



The Abdus Salam
International Centre
for Theoretical Physics



2354-30

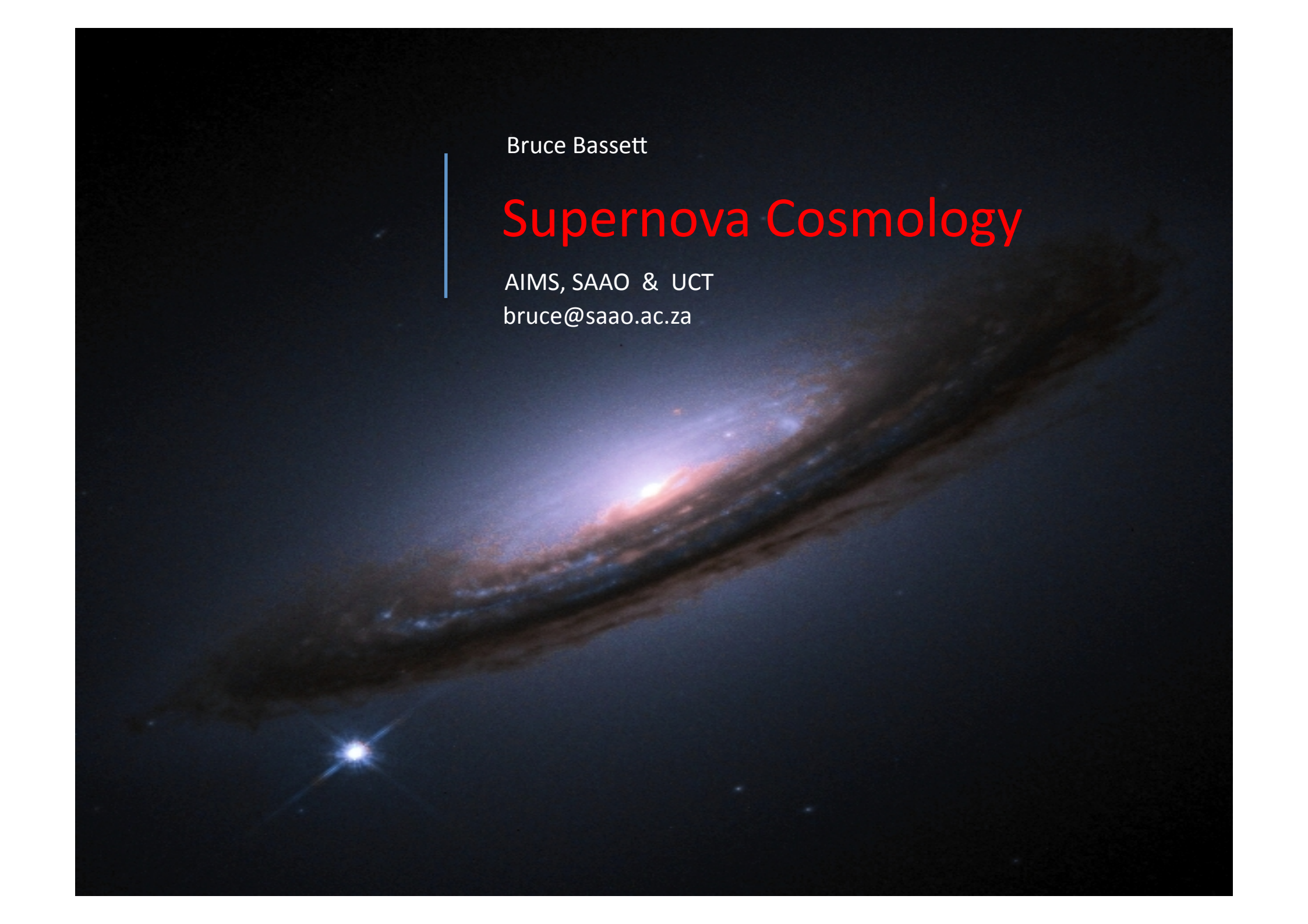
Summer School on Cosmology

16 - 27 July 2012

Supernova Cosmology

B. Bassett

AIMS, SAAO & UCT



Bruce Bassett

Supernova Cosmology

AIMS, SAAO & UCT

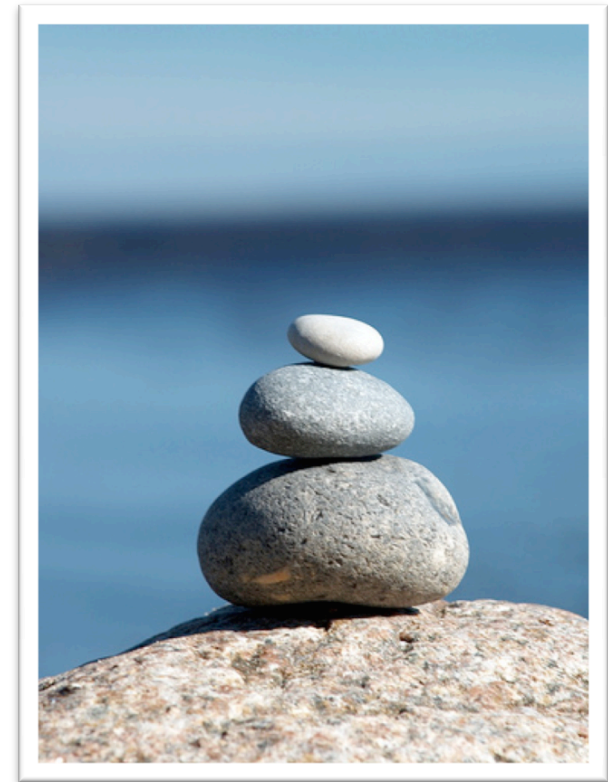
bruce@sao.ac.za

Parkinson Challenge

- Download the Union SNIa data online, write an **MCMC** algorithm in your favourite language and derive constraints on Ω_m , Ω_λ marginalising over H_0 .
- **Time frame**: finish before I stop talking.

Overview

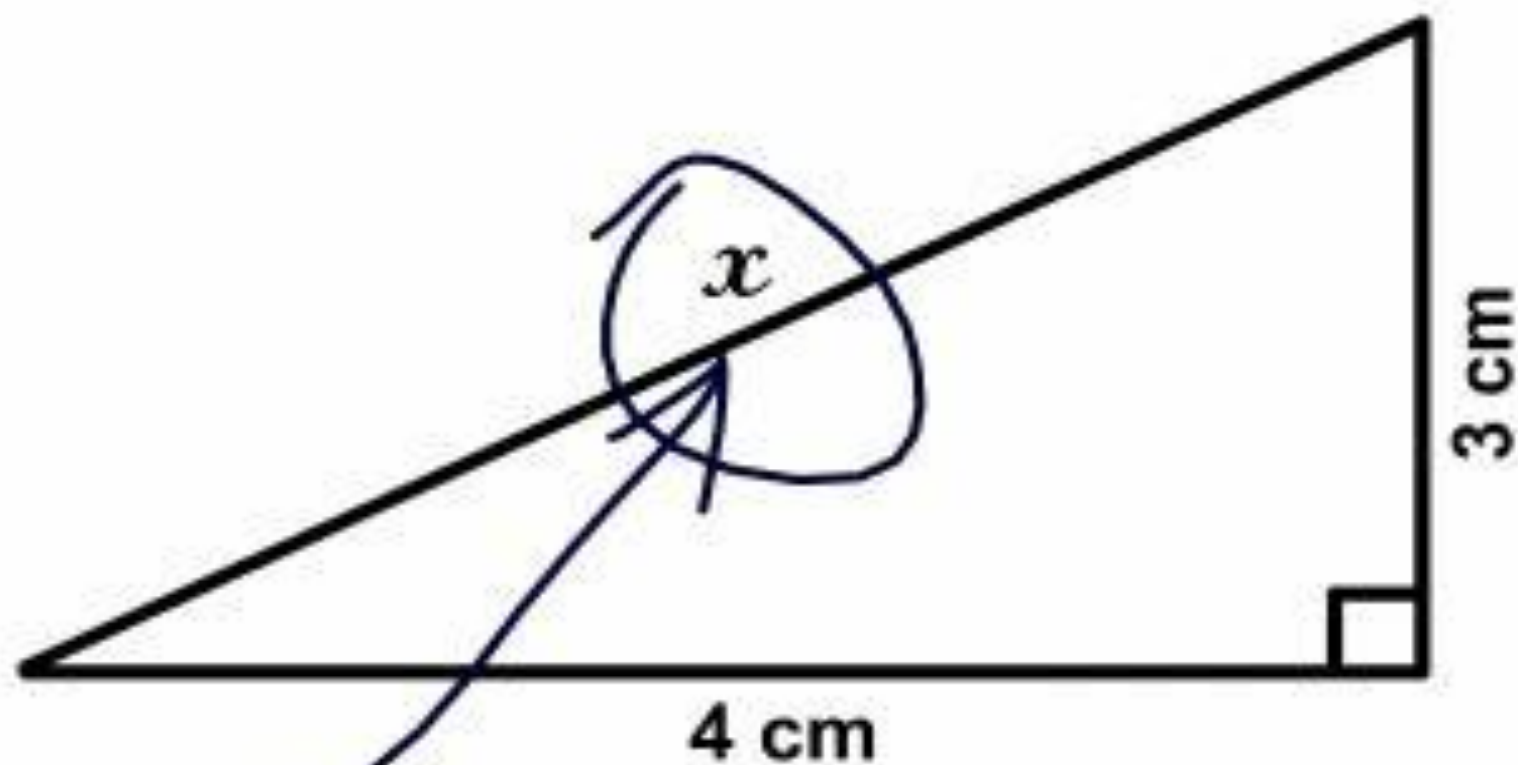
- What distances do
- The Good: SNIa and the Miracle
- The Ugly
- The Future



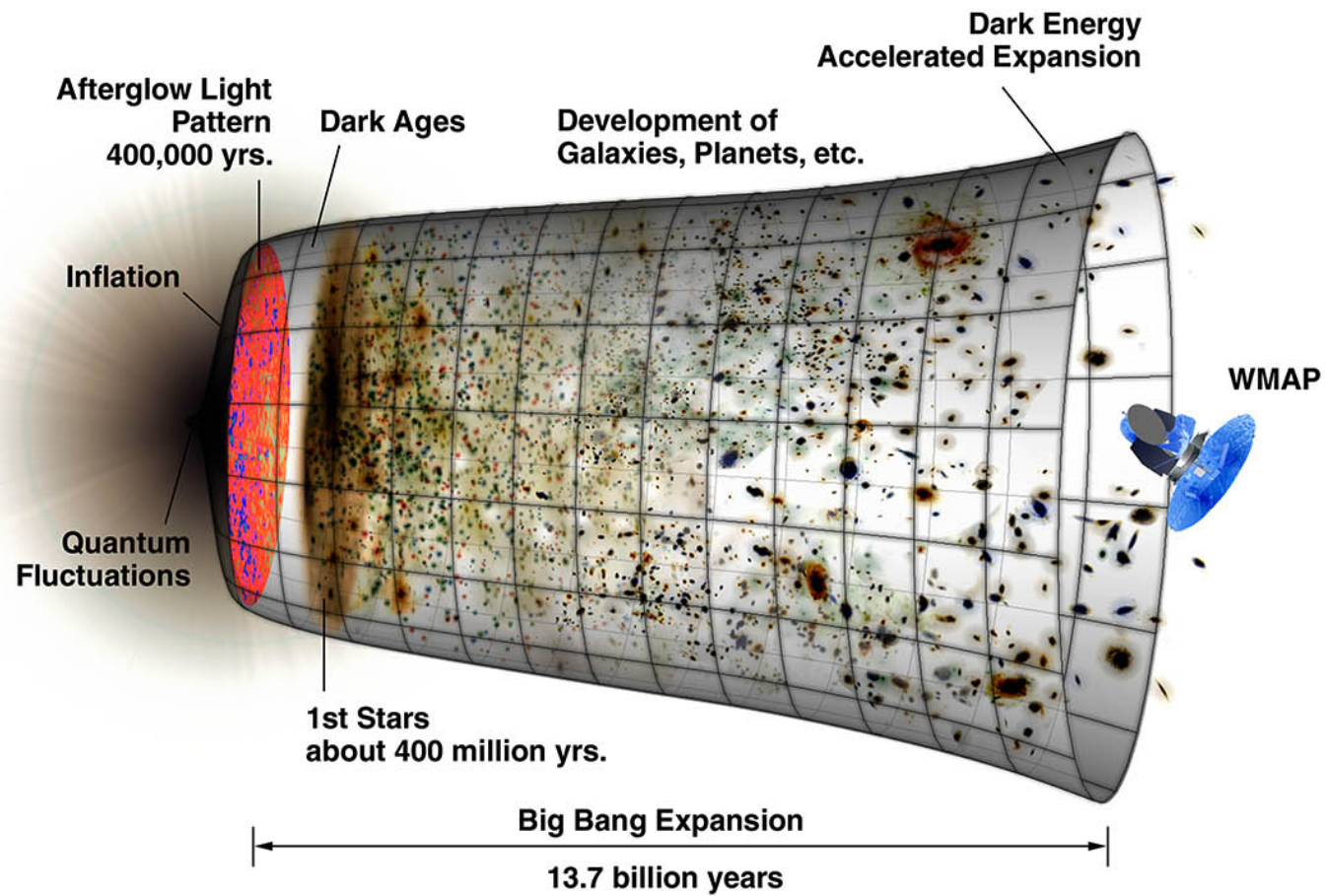
Our life is frittered away by detail.
Simplify, simplify, simplify!

