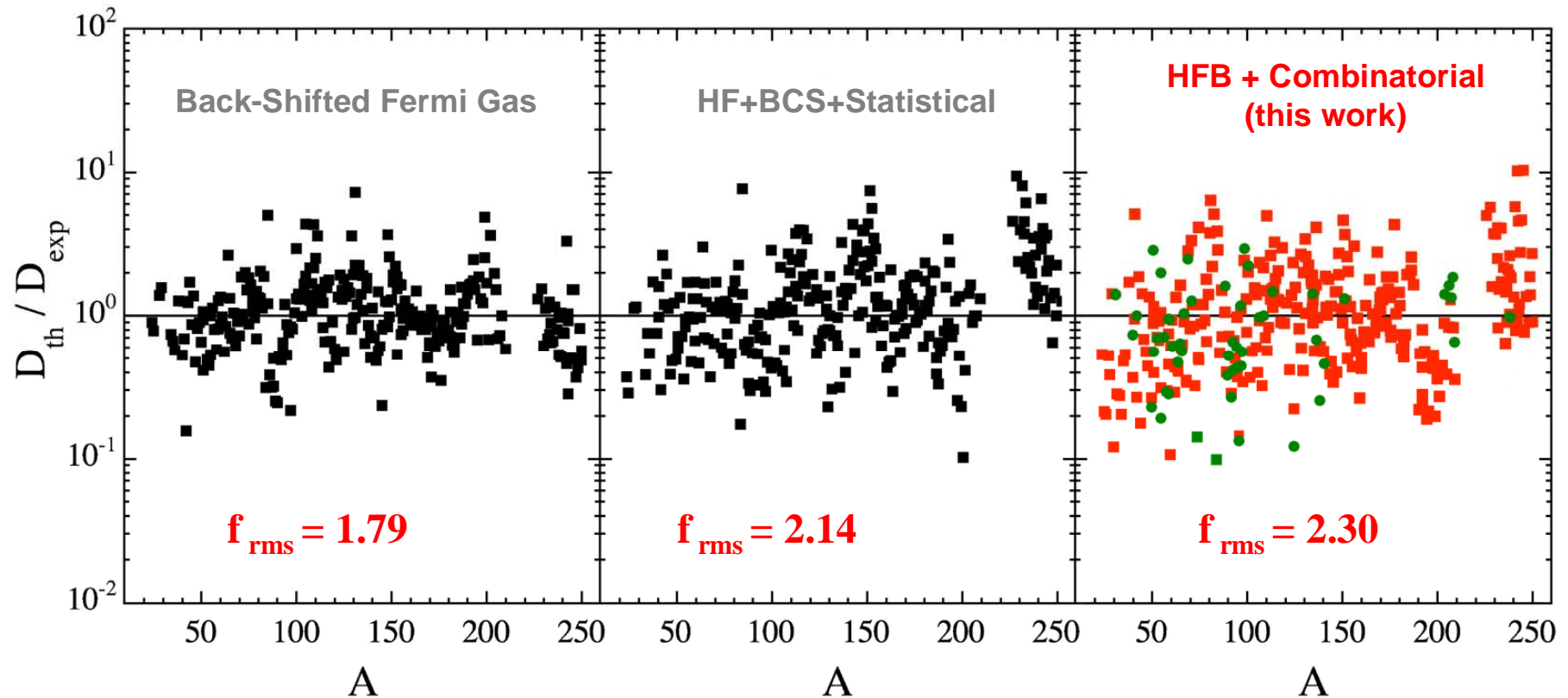
- **Phenomenological** transition for deformed/spherical nucleus

Results for cumulated histograms



→ Structures typical of non-statistical feature

D values (s-waves & p-waves)



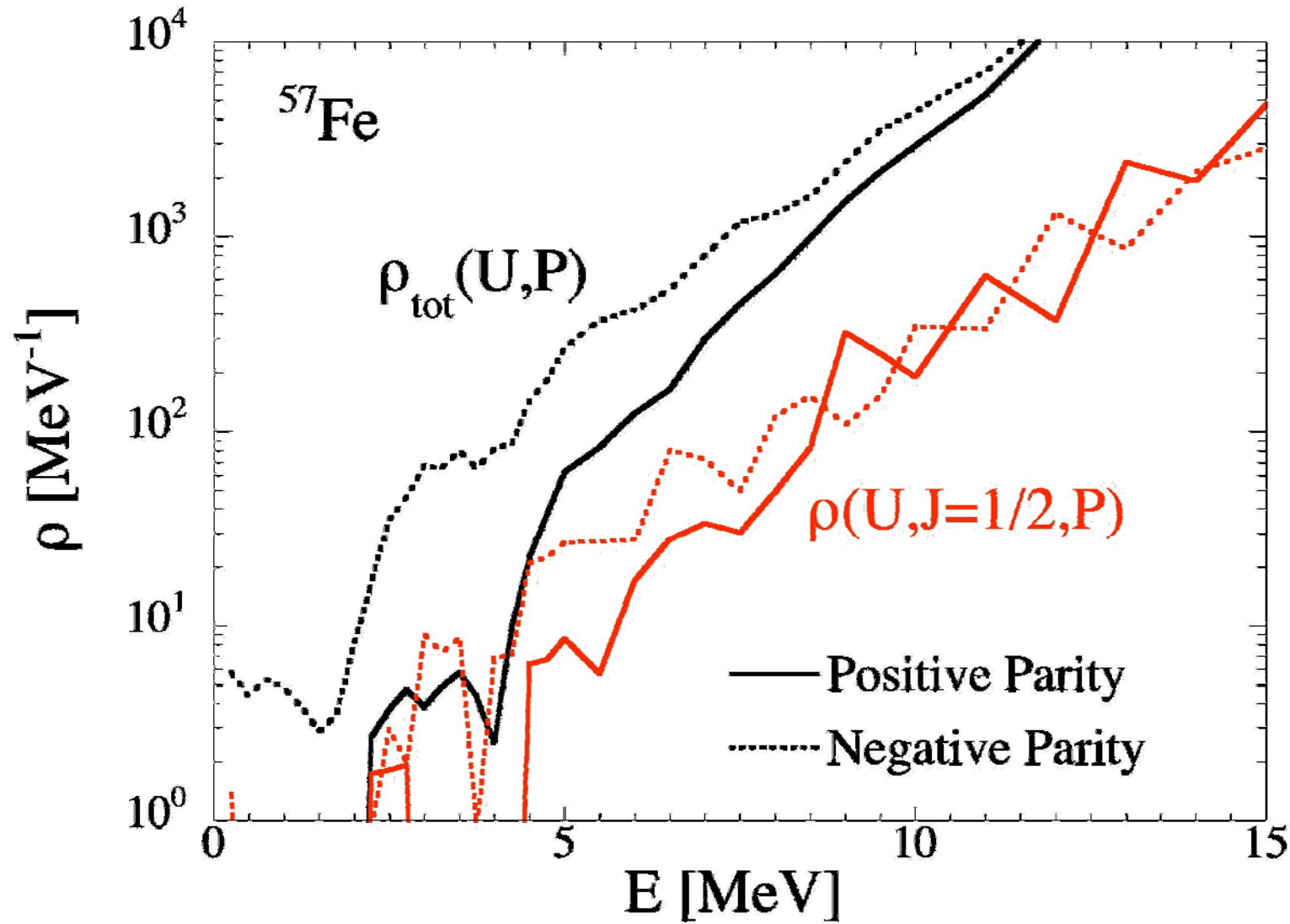
➔ Description similar to that obtained with other **global** approaches



