



*The Abdus Salam*  
**International Centre for Theoretical Physics**



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**Joint ICTP-IAEA School on Synchrotron Applications in Cultural Heritage and  
Environmental Sciences and Multidisciplinary Aspects of Imaging Techniques**

**21 - 25 November 2011**

**X-Ray Absorption Spectroscopy**

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Fondazione Bruno Kessler  
Trento  
ITALY*

# X-Ray Absorption Spectroscopy

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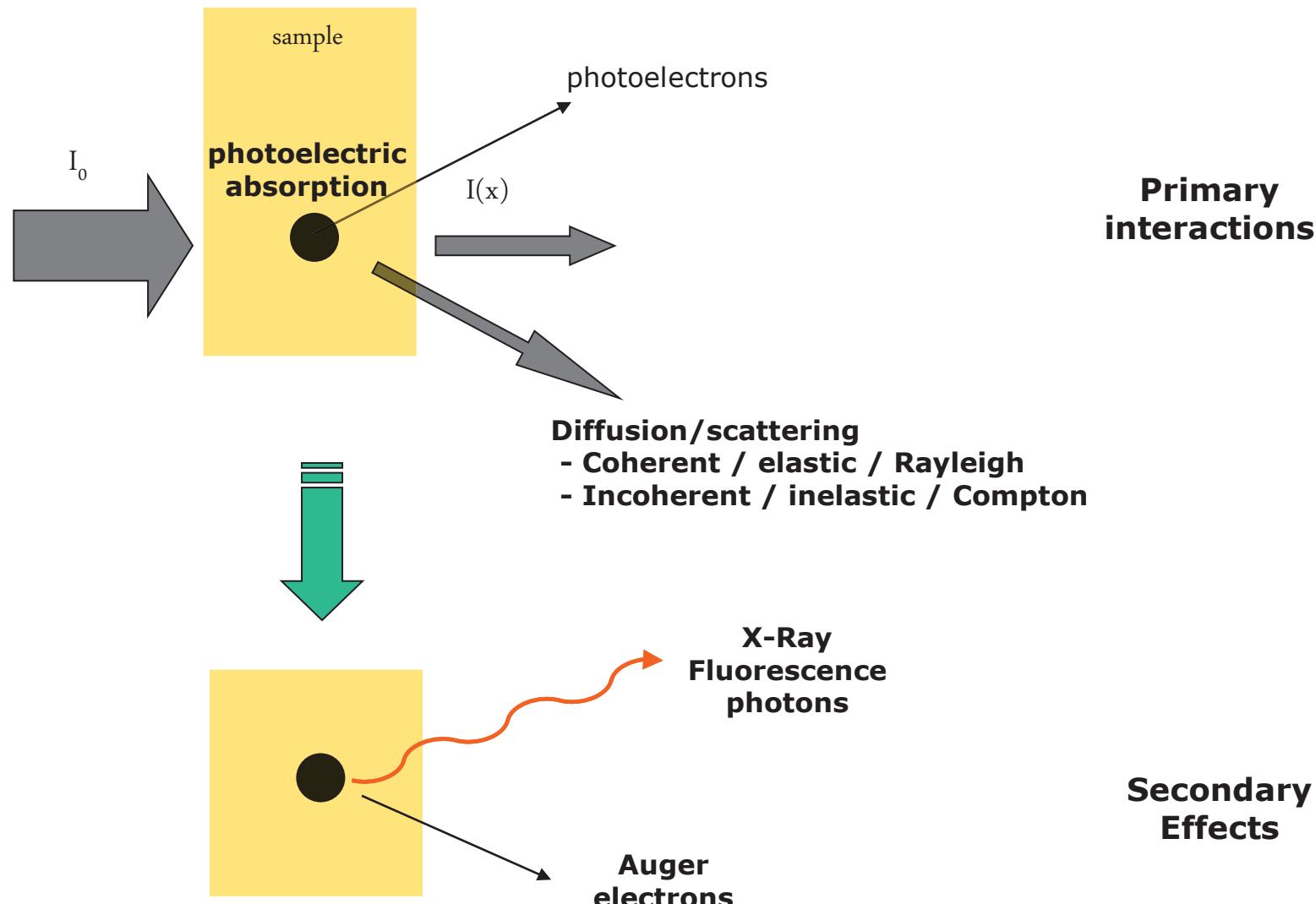
acknowledgement:  
Florian Meirer - FBK  
Paolo Fornasini - UniTN

# **Interaction of X-Rays with matter**

**Photoelectric effect**

**Electron energy levels**

# Interaction of X-Rays with matter

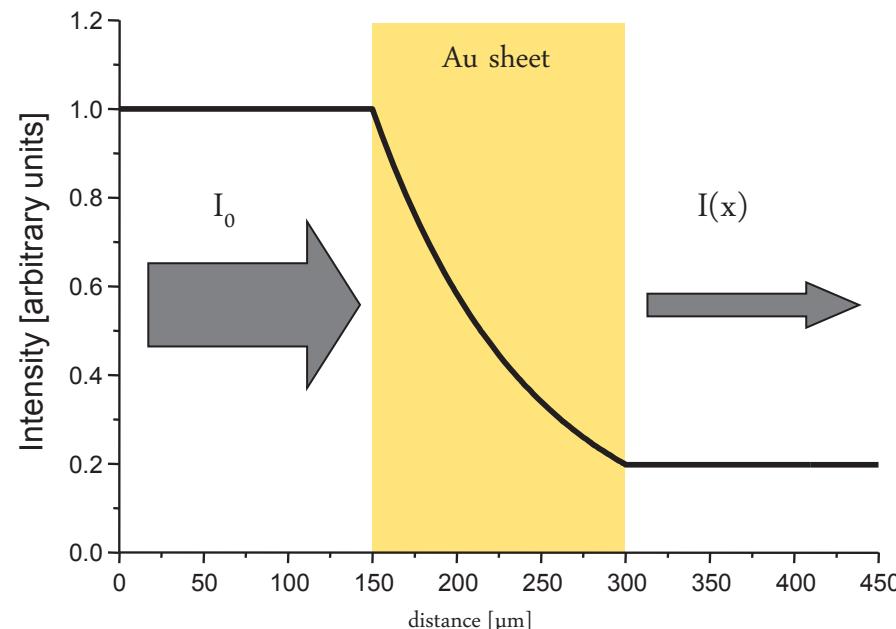


# Beer-Lambert's Law

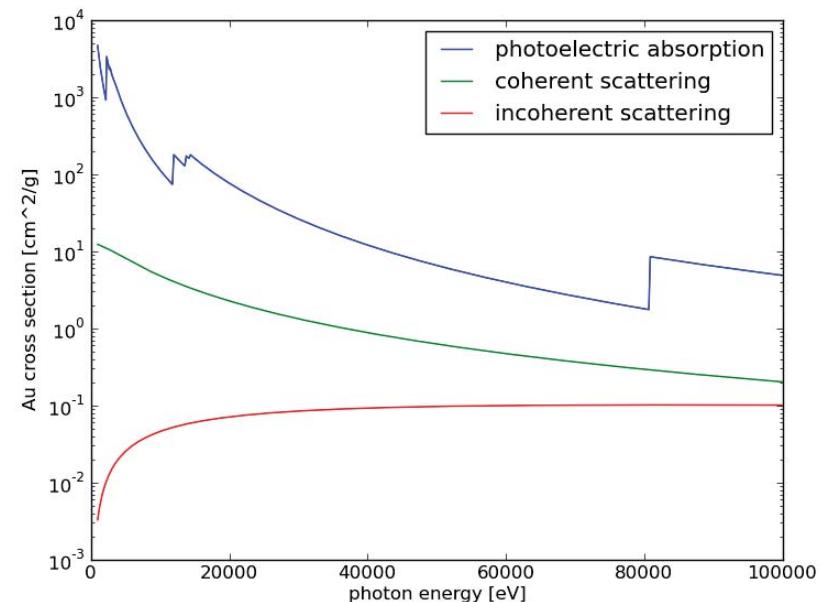
**Beer-Lambert's Law:**

$$I(x) = I_0 \exp(-\mu x)$$

$$\mu = \sigma_c + \sigma_i + \tau$$



Calculated for  $E = 17500\text{eV}$

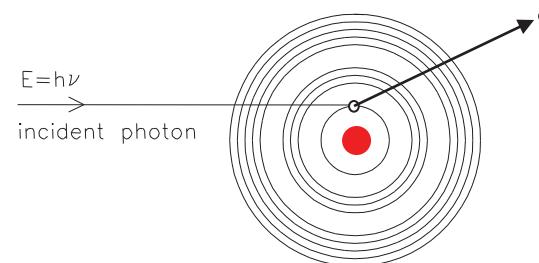


**data from:**

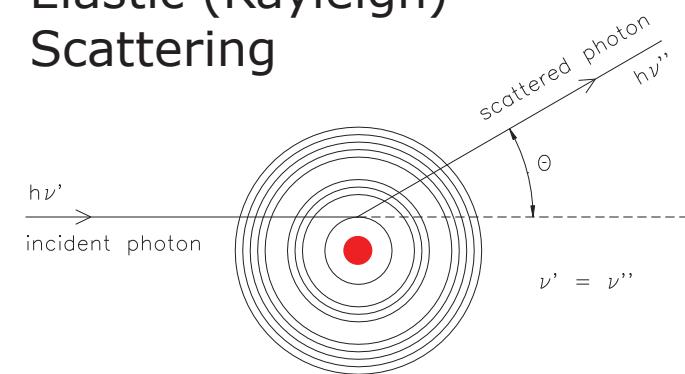
**H. Ebel, R. Svagera, M. F. Ebel, A. Shaltout and J. H. Hubbell, Numerical description of photoelectric absorption coefficients for fundamental parameter programs, X-Ray Spectrometry, 32, 442–451 (2003)**

# Interactions – microscopic view

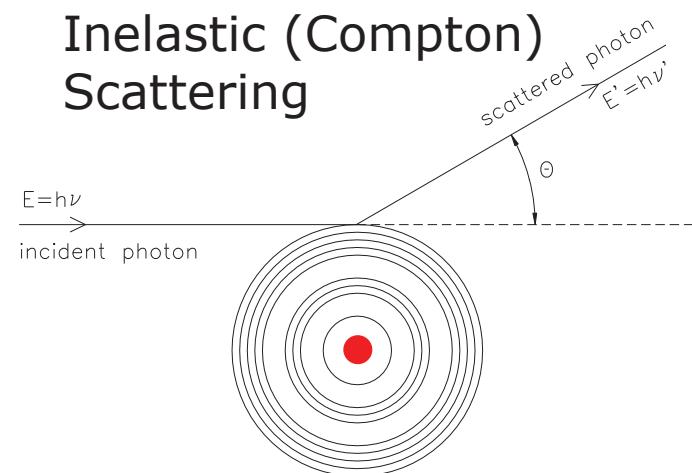
Photoelectric absorption



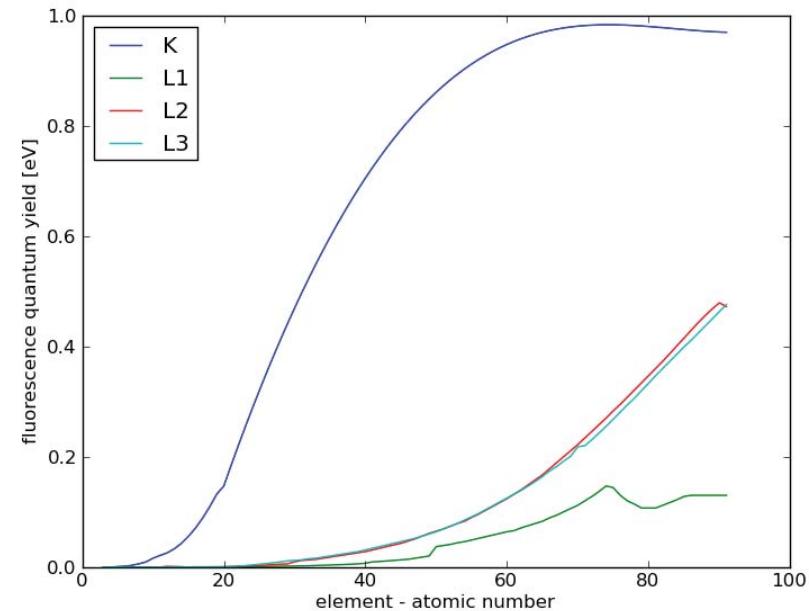
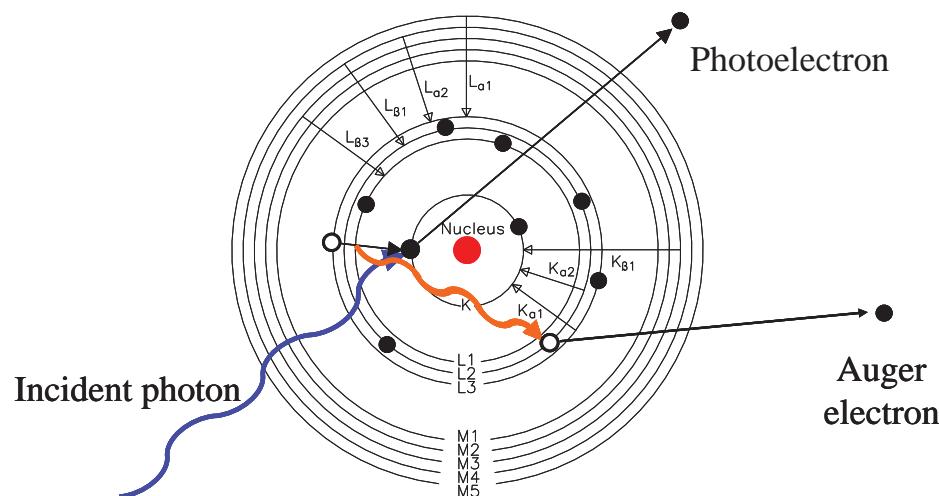
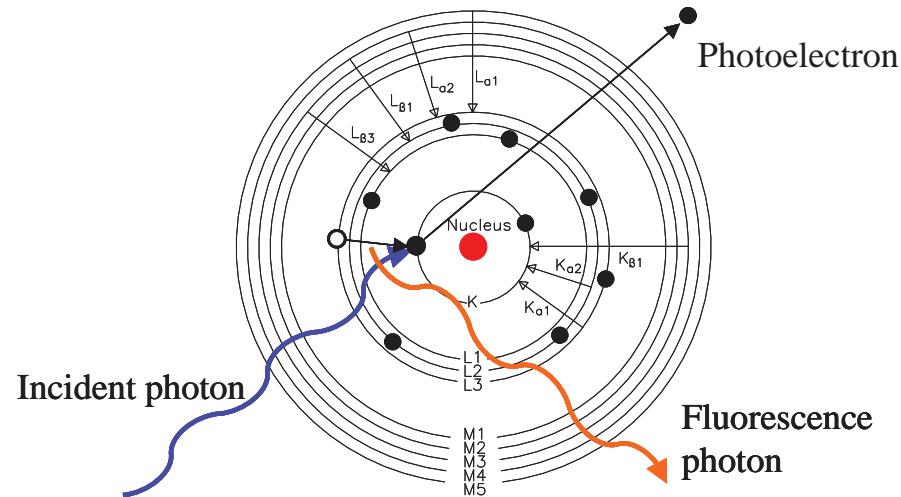
Elastic (Rayleigh) Scattering



Inelastic (Compton) Scattering

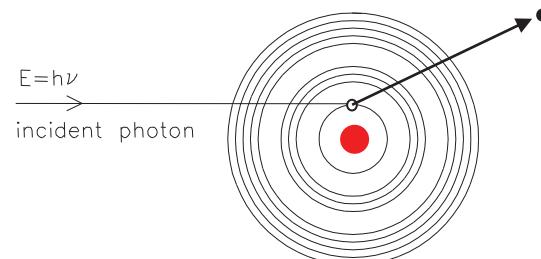


# Secondary effects – Fluorescence vs Auger



**data from:**  
**M. O. Krause,**  
**J. Phys. Chem. Ref. Data 8 (1979) 307**

# Photoelectric effect – energy levels

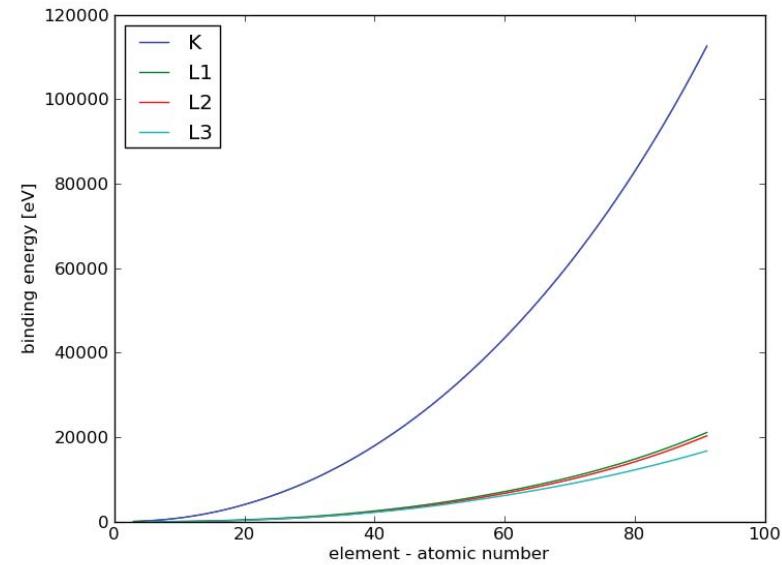
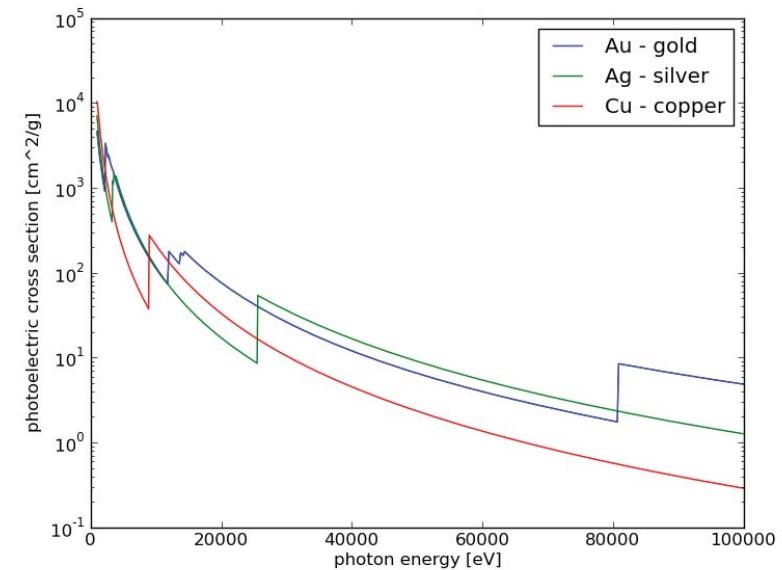


Absorption edges  
 Electron energy levels  
 Shells

shell	n	l	j	spin sign	max number of electrons
K	1	0	0.5	1	2
L1	2	0	0.5	1	2
L2	2	1	0.5	-1	2
L3	2	1	1.5	1	4
M1	3	0	0.5	1	2
M2	3	1	0.5	-1	2
M3	3	1	1.5	1	4
M4	3	2	1.5	-1	4
M5	3	2	2.5	1	6

Z	shell	energy_eV	jump	level_width_eV
79	K	80724.9	4.874	52.1
79	L1	14352.8	1.15567	9.8
79	L2	13733.6	1.4	5.53
79	L3	11918.7	2.55	5.54
79	M1	3424.9	1.04	15.0
79	M2	3147.8	1.058	9.5
79	M3	2743.0	1.15776	8.5
79	M4	2291.1	1.07	2.18
79	M5	2205.7	1.092	2.18

[www.txrf.org/xraydata](http://www.txrf.org/xraydata)

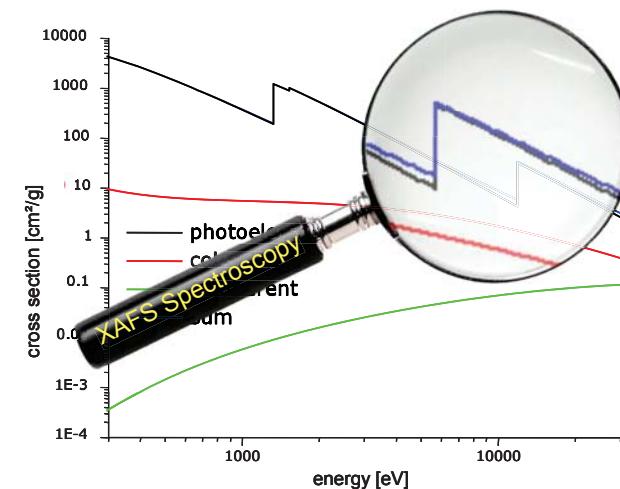
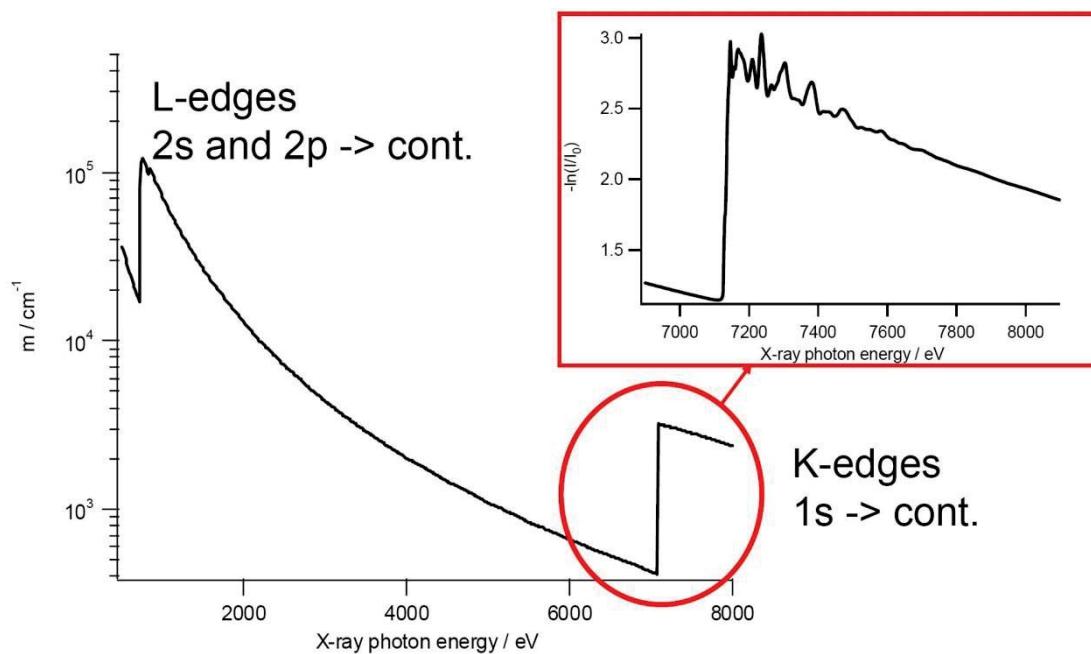