-- Henry David Thoreau

3. Find x .



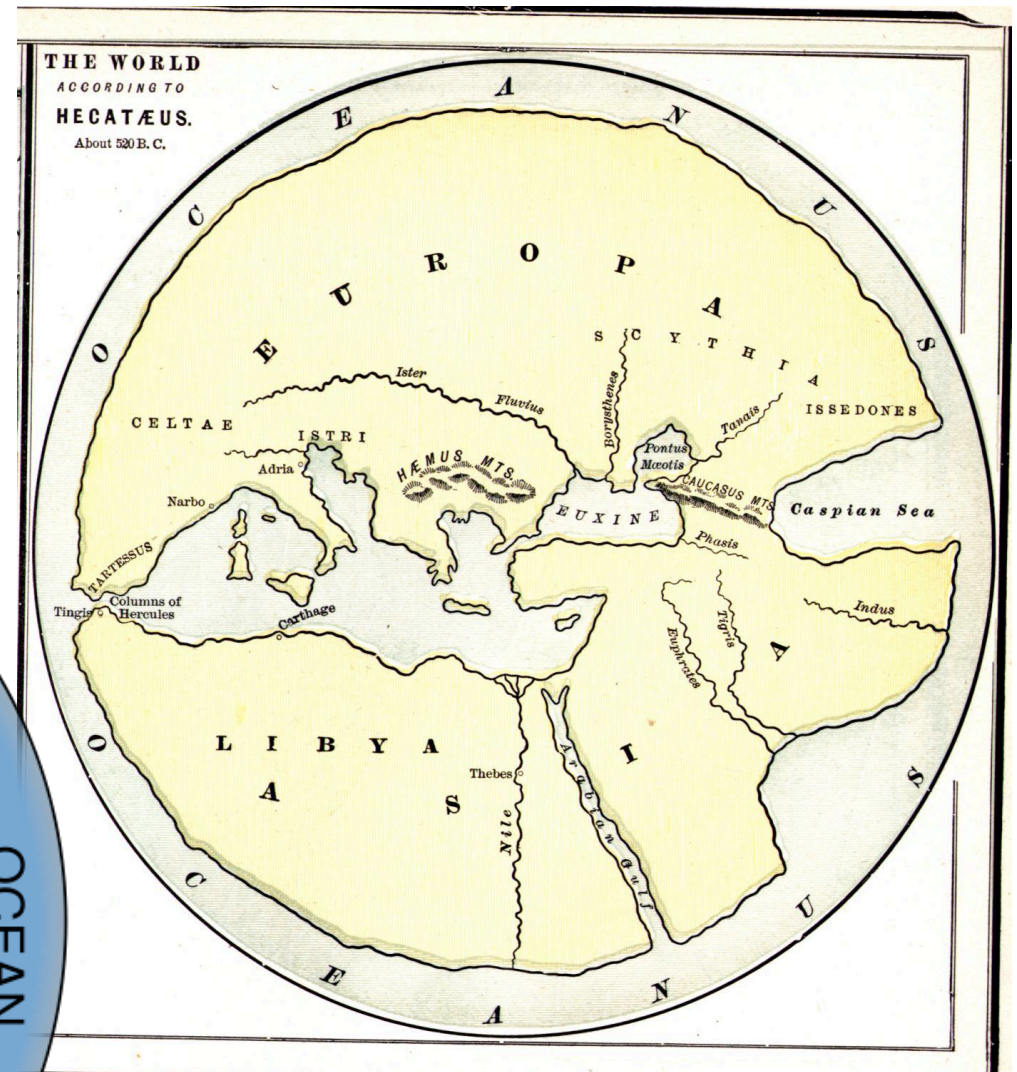
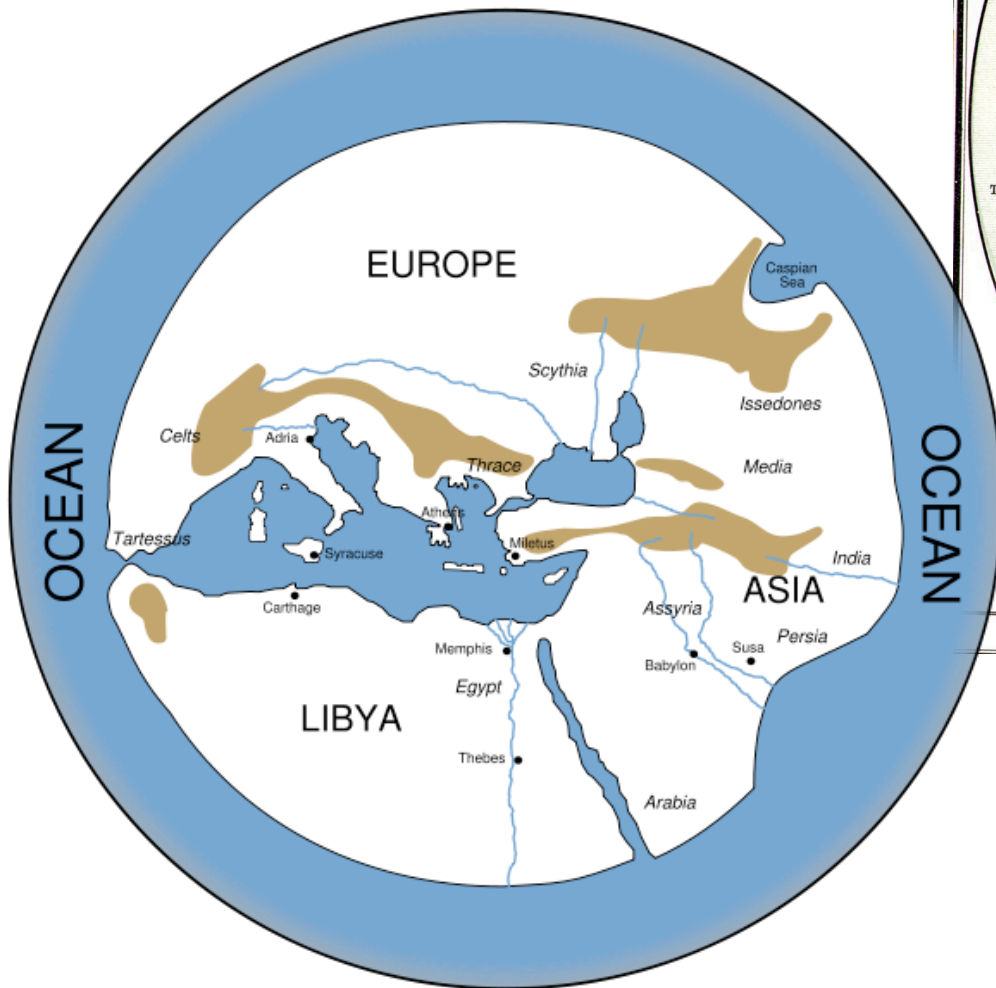
Here it is



Redshift, colours, spectra, fluxes  Dynamics, scale and content of the cosmos

NASA/WMAP Science Team

H_0 , Ω_{DE} , Ω_m , Ω_k etc...



How do we measure distances in everyday life?

“The real benefit is the fact that it’s easier for others to figure out how far away you are... I can say that, after putting two headlighthts on one of my bikes, I had much fewer near-incidence of people stepping out right in front of me”

www.bikeforums.net

A tale of two distances

- Two main distances in cosmology: **Luminosity** distance (d_L) and **angular-diameter** distance (d_A)

Photon Flux $\propto d_L^{-2}$

- They are related by **distance duality**:

$$d_L(z) = (1+z)^2 d_A(z)$$

This is true in any metric theory of gravity with photon conservation and any metric

What do distances tell us?

Let's assume the Universe is flat ($\Omega_k = 0$) to start with.
Then...

$$d_L(z) = (1 + z)$$



Distance we will
Measure somehow



Hubble Scale



Dynamics of the
Universe

What do distances tell us?

Let's assume the Universe is flat ($\Omega_k = 0$) to start with.
Taylor expanding to second order gives...

$$d_L(z) = \frac{c}{H_0} \left[z \right]$$



H_0 sets the overall *scale* for *all*
Distances (Hubble's law
is true for all models)

What do distances tell us?

Let's assume the Universe is flat ($\Omega_k = 0$) to start with.
Taylor expanding to second order gives...

$$d_L(z) = \frac{c}{H_0} \left[z + \frac{1}{4}(1 + 3\Omega_\Lambda)z^2 \dots \right]$$



H_0 sets the overall *scale* for *all* Distances (Hubble's law is true for all models)

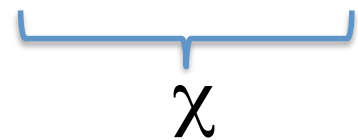


The *shape* of the Hubble Diagram depends on Ω_Λ

What do distances tell us?

What happens when we don't assume flatness?

$$d_L(z) = (1+z) \frac{c}{H_0} \sin \left(\sqrt{-\Omega_K} \int \frac{dz'}{E(z')} \right)$$



Encodes the effect of the curved null geodesics ...

What do distances tell us?

What happens when we don't assume flatness?

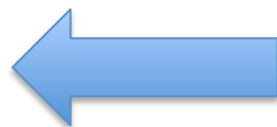
Taylor expanding to second order gives...

$$d_L(z) = \frac{c}{H_0} \left[z + \frac{1}{4} (1 + \Omega_K + 3\Omega_\Lambda) z^2 \dots \right]$$

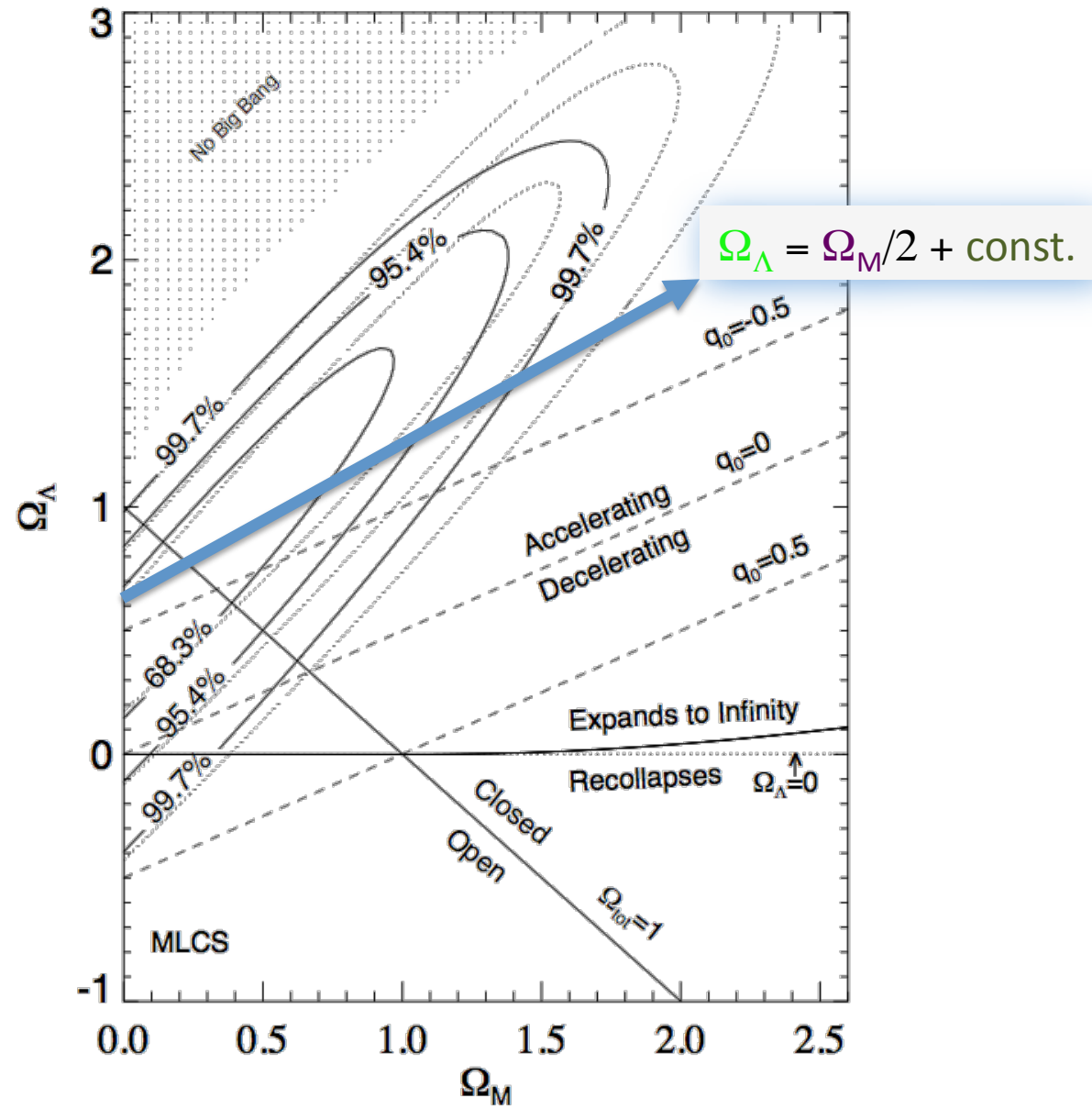


2nd order term is degenerate on the line:

$$\Omega_\Lambda = \Omega_M/2 + \text{const.}$$



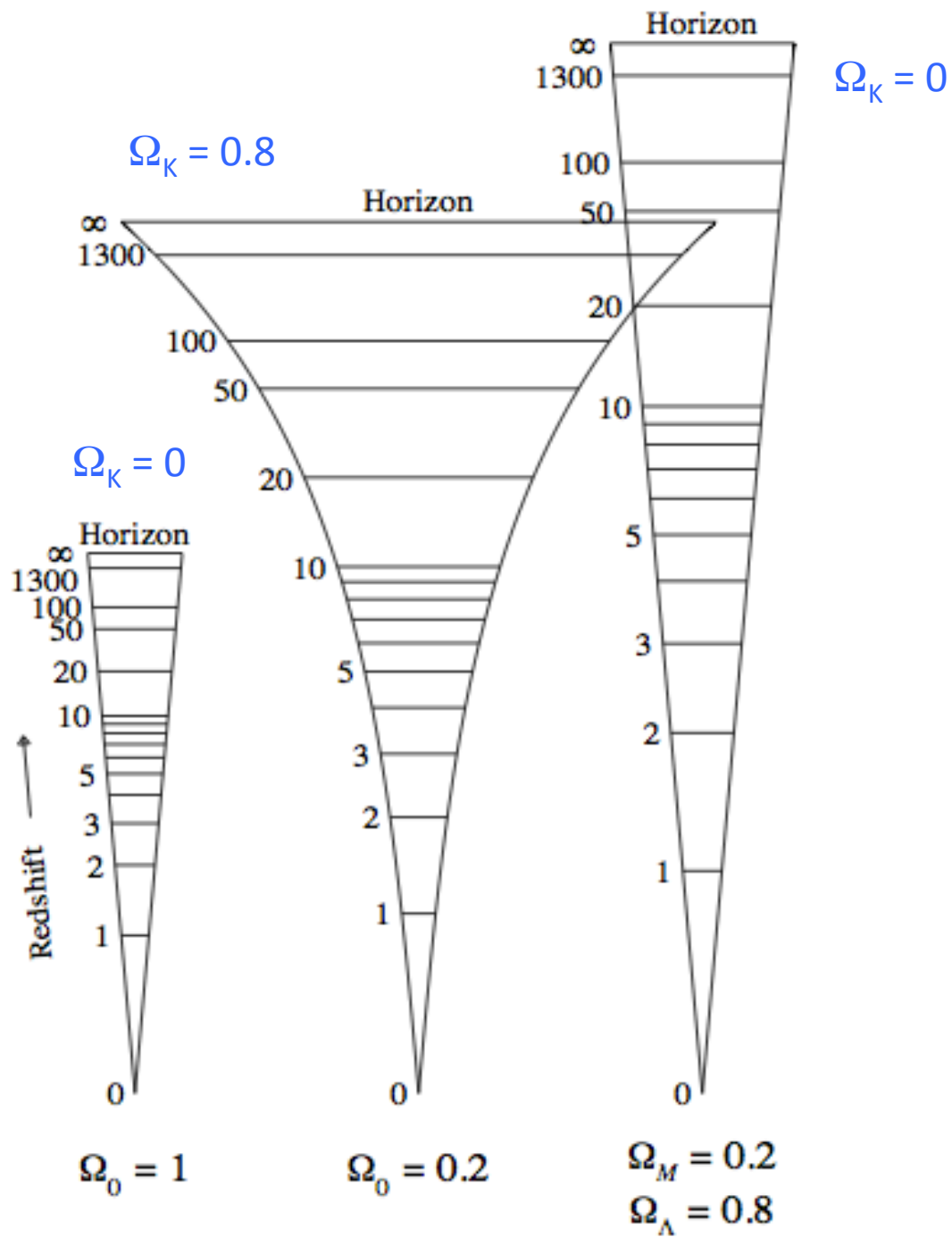
The *curvature* of the Hubble Diagram now depends on the Amount of dark energy **AND** the spatial curvature, Ω_K



