



**The Abdus Salam  
International Centre for Theoretical Physics**



**2164-7**

**Workshop on Nano-Opto-Electro-Mechanical Systems Approaching the  
Quantum Regime**

*6 - 10 September 2010*

**Cavity Optomechanics with Microresonators**

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*EPFL, SB IPEQ UPKIPPE1, PH D2 392 (Batiment PH)*

*Station 3, CH-1015*

*Lausanne*

*SWITZERLAND*



# Cavity Optomechanics with microresonators

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**Tobias J. Kippenberg**

**Laboratory of Photonics and Quantum Measurements**

**Miramar, Italy  
7<sup>th</sup> September 2010**

**Collaborators**

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J. P. Kotthaus

W. Zwerger

I. Wilson-Rae

A. Marx

J. Raedler

R. Holtzwarth (Menlo System)



European Research Council

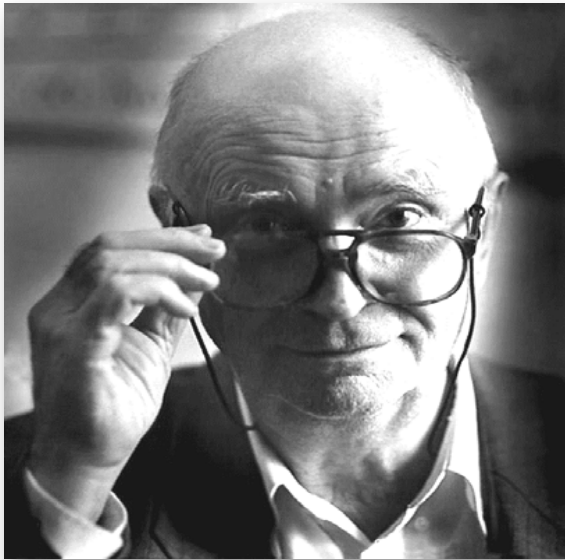




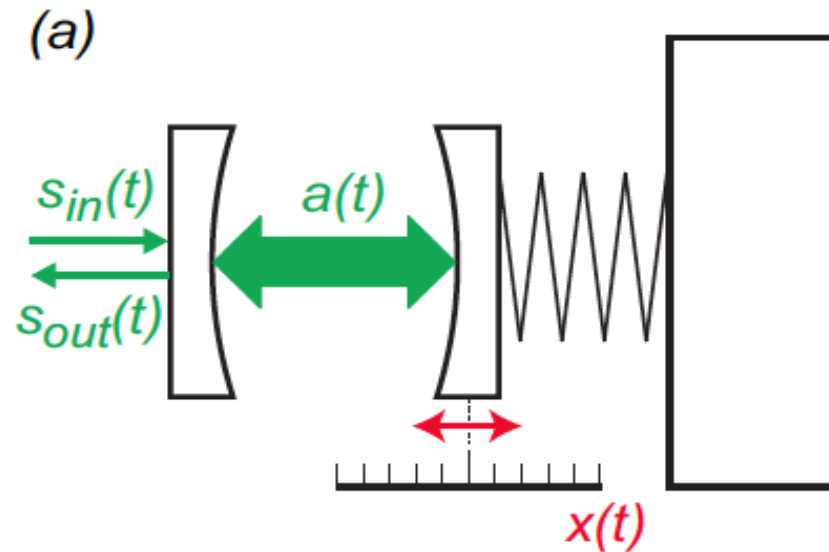
PhD thesis available



# Parametric transducer as canonical optomechanical system



V.B. Braginsky



$$\omega = \omega_0 + Gx(t)$$

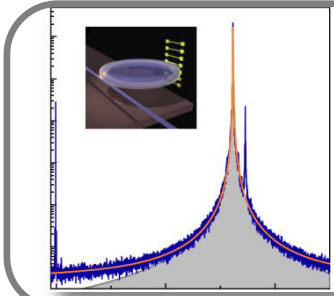
$$\hat{H}_{int} = -\hbar\omega \frac{\hat{x}}{L} (\hat{a}^\dagger \hat{a}) = -\hat{F}_{RP} \hat{x}$$

Braginsky, Manukin: *Measurement of Weak Forces in Physics Experiments* (1977)

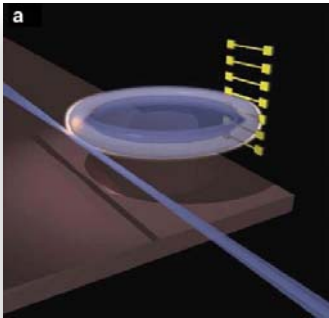
Dykman, M. I., 1978. Heating and cooling of local and quasilocal vibrations by a nonresonance field. *Soviet Physics Solid State* 20, 1306–1311.

Braginskii, V. B., Manukin, A. B., Tikhonov, M. Y., 1970. Investigation of dissipative ponderomotive effects of electromagnetic radiation. *Soviet Physics JETP* 31, 829–830.

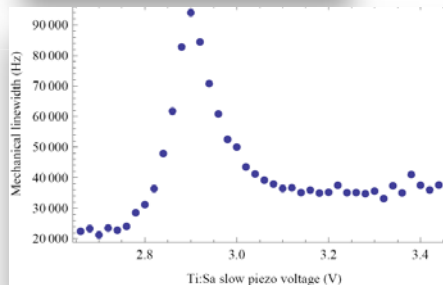
# Outline



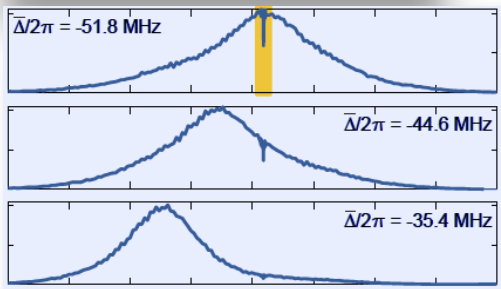
Toroid microresonators and measurement of optomechanical coupling strength



Measurement imprecision below that at The Standard Quantum Limit



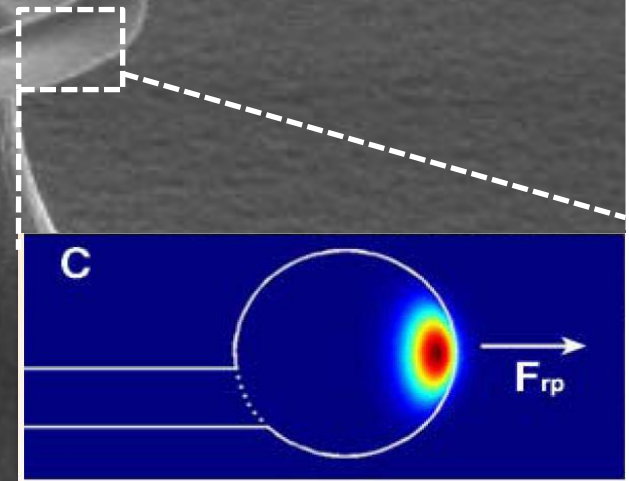
Thermometry of an optomechanically cooled Microresonator in a Helium-3 cryostat



Optomechanically induced transparency

Structures Fabricated at CMI-EPFL

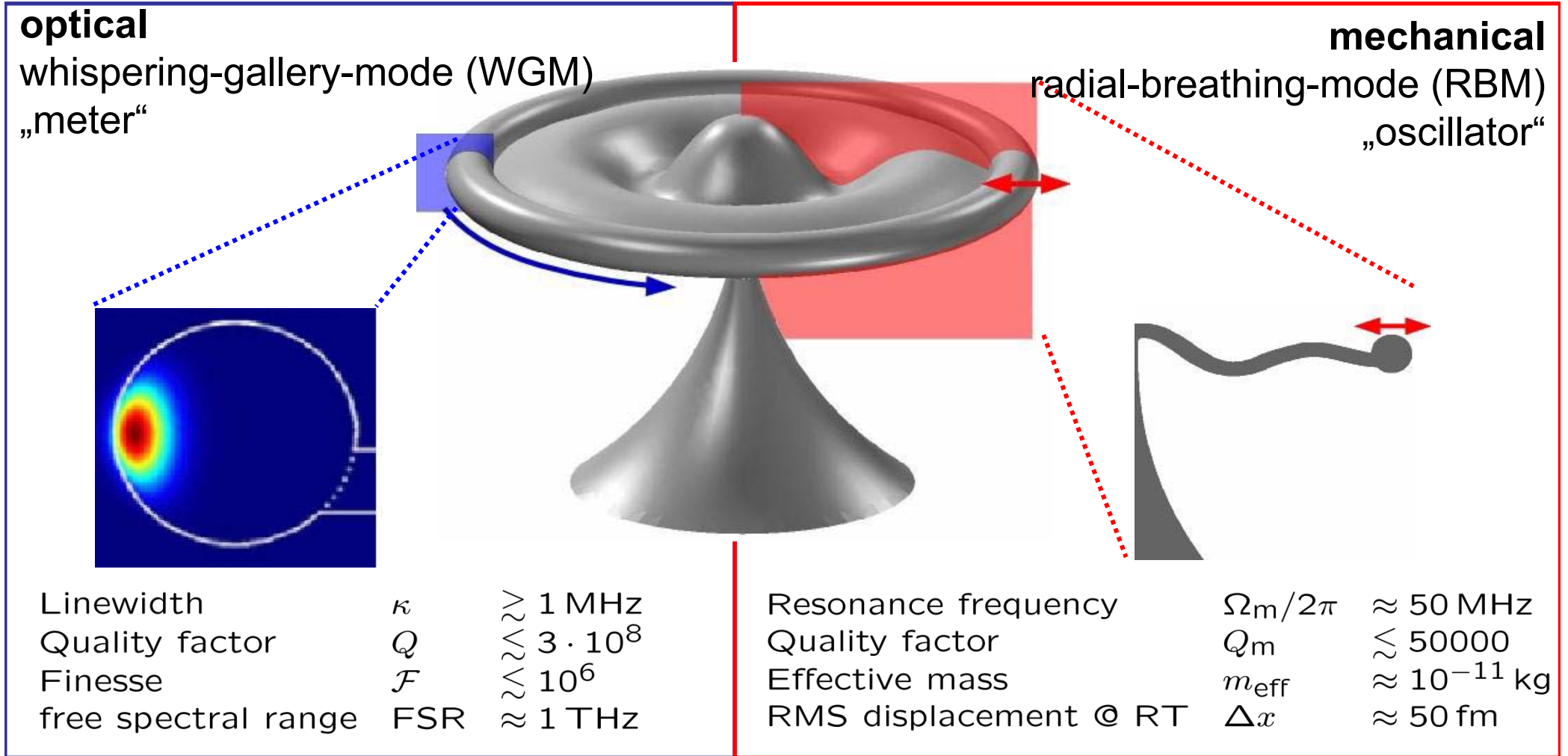
$$Q = \omega \cdot \tau > 10^8$$
$$F > 10^6$$



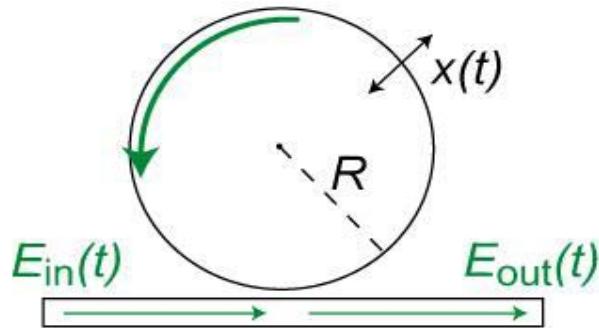
\* D. K. Armani, T. J. Kippenberg, S. M. Spillane, K. J. Vahala. Nature 421, 925-928 (2003).

30.0kV X600 50.0µm

# Natural optomechanical coupling in microresonators



# Optomechanical Coupling

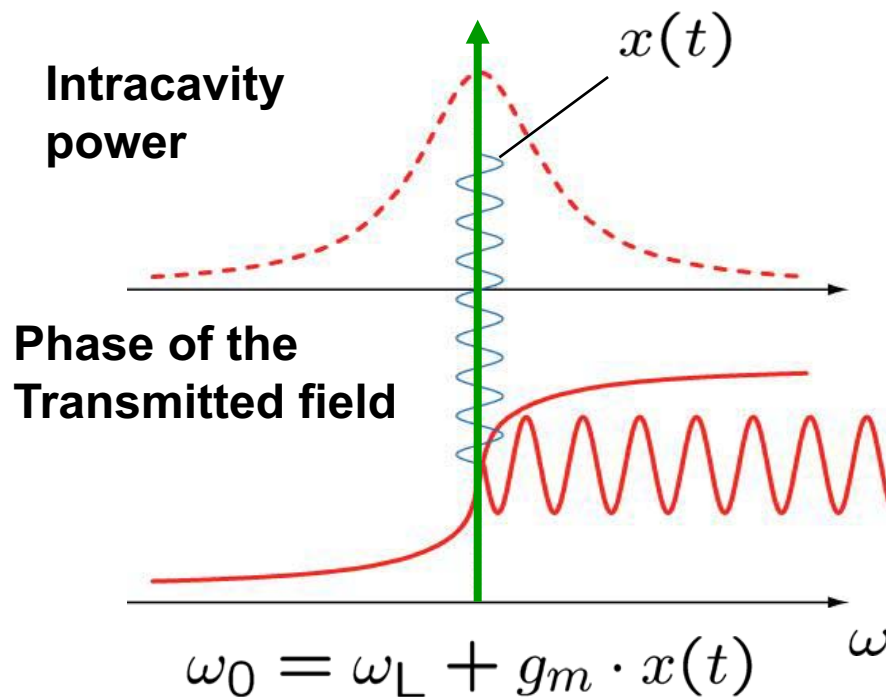


$$H_{int} = \hbar\omega a^\dagger a (1 - g_m x)$$

$$g_m = \frac{\omega}{R} \approx 10\text{GHz/nm}$$

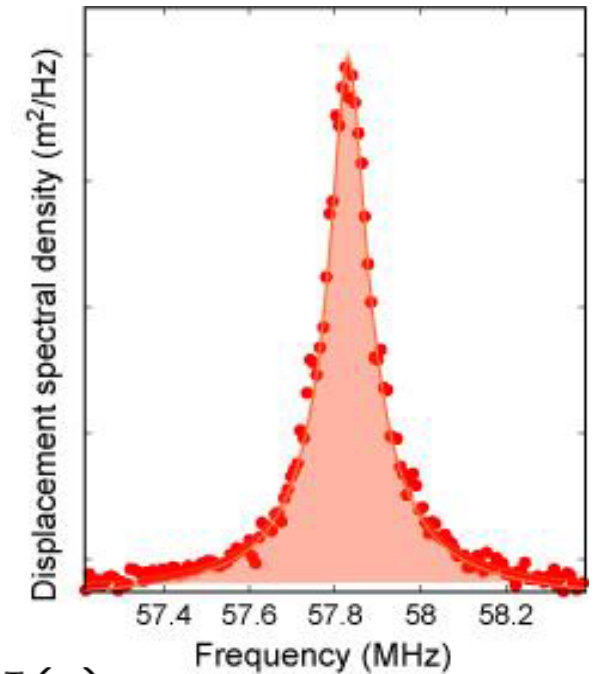
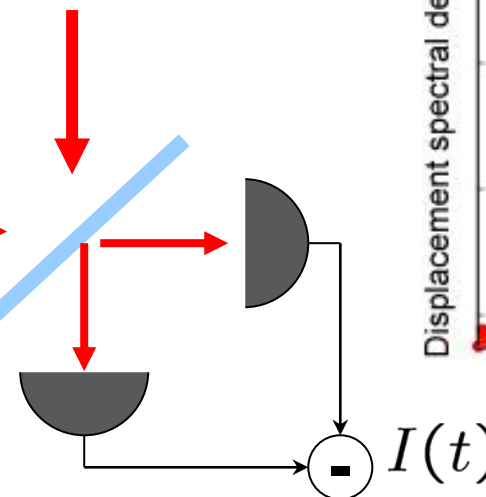
$$F_{rad} = \hbar \cdot g_m$$

$$T_{eff} = \frac{1}{k_B} \int m_{eff} \delta\Omega \cdot x[\Omega]^2 \cdot \Omega^2$$



