



**The Abdus Salam
International Centre for Theoretical Physics**



2234-17

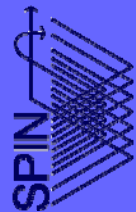
**Meeting of Modern Science and School Physics: College for School
Teachers of Physics in ICTP**

27 April - 3 May, 2011

The century jubilee of superconductivity

Ruggero Vaglio

*SPIN-CNR and University of Naples
Italy*



Istituto SPIN-CNR



Trieste, April 28th 2011



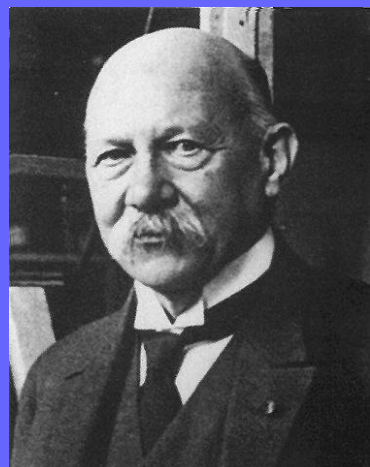
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Ruggero Vaglio

The Century Jubilee of

Superconductivity

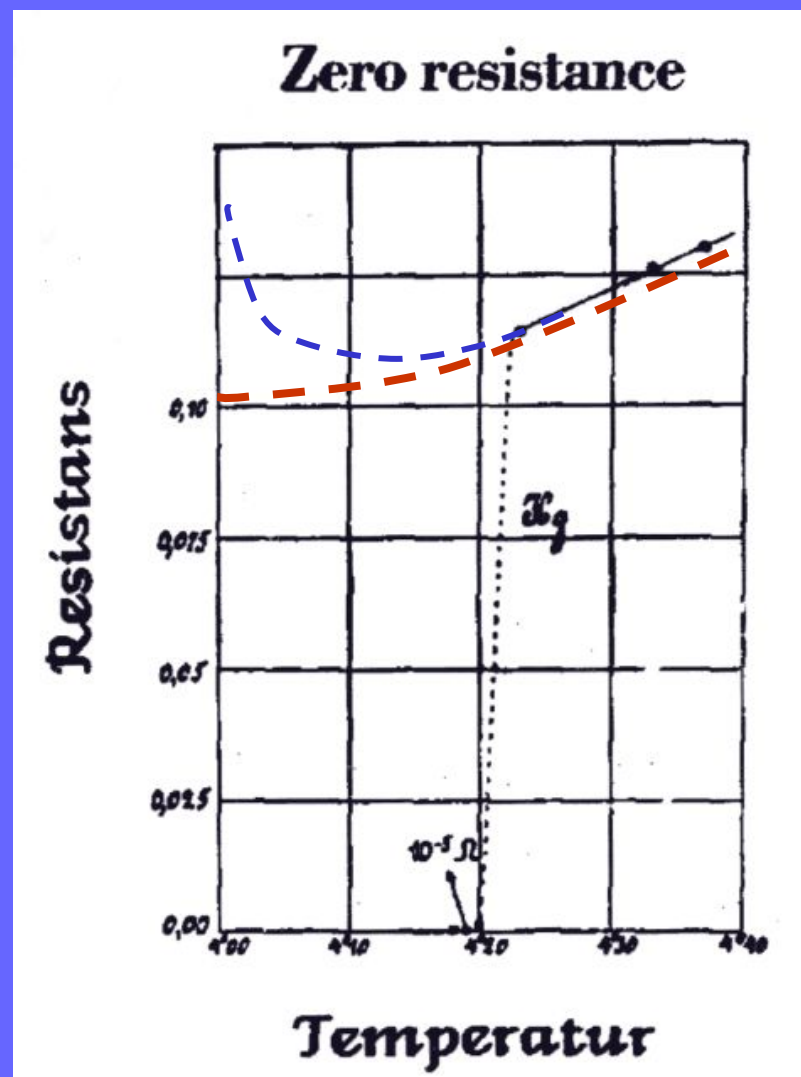
1911: The discovery of superconductivity

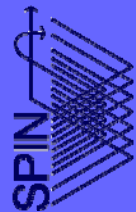


In 1911 Kamerlingh Onnes during his experiments on Helium liquefaction observed the sudden disappearance of mercury electrical resistance.

- In 1912 he discovered that electrical resistance reappeared in presence of intense magnetic fields or currents

Maximum current : J_c Maximum field : B_c





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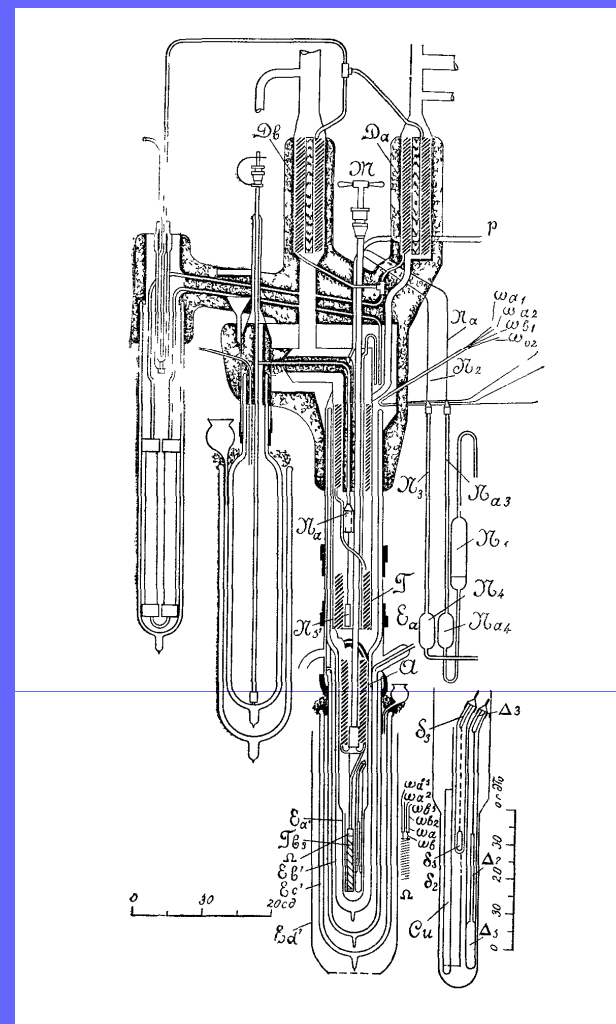


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Kamerlingh Onnes Laboratory



1913



Superconducting elements

H																			He
Li	Be													B	C	N	O	F	Ne
Na	Mg													Al	Si	P	S	Cl	Ar
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn			Ga	Ge	As	Se	Br	Kr
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd			In	Sn	Sb	Te	I	Xe
Cs	Ba	Lu	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg			Tl	Pb	Bi	Po	At	Rn
Fr	Ra																		
			La	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb			
			Ac	Th	Pa	U	Np	Pu											



Supereconductors



Superconductors under high pressure or in thin films



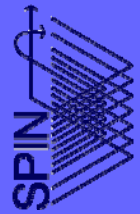
Metallic with Magnetic order



Metallic but not yet found to be superconductors



Non Metallic



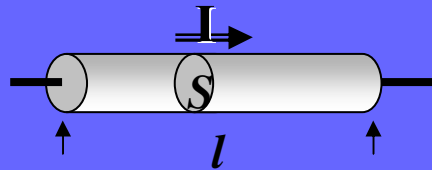
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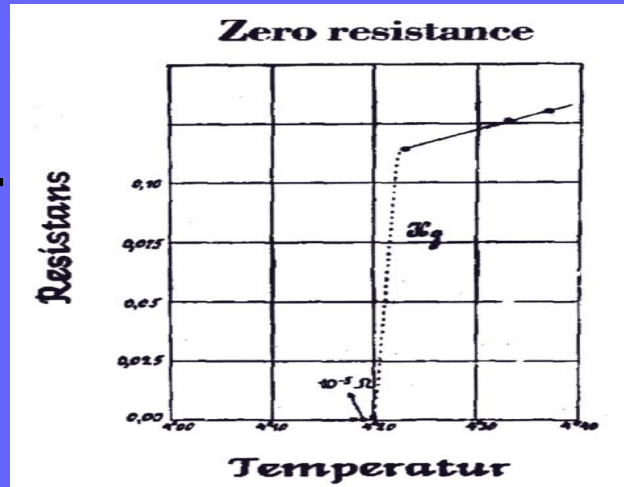
Trieste, April 28th 2011 Zero Resistance



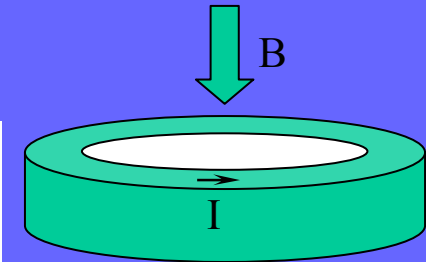
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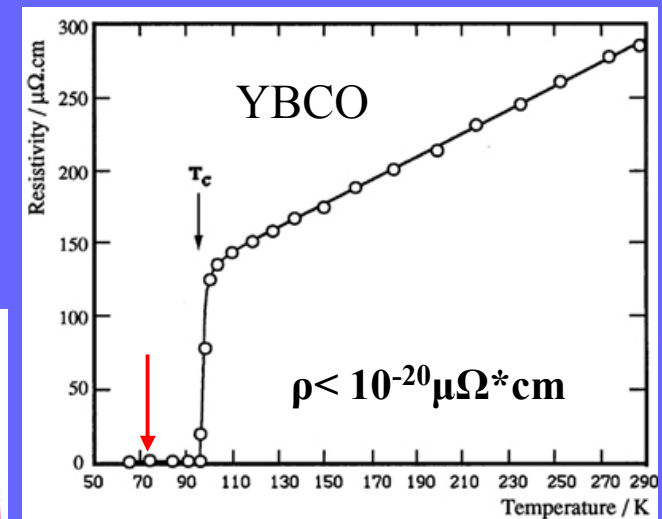
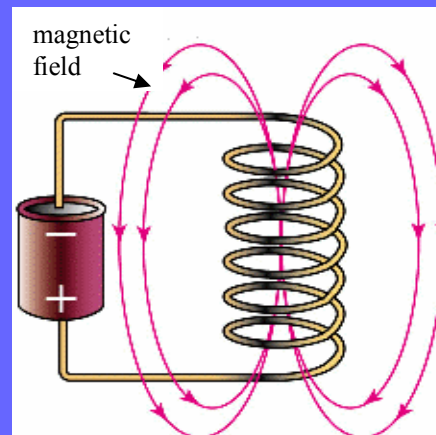
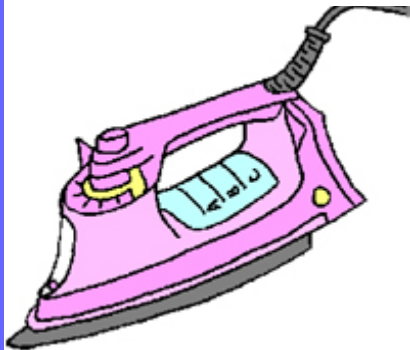
$$V = RI$$
$$R = \rho \frac{l}{S}$$



