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Technology: Analytical Applications**

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**Neutron Self-Shielding in the Thermal Range**

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Things used to be so simple....

# Neutron self-shielding in the thermal range

# Overview

- definition of self-shielding
- principles of cross section dependence on velocity
- self-shielding at 2200 m/s
  - absorption only
  - with scattering
- self-shielding in a thermal neutron spectrum
  - absorption only
  - with scattering
- self-shielding in reality
  - three experiments
- Conclusions

# A definition of 1/v self-shielding

$$f = \frac{R}{V\Sigma_{0,a}\Phi_0} = \frac{R}{N\sigma_{0,a}\Phi_0}$$

where

f is the self-shielding factor

R is the activation rate (captures per second)

V is the sample volume

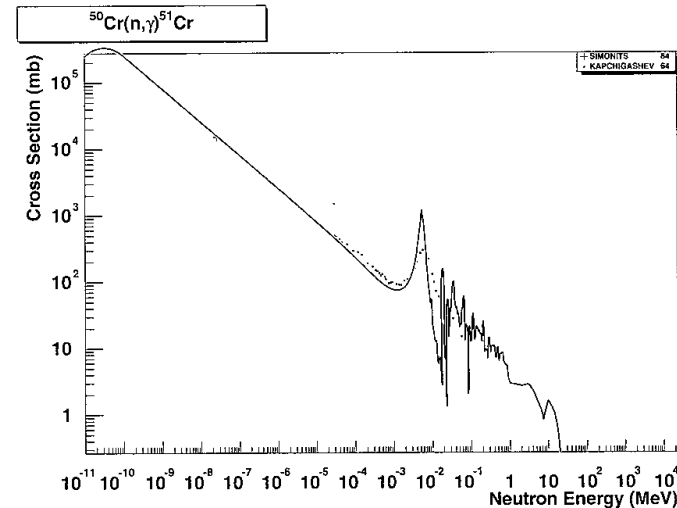
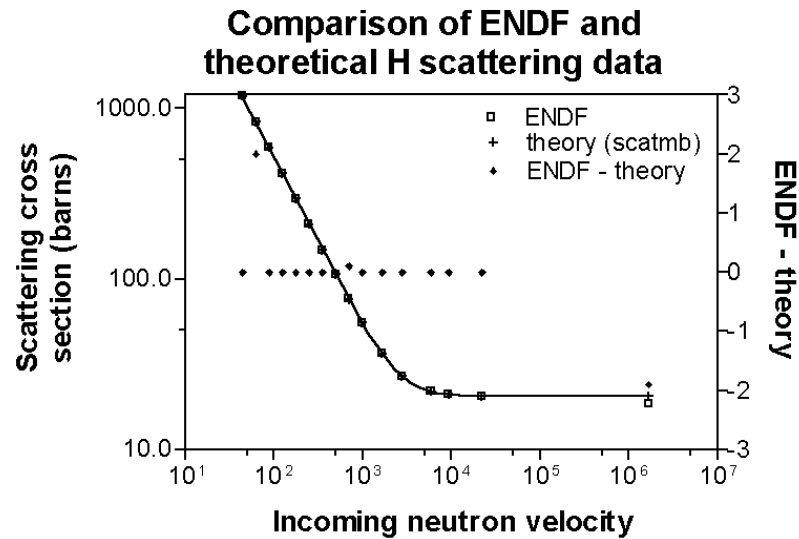
$\Sigma_{0,a}$  is the macroscopic absorption cross section ( $\text{m}^2/\text{m}^3$ )

$\Phi_0$  is the incoming conventional neutron flux  $n v_0$  ( $\text{m}^{-2}\text{s}^{-1}$ )

N is the number of atoms

$\sigma_{0,a}$  is the microscopic absorption cross section ( $\text{m}^2$ )

# Absorption and scattering as a function of neutron velocity

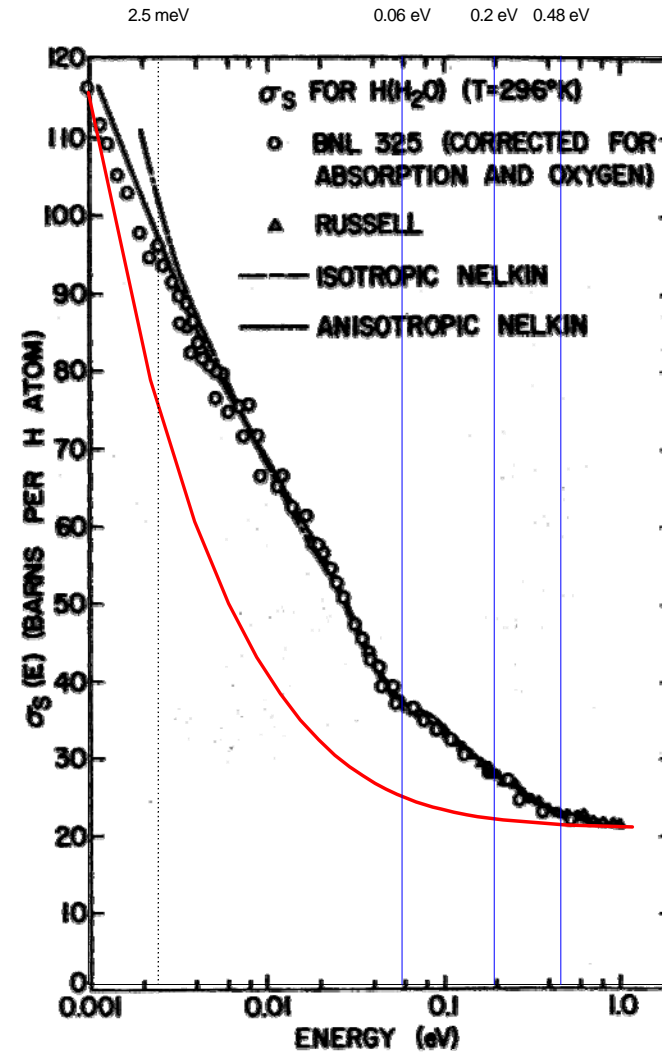


Interaction rate

# Free-gas and the real thing

BNL 325, in  
“Reactor Physics”,  
proc. ANS topical  
meeting, MIT press  
(1966),  
and the same data as  
on the previous sheet  
for the **M=1, free-gas  
model**.