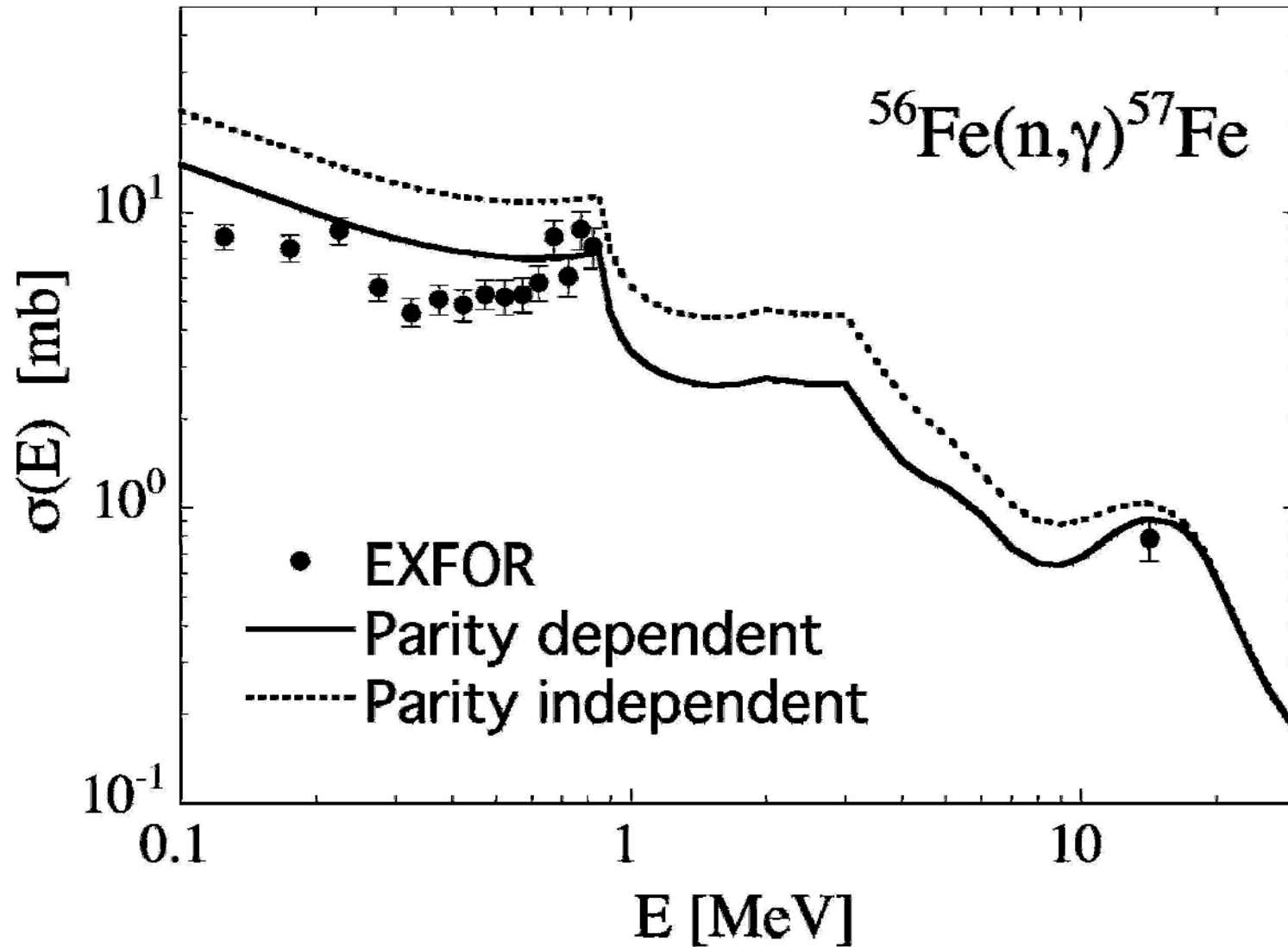
Combinatorial level densities

*Talys deals with realistic (non statistical) **parity** and spin distributions*

Combinatorial level densities



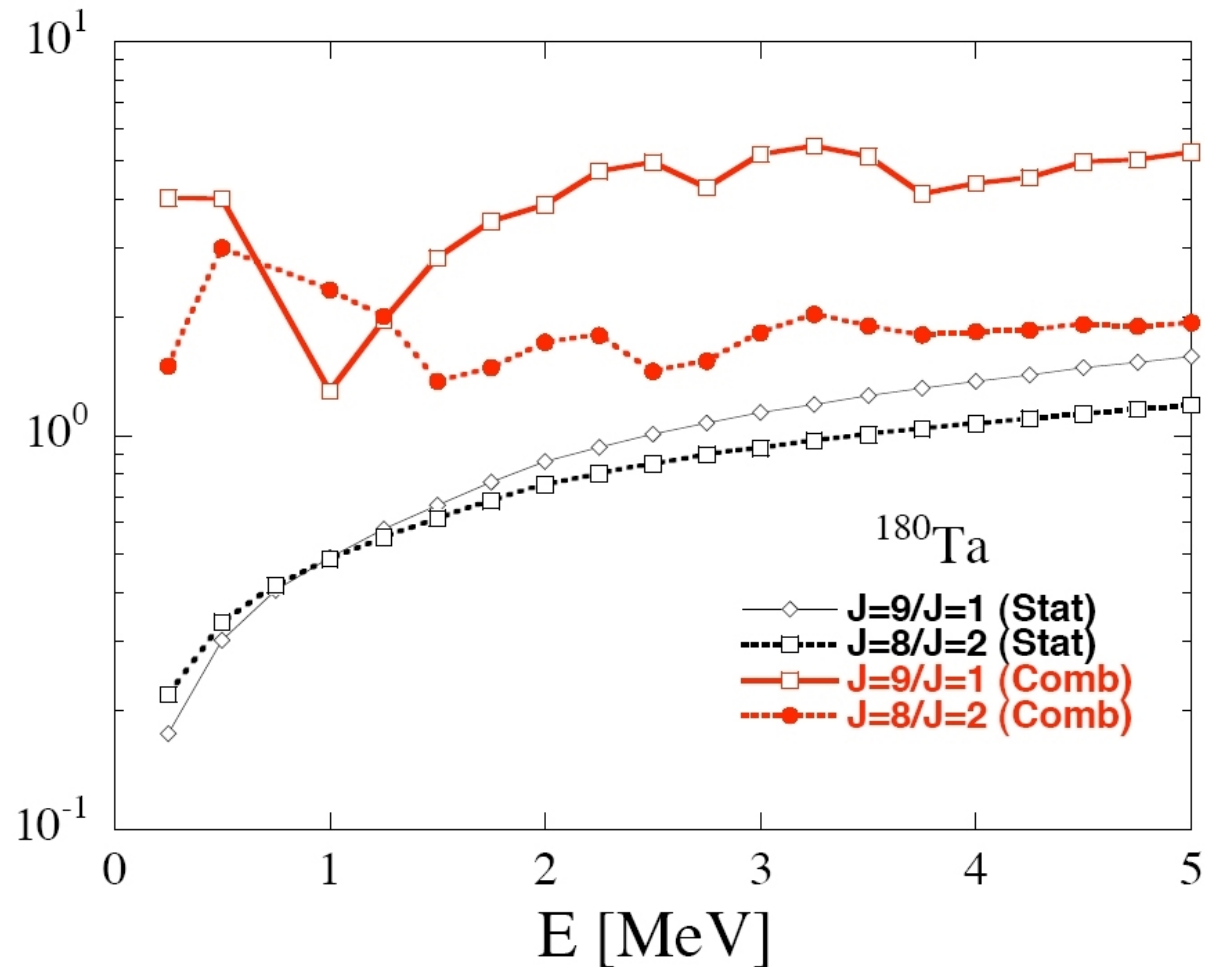
Combinatorial level densities





Combinatorial level densities

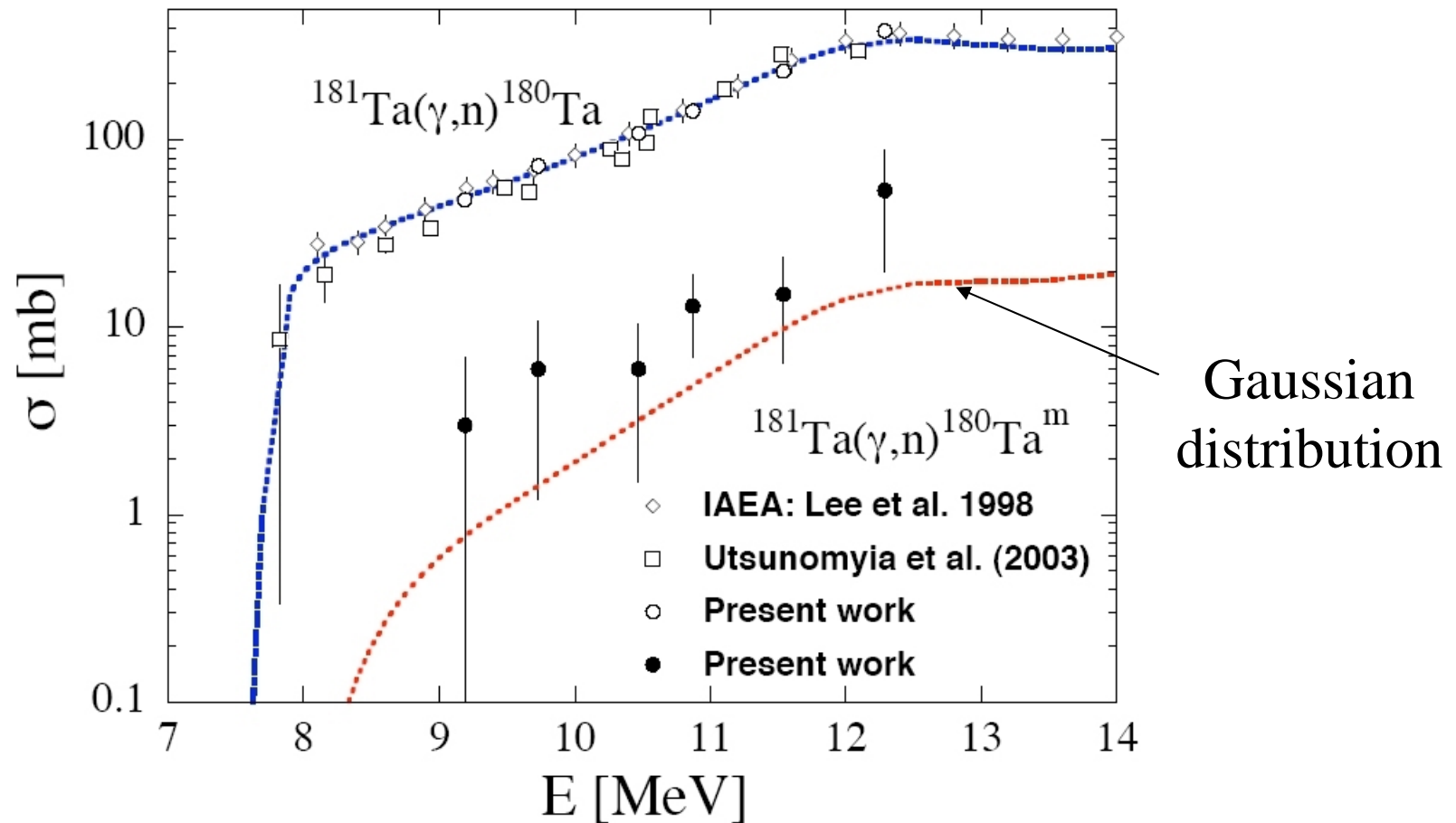
*Talys deals with realistic (non statistical) parity and **spin distributions***



➔ **Non-statistical feature imply significant deviations from the usual gaussian spin dependence**

Combinatorial level densities

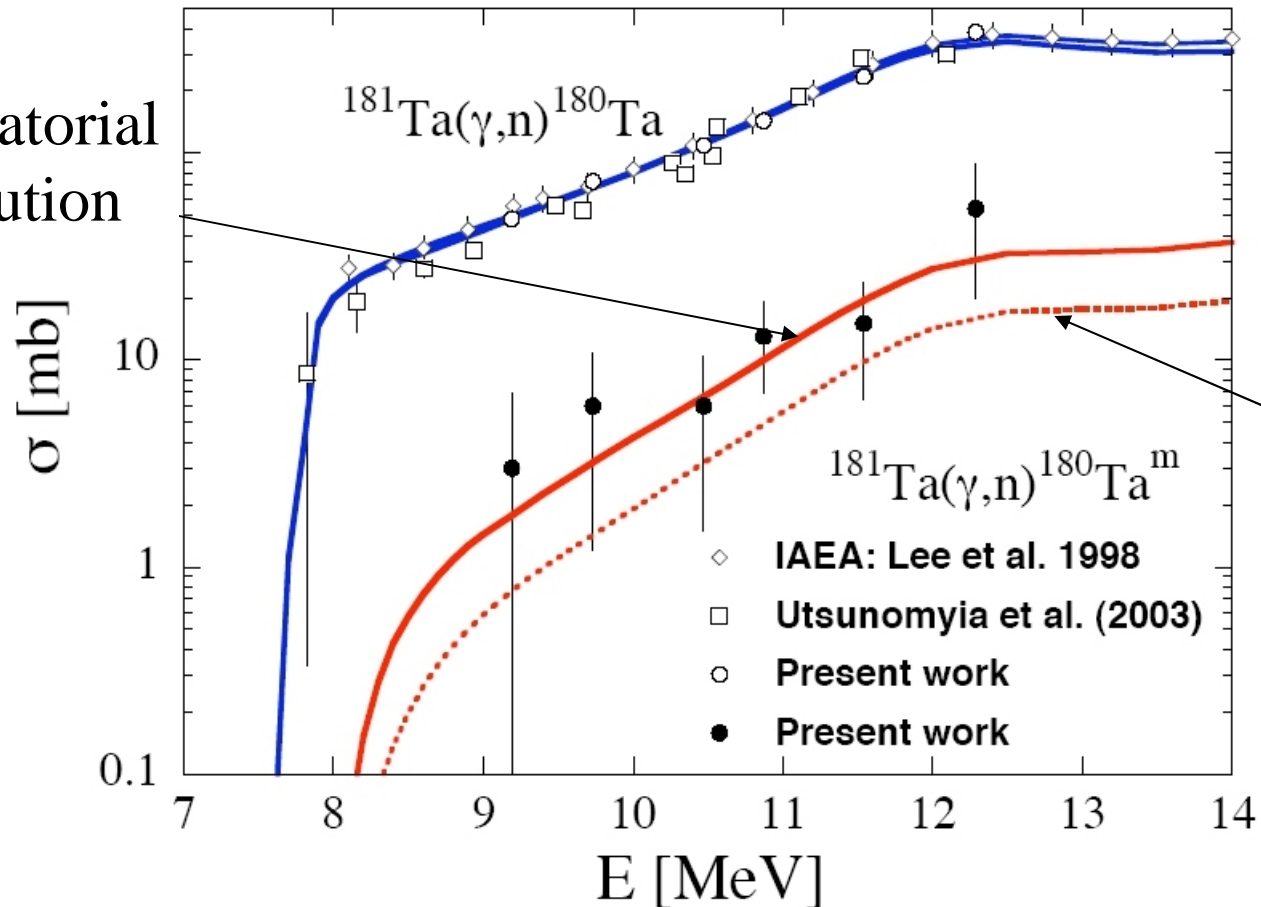
See PRL 96 (2006) 192501 for details



Combinatorial level densities

See PRL 96 (2006) 192501 for details

Combinatorial
distribution

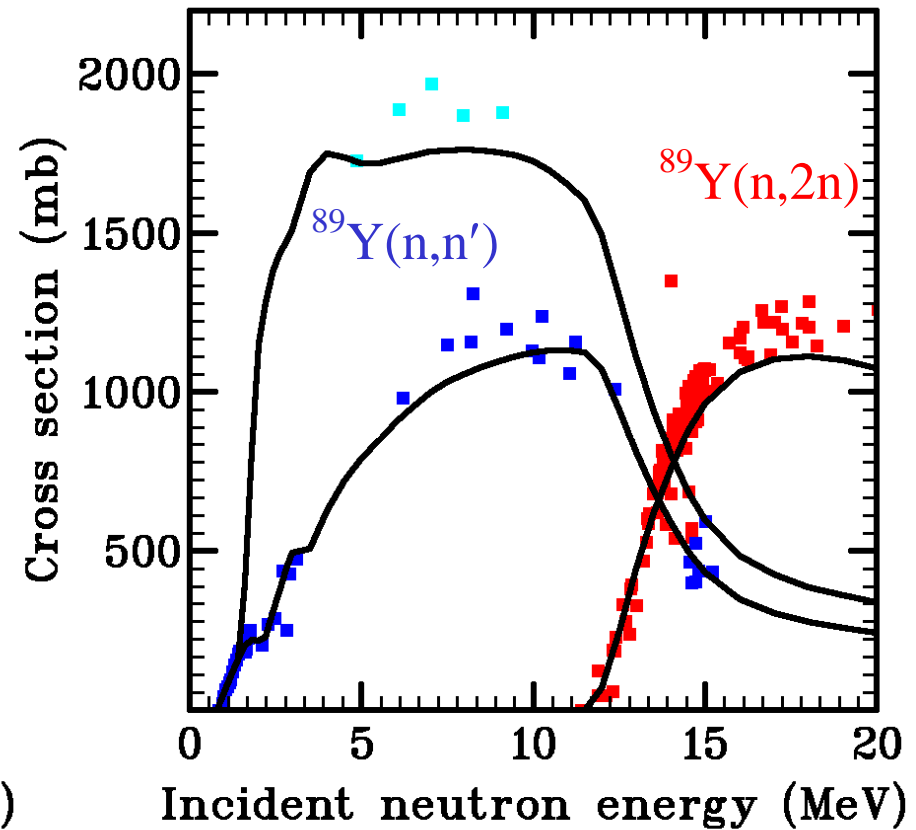
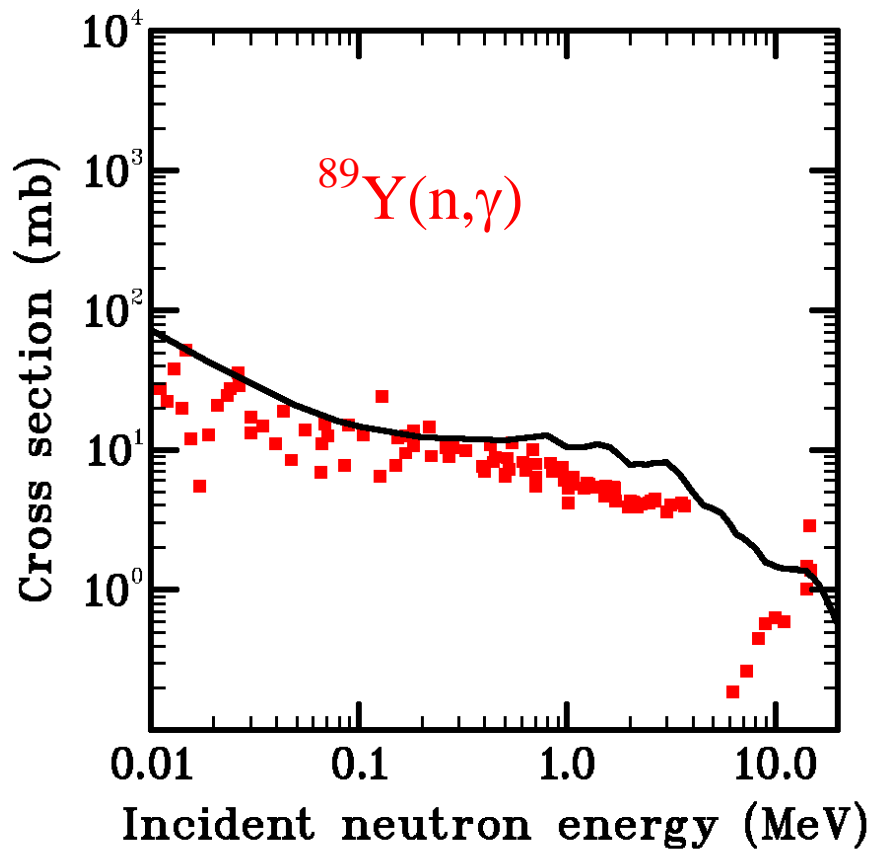