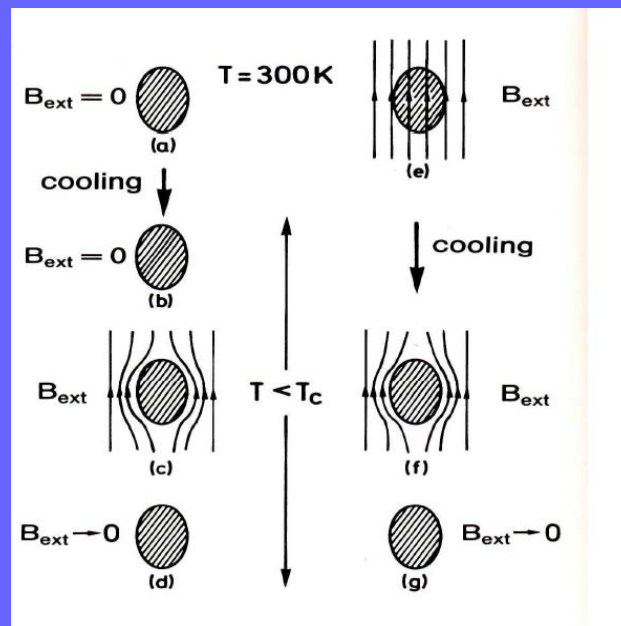
$$I = I_0 e^{-\frac{t}{\tau}}$$
$$\tau = \frac{L}{R}$$



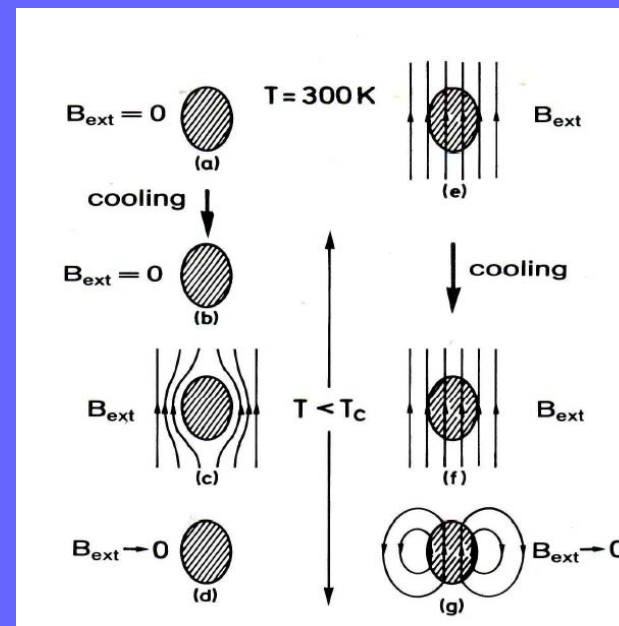
$$P = RI^2$$



1933: The Meissner effect



superconductor



perfect conductor

- In 1933 Meissner discovered that superconducting elements expel the magnetic field, behaving as “perfect diamagnets”

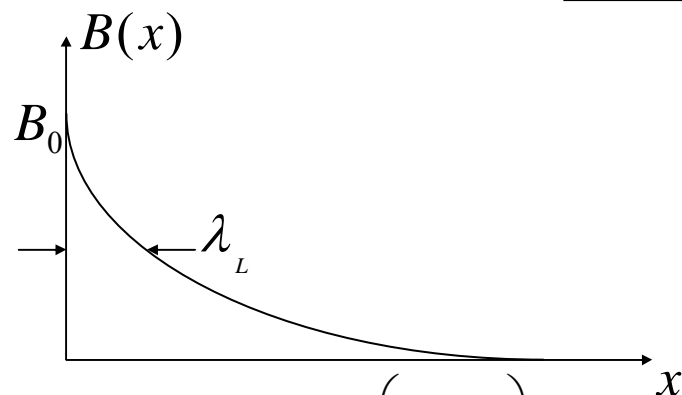


A magnet can “levitate” over a superconductor (and viceversa)

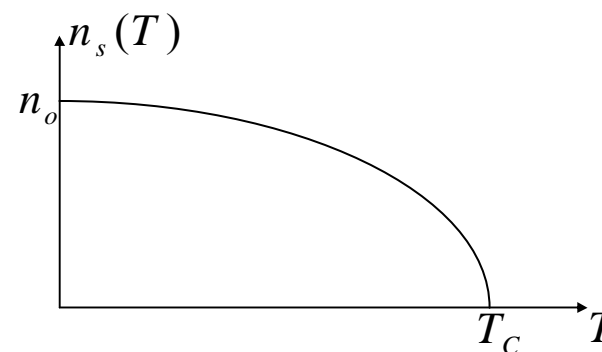
$$\left\{ \begin{array}{l} \vec{J}_n = -n_n e \vec{v}_n \quad \vec{v}_n = -\mu_n \vec{E} \quad \vec{J}_n = n_n e \mu_n \vec{E} \\ \vec{J}_s = n_s q_s \vec{v}_s \quad m_s \frac{d\vec{v}_s}{dt} = q_s \vec{E} \quad \frac{d\vec{J}_s}{dt} = \frac{n_s q_s^2}{m_s} \vec{E} \end{array} \right. + \text{Eq. Maxwell} \left\{ \begin{array}{l} \text{rot } \vec{B} = \mu_o \vec{J}_s \\ \text{rot } \vec{E} = -\frac{\partial \vec{B}}{\partial t} \end{array} \right.$$

$$\left\{ \begin{array}{l} \text{rot rot } \vec{B} = -\nabla^2 \vec{B} = \mu_o \text{rot } \vec{J}_s \\ \frac{\partial \mu_o \text{rot } \vec{J}_s}{\partial t} = \frac{\mu_o n_s q_s^2}{m_s} \text{rot } \vec{E} = -\frac{\mu_o n_s q_s^2}{m_s} \frac{\partial \vec{B}}{\partial t} \end{array} \right. \quad \frac{\partial}{\partial t} \left(\nabla^2 \vec{B} - \frac{1}{\lambda_L^2} \vec{B} \right) = 0$$

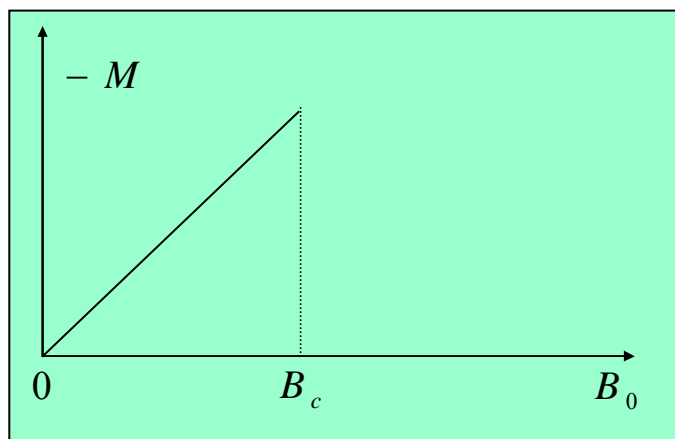
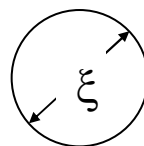
$$\lambda_L = \left(\frac{n_s q_s^2 \mu_o}{m_s} \right)^{-1/2}$$



$$B(x) = B_o \exp\left(-\frac{x}{\lambda_L}\right)$$

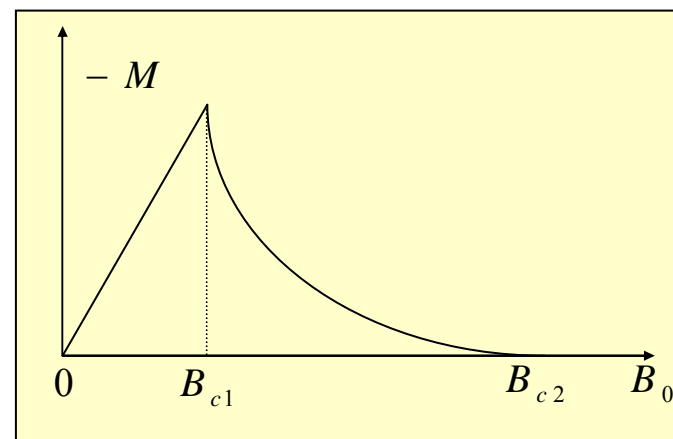


Deviation of values of the penetration depth magnetic from London predictions and anomalous Meissner effect for alloys :



