



The Abdus Salam
International Centre for Theoretical Physics



2234-8

**Meeting of Modern Science and School Physics: College for School
Teachers of Physics in ICTP**

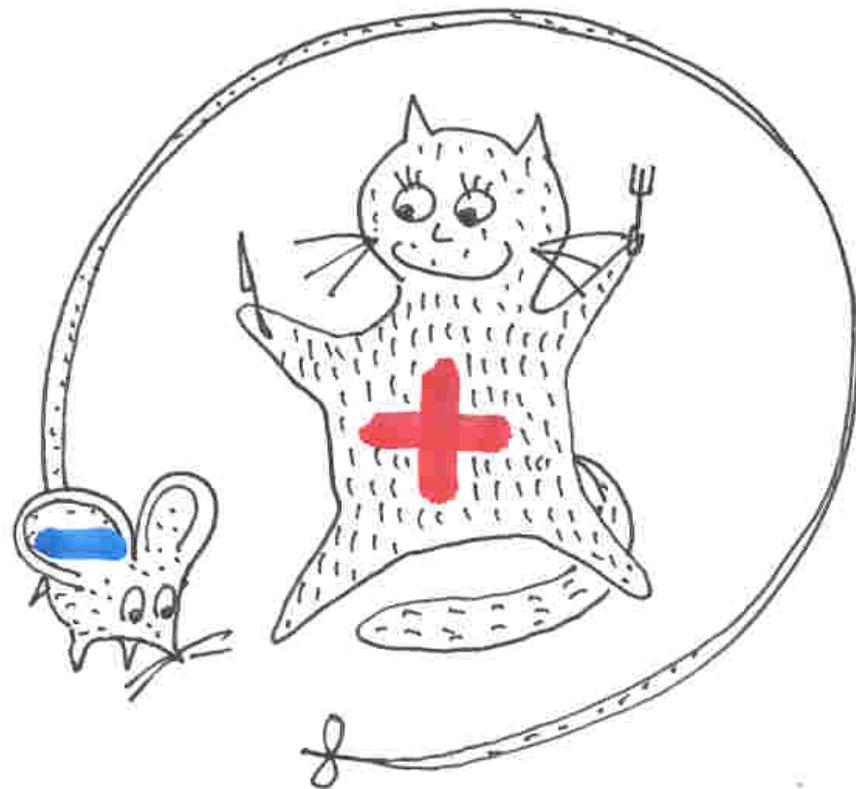
27 April - 3 May, 2011

Excitons: artificial hydrogen atoms in crystals

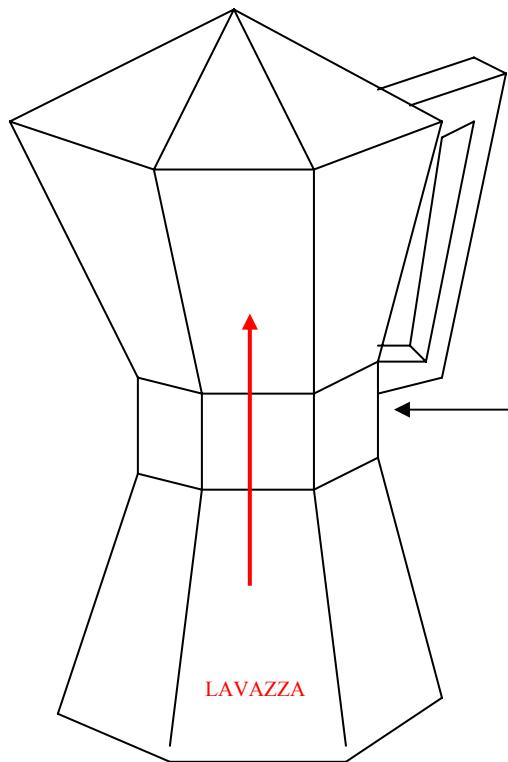
Alexey Kavokin
*University of Southampton
UK*

Excitons: artificial hydrogen atoms in crystals

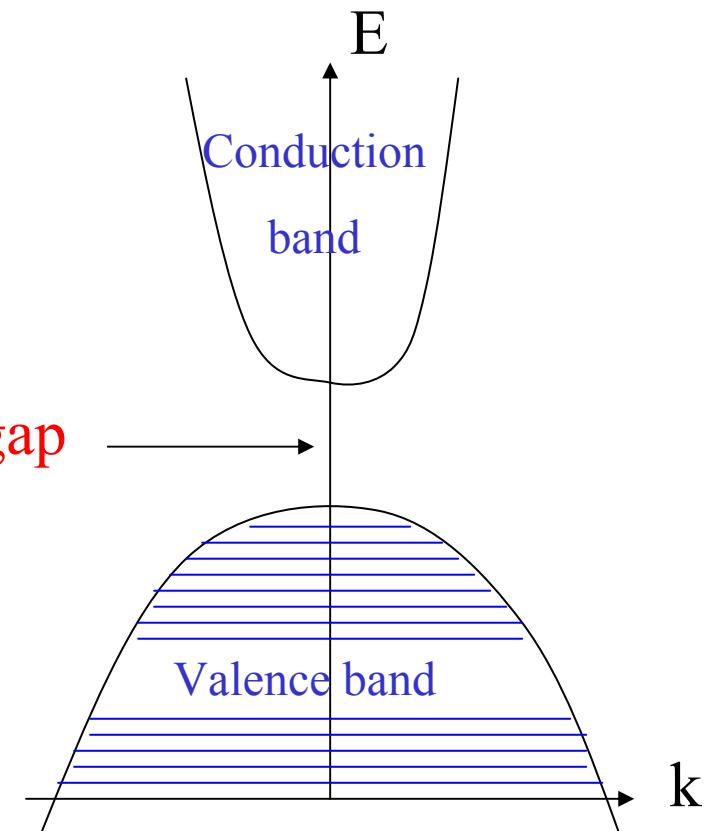
Alexey Kavokin,
University of Southampton



Italian coffee machine

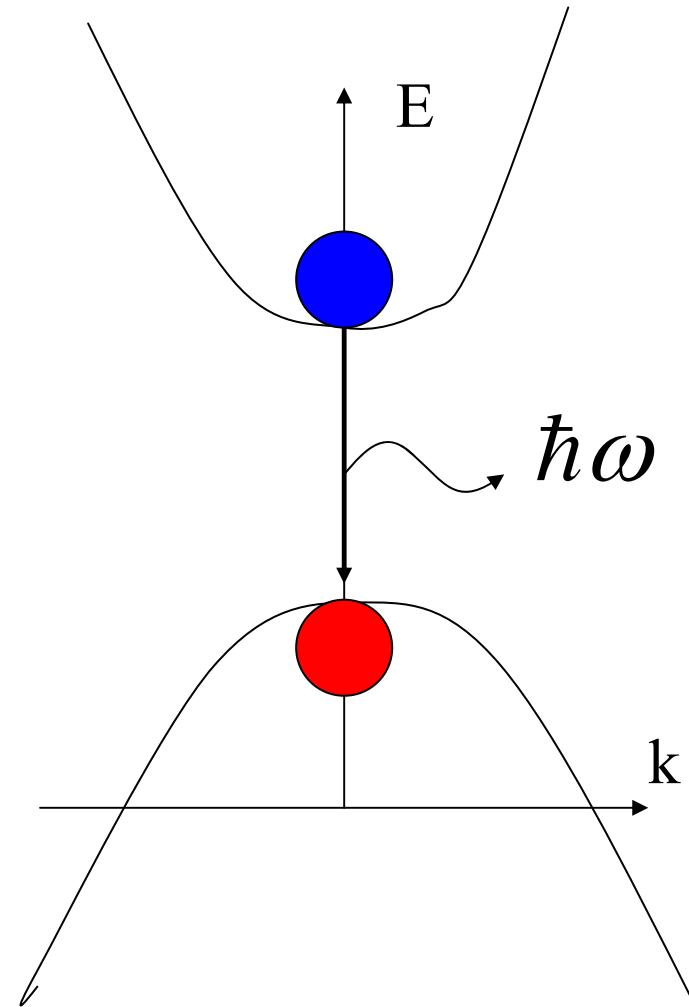
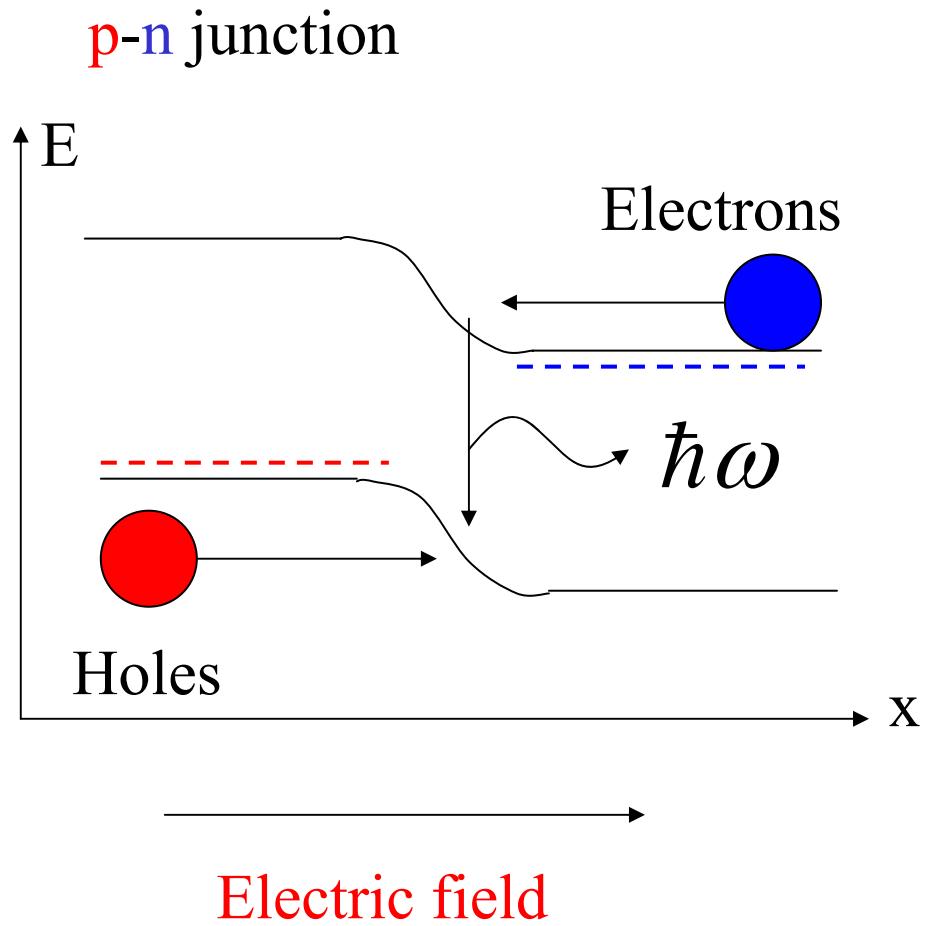