# **The fine Structure**

## **Phenomenological Overview**

# Near Edge Fine Structure

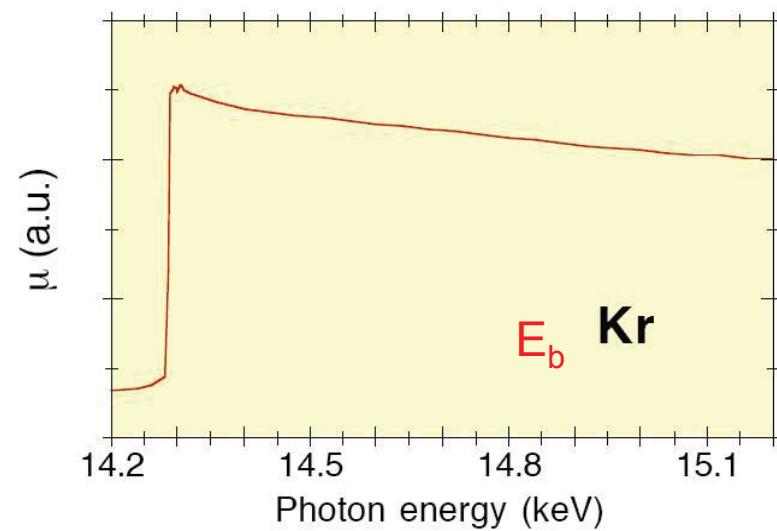
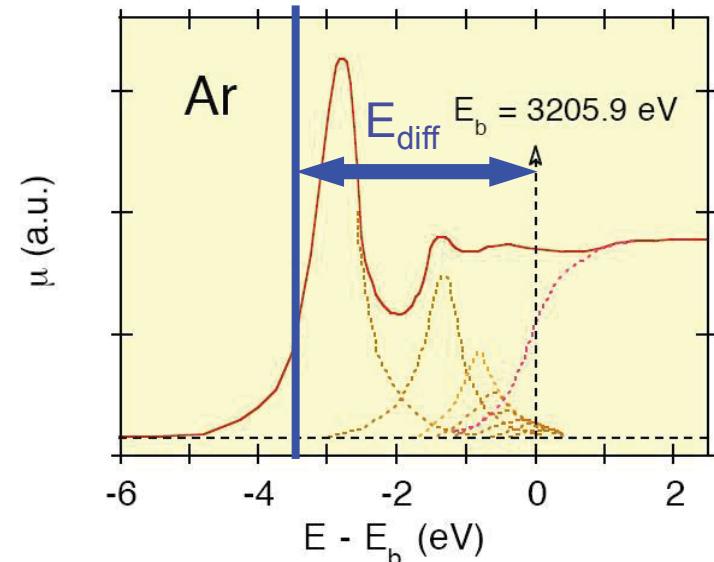
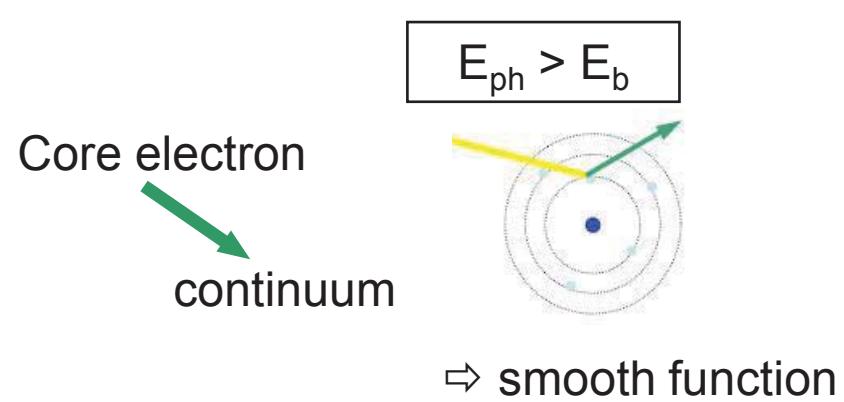
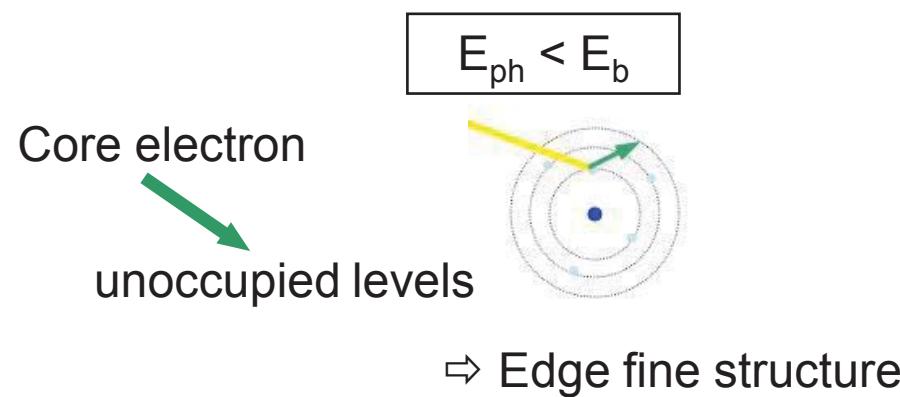
The X-ray Absorption Fine Structure (XAFS) of an iron foil



Different phenomena for:

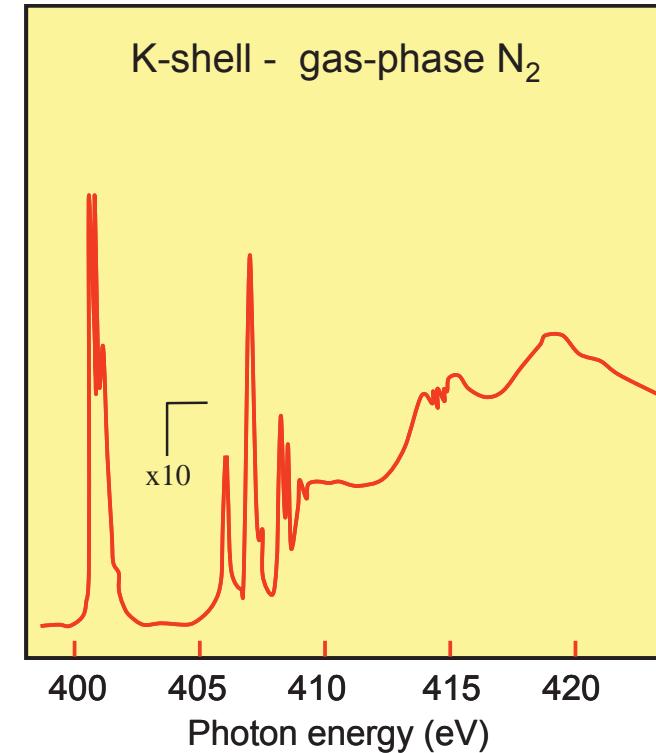
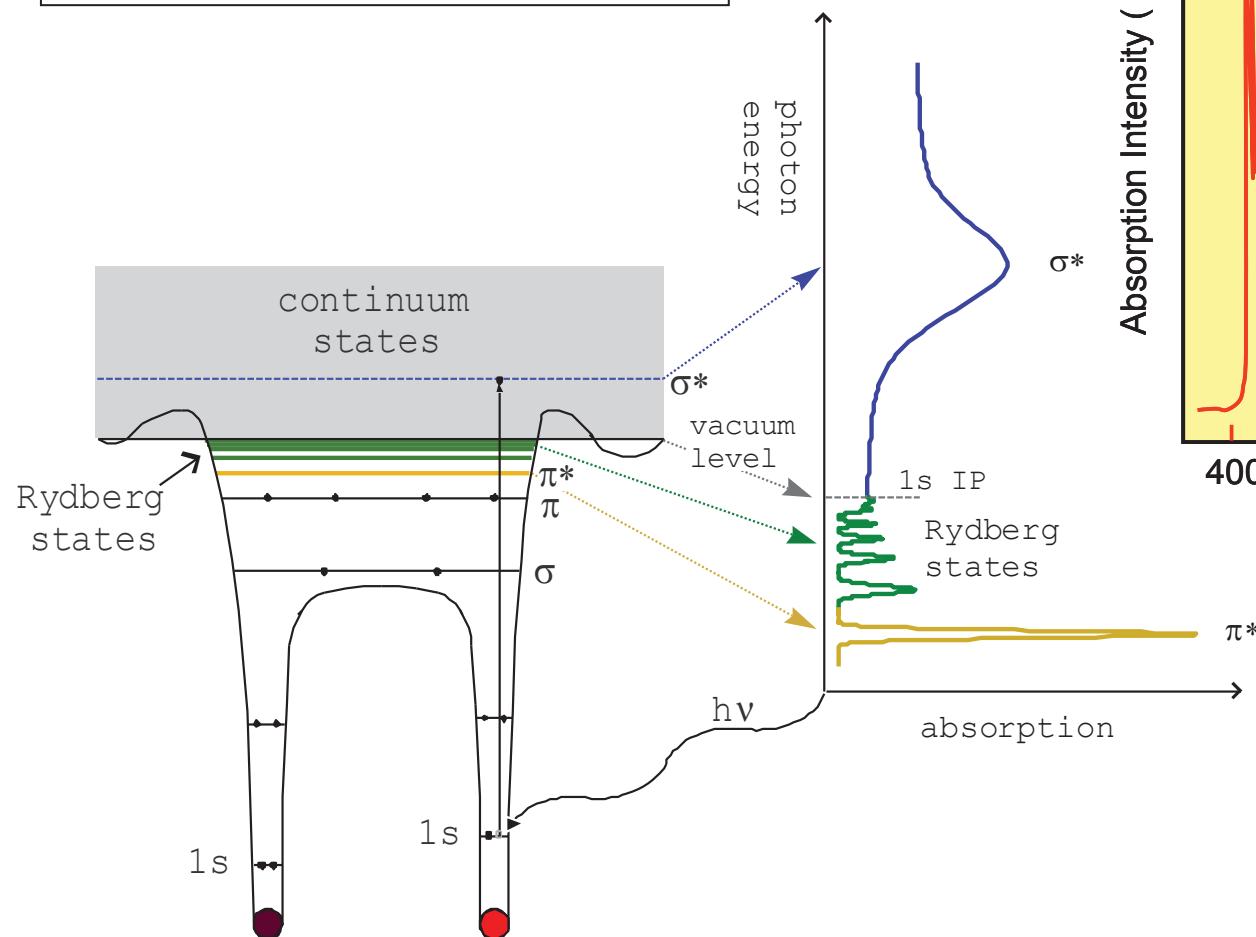
- 'free' atoms
- molecules
- condensed systems

# Near Edge Fine Structure – origin - atoms

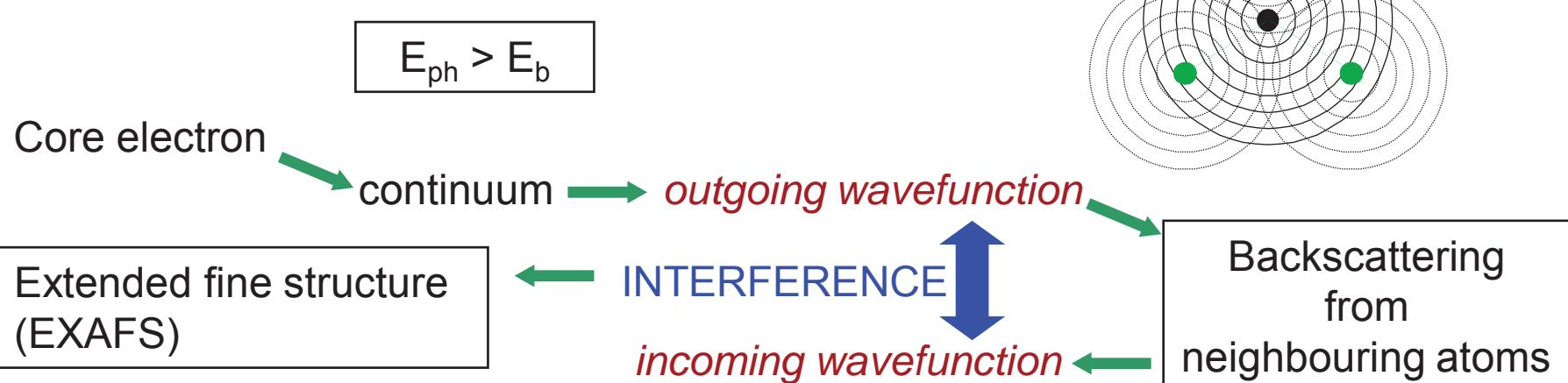
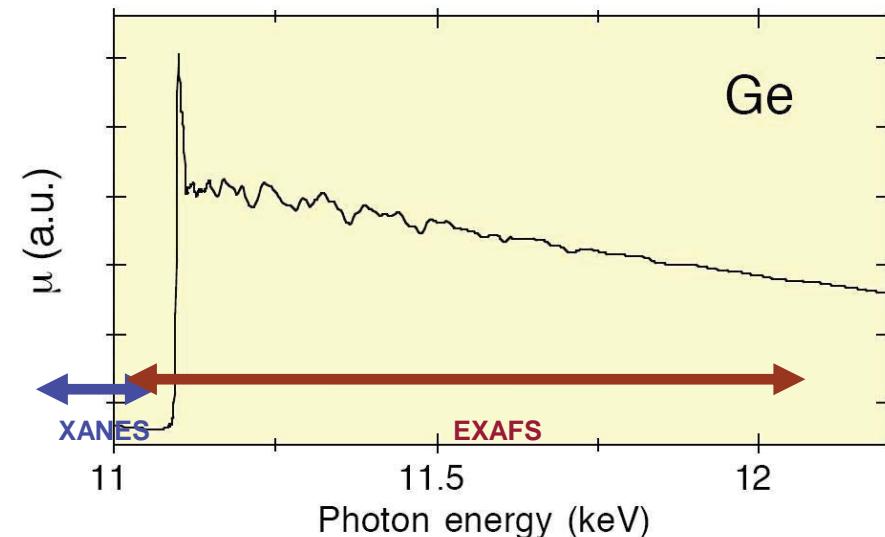
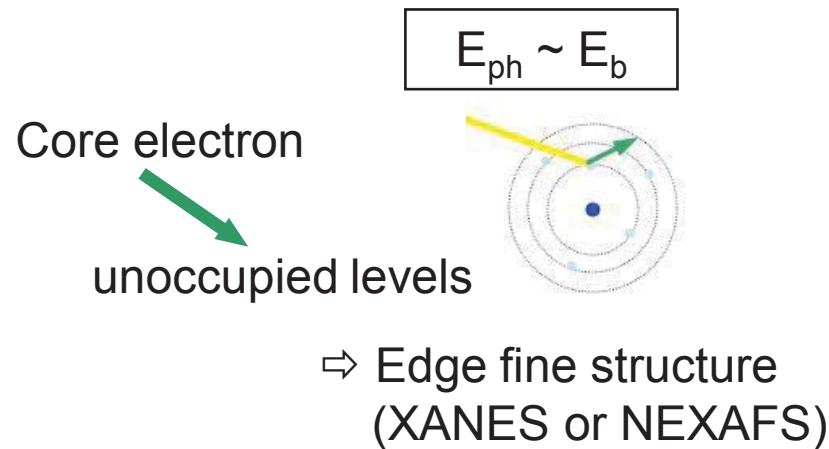


# Near Edge Fine Structure – origin - molecules

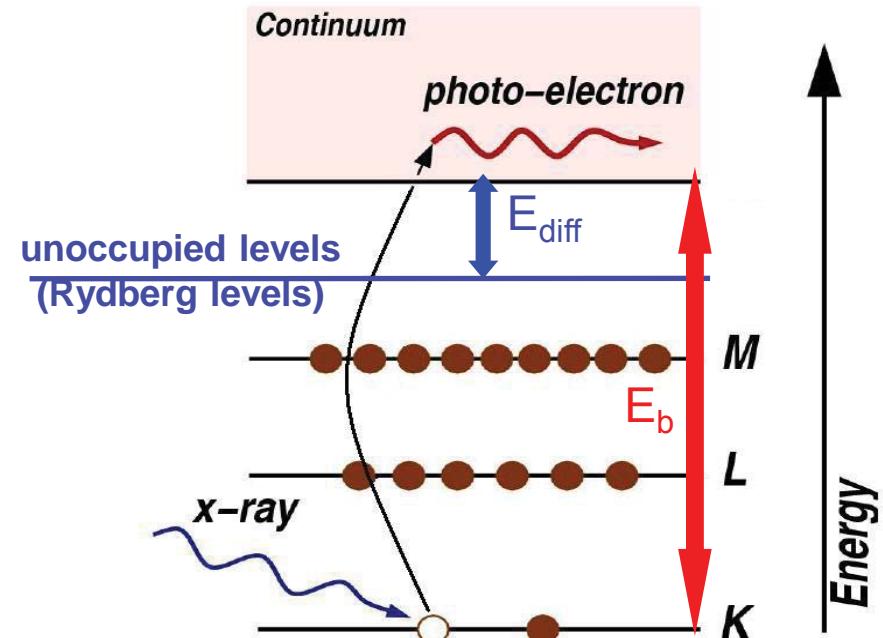
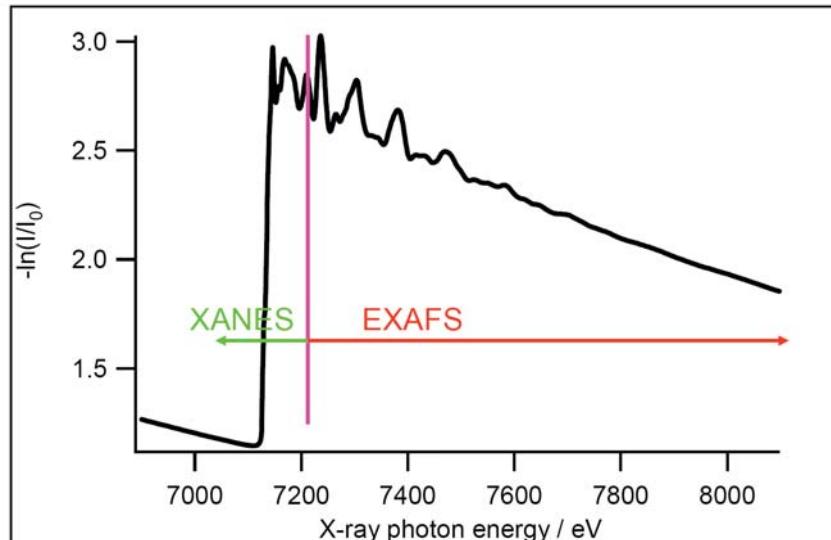
For molecules the predominant features are unfilled molecular orbitals  $\sigma^*$  and  $\pi^*$  (resonances).



# Near Edge Fine Structure – origin Condensed systems



# Near Edge Fine Structure



**XANES:** X-Ray Absorption Near Edge Structure, ends 50-100 eV above the edge

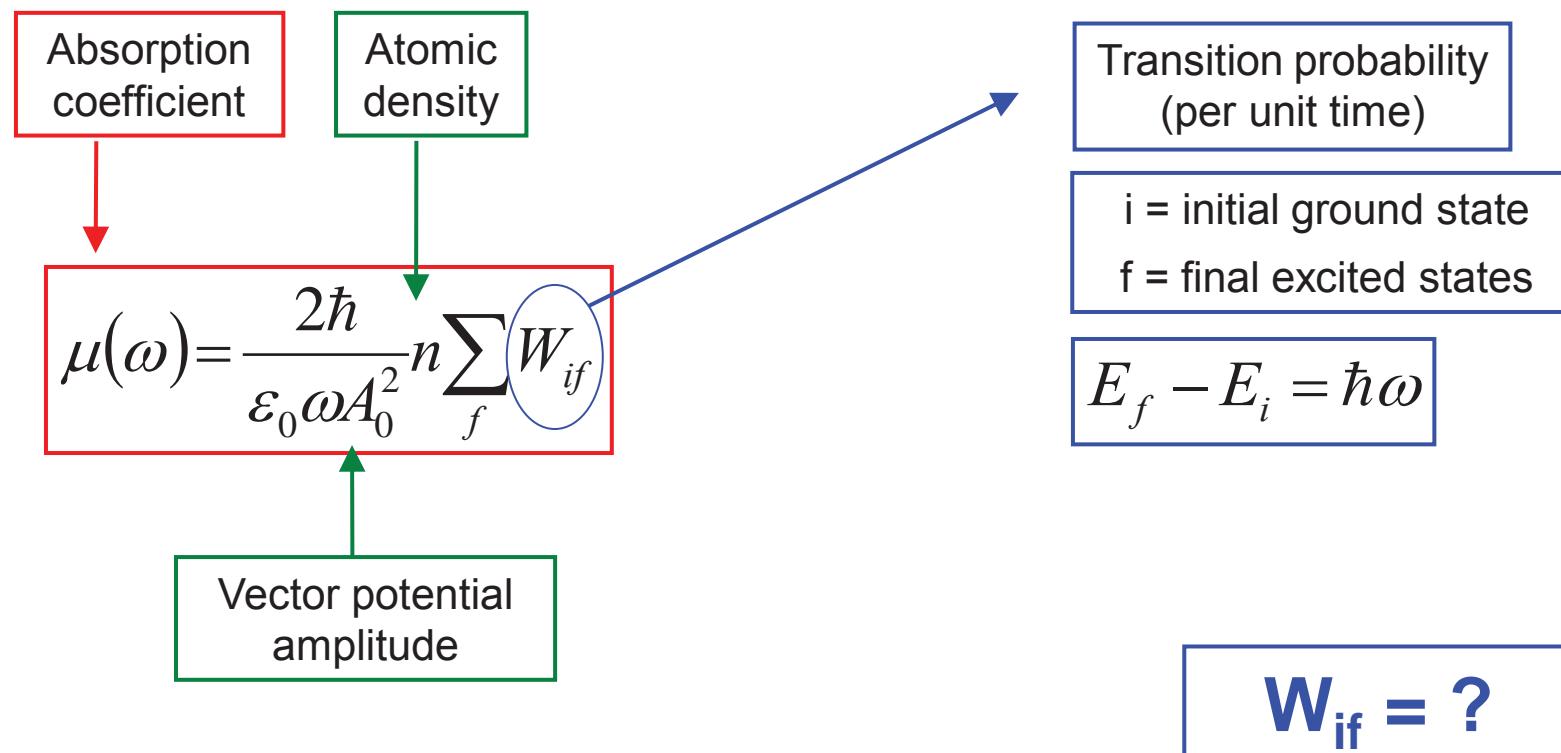
**EXAFS:** Extended X-Ray Absorption Fine Structure, starts 50 - 100 eV above the edge