Type I superconductors
(*Pb, Sn, In, Al*)

$$k \equiv \frac{\lambda}{\xi} < \frac{1}{\sqrt{2}}$$

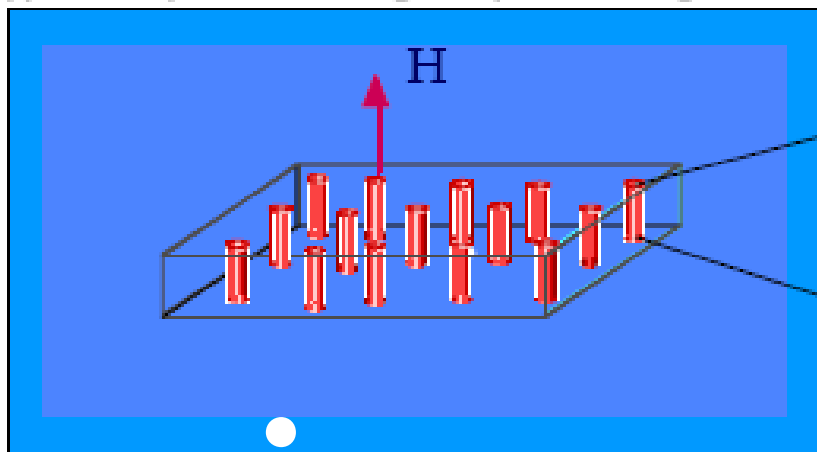


Type II superconductors
(*Nb, alloys*)

$$k \equiv \frac{\lambda}{\xi} > \frac{1}{\sqrt{2}}$$

Explanation of type II Superconductivity : Abrikosov Vortices

Type II superconducting sample in a magnetic field



normal zones

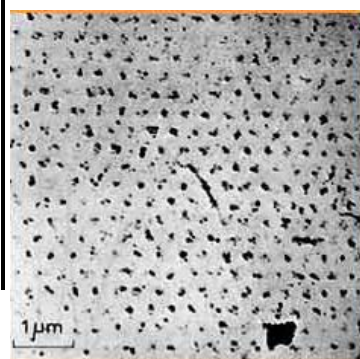
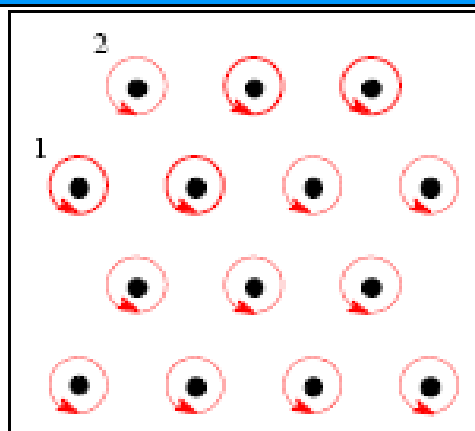


2003

me

$$B_{c2} \xi^2 = \phi_0$$

exagonal vortex lattice



Critical currents are limited by vortex mobility (pinning)

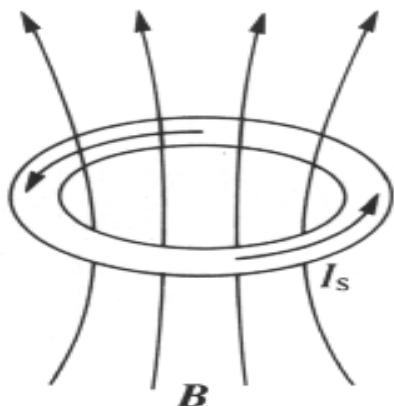
GL theory : quantum interpretation

$$n_s(T) = |\psi|^2 \quad (\Psi \text{ is interpreted as a macroscopic quantum wave-function})$$

$$\psi = \psi_0 e^{i\varphi} = n_s^{1/2} e^{i\varphi}$$

$$\vec{J}_s = n_s q_s \vec{v}_s = \frac{n_s q_s}{m_s} (\hbar \vec{\nabla} \varphi - q_s \vec{A})$$

Flux Quantization



$$\vec{J}_s = 0 \Rightarrow \hbar \vec{\nabla} \varphi = q_s \vec{A}$$

$$\hbar \oint \vec{\nabla} \varphi \cdot d\vec{l} = q_s \oint \vec{A} \cdot d\vec{l} = q_s \int_S \vec{B} \cdot d\vec{S}$$

$$\hbar 2\pi n = q_s \phi$$

$$\phi = n \phi_0$$

$$\phi_0 = \frac{h}{q_s} = 2 \cdot 10^{-7} \text{ gauss} \cdot \text{cm}^2$$

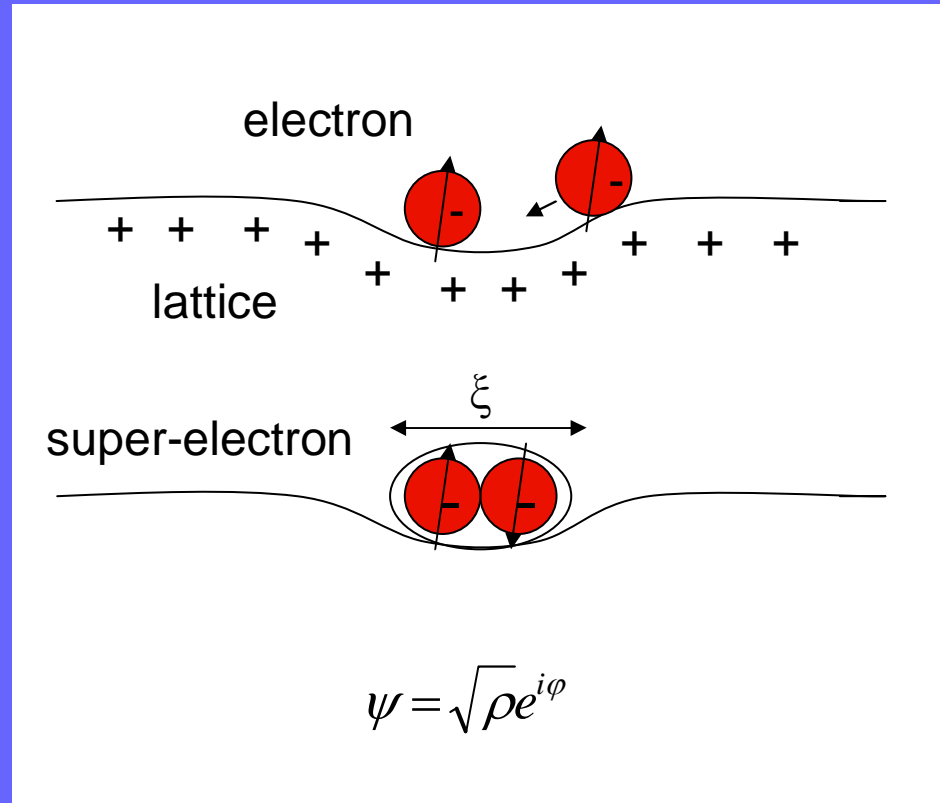
From experimental observation $q_s = 2e$!

1957: The BCS theory

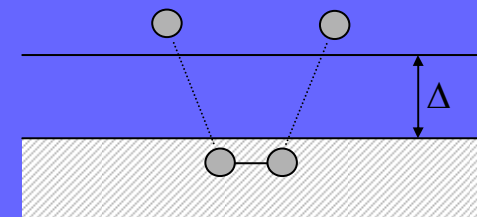


$$T_c = (2\hbar\omega_D\gamma/\pi) \exp[-2/(g\nu)].$$

The BCS theory is based on the interaction of electrons with lattice vibrations (phonons). Its extension (Eliashberg) give correct predictions of critical temperatures (and of the reasons why Cu and Fe are not superconductors) and of the "isotope effect"



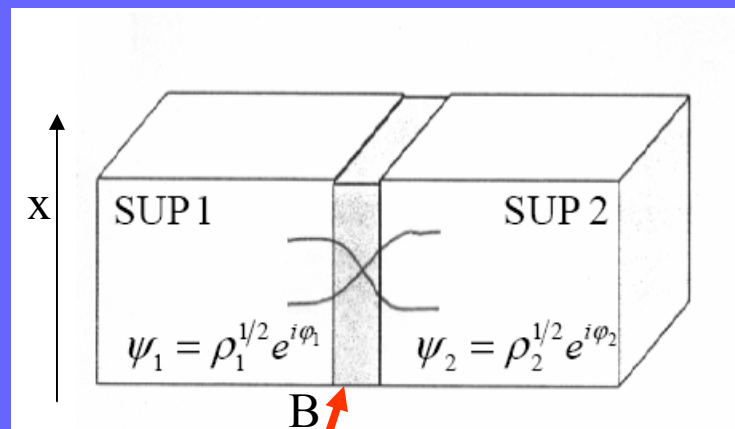
1972



1962: The Josephson effect



Josephson predicted that super-electrons could pass through an insulating barrier by tunnel effect



$$J = \frac{2K}{\hbar} \sqrt{\rho_1 \rho_2} \sin \varphi = J_1 \sin \varphi$$

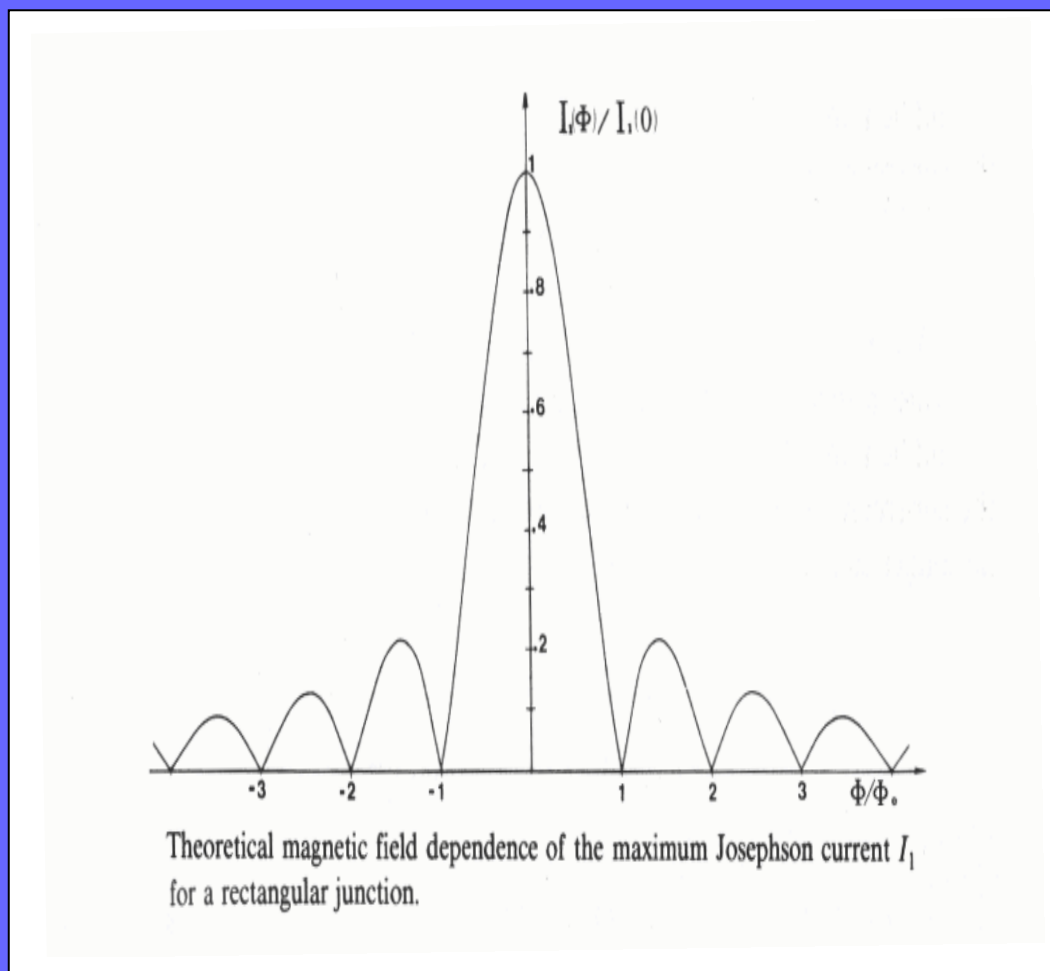
$$\frac{\partial \varphi}{\partial t} = \frac{2e}{\hbar} V$$