The vertical blue  
lines indicate the  $h\nu$   
of  $H_2O$ .



# Beckurts & Wirtz

“Neutron Physics”

by

K.H. Beckurts,

K. Wirtz

(Springer Verlag,  
1958), probably from

E. Fermi, “Sul moto

dei Neutroni nelle

sostanza idrogenate”,

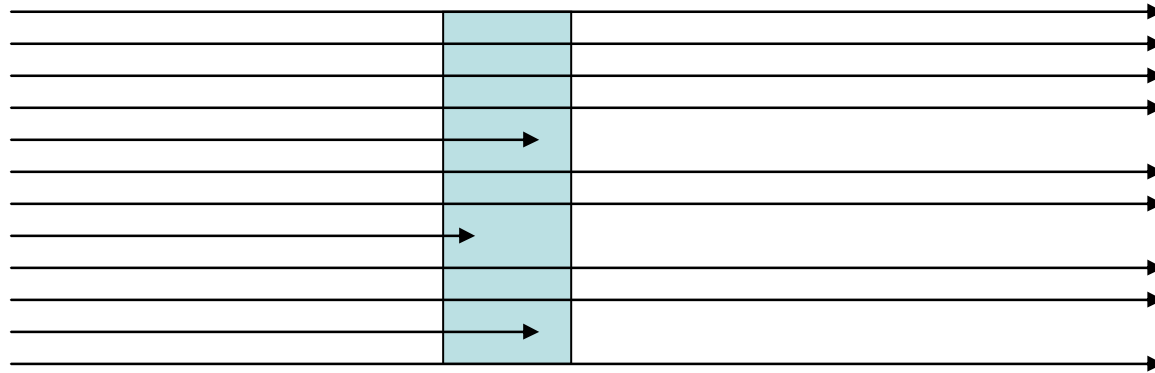
Ricerca Scientifica

7(2) 1936.

*The ratio of the bound atom scattering cross section to the free atom scattering cross section is the square of the ratio of the corresponding reduced masses.*

$$\frac{M_{eff}}{1 + M_{eff}} \frac{1 + M}{M} = \sqrt{\frac{\sigma(M_{eff})}{\sigma_0}}$$

## At 2200 m/s, absorption only, in a beam

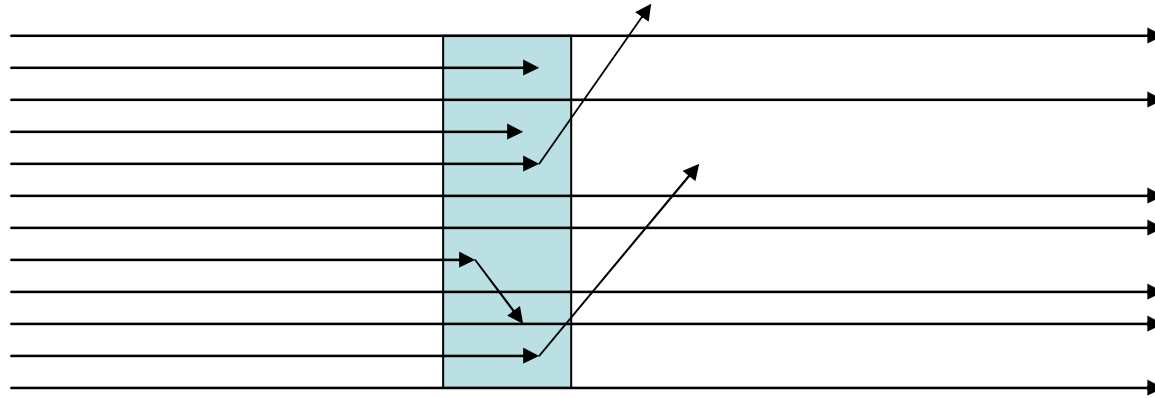


$$f = \frac{(1 - e^{-\Sigma_{0,a}d})}{\Sigma_{0,a}d}$$

This equation for slab only, others available for other shapes



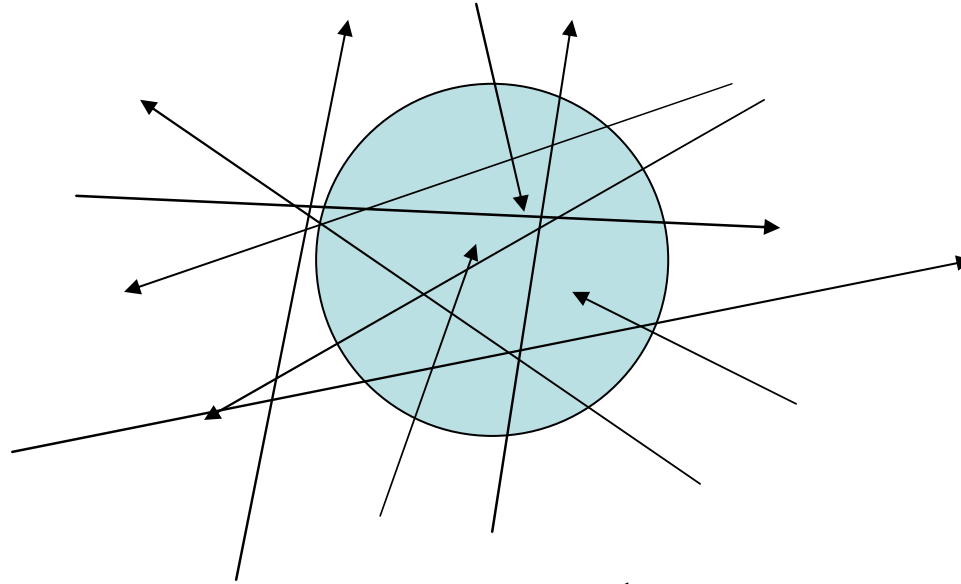
## At 2200 m/s, absorption and scattering, in a beam



Average neutron path length in sample may increase.

