



**The Abdus Salam
International Centre for Theoretical Physics**



2168-4

Joint ICTP-IAEA Workshop on Dense Magnetized Plasma and Plasma Diagnostics

15 - 26 November 2010

Summary of results from POSEJDON Plasma-Focus

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Germany*

Diagnostics and Scaling of Fusion- Produced Neutrons in PF Experiments

Hellmut Schmidt

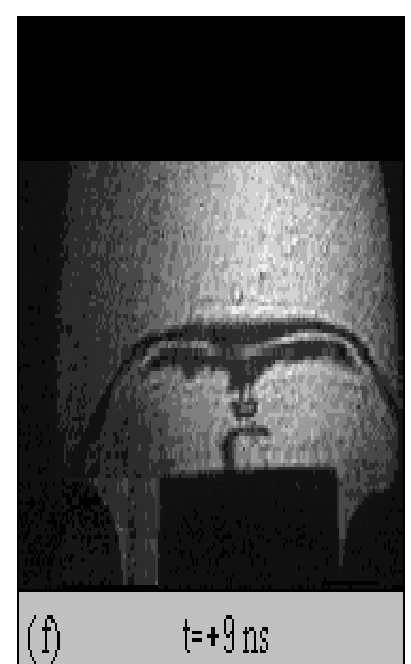
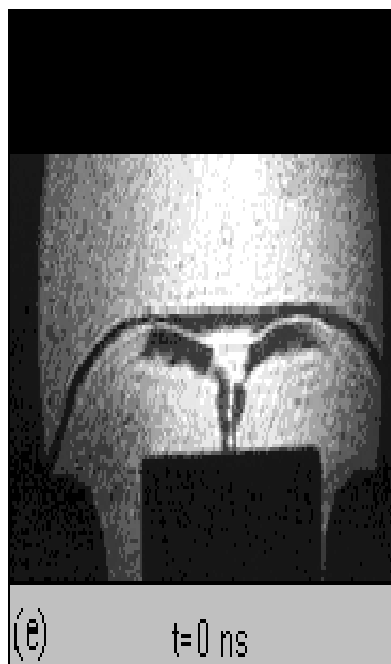
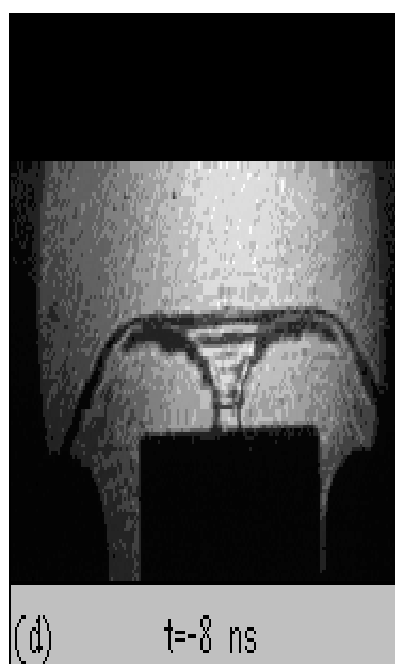
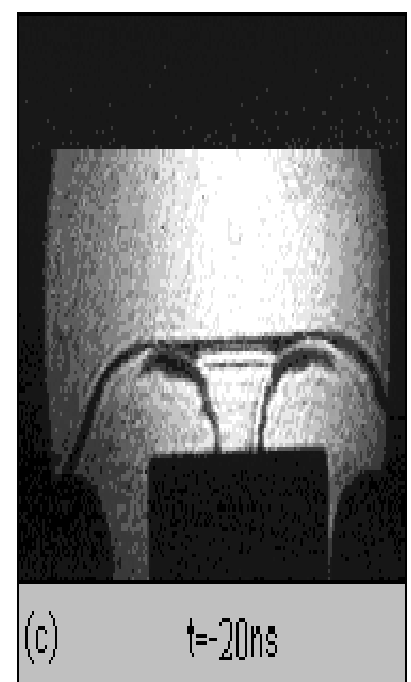
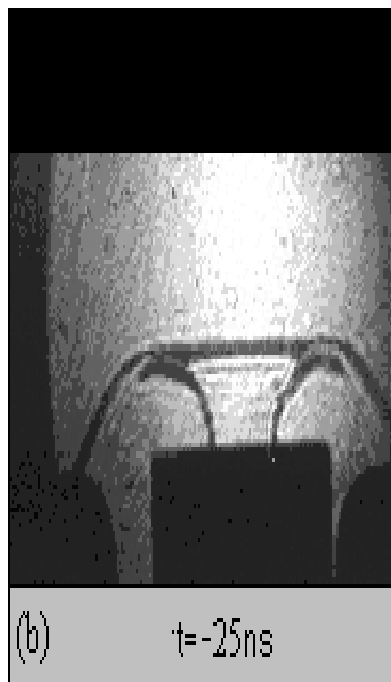
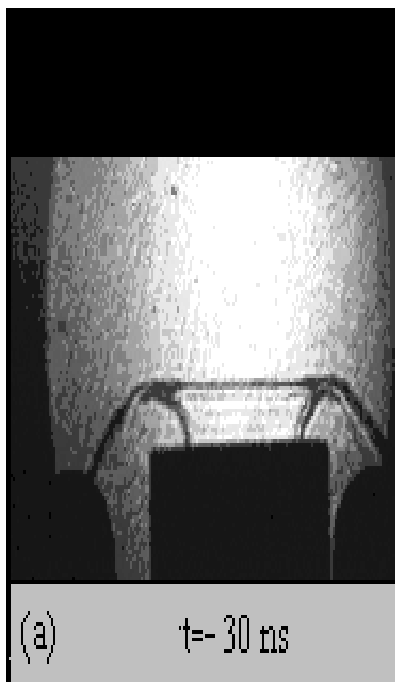
International Centre for
Dense Magnetized Plasmas, Warsaw
and
University of Stuttgart, Germany

Outline

- Introduction
- Neutron Detection
 - Activation
 - Scintillation
 - TLDs and Bubble detectors
- Neutron Diagnostics
 - Fusion Reaction Models
 - Gyration Particle Model
- Scaling of Neutron Yield
 - Beam Target and Thermal Production of Neutrons

Introduction

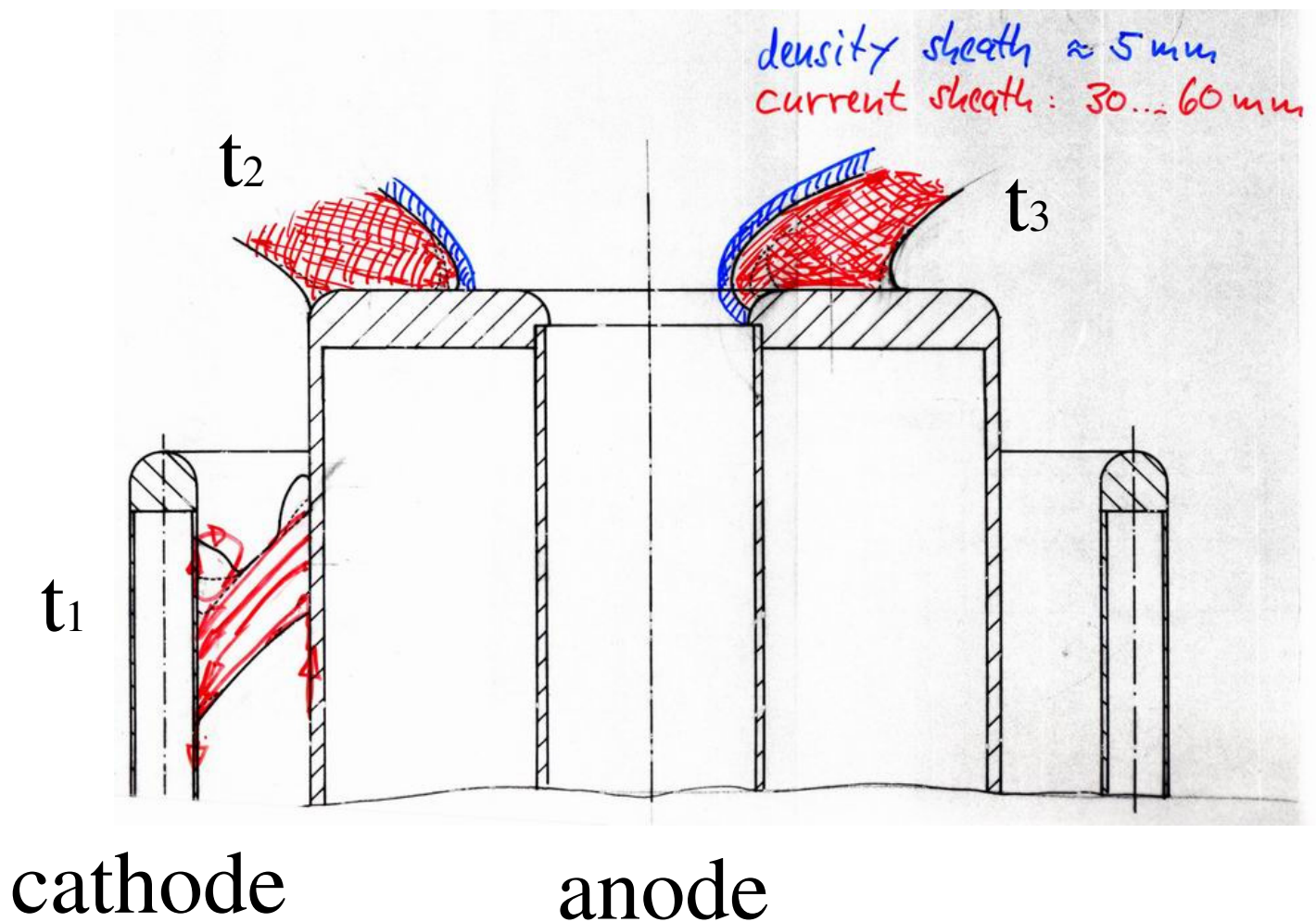
Radial Compression (Pinch) Phase of the Plasma Focus



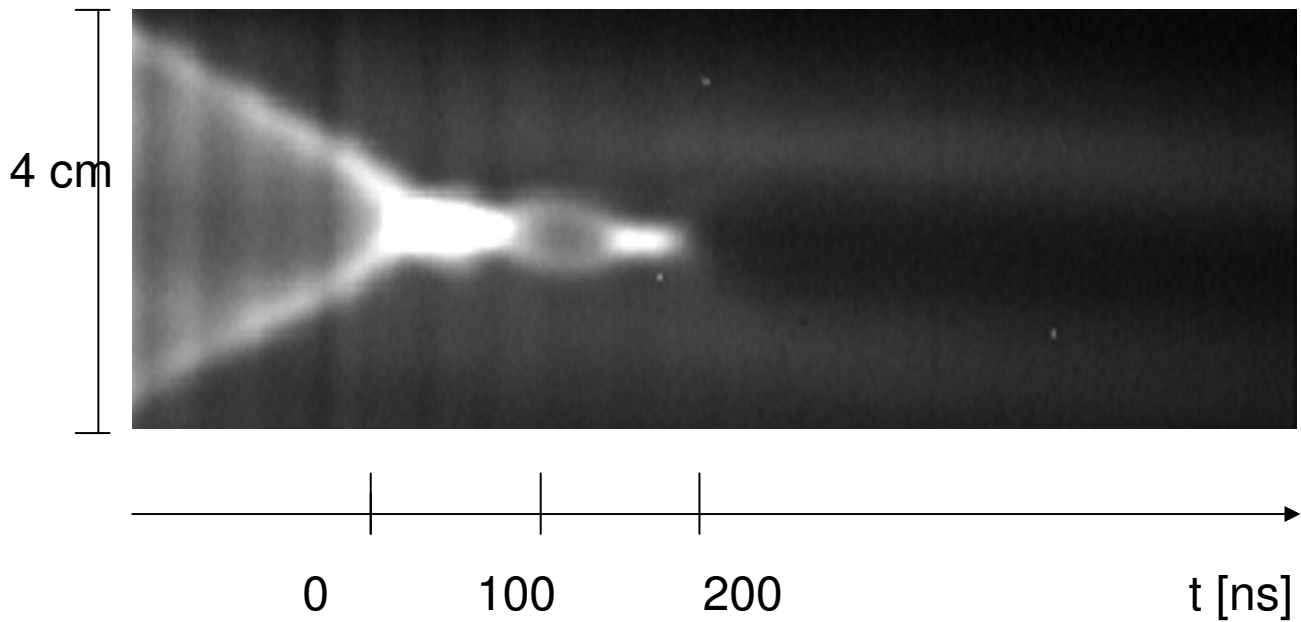


Density sheath of POSEIDON PF Before maximal compression

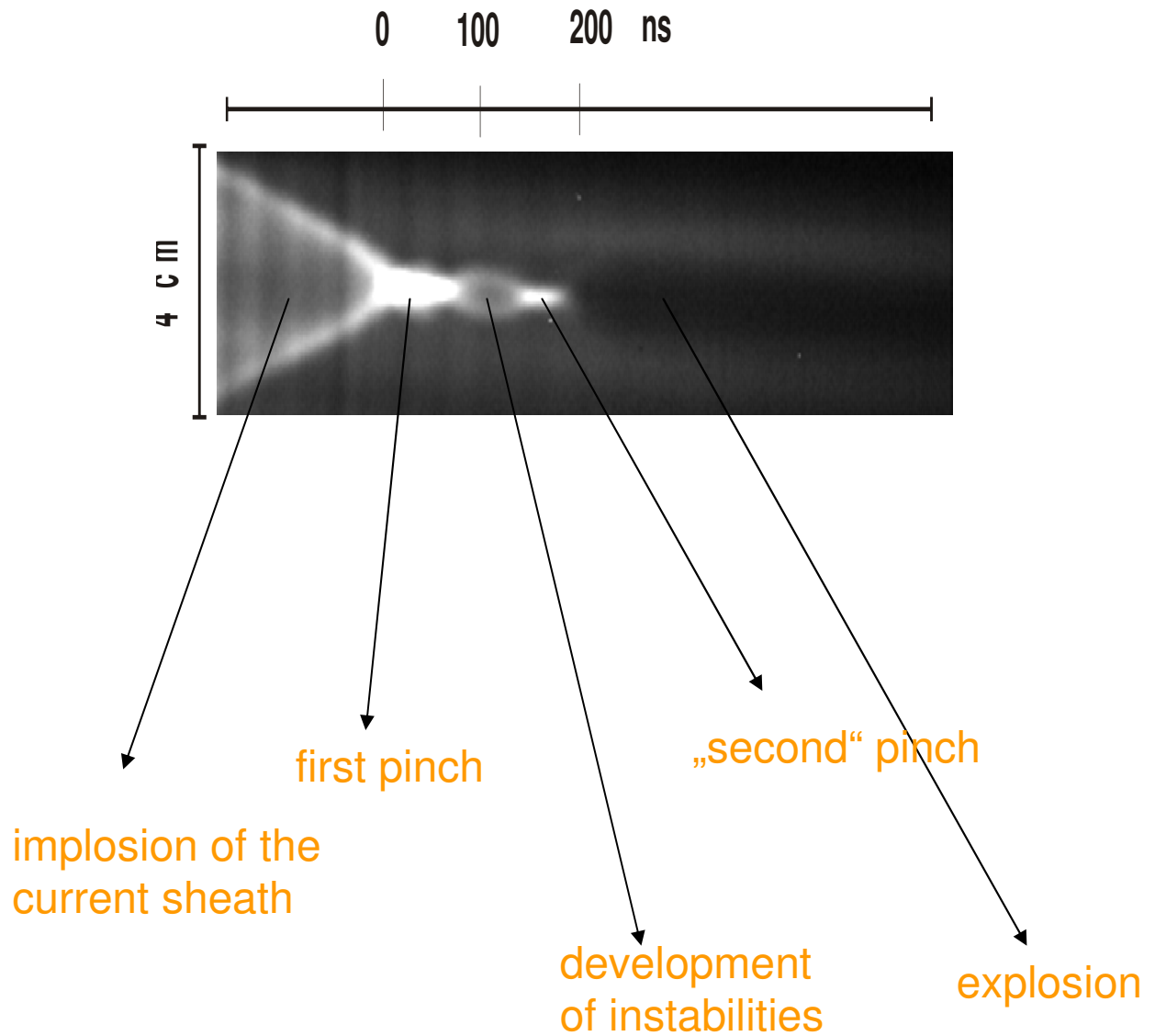
Current and Density Sheaths In a Mather type Plasma Focus



Radial Streak Picture, shot 5055 of PF1000



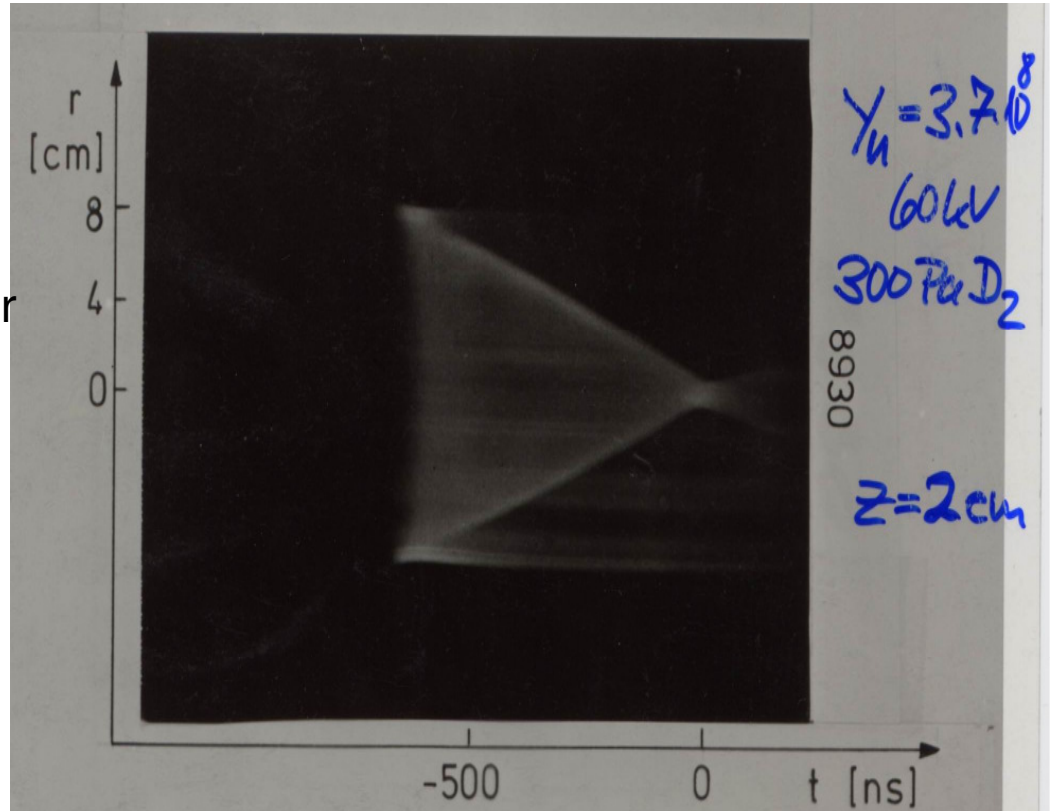
Radial Streak Picture of PF1000 (shot 5055)



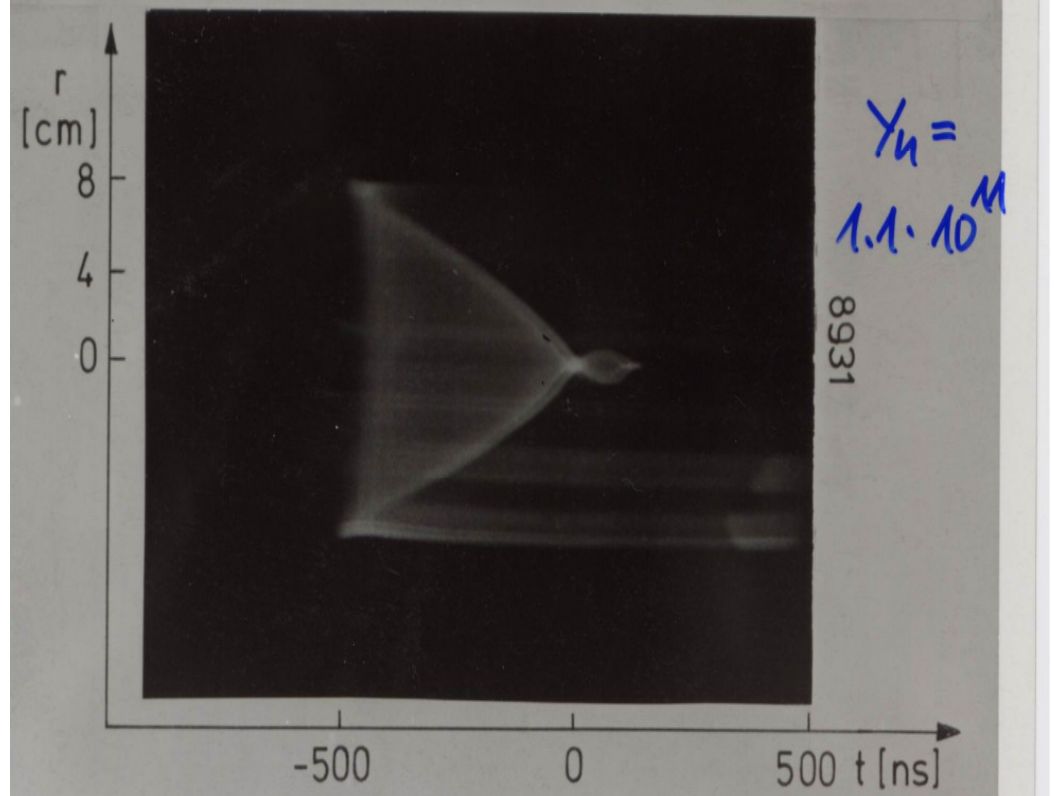
Typical Behaviour

Radial Streak Pictures of POSEIDON (280 kJ)

