



2267-1

Joint ITER-IAEA-ICTP Advanced Workshop on Fusion and Plasma Physics

3 - 14 October 2011

Introduction to Fusion Leading to ITER

SNIPES Joseph Allan

Directorate for Plasma Operation Plasma Operations Group POP, Science Division Building 523/023, Route de Vinon sur Verdon 13115 St Paul lez Durance FRANCE

Introduction to Fusion Leading to ITER

J A Snipes

ITER Organization 13115 St. Paul-lez-Durance, France

Acknowledgements: D J Campbell, many colleagues in the ITER IO, ITER Members

The views and opinions expressed herein do not necessarily reflect those of the ITER Organization.

iter china eu india japan korea russia usa | ICTP Advanced Workshop on Fusion and Plasma Physics, Trieste, Italy 3 – 14 October 2011 | Page 1

- Introduction to thermonuclear fusion
- Basics of magnetic confinement fusion the tokamak
- Some key parameters for magnetic confinement fusion in tokamaks
- How the results from existing tokamaks led to ITER for the next step in fusion research

Fusion – the fundamental principle

 Energy gain from fusion, like fission, is based on Einstein's equation:

$E = \Delta mc^2$

-mass loss for DT reactions corresponds to ~ 0.4%

- As illustrated, energy gain per unit mass is greater for fusion
 - -energy gain/ reaction:

DT fusion: 17.6 MeV U fission: ~200 MeV



Essential Fusion Reactions



+ 20% of Energy (3.5 MeV)

+ 80% of Energy (14.1 MeV)

The D-T fusion reaction is the simplest to achieve under terrestrial conditions:

 $^{2}D + ^{3}T \Rightarrow ^{4}He (3.5 \text{ MeV}) + ^{1}n (14.1 \text{ MeV})$

• Two other important reactions for DT fusion:

 $^{1}n + ^{6}Li \Rightarrow ^{4}He + ^{3}T + 4.8 MeV$

 $^{1}n + ^{7}Li \Rightarrow ^{3}He + ^{3}T + ^{1}n - 2.5 MeV$

- these reactions will allow a fusion reactor to breed tritium

Fusion Power Density vs Temperature

$1 \text{ keV} = 1.16 \times 10^7 \text{ K}$



 High temperatures (~10 keV) are required for significant thermonuclear fusion energy production ⇒ dealing with plasmas!

Basics of Magnetic Confinement Fusion: The Tokamak



Plasma Toroidal Magnetic Confinement

 Magnetic fields cause ions and electrons to spiral around the field lines:

$$F = q(E + v \times B)$$

 in a toroidal configuration plasma particles are lost to the vessel walls by relatively slow diffusion across the field lines



A special version of this torus is called a tokamak:

'toroidal chamber' and 'magnetic coil' (Russian)

Magnetic Confinement in a Tokamak

The Tokamak:

- External coils
 - to produce a toroidal magnetic field
- Transformer with primary winding
 - to produce a toroidal current in the plasma
 - this plasma current creates a poloidal magnetic field
- Finally, poloidal coils
 - to control the position and shape of the plasma



JET: Joint European Torus

- JET is currently the largest tokamak
 - Major/minor radius: 3 m/ 1 m
 - Plasma volume ~100 m³
 - Toroidal field: 3.4 T
 - Plasma Current: 7 MA
- In DT experiments in 1997, a peak fusion power of 16 MW was produced



15m

JET - the largest existing Tokamak Internal View



Magnetic Confinement in a Tokamak



- In configurations with only a toroidal field, ions and electrons drift vertically in opposite directions:
- An additional poloidal field allows particles to follow helical paths, cancelling the drifts
- "Winding number" of helix is an important stability parameter for the system:

$$q_{c} = \frac{aB_{\phi}}{RB_{\theta}} \sim \frac{a^{2}B_{\phi}}{RI_{\rho}}$$

- q_c = "cylindrical" safety factor
- R/a = aspect ratio

Plasma Equilibrium in a Tokamak



• Formal definition of safety factor:

- absolute value of *q* and its variation across the plasma radius are important in plasma stability
- by elongating the plasma, more current can be squeezed into the plasma ring at fixed *q*:

$$\kappa = \frac{b}{a}$$

- *κ* also turns out to have important consequences
 for plasma stability
- Typically the pressure (temperature, density) and current profiles are peaked on the plasma axis:
 - the profile of q is then the inverse, with $q(0) \sim 1$