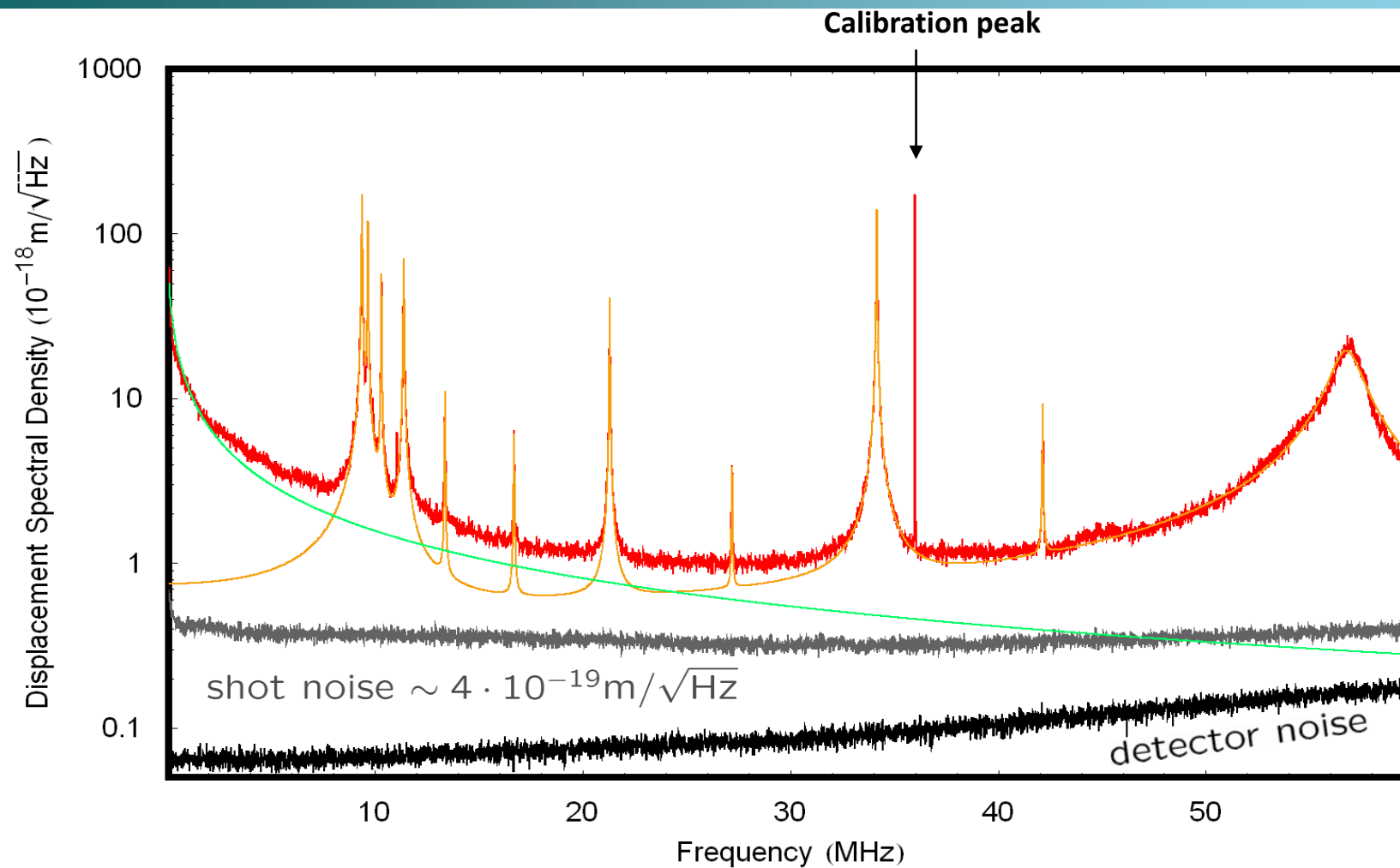
**Quantum limited  
Homodyne Detection**

**LO**





# Quantitative Analysis of Noise Spectra



# Determining the strength of optomechanical coupling

$$\hat{H} = G x_{zpm} \cdot \hat{a}_p^\dagger \hat{a}_p (\hat{a}_m + \hat{a}_m^\dagger)$$

- Optomechanical coupling  $G$  (Hz/nm)  $G \equiv \frac{d\omega(x)}{dx}$   $G = \frac{\omega}{L}$
- The overlap of mechanics and optics is given by: Effective mass ( $m_{eff}$ )
- What is  $L_{eff}$  and what is the amplitude  $x$ ?  
[Eichenfeld et al. Nature]



physically significant parameter is the *vacuum optomechanical coupling parameter*

$$g_0 \equiv G x_{zpm}$$

$$g_0 \equiv G x_{zpm}$$

How to **measure** the *vacuum optomechanical coupling parameter* in an experiment directly?

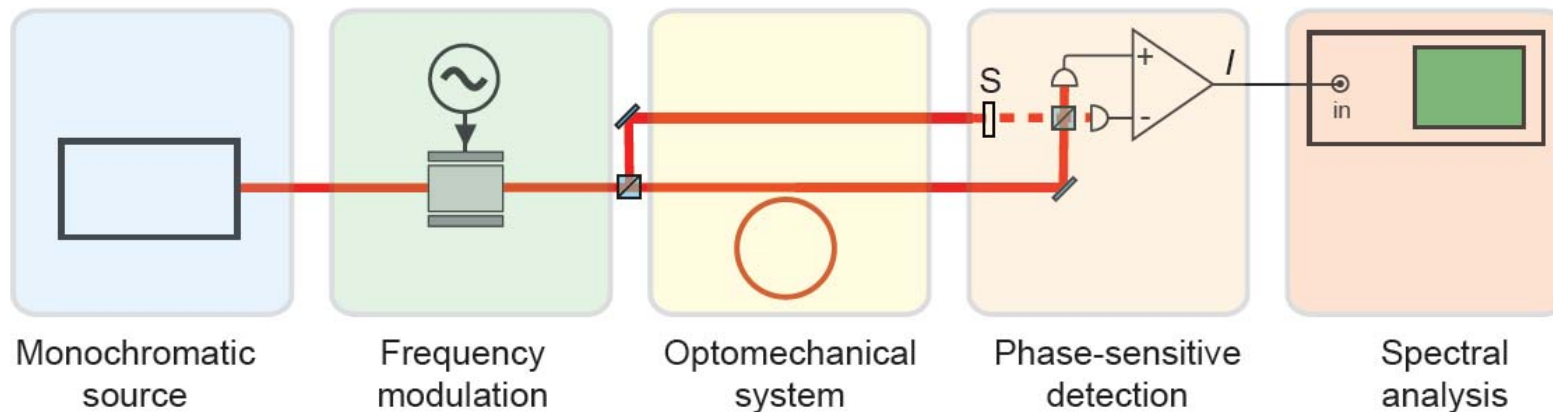
Interesting relation

$$\int_{-\infty}^{\infty} S_{\omega\omega}(\Omega) \frac{d\Omega}{2\pi} = S_{\omega\omega}(\Omega_m) \Gamma_m / 2 = 2 \langle n \rangle g_0^2$$

Cavity frequency noise gives directly the vacuum optomechanical coupling strength. *No finite element simulation needed.*

$$g_0 = \sqrt{\frac{1}{2 \langle n \rangle} \int_{-\infty}^{\infty} S_{\omega\omega}(\Omega) \frac{d\Omega}{2\pi}}$$

# Measuring the cavity frequency noise spectrum



- Scheme to calibrate the cavity frequency (phase) noise

$$S_{\phi\phi} = 2\pi \frac{1}{2} (\delta(\Omega - \Omega_{\text{mod}}) + \delta(\Omega + \Omega_{\text{mod}})) \cdot \frac{\phi_0^2}{2}$$

- Transduction of mechanical motion and phase modulation is the same (\*)

$$S_{II} = K(\Omega, \kappa, \Delta) (S_{\phi\phi}^{\text{mod}}(\Omega) + S_{\phi\phi}^{\text{cav}}(\Omega))$$

Photocurrent  
SD

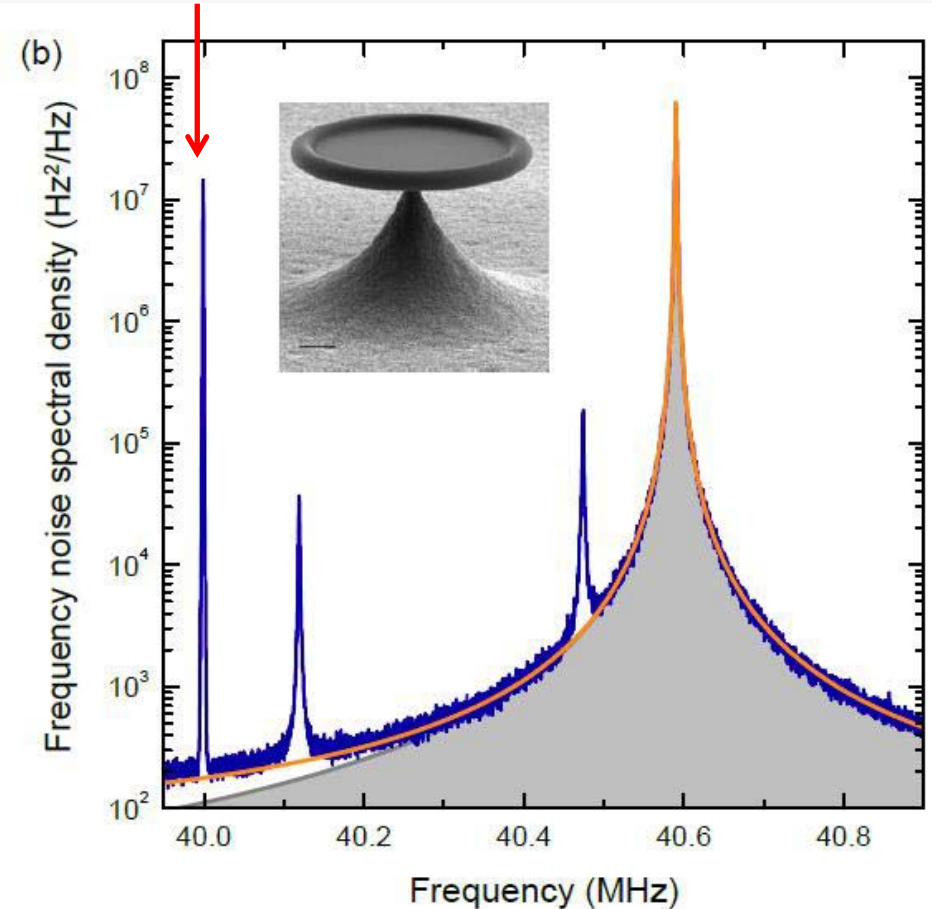
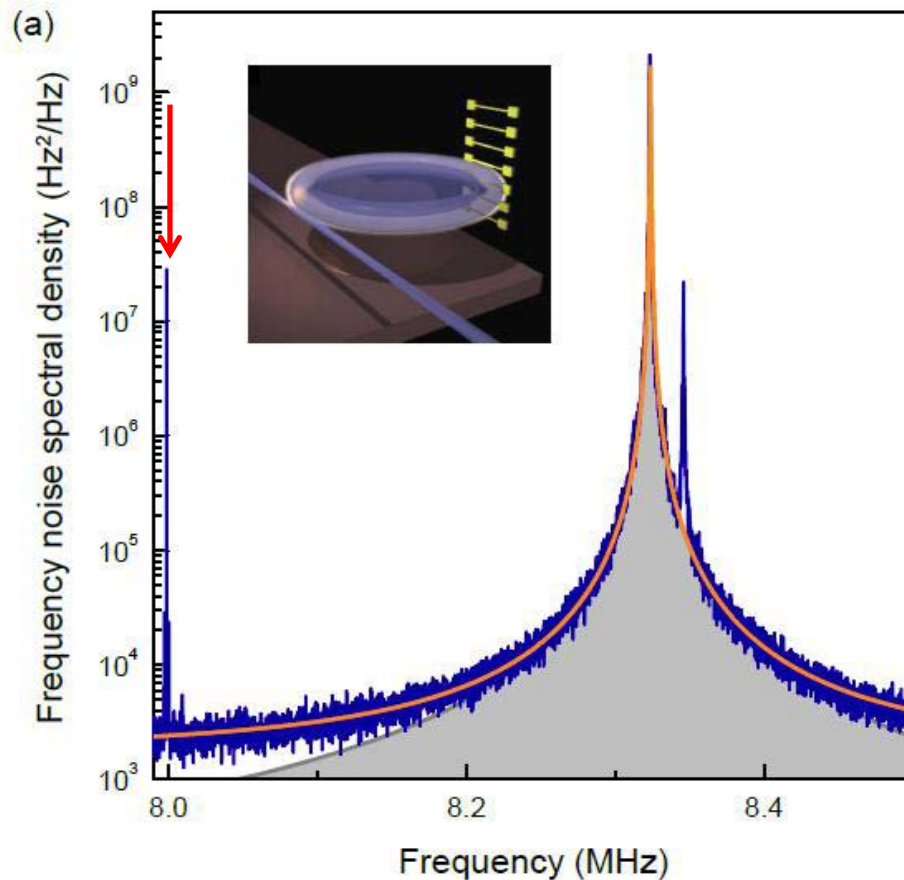
Transduction

Calibration phase  
modulation SD

Phase noise SD of  
the cavity

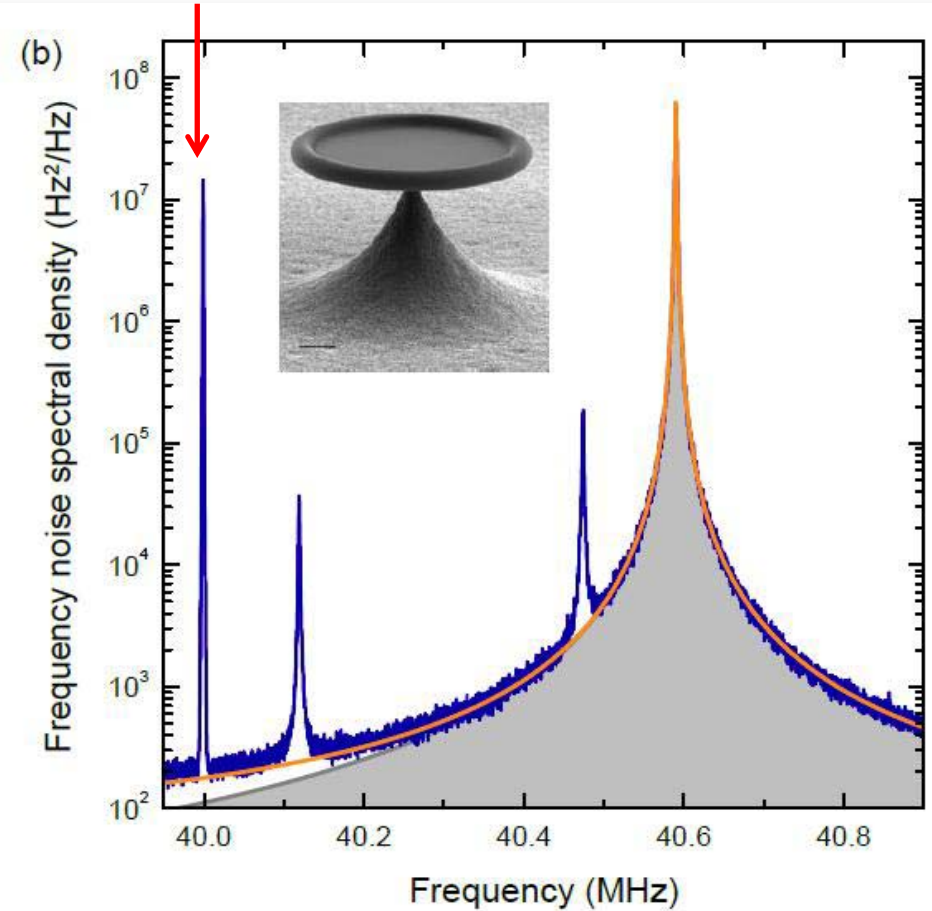
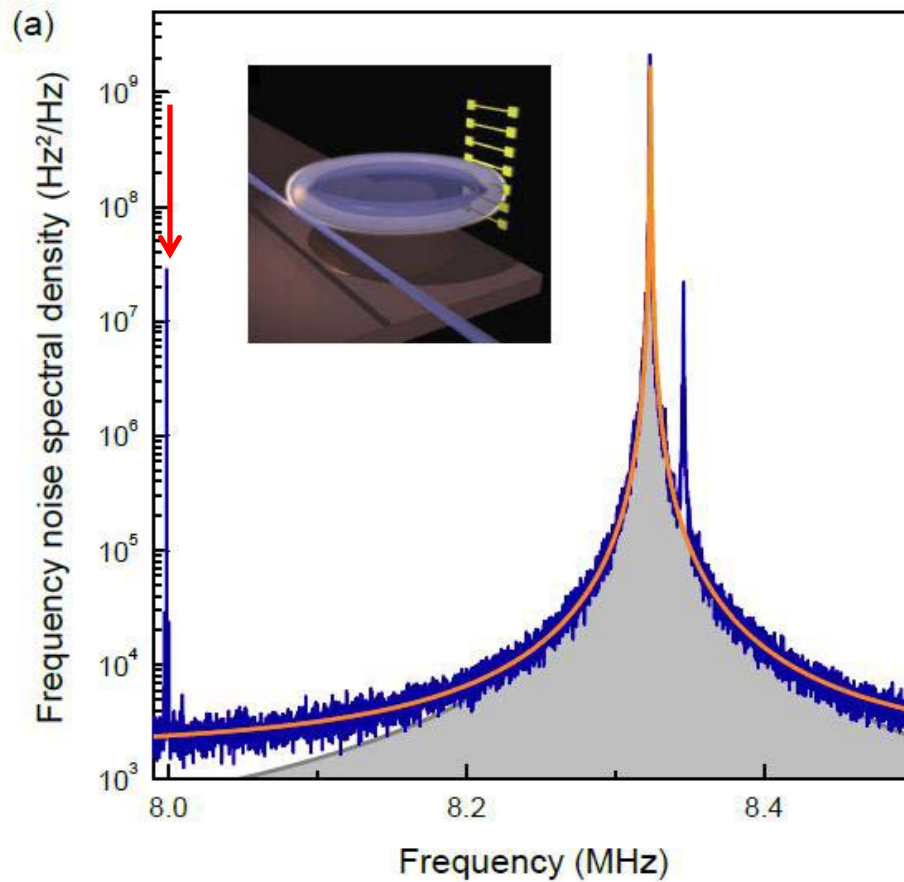
\* Anetsberger, ML Gorodetsky, Schliesser, Kippenberg (to appear: arXiv 2010)

# Determining the strength of optomechanical coupling



$$\langle \delta\omega^2 \rangle = \int_{-\infty}^{\infty} S_{\omega\omega}(\Omega) \frac{d\Omega}{2\pi} = 2\langle n \rangle g_0^2 \quad g_0 \equiv Gx_{zpm}$$

# Determining the strength of optomechanical coupling



$$g_0 = 2\pi \cdot 500\text{Hz}$$

$$g_0 = 2\pi \cdot 5\text{kHz}$$

$$g_0 = 2\pi \cdot 1\text{kHz}$$

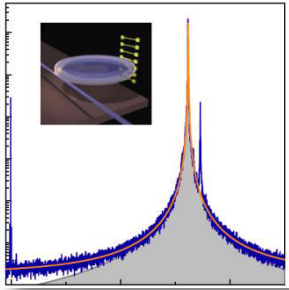
$$g_0 \equiv Gx_{zpm}$$

- Microwave coupled nanobeams  $\sim 10$ - $100$  Hz [ $\sim 1$  kHz Teufel]
- Cantilever based systems  $g_0 \sim 10$ - $100$  Hz
- Toroid microcavities  $g_0 \sim 1$  kHz
- Near field optical beams  $g_0 \sim 1$ - $10$  kHz
- Photonic crystals (predicted)  $g_0 \sim 100$ - $500$  kHz