$$\frac{\partial \varphi}{\partial x} = \frac{2ed}{\hbar c} B$$



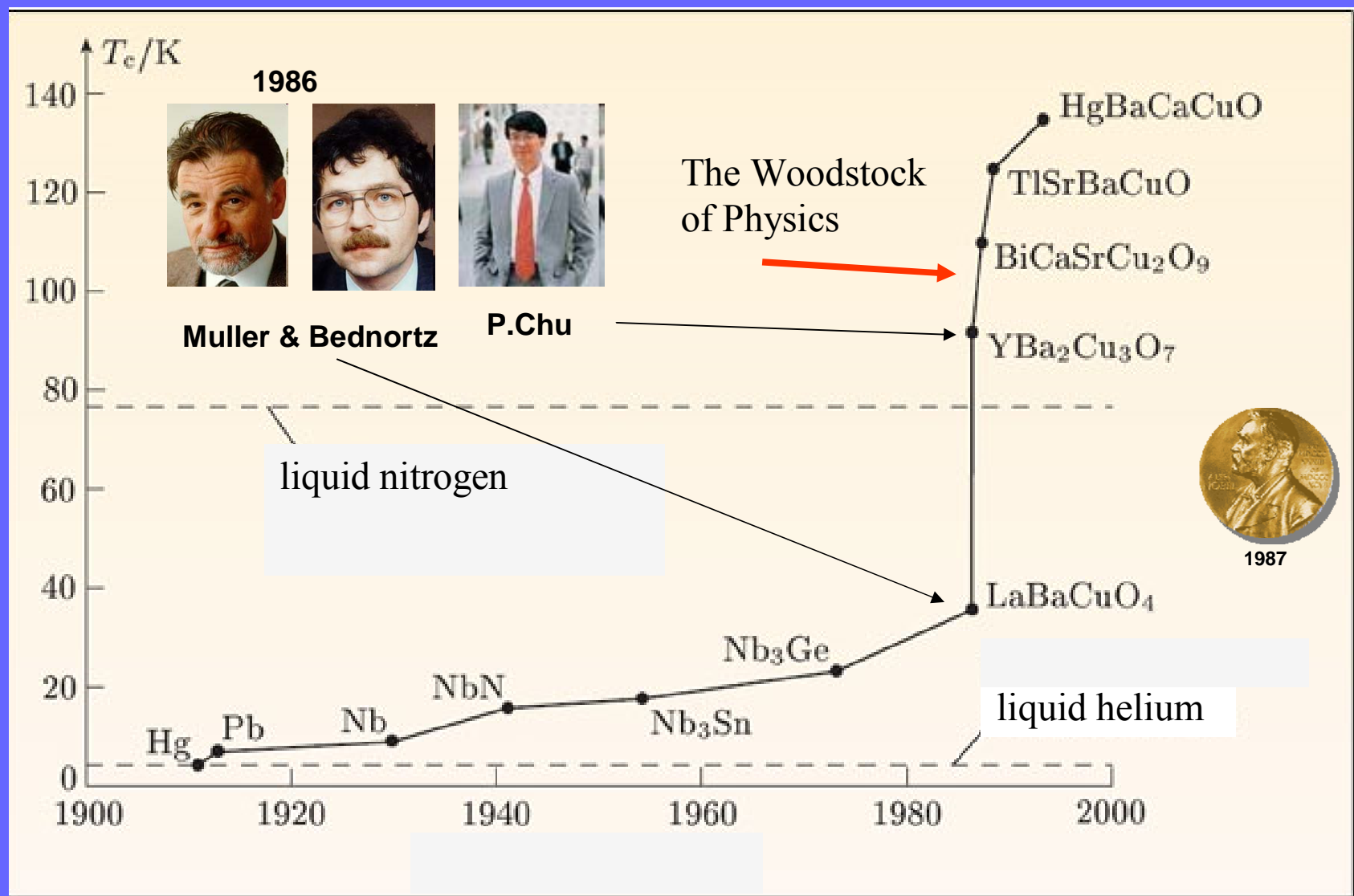
1973

“Theory and Applications of the Josephson Effect”
Barone e Paterno', Wiley 1982

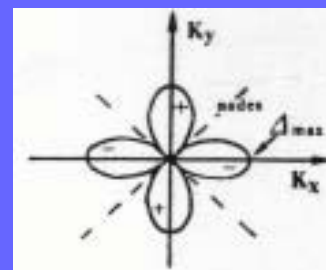
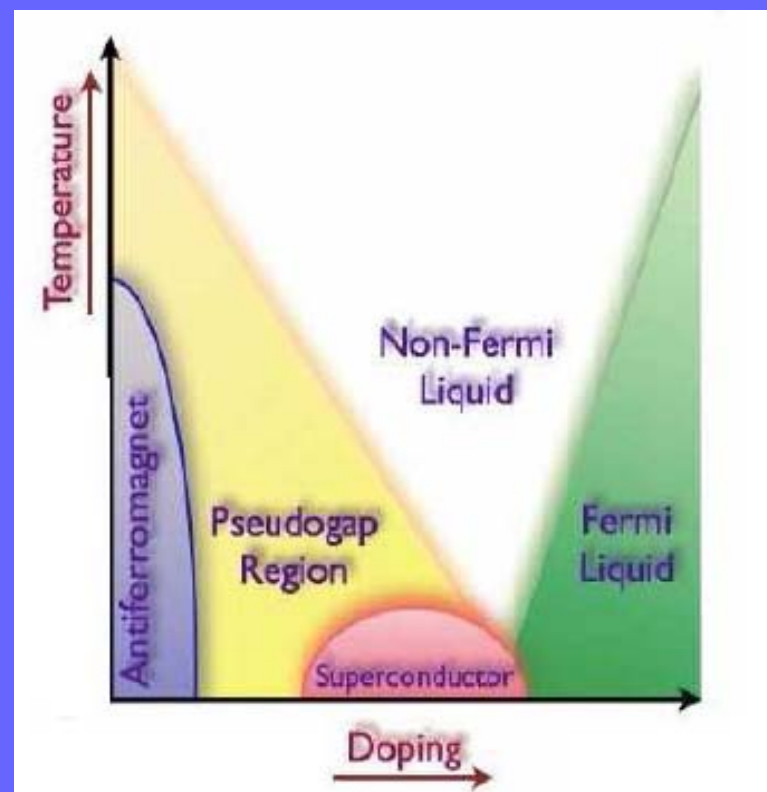
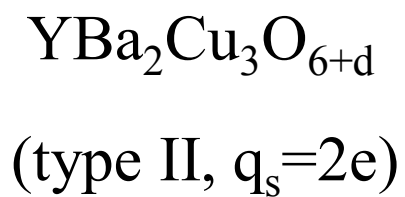
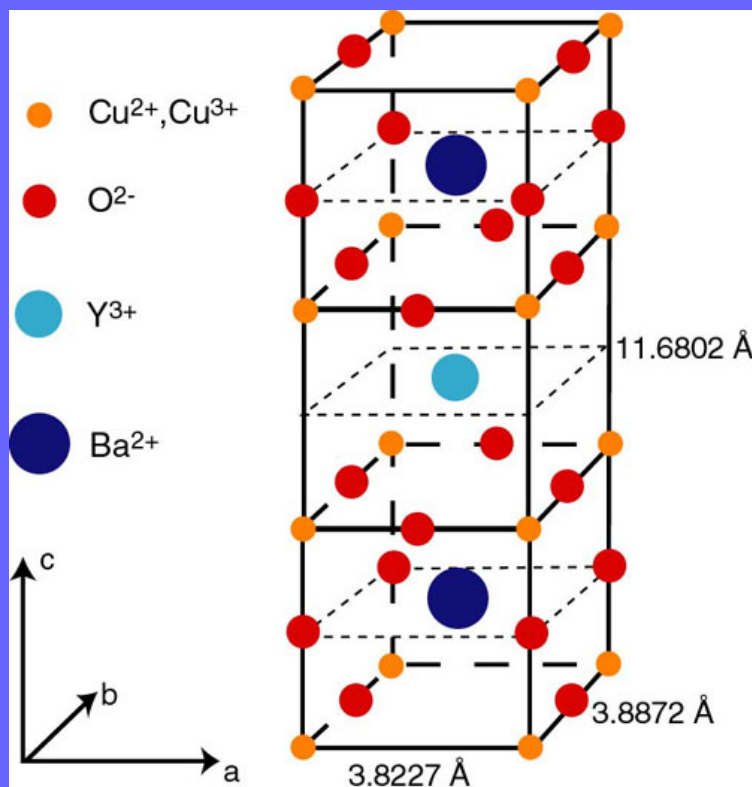


