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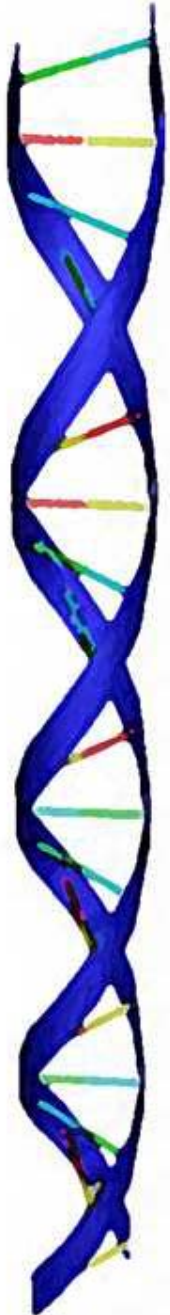
2168-Presentation

**Joint ICTP-IAEA Workshop on Dense Magnetized Plasma and Plasma
Diagnostics**

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Hydromagnetic Instabilities in Magnetized Plasmas

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Hydromagnetic Instabilities in Magnetized Plasmas

by

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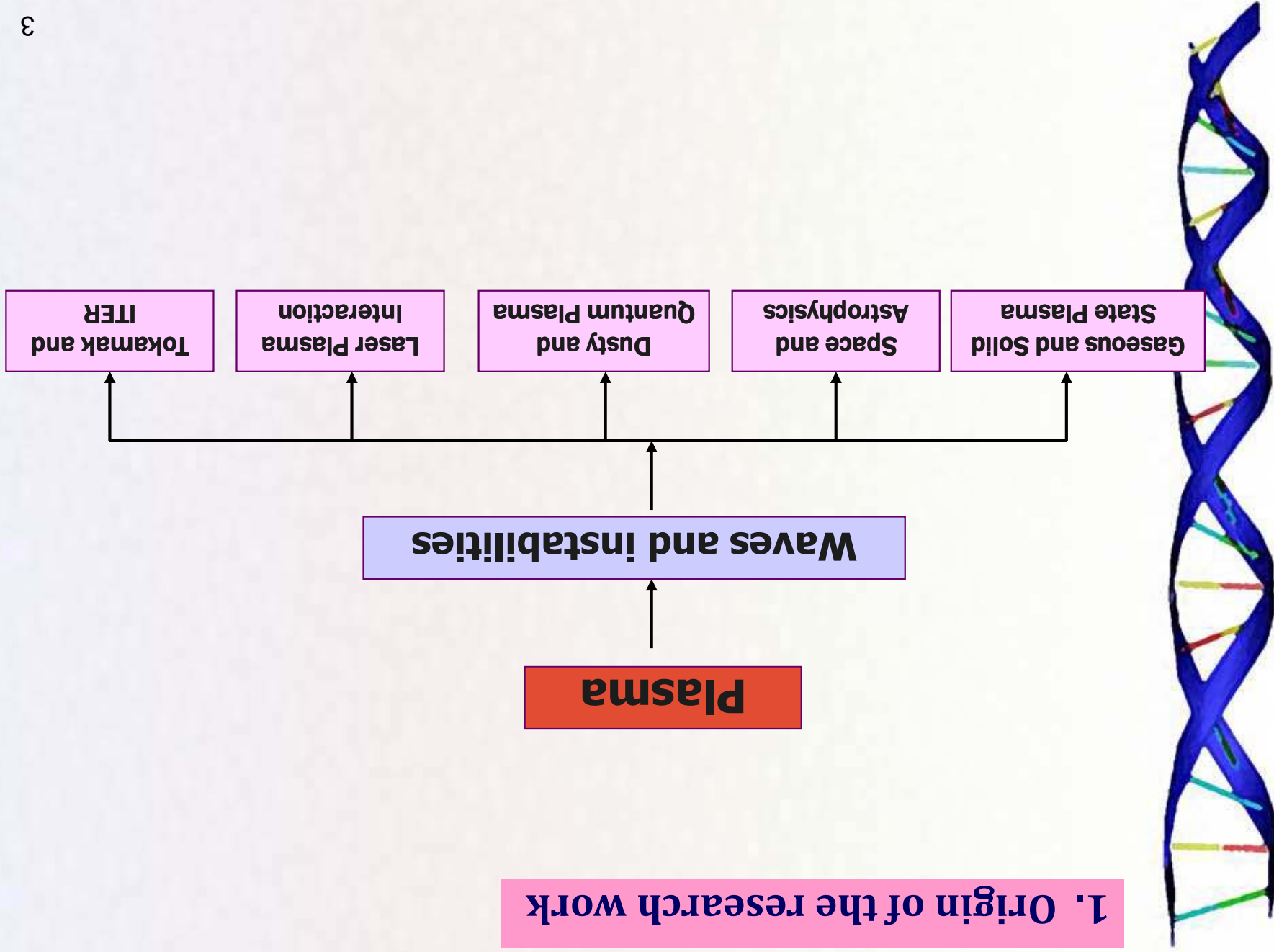
**Joint ICTP/IAEA Workshop on Dense Magnetized Plasma & Plasma Diagnostics,
15-26 November 2010, ICTP Trieste (Italy).**