That means  $f > 1$  !!!!!!!

## At 2200 m/s, absorption only, isotropic



$$\Psi = R\Sigma_a \quad f = \frac{3}{4\Psi^3} \left( \Psi^2 + \left( \frac{1}{2} + \Psi \right) e^{-2\Psi} - \frac{1}{2} \right)$$

This equation for sphere only, others available for other shapes

# Stuart's formula

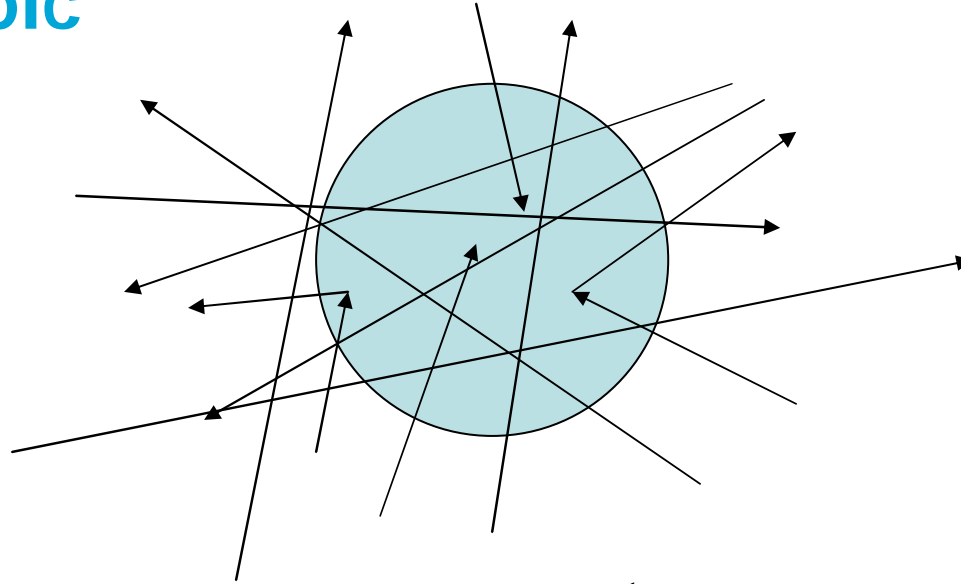
For use in thermal isotropic flux only

$$f = \frac{f_0}{1 - \frac{\sum_s}{\sum_t} (1 - f_0)}$$

$f_0$  to be calculated using the standard formula for the object shape in question and the *total* macroscopic cross section.

M. Blaauw, "The derivation and proper use of Stewart's formula for thermal neutron self-shielding in scattering media", Nuc.Sci.Eng., **124** (1996) 431-435

## At 2200 m/s, absorption and scattering, isotropic



$$\Psi = R\Sigma_t \quad f_0 = \frac{3}{4\Psi^3} \left( \Psi^2 + \left( \frac{1}{2} + \Psi \right) e^{-2\Psi} - \frac{1}{2} \right)$$