Gaussian
distribution

➔ Deviations from the usual gaussian spin dependence can have large impact on isomeric level production cross sections

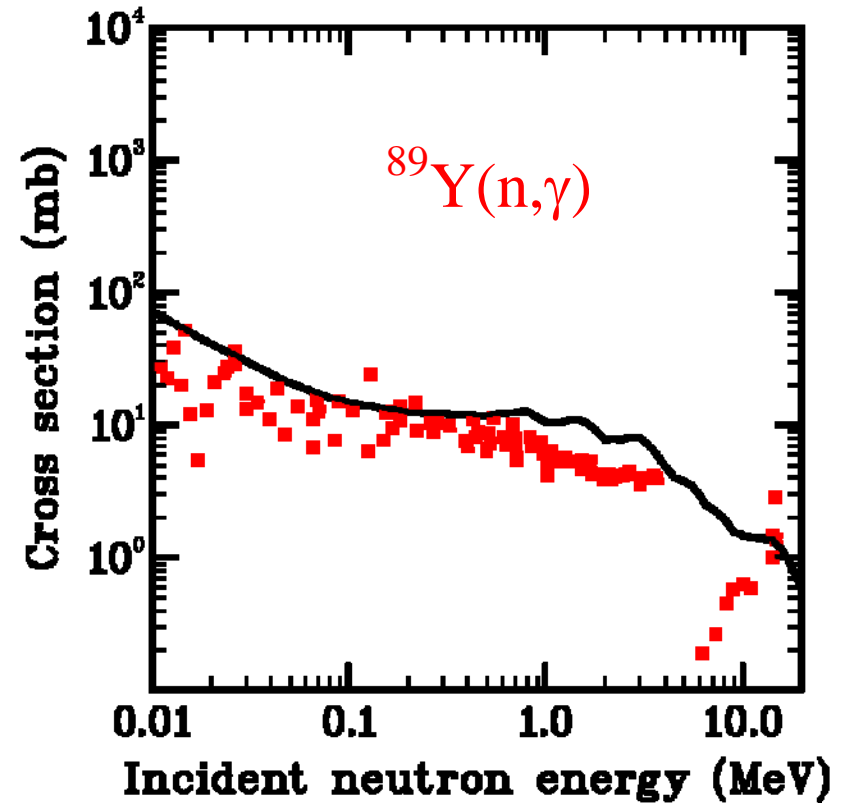
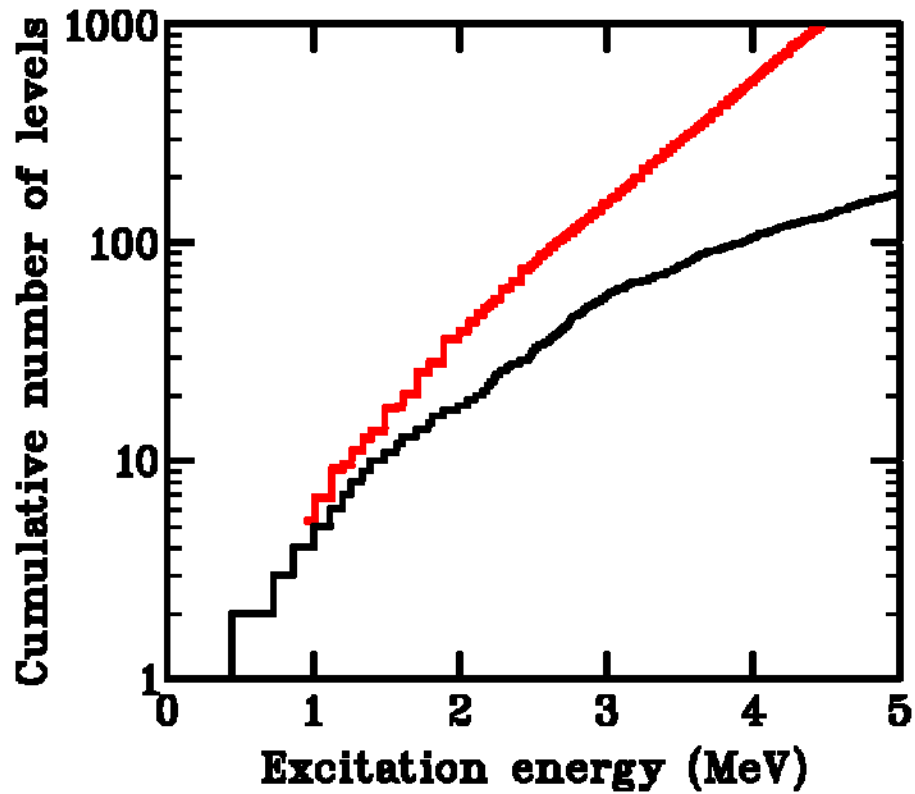
Adjustable level densities : recipe & impact

$$\rho_{\text{renorm}}(U) = e^{\alpha \sqrt{(U - \delta)}} \rho_{\text{global}}(U - \delta)$$



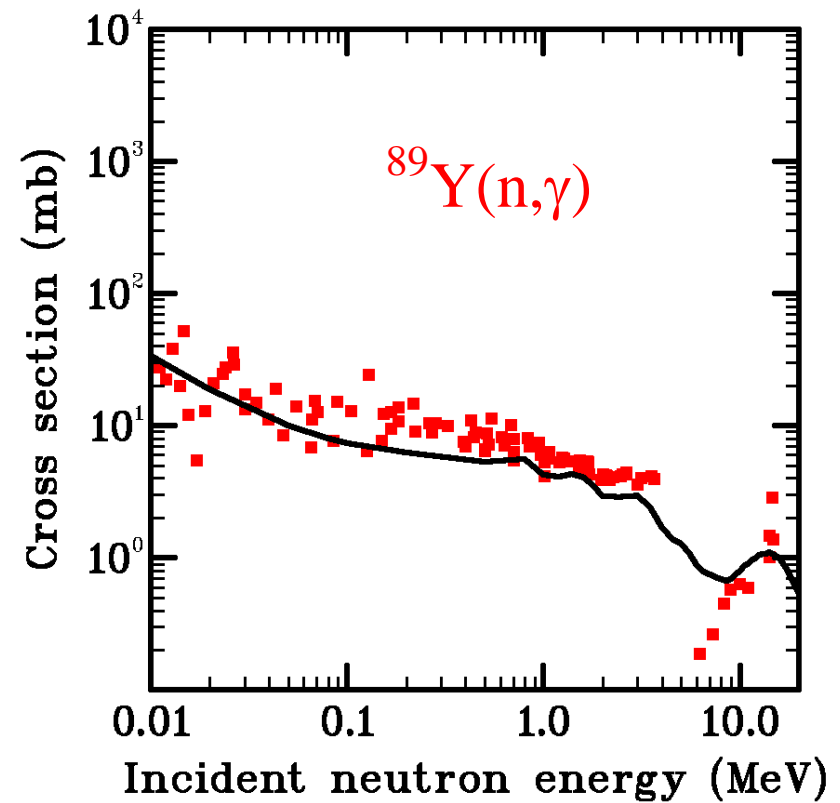
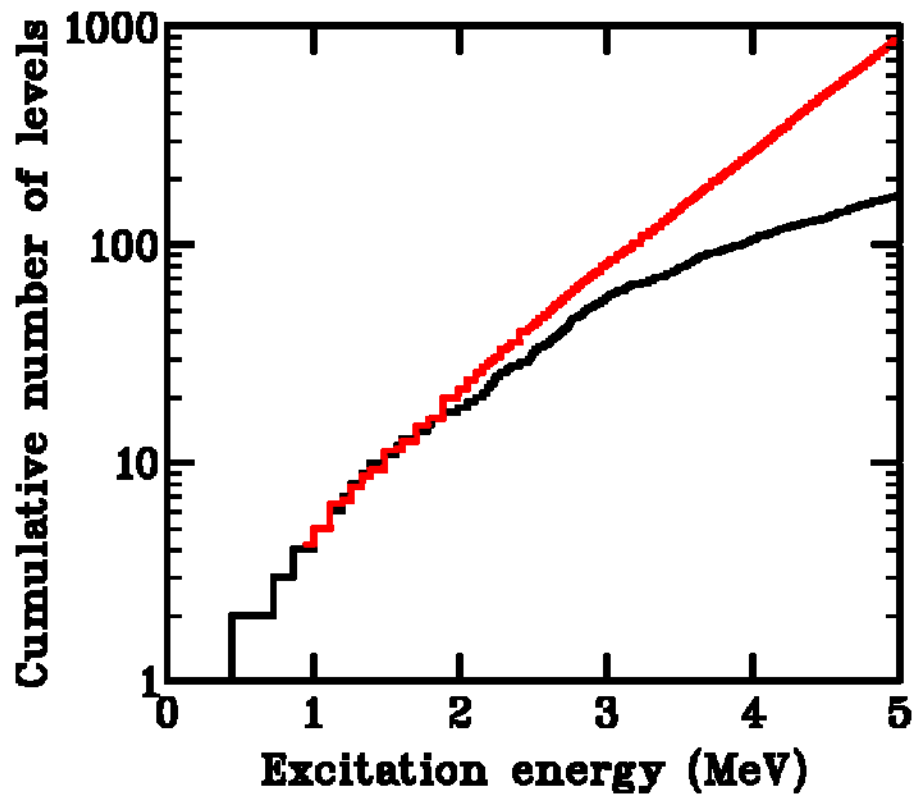
Adjustable level densities : recipe & impact

$$\rho_{\text{renorm}}(U) = e^{\alpha \sqrt{(U - \delta)}} \rho_{\text{global}}(U - \delta)$$