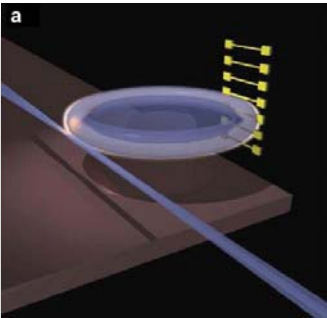
Some relevant formulas:

$$\Gamma_{cool} \approx g_0 \frac{g_0}{\kappa} \bar{n}_c \quad \Gamma_{cool} \approx \frac{4g_0^2}{\Omega_m^2} \frac{P_{in}}{\hbar\omega_L} \quad P_{in}^{SQL} = \frac{\kappa^2}{16g_0^2} \Gamma_m \left[ 1 + \frac{4\Omega_m^2}{\kappa^2} \right]$$

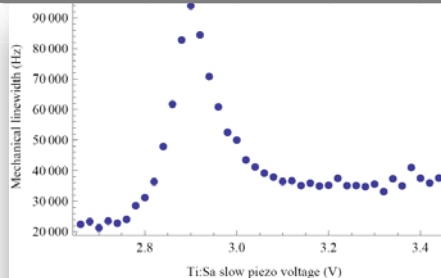
# Outline



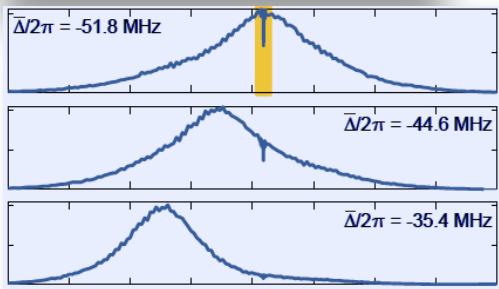
Measuring the vacuum optomechanical coupling strength



Measurement imprecision below that at the Standard Quantum Limit



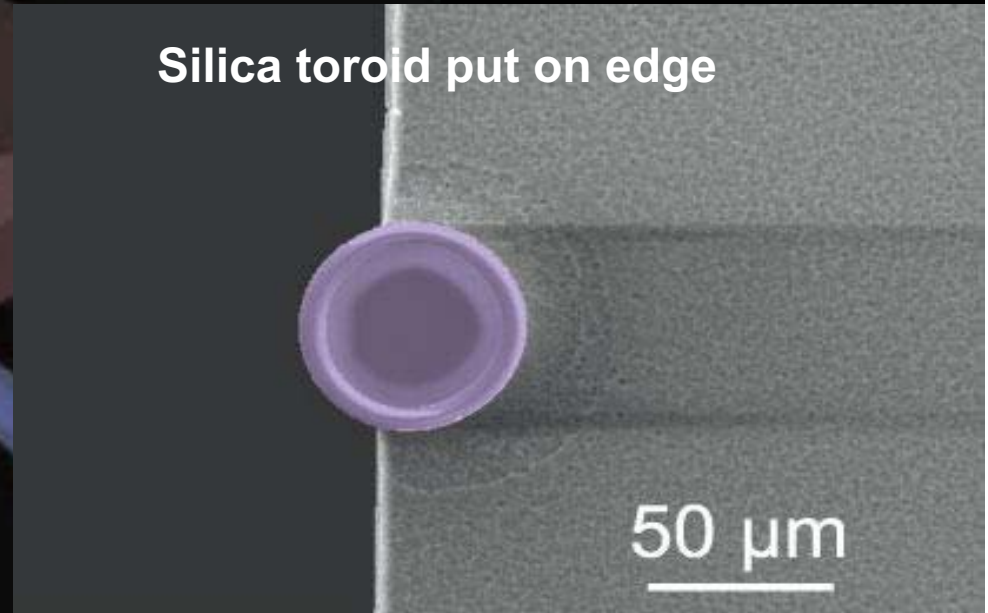
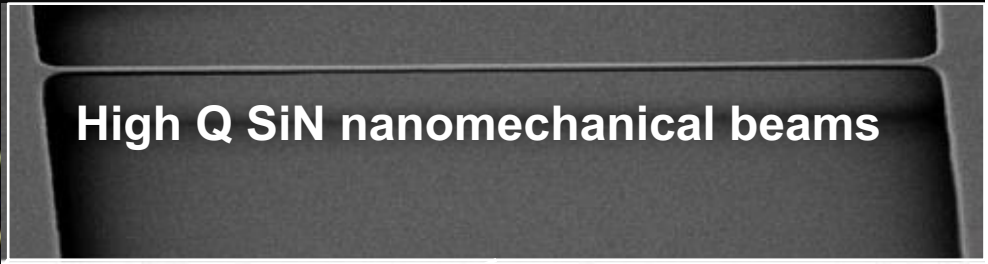
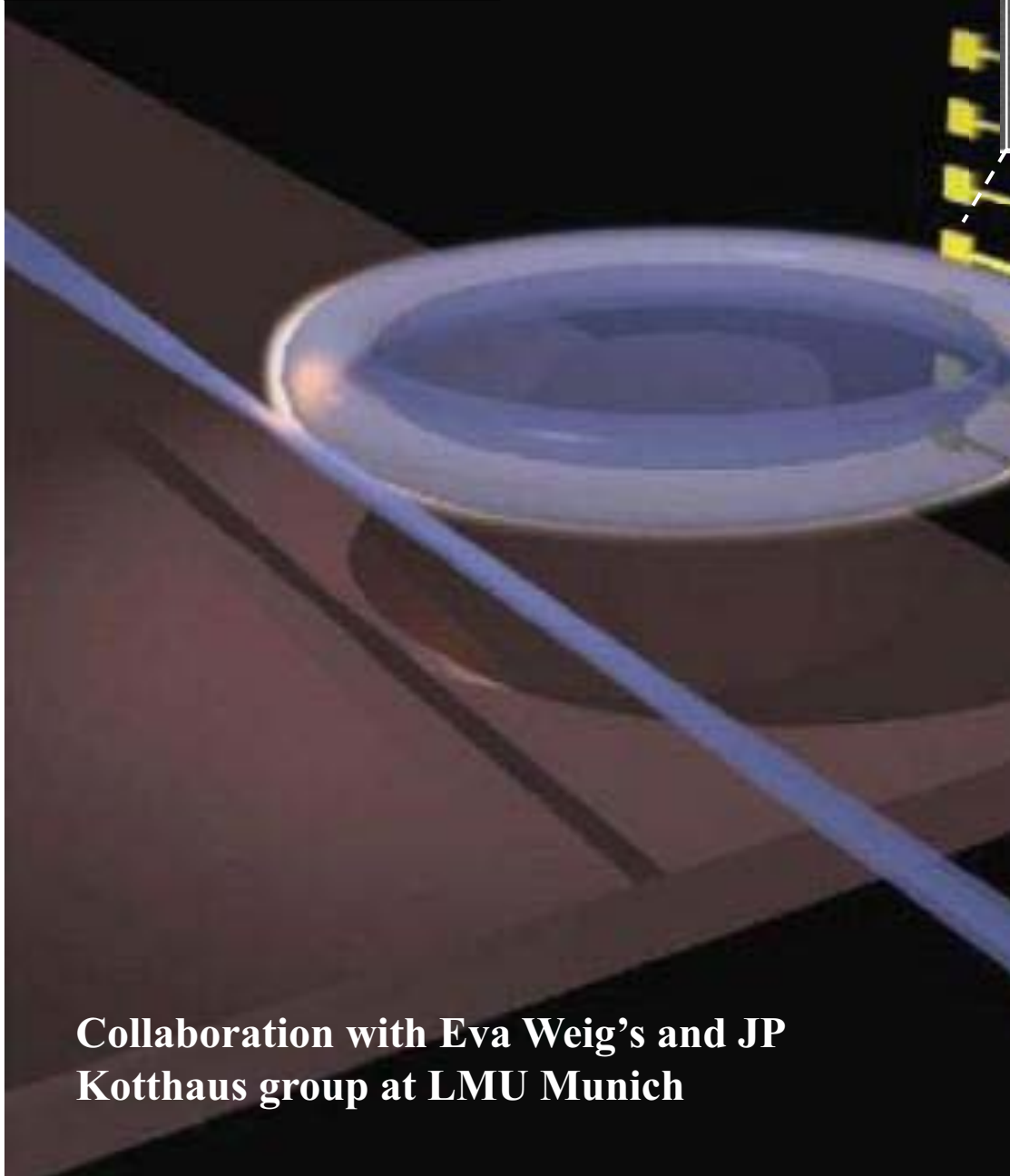
Thermometry of an optomechanically cooled Microresonator in a Helium-3 cryostat



Optomechanically induced transparency

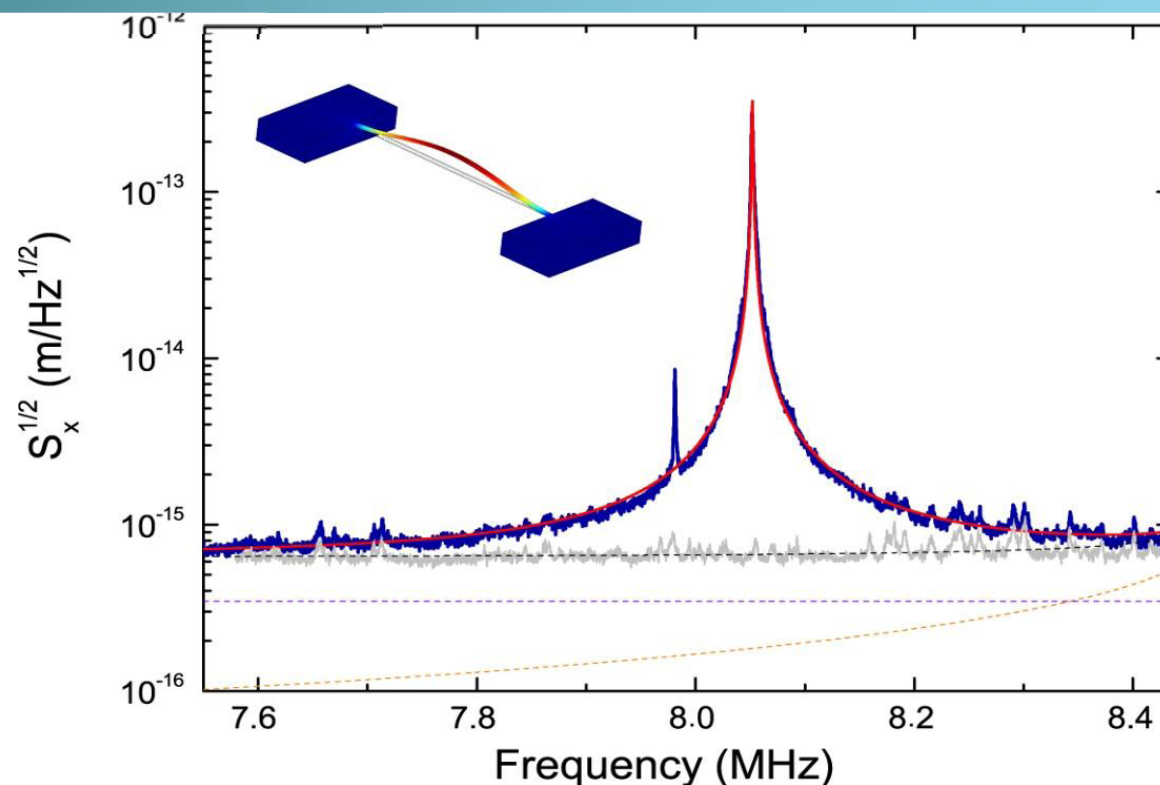
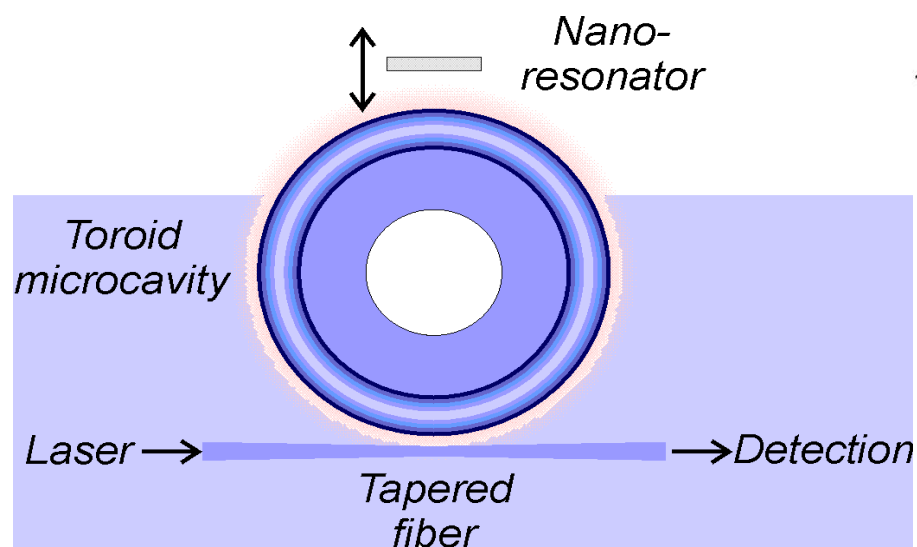
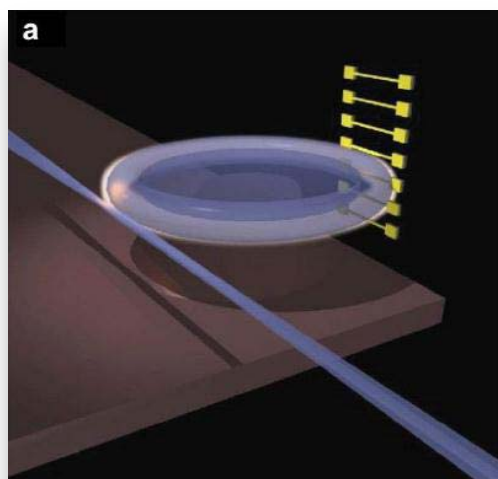