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1. Origin of the research work



2. Instabilities in plasma

Because of free-energy sources in plasma, a very large number of instabilities can develop.

If the involved scale is:

- Comparable to macroscopic size (bulk scale of plasma,.....)
- Macroinstability** (affects plasma globally)
- Comparable to microscopic scale (gyroradius, inertial length,...)
- Microinstability** (affects plasma locally)

Theoretical treatment:

- # Macroinstability** ← Fluid plasma theory (Single and two fluid)
- # Microinstability** ← Kinetic plasma theory



Nonlinear instability

The nonlinearities come from the harmonic generation involving nonlinear Lorentz force, trapping of particles in the wave potential, ponderomotive force etc.

Coherent structures as;
 # Solitary structure
 # Shock waves
 # Vortices

Linear instability

The concept of linear instability arises from the consideration of a linear wave function. Assume any variable (density, magnetic field, etc.) here denoted by A , the fluctuation of which is δA , that can be Fourier decomposed as

$$\delta A = \sum A \exp(ik \cdot x - i\omega t)$$

In general the solution is given as

$$\omega = \omega_r + \gamma$$

For real frequencies disturbances are oscillating waves.

For complex solution
 Amplitude A will grow if $\gamma > 0$

and amplitude A will decay if $\gamma < 0$



3. Methodology and technique

➤ Normal Mode Analysis

Physical quantity = Equilibrium part + Perturbed part

$$\xi = \xi_0 + \xi_1$$

Fourier transform of the perturbed quantity

$$\xi_1 = \xi_1 \exp(\omega t - i k \cdot r)$$

1. If all the values of ω are purely imaginary: System is said to be stable.
2. If some or all the values of ω are real and positive: System is said to be unstable.
3. If some or all the values of ω are complex i.e. $\omega = \omega_r + i \omega_i$

Then for stability $\omega_r < 0$

and for instability $\omega_r > 0$

Routh-Hurwitz criterion for stability analysis:

$$F(\omega) = a_0 \omega^n + a_1 \omega^{n-1} + a_2 \omega^{n-2} + \dots + a_n = 0.$$





➤ Fluid plasma theory

Magnetospheres of planets # stars # Extragalactic jets # Comets tails

The fluid model describes plasma in terms of quantities like density and average velocity around each position.

Single fluid theory

Plasma as single fluid governed by Maxwell's and fluid equations

Two fluid theory

It consists dynamics for each species (electron and ion)

❖ MHD fluid theory (Collisional plasma):

The collisions are so frequent

➤ The MHD theory encloses the regime of scalar pressures for both electrons and ions.

➤ The MHD model can be applied only for low frequency phenomena providing the concept of Alfvén wave and magneto sonic waves.



❖ **CGL fluid theory (Collisionless plasma):**

[Chew et al., Proc. Roy. Soc. A 236, 112 (1956)]

- # Solar wind
- # Interstellar Space
- # Nebula

Multiple Coulomb force dominates over the Charge neutral interaction force

In the absence of collisions the usual scalar pressure is replaced by a pressure tensor due to the presence of strong magnetic field

$$\mathbf{P} = p_{\perp} \mathbf{I} + (p_{\parallel} - p_{\perp}) \mathbf{n} \mathbf{n}$$

where \mathbf{l} is unit dyadic and \mathbf{n} is the unit vector along the magnetic field.

The pressure equations in terms of pressure component, magnetic field and density are given as

$$\frac{d}{dt} \left(p_{\parallel} B^2 \right) = 0 \quad \text{and} \quad \frac{d}{dt} \left(\frac{p_{\perp}}{B} \right) = 0.$$



❖ Polytropic model:

[B. Abraham-Shrauner, Plasma Phys. 15, 375 (1973)]
 [M. Chou and L. N. Hau, Astrophys. J. 611, 1200 (2004)]

MHD approximation



Collision dominated plasma

CGL approximation



Collisionless plasma in presence
 of strong magnetic field

Polytropic model



Transition zone where plasma is
 neither fully collisional nor
 collisionless

$$0 = \frac{d}{dt} \left(\frac{\rho_\epsilon B_\nu}{d_\perp} \right)$$

$$0 = \frac{d}{dt} \left(\frac{\rho_\beta B_\alpha}{d_\parallel} \right)$$

Where α, β, ϵ and ν are the polytropic indices.