Adjustable level densities : recipe & impact

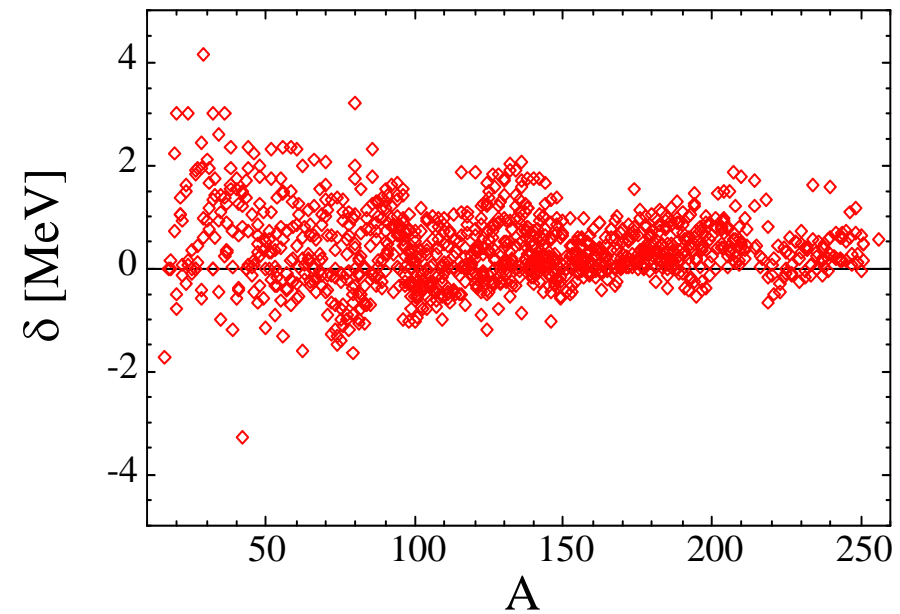
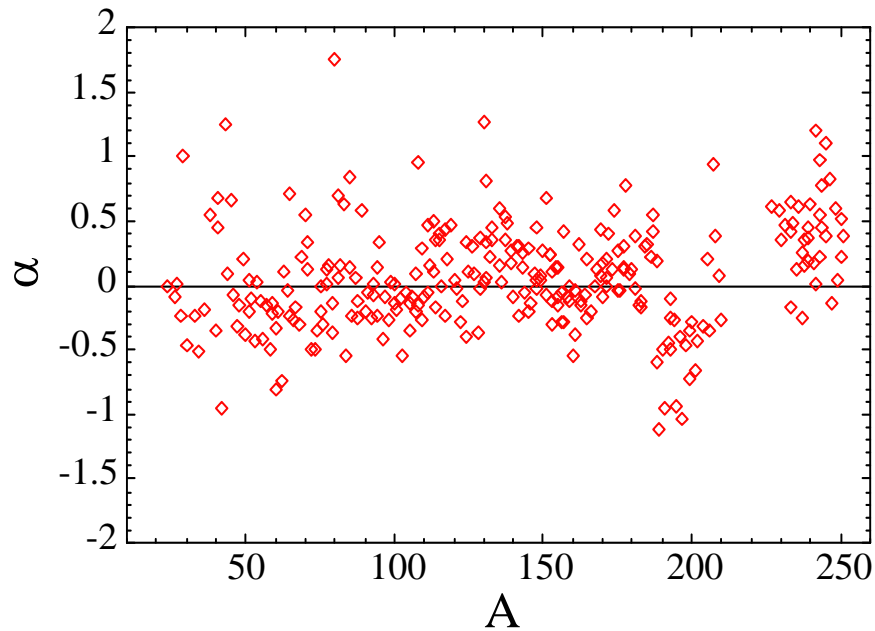
$$\rho_{\text{renorm}}(U) = e^{\alpha \sqrt{(U - \delta)}} \quad \rho_{\text{global}}(U - \delta)$$



Global adjustment

See NPA 810 (2008) 13 for details

α and δ adjusted to fit discrete levels (≈ 1200 nuclei) and D_0 's (≈ 300 nuclei) using the TALYS code



- **Gilbert-Cameron model + Ignatyuk**
 - ⇒ **Default**
- **Back-Shifted Fermi Gas model + Ignatyuk**
 - ⇒ **Default**
- **(Generalized) Superfluid model**
 - ⇒ **More rigorous treatment of pairing correlation at low energy**
 - ⇒ **Fermi gas + Ignatyuk law above some critical energy**
 - ⇒ **Explicit treatment of collective effects**
- **Combinatorial approach**
 - ⇒ **Direct counting method of both partial and total level densities**
 - ⇒ **Access to non statistical effects**

Phenomenological \Rightarrow microscopic predictions ?

Experimentally known (deduced)

Theoretically predicted

Nuclear properties

- Level properties (E , J^π , branching ratios)
- deformations



HFB + ν - ν interaction

Level densities

- Gilbert & Cameron
- Back Shifted Fermi Gas
- Generalized Superfluid Model
- Williams + several corrections (p-h)



Combinatorial method

- Total level densities
- **p-h level densities**

Optical model

- Koning & Delaroche
- Soukhovistkii (actinides)
- Tabulated



Semi-microscopic JLM

γ -strength functions

- Kopecky-Uhl, Brink-Axel



HFBCS or HFB Tables

Fission paths

- Hill-Wheeler

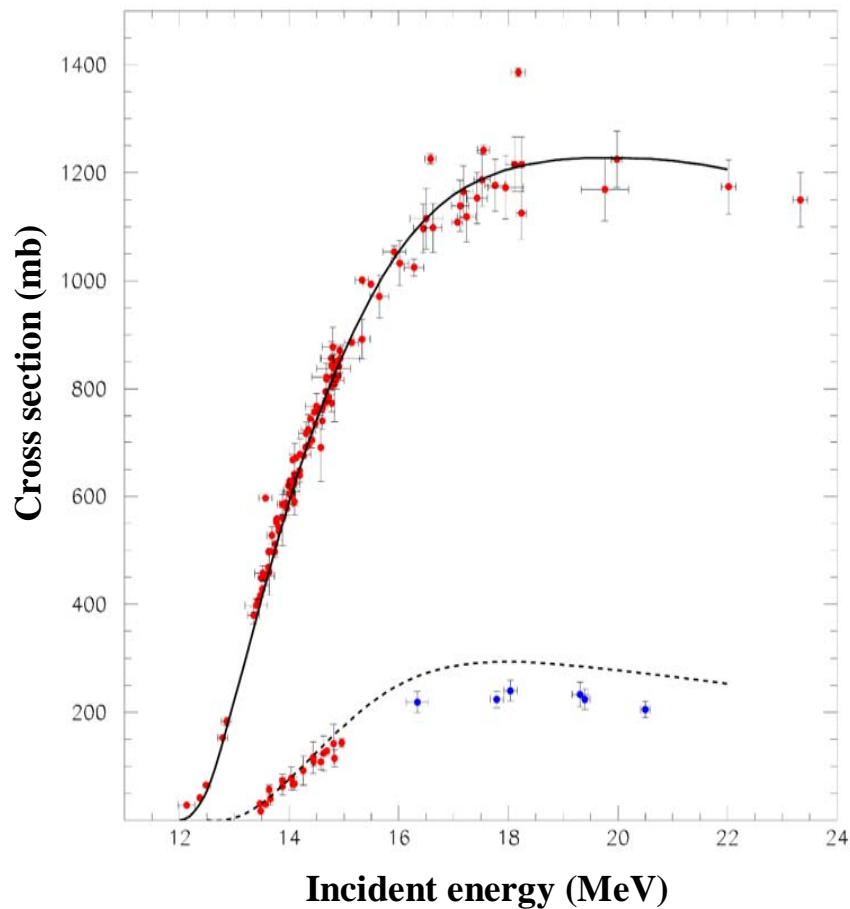


HFB shapes with WKB penetrabilities

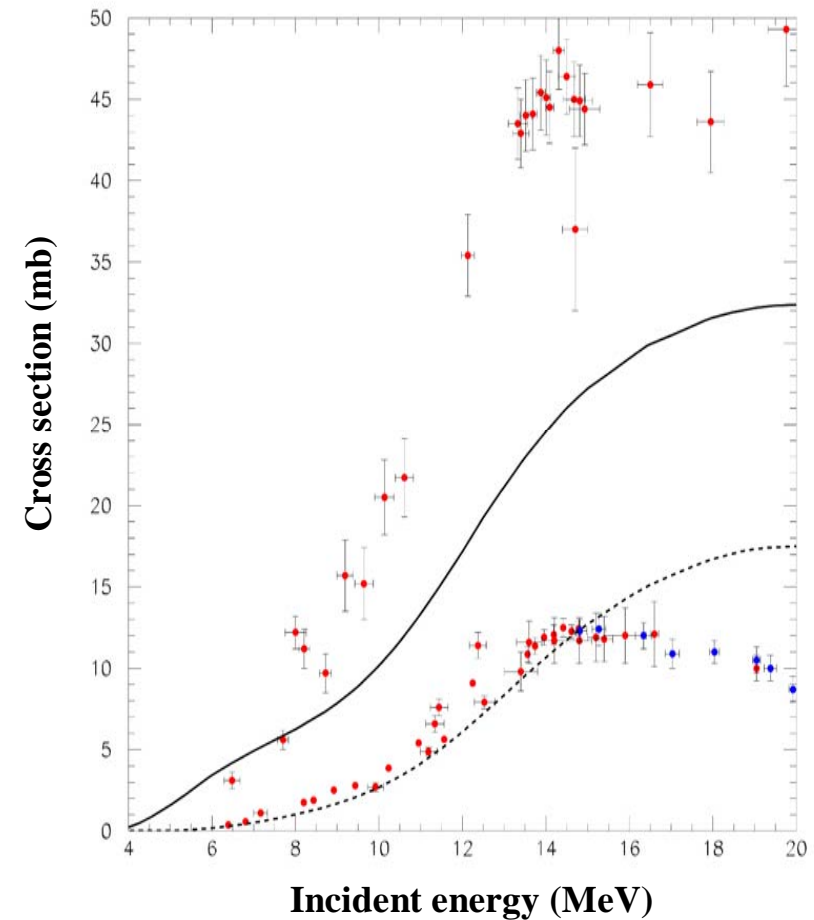
Some TALYS results

Fully microscopic cross section (almost)

$^{90}\text{Zr} (n,2n) ^{89}\text{Zr}$



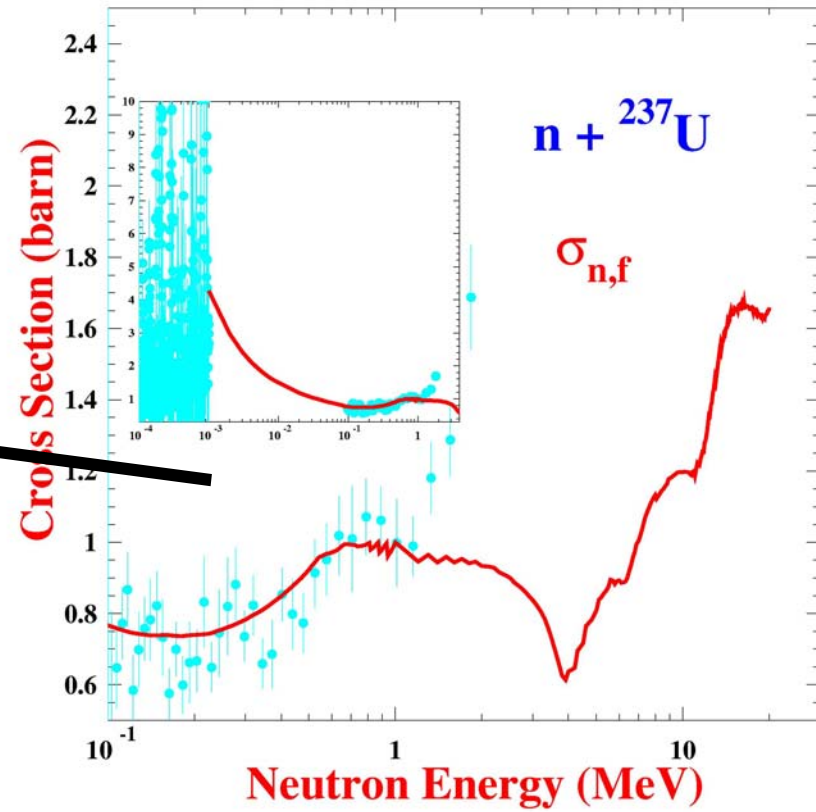
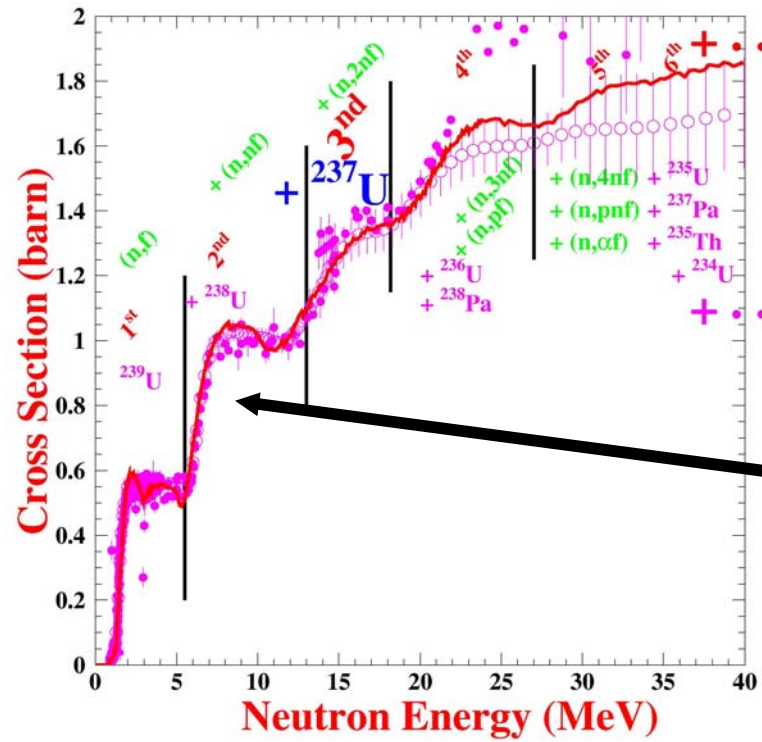
$^{90}\text{Zr} (n,p) ^{90}\text{Y}$