Semiconductor band structure



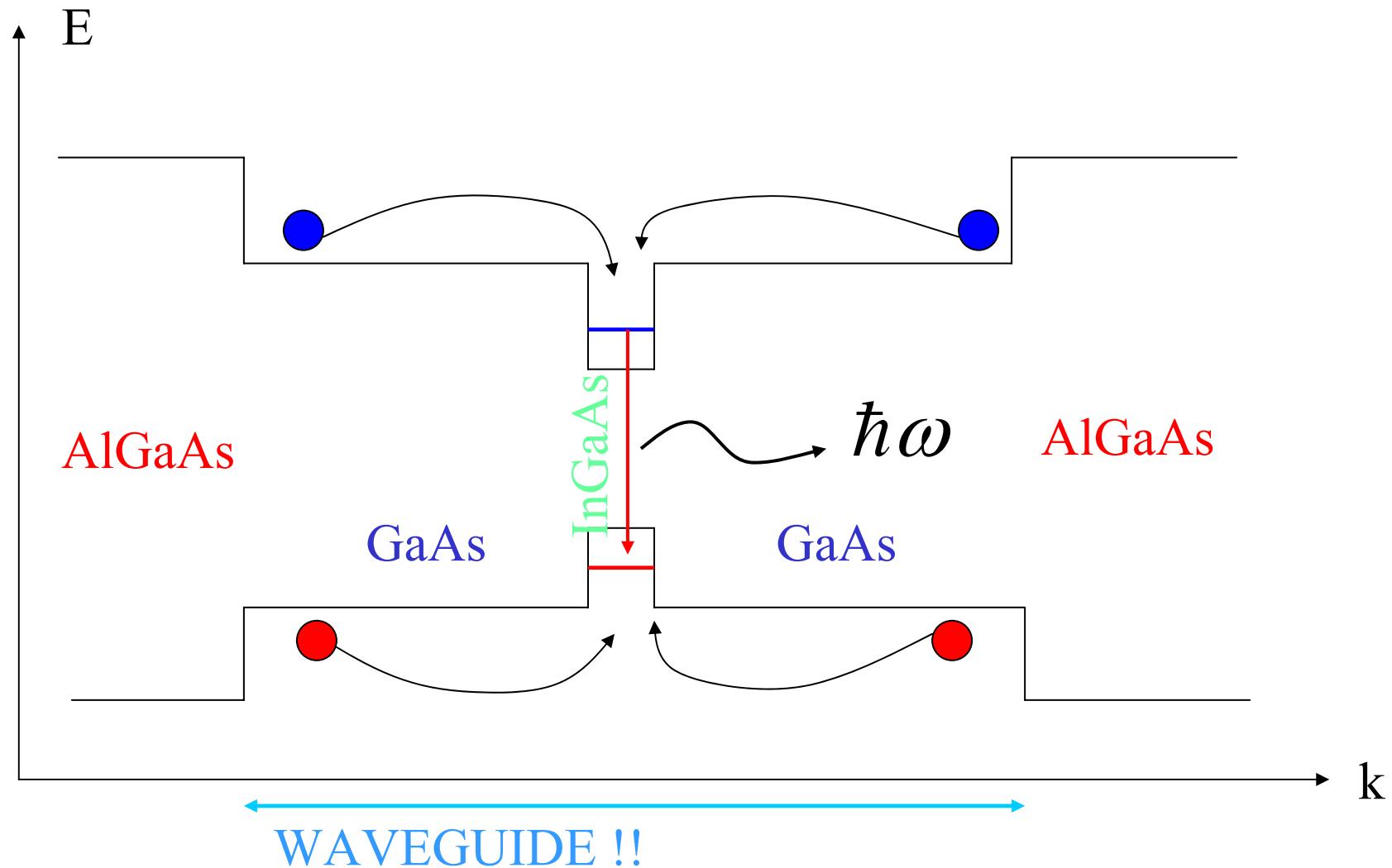
All modern electronics is based on semiconductors...

Semiconductor lasers

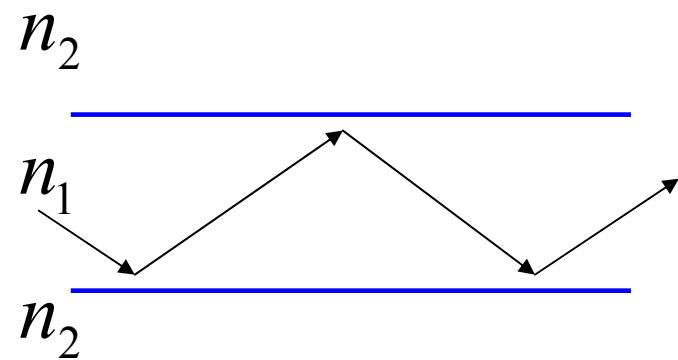


Lasers on heterostructures

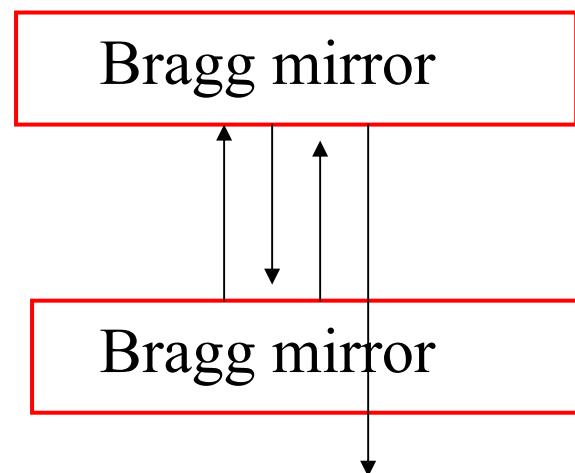
Zh. Alferov, 1972



« Horizontal » and « Vertical » lasers



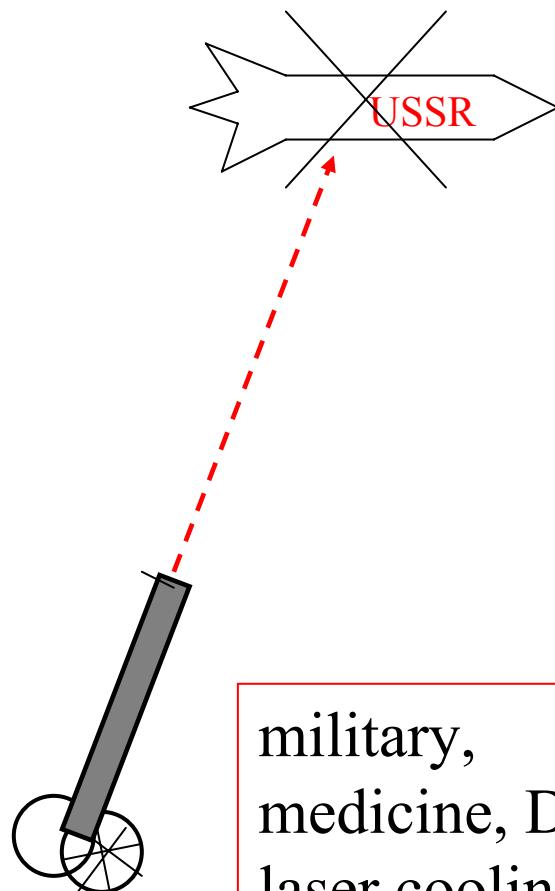
Horizontal lasers emit
in guided modes



Vertical lasers emit
across the mirrors

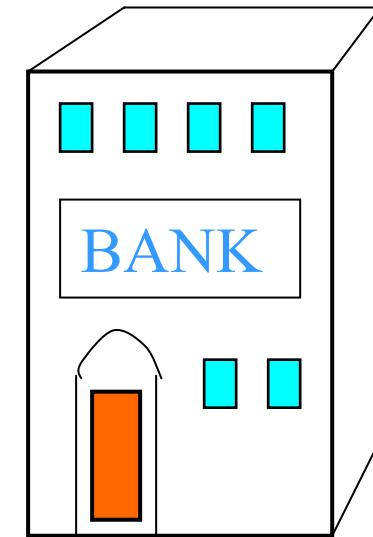
Microlasers, screens...

High power lasers and low power lasers



military,
medicine, DVD,
laser cooling...