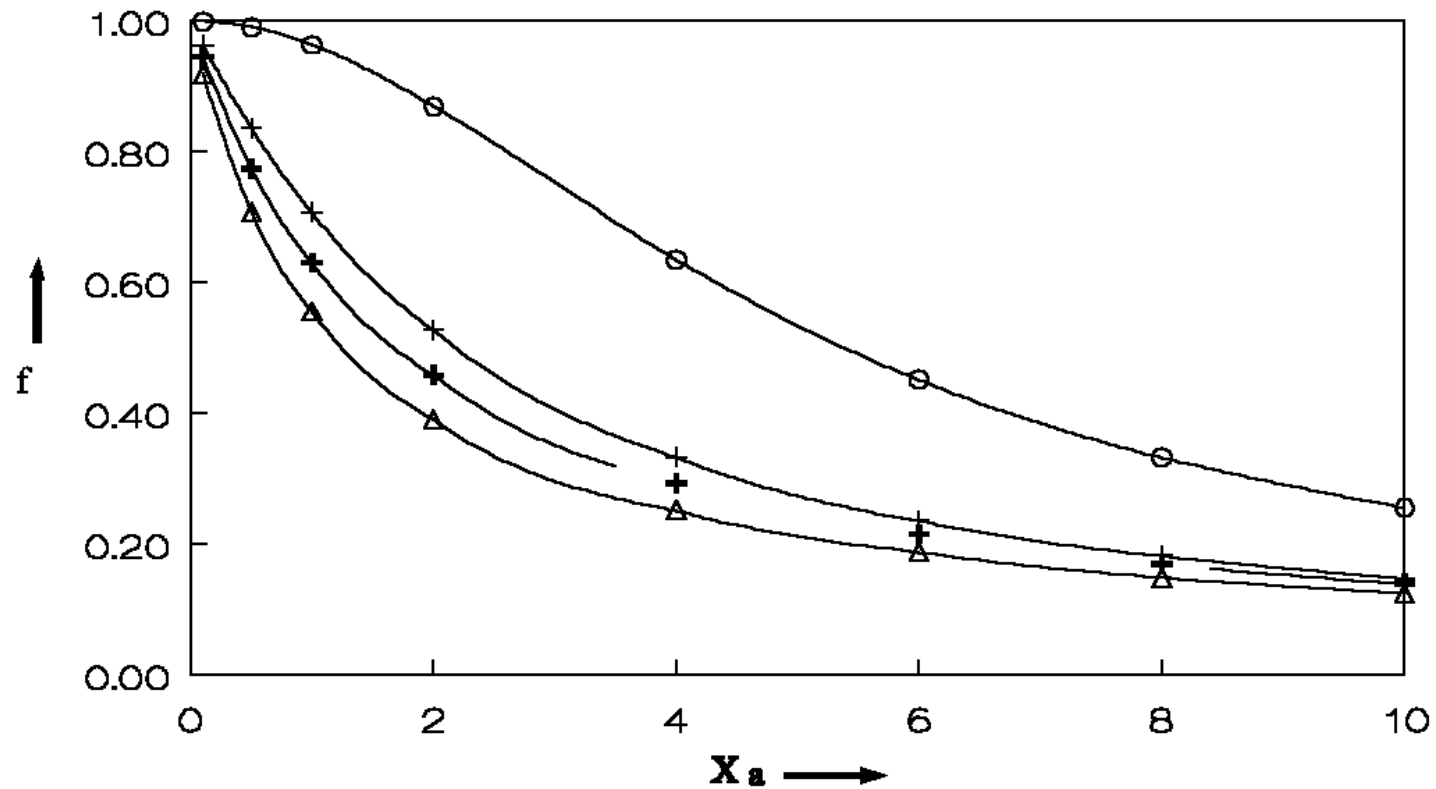
# Stuart, Copley and Blaauw

+  $f(X_a, X_s = 0)$

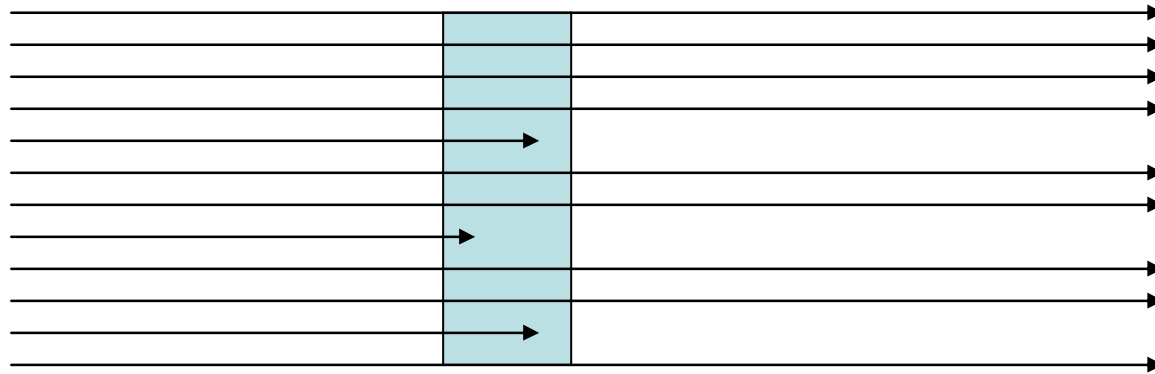
○ Stewart

△ Copley

+ Stewart according to Blaauw



# Thermal flux, absorption only, in a beam



$$f = \frac{(1 - e^{-\langle \Sigma_a \rangle d})}{\langle \Sigma_a \rangle d}$$

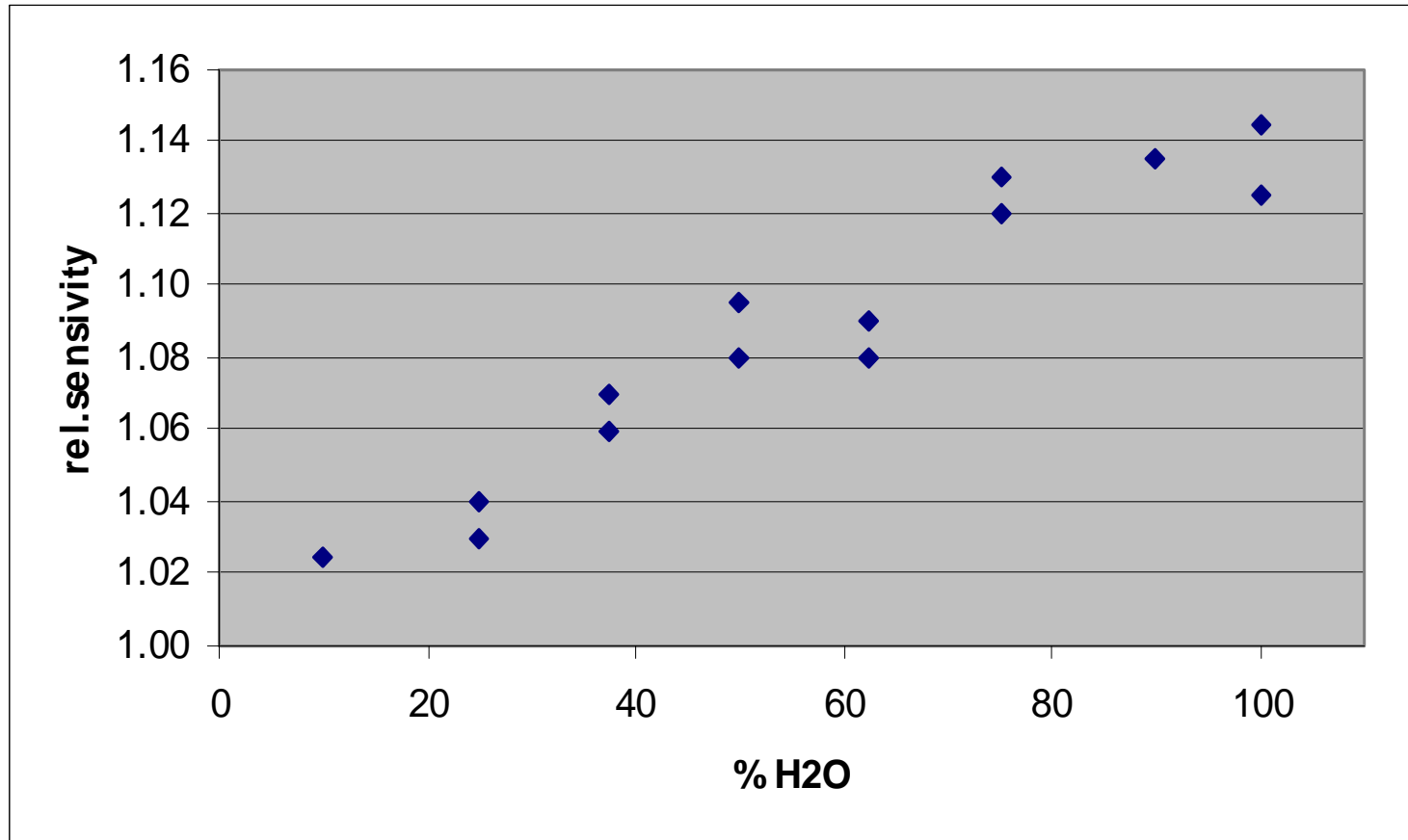
thin:  $\langle \Sigma_a \rangle = \frac{2}{\sqrt{\pi}} \sqrt{\frac{T_0}{T}} \Sigma_{0,a}$