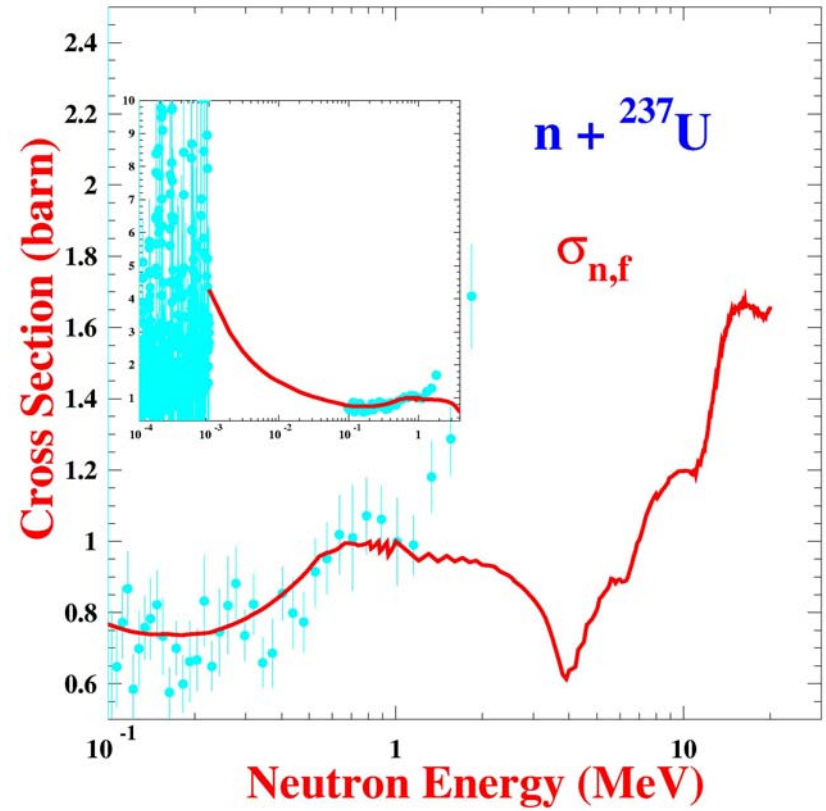
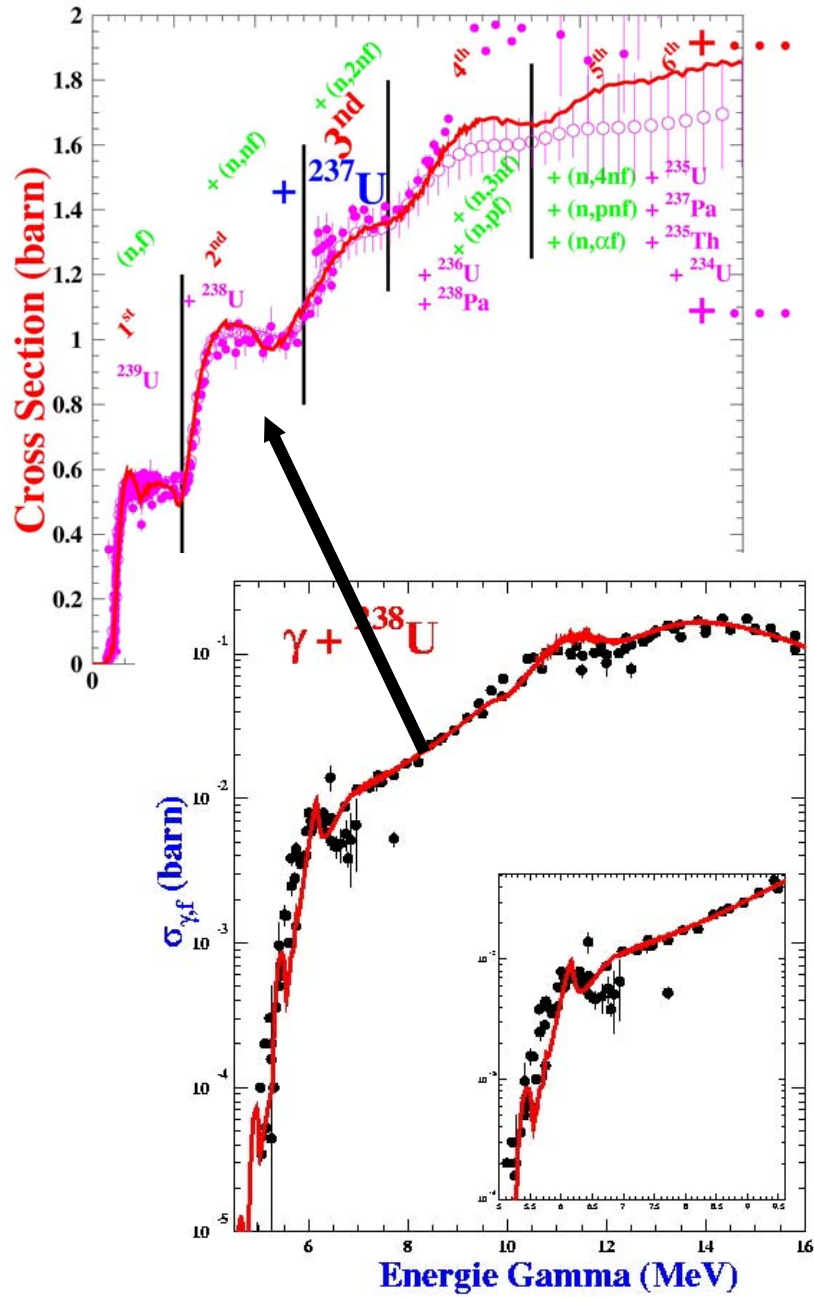


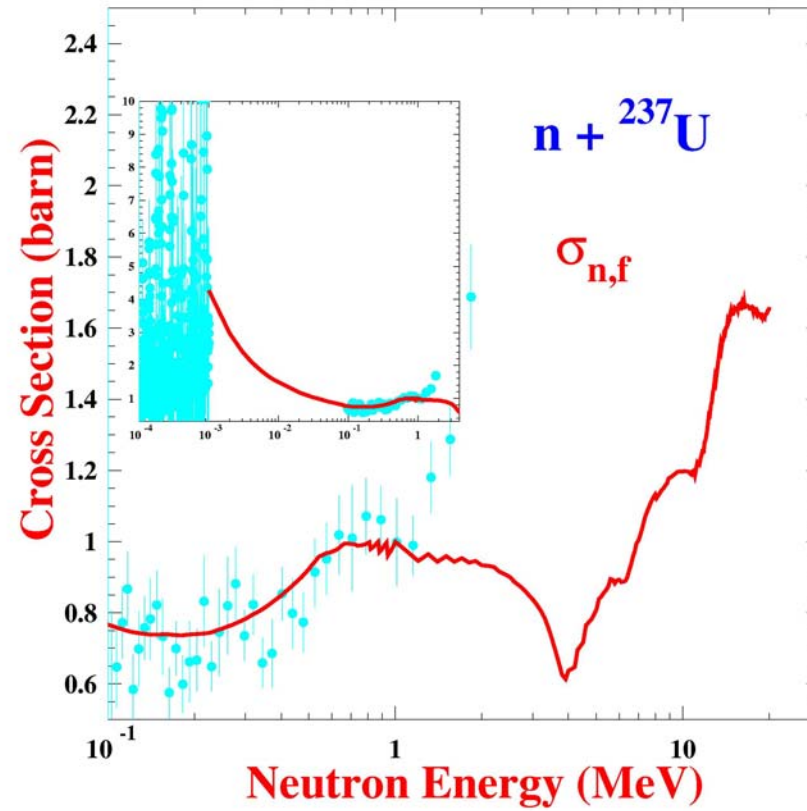
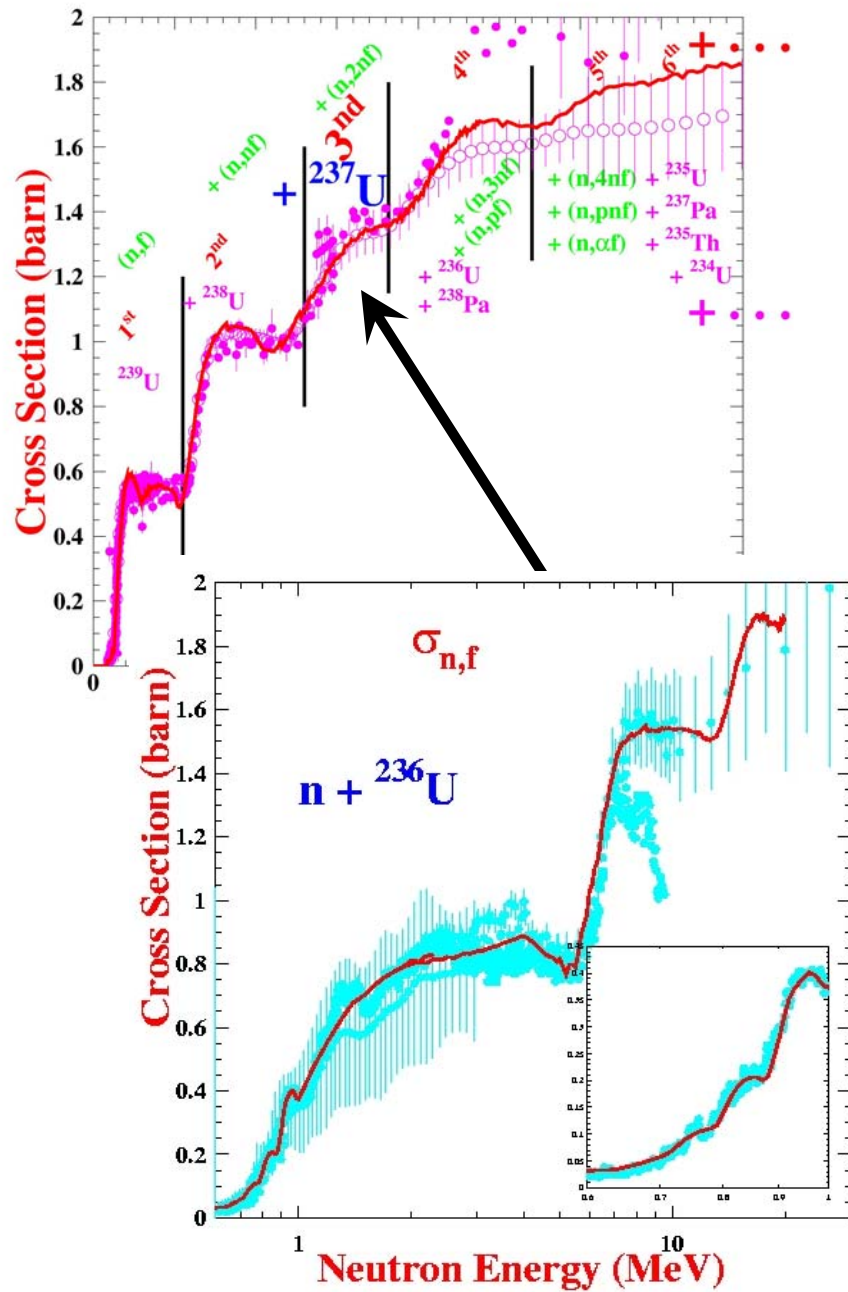
Coherent fission cross sections with phenomenological approach