**NEXAFS:** Near Edge X-Ray Absorption Fine Structure, typically used for molecules and the soft x-ray range

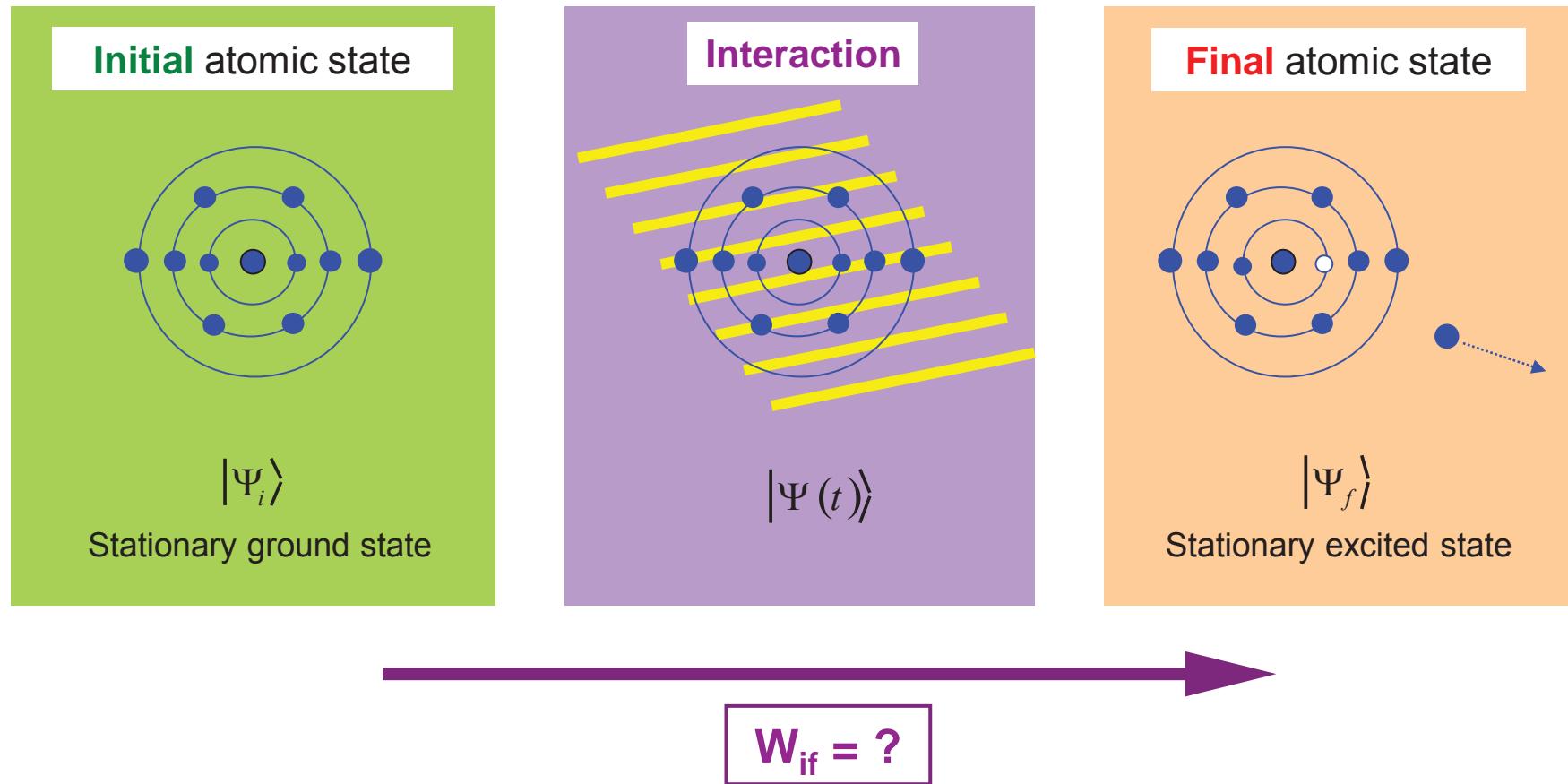
# **Near Edge Fine Structure**

## **Theoretical description**

# XAS theory – absorption coefficient

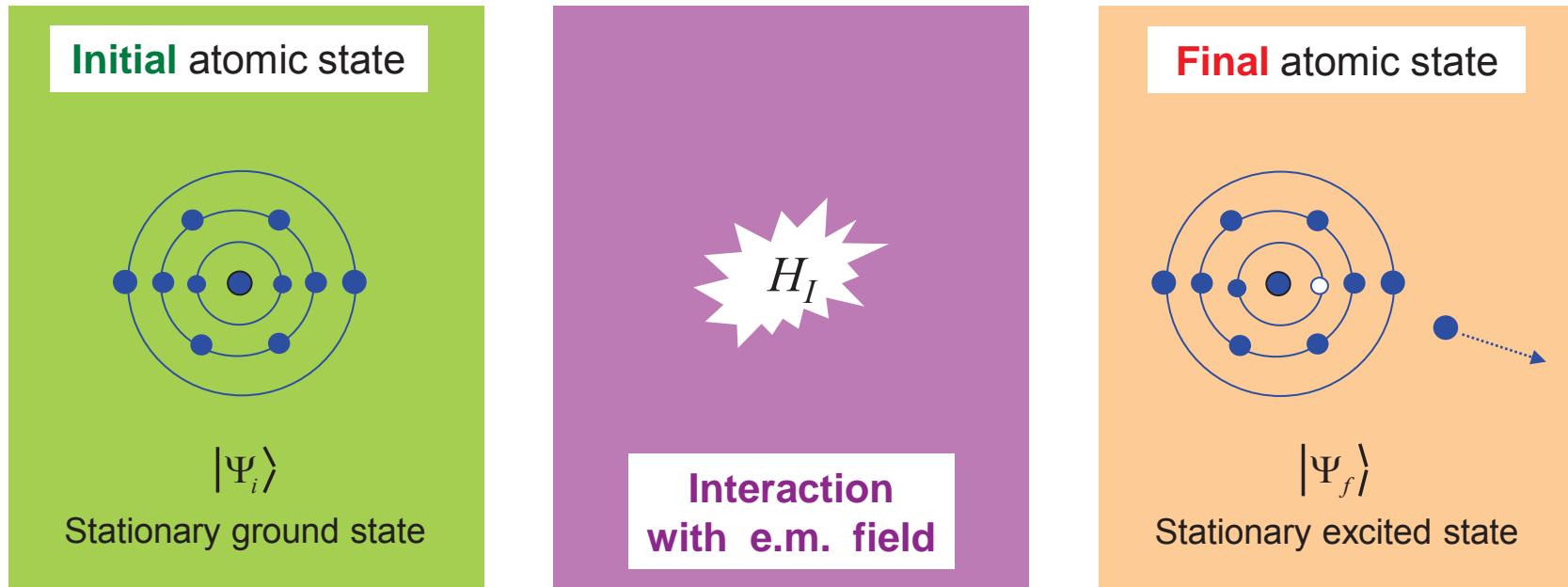


# XAS theory – transition probability



# Fermi Golden Rule

Time-dependent perturbation theory (1st-order)



EXAFS

XANES

$$W_{if} \propto \left| \langle \Psi_i | H_I | \Psi_f \rangle \right|^2 \rho(E_f)$$

matrix element

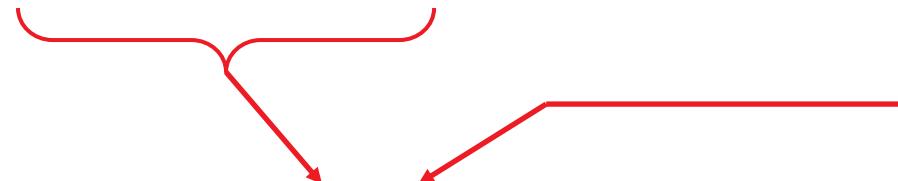
density of  
final states

$$E_f - E_i = \hbar\omega$$

# Interaction Hamiltonian

Sum over electrons

$$\hat{H}_I = \sum_j \left[ \frac{e}{m} \vec{p}_j \cdot \vec{A}(r_j, t) - \frac{e^2}{2m} A^2(r_j, t) \right]$$



$$W_{if} = \frac{\pi \hbar e^2}{m^2} |A_o|^2 \left| \left\langle \Psi_i \left| \sum_j e^{i\vec{k} \cdot \vec{r}_j} \hat{\eta} \cdot \vec{\nabla}_j \right| \Psi_f \right\rangle \right|^2 \rho(E_f)$$

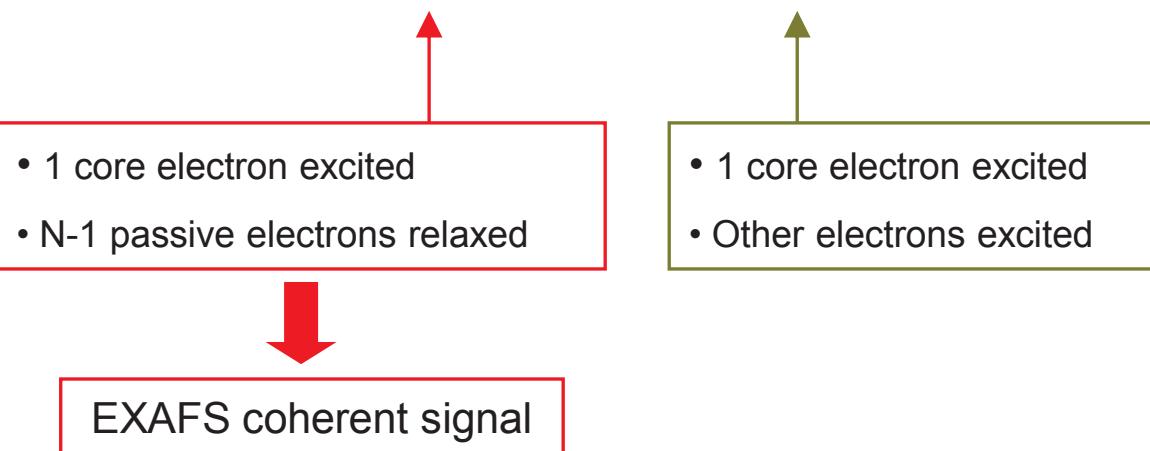
Coulomb gauge

$\vec{\nabla} \cdot \vec{A} = 0$	$\vec{E} = \partial \vec{A} / \partial t$
$\Phi = 0$	$\vec{B} = \vec{\nabla} \times \vec{A}$
$\vec{J} = 0$	$\vec{\nabla} \cdot \vec{A} = \vec{A} \cdot \vec{\nabla}$

$$\vec{A} = A_0 \hat{\eta} e^{i(\vec{k} \cdot \vec{r} - \omega t)} + c.c.$$

# One electron approximation

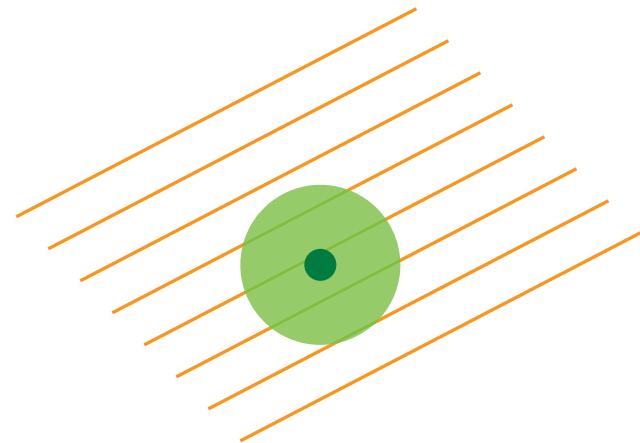
$$\mu_{\text{tot}}(\omega) = \mu_{\text{el}}(\omega) + \mu_{\text{inel}}(\omega)$$



$$\mu_{\text{el}}(\omega) \propto \left| \left\langle \Psi_i^{N-1} \psi_i \right| e^{i\vec{k} \cdot \vec{r}} \hat{\eta} \cdot \vec{p} \left| \psi_f \Psi_f^{N-1} \right\rangle \right|^2 \rho(\epsilon_f)$$

# Electron dipole approximation

$$e^{i\vec{k} \cdot \vec{r}} = 1 + i\vec{k} \cdot \vec{r} - \dots \approx 1$$



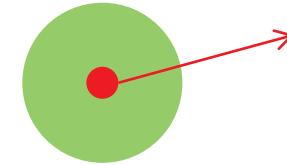
$$H_I \propto e^{i\vec{k} \cdot \vec{r}} \hat{\eta} \cdot \vec{p} \approx \hat{\eta} \cdot \vec{p} = \omega^2 \hat{\eta} \cdot \vec{r}$$



$$\mu_{\text{el}}(\omega) \propto \left| \langle \Psi_i^{N-1} \psi_i | \hat{\eta} \cdot \vec{r} | \psi_f \Psi_f^{N-1} \rangle \right|^2$$

# Sudden approximation

No interaction between **photoelectron** and **passive electrons**



$$|\Psi^{N-1}\psi\rangle = |\Psi^{N-1}\rangle |\psi\rangle$$

1 active electron

N-1 passive electrons

$$\mu_{\text{el}}(\omega) \propto \left| \langle \psi_i | \hat{\eta} \cdot \vec{r} | \psi_f \rangle \right|^2 \rho(\varepsilon_f) |\Psi_i^{N-1}|^2 |\Psi_f^{N-1}|^2$$

$$S_0^2 \approx 0.6 \div 0.9$$

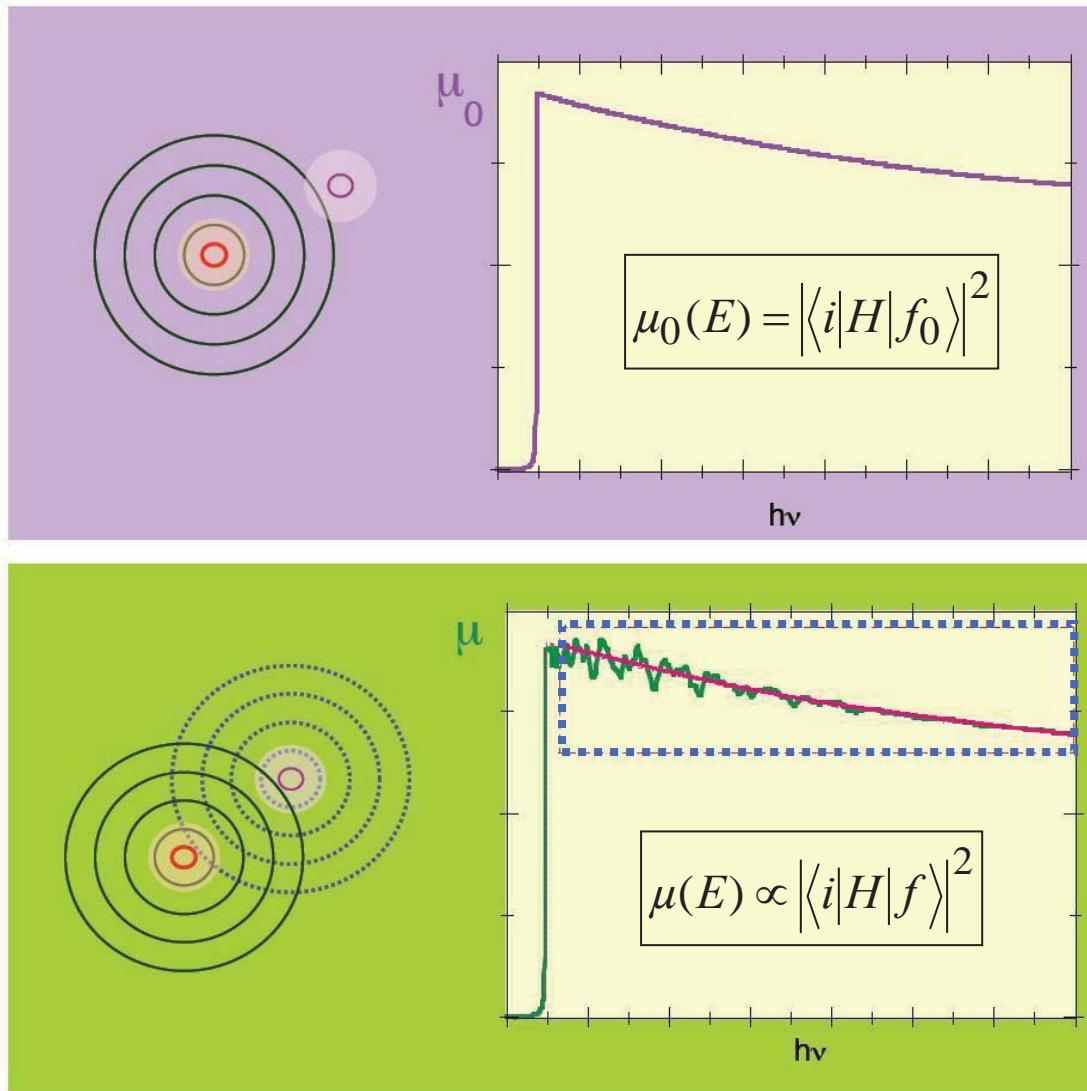
Structural  
information

## X-Ray Absorption Fine Structure

Perturbation of the absorption phenomena

The photoelectron is the real probe

# the EXAFS function



**EXAFS fine-structure function:**

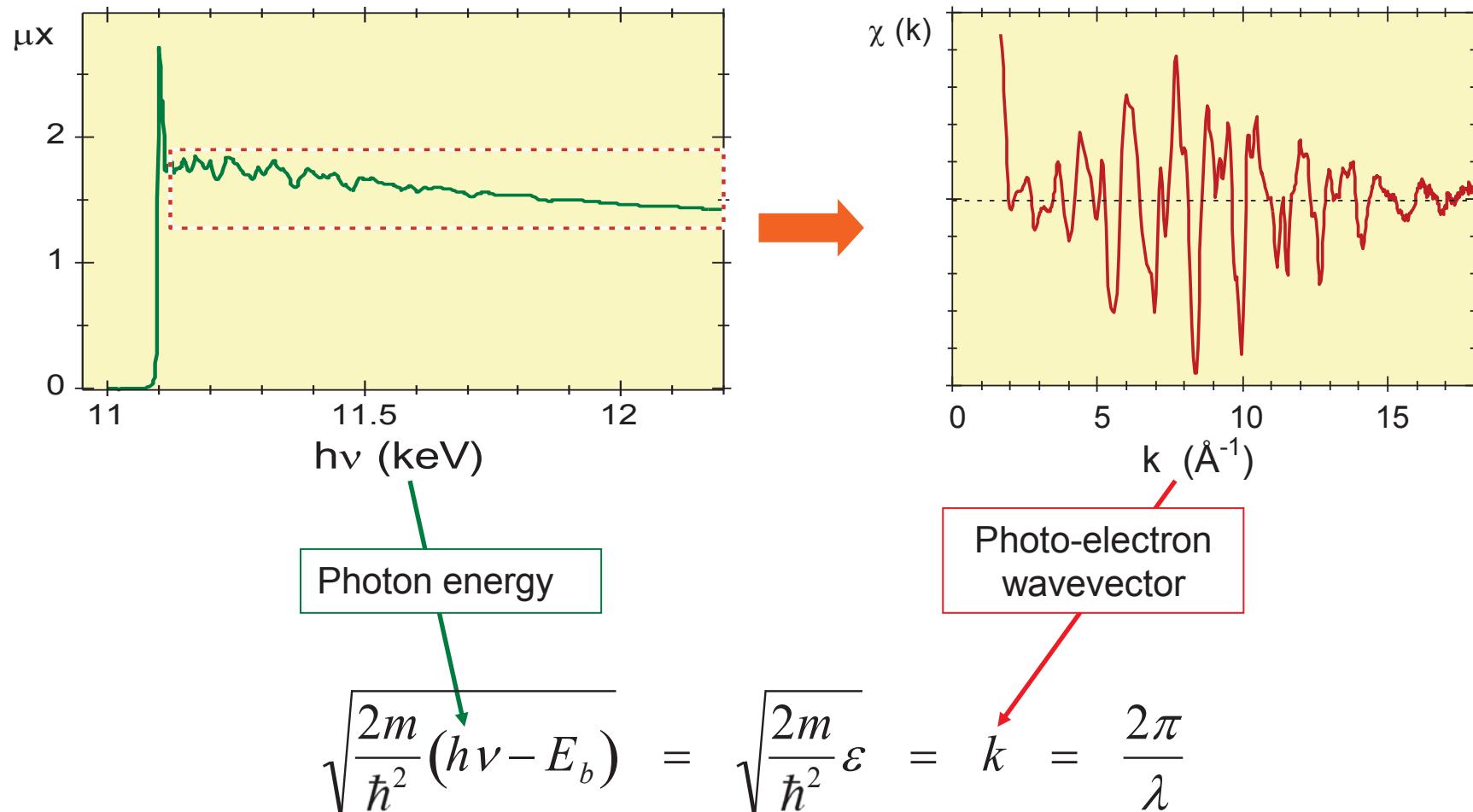
$$\chi(E) = \frac{\mu - \mu_0}{\mu_0}$$

usually transformed to k-space with:

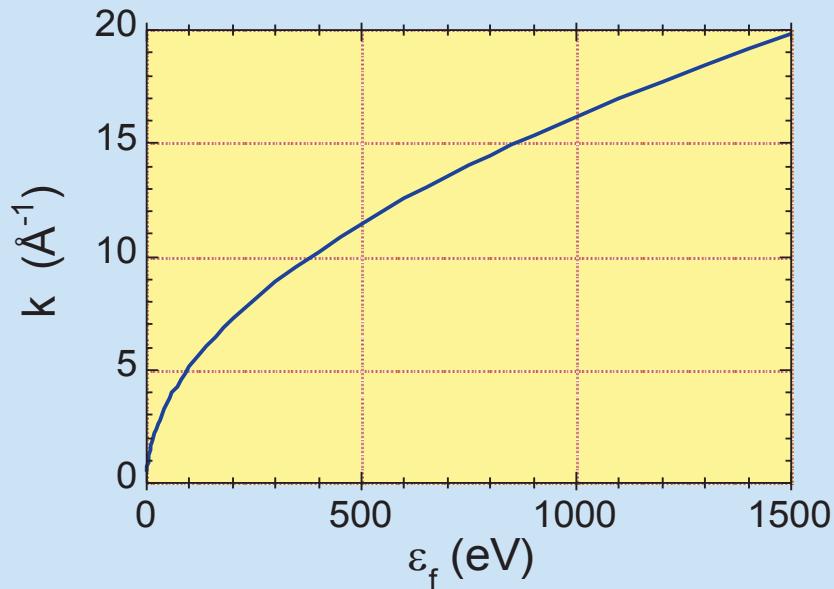
$$k = \sqrt{\frac{2m_e(E - E_0)}{\hbar^2}}$$

k.....photo electron wavenumber  
 m<sub>e</sub>.....mass of electron  
 E<sub>0</sub>.....binding energy