4. Dusty and quantum plasma

A dusty plasma is collection of free electrons and ions with some additional charged dust grains of micron size. A dusty plasma also satisfies the usual quasineutrality condition as

$$n_i^0 = n_e^0 + q^d n^d{}^0$$

In classical plasma

$$\lambda_B \ll \lambda_D$$

$$\lambda_B = \frac{h}{mV^T}$$

i.e. de-Broglie wavelength < Debye length of the system

In case of quantum plasma

$$\lambda_B \gg \lambda_D$$

Quantum statistical effects
(spin magnetization, particles
spin effect etc.)

Quantum Bohm potential
(arises due to low temperature)

- Nanoscale technology [H. G. Craighead, Science 290, 1532 (2000)]
- Microelectronic devices [P.A. Markovich et al., S/C Eqn., Springer-Verlag NY 1990]
- Laser Fusion [S. H. Glenzer et al., Phys. Rev. Letters, 98, 065002 (2007)]
- Dense Astrophysical System [C. Kouveliotou et al., Nature, 393, 235 (1998)]



Publications in dusty and quantum plasma

1. Effect of dust temperature on radiative condensation instability of self-gravitating magnetized dusty plasma
R. P. Prajapati and R. K. Chhajlani, *Physica Scripta* 81, 045501 (2010)

2. Effect of Hall current on Jeans instability of magnetized quantum viscous plasma
R. P. Prajapati and R. K. Chhajlani, *Physica Scripta* 82, 055003 (2010)

3. Effect of magnetic field on Jeans instability of quantum dusty plasma: Application in white dwarf star
R. P. Prajapati and R. K. Chhajlani, *Acta Technica* (2010) [In Press]





5. Some hydromagnetic instabilities

5.1 Jeans (gravitational) instability

Any self-gravitating object opposes the excess gas pressure due to self-gravitation. This causes collapse of the object and triggers an instability called **Self-gravitational or Jeans instability**.

Astrophysics, Star formation, ISM formation, Nebula, Molecular cloud formation, Dwarf star and Neutron star formation etc.

Chandrasekhar (1961) has given the Jeans instability criterion

$$\rho \frac{\partial \mathbf{u}}{\partial t} = -c_s^2 \Delta \delta \rho + \rho \Delta \delta \phi$$

(Momentum conservation)

$$\Delta^2 \delta \phi = -4\pi G \delta \rho$$

(Poisson's equation)

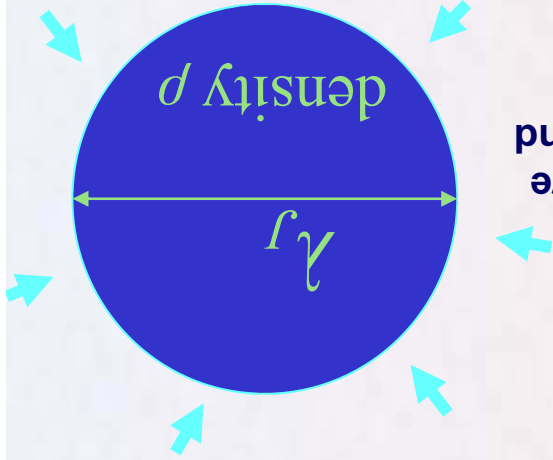
$$\frac{\partial}{\partial t} \delta \rho + \rho \Delta \cdot \mathbf{u} = 0$$

The system will be unstable for all the wavenumbers

$$k > k_j = \frac{(4\pi G \rho)^{1/2}}{c_s}$$

S. Chandrasekhar, *Hydromagnetic and Hydrodynamic Instability*, Oxford (1961).

Mass ≈ 1.4 solar masses



Publications of Jeans instability

1. Self-gravitational instability of rotating anisotropic heat-conducting plasma
R. P. Prajapati et al., *Physics of Plasmas* **15**, 012107 (2008)

2. Self-gravitating rotating anisotropic pressure plasma in presence of Hall current and electrical resistivity with generalized polytropic laws
R. P. Prajapati et al., *Physics of Plasmas* **15**, 062108 (2008)

3. Self-gravitational instability of rotating viscous Hall plasma with arbitrary radiative heat-loss functions and electron inertia
R. P. Prajapati et al., *Astrophysics & Space Science* **327**, 139 (2010)