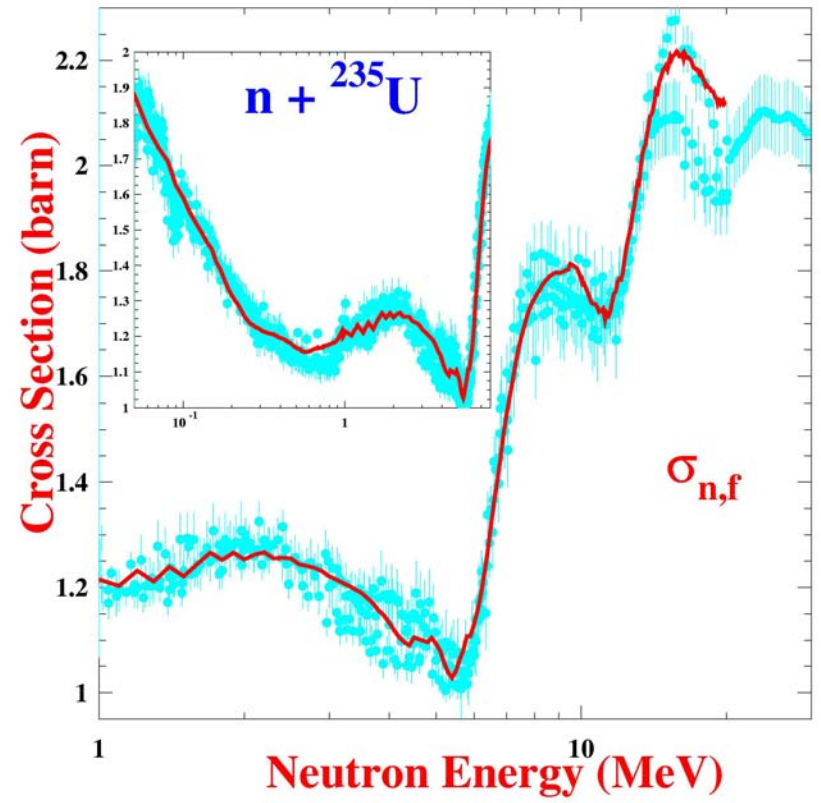
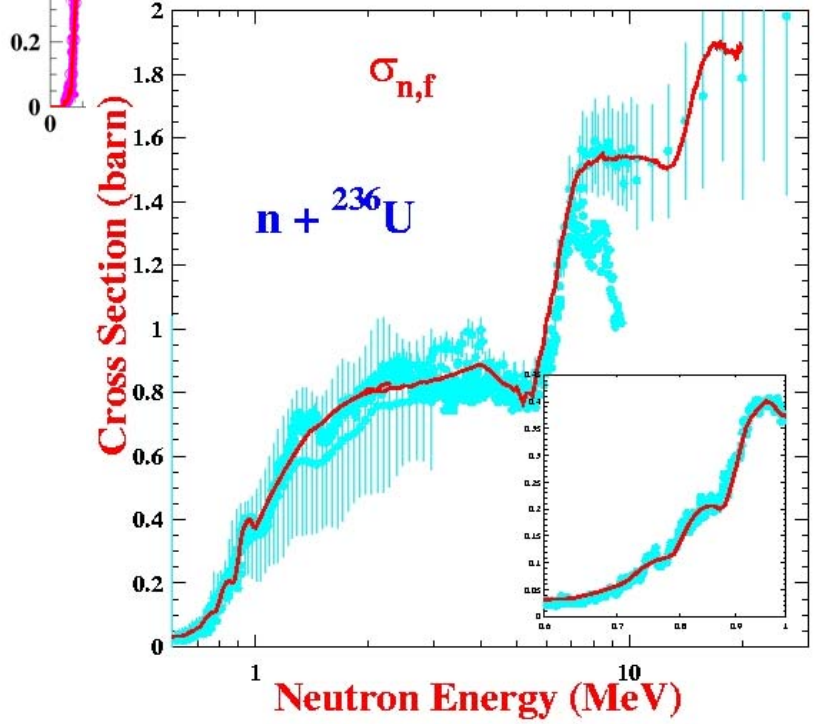
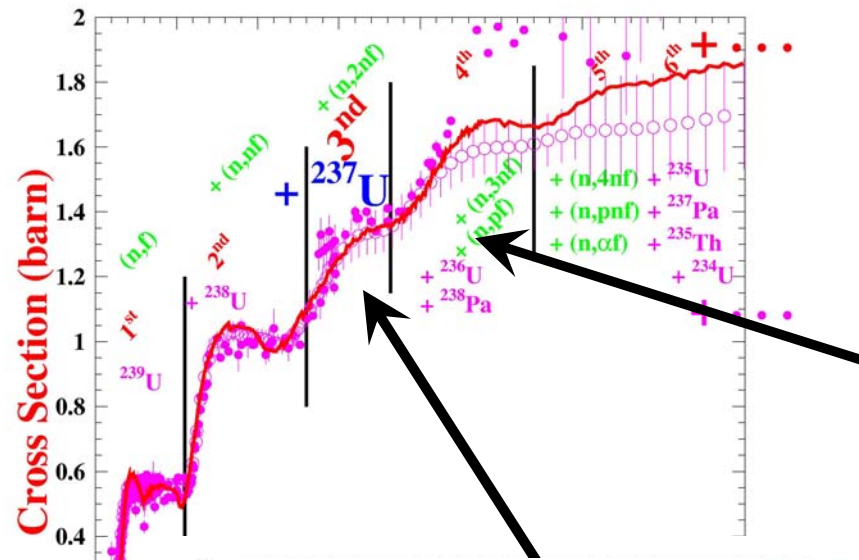
Neutron induced fission on ^{238}U

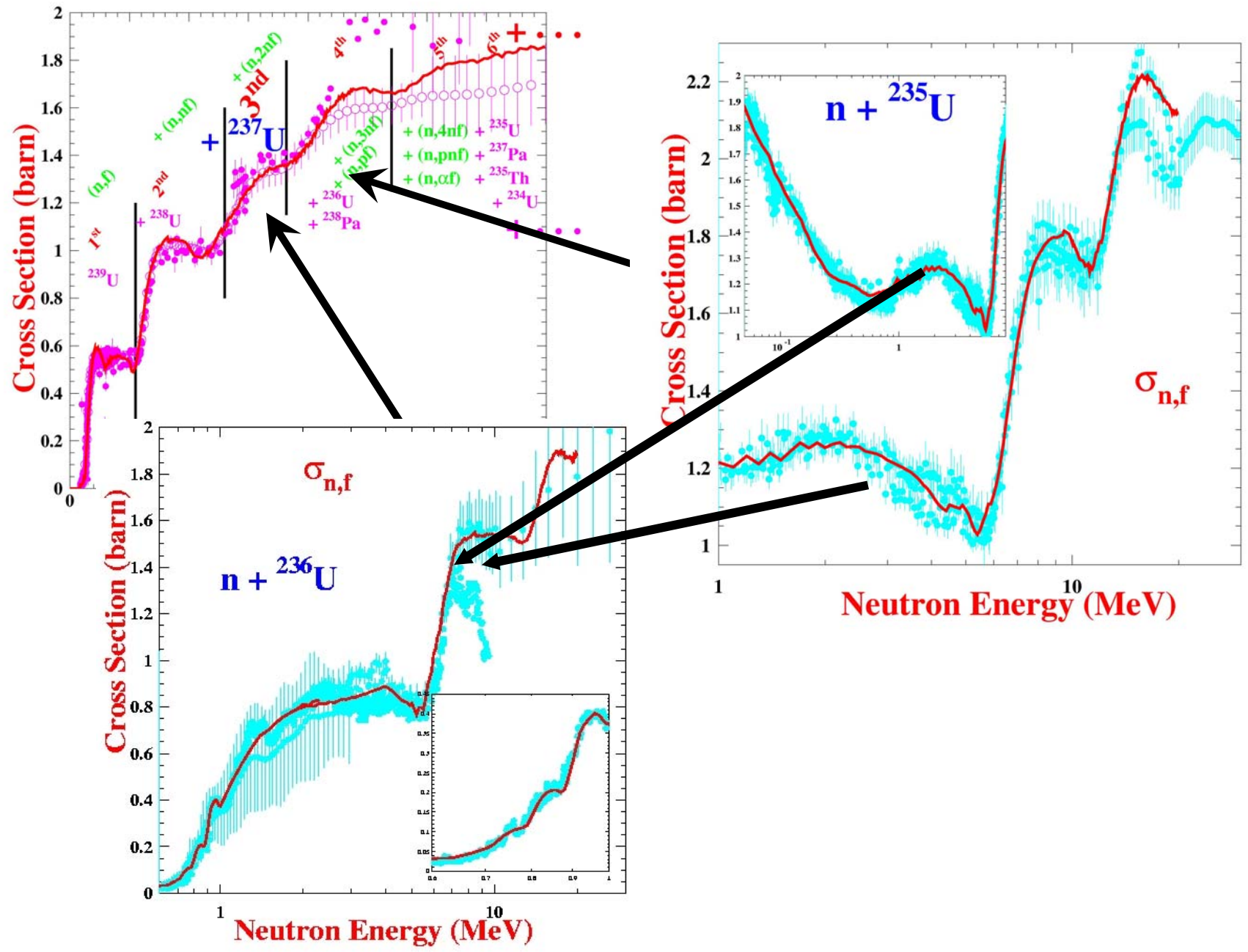
- several hundreds of parameters
- unique set for all fission chances or U targets