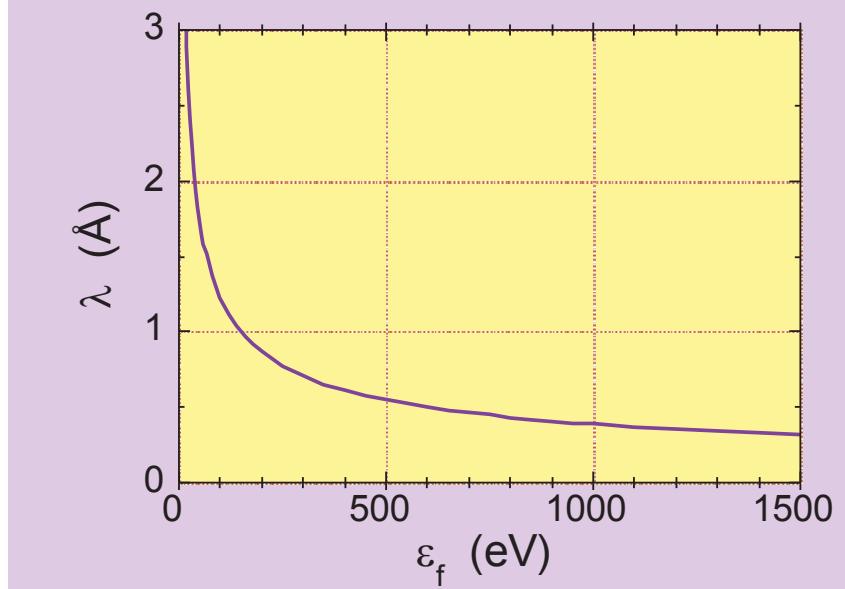
# Photon → photoelectron



Wave-number

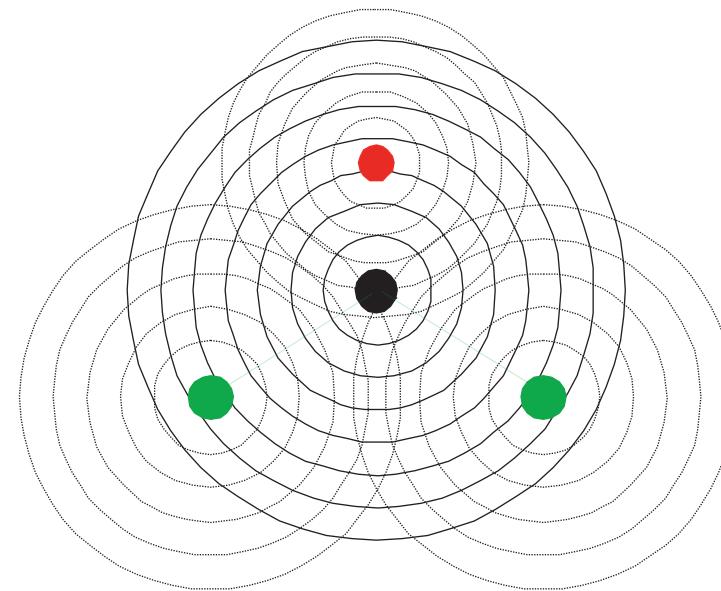


Wave-length



Energy

# the near EDGE region XANES

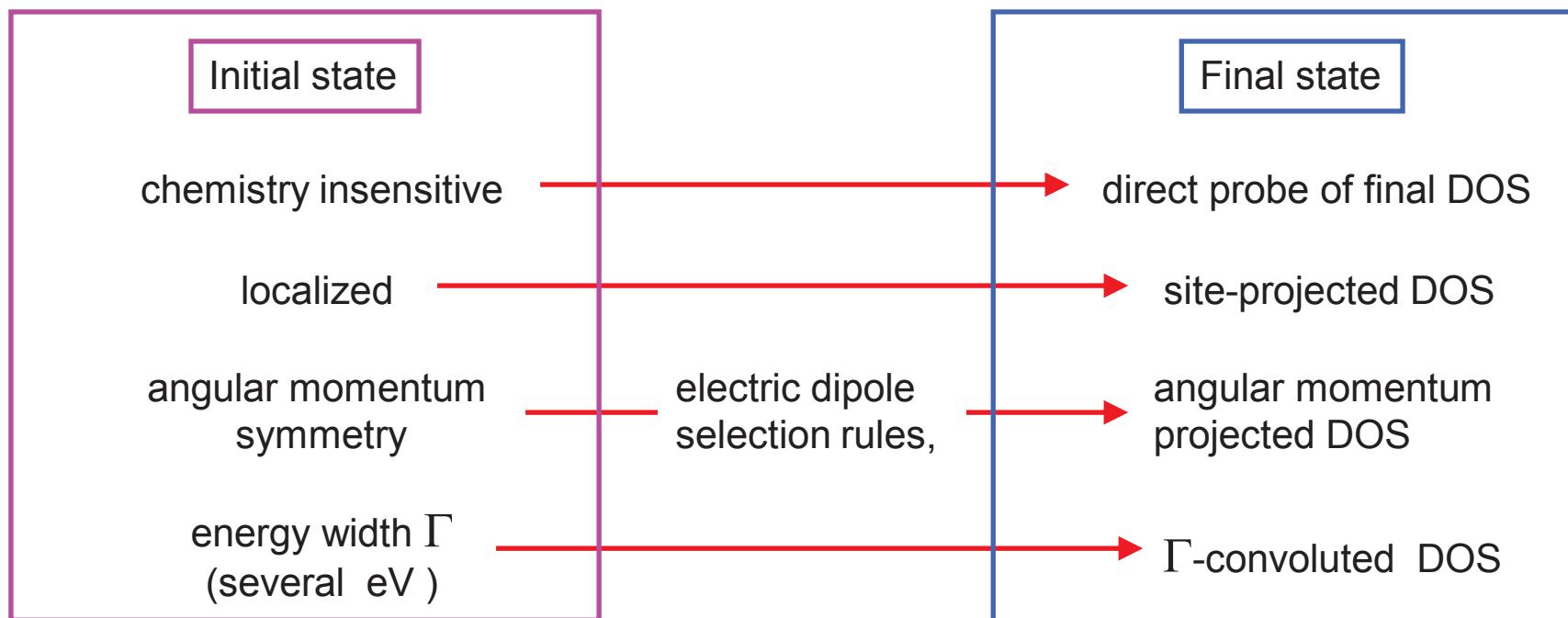


EXAFS

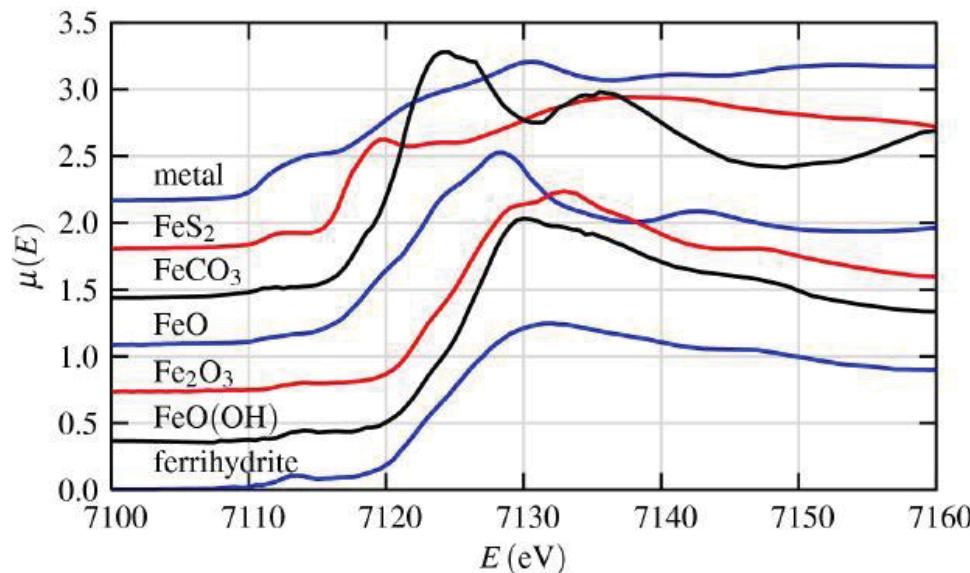
**XANES**

$$W_{if} \propto \left| \langle \Psi_i | H_I | \Psi_f \rangle \right|^2 \rho(E_f)$$

The features of the absorption spectrum is dominated by the the density of final states



- XANES is a much larger signal than EXAFS
  - ⇒ can be done at **lower concentrations**, and **less-than-perfect sample conditions**
- The interpretation of XANES is very complicated
  - ⇒ there is **no simple analytic** (or even physical) **description** of XANES
- the EXAFS equation breaks down at low-k (**1/k term**) and the **increase in the mean-free-path** at very low-k



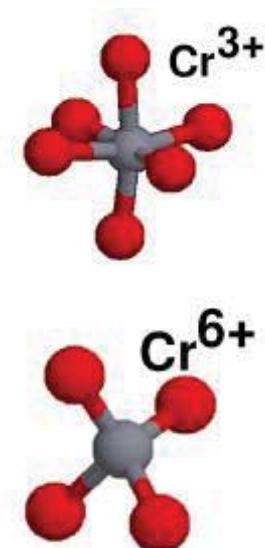
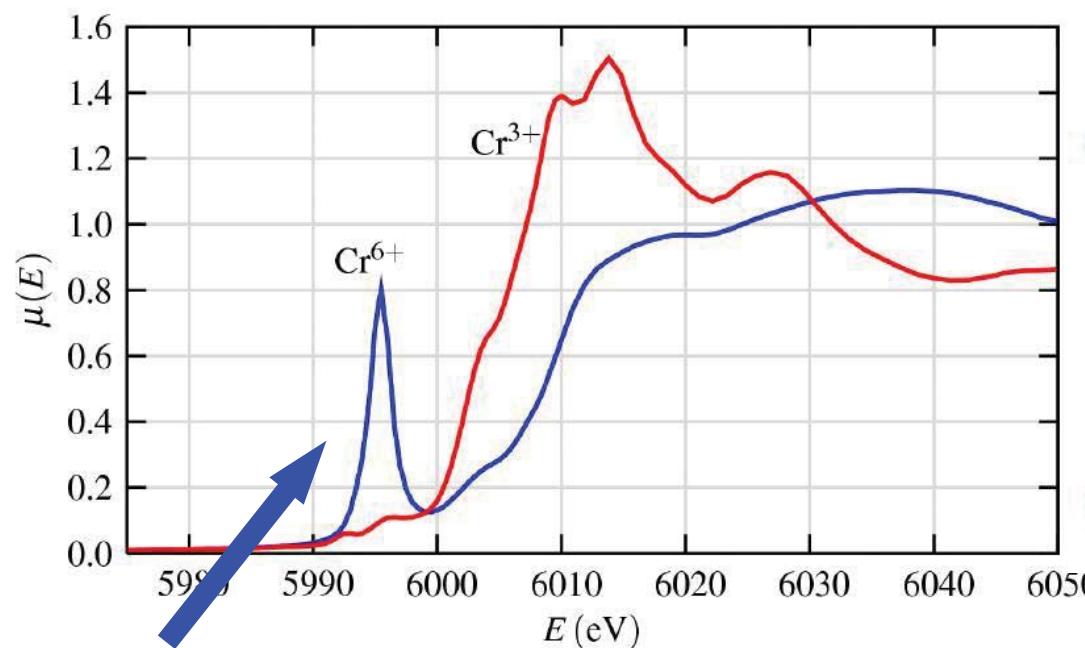
Still, there is much **chemical information** from the XANES region:

- the **edge position and shape** is sensitive to formal valence state, ligand type, and coordination environment
- ⇒ XANES can be used as a **fingerprint method** to identify phases.

## An example: Chromium (Cr)

K edge XANES : edge energy = 5989.2 eV

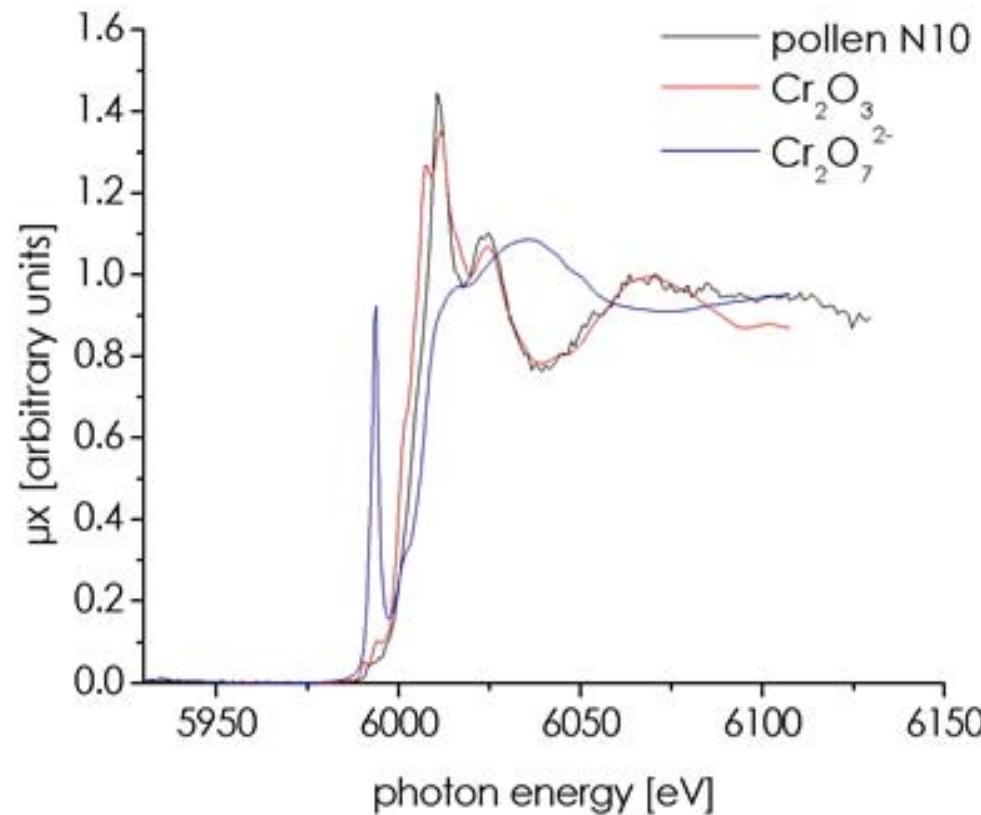
- For ions with unfilled d-electrons bands, the **pd hybridization** is dramatically altered depending on the coordination environment.
- much stronger hybridization for tetrahedral coordination ( $\text{Cr}^{6+}$ ) than for octahedral coordination.



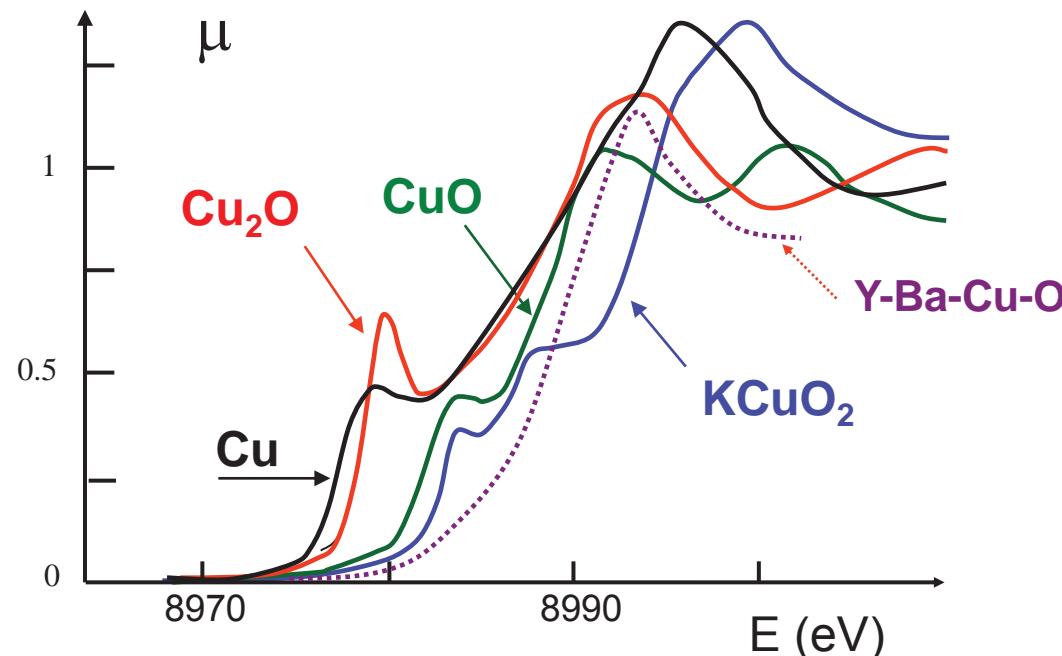
For the case of  $\text{Cr}^{6+}$  pd-hybridization results in a **highly localized molecular orbital state**, giving a well-defined peak below the main absorption edge, indicating a transition to a bound electronic state.

## An example: Chromium (Cr)

K edge XANES : edge energy = 5989.2 eV



Pollen as an indicator of environmental pollution.



(J.B. Boyce et al.  
Phys. Rev. B 1987)

Oxidation Numbers (formal valences)

I Cu<sub>2</sub>O

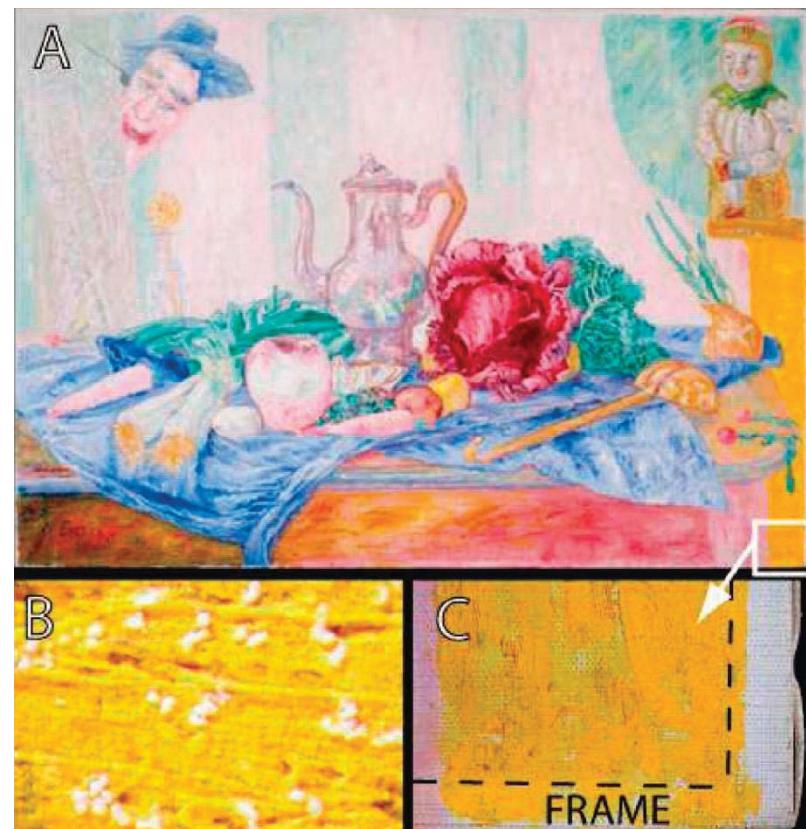
II CuO

III KCuO<sub>2</sub>

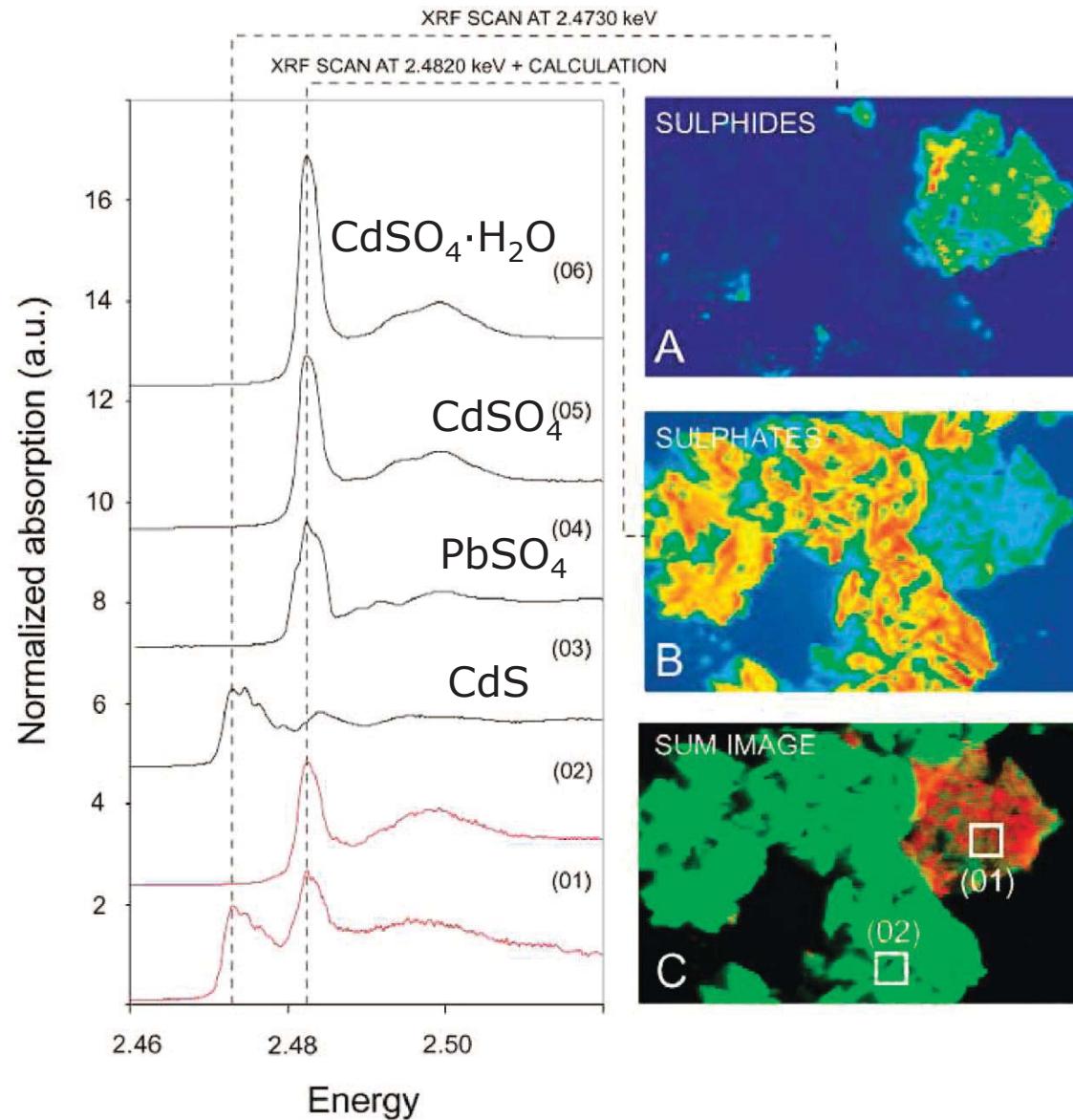
Higher transition energies are expected for higher valence states.

# XANES – example – CdS degradation

**Characterization of a Degraded Cadmium Yellow (CdS) Pigment in an Oil Painting by Means of Synchrotron Radiation Based X-ray Techniques**  
**Geert Van der Snickt et al., Analytical Chemistry, 2009**

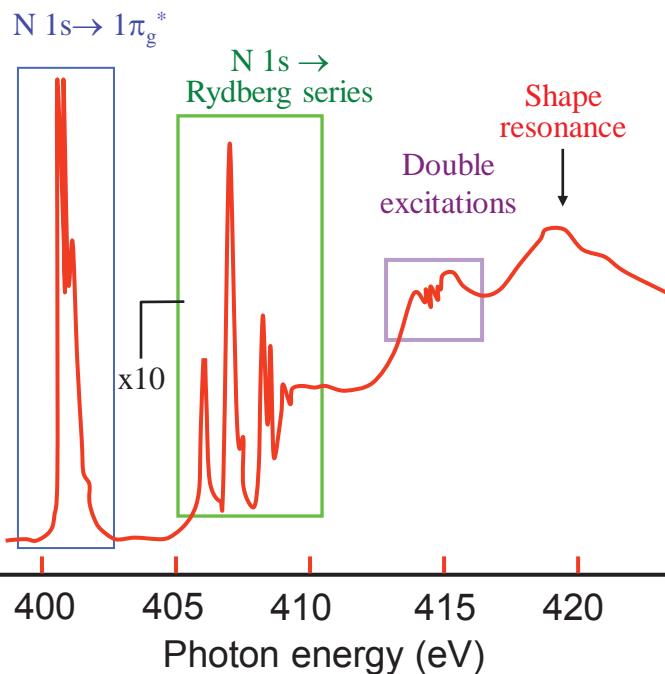


# XANES – example – CdS degradation

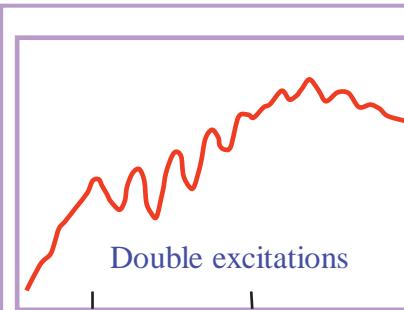


Absorption Intensity ( arb. Units)

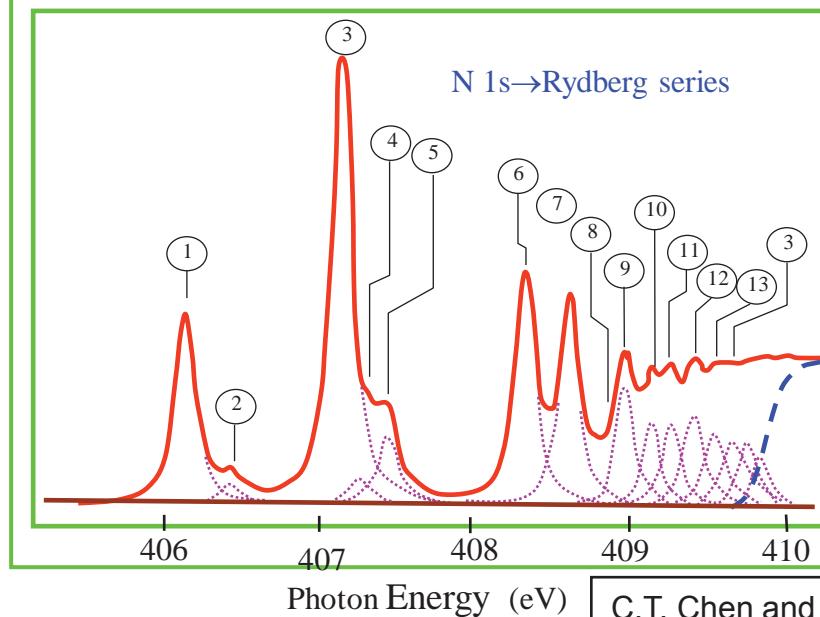
## K-shell - gas-phase N<sub>2</sub>



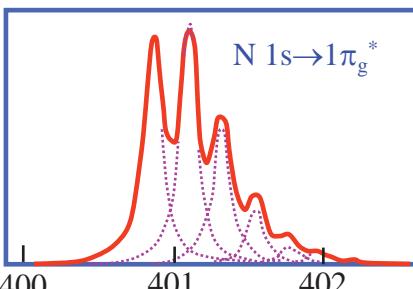
Double excitations associated to the  $1s \rightarrow 1p_g^*$  transition.