The magnetic field sensitivity can reach 10^{-11} gauss

The high Tc race



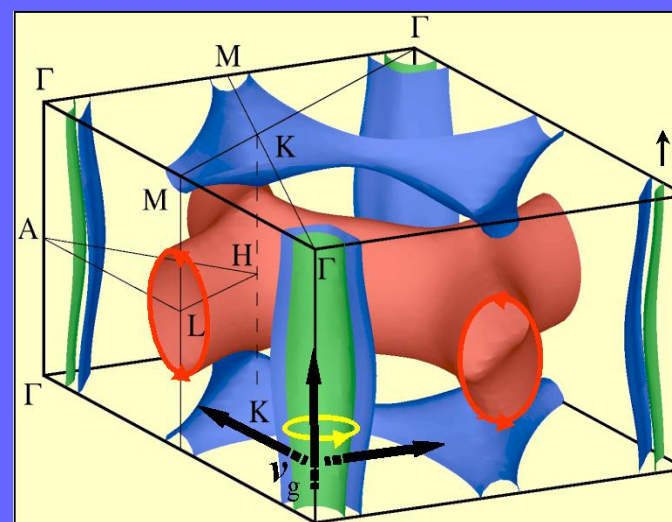
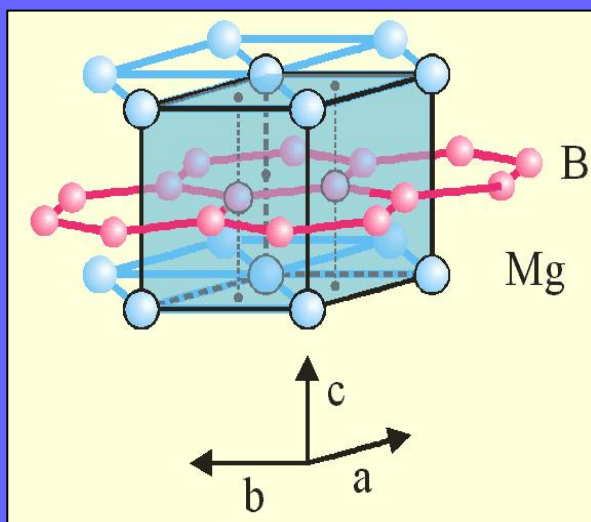
High Tc Superconductors



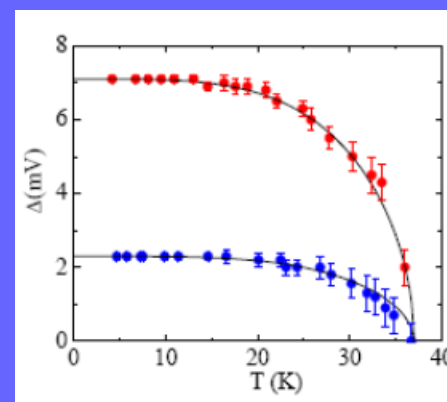
d-wave
 symmetry
 ($d_{x^2-y^2}$)

2001 : MgB_2 $T_c = 40\text{K}$:

the "best" conventional metallic superconductor !

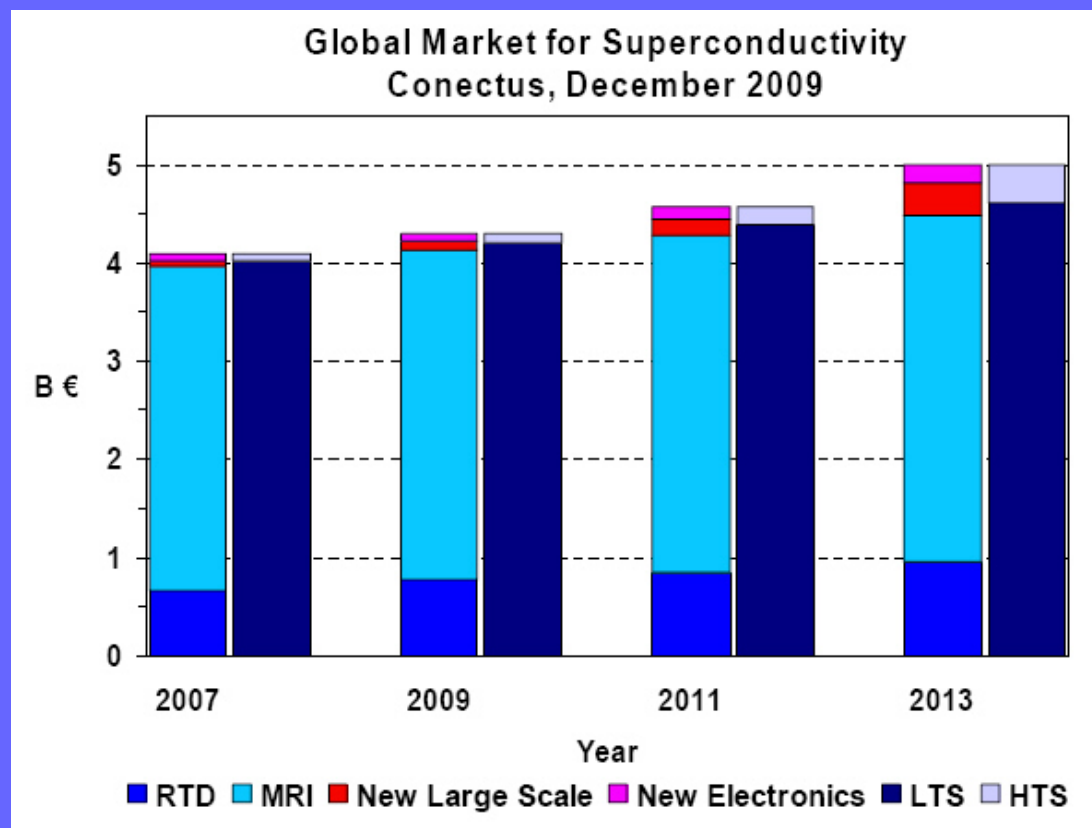


Two-band, two-gaps superconductor !

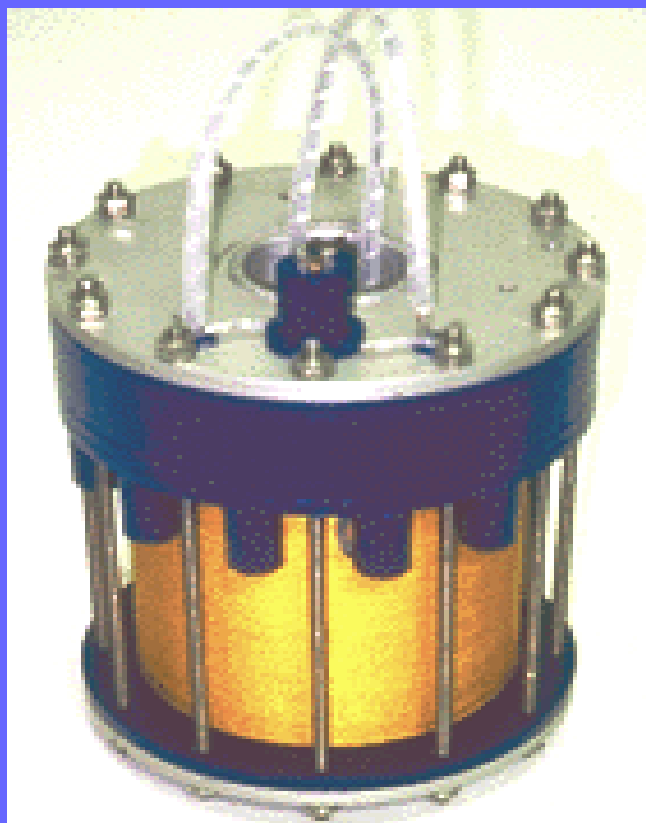


Applications

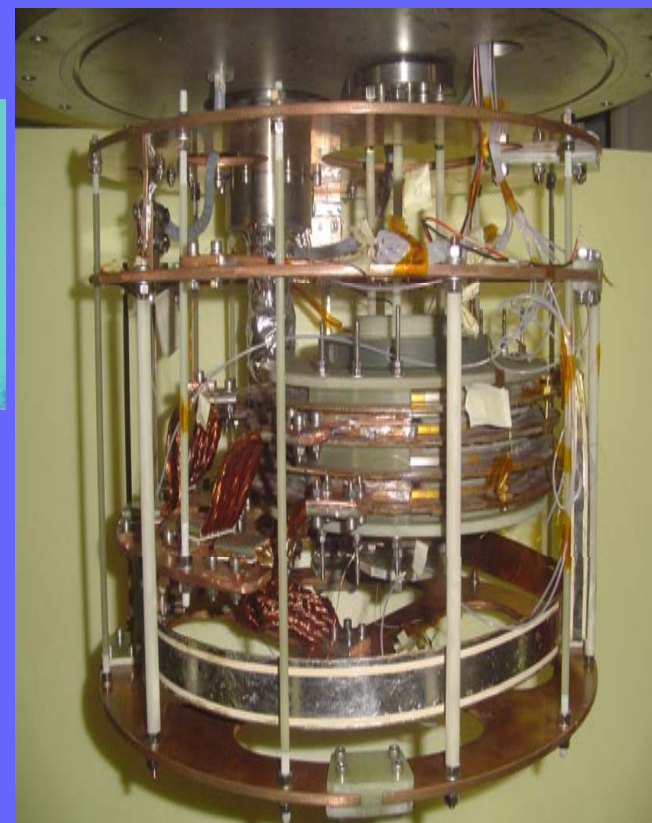
- a) Zero Resistance (magnets and transport)
- b) Josephson effect (magnetic field sensors and electronics)



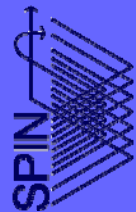
Superconducting magnets for research laboratories