Many Plasma Shapes Have Been Investigated



PBX-M

ASDEX UPGRADE

Plasma shape affects confinement and stability properties

JT-60U

Plasma fusion performance

| <i>Temperature - T_i:</i> | 1-2 × 10 ⁸ K (10-20 keV) (~10 × temperature of sun's core) |
|-------------------------------------|---|
| Density - n _i : | 1 × 10 ²⁰ m ⁻³ (~10 ⁻⁶ of atmospheric particle density) |
| Enorgy confinement time | σ : fow seconds (∞ ourrent \times radius ²) |

Energy confinement time - τ_E : few seconds (\propto current \times radius²) (ITER plasma pulse duration ~1000s)



Fusion Triple Product

- Existing experiments have achieved nTτ values
 ~ 1×10²¹ m⁻³skeV
 ~ Q_{DT} = 1
- JET and TFTR have produced DT fusion powers of >10MW for ~1s
- ITER is designed to a scale which should yield
 Q_{DT} ≥ 10 at a fusion power of 400 500MW for 300-500s



Plasma Heating

- Tokamaks have a built in heating scheme: "Ohmic" heating by the plasma current
 - but plasma resistivity varies as $\rm T_e^{-3/2},$ so heating power declines with increasing $\rm T_e$
 - so Ohmic plasma temperatures of several keV are possible, but additional heating is required to achieve 10-20 keV

• Two basic heating schemes:

- injecting neutral particle beams
- injecting radiofrequency waves because the plasma refractive index depends on density and magnetic fields, several RF options are possible

Each heating technique also provides some current drive

Injection of Neutral Particle Beams

• Neutral beam injection (NBI):

- intense particle beams are accelerated, neutralized and injected into plasma
- $E_b \sim 100 \text{ keV}$, P_b up to 40MW in TFTR
- very effective:
 - heating
 - current drive
 - fuelling
 - rotation drive

• For ITER:

- E_b ~ 1 MeV is required to penetrate plasma/ drive current
- →negative ion source technology
- higher energy ⇒ little fuelling, little rotation drive



Radiofrequency Heating

- Ion Cyclotron Radiofrequency Heating (ICRF):
 - launched at frequencies ~ $\omega_{ci} \Rightarrow$ f ~ 50 MHz
 - technology conventional
 - wave coupling to plasma problematic penetration through edge
- Electron Cyclotron Resonance Heating (ECRH):
 - launched at frequencies ~ $\omega_{ce} \Rightarrow$ f > 100 GHz
 - source technology non-conventional: "gyrotrons"
 - coupling, absorption, space localization very good
- Lower Hybrid Heating/ Current Drive (LHCD):
 - "lower hybrid" a complex wave resonance in plasma: $f \sim 5 \text{ GHz}$
 - technology fairly conventional (source: klystrons)
 - wave coupling to plasma problematic penetration through edge

Current Drive

Current drive provides:

- replacement of the transformer drive \Rightarrow towards steady-state plasma
- manipulation of the current profile to improve confinement/ stability
- direct suppression of plasma instabilities

- Current drive efficiency (η_{CD} = driven current/input power):
 - typically increases with T_e
 - for beams, also increases with E_b

 \Rightarrow favourable for ITER



C Gormezano et al, Nucl Fusion 47 S285 (2007)

Plasma Confinement: H-mode

- It is found that the plasma confinement state (τ_E) can bifurcate:
 - two distinct plasma regimes, a low confinement (L-mode) and a high confinement (H-mode), result
 - this phenomenon has been shown to arise from changes in the plasma flow in a narrow edge region, or pedestal



ITER Physics Basis I

• Predictions of fusion performance in ITER rely essentially on a small number of physics rules:

• H-mode energy confinement scaling (IPB98(y,2)):

$$\begin{aligned} \tau_{\text{E},\text{th}}^{98(\text{y},2)} = 0.144 \, \mathsf{I}^{0.93} \mathsf{B}^{0.15} \mathsf{P}^{-0.69} \mathsf{n}^{0.41} \mathsf{M}^{0.19} \mathsf{R}^{1.97} \varepsilon^{0.58} \kappa^{0.78} \ \text{(s)} \\ \tau_{\mathsf{E}} \propto \mathsf{I} \mathsf{R}^2 \mathsf{P}^{-2/3} \\ \text{NB:} \, \mathsf{H}_{98(\text{y},2)} = \tau_{\mathsf{E},\text{th}}^{\exp} \, / \, \tau_{\mathsf{E},\text{th}}^{98(\text{y},2)} \end{aligned}$$

• H-mode threshold power:

$$P_{LH} = 0.098 M^{-1} B^{0.80} \overline{n}_{20}^{0.72} S^{0.94} \quad (MW)$$

(i.e., a certain level of power needs to flow across the plasma boundary to trigger an H-mode)

Fusion Performance Depends on Confinement

- → Uncertainty in achievable energy $(H_{H98(y,2)})$ and helium particle (τ^*_{He}/τ_E) confinement gives a large uncertainty in resulting fusion performance
- ► Q=10 Inductive Scenario uses $H_{H98(y,2)} = 1$ and $\tau^*_{He}/\tau_E = 5$ based on empirical data from existing tokamaks
- Too much core helium ash accumulation could reduce fusion performance
- Q > 50 is not excluded within the uncertainty



How is ITER scale determined ?