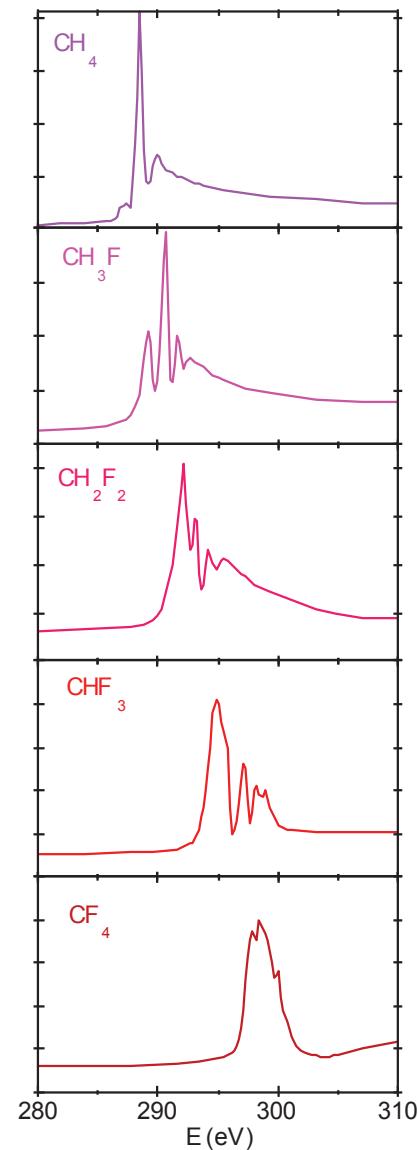
## $1s \rightarrow$ Rydberg series



C.T. Chen and F. Sette,  
Phys. Rev. A 40 (1989)

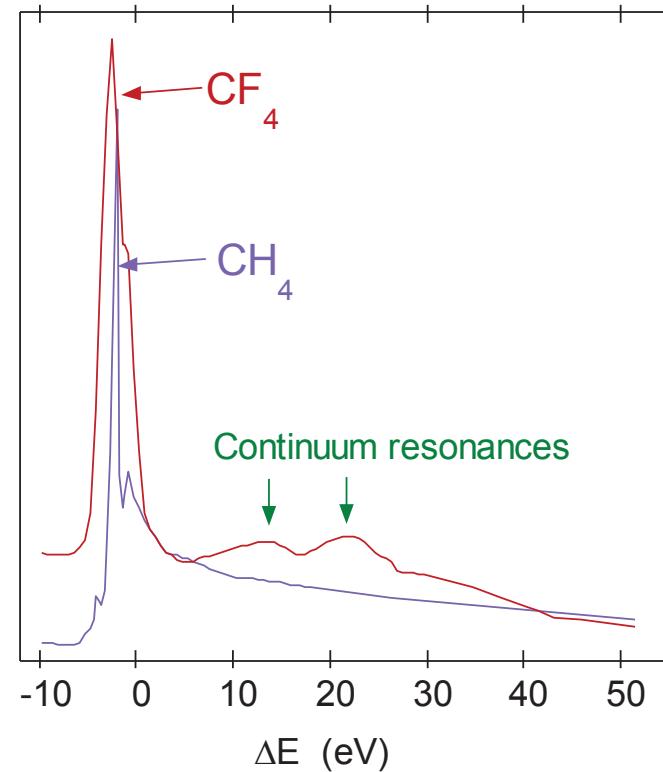


8 vibrational levels  
 $1s \rightarrow 1\pi_g^*$



$\text{CH}_4$

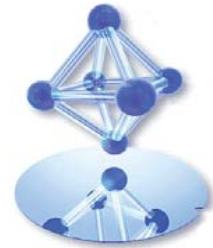
$\text{CF}_4$



Transitions to:

Rydberg states (narrow)  
Valence orbitals (broad)  
Resonances (broad)

# Organic contamination on Si wafers

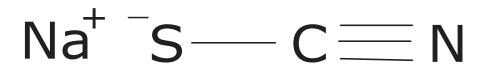
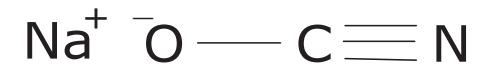
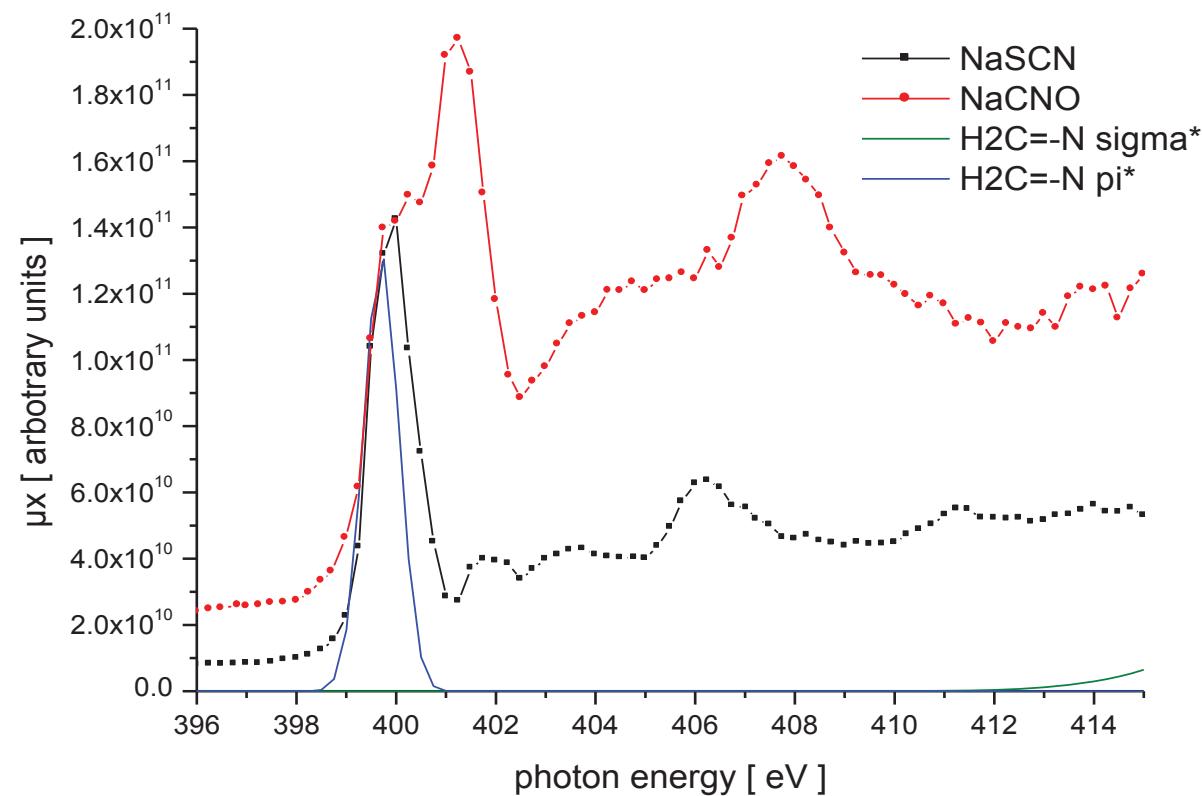


## Motivation

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- Si wafer contamination control is required by the semiconductor industry to produce high quality products with fast decreasing dimensions and prices. Contamination results in quantifiable yield losses.
- the worst are metallic contaminants but:  
ionic and molecular organic contaminants can affect the wettability and the microroughness of the wafer surface, leading either to crystalline defects in the epitaxial layer or even to carbonisation and formation of silicon carbide

# NEXAFS - the systematic of resonance positions

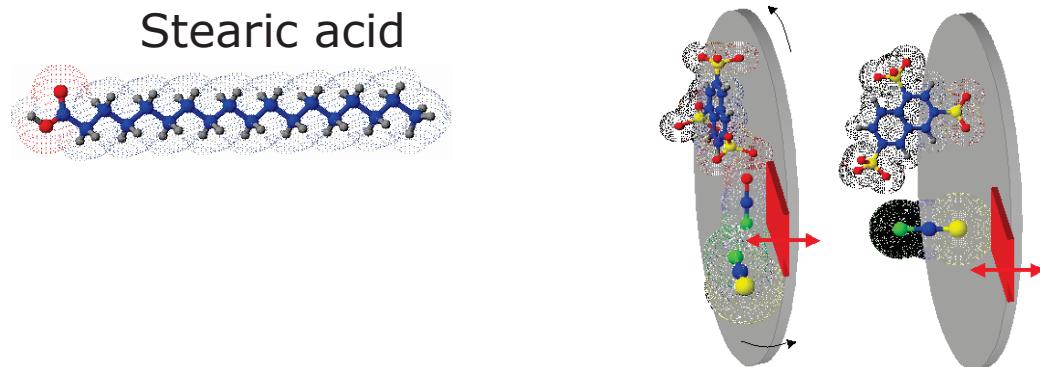
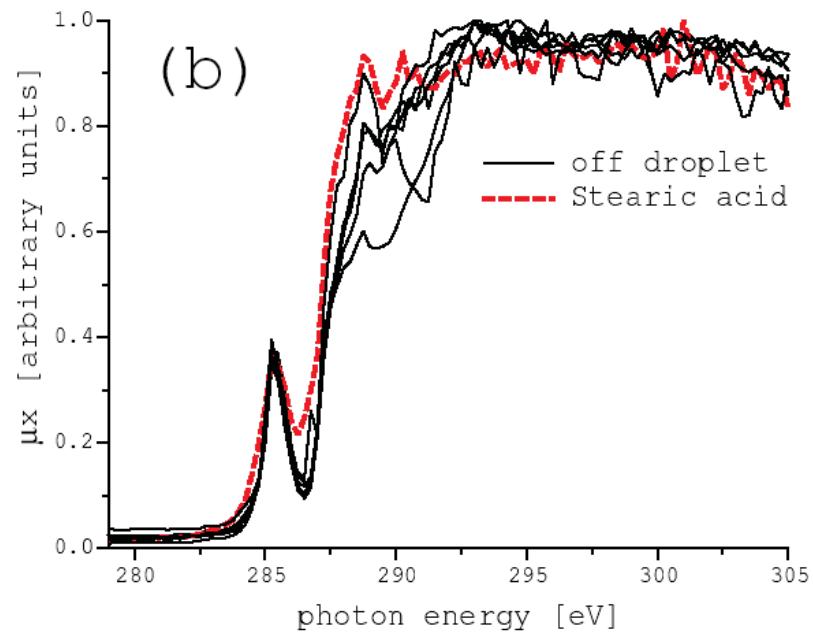
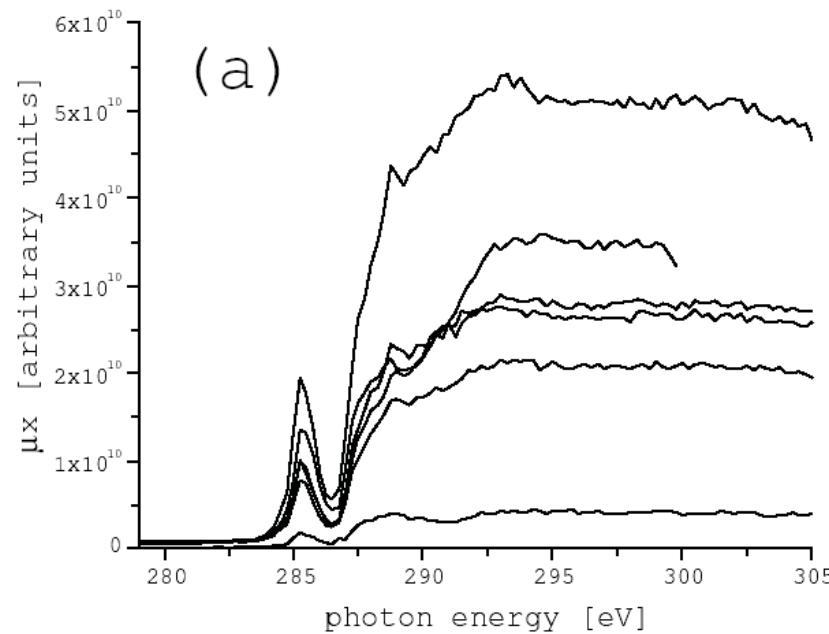


N edge

**from J. Stöhr - NEXAFS Spectroscopy**

bond	energy of resonance [eV]			
carbon	pi*	sigma*		
HC≡CH	285.9	310	nitrogen	
H <sub>2</sub> C=CH <sub>2</sub>	284.7	301	HC≡N	399.7 420.8
H <sub>3</sub> C-CH <sub>3</sub>		291.2	H <sub>3</sub> C-NH <sub>2</sub>	404.6
HC≡N	286.4	307.9	F <sub>3</sub> C-NO	407.5
H <sub>3</sub> C-NH <sub>2</sub>		291.5	N≡N	401 418.9
C≡O	287.3	304	H <sub>2</sub> N-NH <sub>2</sub>	405.1
H <sub>2</sub> C=O	286	300.5	N≡O	399.7 414.5
H <sub>3</sub> C-OH		292	F <sub>3</sub> CN=O	399.1 413.2
H <sub>3</sub> C-F		289.1	N-F <sub>3</sub>	407.1
			oxygen	
			C=O	534.1 550
			H <sub>2</sub> C=O	530.8 544
			H <sub>3</sub> C-OH	537.4
			N≡O	532.7 546.3
			O=O	530.8 540.5
			HO-OH	533
			F <sub>2</sub> -O	534.6

# NEXAFS - organic contamination on wafers



NB.  
Polarise light  
Bond (dipole) orientation

see J. Stöhr - NEXAFS  
Spectroscopy

## Ultratrace speciation of nitrogen compounds in aerosols collected on silicon wafer surfaces by means of TXRF-NEXAFS

S. Török et al., Powder Diffraction **19** (1), March 2004

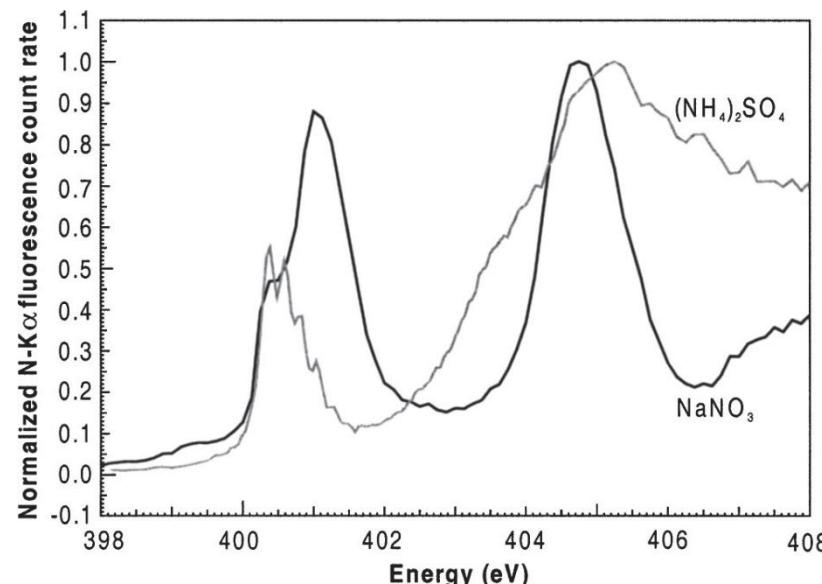


Figure 1. N K-edge TXRF-NEXAFS spectra of ammonium sulfate and sodium nitrate standards prepared on silicon wafers.

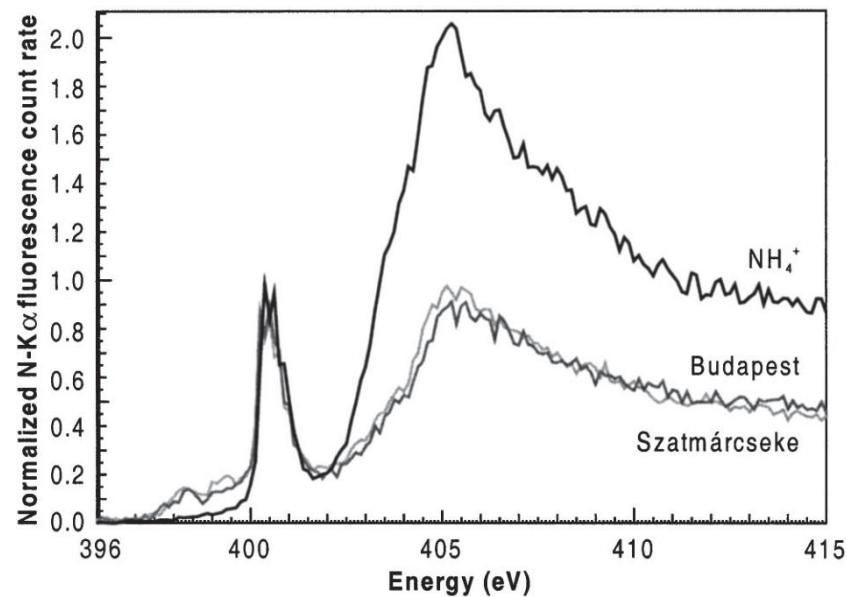


Figure 2. A comparison of N K-edge TXRF-NEXAFS spectra of the submicrometer aerosol samples and of the ammonium sulfate standard.

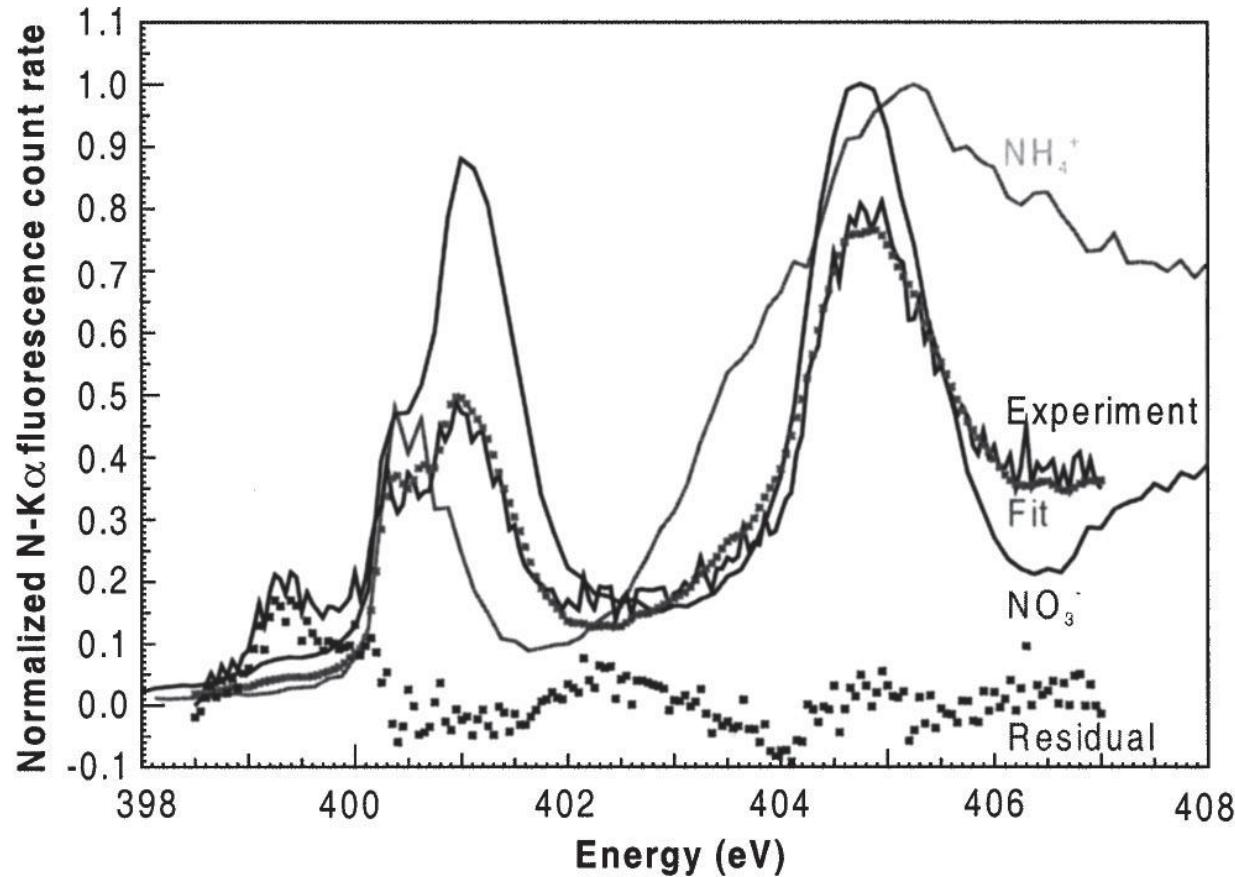
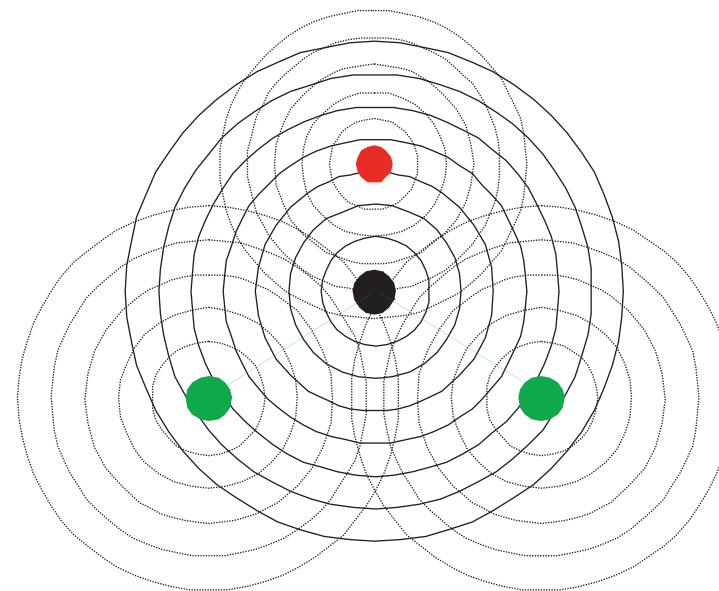


Figure 3. An evaluation of the N K-edge TXRF-NEXAFS spectrum of the 1–2  $\mu\text{m}$  aerosol fraction collected at Szatmarcseke, using a linear combination of the standard spectra of ammonium and nitrate.

# the extended region EXAFS

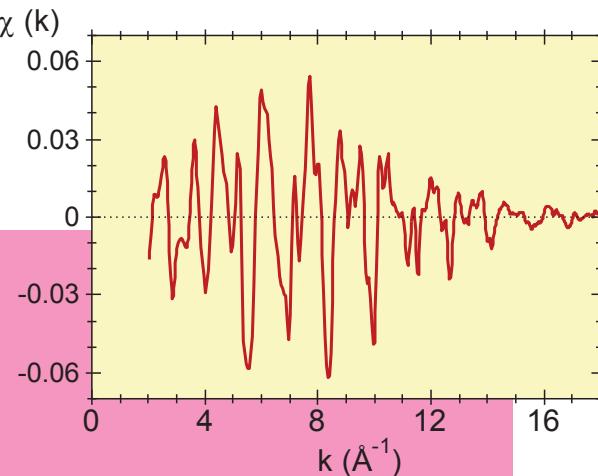


# the EXAFS function

$$\chi(k) = \frac{\mu - \mu_0}{\mu_0}$$

$$\left. \begin{array}{l} \mu_0(\omega) \propto \left| \langle \psi_i | \hat{\eta} \cdot \vec{r} | \psi_f^0 \rangle \right|^2 \\ \mu(\omega) \propto \left| \langle \psi_i | \hat{\eta} \cdot \vec{r} | \psi_f \rangle \right|^2 \end{array} \right\}$$

(weak interaction)



$$|\psi_f\rangle = |\psi_f^0 + \delta\psi_f\rangle$$

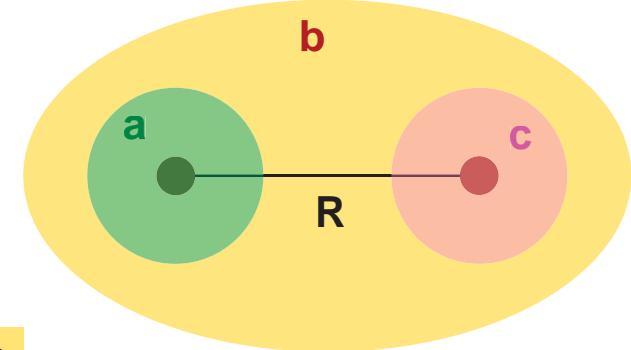
?