A
P
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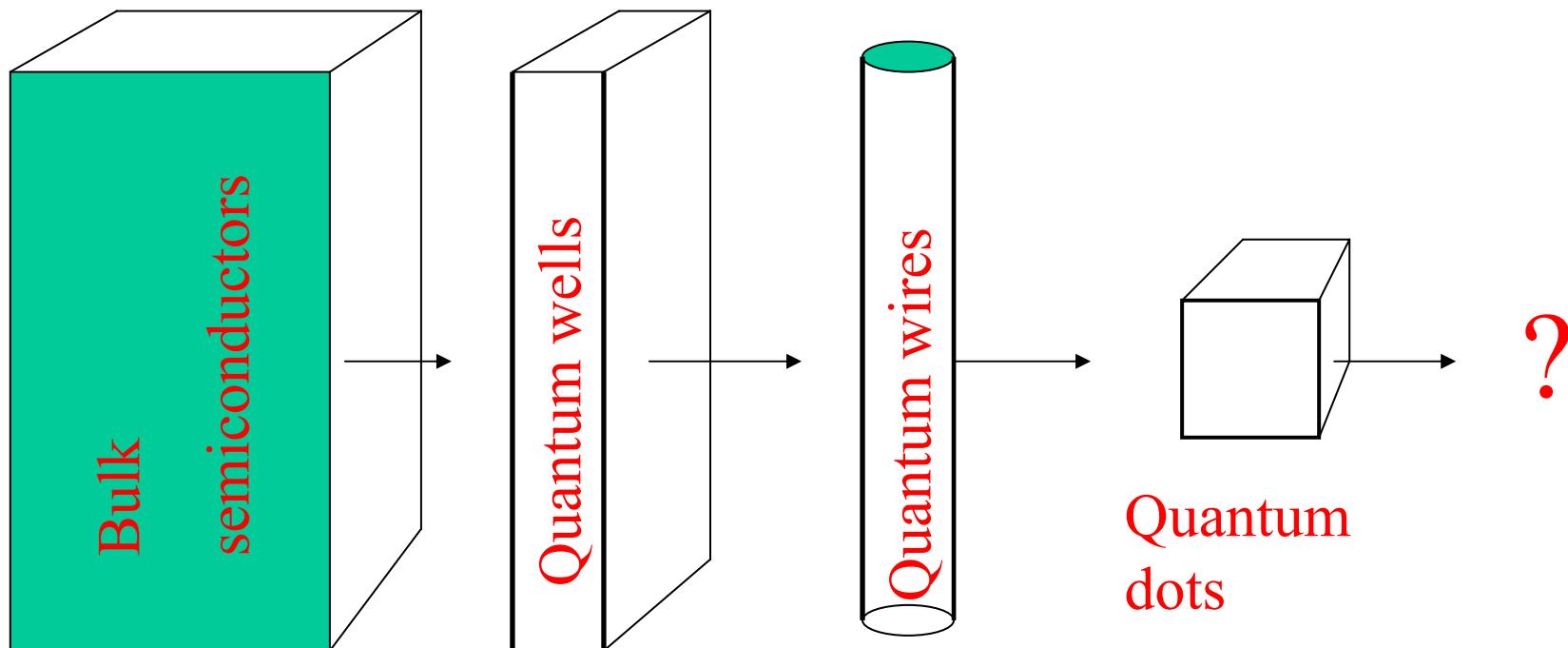


Single photon sources,
quantum cryptography,
quantum computing...

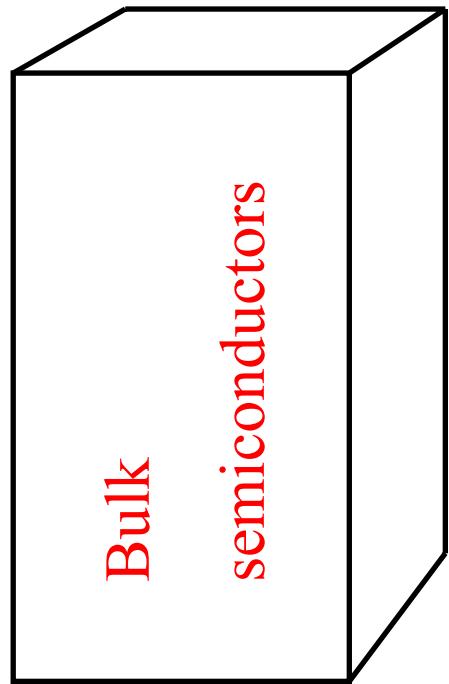
Each laser has a threshold...

Threshold comes from the population inversion...

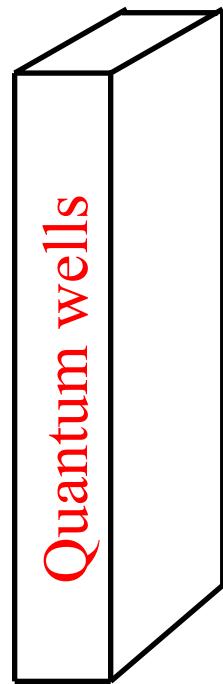
How to make it lower ???



...reducing the density of electronic states !



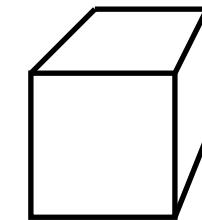
Bulk
semiconductors



Quantum wells



Quantum wires



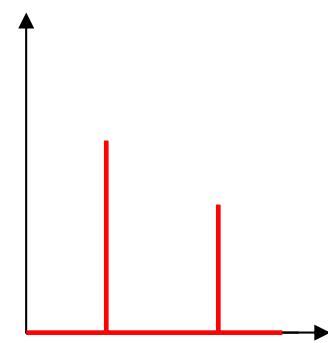
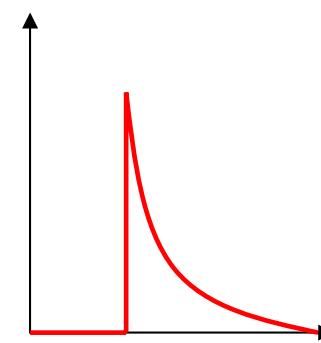
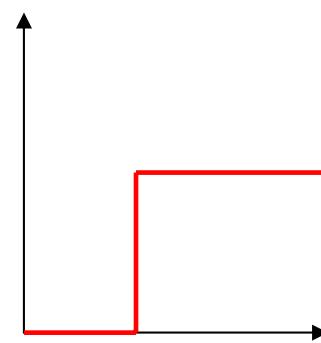
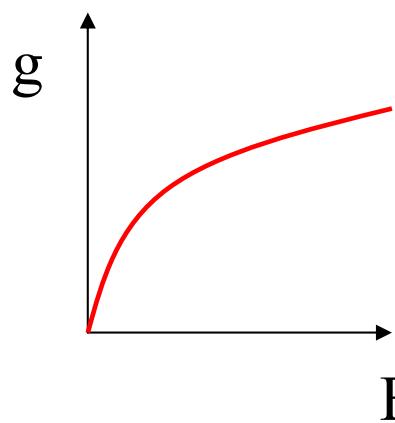
Quantum
dots