# Evanescent sensing



Collaboration with Eva Weig's and JP Kotthaus group at LMU Munich

# Near field displacement transduction



Parameters

$$g_m = \frac{d\omega}{dx} \approx 10 \text{ MHz/nm}$$

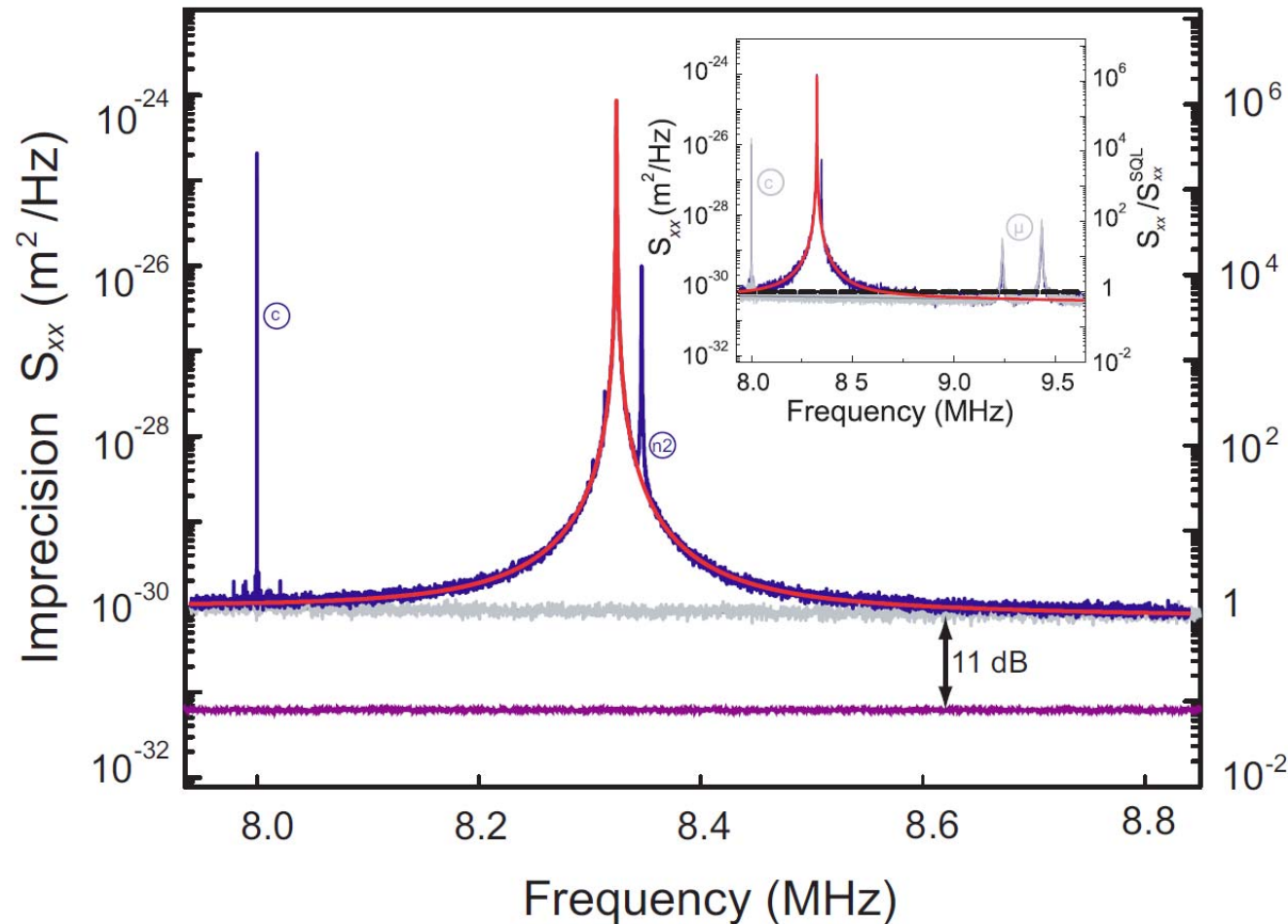
$$m_{\text{eff}} \approx 1 \text{ pg}$$

$$Q_m > 100,000$$

$$\Omega_m \approx 10 \text{ MHz}$$

$$\kappa < 10 \text{ MHz} (F > 250,000)$$

# Near field displacement transduction

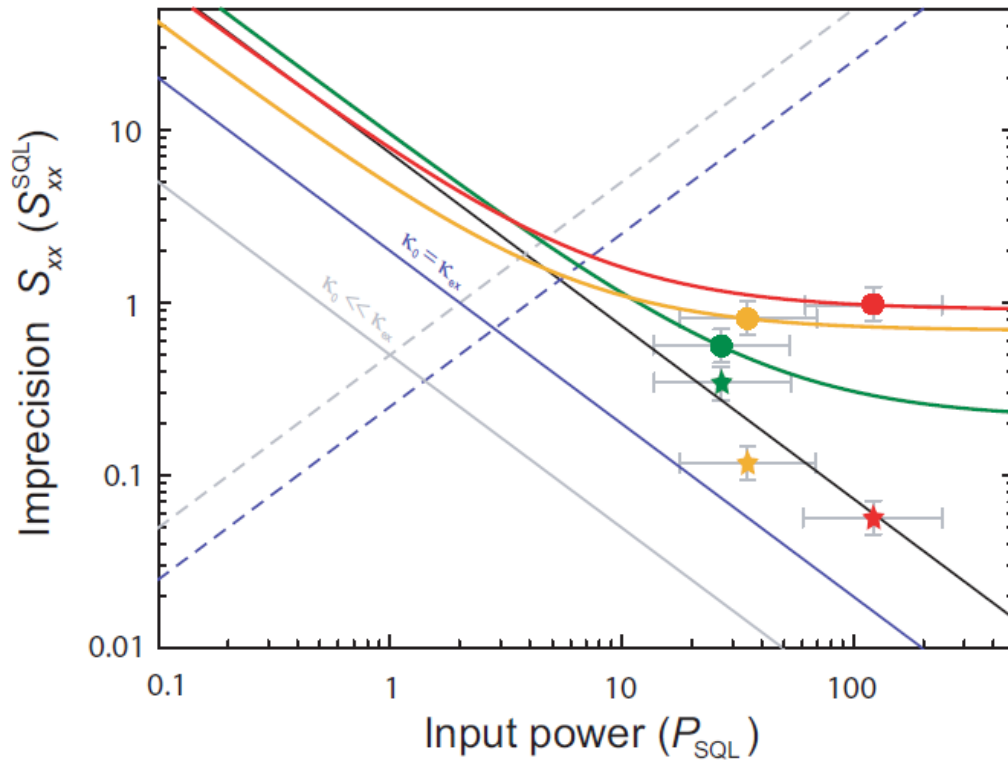


$$S_{xx} = 0.08 \times S_{xx}^{\text{SQL}} \text{ or } (250 \text{ am}/\sqrt{\text{Hz}})^2$$

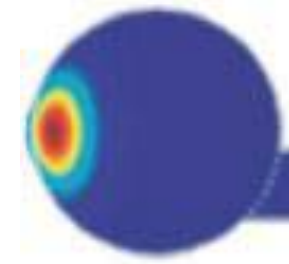
*Optical fields:* Anetsberger et al. <http://arxiv.org/1003.3752f>

*Microwave fields:* Teufel, Donner, Castellanos-Beltran, Harlow, Lehnert Nature Nanotechnology 4, 820 - 823 (2009)

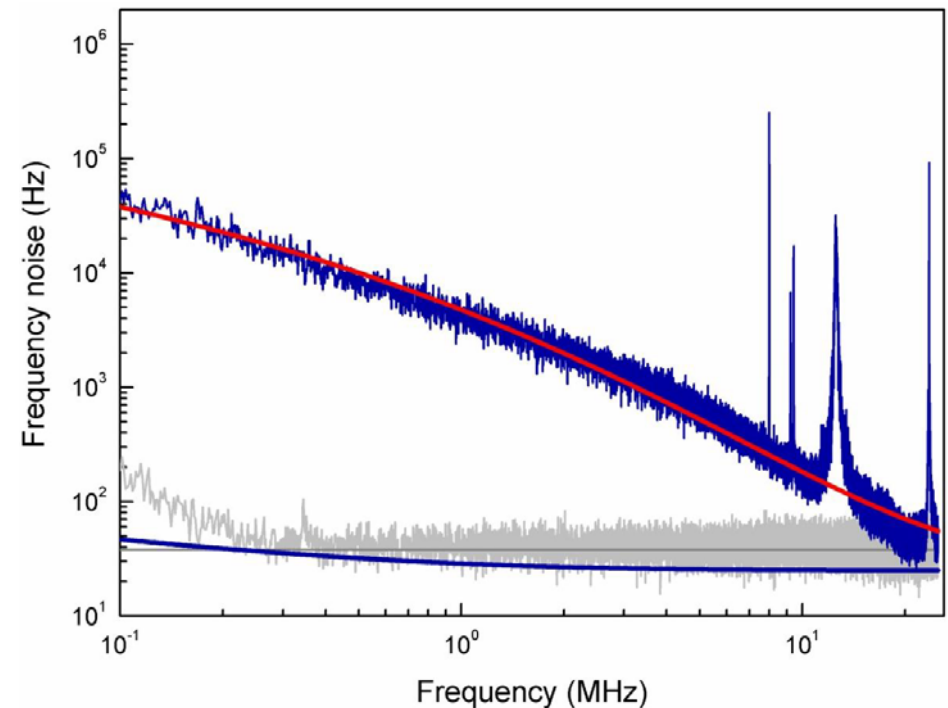
# Near field displacement transduction



$$\langle \Delta T^2 \rangle = \frac{k_B T}{\rho C V}$$



$$S_{\omega}^{\text{thr}}[\Omega] = \omega^2 \left( \frac{dn}{dT} \right)^2 k_B T^2 \frac{\sqrt{l}}{\pi^{3/2} n^2 R^2 \sqrt{\Omega \lambda^2 / D}} \frac{1}{\sqrt{1 - (b/d)^2}} \frac{1}{(1 + (\Omega \tau)^{3/4})}$$



- Imprecision limited by cavity frequency noise.
- Access to QBA dominated regime

- Quantum backaction<sup>[1]</sup> can reach values of unity even at *room temperature*

$$\frac{S^{\text{qba}}}{S^{\text{th}}} = \left( \frac{g/2\pi}{20 \text{ MHz/nm}} \right)^2 \underbrace{\left( \frac{15 \text{ pg}}{m_{\text{eff}}} \right) \left( \frac{Q_m}{10^6} \right) \left( \frac{1 \text{ MHz}}{\Omega_m/2\pi} \right)}_{[2]} \left( \frac{4 \text{ MHz}}{\kappa/2\pi} \right)^2 \left( \frac{P}{100 \mu\text{W}} \right) \left( \frac{\lambda}{780 \text{ nm}} \right) \left( \frac{300 \text{ K}}{T} \right)$$

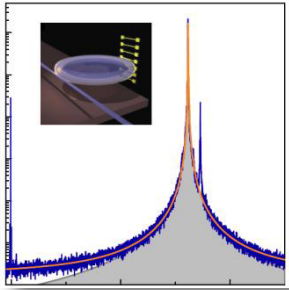
- QBA dominated regime is prerequisite for *room temperature* Quantum Optomechanical Experiments.

[1] Fabre *et al.* PRA 49, 1337 (1994), Heidmann *et al.* Applied Physics B 64, 173 (1997).

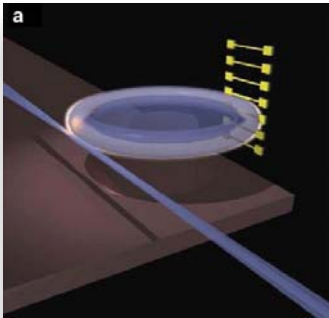
[2] Verbridge *et al.* APL 92, 013112 (2008).

[3] Verlot, P., Tavernarakis, A., Briant, T., Cohadon, P.-F. & Heidmann, Phys. Rev. Lett. 102, 103601 (2009)

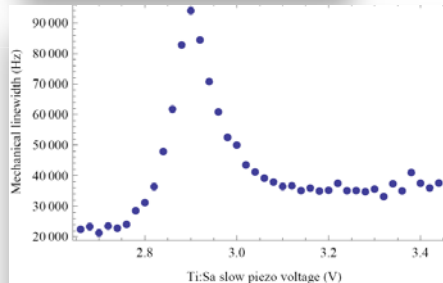
# Outline



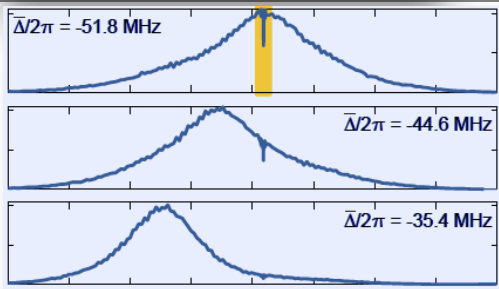
Measuring the vacuum optomechanical coupling strength



Measurement imprecision below that at The Standard Quantum Limit

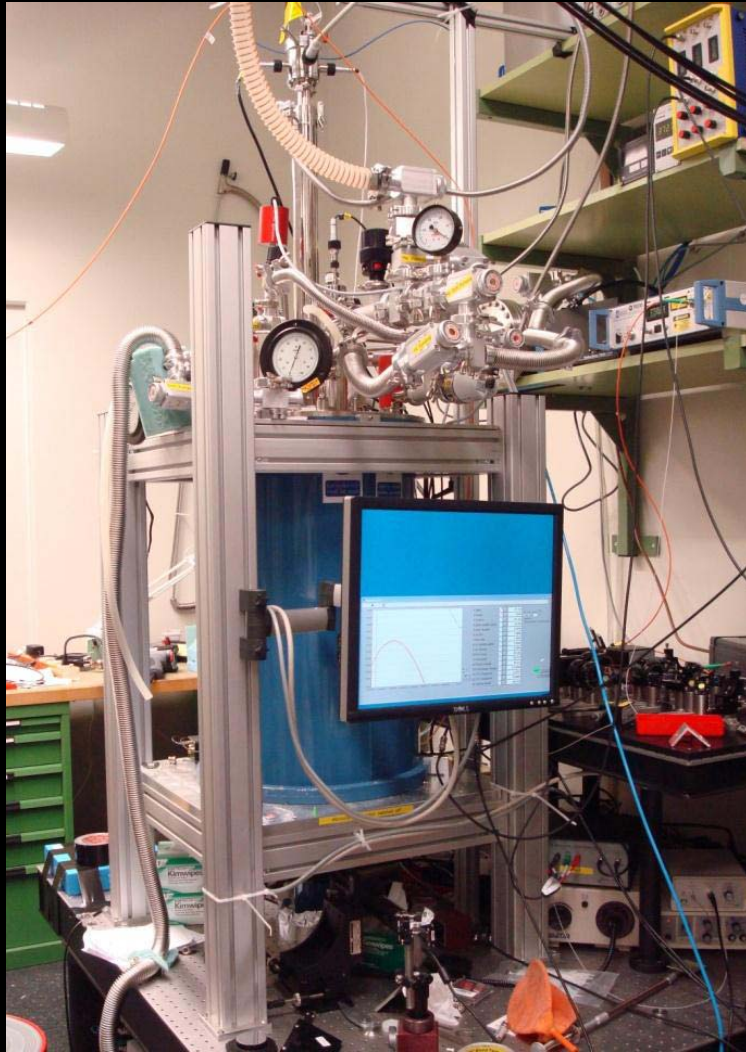


Thermometry of an optomechanically cooled Microresonator in a Helium-3 cryostat



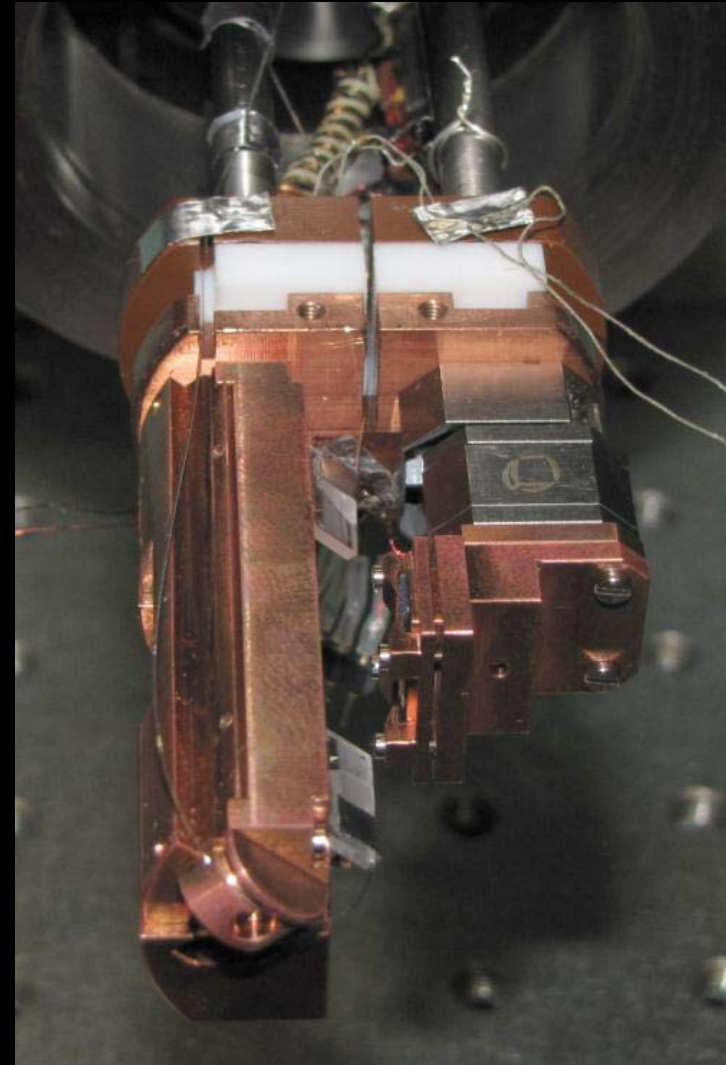
Optomechanically induced transparency

# Optomechanics at Helium-3 Temperatures (600 mK)



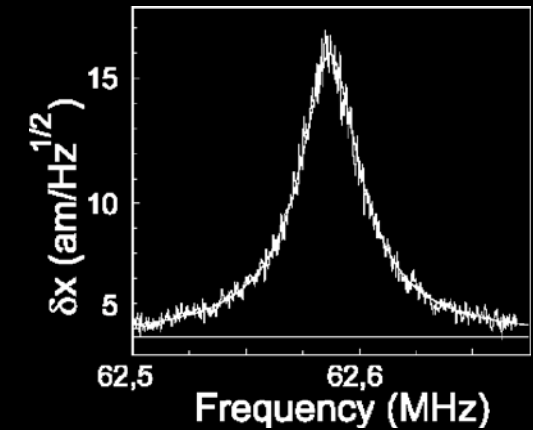
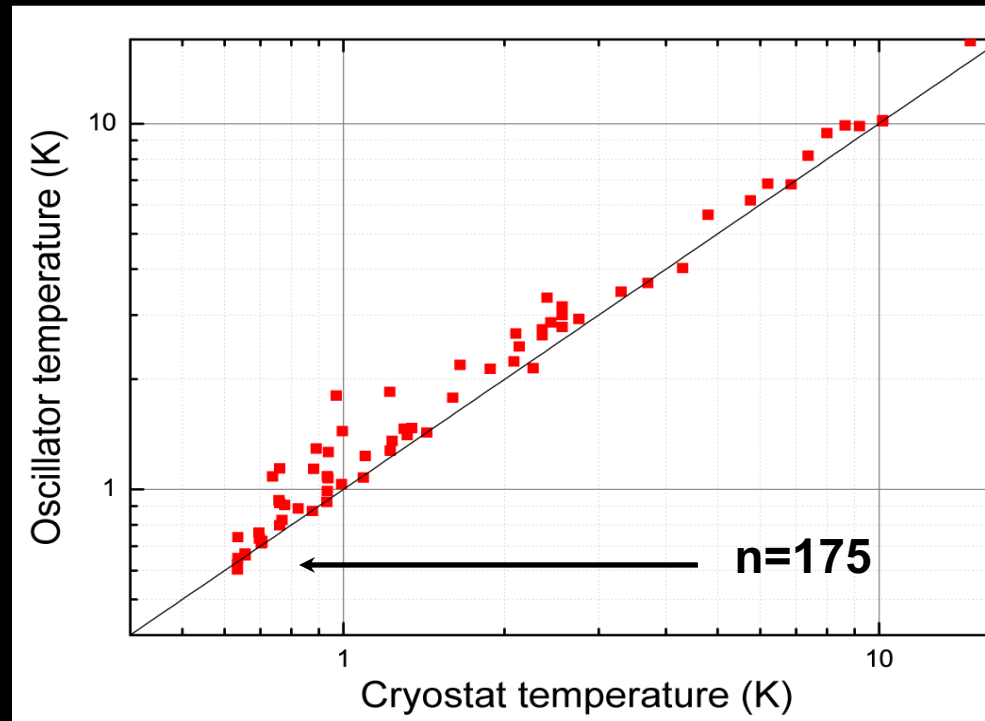
Holder for tapered fiber