NbTi



MgB₂



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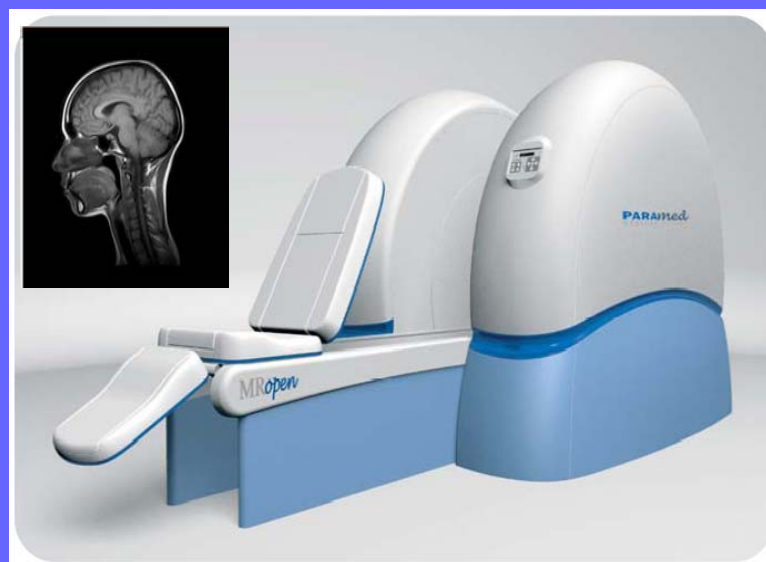


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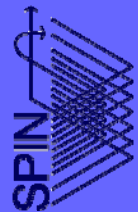
Magnets for MRI



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Columbus Superconductors SPIN- Genova



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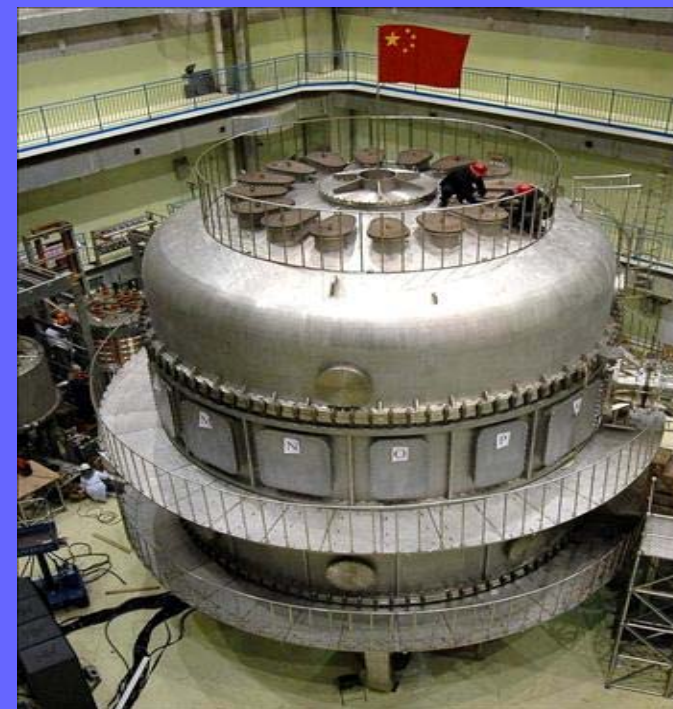
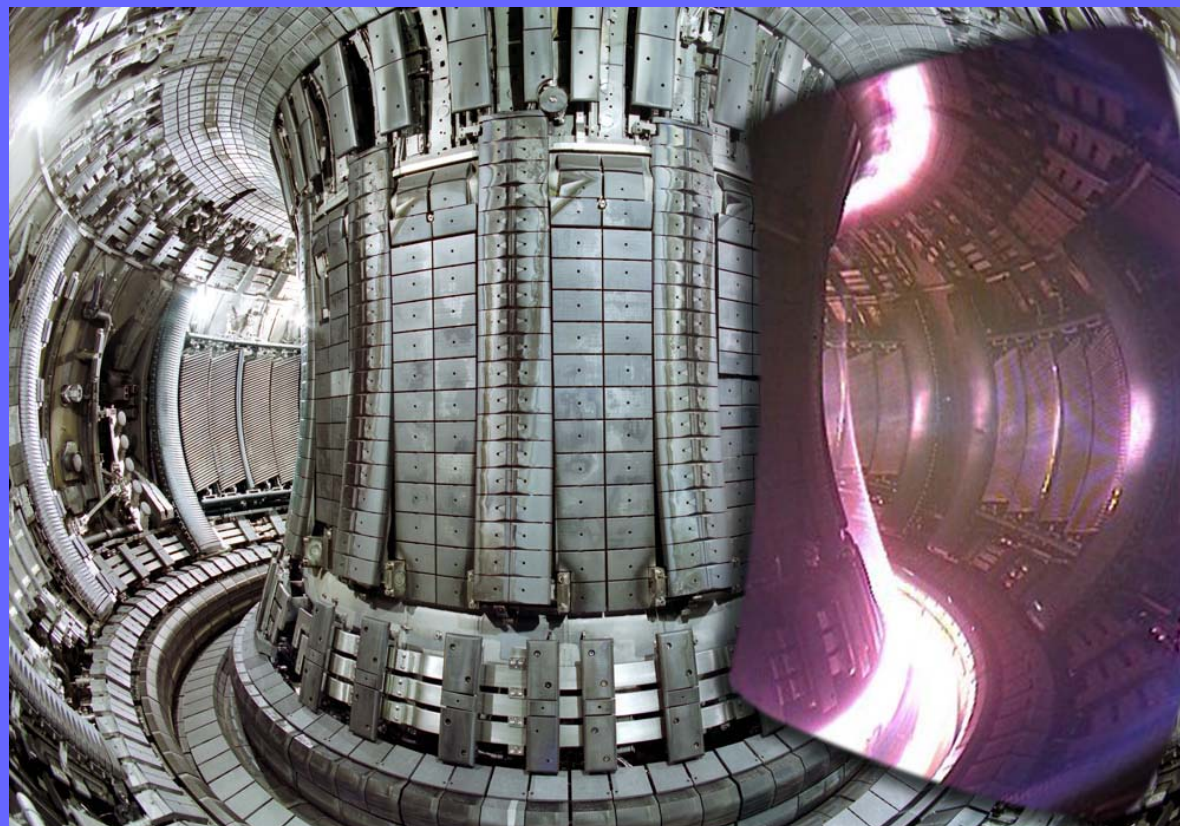


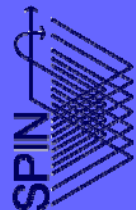
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Magnets for fusion experiments (ITER)





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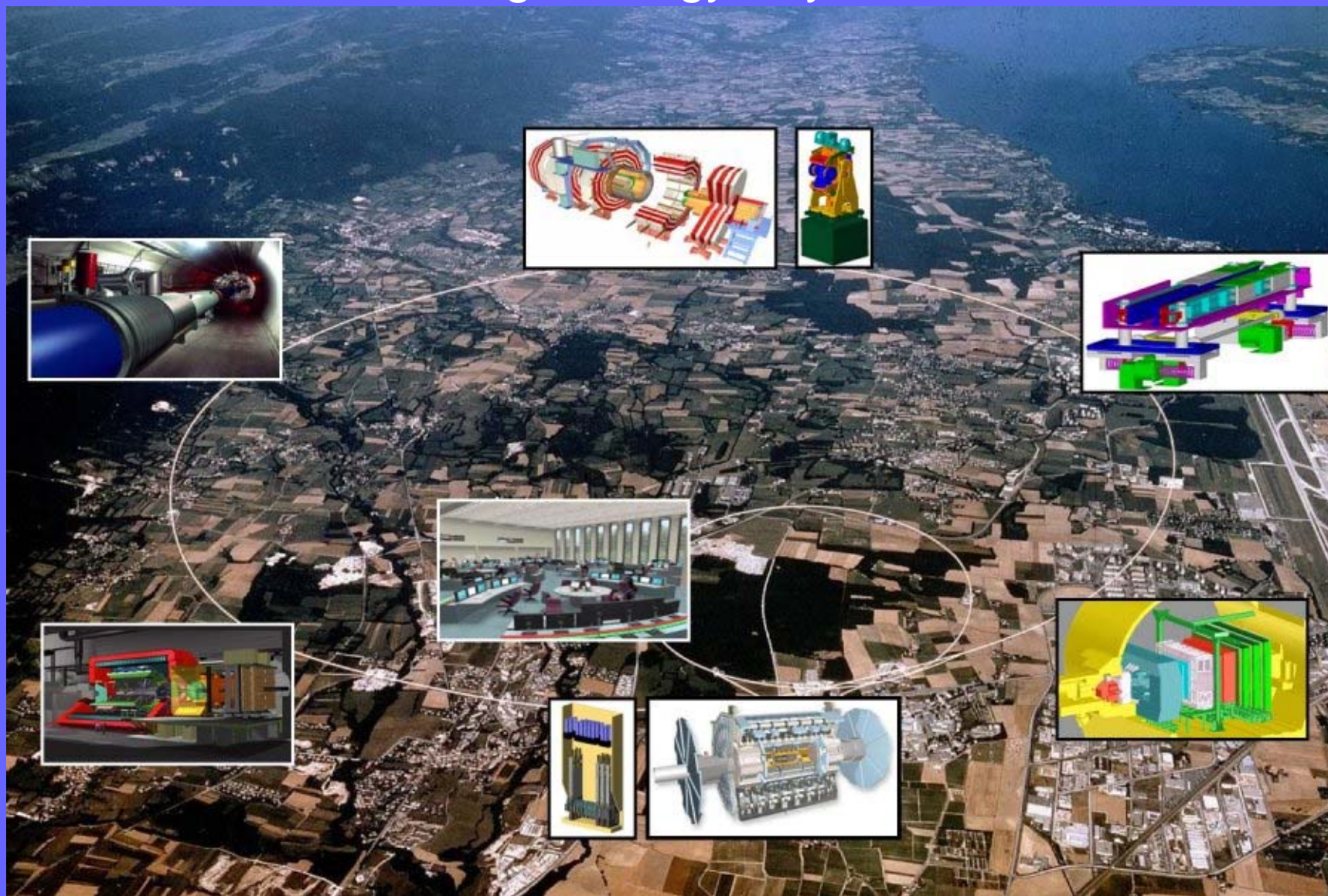