3D

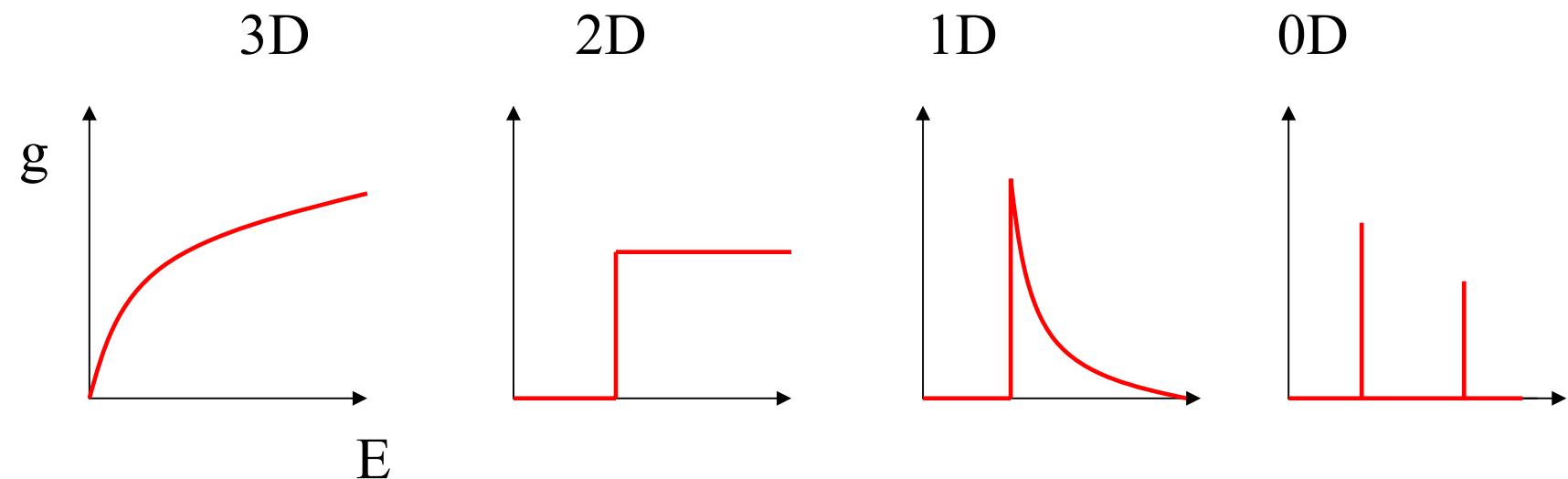
2D

1D

0D

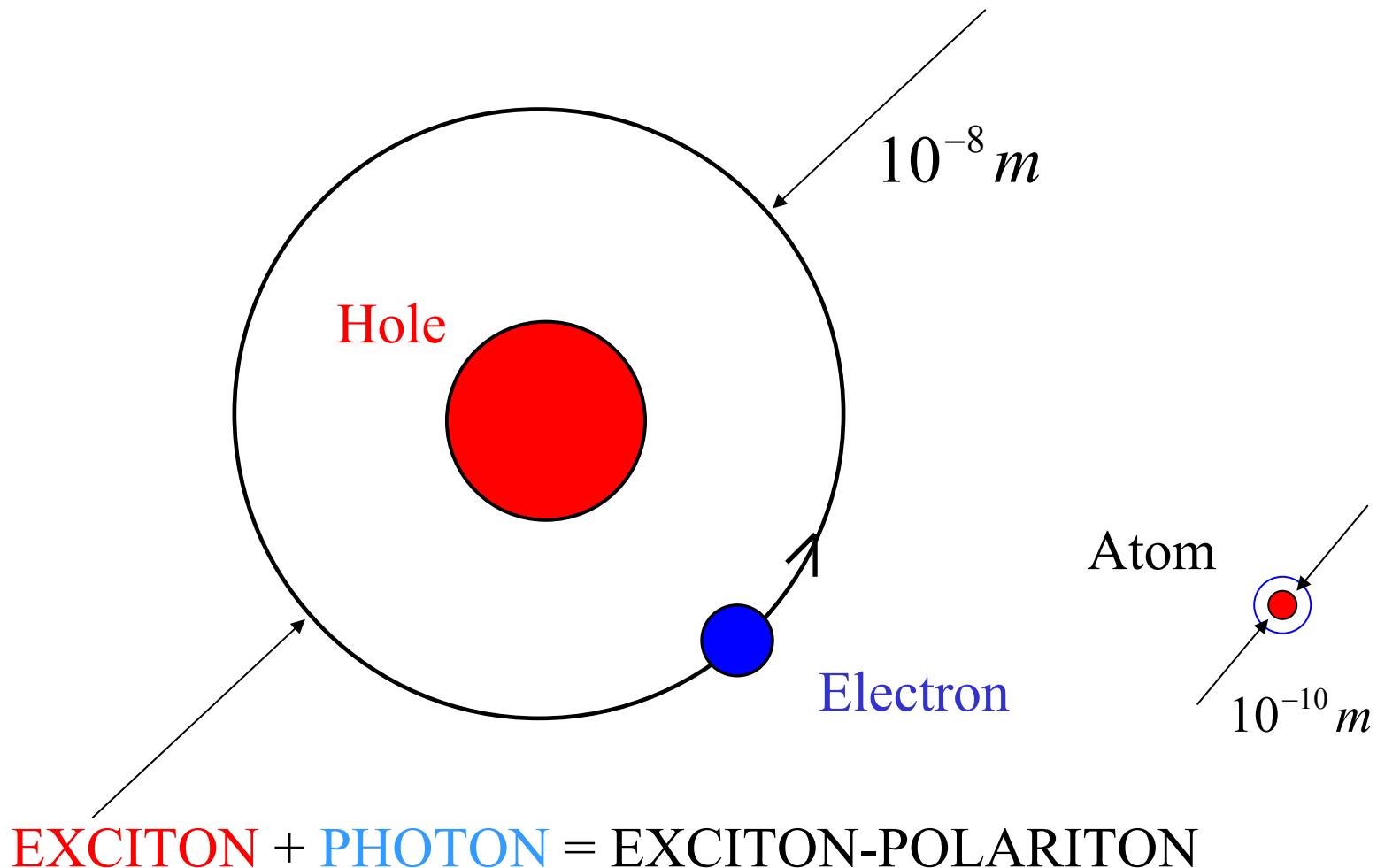


Electron density of states



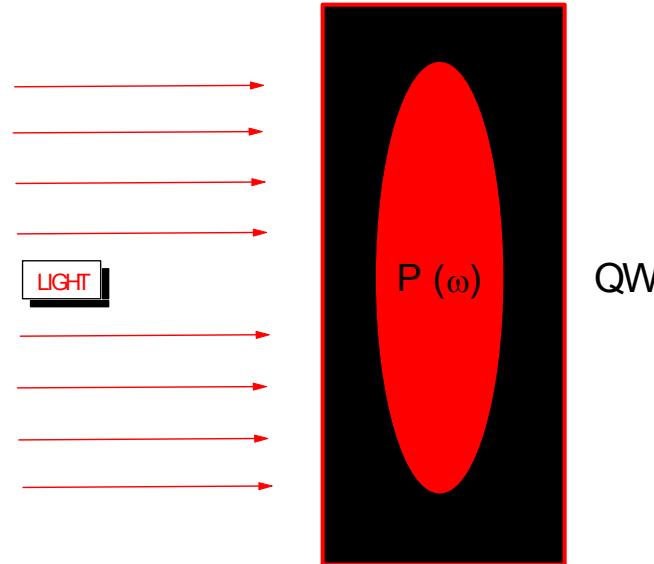
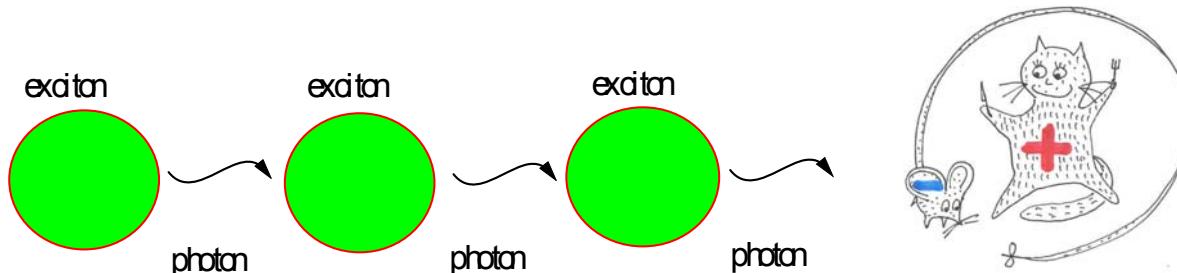
New idea: non-Einstein laser

EXCITON: an artificial ATOM !



Exciton-polaritons: quantum and classical pictures

Virtual chain of
emission-absorption
acts
(quantum
optics)



(semi-classical
approach)

Time- and space-dependent
dielectric polarization

LIGHT-MATTER COUPLING IN SOLIDS

Maxwell equations

$$\nabla \cdot \mathbf{D} = \frac{\rho}{\epsilon_0},$$

$$\nabla \times \mathbf{E} = -\frac{1}{c} \partial_t \mathbf{B},$$

$$\nabla \cdot \mathbf{B} = 0,$$

$$\nabla \times \mathbf{B} = \frac{1}{\epsilon_0 c^2} \partial_t \mathbf{J} + \frac{1}{c^2} \partial_t \mathbf{D}.$$



Electric displacement field $\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} = \epsilon \mathbf{E}$



and Lorentz oscillator model

$$m_0 \ddot{x} + m_0 2\gamma \dot{x} + m_0 \omega_0^2 x = -e E(t)$$

Diagram illustrating the differential equation for the Lorentz oscillator model:

- A horizontal arrow points from the term $m_0 2\gamma \dot{x}$ to the text "damping".
- A vertical arrow points from the term $m_0 \omega_0^2 x$ to the text "potential".
- A diagonal arrow points from the term $-e E(t)$ to the text "driving force".

Solution: $x(t \rightarrow \infty) = \mathcal{A} \cos(\omega t - \phi)$

$$\mathcal{A}(\omega) = \frac{-e E_0}{m_0} \frac{1}{\sqrt{(\omega^2 - \omega_0^2)^2 + (2\gamma\omega)^2}} \quad \phi(\omega) = \arctan \left(\frac{2\gamma\omega}{\omega_0^2 - \omega^2} \right)$$

Hopfield, J. J. (1958). Theory of the contribution of excitons to the complex dielectric constant of crystals, *Phys. Rev.* **112**: 1555.



Hopfield equations

$$\nabla \cdot \mathbf{D} = \frac{\rho}{\epsilon_0},$$

$$\nabla \cdot \mathbf{B} = 0,$$

Maxwell equations

$$\nabla \times \mathbf{E} = -\frac{1}{c} \partial_t \mathbf{B},$$

$$\nabla \times \mathbf{B} = \frac{1}{\epsilon_0 c^2} \partial_t \mathbf{J} + \frac{1}{c^2} \partial_t \mathbf{D}$$

Displacement field

$$\mathbf{D} = \epsilon_B \mathbf{E} + \mathbf{P} = \epsilon \mathbf{E}$$

$$\frac{\epsilon_B}{c^2} \frac{\partial^2}{\partial t^2} \mathbf{E}(\mathbf{r}, t) + \nabla \times \nabla \times \mathbf{E}(\mathbf{r}, t) = -\frac{1}{c^2} \frac{\partial^2}{\partial t^2} \mathbf{P}(\mathbf{r}, t)$$