Bad Behaviour

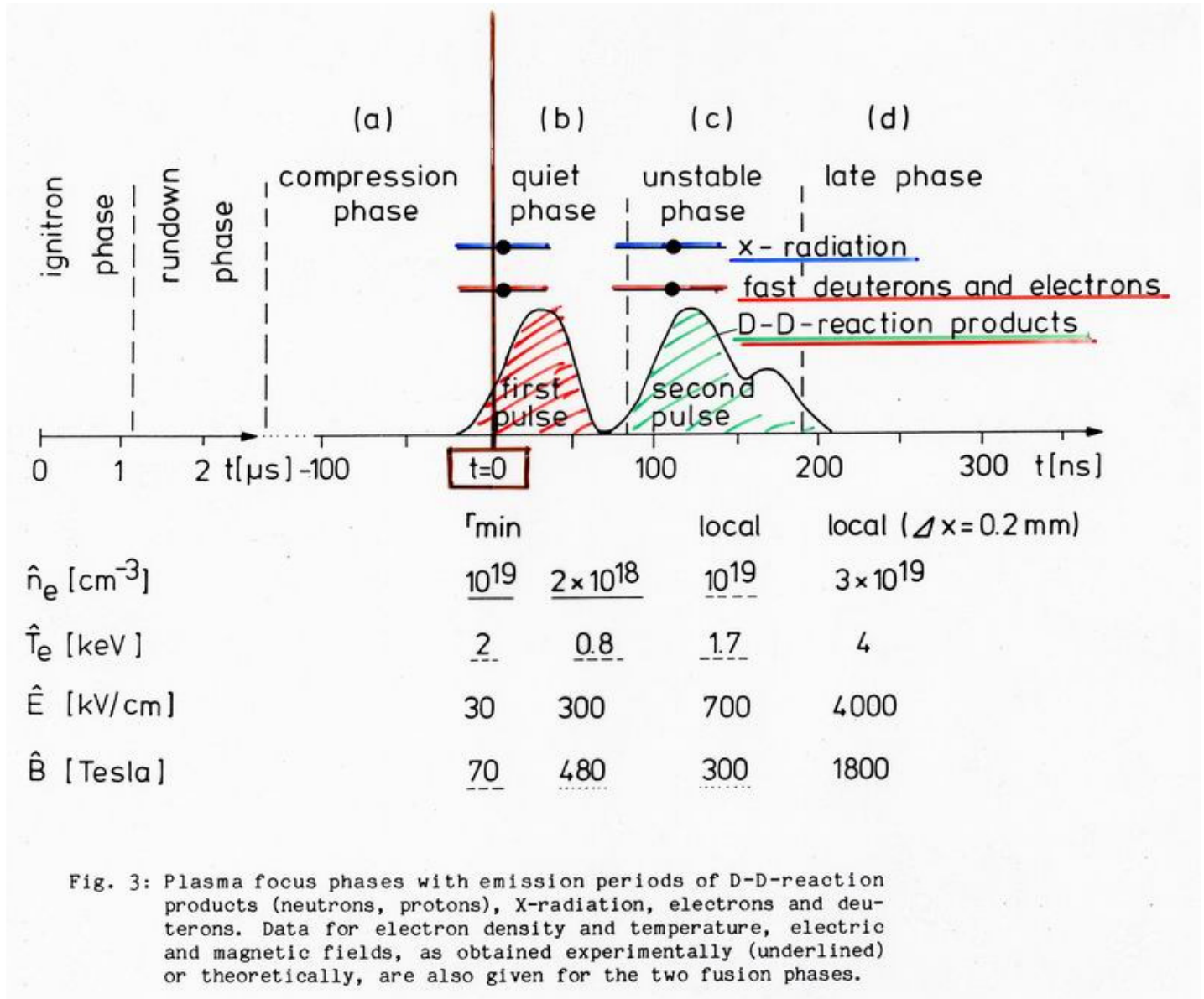


Typical Behaviour

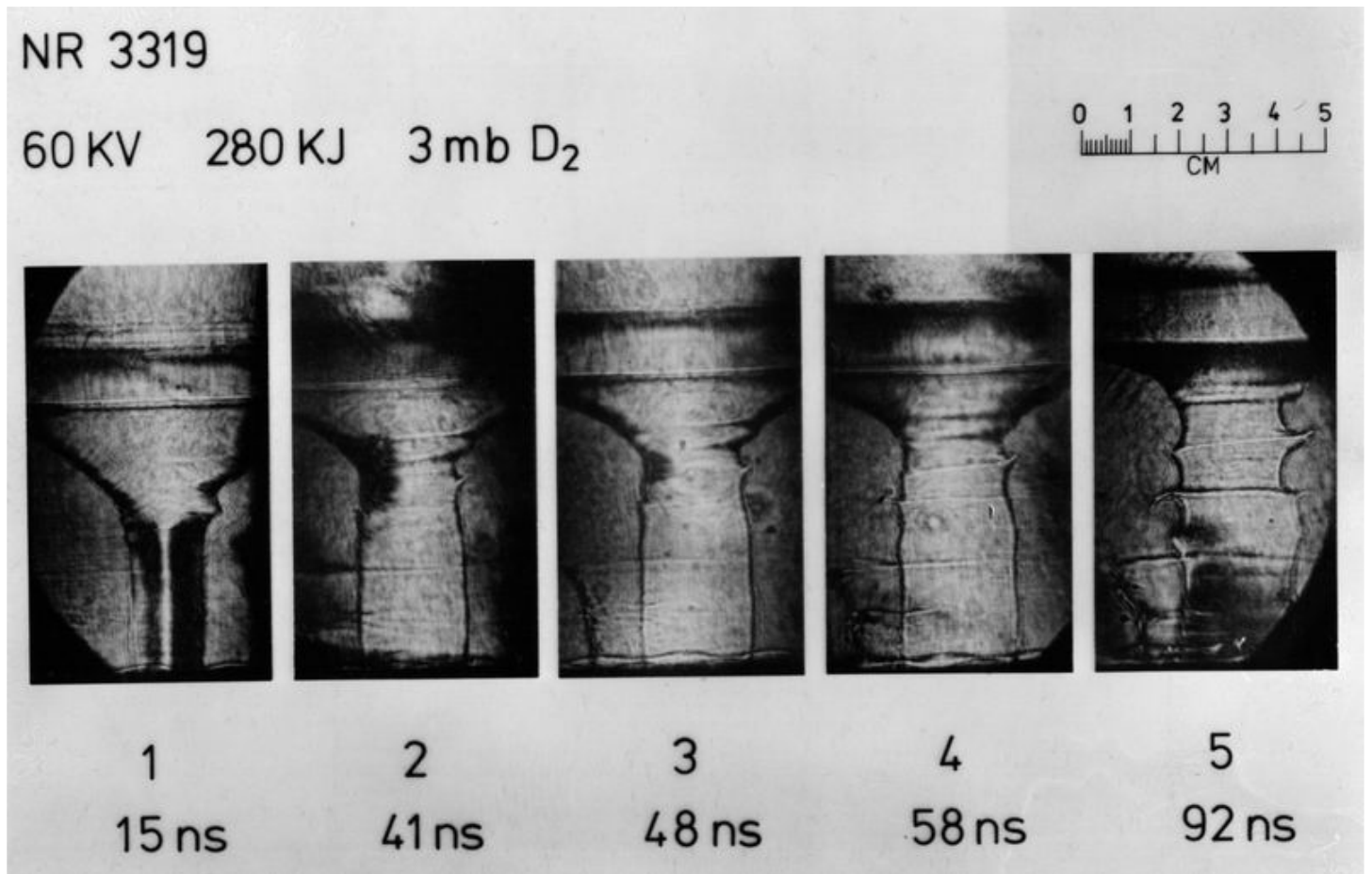


Trieste
November 2010

Plasma Focus Phases

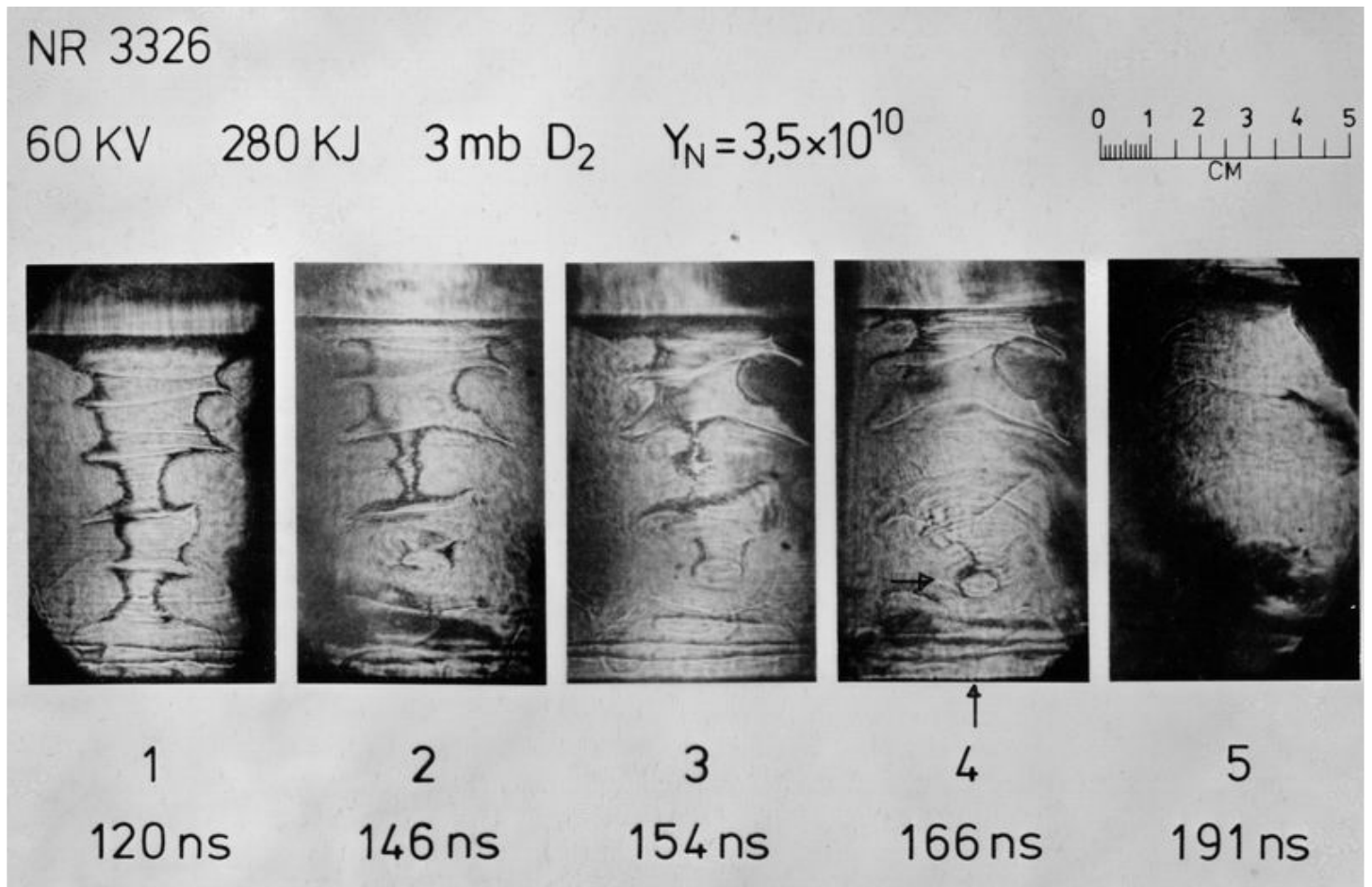


5 Schlieren pictures

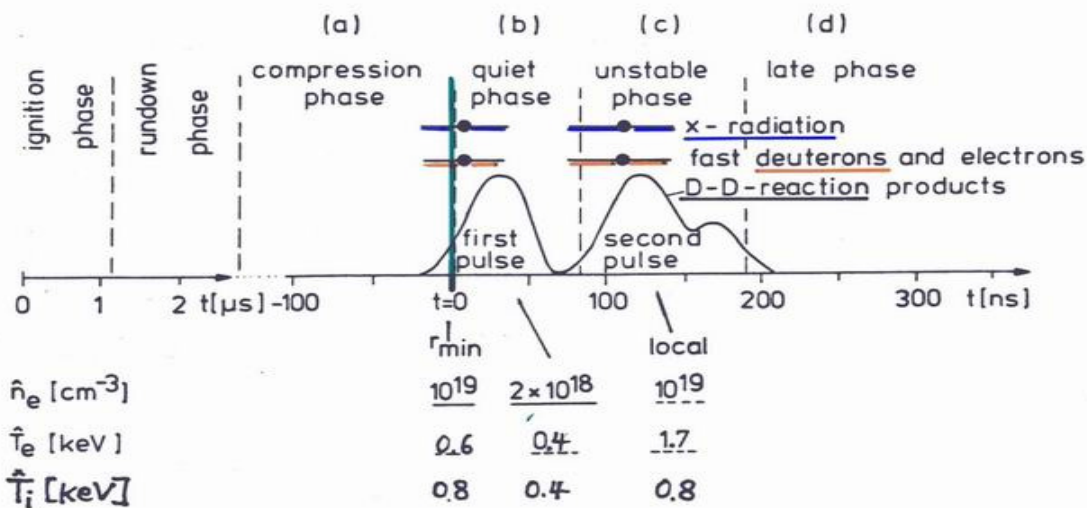


First neutron pulse

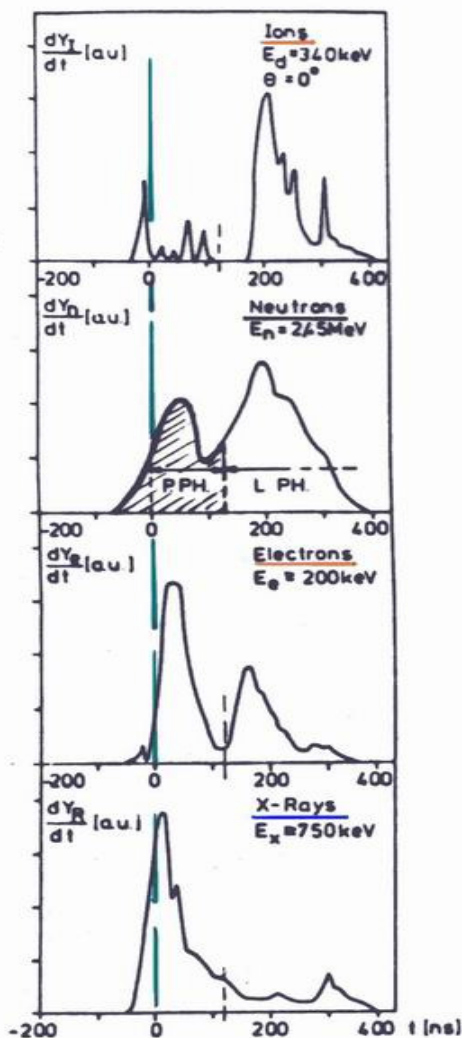
5 Schlieren pictures



Second neutron pulse



Examples of emissions



deuterons
(340 keV)

$\uparrow + z$

neutrons

electrons
(200 keV)

$\downarrow - z$

hard X-rays
(750 keV)

Neutron Detection

Neutron Activation

TABLE I. Comparison of various elements employed for pulsed neutron detection.

Element	Half-life $\tau_{1/2}$ (s)	Reaction	Neutron energy	Operation mode	Particle detected
Arsenic	0.017	$^{75}\text{As}(n, n') ^{75\text{m}}\text{As}$	Fast	Repetitive	γ
Beryllium	0.807	$^9\text{Be}(n, \alpha) ^6\text{He}$	Fast	Repetitive	β
Boron tri fluoride	...	$^9\text{B}(n, \alpha) ^7\text{Li}$	Thermal	Repetitive	α
Helium-3	...	$^6\text{He}(n, p) ^3\text{H}$	Thermal	Repetitive	p
Indium	14.1	$^{115}\text{In}(n, \gamma) ^{116}\text{In}$	Thermal	Single	β
Rhodium	42.3	$^{103}\text{Rh}(n, \gamma) ^{104}\text{Rh}$	Thermal	Single	β
Silver	24.6	$^{109}\text{Ag}(n, \gamma) ^{110}\text{Ag}$	Thermal	Single	β
	142	$^{107}\text{Ag}(n, \gamma) ^{118}\text{Ag}$			

Be-detector

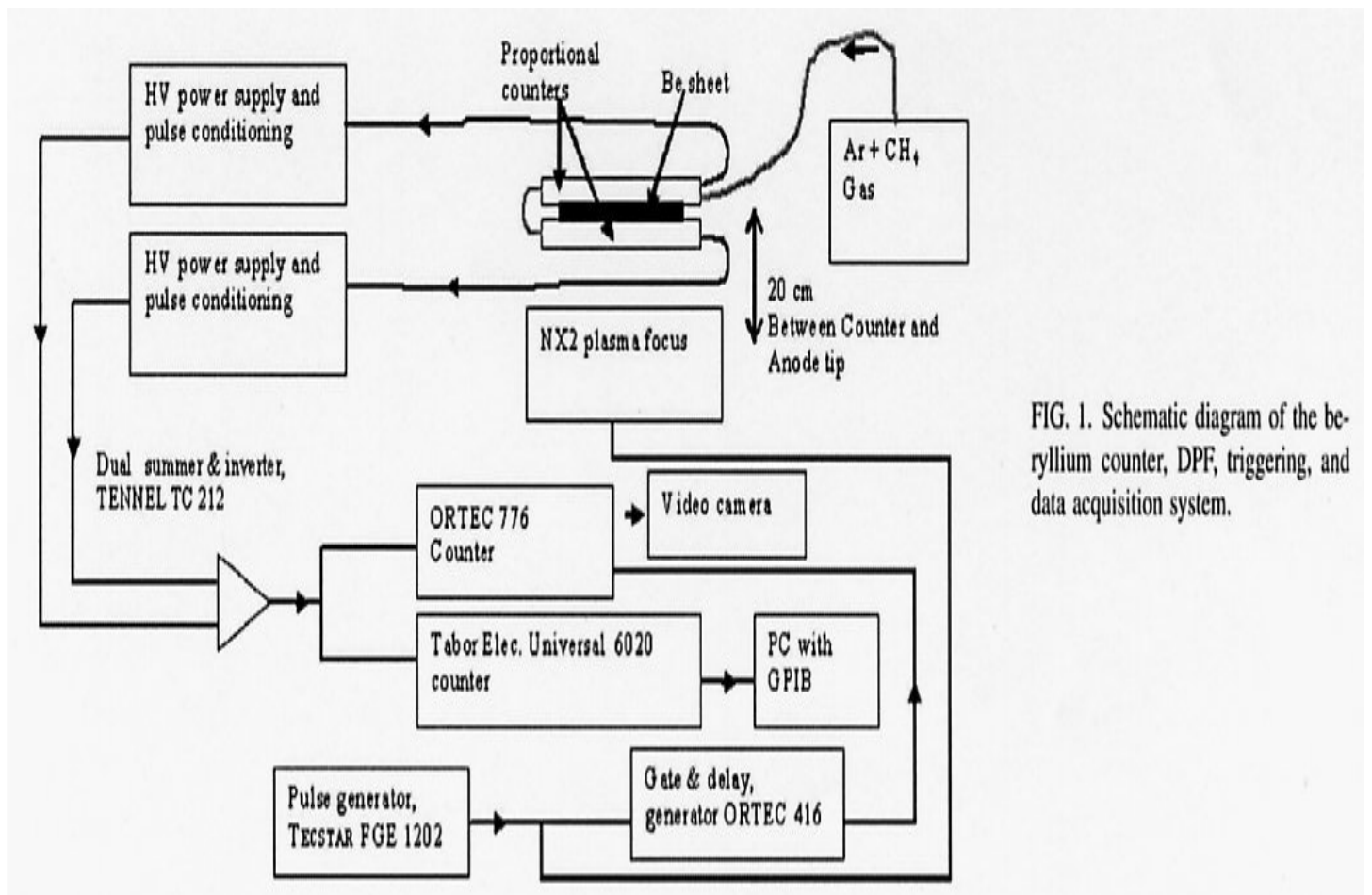
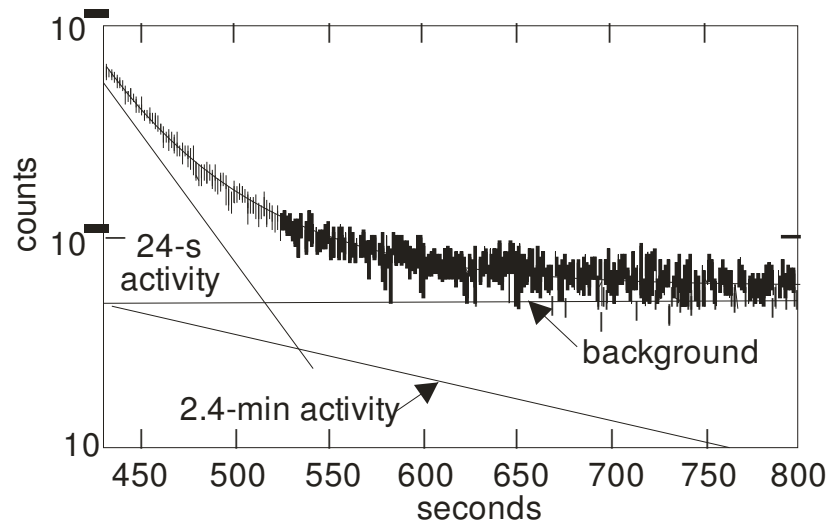
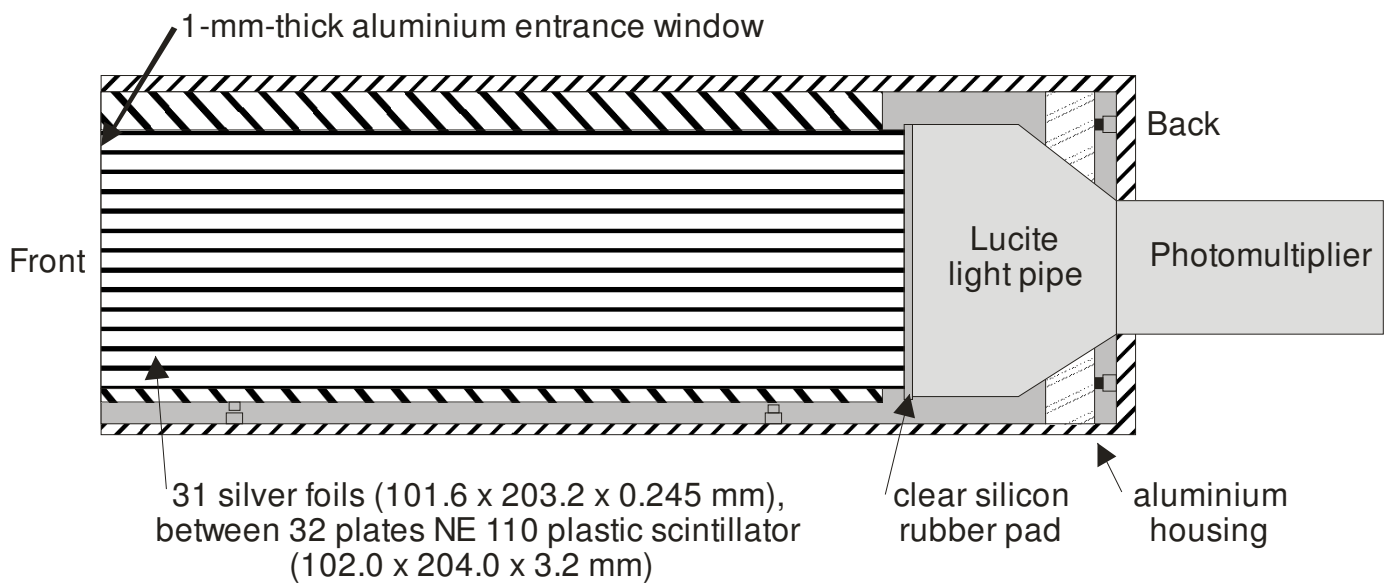


FIG. 1. Schematic diagram of the beryllium counter, DPF, triggering, and data acquisition system.

Silver activation counter

Example of a specific design

When used with fast neutrons, the counter is normally placed within a polyethylene moderator



Activity after insertion in a constant neutron flux

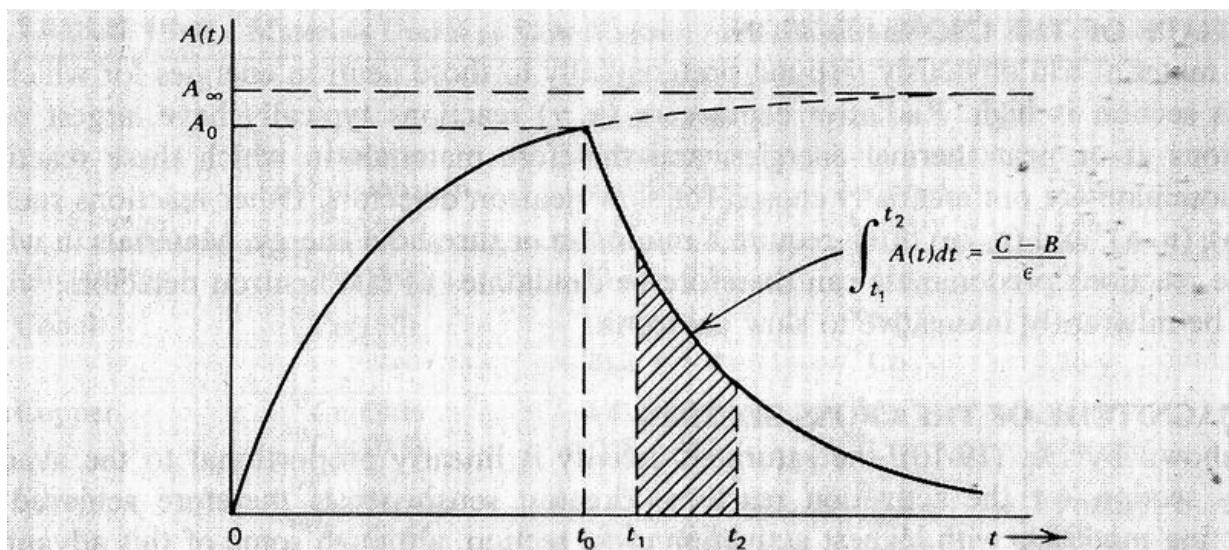
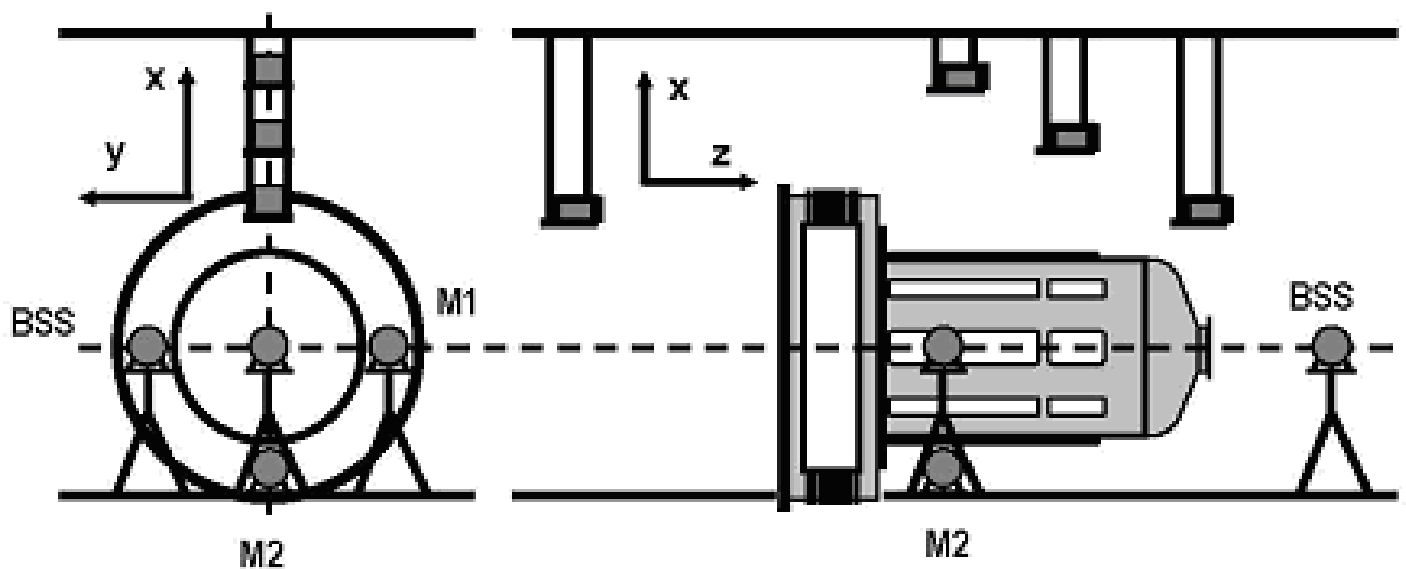


Figure 19-15 The activity of an activator detector after insertion in a constant neutron flux at time = 0 and removal at time = t_0 . The measured number of counts is proportional to the area under the decay curve between t_1 and t_2 .