Riess et al, 1998

$$\underbrace{\frac{c}{H_0} \int \frac{dz'}{E(z')}}_{\chi}$$

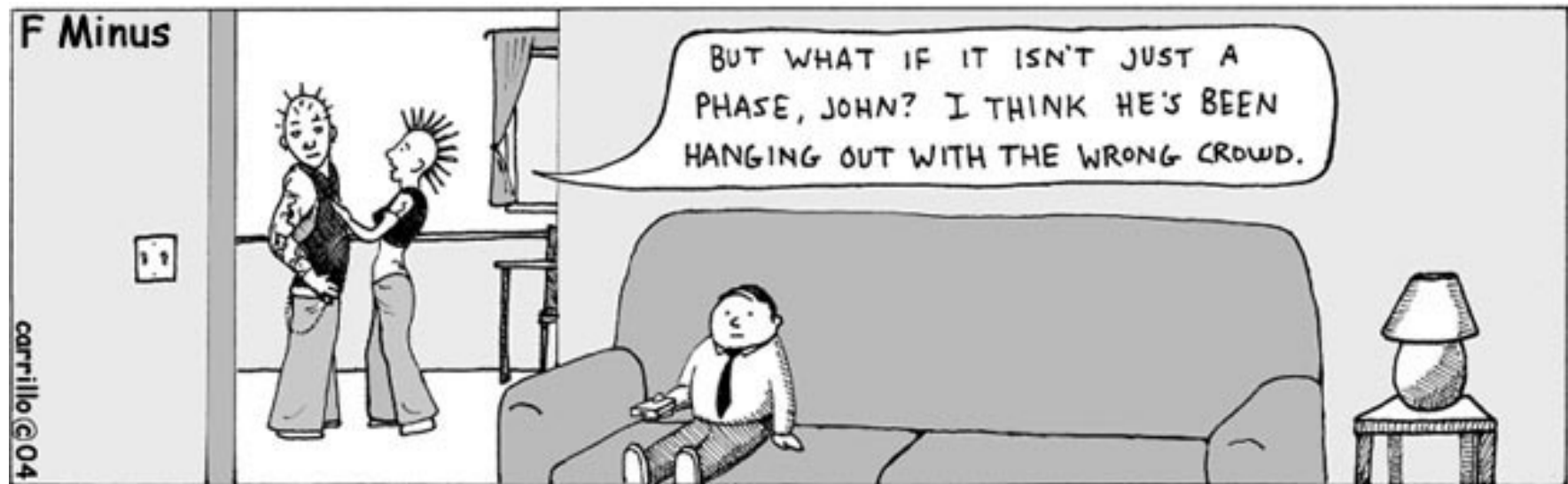
Comoving distance





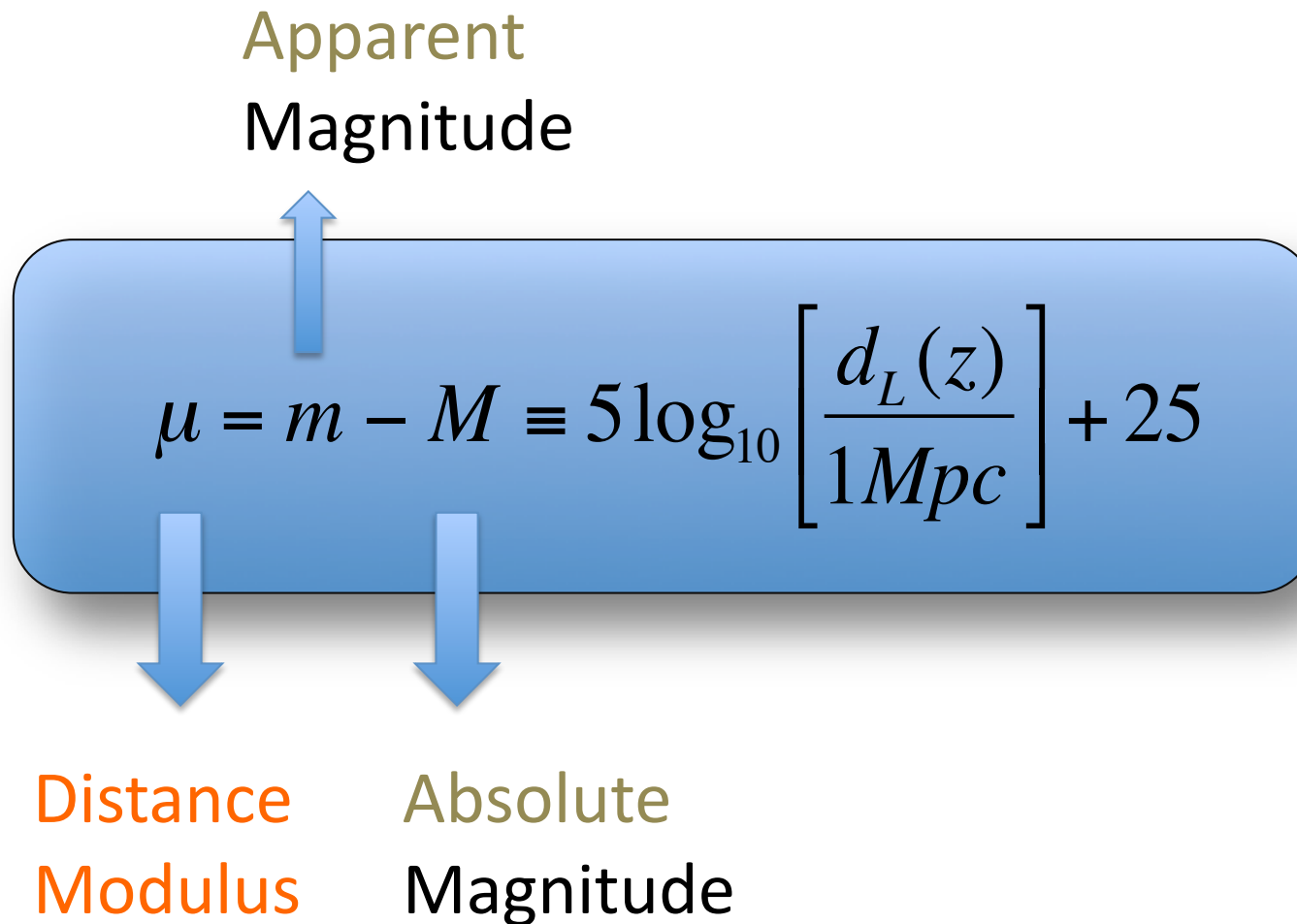
So how do we measure a luminosity distance?

- Or “why cosmologists care about stars”



Distance modulus and magnitudes

Apparent
Magnitude


$$\mu = m - M \equiv 5 \log_{10} \left[\frac{d_L(z)}{1 \text{ Mpc}} \right] + 25$$

Distance
Modulus Absolute
Magnitude

d_L converts $M \rightarrow m$

Common in Supernova cosmology to plot μ or m_B vs redshift

d_L Wish list...

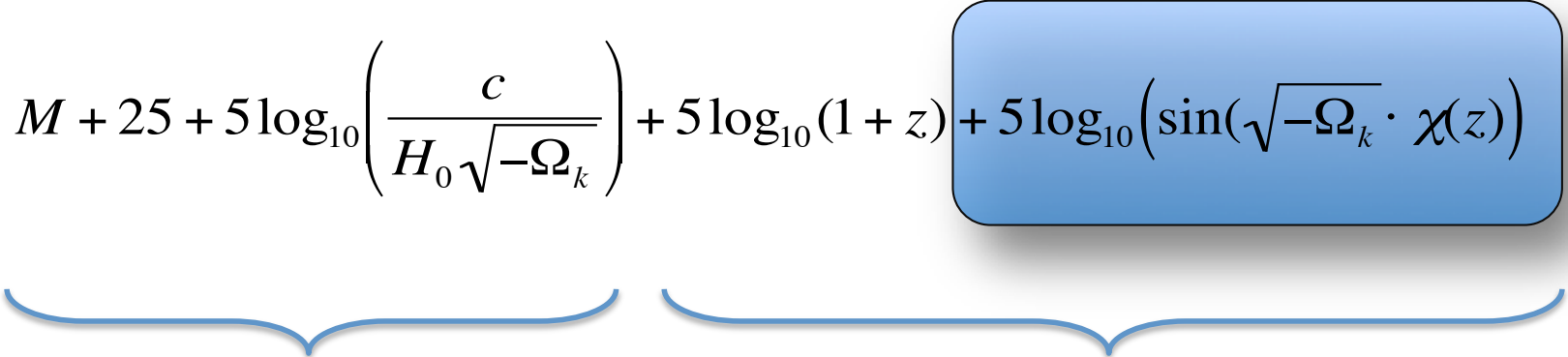
We want objects that are:

- Very Bright...so we can see them across the observable universe
- Known: We know each object's intrinsic luminosity.
- Standard: no environmental dependence of the intrinsic luminosity
- Time-invariant: their intrinsic luminosity doesn't change with redshift
- Easy to find...
- Emit most of their light in the optical
(counter example: binary black holes as GW sirens)

Problem: if we don't know M

What can we do?

$$m = M + 25 + 5 \log_{10} d_L(z)$$
$$= M + 25 + 5 \log_{10} \left(\frac{c}{H_0 \sqrt{-\Omega_k}} \right) + 5 \log_{10} (1+z) + 5 \log_{10} \left(\sin(\sqrt{-\Omega_k} \cdot \chi(z)) \right)$$

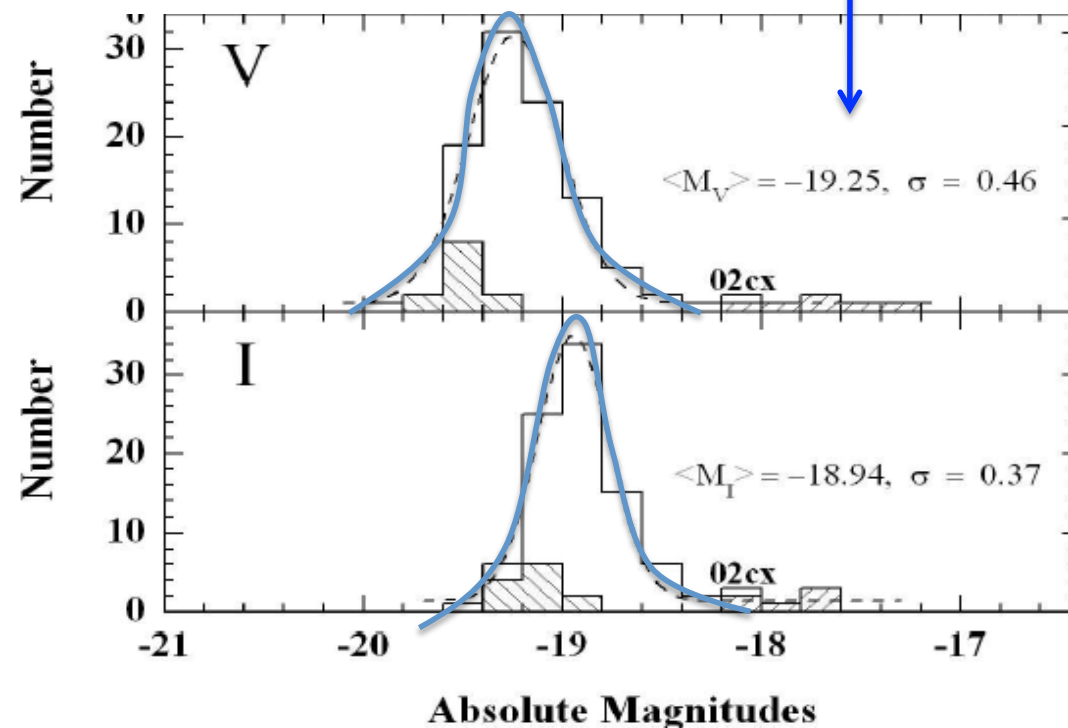


Constants Redshift dependent

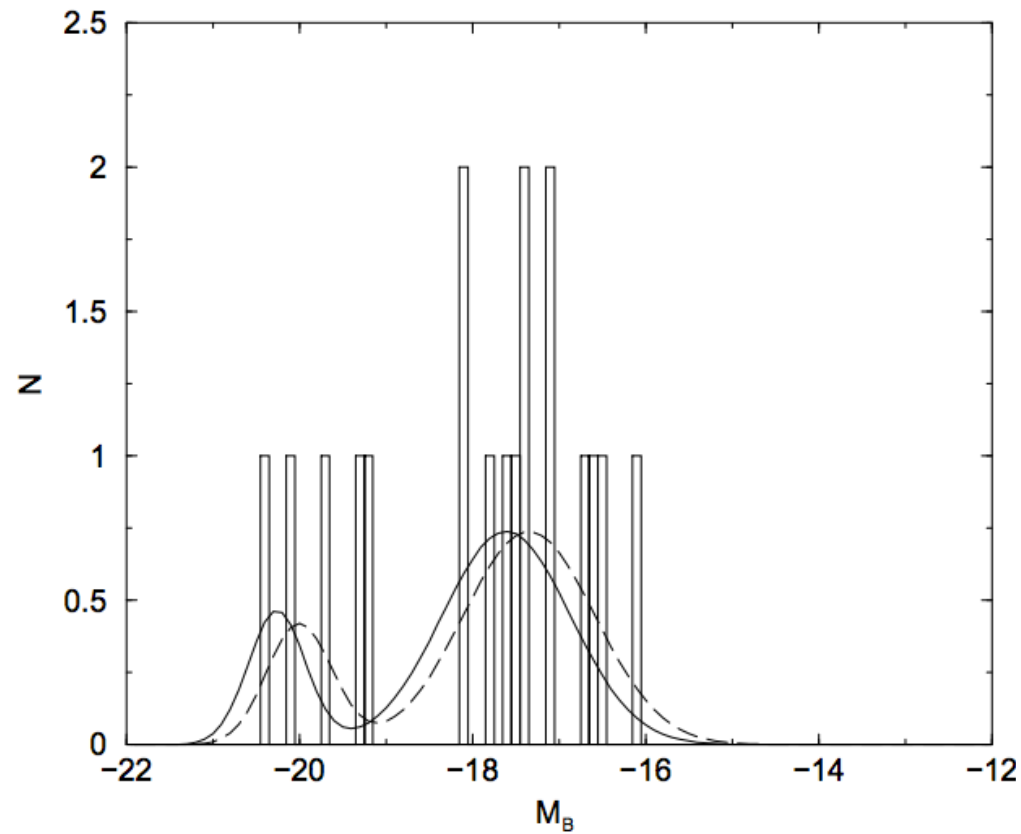
If we subtract the apparent magnitudes of two objects with the same M at different redshifts, all the unknown constants disappear...

Type Ia Supernovae (SNIa)

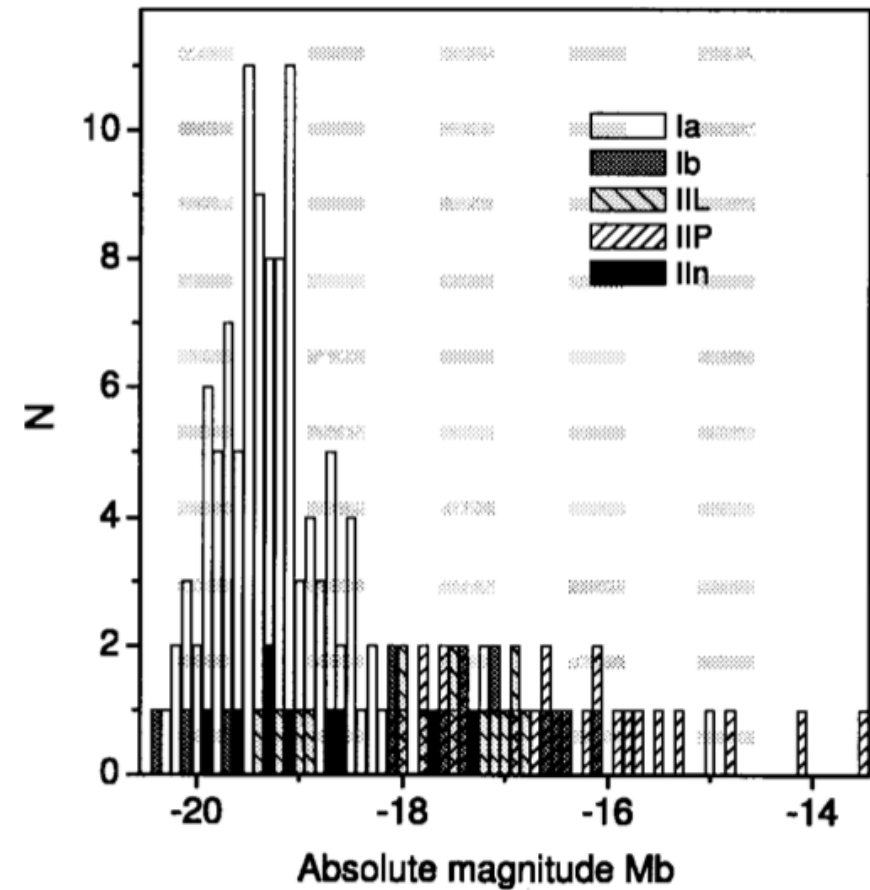
- **Very Bright**...at peak they can outshine their host galaxies
- **Stable**: Don't seem to evolve with redshift
- **Fairly standard**:
 $\sigma = 0.4$ mag
- **Easy to find**...
 ~ 1 per century per galaxy