1st Hopfield equation

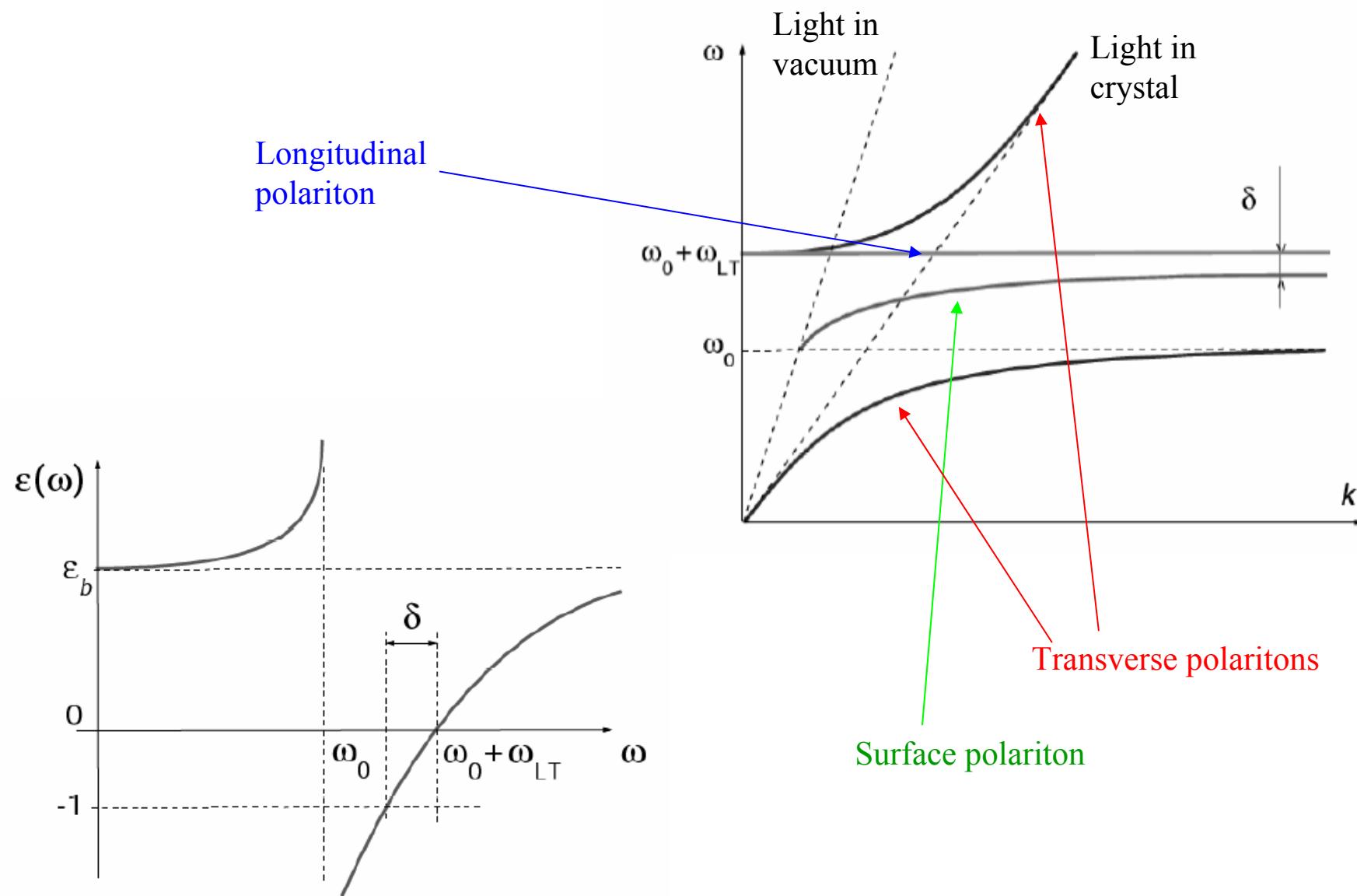
Lorentz oscillator

$$m_0 \ddot{x} + m_0 2\gamma \dot{x} + m_0 \omega_0^2 x = -e E(t)$$

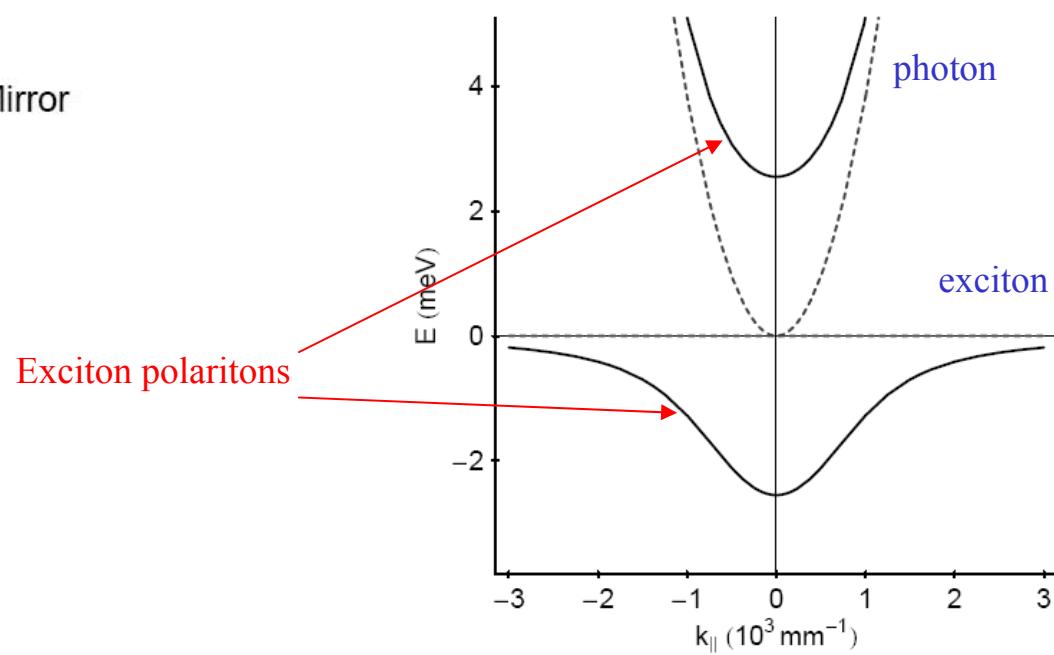
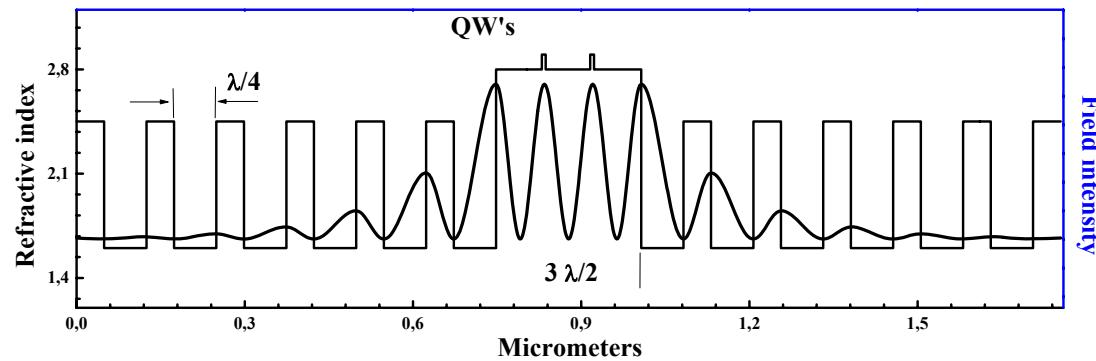
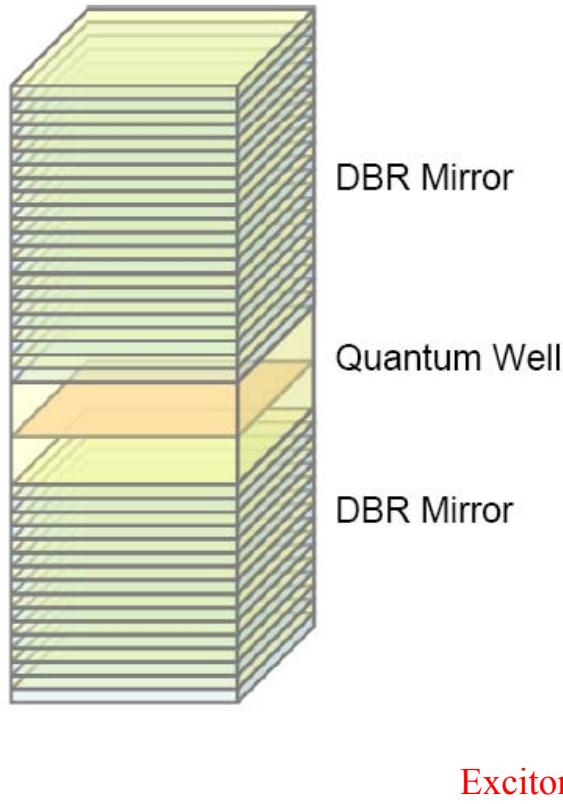
$$\left[\frac{\partial^2}{\partial t^2} + 2\gamma \frac{\partial}{\partial t} + \omega_0^2 - \frac{\hbar \omega_0}{M_x} \nabla^2 \right] \mathbf{P}(\mathbf{r}, t) = \epsilon_B \omega_p^2 \mathbf{E}(\mathbf{r}, t)$$

2nd Hopfield equation

Dispersion of bulk exciton-polaritons in the limit $M_x \rightarrow \infty$ $\gamma \rightarrow 0$

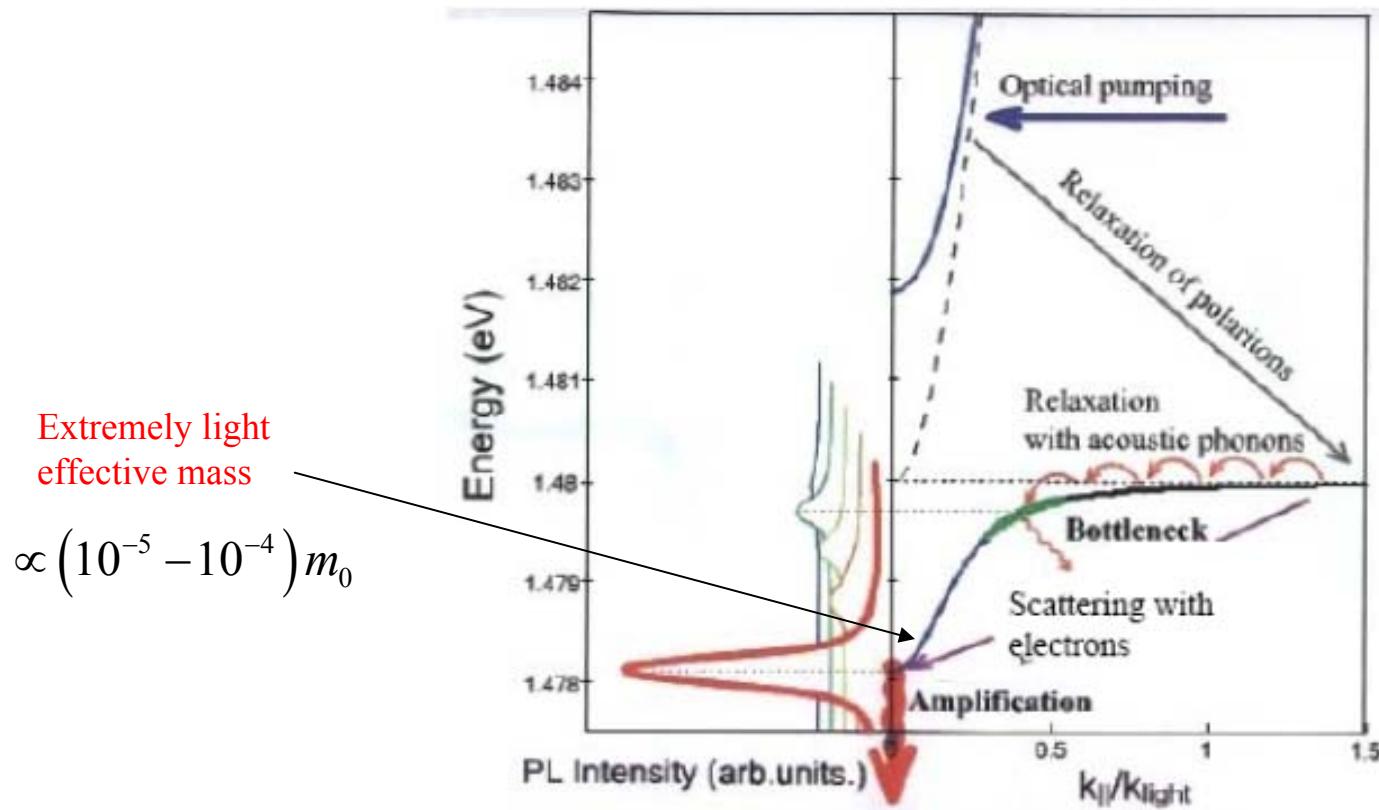


Polaritons in microcavities



Concept of polariton lasing:

A. Imamoglu, et al, Phys. Lett. A 214, 193 (1996).



Photon mode dispersion

$$\frac{\omega}{c}n = \sqrt{\left(\frac{2\pi}{L}\right)^2 + k_{\parallel}^2}$$

Optically or electronically excited exciton-polaritons relax towards the ground state and Bose-condense there. Their relaxation is stimulated by final state population. The condensate emits spontaneously a coherent light

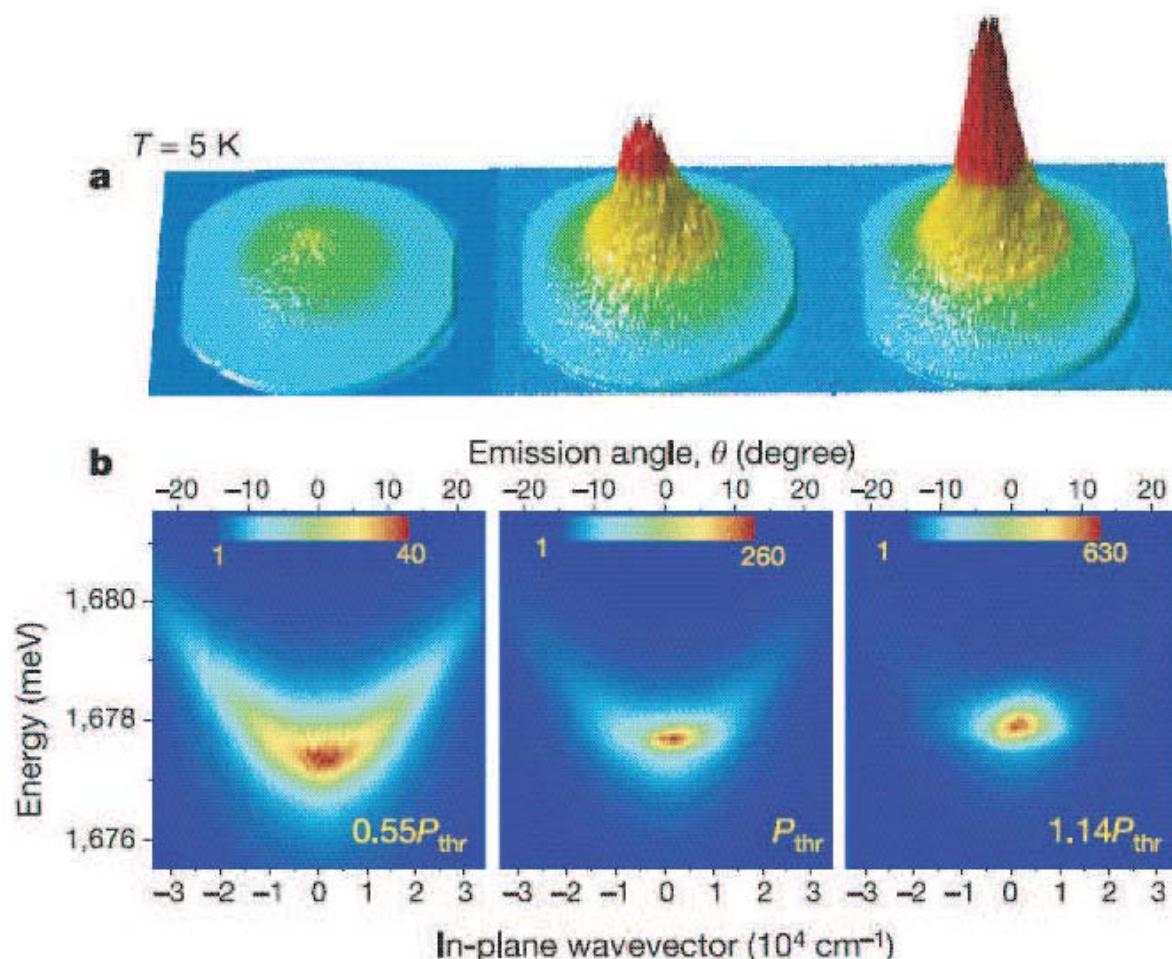
Polariton lasing vs polariton BEC

Criteria	Polariton laser	Polariton BEC
Strong coupling regime	yes	yes
Formation of a macroscopically occupied coherent polariton state	yes	yes
Stimulated scattering	yes	yes
Polaritons with low k at thermal equilibrium	not necessarily	yes
Thermodynamic phase transition	not necessarily	yes
Order parameter build up	not necessarily	yes
Spontaneous symmetry breaking	not necessarily	yes

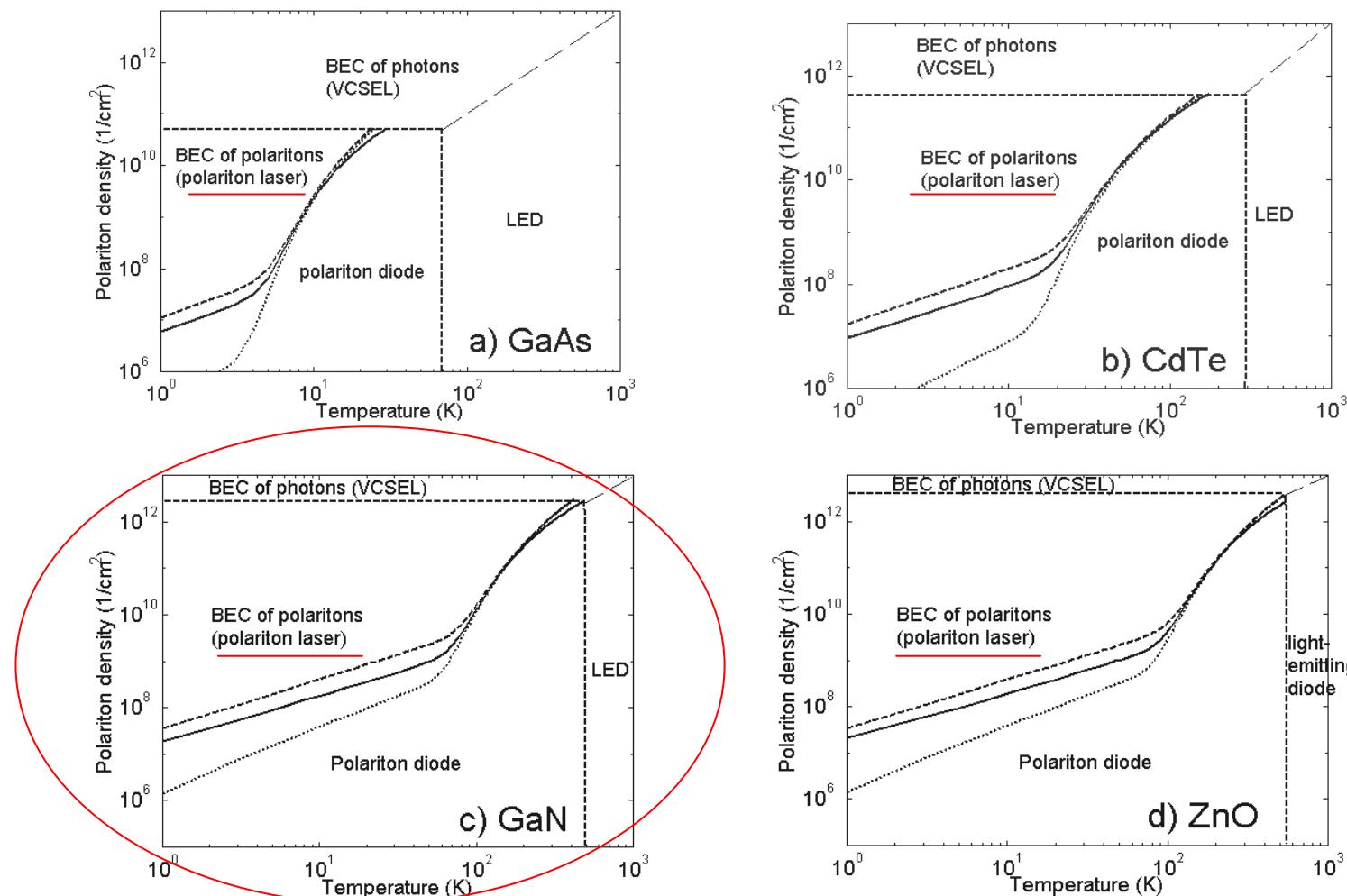
ARTICLES

Bose-Einstein condensation of exciton polaritons

J. Kasprzak¹, M. Richard², S. Kundermann², A. Baas², P. Jeambrun², J. M. J. Keeling³, F. M. Marchetti⁴, M. H. Szymańska⁵, R. André¹, J. L. Staehli², V. Savona², P. B. Littlewood⁴, B. Deveaud² & Le Si Dang¹



Phase diagrams for BEC of exciton-polaritons in different model cavities



Solid red lines show the critical concentration N_c versus temperature of the polariton KT phase transition. Dotted and dashed blue lines show the critical concentration N_c for quasi condensation in $100 \mu\text{m}$ and 1 meter lateral size systems, respectively.

GaN-based microcavities

Exciton binding energy = 28 meV

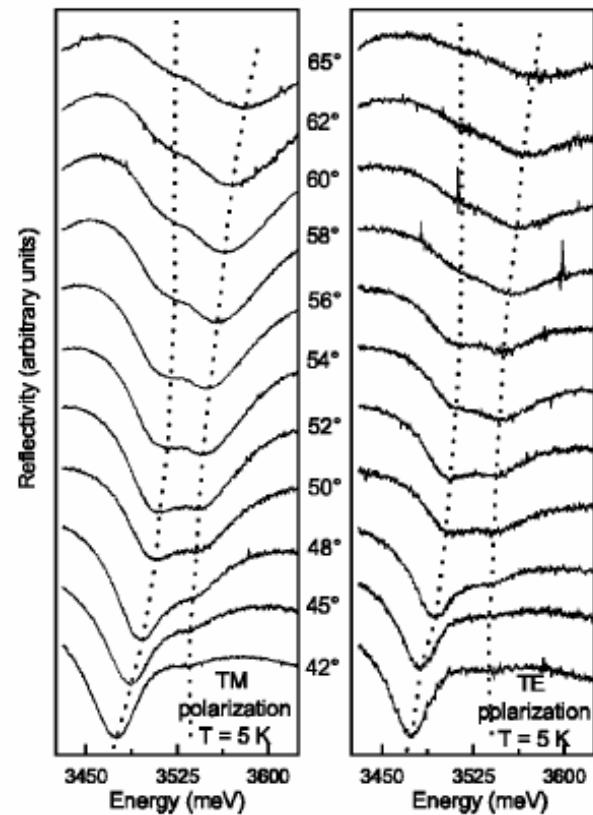


Excitons are stable at room temperature!