thick:  $\langle \Sigma_a \rangle = \frac{\sqrt{\pi}}{2} \sqrt{\frac{T_0}{T}} \Sigma_{0,a}$

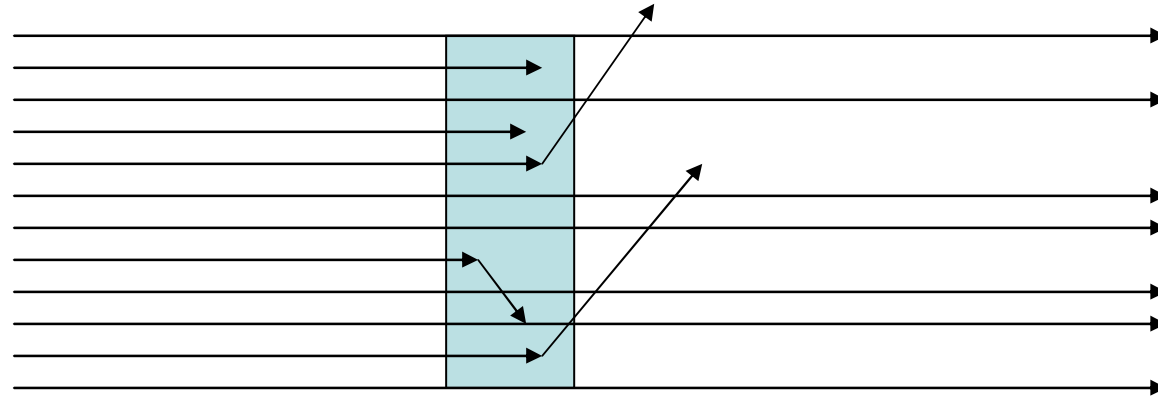
Same for other shapes: Just use the right average

M. Blaauw, "The confusing issue of the neutron capture cross-section to use in thermal neutron self-shielding computations", Nucl.Instr.Meth. A356 (1995) 403-407

# Thermal flux, absorption and scattering, in a beam



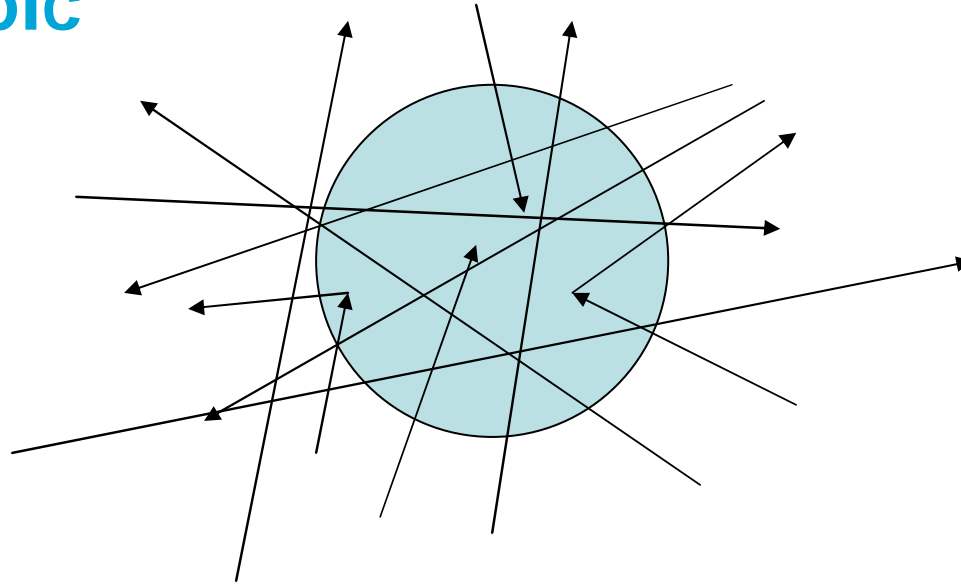
# Thermal flux, absorption and scattering, in a beam



Average neutron path length in sample increases.  
Purely thermal flux: No moderation effect.  
That means  $f > 1$  !!!!!!!



# Thermal flux, absorption and scattering, isotropic



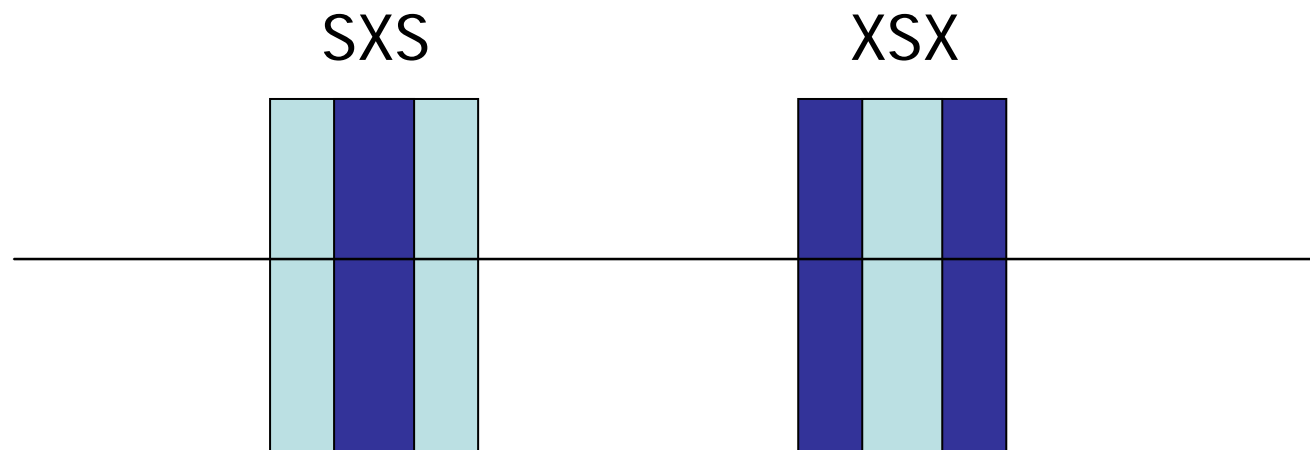
No simple equations known! If flux is purely thermal, 2200 m/s approximation with  $\langle \Sigma_a \rangle$  is good. If epithermal component present, moderation kicks in...