Non-Ia diversity



Richardson et al, 2001



Sveshnikova, 2003

Total incineration of a white dwarf when it exceeds the Chandrasekhar limit of $1.4 M_{\text{sun}}$

 **Bright *and* standard**



www.astroart.org

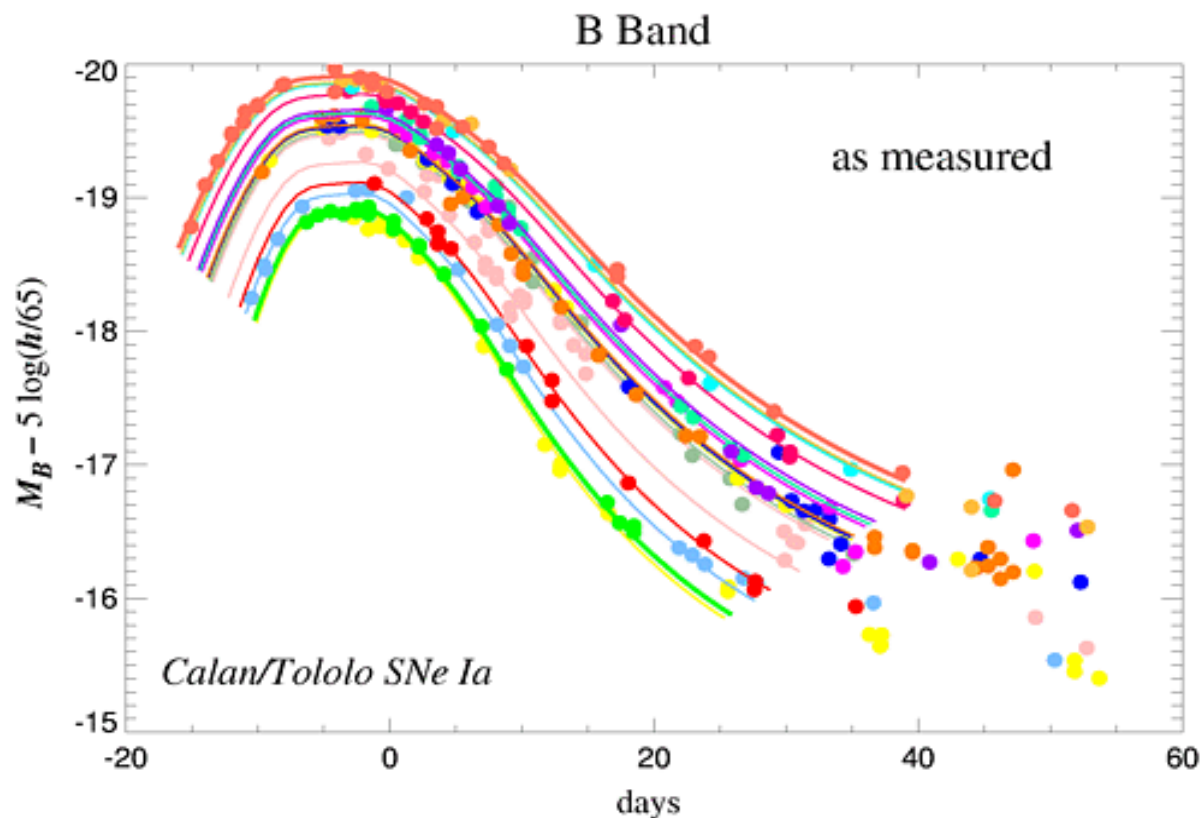
Artist's rendition of a white dwarf accumulating mass from a nearby

But wait, there's more...

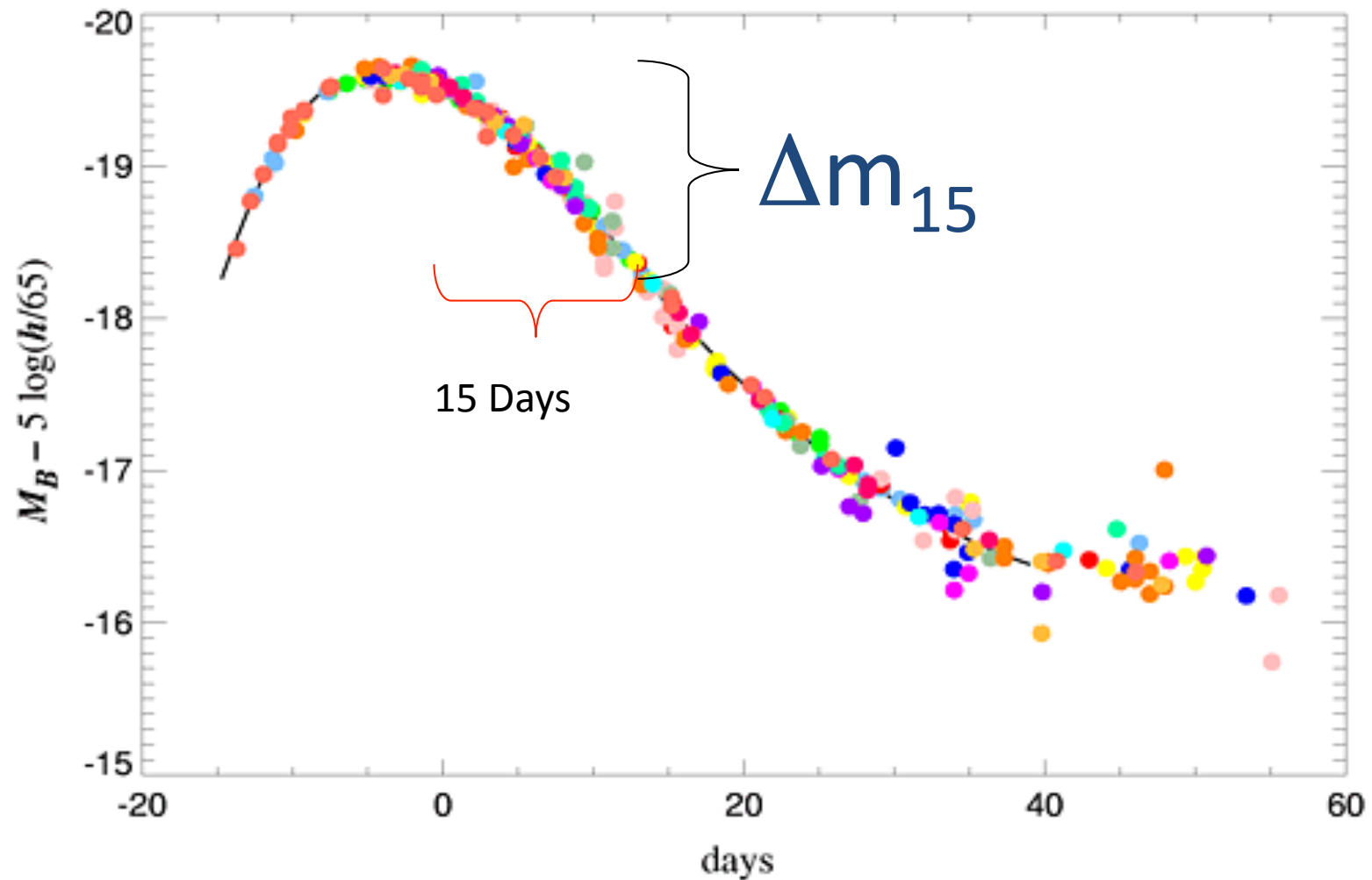
- We can make SNIa even more standard reducing variation in peak brightness to

$$\sigma = 0.12 \text{ mag}$$

- Exploit observed correlation between peak brightness, colours and shape of the lightcurve (Phillips, 1993).



After Standardisation...



Kim, *et al.* (1997)

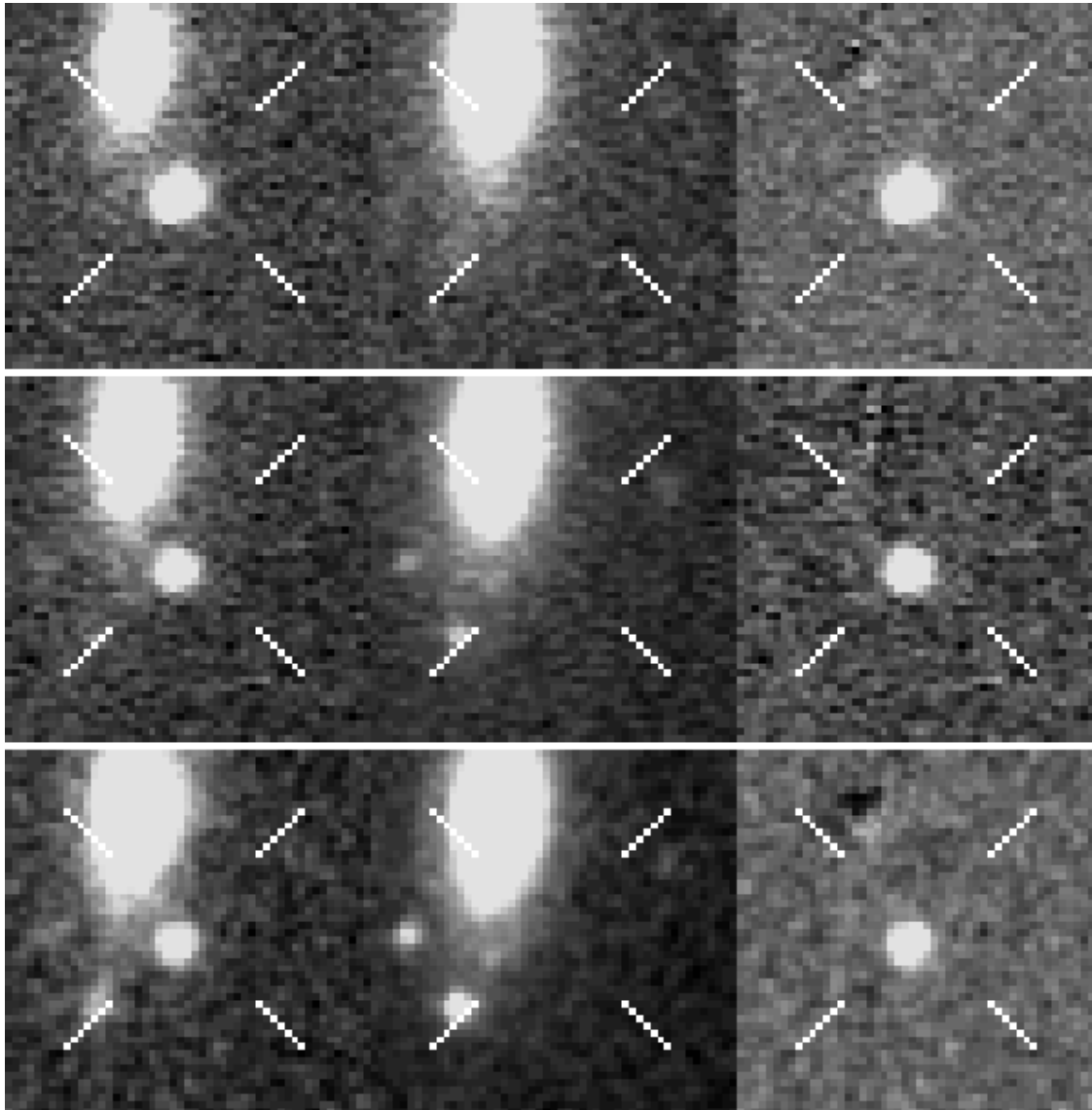
Why does
intrinsically **brighter** → **slower** decline?

How are SNe found these days?