- Energy confinement time is one of many parameters studied in a wide range of tokamak experiments
 - multi-tokamak experimental database provides <u>scaling</u> prediction for ITER energy confinement time, τ



ITER is twice as large as our largest existing experiments



ITER Physics Basis II

• MHD stability:

$$q_{95} = 3$$
 $q_{95} = 2.5 \frac{a^2 B}{RI} f(\varepsilon, \kappa, \delta)$

$$n/n_{GW} \le 1$$
 $n_{GW} (10^{20}) = \frac{I(MA)}{\pi a^2}$

$$\beta_{\rm N} \le 2.5$$
 $\beta_{\rm N} = \beta(\%) \frac{{\rm aB}}{{\rm I(MA)}}$

 $\kappa,~\delta$ determined by control considerations

 β = (plasma kinetic energy)/(plasma magnetic energy)

• Divertor physics:

Peak target power ~ 10MWm⁻²

Helium transport: $\tau_{He}^{*} / \tau_{E} \sim 5$

Impurity content: $n_{Be}/n_e = 0.02$ (+ ~0.1% Ar for radiation)

MHD Stability - Plasma Operational Limits

- The interaction of the plasma fluid and the magnetic field is described by magnetohydrodynamic (MHD) stability theory
 - provides a good qualitative, and to a significant extent quantitative, description of stability limits and the associated instabilities
- There are two basic types of instability:
 - "ideal" instabilities produce field line bending can grow very rapidly
 - "resistive" instabilities cause tearing and reconnection of the magnetic field lines ⇒ formation of "magnetic islands"
- Plasma control techniques are being applied to suppress or avoid the most significant instabilities
 - Neo-classical tearing modes (NTMs)
 - Edge localized modes (ELMs)
 - Disruptions and vertical displacement events
 - Allows access to higher fusion performance



MHD Stability: Disruptions



pical chain of events durin a plasma disruption

- The ultimate stability limit in tokamak plasmas is set by major disruptions: large scale MHD instabilities
 - loss of plasma energy in milliseconds (thermal quench TC)
 - plasma current decays in 10s of milliseconds (current quench QC)
- Produces:
 - very large heat loads on plasma facing surfaces
 - significant electromagnetic forces in vacuum vessel
- Mitigation techniques essential

large runaway electron beam

MHD Stability - Plasma Equilibrium Limits



- I_i-q_a diagram describes stable plasma operating space of internal inductance vs safety factor, limited by disruptions:
 - low I_i typically has to be negotiated during the plasma current ramp-up
 - high-l_i limit typically occurs due to excessive radiation at plasma edge, resulting in cold edge plasma and narrow current channel (e.g., at density limit)

MHD Stability - Density Limits

- Experiments have shown that tokamak plasmas can sustain a maximum density:
 - limit depends on operating regime (ohmic, L-mode, H-mode ...)
 - limit may be determined by edge radiation imbalance or edge transport processes
 - limit can be disruptive or nondisruptive
- Comprehensive theoretical understanding still limited
 - "Greenwald" density:

 $n_{GW} = I(MA)/\pi a^2$

- operational figure of merit



Plasma MHD Stability – Pressure Limit: β

 Maximum value of normalized plasma pressure, β, is limited by MHD instabilities:

$$\beta(\%) = 100 \frac{\langle p \rangle}{B^2 / 2\mu_0}$$
$$\beta_N = \frac{\beta(\%)}{I_p(MA) / aB}$$

• Typically, "Troyon" limit describes tokamak plasmas:

 $\beta_N {\leq} 2.8 {-} 3.5$

• More generally, "no-wall" limit:

$$\beta_N \le 4 \times \mathbf{I}_i$$



Power and Particle Exhaust

Scrape-off layer (SOL) plasma: region of open field lines

Essential problem is:

- handle power produced by plasma with (steady-state) engineering limit for plasma facing surfaces of 10 MWm⁻²
- extract helium from the core plasma to limit concentration below ~6%
- prevent impurities from walls penetrating into plasma core
- ensure plasma facing surfaces survive sufficiently long