Piezo Positioners



38mm

# Optomechanics at Helium-3 Temperatures (600 mK)



$$\Omega_m \approx 50 - 75 \text{ MHz} \quad T = 600 \text{ mK}$$

$$n = \frac{k_B T}{\hbar \Omega} \approx 175 - 250$$

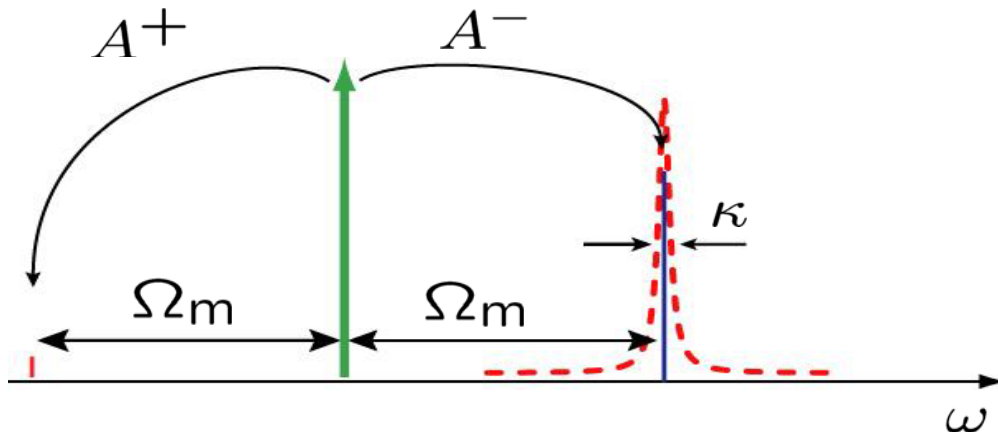
**At low power excellent thermalization observed with the cryostat**



# Resolved sideband cooling in a cryostat

## Resolved Sideband Limit and Cooling:

- Suppresses direct excitation of the optical mode, i.e. reduces resonant heating



$$n_f \approx \frac{\Gamma_m}{\Gamma_{cool}} n_i + \frac{A_+}{A_- - A_+}$$

Reservoir heating

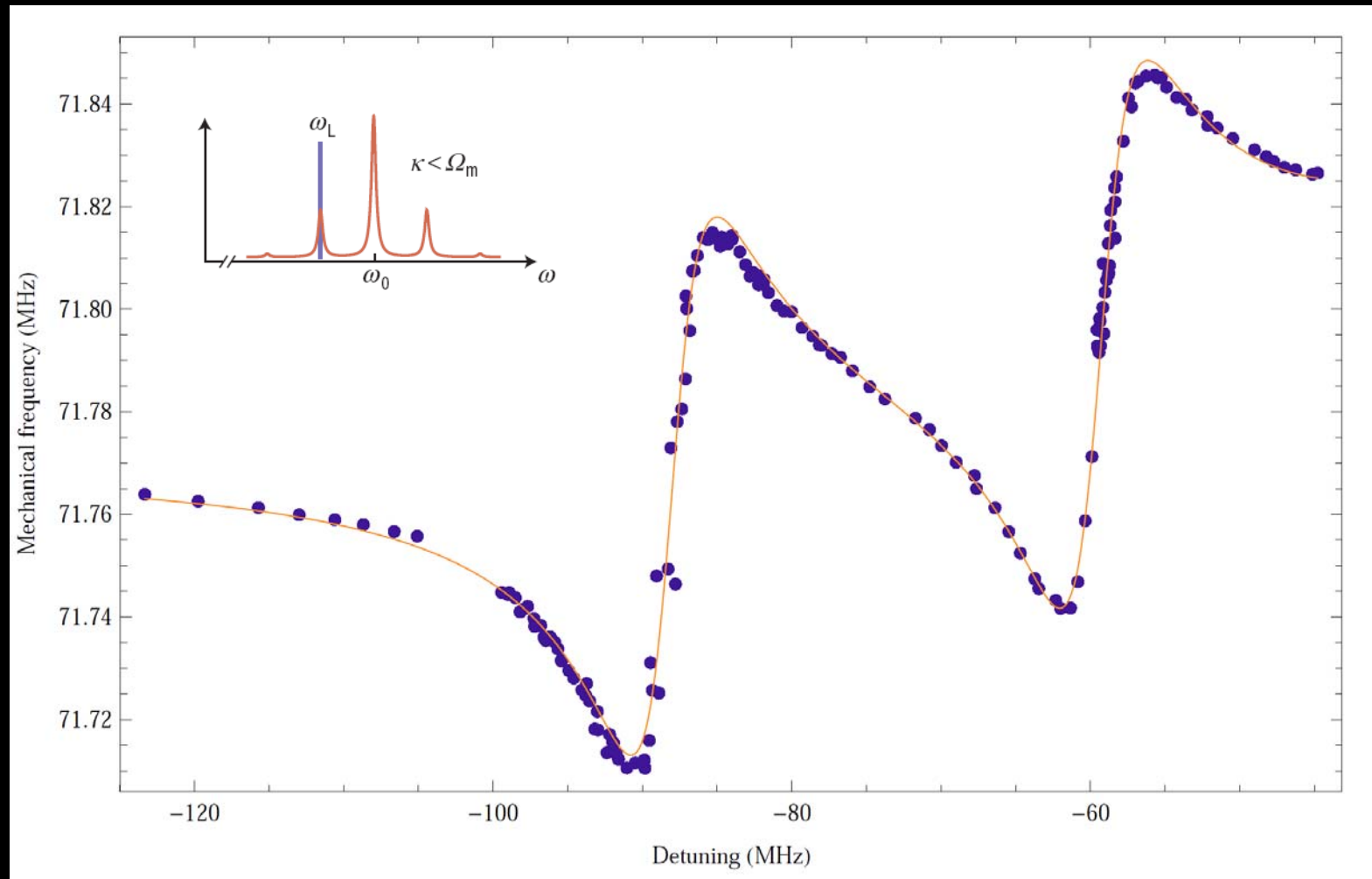
$$\Gamma_{cool} \approx \frac{4g_0^2}{\Omega_m^2} \frac{P_{in}}{\hbar\omega_L}$$

Quantum Backaction

$$n_b \approx \frac{\kappa^2}{16\Omega_m^2} \ll 1$$

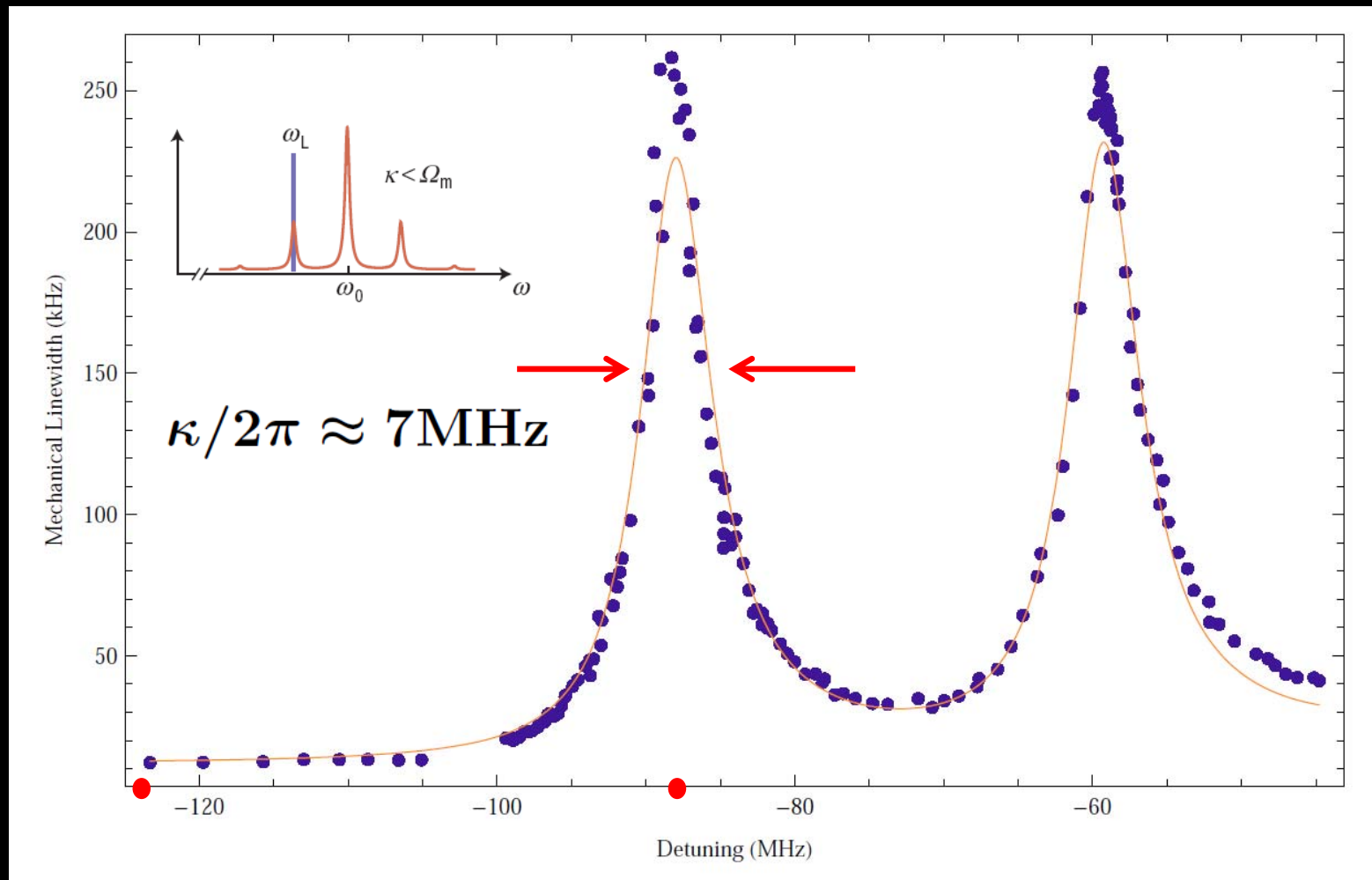
Technical problem: relative detuning of the laser is not precisely known!

# Mechanical frequency shift (in phase radiation pressure)



- Excellent agreement with theory of Dynamical backaction cooling
- Frequency can be measured with high accuracy

# Damping rate (quadrature component of radiation pressure)

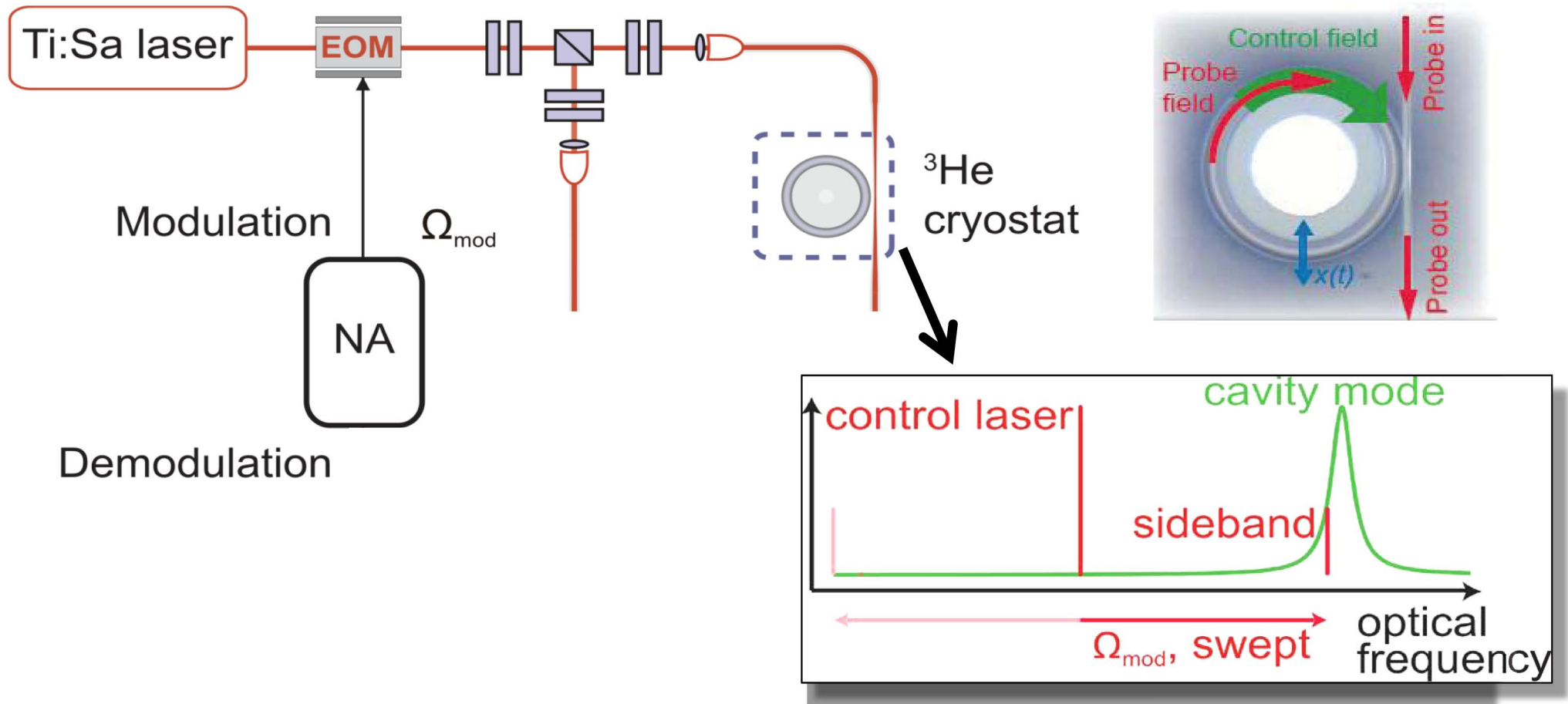


- Maximum cooling rate **> 250 kHz**
- Intrinsic mechanical damping **12 kHz**

$$\frac{\Omega_m}{\kappa} = 10$$

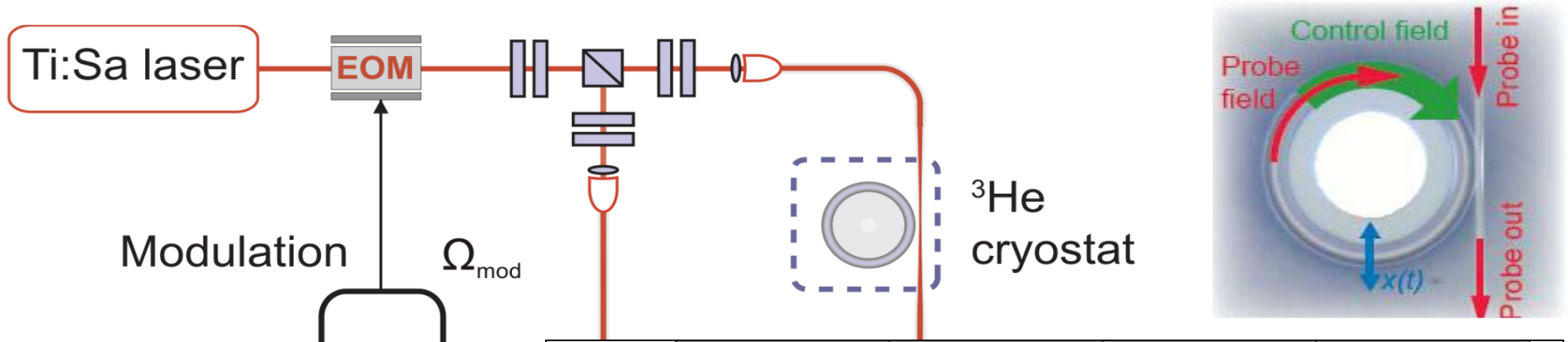
$$\Delta P_{abs} = P_{in} \cdot \frac{4\kappa^2}{\Omega_m^2}$$

# Detuning calibration



- On the fly line-width
- Precise detuning, since centre of Lorentzian is much better known than its width.