NEW Image

OLD Image

Difference



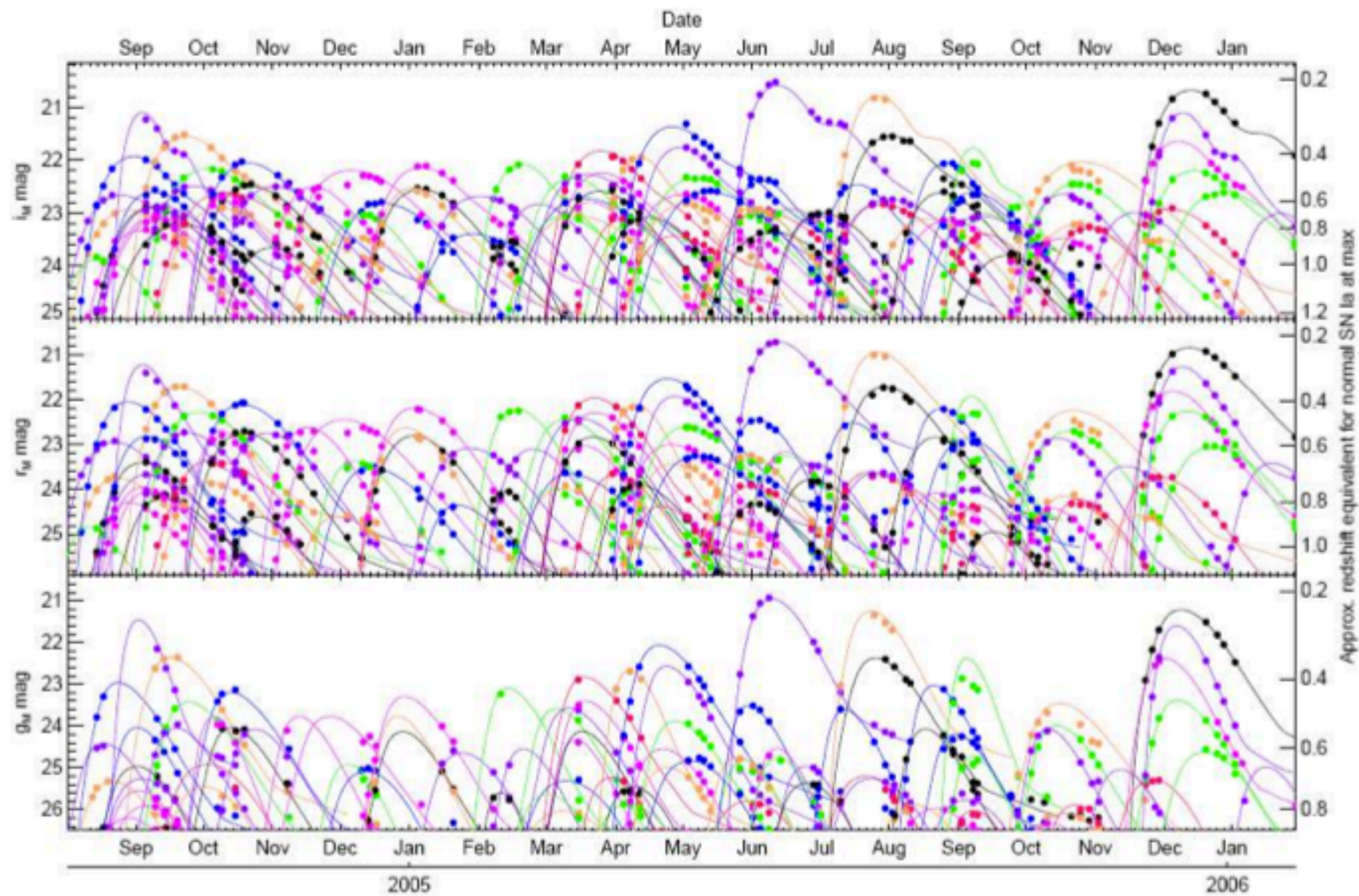
g

r

i



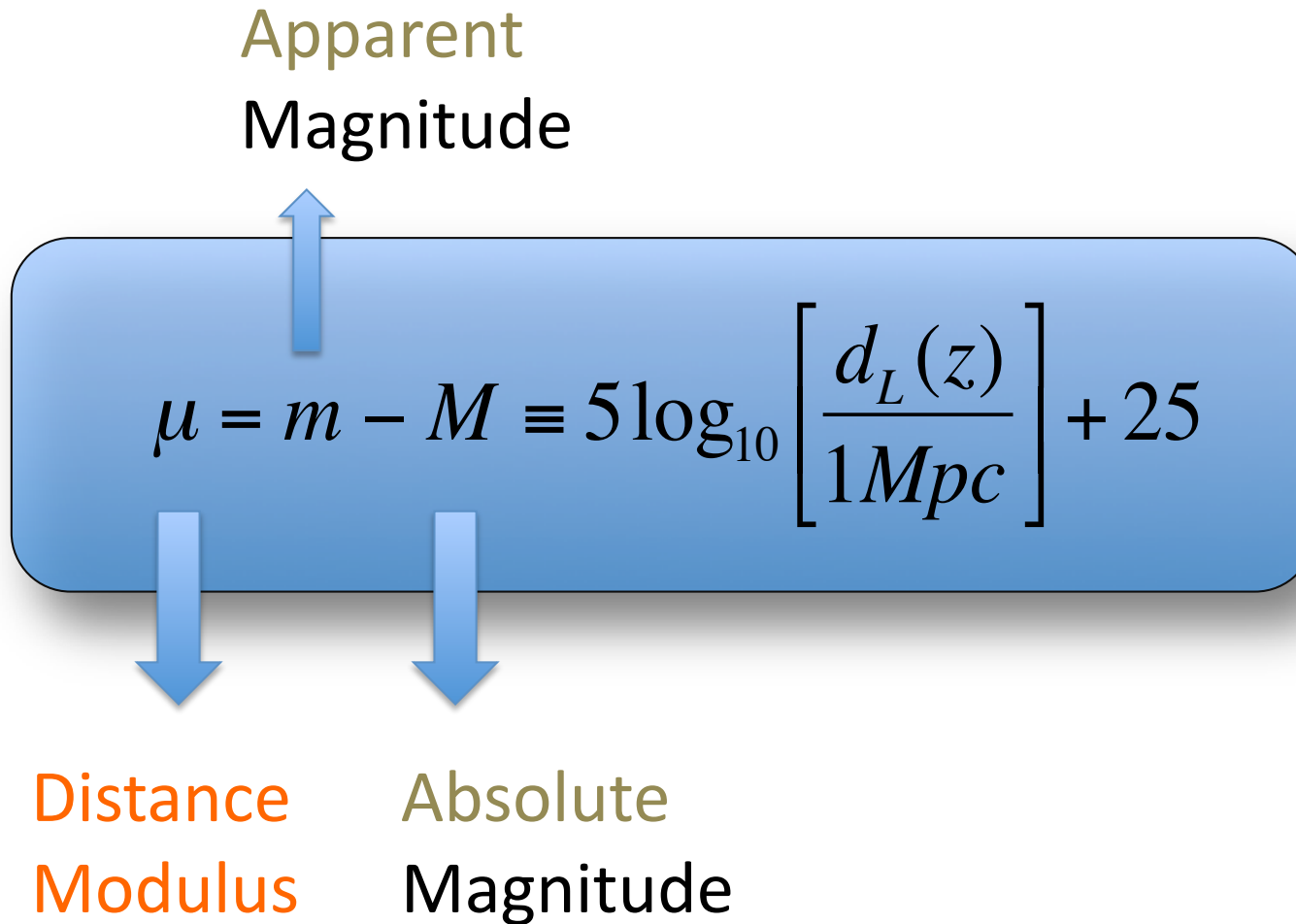
SDSS SN search



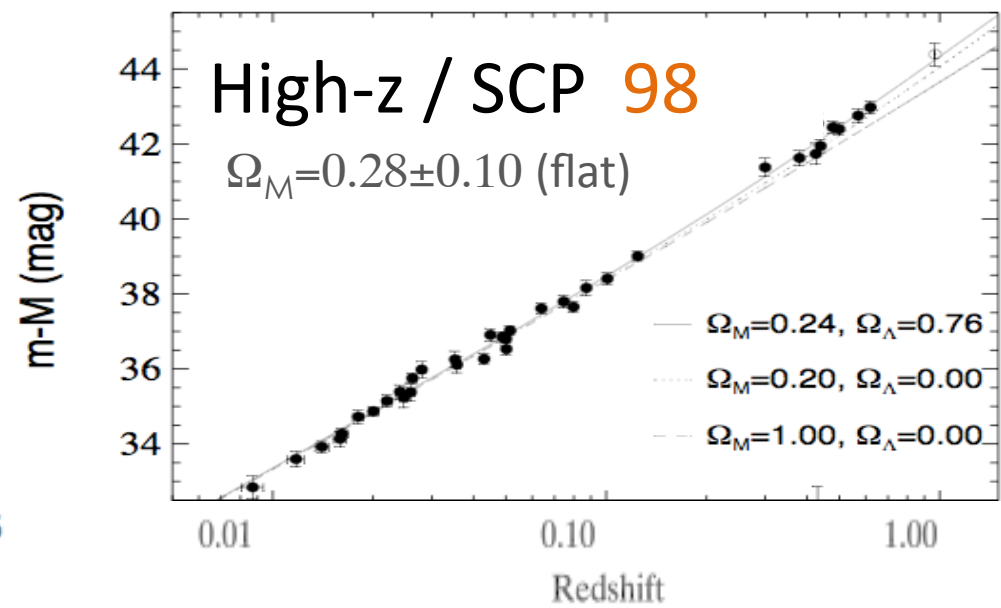
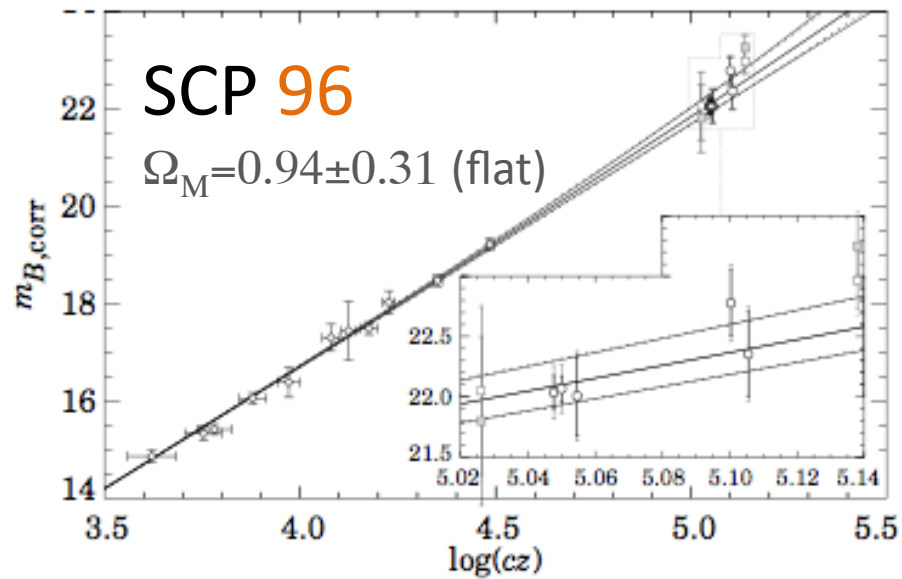
SNLS Rolling Search

Distance modulus and magnitudes

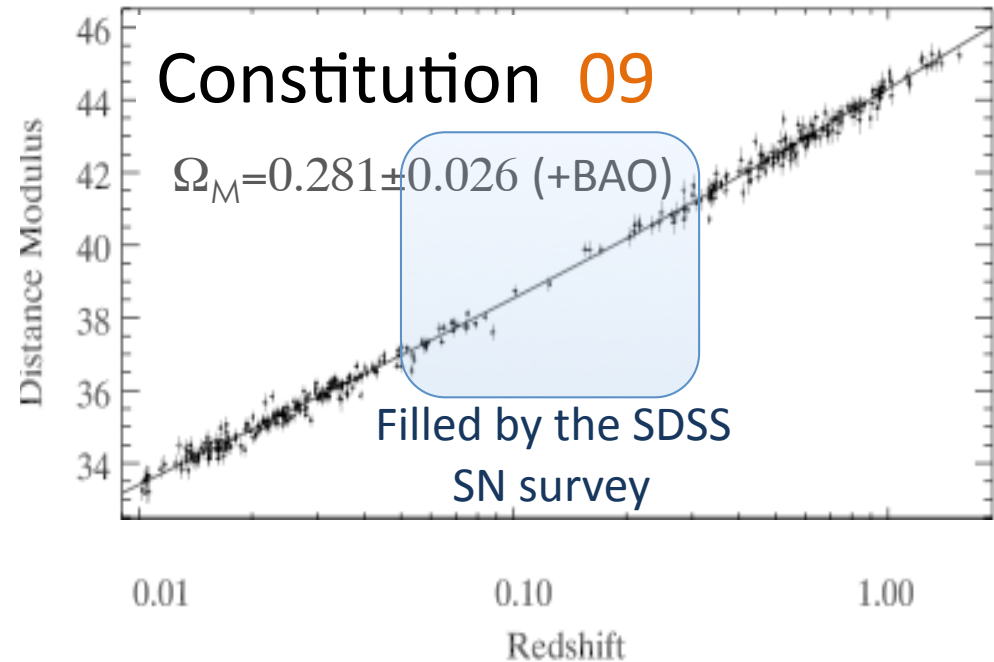
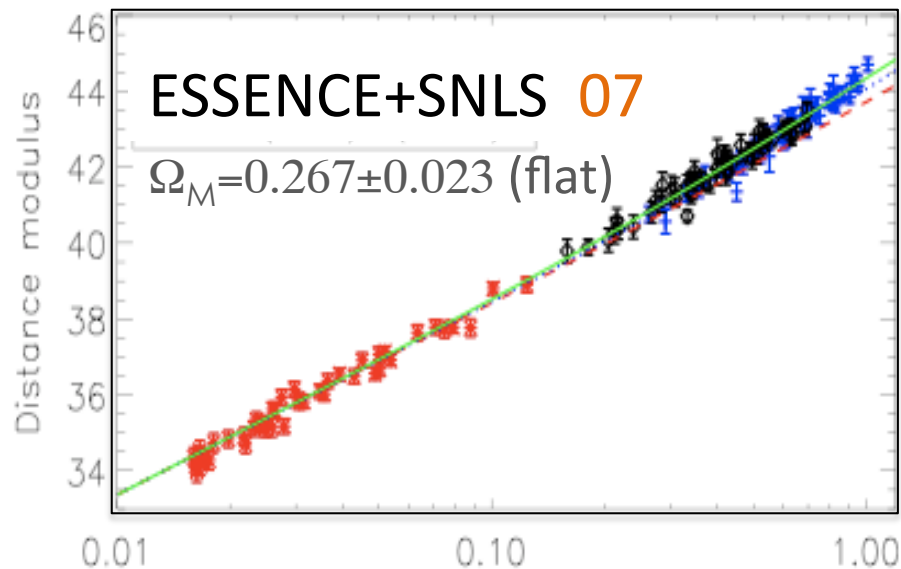
Apparent
Magnitude


$$\mu = m - M \equiv 5 \log_{10} \left[\frac{d_L(z)}{1 \text{ Mpc}} \right] + 25$$