Bonner spheres with TLDs and silver activation detectors PF1000 in Warsaw



Scintillation Detectors

- Scintillators connected to PM-Tubes
- Time resolution in the ns range

Bubble detectors

They integrate the neutron yield

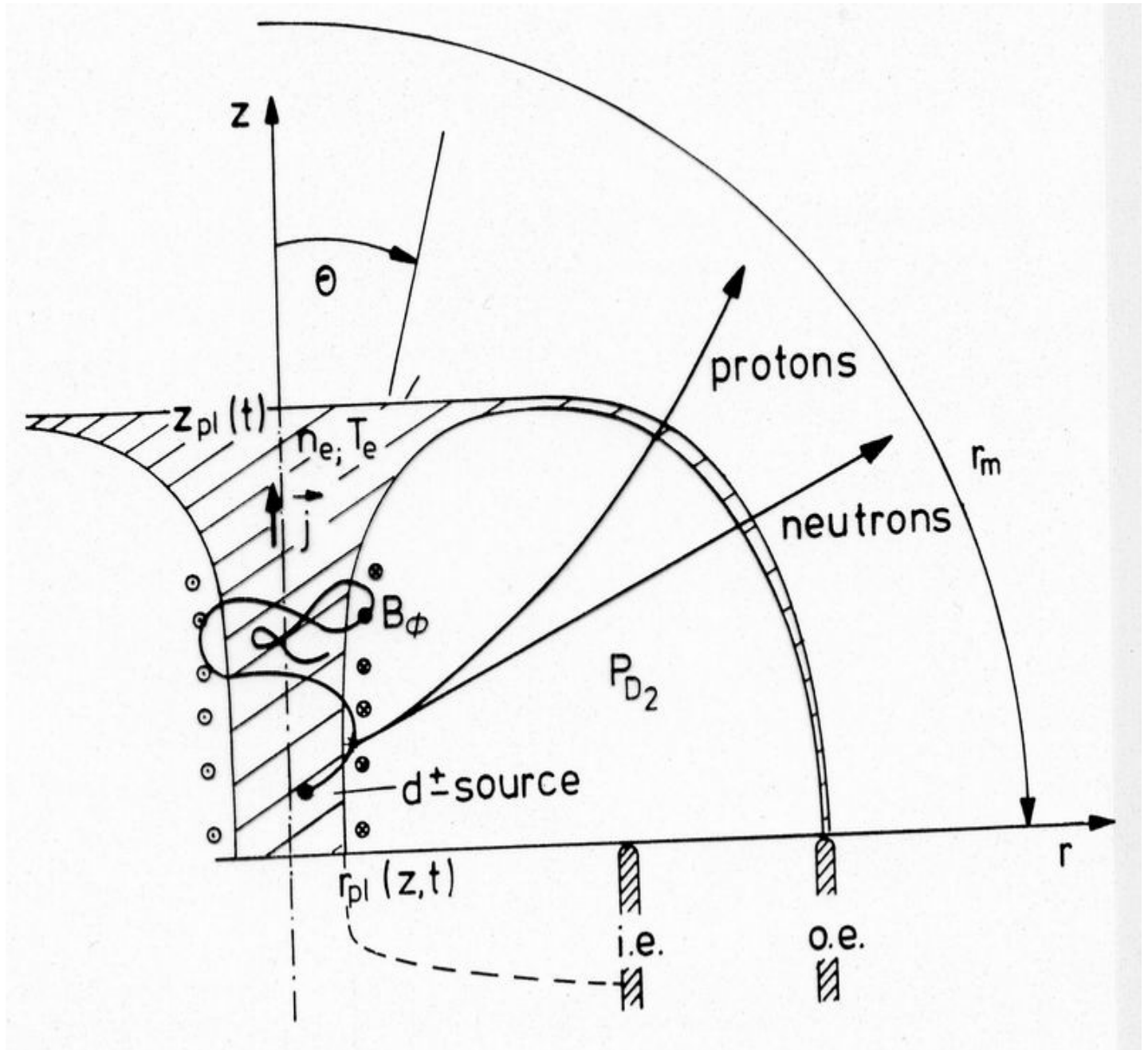
Bubbles have to be counted

Calibration necessary

Neutron Diagnostics

Beam Target Processes and their verification

Ion Trajectories in the Azimuthal Magnetic Field of the Pinch Current



Beam-Target neutron production? ?

statement to be verified:

deuterons of mean energy (50-100 keV)
contribute mainly and essentially to fusion yield

the following processes take place:

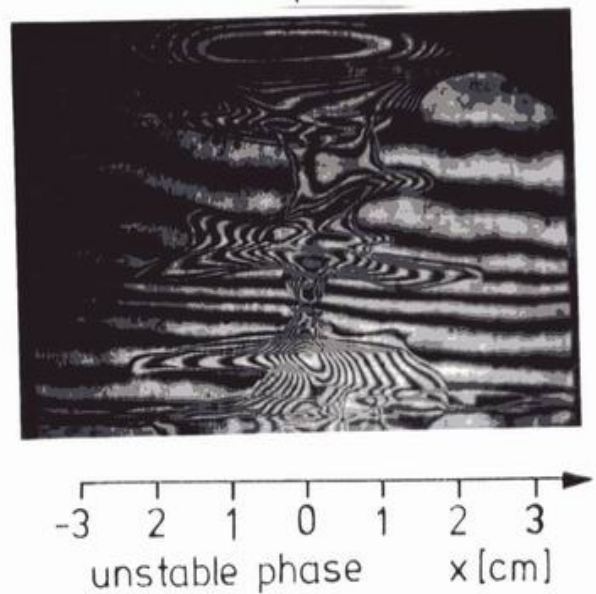
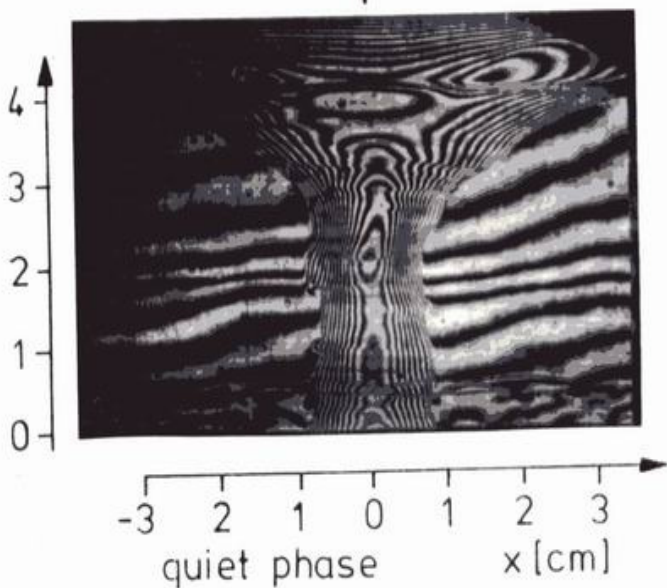
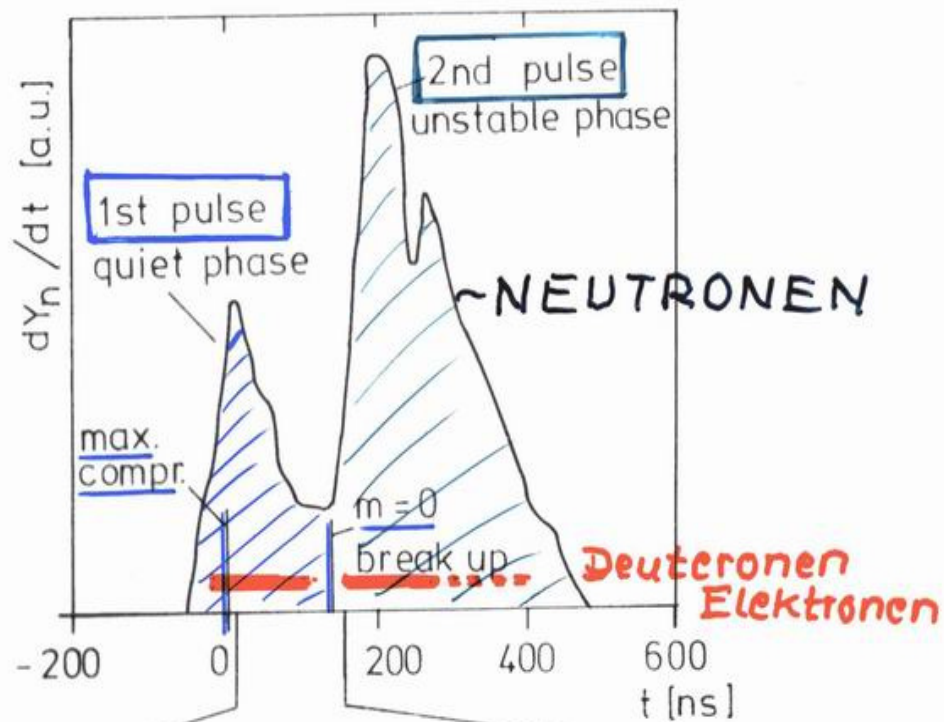
- acceleration of deuterons
by high transient electromagnetic fields
- absorption of those deuterons
in the pinch
- atomic processes in the surrounding
gas/plasma
 - charge exchange
 - scattering

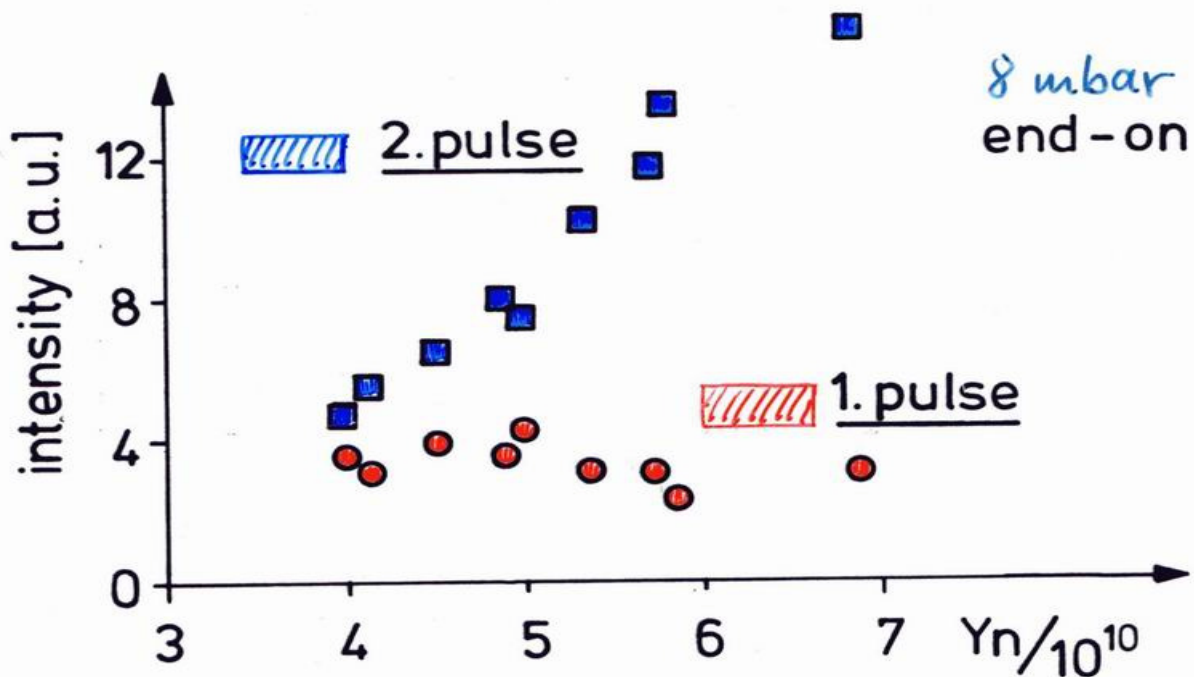
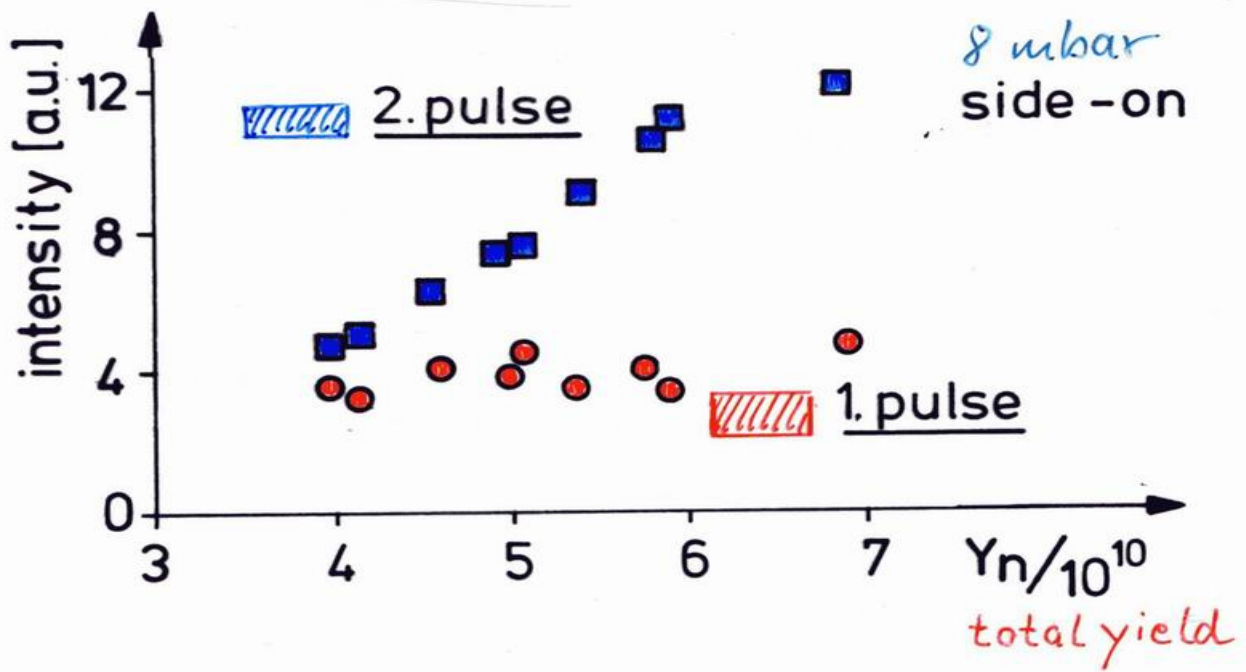
⇒ no direct observation of deuteron
distribution (which is responsible for
fusion reactions)

- indirect methods
- neutron spectroscopy
(time-of-flight, nuclear emulsions)
 - neutron flux anisotropy
 - neutron source location and
intensity distribution

(integral and/or time-resolved
measurements)

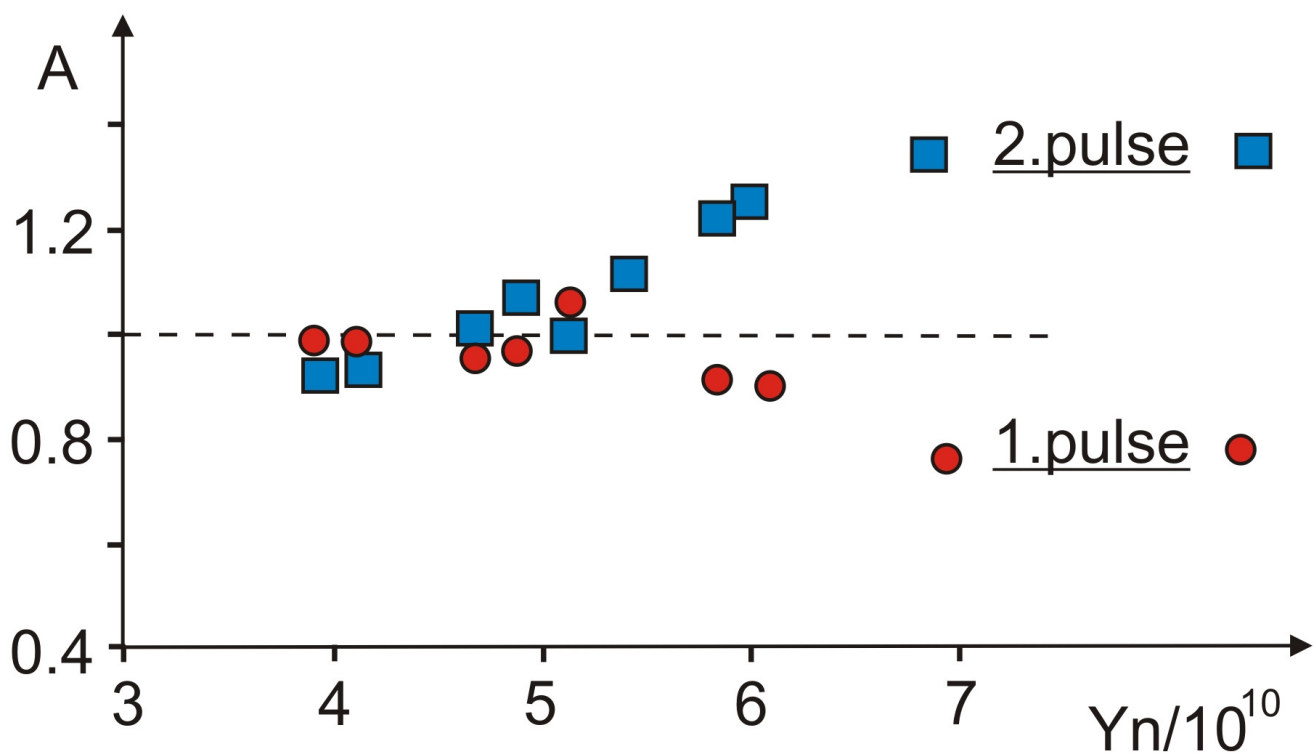
Two Neutron emitting Phases Correlated to Plasma Dynamics





neutron yield of the two phases

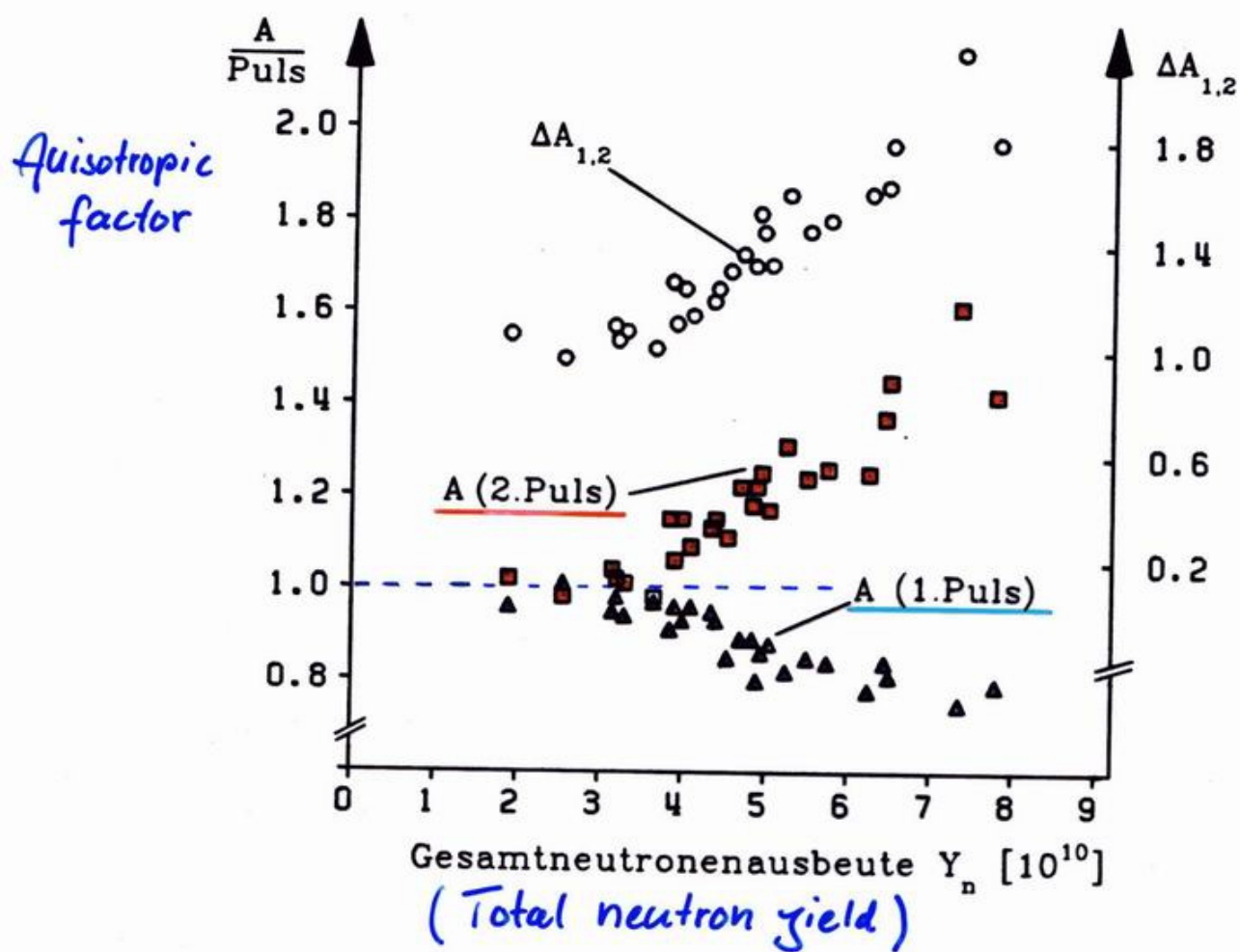
8 mbar



anisotropy factor of the two phases

$$A = \frac{\text{neutron flux end-on}}{\text{neutron flux side-on}}$$

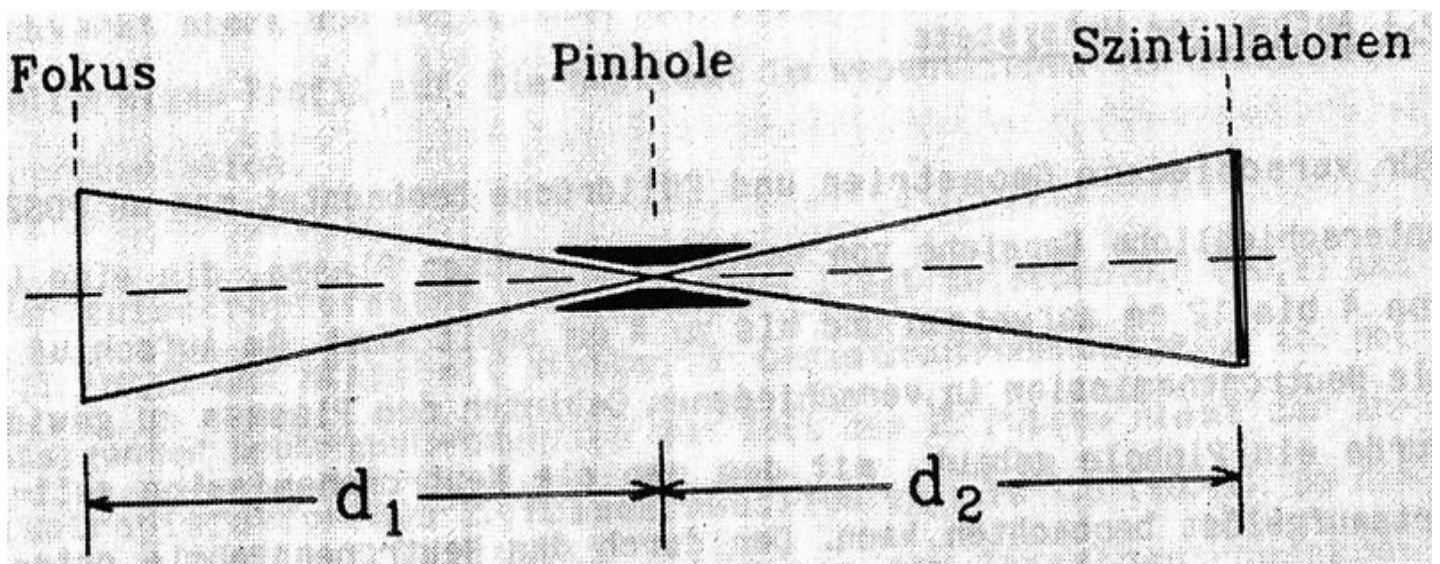
Neutron Anisotropic factor of 1. and 2. Neutron Pulse



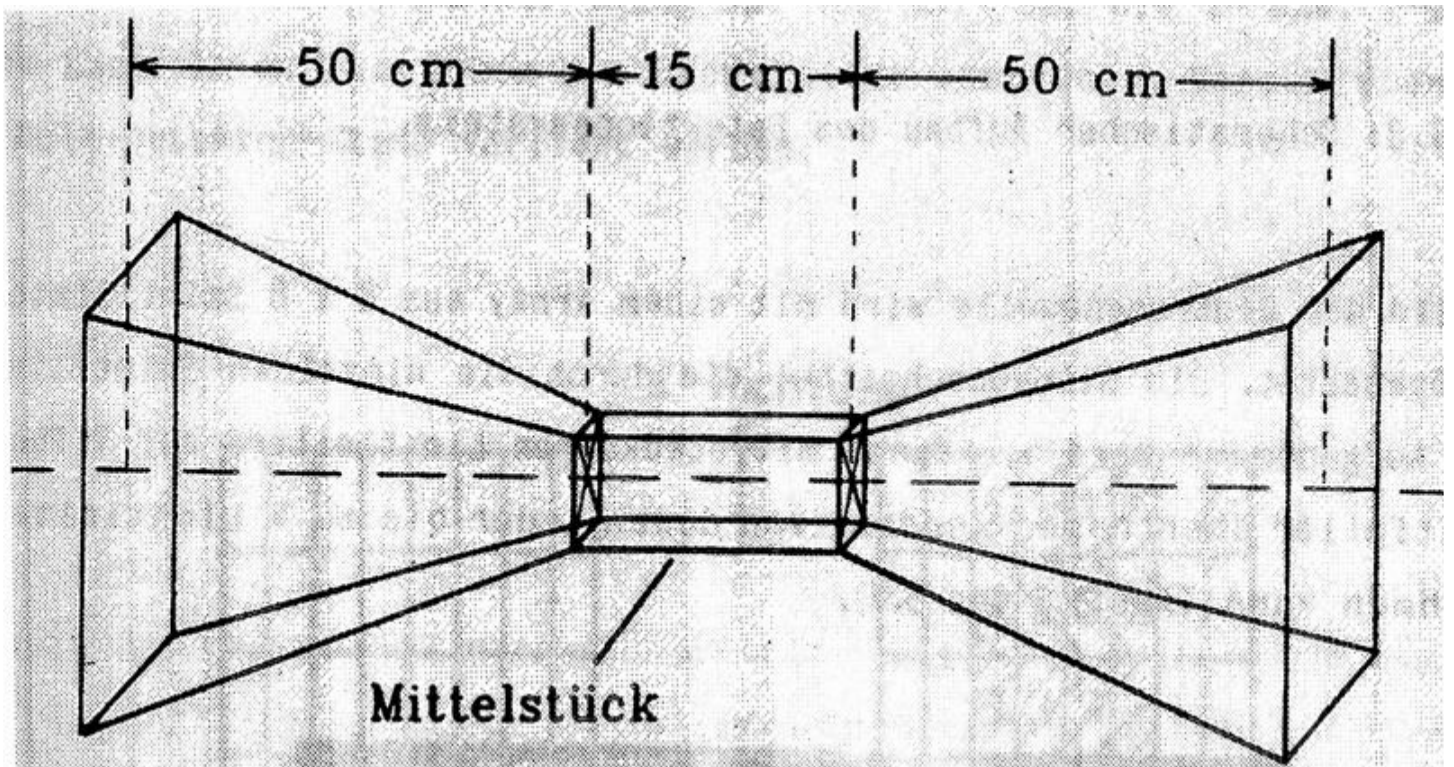
$$\text{Anisotropic factor } A = \frac{Y(\text{"end-on"})}{Y(\text{"side-on"})}$$