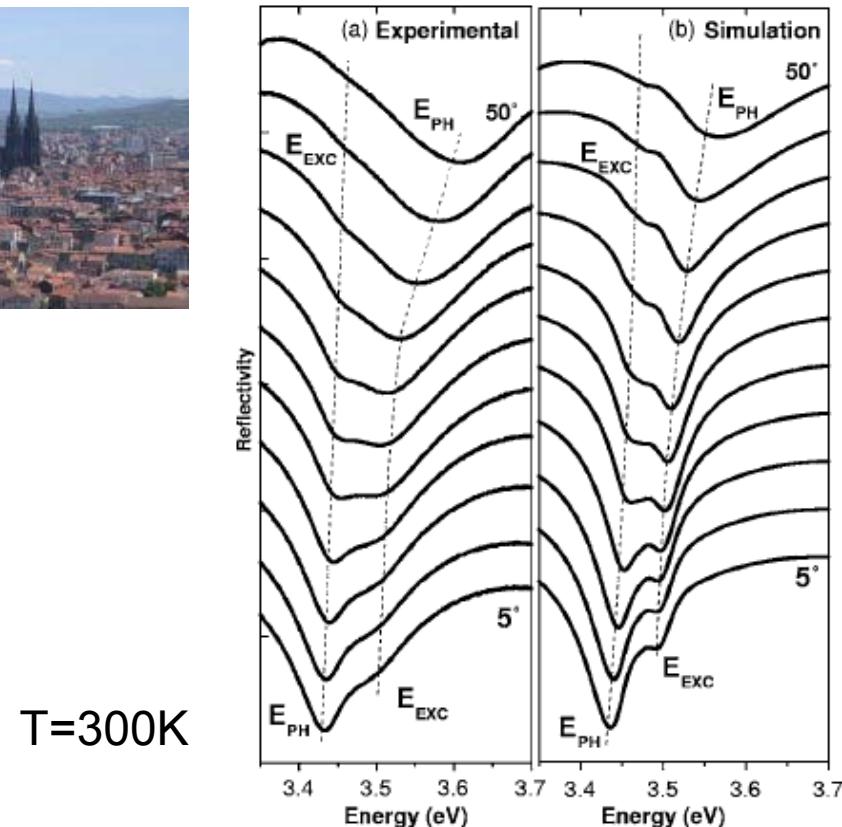


Strong coupling in GaN-based microcavities: experimental realisation

CLERMONT (1999-2003)
CLERMONT2 (2003-2007)
CLERMONT4 (2009-2013)



T=5K

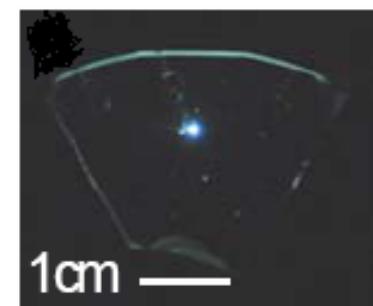
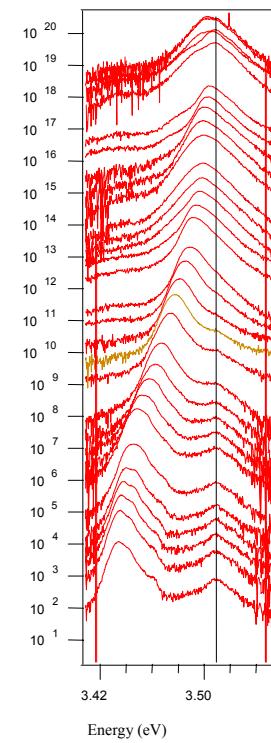
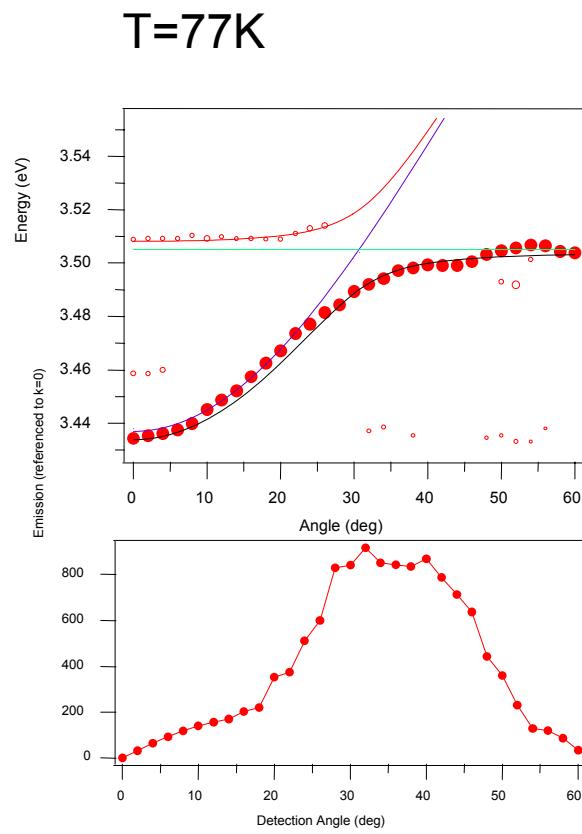


N. Antoin-Vincent et al, PRB 68, 153313 (2003)

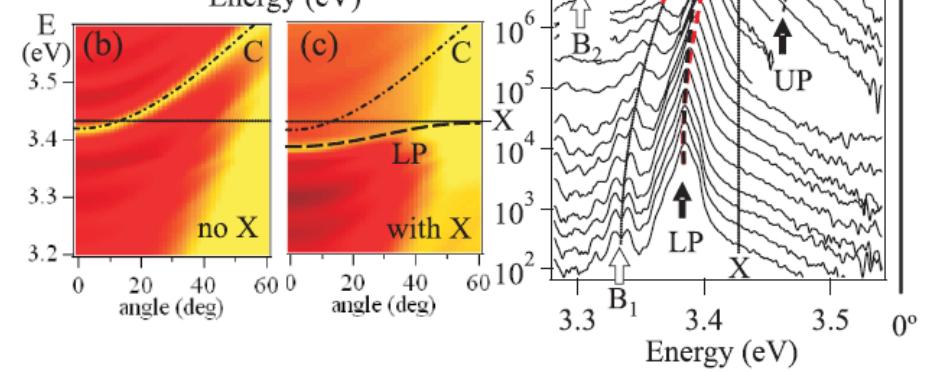
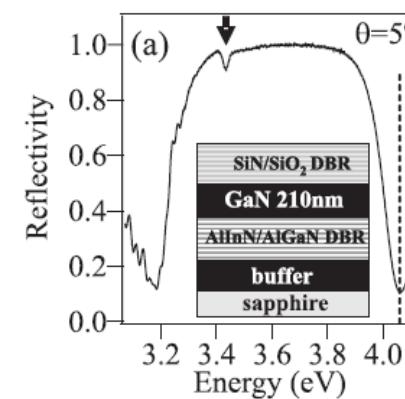
I. Sellers et al, Phys. Rev. B73, 033304 (2006)