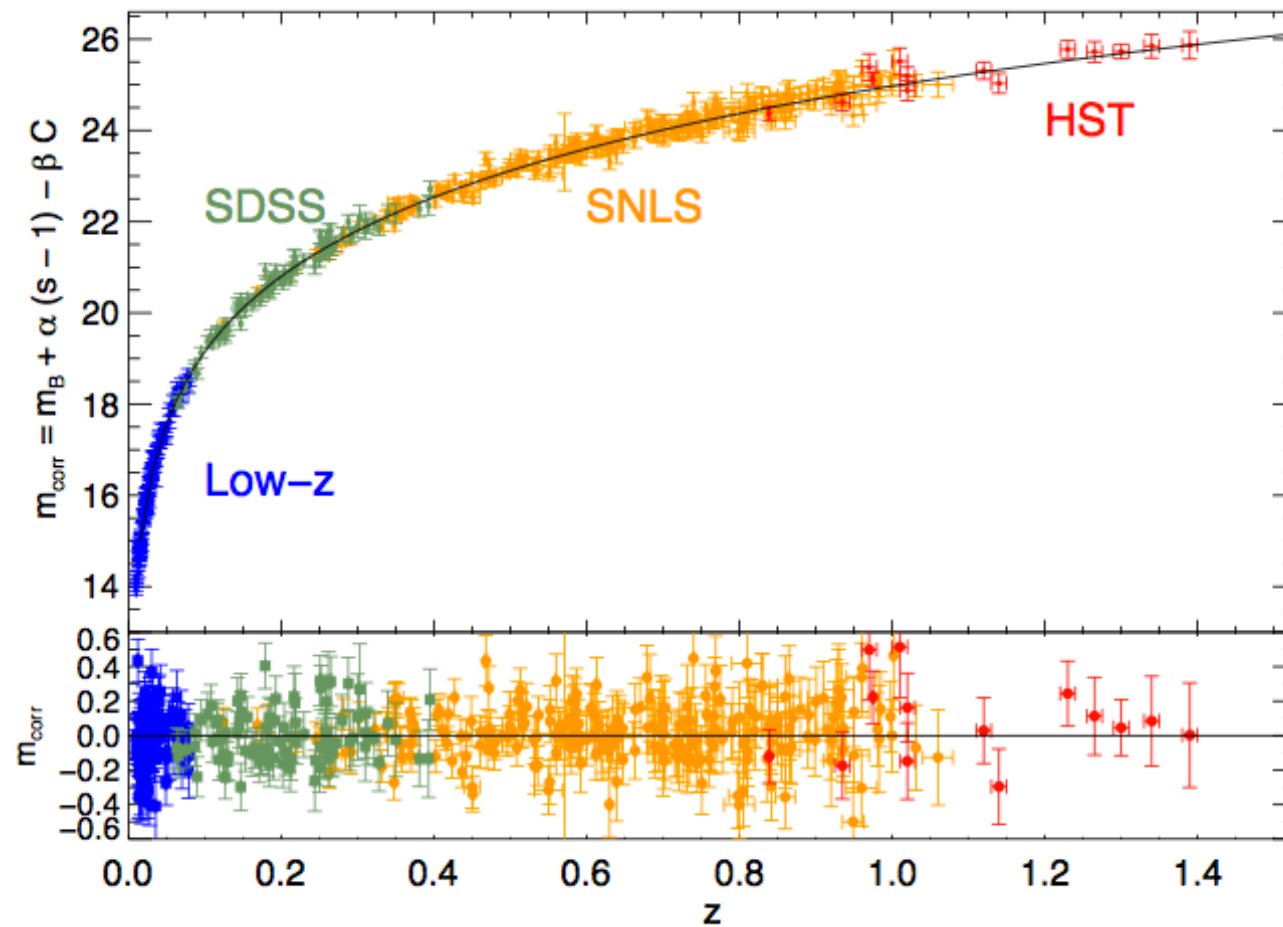
Distance
Modulus Absolute
Magnitude



Evolution of the Hubble Diagram 1996-2009



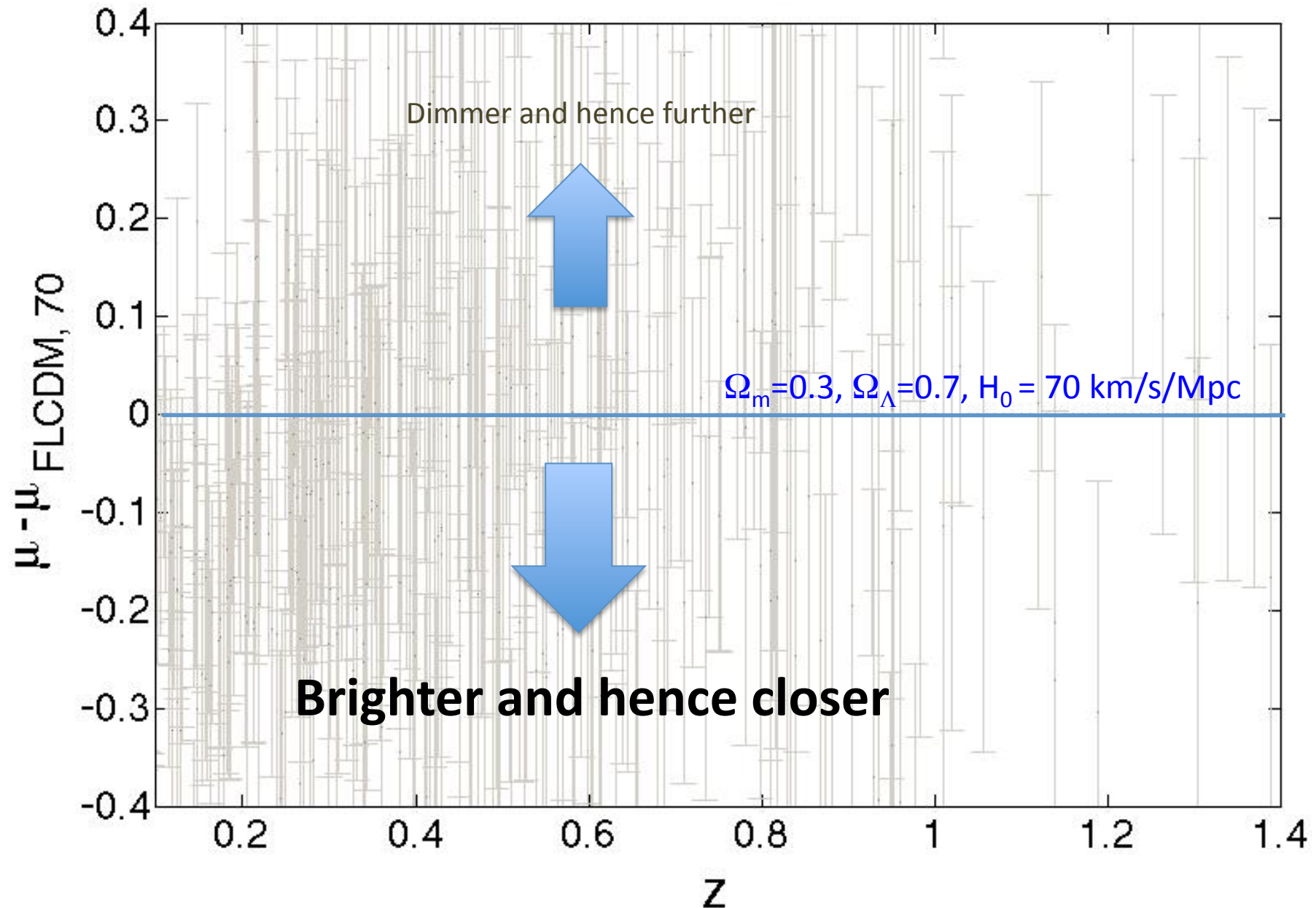
Current state-of-the art

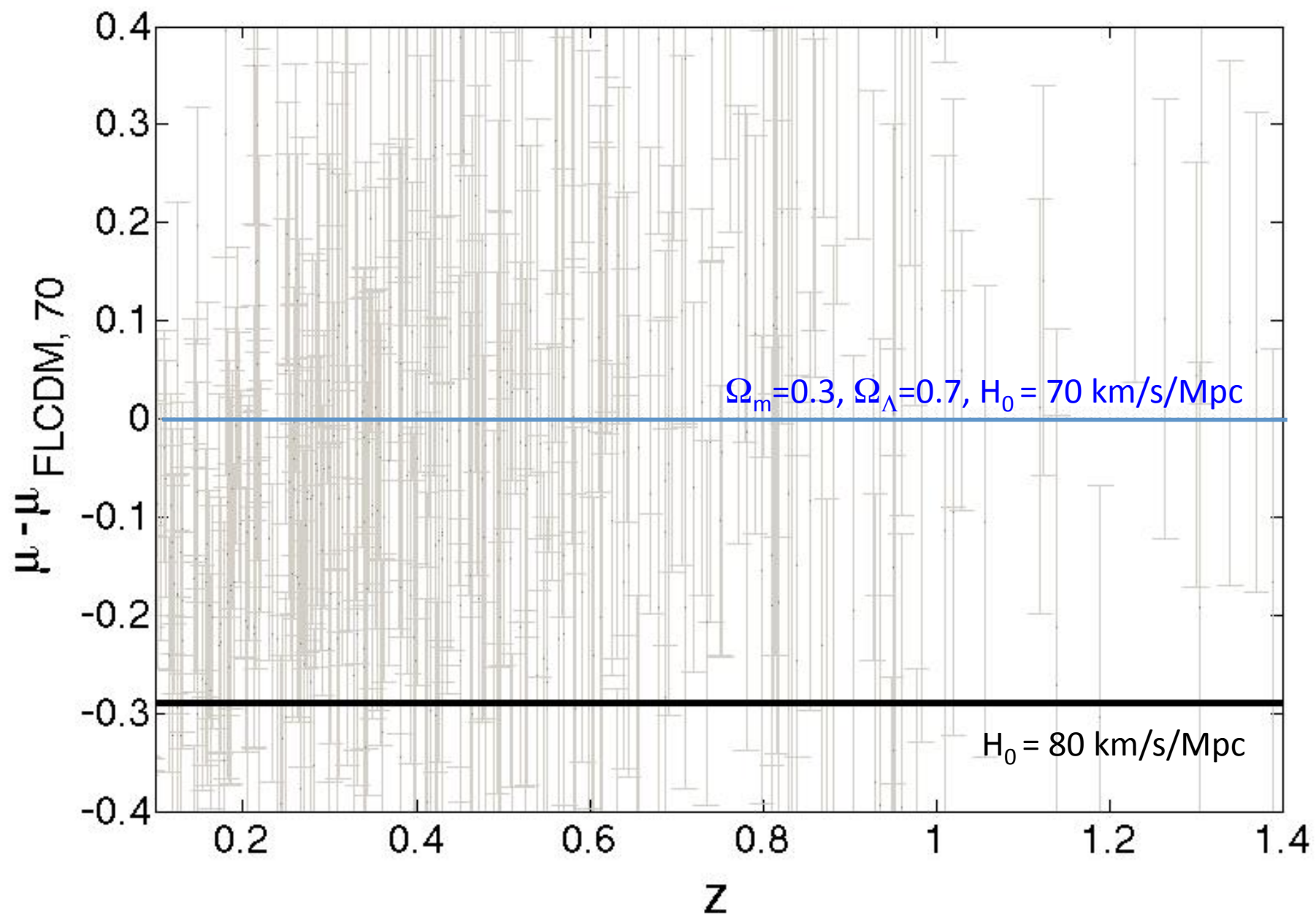


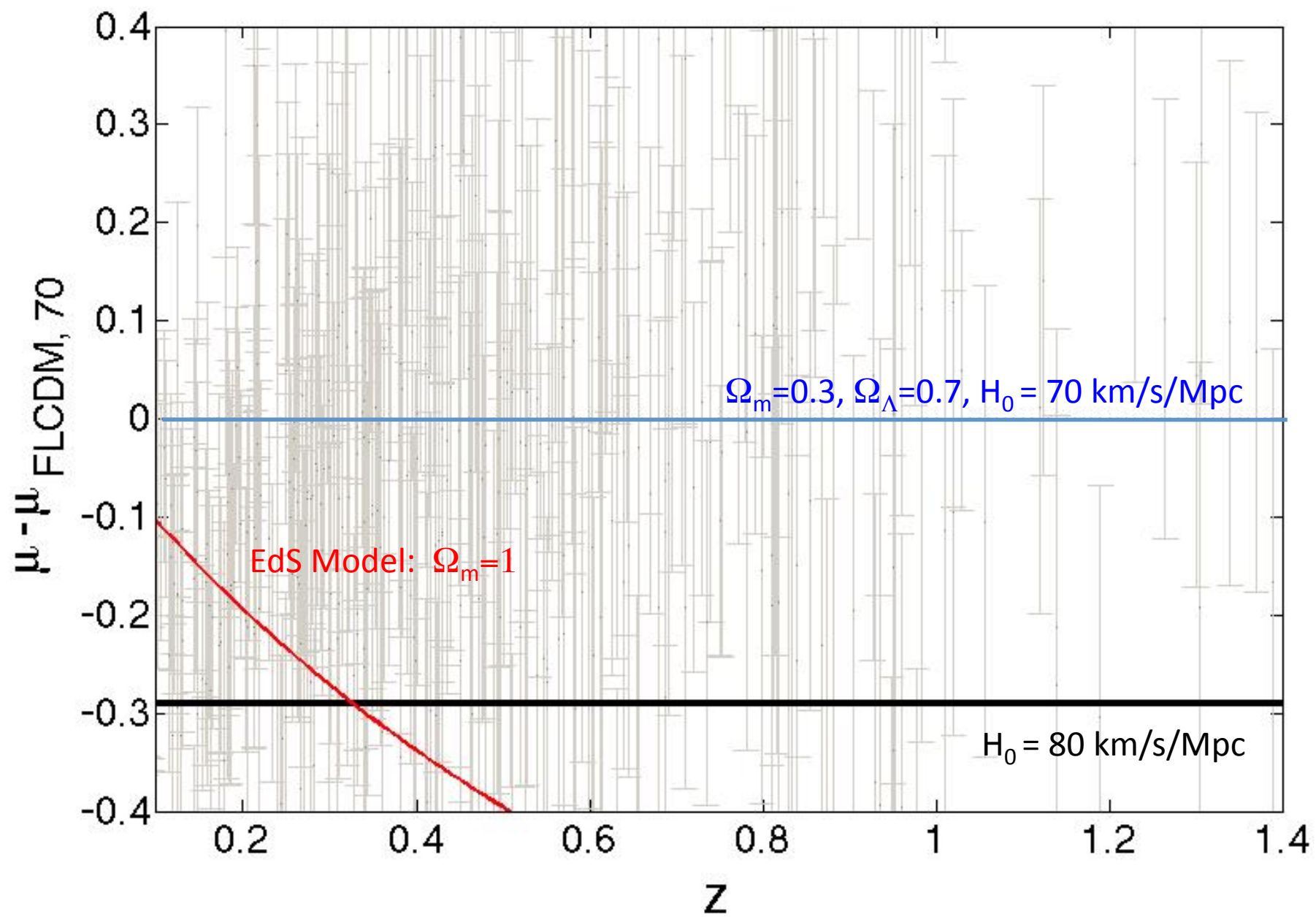
So what can we learn?