Quantum states → wavefunctions

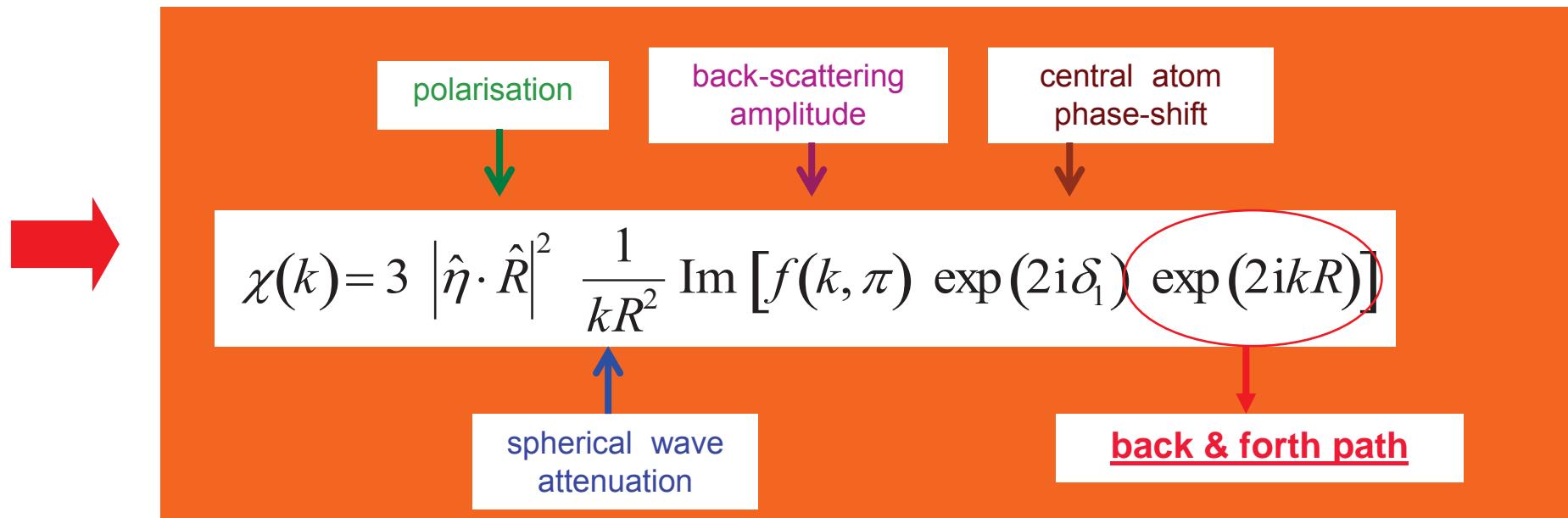
$$\chi(k) = \frac{2 \operatorname{Re} \int d\vec{r} (\psi_i \hat{\eta} \cdot \vec{r} \psi_f^0)^* (\psi_i^* \hat{\eta} \cdot \vec{r} \delta\psi_f)}{\int d\vec{r} |\psi_i^* \hat{\eta} \cdot \vec{r} \psi_f^0|^2}$$

# the EXAFS function

Scattering theory  
in plane-wave approximation



$$\delta\psi_f \propto \psi_f^0 i e^{i\delta} \frac{\exp(ikR)}{2kR} f(k, \pi) \frac{\exp(ikR)}{R} e^{i\delta}$$



# the EXAFS function

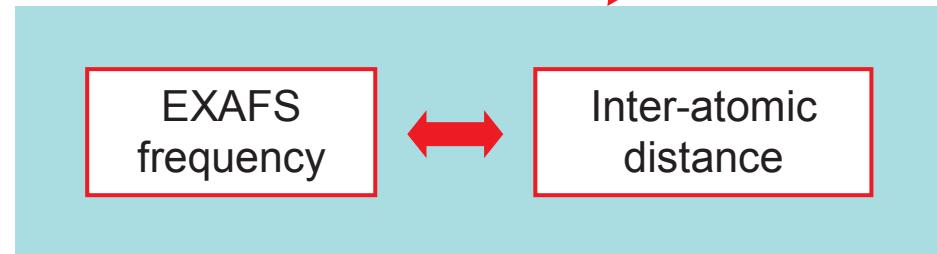
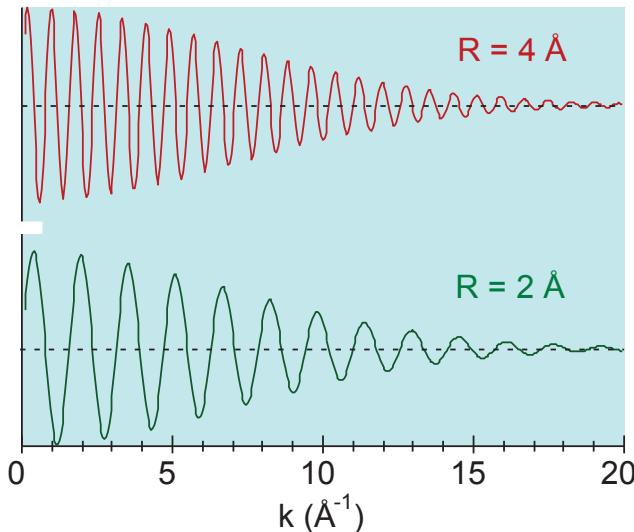
$$\chi(k) = 3 |\hat{\eta} \cdot \vec{R}| \frac{1}{kR^2} \text{Im} [f(k, \pi) \exp(2i\delta_1) \exp(2ikR)]$$



$$f(k, \pi) e^{2i\delta} = |f(k, \pi)| e^{i\phi}$$



$$\chi(k) = 3 |\hat{\eta} \cdot \hat{R}|^2 \frac{1}{kR^2} |f(k, \pi)| \sin [2kR + \phi(k)]$$

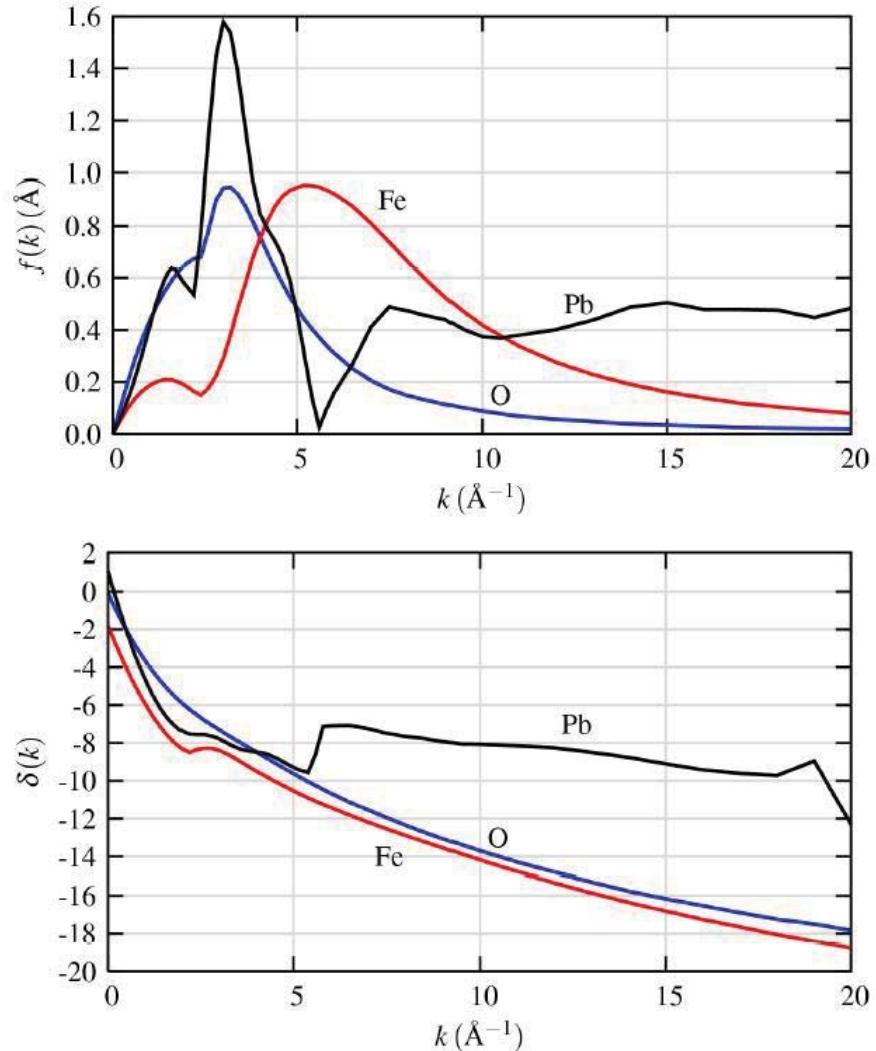


# Scattering amplitude and phase shift

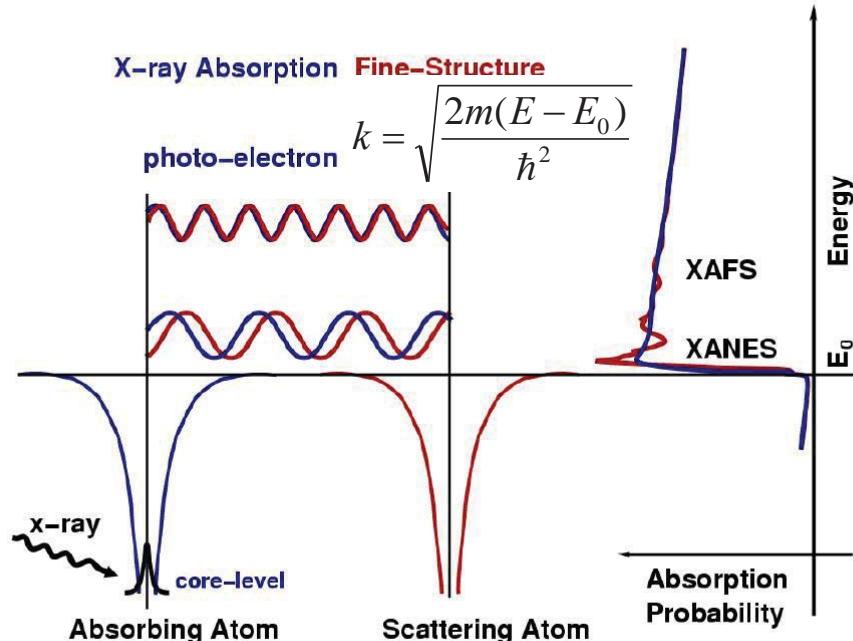
## Examples: scattering properties

$f(k)$  (top) and  $\delta(k)$  (bottom)

for O, Fe, and Pb showing the dependence of these terms on atomic number Z



# the EXAFS function



1. Absorption only takes place if there is an **available state** for the photo-electron:  
i.e. a quantum state at exactly the right energy, and also the right angular momentum state
2. A photo-electron with **wave number  $k$**  is created and propagates away from the atom
3. the photo-electron can **scatter from the electrons of a neighboring atom**, and return to the absorbing atom
4. considering point 1. we see that the presence of the photo-electron scattered back from the neighboring atom **will alter the absorption coefficient**:

This is the origin of XAFS

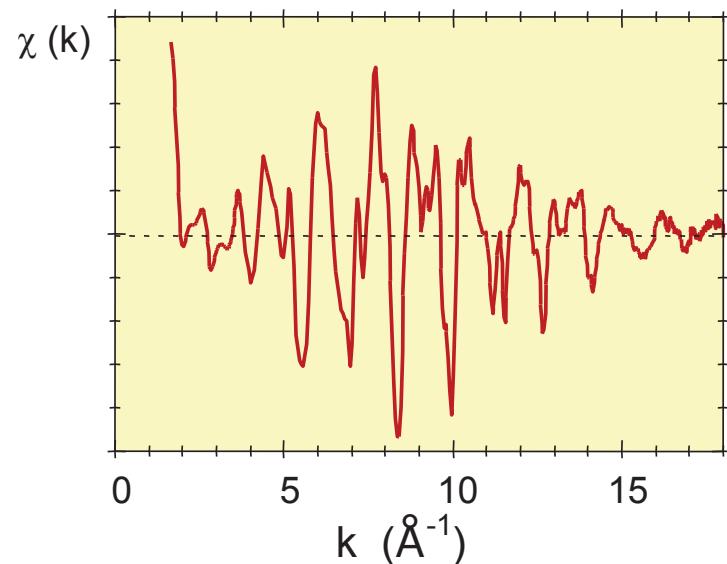
# Inelastic effects – dumping

Intrinsic  
inelastic effects

Photo-electron  
mean-free-path

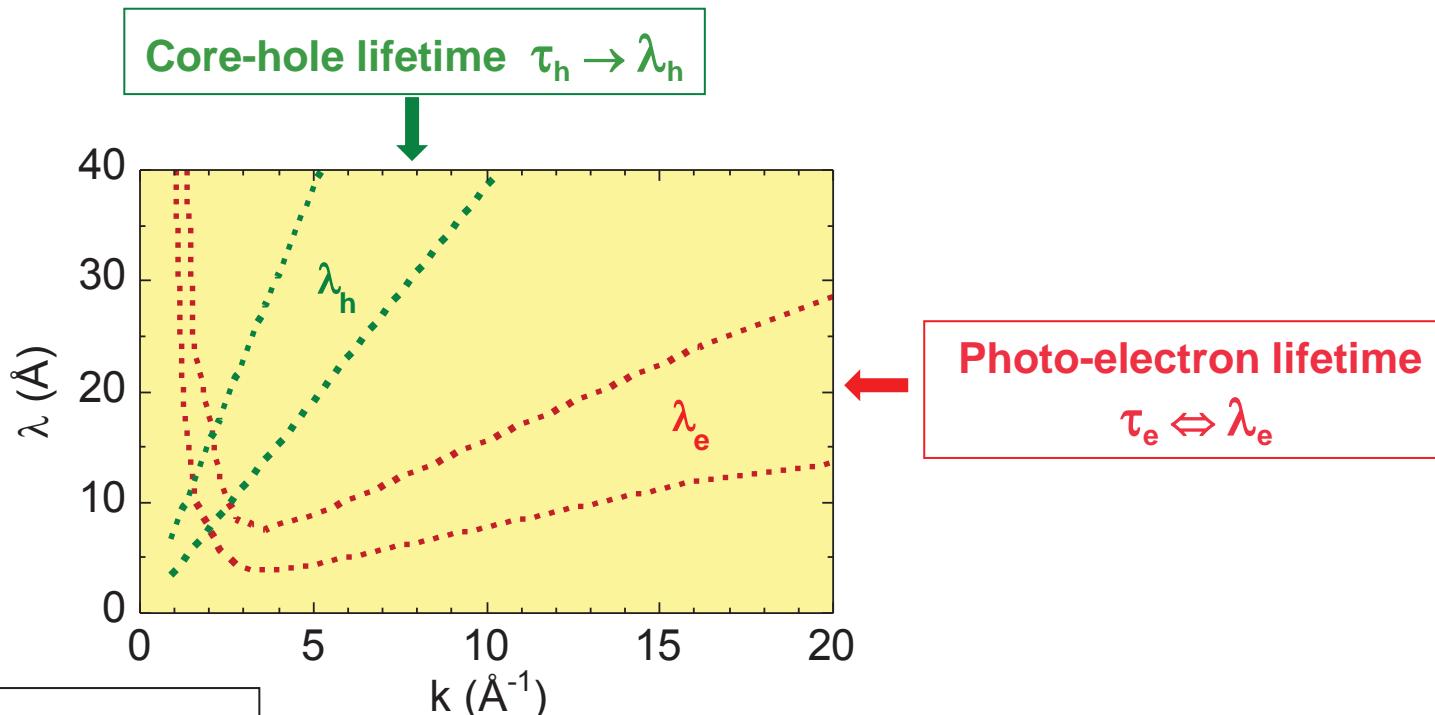
$$\chi(k) = \frac{S_0^2}{k} \sum_{shell} N_s \operatorname{Im} \left[ f_s(k, \pi) e^{2i\delta_l} \frac{e^{-2R_s/\lambda(k)}}{R_s^2} \exp(2ikR_s) \right]$$

**Atoms frozen in equilibrium positions !**



**Disorder effects ?**

# The EXAFS equation – short range



$$\frac{1}{\lambda} = \frac{1}{\lambda_h} + \frac{1}{\lambda_e}$$

$$\psi(k, r) = \frac{e^{ikr} e^{-2r/\lambda(k)}}{kr}$$

Photoelectron  
wavefunction

Attenuation factors

EXAFS Equation

$$\chi(k) = \sum_j \frac{N_j \cdot e^{-2k^2\sigma_j} \cdot e^{-2R_j/\lambda(k)} \cdot f_j(k)}{kR_j^2} \sin[2kR_j + \delta_j(k)]$$

# Disorder

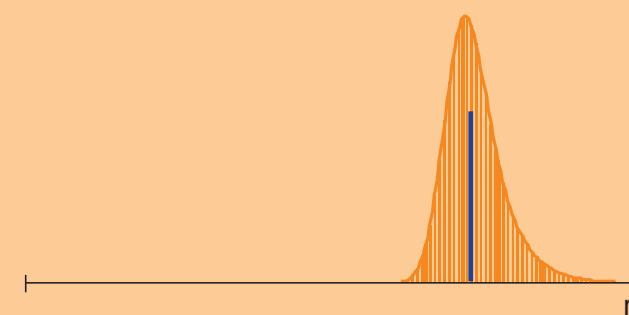
$$\chi(k) = \sum_j \frac{N_j \cdot e^{-2k^2\sigma_j^2} \cdot e^{-2R_j/\lambda(k)} \cdot f_j(k)}{kR_j^2} \sin[2kR_j + \delta_j(k)]$$

## Thermal disorder



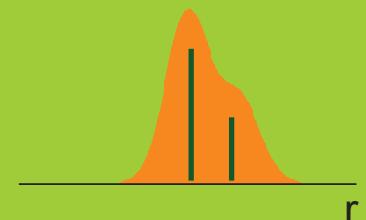
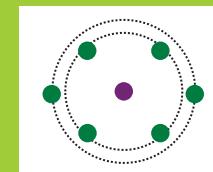
$$\tau_{\text{vib}} \approx 10^{-12} \text{ s}$$

$$\tau_{\text{exafs}} \approx 10^{-15} \text{ s}$$

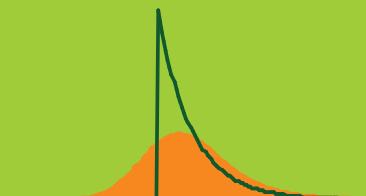
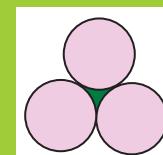


## Structural disorder

- Distorted coordination shells



- Free-volume models



- Sites disorder



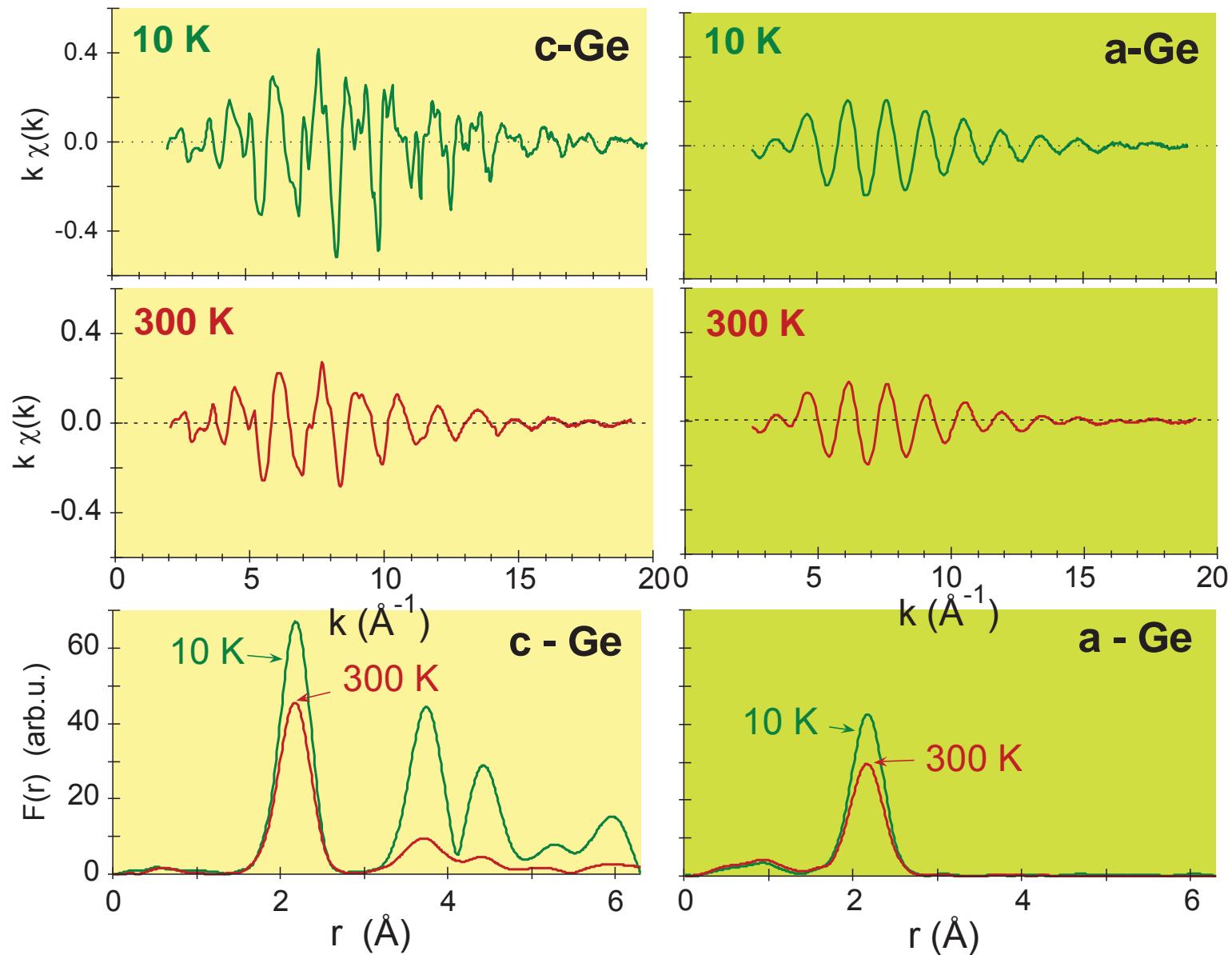
Distance



Distribution of distances

# Disorder

**EXAFS  
signals**

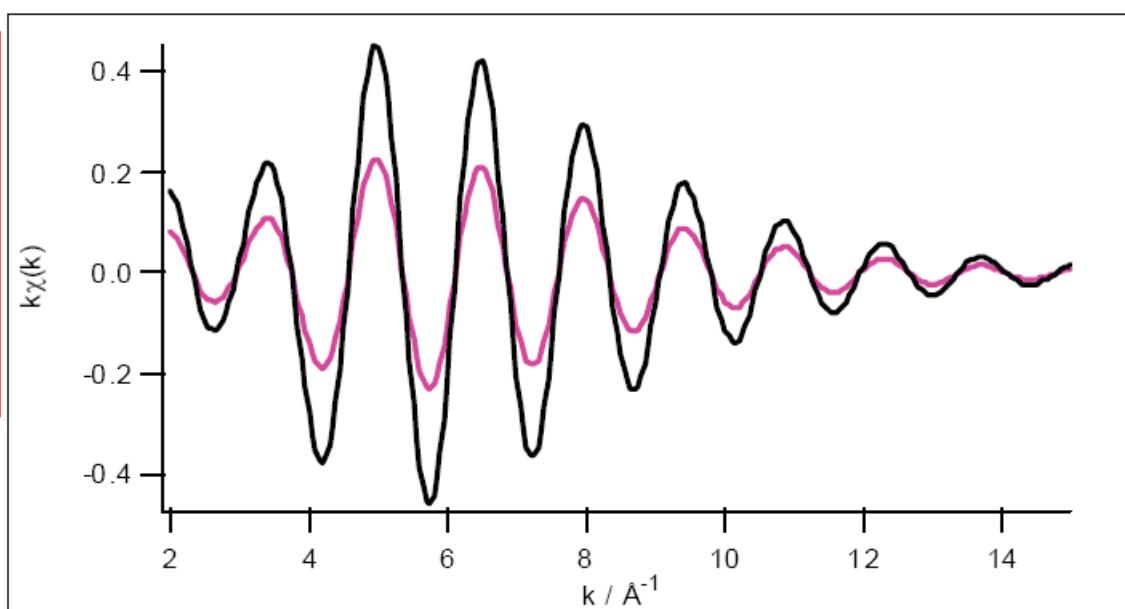


**Fourier  
transforms**

# The EXAFS equation

Theoretical Fe-Fe XAFS,  $R = 2.4855 \text{ \AA}$   $\sigma^2=0.005 \text{ \AA}^2$   
 $N = 4$  compared to  $N = 8$