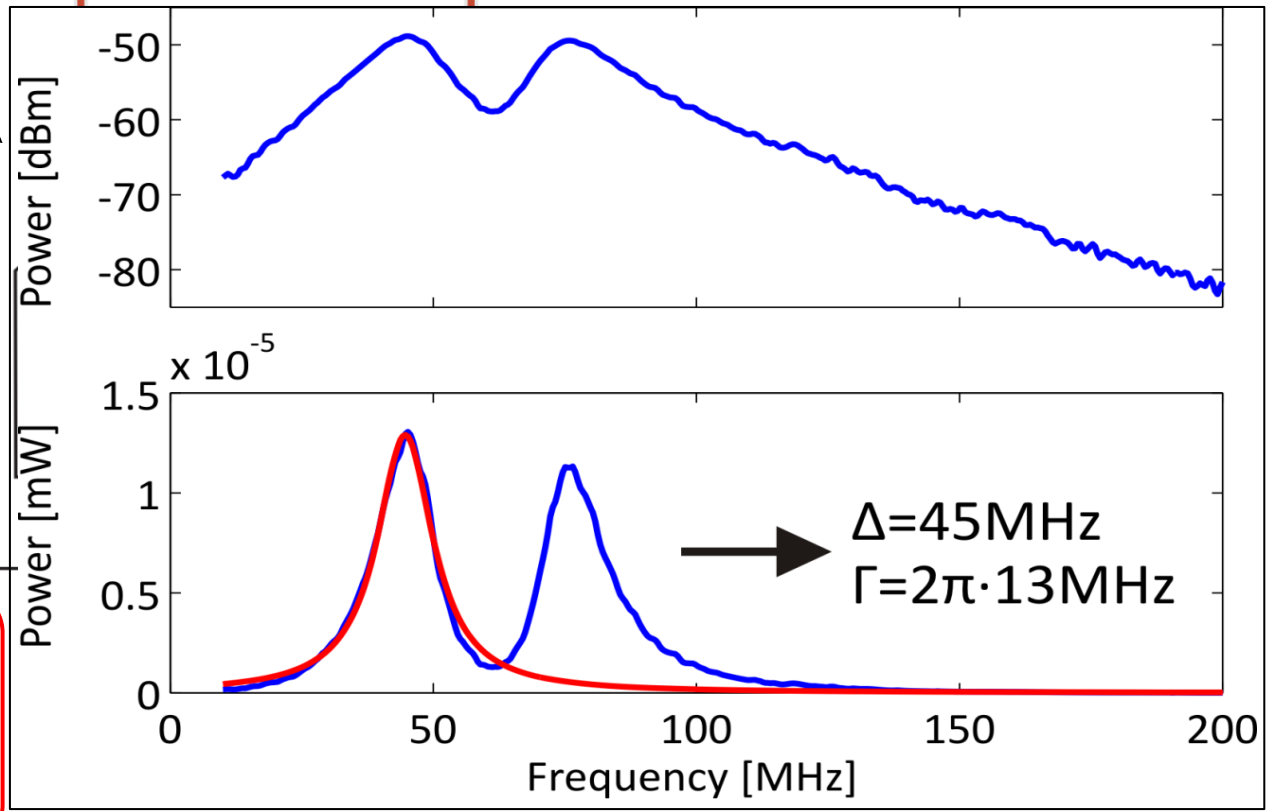
# Detuning calibration



Modulation  $\Omega_{\text{mod}}$

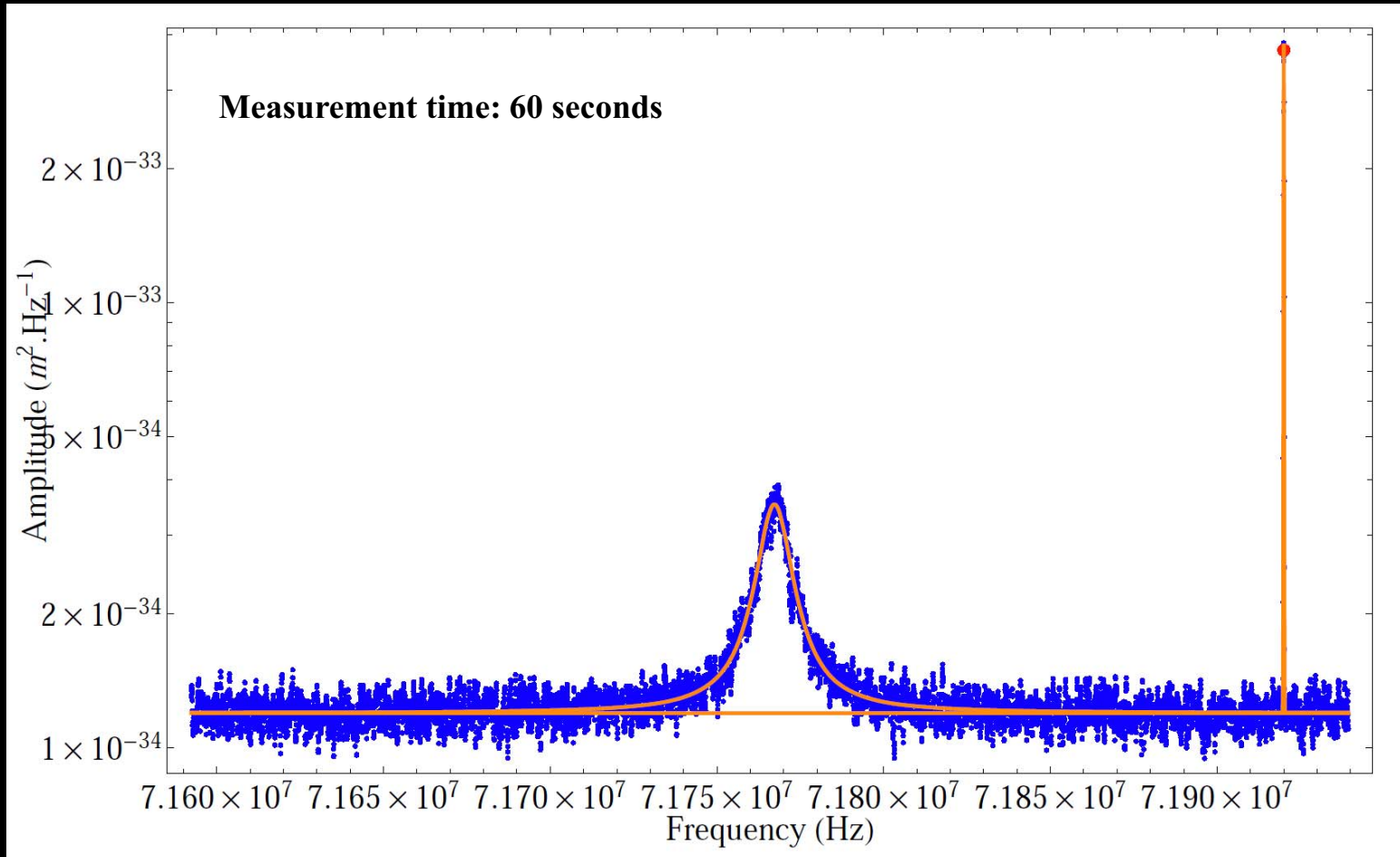
Demodulation

NA



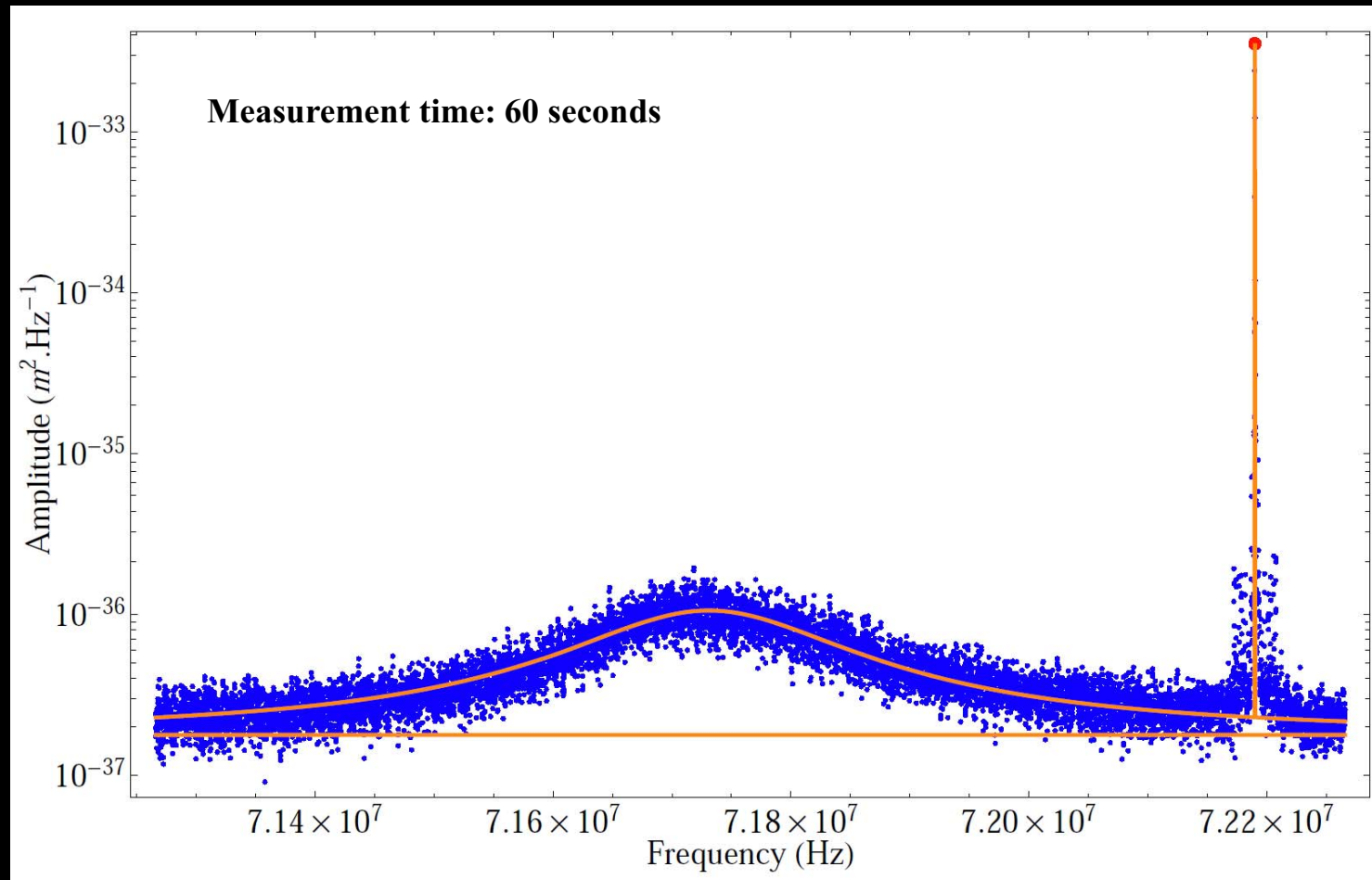
- On the fly line-width
- Precise *detuning*, since centre of Lorentzian is much better known than its width.

# Displacement spectral density



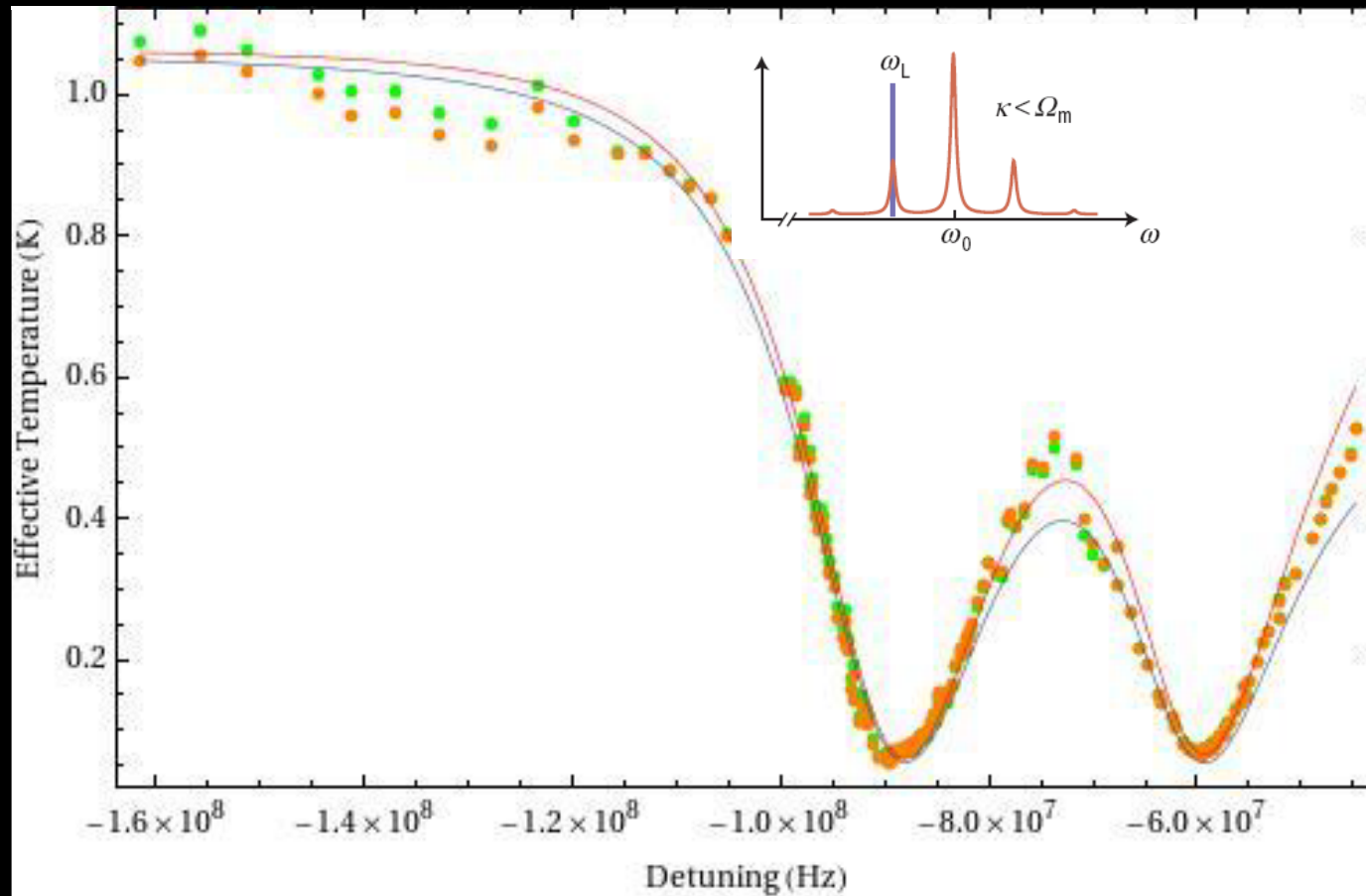
- Excellent agreement with theory of Dynamical backaction cooling  $\Gamma_{eff} = 12\text{kHz}$
- Frequency can be measured with high accuracy
- Measurement time 60 seconds

# Displacement spectral density



- Excellent agreement with theory of Dynamical backaction cooling  $\Gamma_{eff} \approx 250\text{kHz}$
- Frequency can be measured with high accuracy

# Resolved sideband in 3He cryostat



$$T_i = 1100\text{mK} : Q_m \approx 6100, \Omega_m = 71\text{MHz}, \kappa/2\pi = 7\text{MHz}$$

$$\bar{n}_i \approx 320 \quad \bar{n}_f = \bar{n}_i \frac{\Gamma_m}{\Gamma_m + \Gamma_{opt}} \approx 15$$



VOLUME 24, NUMBER 4

PHYSICAL REVIEW LETTERS

26 JANUARY 1970

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ACCELERATION AND TRAPPING OF PARTICLES BY RADIATION PRESSURE

A. Ashkin

Bell Telephone Laboratories, Holmdel, New Jersey 07733

(Received 3 December 1969)



Tapered fiber.

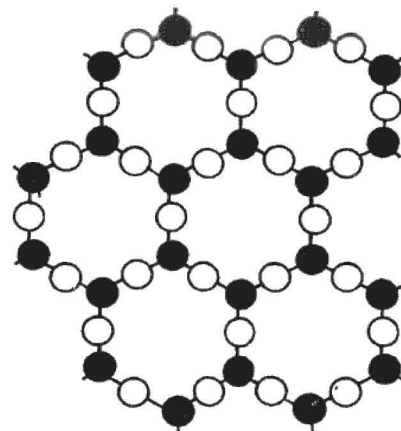
# Low Temperatures: Two level fluctuators (TLS)

## Anomalous Low-temperature Thermal Properties of Glasses and Spin Glasses

By P. W. ANDERSON†, B. I. HALPERIN and C. M. VARMA

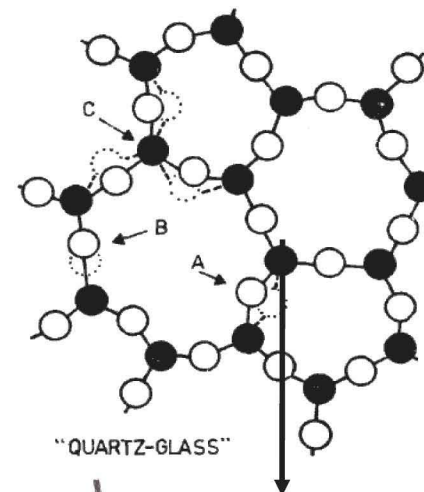
Bell Laboratories, Murray Hill, New Jersey 07974

Crystalline structure



"QUARTZ-CRYSTAL"

[2]

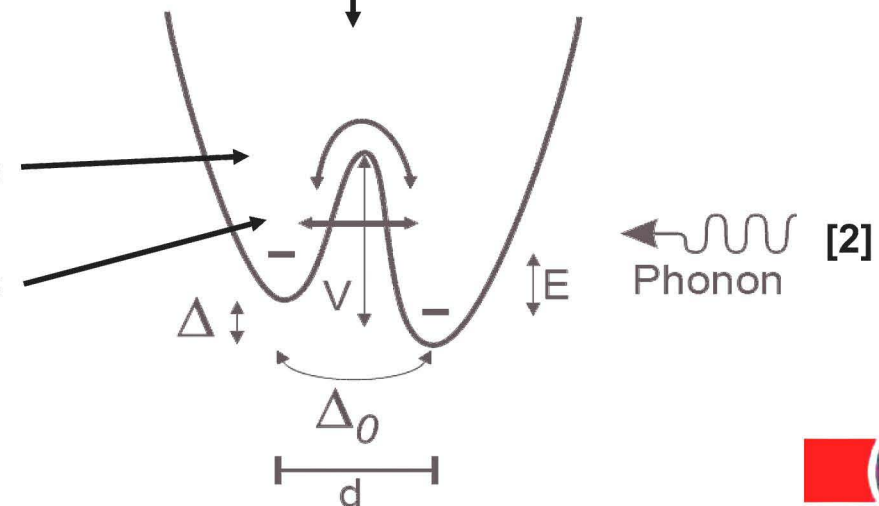


Glassy structure:  
Different possible  
conformations

2-level system approximation [1]