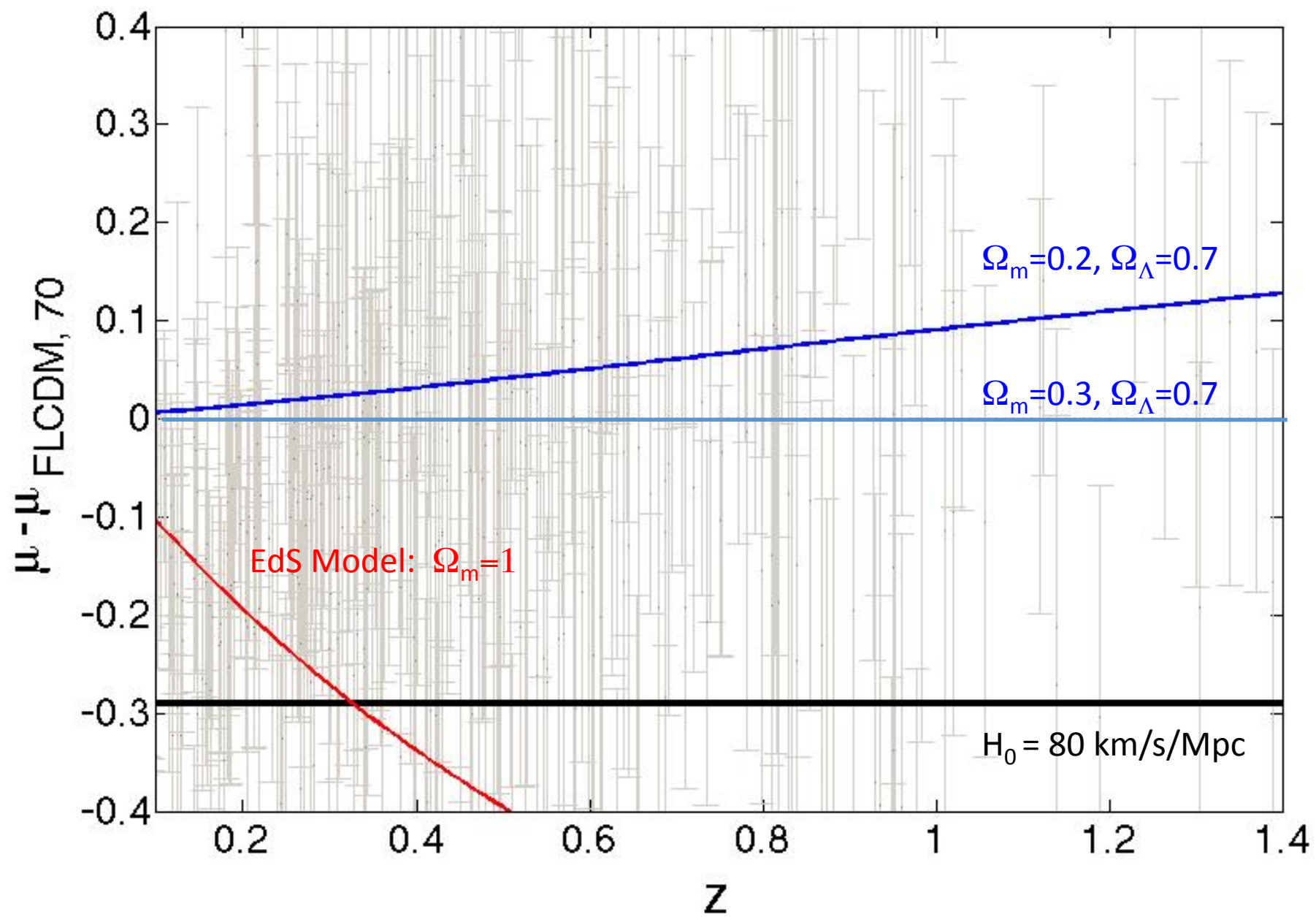


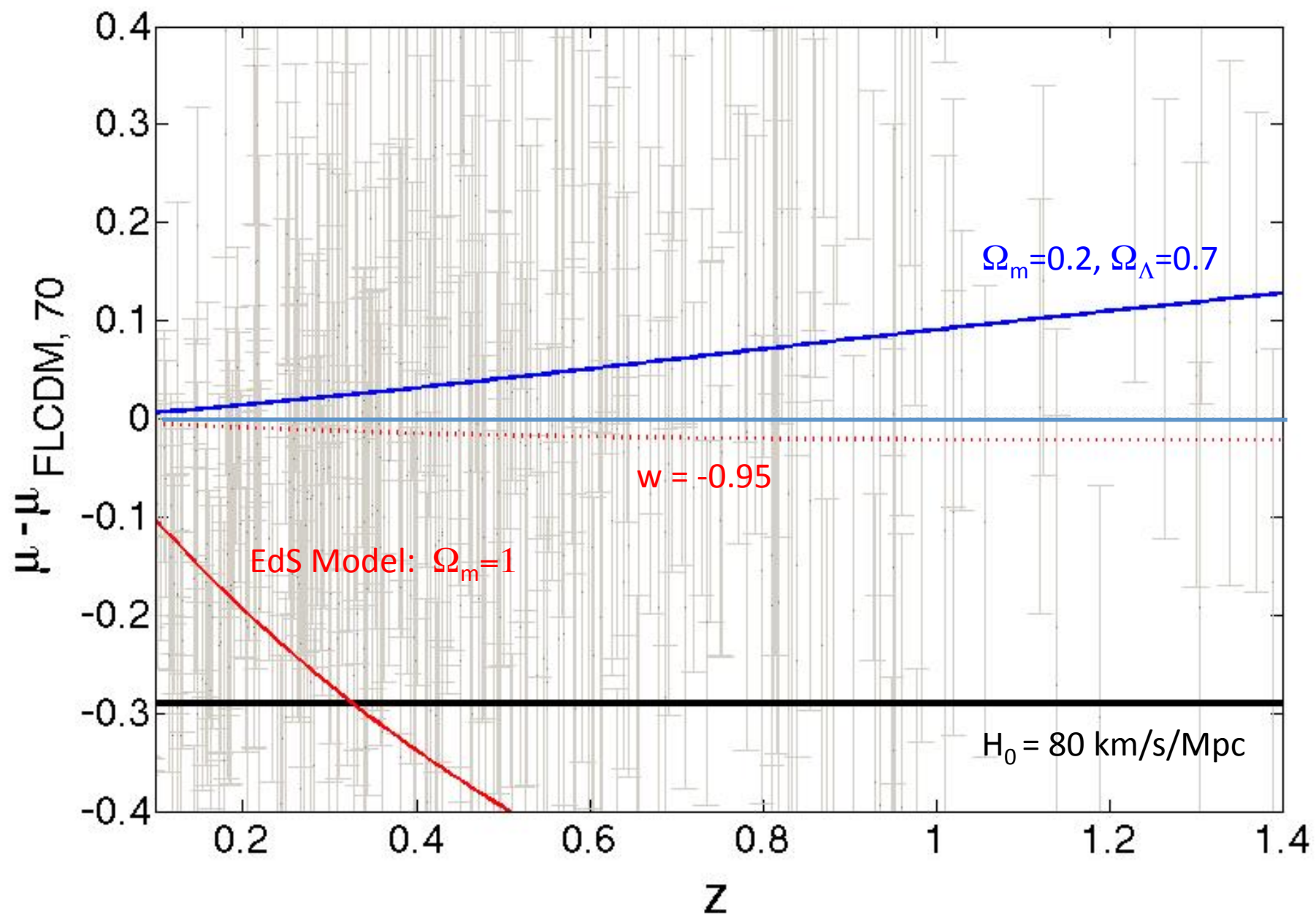
Relative Hubble Diagram

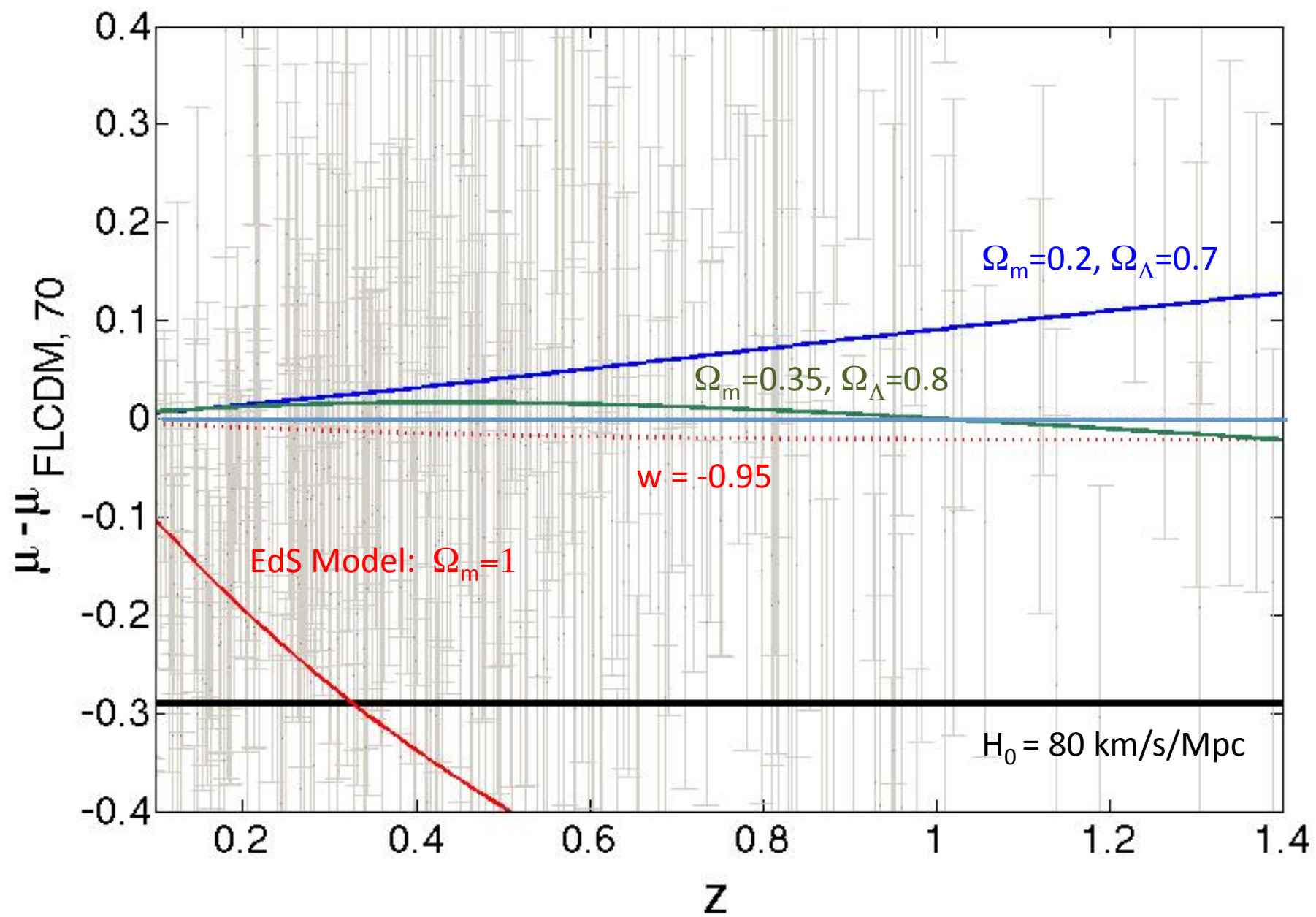


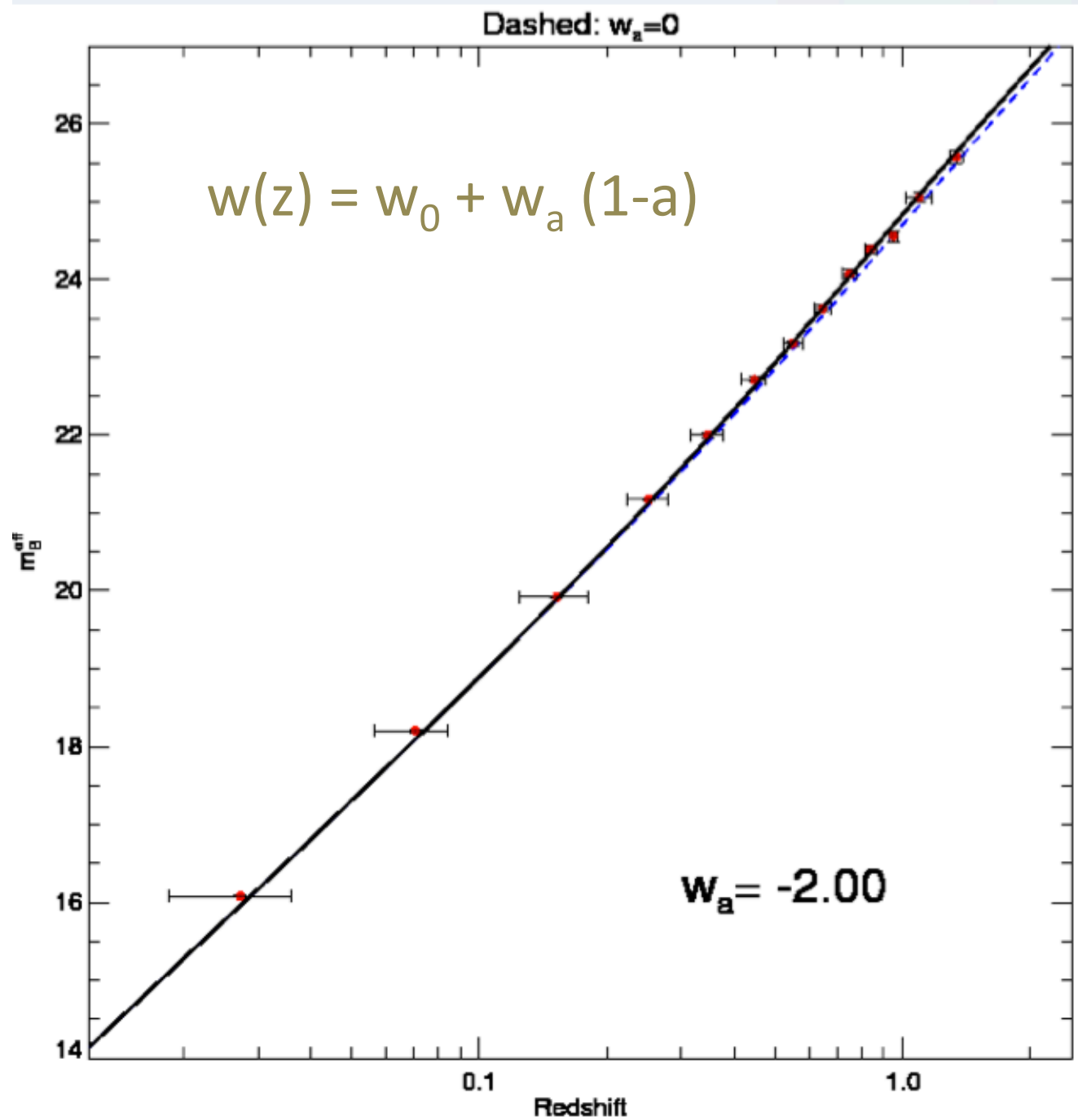






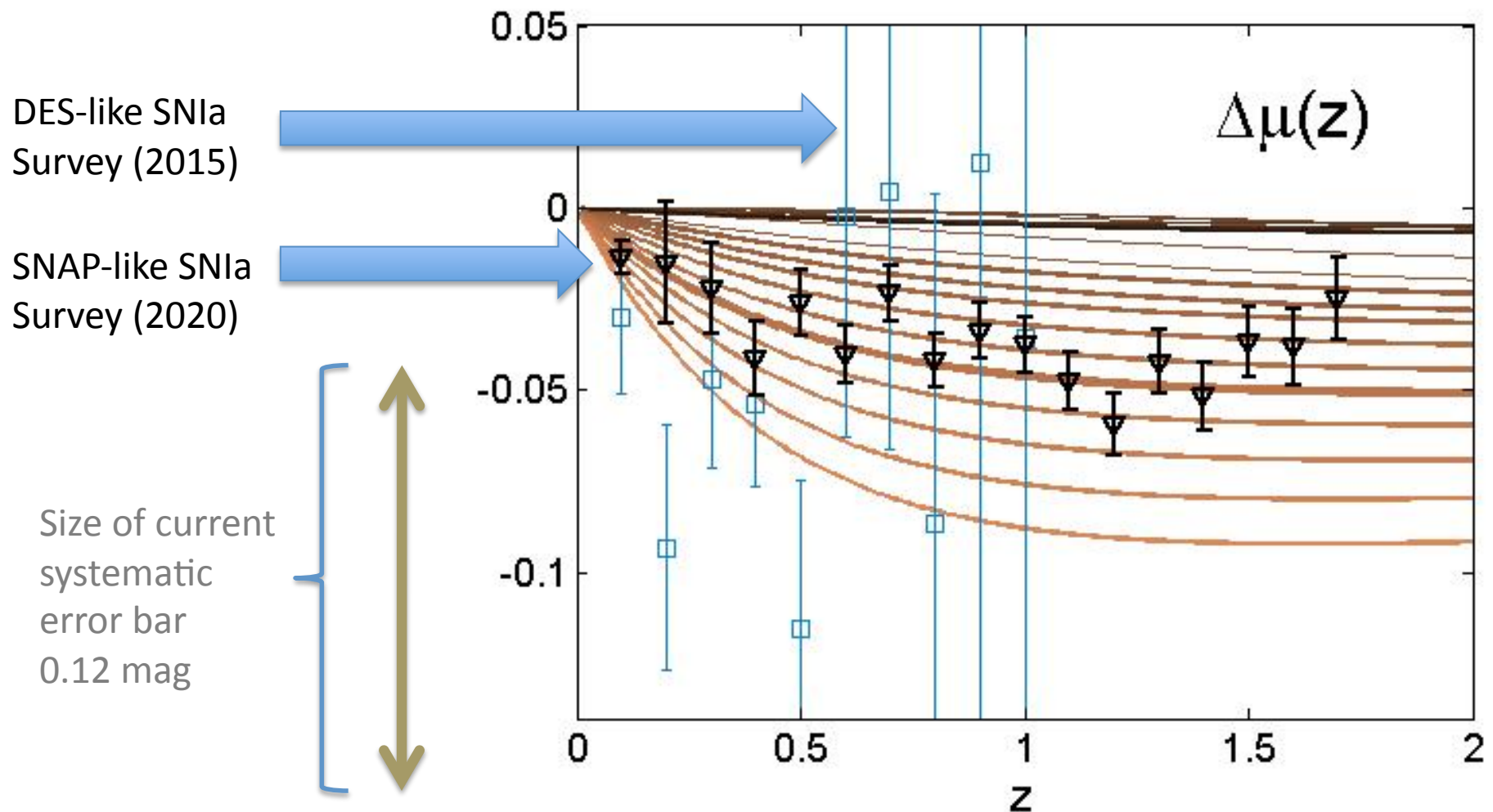






Goobar

It is very hard to distinguish tracking quintessence and Λ CDM...



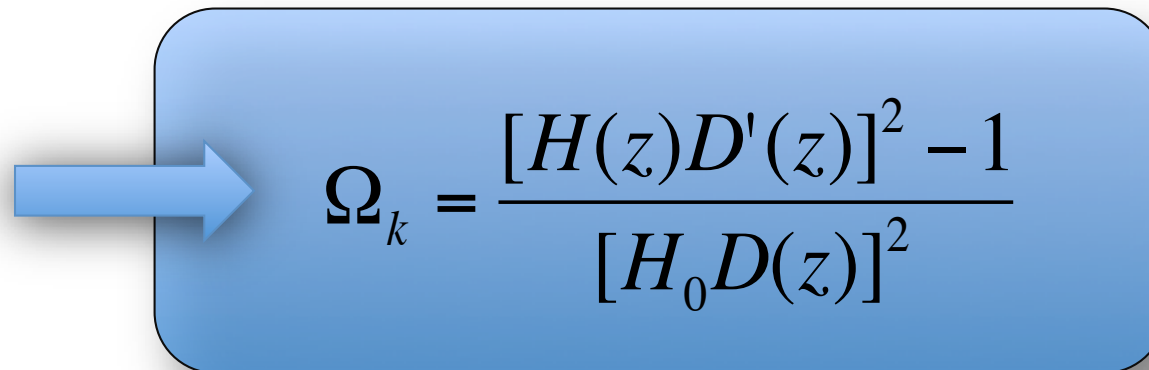
Other difficulties...


- Degeneracies...
- Gravitational lensing...
- Dust obscuration...
- Peculiar objects...
- Host dependence...
- UV evolution...
- The Spectroscopic desert
- The future...