Power and Particle Exhaust

• The divertor is a significant element of the solution

- surfaces for high heat fluxes (10 MWm⁻²)
- cryopumping to extract particles leaving the plasma, including helium



- The divertor is fundamental to High heat exhaust power from a burning plasma:
 - impurities are added to the edge plasma to increase radiation

ITER divertor cassette – 54 cassettes make up the complete toroidal ring

- a large pressure gradient develops along the field lines into the divertor
- > the divertor plasma temperature falls to a few eV
- a large fraction of the plasma exhaust power is redistributed by radiation and ion-neutral collisions

Burning Plasma Physics

- Access to plasmas which are dominated by a-particle heating will open up new areas of fusion physics research, in particular:
 - confinement of α 's in plasma
 - response of plasma to α -heating
 - influence of α -particles on MHD stability
- Experiments in existing tokamaks have already provided some positive evidence
 - "energetic" particles (including α -particles) are well confined in the plasma
 - such particle populations interact with the background plasma and transfer their energy as predicted by theory
 - but energetic particles can induce MHD instabilities (Alfvén eigenmodes) - for ITER parameters at Q=10, the impact is expected to be tolerable

Energetic Ion Confinement

- In existing experiments single particle theory of energetic ion confinement confirmed:
 - − simple estimate, based on banana orbit width shows that $I_p ≥ 3MA$ required for α-particle confinement
- Classical slowing down of fast ions well validated:
 - data range 30keV NBI (ISX-B) to
 3.5MeV α-particles (TFTR)
- Energetic ion heating processes routinely observed in additional heating experiments



W W Heidbrink, G J Sadler, Nucl Fusion 34 535 (1994)

Alfvén Eigenmodes

- In a tokamak plasma, the Alfvén wave continuum splits into a series of bands, with the gaps associated with various features of the equilibrium:
 - a series of discrete frequency <u>Alfvén eigenmodes</u> can exist in these gaps:
 - <u>toroidicity-induced</u> (TAE)
 - ellipticity-induced (EAE)
 - <u>triangularity-induced</u> (NAE)
 - <u>beta-induced</u> (BAE)
 - <u>kinetic toroidal</u> (KTAE)

gap created by toroidicity gap created by elongation gap created by additional noncircular effects

gap created by field compressibility gap created by non-ideal effects such as finite Larmor radius

... and others!

 These modes can be driven unstable by the free energy arising from energetic particle populations with velocities above the Alfvén velocity, eg α-particles

Physics for Fusion Power Plants

- A fusion power plant requires physics parameters that are <u>simultaneously</u> close to the limits of what might be achievable on the basis of our (experimental and theoretical) understanding
- Several key issues in (burning) plasma physics for a tokamak power plant must be developed in the current programme and demonstrated (and extended) in ITER:
 - Operating scenario steady-state ?
 - High confinement at high density and high radiated power fraction
 - High fusion power \Rightarrow high β operation \Rightarrow robust MHD stability
 - Effective disruption avoidance and control
 - Power (and particle) exhaust with relevant PFCs
 - Tritium efficiency
 - α-particle confinement
 - Reactor-relevant auxiliary systems (H&CD, diagnostics, fuelling, control ...)

ITER on the Path to Fusion Energy



References: Tokamak Fusion Physics

Gibson, A. et al. *Deuterium-tritium plasmas in the Joint European Torus (JET): Behavior and implications*, Physics of Plasmas **5**, 1839 (1998).

Greenwald, M, et al., *A new look at density limits in tokamaks*, Nuclear Fusion **28**, 2199 (1988).

Hawryluk, R. J. *Results from deuterium-tritium tokamak confinement experiments*, Reviews of Modern Physics **70**, 537 (1998).

Heidbrink, W. et al., Nuclear Fusion 34, 535 (1994).

Jacquinot, J. et al. Overview of ITER physics deuterium-tritium experiments in JET, Nuclear Fusion **39**, 235 (1999).

Keilhacker, M. et al. *High fusion performance from deuterium-tritium experiments in JET*, Nuclear Fusion **39**, 209 (1999).

Snipes, J. A., et al., Nuclear Fusion 28, 1085 (1988).

Troyon, F. and Gruber, R., Physics Letters A 110, 29 (1985).

ITER Physics Basis, ITER Physics Expert Groups et al, Nucl Fusion **39** 2137-2638 (1999)

Progress in the ITER Physics Basis, ITPA Topical Physics Groups et al, Nucl Fusion **47** S1-S413 (2007)

http://www.iter.org - and associated links