Neutron pinhole measurements

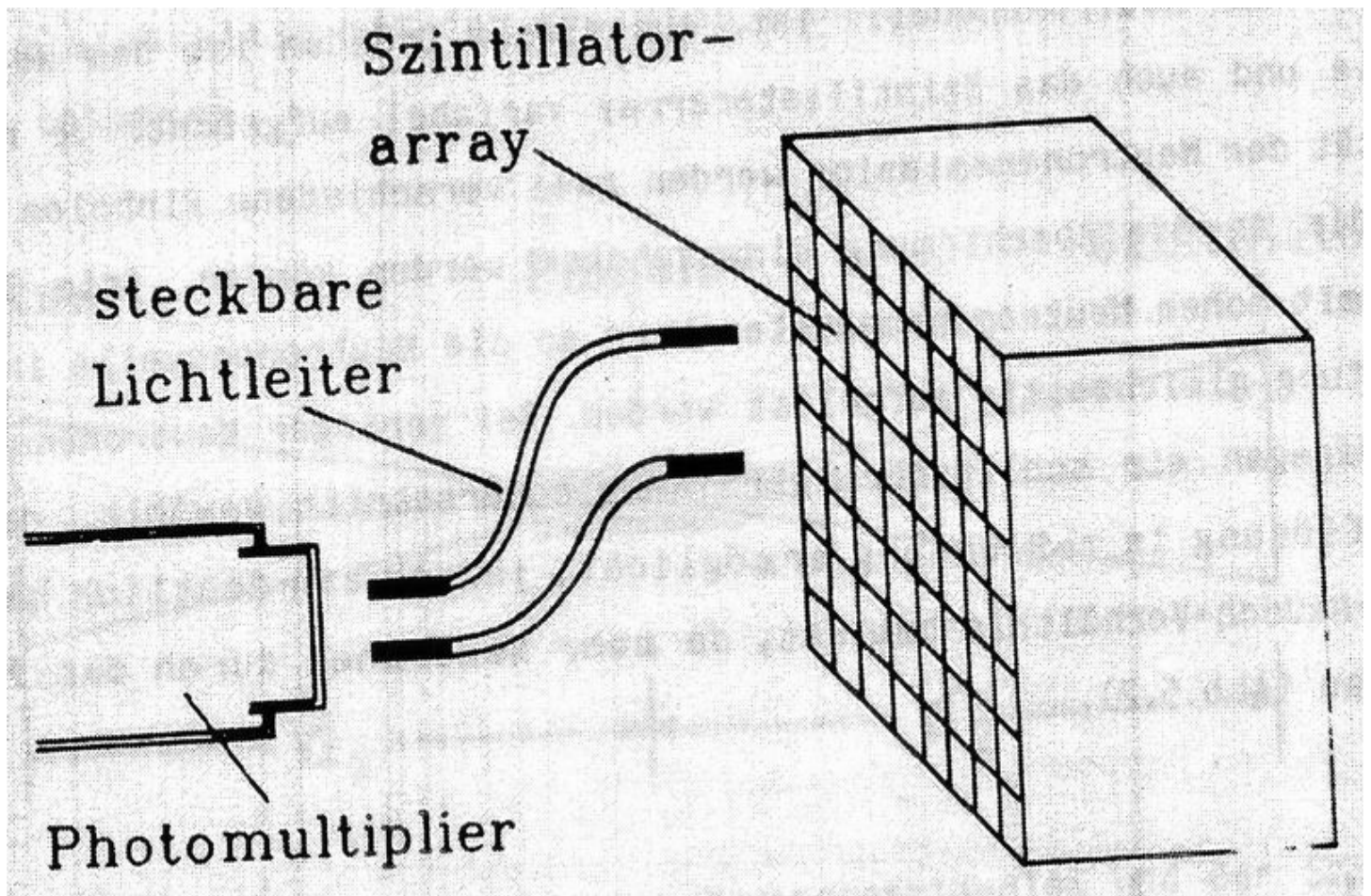
Pinhole Measurements Setup for Neutrons



Pinhole for Neutrons

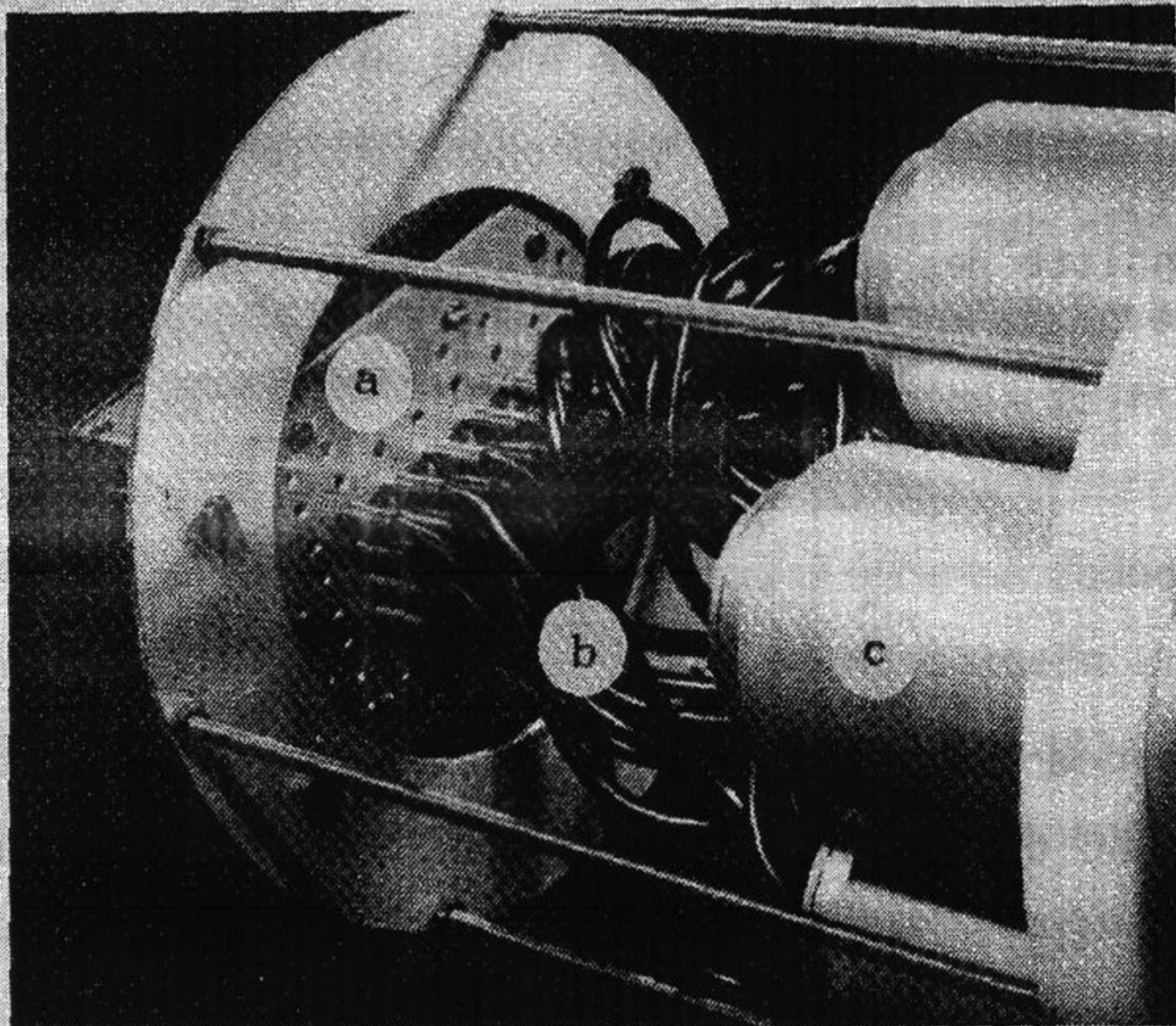


Array of Scintillators Connected via Light Fibres To Photomultipliers



a - Scintillator-array

b - light pipes



7 PM-tubes

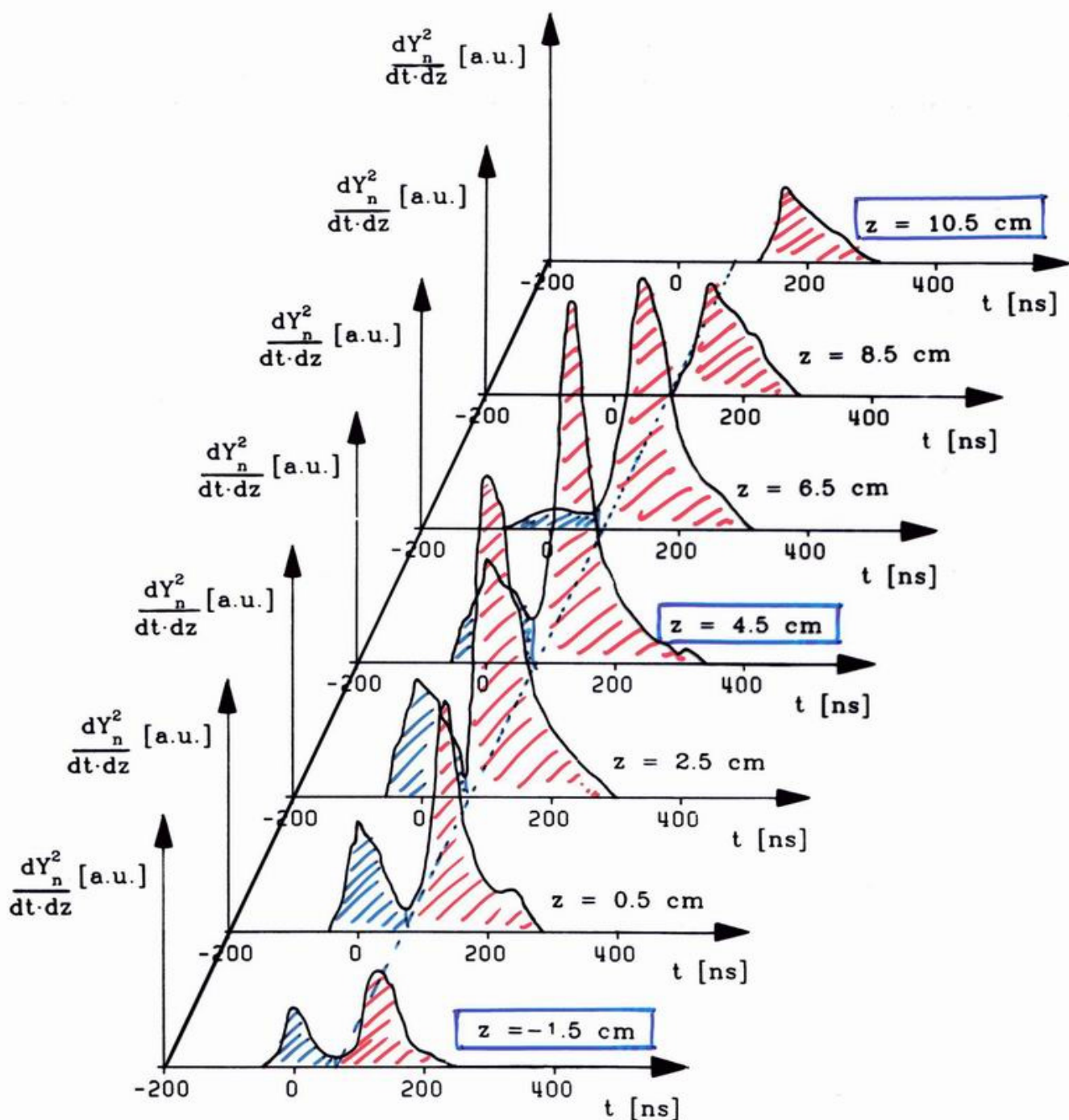
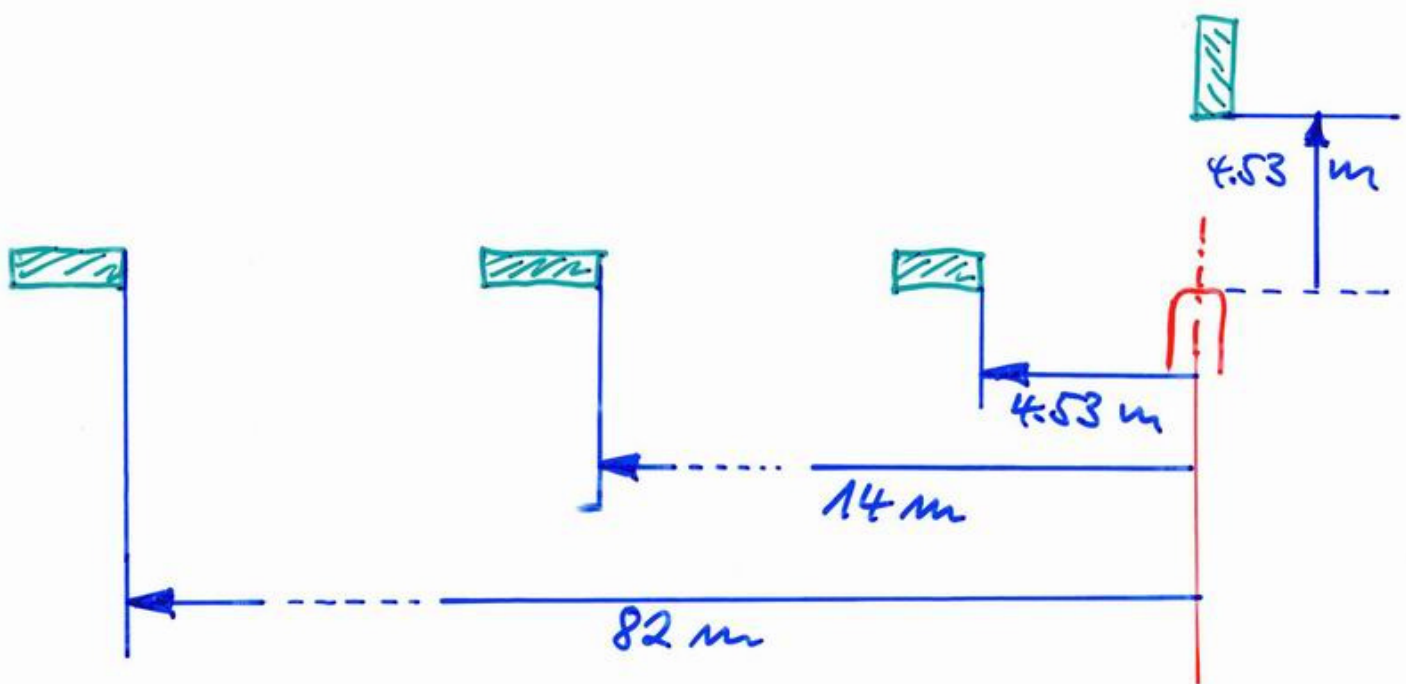


Fig. 9: Result of neutron pinhole measurements. Neutron emission spatially resolved in axial direction (resolution $\Delta z \leq 2$ cm, $\Delta t \leq 20$ ns) on POSEIDON (280 kJ, 60 kV, 500 Pa D₂, $Y_n = 6.6 \cdot 10^{10}$, hollow inner electrode, 131 mm diameter).

Neutron Time of Flight Measurements

Neutron measurement

Time of flight detectors

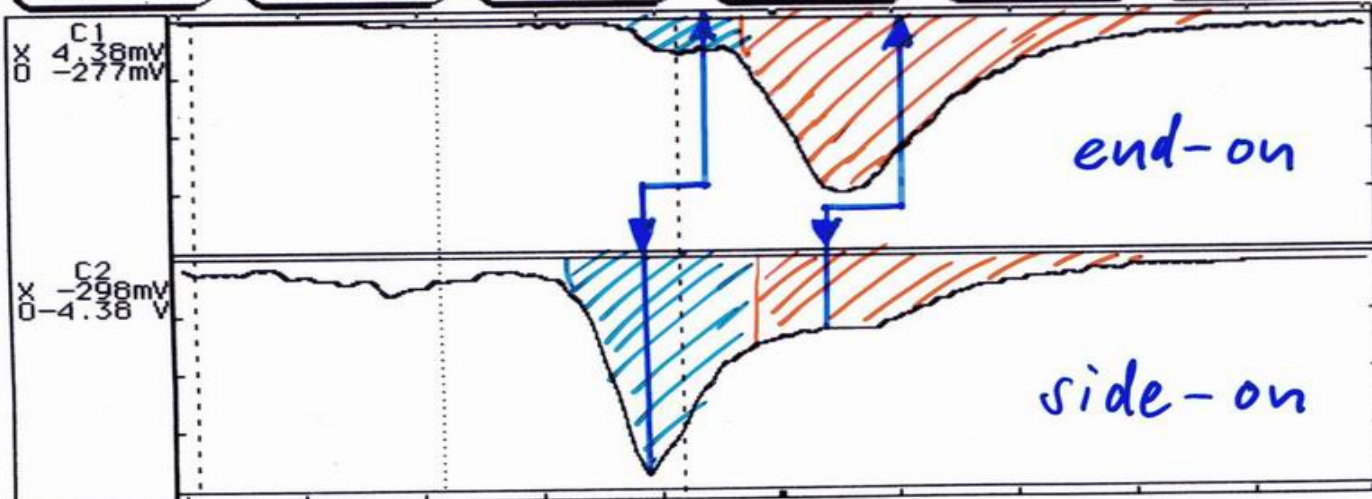


shot 10414

Oscilloscope E	Channel	Autoscale	Cancel	Run
Input C1	V/Div 500 mV	Offset -980 mV	Probe 1:1	Impedance 50 Ohms
s/Div 600 ns	Delay 2.7570 us	Markers On	X to 0 3.1320 us	Trig to X 864.00 ns
				Trig to 0 3.9960 us



Input C2	Period	Freq	Vp_p 5.44 V
s/Div 50.0 ns	Risetime 142.03 ns	+Width	Preshoot 0 %
Delay 142.00 ns	Falltime 28.711 ns	-Width 47.975 ns	Overshoot 1.75 %
Markers On	X to 0 204.00 ns	Trig to X -103.00 ns	Trig to 0 101.00 ns



$$\gamma_n = 8.73 \cdot 10^{10}$$

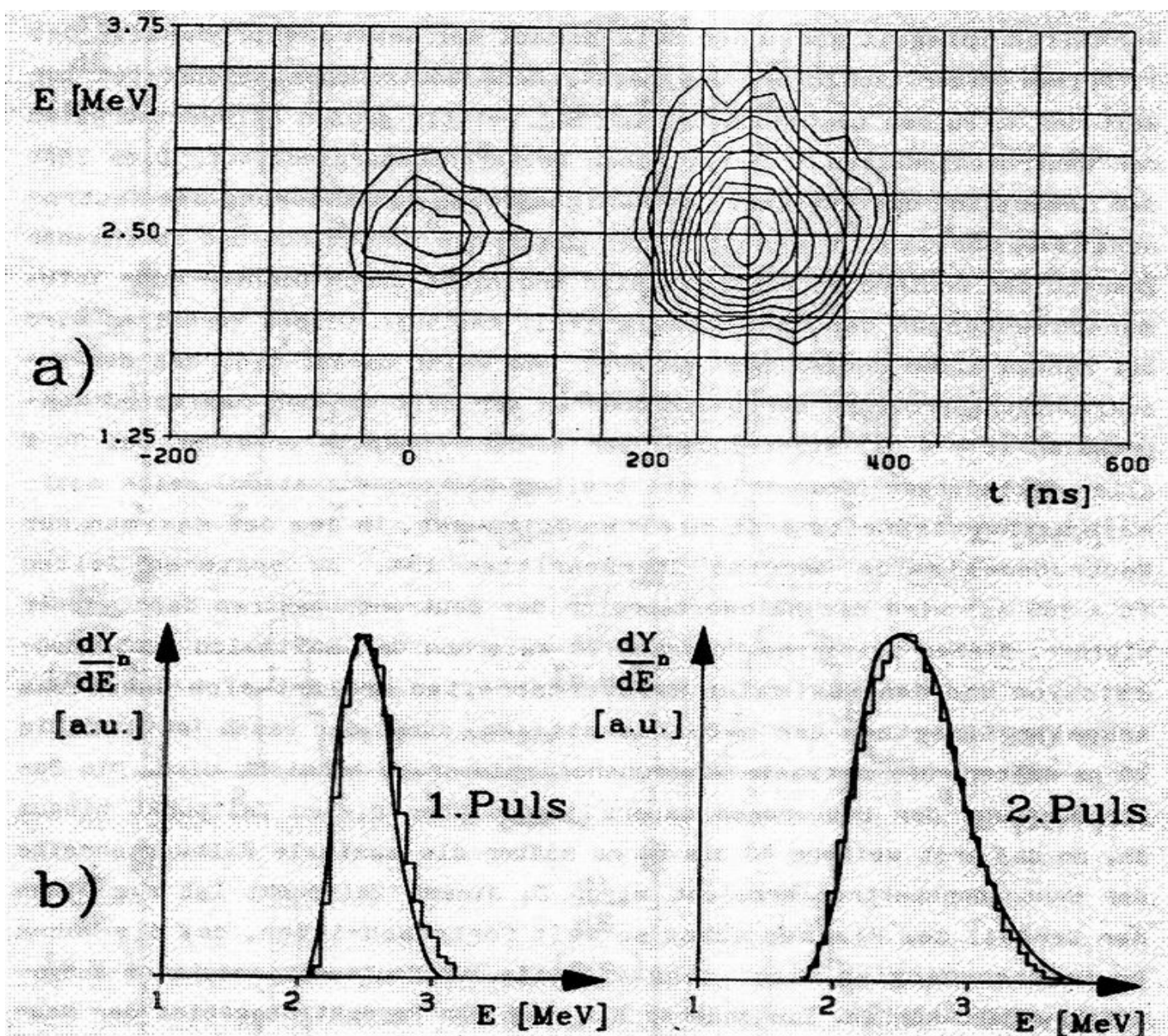


Abb. 6.14: Rekonstruktion einer typischen POSEIDON-Entladung, Schuß Nr. 7522, breiter Elektrodenabstand, Ladespannung 70 kV, Füll-
druck 8 mbar.

a) Höhenlinienbild.

b) Neutronenspektren der Einzelpulse (zeitintegriert): ermittelt aus der rekonstruierten Quellfunktion (Stufenfunktion) und berechnet nach dem verallgemeinerten "Beam-Target"-Modell (glatte Kurve).

Neutronenspektren

1. Puls

2. Puls

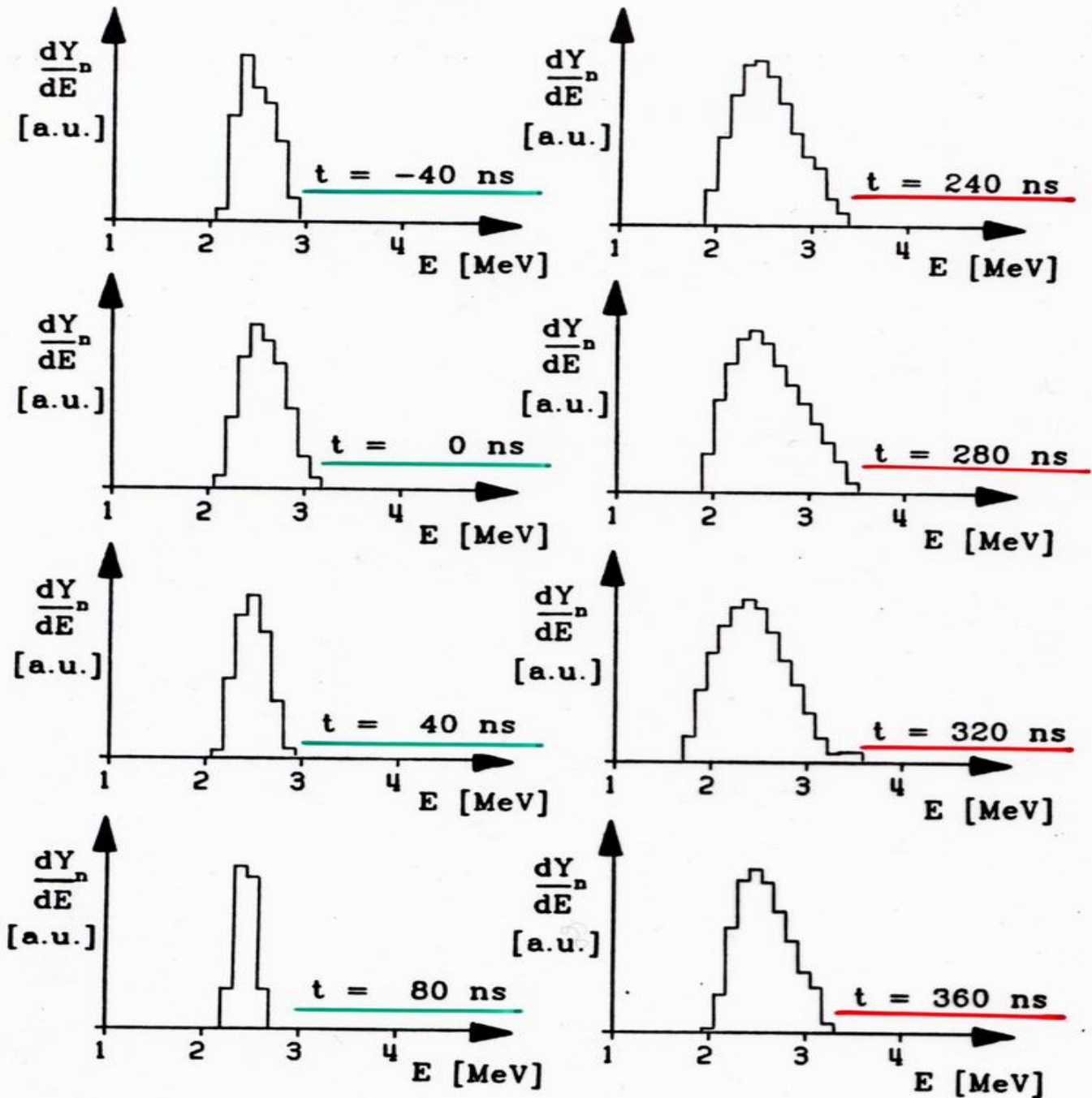


Abb. 6.15: Neutronenspektren einer typischen POSEIDON-Entladung zu verschiedenen Zeiten, Schuß Nr. 7522 (vgl. Abb.6.14a).