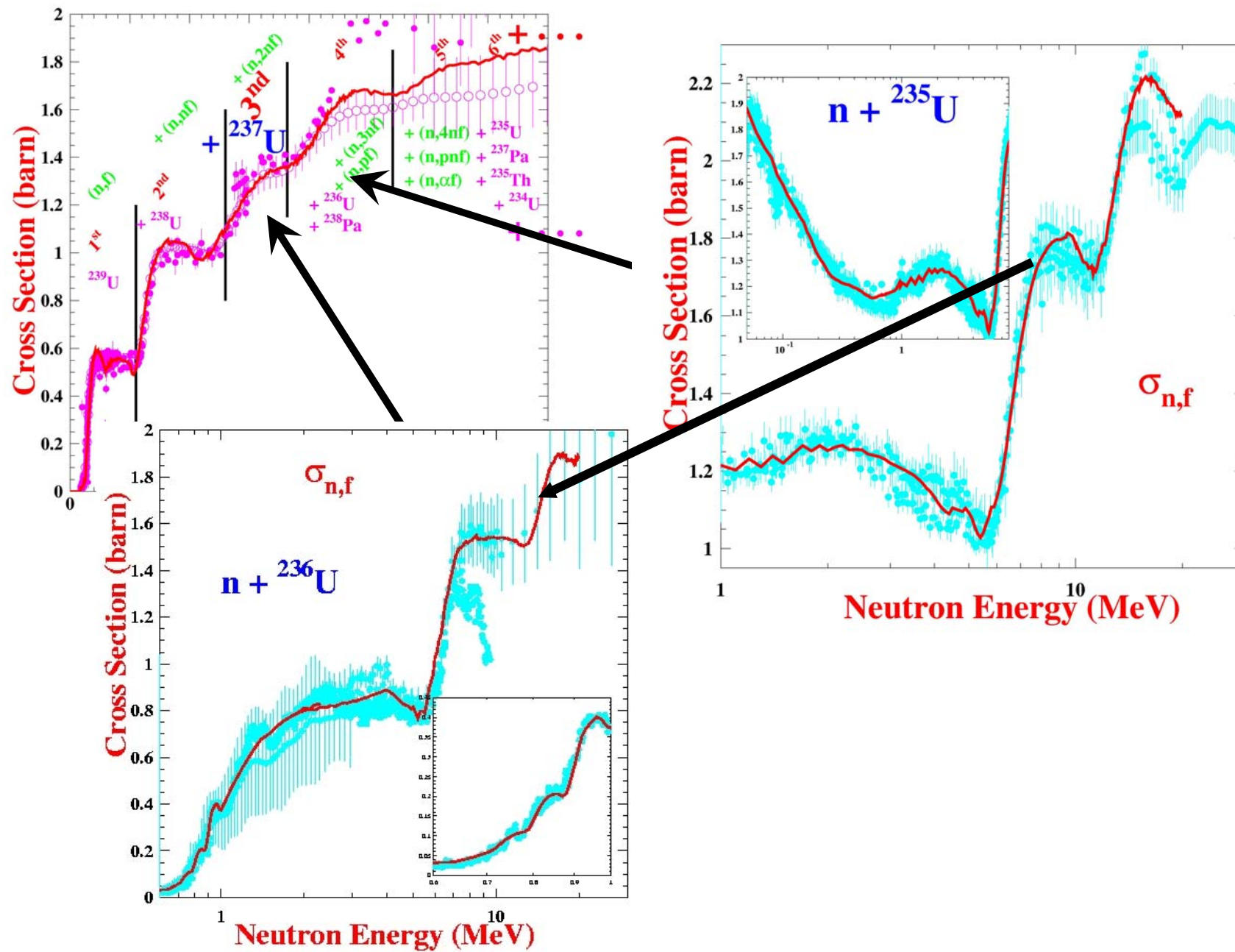


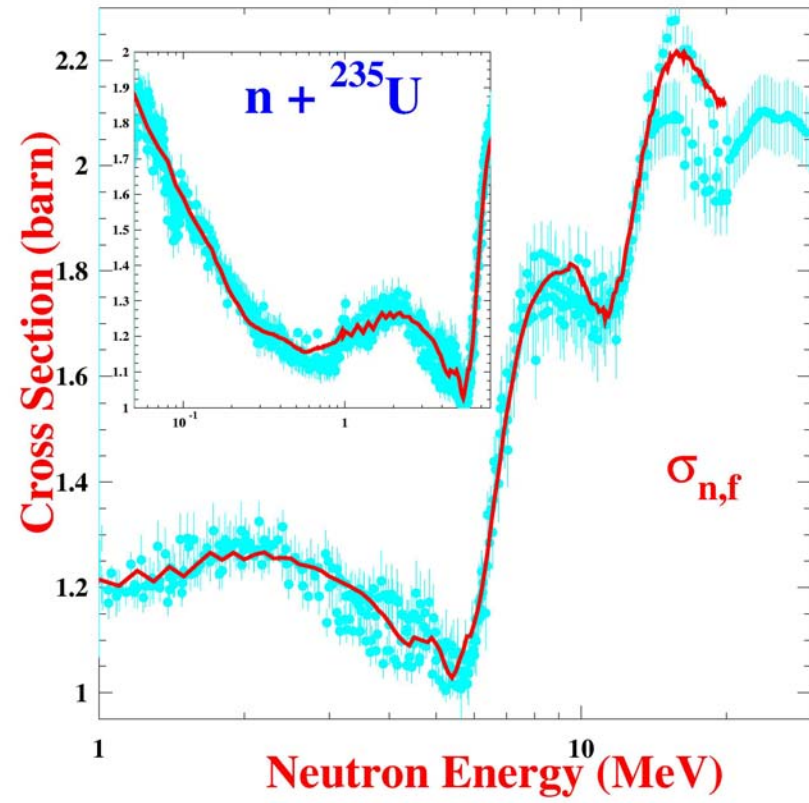
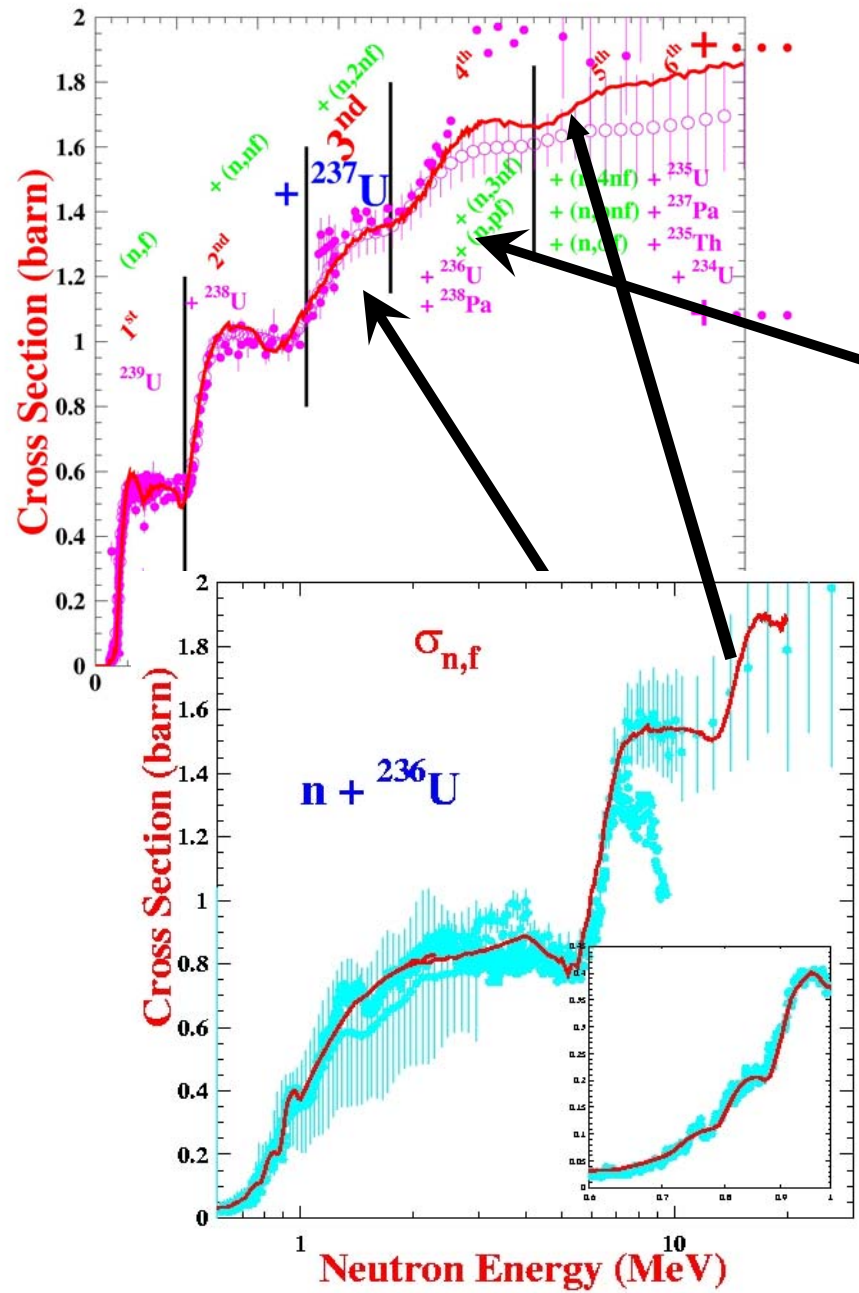


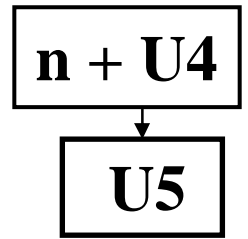
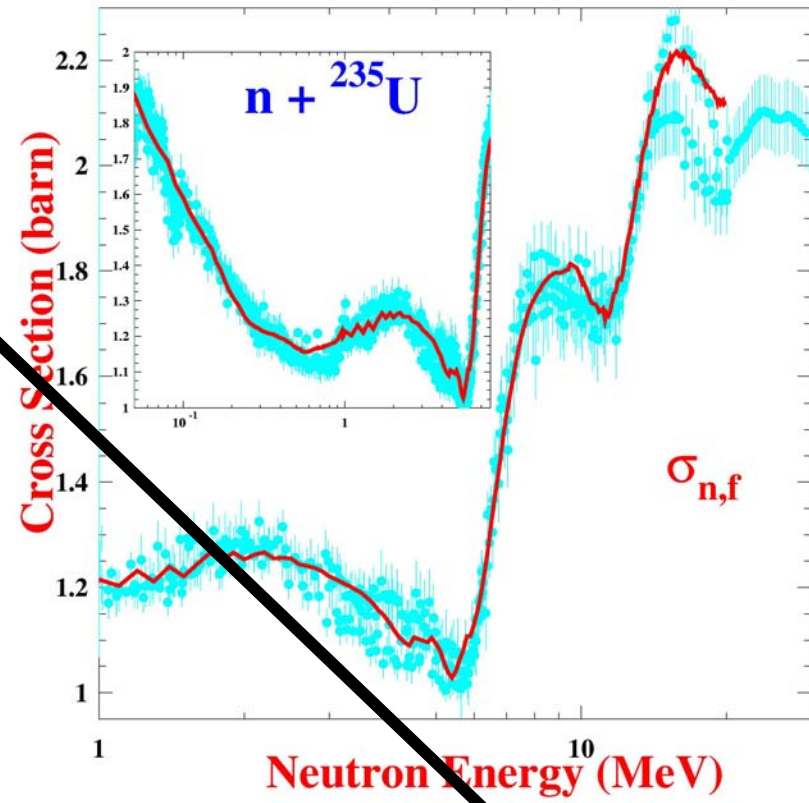
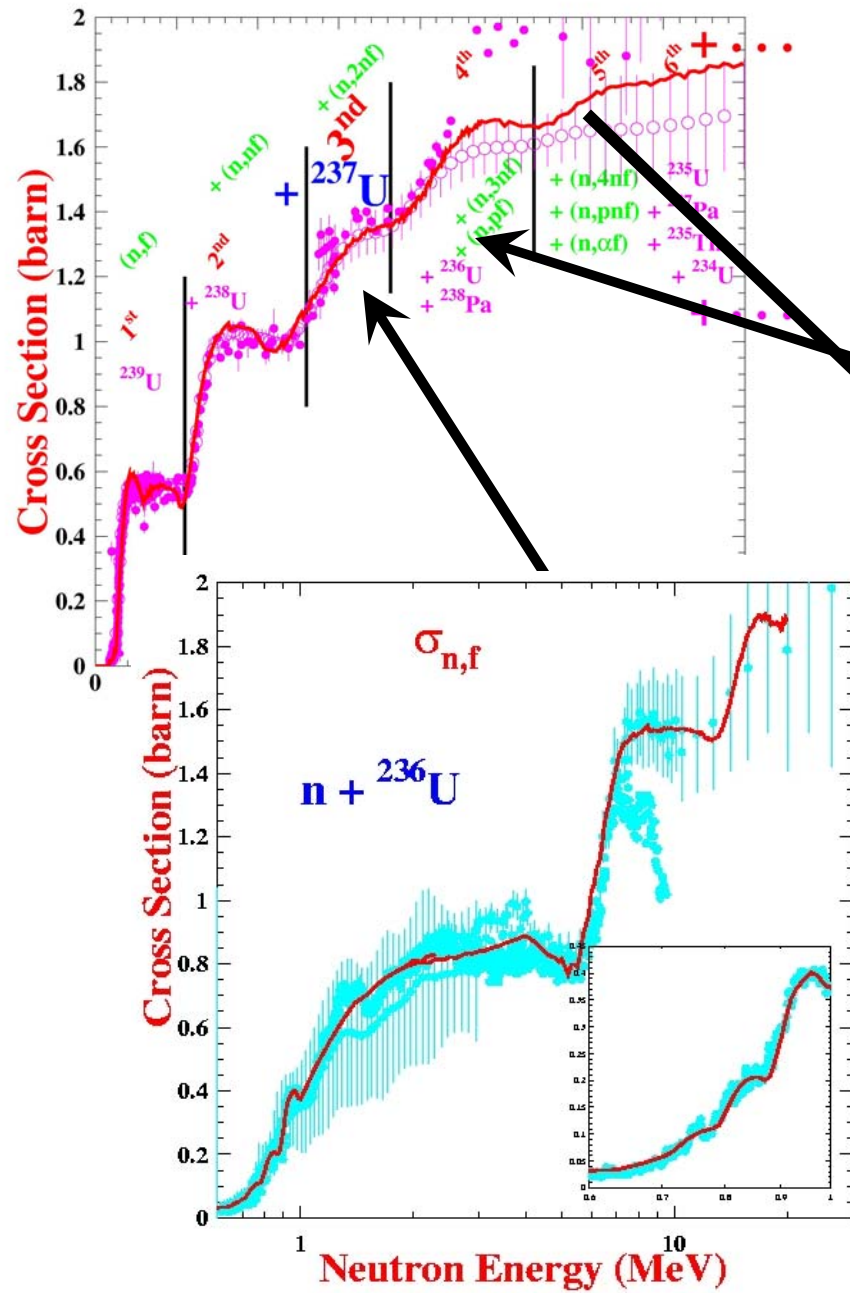












Coherent fission cross sections With microscopic ingredients

HFB-14 predictions of fission barriers and NLD at saddle points, including renormalization (max 5 parameters) of

- fission path height: $B_f'(\beta_2) = B_f(\beta_2) \times v_{corr}$
- NLD at 1st and 2^d saddle points:

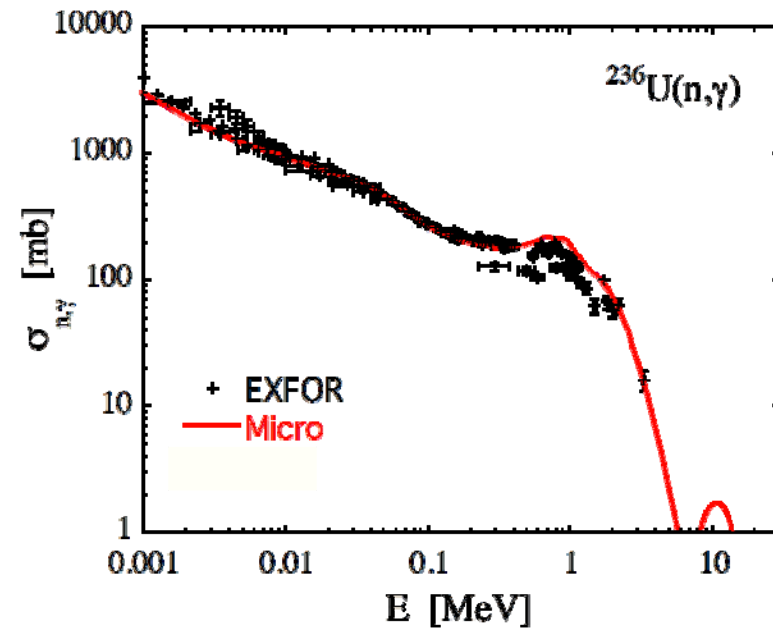
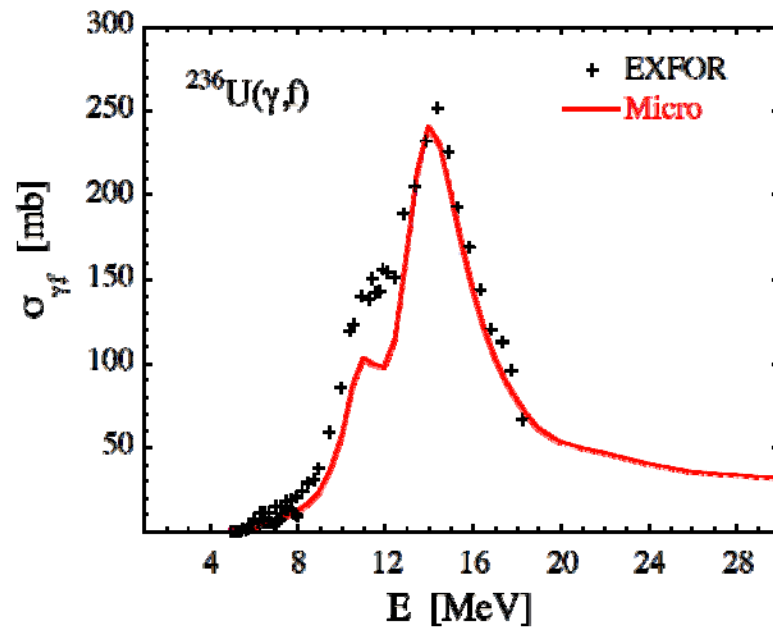
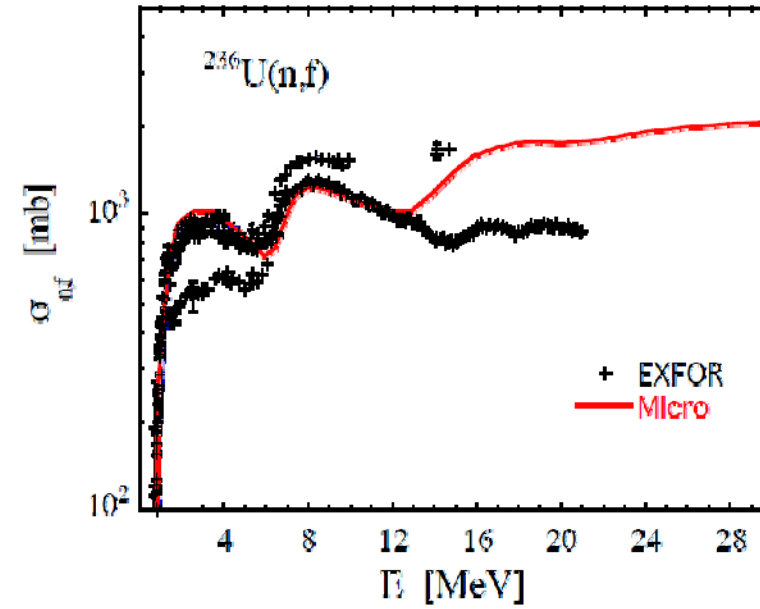
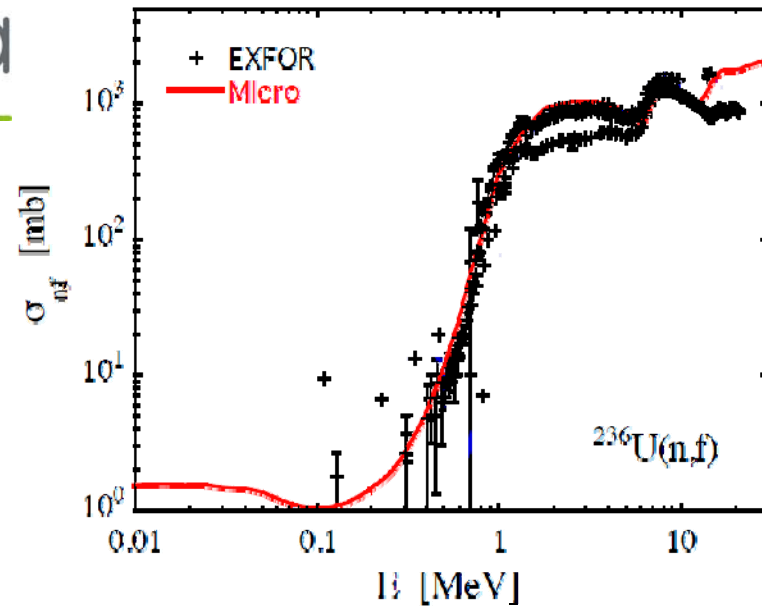
$$\rho'(U, J, P) = \rho(U - \delta, J, P) e^{\alpha\sqrt{U-\delta}}$$

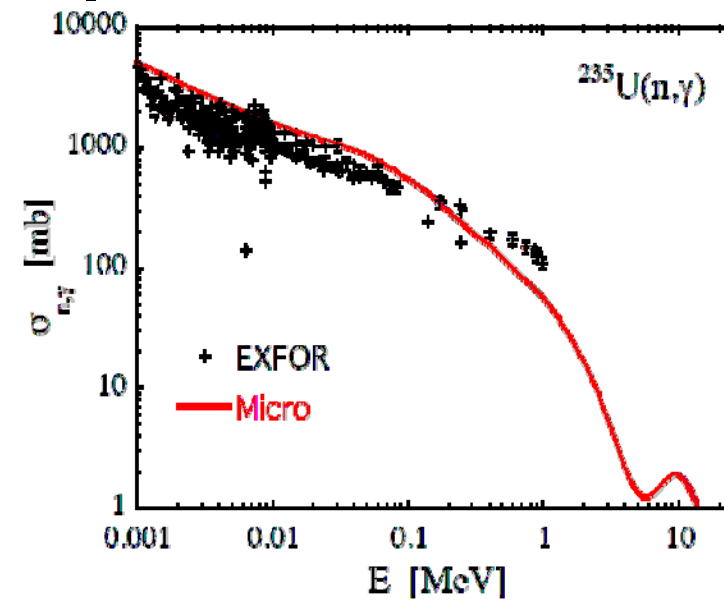
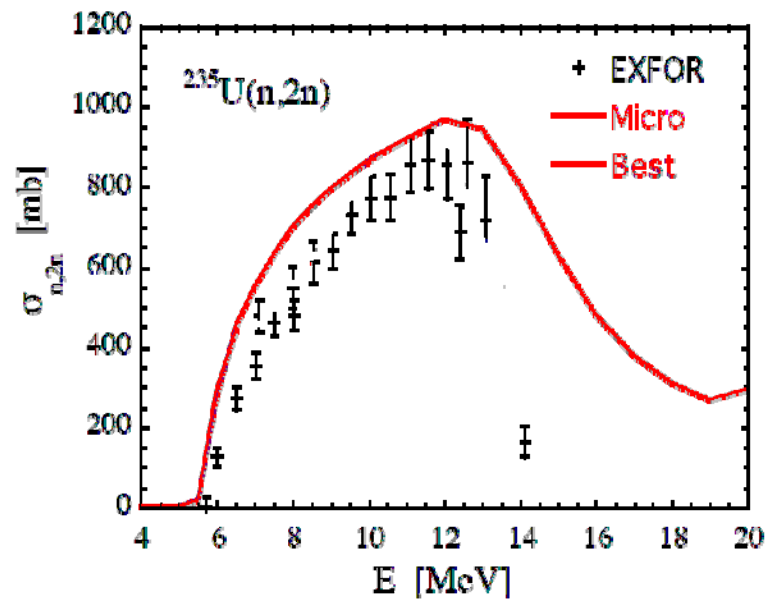
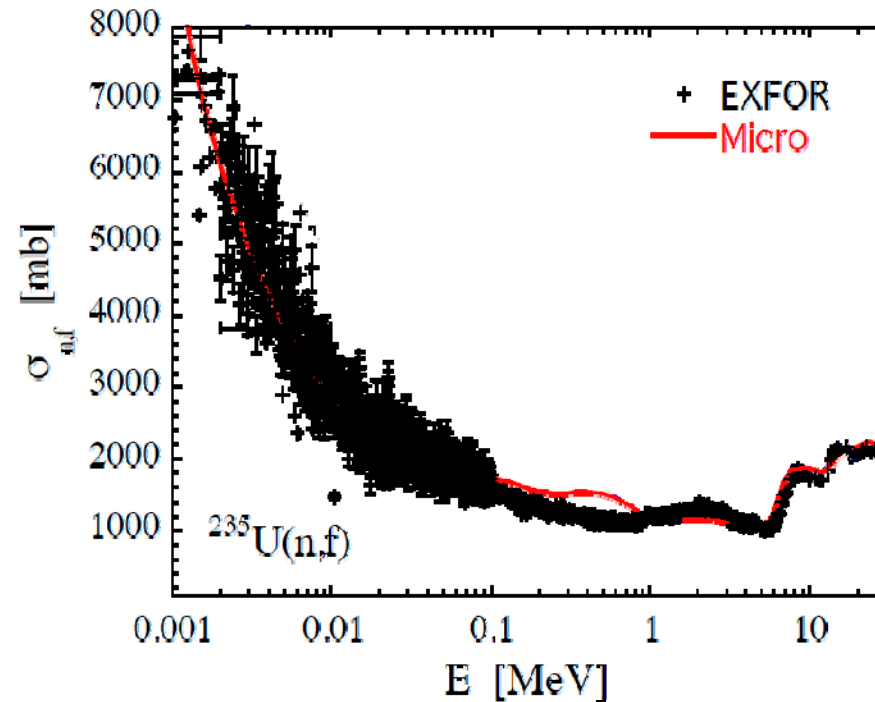
Additional nuclear inputs:

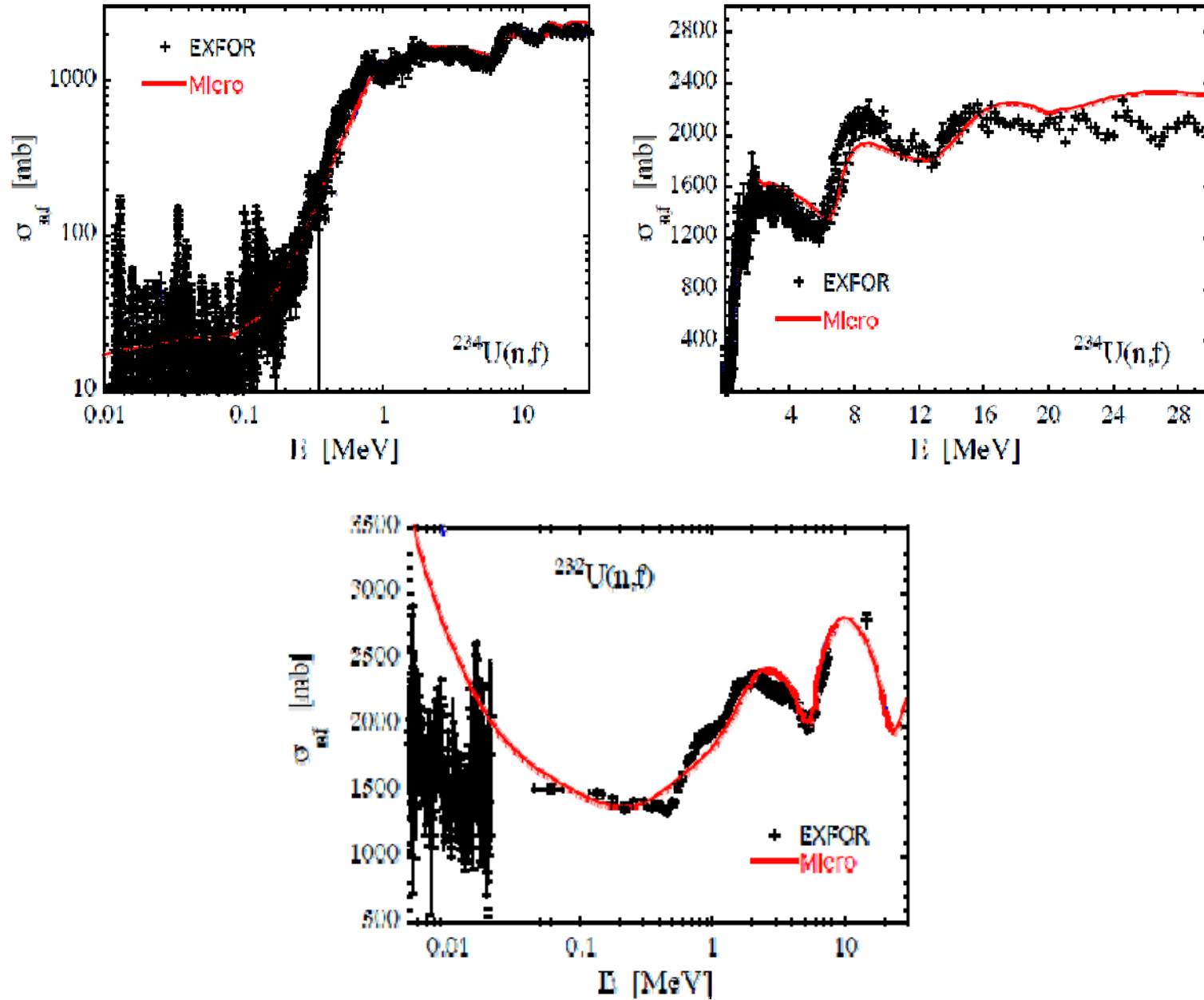
- Nuclear structure properties: HFB-14 (Goriely et al. 2007)
- Optical potential: Soukhovitskii et al. (2004)
- γ -ray strength: Hybrid model (Goriely, 1998)
- NLD: HFB-14 plus combinatorial model (Goriely et al., 2008)
normalized on s-wave spacings and discrete excited levels

Note:

- **1 UNIQUE set of nuclear ingredients for all U isotopes**
- no class 2 states included
- no discrete transition states included







- *Cross section modeling quite easy for non fissile nuclei*

Microscopic or Phenomenological OMP, Γ_γ , LDs

⇒ full microscopic calculation for non fissile nuclei almost possible

- *Difficult cross section modeling for fissile nuclei*

- *Web site opened in October 2006 : www.talys.eu*

⇒ All microscopic ingredients mentioned included in the distribution

- *New level densities for pre-equilibrium (done but not tested)*
- *JLM OMP : spherical (OK) – deformed (soon)*
- *Neutron multiplicities from FF decay (under dev.)*
- *Microscopic ingredients with Gogny instead of Skyrme (under dev.)*