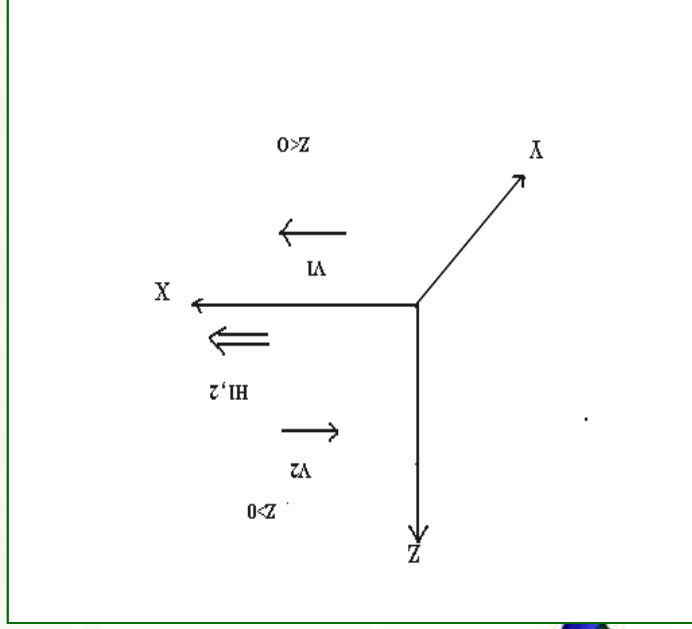


5.2 Kelvin-Helmholtz (K-H) instability

- The instability arising due to the tangential discontinuity of velocities between two plasma streams is called the K-H instability.
- Two superposed fluids are separated by a horizontal boundary which are in relative motion then the instability arising at the interface is called the K-H instability.
- Velocity shear is present.
- There is sufficient velocity difference across the interface between two fluids.

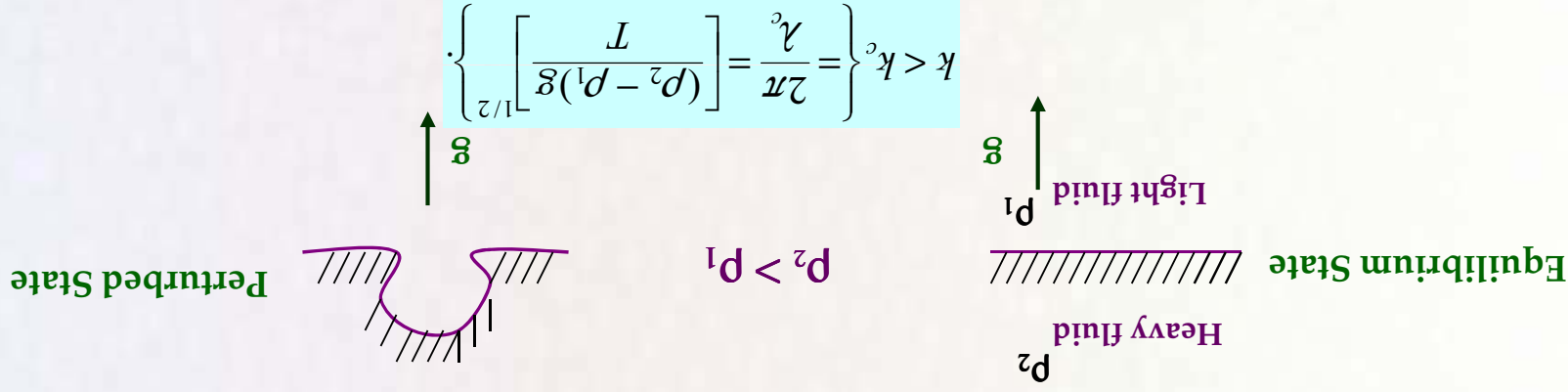
$$k > k_{\min} \left[= \frac{g(\rho_1 - \rho_2)}{\rho_1 \rho_2 (U_1 - U_2)^2} \right]$$

- Space plasma
- Comet tails
- Astrophysical plasmas
- Tokamak & ITER
- Industrial plasma



5.3 Rayleigh-Taylor (R-T) instability

The instability at the interface of two superposed fluids of different densities in which the heavy fluid is supported by the light one is generally known as R-T instability.