# Self-shielding in reality

- Gradient flux effect on scattering, slab-shaped samples
- Effect of scattering samples on gradient
- BISNIS experiences

# Gradient flux effect on scattering, slab-shaped samples

- If  $f$  can be larger than unity in a slab, in a purely thermal beam, what could happen in a strong gradient?



Unpublished work by Lindstrom and Blaauw

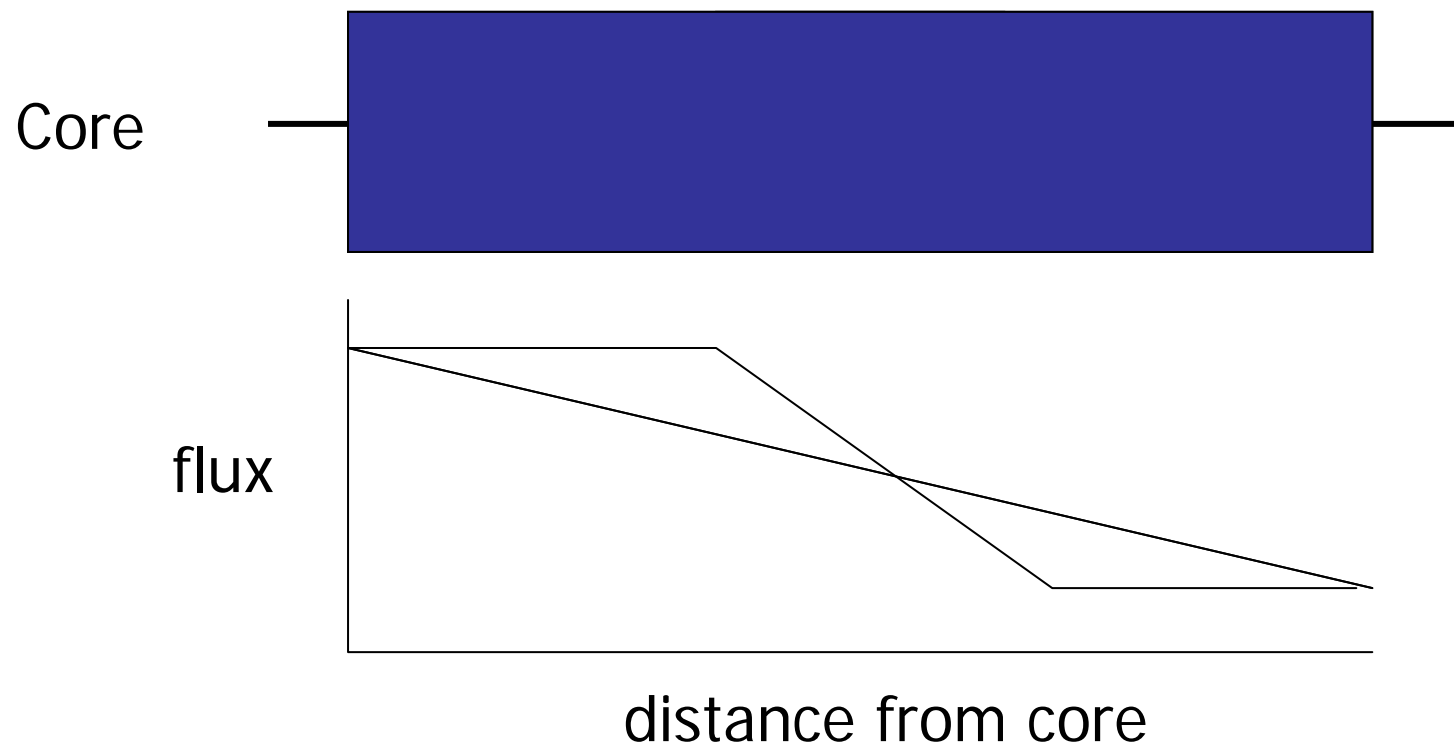
# Results for SXS/XSX ratios

- Experimental
  - In beam:  $1.150 \pm 0.005$
  - In gradient:  $0.993 \pm 0.001$
- Monte Carlo
  - In isotropic:  $0.999 \pm 0.002$
  - In gradient:  $0.998 \pm 0.002$
  - In beam:  $1.118 \pm 0.002$