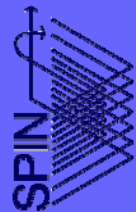
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High Energy Physics





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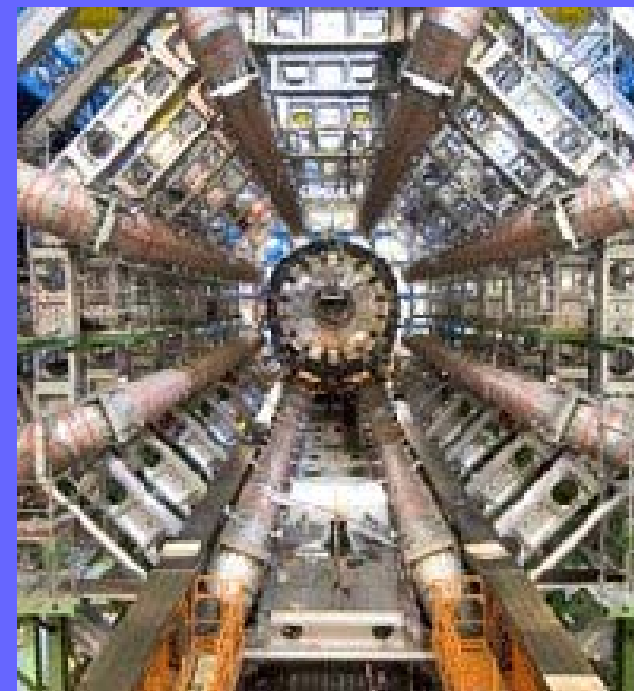


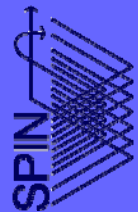
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LHC Superconducting Magnets





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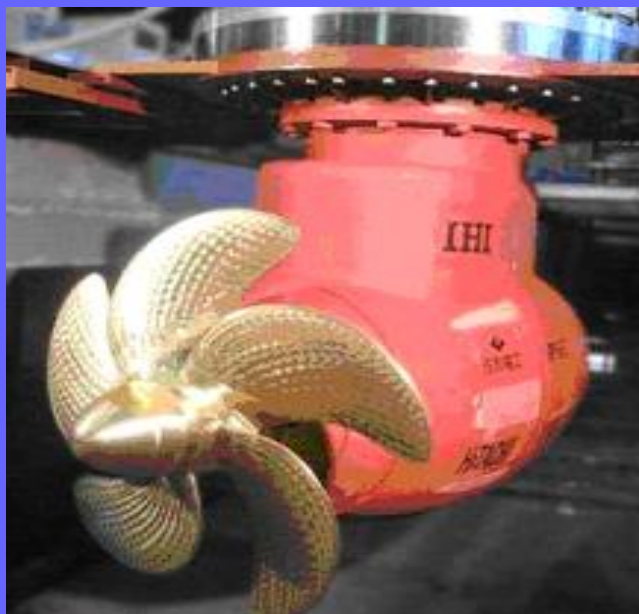
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Superconducting Motors and Generators

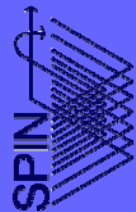
Motore BSCCO



1
maglev
wind
turbine

it powers
750
thousand
homes*

1 maglev
wind turbine
requires less than
100 acres



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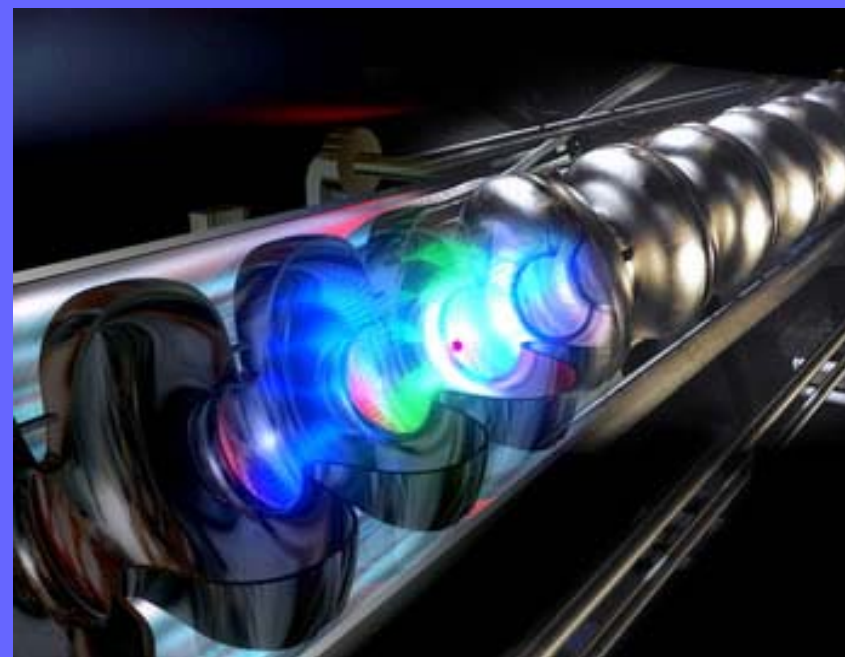


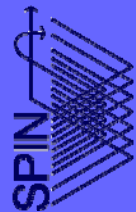
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Superconducting Accelerating Cavities





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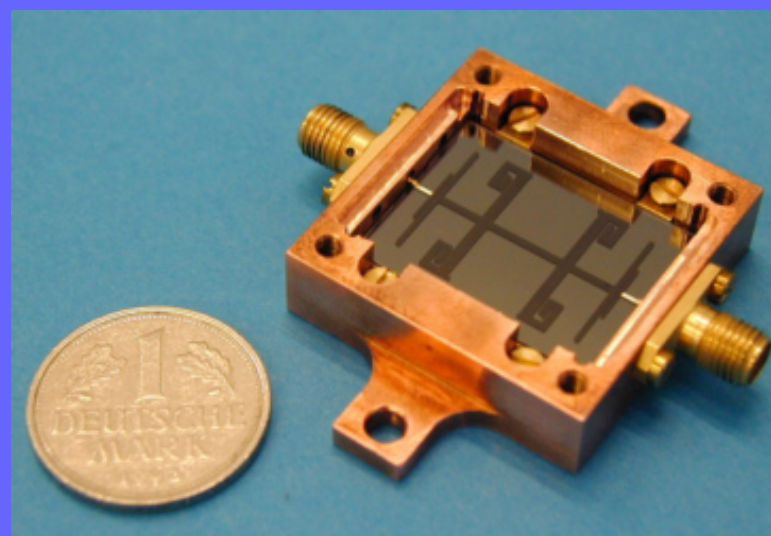


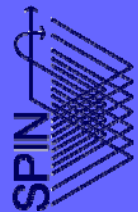
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Passive filters for cellular communications (BST)





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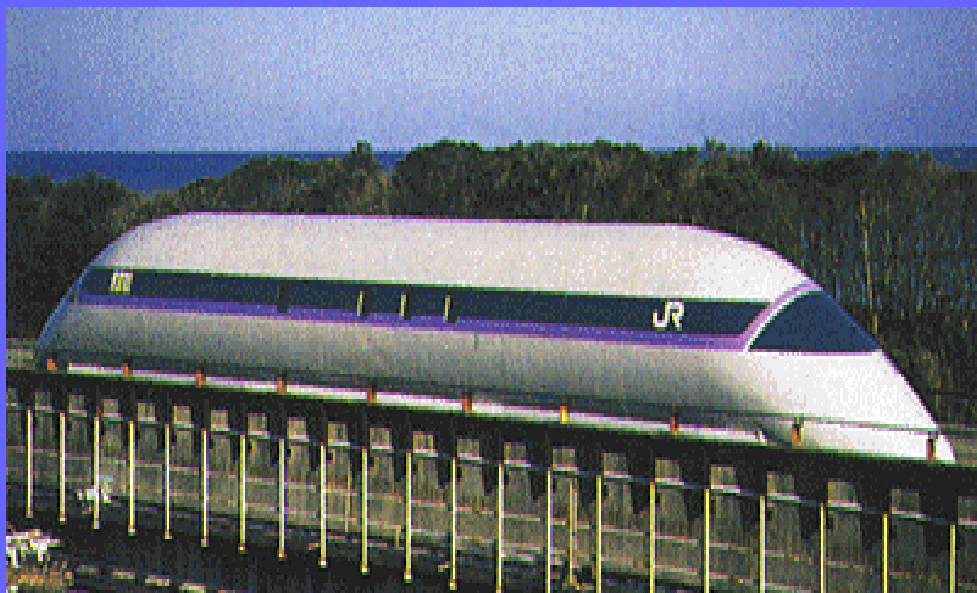


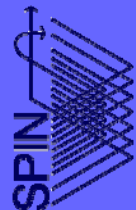
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Magnetic Levitation MAGLEV Trains





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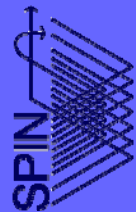
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High Tc superconducting cables for high efficiency connections