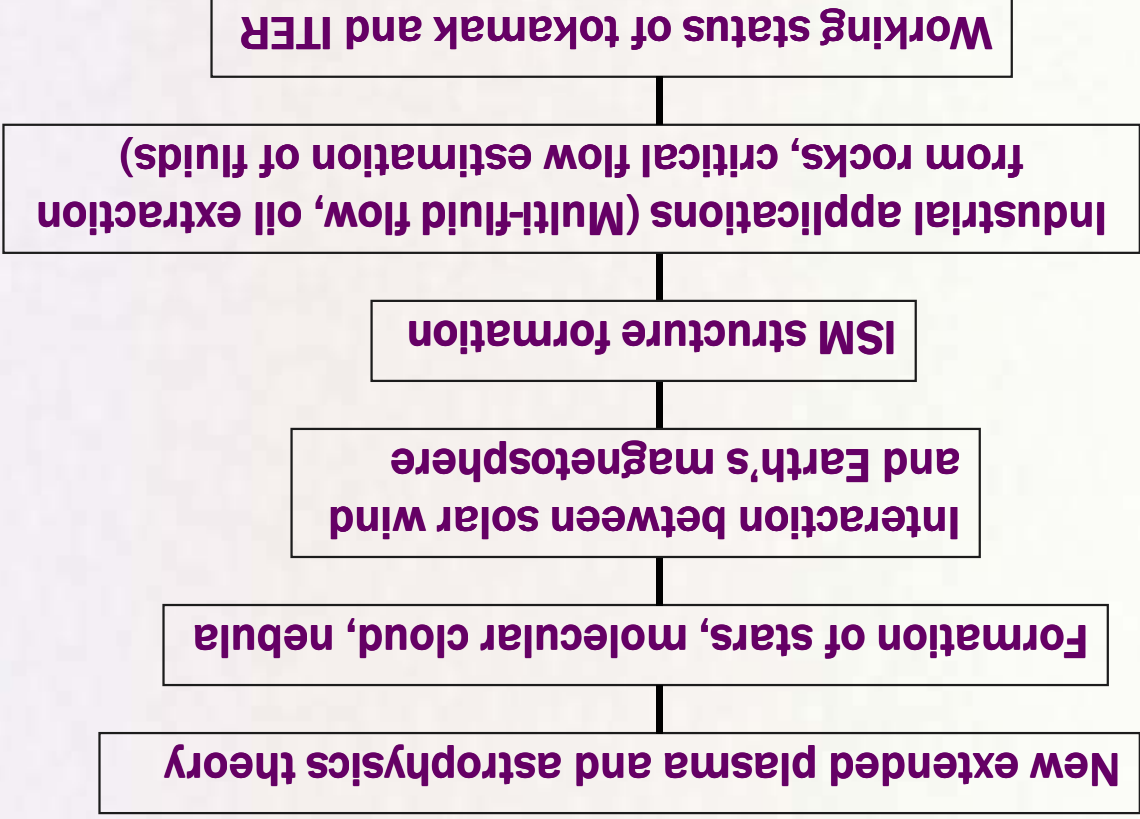
- Inertial Confinement Fusion (ICF)
- Supernova Explosions
- Z-pinches
- Astrophysical plasmas



Publications in K-H and R-T instabilities

1. Kelvin-Helmholtz and Rayleigh-Taylor instability of two superposed magnetized incompressible fluids with suspended dust particles
R. P. Prajapati et al., *Z. Naturforsch A* **64a**, 455 (2009).
2. Kelvin-Helmholtz instability of magnetized plasmas with surface tension and dust particles
R. P. Prajapati and R. K. Chhajani, *J. Physics Conf. Ser.* **208**, 012078 (2010).
3. Kelvin-Helmholtz instability of anisotropic pressure plasma using generalized polytropic laws
R. P. Prajapati et al., *J. Phys. Conf. Ser.* **208**, 012077 (2010).
4. Kelvin-Helmholtz and Rayleigh Taylor instability of two superposed fluids with suspended dust particles flowing through porous media
R. P. Prajapati and R. K. Chhajani, *Journal of Porous Media* **13**, 765 (2010).
5. Rayleigh-Taylor instability of two superposed magnetized fluids with suspended dust particles.
R. P. Prajapati et al., *Thermal Science* **14**, 11 (2010).
6. Effect of surface tension and rotation on Rayleigh-Taylor instability of two superposed fluids with suspended dust particles
P. K. Sharma, R. P. Prajapati and R. K. Chhajani, *Acta Physica Polonica* **118**, 576 (2010).
7. Effect of pressure anisotropy and flow velocity on Kelvin-Helmholtz instability of anisotropic plasma using generalized polytropic laws
R. P. Prajapati and R. K. Chhajani, *Physics of Plasmas* **17**, 1 (2010).

6. Out comes



Thank You