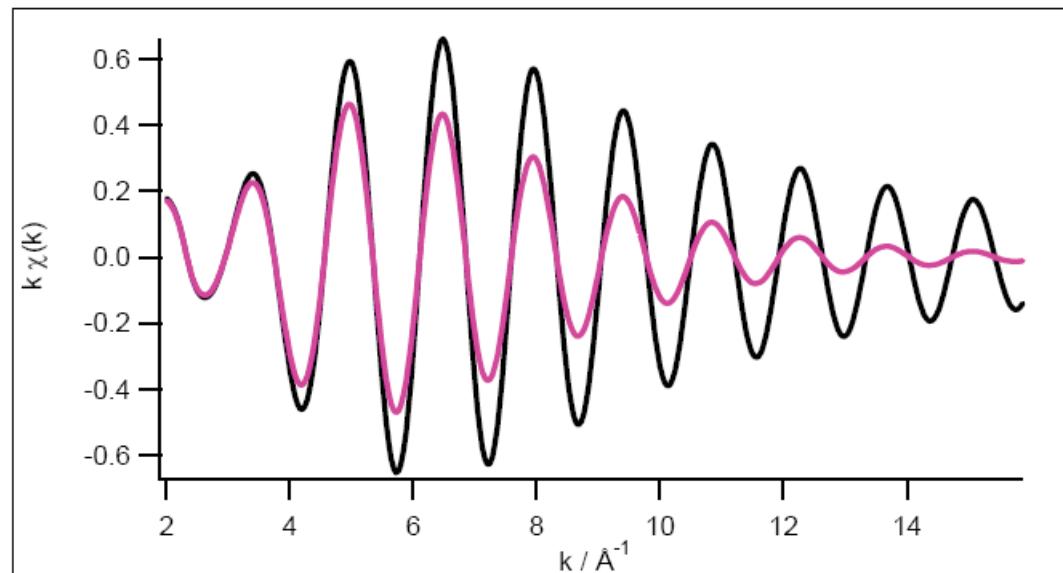
The EXAFS amplitude is proportional to the coordination number



$$\chi(k) = \sum_j^n N_j \times S_i(k) \times F_j(k) \times e^{-2\sigma_j^2 \times k^2} \times e^{-(2r_j / \lambda_j(k))} \times \frac{\sin(2kr_j + \phi_{ij}(k))}{kr_j^2}$$

# The EXAFS equation

**Theoretical Fe-Fe XAFS, N=8 and R = 2.4855 Å  
 $\sigma^2=0.000 \text{ \AA}^2$  compared to 0.005 Å<sup>2</sup>**

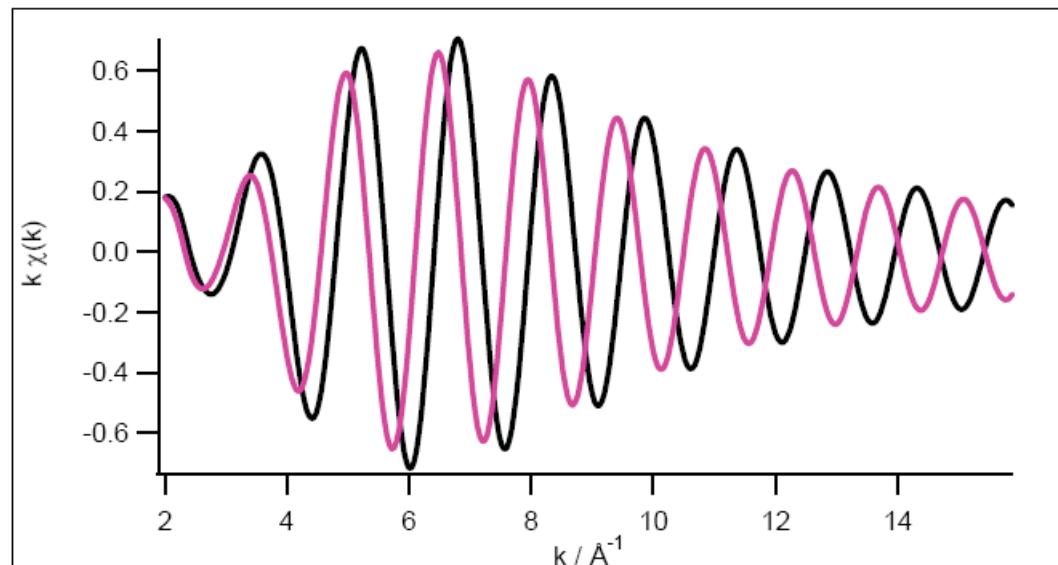


The EXAFS means square displacement dampens the oscillations exponentially

$$\chi(k) = \sum_j^n N_j \times S_i(k) \times F_j(k) \times e^{-2\sigma_j^2 \times k^2} \times e^{-(2r_j / \lambda_j(k))} \times \frac{\sin(2kr_j + \phi_{ij}(k))}{kr_j^2}$$

# The EXAFS equation

**Theoretical Fe-Fe XAFS for N=8 and R=2.3855 compared to R = 2.4855 Å, σ<sup>2</sup>=0.000**



Longer bond distances result in a shorter periodicity of their EXAFS oscillations

$$\chi(k) = \sum_j^n N_j \times S_i(k) \times F_j(k) \times e^{-2\sigma_j^2 \times k^2} \times e^{-(2r_j / \lambda_j(k))} \times \frac{\sin(2kr_j + \phi_{ij}(k))}{kr_j^2}$$

# XAS – experimental acquisition

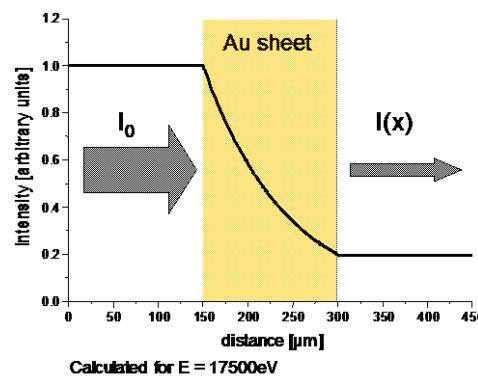
What do we need?

- a “highly” ( $\sim dE/E 10^{-4}$ ) monochromatic but adjustable (in energy) X-ray source (Polychromatic source + monochromator)
- loads of photons

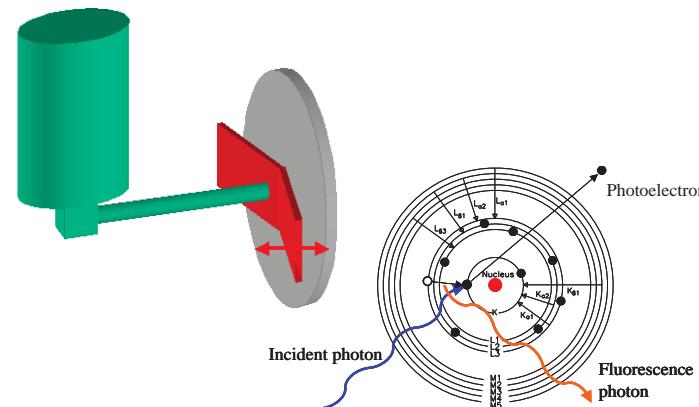


A synchrotron would be good

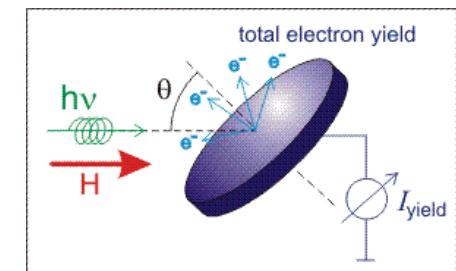
transmitted photons



fluorescence



electrons



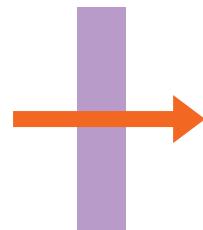
<http://www.fkf.mpg.de/kern/facilities/xmcd.html>

# XAS – experimental acquisition transmission



Bulk information  
(not from surface)  
from:

Thin samples



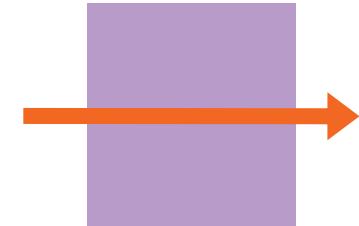
Non-diluted samples

Homogeneous samples

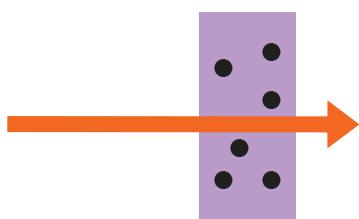
**High accuracy attainable**



Thick samples



Diluted samples



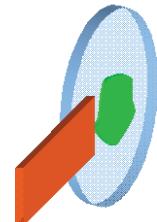
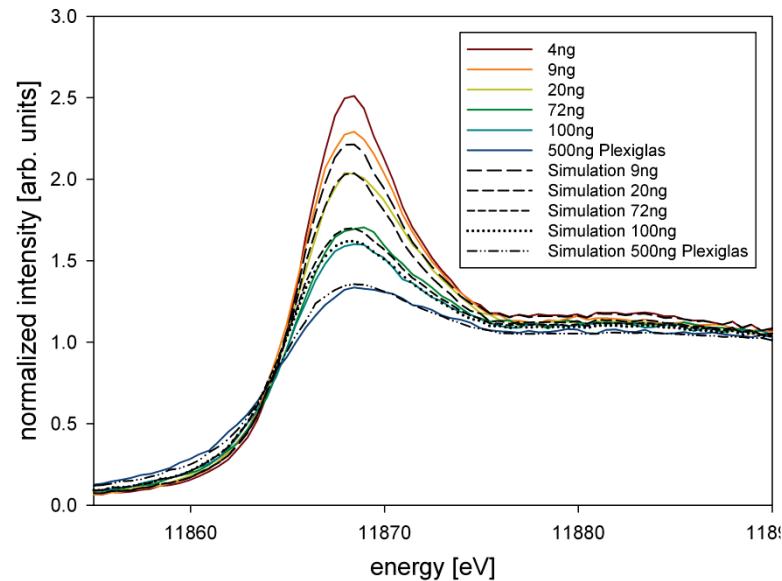
Surface  
information



# XAS – experimental acquisition fluorescence

**Good for diluted samples!**

Arsenic standard solutions  
TXRF geometry

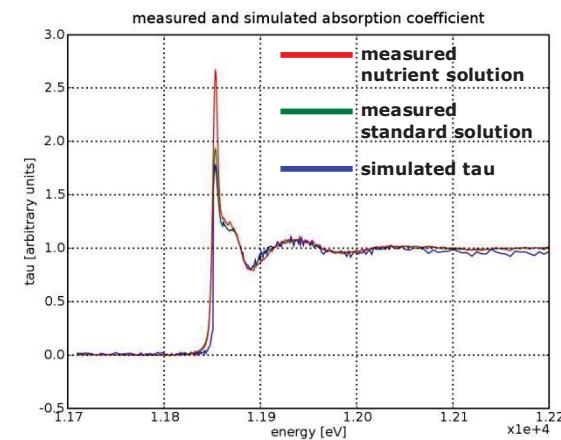
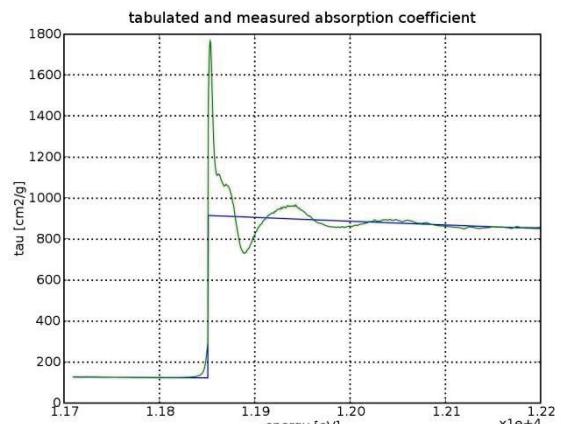


Sample (droplet) size becomes thickness, i.e. path that the primary beam crosses

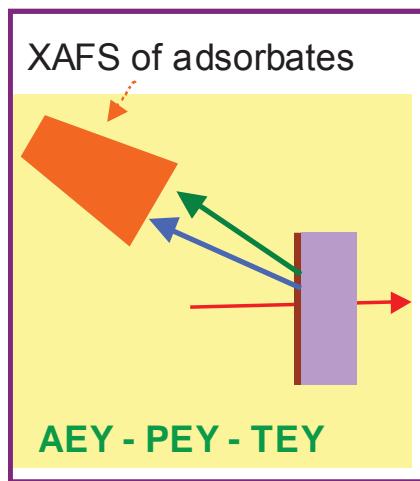
F. Meirer et al., SPECTROCHIMICA ACTA B, vol. 63, n. 12, 2008, pp. 1496-1502

**BUT: watch out for self absorption in “concentrated samples”**

“Simple”  
MONTE  
CARLO  
simulation



# XAS – experimental acquisition electrons



## **AEY = Auger Electron Yield**

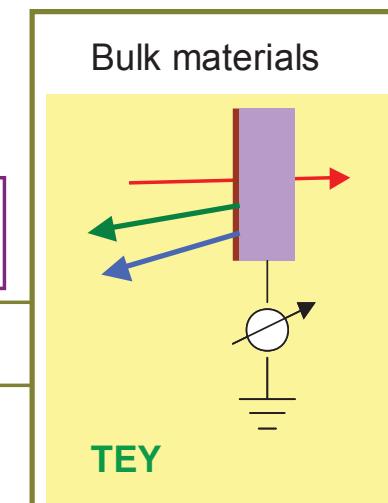
- narrow energy window
- only direct Auger electrons
- spurious structures from photoelectrons

## **PEY = Partial Electron Yield**

- large energy window
- Auger (direct + secondary) = XAFS signal
- Photoel. (direct + secondary) = background

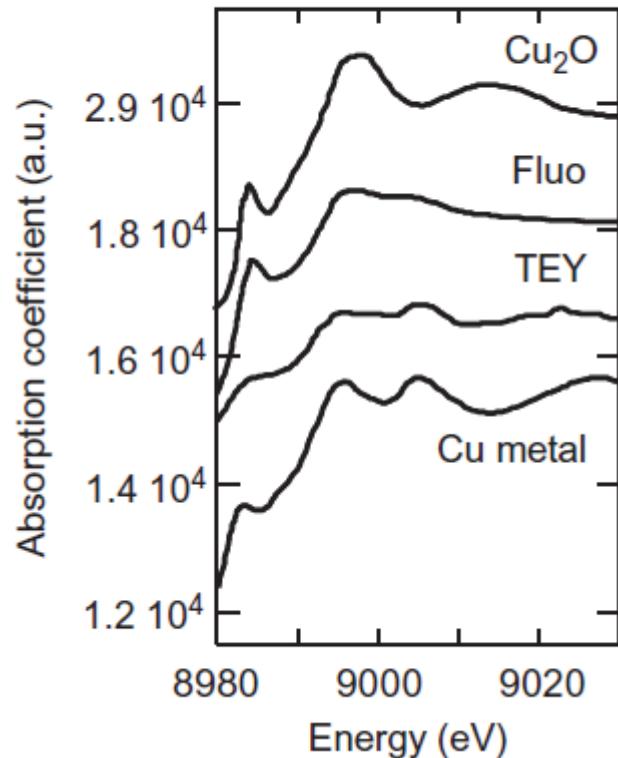
## **TEY = Total Electron Yield**

- all electrons collected
- Auger (direct + second.) = XAFS signal
- Photoel. (direct + second.) = background
- XAFS from Auger and photoel.



# XAS – experimental acquisition Electrons vs. fluorescence

K. Janssens



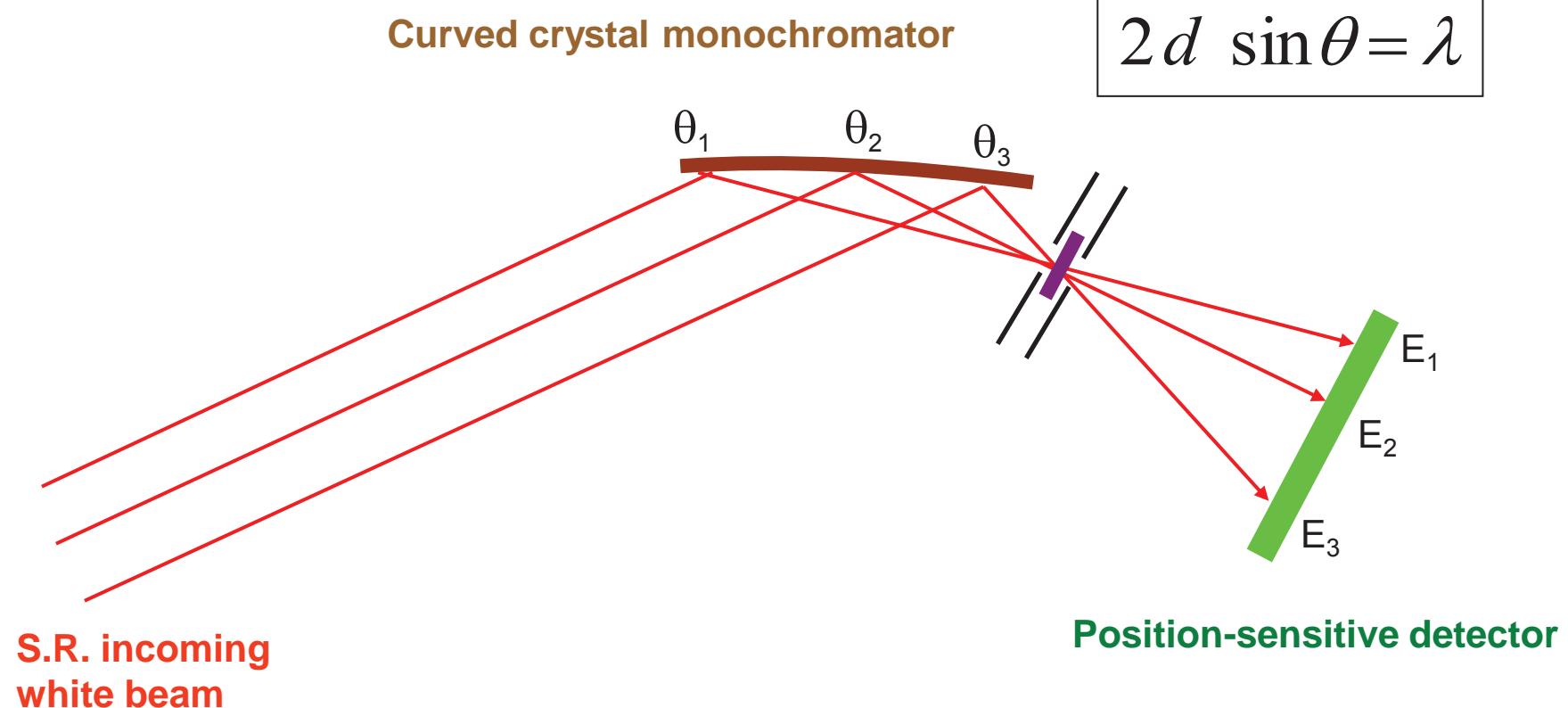
## Copper in lustre

"When XANES data are recorded in the fluorescence mode of detection, where the lustre layer is probed down to a depth of 50 µm, Cu appears to be predominantly present under an oxidized form (i.e., as  $\text{Cu}^+$  in  $\text{Cu}_2\text{O}$ , cuprite). However, if the TEY mode of detection is used, where only the top 100 nm of the same lustre layer is analysed, a different profile is obtained, showing the presence of Cu in the metallic state."

XANES spectra obtained in fluorescence (fluo) and total electron yield (TEY) mode of the red lustre, compared to the XANES profiles of  $\text{Cu}_2\text{O}$  and metallic Cu.

# XAS – experimental acquisition

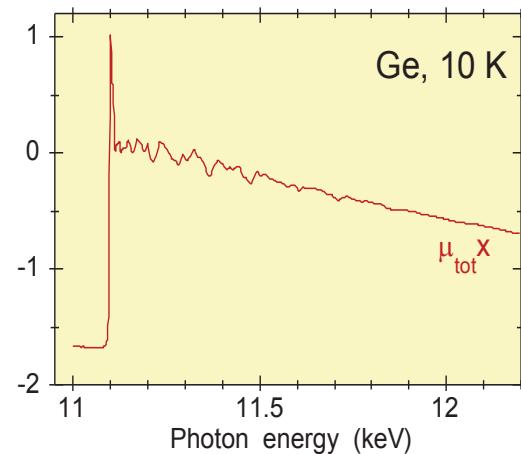
## Dispersive EXAFS



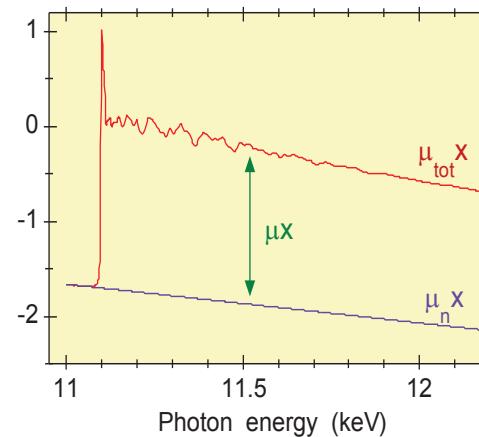
No high resolution but very fast  
Follow dynamics of non equilibrium systems

# XAS – data analysis

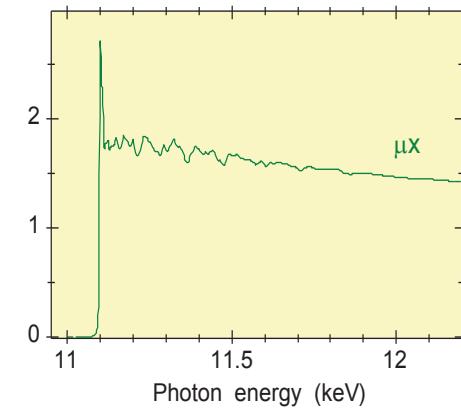
Experimental  
signal



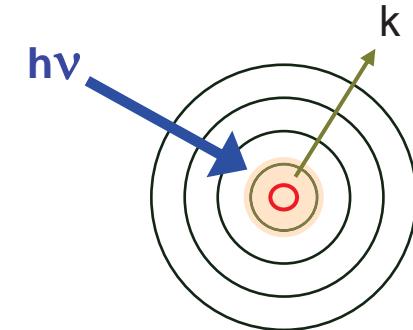
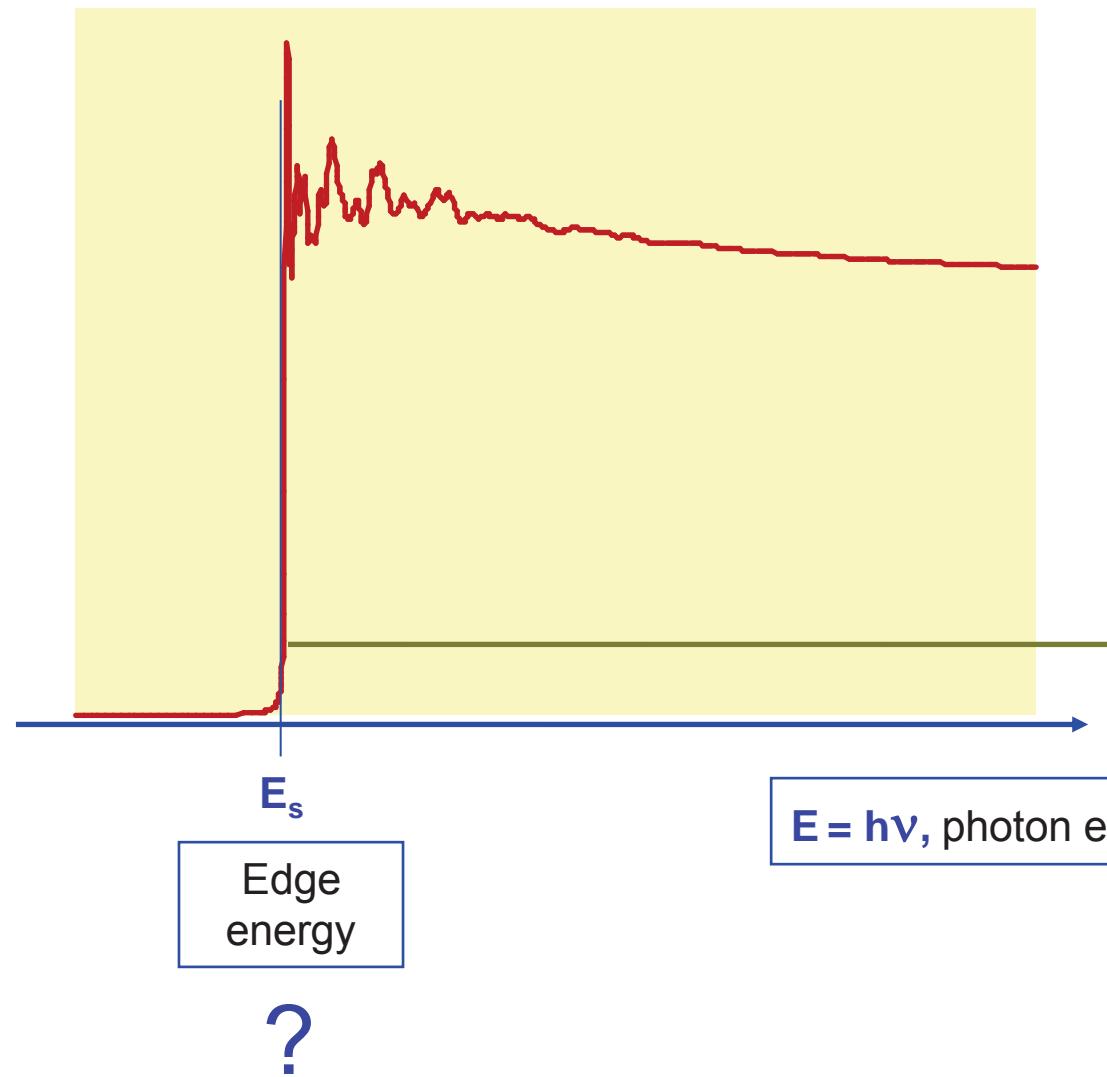
Extrapolation  
of pre-edge  
behaviour