Neutronenspektren

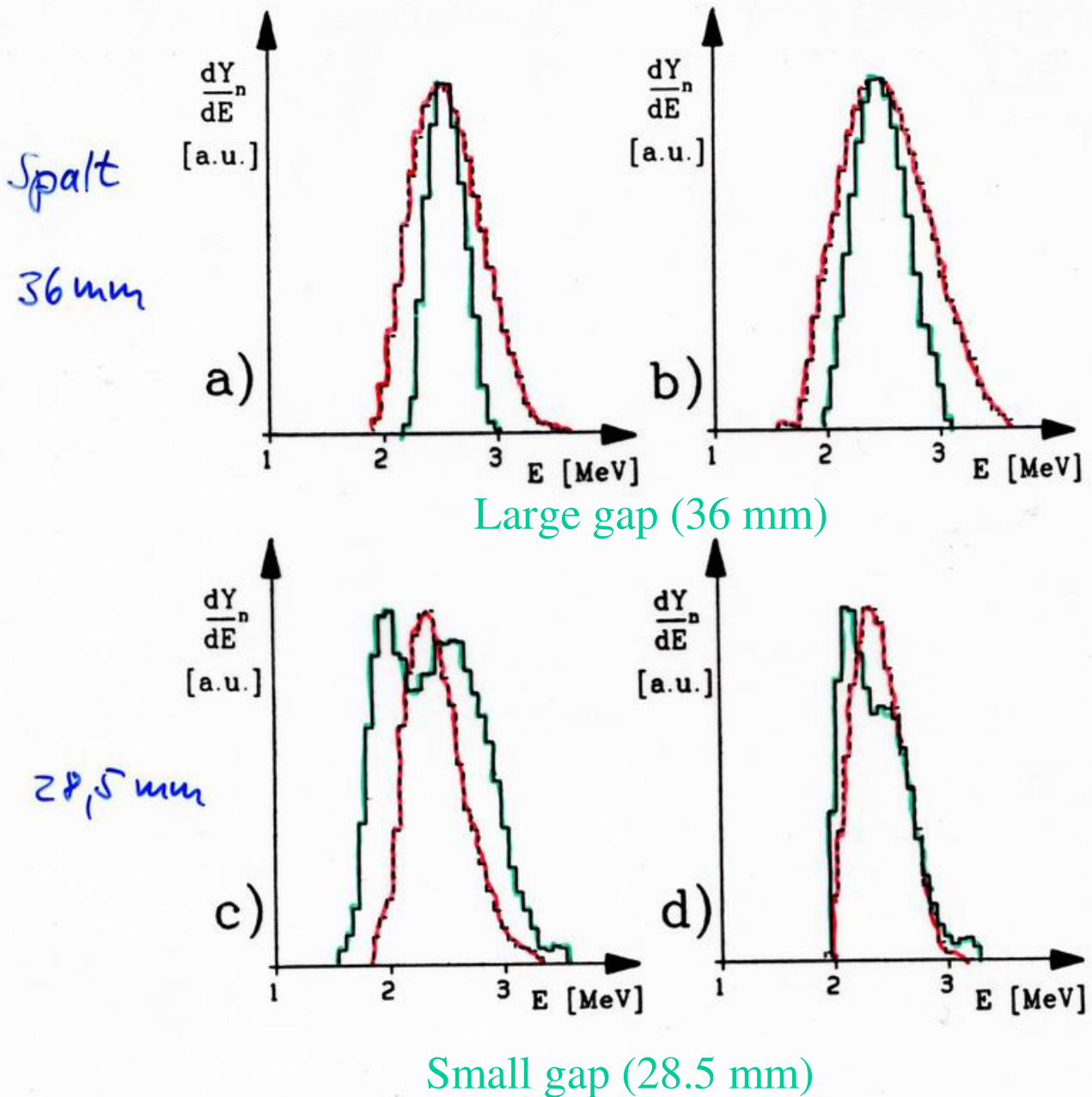


Abb. 6.19: Typische Neutronenspektren des 1. Pulses (—) und des 2. Pulses (---) mit breitem (a,b) und engem (c,d) Elektrodenabstand.

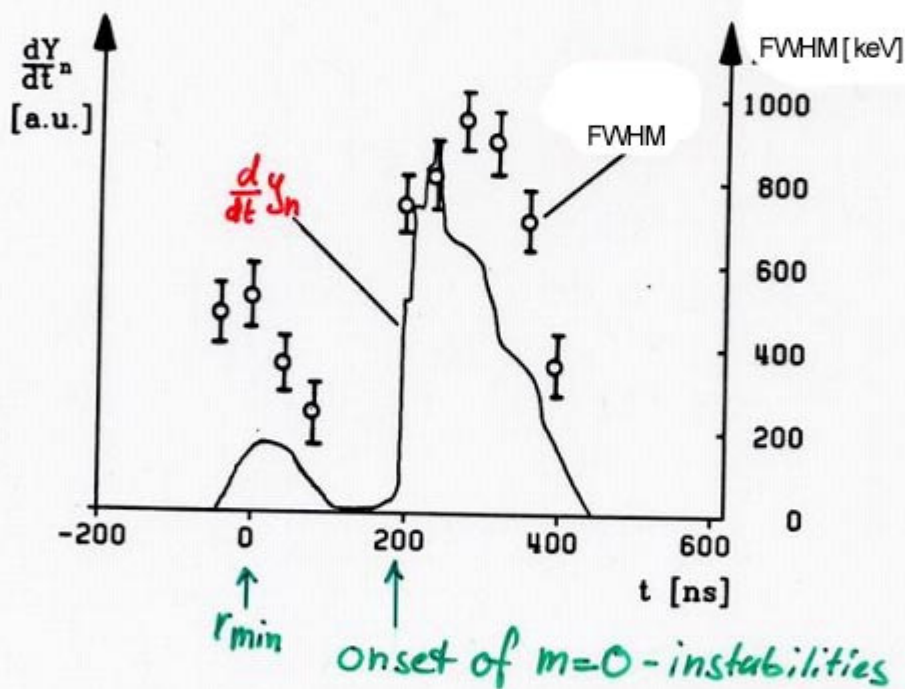
$Y_n/10^{10}$
1.6
6.2
2.7
1.3

a) $Y_n(1.Puls) = 0,3 \cdot 10^{10}$, $Y_n(2.Puls) = 1,3 \cdot 10^{10}$.
b) $Y_n(1.Puls) = 0,6 \cdot 10^{10}$, $Y_n(2.Puls) = 5,6 \cdot 10^{10}$.
c) $Y_n(1.Puls) = 1,9 \cdot 10^{10}$, $Y_n(2.Puls) = 0,8 \cdot 10^{10}$.
d) $Y_n(1.Puls) = 0,5 \cdot 10^{10}$, $Y_n(2.Puls) = 0,8 \cdot 10^{10}$.

Y_2/Y_1
4.3
9.3
0.4
1.6

Neutron Yield and FWHM of Neutron Spectra as a function of time

-> relaxation of deuterons



Typical ROSEIDON shot, voltage 70 kV, filling pressure 8 mbar.

$$Y_n (1. \text{ pulse}) = 0.6 \cdot 10^{10}$$

$$Y_n (2. \text{ pulse}) = 4.1 \cdot 10^{10}$$

Why Reaction Proton Diagnostics?

Beam properties

$$\underline{\underline{E_d(r, z, t, v)}}$$

Target properties

$$\underline{\underline{n_e(r, z, t)}}$$

$$\underline{\underline{T_e(r, z, t)}}$$

Field properties

$$\underline{\underline{B_\theta(r, z, t)}}$$

$$? \underline{\underline{B_z(r, z, t)}}$$

Model
calculations

Neutron emission

$$\underline{\underline{\gamma_n(r, z, t, v)}}$$

$$\underline{\underline{E_n(r, z, t, v)}}$$

Reaction proton
emission

$$\underline{\underline{\gamma_p(r, z, t, v)}}$$

$$\underline{\underline{E_p(r, z, t, v)}}$$

Proton Diagnostics

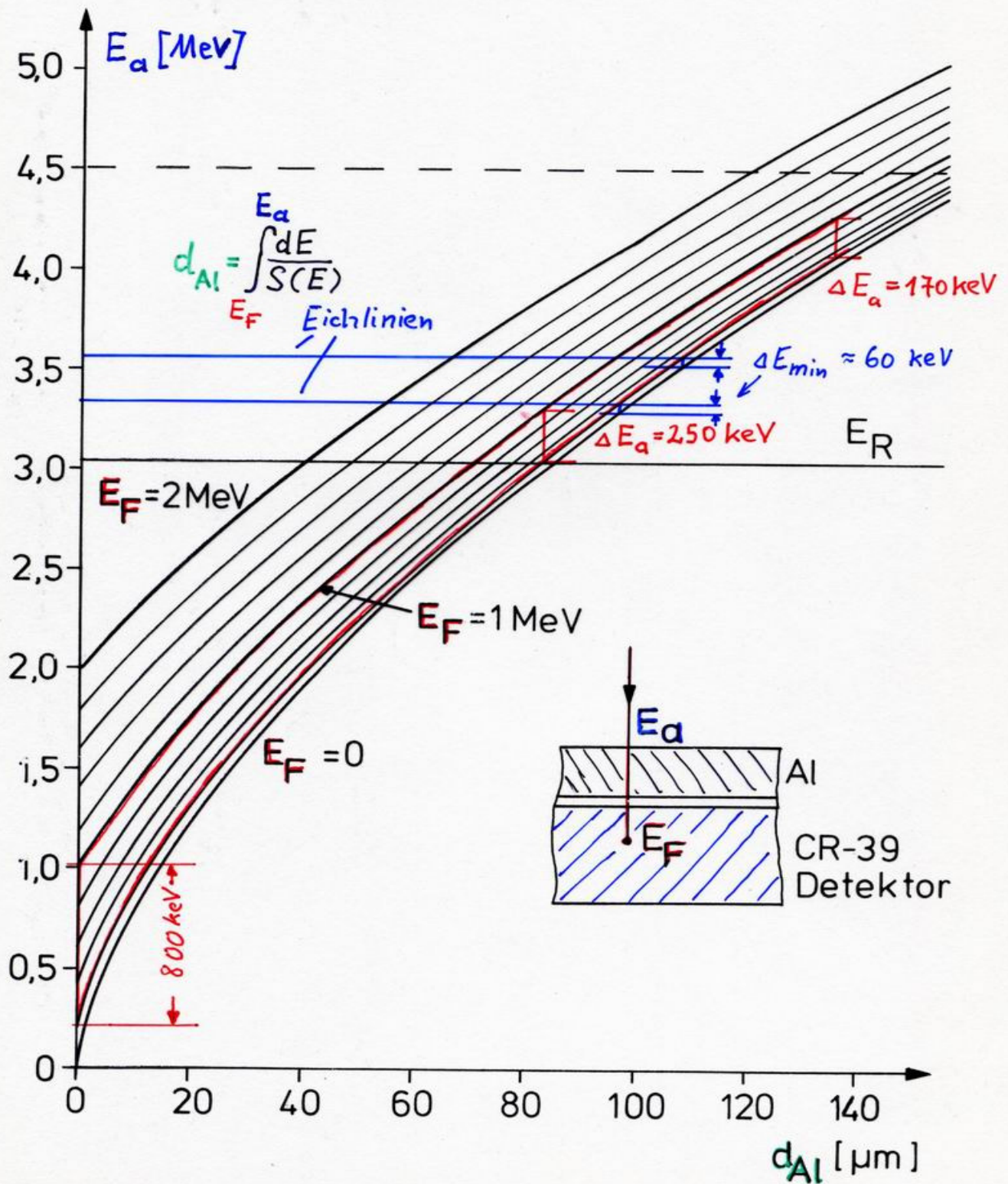
Solid state nuclear track detectors are used for the registration of protons

Individual particles can be registered using an etching process.

Small craters are visible under the Microscope

Spectral information can be obtained using absorbing metal foils

Anfangsenergien für Protonen in Aluminium



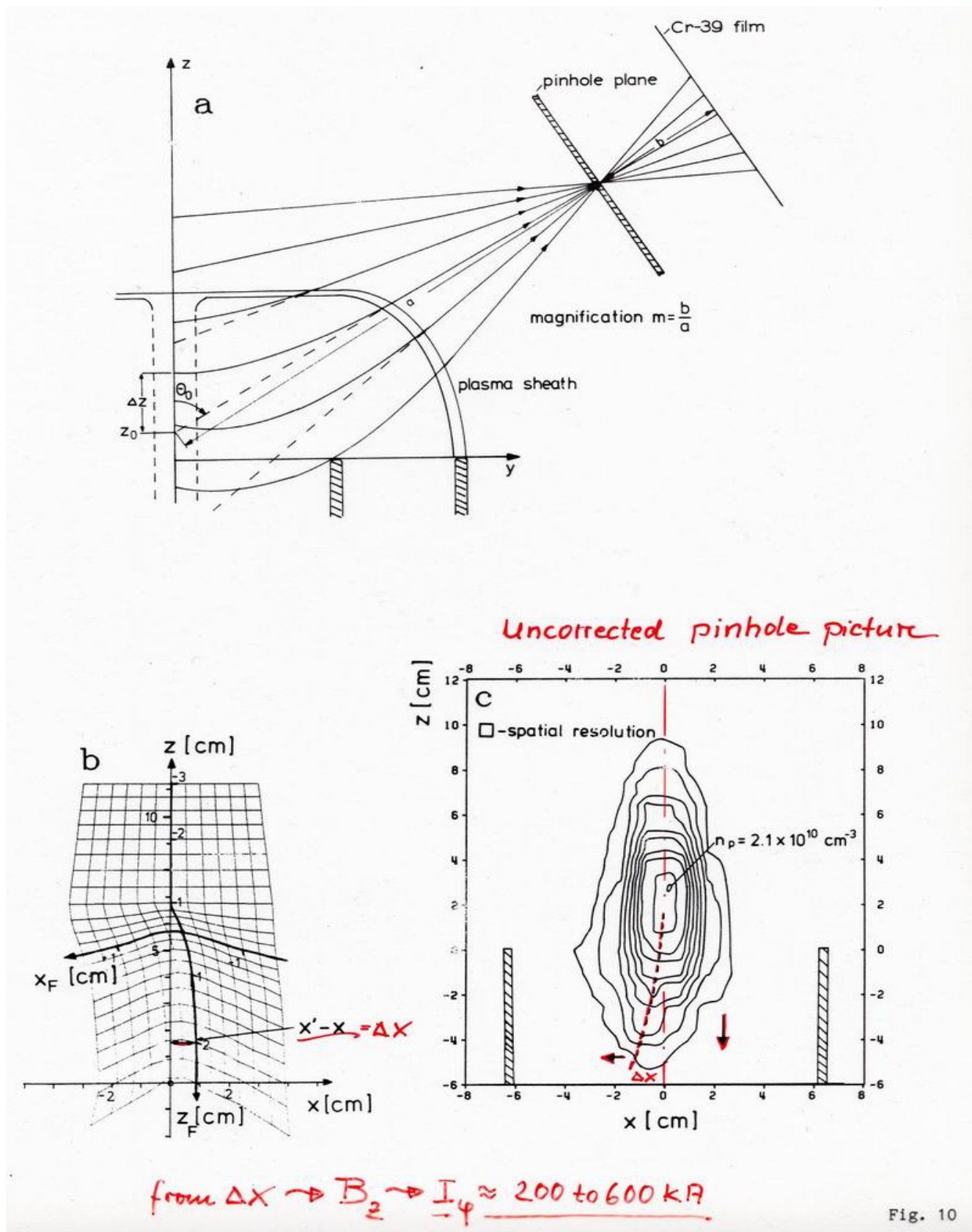
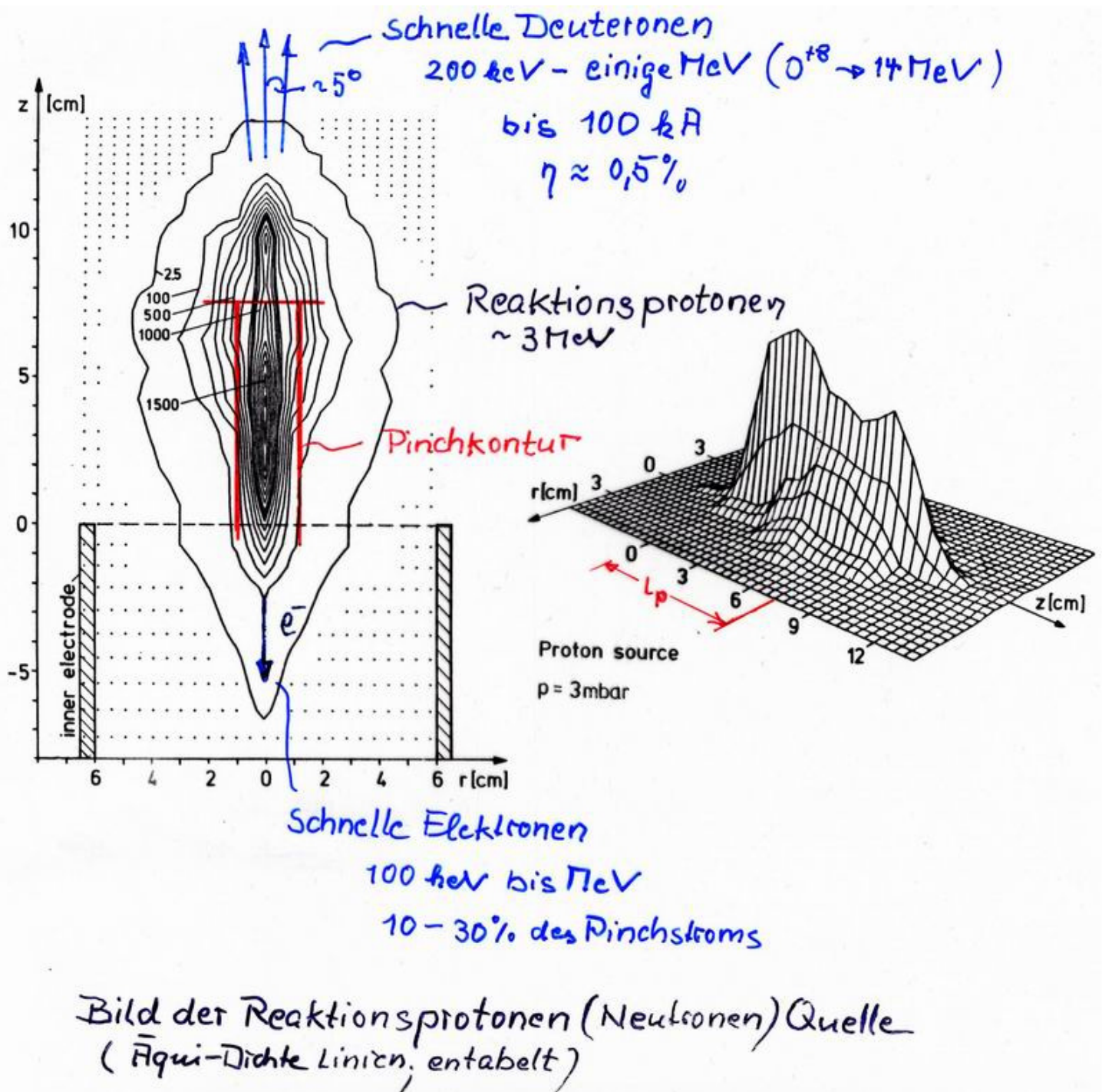


Fig. 10



Pinhole picture of

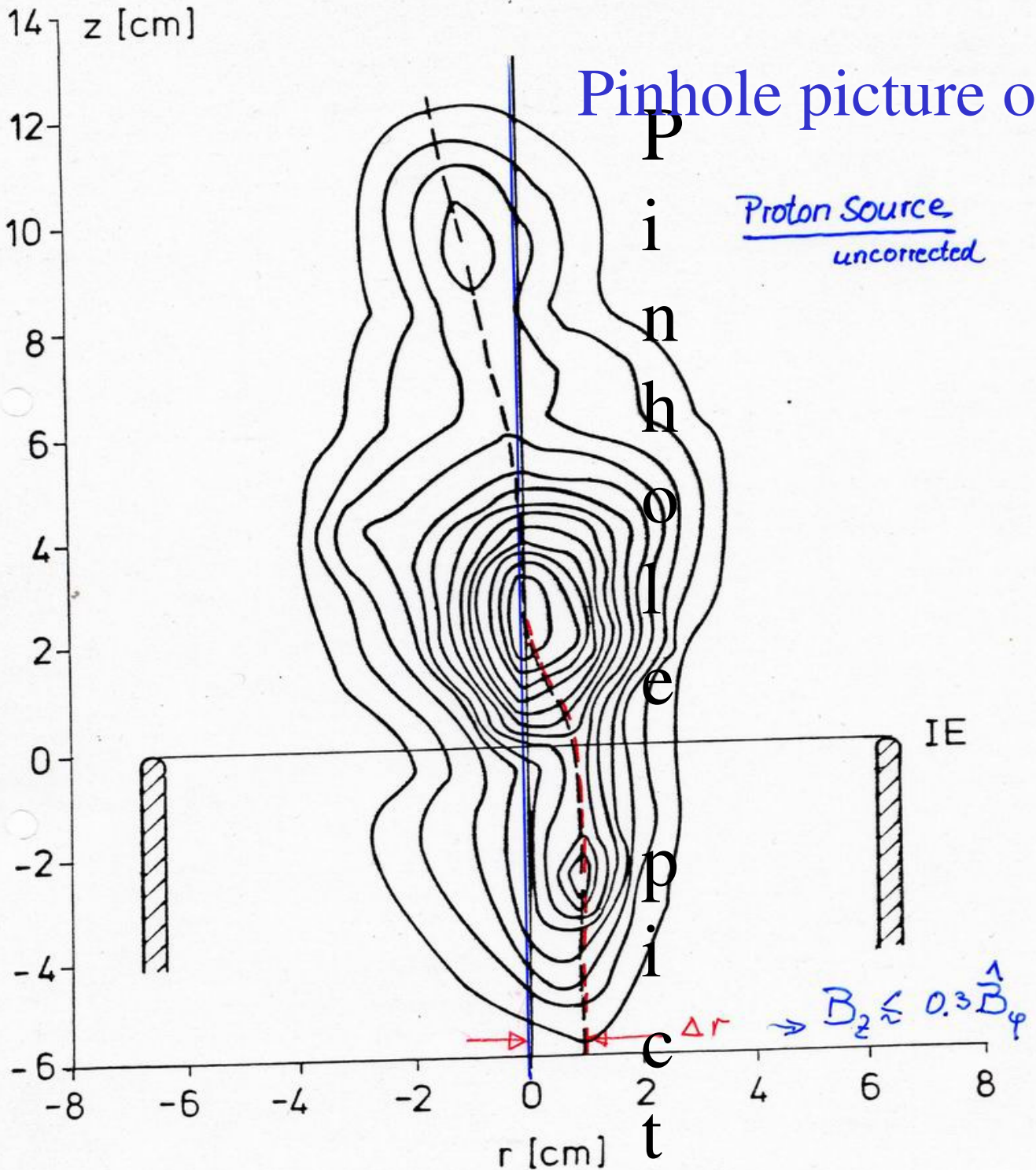


Fig. 7: Pinhole picture of the reaction proton source, uncorrected for beam deviation by focus magnetic fields. The contours are lines of constant proton density on the film. IE = inner electrode. $W_0 = 280$ kJ; $p_0 = 5$ mbar D_2 .

Proton spectra

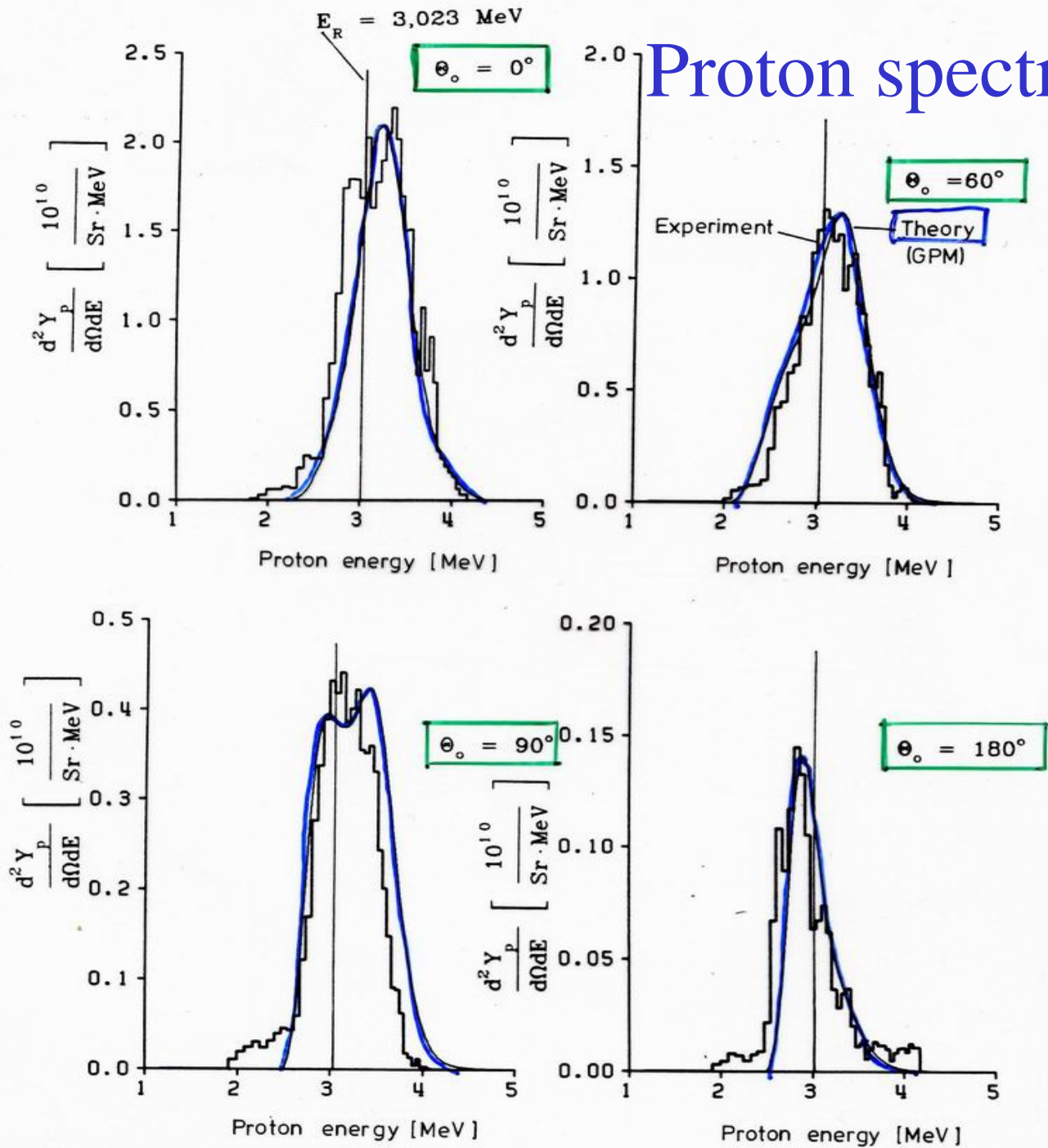
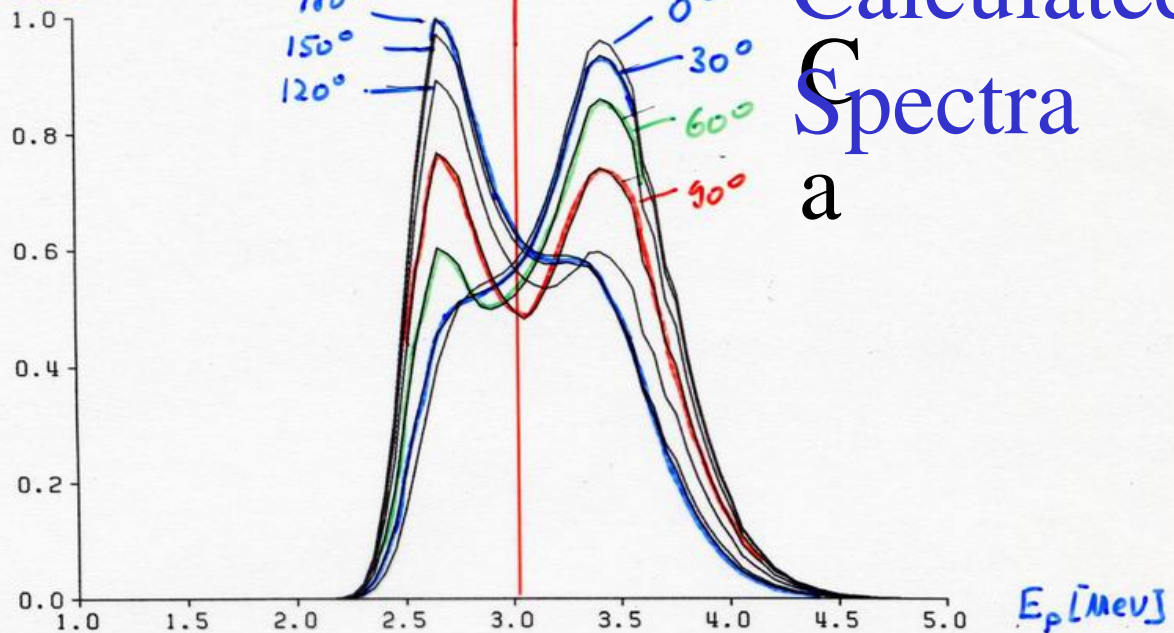