Polariton lasing in GaN microcavities

S. Christopoulos *et al.*, *Phys. Rev. Lett.* **98**, 126405 (2007).

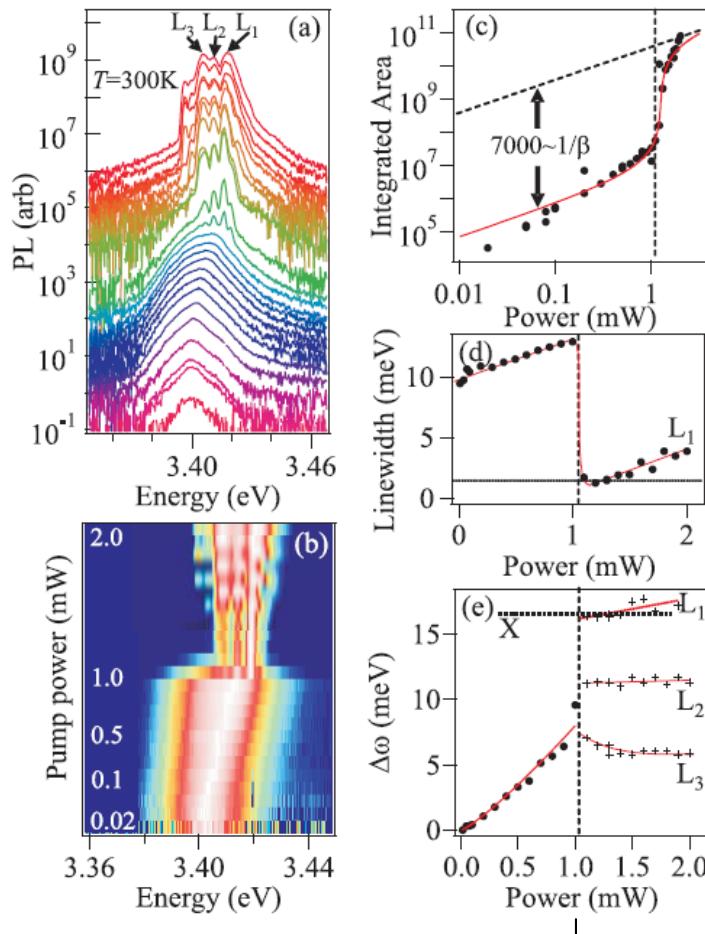


T=300 K



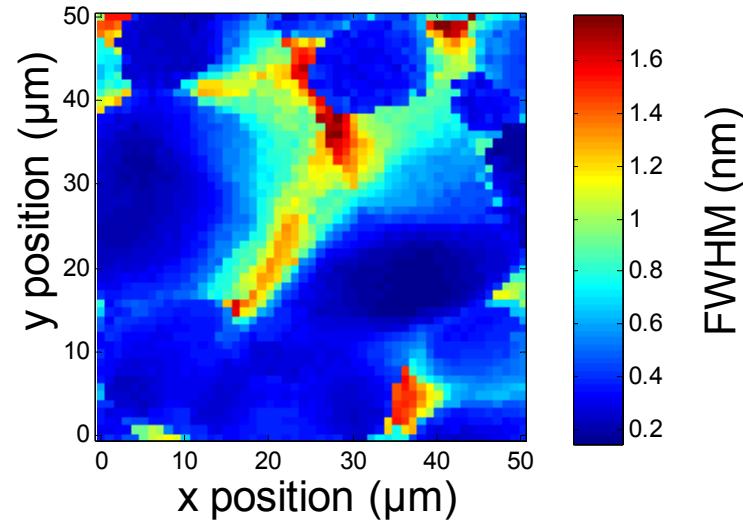
S. Christopoulos, G. Baldassarri von Högerthal, A. J. Grundy, P. G. Lagoudakis, A. V. Kavokin, J. J. Baumberg, G. Christmann, R. Butté, E. Feltin, J.-F. Carlin, and N. Grandjean, *Room-Temperature Polariton Lasing in Semiconductor Microcavities*, **Phys. Rev. Lett.** **98**, 126405 (2007).

Polariton lasing



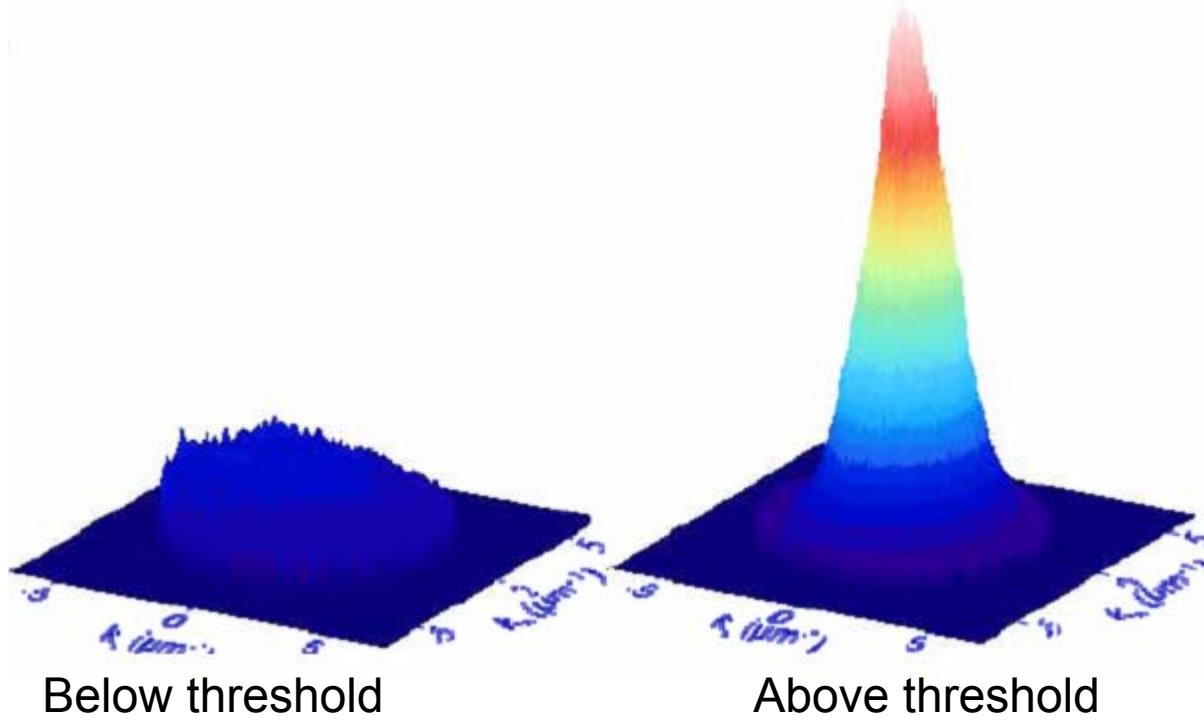
$$I_{\text{th}} = 1.0 \text{ mW} \quad N_{3D} \sim 8 \times 10^{17} \text{ cm}^{-3}$$

Photonic disorder



Micro-transmission maps of the cavity linewidth of a GaN bulk microcavity, with a negatively-detuned cavity of resonant wavelength around 418nm. The spot size is $8\mu\text{m}$.

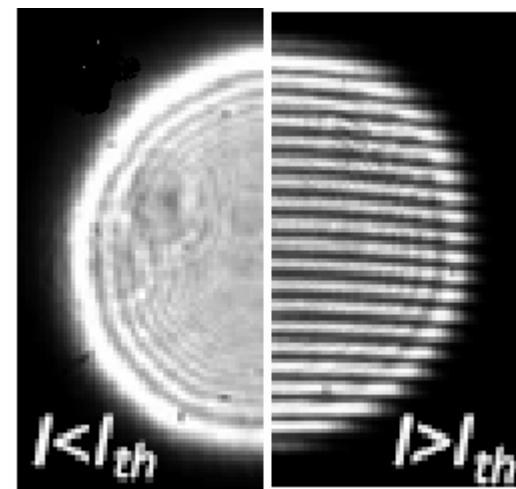
Build-up of the condensate in a GaN microcavity



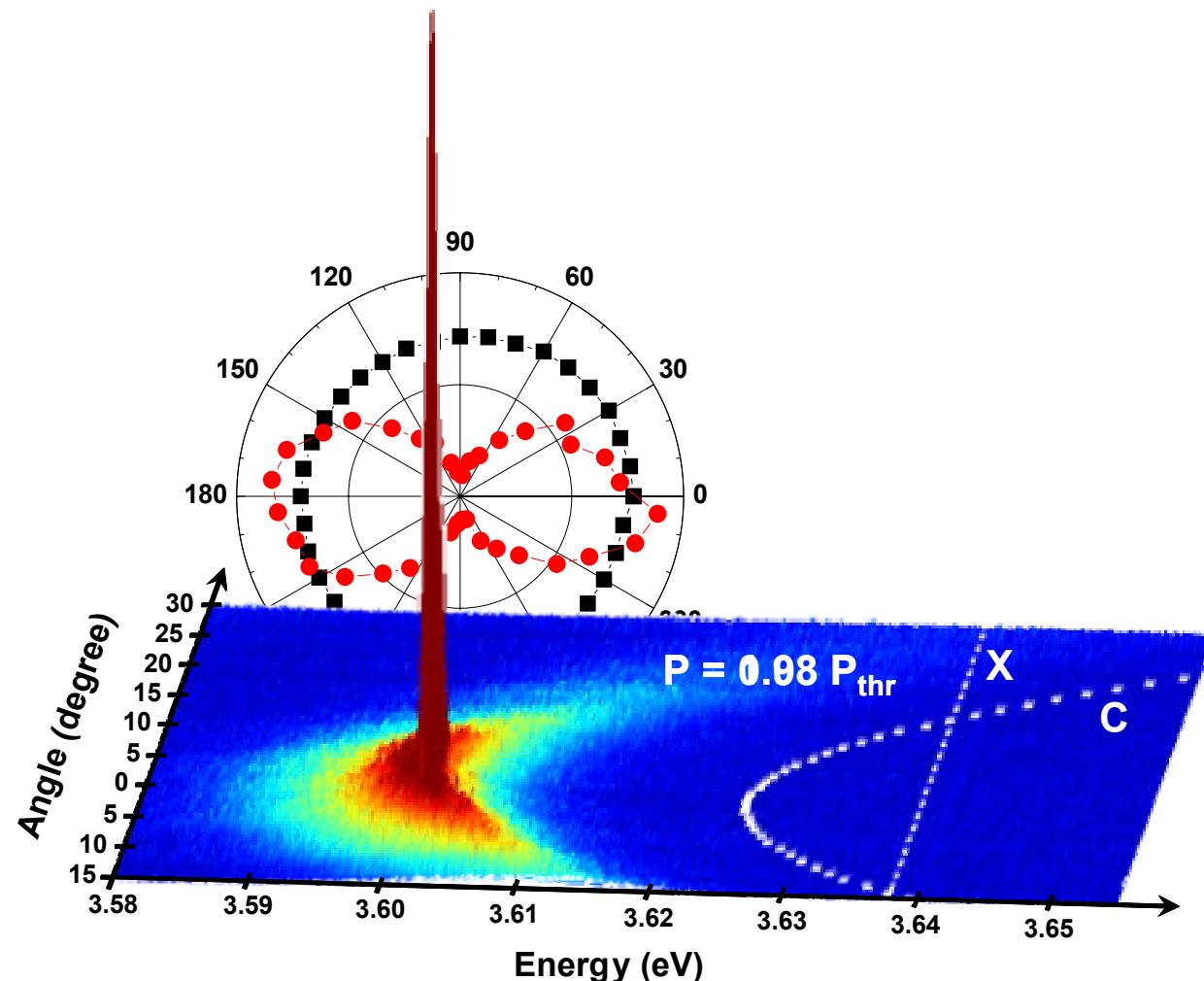
Below threshold

Above threshold

J. J. Baumberg, A. V. Kavokin, S. Christopoulos, A. J. Grundy, R. Butté, G. Christmann, D. D. Solnyshkov, G. Malpuech, G. Baldassarri Höger von Högersthal, E. Feltin, J.-F. Carlin, and N. Grandjean, Spontaneous Polarization Buildup in a Room-Temperature Polariton Laser, **Phys. Rev. Lett.** **101**, 136409 (2008).



Room temperature polariton lasing



**Nonlinear emission
peaked at $k_{\parallel} = 0$**

**Linearly polarized
emission above
threshold (polarization
degree: 80%)**

**Emission still in the
strong coupling
regime**

**Polarization splitting:
~150 μeV**