Specific  
absorption coefficient



# XAS – data analysis



Photoelectron  
wavenumber

$$k = \sqrt{\frac{2m}{\hbar^2} (h\nu - E_s)}$$

$E = h\nu$ , photon energy

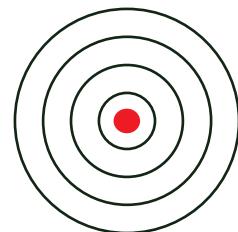
# XAS – data analysis

EXAFS function

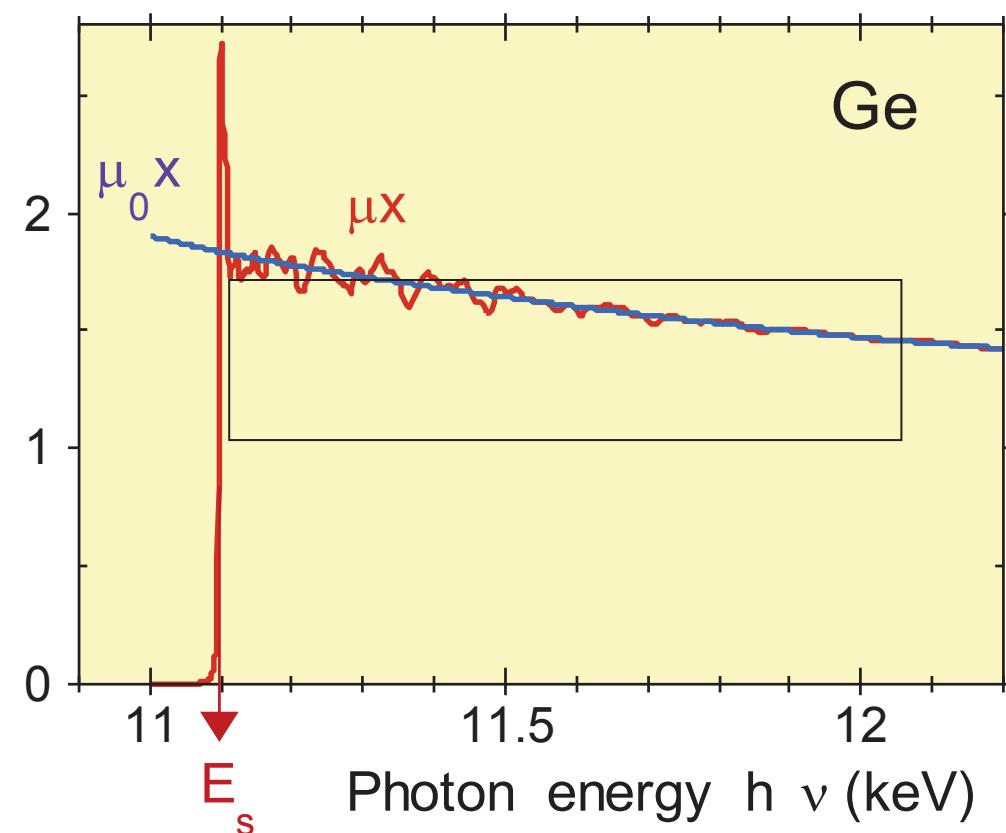
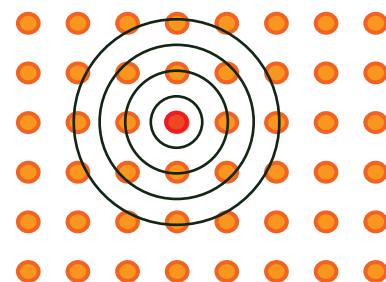
$$\chi(k) = \frac{\mu - \boxed{\mu_0}}{\mu_0}$$

$\boxed{\mu_0 ?}$

Isolated atom



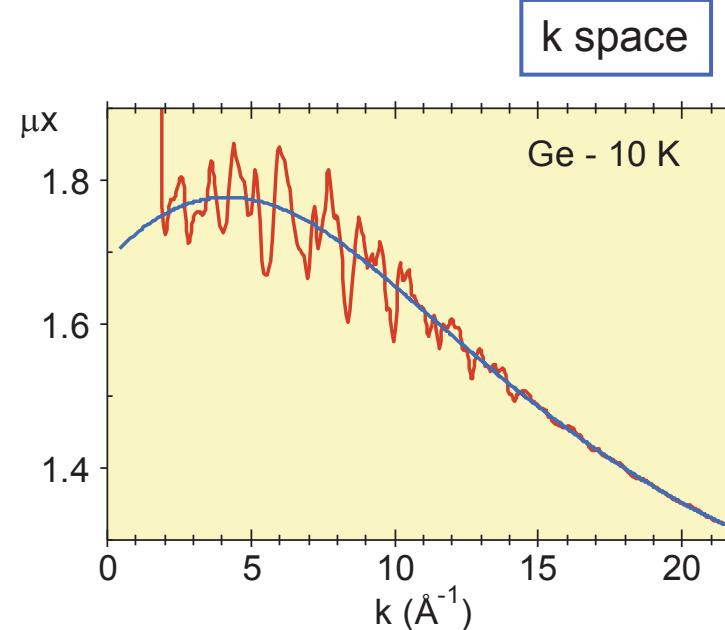
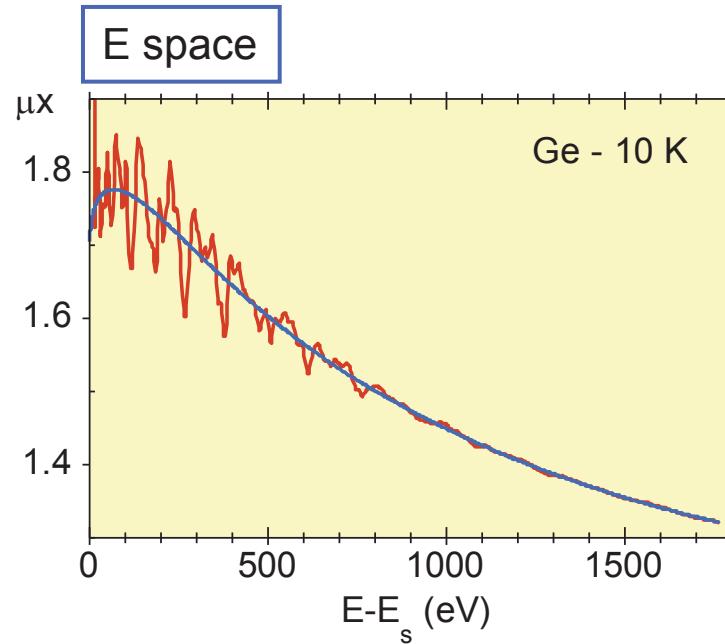
Embedded atom



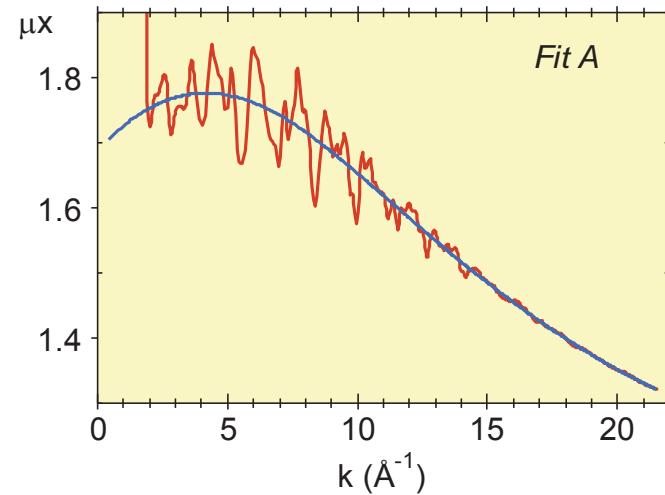
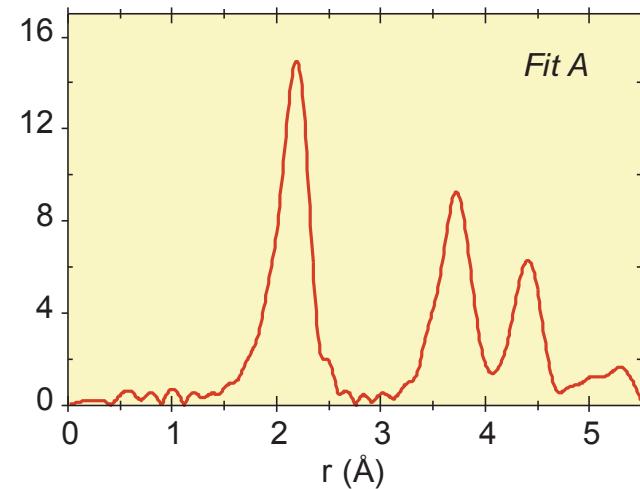
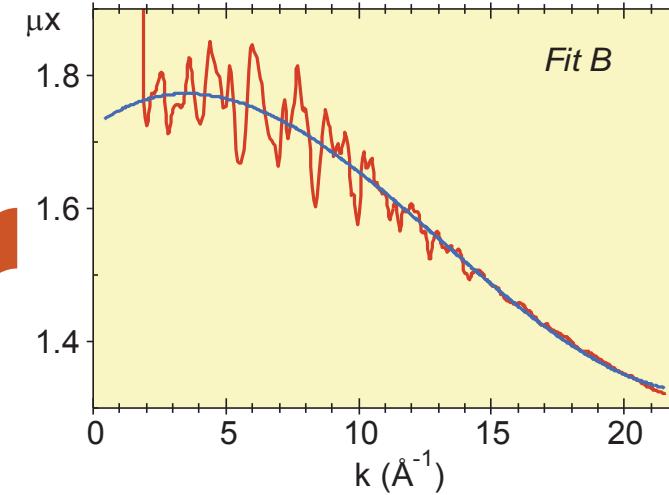
# XAS – data analysis

$$\chi(k) = \frac{\mu - \mu_0}{\mu_0}$$

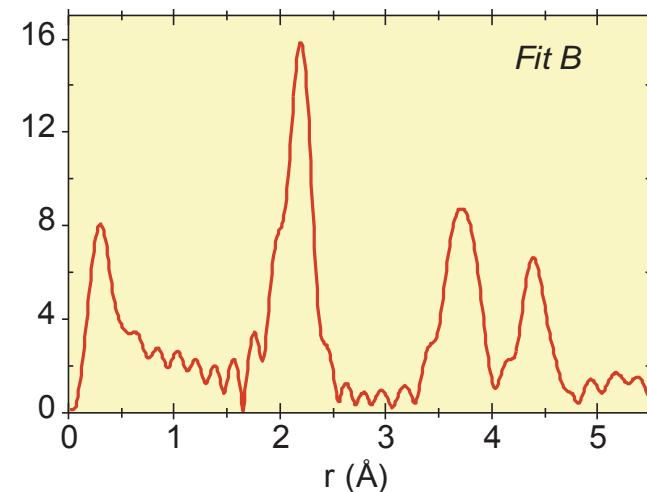
Polynomial spline - best fit



# XAS – data analysis

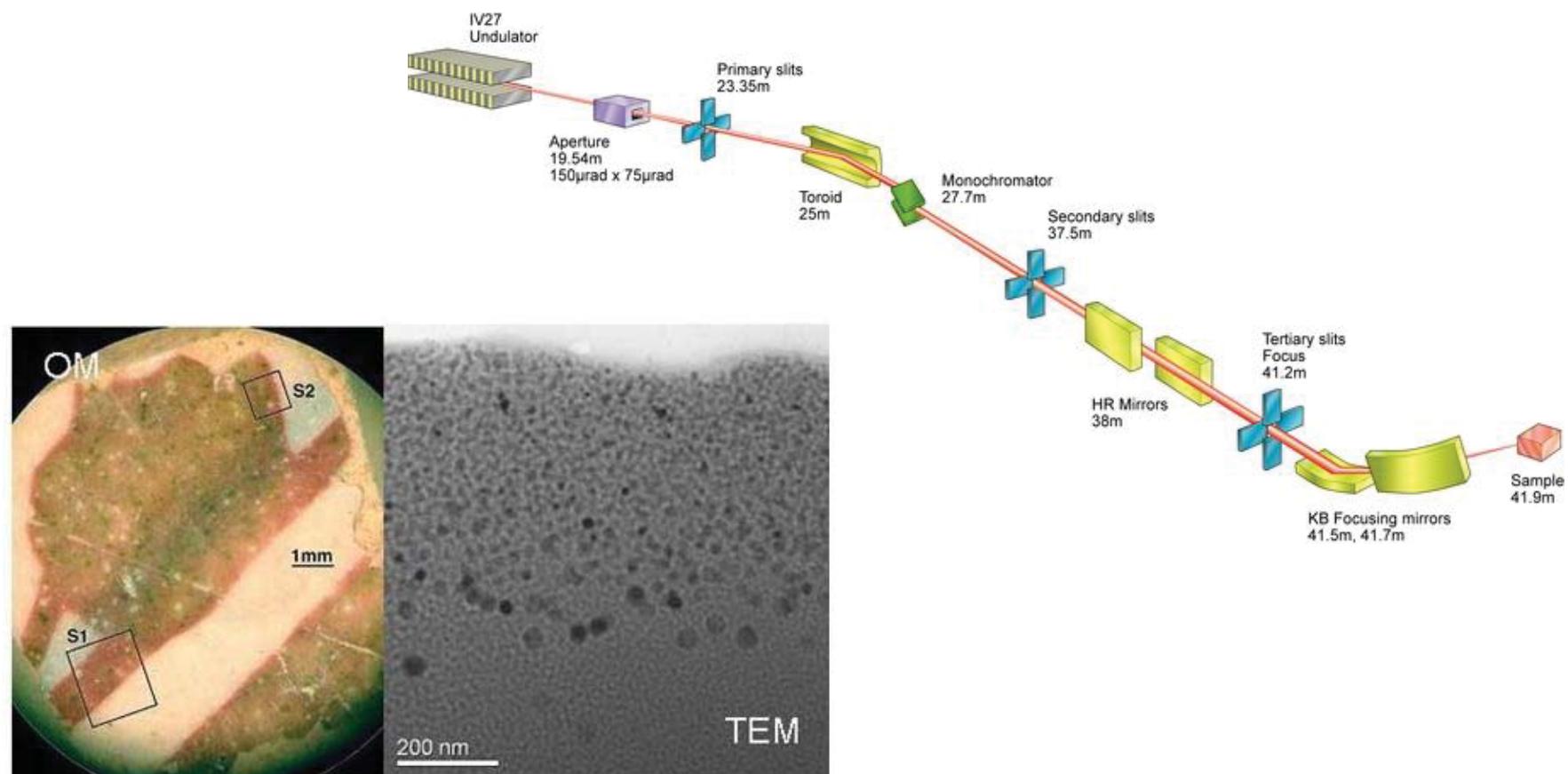

 $\mu - \mu_0$ 


Fourier  
transf.

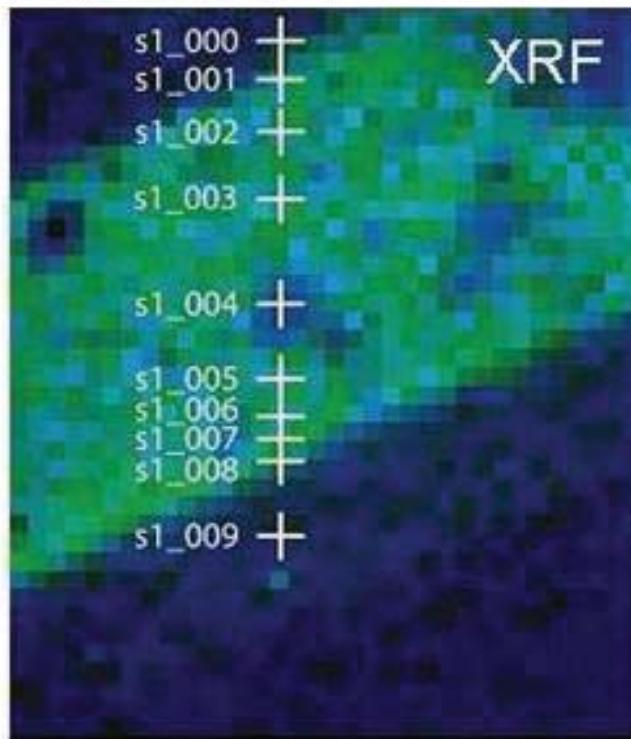
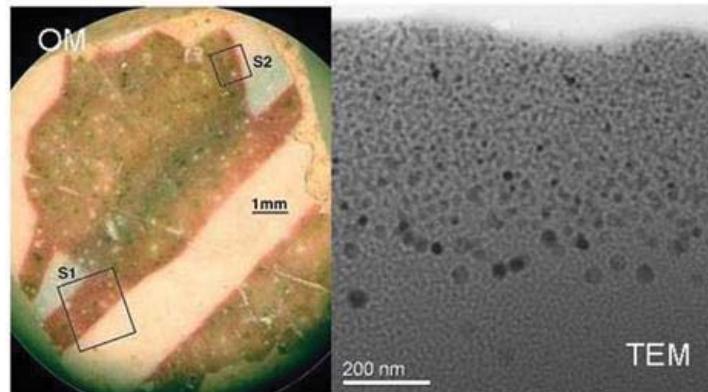


# EXAFS example – Copper in lustre

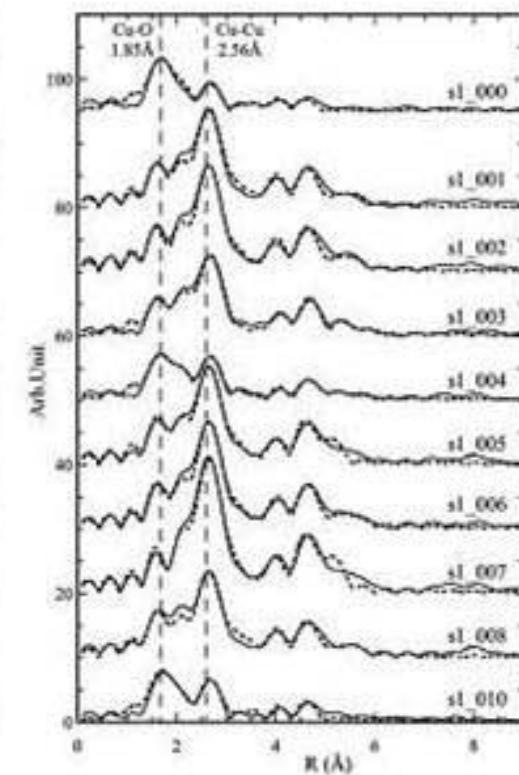
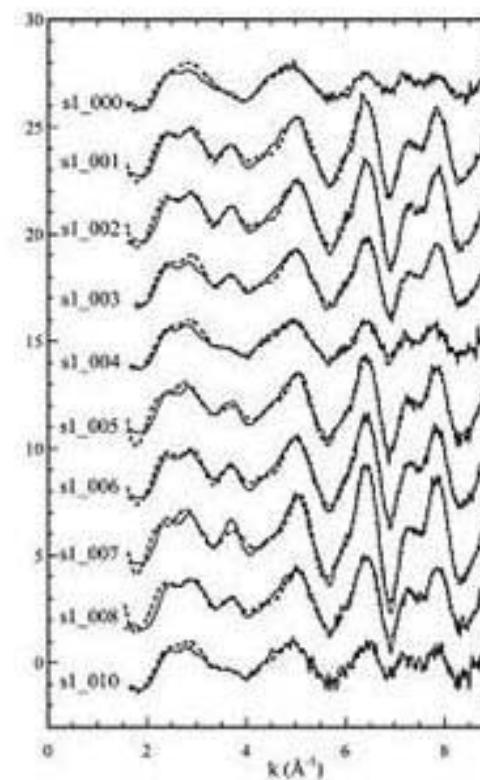
**Microfocus spectroscopy experiments on mineral deposits and in cultural heritage science at Diamond Light Source.**  
*Josep Roque-Rosell1 et al.*



# EXAFS example – Copper in lustre



*Ceramic lustre from Paterna (Spain). Sampled region, TEM pictures showing copper nanoparticles, XRF mapping for Cu Ka, Cu K edge EXAFS spectra obtained across the lustre decoration and magnitude of the Fourier Transform of the EXAFS spectra.*



# XANES ↔ EXAFS

## Multiple scattering

$$\chi(k) = \frac{\mu - \mu_0}{\mu_0}$$

Single Scattering

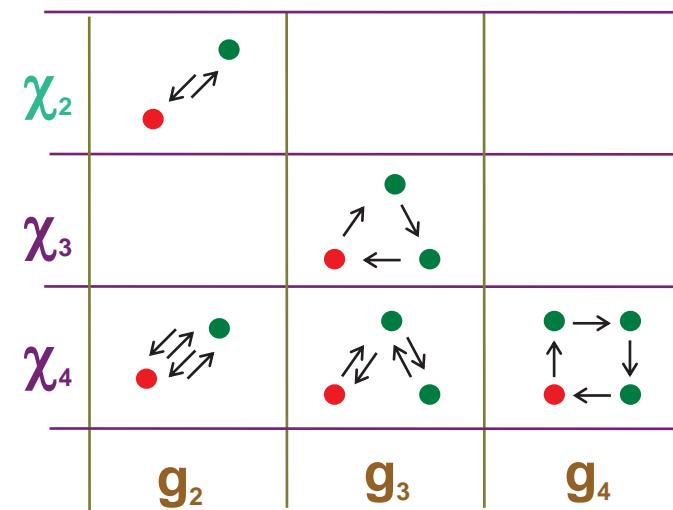
$$\mu(k) = \mu_0(k)[1 + \chi(k)]$$

Multiple Scattering

$$\mu(k) = \mu_0(k)[1 + \chi_2(k) + \chi_3(k) + \chi_4(k) + \dots]$$

$$\chi_n(k)$$

Contribution from all n-order paths



$g_n$  = n-body correlation function

## XAS - References

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