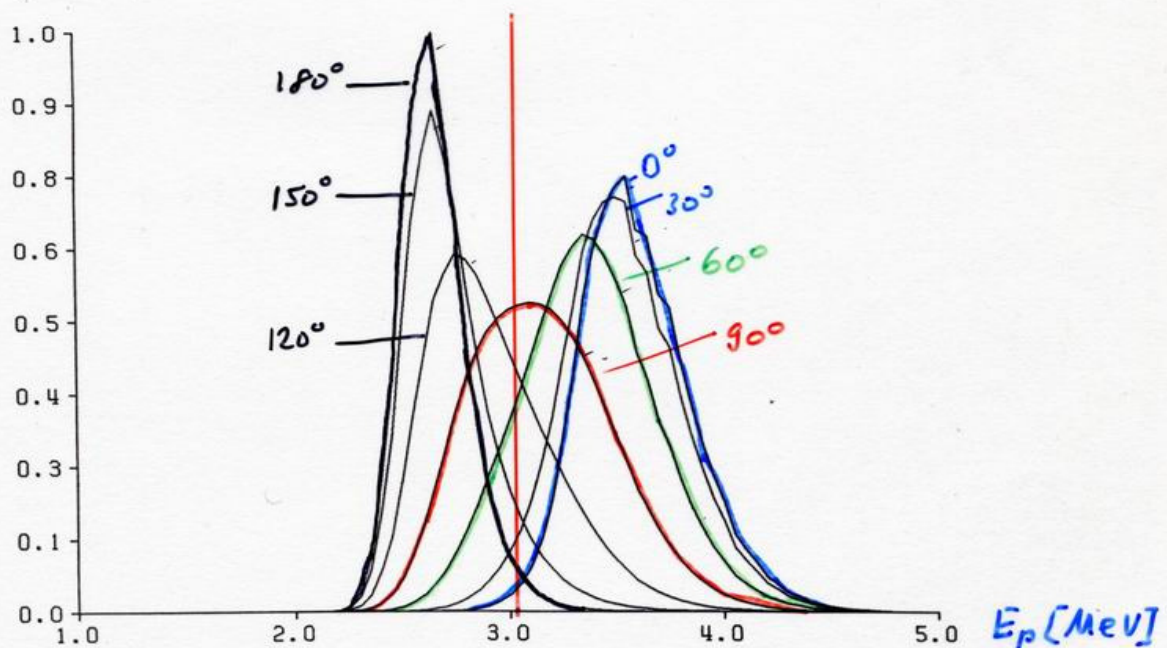
Fig. 14: Reaction proton spectra on POSEIDON (280 kJ, 60 kV, 500 Pa D_2), proton yield $Y_p = 5.7 \cdot 10^{10}$, contribution of the first pulse 34%. For the GPM-calculations [30] of the curves the following parameters were taken: $T_{i1}^* = 75 \text{ keV}$, $N_{b1} = 5 \cdot 10^{16}$, $T_{i2}^* = 200 \text{ keV}$, $N_{b2} = 3 \cdot 10^{14}$, $A_d = 3$, $I_{p1} = 790 \text{ kA}$.

$\frac{dY}{dE d\Omega}$ [a.u.]



Calculated
Spectra
a

Calculated reaction proton spectra ($T_i^* = 100$ keV)
Anisotropy factor: $A_d = 5$



($A_d = 10$; $T_i^* = 100$ keV)

Neutron diagnostics deliver valuable information on various processes which occur in the plasma, such as:

Pinch dynamics

fast ion beams interacting
with a target

fast ion energy distributions

Fusion neutron measurements
should be complemented by
Fusion proton measurements.

NEUTRONS

angular resolution

---> anisotropy
beam target processes

temporal resolution

---> two (three) pulses
various pinch phases

spatial resolution

---> dimensions of source
(better using proton measurements)

axial propagation of source

spectral resolution

---> energy of deuterons

relaxation of deuterons

Gyrating Particle Model

Beam properties

$$\underline{\underline{E_d(r, z, t, v)}}$$

Target properties

$$\underline{\underline{n_e(r, z, t)}}$$

$$\underline{\underline{T_e(r, z, t)}}$$

Field properties

$$\underline{\underline{B_\theta(r, z, t)}}$$

$$? \underline{\underline{B_z(r, z, t)}}$$

Model
calculations

Neutron emission

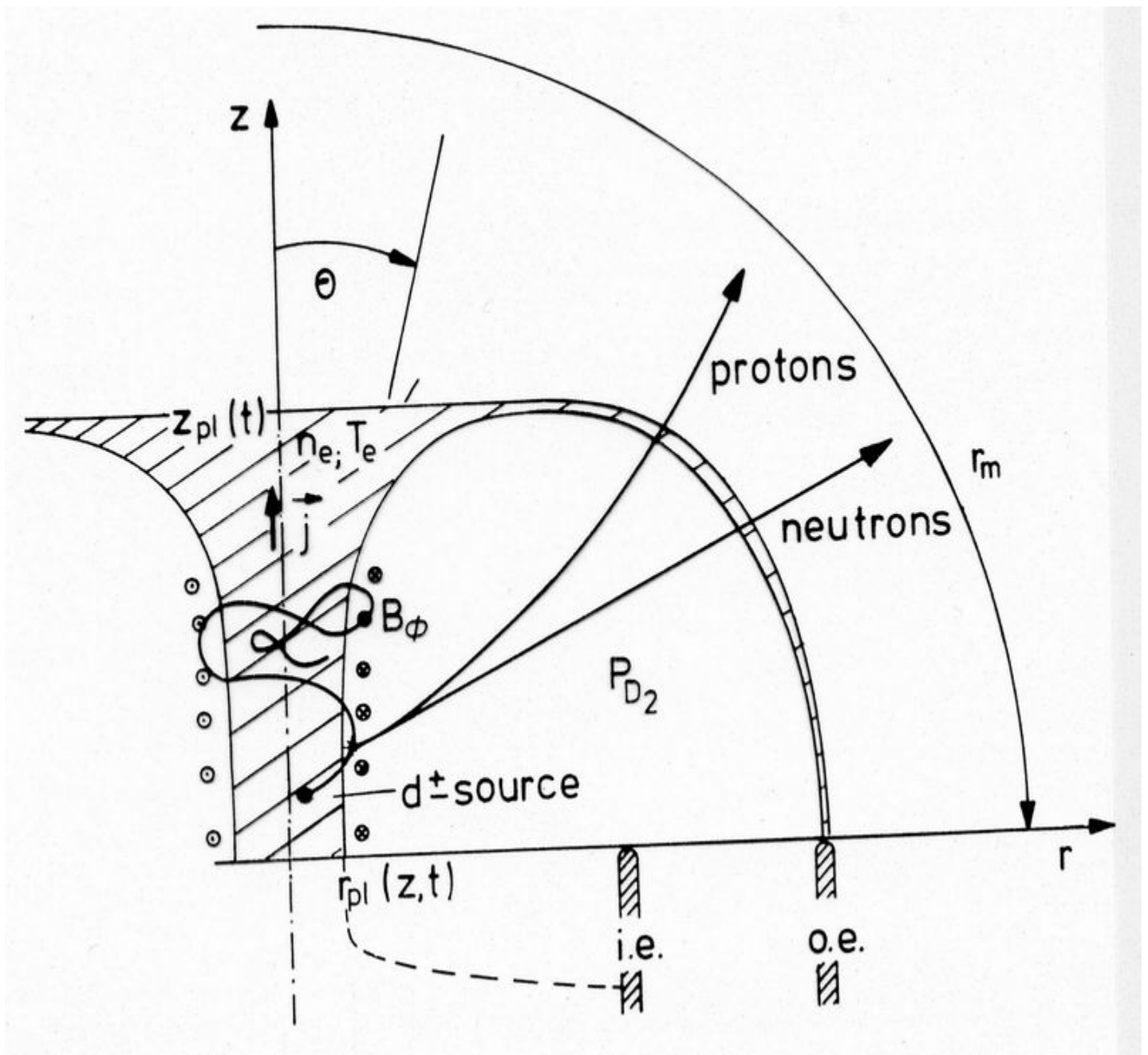
$$\underline{\underline{\gamma_n(r, z, t, v)}}$$

$$\underline{\underline{E_n(r, z, t, v)}}$$

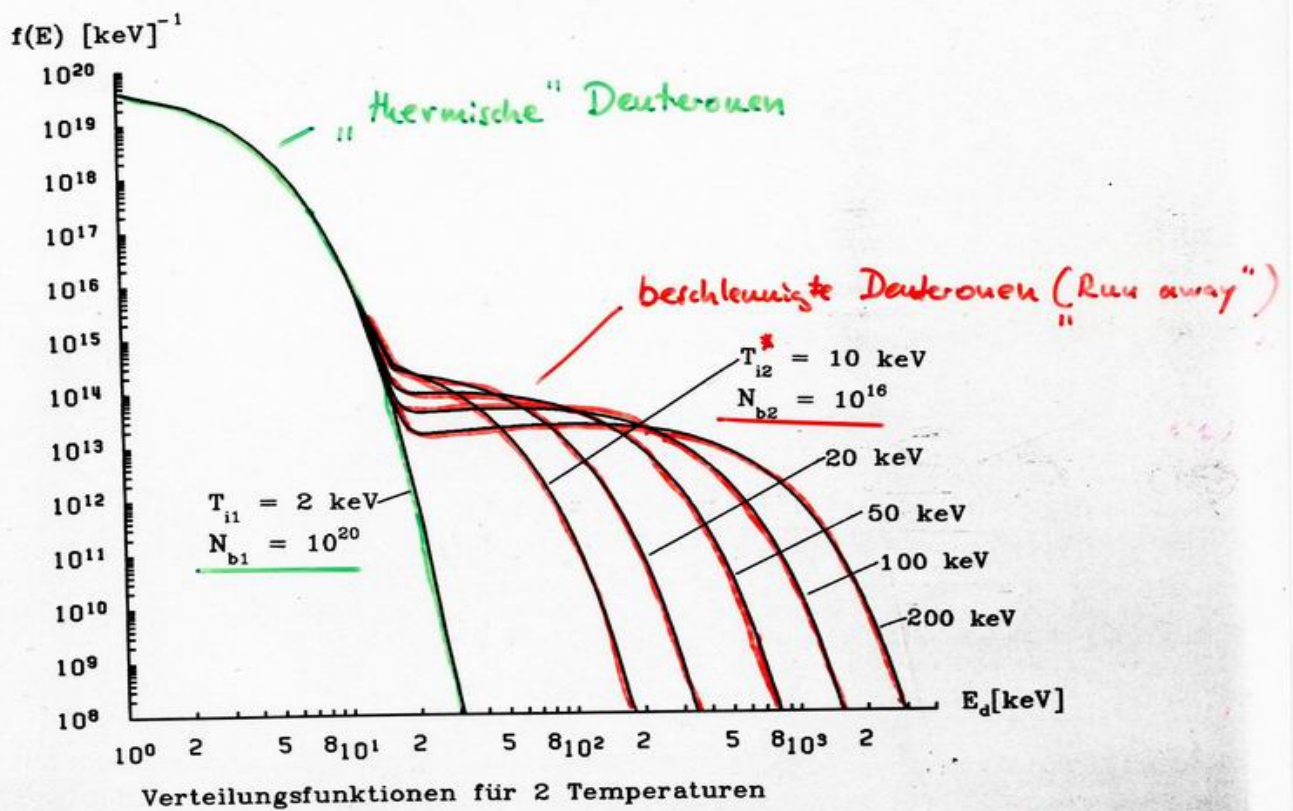
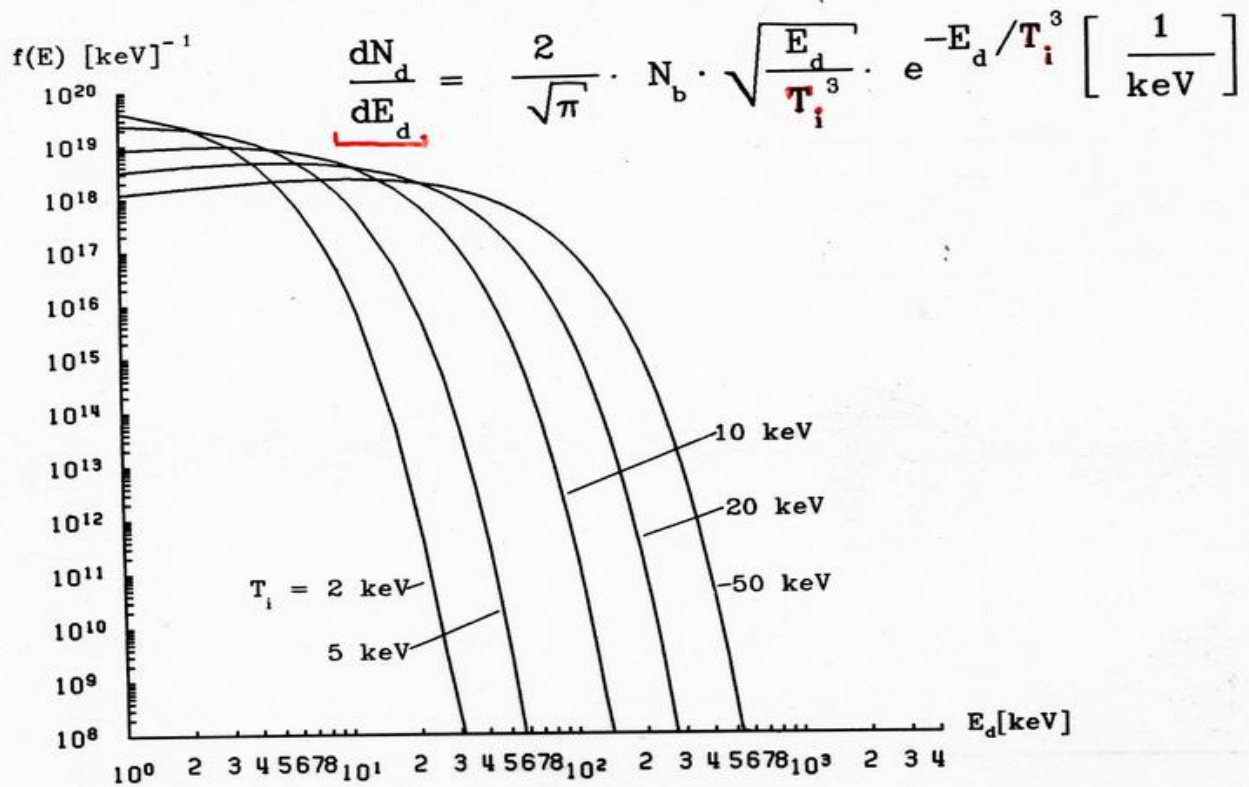
Reaction proton
emission

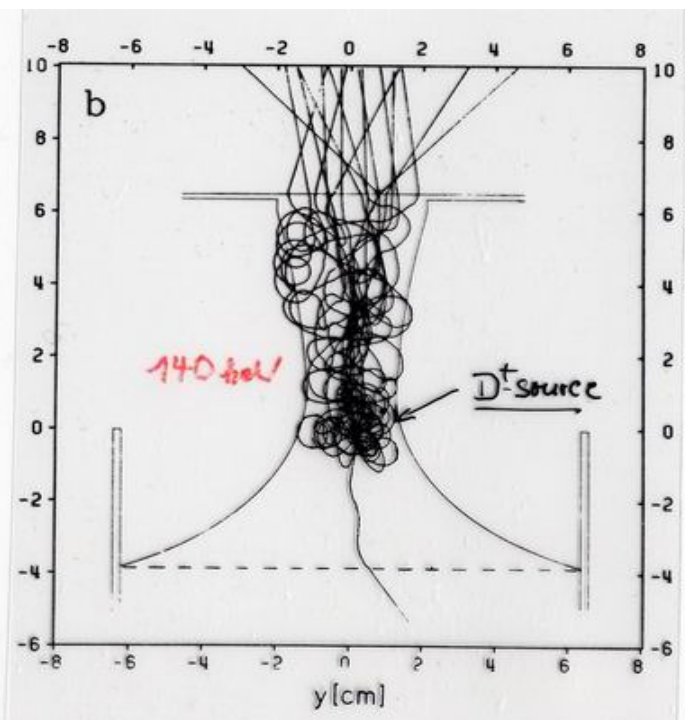
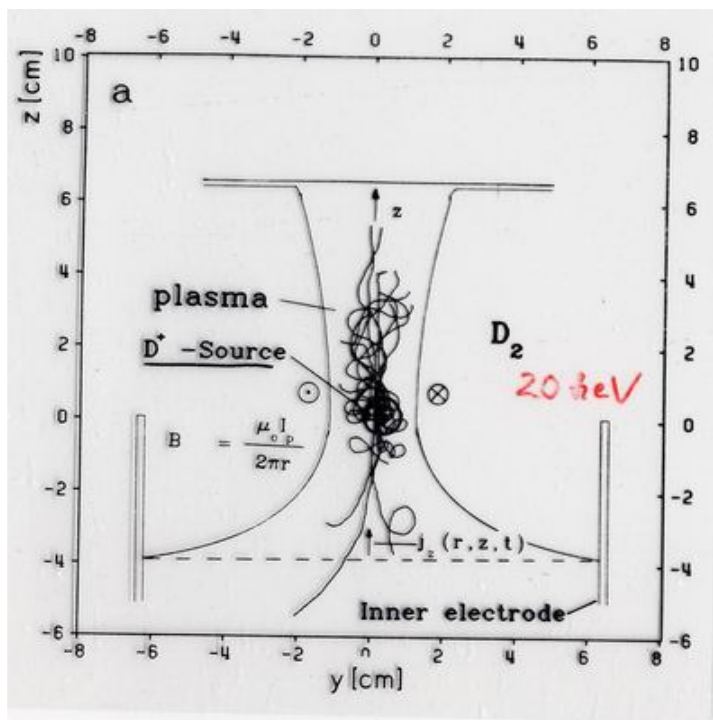
$$\underline{\underline{\gamma_p(r, z, t, v)}}$$

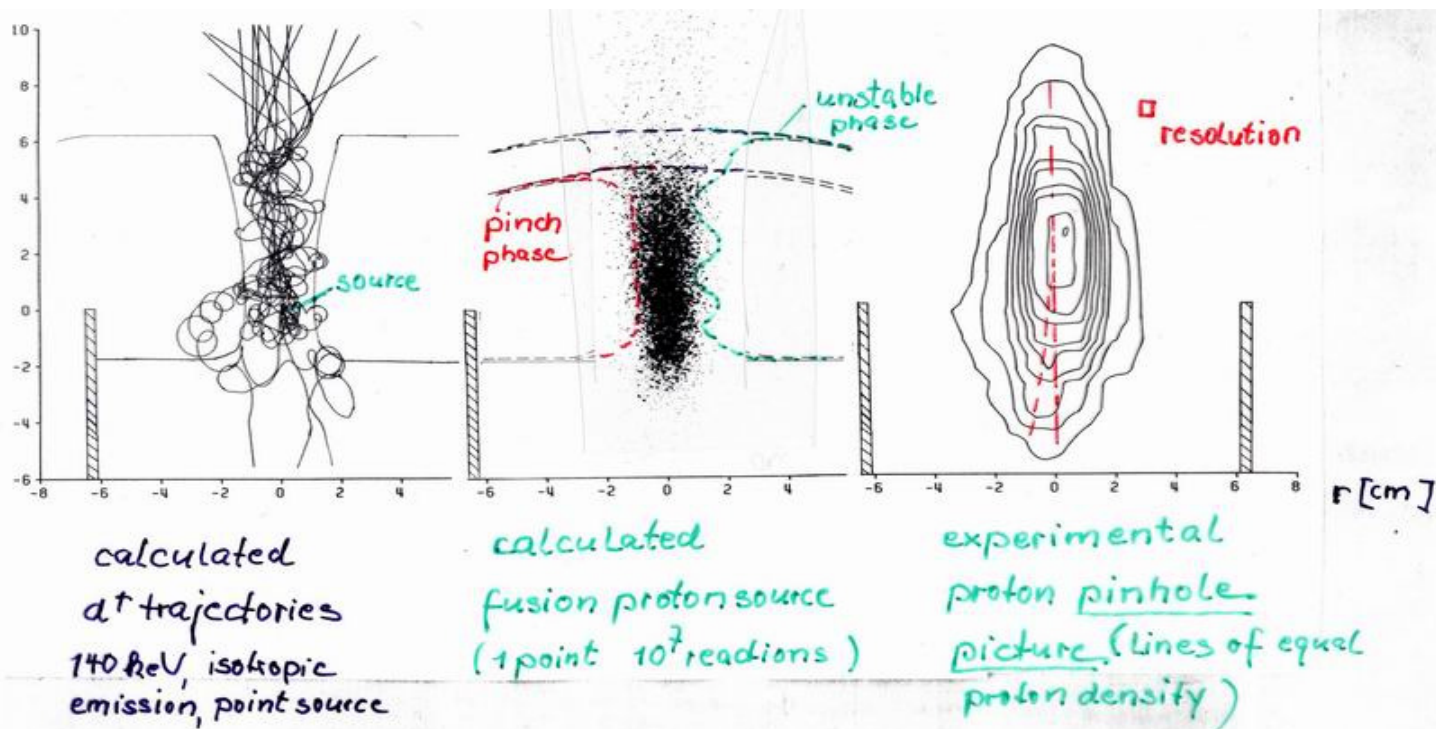
$$\underline{\underline{E_p(r, z, t, v)}}$$

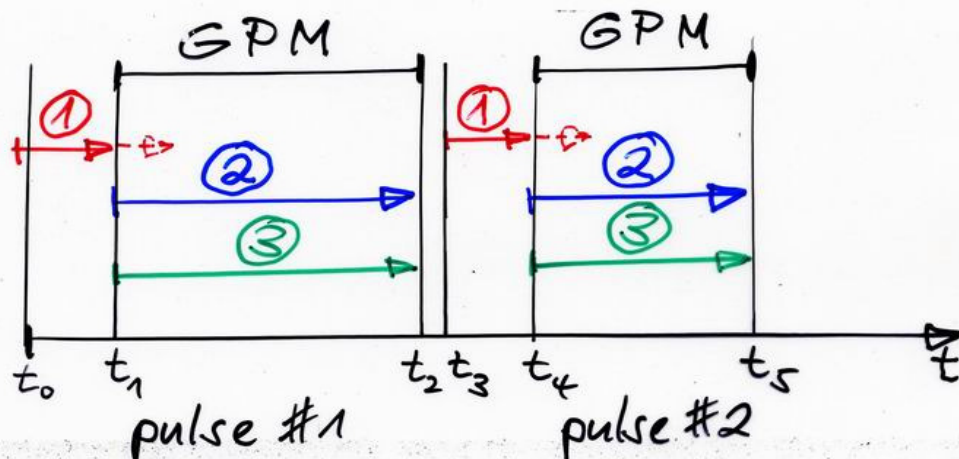


Temperatur-Verteilung für Deuteronen (aus Maxwellverteilung):









- ① ion acceleration in time- and space-dependent electromagnetic fields
 ----> deuteron distribution $f_{D1}(E, r, t, \Theta)$
- ② ion relaxation
 predominantly by Coulomb collisions with electrons $T_e, n_e(r, t), \mathbf{B}(r, t)$
 ----> deuteron distribution $f_{D2}(E, r, t)$
- ③ fusion reactions of the fast deuterons with the target $n_{i,0}(r, t)$
- ④ fusion products $n(E, t)$ and $p(E, t)$ pass through the plasma/gas up to the detectors.
 protons experience deviations in $\mathbf{B}(r, t)$

measured quantities

POSEIDON diagnostics:

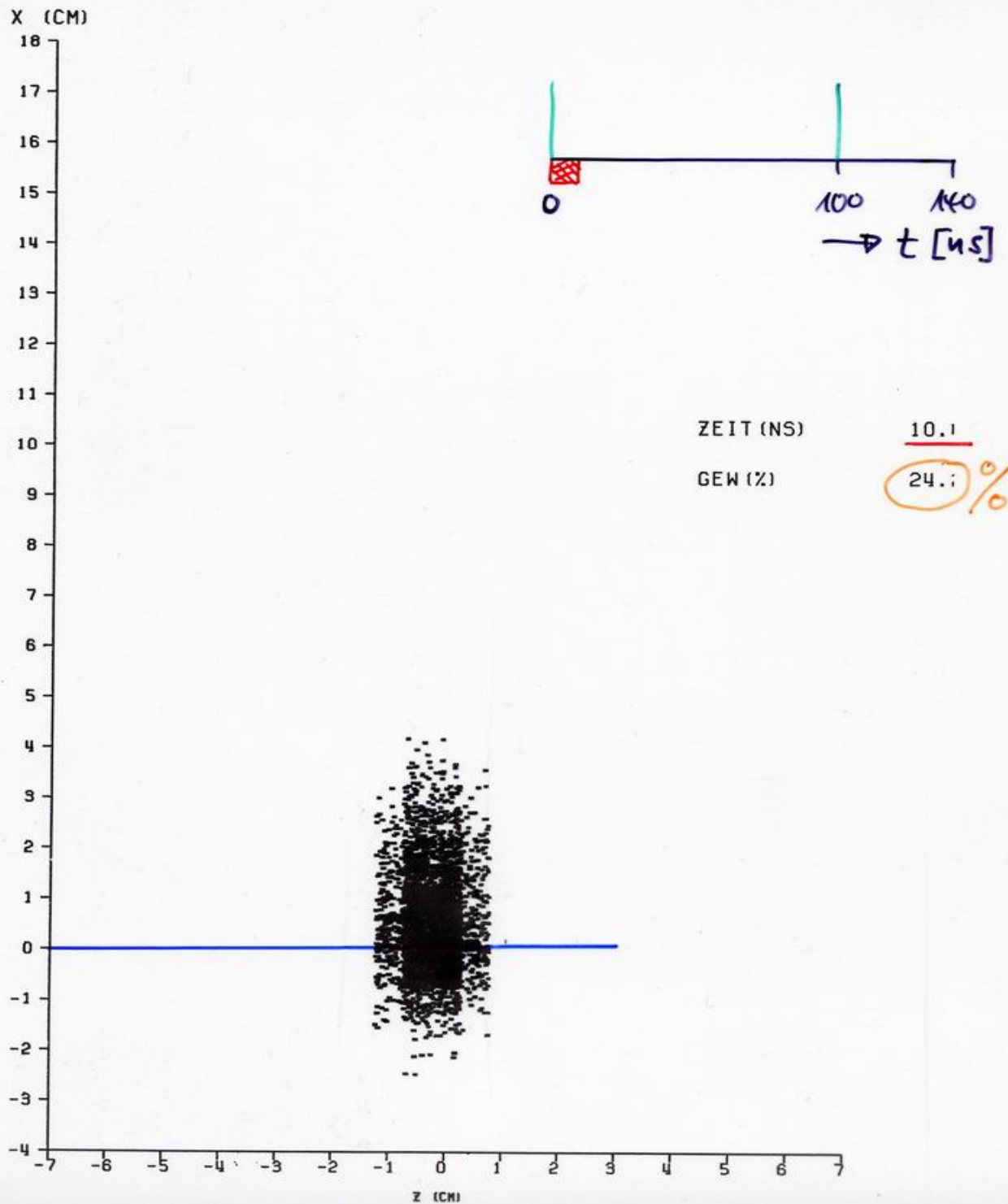
neutrons $n(E, t)$ in one direction (side on)

protons $p(E)$ in 8 directions

ORTSVERTEILUNG NEUTRONEN UND PROTONEN

POSEIDON

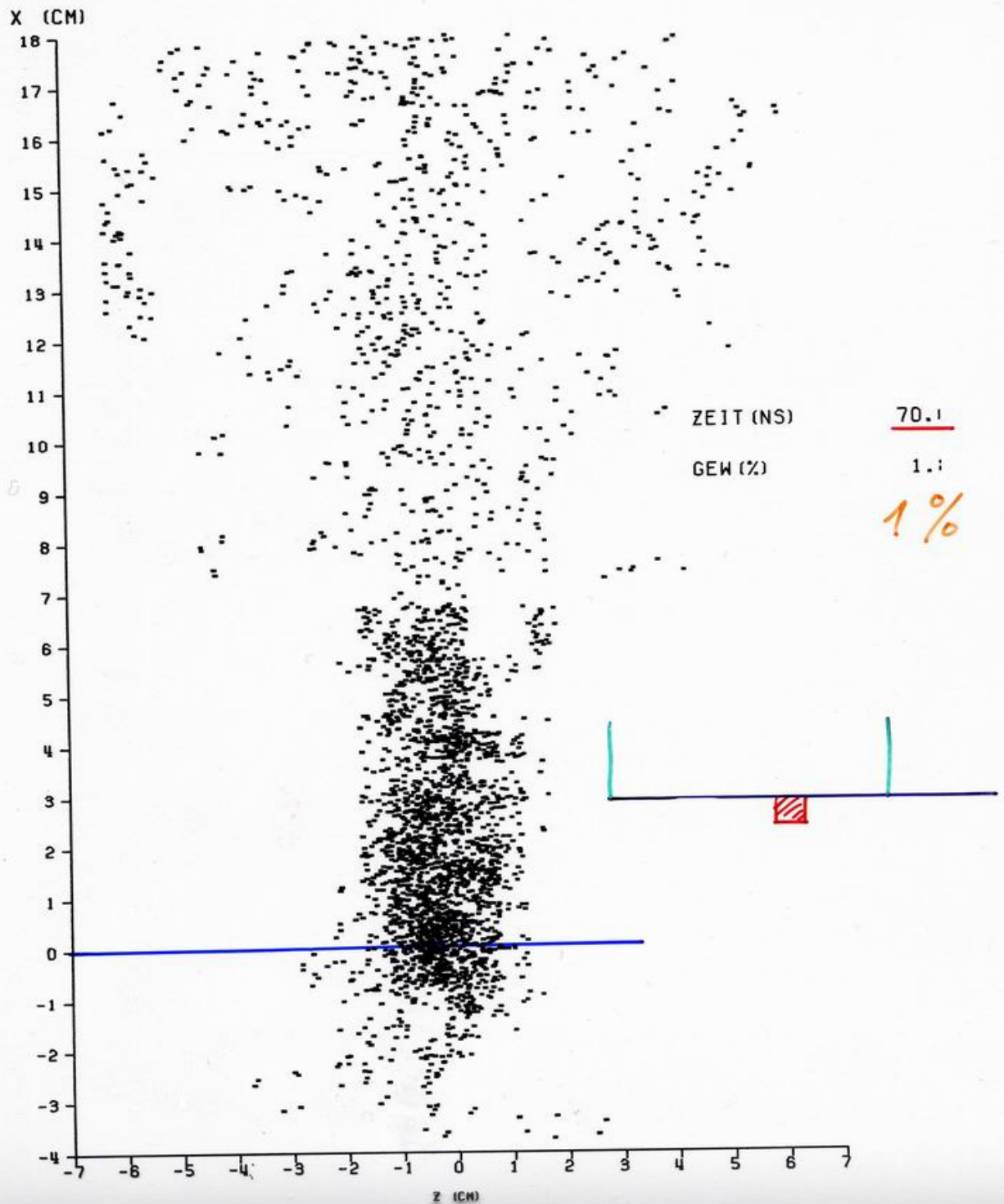
VAR.NR. =	69.00	STRÖM (KA) =	1200.00	STR (FI) =	0.00	DRUCK (MB) =	8.00
D.QU. X0 =	0.02	Z0 (CM) =	0.05	ED (MIN) =	20.00	ED (MAX) =	500.00
TEMAX/KEV =	1.00	TANF (NS) =	0.00	TINSA (NS) =	100.00	TINSE (NS) =	170.00
NI (1) xE19 =	0.80	ND+ (1) xE17xxxxxxxx		ANIS (1) =	1.00	T11 (KEV) =	100.00
NI (2) xE19 =	0.20	ND+ (2) xE17xxxxxxxx		ANIS (2) =	1.00	T12 (KEV) =	400.00



ORTSVERTEILUNG NEUTRONEN UND PROTONEN

POSEIDON

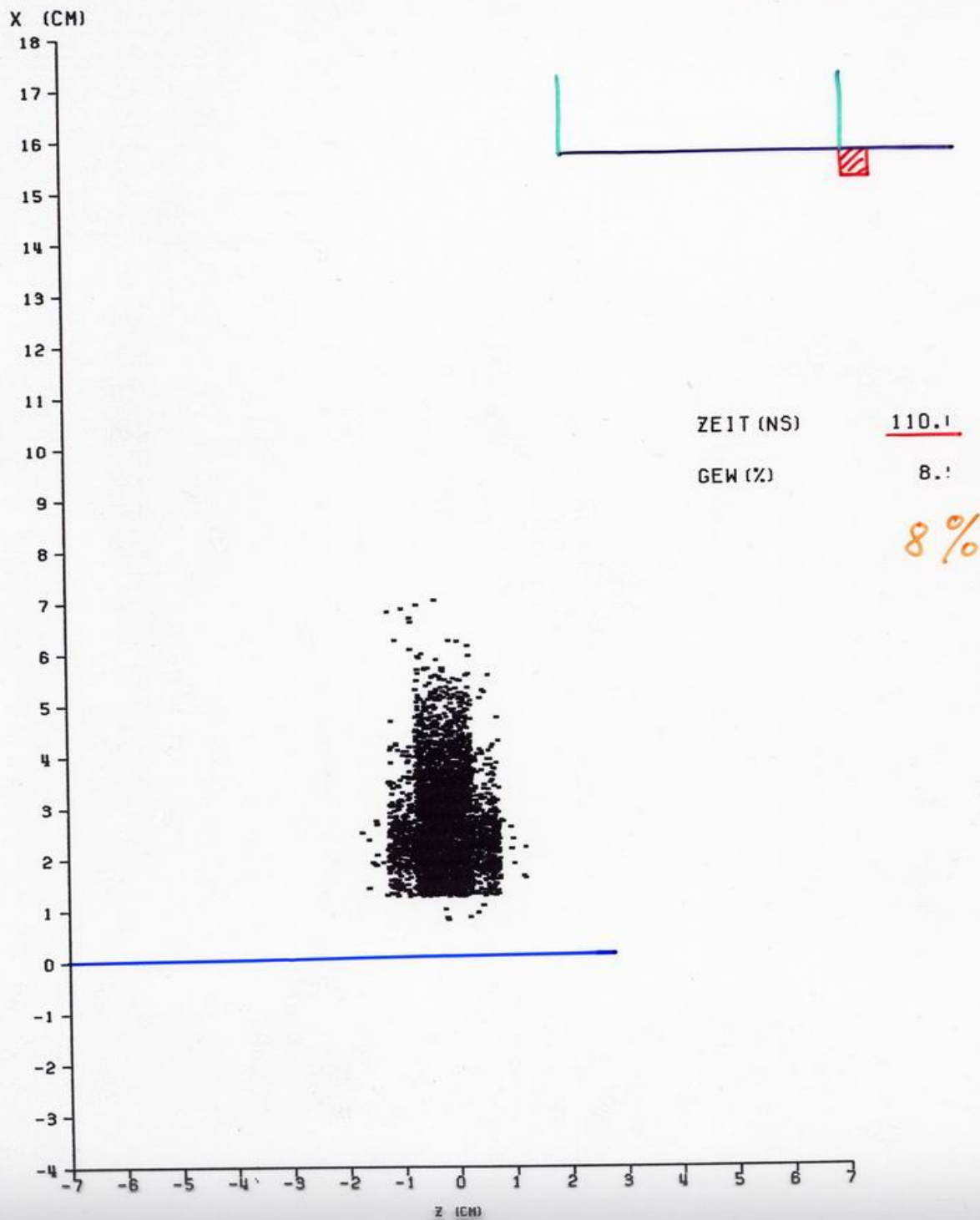
VAR.NR. =	69.00	STROM (KA) =	1200.00	STR (FI) =	0.00	DRUCK (MB) =	8.00
D.QU. XQ=	0.02	ZQ (CM) =	0.05	ED (MIN) =	20.00	ED (MAX) =	500.00
TEMAX/KEV=	1.00	TANF (NS) =	0.00	TINSA (NS) =	100.00	TINSE (NS) =	170.00
NI (1) xE19=	0.80	ND+ (1) xE17xxxxxxxx		ANIS (1) =	1.00	T11 (KEV) =	100.00
NI (2) xE19=	0.20	ND+ (2) xE17xxxxxxxx		ANIS (2) =	1.00	T12 (KEV) =	400.00



ORTSVERTEILUNG NEUTRONEN UND PROTONEN

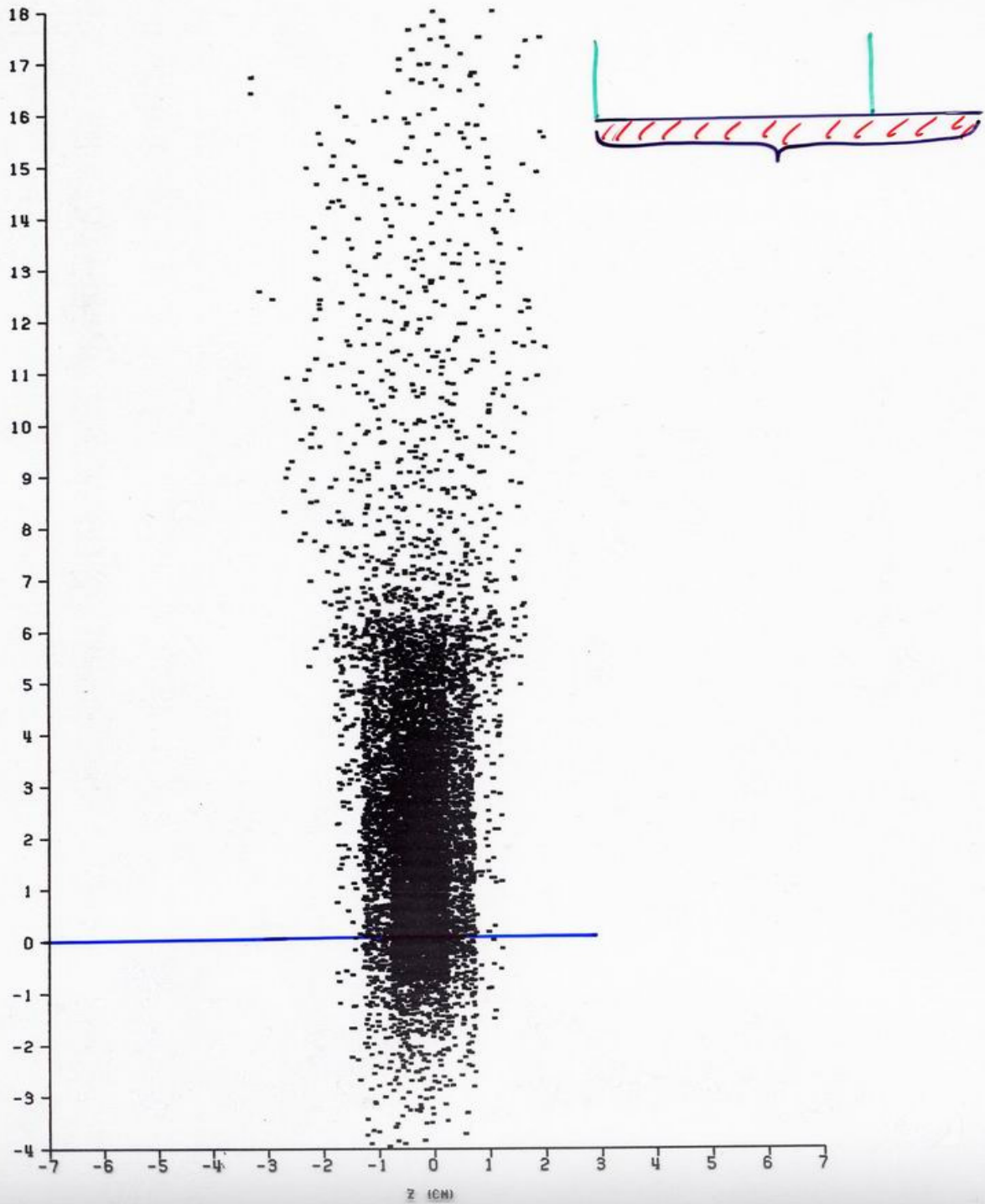
POSEIDON

VAR.NR. =	69.00	STROM (KA) =	1200.00	STR (F1) =	0.00	DRUCK (MB) =	8.00
D.QU. XQ =	0.02	ZQ (CM) =	0.05	ED (MIN) =	20.00	ED (MAX) =	500.00
TEMAX/KEV =	1.00	TANF (NS) =	0.00	TINSA (NS) =	100.00	TINSE (NS) =	170.00
N1 (1) *E19 =	0.80	ND+ (1) *E17 *XXXXXXXXXX		ANIS (1) =	1.00	T11 (KEV) =	100.00
N1 (2) *E19 =	0.20	ND+ (2) *E17 *XXXXXXXXXX		ANIS (2) =	1.00	T12 (KEV) =	400.00

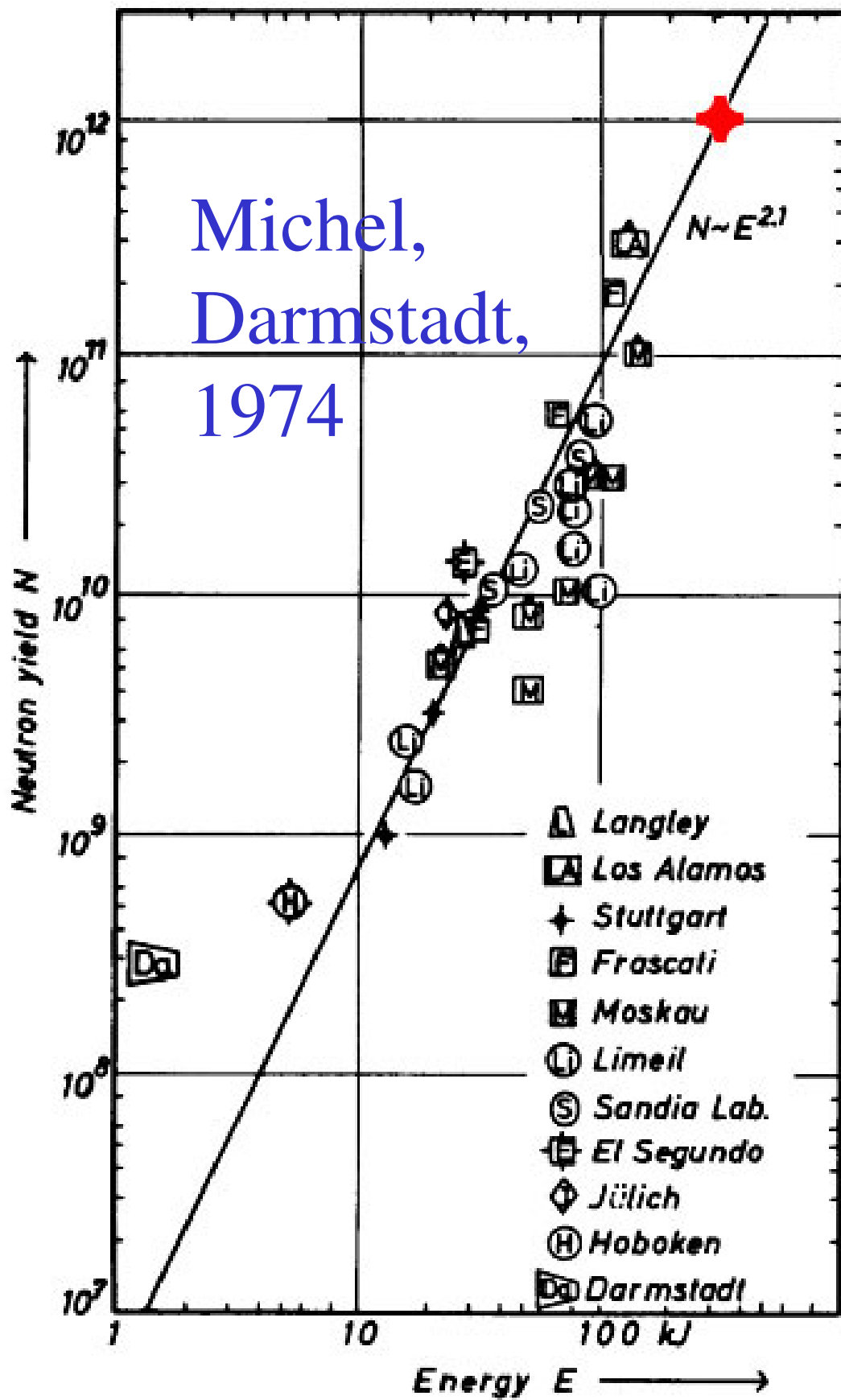


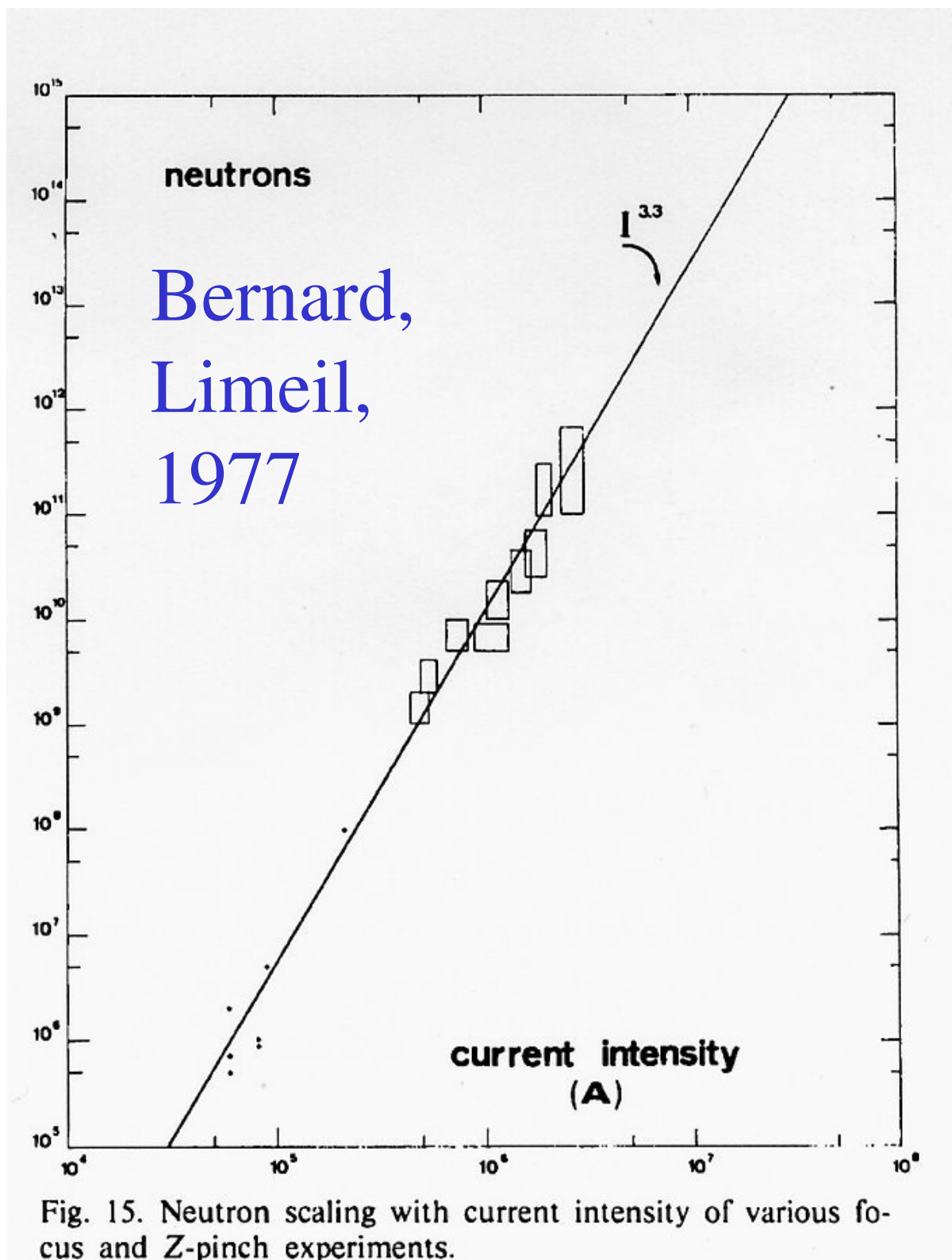
ORTSVERTEILUNG NEUTRONEN UND PROTONEN ZEITINT.