Conclusion: Gradient fields are ok

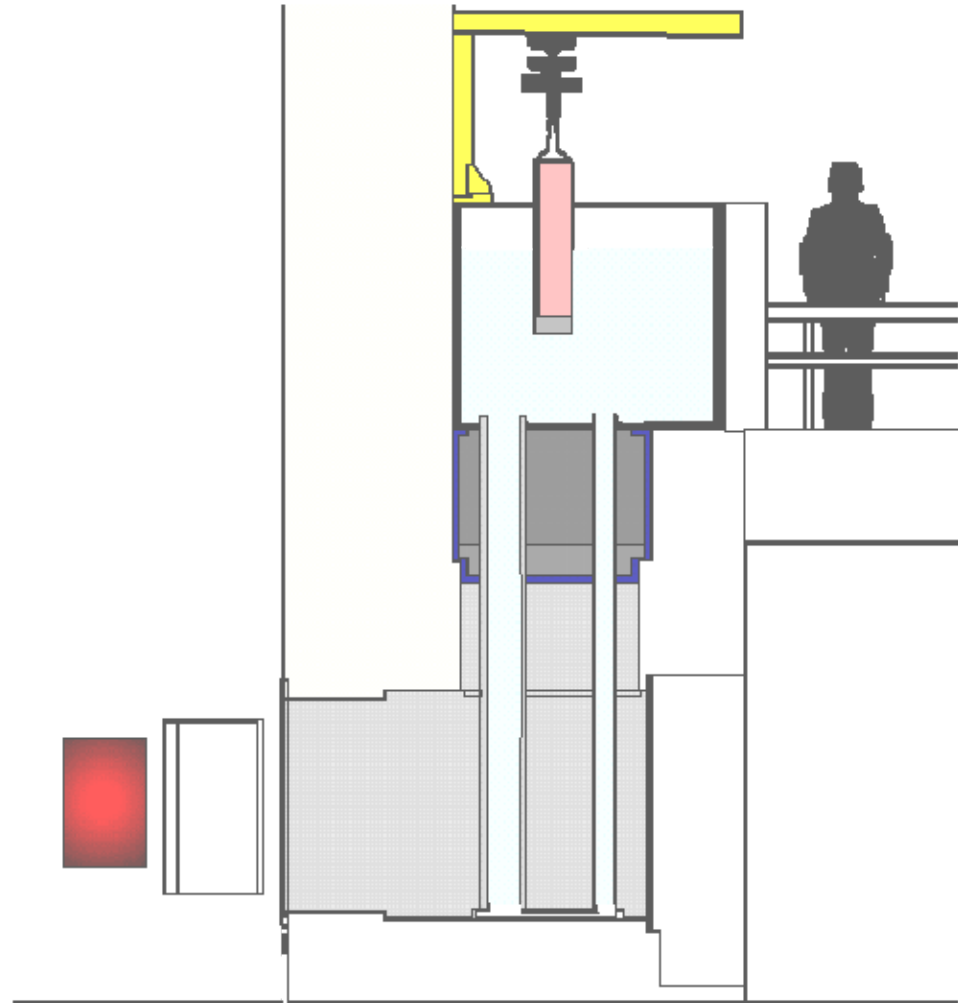
# Effect of scattering samples on gradient