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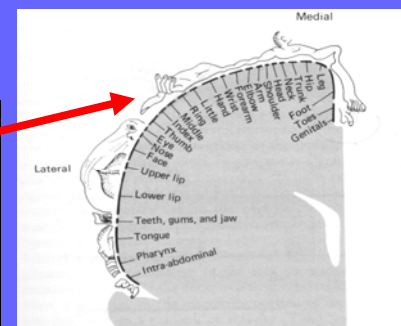
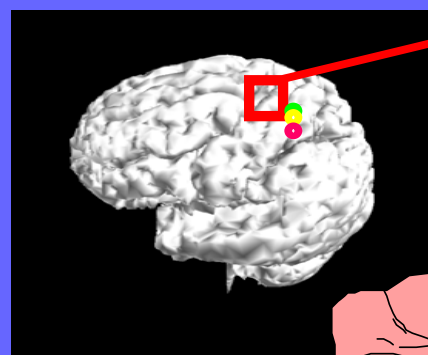
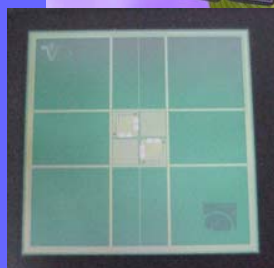
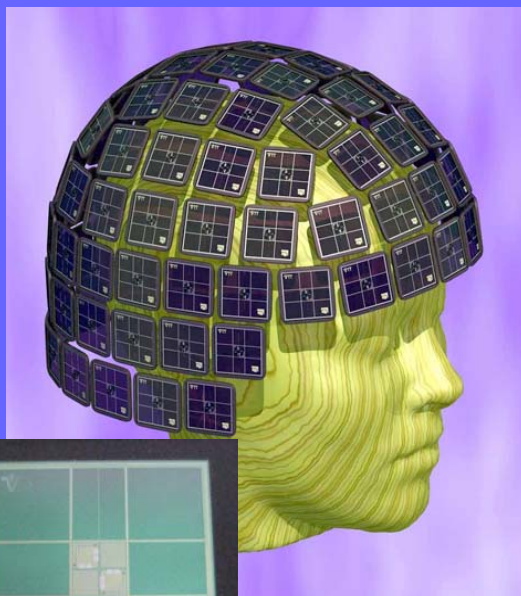


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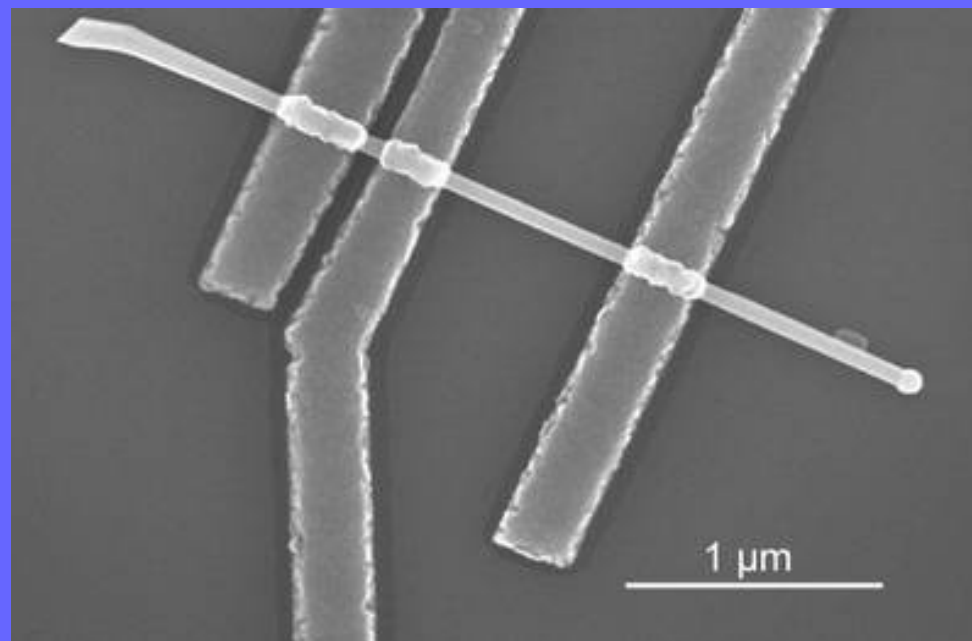
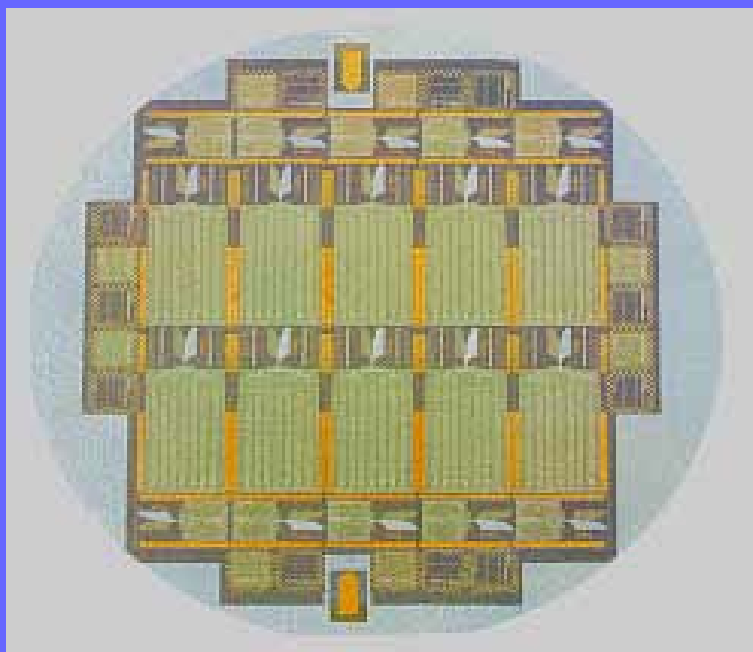


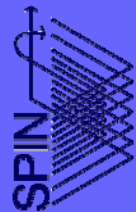
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SQUID for Biomagnetism



Josephson junctions and superconducting circuits for electronics and quantum computation





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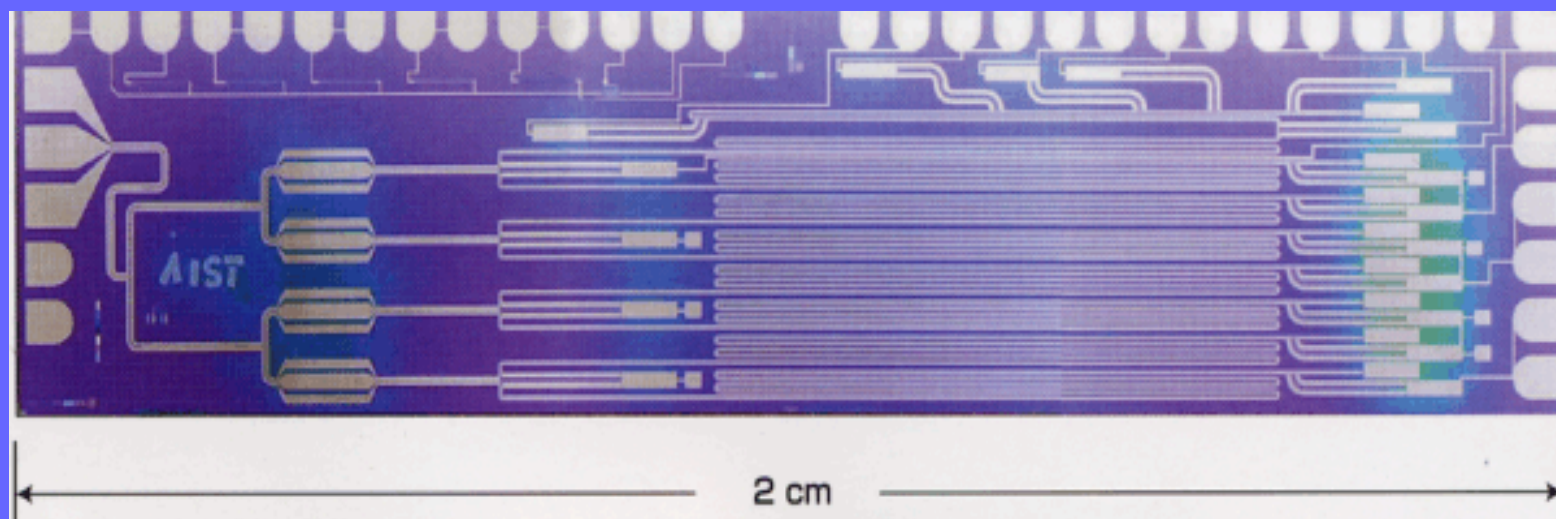


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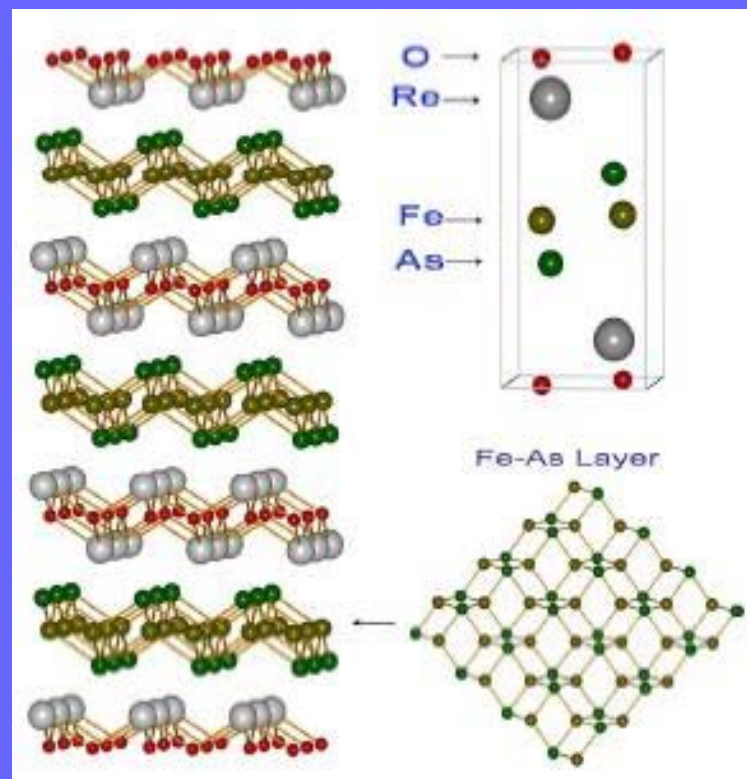


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Voltage standard maintenance

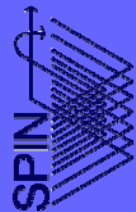


Search for Room Temperature Superconductors !



REOFeAs





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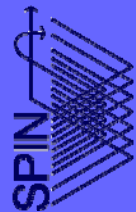
Trieste, April 28th 2011

AVATAR



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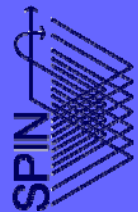
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The need to conquer Pandora is because of room temperature superconductors rocks (based on “unobtainium”, an element with atomic number 120)





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The Superconductive City