Thermally activated process

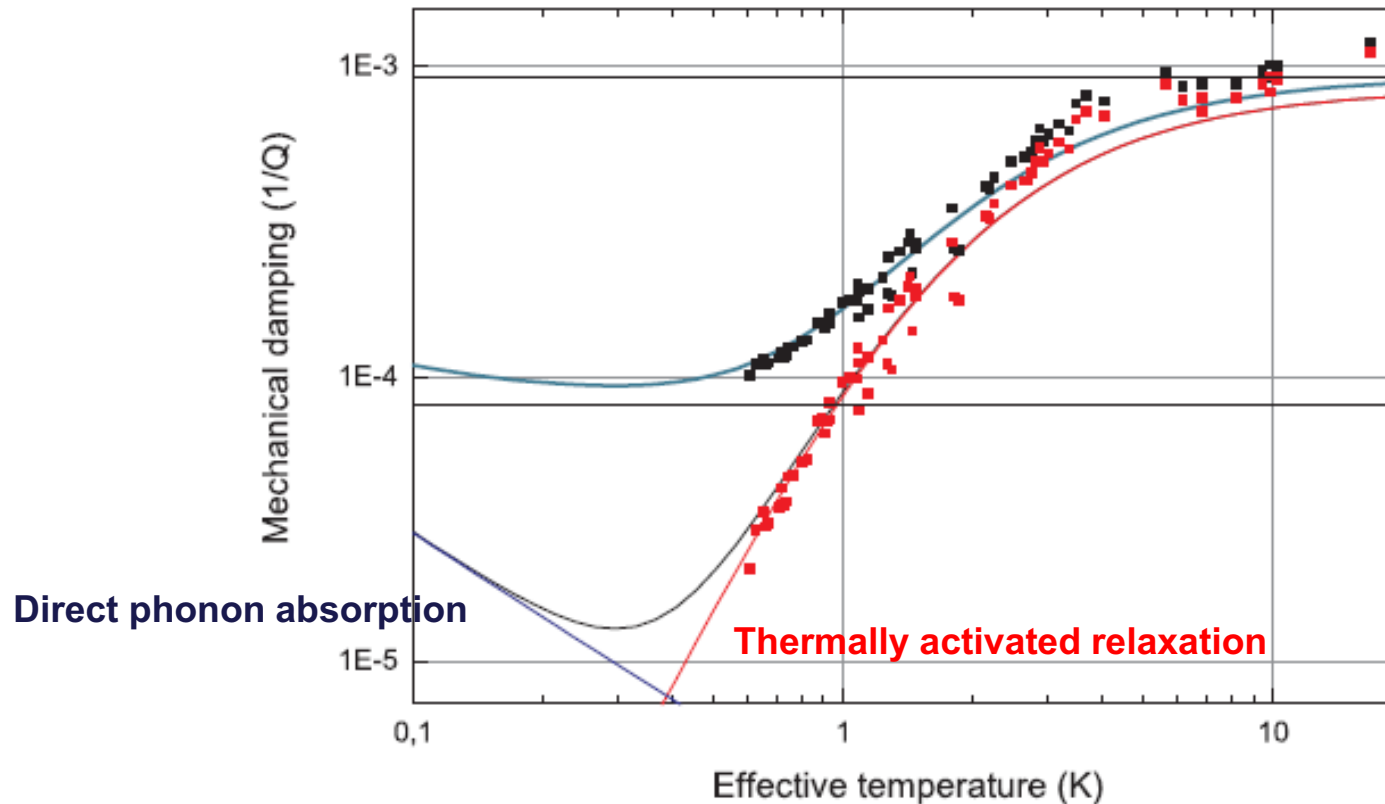
Tunnelling process



[1] Vacher, Courtens, Forêt, PRB 72 214205 (2005)

[2] Jäckle, Piché, Huncklinger, J. Non-Crys. Sol. 20 365 (1976)

# Low Temperatures: Two level fluctuators (TLS)



$$Q_m \propto \frac{\Omega_m}{T^3}$$

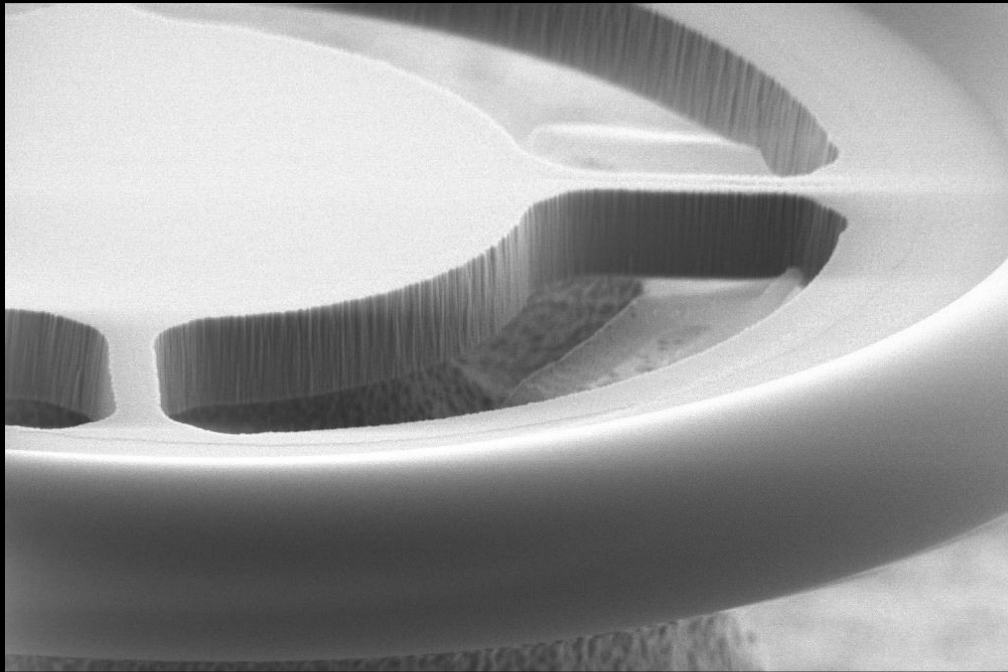
Thermally activated contribution

$$Q^{-1} = C \operatorname{Erf} \left( \frac{\sqrt{2}T}{\Delta_C} \right) \frac{1}{T} \int_0^\infty \left( \frac{V}{V_0} \right)^{-\xi} e^{-\frac{1}{2} \frac{V^2}{V_0^2}} \frac{\Omega \tau_0 e^{V/T}}{1 + \Omega^2 \tau_0^2 e^{2V/T}} dV \quad [3]$$

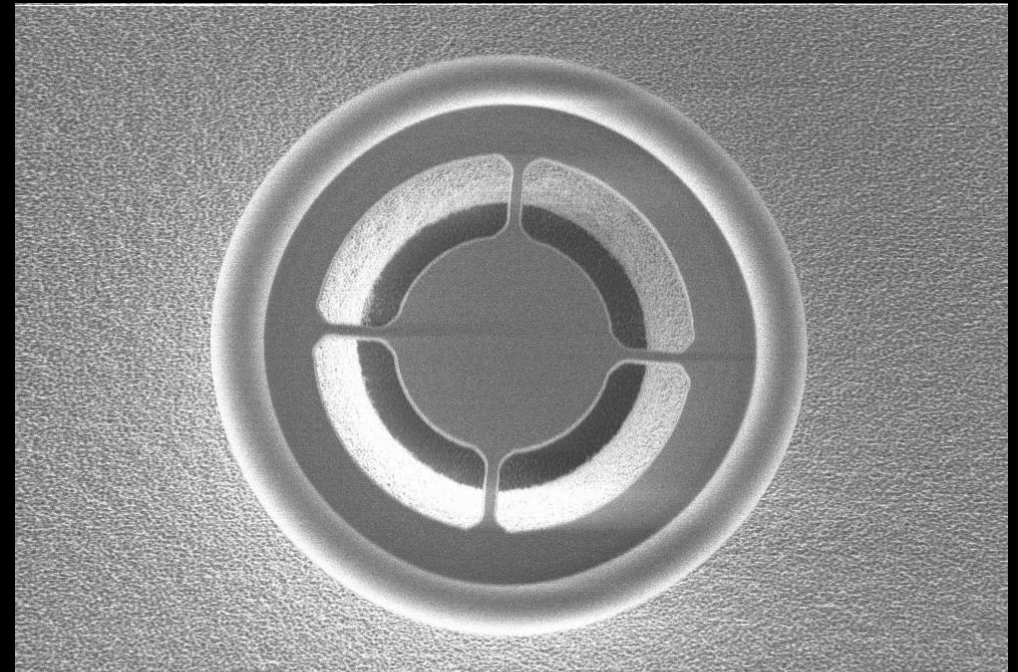
[1] Arcizet, Rivière, Schliesser, Anetsberger, Kippenberg, PRA (2009)

[2] U. Bartell et al., J. Phys. (Paris) Colloq. 43, C9 (1982)11

[3] Vacher, Courtens, Forêt, PRB 72, 214205 (2005)



Mag = 16.50 K X | 2µm | EHT = 3.00 kV | Signal A = SE2 | Date :6 Sep 2010 | EPFL-CMI  
WD = 5 mm | Stage at T = 22.2 ° | File Name = 2\_6.tif



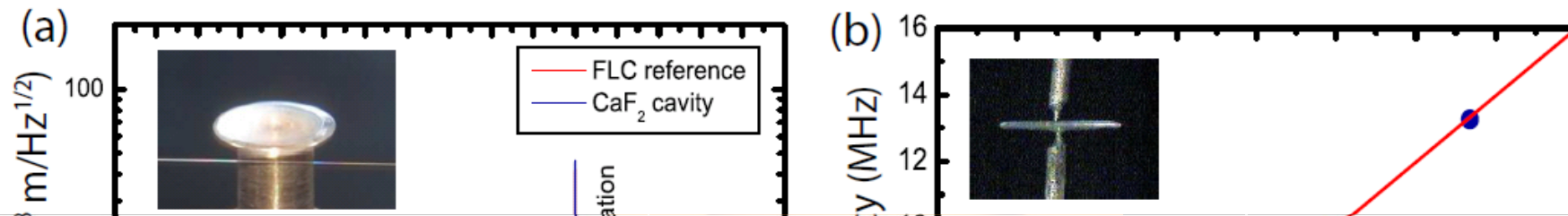
Mag = 4.18 K X | 10µm | EHT = 3.00 kV | Signal A = SE2 | Date :6 Sep 2010 | EPFL-CMI  
WD = 4 mm | Stage at T = 0.0 ° | File Name = 4\_4.tif

## Improved spoke resonators (using RIE processing)

- Reduced clamping losses
  - Increased optomechanical coupling
- (Gavartin, Verhagen et al)

# Optomechanical coupling in crystalline microresonators

High Q mechanical modes observed coupled via radiation pressure



КВАРЦЕВЫЙ РЕЗОНАТОР ЧАСТОТЫ КОЛЕБАНИЙ 1 МГц  
И ДОБРОТНОСТЬЮ  $4,2 \cdot 10^9$  ПРИ ТЕМПЕРАТУРЕ 2 К

А. Г. СМАГИН

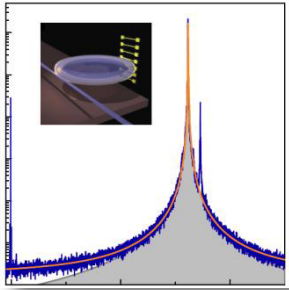
Описаны высокодобротные колебательные системы из бездислокационного кварца. Асимптотические методы обработки поверхности монокристалла и конструкция кристаллодержателя, в котором пьезоэлектрический кварц находится в свободном (не зажатом) состоянии, позволили повысить добротность резонатора до  $4,2 \cdot 10^9$  при температуре 2 К. Обсуждается вопрос о получении макроскопических колебательных систем с добротностью  $\sim 10^{13}$  путем понижения температуры до 0,3 К.

$$\kappa < 20\text{kHz} (Q > 10^{10}) \rightarrow \text{Im}(n) < 10^{-10}$$

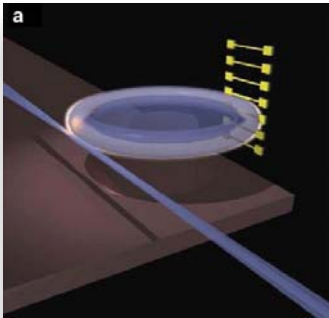
$$\Omega_m / \kappa > 100$$

Pioneering work on crystalline resonators: Maleki, Ilchenko, Matkso et al. JPL  
J. Hofer, A. Schliesser, TJ Kippenberg (arXiv:1003.5922, to appear Phys. Rev. A.)

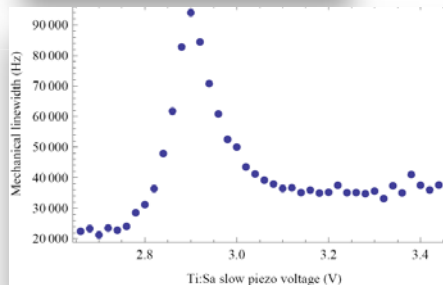
# Outline



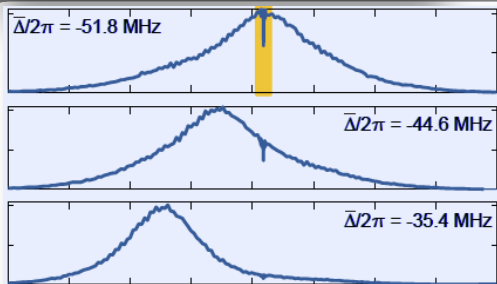
Measuring the vacuum optomechanical coupling strength



Measurement imprecision below that at The Standard Quantum Limit



Thermometry of an optomechanically cooled Microresonator in a Helium-3 cryostat



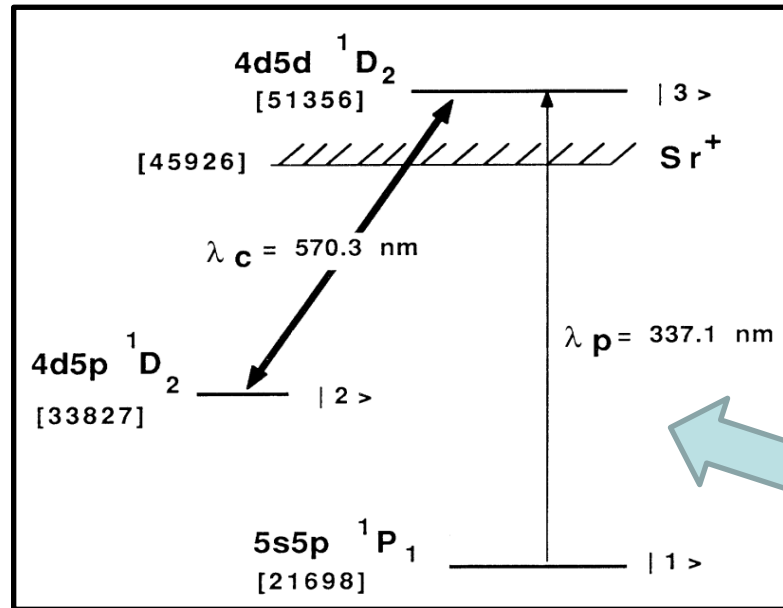
Optomechanically induced transparency

## Observation of Electromagnetically Induced Transparency

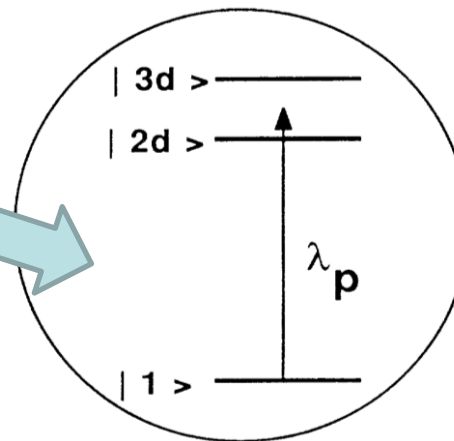
K.-J. Boller, A. Imamoglu, and S. E. Harris

Edward L. Ginzton Laboratory, Stanford University, Stanford, California 94305

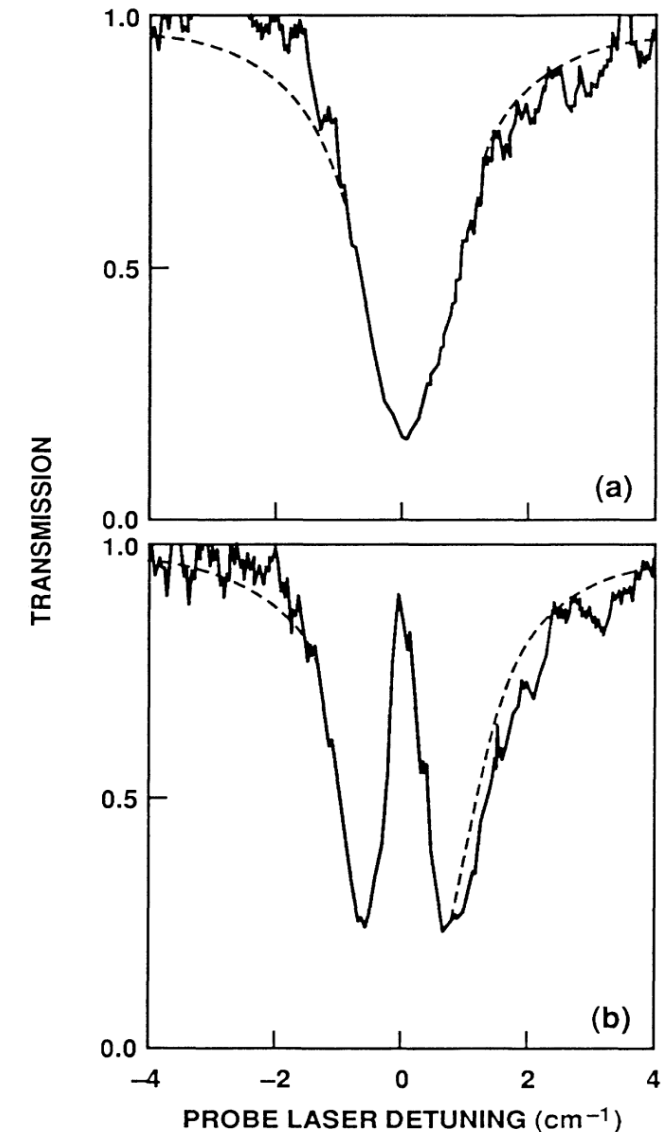
(Received 12 December 1990)