Lindstrom and Blaauw



# BISNIS experiences

Overwater,  
Bode,  
Baas,  
Blaauw



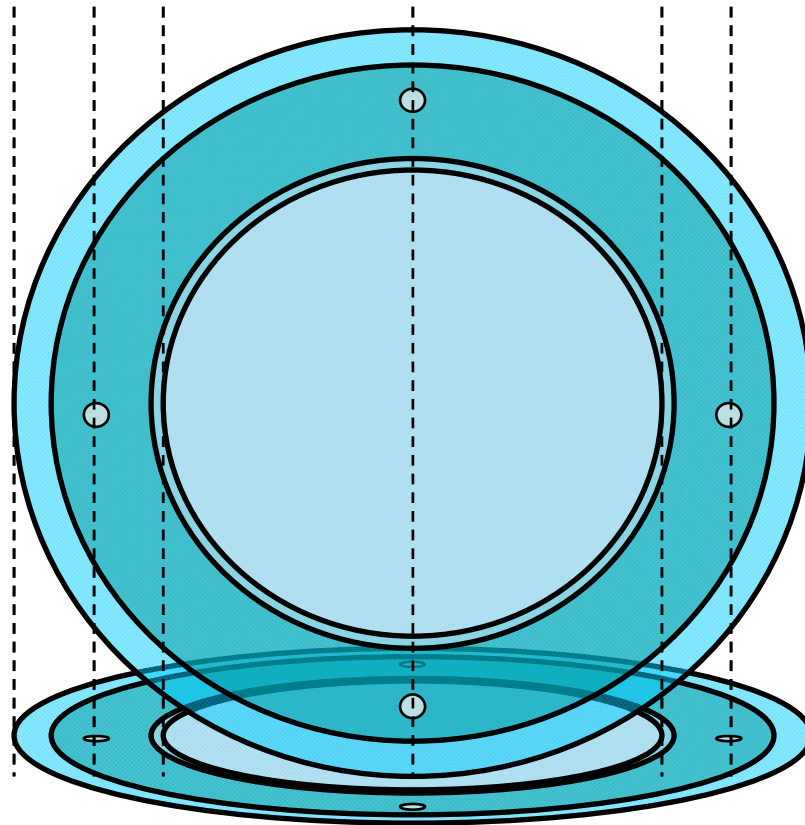
Trieste, Nov 8-10, 2010

# Sample container and sample



Trieste, Nov 8-10, 2010

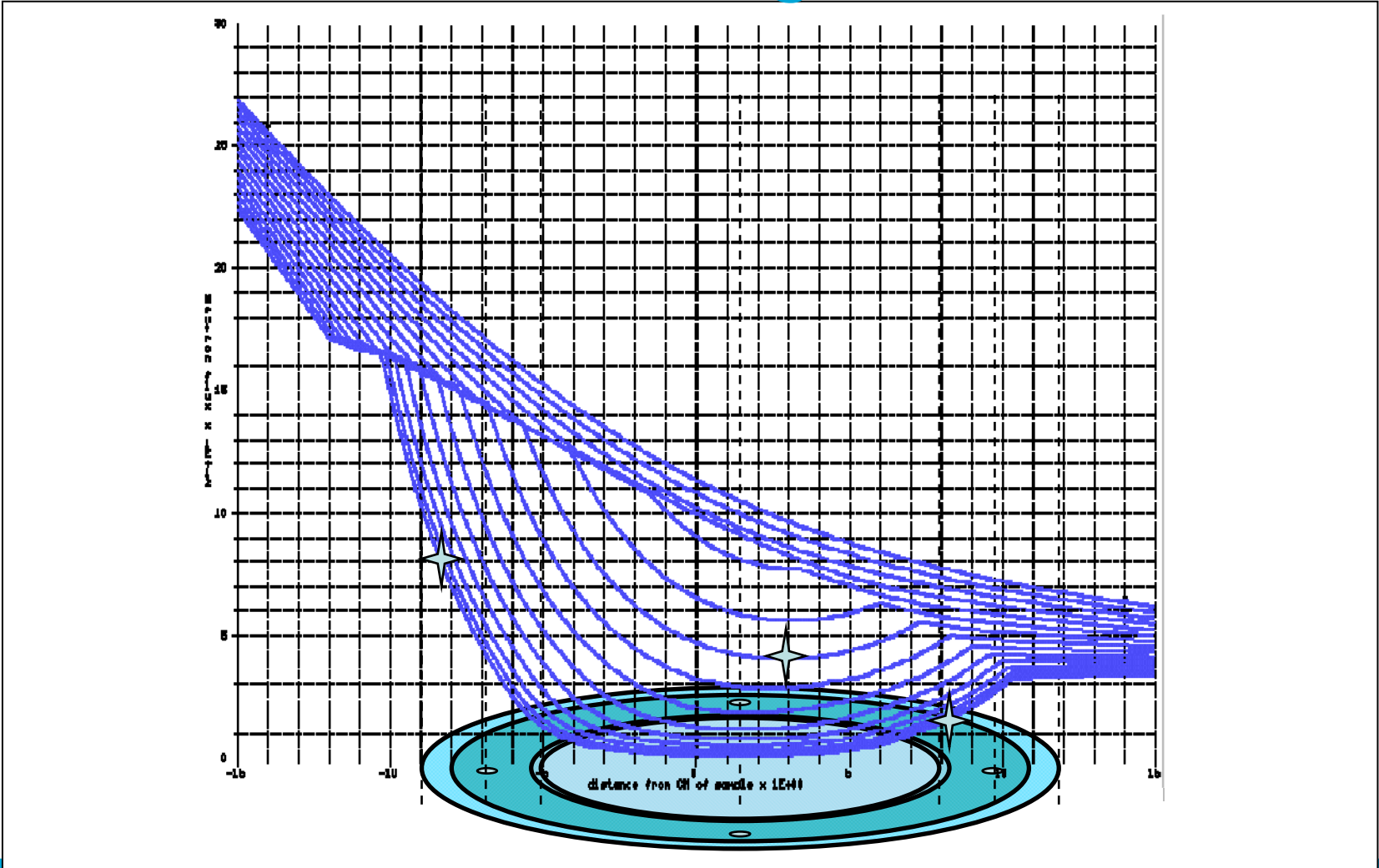
# Neutron self-shielding



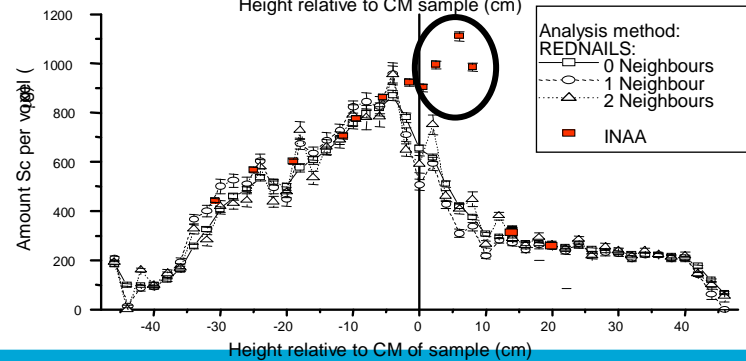
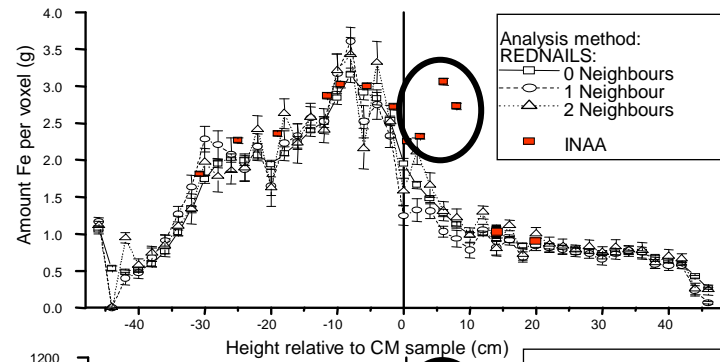
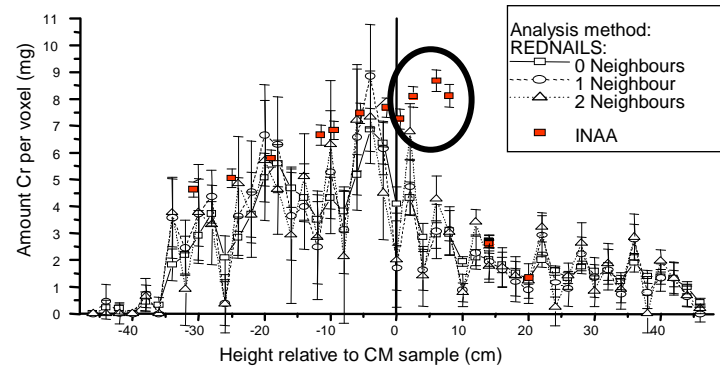
Trieste, Nov 8-10, 2010



# Neutron self-shielding



# Results per layer compared to ordinary INAA



# Conclusions

- Keep your samples small
- Always sandwich between monitors
- If in a beam, either use spherical shape or do internal standardization
- In an isotropic flux, make sure you use the right, temperature-corrected averaged capture cross section.
- For hydrogenous materials, try to use a very thermal flux. Otherwise, Cd-cover may be essential to check for the epithermal-moderated contribution.