Geometry – Dynamics Degeneracy

- Even with perfect **distance** measurements there is a perfect degeneracy between **the curvature** (Ω_k) and **$H(z)$** (Weinberg, '73)

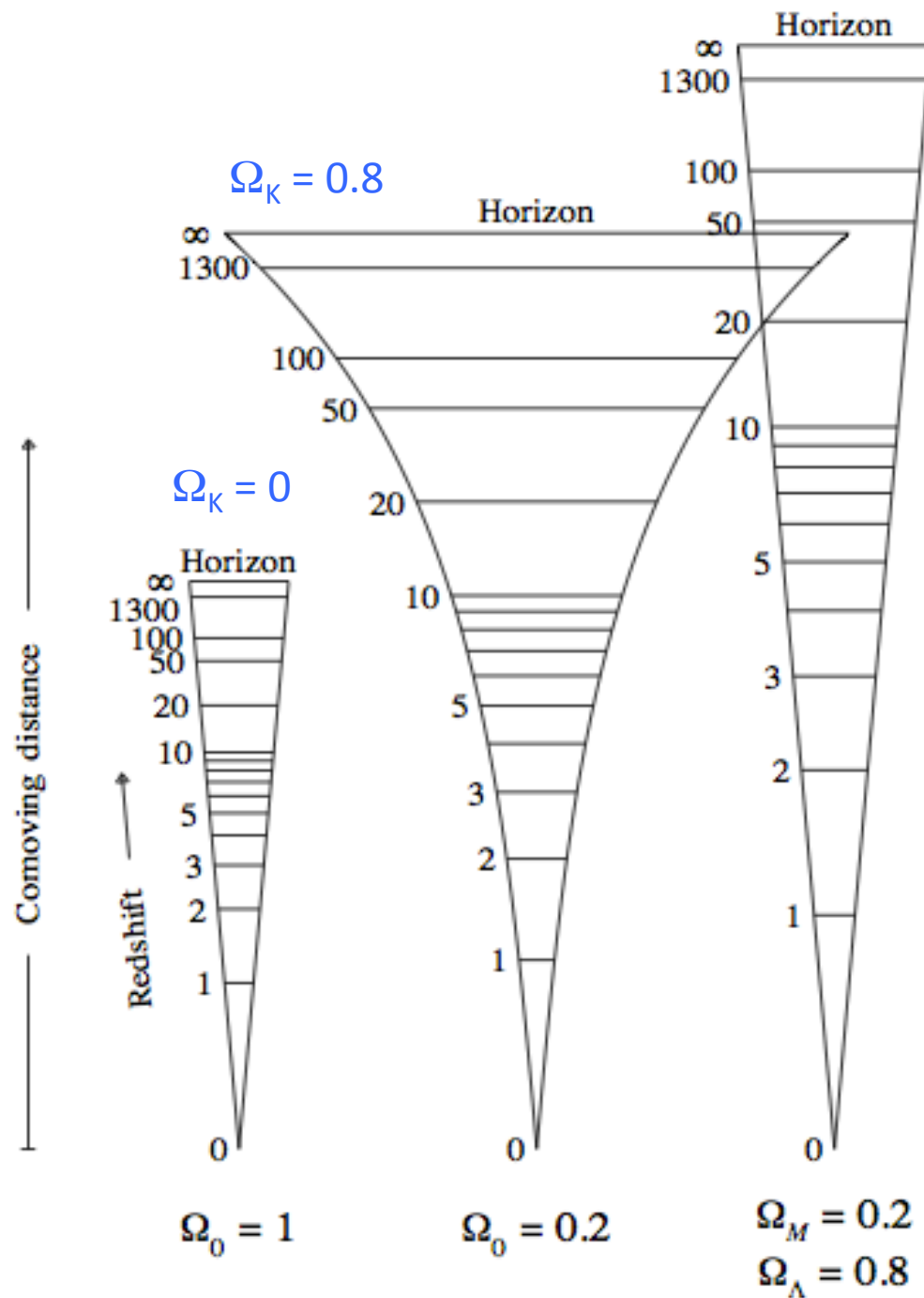
$$d_L(z) = \frac{(1+z)}{H_0 \sqrt{-\Omega_k}} \sin \left(H_0 \sqrt{-\Omega_k} \int \frac{dz'}{H(z')} \right)$$


$$\Omega_k = \frac{[H(z)D'(z)]^2 - 1}{[H_0 D(z)]^2}$$



 Dynamics

$$\chi = \int \frac{dz'}{E(z')}$$



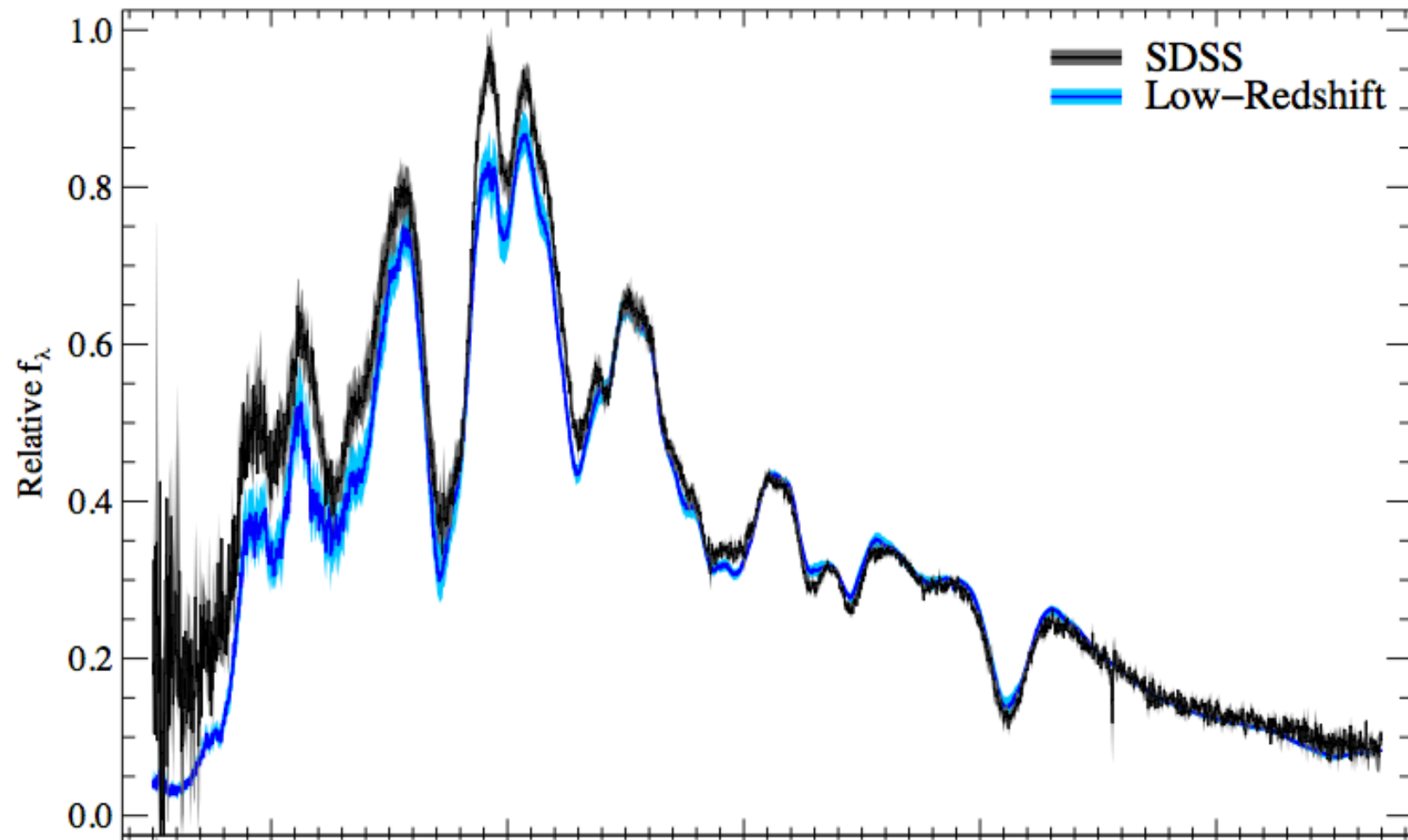
$$\Omega_K = 0.8$$

$$\Omega_K = 0$$

$$\Omega_K = 0$$

The size of an object is not uniquely determined by χ – curvature must also be known

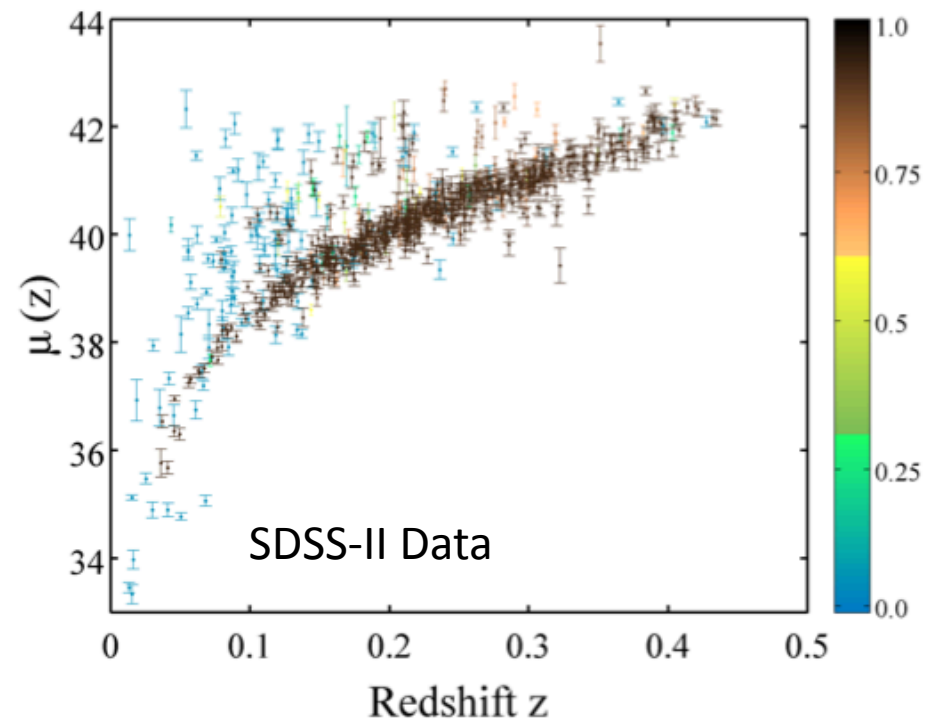
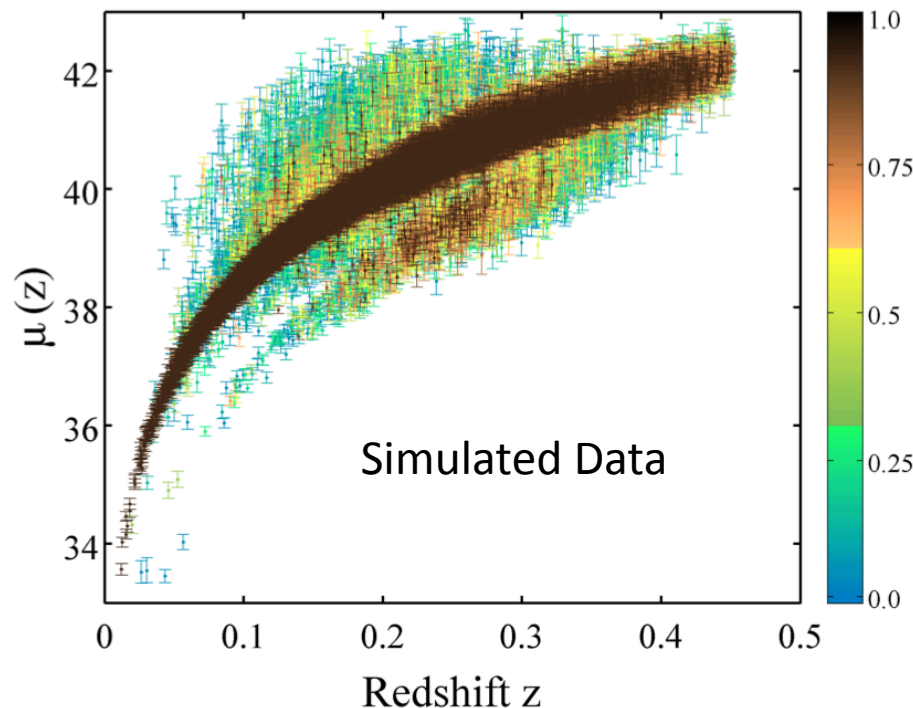
Evolution in the UV?



Foley et al. 2011

Future Surveys

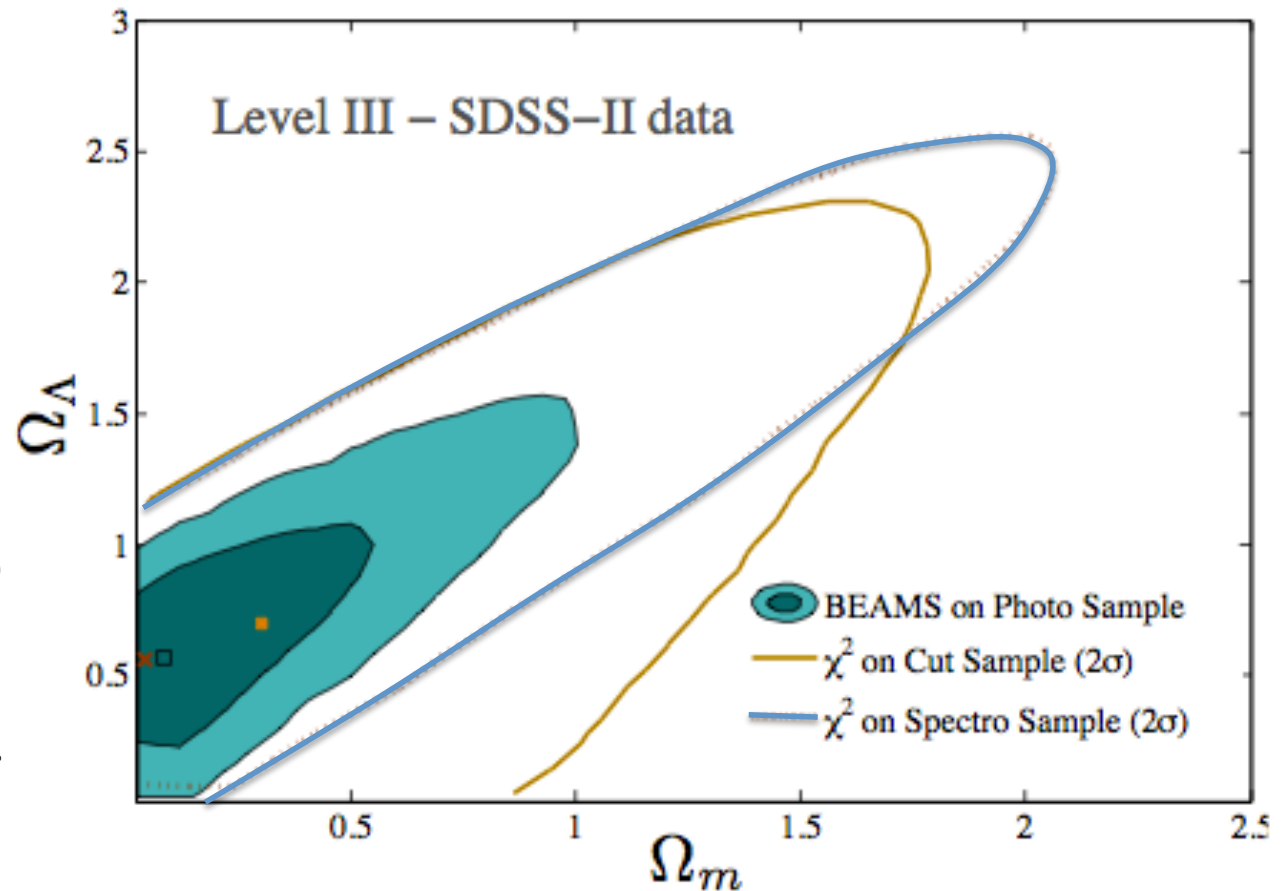
- Current surveys are already spectroscopy limited
- LSST will discover up to 10^6 SNIa...only a fraction will have spectroscopic confirmation.
- Entering the era of photometric supernova cosmology
- Contamination from non-Ias will be **unavoidable**.



Photometric Supernova Cosmology

Basic Idea: use the lightcurves to work out the probability that an object is a SNIa...

Can use Bayesian Methods (BEAMS) to do Cosmological Parameter estimation with these probabilities.



Hlozek et al
Arxiv:1111.5328

Conclusions

- SNIa are arguably still the best evidence for dark energy that we have...
- But distances are fundamentally limited by the curvature-dynamics degeneracy. Requires additional $H(z)$ data to break.



Conclusions II

- As with all methods, SNIa have a hard systematics floor that will need lots of work to break through...
- We are now entering an exciting new era of massive data and a switch to *photometric* supernova cosmology