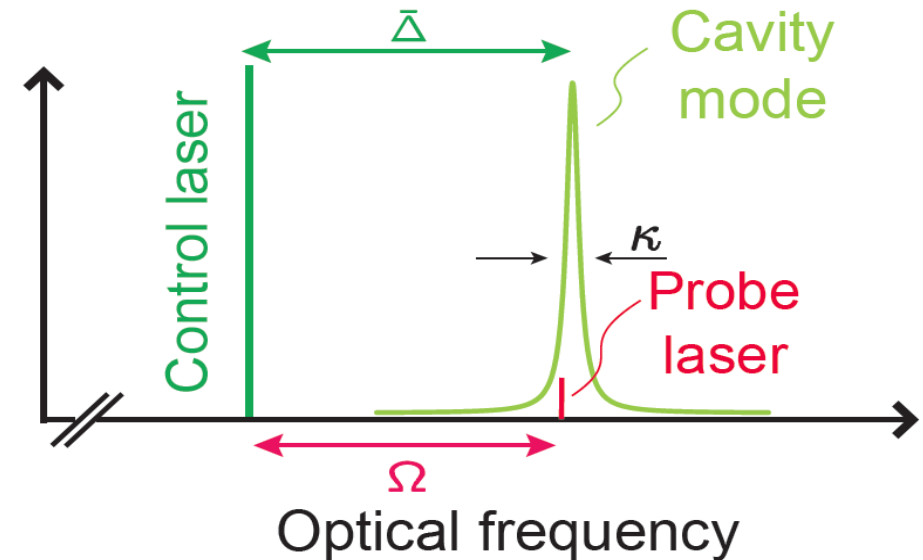
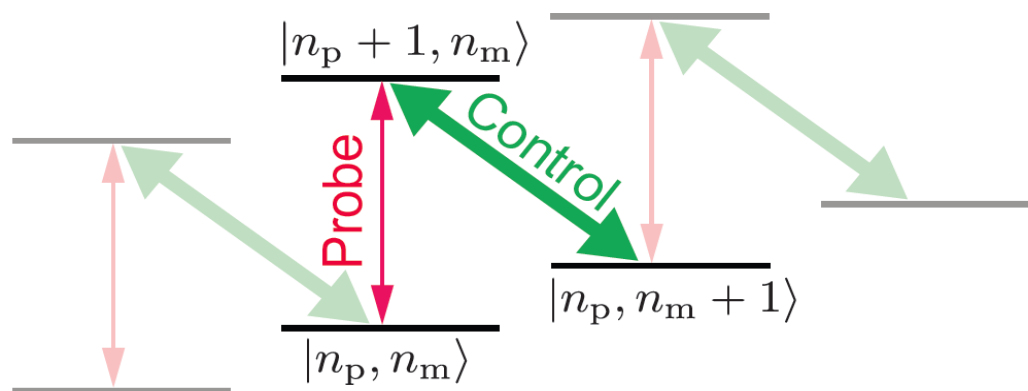
Level Scheme of Sr



Dressed States Picture



- Light storage (Phillips, *PRL*, 86 (2001)); Slow light (Hau, *Nature*, 97 (1999))



$$\begin{aligned} (-i(\bar{\Delta} + \Omega_m + \Delta') + \kappa/2) A^-(\Omega) &= -ig_0 \bar{a} X^-(\Omega) + \sqrt{\eta_c \kappa} s_p \\ 2m_{\text{eff}} \Omega_m (2\Delta' - i\Gamma_m) X^-(\Omega) &= -\hbar g_0 \bar{a} A^-(\Omega). \end{aligned}$$

EIT

projection operator  $\sigma_{13}$  (coherence  $\rho_{13}$ )  
 projection operator  $\sigma_{12}$  (coherence  $\rho_{12}$ )  
 energy difference between ground states  $\hbar\omega_{21}$   
 Rabi frequency  $\mu_{23}\mathcal{E}_c/\hbar$

OMIT

intracavity field amplitude  $A^-$   
 mechanical displacement amplitude  $X$   
 phonon energy  $\hbar\Omega_m$   
 optomechanical coupling rate  $2g_0\bar{a}x_{z\text{pf}}$

Schliesser, PhD thesis (Nov. 2009) <http://edoc.ub.uni-muenchen.de/10940/>

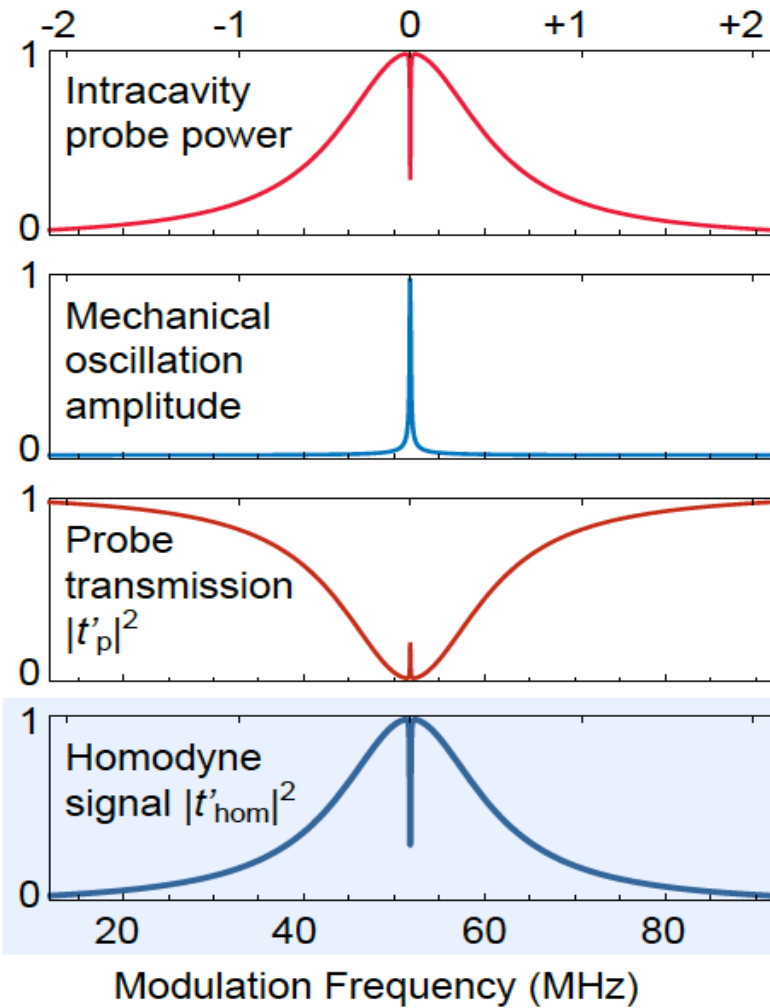
Agarwal Huang Phys. Rev. A 81, 041803(R) (2010)

S. Weis, R. Riviere, S. Deleglise, A. Schliesser & T.J. Kippenberg arXiv:1007.0565

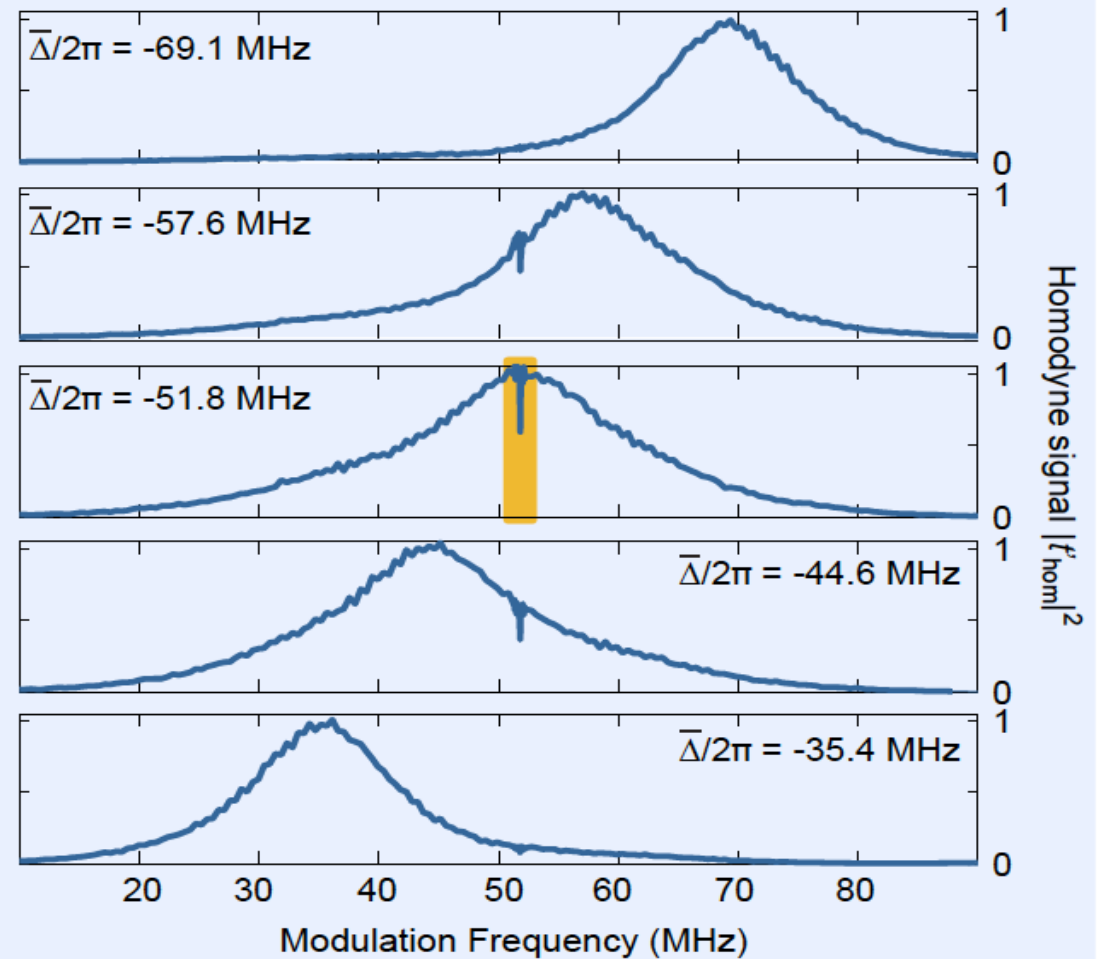


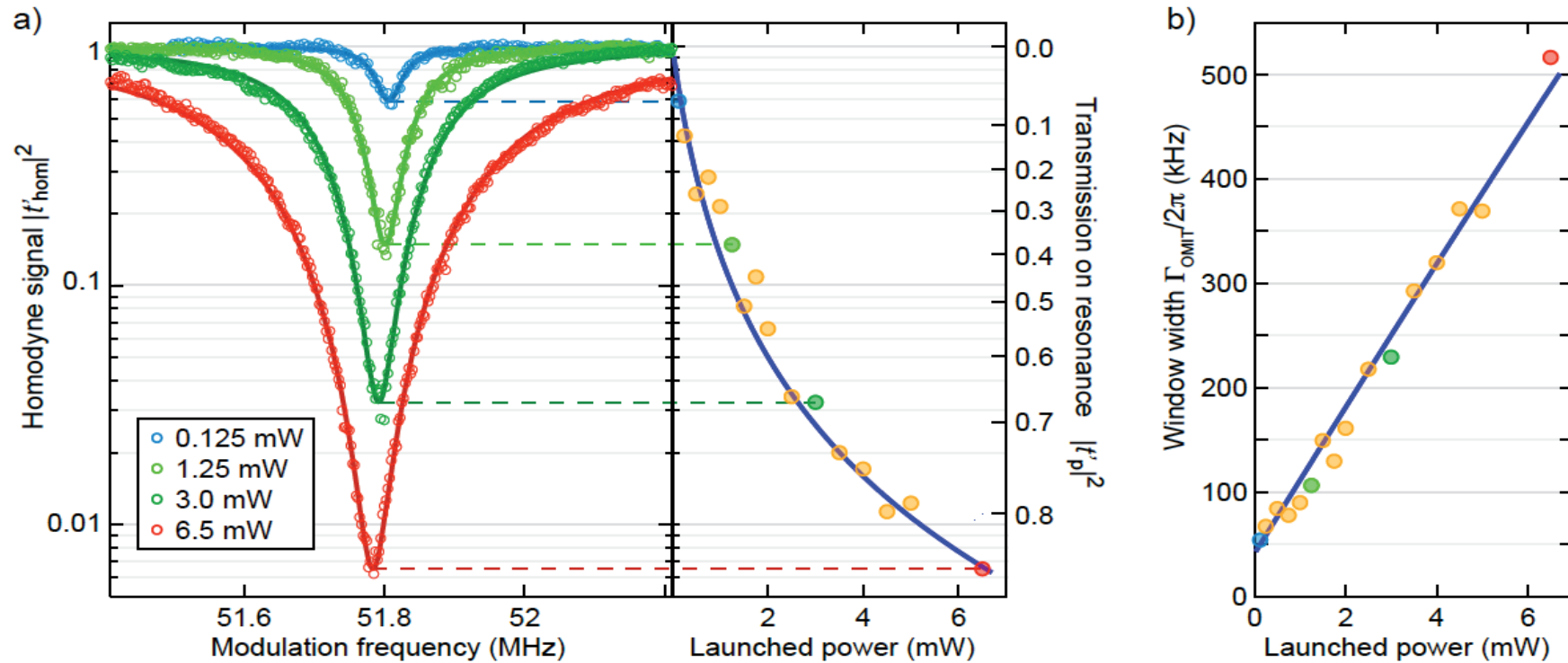
# Optomechanical EIT

a) Normalized probe detuning  $\Delta'/\kappa$



b)

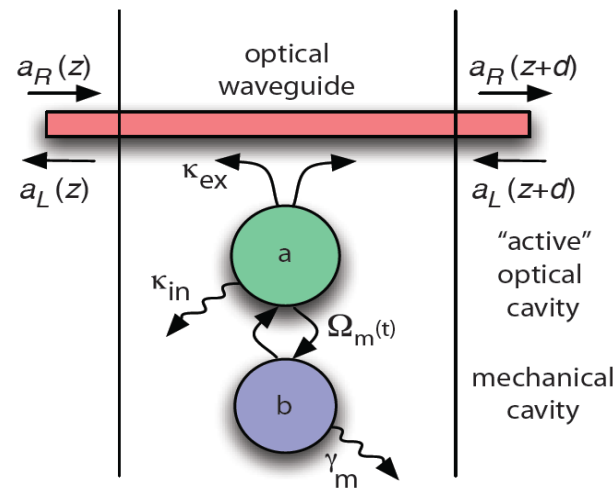




- probe transparency window tunable by control as expected
- window width and depth can be continuously varied
- “transparency” up to 81% achieved

## Stopping and storing light in mechanical systems

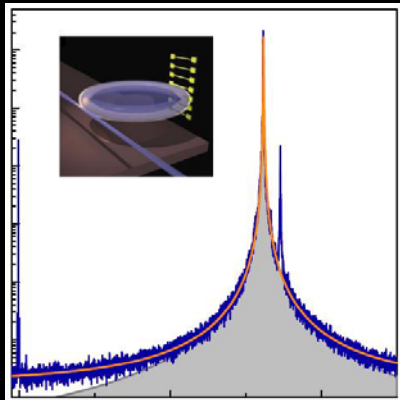
- Requires an array of optomechanical systems
- Tuning of probe transmission while light propagates through the array



Chang, Safavi-Naeini, Hafezi, Painter, arXiv:1006.3829

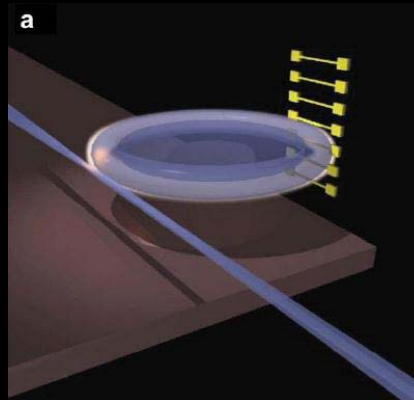
# Summary

## Vacuum optomechanical coupling calibration



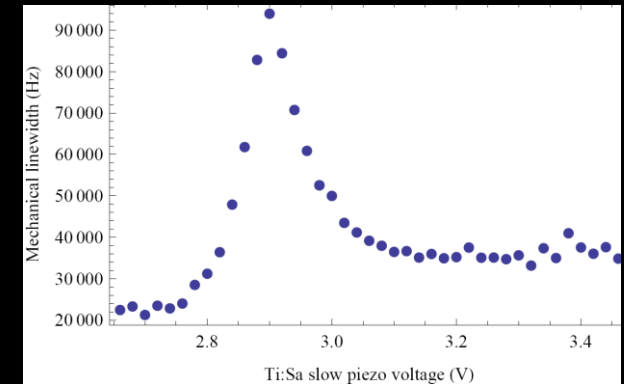
Anetsberger, ML Gorodetsky,  
Schliesser, TJK (to appear: arXiv)

## Imprecision below that at the SQL for nanomechanical beams



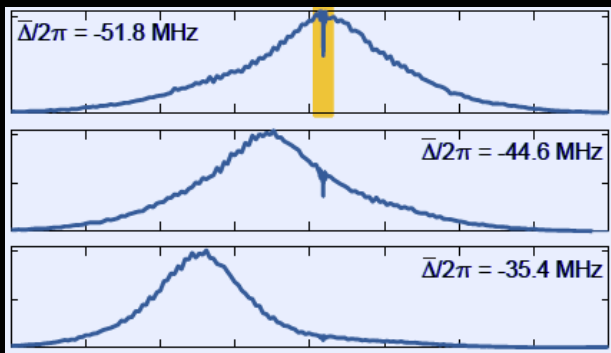
Anetsberger et al. arXiv (2009)

## Optomechanical cooling in a He3 cryostat



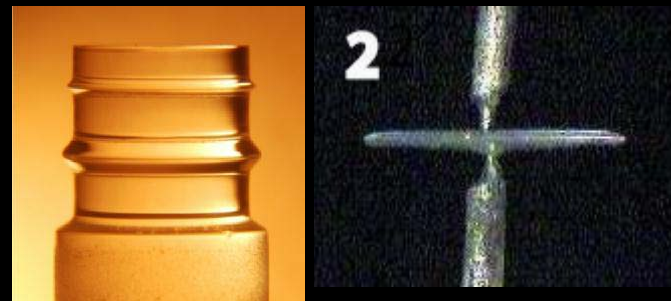
Unpublished (Weis, Riviere,  
Deleglise)

## Observation of Optomechanical EIT



Weis, Riviere, Deleglise, Schliesser, TJK (arXiv)

## Observation of Optomechanical EIT



Hofer et al. arXiv (to appear:  
*Phys. Rev. A.*)

# Acknowledgements



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