VAR.NR. =	69.00	STROM (KA) =	1200.00	STR (F1) =	0.00	DRUCK (MB) =	8.00
D.QU. XQ=	0.02	ZQ (CM) =	0.05	ED (MIN) =	20.00	ED (MAX) =	500.00
TEMAX/KEV=	1.00	TANF (NS) =	0.00	TINSA (NS) =	100.00	TINSE (NS) =	170.00
NI (1) *E19=	0.80	ND+ (1) *E17*****		ANIS (1) =	1.00	T11 (KEV) =	100.00
NI (2) *E19=	0.20	ND+ (2) *E17*****		ANIS (2) =	1.00	T12 (KEV) =	400.00

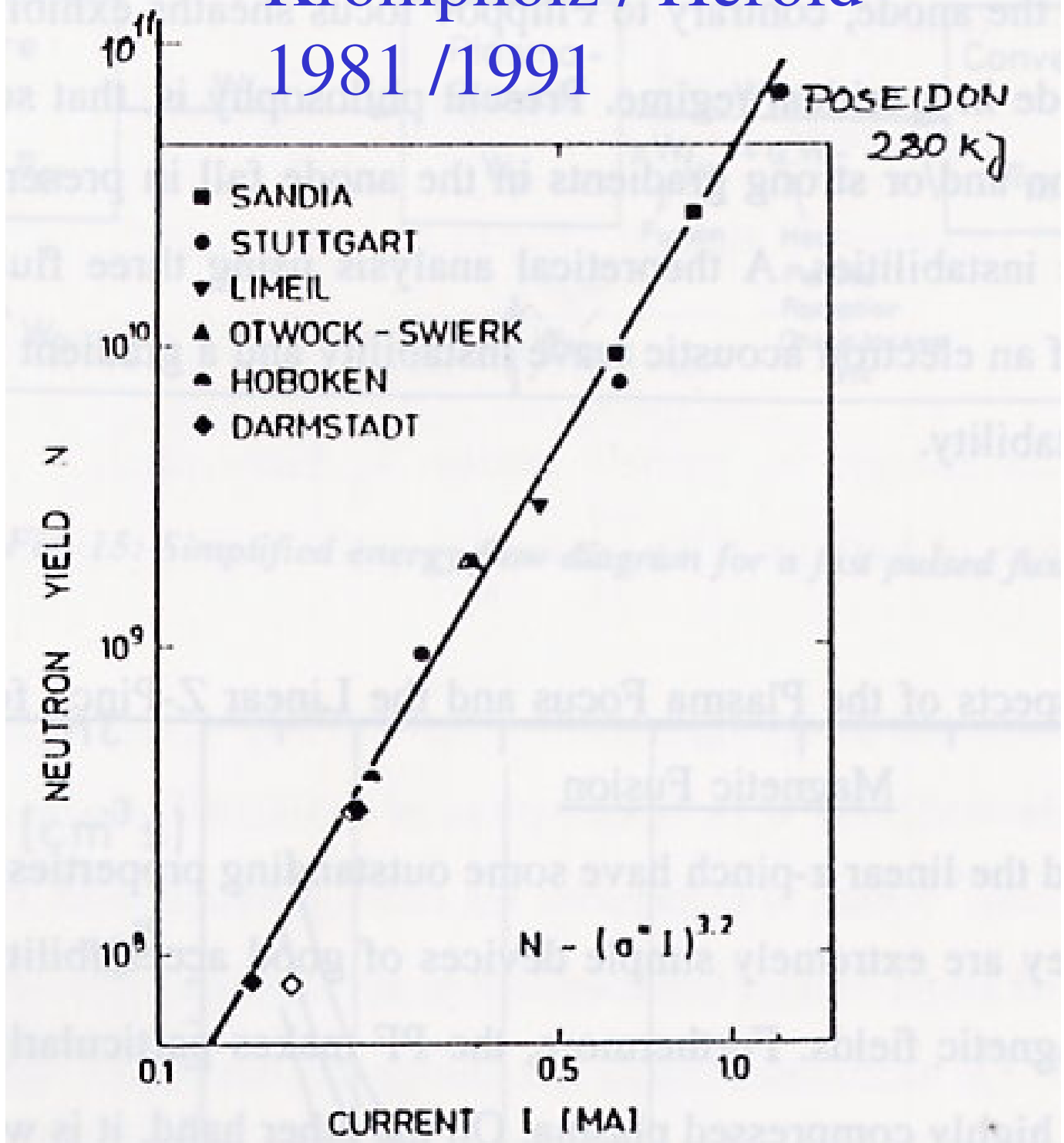


Scaling

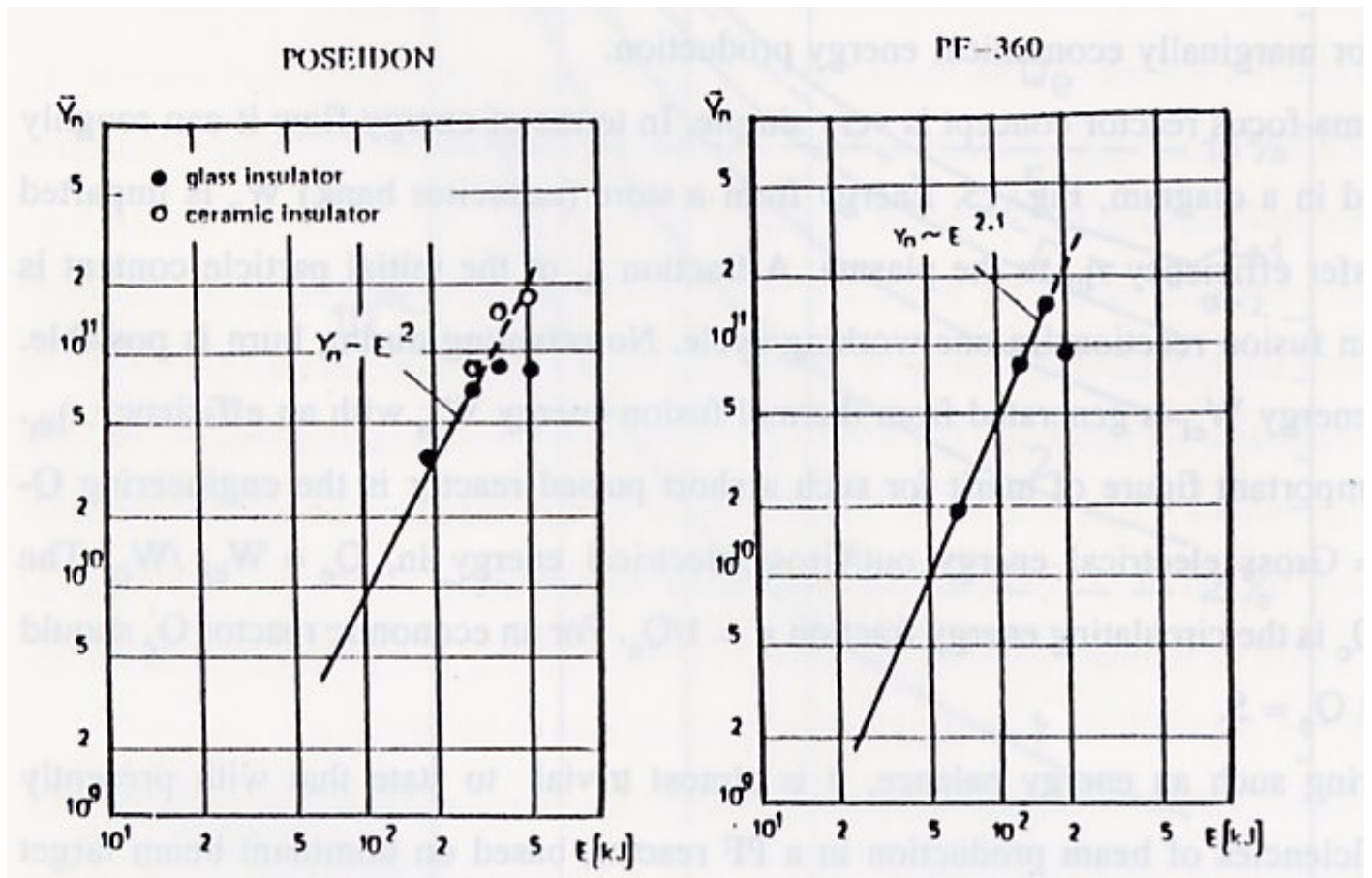




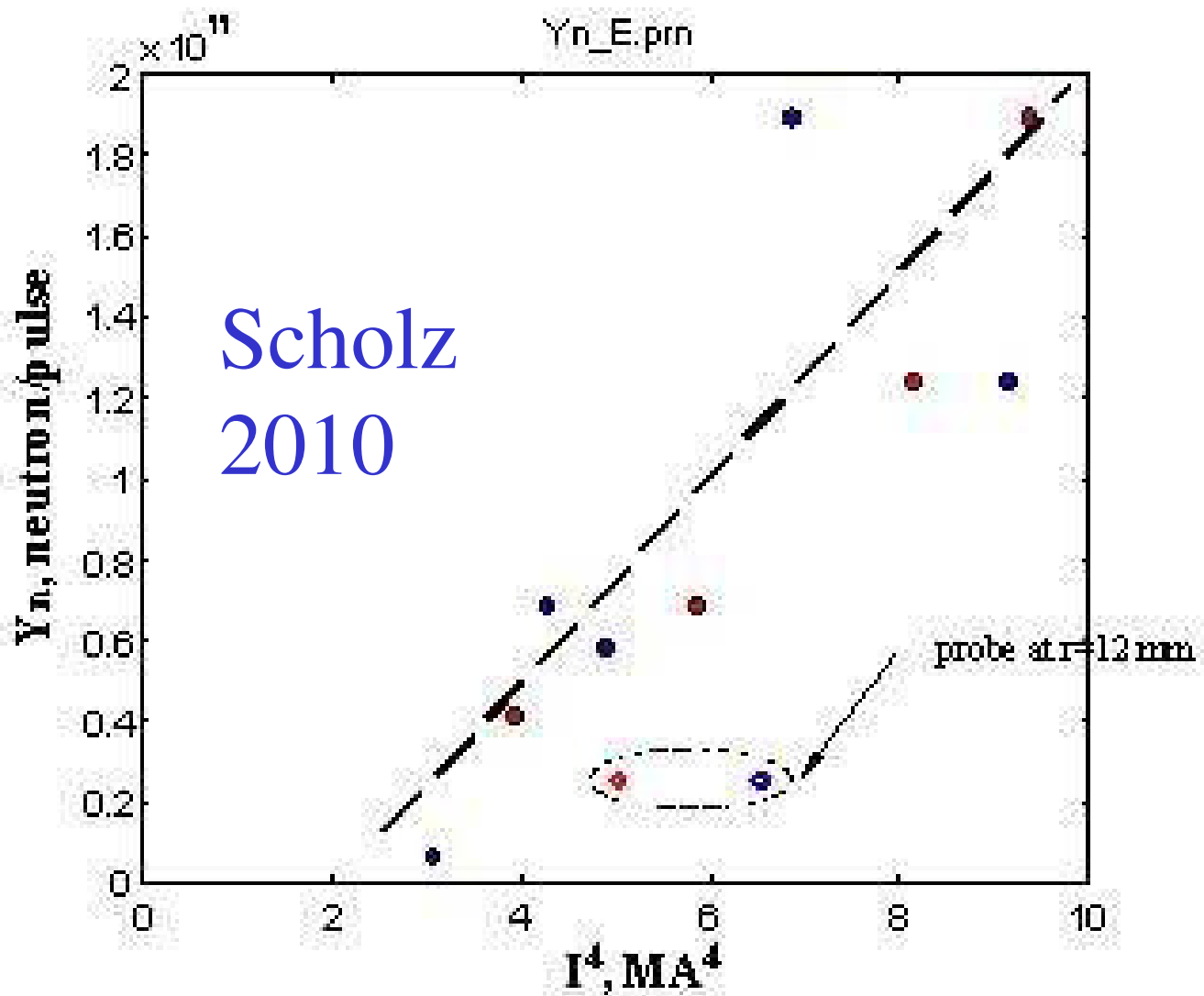
Krompholz / Herold 1981 /1991



Herold 1989



Influence of insulator material!

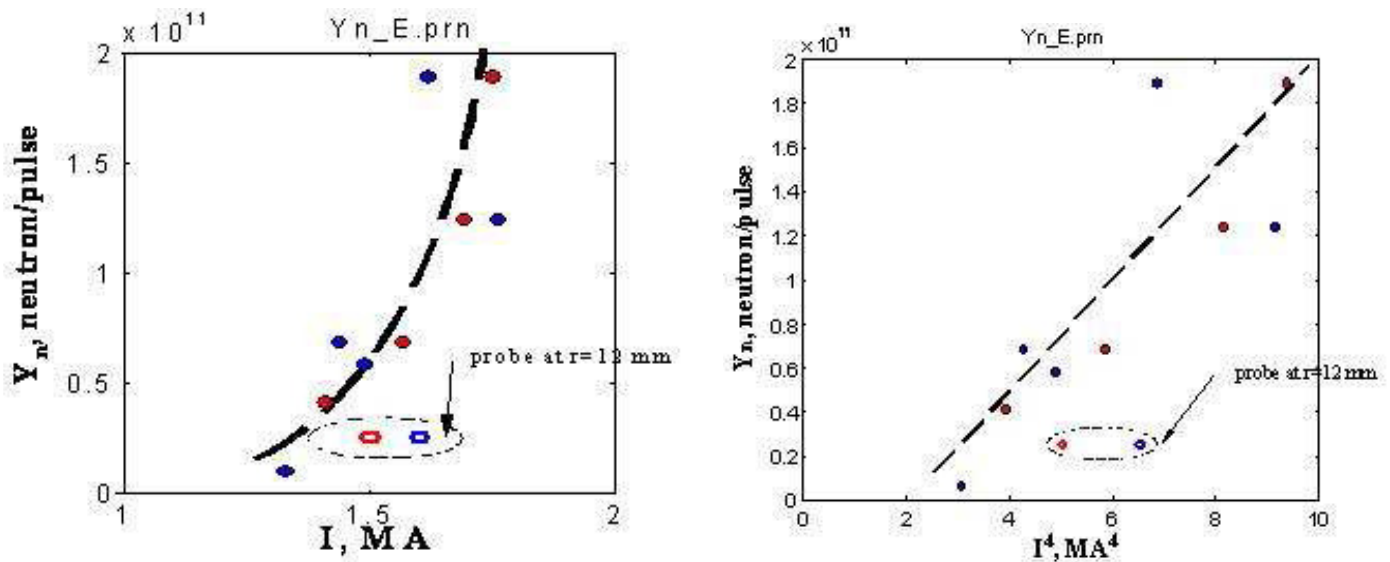


I_c - Current on collector probe, at the max current on ACMP

I_p - Max current on ACMP, measured at 40/12 mm

Scholz 2010

Results for probe position at $r = 40$ mm



I_c - Current on collector probe, at the max current on ACMP

I_p - Max current on ACMP, measured at 40/12 mm

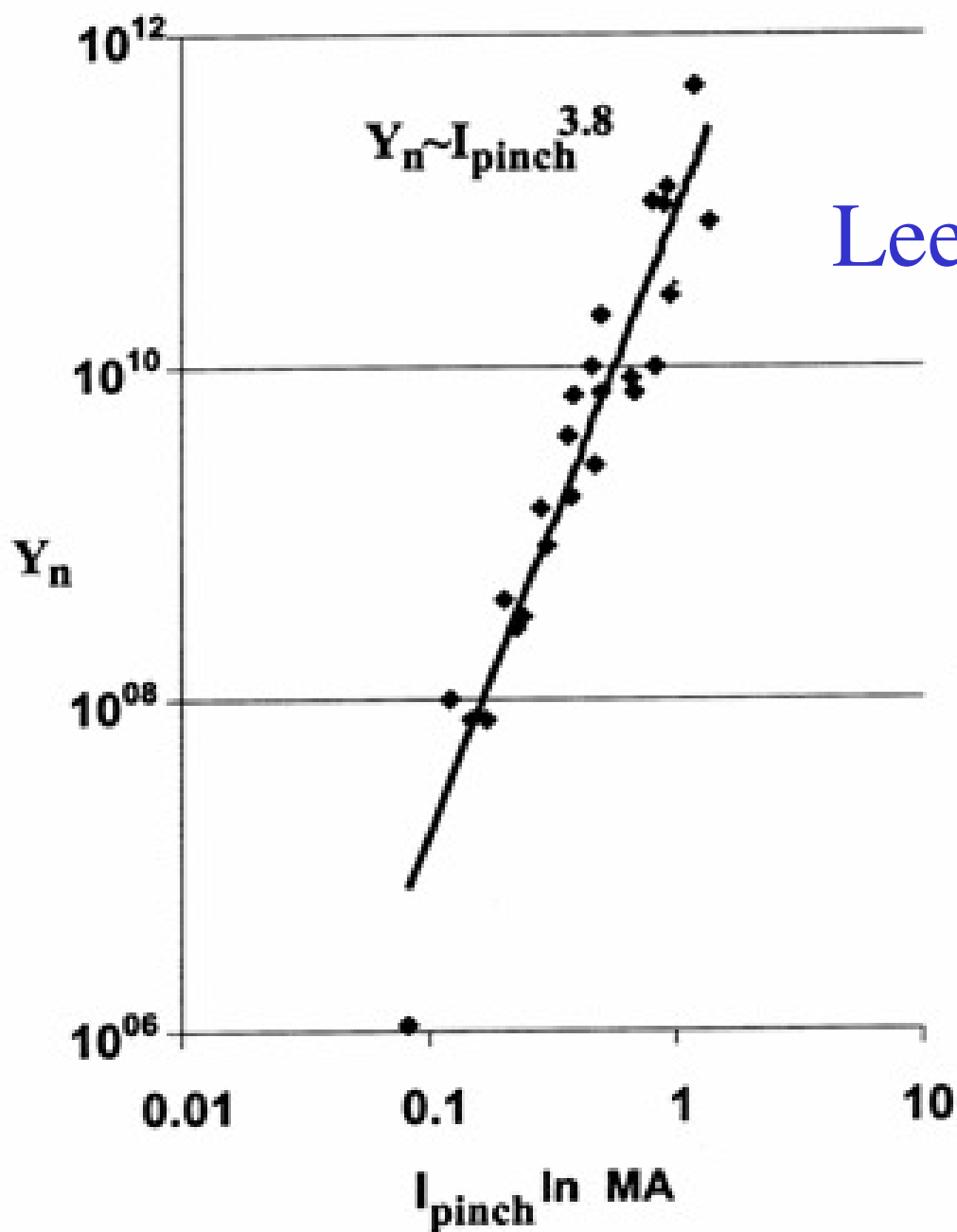
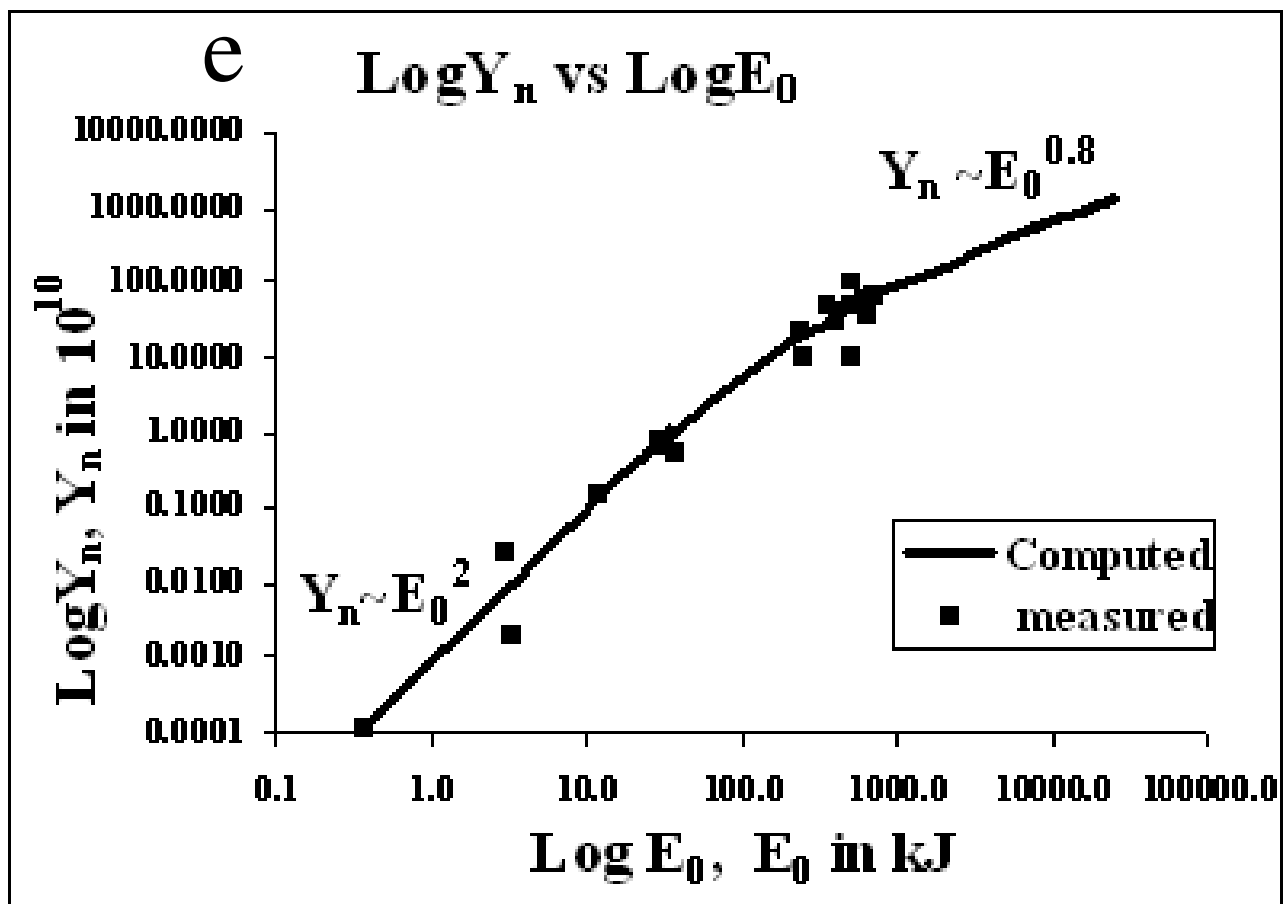
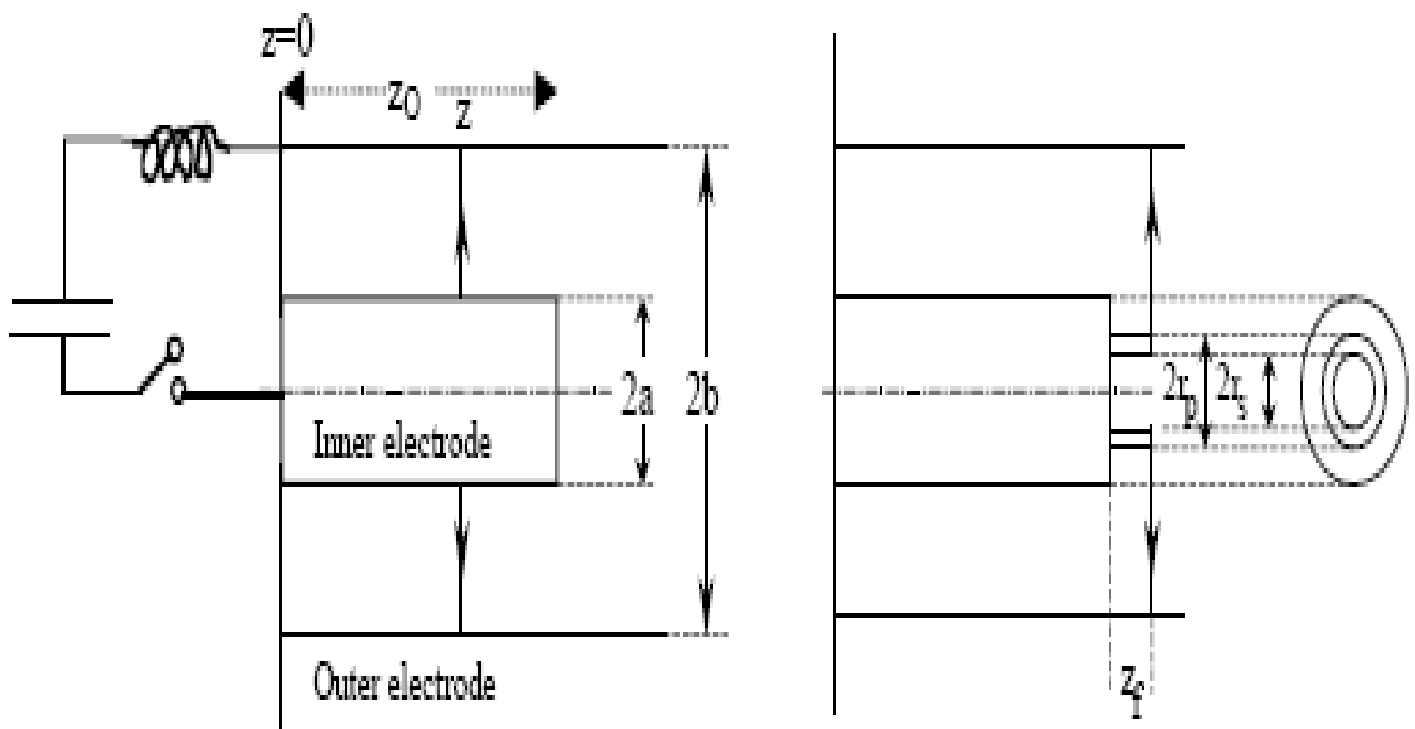


Figure 4. Assembly of experimental data to obtain Y_n scaling with current; loosely termed as the current or pinch current in the literature. This is the experimental curve from which a calibration point is obtained, at 0.5 MA, to calibrate the neutron yield equation (3) for the Lee model code.

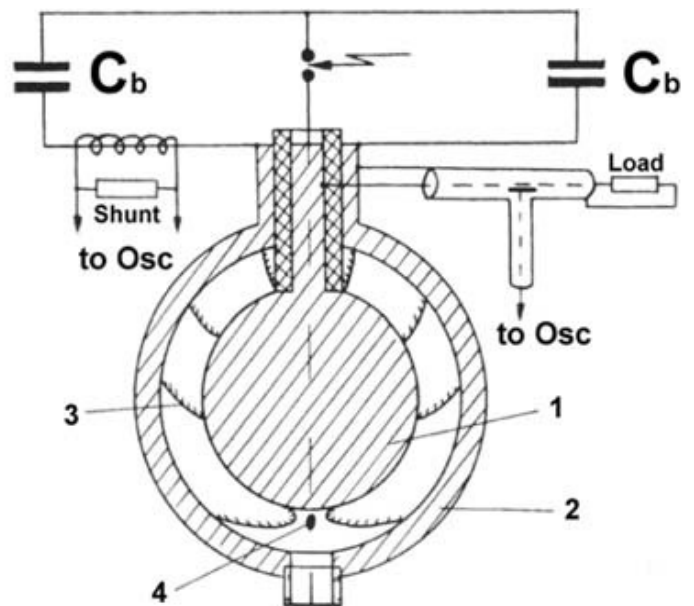
L Lee, 2009
e



Geometry for Lee-Model



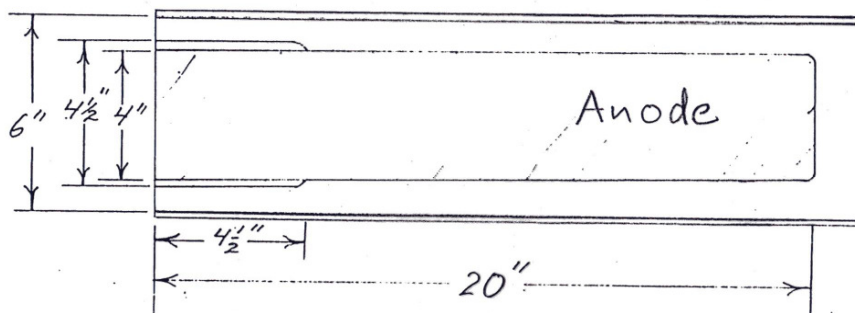
Spherical Geometry – Different results for scaling expected



Kdw

Mather type

Example: DPF-6-1/2 Peak Current (720 kJ Capacitor Bank at 60 kV)



$$\gamma_n \approx 2 \cdot 10^{12}$$

$$C_o = 400 \mu F, L_o = 17.8 mH, R_o = 0.78 m\Omega$$

$$I_{pk} (MA) = \frac{1}{8.75} P_o (Torr D_2)^{1/8} V_o (kV)^{3/4}$$

MEASURED : 3.12 MA AT 50 kV, 20 Torr D_2

CALCULATED, $L(t) = 3.34 MA$

CONST-L = 3.02 MA

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9/16/97 7

ICDMP Visit to Shiva Star Phillips Lab 30.3.1992



Trieste
November 2010

Neutrons in PF
POSEIDON

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Conclusions

- Neutrons from a plasma focus, measured as a function of time, location and direction of emission reveal quite a number of important parameters on fusion reactions occurring in the dense high current phase of the experiment. In addition determination of the energy spectra of the emitted neutrons are important for understanding of the mechanisms taking place for the neutron production. The two main reasons for neutron diagnostics are 1) to determine the fusion yield, i.e. the efficiency of a device and its scaling; and 2) to learn the characteristics of the fusion reactions.
- Correlations of the neutron emission to other properties of a plasma focus such as pinch current, plasma density and temperature as well as electron beam and ion beam emission are of importance as well.
- Methods and results of neutron diagnostics for large experiments such as the former experiment Poseidon in Stuttgart and PF 1000 in Warsaw were presented and discussed.
- Neutron diagnostics presented include nuclear track detectors, plastic scintillators coupled to photomultipliers, activation measurements, time-of-flight methods as well as pinholes for spatial resolution of the neutron source.
- The well known scaling law according to which neutron yield scales roughly with the square of the energy input and the fourth power of the current was commented. Reasons for strong deviations from this law for high energy – known as saturation – are still a subject of debate.

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